

# The $n\Lambda$ scattering length and the $nn\Lambda$ resonance

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## Brief Outline

- What do we know about  $n\Lambda$  scattering?
- Experimental data in hand. ( $p\Lambda$ , CSB)
- An  $nn\Lambda$  bound state? A resonance?
- What is possible theoretically? Our model.
- $nn\Lambda$  resonance properties.
- Summary
- JLab  ${}^3\text{H}(e,e' K^+)nn\Lambda$  tritium experiment.

- Nothing about  $n\Lambda$  scattering has been measured.
- We have inferred something from  $p\Lambda$  scattering plus the binding energies of few-body  $\Lambda$  hypernuclei.
- Starting from the sparse  $p\Lambda$  data we have inferred a Charge Symmetry Breaking (CSB) difference between  $p\Lambda$  and  $n\Lambda$  interactions from the  $A=4$  mirror hypernuclei binding. However, we do **not** know whether the observed CSB arises from the fundamental two-body  $N\Lambda$  interaction or from a possible  $NN\Lambda$  three-body force.
- The HypHI collaboration reported seeing a  ${}^3_{\Lambda}n$  bound state. Given our knowledge of the  $nn$  interaction, a  ${}^3_{\Lambda}n$  bound state would provide a strong  $n\Lambda$  interaction constraint. Moreover, JLab would be an ideal facility to check the claim using the  ${}^3\text{H}(e,e'K^+){}^3_{\Lambda}n$  reaction.

Theoretically, the possibility of a  ${}^3_{\Lambda}n$  bound state seems remote. However, the coming tritium experiment by Tang and collaborators should provide a definitive answer.

Question: How can we model the  $n\Lambda$  interaction, when we have only limited data regarding  $p\Lambda$  scattering?

Few-body hypernuclei:

- $\Lambda$  hypernuclei provide weak constraints
  - ${}^3_{\Lambda}\text{H}$  is weakly bound [ $B_{\Lambda}({}^3_{\Lambda}\text{H}) = 0.13 \pm 0.05$  MeV]; small separation energy implies that  ${}^3_{\Lambda}\text{H}$  is one of our largest halo nuclei.
  - The  $A=4$  isodoublet seems to exhibit significant Charge Symmetry Breaking, some 2-3 times that in the  ${}^3\text{H}$ - ${}^3\text{He}$  isodoublet.
  - The uncertainty in the sparse  $p\Lambda$  data implies a potentially wide range of variation in the  $n\Lambda$  interaction.
- Recent experiments have suggested a decrease in the apparent size of the  $A=4$  CSB.
  - Esser *et al.*, PRL **114**, 232501 (2015) report a value for  $B_{\Lambda}({}^4_{\Lambda}\text{H})$  of  $2.12 \pm 0.01(\text{stat.}) \pm 0.09(\text{syst.})$ .
  - Yamamoto *et al.*, PRL **115**, 222501 (2015) have measured the gamma decay in  ${}^4_{\Lambda}\text{H}$  and obtained a value of  $1.406 \pm 0.02 \pm 0.02$  MeV, indicating that CSB in the  $1^+$  excited states is quite small.

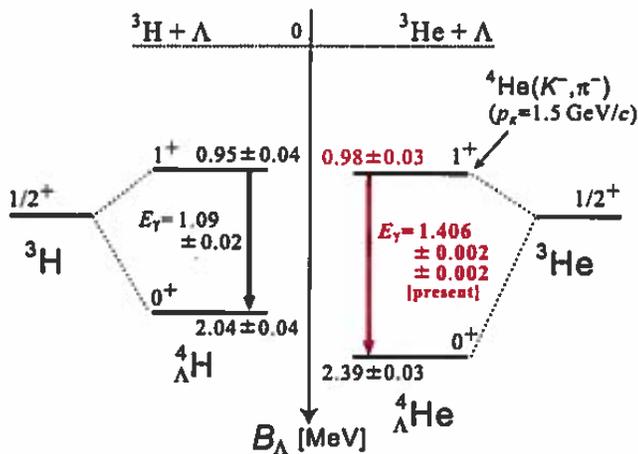
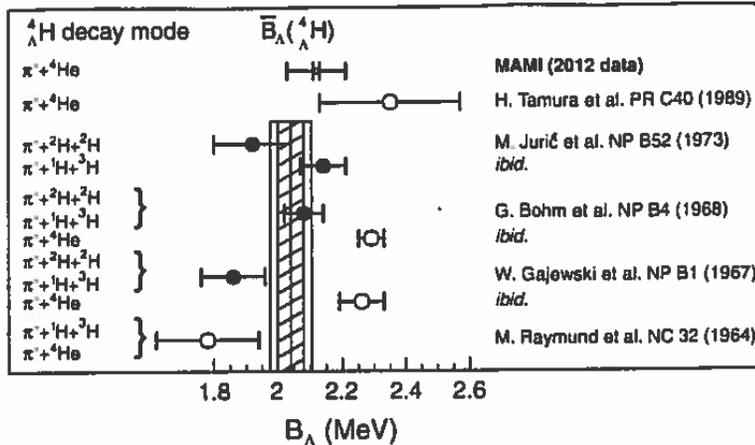


FIG. 1 (color online). Level schemes of the mirror hypernuclei,  ${}^4_{\Lambda}\text{H}$  and  ${}^4_{\Lambda}\text{He}$ .  $\Lambda$  binding energies ( $B_{\Lambda}$ ) of  ${}^4_{\Lambda}\text{H}(0^+)$  and  ${}^4_{\Lambda}\text{He}(0^+)$  are taken from past emulsion experiments [2].  $B_{\Lambda}({}^4_{\Lambda}\text{He}(1^+))$  and  $B_{\Lambda}({}^4_{\Lambda}\text{H}(1^+))$  are obtained using the present data and past  $\gamma$ -ray data [6–8], respectively. Recently,  $B_{\Lambda}({}^4_{\Lambda}\text{H}(0^+)) = 2.12 \pm 0.01(\text{stat}) \pm 0.09(\text{syst}) \text{ MeV}$  was obtained with an independent technique [5].

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Nevertheless, we do not understand CSB in the mirror  $\Lambda$  hypernuclei. Therefore, we cannot accurately model the  $n\Lambda$  interaction.

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A  ${}^3_{\Lambda}\text{n}$  hypernucleus?

- The HypHI collaboration reported a bound  $nn\Lambda$  system ( ${}^3_{\Lambda}\text{n}$ ).
  - C. Rappold *et al.*, Phys. Rev. C **88**, 041001(R) (2013).
  - They observed both two-body and three-body decay modes.
  - ${}^3_{\Lambda}\text{n}$  would be the lightest neutron-rich hypernucleus observed.
  - Such a bound state would provide a significant constraint on the  $n\Lambda$  interaction, because the  $nn$  interaction is well known.
  - Such a bound state could be observed directly in a  ${}^3\text{H}(e,e'\text{K}^+){}^3_{\Lambda}\text{n}$  experiment at JLab, although a weakly bound system would imply a small cross section.
  - Alternative reactions at J-PARC would be  ${}^3\text{H}(\text{K}^-, \pi^0){}^3_{\Lambda}\text{n}$  and  ${}^3\text{He}(\text{K}^-, \pi^+){}^3_{\Lambda}\text{n}$ . The latter, being a double-charge-exchange reaction, suggests a very small cross section.

# A ${}^3_{\Lambda}\text{n}$ hypernucleus?

A  ${}^3_{\Lambda}\text{n}$  bound state has been strongly questioned:

- H. Garcilazo and A. Valcarce, Phys. Rev. C **89** 057001 (2014).
- E. Hiyama *et al.*, Phys. Rev. C **89** 061302 (2014).
- A. Gal and H. Garcilazo, Phys. Lett. **B736**, 93 (2014).

Simple physics suggests that one would not expect a bound state.

- The hypertriton is barely bound and its core is a deuteron.
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Could there instead exist an  $nn\Lambda$  three-body resonance?

If so, then one could still utilize the JLab electro-production reaction (or perhaps the J-PARC strangeness-exchange reaction) to constrain the  $n\Lambda$  interaction.

# Where do we stand?

## The Science

- *No published*  $n\Lambda$  data exist ! Our numerous  $N\Lambda$  potential models have never been tested against  $n\Lambda$  data.
- Were a bound  ${}^3_{\Lambda}\text{n}$  hypernucleus to exist, our knowledge of the  $nn$  interaction would permit us to significantly constrain the low-energy properties of the  $n\Lambda$  system. *Existence of such a bound state has been thoroughly questioned theoretically !*
- However, either a strong sub-threshold resonance or an actual physical resonance in the  $nn\Lambda$  system should exist. The resonance position and the shape of the spectrum in a  $K^+$  electro-production measurement from a tritium target at JLab leading to an  $nn\Lambda$  final state would provide a significant constraint upon the low-energy properties of the heretofore unmeasured  $n\Lambda$  interaction.

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## The Impact

- Understanding  $\Lambda$  hypernuclei (ground states and single-particle spectra) requires knowledge of the  $n\Lambda$  interaction. Scattering data exist only for  $p\Lambda$ .
- Data on an  $nn\Lambda$  (sub-threshold or physical) resonance would constrain the critical low-energy properties of the  $n\Lambda$  interaction.
- $n\Lambda$  knowledge would elucidate Charge Symmetry Breaking in the  $N\Lambda$  interaction, provide a realistic basis for understanding the  $\Lambda$ -hypernuclei data, enhance our calculations of neutron-rich hypernuclei, and constrain our modeling of neutron stars.

# Our $nn\Lambda$ Three-Body Model

- pairwise s-wave interactions of rank one separable form

$$V(k, k') = g(k)Cg(k') ; \quad g(k) = 1/(k^2 + \beta^2)$$

- $nn$  potential strength and range  $\sim$  effective range parameters:

$$a_{nn} = -18.9 \pm 0.4 \text{ fm and } r_{nn} = 2.75 \pm 0.11 \text{ fm}$$

- $n\Lambda$  strength and range fitted to the Nijmegen model D values:

$$a_s = -2.03 \pm 0.32 \text{ fm and } r_s = 3.66 \pm 0.32 \text{ fm}$$

$$a_t = -1.84 \pm 0.10 \text{ fm and } r_t = 3.32 \pm 0.11 \text{ fm}$$

M. M. Nagels, T. A. Rijken, & J. J. deSwart, PRD **15**, 2547 (1977)

- Separable potentials allow us to (simply) analytically continue onto the second sheet of the energy plane.
- We search for resonance poles by examining the eigenvalue spectrum of the kernel of the Faddeev equations for the  $nn\Lambda$  system.
- We previously used such a technique to explore  $\Lambda - d$  scattering:  
I. R. Afnan and B. F. Gibson, PRC **47**, 1000 (1993).

# Searching for Resonances in the $nn\Lambda$ System

We must analytically continue the Faddeev equations onto the second energy sheet.

- For two identical Fermions interacting via Yamaguchi pairwise potentials, the homogeneous integral equation is of the form

$$\lambda_n(E) \phi_{n,k_\alpha}(q, E) = \sum_{k_\beta} \int_0^\infty Z_{k_\alpha, k_\beta}^{JT}(q, q'; E) \tau_{k_\beta}[E - \epsilon_\beta(q')] \\ \times \phi_{n; k_\beta}(q', E) q'^2 dq' .$$

- We analytically continue onto the second energy sheet by utilizing the transformation

$$q \rightarrow q e^{-i\theta} \quad q' \rightarrow q' e^{-i\theta} \quad \text{with} \quad \theta > 0 . \quad (1)$$

- The rotation angle  $\theta$  is limited by kernel singularities . The Born amplitude  $Z_{k_\alpha, k_\beta}^{JT}$  requires that  $\theta < \frac{\pi}{2}$ , which gives us the region  $\Im(E) < 0$  on the second Riemann sheet. The other source of singularities is the quasi-particle propagator  $\tau_{k_\beta}[E - \epsilon_\beta(q')]$ ; because there are no two-body bound states, the rotation is not limited.

# Results of the Eigenvalue Search

Consider the specific example utilizing the  $nn$  and the  ${}^1S_0$  and  ${}^3S_1$   $n\Lambda$  potentials defined previously.

- $nn$  potential strength and range  $\sim$  effective range parameters:

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1. We searched the complex energy plane for the largest eigenvalue of the kernel = **1** and found a pole at:

$$E = -0.154 - 0.753 i \text{ MeV} \quad \text{with eigenvalue} \quad \lambda(E) = 1.0000 - 0.0001 i .$$

2. Because  $\Re(E) < 0$ , this pole corresponds to a sub-threshold resonance, one that lies below the breakup threshold in a region inaccessible by experiment.

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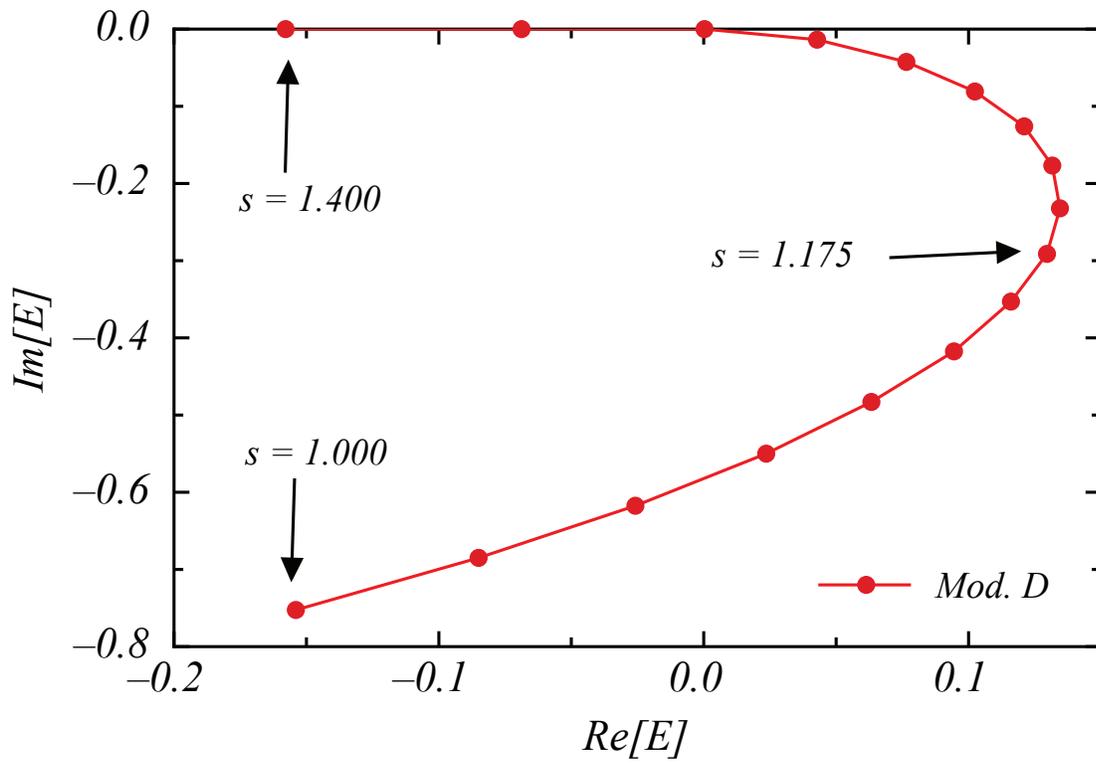
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- Because  $\Re(E) < 0$ , this pole corresponds to a sub-threshold resonance, one that lies below the breakup threshold in a region inaccessible to experiment.

Because the pole lies below the breakup threshold, we ask how easily the pole can be converted into a physical resonance (or bound state).

- We scale the strength of the  $^1S_0$  and  $^3S_1$   $n\Lambda$  potentials by  $s$ .
- We follow the path of the pole as it turns into a "resonance" and then into a bound state.
- An increase in strength of  $\sim 7\%$  produces a physical resonance, one above the three-body breakup threshold.
- An increase of  $35\%$  produces a  $^3_\Lambda n$  bound state.

# Trajectory of the $nn\Lambda$ "Resonance" Pole

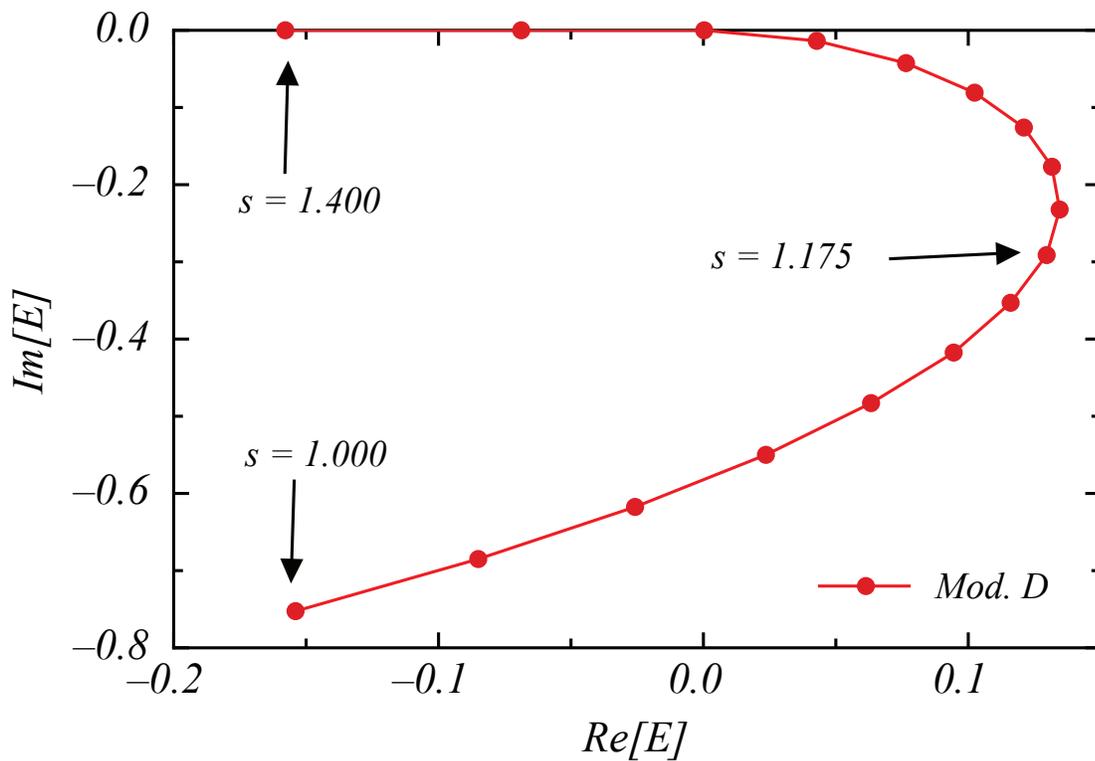


## Trajectory of the $nn\Lambda$ "Resonance" Pole

In the figure one follows the trajectory of the "resonance" pole as the strength  $s$  of the  $n\Lambda$  interaction is increased from a value of 1.0 in increments ( $\Delta s$ ) of 0.025. One starts from a sub-threshold resonance at values of  $s = 1.000$  up to  $s = 1.050$ . For  $s = 1.075$  up to  $s = 1.350$  we obtain a physical resonance; in particular, we obtain a resonance with  $E = 0.129 - 0.291i$  MeV at  $s = 1.175$ . As  $s$  is further increased, we obtain a bound state with energy  $E = -0.069$  MeV at  $s = 1.375$  and  $E = -0.158$  at  $s = 1.400$ . Thus, for this particular model one can see that an  $n\Lambda$  potential whose parameters lie within the uncertainty of the observed low energy  $p\Lambda$  scattering parameters could produce a physical resonance in the  $nn\Lambda$  system.

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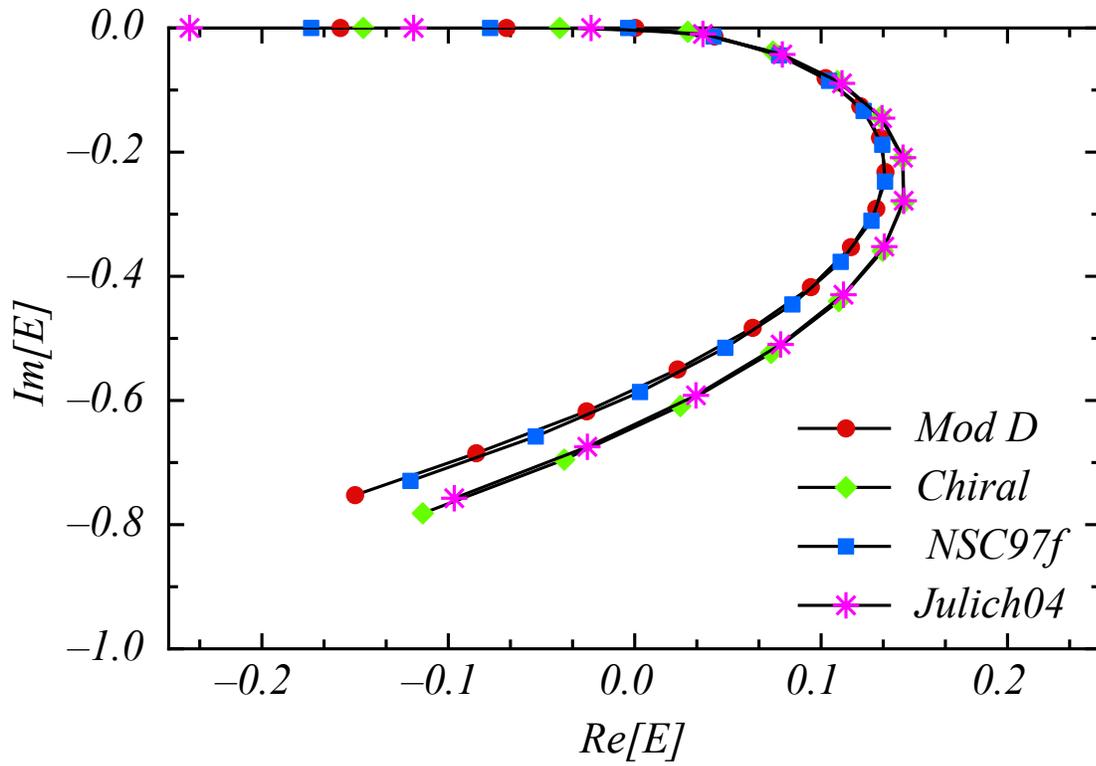
# Trajectory of the $nn\Lambda$ "Resonance" Pole for Four Contemporary Models

To explore the sensitivity to four rather different  $N\Lambda$  potential models

- M. M. Nagels, T. A. Rijken, and J. J. de Swart, "Baryon-baryon scattering in a one-boson-exchange-potential approach, II. Hyperon-nucleon scattering", *Phys. Rev. D* **15**, 2547 (1977). (*ModD*)
- T. A. Rijken, V. G. J. Stoks, Y. Yamamoto, "Soft-core hyperon-nucleon potentials", *Phys. Rev. C* **59**, 21 (1999). (*NSC97f*)
- J. Haidenbauer and U.-G. Meißner, "Jülich hyperon-nucleon model revisited", *Phys. Rev. C* **72**, 044005 (2005). (*Jülich04*)
- J. Haidenbauer, *et al.*, "Hyperon-nucleon interaction at next-to-leading order in chiral effective field theory", *Nucl. Phys. A* **915**, 24 (2013). (*Chiral*)

we repeated the eigenvalue search for each of the potential models. None of the four models produces an observable resonance; each produces a sub-threshold resonance, which lies below the  $nn\Lambda$  threshold. A change in  $s$  of as little as 5% produces a resonance above the three-body threshold, but a change of at least 25% is required to produce a bound  ${}^3_{\Lambda}n$ .

# Trajectory of the $nn\Lambda$ "Resonance" Pole for Four Contemporary Models



# Summary

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*No published  $n\Lambda$  data exist!* Our multiple  $N\Lambda$  potential models have never been tested against  $n\Lambda$  data. Were a bound  ${}^3_{\Lambda}\text{n}$  hypernucleus to exist, our knowledge of the  $nn$  interaction would permit us to significantly constrain the low-energy properties of the  $n\Lambda$  system. *Existence of such a bound state has been strongly questioned theoretically!* However, either a **strong sub-threshold resonance** or an **actual physical resonance** in the  $nn\Lambda$  system should exist. The resonance position and the shape of the spectrum in a  $K^+$  electro-production measurement from a tritium target at JLab leading to a  $n\text{-}n\text{-}\Lambda$  final state would provide a significant constraint upon the low-energy properties of the heretofore unmeasured  $n\Lambda$  interaction.

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## The Impact

Understanding  $\Lambda$  hypernuclei (ground states and single-particle spectra) measured at facilities around the world (JLab, RHIC, Mainz, J-PARC, DaΦne) requires knowledge of the  $n\Lambda$  interaction. Scattering data exist only for  $p\Lambda$ . Data on an  $nn\Lambda$  (sub-threshold or physical) resonance would enhance our knowledge of the critical low-energy properties of the  $n\Lambda$  interaction. Such knowledge would elucidate the important Charge Symmetry Breaking relative to the measured  $p\Lambda$  interaction, provide a realistic basis for understanding the long existing  $\Lambda$  hypernuclei data, enhance our calculations of neutron-rich hypernuclei, and constrain our modeling of neutron stars.

# Immediate Future: JLab Tritium Experiment

*Determining the Unknown  $\Lambda$ -n Interaction by Investigating the  $\Lambda_{nn}$  Resonance*

**Spokespersons:** L. Tang, F. Garibaldi, P.E.C. Markowitz, S.N. Nakamura, J. Reinhold

The availability of a  $T_2$  gas target in Hall A provides a unique opportunity to measure the n-n- $\Lambda$  three-body resonance. Only at JLab with the  $(e,e' K^+)$  reaction can the resonance position (excitation energy) and width be measured with the required precision. It is anticipated that using a 4.4 GeV beam one can achieve an energy resolution of  $\sim 2$  MeV FWHM and an absolute missing mass precision of  $\sim \pm 0.20$  MeV. This data will permit use of a theoretical pole search technique to determine experimentally for the first time the n $\Lambda$  interaction utilizing together the JLab improved p $\Lambda$  scattering data and the  $\Lambda$  hypernuclei Charge Symmetry Breaking data. Scheduled run beginning in mid October.