

(Weak decay of hypernuclei)
focusing on
Pionic decay of hypernuclei

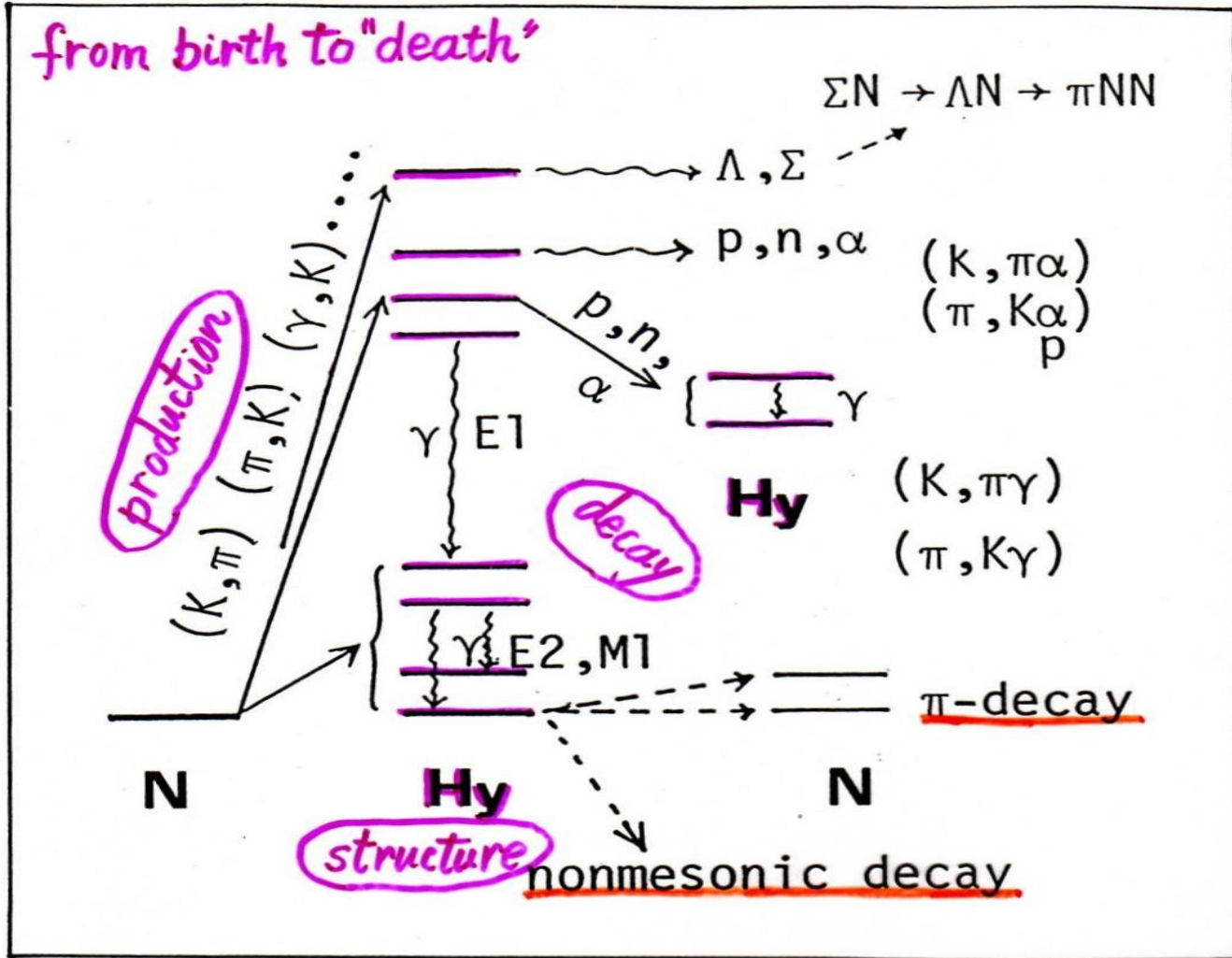
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Hypernuclear Workshop

May 27 - 29, 2014 Jefferson Lab

Life of strangeness many-body systems

Life of the Hypernucleus



Hypernuclear weak decay : 2 modes

Mesonic mode (MWD)

$$q(\pi) = 100 \text{ MeV}/c$$

$$\Lambda \rightarrow p + \pi^- + 37.8 \text{ MeV} \quad (64.2\%),$$

$$\Lambda \rightarrow n + \pi^0 + 41.1 \text{ MeV} \quad (35.8\%).$$

Nonmesonic mode (NMWD)

$$q(\text{NM}) = 400 \text{ MeV}/c$$

$$\Lambda + p \rightarrow n + p + 176 \text{ MeV},$$

$$\Lambda + n \rightarrow n + n + 176 \text{ MeV}.$$

Mesonic Weak Decay

$\Lambda \rightarrow p + \pi^-$	2	$\Delta I = 1/2$ rule	<u>Free Λ decay</u>	
$\Lambda \rightarrow n + \pi^0$	1		$\Lambda \rightarrow p + \pi^-$ ($\sim 2/3$)	$\Lambda \rightarrow n + \pi^0$ ($\sim 1/3$)
$1/2^+$	$1/2^+$	0^-		
\uparrow	\uparrow	L=0	s-wave; PV	
\uparrow	\downarrow	L=1	p-wave; PC	

$$I(\theta) \propto 1 - \alpha P \cos\theta$$

$$\alpha = \frac{2a_s \operatorname{Re} a_p^*}{|a_s|^2 + |a_p|^2}$$

88% s-wave

12% p-wave

Theoretical framework

Pi-mesonic decay interaction is expressed as

$$\mathcal{H}_\pi = ig_w \bar{\psi}_N (1 + \lambda \gamma_5) \tau \begin{pmatrix} 0 \\ \psi_\Lambda \end{pmatrix} \phi_\pi,$$

$$g_w = 0.233 \times 10^{-6} \quad \text{and} \quad \lambda = -6.9$$

Non-relativistic expression is employed:

$$H_\pi = \left[s_\pi X^{(s)}(r) + ip_\pi X^{(p)}(r) \frac{\sigma \cdot \nabla}{q_0} \right] \chi_\pi^{(-)*}(\mathbf{q}; \mathbf{r}),$$

$$s_{\pi^-} = -\frac{\sqrt{2}}{\sqrt{4\pi}} g_w = 0.96 \times 10^{-7}, \quad s_{\pi^0} = -s_{\pi^-} / \sqrt{2},$$

$$p_{\pi^-} = \frac{q_0}{2\sqrt{M_\Lambda M_N}} \lambda s_{\pi^-} = -0.35 \times 10^{-7}, \quad p_{\pi^0} = -p_{\pi^-} / \sqrt{2}.$$

Solve pion distorted waves with the optical potential

$$2\omega U_\pi^{\text{FULL}} = -4\pi[b(r) + B(r)] + 4\pi\nabla \cdot \mathcal{L}(r)[c(r) + C(r)]\nabla \\ - 4\pi \left[\frac{p_1 - 1}{2} \nabla^2 c(r) + \frac{p_2 - 1}{2} \nabla^2 C(r) \right],$$

$$\mathcal{L}(r) = \frac{1}{1 + \frac{4\pi}{3}\lambda[c(r) + C(r)]}$$

$$M_\Lambda + E({}_\Lambda^A Z; J_i T_i) = M_N + E({}^A Z'; J_f T_f) + \frac{q^2}{2M_A} + \omega_q,$$

$$\omega_q = \sqrt{m_\pi^2 + q^2}$$

Pion distorted waves $\chi_{\pi}^{(-)*}(\mathbf{q}; \mathbf{r})$

solve Klein-Gordon Eq.

Optical potential (MSU group)

J.A. Carr et al. P.R. C25(1982)952

effective form:

$$2\omega U_{\pi} = -4\pi [b_{\text{eff}} \rho(r) - c_{\text{eff}} \nabla \rho(r) \nabla + c_{\text{eff}} \frac{\omega}{2M} \nabla^2 \rho(r)]$$

adopted

general form:

⊕ Vertex renormalization

M. Ericson & H. Bandō, P.L. B273(1990)169

$$2\omega U_{\pi} = -4\pi [b(r) + B(r)] + 4\pi \nabla \cdot \{ \underbrace{\mathcal{L}(r)}_{\text{LLEE}} [C(r) + C(r)] \} \nabla - 4\pi \left\{ \frac{P_1 - 1}{2} \nabla^2 c(r) + \frac{P_2 - 1}{2} \nabla^2 C(r) \right\}$$

$$b(r) = p_1 [b_0 \rho(r) - \epsilon_{\pi} b_1 \delta \rho(r)]$$

$$c(r) = \frac{1}{p_1} [c_0 \rho(r) - \epsilon_{\pi} c_1 \delta \rho(r)]$$

$$B(r) = p_2 B_0 \rho(r)^2$$

$$C(r) = \frac{1}{p_2} C_0 \rho(r)^2$$

$$\mathcal{L}(r) = \left\{ 1 + \frac{4\pi}{3} \lambda [C(r) + C(r)] \right\}^{-1}$$

$$\mathcal{L}^{(s)}(r) = 1$$

$$\mathcal{L}^{(p)}(r) = \mathcal{L}(r)$$

Decay rate is expressed with suppression factors $S^{(s)}$ and $S^{(p)}$

$$\Gamma_{\pi}({}^A_{\Lambda}Z; J_i T_i \tau_i) = \sum_f \Gamma_{\pi}(J_i T_i \tau_i \rightarrow J_f T_f \tau_f), \quad \Gamma_{\pi}^{\text{free}} = \frac{2q_0}{1 + (\omega_{q_0}/M_N)} (s_{\pi}^2 + p_{\pi}^2),$$

$$\Gamma_{\pi}(J_i T_i \tau_i \rightarrow J_f T_f \tau_f) = \frac{2q}{1 + (\omega_q/M_A)} \left[s_{\pi}^2 S_{\pi}^{(s)}(i, f; q) + p_{\pi}^2 \left(\frac{q}{q_0}\right)^2 S_{\pi}^{(p)}(i, f; q) \right],$$

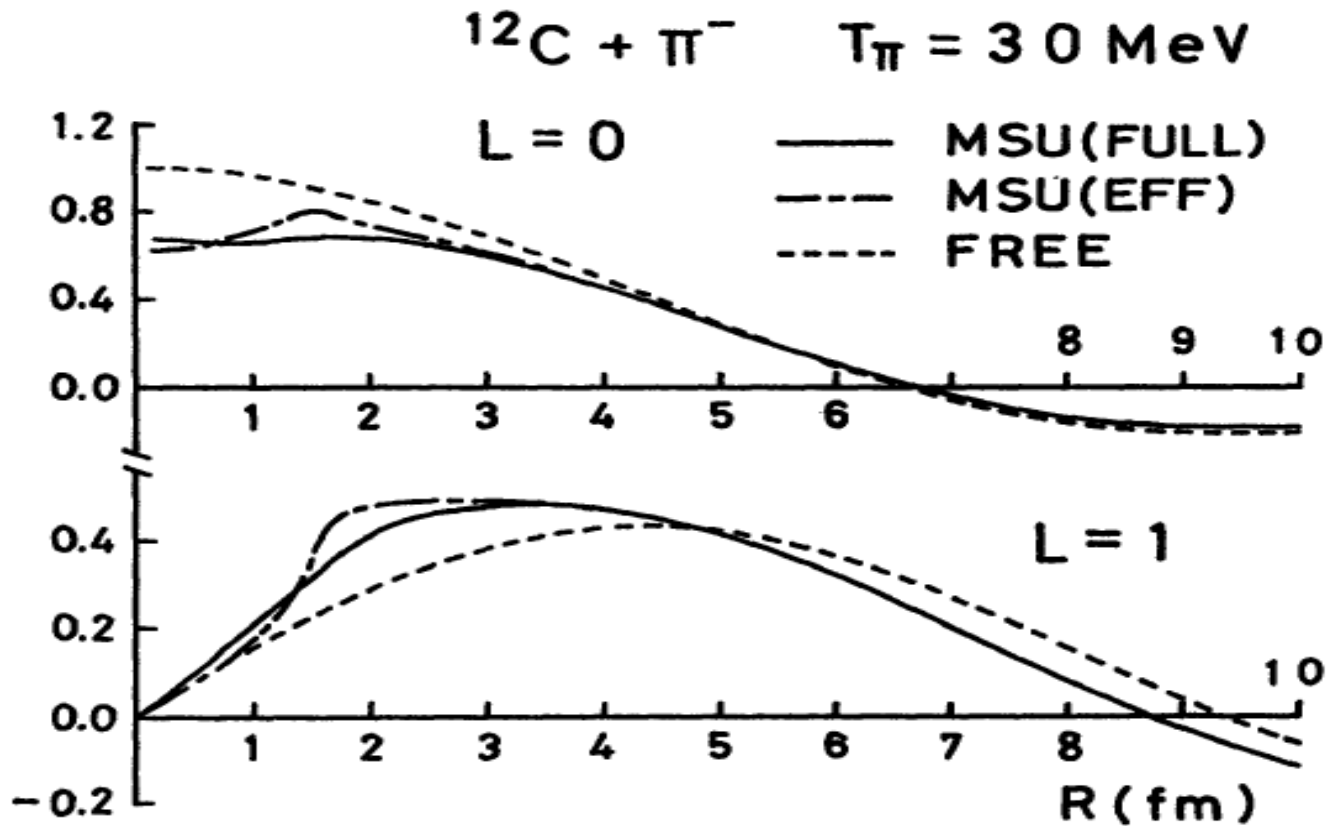
$$S_{\pi}^{(s,p)}(i, f; q) = \frac{4\pi}{2J_i + 1} \sum_K \left| \sqrt{A} \langle \Psi({}^A Z'; J_f T_f \tau_f) \| F_K^{(s,p)} \theta_{\pi} \| \Phi({}^A_{\Lambda} Z; J_i T_i \tau_i) \rangle \right|^2$$

$$F_K^{(s)} = \tilde{j}_K(q; r) X^{(s)}(r) Y_K(\hat{\mathbf{r}}),$$

$$F_K^{(p)} = -i \sum_{k=K \pm 1} (10K0|k0) f_k^K(q; r) X^{(p)}(r) [\sigma \times Y_k(\hat{\mathbf{r}})]_K$$

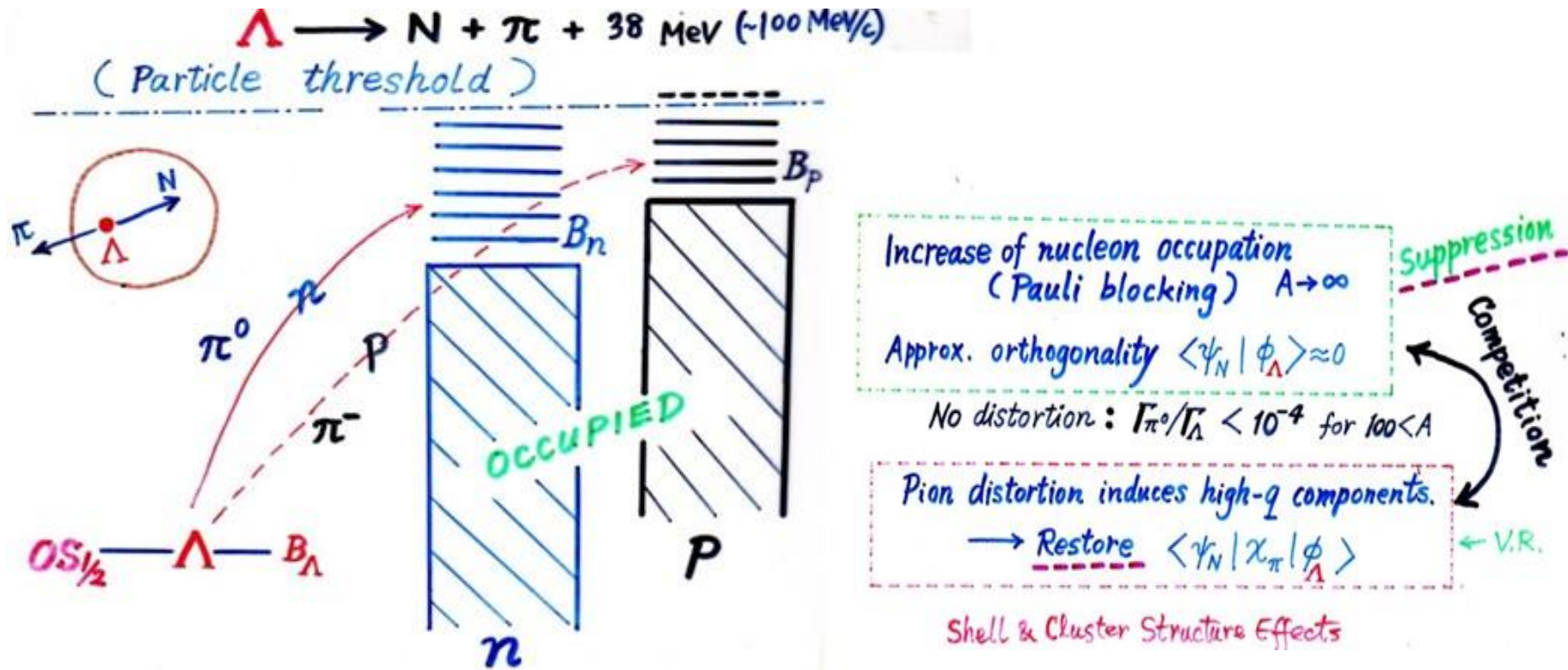
$$f_k^K(q; r) = \frac{\partial \tilde{j}_K(q; r)}{q \partial r} + \left\{ 1 - \frac{k(k+1) - K(K+1)}{2} \right\} \frac{\tilde{j}_K(q; r)}{qr}.$$

Effect of pion wave distortion



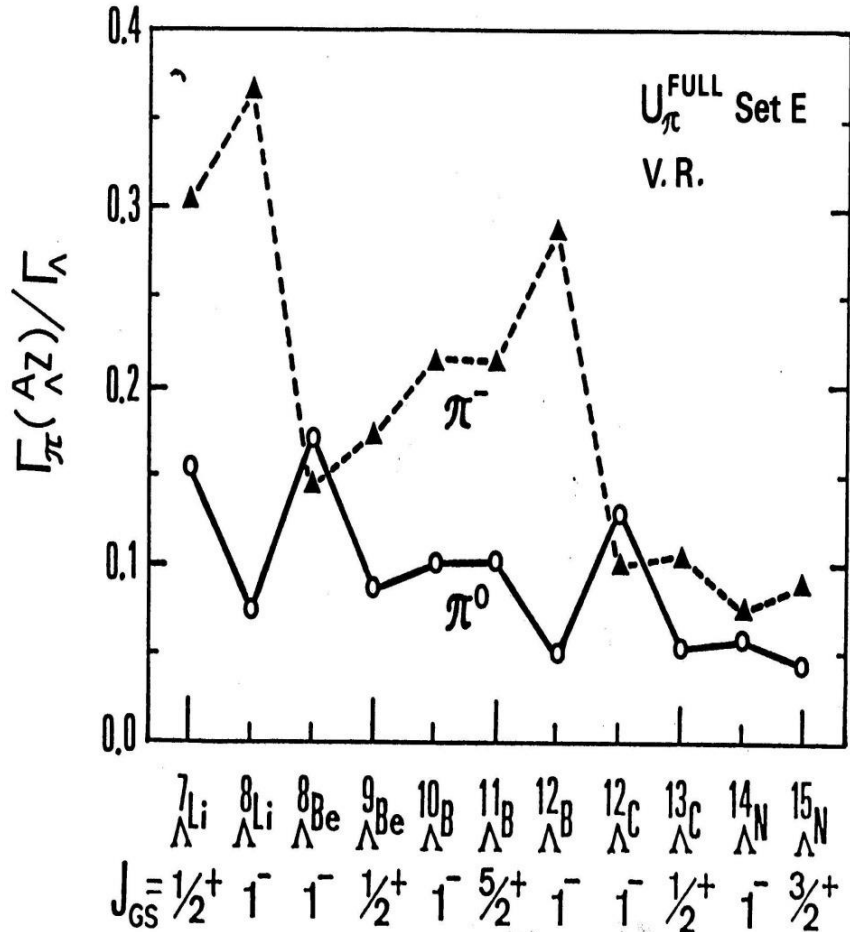
DW vs. PW: $e^{-iq \cdot r}$ if there is no distortion

Actual process (illustrative)

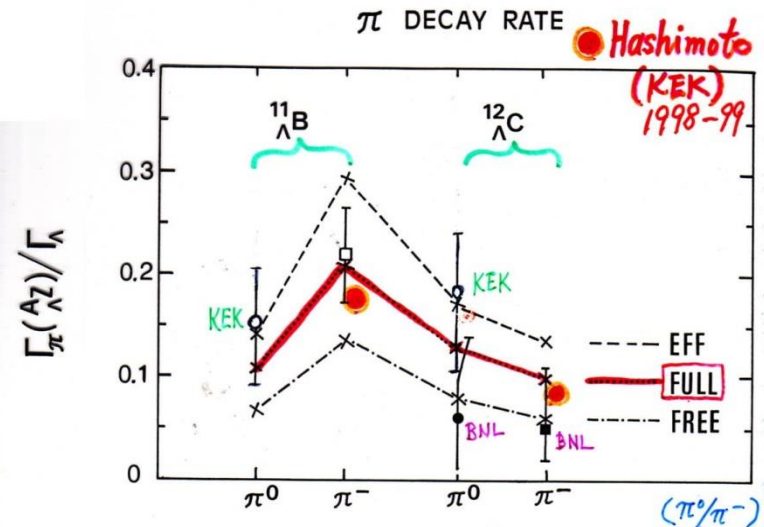


Decay rate and decay pattern are very sensitive to shell structure (spin and binding energy)
 --→ Useful tool to identify the hypernucleus

Very sensitive shell dependence predicted.

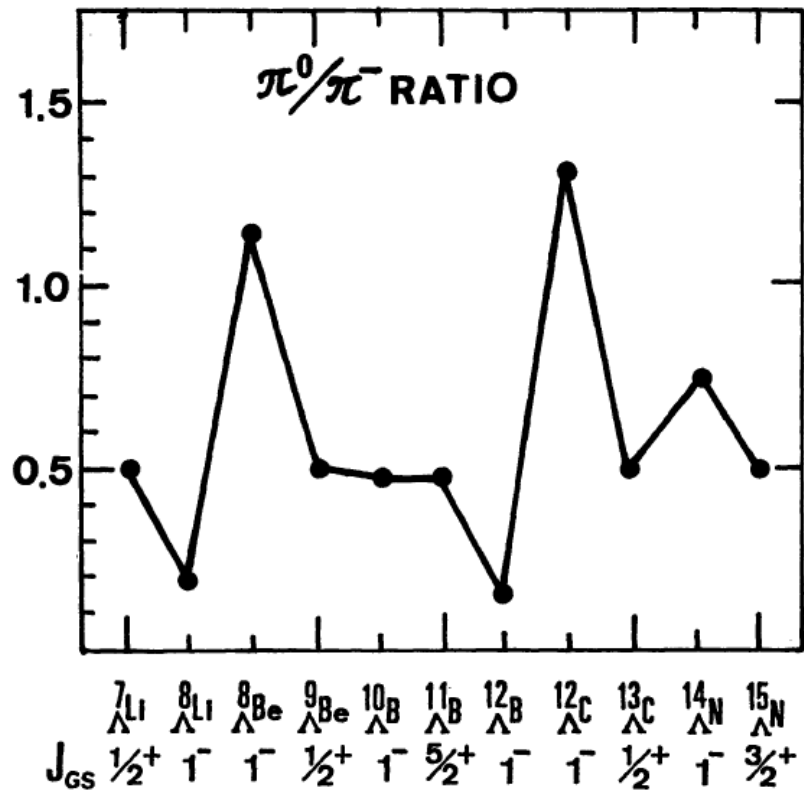
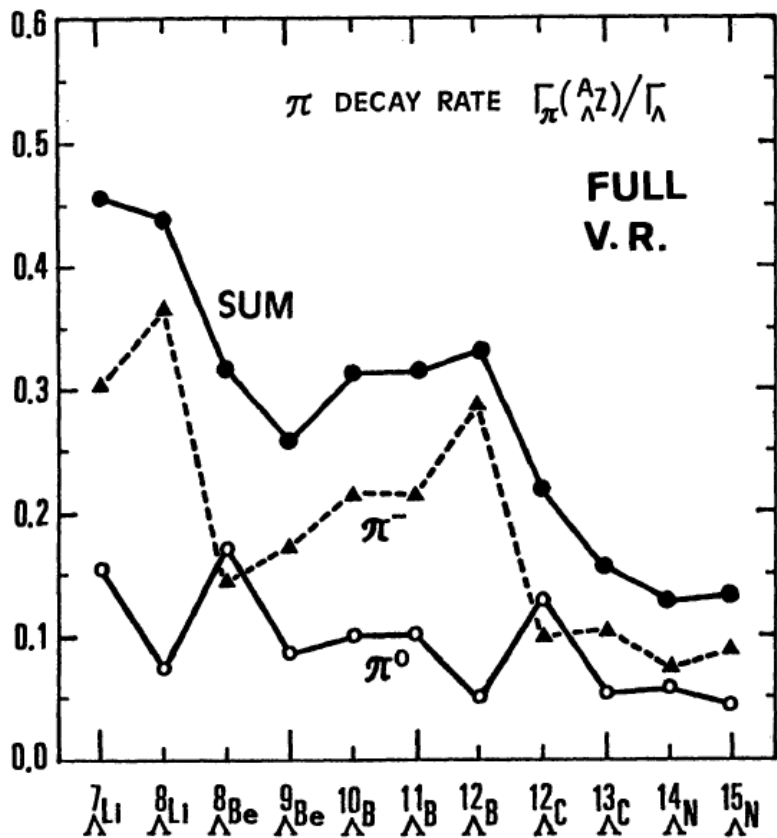


BNL \bullet \blacksquare J.J. Szymanski et al., P.R. C43 (1991)
 KEK \circ A. Sakaguchi et al., P.R. C43 (1991)

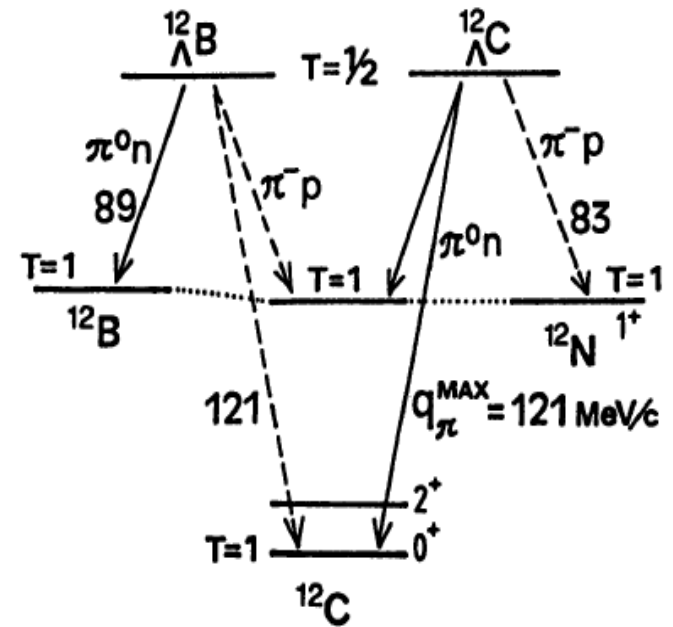
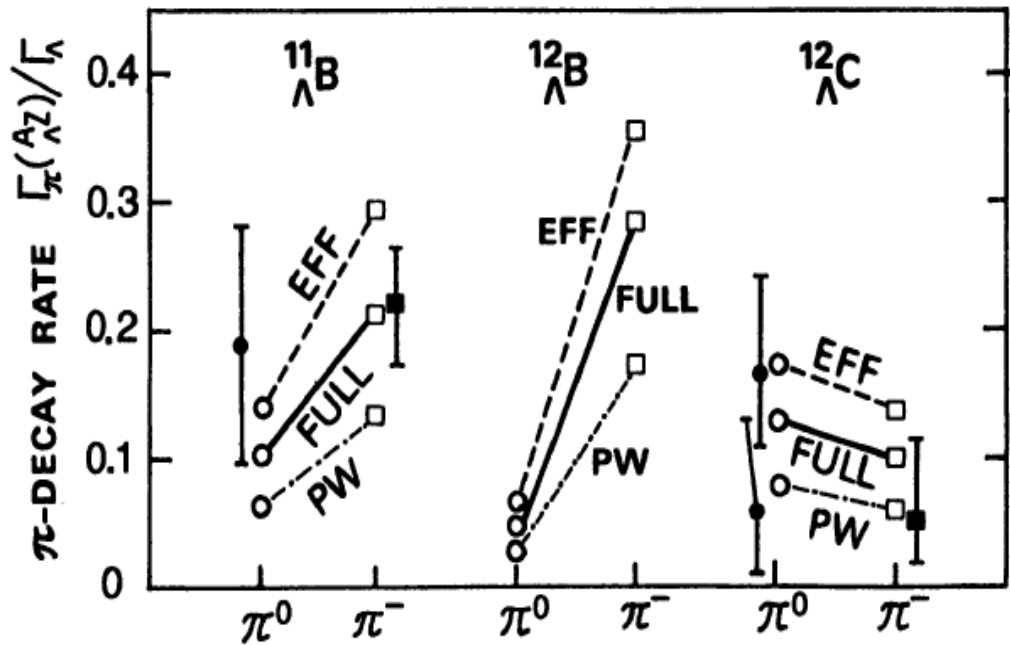


	π^0	π^-	π^0	π^-	(π^0/π^-)
CAL.	0.103	0.213	0.130	0.098	(1.32)
(BNL)			$0.06^{+0.08}_{-0.05}$	$0.05^{+0.06}_{-0.03}$	(1.16)
KEK	0.192 ± 0.090		0.217 ± 0.085		
Hashimoto (1998)		0.163 ± 0.026		0.092 ± 0.020	
Oset et al. (priv. comm.) 1992		0.159		0.086	(1.86)

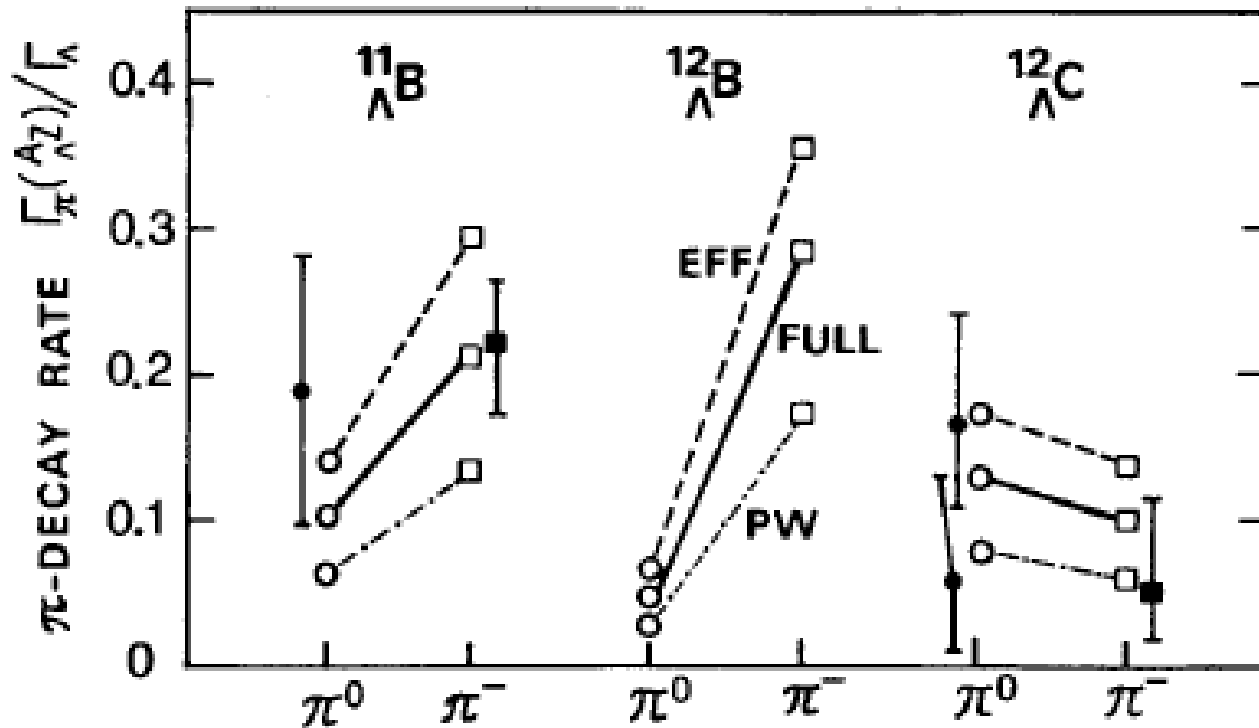
1. Importance of π -DW 30-40% enhancement in p-shell
2. $\Gamma_{\pi^0} > \Gamma_{\pi^-}$ in ${}^{12}_{\Lambda}\text{C}$ has been confirmed by Exp. \rightarrow sensitive to shell structure



Shell dependence comes from energy-momentum available for π

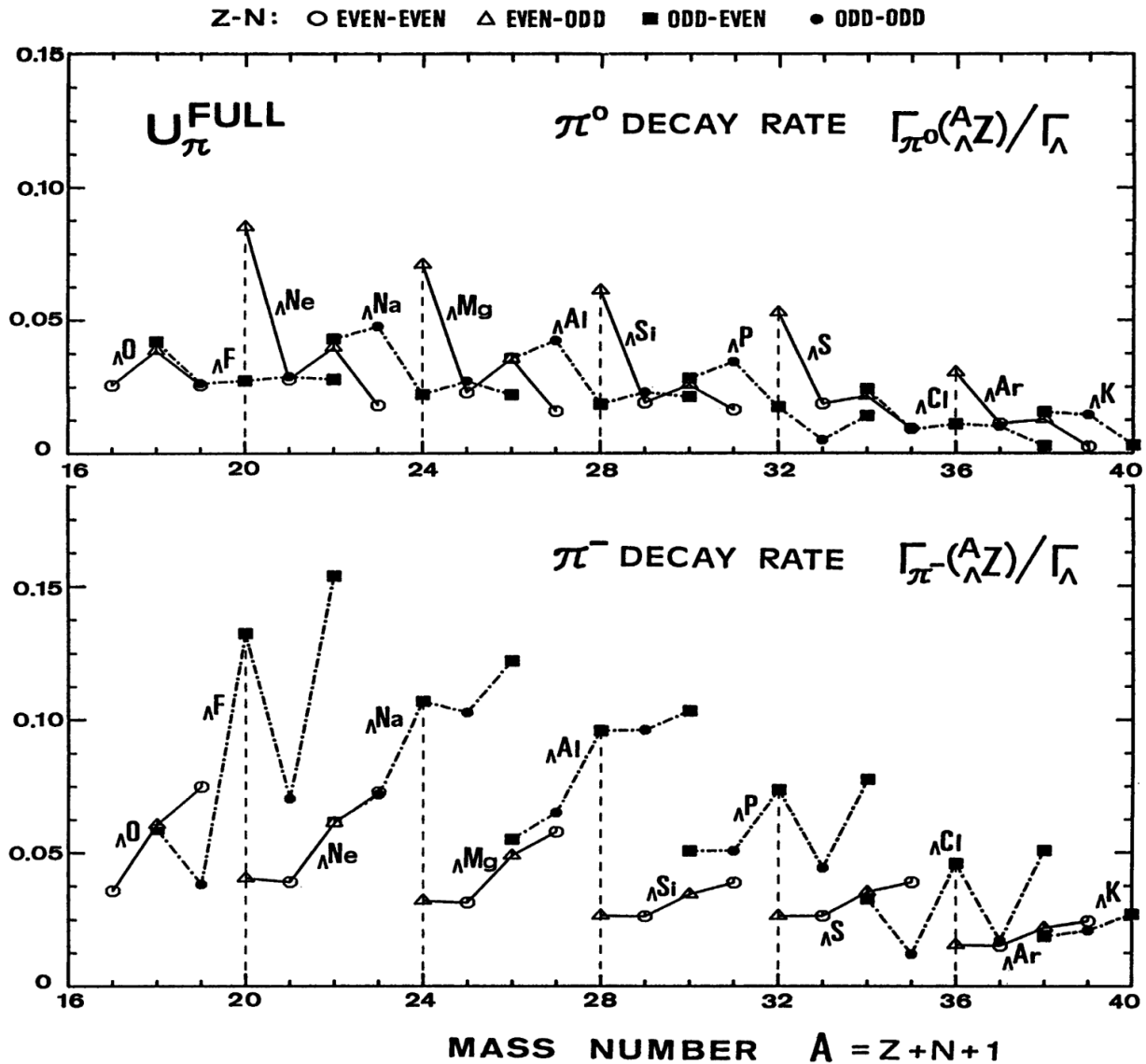


Potential dependence of decay rate

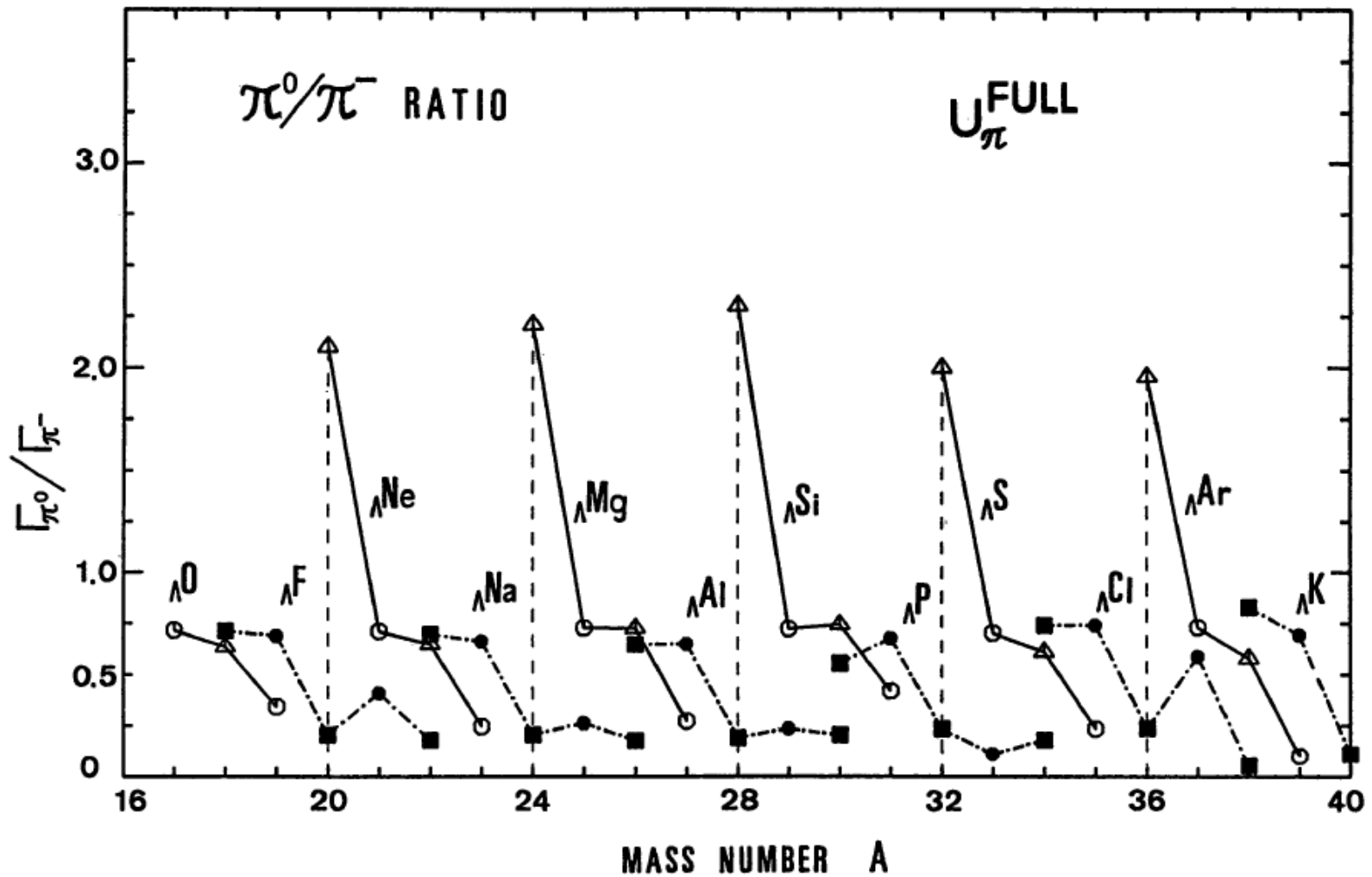


- H. Noumi et al. PRC52, 2936 (1995)
- J. J. Szymanski et al. PRC43, 849 (1991)
- A. Sakaguchi et al. PRC43, 73 (1991)
- T. Motoba et al., PTP Suppl. No.117, 477 (1994)

sd-shell hypernuclear π -decays



Ratio for Mdecay in sd-shell



Very recent assignment of ${}^7_{\Lambda}\text{Li}$ w/ Ge ball & SKS

•KEK-PS E419 experiment

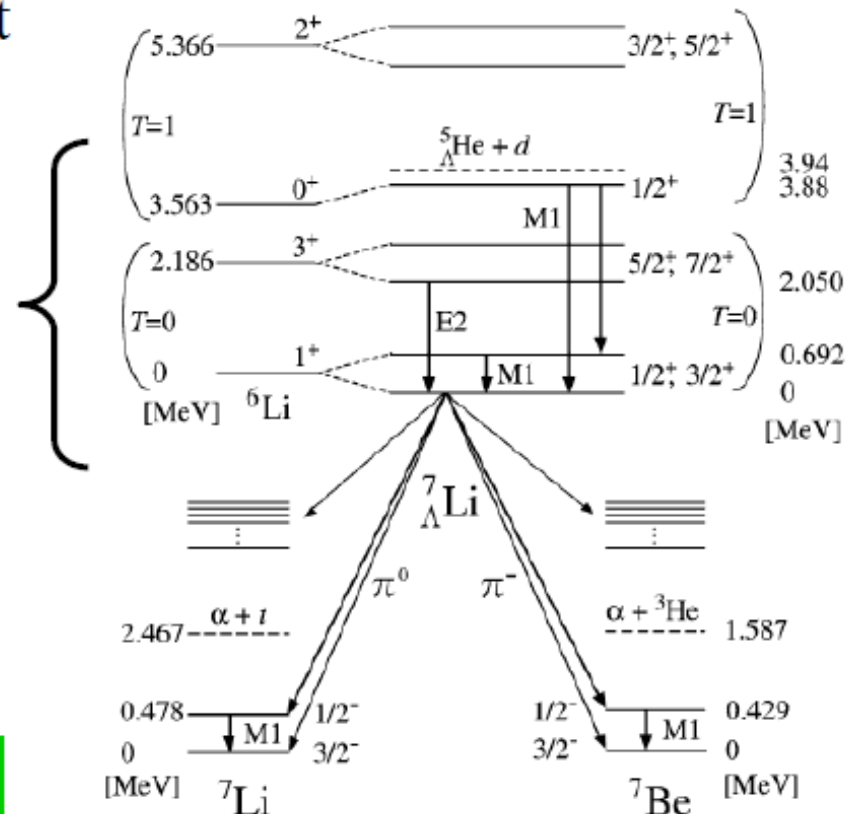
[Phys. Lett. B578 \(2004\) 258-264](#)

Formation of ${}^7_{\Lambda}\text{Li}$ was identified by (π^+, K^+) reaction

Coincidence measurement



M1 429keV



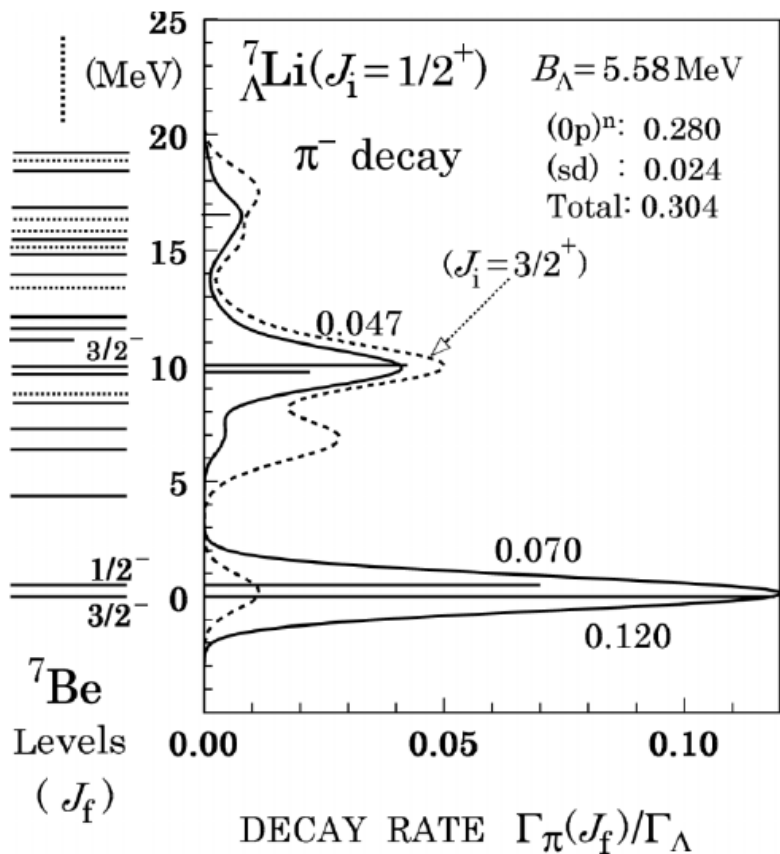


Fig. 6. Calculated π -decay spectrum as a function of ${}^7\text{Be}$ excitation energy.

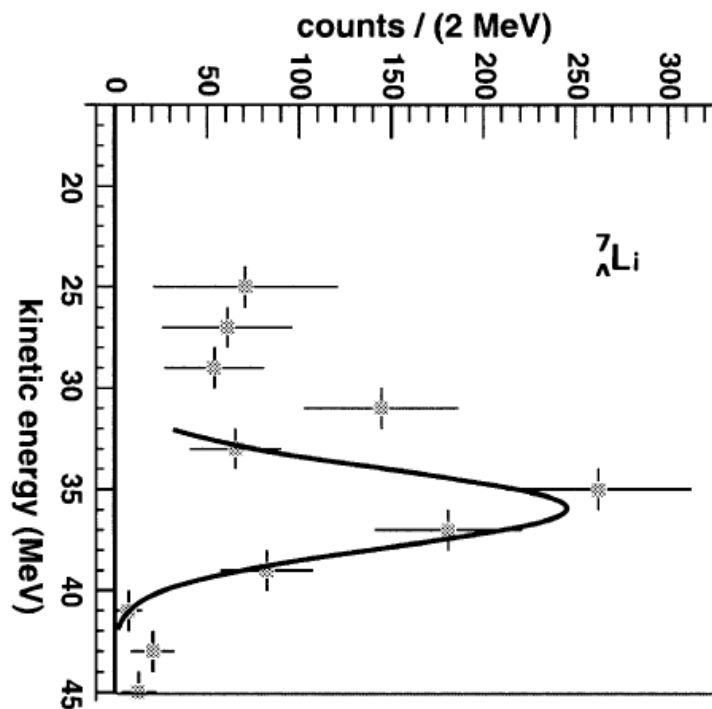


Fig. 7. π^- spectrum as a function of the kinetic energy. The data are taken from Ref. 37) and clipped for comparison.

Γ_{π}/Γ_{nm} and Λ -Nucleus Potential

Central repulsion in α - Λ potential ?

Can be directly checked by the mesonic decay widths

Mesonic decay rate

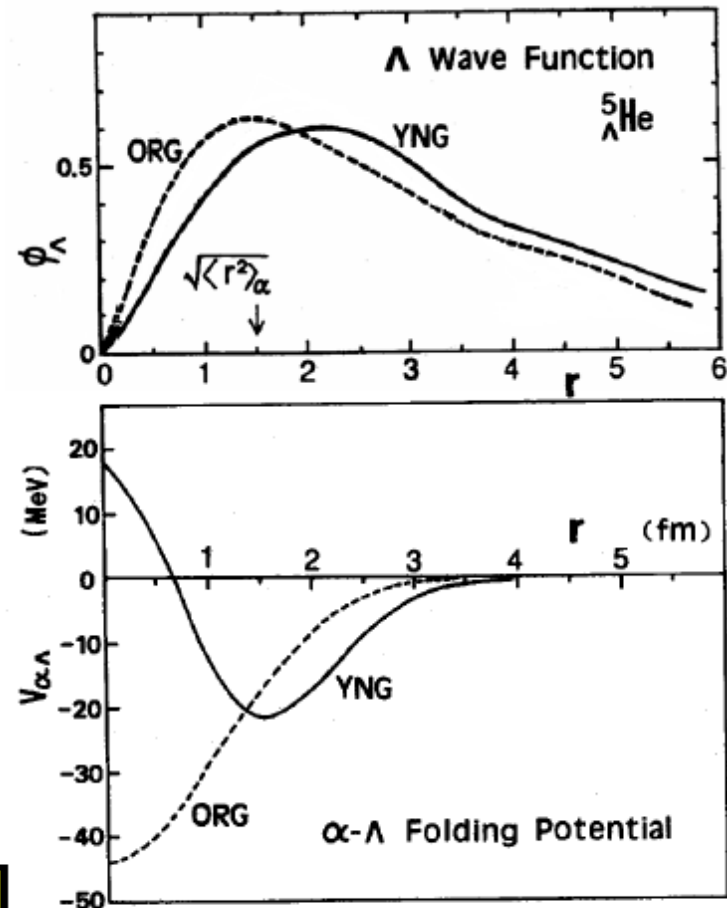
$$\Gamma_{\pi}(YNG) > \Gamma_{\pi}(ORG)$$

Non-mesonic decay rate

$$\Gamma_{nm} \propto \int \frac{\psi_N^2}{\rho_0} \cdot \psi_{\Lambda}^2 d\vec{r} \quad ??$$

$$\Gamma_{nm}(ORG) > \Gamma_{nm}(YNG)$$

Calculation using two types of α - Λ potential ORG and YNG (Motoba *et al.* NPA577)



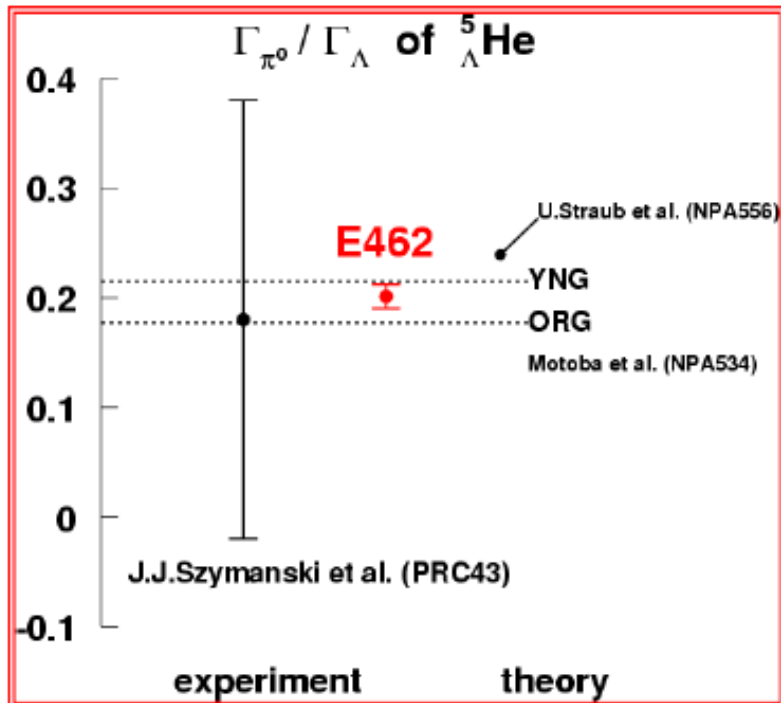
	Refs.	$\Gamma_{tot}/\Gamma_{\Lambda}$	$\Gamma_{\pi^-}/\Gamma_{\Lambda}$	$\Gamma_{\pi^0}/\Gamma_{\Lambda}$	$\Gamma_{nm}/\Gamma_{\Lambda}$
${}^5_{\Lambda}\text{He}$ (exp.)	[2]	1.03 ± 0.08	0.44 ± 0.11	0.18 ± 0.20	0.41 ± 0.14
${}^5_{\Lambda}\text{He}$ (ORG,SG)	[1][5]		0.321(ORG), 0.271(SG)	0.177(ORG), 0.158(SG)	
${}^5_{\Lambda}\text{He}$ (YNG,Isle)	[1][5]		0.393(YNG), 0.354(Isle)	0.215(YNG), 0.205(Isle)	
${}^{12}_{\Lambda}\text{C}$ (exp.)	[3][4]	1.14 ± 0.08	0.113 ± 0.015	0.200 ± 0.068	0.828 ± 0.087
${}^5_{\Lambda}\text{He}$ (exp.)	present	$0.940 \pm 0.040 \pm 0.007$	$0.322 \pm 0.018 \pm 0.003$	$0.207 \pm 0.012 \pm 0.005$	$0.411 \pm 0.023 \pm 0.006$
${}^{12}_{\Lambda}\text{C}$ (exp.)	present	$1.213 \pm 0.034 \pm 0.009$	$0.120 \pm 0.014 \pm 0.005$ [3]	$0.164 \pm 0.008 \pm 0.004$	$0.929 \pm 0.027 \pm 0.016$

Selected comparison to confirm the theory

Results of Γ_{π^0}

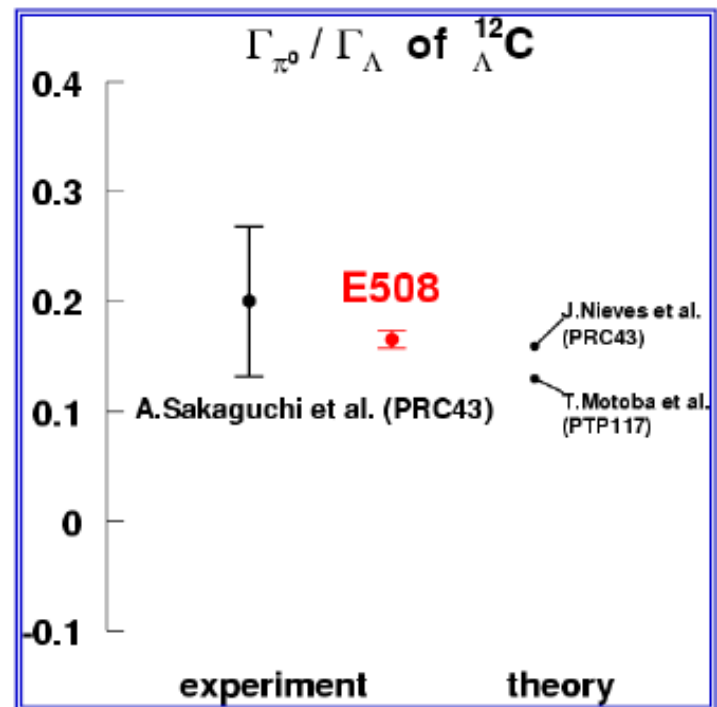
Lifetime : 278^{+11}_{-10} ps (E462)

$${}^5_{\Lambda}\text{He} : \Gamma_{\pi^0} / \Gamma_{\Lambda} = 0.207 \pm 0.013$$



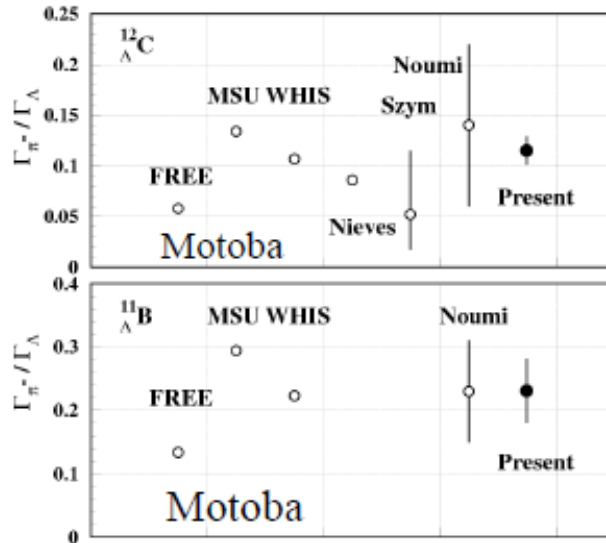
Lifetime : 212^{+7}_{-6} ps (E508)

$${}^{12}_{\Lambda}\text{C} : \Gamma_{\pi^0} / \Gamma_{\Lambda} = 0.164 \pm 0.009$$



Results of π^- Mesonic decay width (Outa slide)

Y.Sato *et al.*
PRC



Branching ratio

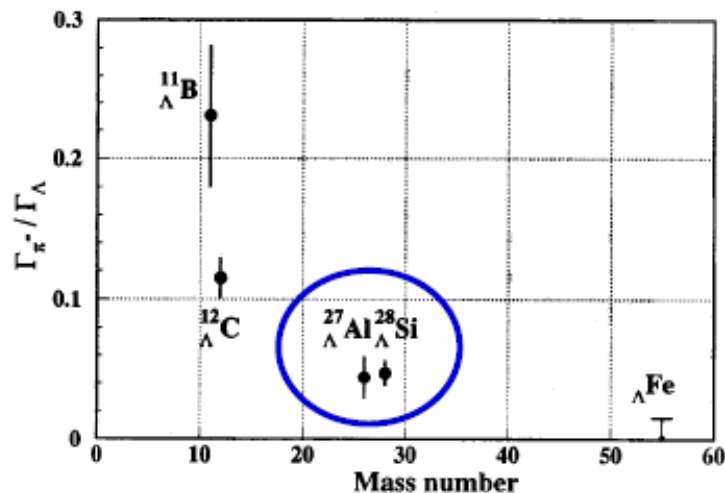
$$B_{\pi^-}(^{12}_{\Lambda}C) = 0.099 \pm 0.011 \pm 0.004$$

$$B_{\pi^-}(^{11}_{\Lambda}B) = 0.170 \pm 0.027 \pm 0.036$$

$$B_{\pi^-}(^{28}_{\Lambda}Si) = 0.036 \pm 0.008 \pm 0.002$$

$$B_{\pi^-}(^{27}_{\Lambda}Al) = 0.032 \pm 0.008 \pm 0.015$$

$$B_{\pi^-}(\Lambda Fe) < 0.012 \quad (90\% \text{ CL})$$



π^- Mesonic decay width

$$\Gamma_{\pi^-}(^{12}_{\Lambda}C) = 0.113 \pm 0.014 \pm 0.005 \Gamma_{\Lambda}$$

$$\Gamma_{\pi^-}(^{11}_{\Lambda}B) = 0.212 \pm 0.036 \pm 0.045 \Gamma_{\Lambda}$$

$$\Gamma_{\pi^-}(^{28}_{\Lambda}Si) = 0.046 \pm 0.011 \pm 0.002 \Gamma_{\Lambda}$$

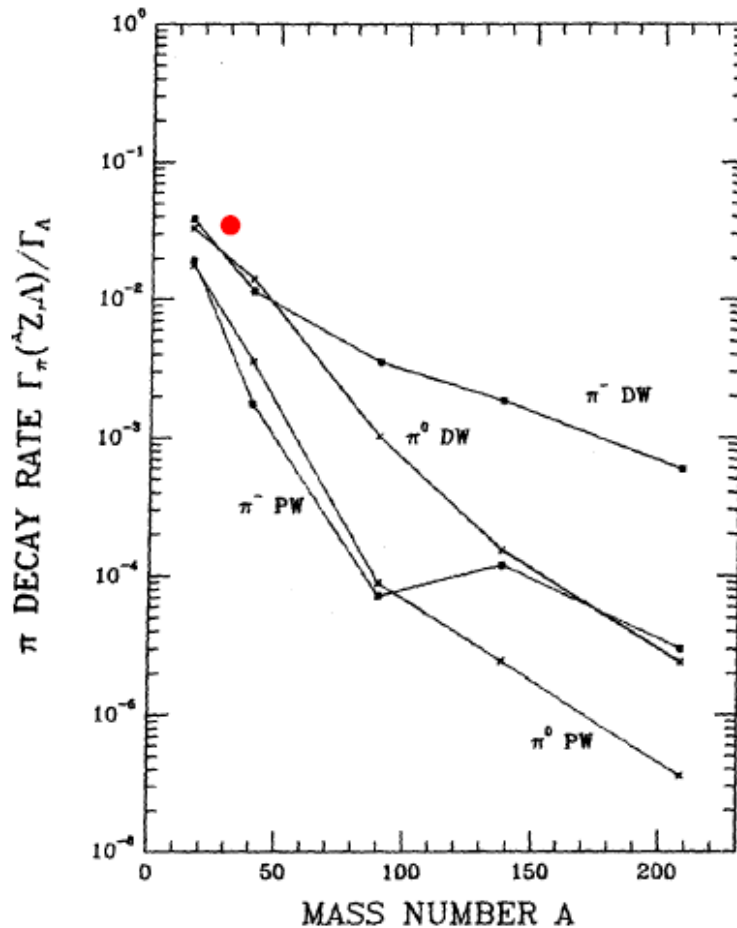
$$\Gamma_{\pi^-}(^{27}_{\Lambda}Al) = 0.041 \pm 0.010 \pm 0.019 \Gamma_{\Lambda}$$

$$\Gamma_{\pi^-}(\Lambda Fe) < 0.015 \Gamma_{\Lambda} \quad (90\% \text{ CL})$$

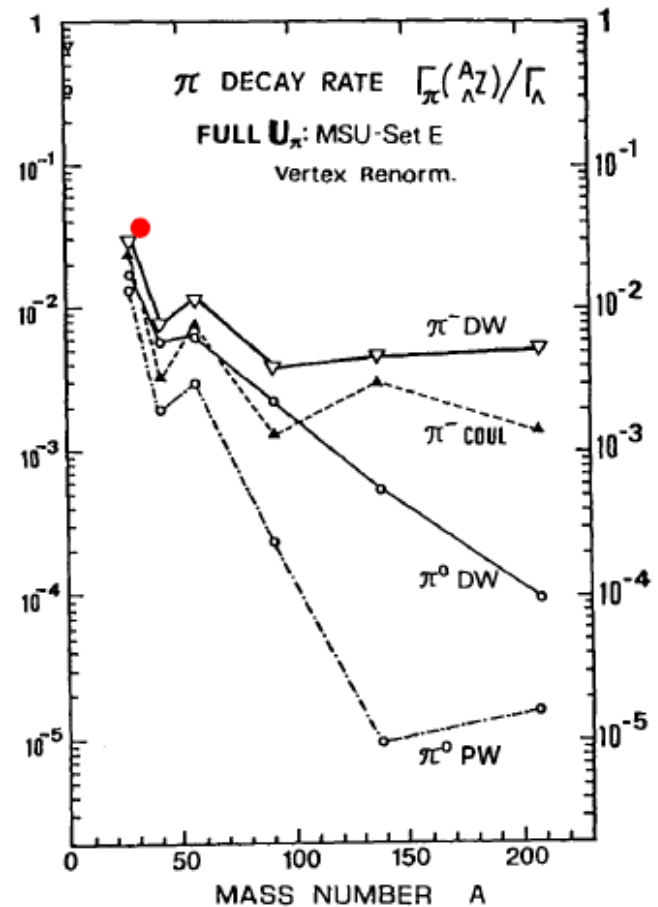
Gross behavior

Gross behavior of hypernuclear π -mesonic decay rate

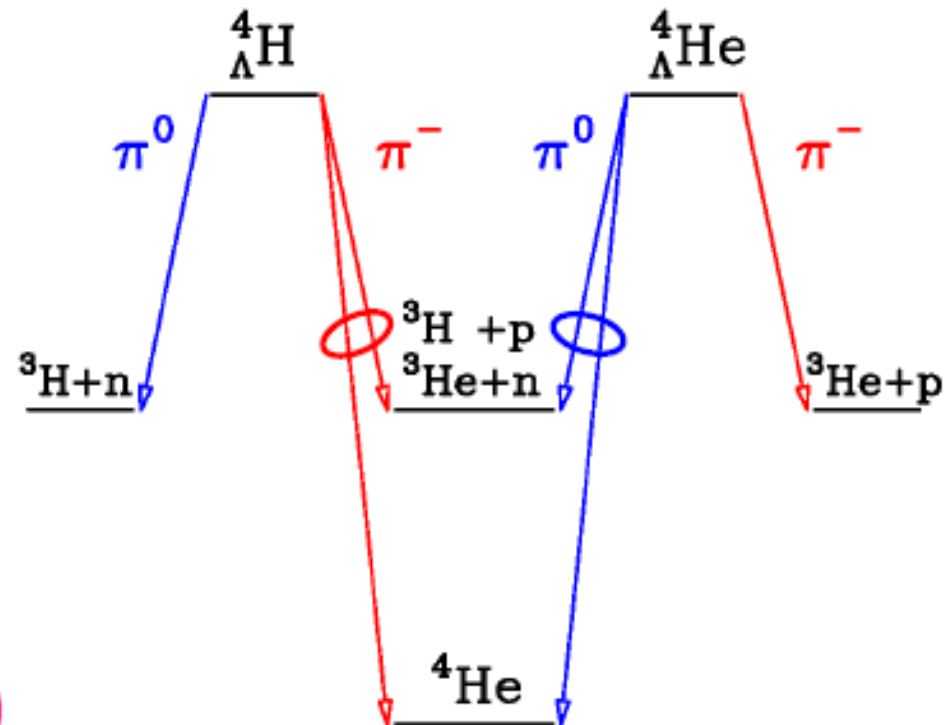
E.Oset et al. Prog. Theo. Phys. Suppl. No.117,461 (1994)



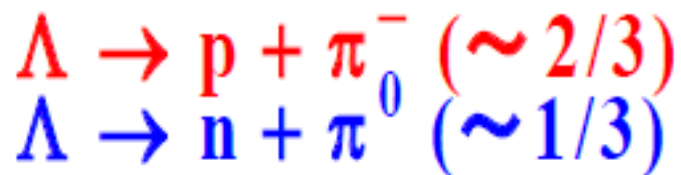
T. Motoba et al. Nucl. Phys. A547, 115c (1992)



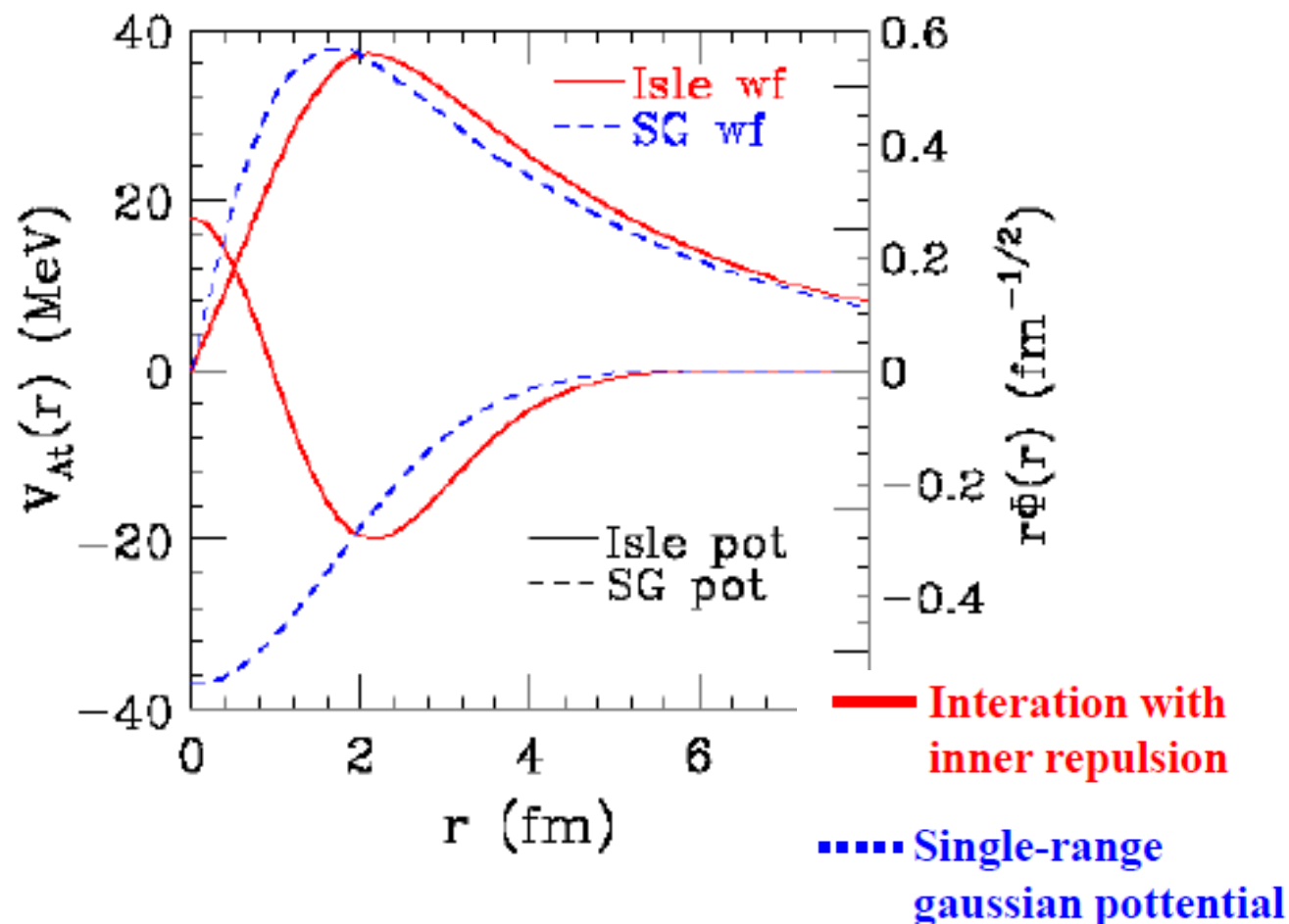
Mesonic decay scheme of ${}^4_{\Lambda}\text{H}$ and ${}^4_{\Lambda}\text{He}$ —“*Mirror*” hypernuclei



Free Λ decay



Test of Λ - nucleus potential for ${}^4_{\Lambda}\text{H}$ and ${}^4_{\Lambda}\text{He}$



Repulsive-core in Λ -nucleus potential?

Table II. Calculated decay rates, Γ_n/Γ_p ratios and intrinsic asymmetry parameters of s - and p -shell and medium-to-heavy hypernuclei. The interaction $V_\pi+V_{2\pi/\rho}+V_{2\pi/\sigma}+V_\omega+V_K+V_{\rho\pi/a_1}+V_{\sigma\pi/a_1}$ is adopted. For p -shell hypernuclei, calculations with configuration-mixed nuclear SM + Λ particle wave functions are shown without parentheses and those with simple nuclear SM + Λ particle wave functions are shown in parentheses. Calculations are done with restriction of $L_0 = 0$ only. The decay rates are given in units of the free Λ decay rate Γ_Λ . See the text.

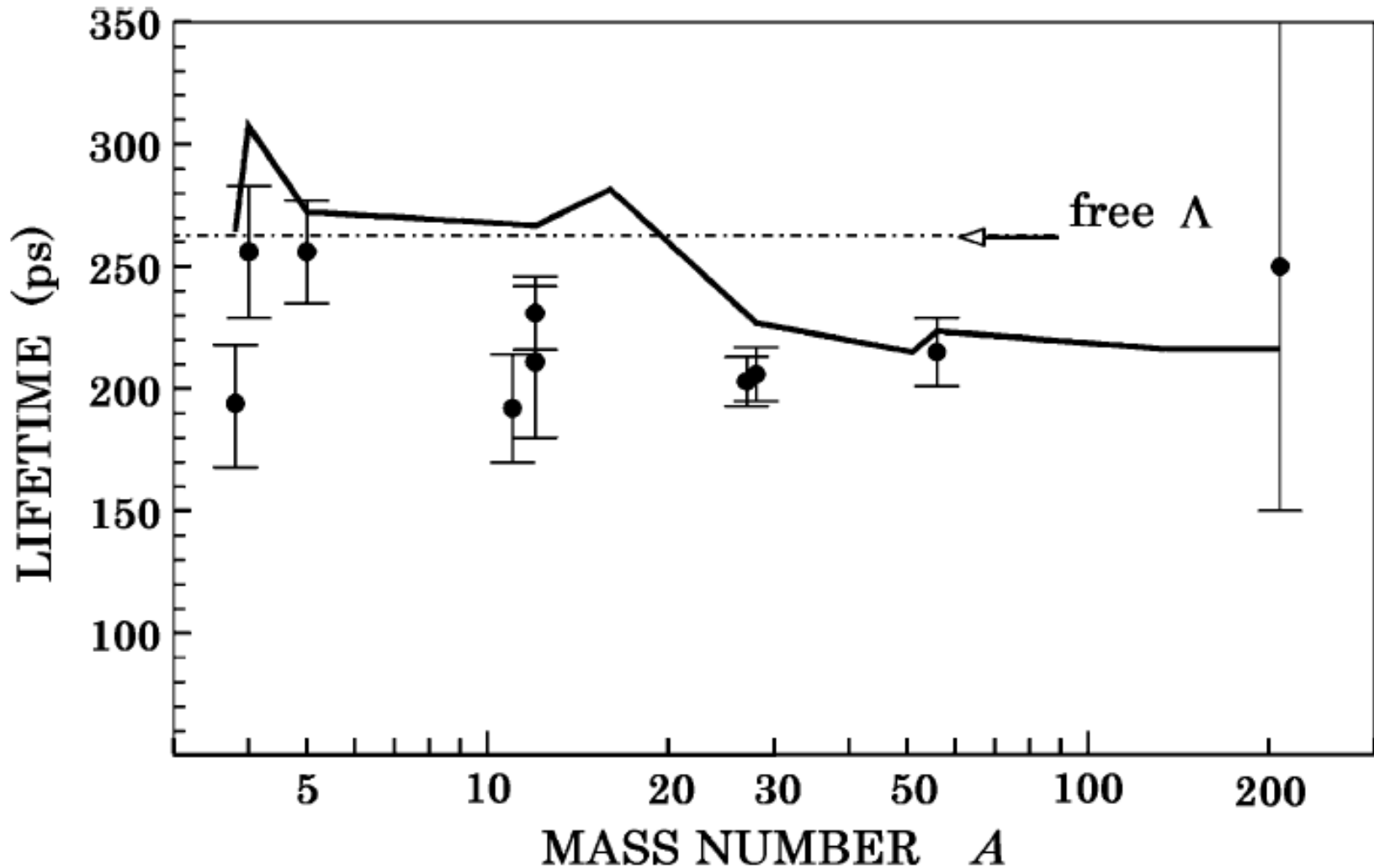
	Γ_p	Γ_n	Γ_{nm}	Γ_n/Γ_p	α_Λ
${}^4_\Lambda\text{H} (0^+)$	0.006	0.099	0.105	16.02	—
${}^4_\Lambda\text{He} (0^+)$	0.185	0.012	0.198	0.066	—
${}^5_\Lambda\text{He} (1/2^+)$	0.237	0.121	0.358	0.508	0.083
${}^7_\Lambda\text{Li} (1/2^+)$	0.297	0.152	0.449 [0.481]	0.512 [0.505]	0.138 [0.093]
${}^8_\Lambda\text{Li} (1^-)$	0.297	0.200	0.498 [0.501]	0.673 [0.657]	0.120 [0.113]
${}^9_\Lambda\text{Be} (1/2^+)$	0.401	0.200	0.601 [0.601]	0.499 [0.499]	0.053 [0.053]
${}^{10}_\Lambda\text{B} (1^-)$	0.442	0.198	0.640 [0.635]	0.448 [0.452]	0.083 [0.086]
${}^{10}_\Lambda\text{B}^* (2^-)$	0.492	0.198	0.690 [0.688]	0.403 [0.412]	0.024 [0.029]
${}^{11}_\Lambda\text{B} (5/2^+)$	0.444	0.223	0.667 [0.660]	0.502 [0.503]	0.072 [0.079]
${}^{12}_\Lambda\text{B} (1^-)$	0.446	0.266	0.711 [0.697]	0.596 [0.617]	0.061 [0.082]
${}^{12}_\Lambda\text{C} (1^-)$	0.535	0.223	0.758 [0.754]	0.418 [0.407]	0.044 [0.045]
${}^{12}_\Lambda\text{C}^* (2^-)$	0.536	0.237	0.773 [0.778]	0.443 [0.451]	0.044 [0.045]
${}^{13}_\Lambda\text{C} (1/2^+)$	0.495	0.247	0.741 [0.740]	0.499 [0.498]	0.023 [0.019]
${}^{14}_\Lambda\text{N} (1^-)$	0.551	0.246	0.797 [0.803]	0.446 [0.451]	0.034 [0.045]
${}^{15}_\Lambda\text{N} (3/2^+)$	0.555	0.278	0.833 [0.832]	0.502 [0.503]	0.037 [0.049]
${}^{16}_\Lambda\text{N} (1^-)$	0.519	0.298	0.816 [0.818]	0.574 [0.571]	0.038 [0.038]
${}^{16}_\Lambda\text{O} (0^-)$	0.586	0.275	0.860 [0.861]	0.469 [0.470]	—
${}^{16}_\Lambda\text{O}^* (1^-)$	0.586	0.264	0.850 [0.850]	0.451 [0.452]	0.037 [0.035]
${}^{28}_\Lambda\text{Si} (2^+)$	0.735	0.336	1.071	0.456	0.024
${}^{51}_\Lambda\text{V} (11/2^+)$	0.792	0.412	1.203	0.520	0.017
${}^{56}_\Lambda\text{Fe} (1^-)$	0.764	0.392	1.156	0.513	0.008
${}^{89}_\Lambda\text{Y} (7/2^-)$	0.774	0.418	1.192	0.540	-0.003
${}^{139}_\Lambda\text{La} (9/2^+)$	0.757	0.455	1.212	0.601	-0.007
${}^{209}_\Lambda\text{Bi} (9/2^+)$	0.733	0.477	1.210	0.651	-0.016

Calculated MWD and NMWD rates

Table VIII. Calculated hypernuclear weak decay rates and lifetimes for s -shell, p -shell, medium and heavy hypernuclei. The decay rates are given in units of the free Λ decay rate Γ_Λ . The lifetimes are given in units of picosecond. Experimental data of nonmesonic decay rates and lifetimes are listed for comparison. See the text.

	Γ_π	Γ_{nm}	$\Gamma_\pi + \Gamma_{nm}$	τ	Γ_{nm}^{exp}	τ^{exp}
${}^4_\Lambda\text{H}$	0.891	0.105	0.996	264.1	0.17 ± 0.11 [Ref. 62)]	194^{+24}_{-26} [Ref. 62)]
${}^4_\Lambda\text{He}$	0.658	0.198	0.856	307.2	0.17 ± 0.05 [Ref. 62)] 0.20 ± 0.03 [Ref. 63)]	256 ± 27 [Ref. 62)]
${}^5_\Lambda\text{He}$	0.608	0.358	0.966	272.3		256 ± 21 [Ref. 64)]
${}^{11}_\Lambda\text{B}$	0.316	0.667	0.983	267.5	$0.95 \pm 0.13 \pm 0.04$ [Ref. 67)] $0.861 \pm 0.063 \pm 0.073$ [Ref. 51)]	192 ± 22 [Ref. 65)] 211 ± 13 [Ref. 66)]
${}^{12}_\Lambda\text{C}$	0.228	0.758	0.986	266.7		211 ± 31 [Ref. 65)] 231 ± 15 [Ref. 66)]
${}^{16}_\Lambda\text{O} (0^-)$	0.074*	0.860	0.934	281.6		86^{+33}_{-26} [Ref. 68)]
${}^{28}_\Lambda\text{Si}$	0.088	1.071	1.159	226.9	$1.125 \pm 0.067 \pm 0.106$ [Ref. 51)]	206 ± 12 [Ref. 66)]
${}^{51}_\Lambda\text{V}$	0.02*	1.203	1.223	215.0		
${}^{56}_\Lambda\text{Fe}$	0.02*	1.156	1.176	223.6	1.21 ± 0.08 [Ref. 51)]	$215 \pm 14^{**}$ [Ref. 66)]
${}^{89}_\Lambda\text{Y}$	0.005*	1.192	1.197	219.7		
${}^{139}_\Lambda\text{La}$	0.005*	1.212	1.217	216.1		
${}^{209}_\Lambda\text{Bi}$	0.005*	1.210	1.215	216.5		250^{+250}_{-100} [Ref. 69)]

Lifetimes (light to heavy hypernuclei)



MWD rates for p-shell hypernuclei

Table IX. π^- -decay rates of p -shell hypernuclei in units of the free Λ decay rate Γ_Λ .

	${}^7_\Lambda\text{Li}$	${}^8_\Lambda\text{Li}$	${}^8_\Lambda\text{Be}$	${}^9_\Lambda\text{Be}$	${}^{10}_\Lambda\text{B}$	${}^{11}_\Lambda\text{B}$	${}^{12}_\Lambda\text{B}$	${}^{12}_\Lambda\text{C}$	${}^{13}_\Lambda\text{C}$	${}^{14}_\Lambda\text{N}$	${}^{15}_\Lambda\text{N}$
J_{gs}	$1/2^+$	1^-	1^-	$1/2^+$	1^-	$5/2^+$	1^-	1^-	$1/2^+$	1^-	$3/2^+$
$(0p)^n$	0.280	0.344	0.130	0.144	0.197	0.195	0.256	0.077	0.079	0.053	0.067
$(1s0d)$	0.024	0.024	0.019	0.029	0.018	0.018	0.030	0.022	0.025	0.020	0.023
$\Gamma^{(\text{Sum})}$	0.304	0.368	0.149	0.172	0.215	0.213	0.286	0.099	0.105	0.073	0.090
$\Gamma_{\pi^-}^{(\text{AG})}$	0.332			0.157		0.178					0.066
Exp ³⁷⁾	0.353			0.178		0.249					0.108
\pm error	± 0.059			± 0.050		± 0.051					± 0.038
Exp ^{65), 82)}						0.218		0.05			
\pm error						± 0.046		$+0.06$ -0.03			

Comparison with FINUDA data

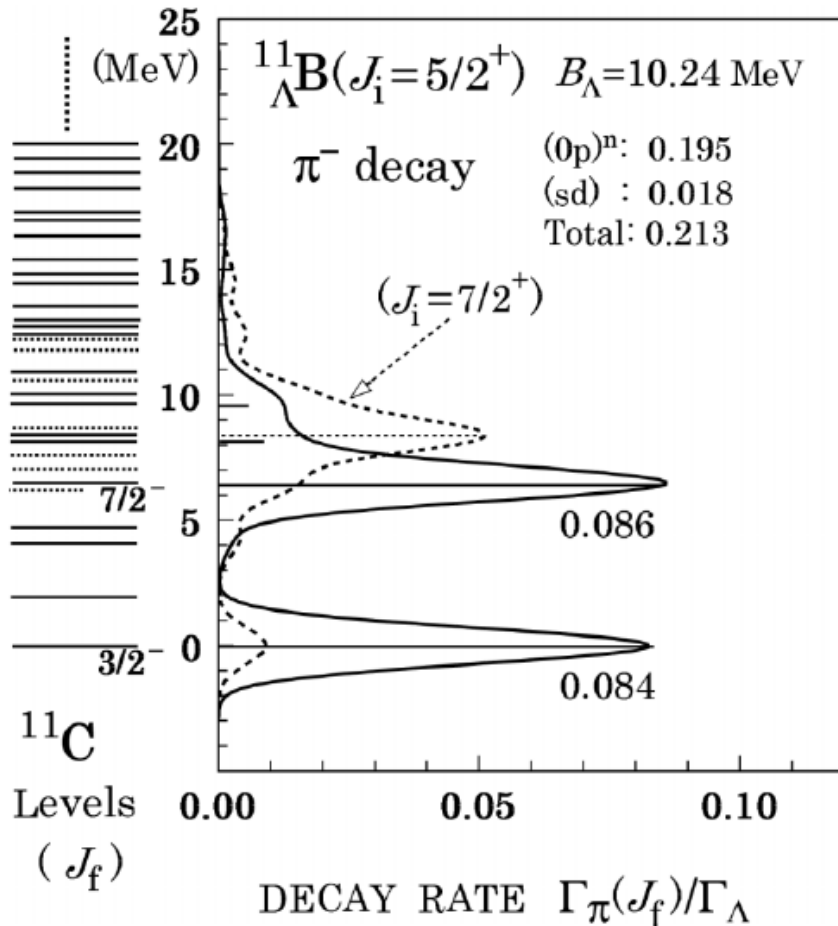


Fig. 10. Calculated π^- -decay spectrum as a function of ^{11}C excitation energy.

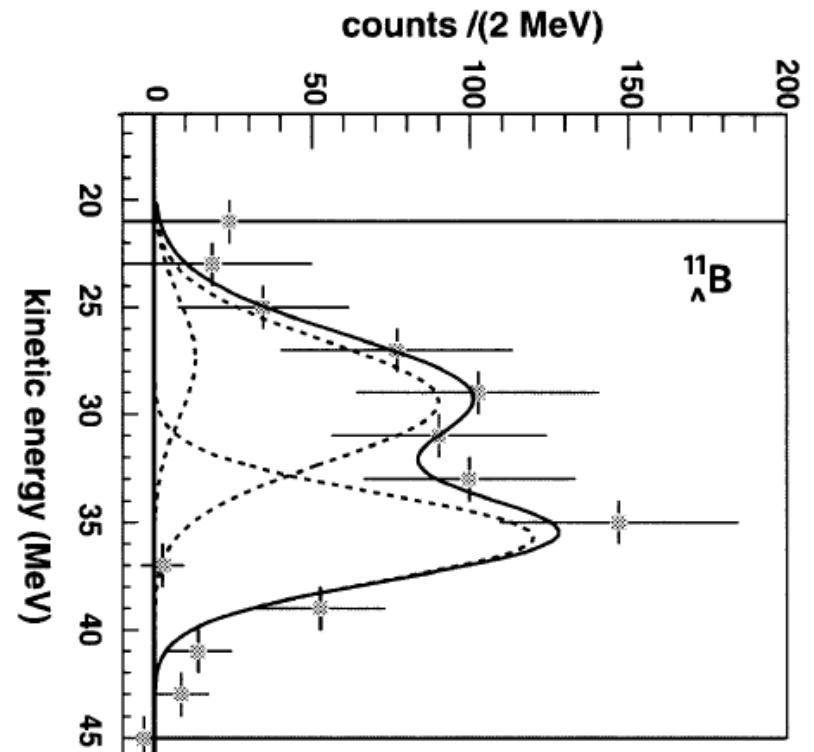


Fig. 11. π^- spectrum as a function of the kinetic energy. The data are taken from Ref. 37) and clipped for comparison.

Comparison with recent FINUDA exp.

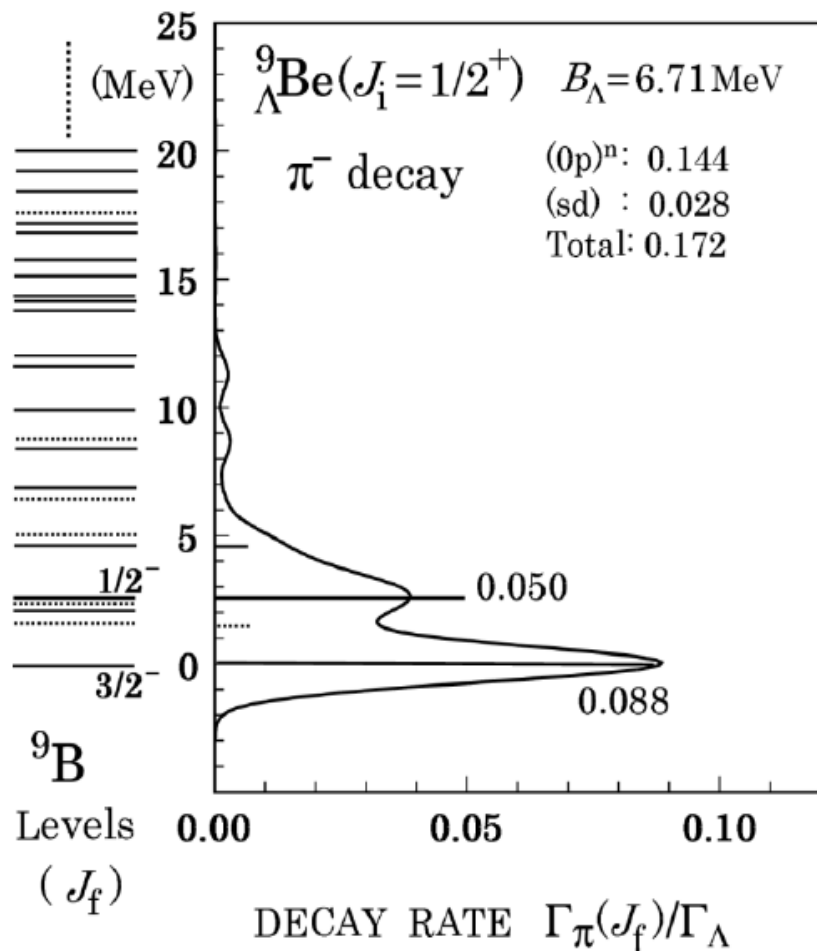


Fig. 8. Calculated π -decay spectrum as a function of ${}^9\text{B}$ excitation energy.

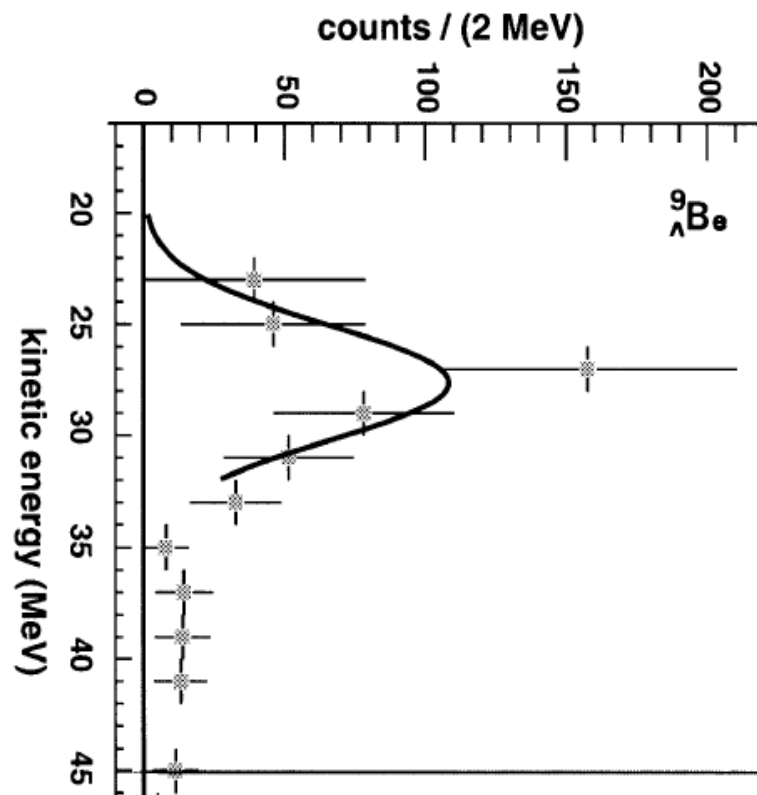
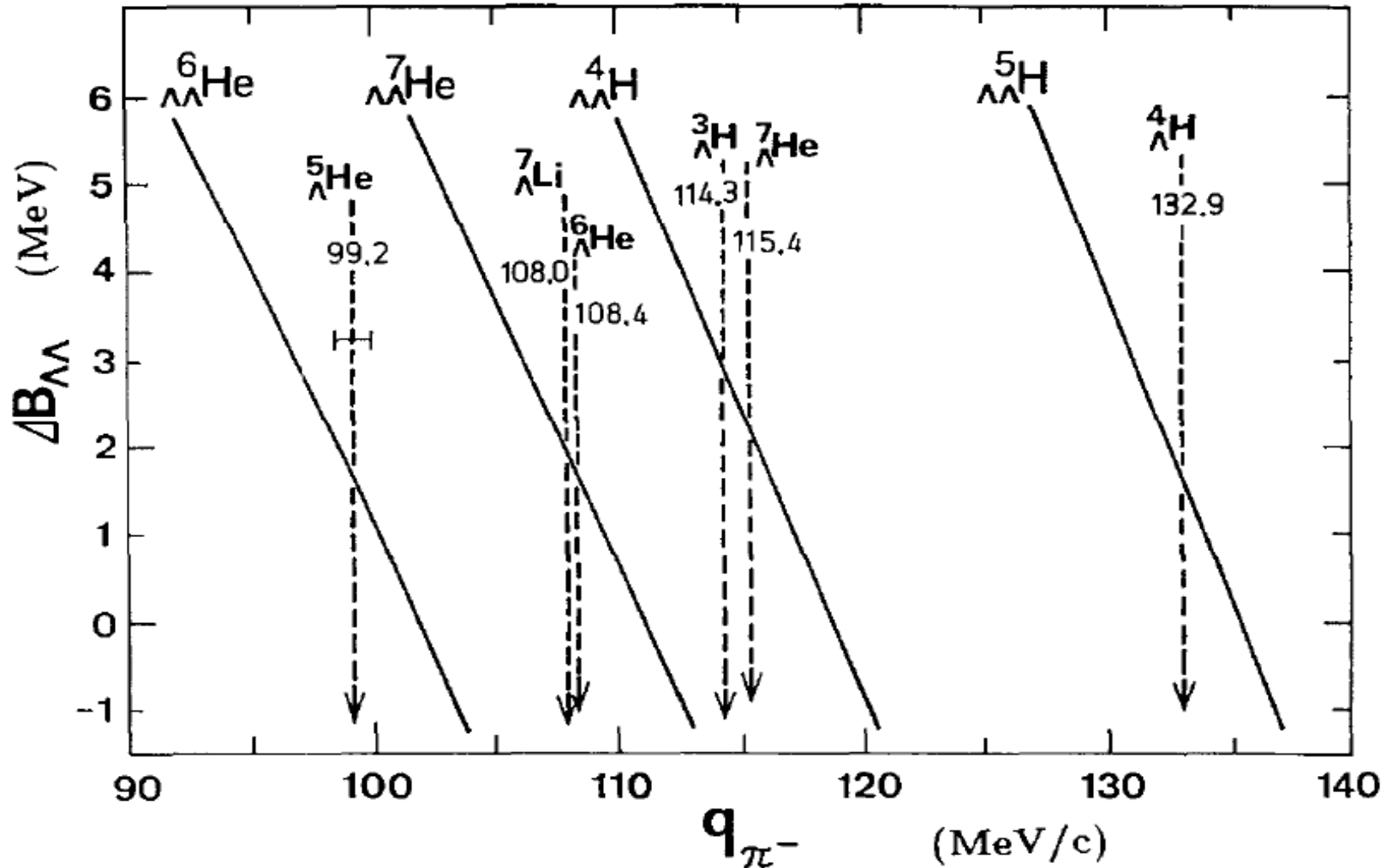


Fig. 9. π^- spectrum as a function of the kinetic energy. The data are taken from Ref. 37) and clipped for comparison.

Characteristic pion momentum helps to identify the hypernucleus



Case of 2-body and 3-body decays

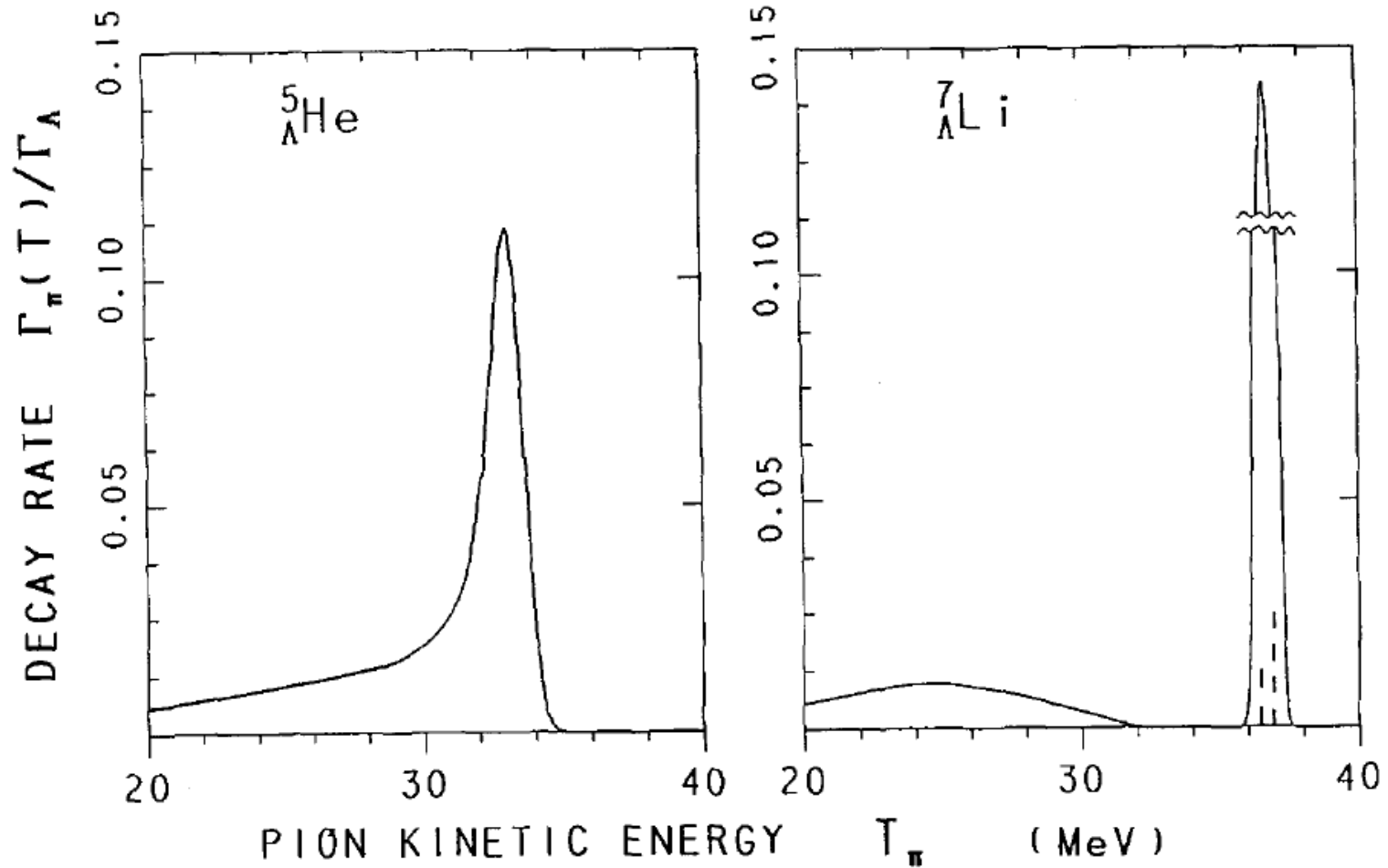


Fig. 10. The continuum pion spectrum calculated for the weak decay ${}^5_{\Lambda}\text{He} \rightarrow {}^4\text{He} + p + \pi^{-}$ (left) and that for the two-body and three-body π^{-} decay of ${}^7_{\Lambda}\text{Li}$ (right).

2-body and 3-body decays

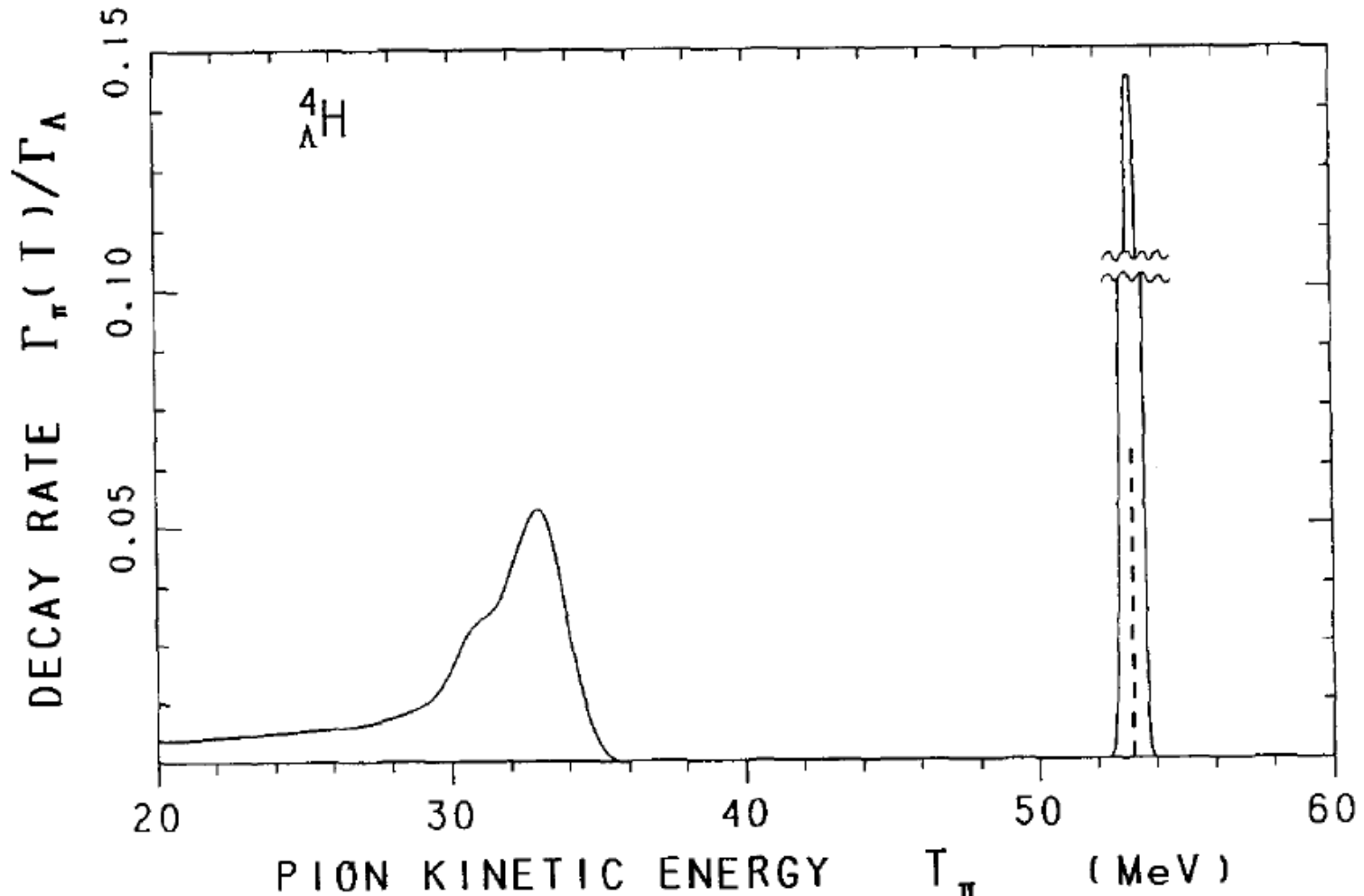


Fig. 11. The calculated π^- spectrum from the weak decay of ${}^4_{\Lambda}\text{H}$, which consists of the monochromatic peak for the two-body decay (${}^4\text{He} + \pi^-$) and the continuum part for the three-body decays (${}^3\text{H} + p + \pi^-$).

Pion spectrum helps

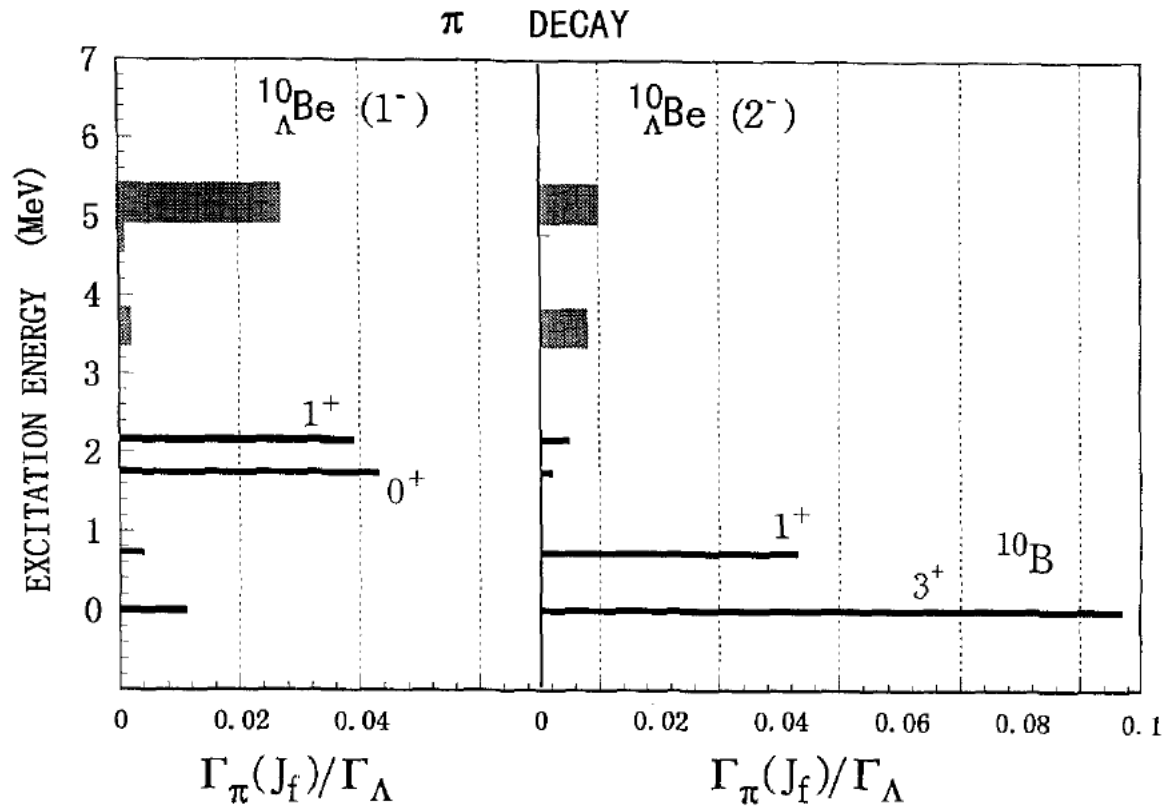


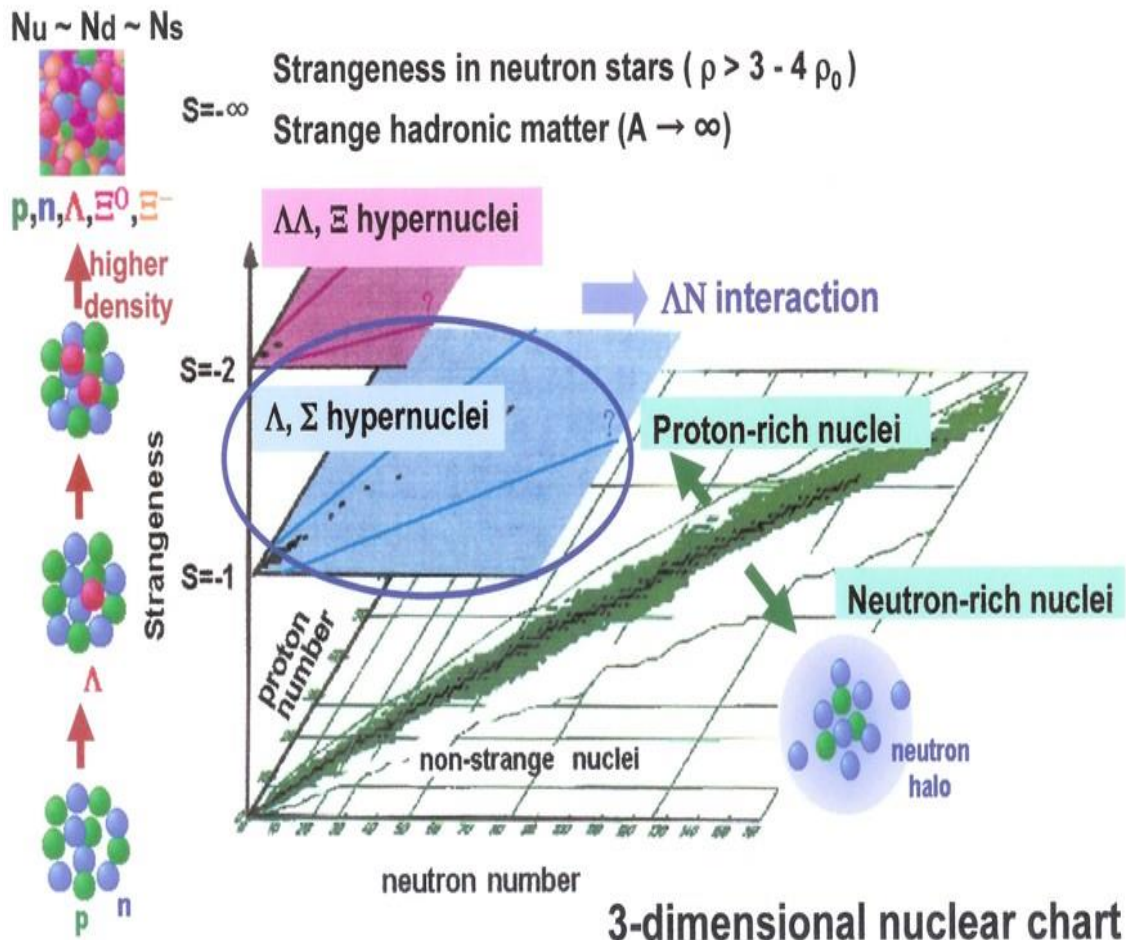
Fig. 18. The calculated π^- spectrum from the weak decay of $^{10}_{\Lambda}\text{Be}(1^-, 2^-)$. Comment as for Fig. 13.

Summary

- 1) Pionic decay rates and spectra have been extensively calculated, using pion optical potential which reproduces nucleus-pion scattering at low energies.
- 2) Sensitivity of pi-decay momentum helps to identify the hypernucleus unambiguously.
- 3) Pionic decay spectrum also helps to determine the Ground-state spin.
- 4) Coincidence measurements are promising

Three-dimensional nuclear chart

(from H. Tamura)



(First 3D illustration

appeared in literature:

Genshikaku Kenkyu **32**

(1987) 97.

See P.T.P. S.185 (2010) 1.)

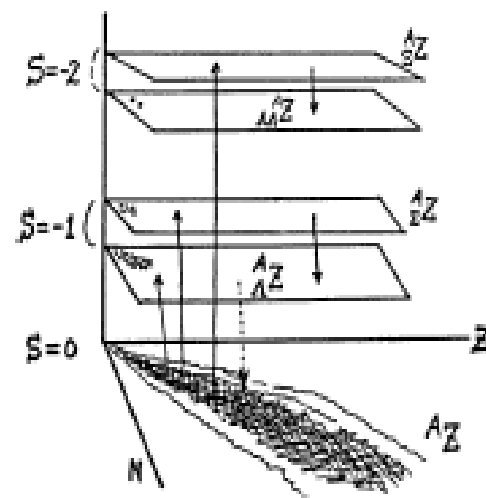


Fig. 3. 3D-illustration of hypernuclear chart drawn for the first time with the strangeness axis added (1986).⁶¹⁾