

$|V_{us}|$ FROM INCLUSIVE HADRONIC τ DECAYS

An example of the utility of combined lattice/continuum approaches

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A New Era for Hadro-Particle Physics

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OUTLINE

- *The inclusive flavor-breaking FESR V_{us} determination*
 - *The conventional implementation and the $> 3\sigma$ low $|V_{us}|$ puzzle*
 - *Continuum, lattice input: conventional implementation systematics and a new implementation strategy*
 - *Results of the new implementation and a resolution of the $> 3\sigma$ low $|V_{us}|$ puzzle*
- *A new lattice+inclusive us $V+A$ τ data approach*

BASICS: HADRONIC τ DECAYS IN THE SM

- $\Pi_{ij;V/A}^{(J)}(Q^2)$, $\rho_{ij;V/A}^{(J)}(s)$: the flavor $ij = ud, us$, spin $J = 0, 1$ V, A current-current 2-pt function polarizations and spectral functions. Generically:

$$\begin{aligned}\Pi_{\mu\nu}(q) &= i \int d^4x e^{iq \cdot x} \langle 0 | T \{ J_\mu(x) J_\nu^\dagger(0) \} | 0 \rangle \\ &= (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi^{(1)}(q^2) + q_\mu q_\nu \Pi^{(0)}(q^2) \\ \rho^{(J)}(s) &= \frac{1}{\pi} \text{Im} \Pi^{(J)}(s + i\epsilon)\end{aligned}$$

- $R_{ij;V/A} \equiv \Gamma[\tau \rightarrow \nu_\tau \text{ hadrons}_{ij;V/A}(\gamma)] / \Gamma[\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e(\gamma)]$

- In SM, with $y_\tau \equiv s/m_\tau^2$, $w_\tau(y) = (1 - y)^2(1 + 2y)$,
 $w_L(y) = -2y(1 - y)^2$ [Tsai PRD4 (1971) 2821]

$$\frac{dR_{ij;V/A}}{ds} = \frac{12\pi^2 |V_{ij}|^2 S_{EW}}{m_\tau^2} \left[w_\tau(y_\tau) \rho_{ij;V/A}^{(0+1)}(s) + w_L(y_\tau) \rho_{ij;V/A}^{(0)}(s) \right]$$

- $\rho_{ud,us;V/A}^{(1)}(s)$; π , K pole contributions to $\rho_{ud,us;A}^{(0)}(s)$ all chirally unsuppressed
- Continuum $\rho_{ij;V/A}^{(0)}(s) \propto (m_i \mp m_j)^2$

THE INCLUSIVE FB τ $|V_{us}|$ DETERMINATION

- **Context:** (conventional implementation) inclusive FB τ result c.f. K physics, 3-family unitarity

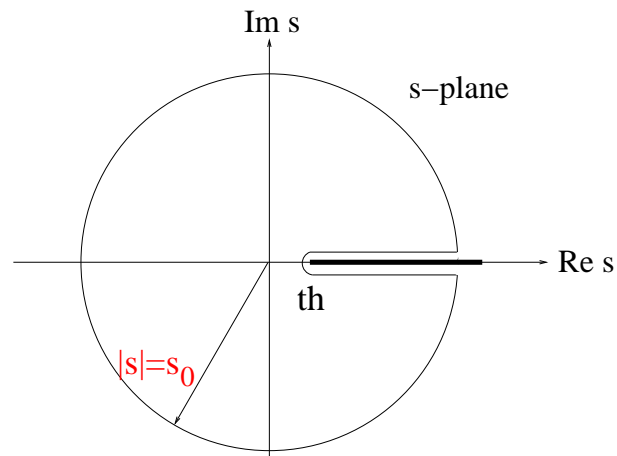
$ V_{us} $	Source
0.2258(9)(?)	3-family unitarity, HT14 $ V_{ud} $
0.2231(4) $_{exp(7)latt}$	$K_{\ell 3}$, 2+1+1 lattice $f_+(0)$
0.2250(4) $_{exp(9)latt}$	$\Gamma[K_{\mu 2}]/\Gamma[\pi_{\mu 2}]$, lattice f_K/f_π
0.2176(19) $_{exp(10?)th}$	Inclusive FB τ FESR (Passemar CKM14)

- Inclusive FB τ result $> 3\sigma$ low c.f. others: interesting **if real**, but theory systematics?

THE CONVENTIONAL IMPLEMENTATION

- **Basic tool:** FESRs (*Cauchy's Thm*), valid for any s_0 , analytic $w(s)$, kinematic-singularity-free Π :

$$\int_{s_{th}}^{s_0} ds w(s) \rho(s) = \frac{-1}{2\pi i} \oint_{|s|=s_0} ds w(s) \Pi(s)$$



- General inclusive FB τ sum rules for $|V_{us}|$
 - * FESRs involving FB polarization, spectral function combinations $\Pi_{ud-us;V+A}^{(J)}(Q^2)$, $\rho_{ud-us;V+A}^{(J)}(s)$
 - * Experimental (spectral) input: $|V_{ij}|^2 \rho_{ij;V/A}^{(J)}(s)$ from $dR_{ij;V/A}/ds$, $R_{ij;V/A} \equiv \frac{\Gamma[\tau \rightarrow \nu_\tau \text{ hadrons}_{ij;V/A}(\gamma)]}{\Gamma[\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e(\gamma)]}$
(small external us $J=0$ “subtraction” required)

[SM “kinematic weight” w_τ]

- * $R_{ij;V/A}^w(s_0)$: Re-weighted $R_{ij;V/A}$ analogue, integrated to variable upper endpoint s_0 in spectrum

$$R_{ij;V/A}^w(s_0) \sim \int_{th}^{s_0} ds \frac{dR_{ij;V/A}}{ds} \frac{w(s/s_0)}{w_\tau(s/m_\tau^2)}$$

* FESR for $\delta R^w \equiv \frac{R_{ud;V+A}^w}{|V_{ud}|^2} - \frac{R_{us;V+A}^w}{|V_{us}|^2}$ yields

$$|V_{us}| = \sqrt{\frac{R_{us;V+A}^w(s_0)}{\frac{R_{ud;V+A}^w(s_0)}{|V_{ud}|^2} - [\delta R^w(s_0)]^{OPE}}}$$

valid for arbitrary analytic $w(y)$, $s_0 \leq m_\tau^2$

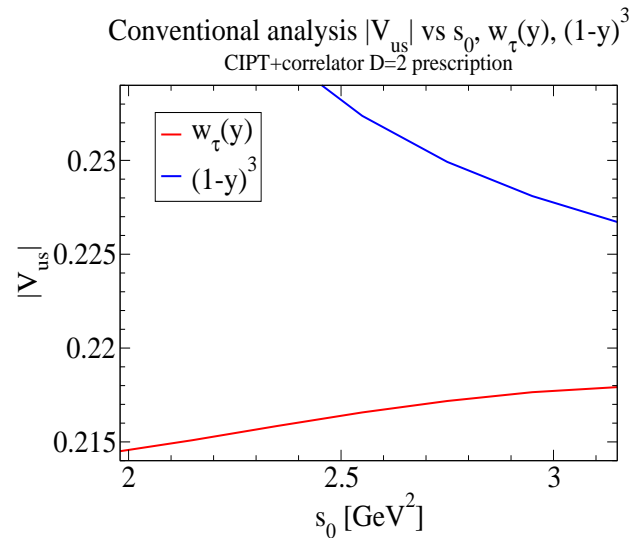
* “Self-consistency tests”:

- $|V_{us}|$ independent of s_0 , w (tests control, understanding of OPE, experimental systematics)
- E.g., **integrated $D = 2k + 2$ OPE $\sim 1/s_0^k$** : errors in higher D treatment $\leftrightarrow s_0$ -instability

• **Below: FESRs with variable s_0 , w , $\Delta\Pi_\tau \equiv \Pi_{ud-us;V+A}^{(0+1)}$**

- The conventional implementation [Gamiz et al. JHEP03(2003)060]
 - $s_0 = m_\tau^2$, $w = w_\tau$ only (degree 3, OPE to $D = 8$)
 - $D = 2$, $D = 4$ known, $D = 6$ OPE estimated with “VSA” (very small), $D = 8$ neglected
 - Resulting $\left[\delta R^{w_\tau}(m_\tau^2) \right]_{D \geq 2}^{OPE}$ estimate a few to several % of individual $R_{ud,us;V+A}$
 - \Rightarrow modest OPE accuracy enough for precision $|V_{us}|$
 - Rationale for $w = w_\tau$, $s_0 = m_\tau^2$ choice: $R_{ud,us;V+A}$ from inclusive ud , us BFs ($dR_{ij;V/A}/ds$ not needed)
 - [Downside: unquantified systematic error from absence of $w(s)$ -, s_0 -independence tests]

- Conventional implementation tests [KM et al. arXiv:1511.08514]
 - Variable $s_0 \leq m_\tau^2$, $|V_{us}|$ s_0 -stability checks
 - Targeted $|V_{us}|$ w -independence test: $y = (s/s_0)$,
 $w = w_\tau(y) = 1 - 3y^2 + 2y^3$ c.f. $\hat{w}(y) = 1 - 3y + 3y^2 - y^3$
 - * $D = 6/8$ \hat{w} OPE contributions $-1/-1/2$ \times $D = 6/8$ w_τ OPE contributions
 - * IF $D > 4$ assumptions OK for w_τ FESR \Rightarrow also OK for \hat{w} FESR $\Rightarrow |V_{us}|$ agreement, s_0 -stability
 - * $D > 4$ assumptions not OK \Rightarrow opposite-sign w_τ , \hat{w} s_0 -instabilities, both decreasing with s_0



- Candidate self-consistency problem sources:
 - * **Experiment:** The less-well-known us distribution
 - * **Theory:** Conventional $D > 4$ assumptions (w_τ , \hat{w} comparison); slow $D = 2$ OPE convergence

- The $D = 2$ series and slow $D = 2$ convergence

$$\left[\Delta \Pi_\tau(Q^2) \right]_{D=2}^{OPE} = \frac{3}{2\pi^2} \frac{\bar{m}_s^2}{Q^2} \left[1 + 2.333\bar{a} + 19.933\bar{a}^2 + 208.746\bar{a}^3 + \dots \right]$$

with running $\overline{MS} \bar{a} = \frac{\alpha_s(Q^2)}{\pi}$, $\bar{m}_s = m_s(Q^2)$

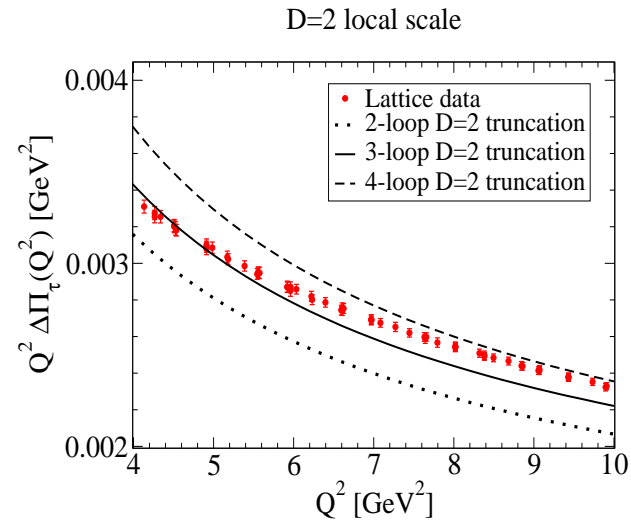
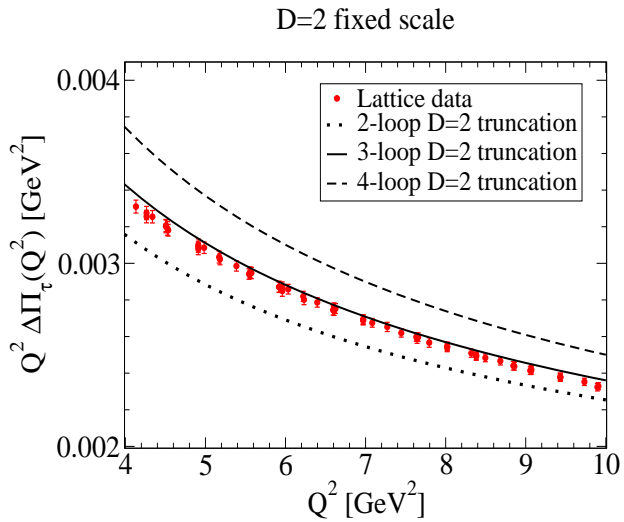
$\bar{a}(m_\tau^2) > 0.1 \Rightarrow$ slow convergence at ALL τ decay scales

- w_τ , \hat{w} comparison: $D > 4$ assumption problem(s) likely

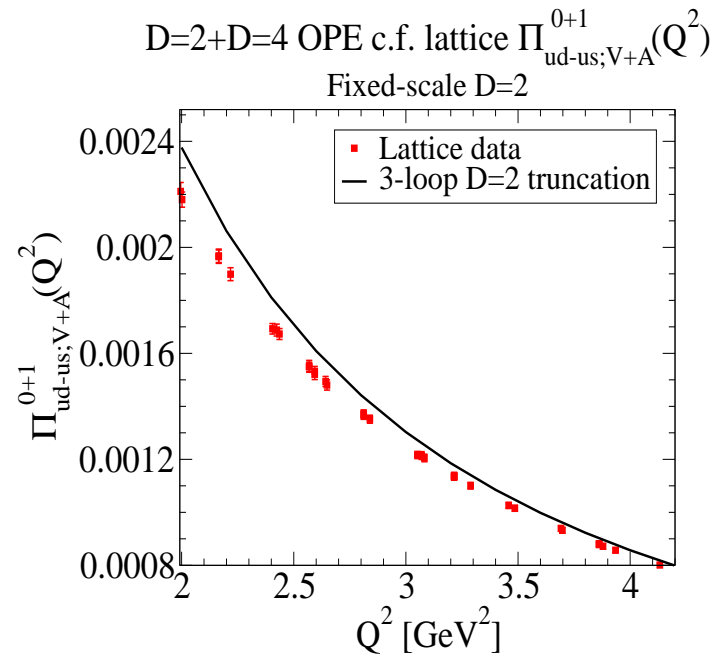
Note: VSA VERY crude; channel-dependent ud V,A breaking; caution re strong double cancellation (in ud , us V+A sums, AND resulting $ud - us$ difference)

LATTICE RE $D = 2$ SERIES, $D > 4$ OPE ISSUES

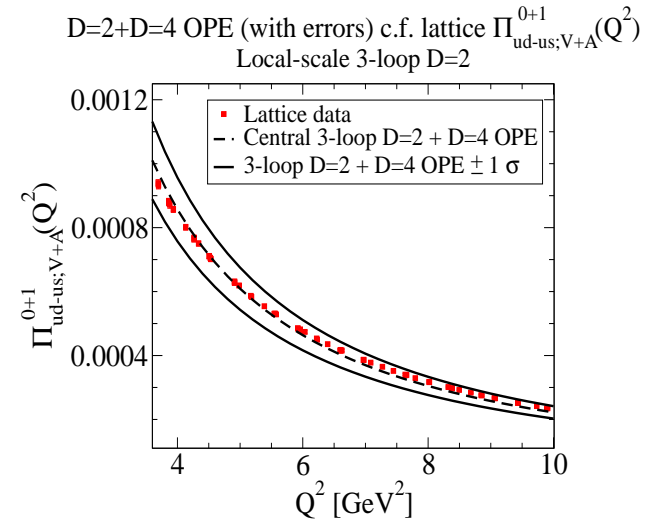
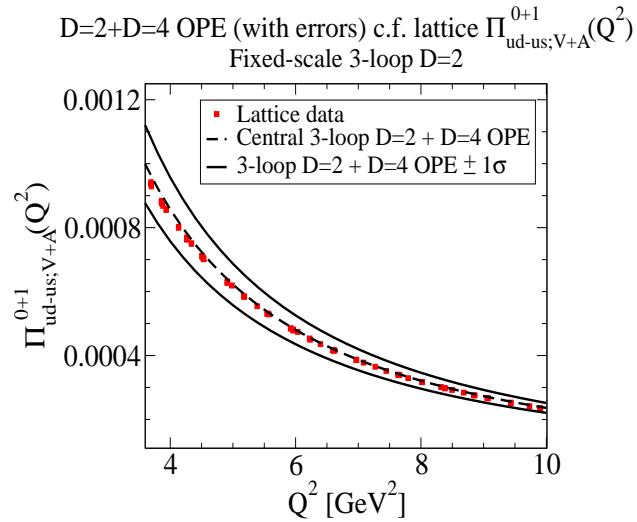
- OPE c.f. RBC/UKQCD lattice $\Delta\Pi_\tau(Q^2)$
 - KM with J. Hudspith, R. Lewis, J. Zanotti
 - Here: lightest m_π , fine ($1/a = 2.38 \text{ GeV}$) $32^3 \times 64$
 $2 + 1$ ensemble ($m_\pi \sim 300 \text{ MeV}$, $m_\pi L \sim 4.1$)
 - Tight cylinder cut for continuum correlator behavior
(width from stabilization in ud V channel α_s study)
 - Large Q^2 : $D = 2 + 4$ OPE exploration; lower Q^2 : re
possible non-negligible $D > 4$ at τ decay scales
 - * Fixed- or local-scale $D = 2$? [c.f. FESR FOPT/CIPT]



- Higher Q^2 : best (excellent) lattice vs $D = 2 + 4$ OPE match for 3-loop-truncated, fixed-scale $D = 2$
- Fixed scale suggests FOPT for FESR $D = 2$



- Onset of $D > 4$ contributions below $\sim 4 \text{ GeV}^2$



- Standard $D = 2$, $D = 4$ error estimates VERY conservative, despite slow $D = 2$ series convergence

A NEW IMPLEMENTATION STRATEGY BASED ON LATTICE, CONTINUUM LESSONS

- No assumptions re $D > 4$; include contributions and fit effective condensates C_D as part of analysis
- 3-loop-truncated FOPT for $D = 2$ contribution; standard error estimates for $D = 2 + 4$
- Range of s_0 to fit both $|V_{us}|$, C_D
- Fits simplest for $w(y) = w_N(y) = 1 - \frac{y}{N-1} + \frac{y^N}{N-1}$ (C_{2N+2} , $|V_{us}|$ as only fit parameters)
- Self-consistency check: $|V_{us}|$ from different w_N

RESULTS OF THE NEW ANALYSES

- For brevity: results for $K\pi$ normalization using preliminary BaBar $B[\tau \rightarrow K^- \pi^0 \nu_\tau]$ update (Adametz thesis)
- Fitted $|V_{us}|$ (as expected) between (s_0 -unstable) w_τ and \hat{w} conventional implementation results
- Fitted C_D show FB cancellation in comparison to fitted ud $V+A$ analogues (qualitative self-consistency test)
- Excellent $|V_{us}|$, C_{2N+2} stability wrt fit window variation; excellent agreement of $|V_{us}|$ from different w_N

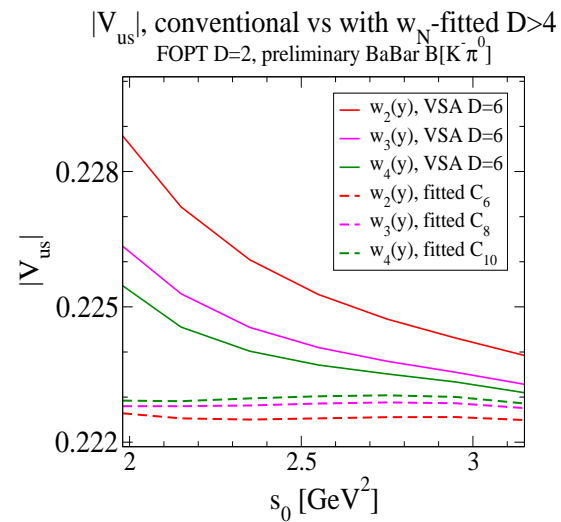
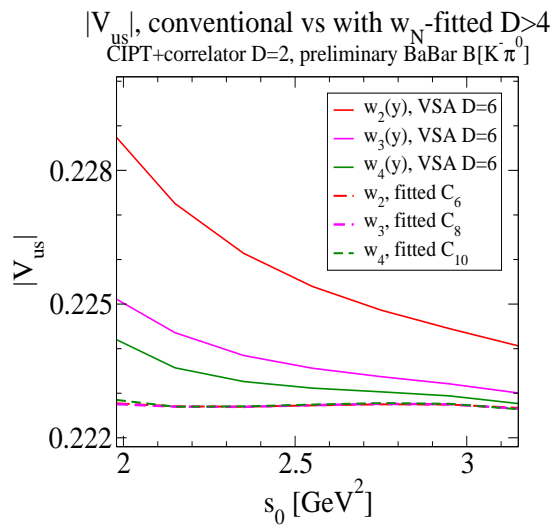
- Results using preliminary $K\pi$ BF update

Weight	$ V_{us} $	$ V_{us} $
	CIPT+corr $D = 2$	FOPT $D = 2$
w_2	0.22271(228)	0.22252(228)
w_3	0.22271(228)	0.22282(228)
w_4	0.22271(229)	0.22296(229)

Error budget, 3-loop-truncated FOPT $D = 2$ fits

Source	$\delta V_{us} $	$\delta V_{us} $	$\delta V_{us} $
	(w_2 FESR)	(w_3 FESR)	(w_4 FESR)
$\delta\alpha_s$	0.00001	0.00004	0.00004
$\delta m_s(2 \text{ GeV})$	0.00017	0.00019	0.00019
$\delta\langle m_s \bar{s}s \rangle$	0.00035	0.00035	0.00035
$\delta(\text{long corr})$	0.00009	0.00009	0.00009
ud exp	0.00027	0.00028	0.00028
us exp	0.00226	0.00227	0.00227

Fit quality: s_0 -stability for conventional implementation, except with fitted C_{2N+2}



Solid/dashed lines: with conventional implementation/new implementation (fit) treatments of $D > 4$ contributions

SUMMARY, CONVENTIONAL FB FESR ANALYSIS

- Continuum, lattice \Rightarrow conventional $D > 4$ assumptions untenable, must fit $D > 4$ effective condensates

[Additional $\sim 0.0020 - 0.0030$ systematic from breakdown of $D > 4$ assumptions]

- Fitting $D > 4$ resolves s_0 -, $w(y)$ -dependence problems, raises $|V_{us}|$ by ~ 0.0020
- $D > 4$ fit needs variable $s_0 \Rightarrow$ **analysis using only inclusive non-strange, strange BFs NOT possible**
- Lattice \Rightarrow 3-loop FOPT for $D = 2$, $D = 2 + 4$ OPE error estimate conservative

- $|V_{us}|$ errors strongly dominated by (improvable) us experimental uncertainties (currently BF dominated)
- New preliminary $K^-\pi^0$ BF normalization, 3-weight combined fit result $|V_{us}| = 0.2228(23)_{exp}(5)_{th}$ in agreement with other determinations (especially $K_{\ell 3}$)
(c.f. $0.2200(23)_{exp}(5)_{th}$ with 2014 HFAG $B[K^-\pi^0]$)
- Theory error (conservative) ~ 0.0005 c.f. 0.0009 for $K_{\ell 3}$, $\Gamma[K_{\mu 2}]/\Gamma[\pi_{\mu 2}]$ (lattice crucial here)
- \Rightarrow **FB FESR method competitive with sufficient improvement to us experimental errors**

- However sub-0.5% $|V_{us}|$ needs sub-% $R_{us;V+A}^w$ error

Relative exclusive mode $R_{us;V+A}^w$ contributions

Wt	s_0 [GeV ²]	K	$K\pi$	$K\pi\pi$ (B-factory)	Other
w_2	2.15	0.496	0.426	0.062	0.010
	3.15	0.360	0.414	0.162	0.065
w_3	2.15	0.461	0.446	0.073	0.019
	3.15	0.331	0.415	0.182	0.074
w_4	2.15	0.441	0.456	0.082	0.021
	3.15	0.314	0.411	0.194	0.081

- “Other”: 1999 ALEPH data/MC, $\sim 25\%$ error

\Rightarrow “sufficient improvement” includes experimentally (much) more challenging higher-multiplicity modes

A PROMISING τ -BASED ALTERNATIVE

- Work with J. Hudspith, T. Izubuchi, R. Lewis, H. Ohki, C. Lehner + RBC/UKQCD
- Basic idea: generalized dispersion relations for products of combination $\tilde{\Pi}$ of $J = 0, 1$ us $V+A$ polarizations with weights having poles at Euclidean Q^2
 - $\tilde{\Pi}(Q^2)$ such that inclusive $dR_{us;V+A}/ds$ distribution directly determines related spectral function $\tilde{\rho}(s)$
 - Theory: Lattice us 2-point function data (no OPE)
 - Weights tunable, allow suppression of larger-error, higher-multiplicity us spectral contributions

More on the lattice-inclusive us τ approach

- Experimental $dR_{us;V+A}/ds$ distribution yields $|V_{us}|^2 \tilde{\rho}(s)$,

$$\tilde{\rho}_{us;V+A}(s) \equiv \left(1 + 2\frac{s}{m_\tau^2}\right) \rho_{us;V+A}^{(J=1)}(s) + \rho_{us;V+A}^{(J=0)}(s)$$

(no continuum us $J = 0$ subtraction needed)

- Associated (kinematic-singularity-free) polarization

$$\tilde{\Pi}_{us;V+A}(Q^2) \equiv \left(1 - 2\frac{Q^2}{m_\tau^2}\right) \Pi_{us;V+A}^{(J=1)}(Q^2) + \Pi_{us;V+A}^{(J=0)}(Q^2)$$

Reminder: strange hadronic τ decays in the SM

With $y_\tau \equiv s/m_\tau^2$ [Tsai PRD4 (1971) 2821]

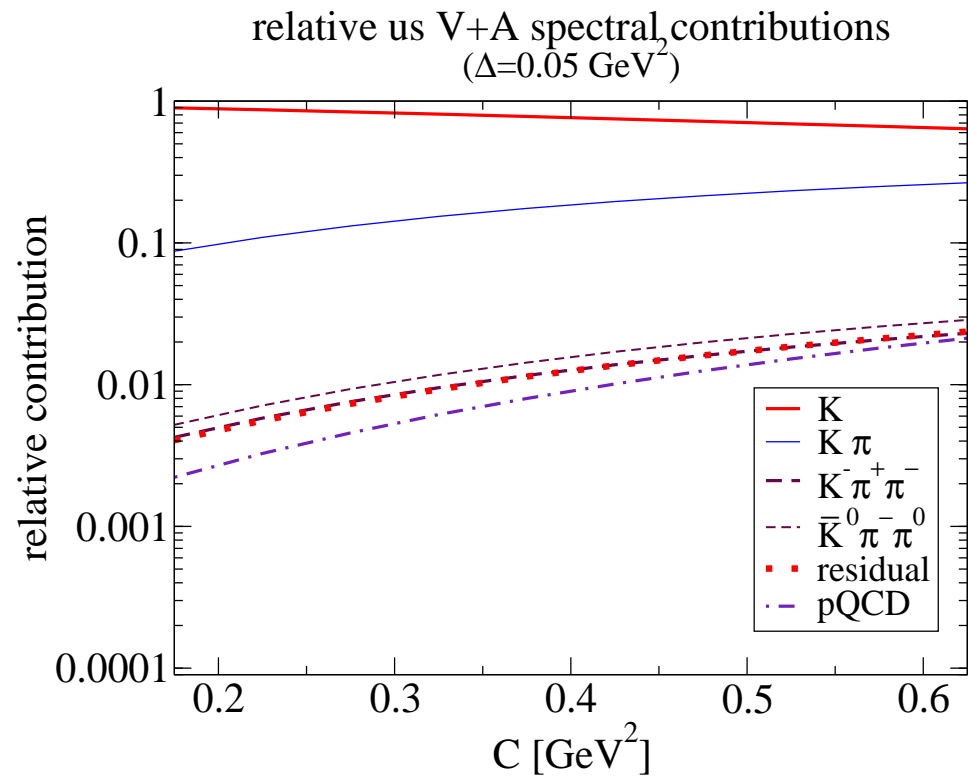
$$\begin{aligned} \frac{dR_{us;V+A}}{ds} &= \frac{12\pi^2 |V_{us}|^2 S_{EW}}{m_\tau^2} [1 - y_\tau]^2 \\ &\quad \cdot \left[(1 + 2y_\tau) \rho_{us;V+A}^{(J=1)}(s) + \rho_{us;V+A}^{(J=0)}(s) \right] \\ &= \frac{12\pi^2 |V_{us}|^2 S_{EW}}{m_\tau^2} [1 - y_\tau]^2 \tilde{\rho}_{us;V+A}(s) \end{aligned}$$

\Rightarrow Division by known factors yields $\tilde{\rho}_{us;V+A}(s)$ with no need for continuum $\rho_{us;V/A}^{(0)}(s)$ subtraction

- For weights $w_N(s) \equiv \frac{1}{\prod_{k=1}^N (s+Q_k^2)}$, $N \geq 3$, obtain convergent, unsubtracted 'dispersion relation'

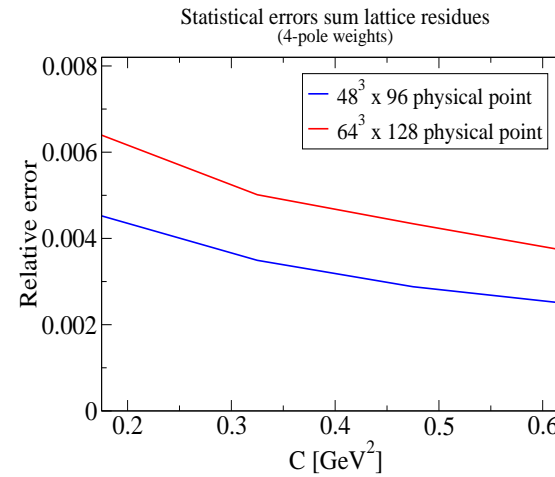
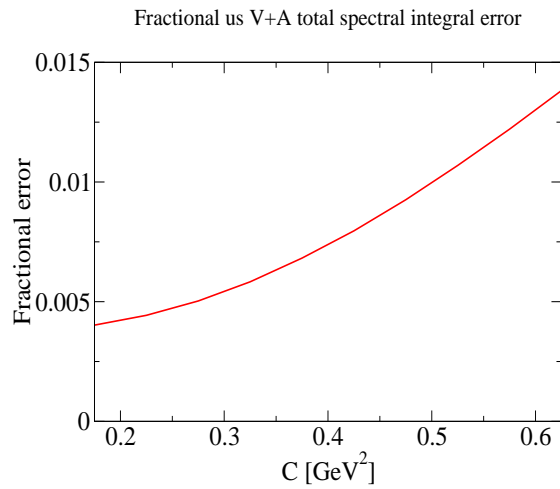
$$\int_{th}^{\infty} ds w_N(s) \tilde{\rho}_{us;V+A}(s) = \sum_{k=1}^N \frac{\tilde{\Pi}_{us;V+A}(Q_k^2)}{\prod_{j \neq k} (Q_j^2 - Q_k^2)}$$

- Lattice data for $\tilde{\Pi}_{us;V+A}(Q_k^2)$ on RHS
- LHS from experimental $dR_{us;V+A}/ds$, up to $|V_{us}|^2$
- $w_N(s)$: rapid fall-off if all $Q_k^2 < 1 \text{ GeV}^2 \Rightarrow K, K\pi$ **dominate LHS, near-endpoint multi-particle, $s > m_\tau^2$ contributions strongly suppressed [FIG]**
- Optimization: increasing $\{Q_k^2\}$ decreases RHS lattice error, increases LHS experimental error



4-pole weights, pole spacing Δ , pole centroid C

- An illustration of expected precision: $N = 4$, $C = 0.325 \text{ GeV}^2$ [$\{Q_k^2\} = \{0.25, 0.30, 0.35, 0.40 \text{ GeV}^2\}$]
 - Spectral integral error $\sim 0.4\%$
 - Statistical error, sum of residues, near-physical-point $n_f = 2 + 1$ RBC/UKQCD ensembles:
 - $\sim 0.4\%$ for $48^3 \times 96$, $1/a = 1.73 \text{ GeV}$
 - $\sim 0.5\%$ for $64^3 \times 128$, $1/a = 2.36 \text{ GeV}$
 - Quadrature sum $\Rightarrow \delta|V_{us}| \sim 0.0007$, ~ 0.0008 for near-physical-point 48^3 , 64^3 ensembles, respectively
- Optimization/systematics (FV, continuum limit, correcting for small m_ℓ , m_s mis-tunings) under study



Relative errors as a function of C , 4-pole weights, pole spacing $\Delta = 0.05 \text{ GeV}^2$, pole interval centroid C

OVERALL SUMMARY

- Old 3σ low inclusive FB τ FESR $|V_{us}|$ problem resolved
 - Alternate, no-assumptions implementation: $|V_{us}|$ higher by ~ 0.0020 , compatible with other determinations
 - Near-term improvements feasible through improvements in us exclusive mode BFs
 - Highly favorable theoretical error situation
 - However, for competitive $|V_{us}|$ need improvements to old ALEPH higher-multiplicity, low-statistics data [unlikely in the near-term]

- Advantage of new lattice-inclusive $us V + A \tau$ approach

- Theory:

- * Lattice in place of OPE; no (albeit mildly model-dependent) $us J = 0$ continuum subtraction; improvement through increased statistics
 - * Easily parasitic on lattice a_μ effort (major effort being pursued by many lattice collaborations)

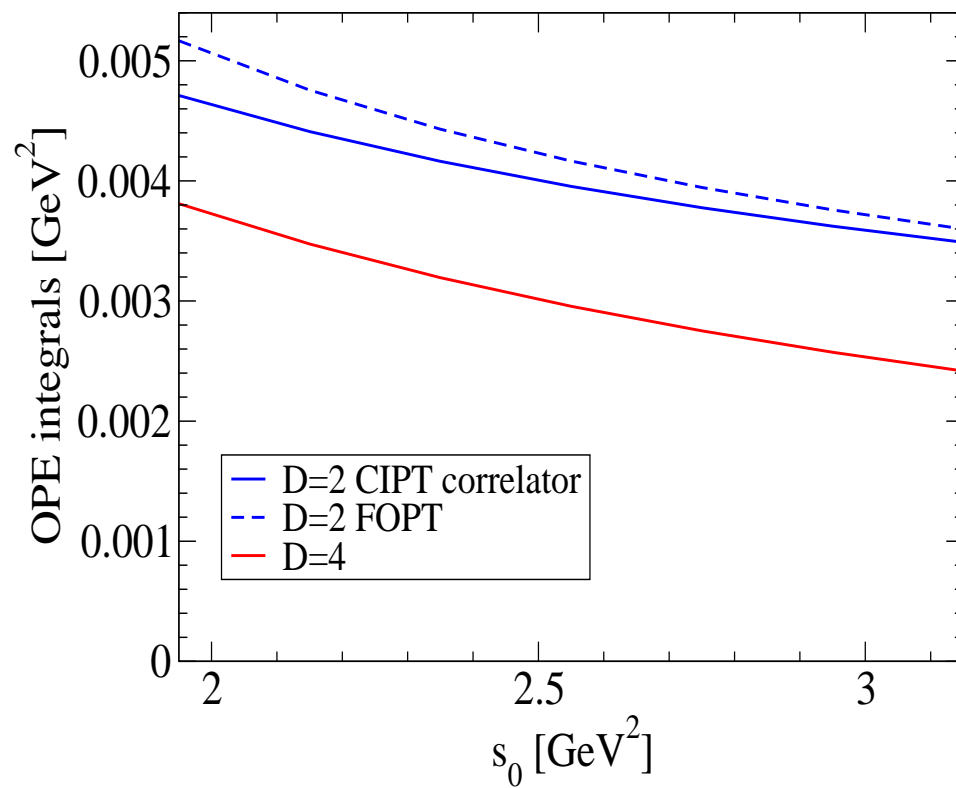
- Spectral integrals:

- * Theory errors still small for weights strongly suppressing higher multiplicity contributions
 - * Strong $K, K\pi$ dominance of spectral integral
 - * Significant experimental improvements possible through just improved $K\pi$ BFs, distributions

BACKUP SLIDES

- Are $D = 6, 8$ OPE contributions likely to be small for the conventional inclusive FB FESRs?
 - $D = 4 \ll D = 2$ for $w_\tau(y) = 1 - 3y^2 + 2y^3$, $y = s/s_0$, “accidental” [$O(\alpha_s^2)$ suppression due to absence of term linear in y in $w_\tau(y)$]
 - Comparison of $D = 2, 4$ OPE contributions for $w(y) = (1 - y)^2$ (a case without this suppression) to see natural relative sizes

D=2, 4 OPE integrals
for weight $w(y)=(1-y)^2 (y=s/s_0)$



OPE, SPECTRAL INPUT

- PDG, FLAG, HPQCD input for $D = 2, 4$ OPE
- ud V+A spectral data from ALEPH 2013
- us V+A spectral data from sum over exclusive modes [$> 90\%$ of B_{us}^{TOT} from $K_{\ell 2}$, Belle, BaBar $K\pi$, $K\pi\pi$, $3K$ results; residual: 1999 ALEPH]
- Here, for brevity, $K\pi$ normalization including preliminary BaBar $B[\tau \rightarrow K^- \pi^0 \nu_\tau]$ update (Adametz thesis)

MORE ON THE us DATA

- K pole via $f_K|V_{us}|$ from $K_{\ell 2}$
- Rather precise unit-normalized $K^-\pi^0$, $\bar{K}^0\pi^-$, $K^-\pi^+\pi^-$, $\bar{K}^0\pi^-\pi^0$, $3K$ distributions from Belle, BaBar (main uncertainties from BFs)
- K , B-factory modes over 90% of B_{us}^{TOT}
- Residual us exclusive mode contributions from 1999 ALEPH data, covariances

THE EXPERIMENTAL $K\pi$ BF SITUATION

- HFAG 2014 $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau] = 0.0126$
- HFAG 2014 $B[K^-\pi^0\nu_\tau] = 0.00431(15)$ value \rightarrow preliminary BaBar (Adametz thesis) result $0.00500(15)$ yields $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau] = 0.0133$
- Central $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau]$ from $K_{\ell 3}$, dispersive analysis expectations [ACLP13] also 0.0133
- 0.07% difference “small” but represents $\sim 2.4\%$ of B_{us}^{TOT} , hence $\sim 1.2\%$ increase in $|V_{us}|$

Results for $|V_{us}|$ for current HFAG 2014 $K\pi$ BFs

Weight	$ V_{us} $ CIPT+corr $D = 2$	$ V_{us} $ FOPT $D = 2$
w_2	0.21985(230)	0.21966(230)
w_3	0.21985(231)	0.21966(231)
w_4	0.21985(231)	0.22009(231)

Error budget, existing $K\pi$ BFs

Source	$\delta V_{us} $ (w_2 FESR)	$\delta V_{us} $ (w_3 FESR)	$\delta V_{us} $ (w_4 FESR)
$\delta\alpha_s$	0.00001	0.00003	0.00005
$\delta m_s(2 \text{ GeV})$	0.00017	0.00018	0.00020
$\delta\langle m_s \bar{s}s \rangle$	0.00034	0.00034	0.00034
$\delta(\text{long corr})$	0.00009	0.00009	0.00009
ud exp	0.00027	0.00027	0.00027
us exp	0.00229	0.00229	0.00230

Stability of $|V_{us}|$ with fitted C_{2N+2} input, existing $K\pi$ BF normalization

