

Numerical Investigation of the 2-D $\mathcal{N}=2$ Wess-Zumino Model

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Motivation (Physics)

- The lattice breaks supersymmetry explicitly.
- No spontaneous supersymmetry breaking of the continuum model expected.
⇒ **Supersymmetry restoration** in continuum limit can be analyzed.
- In former works (M. Beccaria et al. (1998), S. Catterall and S. Karamov (2003)) only Wilson fermions with **Nicolai improved action** were used. Problems at stronger couplings.
- Effects of Nicolai improvement?

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2. Motivation (Algorithms)

- Explicit investigation and improvement of the used algorithms, cf. e.g. **Bergner et al. (2007)** for WZ model in $1d$ with different discretizations.
- High precision measurements available in lower dimensions.

- The continuum action

$$S_{\text{cont}} = \int d^2x \left(2\bar{\partial}\bar{\varphi}\partial\varphi + \frac{1}{2}|W'(\varphi)|^2 + \bar{\psi}M\psi \right),$$

$$M = \gamma^z\partial + \gamma^{\bar{z}}\bar{\partial} + W''P_+ + \overline{W}''P_-$$

allows for **4 real supersymmetries**, $\varphi = \varphi_1 + i\varphi_2$.

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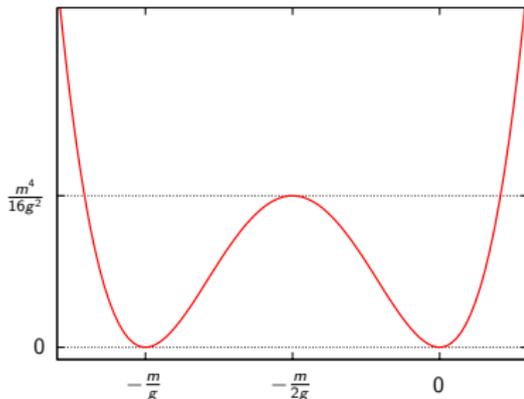
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allows for **4 real supersymmetries**, $\varphi = \varphi_1 + i\varphi_2$.

- We use $W(\varphi) = \frac{1}{2}m\varphi^2 + \frac{1}{3}g\varphi^3$ with dimensionless coupling $\lambda = \frac{g}{m}$.

Classical potential $|W'(\varphi_1)|^2$:



- $\lambda = 0$ corresponds to free theory
 \Rightarrow **perturbative expansion** in λ possible.

Using the Nicolai variable $\xi_x = 2(\bar{\partial}\bar{\varphi})_x + W_x$ an action on the lattice preserving **one supersymmetry** is given by

$$S = \frac{1}{2} \sum_x \bar{\xi}_x \xi_x + \sum_{xy} \bar{\psi}_x M_{xy} \psi_y$$

with $W_x = W'(\varphi_x)$, $W_{xy} := \partial W_x / \partial \varphi_y$ and

$$M_{xy} = \begin{pmatrix} W_{xy} & 2\bar{\partial}_{xy} \\ 2\partial_{xy} & W_{xy} \end{pmatrix} = \begin{pmatrix} \frac{\partial \xi_x}{\partial \varphi_y} & \frac{\partial \xi_x}{\partial \bar{\varphi}_y} \\ \frac{\partial \xi_x}{\partial \varphi_y} & \frac{\partial \xi_x}{\partial \bar{\varphi}_y} \end{pmatrix}.$$

The model

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In terms of the original fields the action reads

$$S = \sum_x \left(2(\bar{\partial}\bar{\varphi})_x (\partial\varphi)_x + \frac{1}{2} |W_x|^2 + W_x (\partial\varphi)_x + \bar{W}_x (\bar{\partial}\bar{\varphi})_x \right) + \sum_{xy} \bar{\psi}_x M_{xy} \psi_y.$$

The difference to a straightforward discretization is given by **surface terms**

$$\Delta S = \sum_x \left(W_x (\partial\varphi)_x + \bar{W}_x (\bar{\partial}\bar{\varphi})_x \right).$$

The model

The lattice discretization

We use different lattice derivatives (*the same for bosonic and fermionic degrees of freedom*):

- Symmetric derivative $(\partial_\mu^S)_{xy} = \frac{1}{2}(\delta_{x+\hat{\mu},y} - \delta_{x-\hat{\mu},y})$ with **standard Wilson** term $W_x = W'(\varphi_x) - \frac{r}{2}(\Delta\varphi)_x$ using ($r = 1$).

$$M_{xy} = \begin{pmatrix} W'''(\phi_x)\delta_{xy} & 2\bar{\partial}_{xy} \\ 2\partial_{xy} & W''(\phi_x)\delta_{xy} \end{pmatrix} - \frac{r}{2}\Delta_{xy}$$

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The choice $r = 2/\sqrt{3}$ renders the mass of the free theory exact up to $\mathcal{O}(a^4)$.

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- **SLAC derivative** $\partial_{x \neq y} = (-1)^{x-y} \frac{\pi/N}{\sin(\pi(x-y)/N)}$, $\partial_{xx} = 0$ with

M_{xy} unchanged.

⇒ Simulate the (un)improved model with these different discretizations!

We use a **combination of fourier accelerated HMC with higher-order integrators**.

Limitations of improvement

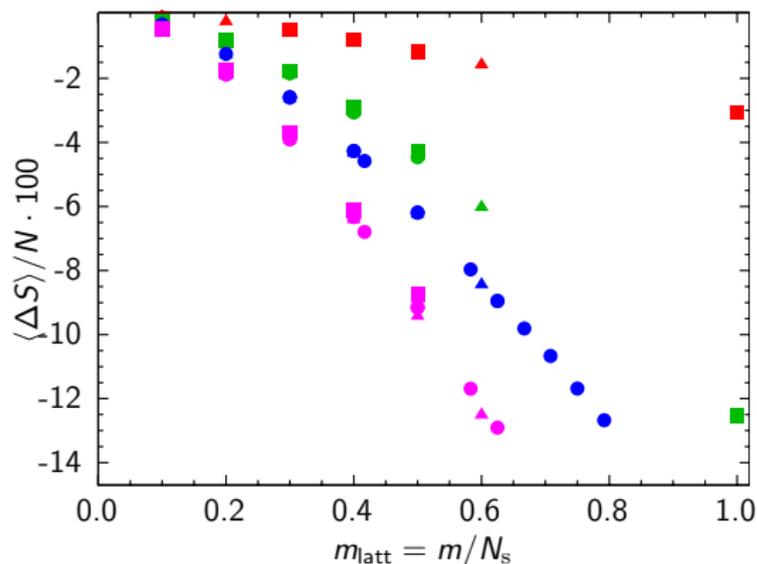
For **dynamical simulations** of the **improved model** the bosonic action is fixed to $\langle S_B \rangle = N = \#$ lattice points.

With SLAC fermions at different coupling strenghts we observe the improvement term $\Delta S = \sum_x \left(W_x (\partial\varphi)_x + \overline{W}_x (\partial\overline{\varphi})_x \right)$:

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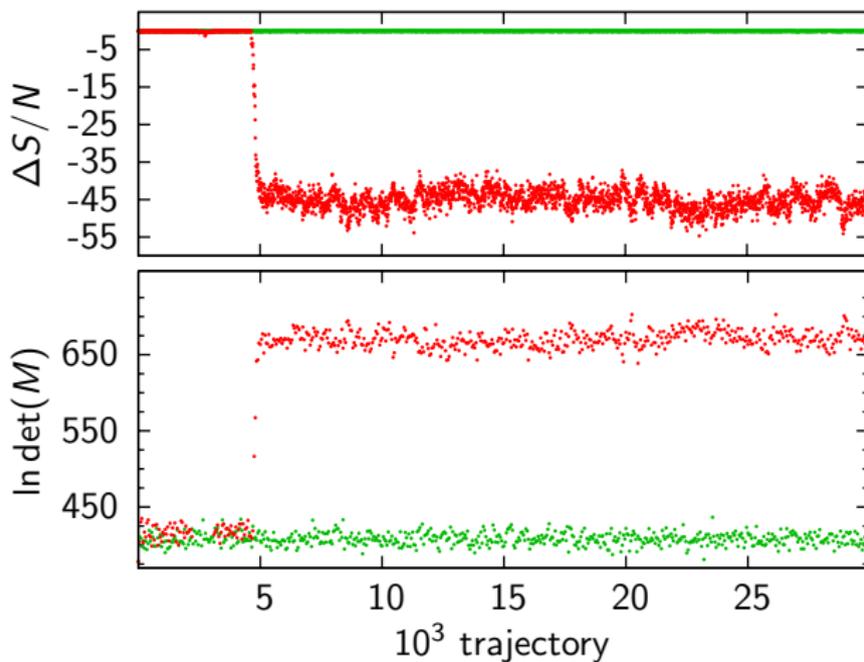
$N \in \{9 \times 9, 15 \times 15, 25 \times 25\}$

$\lambda = 0.8, 1.0, 1.2, 1.5$

Simulations **break down** when $\langle \Delta S \rangle / N$ exceeds 14%.

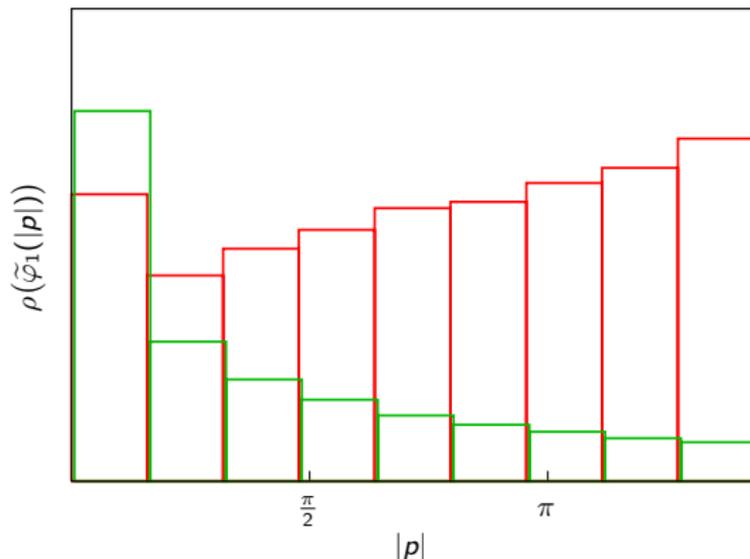
Limitations of improvement

MC history of the improvement term and the fermion determinant at $\lambda = 1.4$ and $\lambda = 1.7$ ($m_{\text{latt}} = 0.6$, $N = 15 \times 15$), $\langle S_B \rangle \approx N$ in each run:



Limitations of improvement

Analyzing the distribution of the fields in momentum space at $\lambda = 1.4$ and $\lambda = 1.7$:



\Rightarrow For too large couplings λ (or lattice masses m_{latt}) the simulation samples only **unphysical UV dominated** configurations.

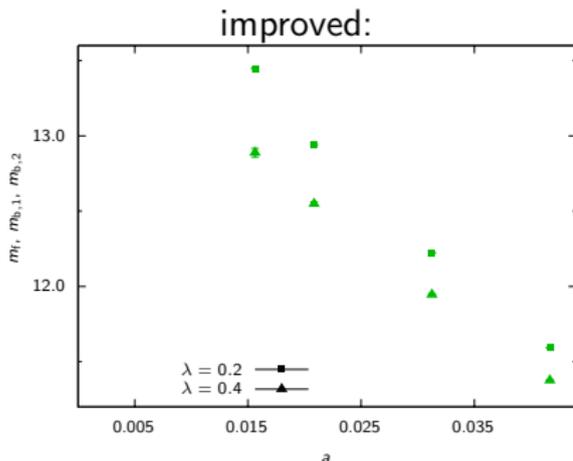
\Rightarrow At larger couplings a **careful analysis of the improvement term** during the simulation must be ensured.

Weak coupling results

Bosons vs. fermions

With Wilson fermions we test for supersymmetry breaking effects on the lattice at different lattice spacings for $\lambda \in \{0.2, 0.4\}$, $m = 15$.

Masses for bosons (φ_1, φ_2 , statistics 10^6 – 10^7 configs)
and fermions (statistics 10^4 configs)

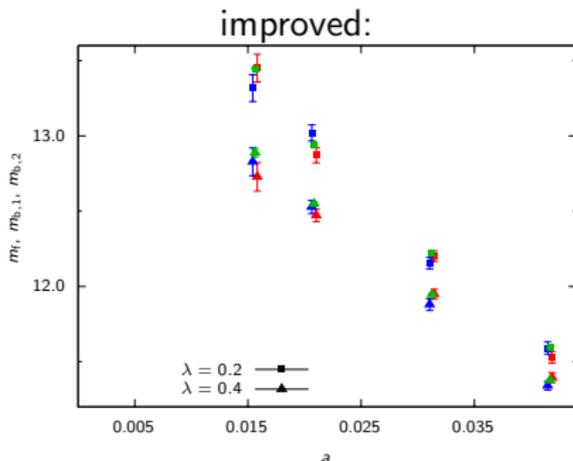


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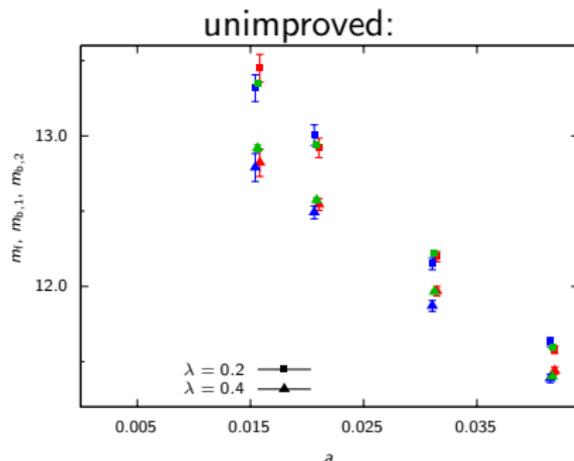
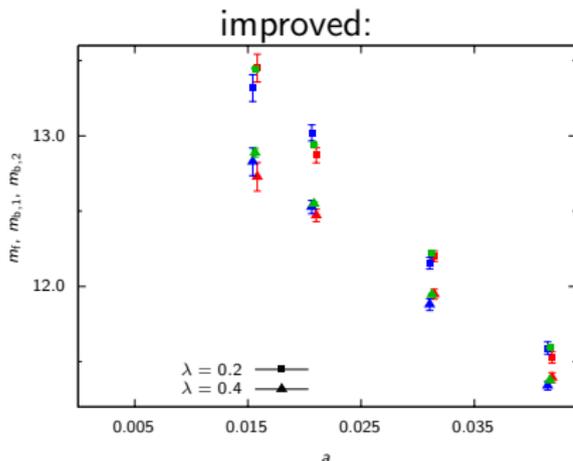


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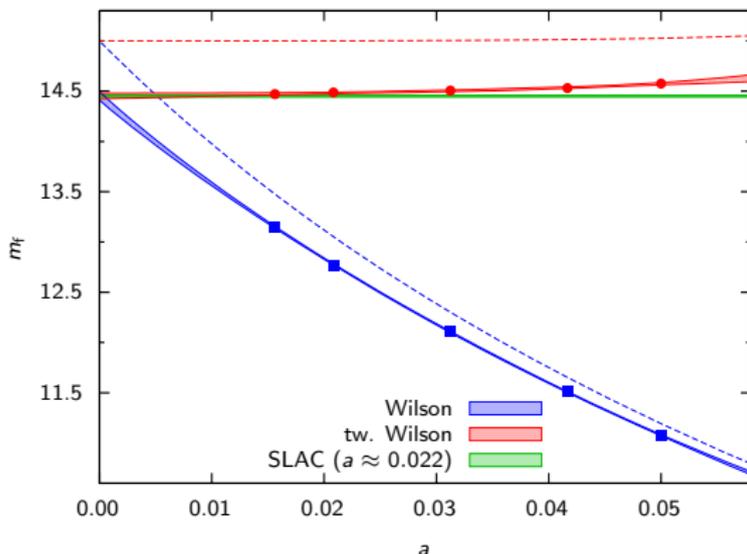
⇒ Improved and unimproved model can **not be distinguished** even with that high statistics.

⇒ Bosonic and fermionic masses coincide.

Weak coupling results

Continuum extrapolation

Extrapolation from finite lattice spacing to the continuum using Wilson and twisted Wilson fermions for the improved model ($m = 15$, $\lambda = 0.3$):

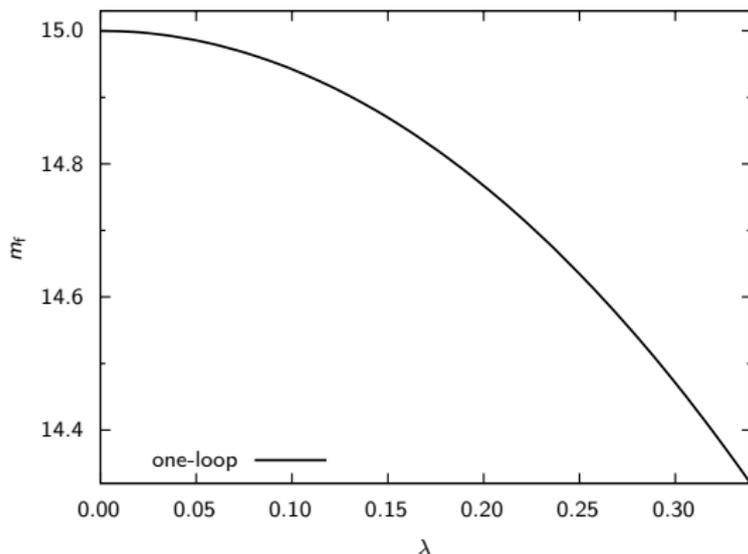


⇒ All formulations yield the same continuum result.

Weak coupling results

Comparing with perturbation theory

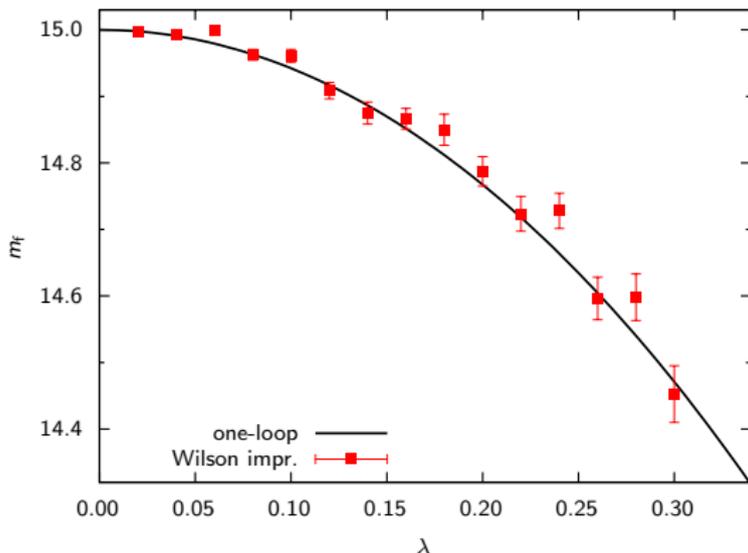
The perturbative one-loop result $m_{\text{ren}}^2 = m^2 \left(1 - \frac{4\lambda^2}{3\sqrt{3}}\right) + \mathcal{O}(\lambda^4)$ can be compared to the continuum extrapolation of the lattice data:



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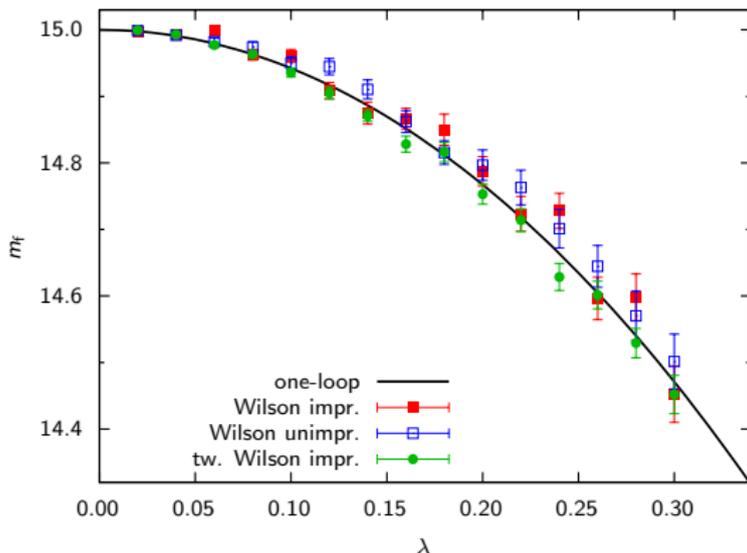
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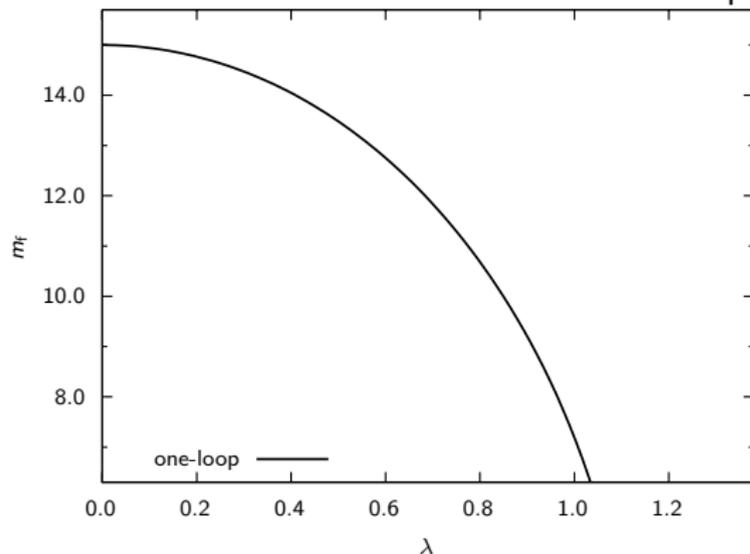


⇒ All different formulations coincide with perturbation theory.

⇒ The supersymmetric continuum limit is reached.

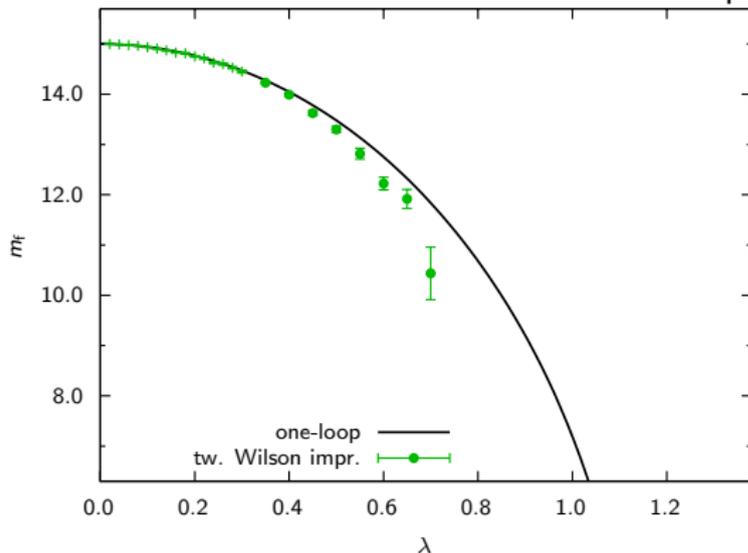
Intermediate coupling results

We checked for the limitations of the one-loop calculation using $\lambda \in [0, 1.2]$:



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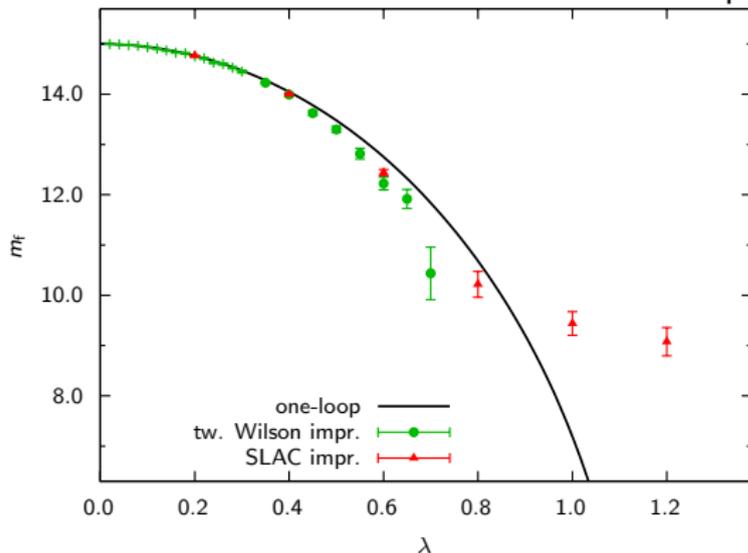
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impr. tw. Wilson (cont.)

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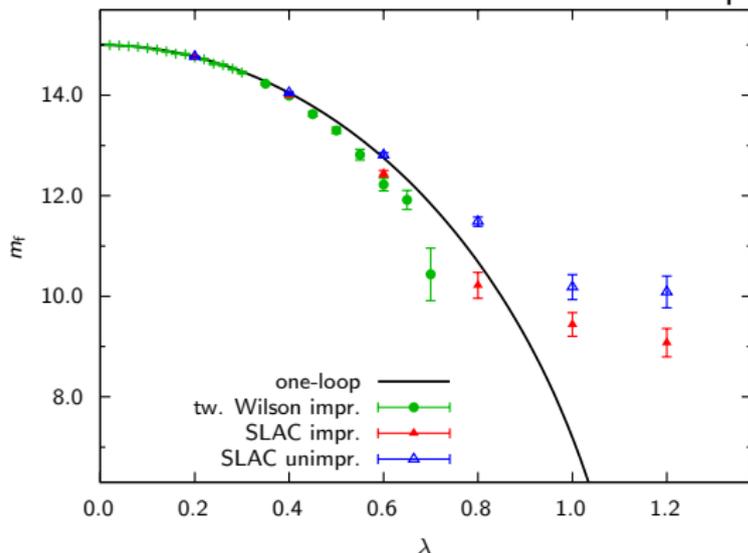
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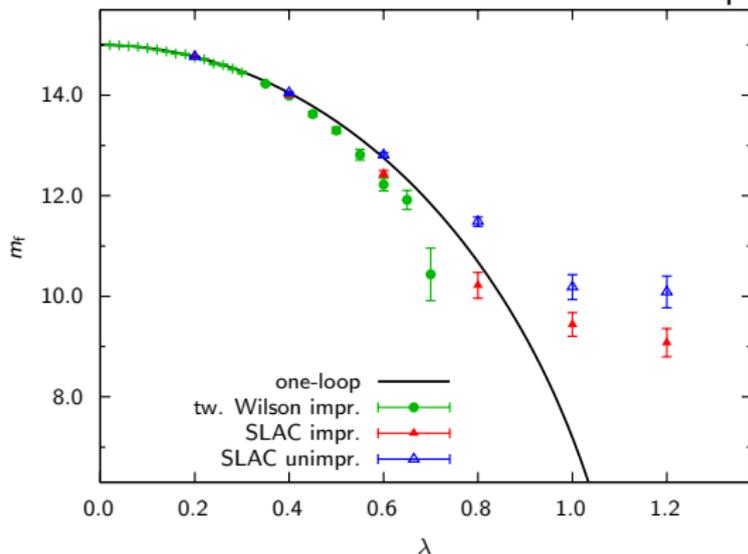
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N_s	improved	unimproved
45	10.22(26)	11.49(9)
63	10.54(15)	10.70(19)

\Rightarrow The correct continuum limit is reached for both models, where the improved SLAC model is closer to the continuum limit.

Results

- With very high statistics bosonic and fermionic masses can **not be distinguished** in the weak to intermediate coupling region for both improved and unimproved formulation.
- For intermediate coupling the improved action is closer to the continuum limit (at least for SLAC fermions).
- The “Nicolai improvement” introduces new problems due to the sampling of unphysical (high-momentum) states. (**no real improvement?**)
- Even without improvement the **correct continuum limit** is reached.

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Outlook

- Further algorithmic improvements (PHMC, multiple r.h.s. solvers) to obtain results for strong coupling ($\lambda > 1.5$).
- Use the elaborate algorithms to explore the $\mathcal{N} = 1$ WZ model in $d = 2$ (SUSY breaking expected).

More on this model including
finite size effects, discussion of the sign problem and technical details
can be found under
[arXiv:0807.1905 \[hep-lat\]](https://arxiv.org/abs/0807.1905)

Thank you!

... and please be cautious when using Nicolai improved actions.