

# Exploring the phase diagram of sextet QCD

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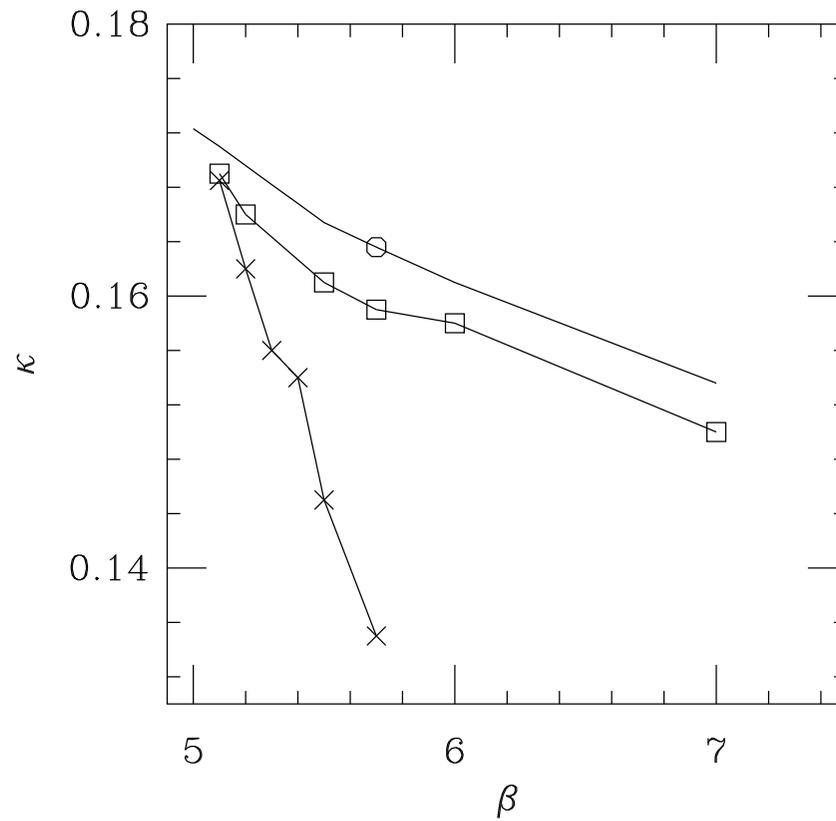
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## Overview

- Follow up to Schrodinger functional study (Svetitsky's talk)
- What we did
- What we saw
  - Near the  $N_t = 8$  deconfinement line
  - Near the  $N_t = 8$  chiral restoration line
  - At small quark mass
- Trying to sum up

Disclaimer – this is exploratory, people have not been doing simulations on this model since 1981!

# The Big Picture



## What we did

- $N_c = 3$ ,  $N_f = 2$ , symmetric rep clover fermions
- Scan parameter space, try to find lines of constant physics
- Use finite size as “marker” for scales (as in SF)
- Bulk of simulations
  - $8^3 \times 12$  and  $12^4$  for spectroscopy, potential
  - $(12 \times 8^2) \times 8$  for chiral and deconfinement crossovers
  - $12^3 \times 8$  for some deconfinement
  - $12^4$  search for scaling
- “P+A” trick
- Trade  $\kappa$ ,  $\kappa_c$  for AWI quark mass

$$\partial_t \langle A_0(t) X(0) \rangle = 2m_q \langle P(t) X(0) \rangle \quad (1)$$

## Connection to conformal physics?

IRFP, if there, would appear as peculiar behavior

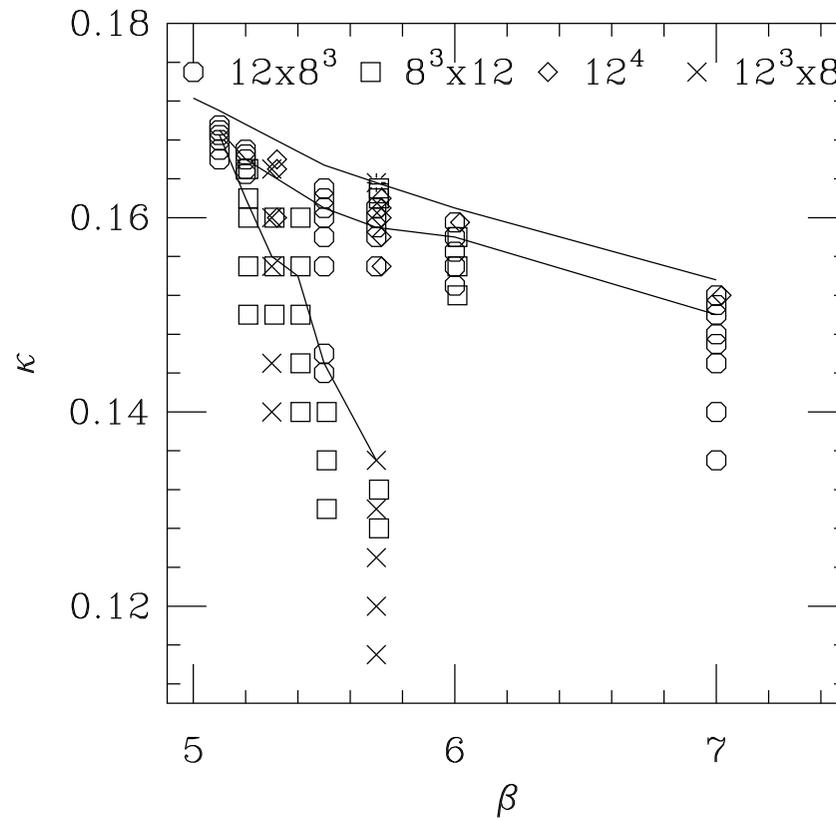
- No confinement, Coulombic potential
- No chiral symmetry breaking,  $f_\pi \rightarrow 0$
- ???

BUT

- Discretization effects break conformal symmetry
- Finite volume breaks conformal symmetry
- Nonzero mass breaks conformal symmetry

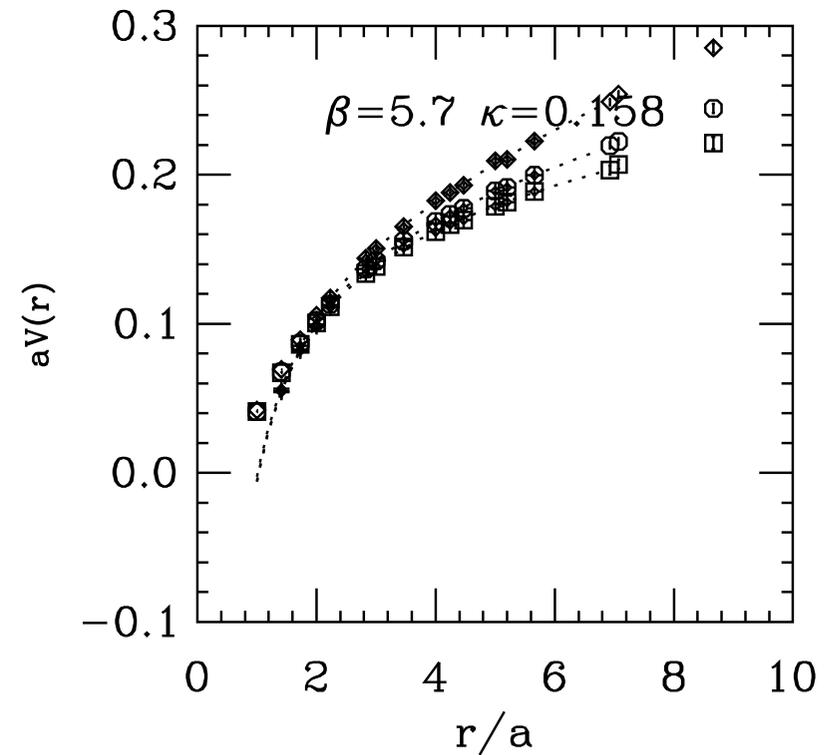
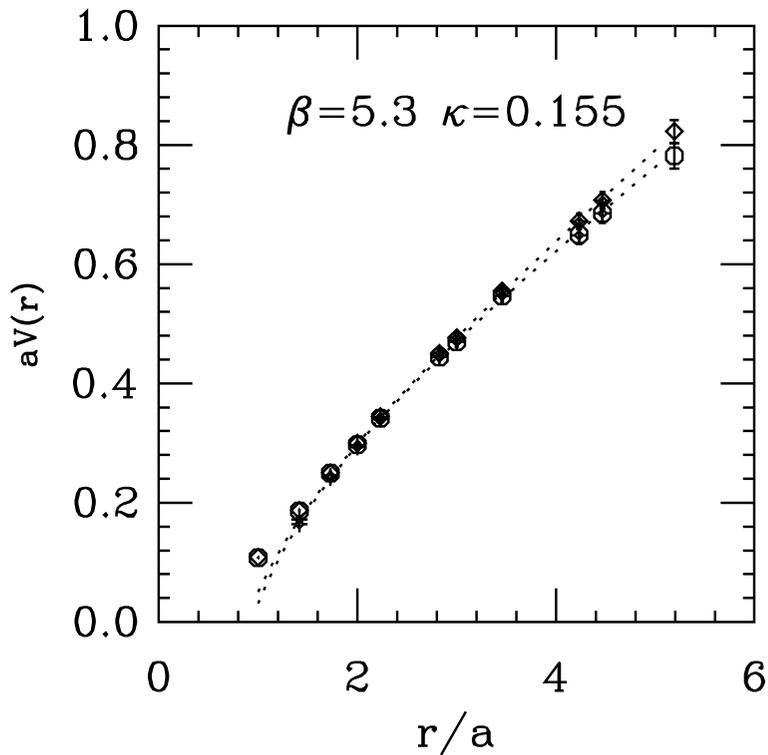
## Where we simulated

- $8^4$ ,  $(12 \times 8^2) \times 8$  as rounded finite temperature ( $N_t = 8$ )
- $8^3 \times 12$  when confinement physics sets in, for  $V(r)$
- $12^3 \times 8$  slightly better finite temperature
- $12^4$  slightly better zero temperature



## Along the $N_t = 8$ deconfinement line

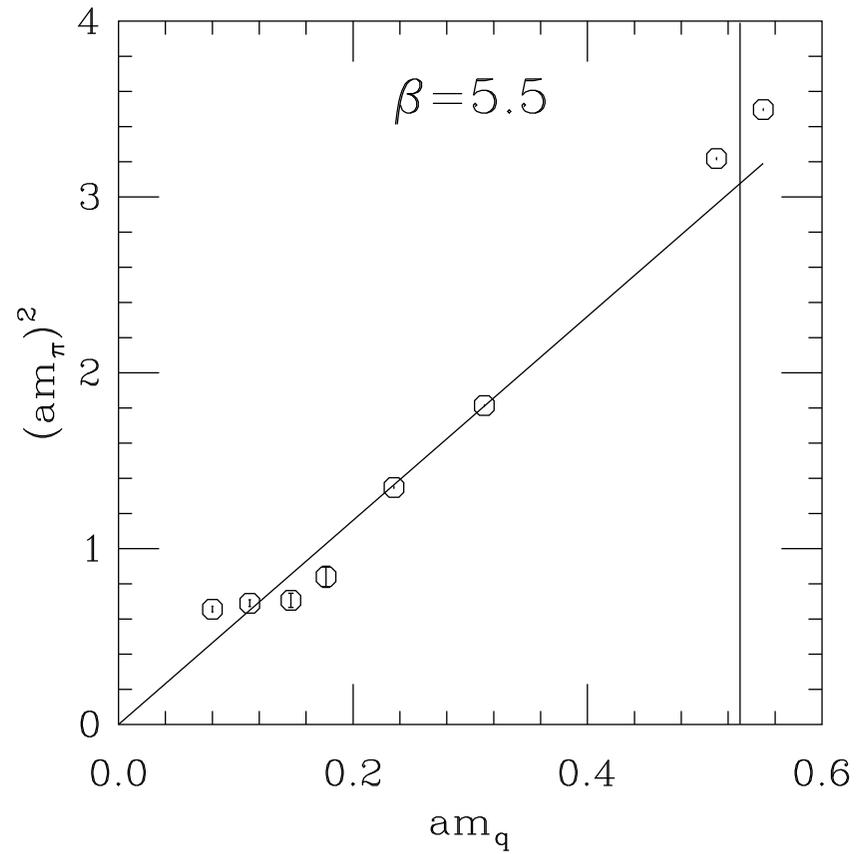
- Observe order-disorder transition in Polyakov loop
- “Conventional”  $V(r)$  measurements show rapid increase in  $r_0/a$ , rapid fall in  $a^2\sigma$  as  $\beta$  rises,  $m_q$  falls



Confinement length (as seen by  $T_c$  and  $V(r)$ ) is squeezed out of the volume

# Scale separation - I

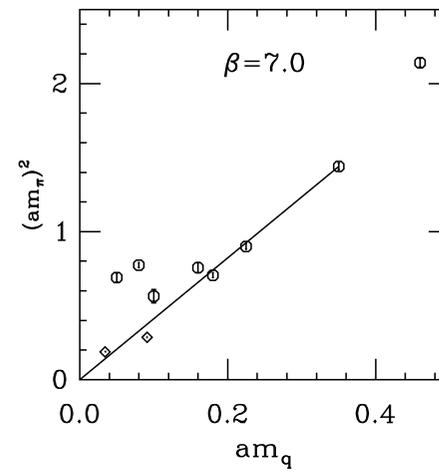
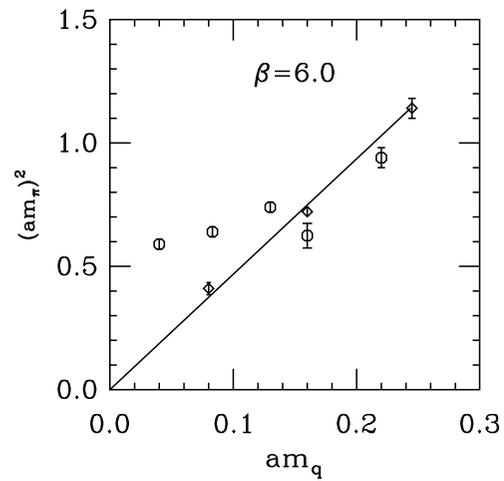
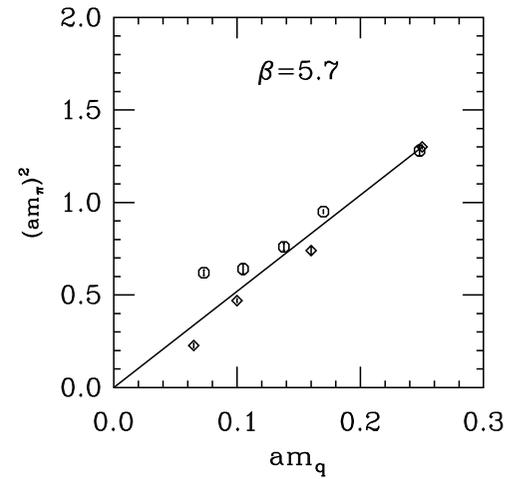
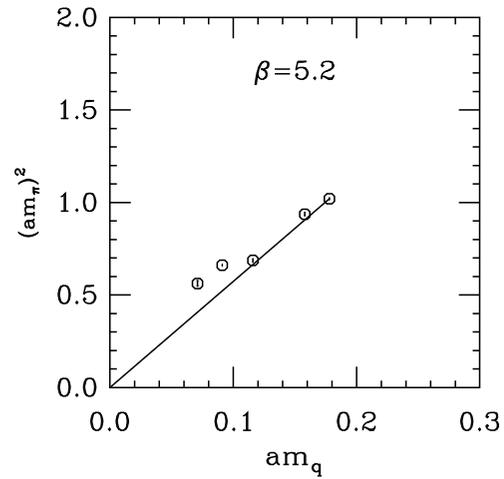
Deconfinement transition has no apparent effect on pion mass (there is a deconfined, chirally broken phase)



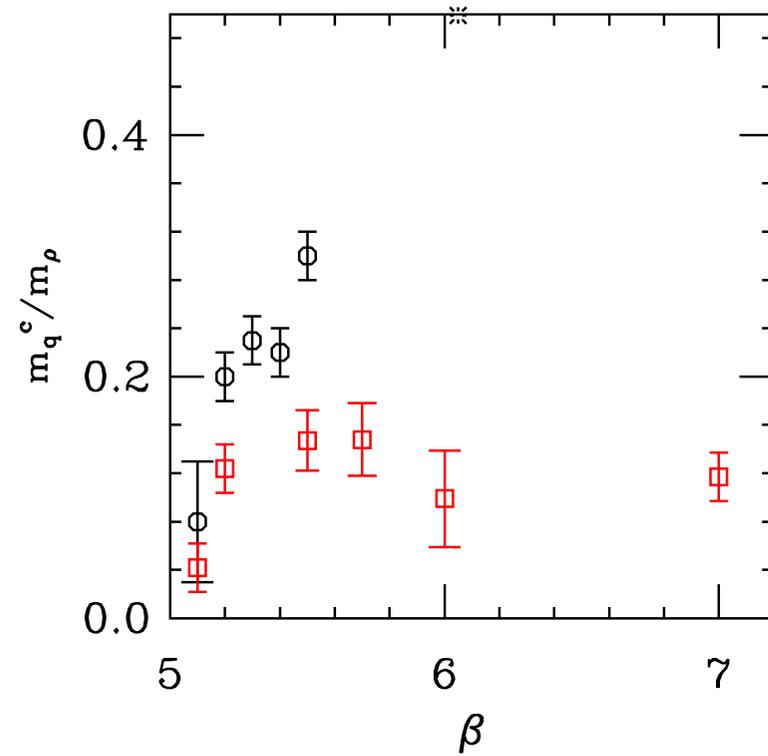
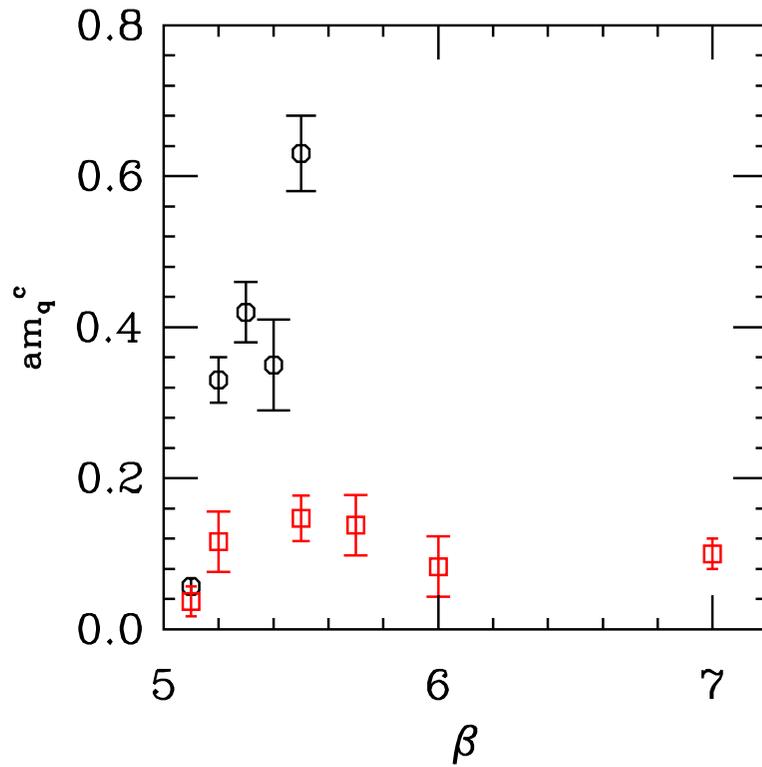
$\beta = 5.5$

## Along the $N_t = 8$ chiral restoration line

With  $N_t = 8$  the fermions see AP b.c.'s at  $m_\pi^2 \sim (2(\pi/8))^2$ :  $m_\pi^2$  no longer  $\sim m_q$



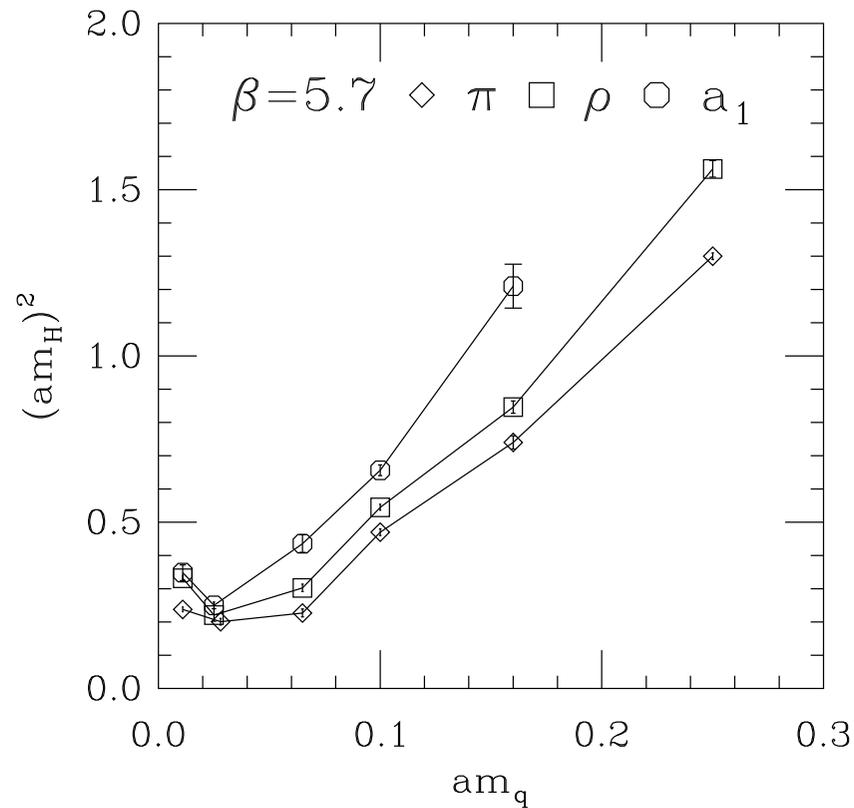
# The two transitions



Octagons –  $N_t = 8$  deconfinement, squares –  $N_t = 8$  chiral restoration

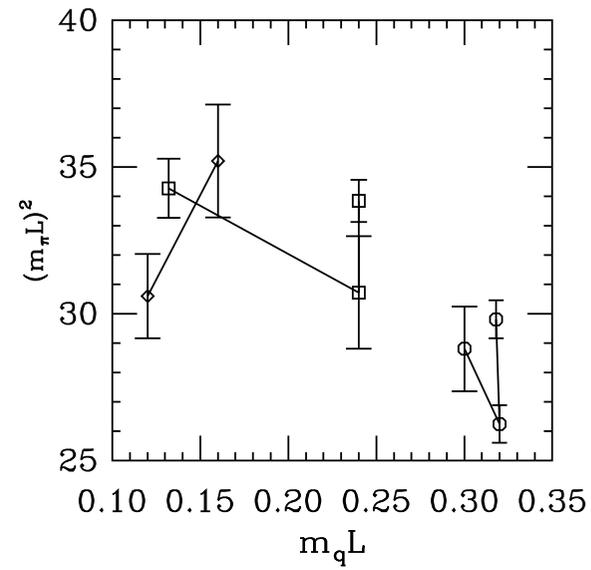
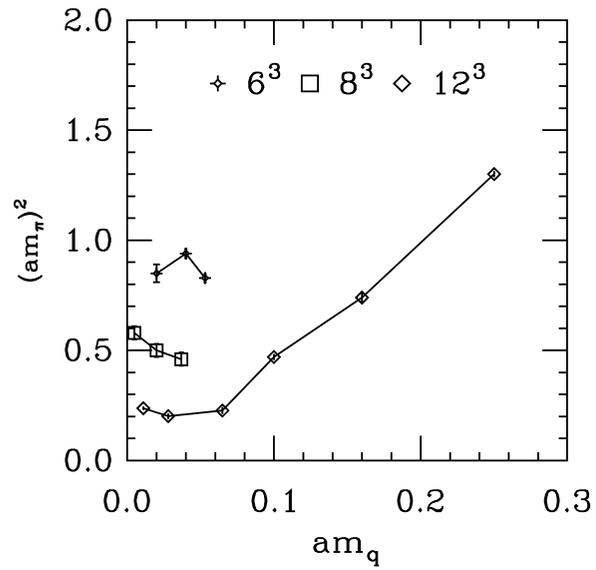
## In the deconfined, chirally broken phase, life is strange

- $m_\pi^2/m_q \sim \text{constant}$
- Never get small  $m_\pi/m_\rho$
- $m_{a_1}/m_\rho$  quite close to 1
- $f_\pi$  quite  $m_q$  - dependent, drops a lot as  $m_q \rightarrow 0$ , can't tell if it vanishes at  $am_q = 0$
- $am_\rho(\beta, am_q) = f(am_q)$ , nearly  $\beta$  independent for  $5.2 < \beta < 7$



# Very small quark mass

At even smaller quark masses, finite volume is everything,  $mL \sim \text{constant}$

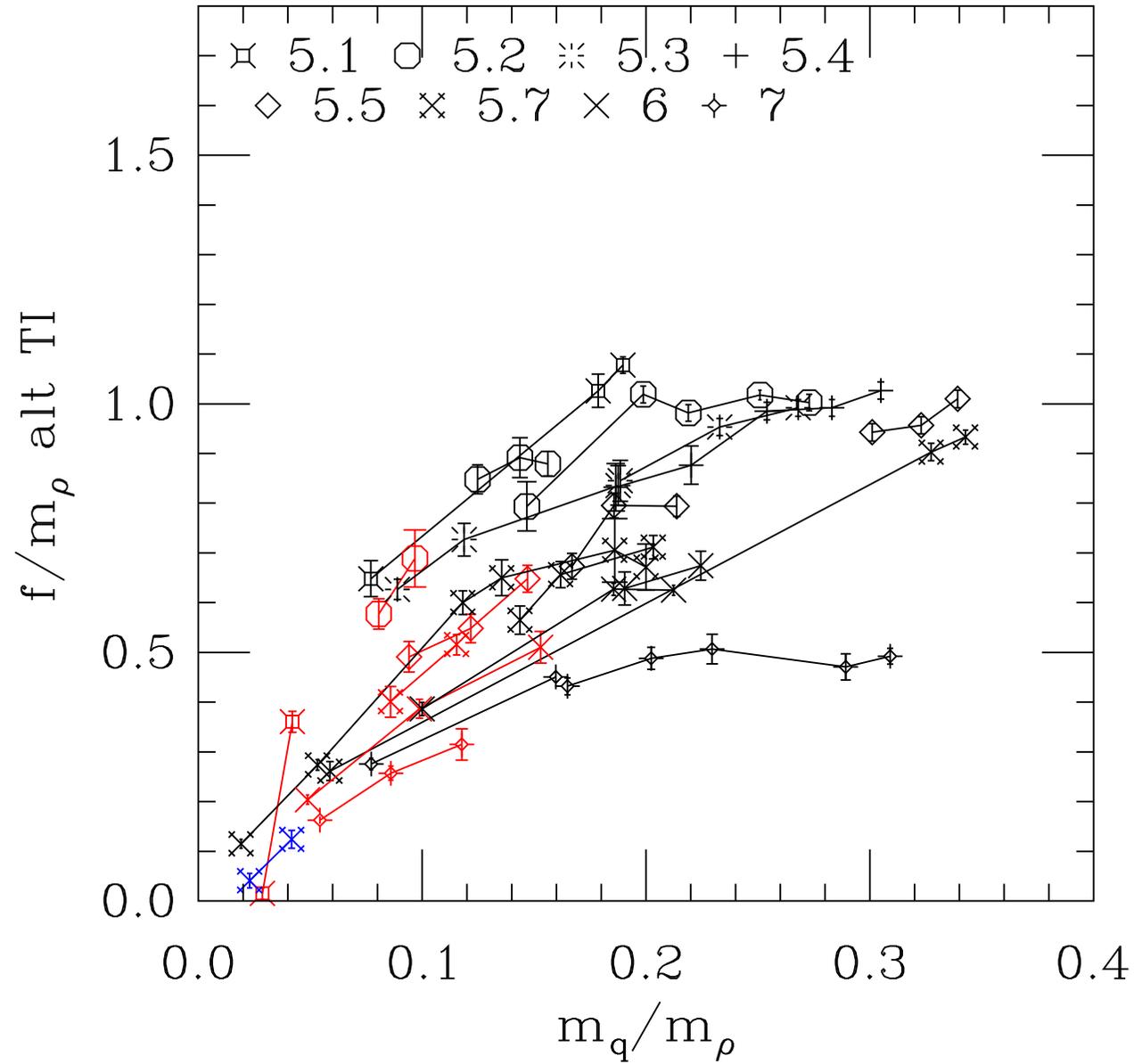


## $f_\pi$ - A crucial diagnostic?

- In technicolor models  $f_\pi = \text{Higgs VEV (246 GeV)}$  sets the scale
- $f_\pi/m_\rho \rightarrow m_\rho$  - New physics, the technirho mass
- In large  $N_c$ , sextet  $f_\pi \sim N_c$ , vs fundamental rep  $f_\pi \sim \sqrt{N_c}$
- With an IRFP, no chiral symmetry breaking,  $f_\pi \rightarrow 0$  as  $m_q \rightarrow 0$

But, simulational dirt can mar the result

- Lattice to continuum conversion
  - $a f_\pi^L = Z_A \langle 0 | A_0(x) | \pi \rangle$
  - $Z_A = (1 + \frac{g^2}{16\pi^2} C(R)W) (1 - \frac{6\kappa}{8\kappa_C})$
- Finite volume effectively restores chiral symmetry
  - $m_\pi L < 1, f_\pi L > 1$  - epsilon regime
  - $f_\pi L < 1$  - hell regime



$f_\pi/m_\rho$  falls to zero as  $m_q$  goes to zero, even if volume too big for chiral transition

## Conclusions

$N_f = 2$  sextet QCD is not like low- $N_f$  fundamental rep QCD!

- Scale separation:  $T_{th} \ll T_{ch}$
- Funny “chirally broken” phase
- Still to do: can “slow running” of  $g^2$  be seen in  $aM$ ?
- Direct search for IRFP in spectroscopy seems difficult, as hard as seeing asymptotic freedom in spectroscopy was in QCD
- But never mind! Discoveries await, even though existing language may be inadequate...

## “P+A” trick

- Add quark propagators with antiperiodic (A) and periodic (P) temporal b.c.'s
- $C(t) = Z(\exp(-mt) \pm \exp(-m(2N_t - t)))$

