

**MeAsurement of F_2^n/F_2^p , d/u RATios and $A=3$
EMC Effect in Deep Inelastic Electron Scattering
Off the Tritium and Helium MirrOr Nuclei**

JLab MARATHON Experiment

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Hall A / C Collaboration

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Overview and Goals of the Experiment

- MARATHON will use deep inelastic scattering to measure the cross section ratio of the ^3H and ^3He mirror nuclei using the Hall A spectrometer facility of JLab.
- Using a new novel method that exploits the isospin symmetry of the $A=3$ nuclei, will measure ratio of the neutron to proton structure functions F_2^n / F_2^p , and the ratio of the proton d/u quark distribution functions.
- New data will be almost free of nuclear structure theoretical uncertainties that dominate the previous SLAC data, extracted from proton and deuteron deep inelastic scattering (DIS) measurements.
- Expt. will also measure the EMC effect for ^3H and ^3He . $A=3$ EMC data are considered crucial for the effect understanding.

Deep Inelastic Scattering and Quark Parton Model

- DIS cross section - Nucleon structure functions F_1 and F_2 :

$$\frac{d\sigma}{d\Omega dE'} = \frac{\alpha^2}{4E^2 \sin^4\left(\frac{\theta}{2}\right)} \left[\frac{F_2(\nu, Q^2)}{\nu} \cos^2\left(\frac{\theta}{2}\right) + \frac{2F_1(\nu, Q^2)}{M} \sin^2\left(\frac{\theta}{2}\right) \right]$$

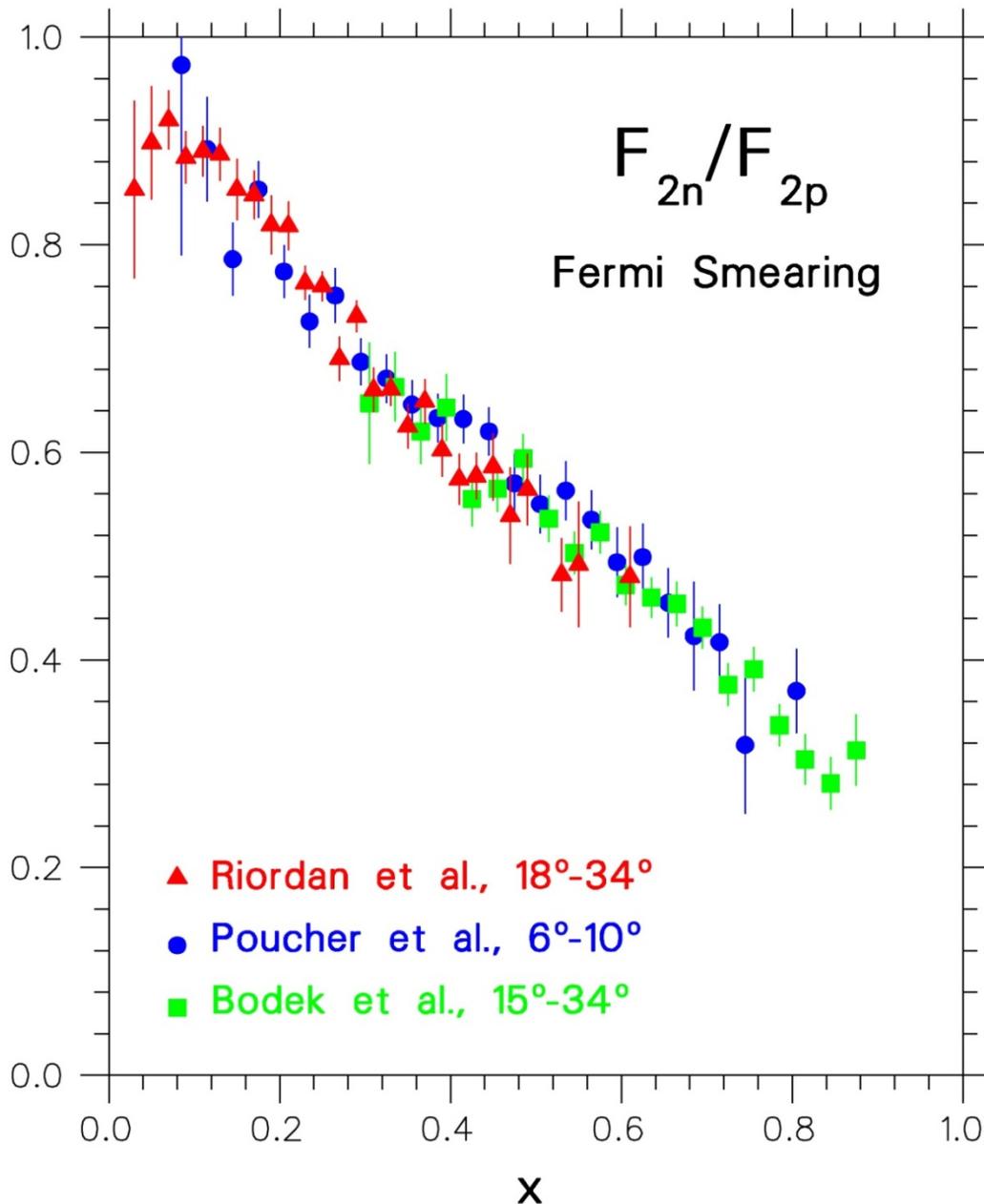
$$R = \frac{\sigma_L}{\sigma_T} = \frac{F_2 M}{F_1 \nu} \left(1 + \frac{\nu^2}{Q^2} \right) - 1 \quad \begin{array}{l} \nu = E - E' \\ Q^2 = 4EE' \sin^2(\theta/2) \end{array}$$

- QPM interpretation in terms of quark momentum probability distributions $q_i(x)$ (large Q^2 and ν , fixed x):

$$F_1(x) = \frac{1}{2} \sum_i e_i^2 q_i(x) \quad F_2(x) = x \sum_i e_i^2 q_i(x)$$

- Bjorken x : fraction of nucleon momentum carried by struck quark:

$$x = Q^2 / 2M\nu$$



SLAC 1968-1972

Friedman, Kendall, Taylor
Nobel 1991

F_2^n / F_2^p extracted from p and d DIS
using a Fermi-smearing model and
a non-relativistic N - N potential

Data in disagreement with
 $SU(6)$ prediction: $2/3=0.67!$

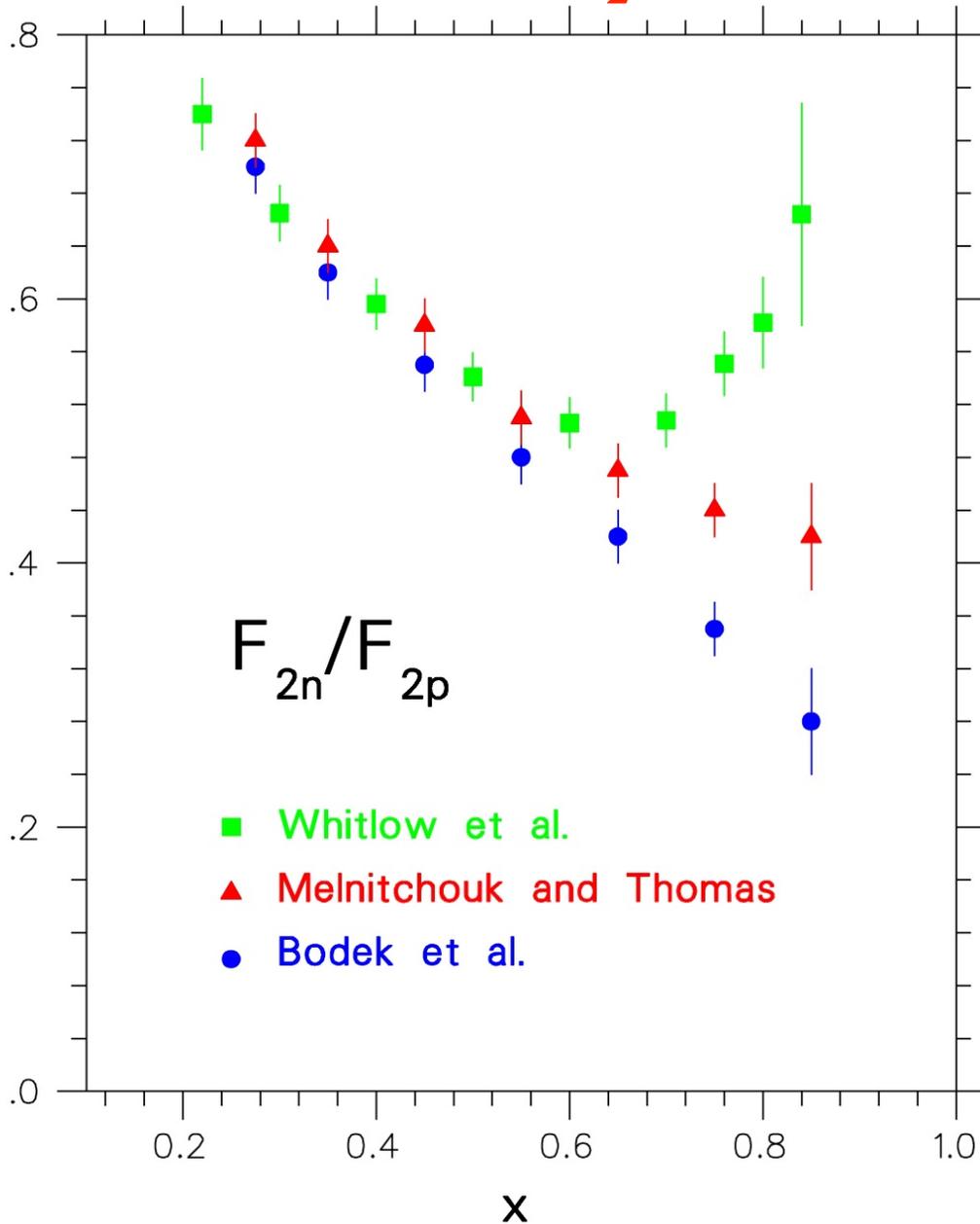
Data consistent with di-quark
model by Feynman and others

High momentum quarks in
 $p(n)$ are $u(d)$ valence quarks

There are no high momentum
strange quarks in p and n

Sea quarks dominate at small x

Nucleon F_2 Ratio Extraction Revisited



← SU(6)

$$F_{2n}^n / F_{2p}^p \rightarrow \frac{2}{3} \quad d/u \rightarrow \frac{1}{2}$$

← pQCD / Quark Counting Rules

$$F_{2n}^n / F_{2p}^p \rightarrow \frac{3}{7} \quad d/u \rightarrow \frac{1}{5}$$

← Diquark / Quark Model

$$F_{2n}^n / F_{2p}^p \rightarrow \frac{1}{4} \quad d/u \rightarrow 0$$

Nucleon F_2 Ratio Extraction from ${}^3\text{He}/{}^3\text{H}$

- Binding of nucleons in the two nuclei is of same nature. Differences between bound and free nucleons in the two nuclei is calculable, summarized, for their ratio, by some parameter R^* (W. Melnitchouk *et al.*).

$$R^* = \frac{R({}^3\text{He})}{R({}^3\text{H})} \quad \text{where} \quad R({}^3\text{He}) = \frac{F_2^{3\text{He}}}{2F_2^p + F_2^n} \quad R({}^3\text{H}) = \frac{F_2^{3\text{H}}}{F_2^p + 2F_2^n}$$

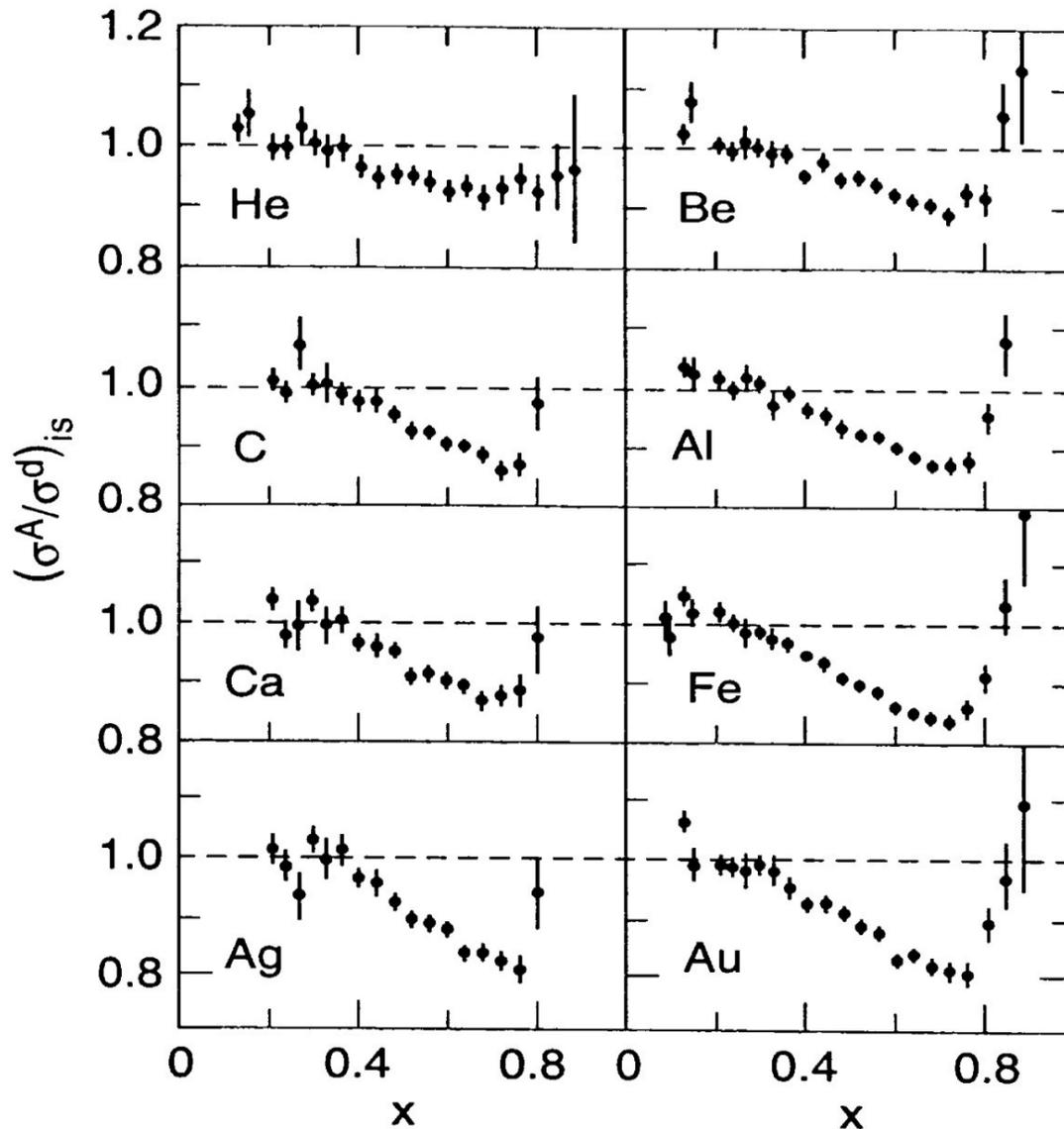
- If $R = \sigma_L / \sigma_T$ is the same for ${}^3\text{He}$ and ${}^3\text{H}$, measured DIS cross section ratio must be equal to the F_2 structure function ratio as calculated using R^* :

$$\frac{\sigma^{3\text{He}}}{\sigma^{3\text{H}}} = \frac{F_2^{3\text{He}}}{F_2^{3\text{H}}} = R^* \frac{2F_2^p + F_2^n}{F_2^p + 2F_2^n}$$

- Determine nucleon F_2 ratio using $A=3$ DIS cross section data and $R^*(\approx 1)$ from theory:

$$\frac{F_2^n}{F_2^p} = \frac{2R^* - F_2^{3\text{He}} / F_2^{3\text{H}}}{2F_2^{3\text{He}} / F_2^{3\text{H}} - R^*}$$

EMC Effect



- Nuclear F_2 structure function per nucleon is different than that of deuterium: large dependence on Bjorken x and mass A .
- No universally accepted theory for the effect explanation.
- $A=3$ data will be pivotal for understanding the EMC effect.

SLAC E-139, 1984
J. Gomez et al.

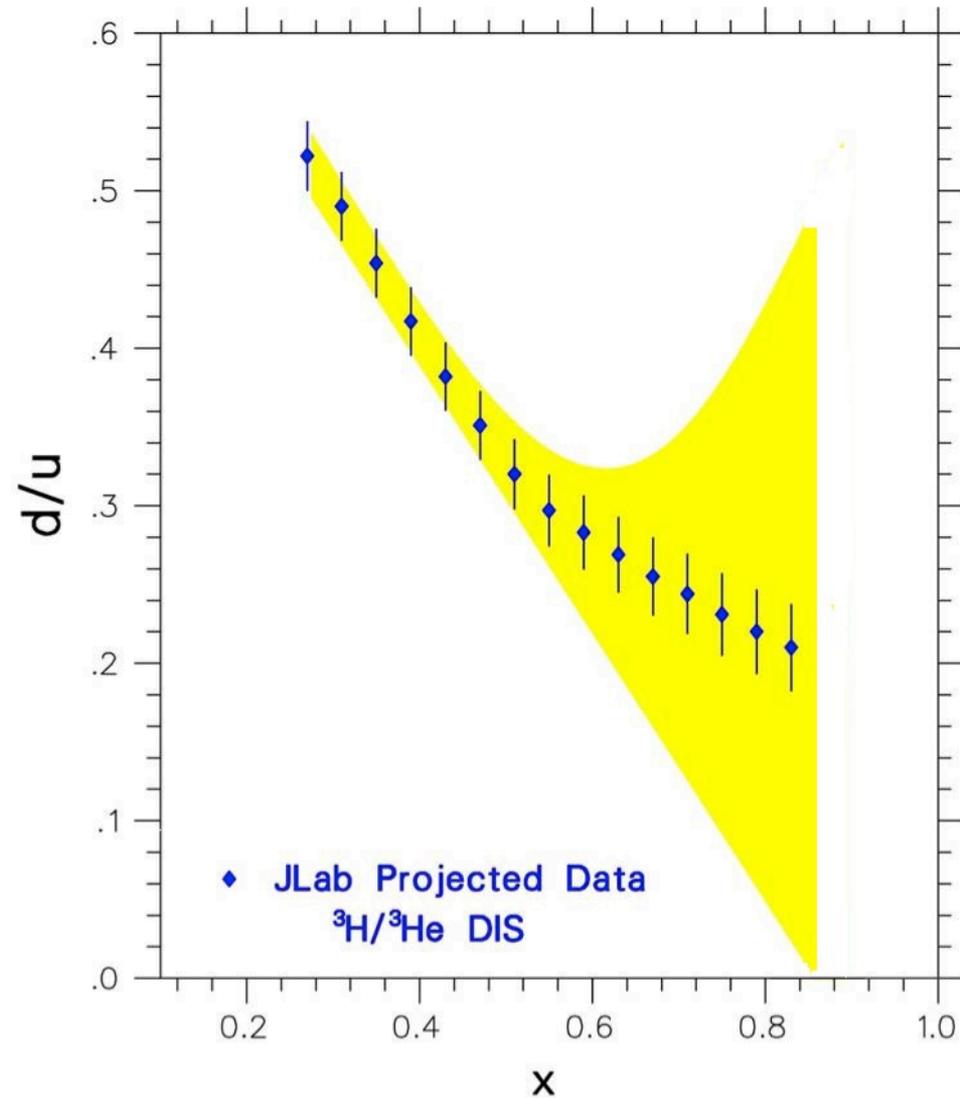
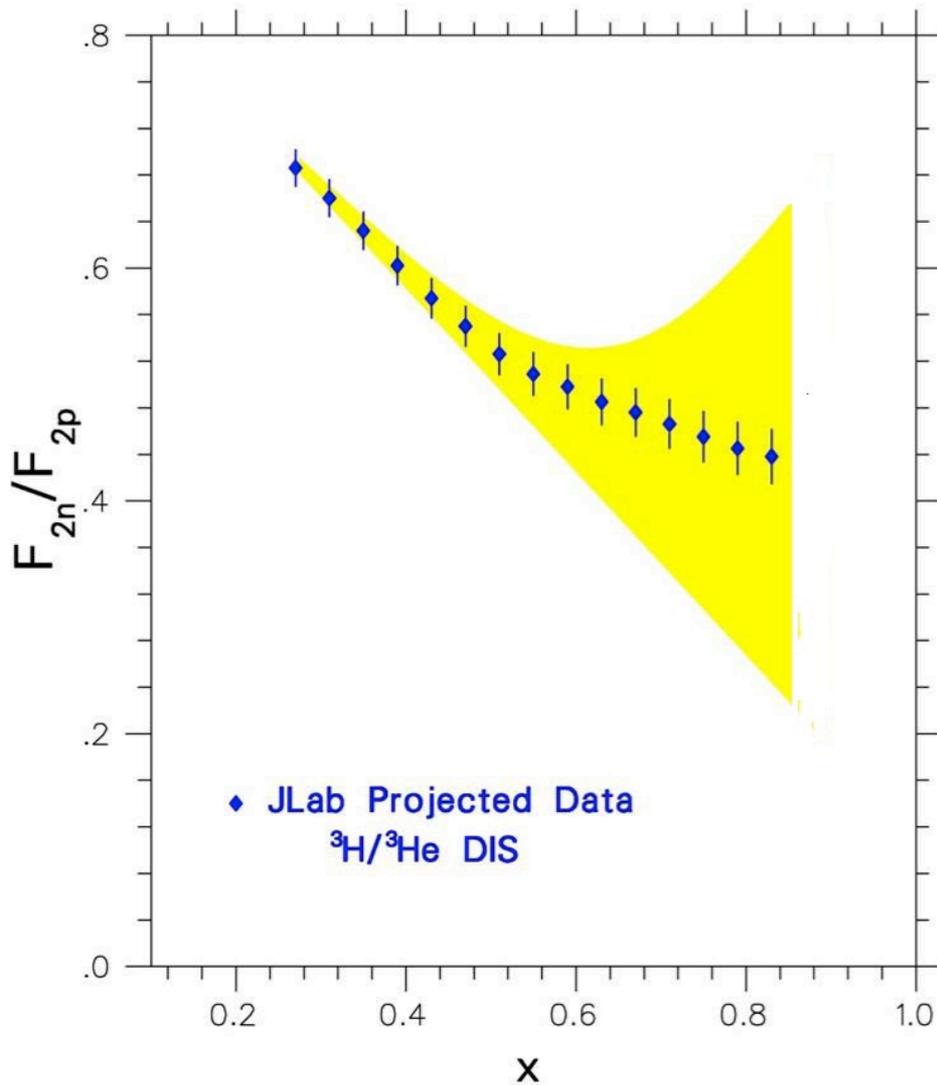
Experimental Plan and Requirements (I)

- Proposal was based on using the Left High Resolution Spectrometer (HRS) for low- x measurements and the Big Bite Spectrometer (BBS) for high- x measurements.
- *Change of plans after successful final target safety review that took place in the fall of 2015!*
- For safety reasons, the target review committee had to impose severe:
 - Restrictions in the BBS movement (change of angle)
 - Limitations in accessing the BBS components and detectors for checkout & repair purposes. Access from the front was to be totally forbidden!
- MARATHON requested to take the data with both the HRS spectrometers (instead of the Left HRS and BBS) at the expense of:
 - Eliminating the highest- x point ($x=0.87$) in the resonance region
 - Reducing the W^2 value of the highest- x DIS point ($x=0.83$) from 4.0 to about 3.5 GeV².
- Request has been approved!

Experimental Plan and Requirements (II)

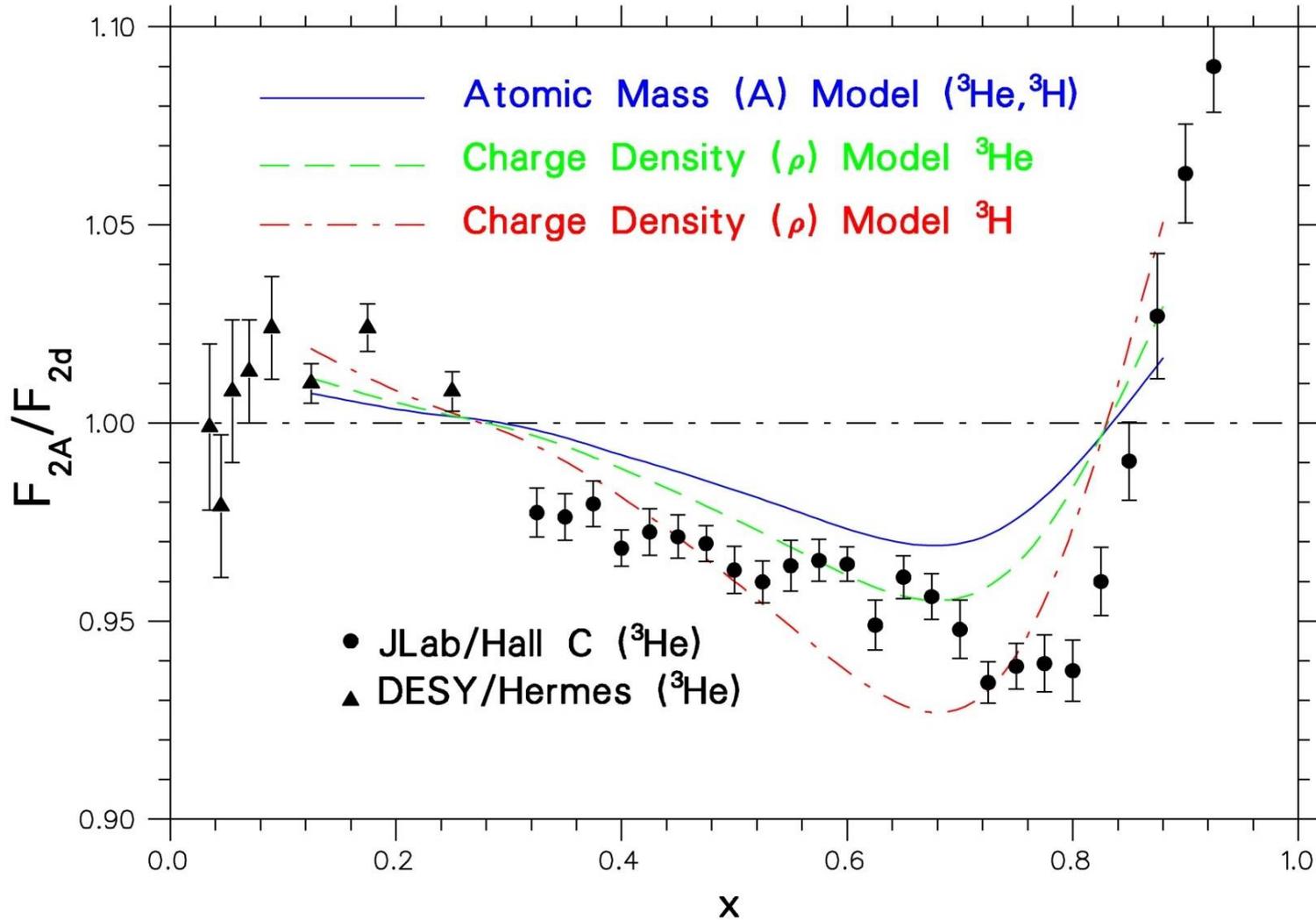
- Experiment requires (unpolarized) electron beam with the highest energy available (11 GeV). Beam current will be limited to 20 μA for target safety reasons.
- Target system is under construction. Ladder structure will have helium, tritium, deuterium, and hydrogen high pressure cells [25 cm long, 1.25 cm diameter, 14 atm (^3H), 30 atm (^1H , ^2H , ^3He)].
- HRS Systems will operate in angular range between 15° and 60° , and in momentum range between 0.5 and 4.0 GeV.
- Experiment will use the existing detector packages of the HRS systems, and standard Hall A data acquisition and online analysis and monitoring software.

Projected JLab Hall A Data for F_2^n/F_2^p and d/u Ratios



11 GeV beam; Left and Right High Resolution Hall A Spectrometers
3H and 3He gas target system; 42 days of data taking

EMC Effect for $A=3$ Mirror Nuclei



Hall A data on ^3H , ^3He will be of similar precision to Hall C data

Summary

- JLab Hall A MARATHON experiment target system has passed safety review.
- Experiment passed recently readiness review.
- Will use the two HRS's as electron detection systems.
- Is scheduled to take data in the spring of 2017.
- Will provide data on F_2^n/F_2^p and d/u with minimal theoretical uncertainties, and input for parton distribution parametrizations.
- Will provide precision data on the EMC effect for the $A=3$ nuclear systems.

Thanks!

Back Up Slides

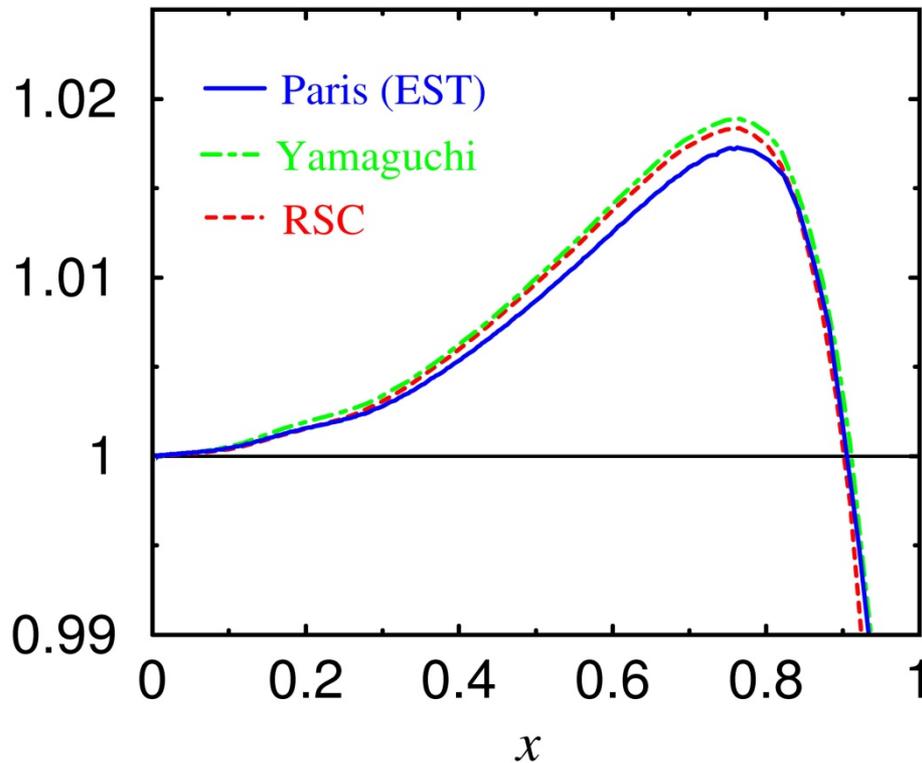
SuperRatio $R^* = R(^3\text{He})/R(^3\text{H})$ has been calculated by three expert groups to deviate from 1 only up to $\sim 1.5\%$ taking into account all possible effects:

* I. Afnan et al., Phys. Lett. B493, 36 (2000); Phys. Rev. C68, 035201 (2003)

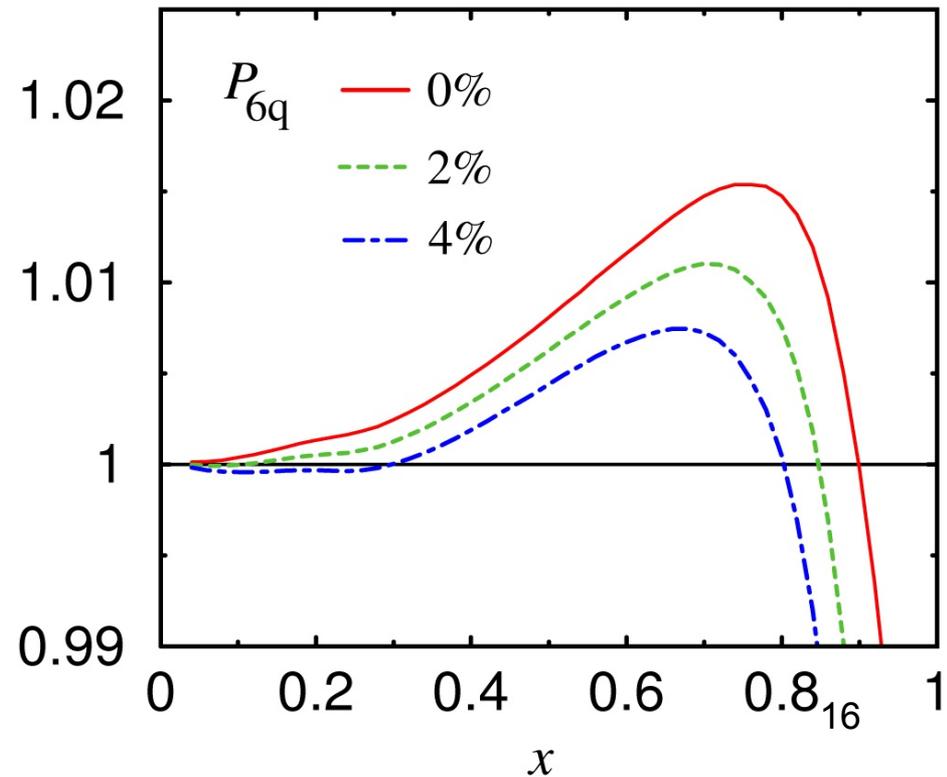
* E. Pace, G Salme, S. Scopetta, A. Kievsky, Phys. Rev. C64, 055203 (2001)

* M. Sargsian, S. Simula, M. Strikman, Phys. Rev. C66, 024001 (2002)

$R(^3\text{He}) / R(^3\text{H})$



$R(^3\text{He}) / R(^3\text{H})$



F_2^n / F_2^p in Quark Parton Model

- Assume isospin symmetry:

$$\begin{aligned} u^p(x) &\equiv d^n(x) \equiv u(x) & \bar{u}^p(x) &\equiv \bar{d}^n(x) \equiv \bar{u}(x) \\ d^p(x) &\equiv u^n(x) \equiv d(x) & \bar{d}^p(x) &\equiv \bar{u}^n(x) \equiv \bar{d}(x) \\ s^p(x) &\equiv s^n(x) \equiv s(x) & \bar{s}^p(x) &\equiv \bar{s}^n(x) \equiv \bar{s}(x) \end{aligned}$$

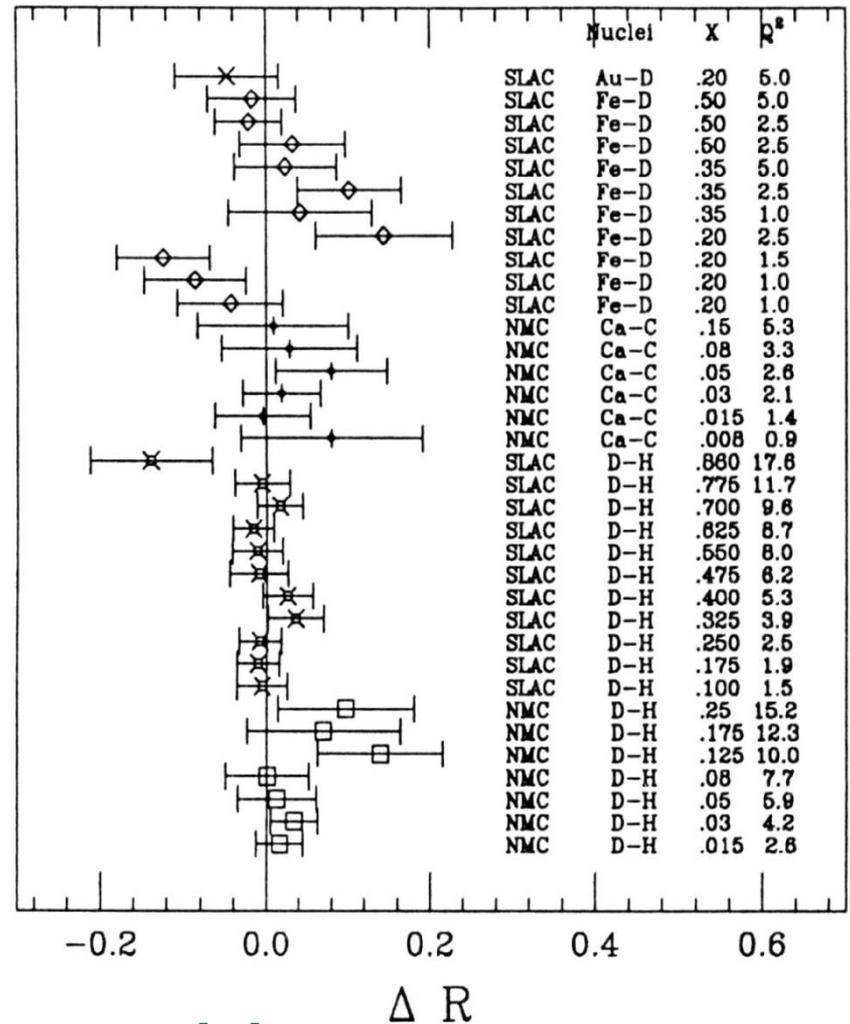
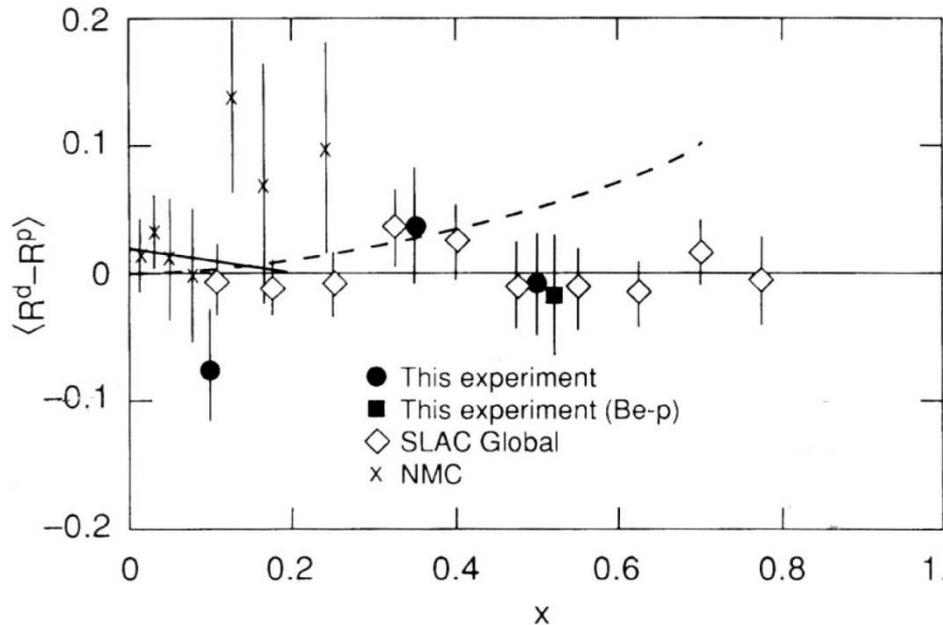
- Proton and neutron structure functions:

$$\begin{aligned} F_2^p &= x \left[\frac{4}{9} (u + \bar{u}) + \frac{1}{9} (d + \bar{d}) + \frac{1}{9} (s + \bar{s}) \right] \\ F_2^n &= x \left[\frac{4}{9} (d + \bar{d}) + \frac{1}{9} (u + \bar{u}) + \frac{1}{9} (s + \bar{s}) \right] \end{aligned}$$

- Nachtmann inequality: $1/4 \leq F_2^n / F_2^p \leq 4$

$R = \sigma_L / \sigma_T$ Measurements

SLAC/CERN data show that the ratio $R = \sigma_L / \sigma_T$ is the same for all nuclei within experimental errors



$^3\text{He}/^3\text{H}$ JLab data will be of better precision than SLAC data [wider angular range!]

F_2^n / F_2^p Deuteron Extraction Uncertainties

- Theoretically, F_2^d is a convolution of the “free nucleon” F_2^p and F_2^n functions with the deuteron wave function. Uncertainty in knowledge of the nucleon-nucleon potential introduces a large uncertainty in the determination of F_2^n (and F_2^n / F_2^p).
- Realistic relativistic calculations also show a large sensitivity in the F_2^n extraction on the nucleon binding effects in the deuteron (W. Melnitchouk and A. Thomas).
- Existence of an EMC effect in the deuteron (if the nucleon $q_i(x)$ in the deuteron are different than those of the free nucleon) would alter significantly the extracted F_2^n value. Note: recent JLab data support the existence of the EMC effect in the deuteron!

F_2^n / F_2^p , d/u Ratios and A_1 Limits for $x \rightarrow 1$

	F_2^n / F_2^p	d/u	A_1^n	A_1^p
SU(6)	2/3	1/2	0	5/9
Diquark/Feynman	1/4	0	1	1
Quark Model/Isgur	1/4	0	1	1
Perturbative QCD	3/7	1/5	1	1
Quark Counting Rules	3/7	1/5	1	1

A_1 : Asymmetry measured with polarized electrons and nucleons.
 QPM probability that the quark spins are aligned with the nucleon spin.

A_1^p, A_1^n : Extensive experimental programs at CERN, SLAC, DESY and JLab (6 GeV and 12 GeV Programs)

EMC Effect

- Nuclear F_2 structure function per nucleon is different than that of deuterium: large dependence on Bjorken x and mass A .
- Quark distribution functions modified in the nuclear medium.
- Possible explanations include:
 - Binding effects beyond nucleon Fermi motion
 - Enhancement of pion field with increasing A
 - Influence of possible multi-quark clusters
 - Change in the quark confinement scale in nuclei
 - Local environment and density dependence
- No universally accepted theory for the effect explanation.
- $A=3$ data will be pivotal for understanding the EMC effect.
- **Theorists:** Ratio of EMC effect for ${}^3\text{H}$ and ${}^3\text{He}$ is the best quantity for quantitative check of the theory, free of most uncertainties.

Deep Inelastic Scattering and Quark Parton Model

- DIS cross section - Nucleon structure functions F_1 and F_2 :

$$\frac{d\sigma}{d\Omega dE'} = \frac{4\alpha^2}{Q^4} \left[\frac{F_2(\nu, Q^2)}{\nu} \cos^2\left(\frac{\theta}{2}\right) + \frac{2F_1(\nu, Q^2)}{M} \sin^2\left(\frac{\theta}{2}\right) \right]$$

$$R = \frac{\sigma_L}{\sigma_T} = \frac{F_2 M}{F_1 \nu} \left(1 + \frac{\nu^2}{Q^2} \right) - 1$$

$$\frac{d\sigma}{d\Omega dE'} = \frac{4\alpha^2}{Q^4} \cos^2\left(\frac{\theta}{2}\right) F_2(\nu, Q^2) \left[\frac{1}{\nu} + \frac{(1 + Q^2/\nu^2)}{xM(1+R)} \tan^2\left(\frac{\theta}{2}\right) \right]$$

$$\begin{aligned} \nu &= E - E' \\ Q^2 &= 4EE' \sin^2(\theta/2) \\ x &= Q^2 / 2M\nu \end{aligned}$$

- QPM interpretation in terms of quark momentum probability distributions $q_i(x)$ (large Q^2 and ν , fixed x):

$$F_1(x) = \frac{1}{2} \sum_i e_i^2 q_i(x) \quad F_2(x) = x \sum_i e_i^2 q_i(x)$$

SLAC/CERN Data Interpretation in QPM

- Nachtmann inequality satisfied: $1/4 \leq F_2^n / F_2^p \leq 4$

- For $x \rightarrow 0 : F_2^n / F_2^p \rightarrow 1$: Sea quarks dominate with:

$$u + \bar{u} = d + \bar{d} = s + \bar{s}$$

- For $x \rightarrow 1 : F_2^n / F_2^p \rightarrow 1/4$: High momentum partons in **proton** (**neutron**) are **up** (**down**) quarks, and:

$$s + \bar{s} = 0$$

- For medium and high x , safe to assume that (with d and u denoting now quark plus antiquark distributions):

$$\frac{F_2^n}{F_2^p} = \frac{[1 + 4(d/u)]}{[4 + (d/u)]}$$

Theory Overview (I)

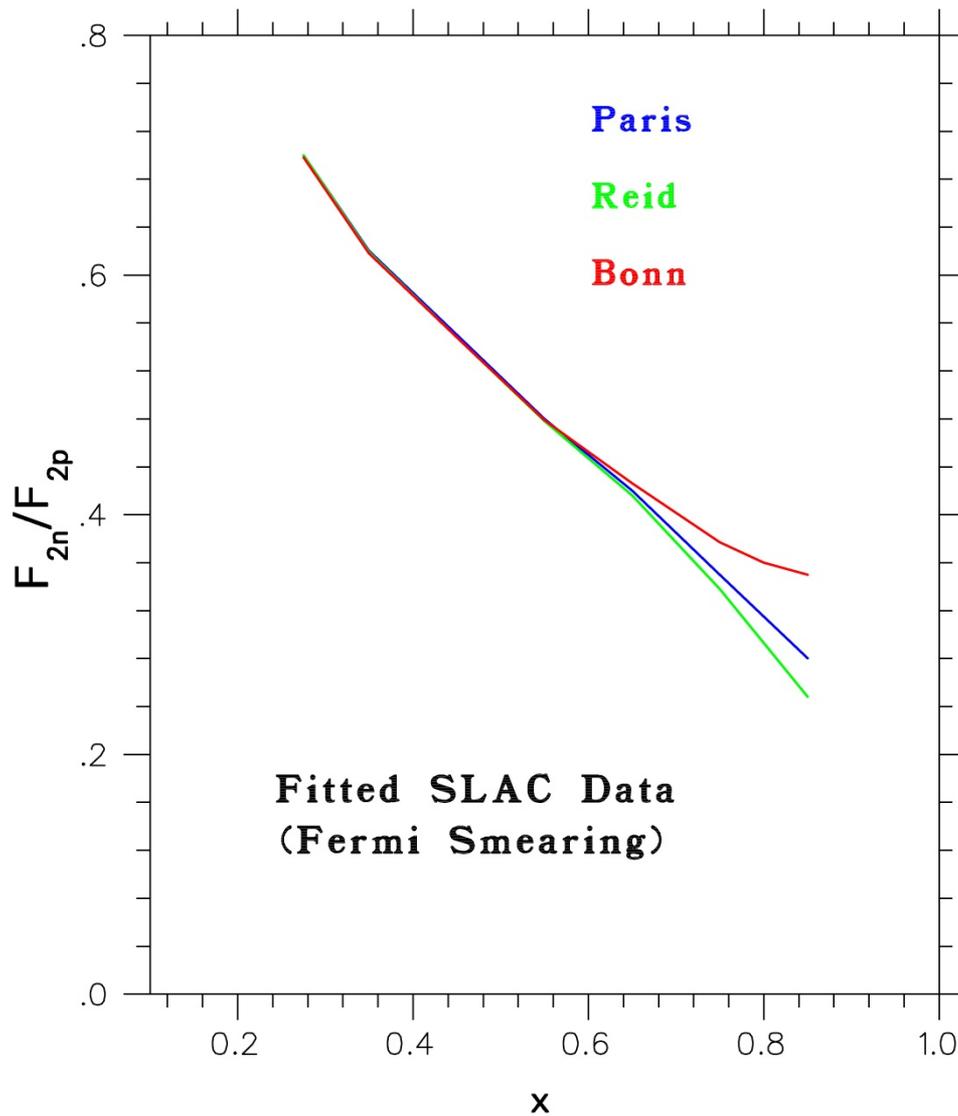
- SU(6)-breaking Diquark Model (Close/Feynman, 1970's)
 - Phenomenological explanation: fall off of F_2^n/F_2^p and its $x=1$ limit easily explainable using SU(6) nucleon wave function (Close, 1973) or Regge theory (Feynman, 1975) and assuming that the spectator diquark configuration with total spin $s=1$ is suppressed relative to that with $s=0$.
 - Limit at $x=1$ in agreement with the SLAC/CERN data.
- Constituent Quark Model (Isgur and Carl, 1980-2000)
 - Suppression fully understood in the hyperfine-perturbed quark model: color-magnetic hyperfine interaction generated by gluon exchange is proportional to quark spin dot products $\mathbf{s}_i \cdot \mathbf{s}_j$. Force is repulsive(attractive) for parallel(antiparallel) spins and allows the $u(x)$ and $d(x)$ quark momentum distributions to be different.
 - Limit for $x=1$ in agreement with the diquark model and data.

Theory Overview (II)

- Perturbative QCD (Farrar and Jackson, 1975)
 - Simple consideration of isospin and helicity structure of nucleon's quark wave function within perturbative QCD dictates that for diquark spin alignment ($s=1$) only exchange of longitudinal gluons is permitted, resulting in suppression of Compton scattering amplitude, which causes the F_2^n/F_2^p fall off.
 - Problem: pQCD $x=1$ limit different than that of quark/diquark model!
- QCD Counting Rules (Brodsky *et al.*, 1995)
 - Early pQCD findings substantiated within first-principles calculations of Counting Rules of QCD predicting the high- x behavior of quark distributions: $q(x) \sim (1-x)^{2n-1+2\Delta h}$, where n is the minimum number of not interacting (spectator) quarks and Δh is the difference in helicity between the struck quark and the nucleon.
 - Same $x=1$ limit as pQCD.

The Tritium Target System

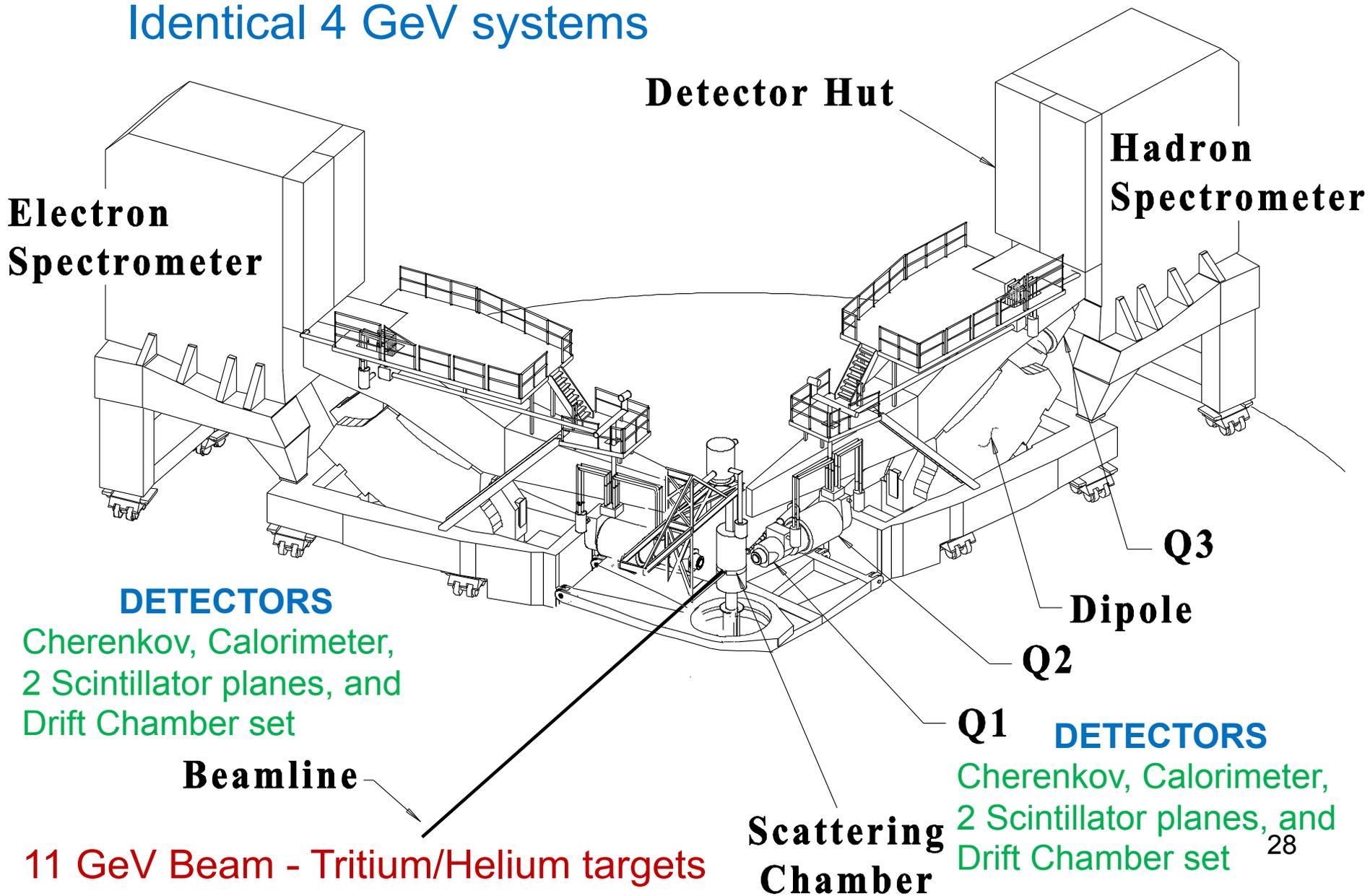
- Five-cell target structure (^1H , ^2H , ^3H , ^3He , Al Dummy):
 - All cells: 25 cm long, 1.25 cm diameter, 20 μA max current
 - ^3H cell: 10 atm, 1100 Ci activity
 - ^1H , ^2H , ^3He cells: 25 atm
- Cells will be filled at the Savannah River National Laboratory.
- Target will be shipped to JLab using commercial services.
- System will be enclosed in secondary “sealed” containment vessel (no reliance on fast valves).
- A catastrophic leak will be vented through stack above Hall A.
- Resulting emission will keep tritium levels at the JLab site boundary below regulatory limits.
- Tritium target project has passes a safety review.

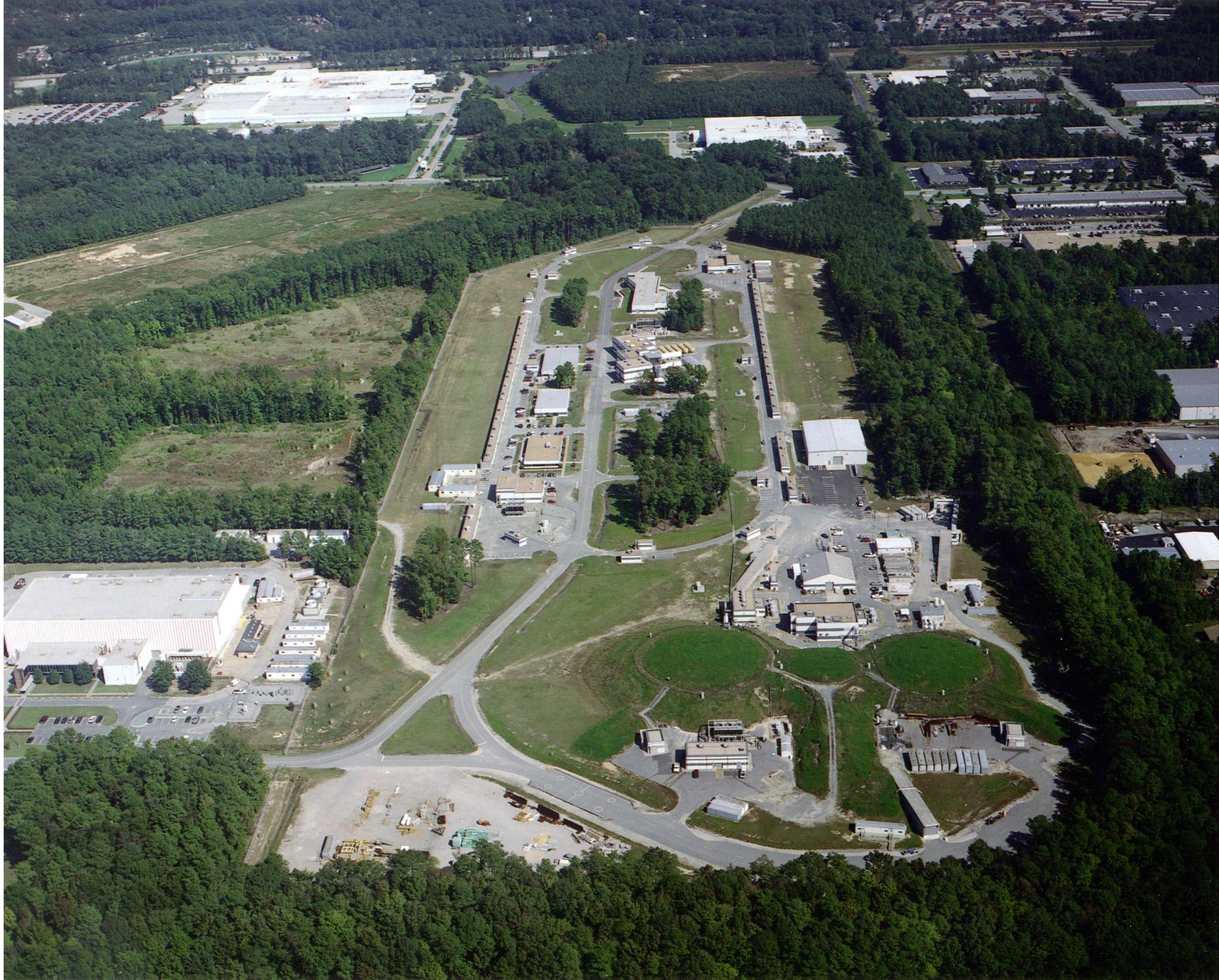


Standard analysis of F_2^n / F_2^p from proton and deuteron data shows large uncertainty at high-x values depending on nucleon-nucleon potential used.

JLab Hall A High Resolution Spectrometers

Identical 4 GeV systems





^3He and ^3H Structure Functions

- Nucleon structure function in nucleus A in the Impulse Approximation: extension of deuteron model (Melnitchouk and Thomas):

$$F_2^{NA}(x) = \int dy [f(y)]_{N/A} F_2^N(x/y) \equiv f_{N/A} \otimes F_2^N$$

- For ^3He :
$$F_2^{^3\text{He}} = 2 f_{p/^3\text{He}} \otimes F_2^p + f_{n/^3\text{He}} \otimes F_2^n$$

- With isospin symmetry:
$$f_{n/^3\text{H}} = f_{p/^3\text{He}} \equiv f_p$$

$$f_{p/^3\text{H}} = f_{n/^3\text{He}} \equiv f_n$$

- Then for ^3H :
$$F_2^{^3\text{H}} = f_n \otimes F_2^p + 2 f_p \otimes F_2^n$$

- Charge symmetry breaking effects are small!