Testing the Structure of Light Scalar Mesons in the Decays of $\overline{B}{}^0 \to D^0 \pi^+ \pi^-$, $D^0 K^+ K^-$, $D^0 \pi^0 \eta$ and $\overline{B}{}^0_s \to D^0 \pi^- K^+$

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• The nature of the light scalar mesons is a topic of longstanding debate.

In the Quark Model: meson $\implies q\overline{q}$

Comparing with the Quark Model's predictions, there exist much more light scalar mesons.

Scalar mesons below 1 GeV:

 $f_0(500) \text{ (or } \sigma), f_0(980), a_0(980), \kappa(800)$

Possible structure:

normal meson $(q\overline{q})$, tetraquark $[(q)^2(\overline{q})^2]$, molecular state $[(q\overline{q})(q\overline{q})]$, glueball (gg, ggg), hybrid $(q\overline{q}g)$,.....

• In the Chiral Unitary Approach:

 $f_0(500), f_0(980), a_0(980)$ and κ are dynamically generated from the interaction of pseudoscalar mesons, and could be interpreted as a kind of molecular states of meson-meson.

 $f_0(500)$ couples mostly to $\pi\pi$; $f_0(980)$ couples mostly to $K\overline{K}$; and $a_0(980)$ mostly to $\pi\eta$ and $K\overline{K}$.

$$f_0(500) \longrightarrow \pi\pi$$
 resonance $f_0(980) \longrightarrow K\overline{K}$ molecule
 $a_0(980) \longrightarrow \pi\eta + K\overline{K}$ molecule

[Oller and Oset, NPA620(1997)438]

• The weak decay of *B* mesons is a good place for testing the structure of the light scalar mesons.

Experimental results:

 $\overline{B}_{s}^{0} \rightarrow J/\psi \pi^{+}\pi^{-}$ decay : A clear peak is observed for $f_{0}(980)$ production ; $f_{0}(500)$ production is not seen;

 $\overline{B}^0 \to J/\psi \pi^+ \pi^-$ decay : A signal is seen for $f_0(500)$ production ;

Only a very small fraction is observed for $f_0(980)$ production.

LHCb: PLB698(2011)115; PRD86(2012)052006; PRD87(2013)052001; PRD89(2014)092006; PRD90(2014)012003; Belle, PRL106(2011)121802; CDF, PRD84(2011)052012; 5 D0, PRD85(2012)011103.



[LHCb, PRL109(2012)131801] 6

In this work: Using these experimental information, we can investigate the production of the light scalar mesons in

 $\overline{B}^{0} \to D^{0} \pi^{+} \pi^{-}, D^{0} K^{+} K^{-}, D^{0} \pi^{0} \eta$ $\overline{B}^{0}_{s} \to D^{0} \pi^{-} K^{+}.$

Considering the final state interactions between the light pseudoscalar meson pairs in chiral unitary approach, $f_0(500)$, $f_0(980)$, $a_0(980)$ and $\kappa(800)$ can be produced dynamically.

 $\pi^+\pi^-, K^+K^-, \pi^0\eta, \pi^-K^+$ invariant mass distribution



• The Cabibbo favored dominant mechanism:



d and \overline{s} act as spectators. \implies Amplitudes for these processes are identical.

Assume: A $q\bar{q}$ nature for vector mesons;

 $f_0(500), f_0(980), a_0(980)$ and κ generated dynamically from PS-PS interaction.



• $\overline{B}^{0}(\overline{B}_{s}^{0}) \rightarrow D^{0}M_{1}M_{2}$ a pair of PS mesons

The $q\bar{q}$ pair is allowed to hadronize into two PS mesons.





The $q\bar{q}$ matrix :

$$M = \begin{pmatrix} u\bar{u} & u\bar{d} & u\bar{s} \\ d\bar{u} & d\bar{d} & d\bar{s} \\ s\bar{u} & s\bar{d} & s\bar{s} \end{pmatrix}$$

with the property $M \cdot M = M \times (\bar{u}u + \bar{d}d + \bar{s}s)$.

Write the matrix *M* in terms of PS mesons :

$$\Phi = \begin{pmatrix} \frac{1}{\sqrt{2}}\pi^{0} + \frac{1}{\sqrt{3}}\eta + \frac{1}{\sqrt{6}}\eta' & \pi^{+} & K^{+} \\ \pi^{-} & -\frac{1}{\sqrt{2}}\pi^{0} + \frac{1}{\sqrt{3}}\eta + \frac{1}{\sqrt{6}}\eta' & K^{0} \\ & K^{-} & \bar{K}^{0} & -\frac{1}{\sqrt{3}}\eta + \sqrt{\frac{2}{3}}\eta' \end{pmatrix}$$

♦ Formalism

$$d\bar{d}(\bar{u}u + \bar{d}d + \bar{s}s) \to (\Phi \cdot \Phi)_{22} = \pi^{-}\pi^{+} + \frac{1}{2}\pi^{0}\pi^{0} + \frac{1}{3}\eta\eta - \sqrt{\frac{2}{3}}\pi^{0}\eta + K^{0}\bar{K}^{0},$$

$$s\bar{d}(\bar{u}u + \bar{d}d + \bar{s}s) \to (\Phi \cdot \Phi)_{23} = \pi^- K^+ - \frac{1}{\sqrt{2}}\pi^0 K^0,$$

Let these PS mesons interact:



(direct mechanism)

(rescattering mechanism)

The amplitude for $\pi^+\pi^-$ production :

$$t(\bar{B}^{0} \to D^{0}\pi^{-}\pi^{+}) = V_{P}(1 + G_{\pi^{-}\pi^{+}}t_{\pi^{-}\pi^{+}\to\pi^{-}\pi^{+}} + \frac{1}{2}G_{\pi^{0}\pi^{0}}t_{\pi^{0}\pi^{0}\to\pi^{-}\pi^{+}} + \frac{1}{3}G_{\eta\eta}t_{\eta\eta\to\pi^{-}\pi^{+}} + G_{K^{0}\bar{K}^{0}}t_{K^{0}\bar{K}^{0}\to\pi^{-}\pi^{+}}),$$

 V_P : *a* common factor to all $\overline{B}^0(\overline{B}^0_s) \to D^0 MM$ decays;

 G_i : the loop function of two PS meson propagators.

$$G_i(s) = i \int \frac{d^4q}{(2\pi)^4} \frac{1}{(P-q)^2 - m_1^2 + i\varepsilon} \frac{1}{q^2 - m_2^2 + i\varepsilon}, \qquad s = \mathbf{P}^2$$

 $t_{i \rightarrow j}$: scattering matrix for $i \rightarrow j$, calculated in the chiral unitary approach, following the Bethe-Salpeter Equation in coupled channels:

$$t = V + VGt$$
, or $t = [1 - VG]^{-1}V$,

[Oller and Oset, NPA620(1997)438] [F.K. Guo et al., NPA773(2006)78]



$$\pi^{+}\pi^{-} \text{ invariant mass } (M_{\text{inv}}) \text{ distributi on :}$$
$$\frac{d\Gamma}{dM_{\text{inv}}} = \frac{1}{(2\pi)^{3}} \frac{p_{D}\tilde{p}_{\pi}}{4M_{\bar{B}^{0}}^{2}} |t(\bar{B}^{0} \to D^{0}\pi^{-}\pi^{+})|^{2},$$

Similar formulas for $\overline{B}{}^{0} \to D^{0}K^{-}K^{+}$, $\overline{B}{}^{0} \to D^{0}\pi^{0}\eta$, $\overline{B}{}^{0} \to D^{0}\pi^{-}K^{+}$.



•
$$\overline{B}^{0}(\overline{B}_{s}^{0}) \rightarrow D^{0}V_{i}$$

$$|\rho^{0}\rangle = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}); \ |\omega\rangle = \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d}); \ |K^{*0}\rangle = d\bar{s}.$$

$$J^P: \mathbf{0}^- \rightarrow \mathbf{0}^- \mathbf{1}^-, (L=1, P-\text{wave})$$

The amplitudes:

$$\begin{split} t_{\bar{B}^{0}\to D^{0}\rho^{0}} &= -\frac{1}{\sqrt{2}} V'_{P} p_{D}, \qquad t_{\bar{B}^{0}\to D^{0}\phi} = 0, \\ t_{\bar{B}^{0}\to D^{0}\omega} &= \frac{1}{\sqrt{2}} V'_{P} p_{D}, \qquad t_{\bar{B}^{0}_{s}\to D^{0}K^{*0}} = V'_{P} p_{D}, \end{split}$$

 V'_P : *a* common factor to all $\overline{B}^0(\overline{B}^0_s) \to D^0 V_i$ decays; P_D : the momentum of D^0 in the rest frame of \overline{B}^0 or \overline{B}^0_s .



The decay width:

$$\Gamma_{\bar{B}_{i}^{0} \to D^{0}V_{i}} = \frac{1}{8\pi M_{\bar{B}_{i}^{0}}^{2}} |t_{\bar{B}_{i}^{0} \to D^{0}V_{i}}|^{2} p_{D}.$$

Convert the total date for vector production into a mass distibution:

$$\frac{d\Gamma_{\bar{B}^0 \to D^0 \rho^0 \to D^0 \pi^+ \pi^-}}{dM_{\rm inv}} = -\frac{2m_\rho}{\pi} \times \\ \operatorname{Im}\left[\frac{1}{M_{\rm inv}^2 - m_\rho^2 + im_\rho \Gamma_\rho(M_{\rm inv})}\right] \tilde{\Gamma}_{\bar{B}^0 \to D^0 \rho^0},$$

 $M_{\rm inv}$ is the $\pi^+\pi^-$ invariant mass.

Similar fomula for $\overline{B}_{s}^{0} \rightarrow D^{0}\pi^{-}K^{+}$ decay.

Determine the common factors V_P and V'_P .

Exp. information from Belle:

[Belle, PRD76(2007)012006]

$$\begin{split} Br[\bar{B}^0 \to D^0 f_0(500)] \cdot Br[f_0(500) \to \pi^+ \pi^-] \\ &= (0.68 \pm 0.08) \times 10^{-4}, \end{split}$$

$$Br(B^0 \to \bar{D}^0 \rho^0) = (3.2 \pm 0.5) \times 10^{-4},$$

$$\tilde{V}_P = (8.8 \pm 0.5) \times 10^{-2} \text{ MeV}^{-1/2},$$
$$\tilde{V}_P' = (6.8 \pm 0.5) \times 10^{-3} \text{ MeV}^{-1/2}.$$

Make some predictions with no free parameters.

• $f_0(500)$ and ρ^0 production s in $\overline{B}{}^0 \to D^0 \pi^+ \pi^-$ decay



Invariant mass distribution for $\pi^+\pi^-$ in $\overline{B}{}^0 \to D^0 \pi^+\pi^-$ decay.

 $f_0(500)$ is clearly vi sible in the region of 400-600 MeV. A larger contributi on from ρ^0 . The shape of $M_{\pi\pi}$ distribution is similar to that of Belle.

• **Production of the scalar resonances (in S-wave)**



Invariant mass distribution for the $\pi^+\pi^-, K^+K^-, \pi^0\eta, \pi^-K^+$ in $\overline{B}{}^0 \to D^0 \pi^+\pi^-, D^0K^+K^-, D^0\pi^0\eta$ and $\overline{B}{}^0_s \to D^0\pi^-K^+$ decays.

The relative weight of the distributions are predicted.

Test some of our predictions:

$$\frac{Br[\bar{B}^0 \to D^0 f_0(980)] \cdot Br[f_0(980) \to \pi^+\pi^-]}{Br[\bar{B}^0 \to D^0 f_0(500)] \cdot Br[f_0(500) \to \pi^+\pi^-]}\Big|_{\text{Exp.}} = 0.12 \pm 0.06.$$

[Belle, PRD76(2007)012006]

$$\frac{Br[\bar{B}^0 \to D^0 f_0(980)] \cdot Br[f_0(980) \to \pi^+ \pi^-]}{Br[\bar{B}^0 \to D^0 f_0(500)] \cdot Br[f_0(500) \to \pi^+ \pi^-]}\Big|_{\text{Theo}} = 0.08,$$

with an estimated error of about 10%.

A good agreement within errors.

• The strength of K^{*0} production versus $\kappa(800)$ production in $\overline{B}_s^0 \to D^0 \pi^- K^+$ decay.



Invariant mass distribution for $\pi^- K^+$ in $\overline{B}_s^0 \to D^0 \pi^- K^+$ decay.

$$\frac{\Gamma(B^0 \to \bar{D}^0 K^+ K^-)}{\Gamma(B^0 \to \bar{D}^0 \pi^+ \pi^-)} \Big|_{\text{Exp.}} = 0.056 \pm 0.011 \pm 0.007$$

[LHCb, PRL 109(2012)131801]

$$\frac{\Gamma(B^0 \to \bar{D}^0 K^+ K^-)}{\Gamma(B^0 \to \bar{D}^0 \pi^+ \pi^-)} \Big|_{\text{Theo.}} = 0.03 \sim 0.06.$$

A qualitative agreement.

Conclusion:

Our results agree with the experimental results. This gives a strong support to the idea of the low lying scalar mesons as being formed from the interaction of pairs of pseudoscalar mesons.



Using the chiral unitary theory, we investigate the Decays of $\overline{B}{}^{0} \rightarrow D^{0} \pi^{+} \pi^{-}$, $D^{0} K^{+} K^{-}$, $D^{0} \pi^{0} \eta$ and $\overline{B}{}^{0}_{s} \rightarrow D^{0} \pi^{-} K^{+}$, and find :

The results agree with Exp., surporting $f_0(500)$, $f_0(980)$, $a_0(980)$, κ are dynamically generated from the PS-PS interaction. $f_0(500)$ is a $\pi\pi$ resonance, $f_0(980)$ is a $K\overline{K}$ molecule, $a_0(980)$ is a $\pi\eta$ molecule, κ is a π K molecule.

- The production rates of some lightscalar mesons in the weak decays of *B* mesons are predicted with no free parameters.
- More Exp. data from LHCb or other facilities are expected to test the structure of the light scalar mesons.

Thanks for your attention!