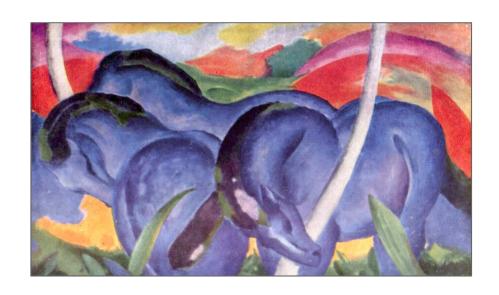
THREE-BODY INTERACTION IN ISOBAR FORMALISM



Maxim Mai The George Washington University



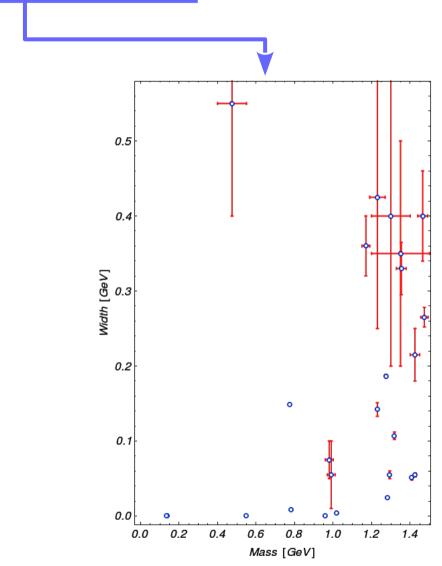
INTRODUCTION

→ mass generation & confinement

Non-perturbative dynamics

→ mass generation & confinement

→ rich spectrum of excited states



→ mass generation & confinement

Non-perturbative dynamics

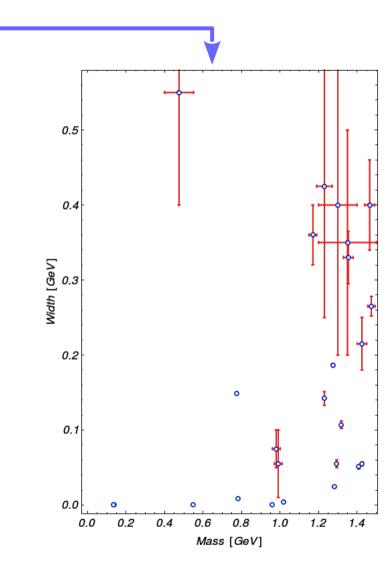
→ rich spectrum of excited states

Q1: how many are there?

→ missing resonance problem

Q2: what are they?

- quark-antiquark
- gluons
- hadron-hardon dynamics



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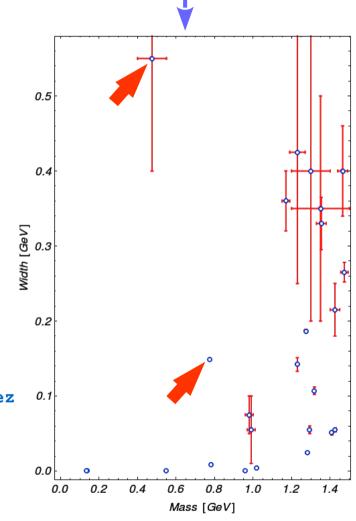
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EXAMPLES:

 $-f0(500), \varrho(770), \kappa(800)...$ couple dominantly to $2\pi, K\pi,...$ \rightarrow talk by Jose R. Pelaez



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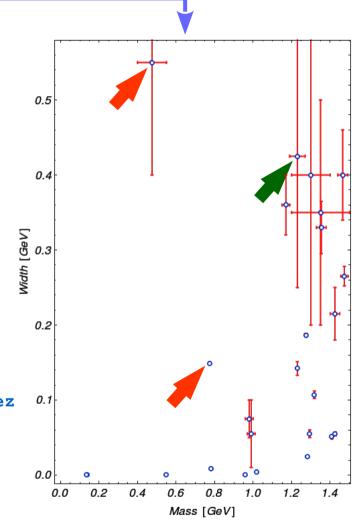
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EXAMPLES:

 $-f0(500), \varrho(770), \kappa(800)...$ couple dominantly to $2\pi, K\pi, \dots$ $\rightarrow \text{talk by Jose R. Pelaez}$

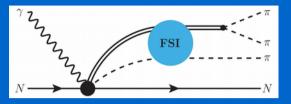
-a1(1260), K1(1400)couple dominantly to 3π , $K\pi\pi$,...



Experiment

- Search for QCD exotics @ GlueX

* a1(1260)

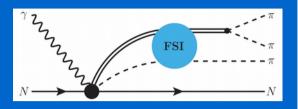


- KL Beam @ GlueX
 - * $K^*(892)$ signature in $KN \rightarrow K\pi N$
 - * $K\pi\pi$ channels(?)
- Further applications:
 - * Roper puzzle $(\pi\pi N)$
 - * X(3872)

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Lattice QCD

Ab-initio numerical calculations

- Euclidean ST
- finite lattice spacing
- finite volume effects
 - → 2-body Quantization Condition

```
[Lüscher (1986)]
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- → talk by Morningstar
- → 3-body **QC** not **yet** established

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[Rusetsky, Polejaeva,
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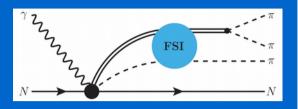
Davoudi, Guo

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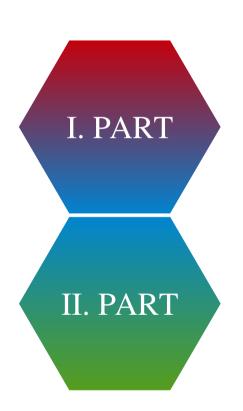
[Rusetsky, Polejaeva, Sharpe, Hansen, Briceno, Davoudi, Guo MM, Doring, ...]

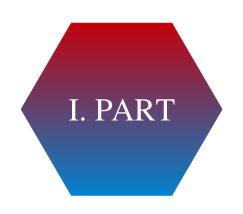
THIS TALK: 3-BODY SCATTERING AMPLITUDE IN ISOBAR-FORMULATION

UNITARITY OF S-MATRIX

IMAGINARY PARTS (INF. VOL.)

POWER LAW FIN. VOL. EFFECTS

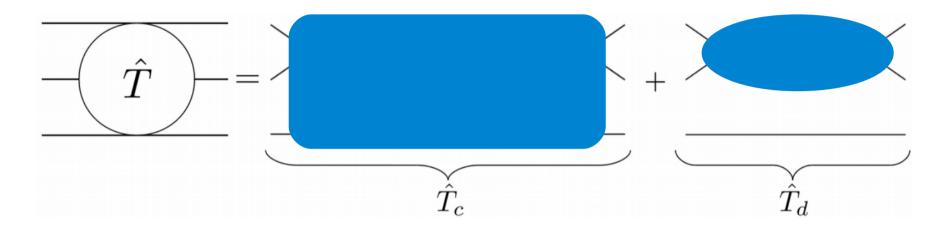




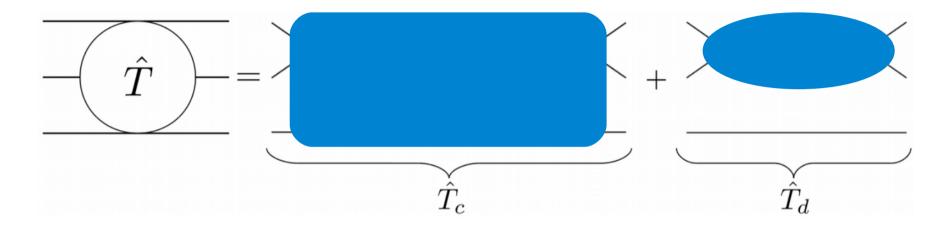
3→3 SCATTERING AMPLITUDE IN INFINITE VOLUME

• 3 asymptotic states (scalar particles of equal mass (m))

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- Connectedness structure of matrix elements: (all permutations considered)



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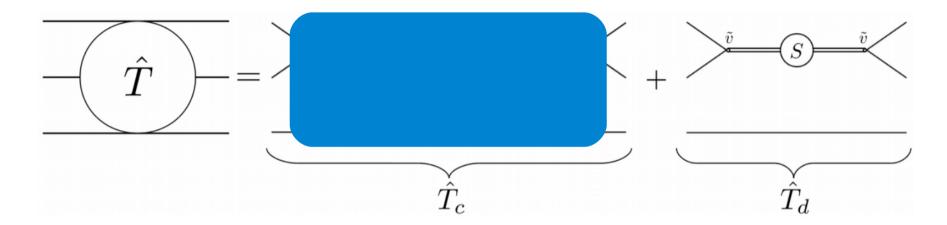


• isobar-parametrization of two-body amplitude

[Bedaque, Griesshammer (1999)

 \rightarrow "isobars" $\sim S(M_{inv})$ for definite QN & correct r.h.-singularities

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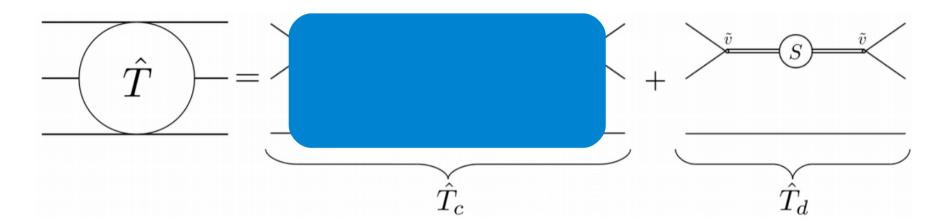


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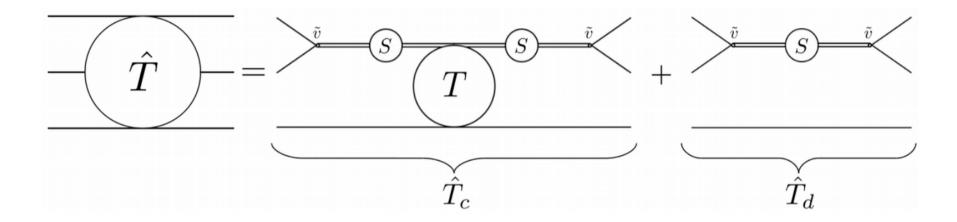
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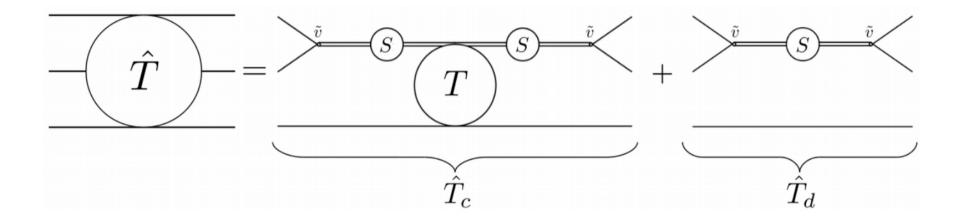
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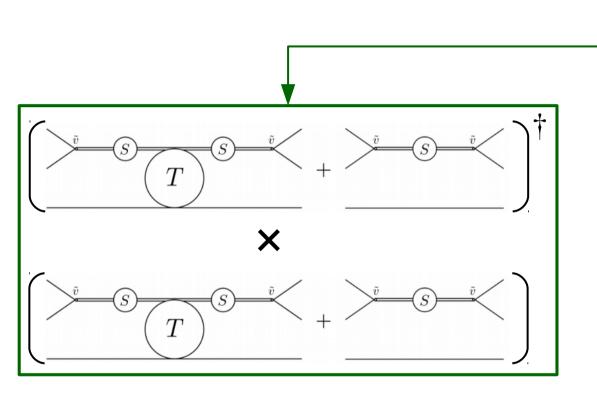


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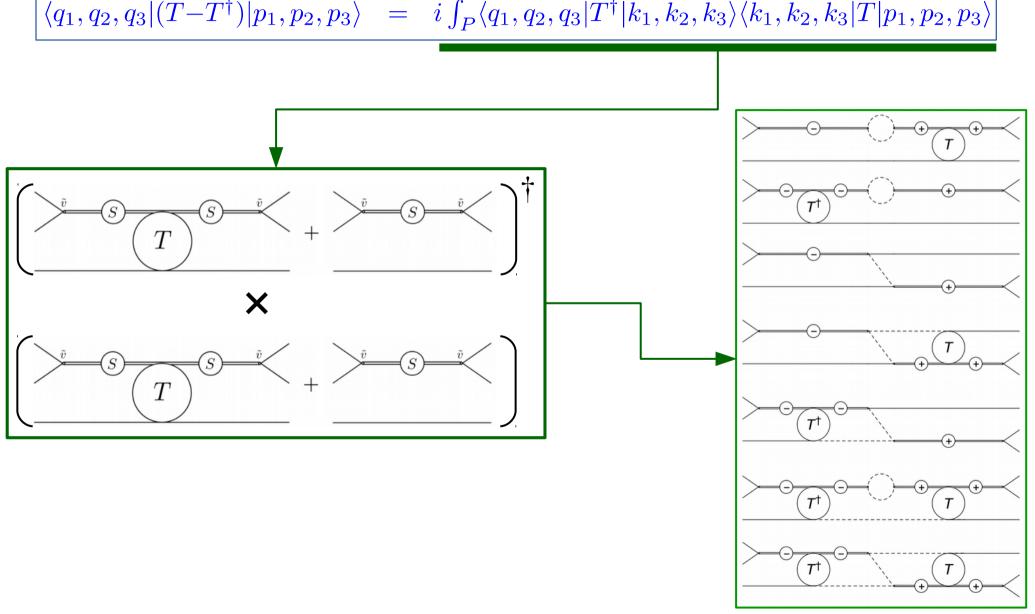
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- Connected part: due to isobar-spectator interaction $\rightarrow T(q_{in}, q_{out}, s)$
 - → 3 unknown functions
 - → 8 kinematic variables

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^{\dagger}) | p_1, p_2, p_3 \rangle \quad = \quad i \int_P \langle q_1, q_2, q_3 | \hat{T}^{\dagger} | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

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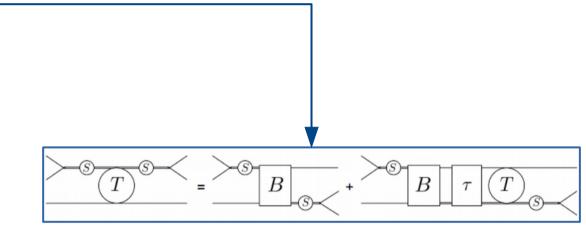
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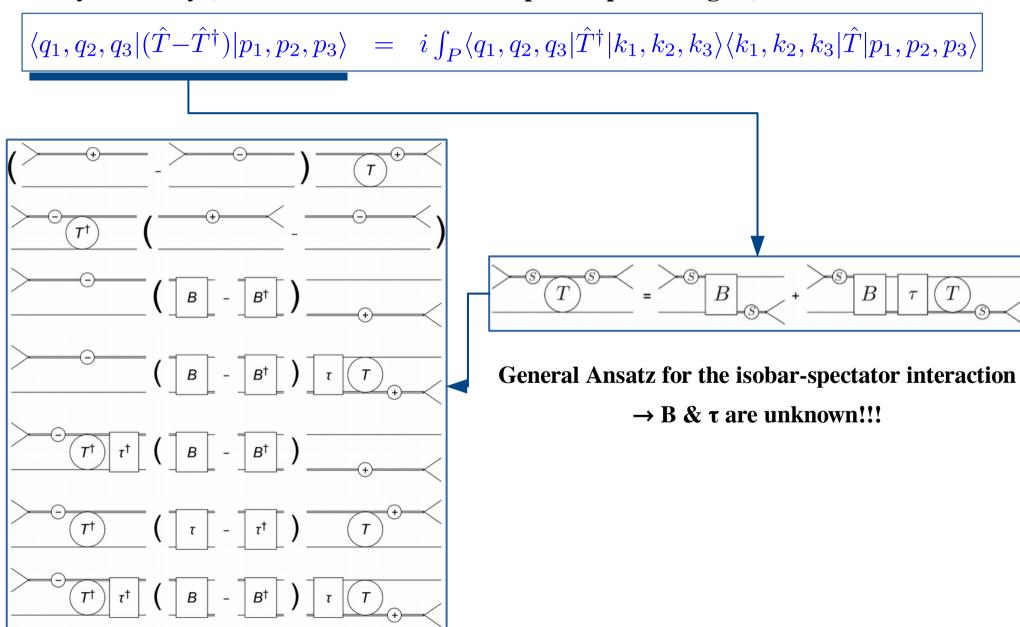
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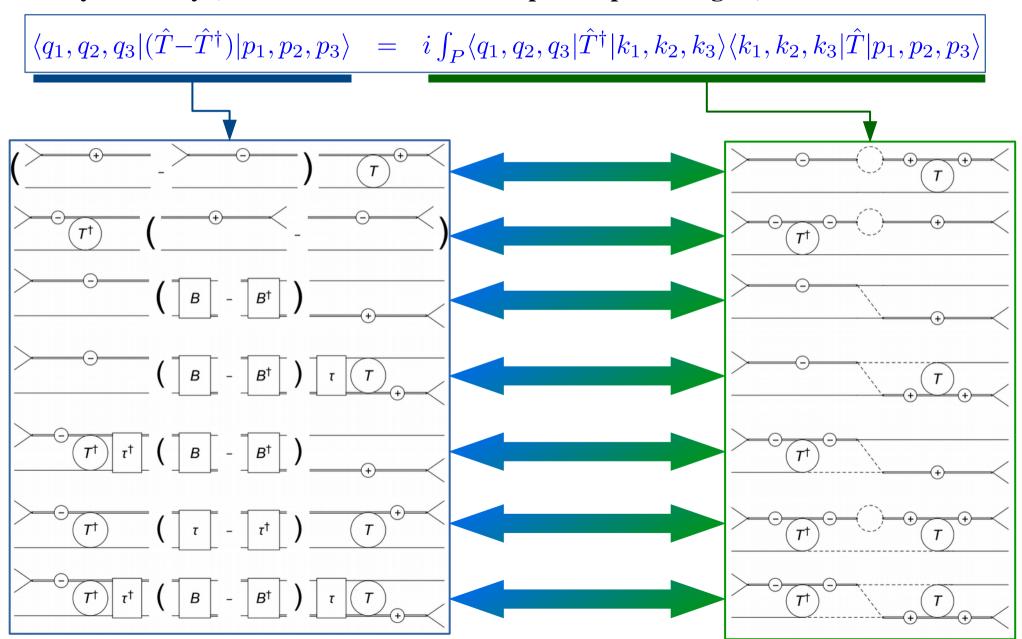
3-body Unitarity (normalization condition ↔ phase space integral)

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^{\dagger}) | p_1, p_2, p_3 \rangle \quad = \quad i \int_P \langle q_1, q_2, q_3 | \hat{T}^{\dagger} | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

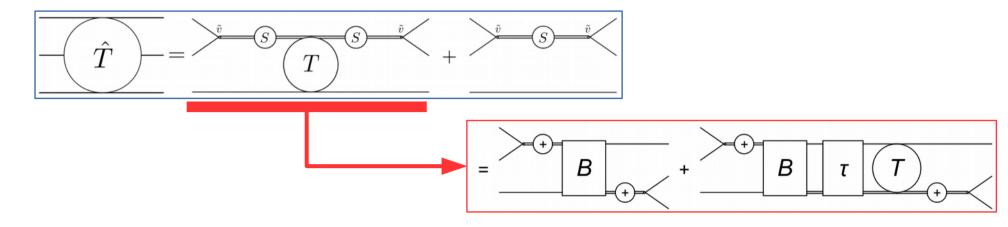


General Ansatz for the isobar-spectator interaction \rightarrow B & τ are unknown!!!

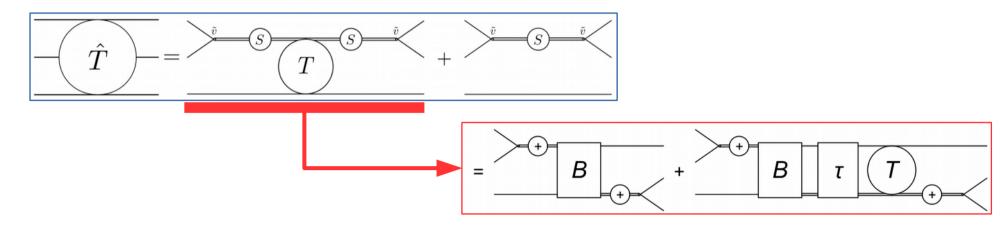




$3 \rightarrow 3$ scattering amplitude as a 3-dimensional integral equation



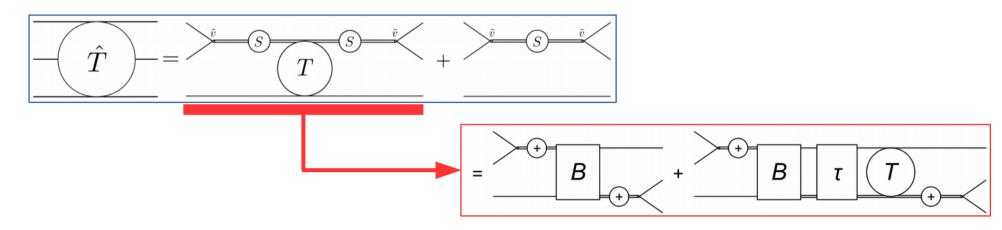
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Unitarity/matching

Disc
$$B(u) = 2\pi i \frac{\delta \left(E_Q - \sqrt{m^2 + \mathbf{Q}^2} \right)}{2\sqrt{m^2 + \mathbf{Q}^2}} v^2$$

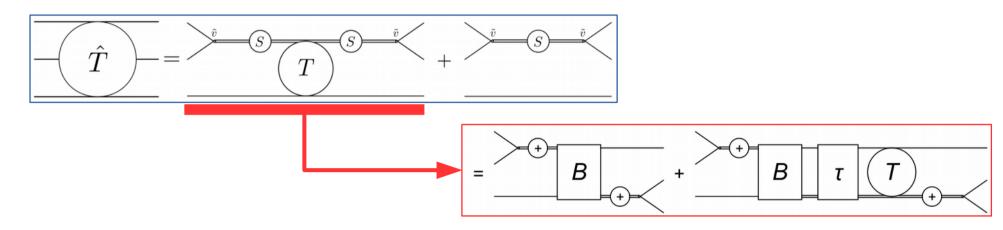
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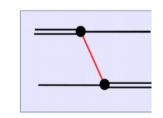
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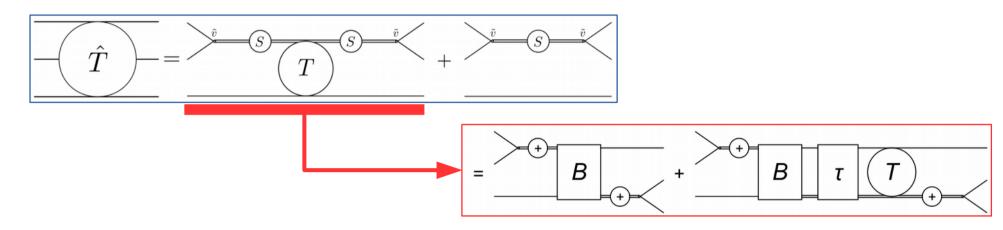


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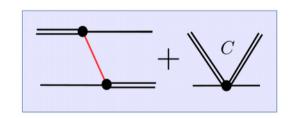


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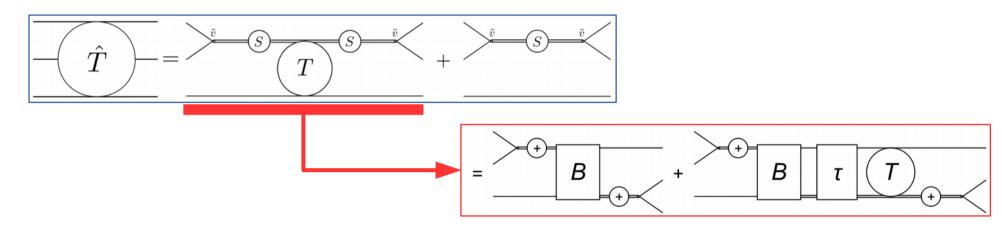


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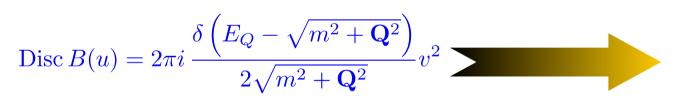
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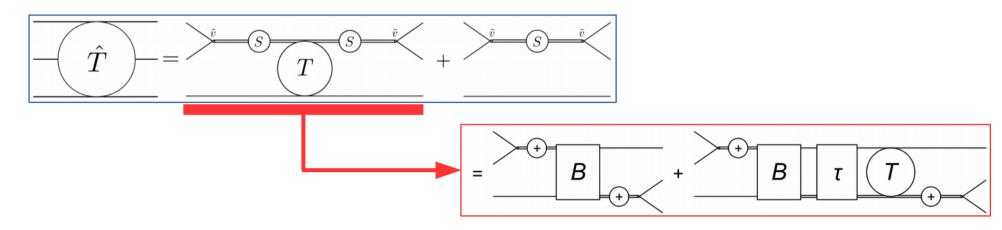


Unitarity/matching



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$$\frac{1}{S(\sigma(k))} = \frac{-i}{64\pi^2 K_{\text{cm}}} \int d^3 \mathbf{\bar{K}} \frac{\delta \left(|\mathbf{\bar{K}}| - K_{\text{cm}}\right)}{\sqrt{(\mathbf{\bar{K}})^2 + m^2}} v^2$$

 $3 \rightarrow 3$ scattering amplitude as a 3-dimensional integral equation



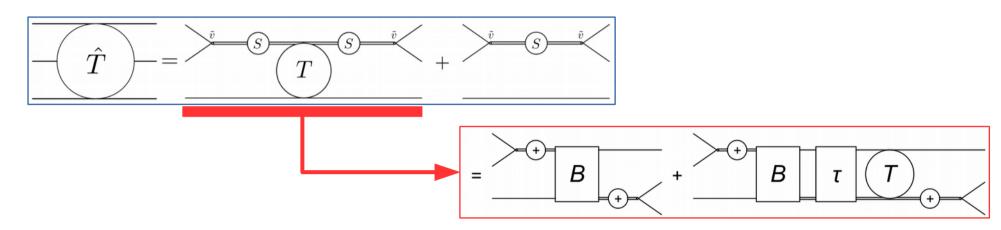
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 $3 \rightarrow 3$ scattering amplitude as a 3-dimensional integral equation



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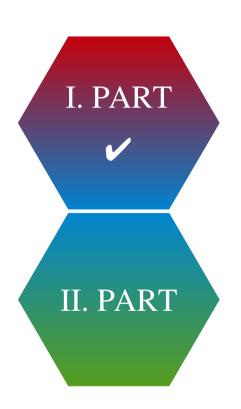
Dispersion relation

THE ONLY IMAGINARY PARTS REQUIRED BY 3b UNITARITY

UNITARITY OF S-MATRIX

► ✓ IMAGINARY PARTS (INF. VOL.)

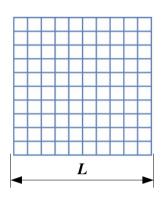
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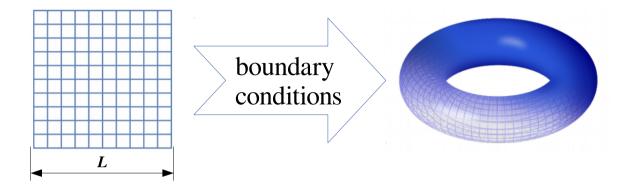
[MM & Döring EPJ A53 (2017)]

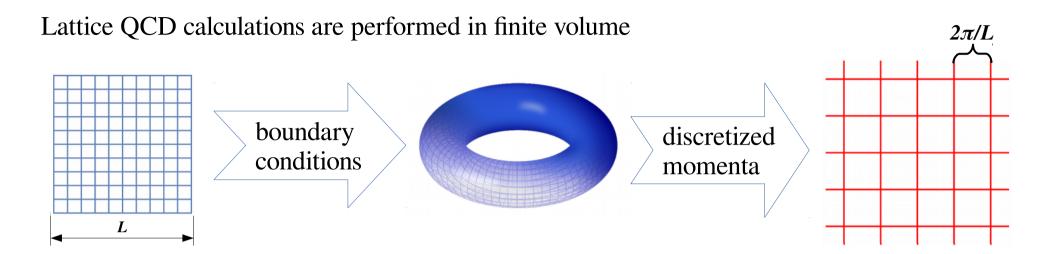
LATTICE QCD SETUP

Lattice QCD calculations are performed in finite volume

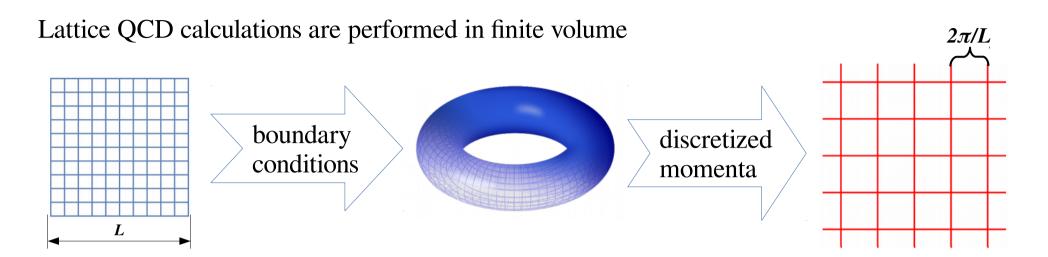


Lattice QCD calculations are performed in finite volume



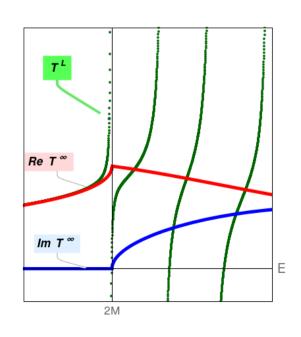


momenta & spectra are discretized

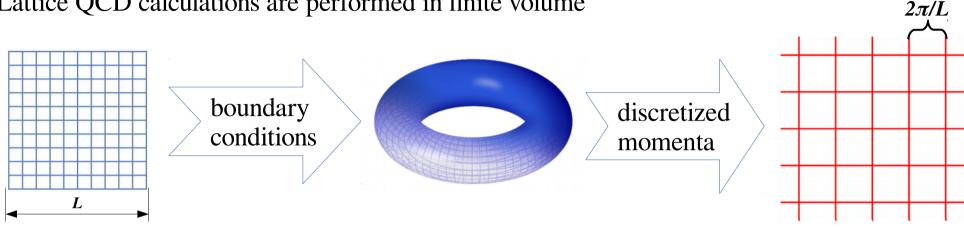


momenta & spectra are discretized

– LSZ formalism relates Greens fct. & S-matrix



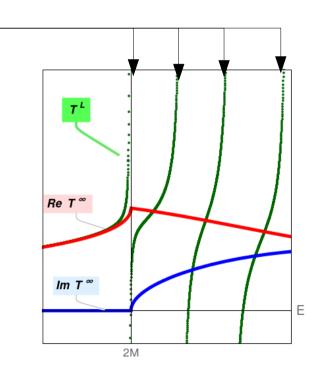
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$$\rightarrow T(E^*) = \infty \leftrightarrow E^* \in Energy-Eigenvalues$$



Lattice QCD calculations are performed in finite volume

2π/L

boundary

conditions

discretized

momenta

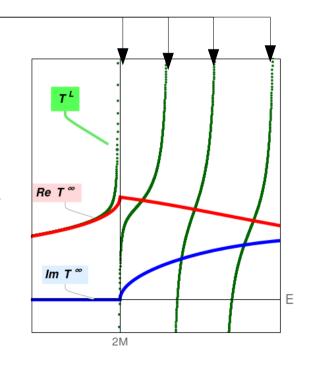


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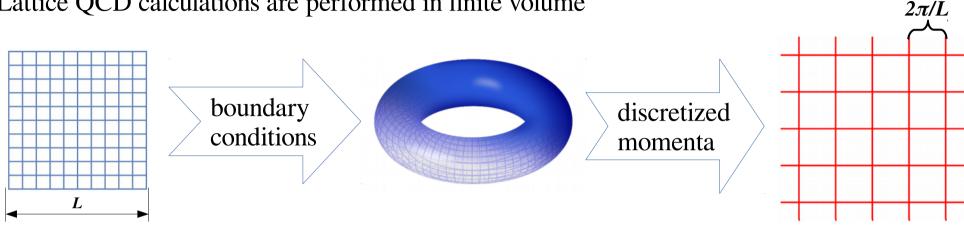
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– well established in 2-body

[Lüscher (1986)]



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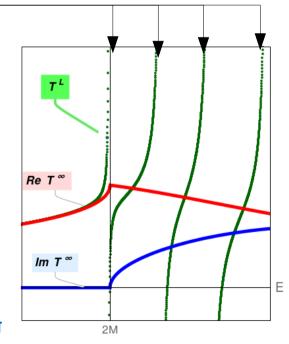
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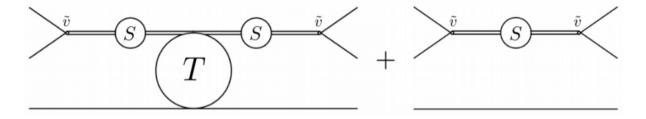
[Lüscher (1986)]

- 3-body analog under investigation Sharpe, Rusetsky, Hansen,

Polejaeva, Briceno, Davoudi, Guo, Pang, MM, Doring

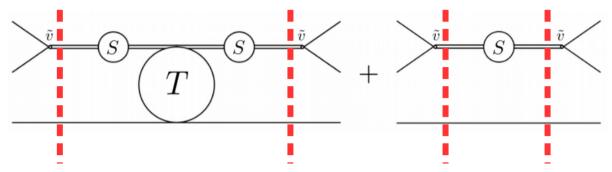


Discretize 3b-scattering amplitude → **3b Quantization Condition**



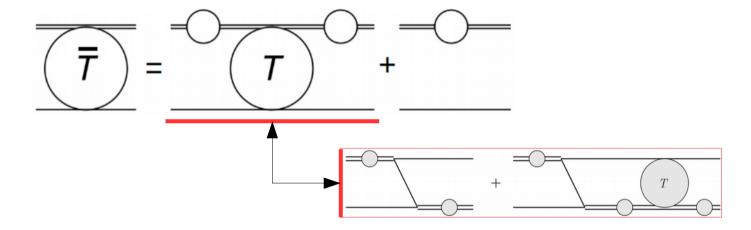
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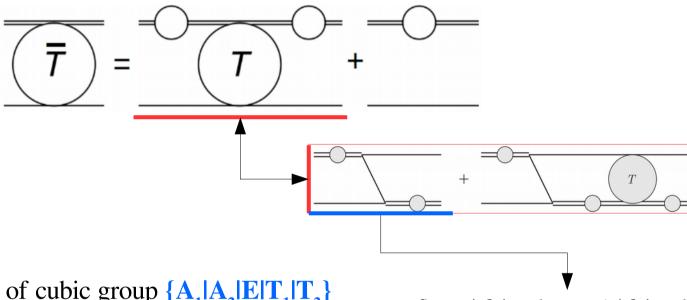
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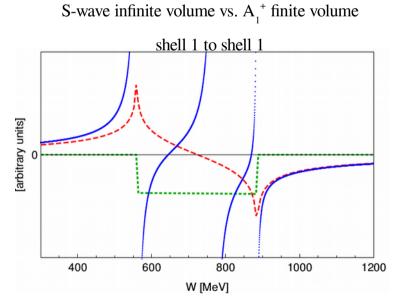


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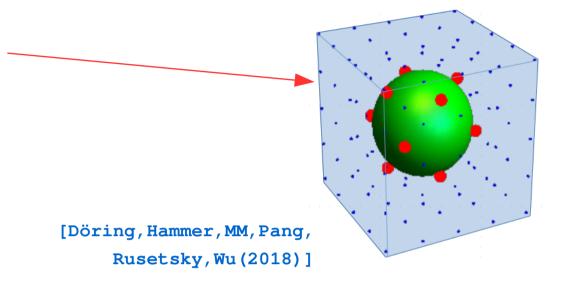


- Project to irreps of cubic group $\{A_1|A_2|E|T_1|T_2\}$
 - reduce dimensionality
 - -B (ope potential) is singular!



1) Separation of variables

- shells = sets of points related by O_h
- inf. vol. analog: radii and angles



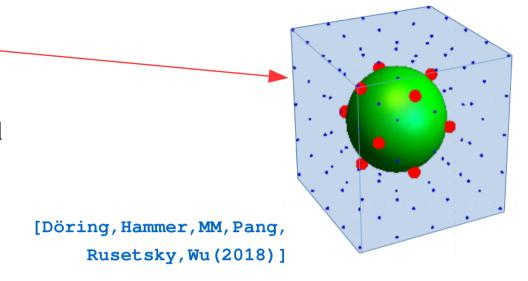
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- inf. vol. analog: radii and angles

2) Find the ONB of functions on each shell

$$- f^{s}(\hat{\mathbf{p}}_{j}) = \sqrt{4\pi} \sum_{\Gamma \alpha} \sum_{u} f_{u}^{\Gamma \alpha s} \chi_{u}^{\Gamma \alpha s}(\hat{\mathbf{p}}_{j})$$

- inf. vol. analog: PWA



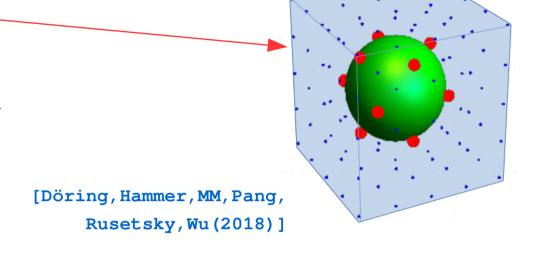
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Projection of 3-body-Quantization-Condition = FINAL RESULT

$$\operatorname{Det}\left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\mathbf{\Gamma}\mathbf{s}\mathbf{s}'}(\mathbf{W}^{2}) + \frac{2\mathbf{E}_{\mathbf{s}}\mathbf{L}^{3}}{\vartheta(\mathbf{s})}\tau_{\mathbf{s}}(\mathbf{W}^{2})^{-1}\delta_{\mathbf{s}\mathbf{s}'}\delta_{\mathbf{u}\mathbf{u}'}\right) = \mathbf{0}$$

[MM, Döring]

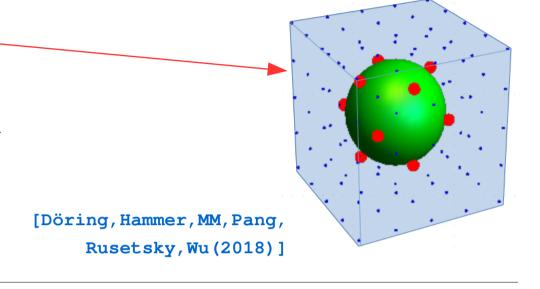
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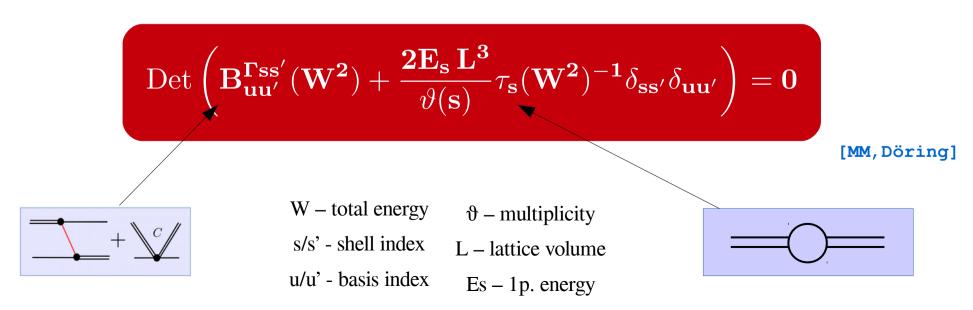
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$$- f^{s}(\hat{\mathbf{p}}_{j}) = \sqrt{4\pi} \sum_{\Gamma \alpha} \sum_{u} f_{u}^{\Gamma \alpha s} \chi_{u}^{\Gamma \alpha s}(\hat{\mathbf{p}}_{j})$$

- inf. vol. analog: PWA



Projection of 3-body-Quantization-Condition = FINAL RESULT



$$\operatorname{Det}\left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\mathbf{\Gamma}\mathbf{s}\mathbf{s}'}(\mathbf{W^2}) + \frac{2\mathbf{E_s}\,\mathbf{L^3}}{\vartheta(\mathbf{s})}\tau_{\mathbf{s}}(\mathbf{W^2})^{-1}\delta_{\mathbf{s}\mathbf{s}'}\delta_{\mathbf{u}\mathbf{u}'}\right) = \mathbf{0}$$

$$\operatorname{Det}\left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\mathbf{\Gamma}\mathbf{s}\mathbf{s}'}(\mathbf{W}^{2}) + \frac{2\mathbf{E}_{\mathbf{s}}\mathbf{L}^{3}}{\vartheta(\mathbf{s})}\tau_{\mathbf{s}}(\mathbf{W}^{2})^{-1}\delta_{\mathbf{s}\mathbf{s}'}\delta_{\mathbf{u}\mathbf{u}'}\right) = \mathbf{0}$$

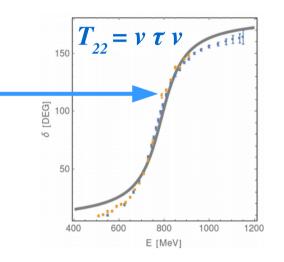
• 3 particles in finite volume: m=138 MeV, L=3 fm

$$\operatorname{Det}\left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\mathbf{\Gamma}\mathbf{s}\mathbf{s}'}(\mathbf{W}^{2}) + \frac{2\mathbf{E}_{\mathbf{s}}\mathbf{L}^{3}}{\vartheta(\mathbf{s})}\tau_{\mathbf{s}}(\mathbf{W}^{2})^{-1}\delta_{\mathbf{s}\mathbf{s}'}\delta_{\mathbf{u}\mathbf{u}'}\right) = \mathbf{0}$$

- 3 particles in finite volume: m=138 MeV, L=3 fm
- one S-wave isobar → two unknowns:
 - vertex(Isobar→2 stable particles)
 - subtraction constant (~mass)

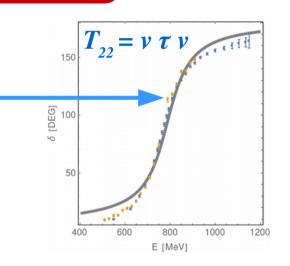
$$\operatorname{Det}\left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\mathbf{\Gamma}\mathbf{s}\mathbf{s}'}(\mathbf{W^2}) + \frac{2\mathbf{E_s}\,\mathbf{L^3}}{\vartheta(\mathbf{s})}\tau_{\mathbf{s}}(\mathbf{W^2})^{-1}\delta_{\mathbf{s}\mathbf{s}'}\delta_{\mathbf{u}\mathbf{u}'}\right) = \mathbf{0}$$

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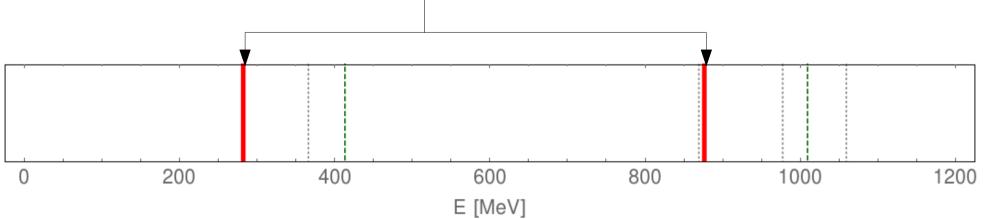
- 3 particles in finite volume: m=138 MeV, L=3 fm
- one S-wave isobar \rightarrow two unknowns:
 - vertex(Isobar→2 stable particles)
 - subtraction constant (~mass)
- Project to $\Gamma = A^{1+}$
 - → prediction of 3body energy-eigenlevels



$$\operatorname{Det}\left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\mathbf{\Gamma}\mathbf{s}\mathbf{s}'}(\mathbf{W^2}) + \frac{2\mathbf{E_s}\,\mathbf{L^3}}{\vartheta(\mathbf{s})}\tau_{\mathbf{s}}(\mathbf{W^2})^{-1}\delta_{\mathbf{s}\mathbf{s}'}\delta_{\mathbf{u}\mathbf{u}'}\right) = \mathbf{0}$$

- 3 particles in finite volume: m=138 MeV, L=3 fm
- one S-wave isobar \rightarrow two unknowns:
 - vertex(Isobar→2 stable particles)
 - subtraction constant (~mass)
- Project to $\Gamma = A^{1+}$





1000

E [MeV]

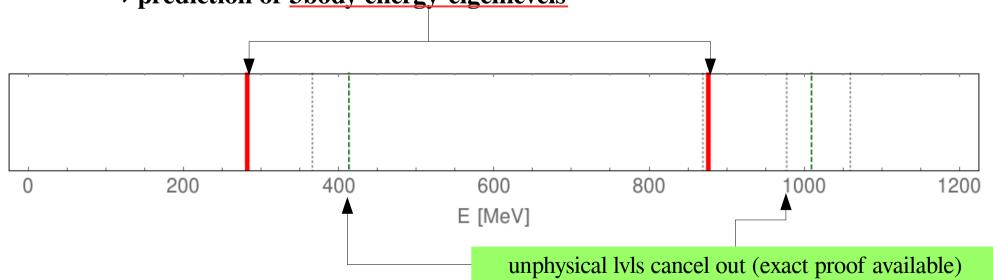
1200

δ [DEG]

$$\operatorname{Det}\left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\mathbf{\Gamma}\mathbf{s}\mathbf{s}'}(\mathbf{W^2}) + \frac{2\mathbf{E_s}\,\mathbf{L^3}}{\vartheta(\mathbf{s})}\tau_{\mathbf{s}}(\mathbf{W^2})^{-1}\delta_{\mathbf{s}\mathbf{s}'}\delta_{\mathbf{u}\mathbf{u}'}\right) = \mathbf{0}$$

- 3 particles in finite volume: m=138 MeV, L=3 fm
- one S-wave isobar \rightarrow two unknowns:
 - vertex(Isobar→2 stable particles)
 - subtraction constant (~mass)
- Project to $\Gamma = A^{1+}$





1000

E [MeV]

1200

SUMMARY/OUTLOOK

3-body scattering amplitude derived from 2&3 body Unitarity

- interaction kernel = one-particle-exchange
- flexible parametrization: real contributions can be added to the kernel

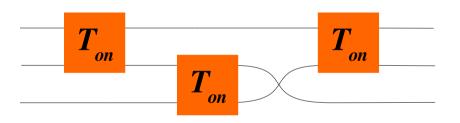
TBD: analysis of physical systems

3-body Quantization Condition in fin. vol. derived

- cancellations of unphysical poles revealed
- projection to irreps done
- technical feasibility on a numerical example
- the only approximation = number of isobars

TBD: multiple channels

TBD: inclusion of isospin & angular momentum



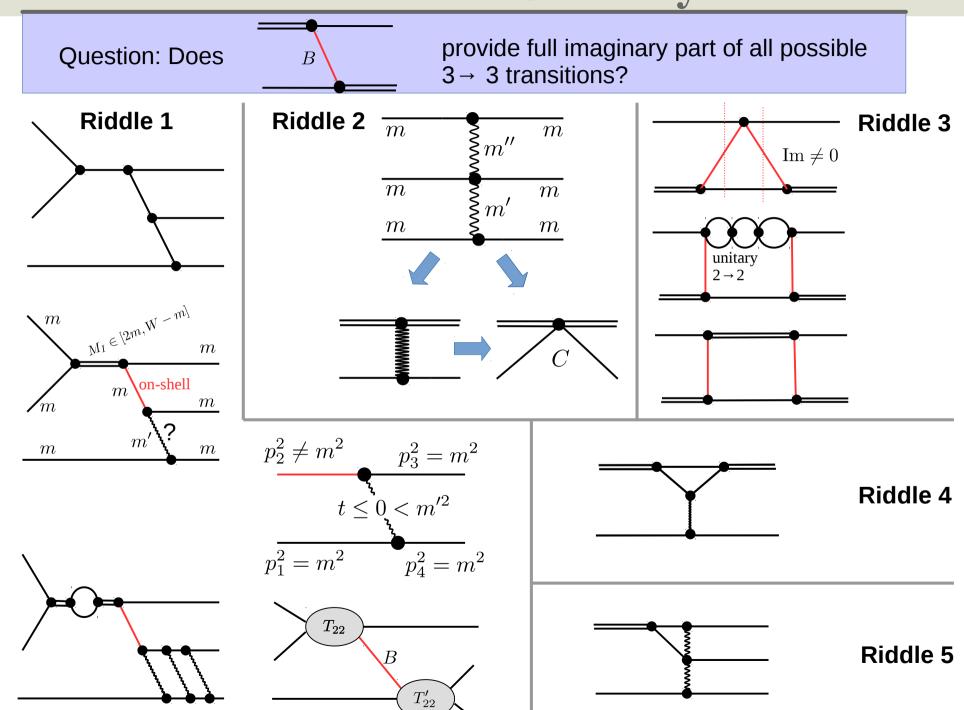
"power of Unitarity"

THANK YOU!



SPARES

The Power of Unitarity



• Projection of T

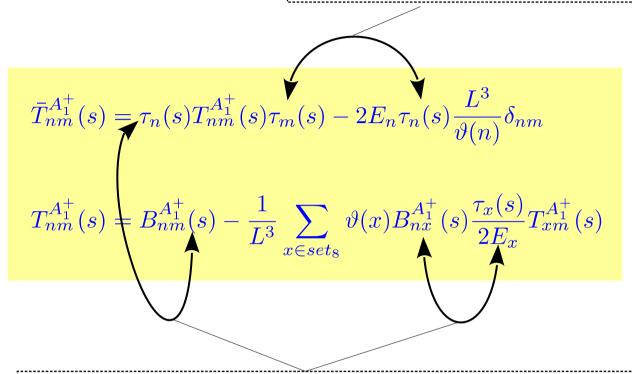
$$T^{ss'}(\hat{\mathbf{p}}_{j}, \hat{\mathbf{p}}_{j'}) = 4\pi \sum_{\Gamma\alpha} \sum_{uu'} \chi_{u}^{\Gamma\alpha s}(\hat{\mathbf{p}}_{j}) T_{uu'}^{\Gamma ss'} \chi_{u'}^{\Gamma\alpha s'}(\hat{\mathbf{p}}_{j'}),$$

$$T_{uu'}^{\Gamma ss'} = \frac{4\pi}{\vartheta(s)\vartheta(s')} \sum_{j=1}^{\vartheta(s)} \sum_{j'=1}^{\vartheta(s')} \chi_{u}^{\Gamma\alpha s}(\hat{\mathbf{p}}_{j}) T^{ss'}(\hat{\mathbf{p}}_{j}, \hat{\mathbf{p}}_{j'}) \chi_{u'}^{\Gamma\alpha s'}(\hat{\mathbf{p}}_{j'})$$

QUANTIZATION CONDITION

Cancellations:

 \rightarrow fin. vol. normalization of δ -distribution!



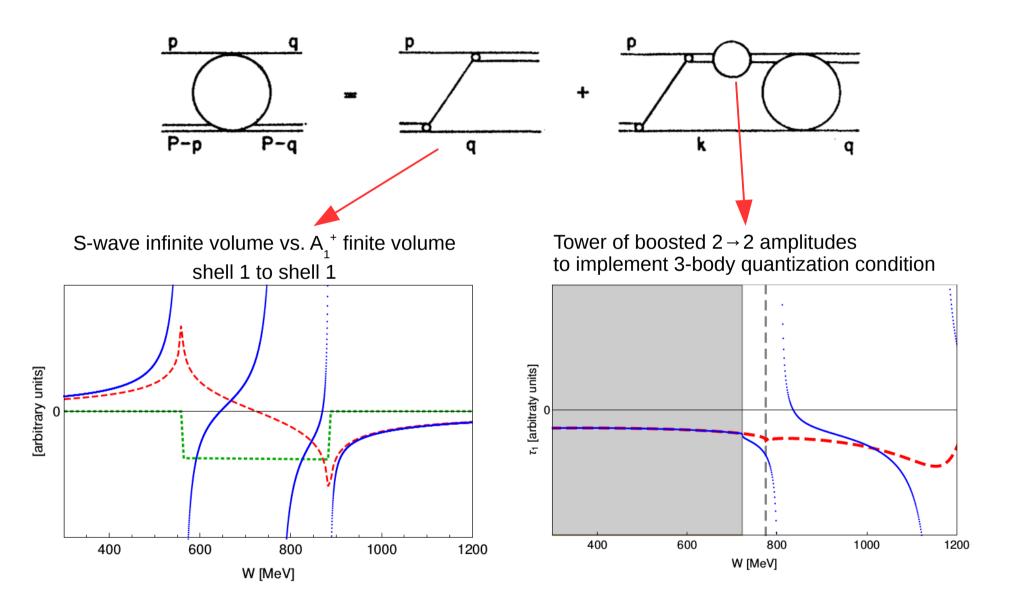
$$B^{A_1^+}$$
 singular at $W^+ = E_m + E_n + E(\boldsymbol{q}_{nj} + \boldsymbol{p}_{mi})$

$$\tau_m^{-1}$$
 singular at $W^{\pm\pm} = E_m \pm E((2\pi/L)\boldsymbol{y}) \pm E((2\pi/L)\boldsymbol{y} + \boldsymbol{p}_{mi})$ for $\boldsymbol{y} \in \mathbb{Z}^3$

- when isobar-momenta are discretized in the 3-body cms momenta

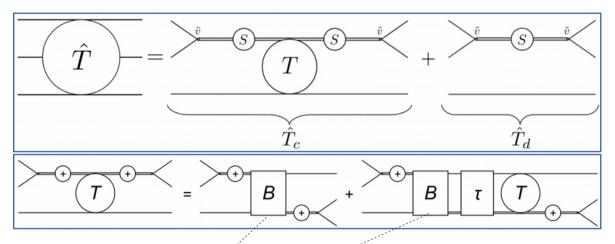
$$\tau = \sigma(k) - M_0^2 - \frac{1}{(2\pi)^3} \int d^3 \ell \frac{\lambda^2}{2E_{\ell}(\sigma(k) - 4E_{\ell}^2 + i\epsilon)}$$

Power-law finite-volume effects dictated by three-body unitarity



SCATTERING AMPLITUDE

 $3 \rightarrow 3$ scattering amplitude is a 3-dimensional integral equation



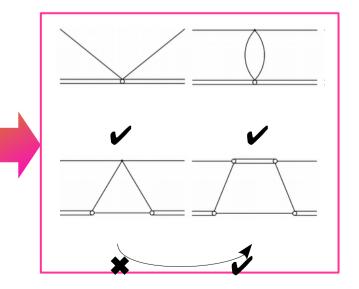
- Imaginary parts of **B**, **S** are fixed by **unitarity/matching**
- For simplicity $v=\lambda$ (full relations available)

Disc
$$B(u) = 2\pi i \lambda^2 \frac{\delta \left(E_Q - \sqrt{m^2 + \mathbf{Q}^2} \right)}{2\sqrt{m^2 + \mathbf{Q}^2}}$$

• un-subtracted dispersion relation

$$\langle q|B(s)|p\rangle = -\frac{\lambda^2}{2\sqrt{m^2+\mathbf{Q}^2}\left(E_Q-\sqrt{m^2+\mathbf{Q}^2}+i\epsilon\right)}$$

• one- π exchange in TOPT $\rightarrow RESULT!$



Unitarity & Matching

• 3-body Unitarity (normalization condition ↔ phase space integral)

