

# Binding, Bonding and Charge Symmetry Breaking in $\Lambda$ -Hypernuclei

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# Outline of the Talk

- Introduction
- A single Mass formula for Normal nuclei, Strange & Multi-Strange Hypernuclei.
- The  $\Lambda$ ,  $\Lambda\Lambda$  and  $\Xi^-$  Binding Energy.
- Lambda-Lambda Bonding Energy.
- Charge Symmetry Breaking
- Summary

# How to find the binding energy of a hypernucleus?

- Experiment
- Microscopic Theory
  - RMF
  - QMF
  - W-S Potential
- Phenomenology
  - Mass Formula
    - Provides a quick check of the microscopic calculations.
    - Can extrapolate to a wider mass region - beyond the domain of the microscopic calculations.
    - Can provide an estimation of the hyperon(s) binding energy to guide the experiments.
    - Can be used as an input to calculate hypernuclear production cross section(s).

# Generalized Mass Formula For Non-Strange & Strange nuclei

*C. Samanta et al., J. Phys. G 32 (2006) 363*

A systematic search using experimental separation energy ( $S_Y$ ) for  $\Lambda^0$ ,  $\Lambda\Lambda$ ,  $\Sigma^+$  and  $\Xi^-$ -hypernuclei leads to a generalized mass formula that is valid for normal nuclei ( $n_Y=0$ ) and Hypernuclei ( $n_Y \neq 0$ ) of all kind, having different mass and Strangeness.

$$B(A, Z) = 15.777A - 18.34A^{2/3} - 0.71 \frac{Z(Z-1)}{A^{1/3}} - \frac{23.21(N-Z_c)^2}{[(1+e^{-A/17})A]} + (1 - e^{-A/30})\delta + n_Y [0.0335(m_Y) - 27.8 - 48.7|S|A^{-2/3}],$$

$A = N + Z_c + n_Y =$  total no. of baryons,  
 $Z_c =$  number of protons,  
 $Z = Z_c + n_Y q_Y =$  net charge number,  
 $q_Y =$  charge no. of Hyperon with proper sign.  
 (viz.,  $q = -1, 0, 1$ ).

Hyperon Separation Energy:

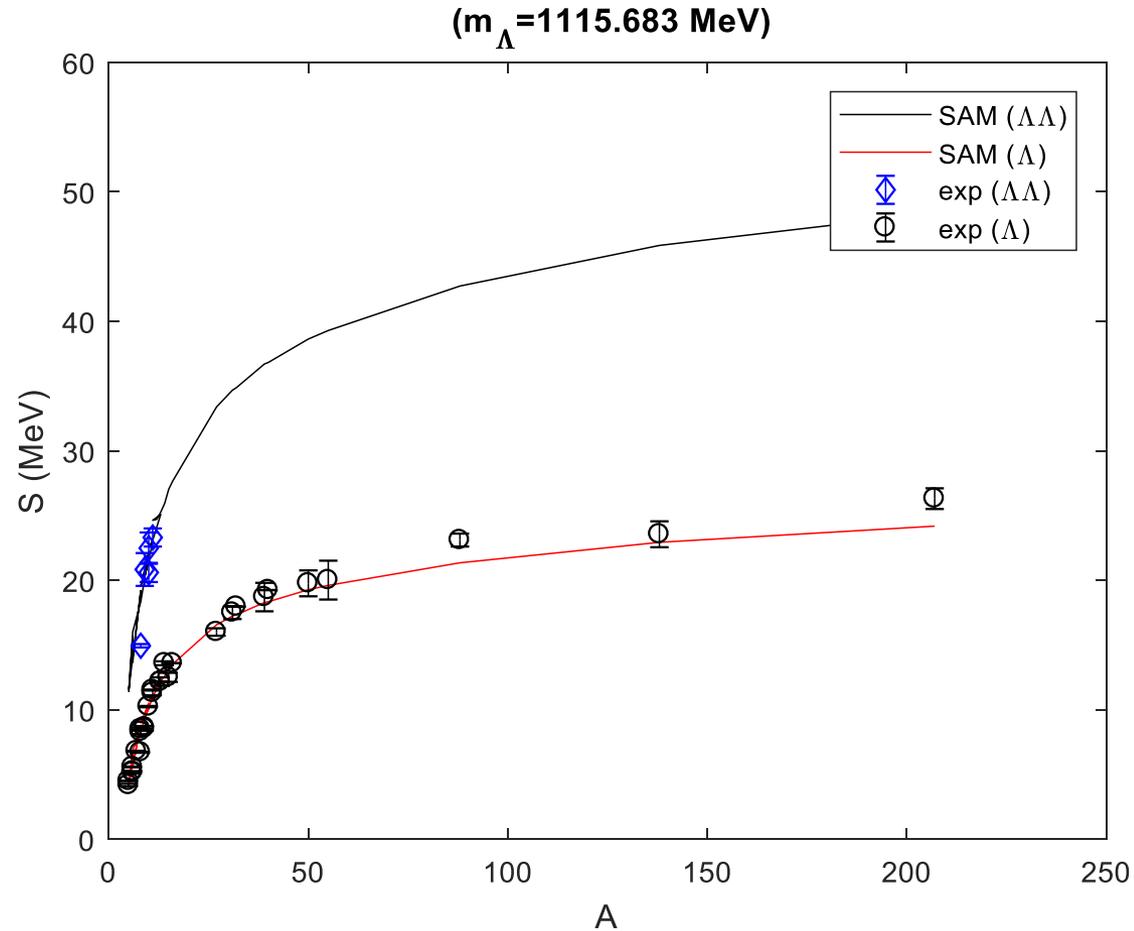
$$S_Y = B(A, Z)_{\text{Hyper}} - B(A - n_Y, Z_c)_{\text{Core}}$$

Explicit dependence on:

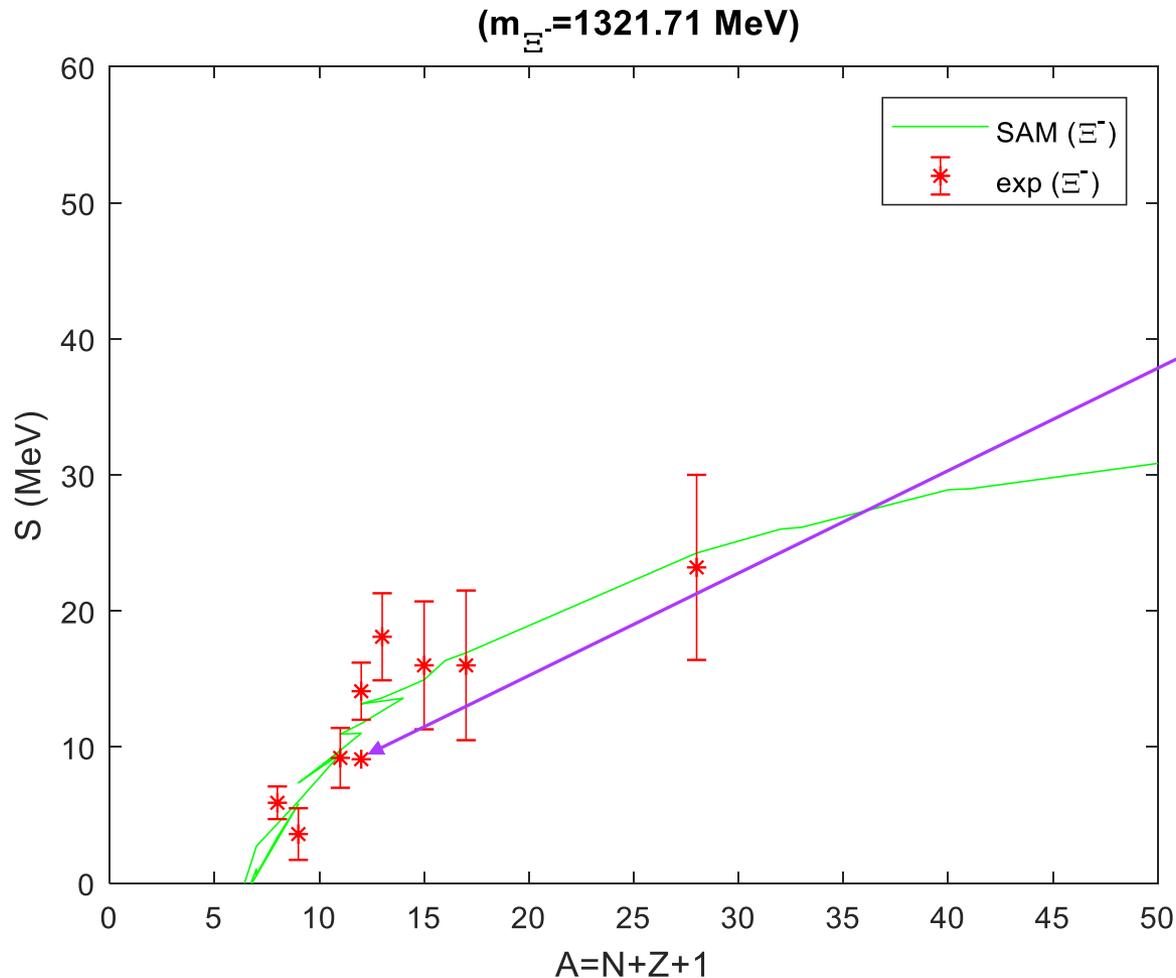
- $n_Y =$  no. of hyperons in a nucleus,
- $m_Y =$  mass of hyperon in MeV,
- $S =$  strangeness of the hyperon.

Hyperon	S	$n_Y$	$q_Y$
$\Lambda^0$	-1	1	0
$\Lambda\Lambda$	-2	2	0
$\Xi^-$	-2	1	-1
Normal	0	0	0

# Variation of Single- $\Lambda$ and Double- $\Lambda$ Separation Energies with $A=N+Z$



# Variation of Single $\Xi^-$ Separation Energies with $A=N+Z+1$



Old Emulsion data

$^{12}_{\Xi^-}\text{Be}$ :

BE(Exp)  $\geq$  9.1 MeV

From the talk of

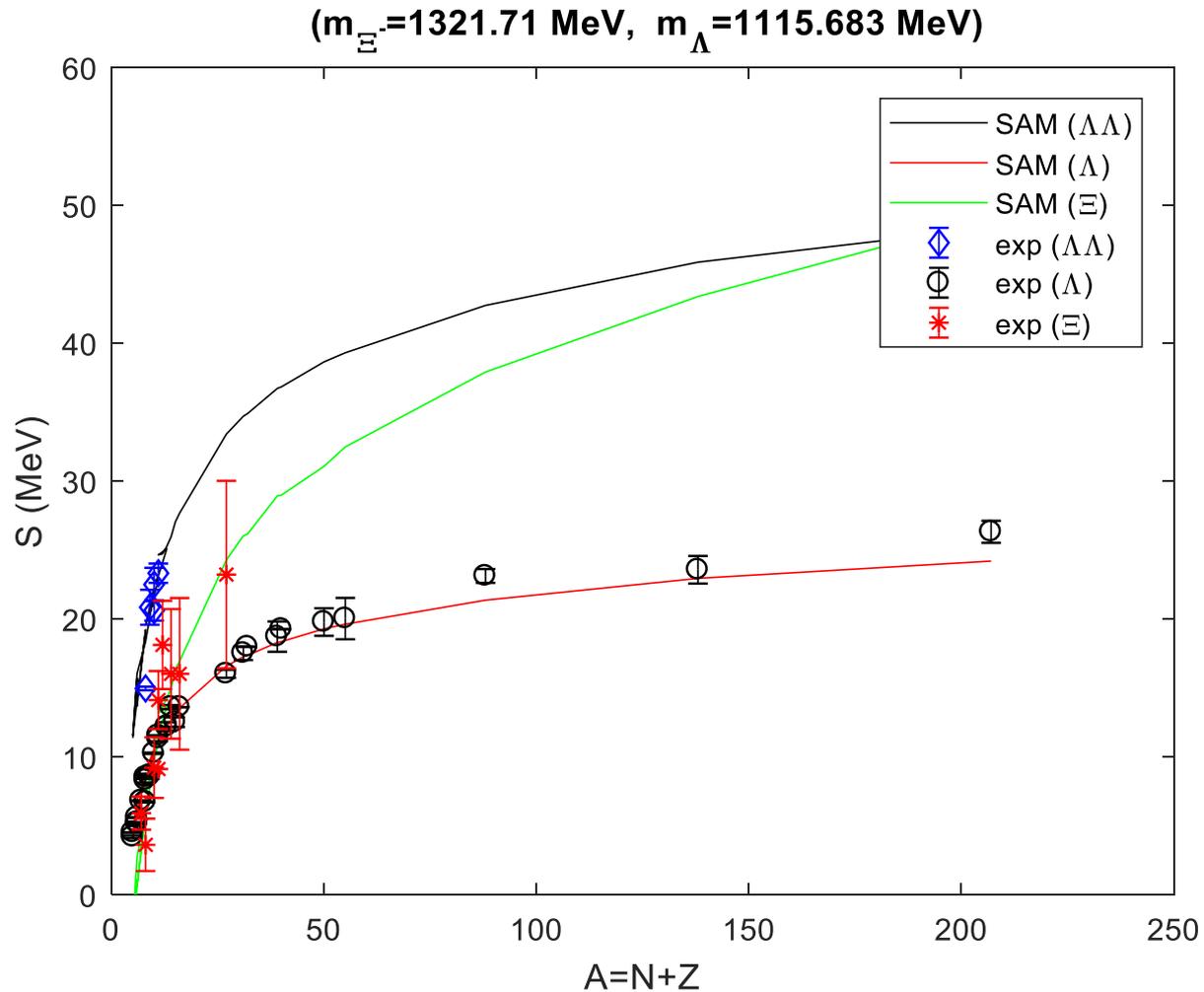
T. Nagae

BE(This work):

$^{12}_{\Xi^-}\text{Be} = ^{11}\text{B} + 1\Xi^-$

=11.01 MeV

# Variation of $\Lambda$ , $\Xi^-$ & $\Lambda\Lambda$ Separation Energies with $A=N+Z$



# Effect of Addition of Lambda

A Lambda-particle makes a nucleus more bound, and can change the neutron-, proton-drip points by creating bound nucleus beyond the normal drip lines.

**Table 1.** One-nucleon separation energies (in MeV) on drip lines for each element with the lowest and highest numbers of bound neutrons in normal and  $\Lambda$ -hypernuclei.

Symbol	Normal p-drip ${}^AZ, S_p$	Normal n-drip ${}^AZ, S_n$	Hyper p-drip ${}^AZ, S_p$	Hyper n-drip ${}^AZ, S_n$
Li	${}^5\text{Li}$ , 3.36	${}^{11}\text{Li}$ , 0.58	${}^5_{\Lambda}\text{Li}$ , 1.22	${}^{12}_{\Lambda}\text{Li}$ , 1.87
Be	${}^6\text{Be}$ , 1.11	${}^{14}\text{Be}$ , 0.90	${}^7_{\Lambda}\text{Be}$ , 3.79	${}^{15}_{\Lambda}\text{Be}$ , 1.84
B	${}^8\text{B}$ , 0.74	${}^{17}\text{B}$ , 0.99	${}^9_{\Lambda}\text{B}$ , 2.39	${}^{18}_{\Lambda}\text{B}$ , 1.73
C	${}^9\text{C}$ , 0.17	${}^{20}\text{C}$ , 1.01	${}^{10}_{\Lambda}\text{C}$ , 1.80	${}^{21}_{\Lambda}\text{C}$ , 1.63
N	${}^{12}\text{N}$ , 1.98	${}^{23}\text{N}$ , 0.97	${}^{12}_{\Lambda}\text{N}$ , 0.24	${}^{24}_{\Lambda}\text{N}$ , 1.50
O	${}^{13}\text{O}$ , 1.98	${}^{26}\text{O}$ , 0.94	${}^{13}_{\Lambda}\text{O}$ , 0.44	${}^{27}_{\Lambda}\text{O}$ , 1.40
F	${}^{15}\text{F}$ , 0.09	${}^{29}\text{F}$ , 0.89	${}^{16}_{\Lambda}\text{F}$ , 0.81	${}^{32}_{\Lambda}\text{F}$ , 0.01
Ne	${}^{16}\text{Ne}$ , 0.64	${}^{32}\text{Ne}$ , 0.87	${}^{17}_{\Lambda}\text{Ne}$ , 1.39	${}^{35}_{\Lambda}\text{Ne}$ , 0.04
Na	${}^{19}\text{Na}$ , 0.54	${}^{35}\text{Na}$ , 0.84	${}^{20}_{\Lambda}\text{Na}$ , 1.05	${}^{38}_{\Lambda}\text{Na}$ , 0.08
Mg	${}^{20}\text{Mg}$ , 1.37	${}^{38}\text{Mg}$ , 0.84	${}^{20}_{\Lambda}\text{Mg}$ , 0.05	${}^{41}_{\Lambda}\text{Mg}$ , 0.13

Normal drip nucleus

is  ${}^{35}\text{Na}$ .

Neutron-drip

nucleus is  ${}^{38}_{\Lambda}\text{Na}$

instead of  ${}^{36}_{\Lambda}\text{Na}$  i.e.,

2 extra neutrons can

be accommodated

due to the addition

of a  $\Lambda$ .

The exact drip-point can change due to the deformation and other parameters near the drip line. Need microscopic calculations.

# Generalized Mass Formula for Non-strange-, Strange- and Multiply-strange Nuclear Systems

C. Samanta, JPG 37 (2010) 075104

$\mathbf{A} = \mathbf{n} + \mathbf{z}_c + \mathbf{n}_Y$  :  $\mathbf{n}$  = no. of neutrons,  $\mathbf{z}_c$  = no. of protons,  $\mathbf{n}_Y$  = no. of hyperons

Binding Energy:  $B(A, Z) = m_A - z_c \cdot m_p - n \cdot m_n - n_Y \cdot m_Y = \text{Negative}$

$$\begin{aligned}
 -B(A, Z) = & 15.777A - 18.34A^{2/3} - 0.71Z(Z-1)/A^{1/3} - 23.21(n - z_c)^2 / [(1 + e^{-A/17})A] + (1 - e^{-A/30})\delta \\
 & + \sum_Y n_Y [0.0335(m_Y) - 26.7 - 48.7 |S| / A^{2/3} \\
 & - a_Y \{ (n_\Lambda + n_{\Xi^0} + n_{\Xi^-} - z_c)^2 + (n_\Lambda + n_{\Xi^0} + n_{\Xi^-} - n)^2 \} / \{ (1 + e^{-A/17})A \}].
 \end{aligned}$$

$m_Y$  = mass of hyperon in MeV

$S$  = strangeness no. of the hyperon

$Z$  = Net Charge

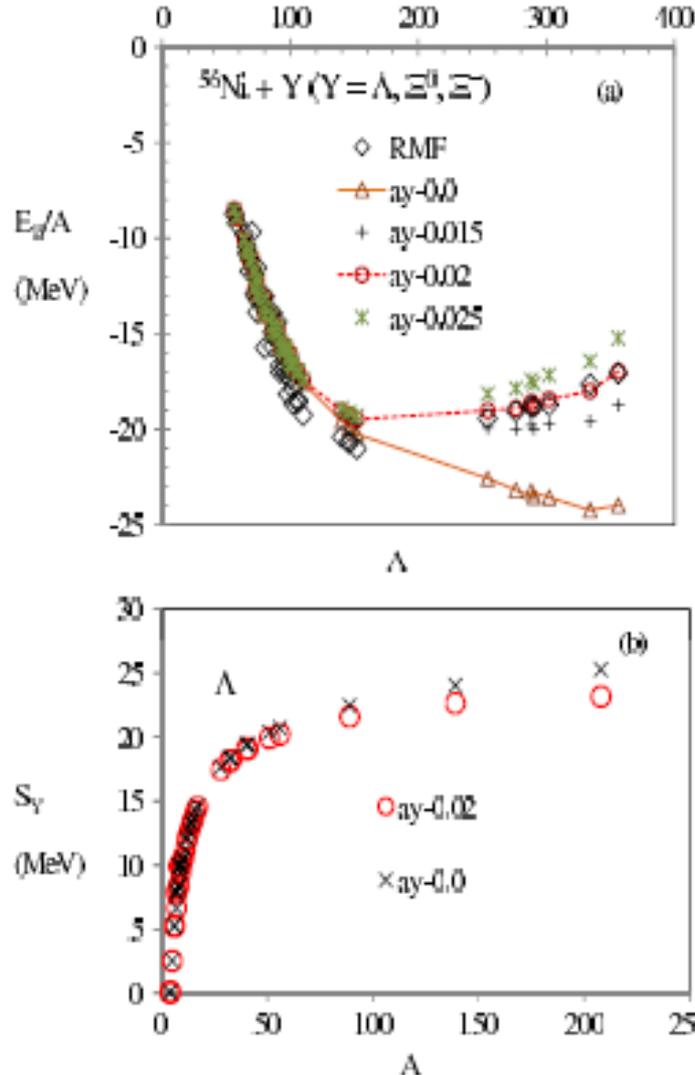
= Proton charges + Total hyperon charges

$q = -1, 0, 1$  depending on the hyperon type

Note: the net charge of a nucleus can be **negative** if the hyperon number is larger than the proton number and the hyperon has a negative charge.

# Multi-Strange Nuclei: Model-2 with Y-Y interaction

C. Samanta, JPG 37 (2010) 075104



**Effect of the new hyperon-asymmetry term: choice of different  $a_Y$  value**

Figures:  $E_B/A$  vs.  $A$  plots for stable multi-strange systems.

(a) The  $a_Y = 0.02$  gives the best fit to the relativistic mean field (RMF) calculations of Schaffner et al. (Ref.1) that is based on  $^{56}\text{Ni}$  nuclear cores with YY interaction (model 2).

(b) Single lambda-hyperon separation energy  $S_Y$  vs.  $A$  for different elements. With  $a_Y = 0.0$  and  $0.2$ .

Ref.1: J. Schaffner J, C.B. Dover, A. Gal, C. Greiner, D.J. Millener and H. Stöcker, Ann. Phys. NY 235 (1994) 35

# Multi-Strange Nuclei: Model-1 without Y-Y interaction

C.Samanta, *JPG 37 (2010) 075104*

(Current wisdom: Y-Y interaction is weak.)

**Model-1=No Y-Y interaction**

**Model-2=Strong Y-Y interaction**

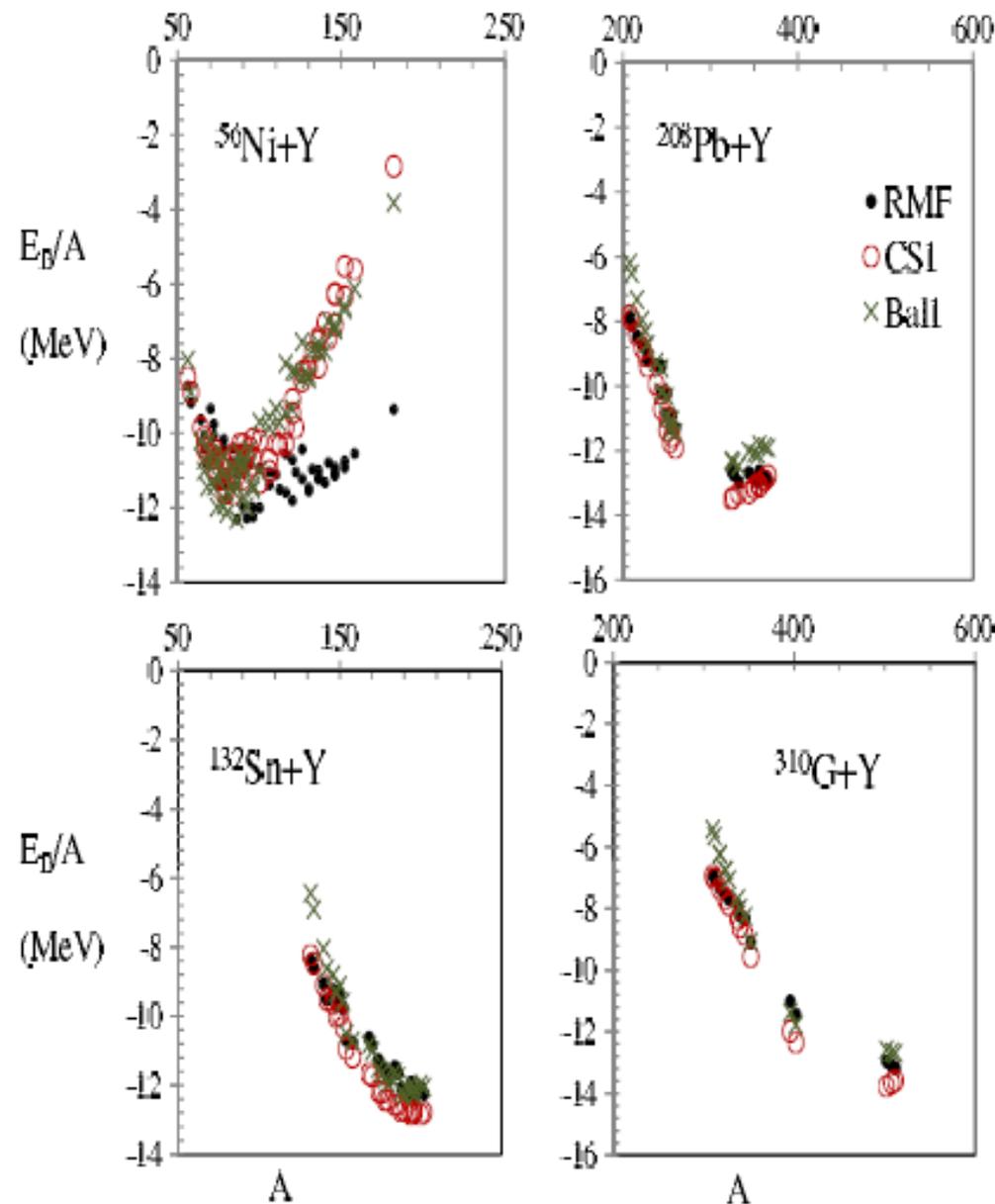
CS1(BE,Model1)=CS2(BE,Model2) -Cr

$$C_r = 12.0 A f_s (f_\Lambda + f_{\Xi^0} + f_{\Xi^-})$$

$$f_s = \sum_Y n_Y |S| / A.$$

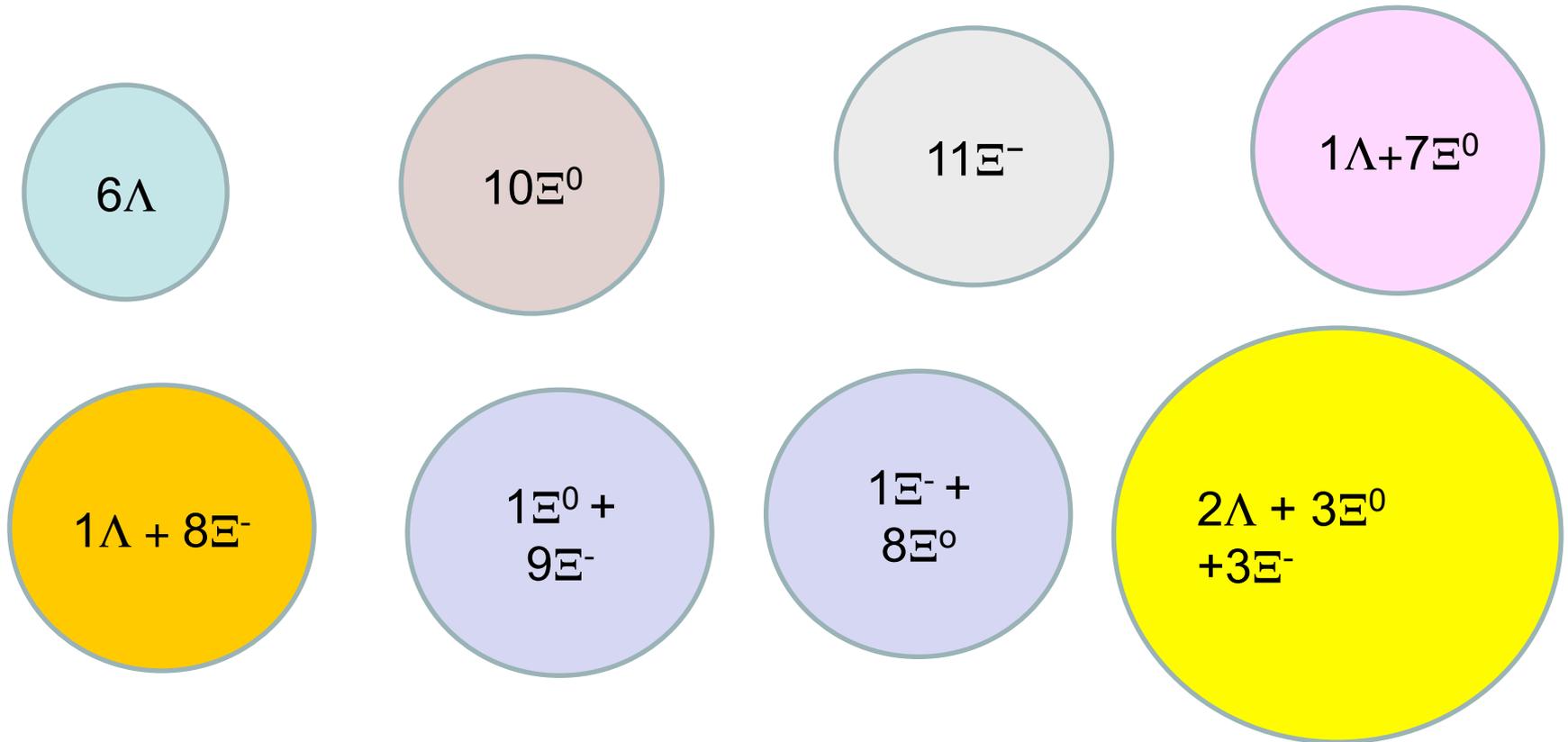
Stable-multiply-strange systems in RMF calculations, based on  $^{56}\text{Ni}$ ,  $^{132}\text{Sn}$ ,  $^{208}\text{Pb}$  and  $^{310}\text{G}$  ( $Z_c=126$ ,  $n=184$ ) core, by:

1. Model-1 of Schaffner et al.,
2. SET-I of Balberg et al.,
3. this work CS1.



# Lightest Bound Nuclei without any Neutrons and Protons!

C. Samanta, Jour. Phys. G: Nucl. Part. Phys. 37 (2010)075104



Pure Hyperonic Systems can be bound (predicted by CS2, with Y-Y interaction)

No bound pure-hyperonic matter is possible by Model-1 (no Y-Y interaction) 12

# Production of hypernuclei in multifragmentation in H.I. collision

In Statistical Multifragmentation model (SMM), mass formula has been used as an input to calculate the relative yield of hypernuclei.

Two different mass formulae have been used:

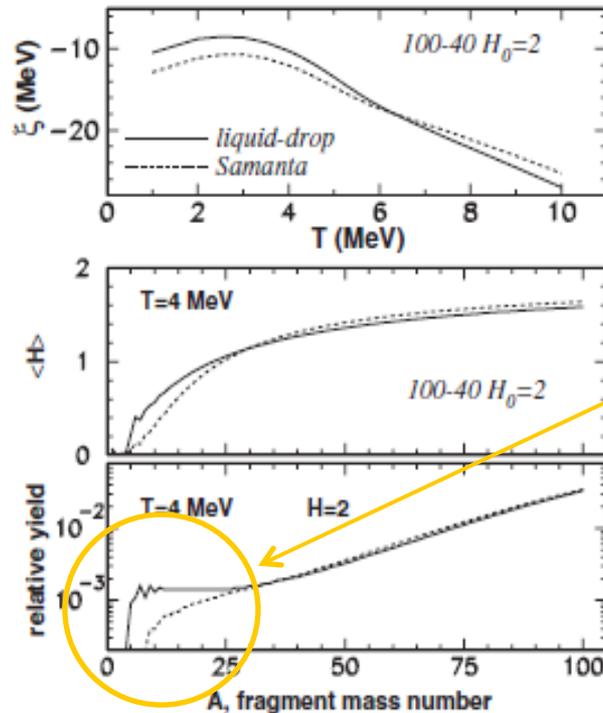
1. Liquid drop model (applicable for  $\Lambda$  hypernuclei only).
2. Generalized mass formula (applicable for all hypernuclei)

A. Botvina and J. Pochodzalla PRC 76 (2007) 024909

# Production of hypernuclei in multifragmentation of nuclear spectator matter

A.Botvina and J. Pochodzalla PRC 76 (2007) 024909

PHYSICAL REVIEW C 76, 024909 (2007)



The SMM calculations show a discrepancy between the generalized mass formula and the Liquid drop mass formula at the low fragment mass number ( $A$ ) region.

FIG. 3. Comparison of SMM calculations with the liquid-drop and Samanta descriptions of hyper terms in the mass formula, for the same sources as in Fig. 2. Top panel – the strangeness chemical potential  $\xi$  versus temperature  $T$ . Middle panel – average number of  $\Lambda$  hyperons in fragments, and bottom panel – yields of fragments with two  $\Lambda$ , at  $T = 4$  MeV.

# Liquid-drop Mass formula (valid for $\Lambda$ -hypernuclei only)

Binding Energy:  $B(A,Z) = m_A - z_c \cdot m_p - n \cdot m_n - n_Y \cdot m_Y = \text{Negative}$

$n$  = no. of neutrons,  $z_c$  = no. of protons,

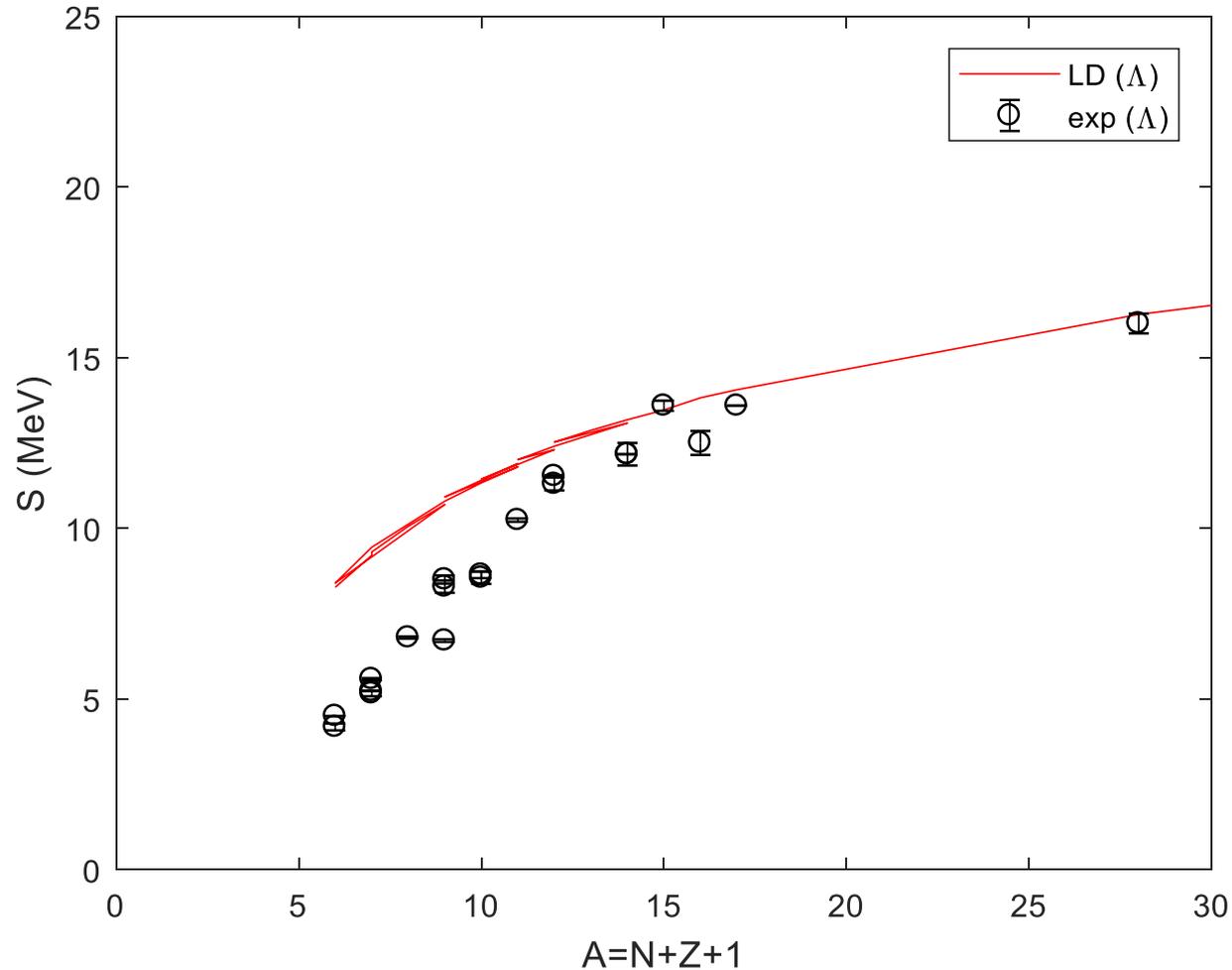
$$A = n + z_c + n_Y$$

$n_Y$  = no. of  $\Lambda$ -hyperons

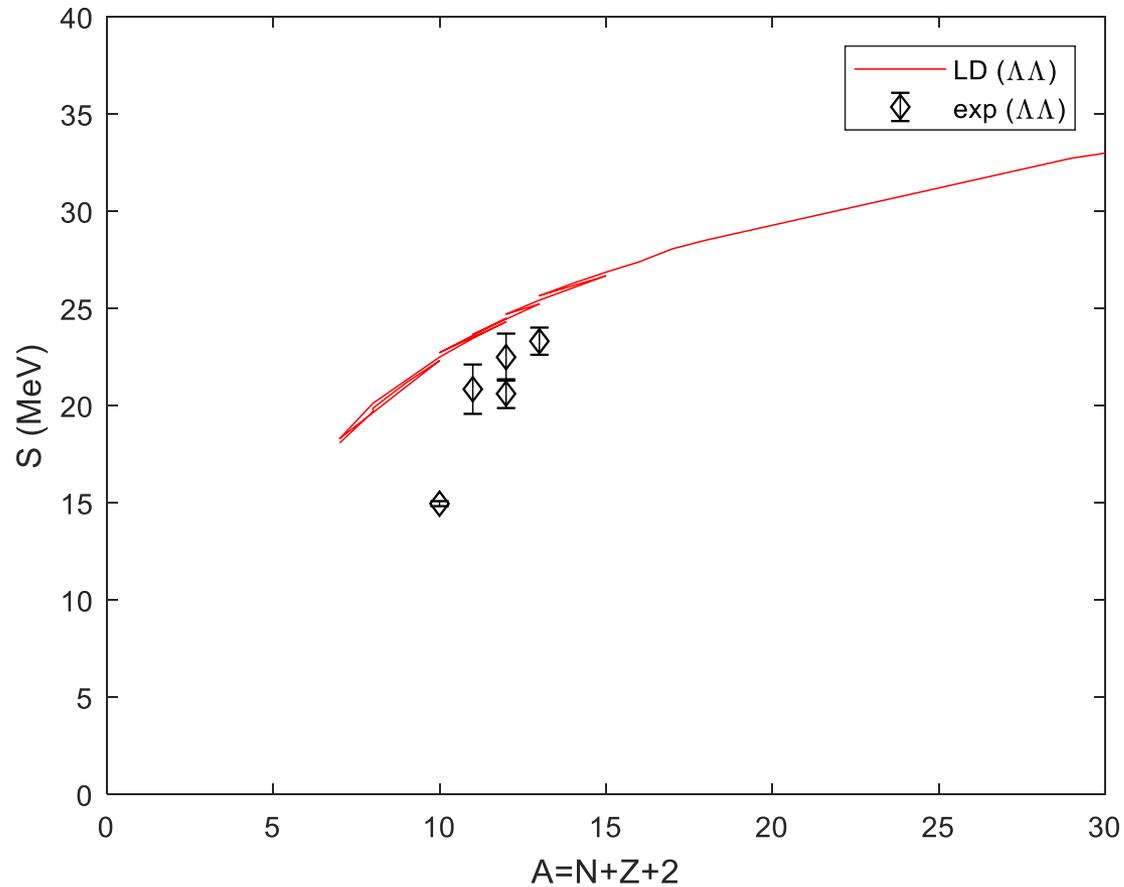
$$B(A,Z) = -16A + 18A^{2/3} + 0.72z_c^2/A^{1/3} + 25(N - z_c)^2/(A - n_Y) - (n_Y/A) [10.68A + 21.27 A^{2/3}]$$

A.S. Botvina, J. Pochodzalla, PRC 76 (2007) 024909

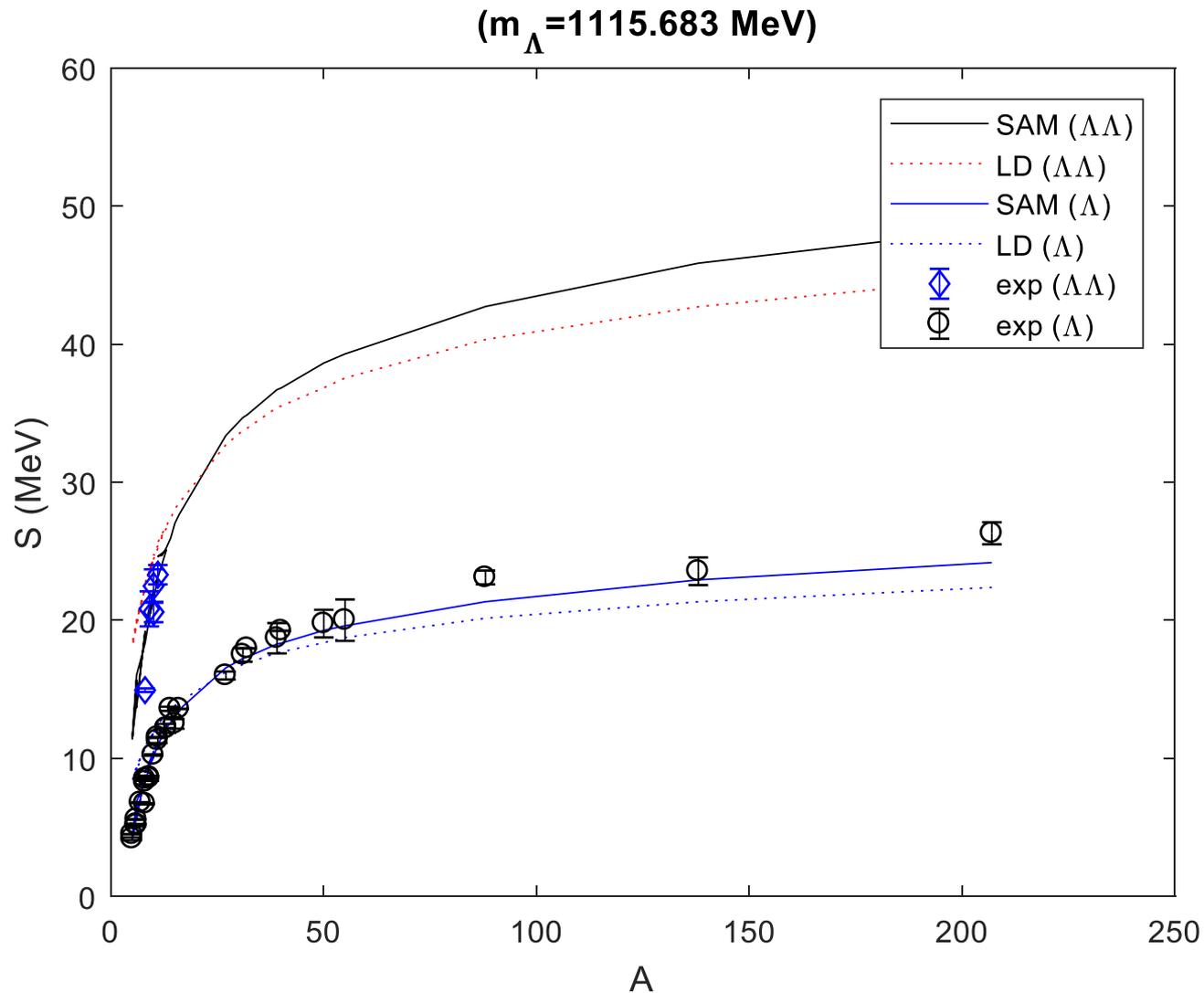
# Variation of Single- $\Lambda$ Separation Energies with Mass Number (A)

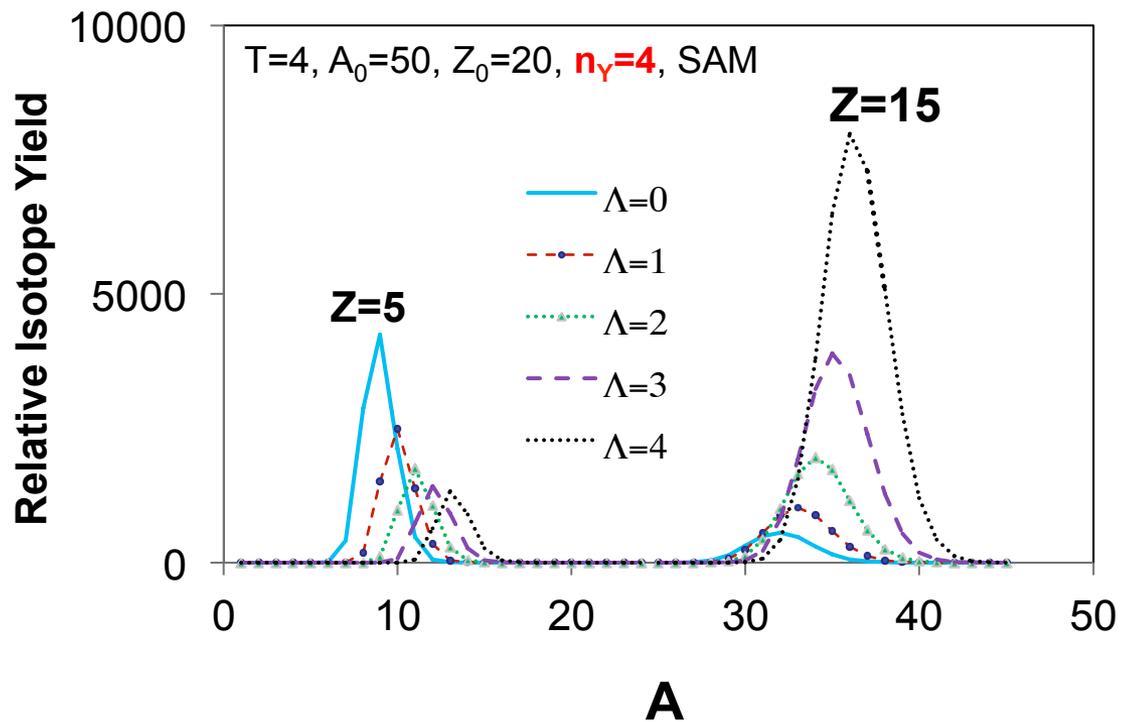
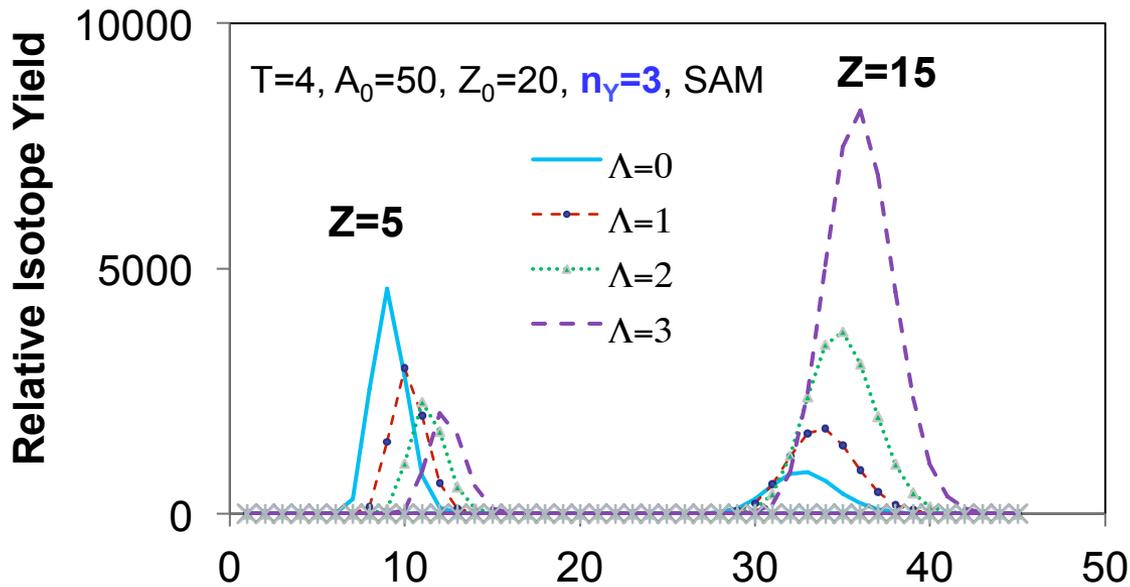


# Variation of Double- $\Lambda$ Separation Energies with Mass Number (A)



# Variation of Single- $\Lambda$ and Double- $\Lambda$ Separation Energies with $A = N + Z$





C. Samanta,  
 A. Botvina,  
 I. Mishustin,  
 & W. Greiner

Relative yields of  $1\Lambda, 2\Lambda, 3\Lambda$  and  $4\Lambda$  was calculated using Statistical Multifragmentation Model (SMM) and the generalized mass formula.

It indicates production yield of  $2\Lambda, 3\Lambda$  and  $4\Lambda$ -hypernuclei to be higher than normal nuclei ( $\Lambda=0$ ) at high temperature for fragments with larger mass number.

# $\Lambda\Lambda$ -Bond Energy

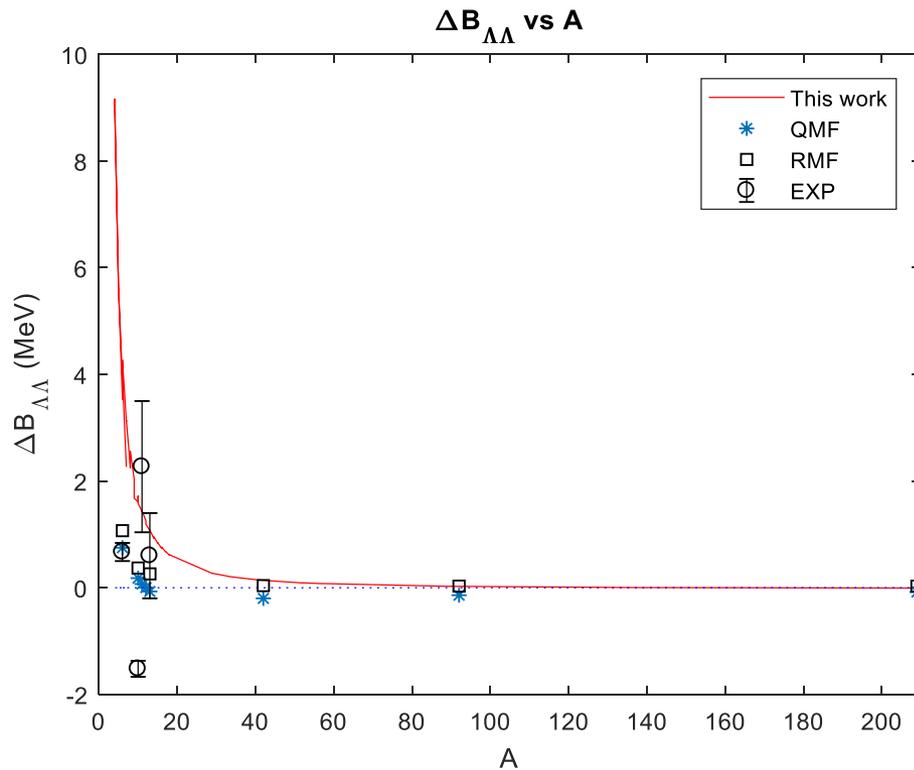
$$B_{\Lambda}({}^A_{\Lambda}Z) = M({}^{A-1}Z) + M(\Lambda) - M({}^A_{\Lambda}Z) \quad (1)$$

$$B_{\Lambda\Lambda}({}^A_{\Lambda\Lambda}Z) = M({}^{A-2}Z) + 2M(\Lambda) - M({}^A_{\Lambda\Lambda}Z) \quad (2)$$

$$\Delta B_{\Lambda\Lambda}({}^A_{\Lambda\Lambda}Z) = B_{\Lambda\Lambda}({}^A_{\Lambda\Lambda}Z) - 2B_{\Lambda}({}^{A-1}Z) \quad (3)$$

where  $M(\Lambda)$  is the mass of the  $\Lambda$  hyperon,  $M({}^{A-1}Z)$  is the mass of the core nucleus of the single- $\Lambda$  nucleus of mass  $M({}^A_{\Lambda}Z)$ , and  $M({}^{A-2}Z)$  is the mass of the core nucleus of the double- $\Lambda$  nucleus of mass  $M({}^A_{\Lambda\Lambda}Z)$ . Here the mass number ( $A$ ) of the hypernuclei is the total number of baryons (neutrons, protons and hyperons) in the nuclei. The bond energies ( $\Delta B_{\Lambda\Lambda}$ ) are the measures of the energy released when the  $\Lambda\Lambda$  bond is broken. These energies assist in understanding the nature of the in-medium strength of the  $\Lambda$ - $\Lambda$  interaction.

# $\Lambda\Lambda$ -Bond Energy



**Table 2.**  $\Lambda\Lambda$  binding and bond energies,  $B_{\Lambda\Lambda}$  and  $\Delta B_{\Lambda\Lambda}$  (in MeV) respectively for double- $\Lambda$  hypernuclei for the QMF model, the RMF model, and the available experimental data.

${}^A_{\Lambda\Lambda}Z$	$B_{\Lambda\Lambda}$			$\Delta B_{\Lambda\Lambda}$		
	Exp.	QMF	RMF	Exp.	QMF	RMF
${}^6_{\Lambda\Lambda}\text{He}$	$6.91 \pm 0.16$	7.94	5.52	$0.67 \pm 0.17$	0.75	1.07
${}^{10}_{\Lambda\Lambda}\text{Be}$	$11.90 \pm 0.13$	17.61	16.34	$-1.52 \pm 0.15$	0.18	0.37
${}^{11}_{\Lambda\Lambda}\text{Be}$	$20.49 \pm 1.15$	19.46		$2.27 \pm 1.23$	0.07	
${}^{12}_{\Lambda\Lambda}\text{Be}$	$22.23 \pm 1.15$	21.00			0.002	
${}^{13}_{\Lambda\Lambda}\text{B}$	$23.30 \pm 0.70$	22.41	22.14	$0.60 \pm 0.80$	-0.07	0.26
${}^{42}_{\Lambda\Lambda}\text{Ca}$		37.25	38.15		-0.20	0.04
${}^{92}_{\Lambda\Lambda}\text{Zr}$		45.24	47.11		-0.14	0.03
${}^{208}_{\Lambda\Lambda}\text{Pb}$		50.73	52.19		-0.08	0.03

## Quark mean-field model for single and double $\Lambda$ and $\Xi$ hypernuclei

J. N. Hu<sup>1</sup>, A. Li<sup>2,3,4</sup>, H. Shen<sup>4</sup>, and H. Toki<sup>5</sup>

# Charge Symmetry Breaking (CSB)

(Binding energy difference in mirror nuclei)

Mirror pair	$\Delta B_\Lambda$ Exp	$\Delta B_\Lambda$ T.W.	CSB Ref.
${}^{40}_{\Lambda}\text{Ca} - {}^{40}_{\Lambda}\text{K}$	-	0.07	
${}^{32}_{\Lambda}\text{S} - {}^{32}_{\Lambda}\text{P}$	-	0.08	
${}^{28}_{\Lambda}\text{Si} - {}^{28}_{\Lambda}\text{Al}$	-	0.09	
${}^{16}_{\Lambda}\text{O} - {}^{16}_{\Lambda}\text{N}$	$-0.74 \pm 0.43$	0.09	$-0.36 \pm 0.43$
${}^{14}_{\Lambda}\text{N} - {}^{14}_{\Lambda}\text{C}$	$0 \pm 0.33$	0.09	
${}^{12}_{\Lambda}\text{C} - {}^{12}_{\Lambda}\text{B}$	$-0.764 \pm 0.20$	0.09	$-0.23 \pm 0.19$
${}^{10}_{\Lambda}\text{B} - {}^{10}_{\Lambda}\text{Be}$	$-0.60 \pm 0.36$	0.09	$-0.22 \pm 0.250$
${}^8_{\Lambda}\text{Be} - {}^8_{\Lambda}\text{Li}$	$0.04 \pm 0.08$	0.10	$0.04 \pm 0.060$

Ref: E. Botta et al, NPA960(2017)165

The  $\Delta B_\Lambda$  values obtained in this work (T.W.) is basically the Coulomb energy difference of the mirror nuclei, **not the CSB effect**.

Since the binding energies of this mass formula well reproduces the experimental binding energies, the  $\Delta B_\Lambda$  values of this work can be taken out from the experimental  $\Delta B_\Lambda$  values to estimate the Coulomb deducted CSB effect.

# Summary

- ❖ A single generalized mass formula is prescribed for all hypernuclei and normal nuclei.
- ❖ It reproduces binding energies obtained from the experimental data, and microscopic (RMF, QMF) calculations available on limited nuclei.
- ❖ It can guide future experiments by giving an estimation of the possible binding energy location.  

For example, the  $\Xi^-$  - separation energy of  $^{12}_{\Xi^-}\text{Be}$ , is predicted to be  $\sim 11.01$  MeV.
- ❖ This mass formula indicates shifts in n, p drip lines, and predicts possible existence of bound nuclei without neutrons and protons.
- ❖ Relative yields of  $1\Lambda$ ,  $2\Lambda$ ,  $3\Lambda$  and  $4\Lambda$  calculated using Statistical Multifragmentation Model (SMM) indicate production yield of  $2\Lambda$ ,  $3\Lambda$  and  $4\Lambda$ -hypernuclei to be higher than normal nuclei at high temperature for fragments with larger mass number.
- ❖ The  $\Lambda\Lambda$  bond energy rapidly drops off with the increasing mass number in agreement with RMF & QMF.
- ❖ The Coulomb energy difference of mirror nuclei can be deduced from this mass formula to extract the CSB effect.
- ❖ More Experimental data are needed.

**Thank You**