

Structure and reactions of Θ^+

Atsushi Hosaka (RCNP, Osaka Univ)

- [hep-ph/0507105](#)
E. Hiyama, Kamimura, Yahiro, A.H., and Toki
- [hep-ph/0505134](#), [hep-ph/0503149](#) to appear PRD
S.I. Nam, A. *Hosaka* and H.-Ch. *Kim*

Outline

1. Full 5-body calculation

2. Photoproduction reconsidered

pn asymmetry when $J = 3/2$

$\gamma n \rightarrow K^- \Theta^+$ vs. $\gamma p \rightarrow \bar{K}^0 \Theta^+$ [Θ^+ or $\Lambda(1520)$]

K^* production

1. Full 5-body calculation

Hiyama-Kamimura-Yahiro-Hosaka-Toki

hep-ph/0507105

Need to handle the (at least) 5-body system

- So far calculations were only approximate and only for bound state
- Better method with scattering states included

Method available developed in nuclear-physics

Assumption: NRQM

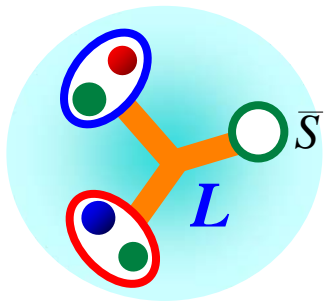
- Validity of the use of the Schrodinger picture
- What is the effective hamiltonian
e.g. *const* is not known => *confinement?*

Clarify what this hamiltonian tells for 5-body system
Then improve this method or choose others?

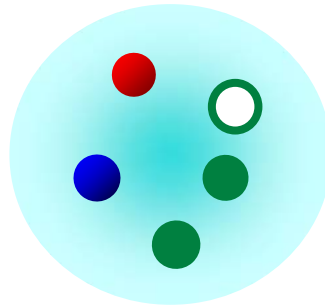
Decay (fall-apart) is very sensitive to WFs

Hadronic (color-singlet) or colored correlations?

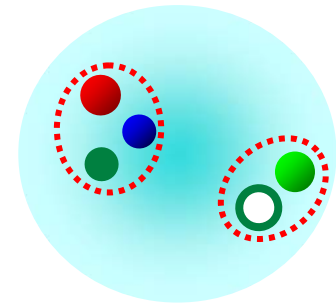
SU(3) qqq or $qq\bar{q}$ are enough to make color singlets



Diquark correl.



Five quarks
random



Hadronic correl.

Difficult to decay

Easy to decay



Dependence on J^P

$J^P = 1/2^-$: $l = 0$ (ground state)

~ KN scattering => can not be narrow

Excited or complicated state may be a narrow res.

$J^P = 1/2^+$: $l = 1$

($3/2^+$) Depends much on the configuration

$J^P = 3/2^-$: $l = 0$

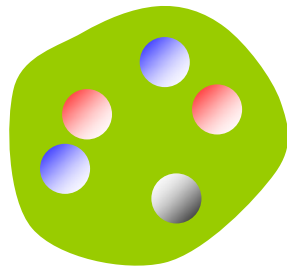
D-wave KN decay is forbidden, can be narrow

Seems consistent with phenomenology

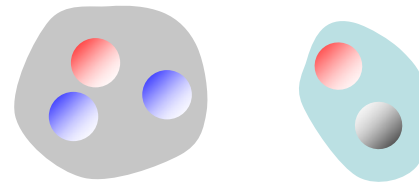
=> Hyodo, PRD, [hep-ph/0509104](#)

[hys.Rev.D71:054017,2005](#) , [hep-ph/0502093](#)

Method **Most serious calculation** for **5-body system**
 with scattering states included
 Gaussian expansion method

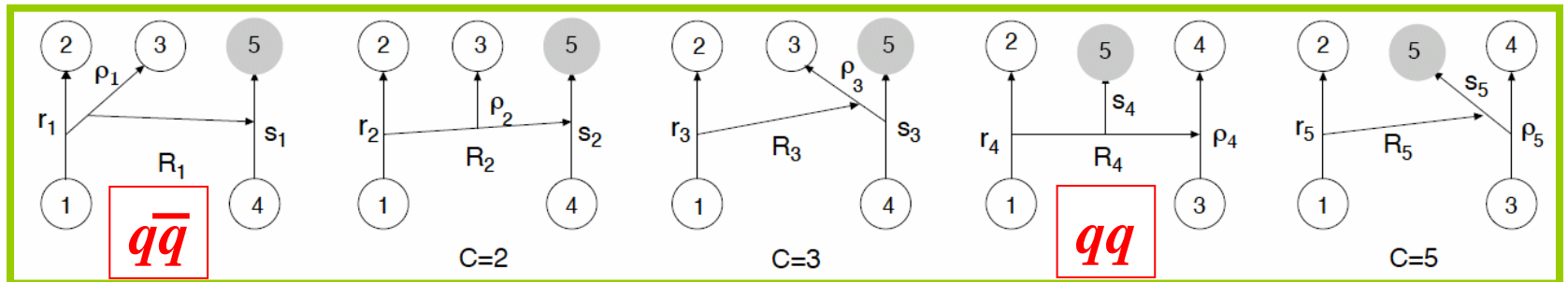


Θ^+ -confined



NK-scattering

Compute phase shifts



Hamiltonian

NR quark model of Isgur-Karl

$$H = \sum_i \left(m_i + \frac{\mathbf{P}_i^2}{2m_i} \right) - T_G + V_{\text{Conf}} + V_{\text{CM}}$$

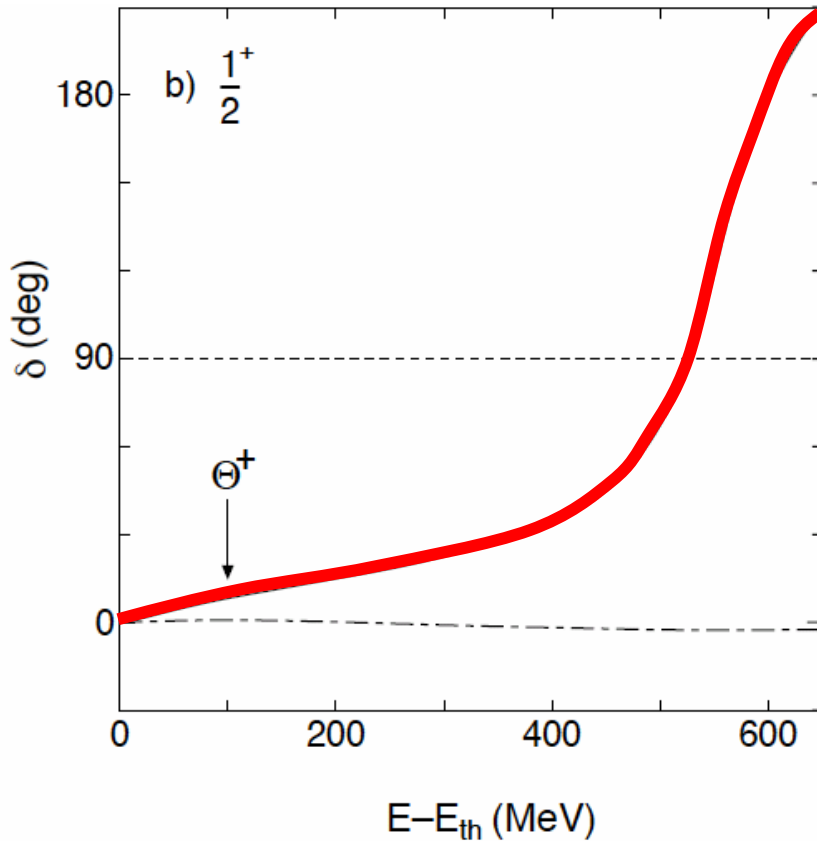
$$V_{\text{Conf}} = - \sum_{i < j} \sum_{\alpha=1}^8 \frac{\lambda_i^\alpha}{2} \frac{\lambda_j^\alpha}{2} \left[\frac{k}{2} (\mathbf{x}_i - \mathbf{x}_j)^2 + v_0 \right]$$

$$V_{\text{CM}} = \sum_{i < j} \sum_{\alpha=1}^8 \frac{\lambda_i^\alpha}{2} \frac{\lambda_j^\alpha}{2} \frac{\xi_\sigma}{m_i m_j} e^{-(\mathbf{x}_i - \mathbf{x}_j)^2 / \beta^2} \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$$

Good for conventional baryons

	Mass		Magnetic moments		Charge radii	
	Cal.	Exp.	Cal.	Exp.	Cal.	Exp.
	(MeV)		(nm)		(fm ²)	
p	939	939	2.7737	2.7828	(0.60) ²	(0.87) ²
n	939	939	-1.826	-1.913	-0.04	-0.12
Λ	1058	1115	-0.602	-0.613	-0.01	-
Σ^+	1119	1189	2.691	2.458	0.35	-
Σ^0	1119	1192	0.819	-	0.03	-
Σ^-	1119	1197	-1.054	-1.160	-0.30	-
Ξ^0	1309	1314	-1.414	-1.250	-0.01	-
Ξ^-	1309	1321	-0.507	-0.651	-0.28	-
Δ^Q	1232	1232	2.843 Q	-	0.41 Q	-
Σ^{*+}	1320	1384	3.18	-	0.64	-
Σ^{*0}	1320	1384	0.33	-	0.12	-
Σ^{*-}	1320	1384	-2.51	-	-0.38	-
Ξ^{*0}	1512	1533	0.67	-	0.03	-
Ξ^{*-}	1512	1533	-2.17	-	-0.35	-
Ω	1506	1672	-1.840	-2.02	-0.32	-

KN-phase shifts $1/2^+$



$E_{res} \sim 530 \text{ MeV}$
 $\Gamma_{res} \sim 110 \text{ MeV}$

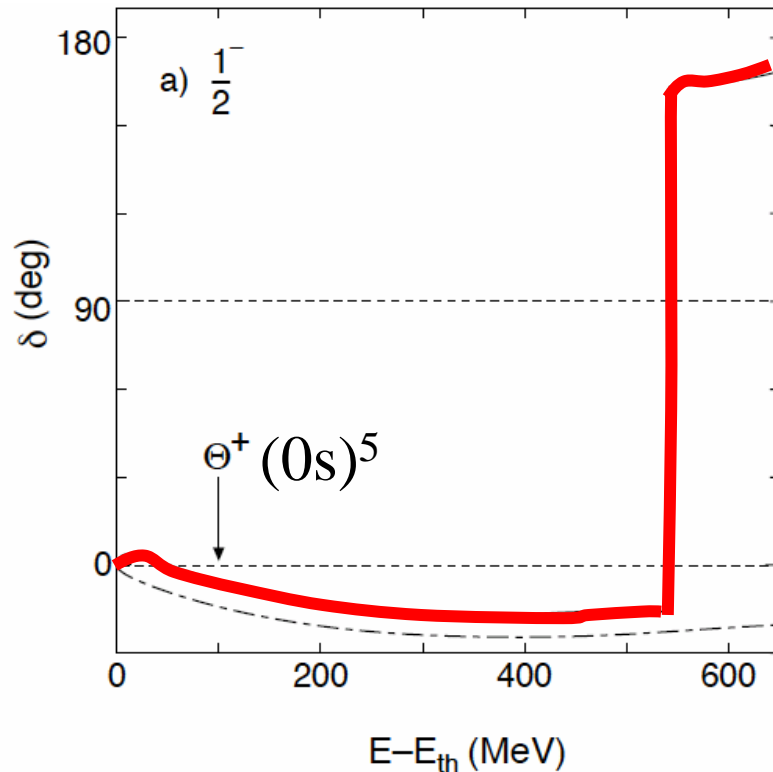
$(0s)^4 1p$

- Strong $q\bar{s}$ correlation favors KN
Rather than $[ud][ud]\bar{s}$

\Rightarrow

- Formation of the JW type conf.
is a dynamical problem

KN-phase shifts $1/2^-$



$(0s)^5$: KN scattering state
Likely to be $1s(0s)^4$

$$\Gamma < 1 \text{ MeV}$$

The nature of
the narrow resonance
is interesting to analyze

We have seen:

- **5-body calculation** of the Isgur-Karl quark model
Two states at ~ 500 MeV above the KN threshold
 $\Gamma(1/2^-) \sim$ Very narrow, < 1 MeV
 $\Gamma(1/2^+) \sim 100$ MeV

When the same *const*-parameter is used
as for conventional baryons

- The ground $(0s)^5$ configuration melts into the continuum
- The $1/2^+$ state is dominated by *qqq-qq* configuration

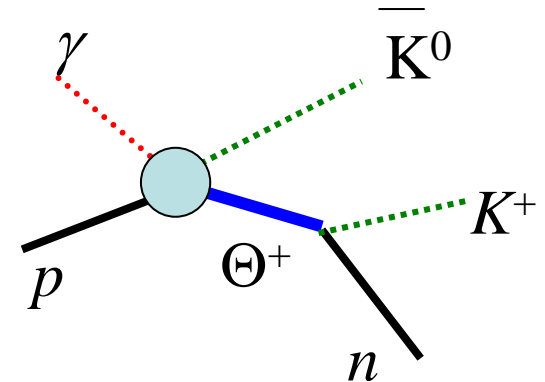
2. Photoproductions

1. K-production
2. K*-production

(1) **K-production** with new J-Lab data

$$\gamma p \rightarrow n K^+ K^0$$

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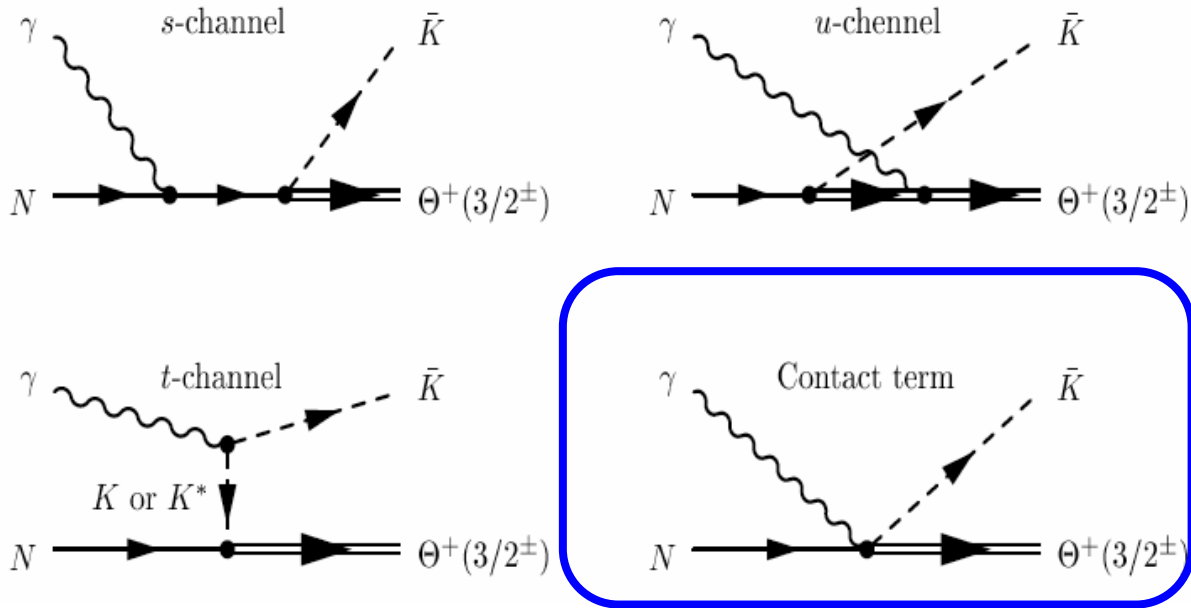


Teken from DeVita's
talk at spring APS meeting

This is serious, but leads immediately to the absence of Θ^+ ?

Effective Lagrangian approach

hep-ph/0505134 Nam-Hosaka-Kim



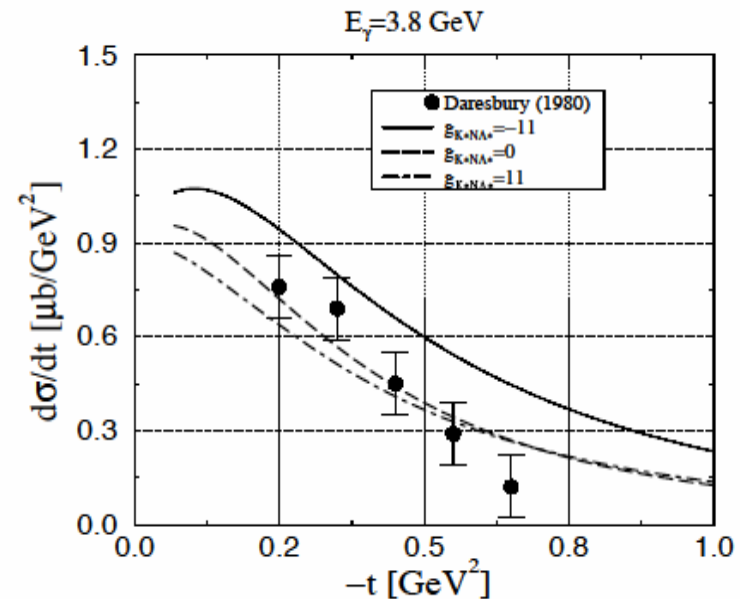
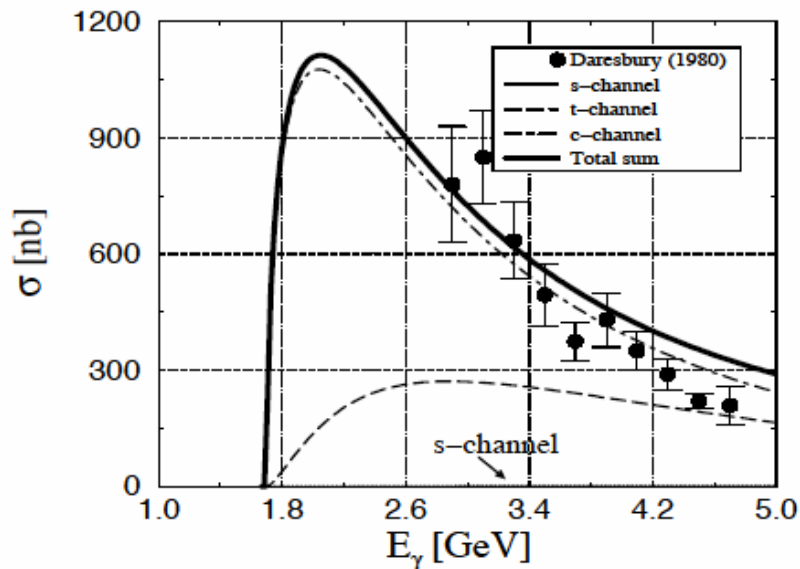
appears only for n-target

For $J = 3/2$, only **PV** scheme is possible

Before the Θ -production

$\gamma n \rightarrow K^- \Lambda(1520)$ and $\gamma p \rightarrow \bar{K}^0 \Lambda(1520)$
 was studied and large **pn** asymmetry was found

Nam-Hosaka-Kim, hep-ph/0503149 to appear PRD



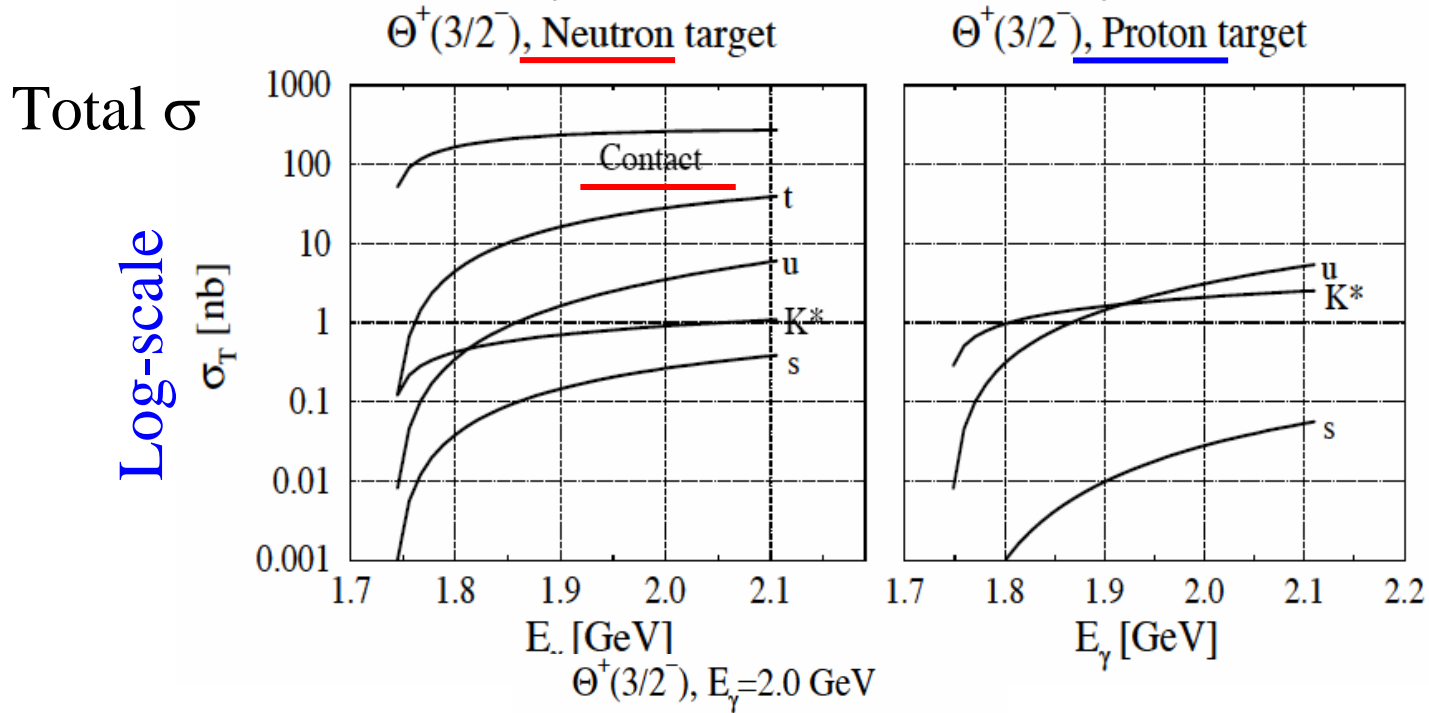
Comparison

Form factor	F_1	
Reactions	<u>$\gamma p \rightarrow K^+ \Lambda^*$</u>	<u>$\gamma n \rightarrow K^0 \Lambda^*$</u>
σ	$\sim 900 \text{ nb}$	$\sim 30 \text{ nb}$
$d\sigma/d(\cos \theta)$	Forward peak	Peak at $\sim 45^\circ$
$d\sigma/dt$	Good	No data

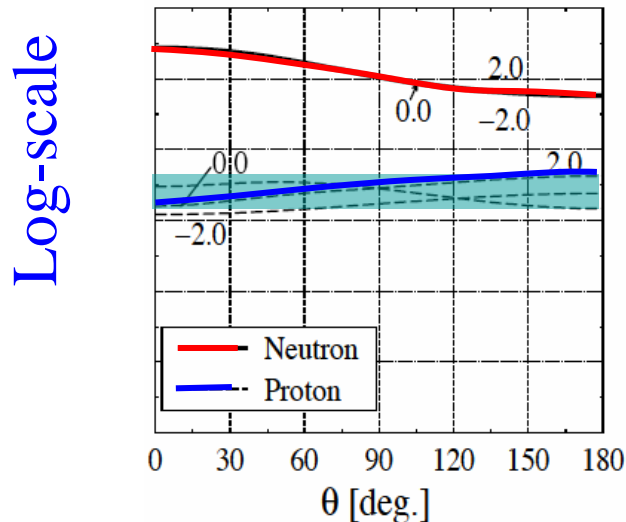
The presence (for p) or absence (for n) contact term is important

LEPS data seems to support this result

Theta production, $J^P = 3/2^-$



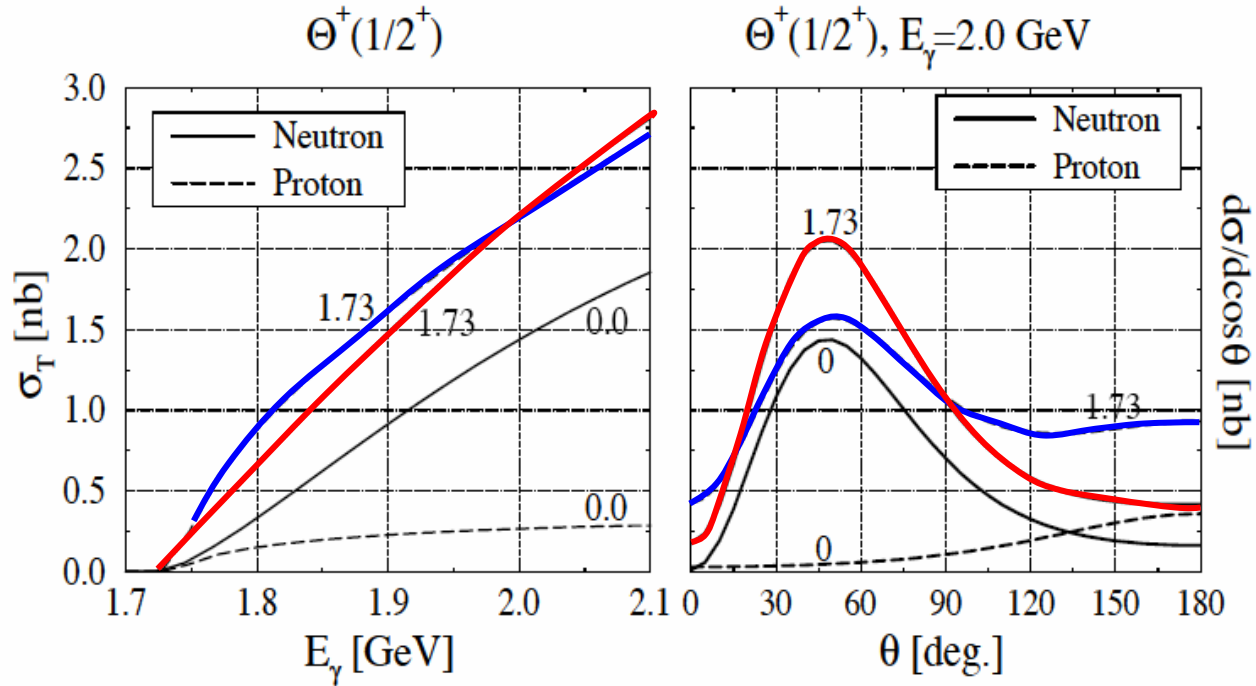
Angular dist



neutron ~ forward peak
Contact term

proton ~ rather flat

$$J^P = 1/2^+$$



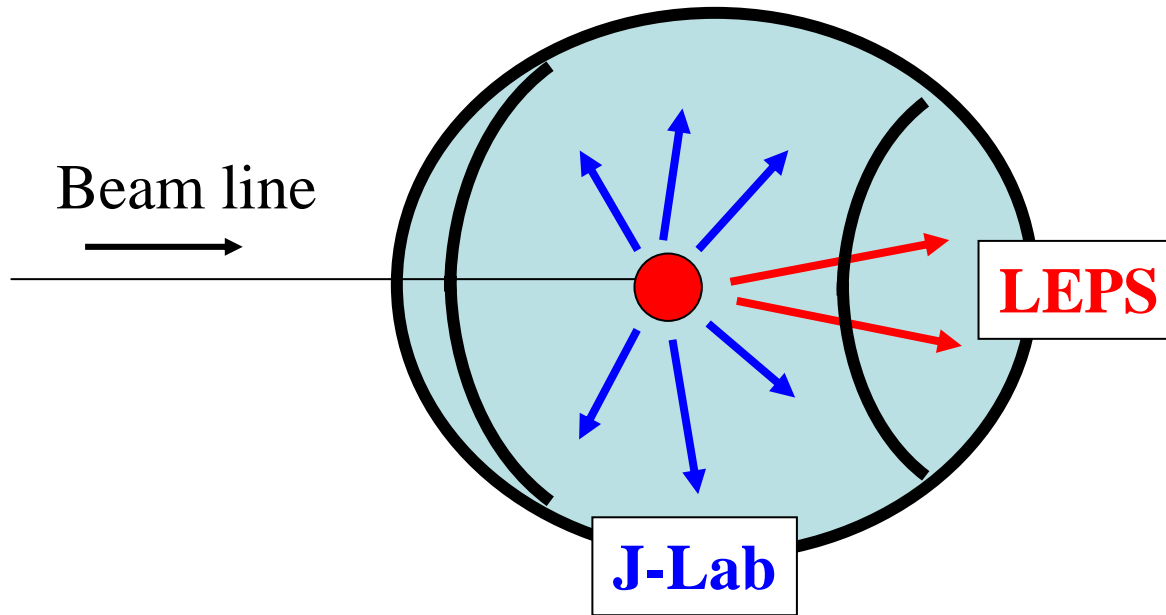
The role of the contact term is more important for $J^P = 3/2^-$ than $1/2^+$

Predictions

J^P	$3/2^+$		$3/2^-$		$1/2^+$	
$g_{KN\Theta}$	0.53		4.22		1.0	
$g_{K^*N\Theta}$	± 0.91		± 2		± 1.73	
Target	n	p	n	p	n	p
σ	~ 25 nb	~ 1 nb	~ 200 nb	~ 4 nb	~ 1 nb	~ 1 nb
$\frac{d\sigma}{d\cos\theta}$	Forward	$\sim 60^\circ$	Forward	–	$\sim 45^\circ$	$\sim 45^\circ$

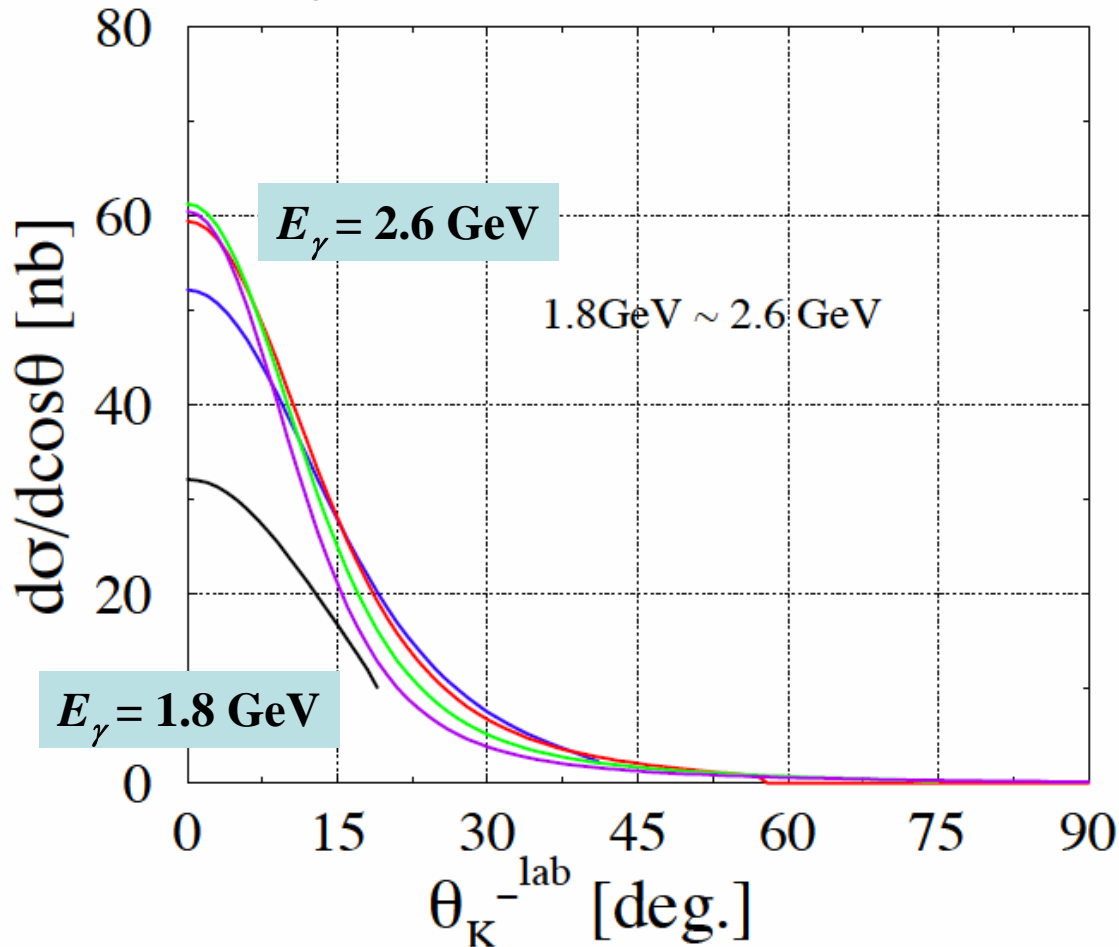
- We see a large asymmetry between pn targets
- Cross section for proton \sim few nb is consistent with the upper limit estimated by CLAS

Different exp. config.



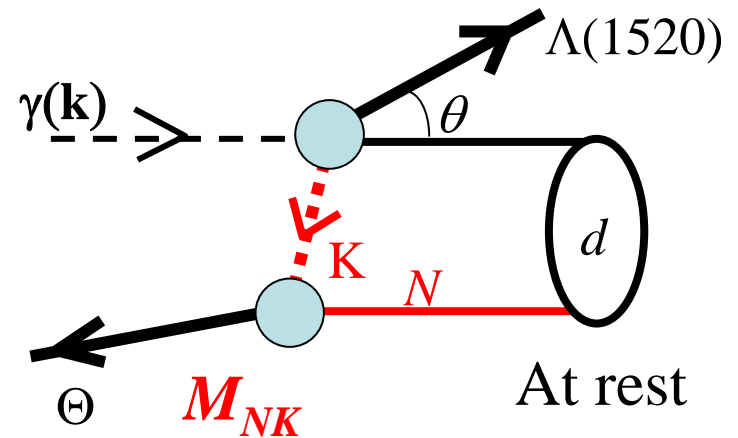
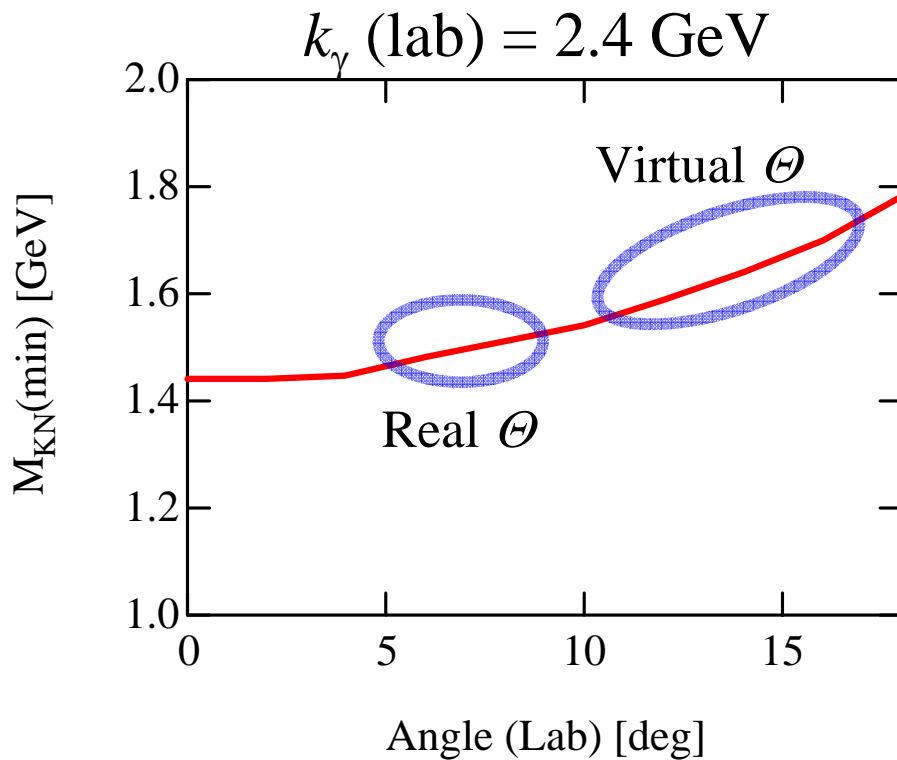
Angular dist. in lab frame

$\gamma n \rightarrow K^- \Theta^+(3/2^+)$ in the lab frame



At $E_\gamma = 2$ GeV
 $\sigma(\theta < 25^\circ) \sim 5\sigma(\theta > 25^\circ)$

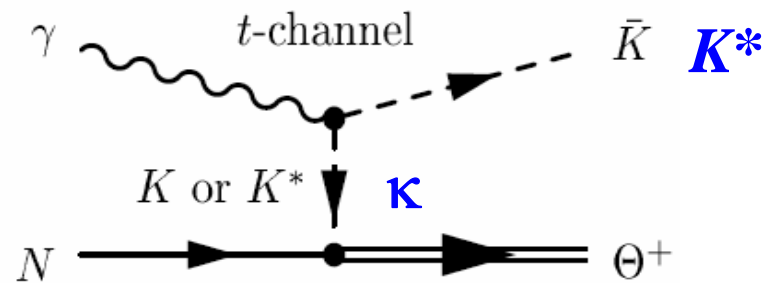
Special kinematics



(2) $K^* (1^-)$ production

- Physics in the t-channel

Now $\kappa (0^-)$ is allowed to be exchanged



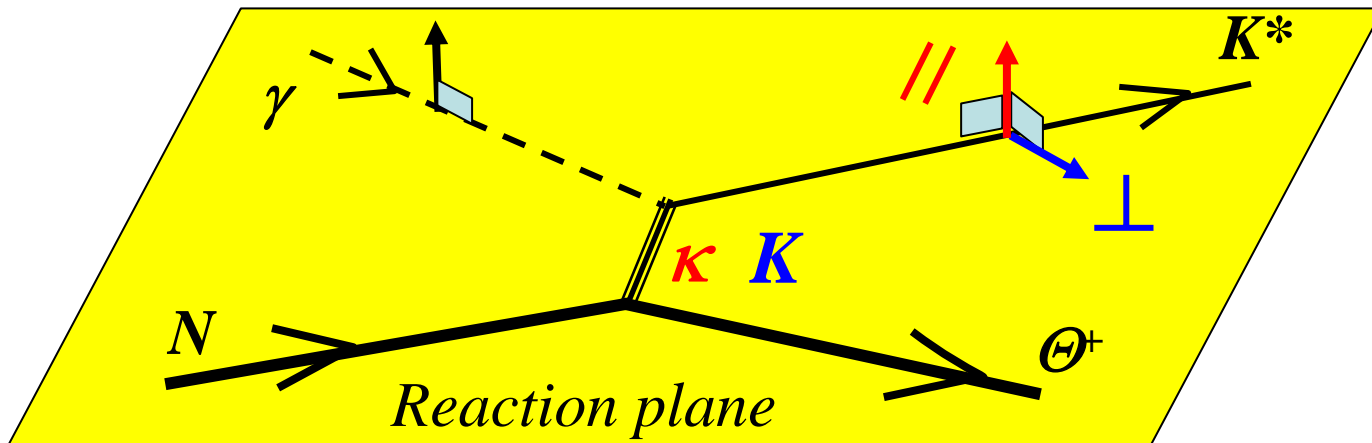
Exotic tetraquark κ may couple strongly to Θ^+

D.P. Roy, J. Phys. G30, R113 (2004)

- Using polarizations of γ and K^* , we can distinguish the exchanged particles

Polarizations as a particle filter

Pol. of γ perp. to react. plane



If parallel [$//$], only κ is exchanged

If perpendicular [\perp], only K is exchanged

Summary

- **5-body calc.**

$1/2^-$ $E \sim 2 \text{ GeV}$, $\Gamma \sim 1 \text{ MeV}$

$1/2^+$ $E \sim 2 \text{ GeV}$, $\Gamma \sim 100 \text{ MeV}$ Configurations mix

(cf: quark model calc. Hosaka-Oka-Shinozaki

[hep-ph/0409102](#), PRD71: 074021 (2005))

- **Photoproduction, revised**

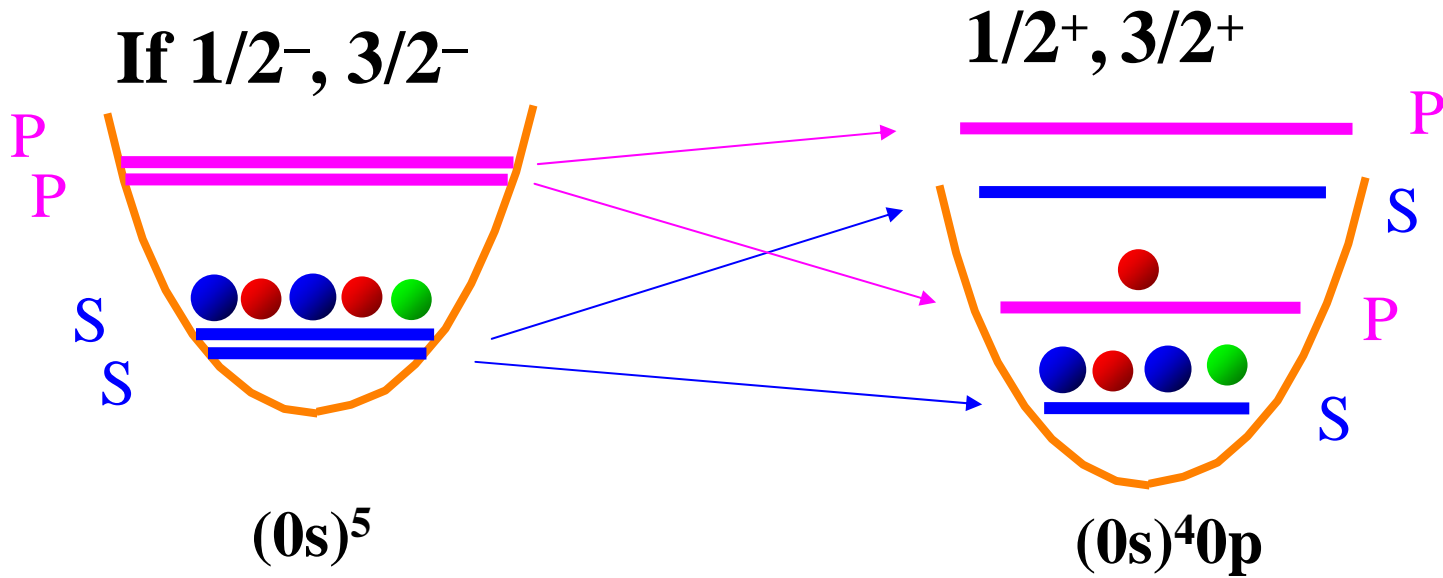
*There is a large *pn asymmetry*, especially for $J = 3/2$

*No signal from the CLAS does not lead immediately to the absence of Θ^+

**Kinematics* at LEPS is very interesting

**K* can be used as a *particle (t-channel) filter*

Interpretation of results



Diquark }
 Tri-quark } correlations
 Chiral }