

Transversity GPDs in the large N_c limit

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work in progress in collaboration with Christian Weiss

Overview:

- Limit of large number of colors N_c in QCD
- What is it useful for? What can it predict?
- Applications to nucleon structure
- Chiral odd GPDs
- Conclusions

quark correlator: $\mathcal{M}(\Gamma)_{\lambda'\lambda} = \int \frac{dP^+ z^-}{2\pi} e^{ixP^+ z^-} \langle p', \lambda' | \bar{\psi}_q(-\frac{z}{2}) \Gamma [-\frac{z}{2}, \frac{z}{2}] \psi_q(\frac{z}{2}) | p, \lambda \rangle \Big|_{z^+ = z_T^j = 0}$

kinematic variables: $P = \frac{1}{2}(p' + p)$, $\Delta = (p' - p)$, $\xi = -\frac{\Delta^+}{P^+}$, $t = \Delta^2$

quark GPDs

$$\mathcal{M}(\gamma^+) = \bar{u}(p', \lambda') \left[\gamma^+ H^q + \frac{i\sigma^{+\mu}\Delta_\mu}{2M_N} E^q \right] u(p, \lambda)$$

$$\mathcal{M}(\gamma^+ \gamma_5) = \bar{u}(p', \lambda') \left[\gamma^+ \gamma_5 \tilde{H}^q + \frac{i\gamma_5 \Delta^+}{2M_N} \tilde{E}^q \right] u(p, \lambda)$$

$$\begin{aligned} \mathcal{M}(i\sigma^{+j}) = \bar{u}(p', \lambda') & \left[\sigma^{+j} H_T^q + \frac{\gamma^+ \Delta^j - \Delta^+ \gamma^j}{M_N^2} E_T^q \right. \\ & \left. + \frac{P^+ \Delta^j - \Delta^+ P^j}{M_N^2} \tilde{H}_T^q + \frac{\gamma^+ P^j - P^+ \gamma^j}{M_N^2} \tilde{E}_T^q \right] u(p, \lambda) \end{aligned}$$

with $\text{GPD}^q = \text{GPD}^q(x, \xi, t; \mu^2)$

gluon GPDs analog (discussed elsewhere)

non-perturbative properties of GPDs

general constraints

- time-reversal & hermiticity: $\text{GPD}(x, -\xi, t) = \pm \text{GPD}(x, \xi, t) = \begin{cases} H, E, \tilde{H}, \tilde{E}, H_T, E_T, \tilde{H}_T \\ \tilde{E}_T \end{cases}$
- Lorentz-invariance: polynomiality $\int dx x^N \text{GPD}(x, \xi, t) = a_0(t) + a_2(t) \xi^2 + \dots + a_N \xi^N$
- inequalities (\dots , Poblitsa, Poblitsa & Kirch)

first principles

- lattice QCD (valuable information on Mellin moments, form factors)
- chiral perturbation theory (Diehl, Manashov, Schäfer; Kivel, Polyakov, Vladimirov)
- large N_c (?) \Leftarrow this talk

models

- bag (Ji et al), chiral quark soliton (Petrov et al; \dots),
light-cone approaches (Pasquini et al; Lorcé), \dots

What is large N_c limit?

observation

- in nature $N_c = 3$ experimentally pretty well established

nevertheless

- in theory limit $N_c \rightarrow \infty$ not only fascinating, but also powerful and useful tool

statement

- “ $\frac{1}{N_c}$ = only small parameter in QCD at *all* energies” (Sidney Coleman, in “Aspects of Symmetry”)

technicalities

- $g_s^2 N_c =$ fixed as limit $N_c \rightarrow \infty$ is taken 't Hooft limit
- planar diagrams dominate; in 1+1 possible to resum 't Hooft NPB 72 & 75 (1974) 461
- 3+1 not solved. But ...

conclusions in large N_c

- shown that meson & glueballs are free, stable, non-interacting particles, decay widths $\sim 1/\sqrt{N_c}$, coupling constants $\sim 1/N_c$ ('t Hooft; Veneziano; Callan et al; ... 1970s)
- mesons = pure $\bar{q}q$ states
(and not $\bar{q}q\bar{q}q$ etc, which are hard to find in nature; implying that our real-world $N_c \gg 1$?)
- Zweig rule = exact, if $N_c \rightarrow \infty$ (i.e. Zweig rule explained in the large N_c limit!?)
- baryons = solitons of meson fields Witten, NPB 160 (1979) 57
solitonic field not known, but it's symmetries are known \Leftarrow explore!

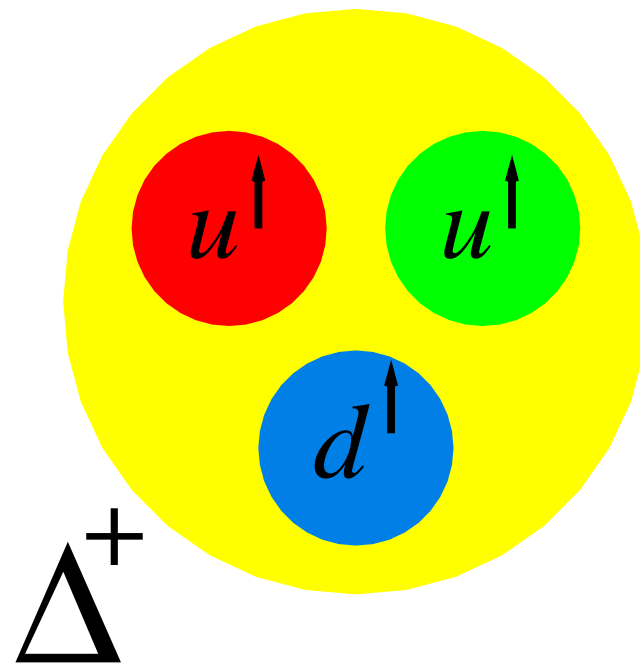
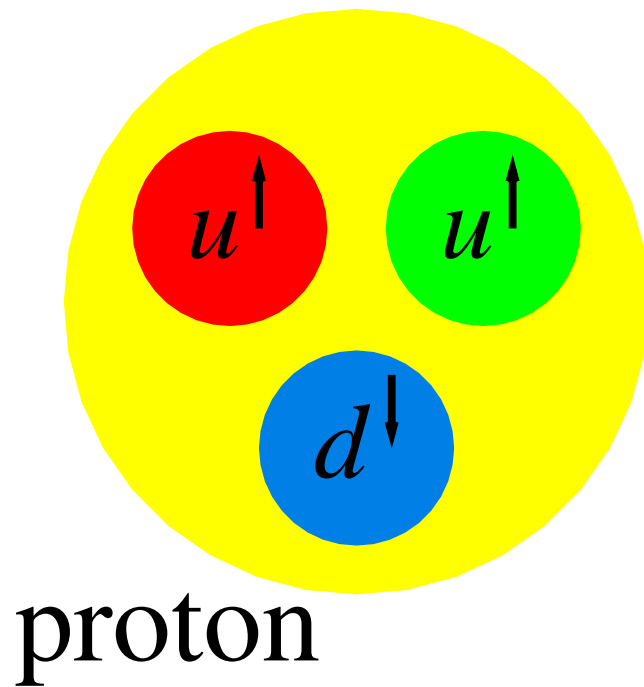
Important to keep in mind:

the mere “counting of colors N_c ” is independent of dynamics!

many generic results are the same — in QCD gauge theory, Skyrme model, quark models.

In particular: quark-model-picture provides useful intuition (and supports rigorous derivations)

for $N_c = 3$: baryon = 3q in antisymmetric color state



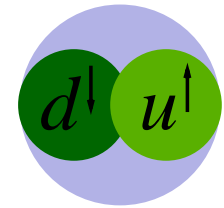
How would nucleon & Δ look for $N_c = 3, 5, 7, \dots$?

$$|\text{nucleon}\rangle = |S_N = \frac{1}{2}, T_N = \frac{1}{2}\rangle, |\Delta\rangle = |S_\Delta = \frac{3}{2}, T_\Delta = \frac{3}{2}\rangle$$

construction of nucleon wave-function in large N_c

(consideration in quark model, result generic)

must preserve spin and isospin(!) \rightarrow add light quarks in scalar-isoscalar pairs
 \rightarrow only in this way S, T unchanged!



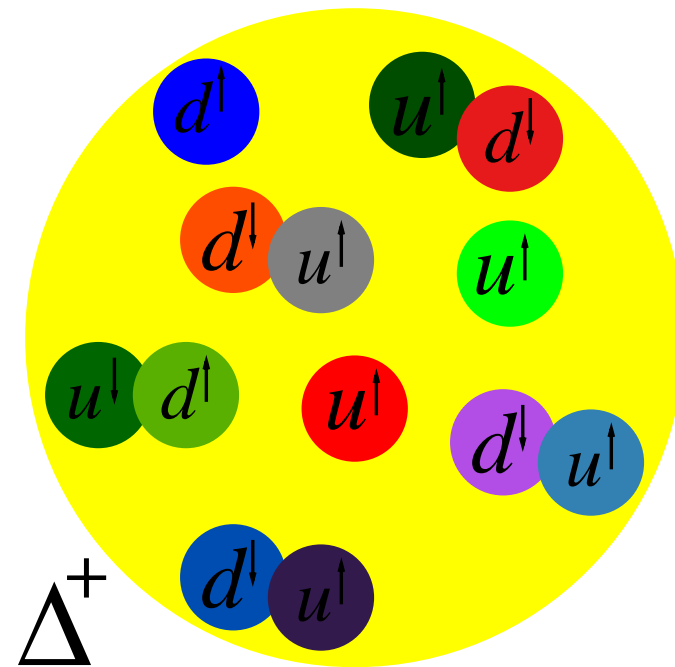
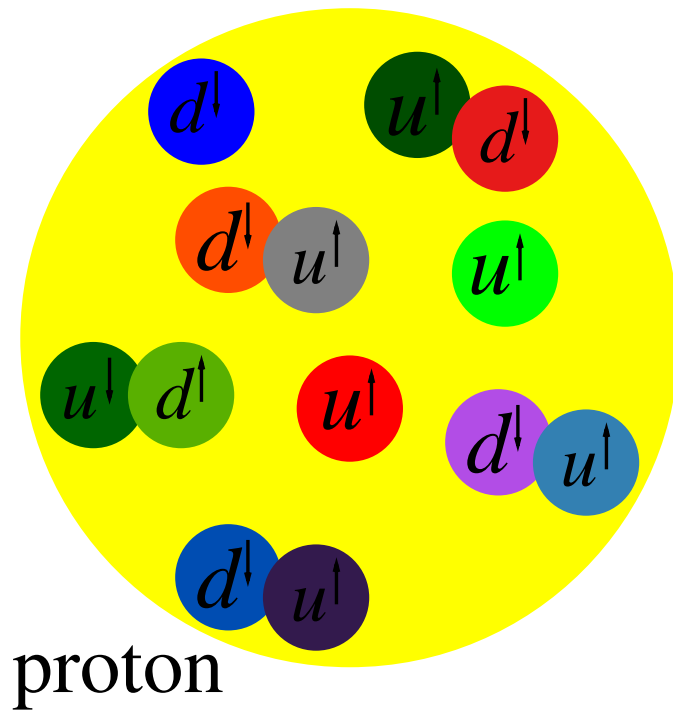
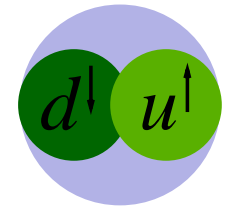
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(in quark model, but generic result)

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Observations:

- not at all surprising that we have a static mean field (“soliton”)!
- total nucleon spin = result of enormous cancellations
- if for some reason \exists tendency for spin of one flavor to be (anti-)aligned) with nucleon spin \Rightarrow the other other flavor automatically exhibits opposite polarization

(does not explain *why* there should be a correlation between flavor- and nucleon-polarization, \rightarrow need dynamics e.g. chiral dynamics)

technical description

- $\langle B', p' | \hat{O}_{\text{QCD}} | B, p \rangle = 2M_N \int d^3 X e^{i(\vec{p}-\vec{p}')\vec{X}} \int dR \phi_{B'}^*(R) O(R)_{\text{eff}} \phi_B(R) + \dots$
 $R \in \text{SU}(2)$ describes collective coordinates in 3-space & flavor-space
 $\phi_B(R)$ = rotational wave-function (defined in terms of Wigner matrices)
 $O(R)_{\text{eff}}$ is the effective mean-field function associated with the operator \hat{O}_{QCD}
- $O(R)_{\text{eff}}$ is unknown (because 3+1 QCD in large N_c not solved)
but it's symmetry properties are known!
 \Rightarrow sufficient!

results (static properties):

- baryons are massive objects $M_B \sim N_c$
- nucleon and Δ different rotational states of the *same* soliton solution
- $M_\Delta - M_N \sim \mathcal{O}(1/N_c)$
 $\rightarrow \frac{M_\Delta - M_N}{M_\Delta + M_N} \sim \mathcal{O}(1/N_c^2)$ vs in nature $\frac{M_\Delta - M_N}{M_\Delta + M_N} \sim \frac{1}{7}$ (reasonable)
- magnetic moments $\mu_p - \mu_n \sim \mathcal{O}(N_c)$ vs $\mu_p + \mu_n \sim \mathcal{O}(N_c^0)$
vs in nature $\mu_p - \mu_n = 4.71 \gg \mu_p + \mu_n = 0.88$ works!
(\exists many more examples)

partonic (dynamical) properties:

- heavy nucleon moves slowly: $\vec{p}, \vec{p}' \sim N_c^0$ i.e. $|t| \sim N_c^0 \ll M_N^2 \sim N_c^2$
- $xN_c = \text{fixed}$ in the limit $N_c \rightarrow \infty$ limit (address “non-exceptional” values of x)
- $\xi N_c = \text{fixed}$ in the limit $N_c \rightarrow \infty$ limit (address “non-exceptional” values of ξ)
- normalization $\langle B', p' | B, p \rangle = 2P^0 (2\pi)^3 \delta^{(3)}(\vec{p} - \vec{p}')$ with $P^0 \sim N_c$
- operator is color-singlet: $\bar{\psi} \dots \psi \sim N_c$
- generic expectation for correlator $\langle B', p' | \bar{\psi}_q \Gamma \psi_q | B, p, \rangle \sim N_c^2$ or (due to symmetry) less
- predictions $\text{PDF}_i(x) = N_c^{A_i} D_i(xN_c)$ with $A_i = 2$ or 1 (Diakonov et al 1996, ...)

$$(f_1^u + f_1^d)(x) = N_c^2 D_{f_1}^{u+d}(xN_c) \gg (f_1^u - f_1^d)(x) = N_c D_{f_1}^{u-d}(xN_c) \quad \checkmark$$

$$(g_1^u - g_1^d)(x) = N_c^2 D_{g_1}^{u-d}(xN_c) \gg (g_1^u + g_1^d)(x) = N_c D_{g_1}^{u+d}(xN_c) \quad \checkmark$$

$$(h_1^u - h_1^d)(x) = N_c^2 D_{h_1}^{u-1}(xN_c) \gg (h_1^u + h_1^d)(x) = N_c D_{h_1}^{u+d}(xN_c) \quad \checkmark (\text{prediction!})$$

helpful predictions:

- analog predictions for $\text{TMD}_i(x, p_T) = N_c^{Bi} D_i(x N_c, p_T)$ (Pobylitsa hep-ph/0301236)
- Sivers function $(u - d) \gg (u + d)$ ✓ (used in early fits PLB 612 (2005) 233, PRD 73 (2006) 014021)
- other TMDs to be tested
- also for gluon distributions large N_c counting possible

$$\Delta G(x)/G(x) = \mathcal{O}(1/N_c) \text{ Efremov, Goeke, Pobylitsa PLB 488 (2000) 182}$$

$$(\text{gluon Sivers})/(\text{quark Sivers}) = \mathcal{O}(1/N_c) \text{ Efremov et al, PLB 612 (2005) 233}$$

- extremely helpful, if we do not know anything otherwise!

Remark: Sivers function $\sim N_c^{2+1}$

while $f_1 \sim N_c^2$ (?), does Sivers asymmetry grow with N_c ?

No! Trivial additional power of N_c due to correlatorⁱ $\sim \frac{p_T^i}{M_N} f_{1T}^\perp(x, p_T)$

correlators of PDFs, TMDs, GPDs $\sim N_c^2$ at most (same story with GPDs below)

chiral-even GPDs (Petrov et al 1998; Pentinnen et al; Goeke, Polyakov, Vanderhaeghen)

- $(H^u + H^d)(x, \xi, t) = N_c^2 D_H^{u+d}(xN_c, \xi N_c, t) \gg (u - d) = \text{smaller}$
- $(E^u - E^d)(x, \xi, t) = N_c^3 D_E^{u-d}(xN_c, \xi N_c, t) \gg (u + d) = \text{smaller}$
- $(\tilde{H}^u - \tilde{H}^d)(x, \xi, t) = N_c^2 D_{\tilde{H}}^{u-d}(xN_c, \xi N_c, t) \gg (u + d) = \text{smaller}$
- $(\tilde{E}^u - \tilde{E}^d)(x, \xi, t) = N_c^4 D_{\tilde{E}}^{u-d}(xN_c, \xi N_c, t) \gg (u + d) = \text{smaller}$

predictions being tested (so far, so good)

chiral-odd GPDs

- marginally mentioned in literature (works & proceedings by Matthias Burkardt)
- no systematic study so far ← **ongoing study**

chiral-odd GPDs (first results of ongoing study)

- $(H_T^u - H_T^d)(x, \xi, t) = N_c^2 D_{H_T}^{u-d}(xN_c, \xi N_c, t) \gg (u + d) = \text{smaller}$ (cf. transversity ✓)
- $(E_T^u - E_T^d)(x, \xi, t) = N_c^3 D_{E_T}^{u-d}(xN_c, \xi N_c, t) \gg (u + d) = \text{smaller}$
- $(\tilde{H}_T^u + \tilde{H}_T^d)(x, \xi, t) = N_c^3 D_{\tilde{H}_T}^{u-d}(xN_c, \xi N_c, t) \gg (u - d) = \text{smaller}$
- $(\tilde{E}_T^u - \tilde{E}_T^d)(x, \xi, t) = N_c^3 D_{\tilde{E}_T}^{u-d}(xN_c, \xi N_c, t) \gg (u + d) = \text{smaller}$

compatible with lattice, models

experiment (?) → JLab, PRC 90 (2014) 039901 + preliminary, see talk by V. Kubarovsky

data indicate (important: at small $|t| \sim N_c^0 \ll M_N^2 \sim N_c^2$)

(i) $\langle H_T^u \rangle \approx - \langle H_T^d \rangle$

compatible with large N_c ✓ (work in progress)

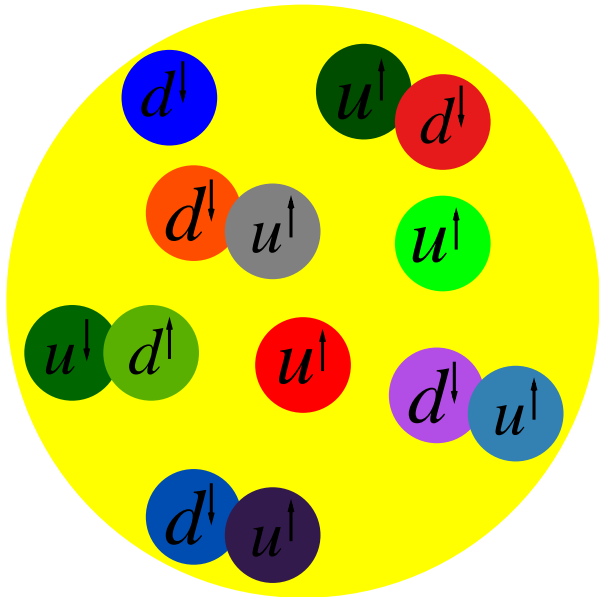
(ii) $\langle \bar{E}_T^u \rangle \approx \langle \bar{E}_T^d \rangle$ where $\bar{E}_T = 2\tilde{H}_T + E_T$

in large N_c we found $\tilde{H}_T^u \approx \tilde{H}_T^d$ ✓ but opposite for E_T (?)
is E_T small (?) need studies in dynamical approaches

- in any case: same-sign flavors for chiral-odd GPDs must involve \tilde{H}_T (and does! ✓)
- and it even enters with factor 2 enhancement (this is numerically almost factor $N_c \dots$)
- overall claim (with all reservations): it's encouraging! ✓
- more work needed! (much motivated!!)

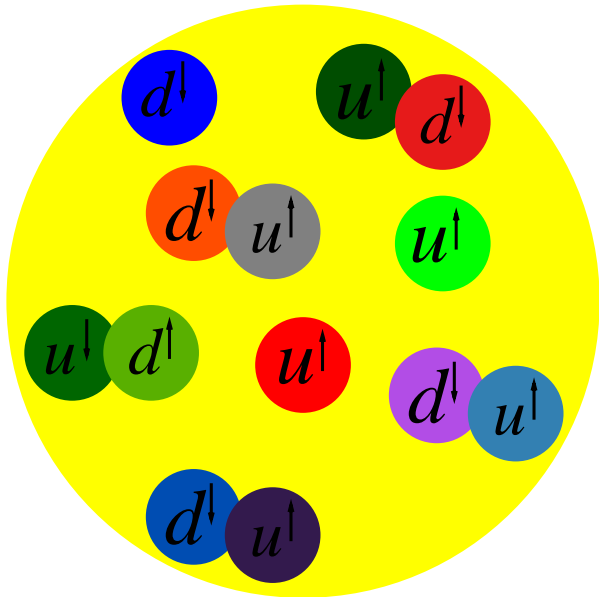
Conclusions

- **large N_c** : powerful, model-independent theoretical tool
- it provides insights & guidelines, **works** phenomenologically
- report from study in progress of **chiral odd GPDs** in large N_c
- first results are in **encouraging agreement** with first data
- more work needed, underway, impact of **chiral dynamics**
- looking forward to **exciting results**



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Thank you !!!