

Final State Nucleons in Neutrino-Nucleus Interactions

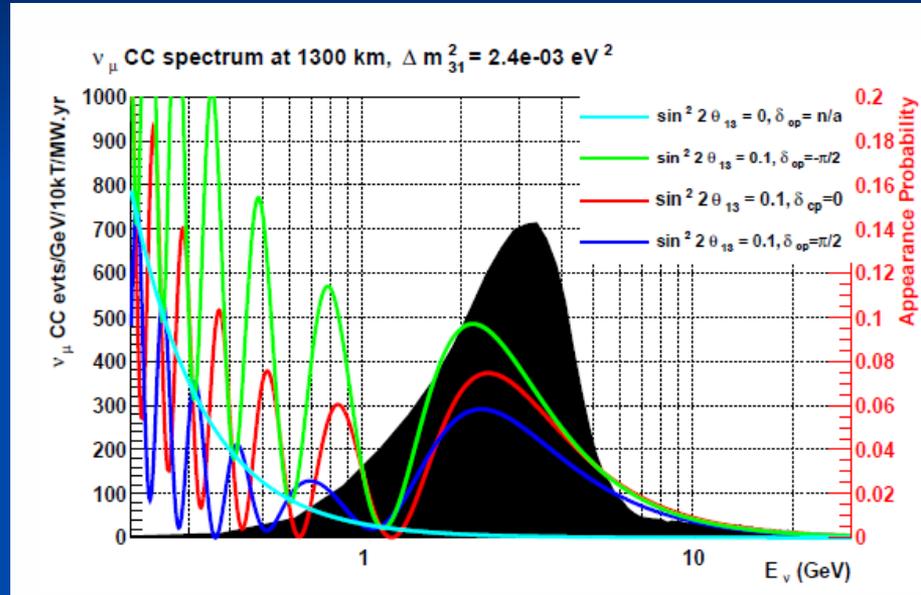
Ulrich Mosel



**Institut für
Theoretische Physik**



DUNE, δ_{CP} Sensitivity



Appearance probability:
 $P_{\mu \rightarrow e}$

Need energy to distinguish
 between different δ_{CP}

Need to know neutrino energy to better than about 100 MeV

FSI and Transport Theory

- Need to reconstruct energy from final state observables
- All modern experiments use nuclear targets

Need to describe

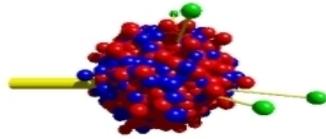
$\nu + A \rightarrow X + \text{many hadrons}$



Nuclear Theory

- Necessary Ingredients
 - Nuclear groundstate (correlations, spectral functions)
 - Nuclear reaction mechanisms (IA vs. 2p2h, coll. excit.)
 - Electroweak interaction vertices, in medium
 - Particle production, also in secondary colls
 - *Propagation of all particles to final state, incl fsi*
- Only capable method: Transport Theory, guidance from QGP generators (MC is poor man's transport theory)

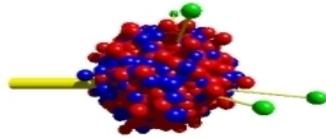




- **GiBUU** describes (within the same unified theory and code)
 - heavy ion reactions, particle production and flow
 - pion and proton induced reactions
 - low and high energy photon and electron induced reactions
 - **neutrino induced reactions**

using the same physics input! And the same code!
Most widely tested generator available





- **GiBUU : Theory and Event Generator**
based on a BM solution of Kadanoff-Baym equation
- Code from *gibuu.hepforge.org*
- Physics content and details of implementation in:
Buss et al, Phys. Rept. 512 (2012) 1- 124
Mine of information on theoretical treatment of potentials, collision terms, spectral functions and cross sections, useful for any generator



Transport Equation

Collision term

$$\mathcal{D}F(x, p) + \text{tr} \left\{ \text{Re} \tilde{S}^{\text{ret}}(x, p), -i \tilde{\Sigma}^<(x, p) \right\}_{\text{pb}} = C(x, p).$$

Drift term

$$\left[\left(1 - \frac{\partial H}{\partial p_0} \right) \frac{\partial}{\partial t} + \frac{\partial H}{\partial \mathbf{p}} \frac{\partial}{\partial \mathbf{x}} - \frac{\partial H}{\partial \mathbf{x}} \frac{\partial}{\partial \mathbf{p}} + \frac{\partial H}{\partial t} \frac{\partial}{\partial p^0} + \text{KB term} \right] F(x, p) = - \text{loss term} + \text{gain term}$$

$$F(x, p) = 2\pi g f(x, p) A(x, p).$$

Spectral function

GiBUU Ingredients

- In-medium corrected primary interaction cross sections for QE, 2p2h, N^* resonances, DIS, boosted to rest frame of bound nucleon, moving in *bound local Fermigas*
- Includes **spectral functions for baryons and mesons** (binding + collision broadening)
- **Hadronic couplings** for FSI taken from PDG
- **Vector couplings** taken from electro-production (MAID)
- **Axial couplings** modeled with PCAC and dipole FF



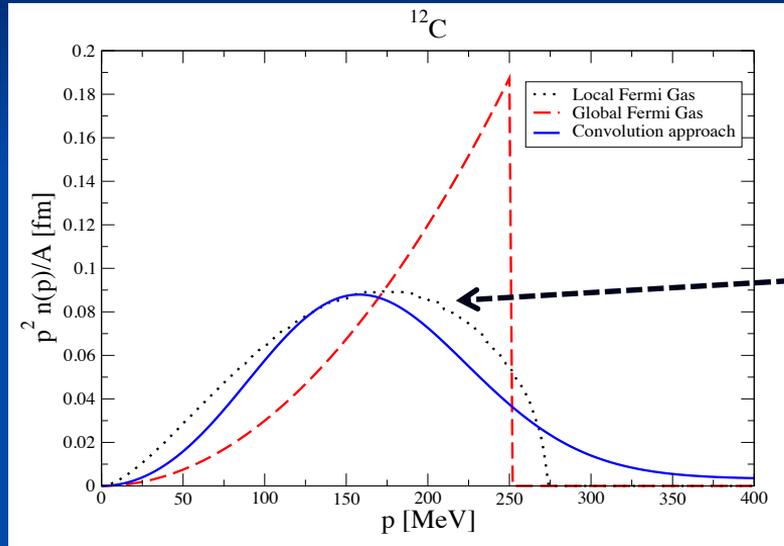
GiBUU Output

- GiBUU produces

1. Full event file: four-vectors of all particles in final state
2. Hundreds of cross sections, directly in Gnuplotable format: energy and angle differential
3. Energy and Q^2 reconstruction files incl migration files



Momentum Distribution GiBUU



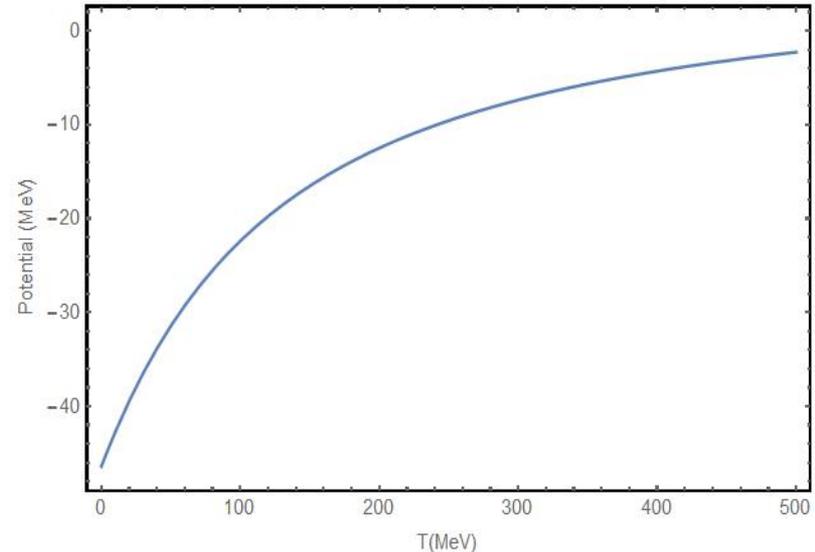
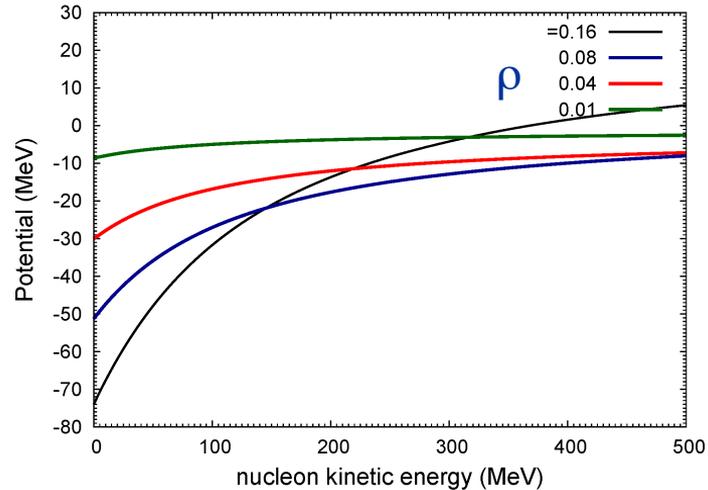
GiBUU uses
Local Fermi Gas
Energy-distribution smooth
because of r-dependent
potential

$$P_h(\vec{p}, E) \propto \int d^3r [\Theta(p_F(\vec{r}) - p) \delta(E + T_p + V(\vec{r}, \vec{p}))]$$

FSI Potential

density and momentum dependent

density-averaged potential for ^{12}C



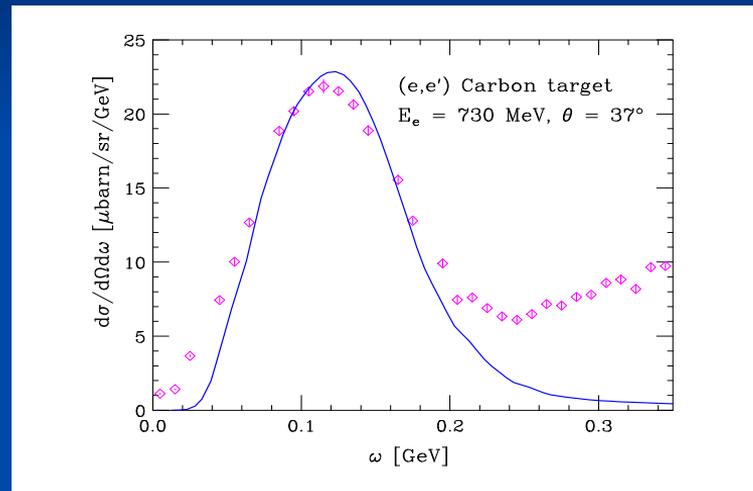
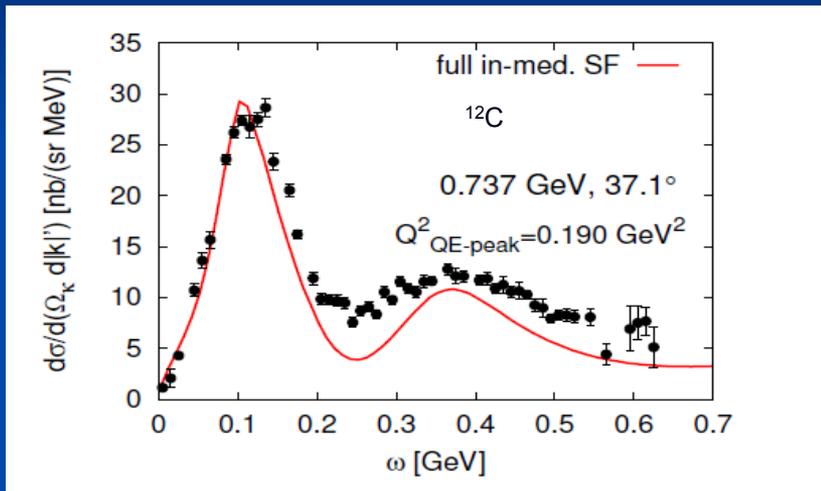
Consistent with binding potential for all nucleons, also for fsi

Test with electron data



Electrons as Benchmark for GiBUU

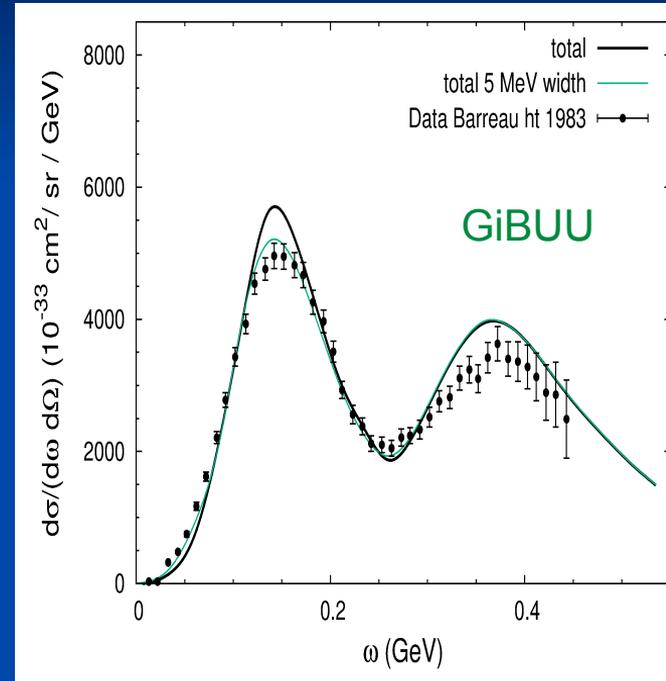
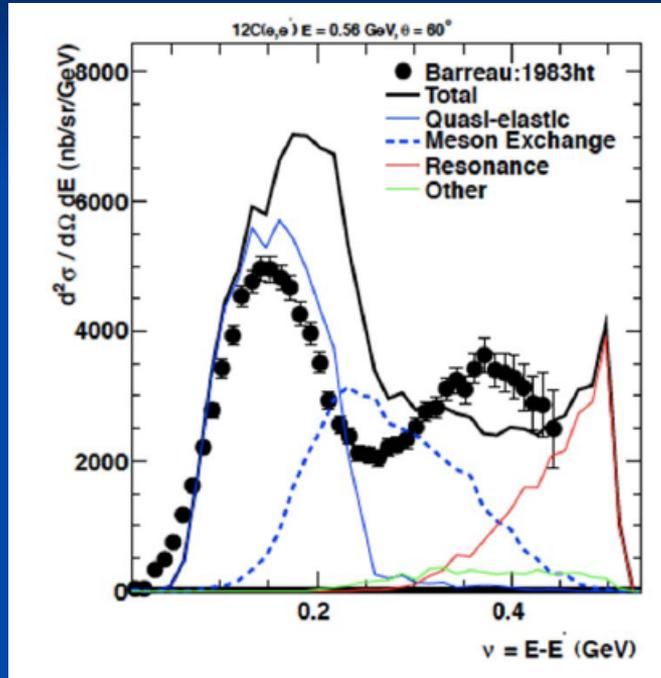
Trouble for neutrinos: ω must be reconstructed



No free parameters!
no 2p-2h, contributes
in dip region and under Δ

O. Benhar, spectral fctn

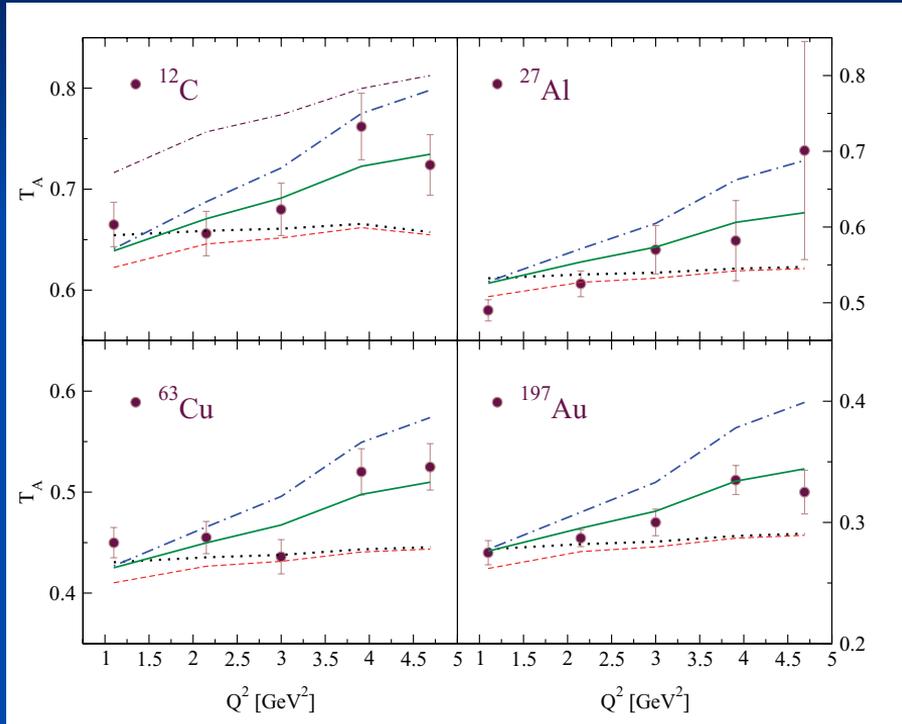
GENIE vs GiBUU



GENIE, from S.Dytman, BNL meet, Febr. 2015



JLAB Pion Production



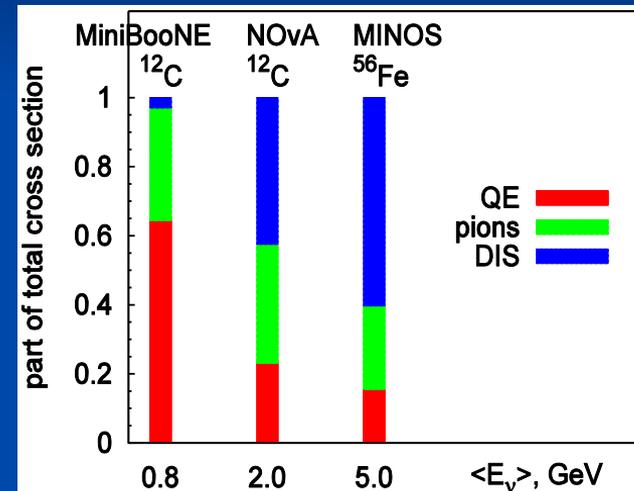
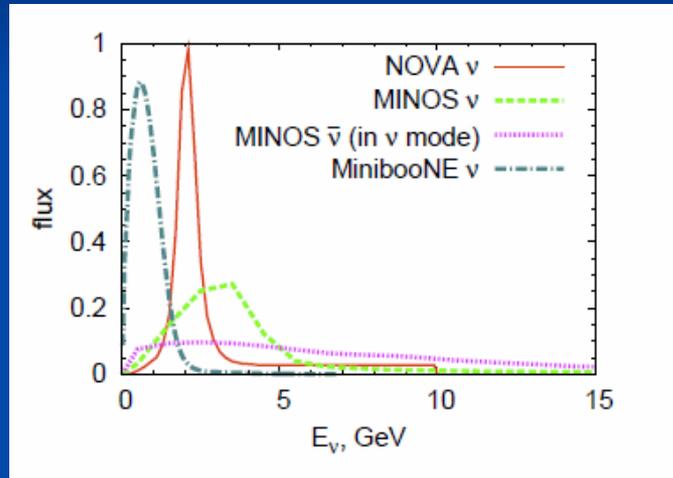
Exp: B. Clasie et al.
Phys. Rev. Lett. 99, 242502 (2007)

GiBUU: Kaskulov et al,
Phys.Rev. C79 (2009) 015207



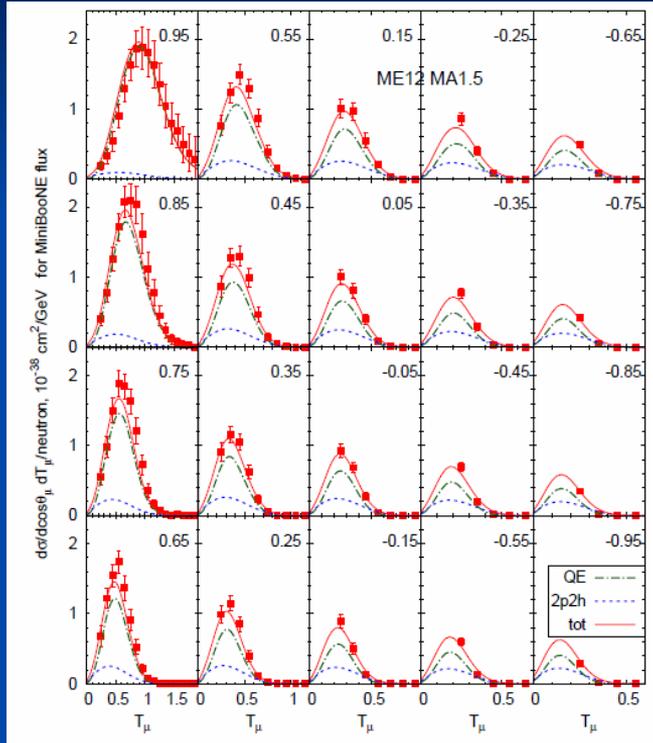
Neutrino Beams

- Neutrinos do not have fixed energy nor just one reaction mechanism



Have to reconstruct energy from final state of reaction
Different processes are entangled

The MiniBooNE QE Puzzle Explanations



ME12, MB flux averaged

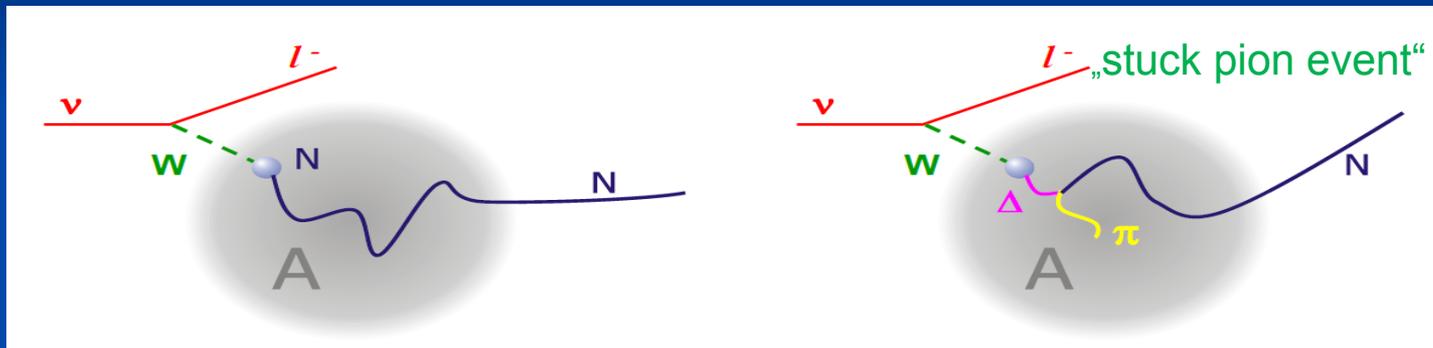
Data corrected
for stuck-pion events!

$$W^{\mu\nu} \sim P_T^{\mu\nu}(q) F(Q^2), \text{ educated guess}$$

Inclusive double-differential
X-sections fairly insensitive to
details of interaction



Final State Nucleons from Nuclear Targets

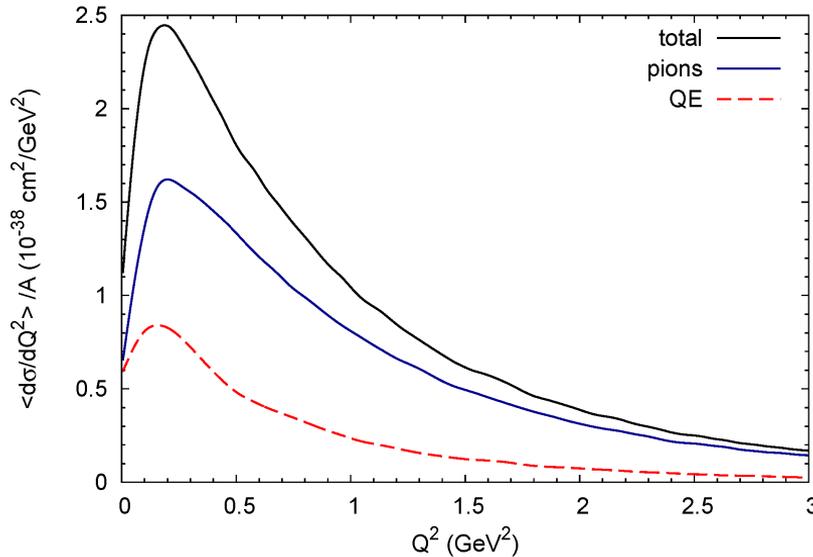


Complication to identify QE, entangled with π production

QE data can never be better than the pion subtraction method

QE vs. Pion Production at DUNE

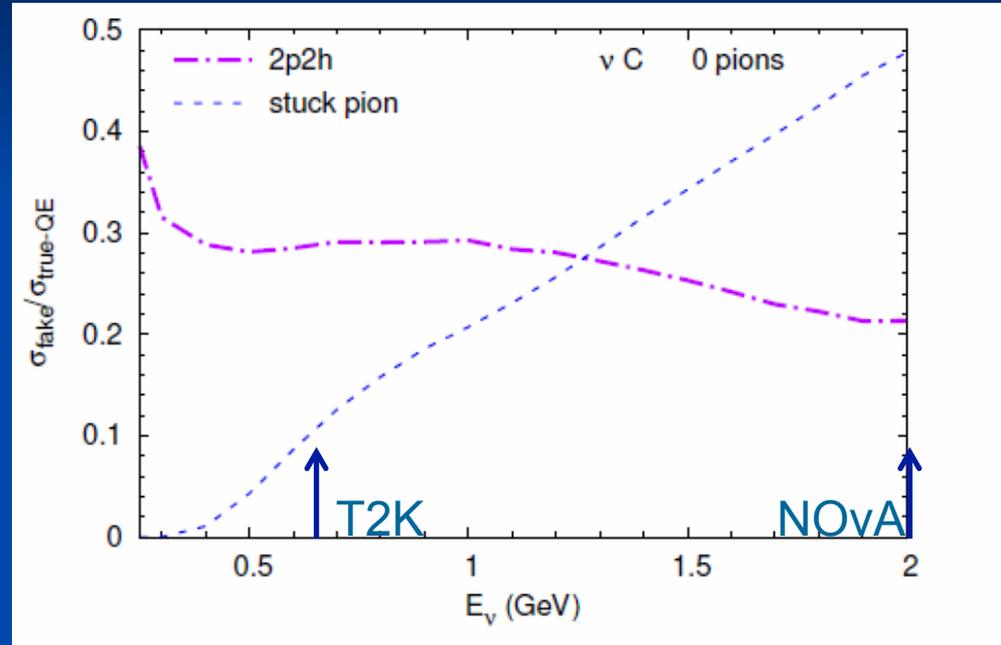
Target:
 ^{40}Ar



Pions: Resonance + DIS
QE: 'true' QE + 2p2h

QE \cong 1/3 total
Pions \cong 2/3 total

2p-2h vs. ‚stuck pions‘



Neutrino QE Data

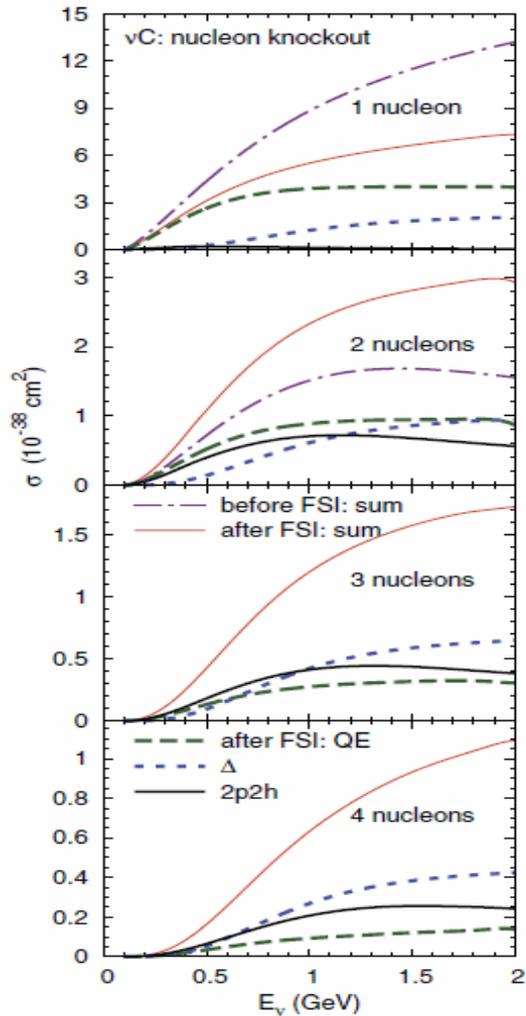
- Neutrino QE Data can be obtained from experiment only if pion production is under control
- Neutrino QE Data can never be better than the description of neutrino-induced pion production



Effects of FSI on nucleon multiplicities

Multiplicities do not distinguish between primary reaction mechanisms:

- 2 nucleons come to equal parts from
- True (1-body) QE
 - 2p2h
 - Delta ($\Delta + N \rightarrow N + N$)



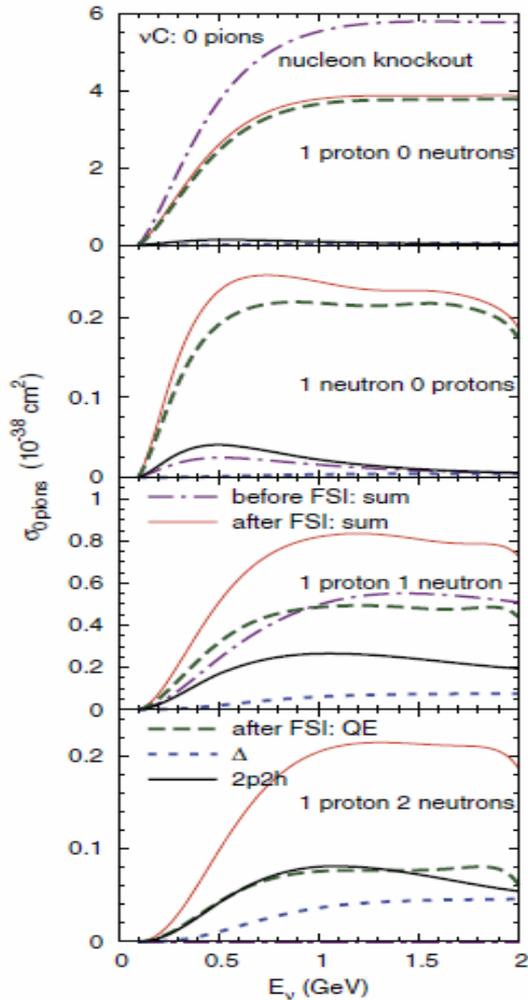
QE + 2p2h Identification with 0 pion events

1p Xn X π : fairly clean QE event

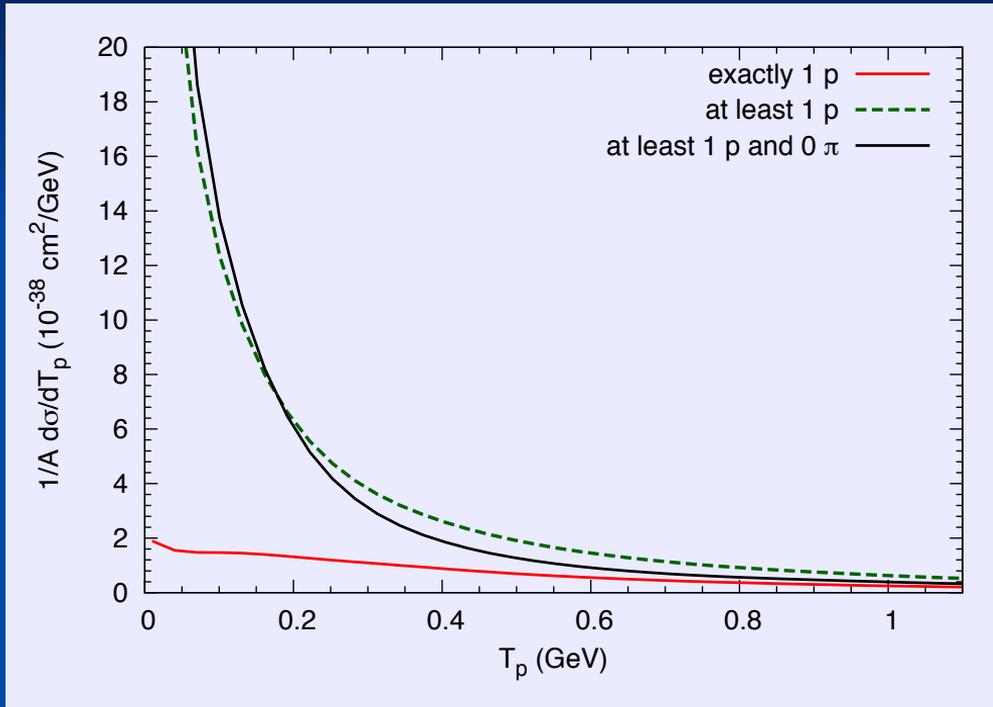
1p 0n 0 π : very clean QE event

No clean signal for 2p-2h
because of FSI

True QE can be rather well
identified by 1p, Xn, 0pi
exploited by MINERvA



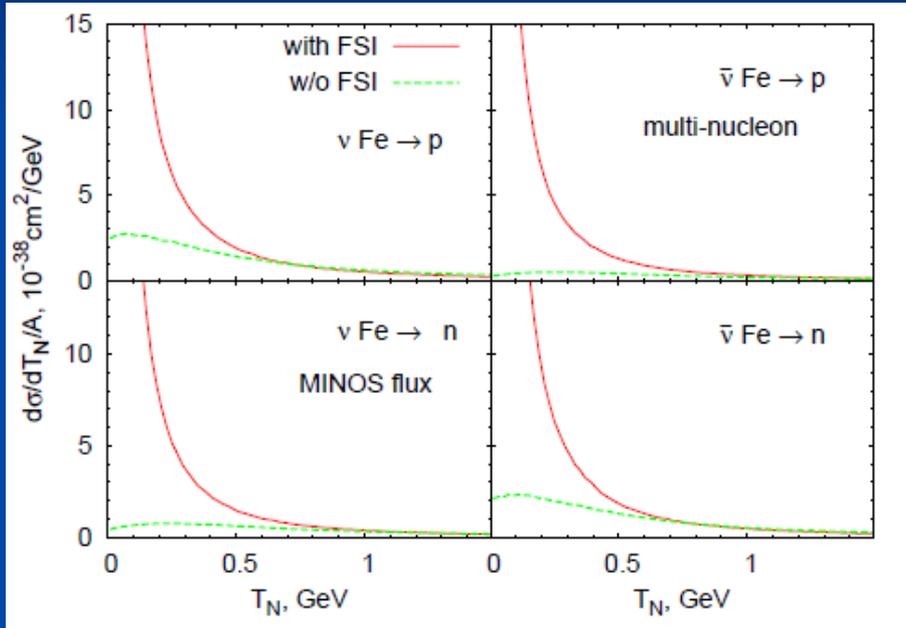
Protons at MINERvA



Strong fsi effect

Experiments at higher energies

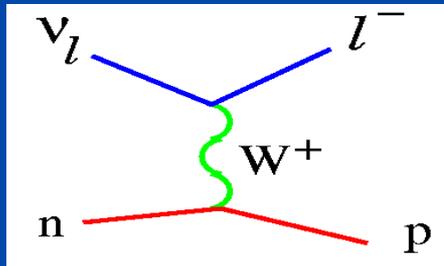
Knock-out Nucleons



Strong rise below 300 MeV,
dominated by FSI:
Avalanche Effect: one initial
nucleon kicks out many others

Energy Reconstruction by QE

- In QE scattering on nucleon at rest, only $l + p$, 0π , is outgoing. lepton determines neutrino energy:

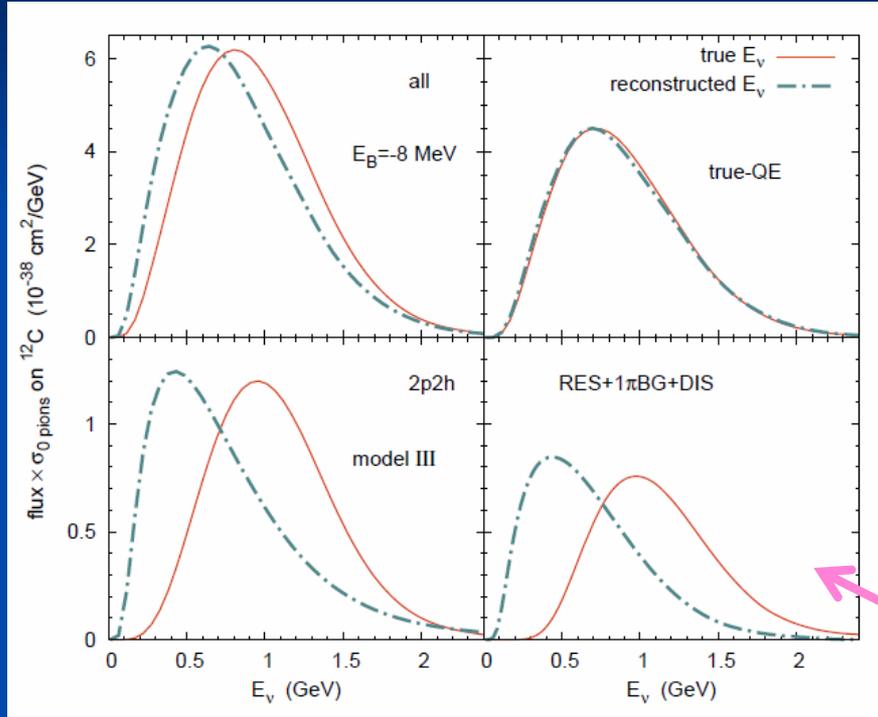


$$E_\nu = \frac{2M_N E_\mu - m_\mu^2}{2(M_N - E_\mu + p_\mu \cos \theta_\mu)}$$

- **Trouble:** all presently running expts use nuclear targets
 1. Nucleons are Fermi-moving
 2. Final state interactions may hinder correct event identification

Energy reconstruction in MB

Event rates = flux x crosssection



Reconstructed energy shifted to lower energies for all processes beyond QE

Reconstruction must be done for 0 pion events

Not only 2p-2h important

NOT contained in Nieves model

MiniBooNE flux

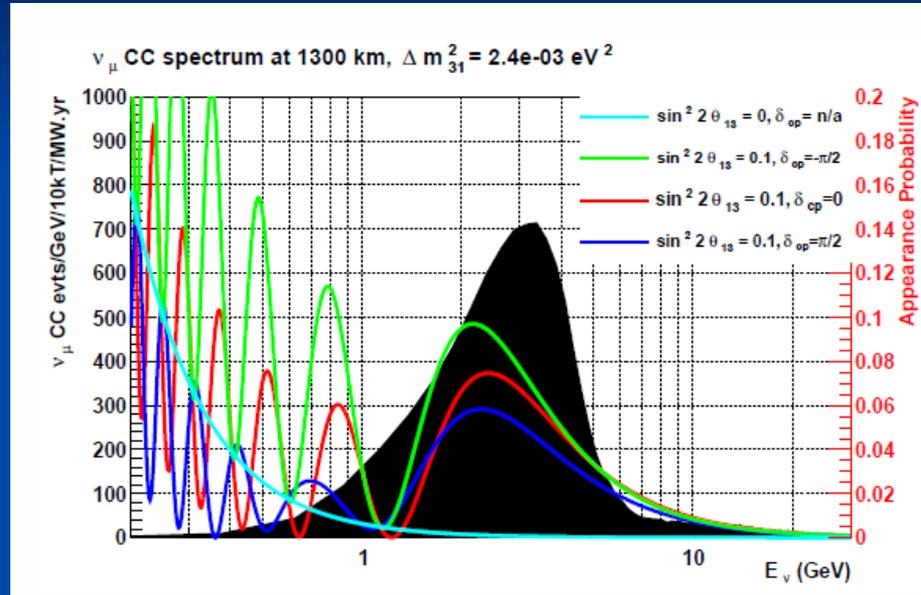
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DUNE, δ_{CP} Sensitivity

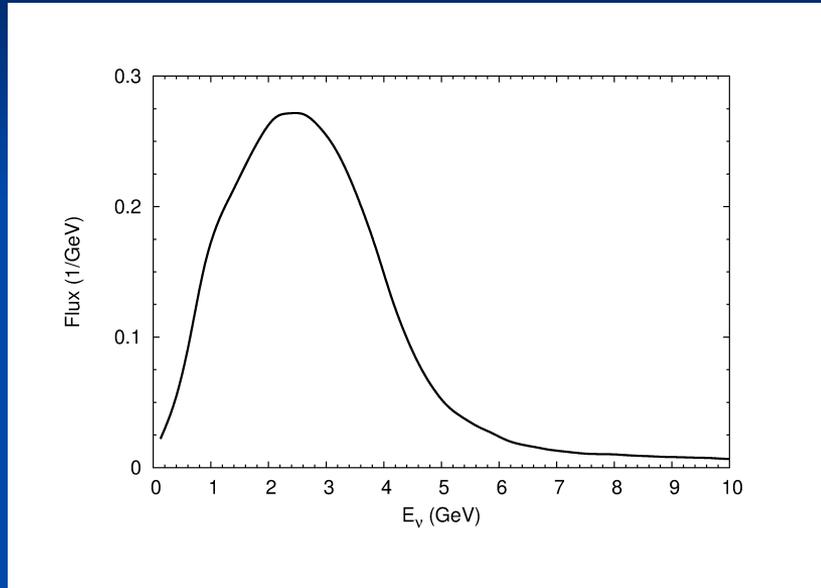


Appearance probability:
 $P_{\mu \rightarrow e}$

Need energy to distinguish
between different δ_{CP}

Need to know neutrino energy to better than about 100 MeV

Energy Reconstruction for DUNE

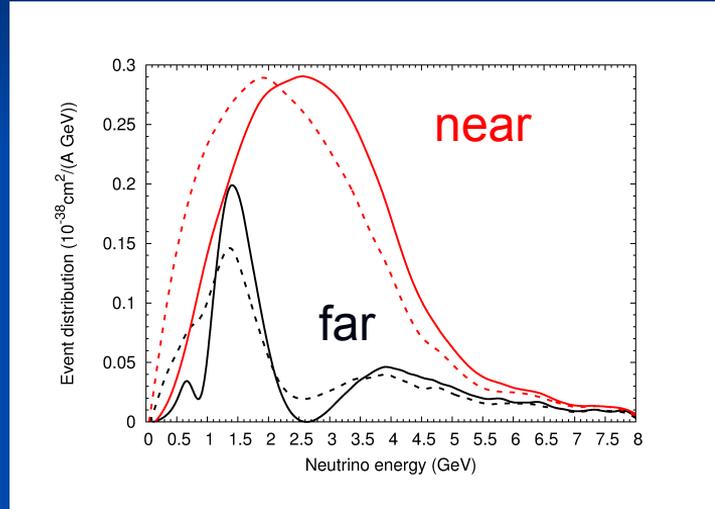


Flux of incoming neutrinos (energy distribution)

QE Energy Reconstruction for DUNE

Muon survival in 0 pion sample

Target: ^{40}Ar



Dashed: reconstructed,
solid: true energy

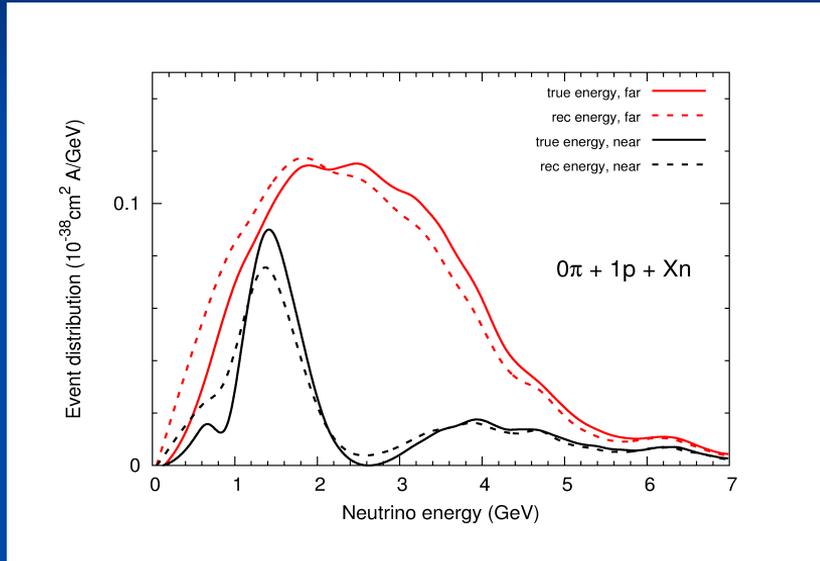
All calculations from GiBUU

Mosel et al.,
Phys.Rev.Lett. 112 (2014) 151802

Nearly 500 MeV difference between true and reconstructed event distributions → not a useful method

QE Energy Reconstruction for DUNE

Muon survival in $0\pi + 1p + Xn$ sample

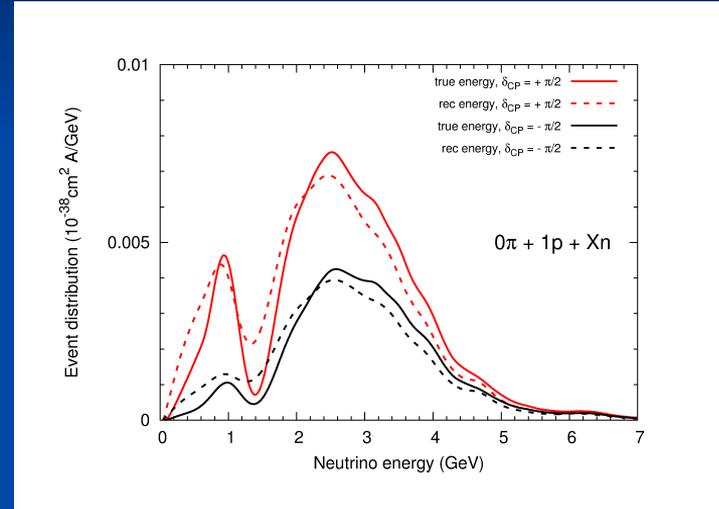
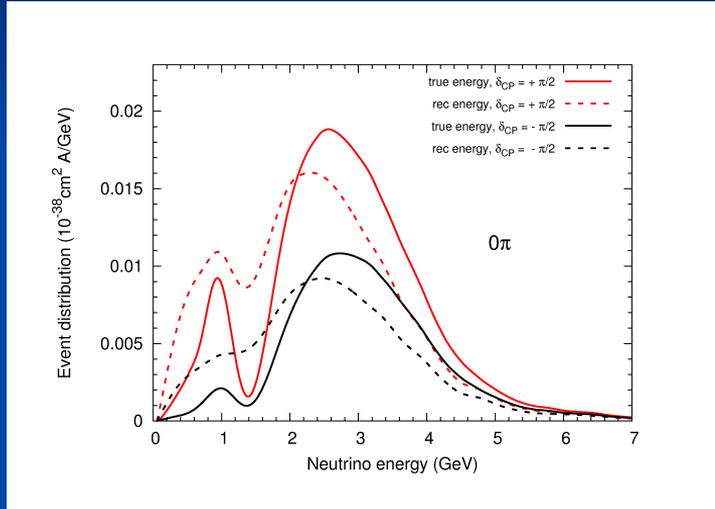


Dashed: reconstructed,
solid: true energy

Dramatic improvement in $0\pi, 1p, Xn$ sample, down by only factor 3

DUNE e-appearance

Sensitivity to δ_{CP}



Dramatic improvement in 0π , $1p$, Xn sample, down by only factor 3

Summary

- QE scattering and pion production are always entangled in neutrino reactions on nuclei
- QE data can never be better than description of pion production
- Full event analysis is needed to help to identify QE
- Energy reconstruction (and neutrino oscillation parameter extraction) depends on a correct modeling of the final state



Backups



Theoretical Basis of GiBUU

Simplicity

- Kadanoff-Baym equation (1960s)
 - full equation can not be solved yet
 - not (yet) feasible for real world problems
- Boltzmann-Uehling-Uhlenbeck (BUU) models
 - Boltzmann equation as gradient expansion of Kadanoff-Baym equations, in Botermans-Malfliet representation (1990s): **GiBUU**
- Cascade models (typical event generators, NUANCE, GENIE, NEUT,..)
 - no mean-fields, primary interactions and FSI not consistent



Factorization of GiBUU

- GiBUU is approximately factorized into
 - initial, first interaction
 - final state interactions (2nd, 3rd, ... coll.)

- Particular strength: FSI treatment

- Detailed infos from

Buss et al, Phys. Rept. 512 (2012) 1- 124

and website

gibuu.hepforge.org



Theoretical Basis: GiBUU

Time evolution of spectral phase space density (for $i = N, \Delta, \pi, \rho, \dots$) given by KB equation in Botermans-Malfliet form:

$$\left[\left(1 - \frac{\partial H}{\partial p_0} \right) \frac{\partial}{\partial t} + \frac{\partial H}{\partial p} \frac{\partial}{\partial x} - \frac{\partial H}{\partial x} \frac{\partial}{\partial p} + \frac{\partial H}{\partial t} \frac{\partial}{\partial p_0} \right] F_i(x, p) = C[F_i(x, p), F_j(x, p)]$$

Hamiltonian H includes off-shell propagation correction and potentials

8D-Spectral phase space density

Collision term

Off shell transport of collision-broadened hadrons included with proper asymptotic free spectral functions



Transport Equation

- Kadanoff-Baym equation for space-time development of one particle spectral phase space density F after gradient expansion in Wigner repres.:

$$\mathcal{D}F(\boldsymbol{x}, p) + \text{tr} \left\{ \text{Re} \tilde{S}^{\text{ret}}(\boldsymbol{x}, p), -i \tilde{\Sigma}^<(\boldsymbol{x}, p) \right\}_{\text{pb}} = C(\boldsymbol{x}, p).$$

F = spectral phase-space density: $F(\boldsymbol{x}, p) = -2f(\boldsymbol{x}, p) \text{tr}[\text{Im}(\tilde{S}^{\text{ret}}(\boldsymbol{x}, p))\gamma^0]$,

$$\mathcal{D}F = \{p_0 - H, F\}_{\text{pb}} \quad \text{with} \quad H = E^*(\boldsymbol{x}, p) - \text{Re} \tilde{\Sigma}_V^0(\boldsymbol{x}, p).$$

Collision term

$$\begin{aligned} C^{(2)}(x, p_1) &= C_{\text{gain}}^{(2)}(x, p_1) - C_{\text{loss}}^{(2)}(x, p_1) = \frac{\mathcal{S}_{1'2'}}{2p_1^0 g_{1'} g_{2'}} \int \frac{d^4 p_2}{(2\pi)^4 2p_2^0} \int \frac{d^4 p_{1'}}{(2\pi)^4 2p_{1'}^0} \int \frac{d^4 p_{2'}}{(2\pi)^4 2p_{2'}^0} \\ &\times (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_{1'} - p_{2'}) |\overline{\mathcal{M}}_{12 \rightarrow 1'2'}|^2 [F_{1'}(x, p_{1'}) F_{2'}(x, p_{2'}) \overline{F}_1(x, p_1) \\ &\times \overline{F}_2(x, p_2) - F_1(x, p_1) F_2(x, p_2) \overline{F}_{1'}(x, p_{1'}) \overline{F}_{2'}(x, p_{2'})] \end{aligned}$$

with

$$\begin{aligned} F(x, p) &= 2\pi g A(x, p) f(x, p) \\ \overline{F}(x, p) &= 2\pi g A(x, p) [1 - f(x, p)] \end{aligned}$$



GiBUU: numerical implementation

- Hadrons feel potentials (nuclear + Coulomb): essential for nucleon spectral functs
- Wigner functions represented by testparticles (100 – 1000 per nucleon). Off-shell transport of hadrons (spectral functions) with proper asymptotics
- Fully relativistic treatment of interaction vertices and propagation. Collision criterion respects relativity (as far as possible), not just $\sigma = \pi r^2$ prescription.



Practical Basis: GiBUU

- one transport equation for each particle species (61 baryons, 21 mesons)
- coupled through the potential in H and the collision integral C
- $W < 2.5$ GeV: Cross sections from resonance model (PDG and MAID couplings), consistent with electronuclear physics
- $W > 2.5$ GeV: particle production through string fragmentation (PYTHIA)



GiBUU Ingredients

- Various options in code are controlled by extended job card with all relevant switches.
- Output are (many) cross sections, reconstructed and true event distributions, full final state with four-vectors of all particles in Les Houches or ROOT format.
- Website *gibuu.hepforge.org* contains extensive documentation for code and explanation of output



GiBUU

- Code can be obtained from gibuu.hepforge.org
- Inclusive X-section needs only initial interaction:
Running time \approx 1 hour on PC
- Full event (incl all semi-incl. X-sections) needs running time of order weeks for reasonable statistics, statistics can be obtained by several parallel runs
- Code is open source, users are encouraged to find bugs, improve code, implement new features



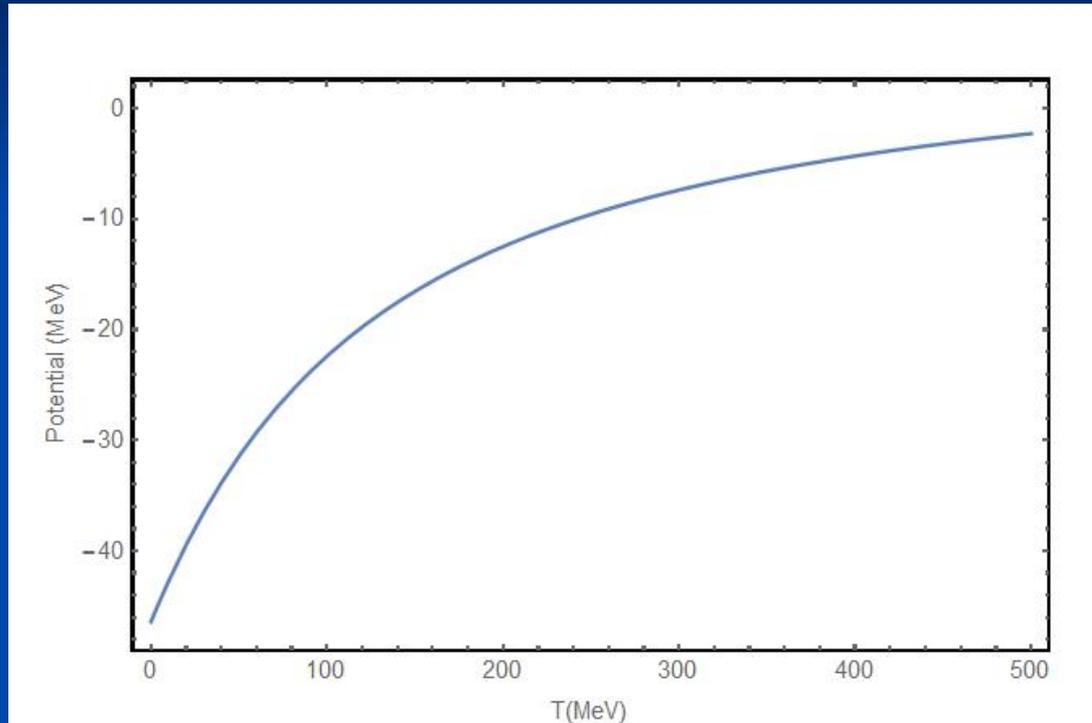
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FSI Potential for 1st Interaction



Welke Potential averaged over density for C12, fitted to p+A data and nuclear EOS

used not only for primary reaction, but also for final state propagation

Model Ingr

Local Fermi
momentum
smoothes E-p
distributions

ISI

Potential
smoothes E-p
distributions

- Hole spectral function (local TF)

$$P_h(\vec{p}, E) \sim \int d^3r [\Theta(p_F(r) - p) \delta(E + T_p + V(r))]$$

- Particle spectral function: collisional broadening

$$P_p(\vec{p}, E) \sim \frac{-\Im\Sigma(\vec{p}, E)}{(p^2 - M_0^2 - \Re\Sigma(\vec{p}, E))^2 + (\Im\Sigma(\vec{p}, E))^2}$$

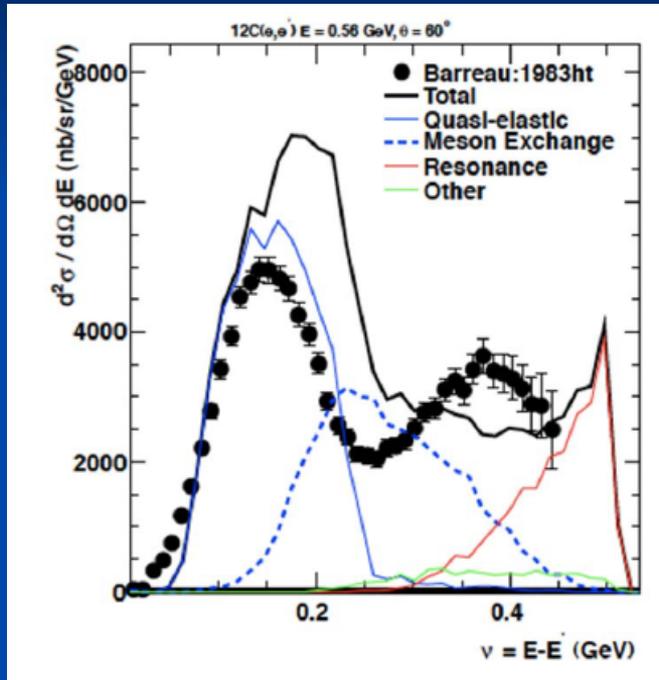
Consistency: $\text{Im}(\Sigma) \sim \sqrt{s} \rho \sigma v$

- Inclusive cross section

$$d\sigma_{\text{tot}}^{lA \rightarrow l'X} = g \int dE \int \frac{d^3p}{(2\pi)^3} P_h(\vec{p}, E) \frac{k \cdot p}{k^0 p^0} d\sigma_{\text{tot}}^{lN} P_{\text{PB}}(\vec{p}, E)$$



Generators must be tested and improved



QE reasonable

MEC bad

Resonance miserable

NOTE: GENIE still uses Rein-Sehgal
for resonance excitations !!

Shortcomings of RS known since ~ NUINT09

GENIE, from S.Dytman, BNL meet, Febr. 2015

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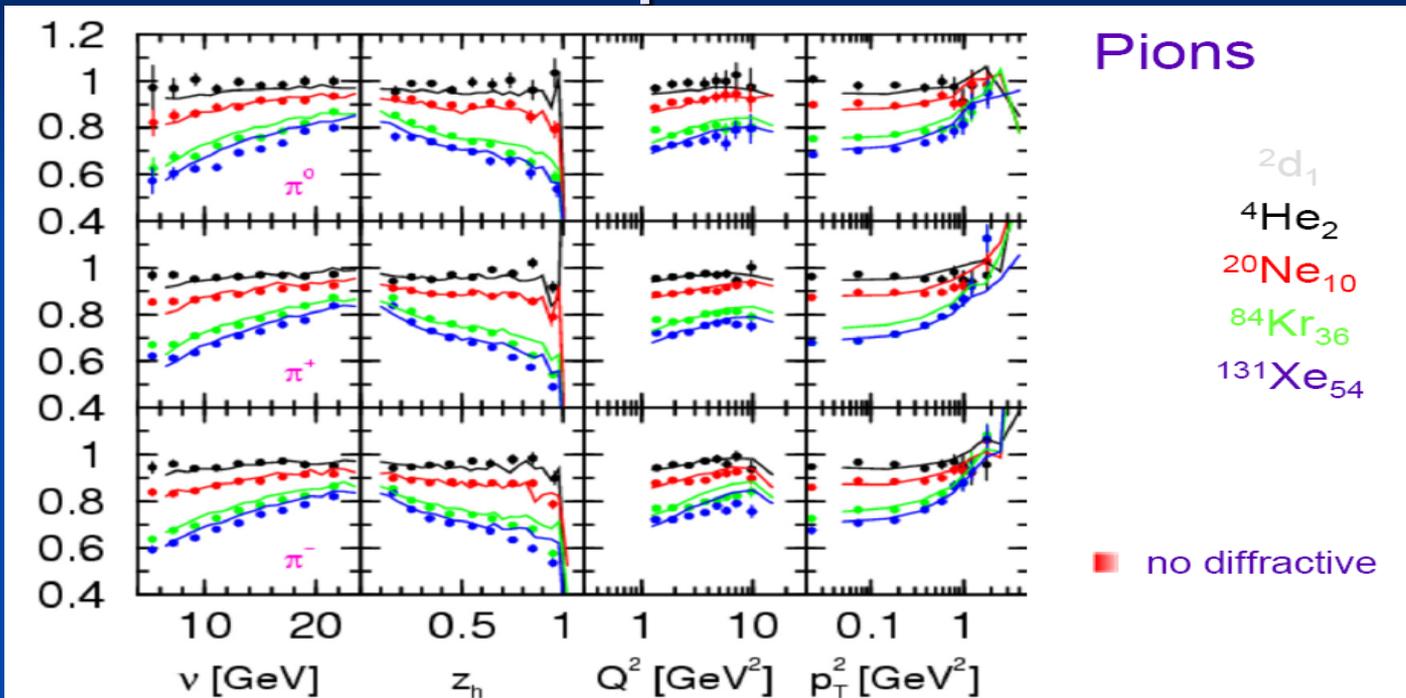


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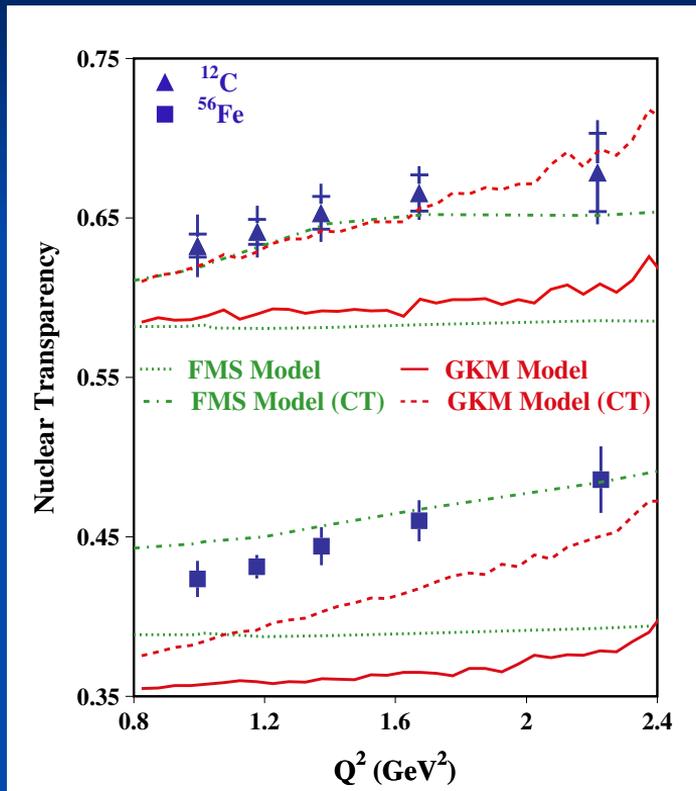
JUSTUS-LIEBIG-
UNIVERSITÄT
GIESSEN

HERMES@27 GeV and GiBUU

Airapetian et al.



JLAB Rho Production



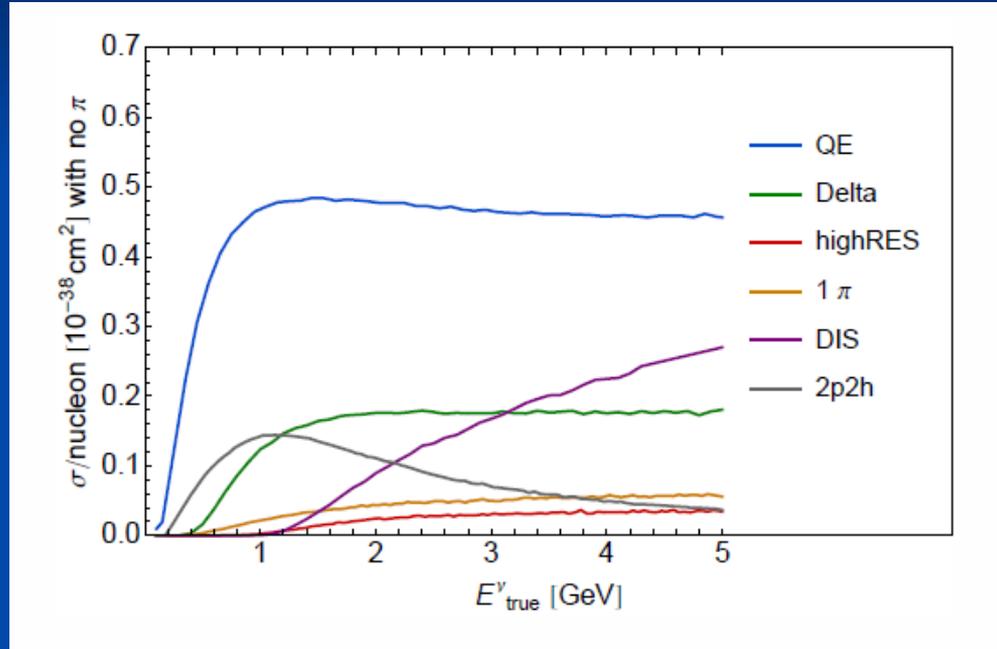
Exp: Hafidi et al,
Phys.Lett. B712 (2012) 326-330

GiBUU: Gallmeister et al.
Phys.Rev. C83 (2011)



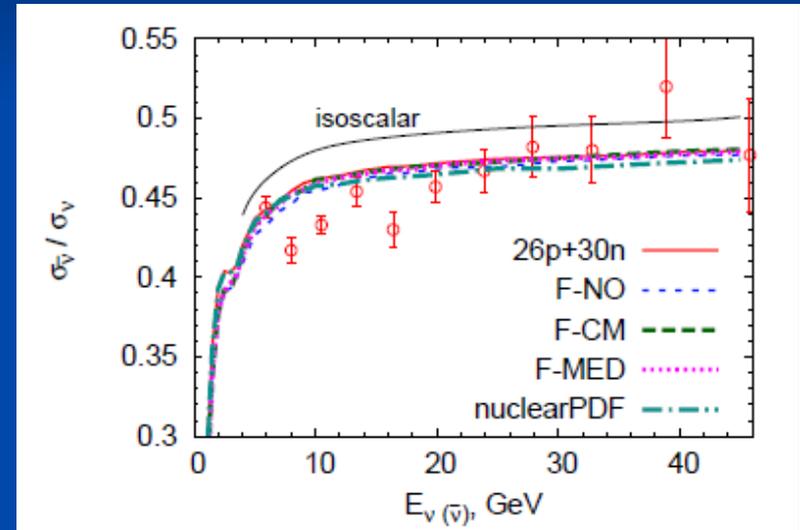
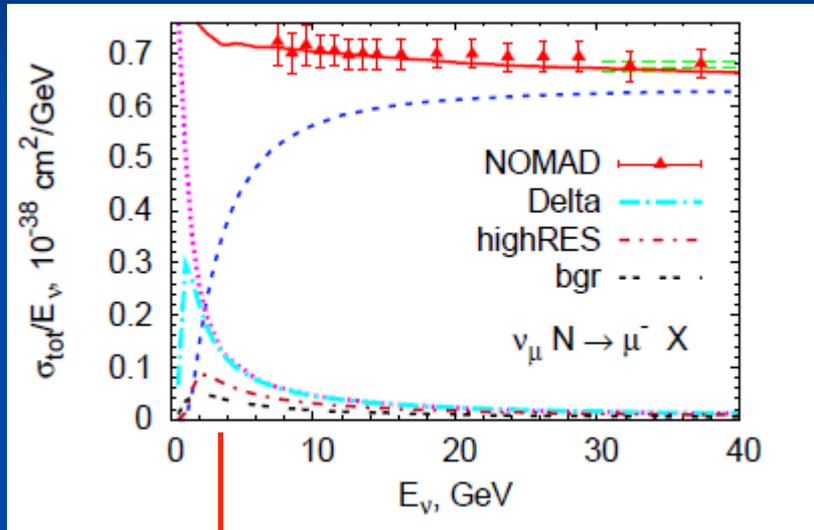
0 Pion Events from GiBUU

From Coloma & Huber: arXiv:1307.1243v1



Experiments at higher energies

Phys. Rev. C86 (2012) 014607



Shallow Inelastic Region, very sensitive to interplay of different reaction mechanisms

