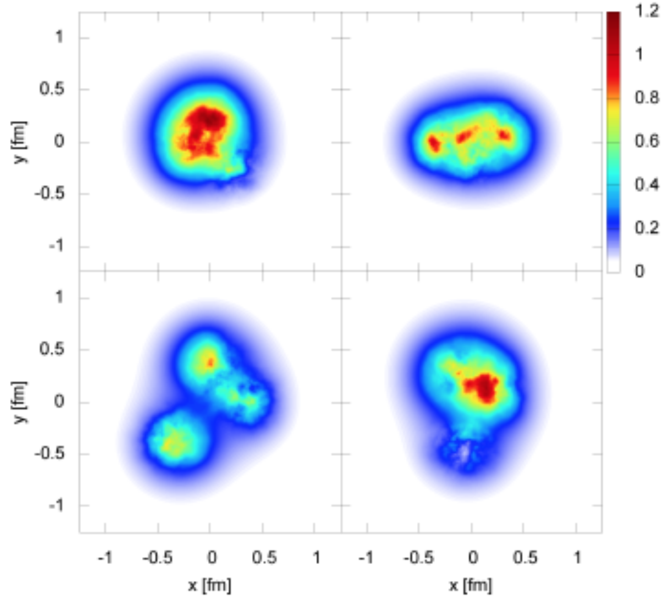
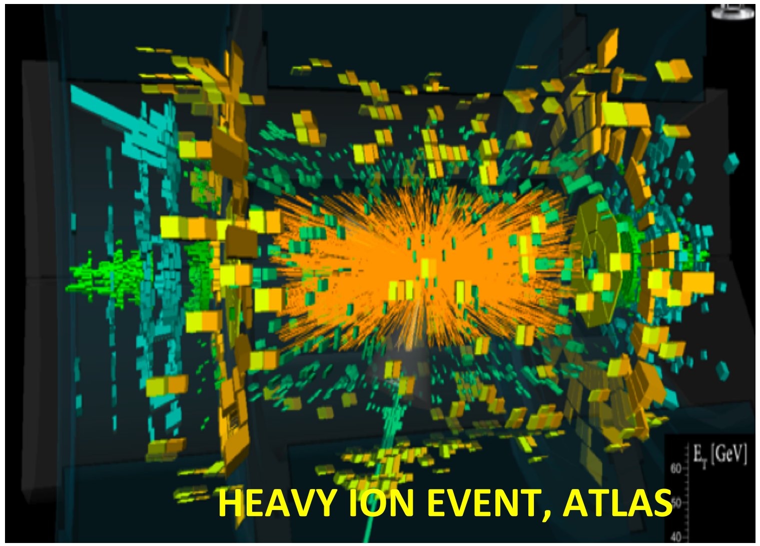
What does the 3-D structure of the proton teach us about matter in heavy-ion collisions?

Raju Venugopalan

Brookhaven National Laboratory



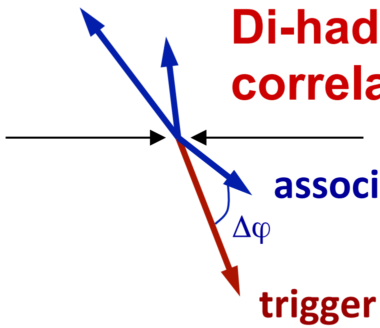
Mantysaari,Schenke,arXiv:1603.043049



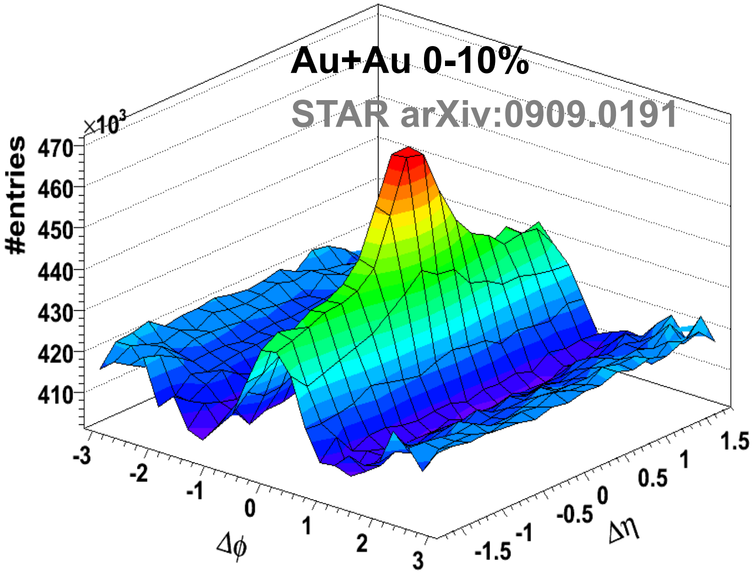
QCD evolution workshop, Jlab, May22-26, 2017

Azimuthal correlations in A+A collisions: the ridge

Di-hadron correlations



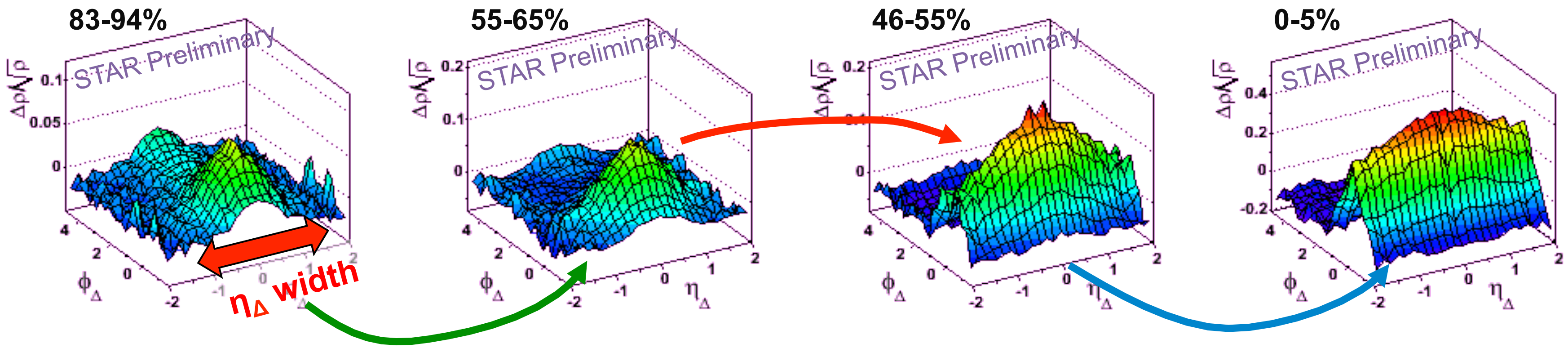
$3 < p_{t,trigger} < 4$ GeV
$p_{t,assoc.} > 2$ GeV



Novel long range "near-side" (Δφ≈0) collimation

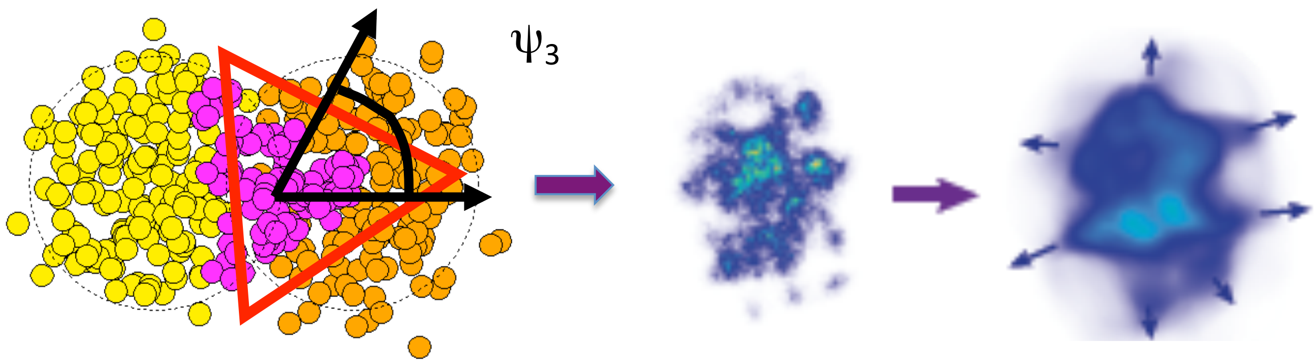
Long range "away-side" (Δφ≈π) correlations suppressed…jet quenching

Also seen in untriggered correlations



First seen by RHIC Au+Au experiments: STAR, PHOBOS, PHENIX

The A+A ridge and "collectivity"



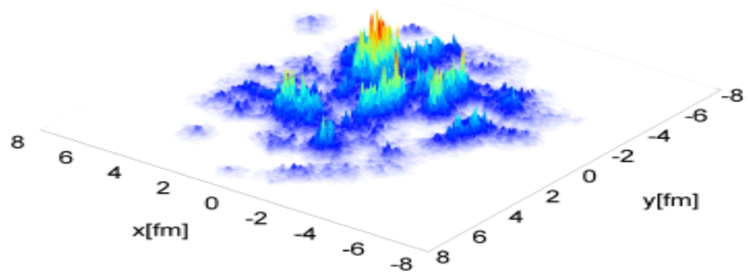
Alver,Roland, PRC81(2010)054905
Alver, Gombeaud, Luzum, Ollitrault, PRC82 (2010) 03491

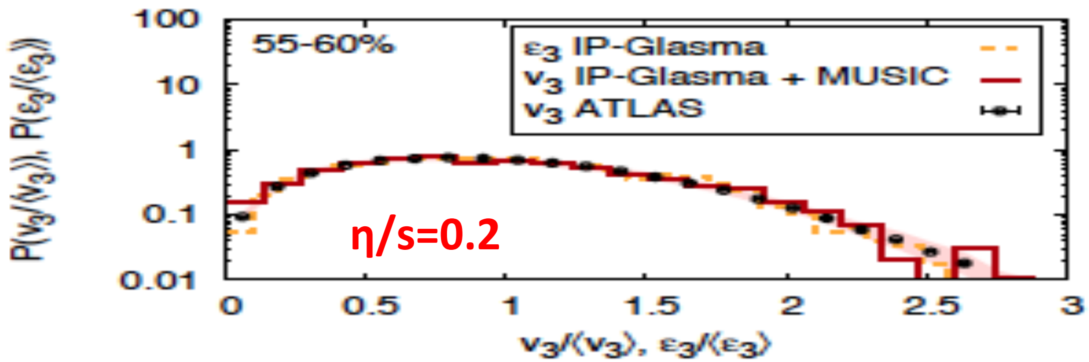
Structure of ridge correlations: hydrodynamic flow driven by event-by-event ``eccentricity" fluctuations in nucleon positions

In hydro: multi-particle correlations factorize into product of single particle distributions that are commonly correlated to an "event plane"

$$\frac{1}{N_{\rm trig} N_{\rm assoc.}} \frac{d^2N}{d\Delta\Phi} = 1 + V_1\, Cos(\Delta\Phi) + V_2\, Cos(2\Delta\Phi) + \cdots$$

IP-Glasma

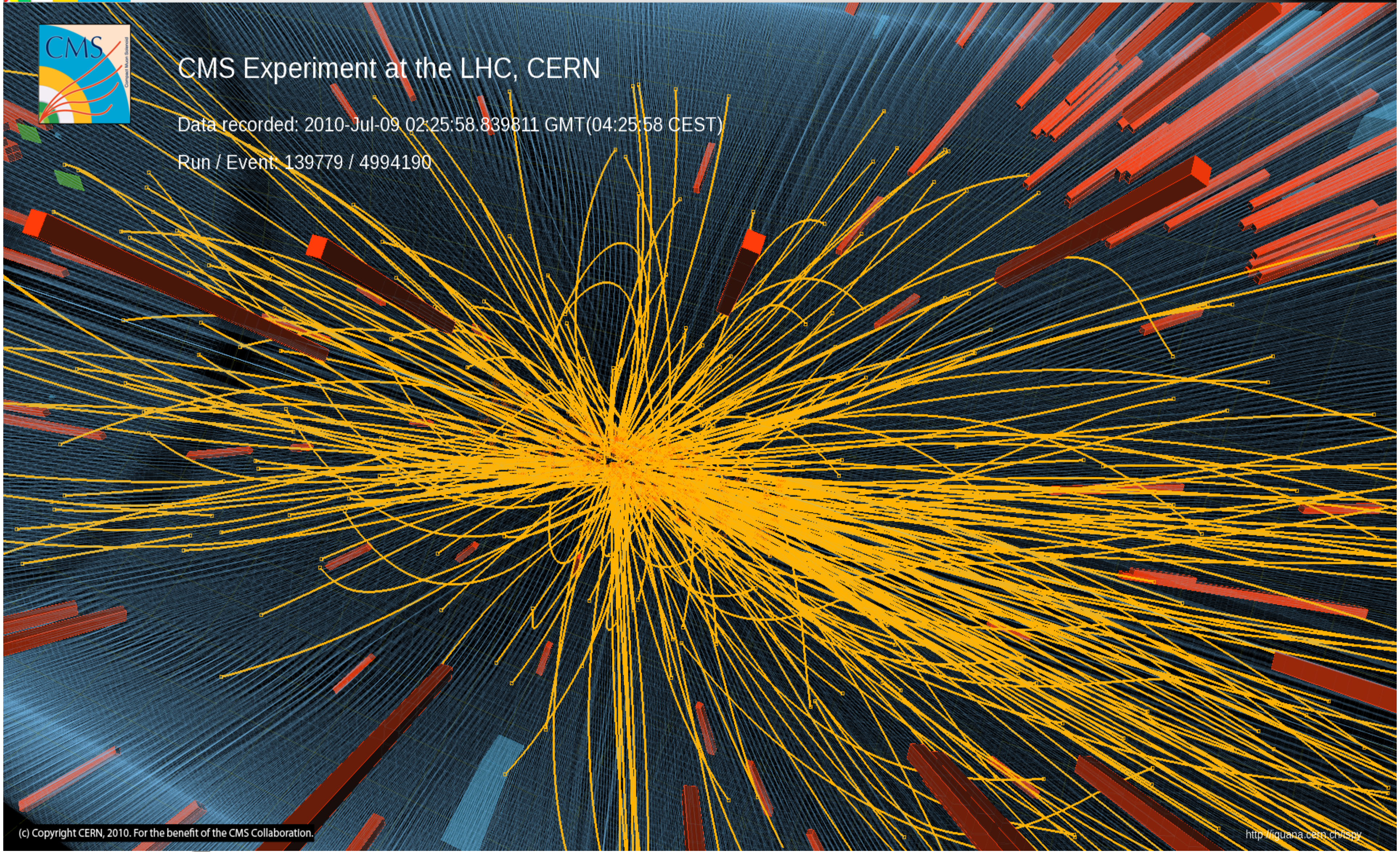




Some evidence of sensitivity of data to sub-nucleon scale fluctuations

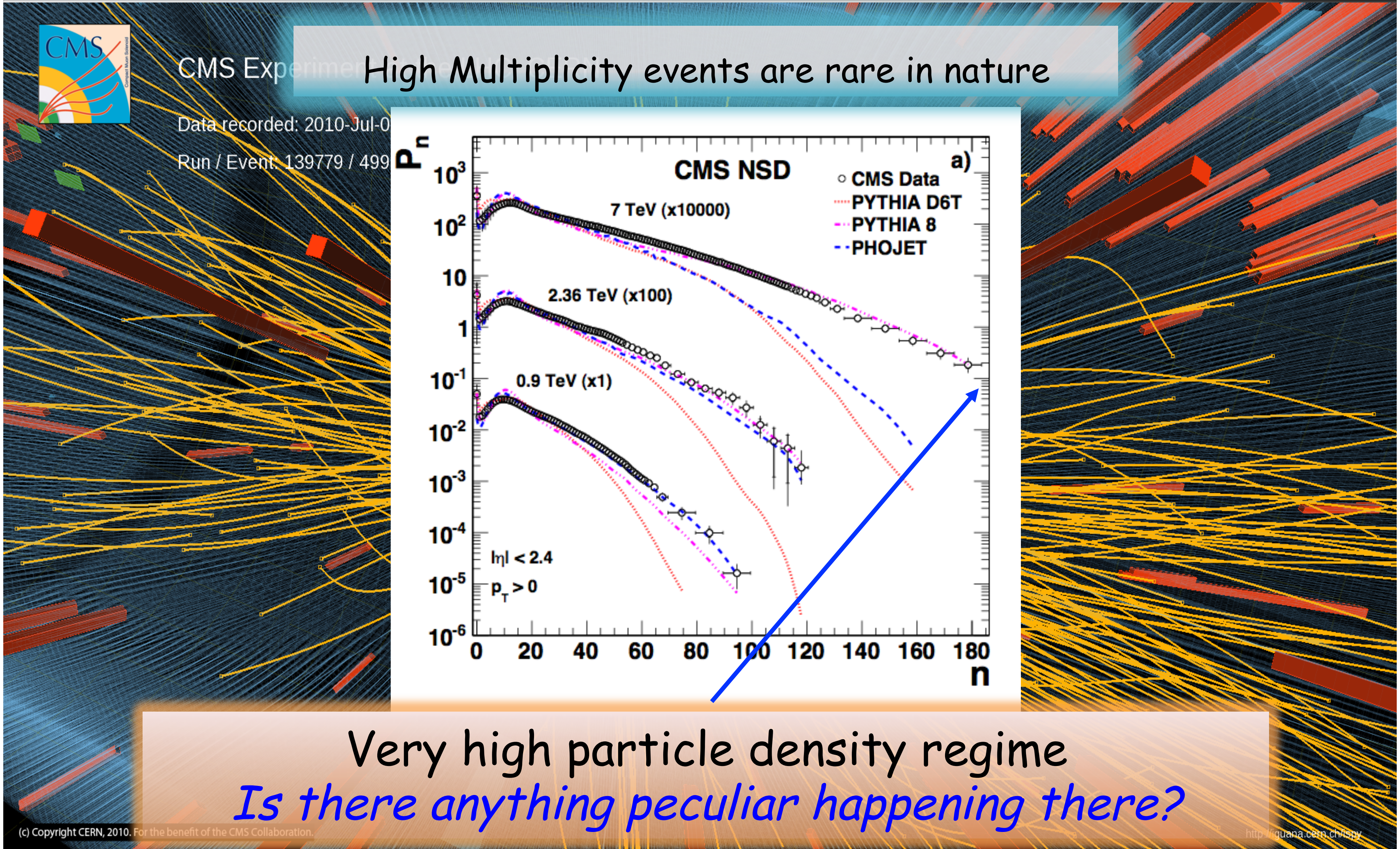
Gale,Jeon,Schenke,Tribedy,Venugopalan, PRL110 (2013) 012302

High Multiplicity pp collisions



High Multiplicity pp collisions

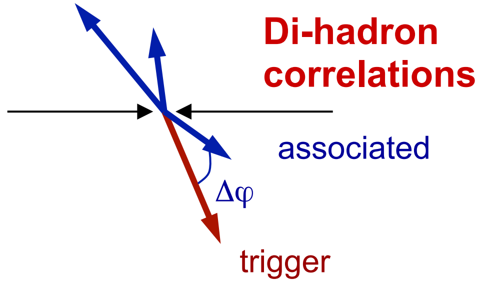
High Multiplicity events are rare in nature

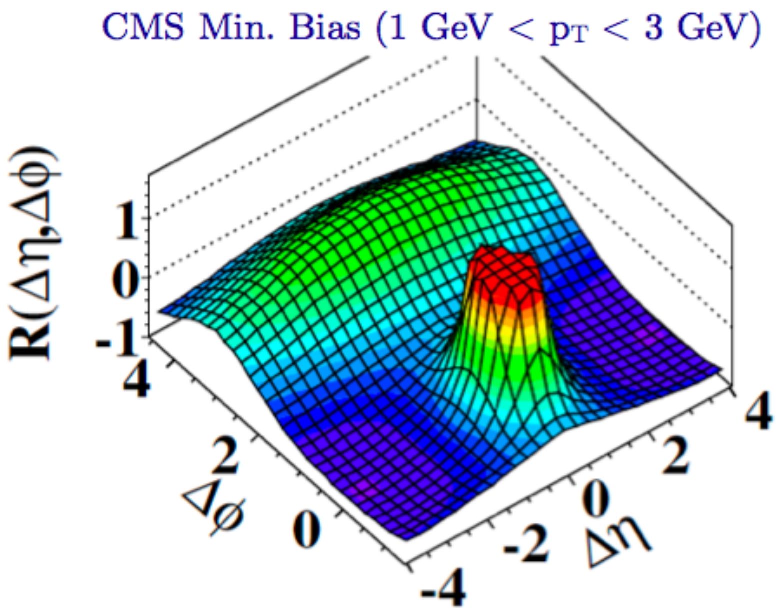


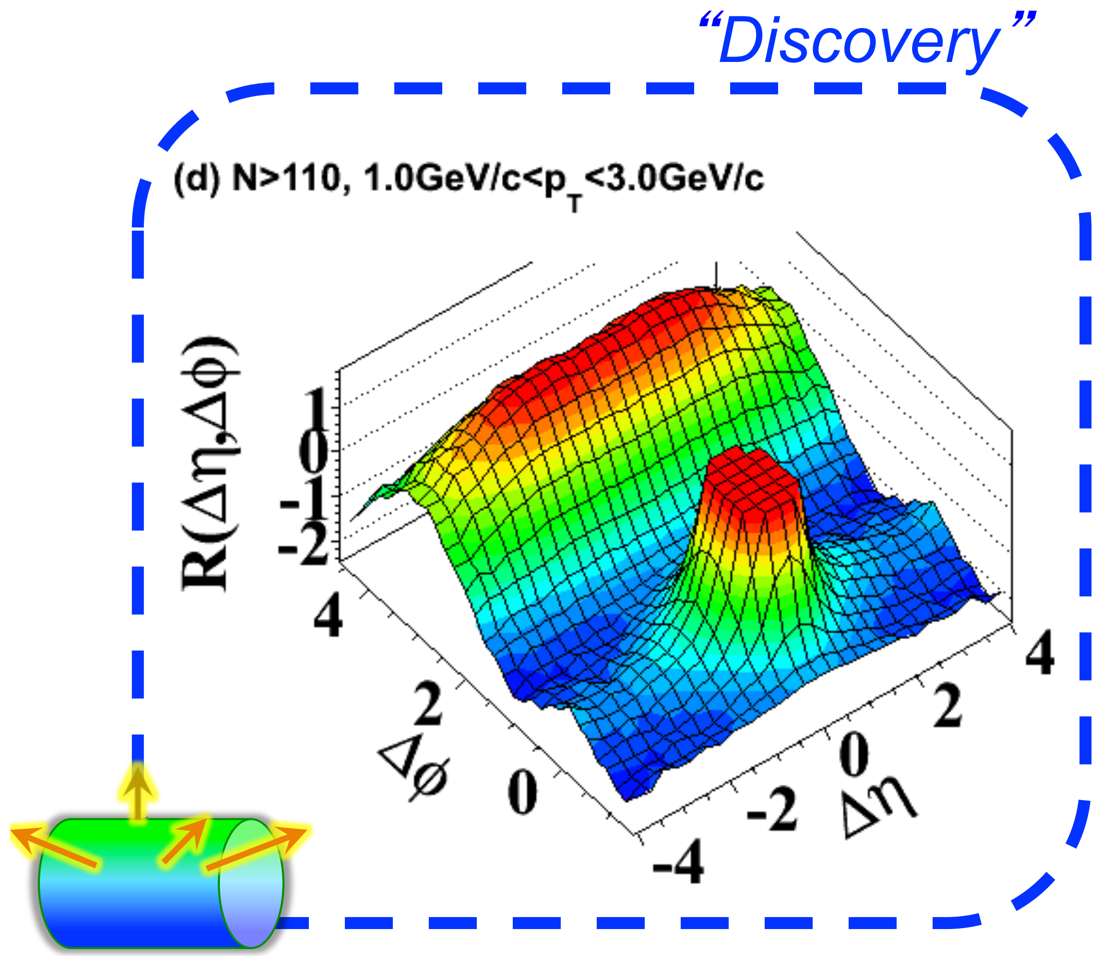
Very high particle density regime

Is there anything peculiar happening there?

Two particle correlations: CMS results





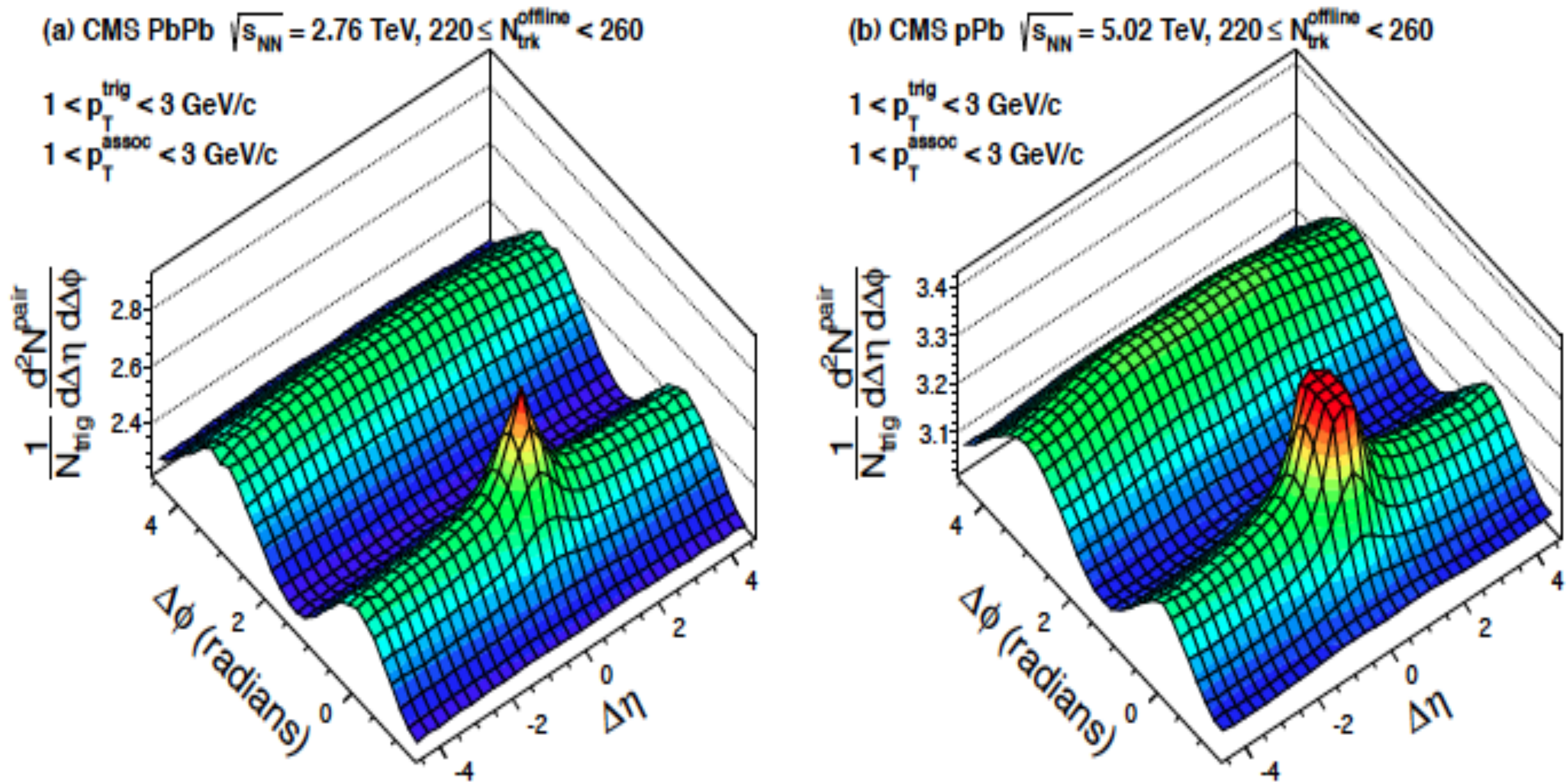


Observation of Long-Range Near-Side Angular Correlations in Proton-Proton Collisions at the LHC CMS Collaboration (Vardan Khachatryan (Yerevan Phys. Inst.) *et al.).* **JHEP 1009 (2010) 091**

Cited by 564 records

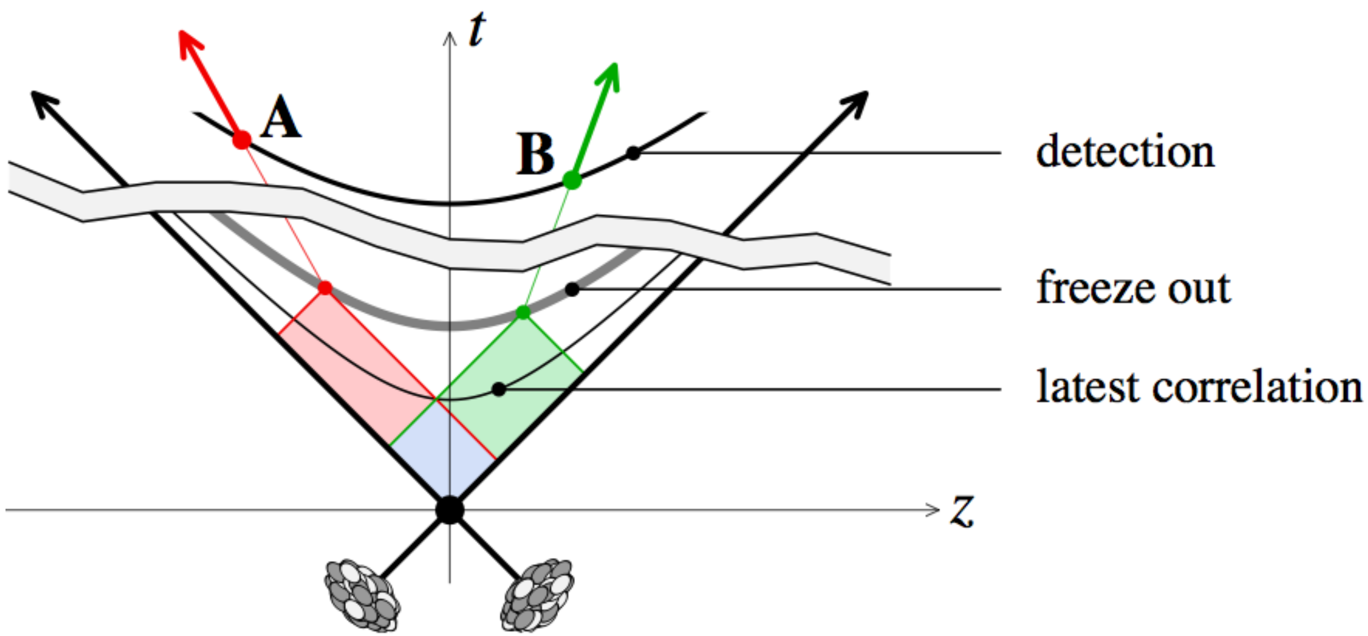
5th most cited CMS physics paper to date!

Striking results from LHC p+A collisions



p+A ridge much large than p+p at same multiplicity and nearly as large as that in peripheral Pb+Pb collisions

Long range rapidity correlations as a chronometer



$$\tau \leq \tau_{\text{frz-out}} \exp\left(-\frac{1}{2}|y_A - y_B|\right)$$

❖ **Long range correlations sensitive to very early time (fractions of a femtometer ~ 10^{-24} seconds) dynamics in collisions**

What's the underlying QCD dynamics?

- **Is it final state collective flow of the world's smallest droplets?**
- **Is it the initial state dynamics arising from rare configurations in the hadron wavefunctions?**
- **Or, is it some combination, where there is a smooth transition from one description to the other?**

Option 1 stretches to the limit -- the applicability of thermodynamic and hydrodynamic concepts in high energy physics

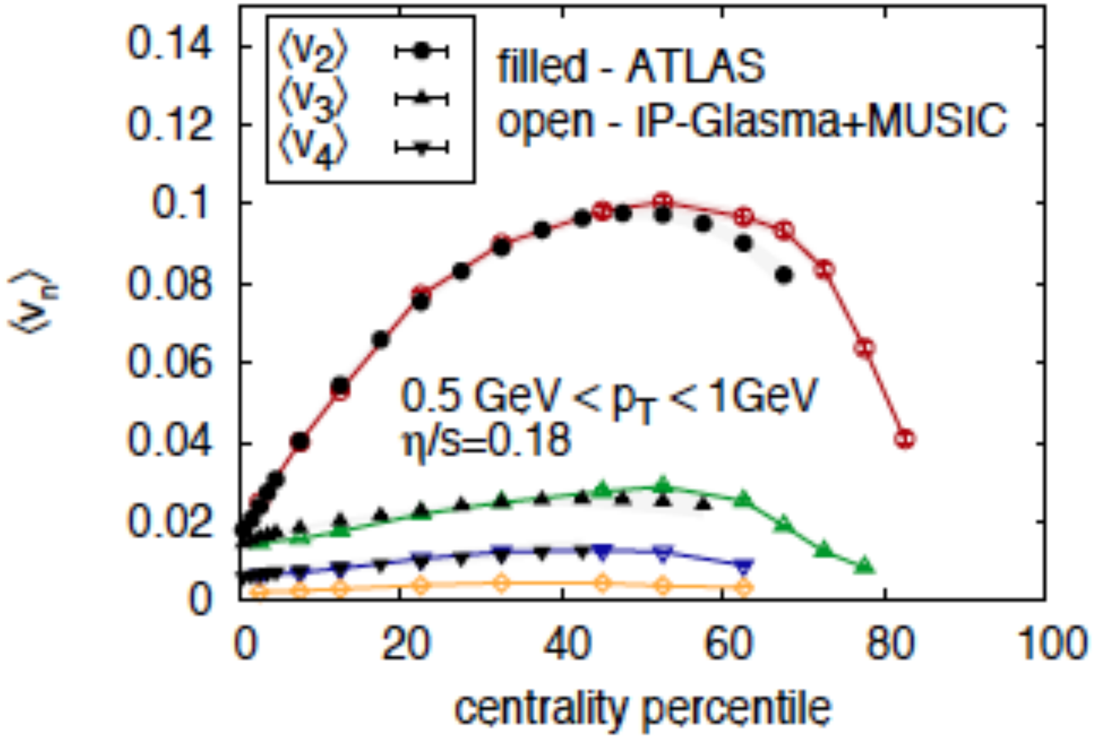
Option 2 stretches to the limit -- our understanding of the quark-gluon sub-structure of hadrons

What's the smallest sized QGP droplet?

IP-Glasma: Initial state from Yang-Mills evolution of two lumpy light cone sources

MUSIC: Event-by-event relativistic viscous hydrodynamics

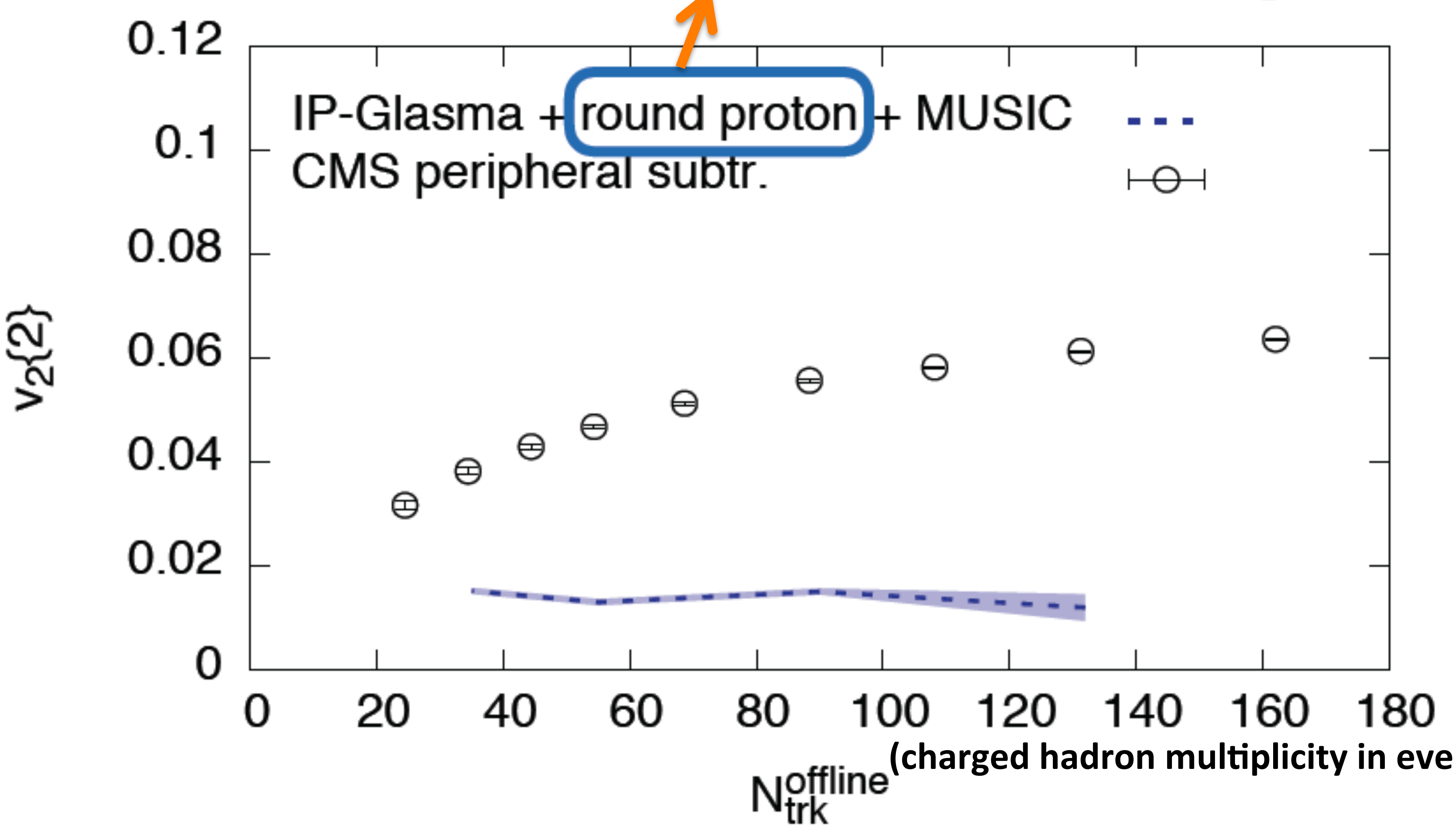
Schenke, Venugopalan, PRL 113 (2014) 102301



Where does the hydro paradigm break down?

Shape fluctuations essential to generate flow-I

Gaussian impact parameter profile: $T_G(b) = \frac{1}{2\pi B_G} \exp\left(-b^2/2B_G\right)$

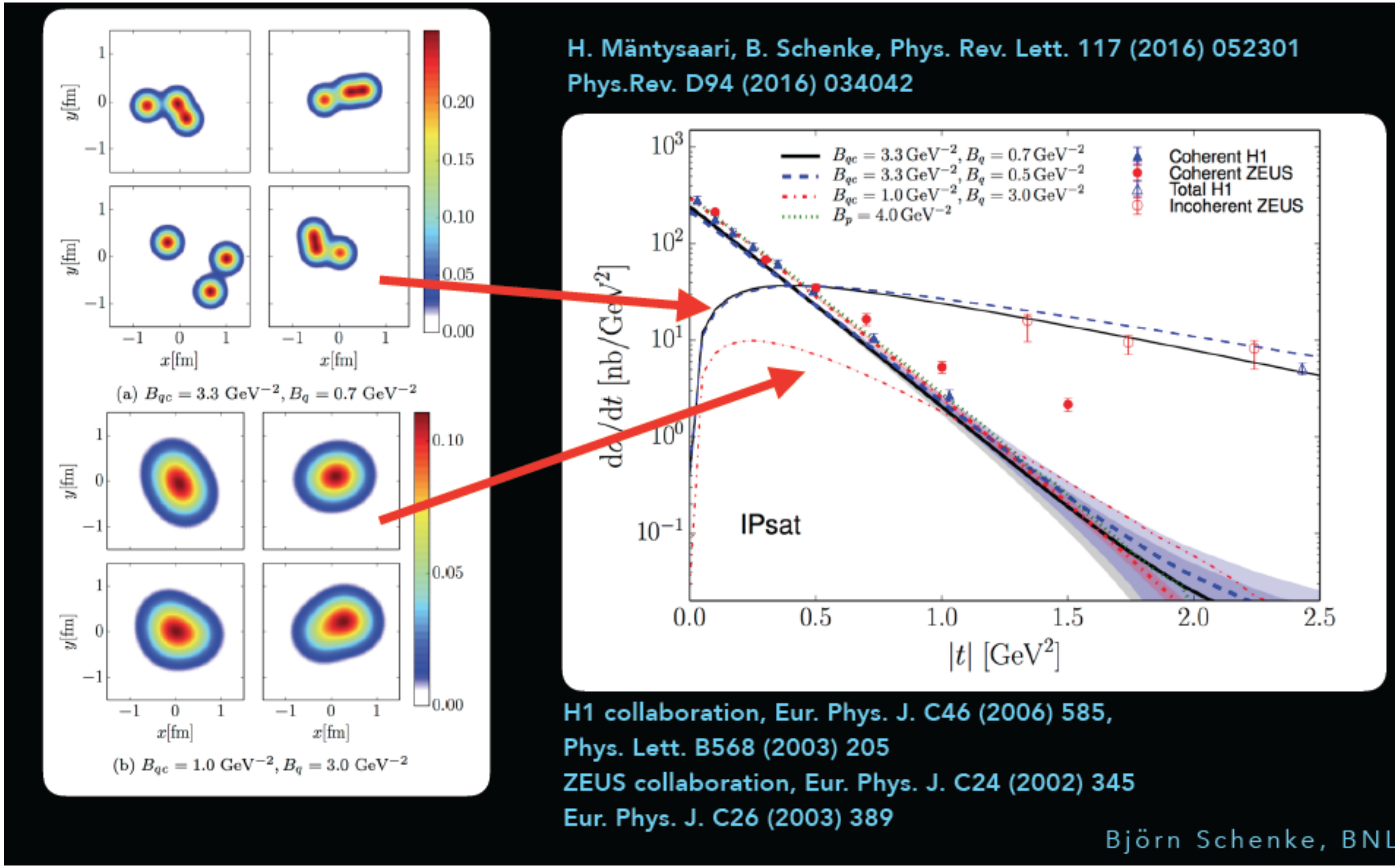


(charged hadron multiplicity in event)

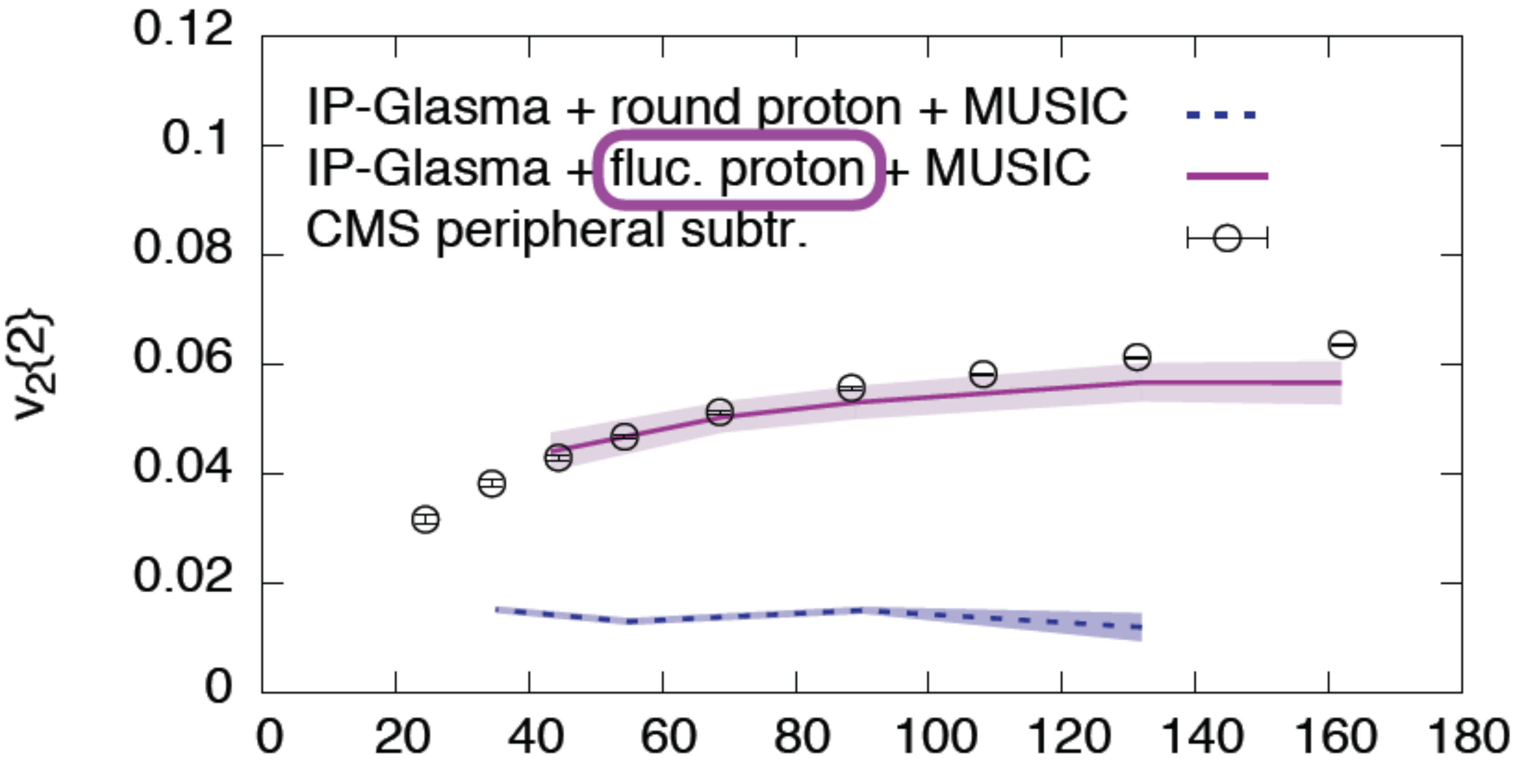
Schenke,Venugopalan,PRL113 (2014)102301

Shape fluctuations essential to generate flow-I

Incoherent exclusive vector meson (J/ψ) production is sensitive to fluctuations in transverse spatial profile of proton – compute in dipole picture



Shape fluctuations essential to generate flow-III



$\eta/s=0.2 + \zeta/s(T)$

$N_{trk}^{offline}$ **(charged hadron multiplicity in event)**

Mantysaari,Schenke,Shen,Tribedy,arXiv:1705.03177

Sensitive measure of collectivity: azimuthal cumulants

2m-particle azimuthal cumulants

Borghini,Dinh,Ollitrault, nucl-th/0105040

$$c_n\,\{2m\} = \langle\langle e^{in(\phi_1+\cdots\phi_m-\phi_{m+1}-\cdots\phi_{2m})}\rangle\rangle$$

If cumulants factorize into product of correlations relative to a reaction plane, define flow coefficients:

$$v_n\{2\}^2 \equiv c_n\{2\} \quad v_n\{4\}^4 \equiv -c_n\{4\} \quad v_n\{6\}^6 \equiv c_n\{6\}/4$$

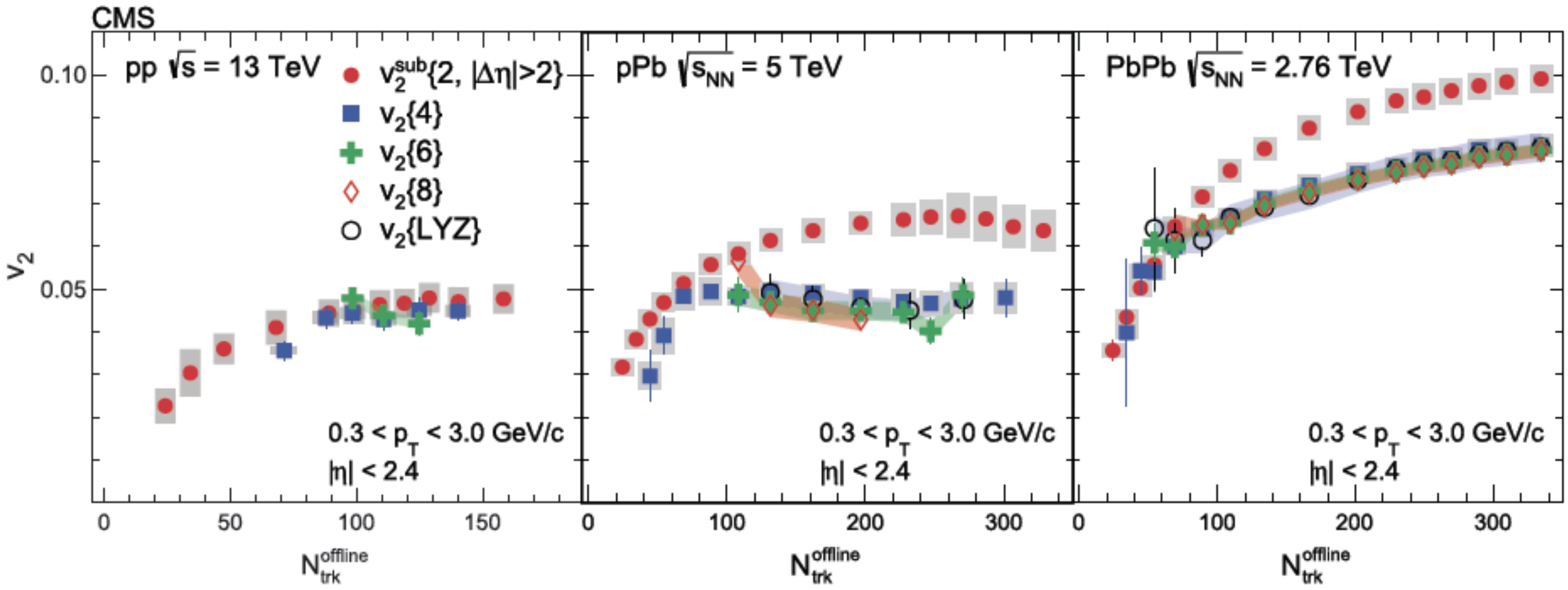
Spatial eccentricities: $\epsilon_n = \frac{1}{< r_\perp^n >}\int d^2r_\perp\, e^{in\phi_r}\, r_\perp^n\, \frac{dN}{dyd^2r_\perp}$

A number of simple "Gaussian" models give $\epsilon_n\{2\} > \epsilon\{4\} = \epsilon\{6\} = \cdots$

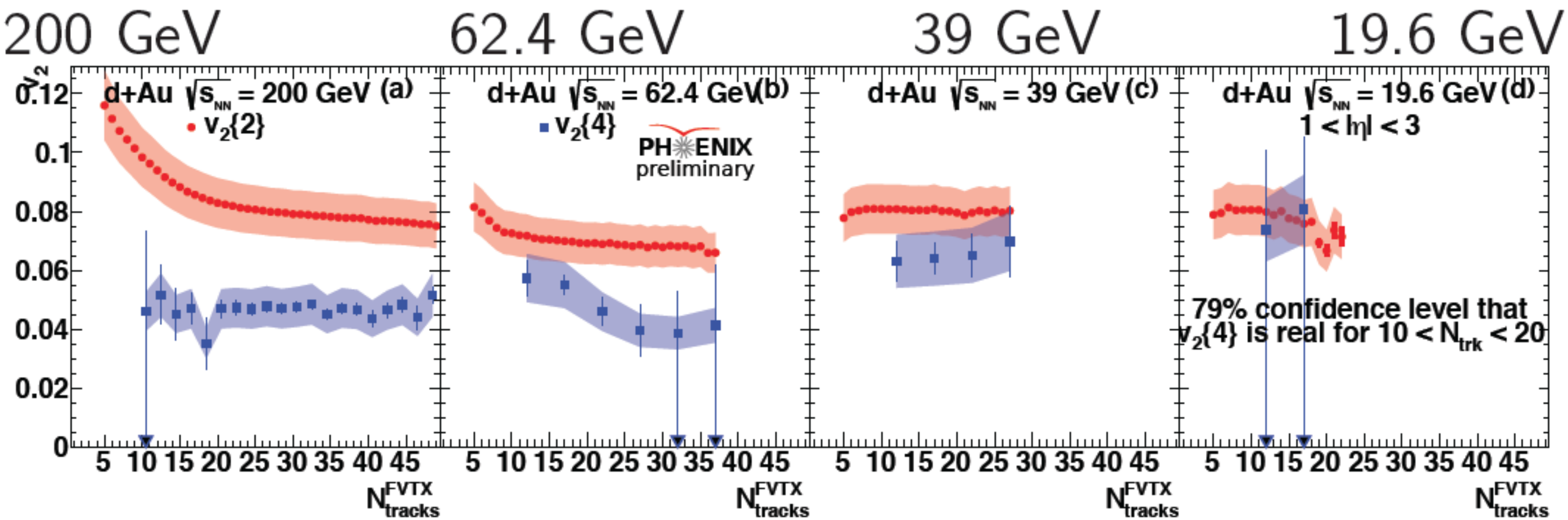
Hydro linear response: $v_n\{m\} \approx c_n\varepsilon_n\{m\}$

Gardim,Grassi,Luzum,Ollitrault, PRC (2012)024908; Niemi,Denicol,Holopainen,Huovinen, PRC87 (2013)054901
Bzdak,Bozek,McLerran, arXiv:1311.7325, Bzdak, Skokov, arXiv: 1312.7349
Yan, Ollitrault, arXiv:1312.6555, Basar,Teaney, arXiv:1312.6770

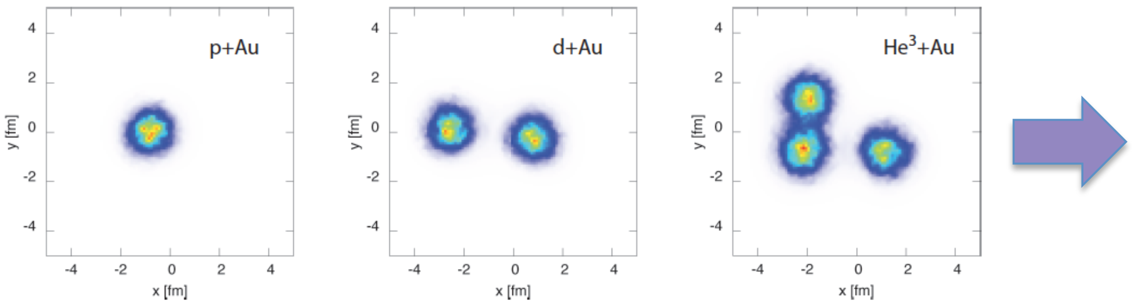
Collectivity across system size

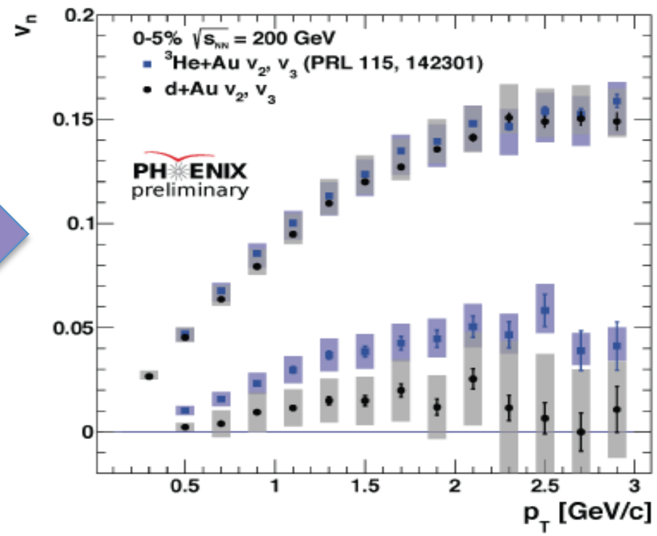


Collectivity across wide energy scales



Schenke, RV:1407.7557





Panta Rhei?

Heraclitus of Ephesus
535-475 BC

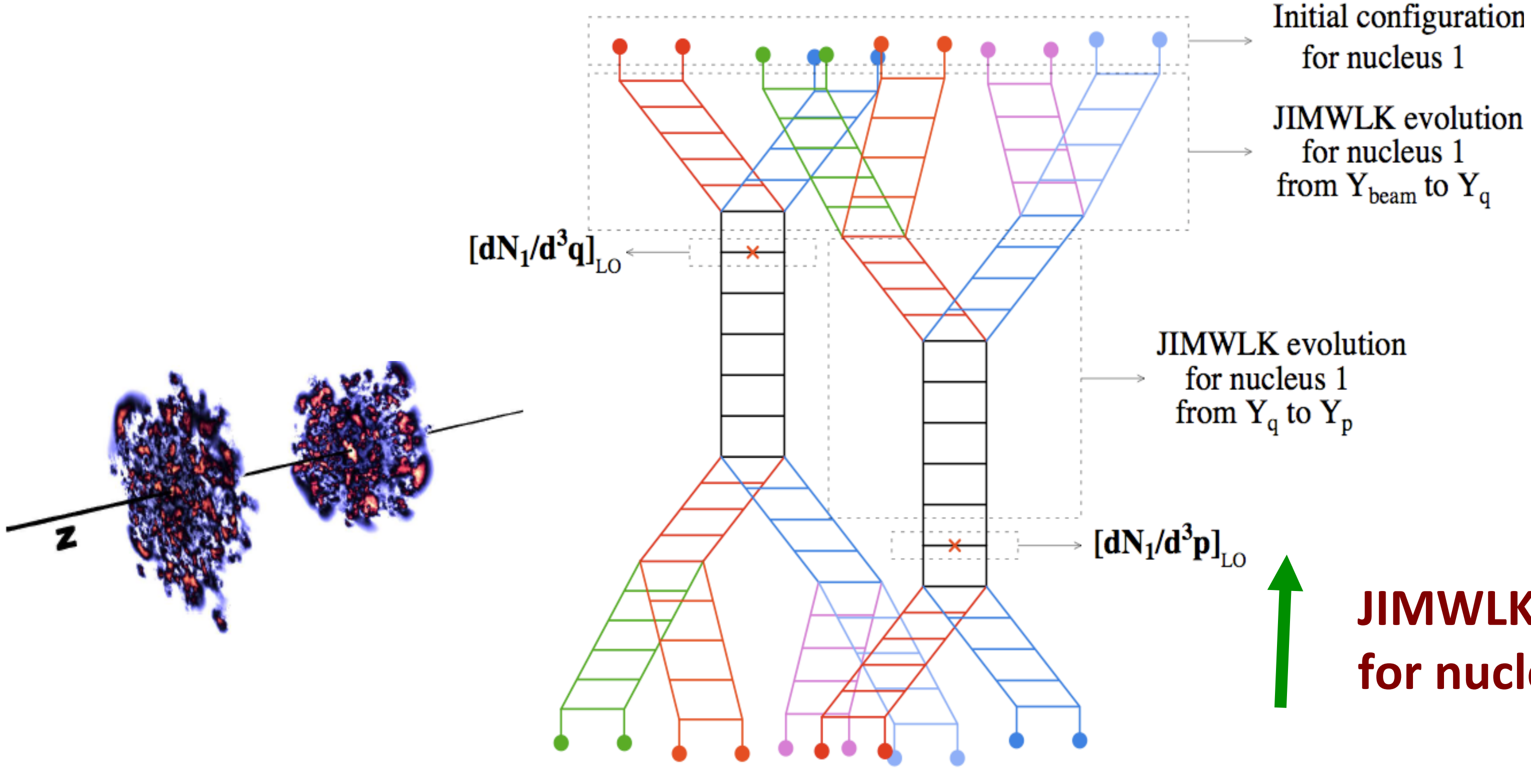
Natural in hydro – yet, very few ab initio hydro computations of 4-particle cumulants for p+A -- none for p+p

Problems with hydro interpretation for small systems:
i) Absence of “jet” quenching
ii) Lack of convergence of hydro expansion (large Knudsen numbers)
iii) Effects seen for small multiplicity and high p_T (9 GeV)

Can we understand multiparticle correlations in an *ab initio* approach ?

Review: Dusling,Li,Schenke, arXiv:1509.07939

Two-parton azimuthal correlations in the CGC



JIMWLK evolution for nucleus 2

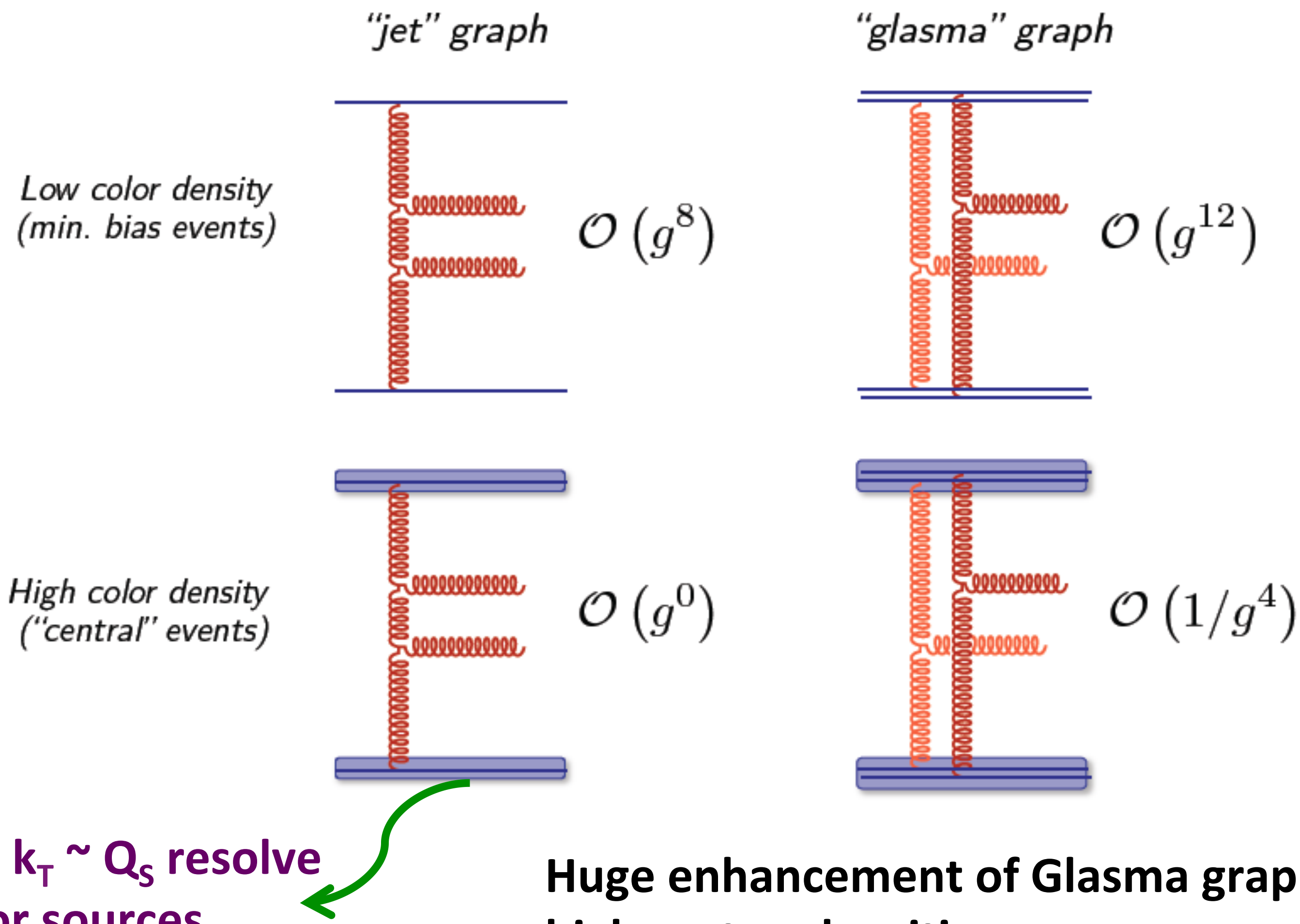
$$\langle \frac{dN_2}{d^3p\, d^3q} \rangle_{\rm LLogs} = \int [d\rho_1][d\rho_2] W_{Y_1}[\rho_1]\, W_{Y_2}[\rho_2]\, \frac{dN}{d^3p}|_{\rm LO}\, \frac{dN}{d^3q}|_{\rm LO}$$

JIMWLK evolution

Solution of Yang-Mills for dense-dense systems (as in A+A)

Dumitru,Gelis,McLerran,Venugopalan: 0804.3858
Gelis, Lappi, Venugopalan, arXiv: 0807.1306
Dusling,Gelis,Lappi,Venugopalan, arXiv:0911.2720

Glasma graph approximation: power counting

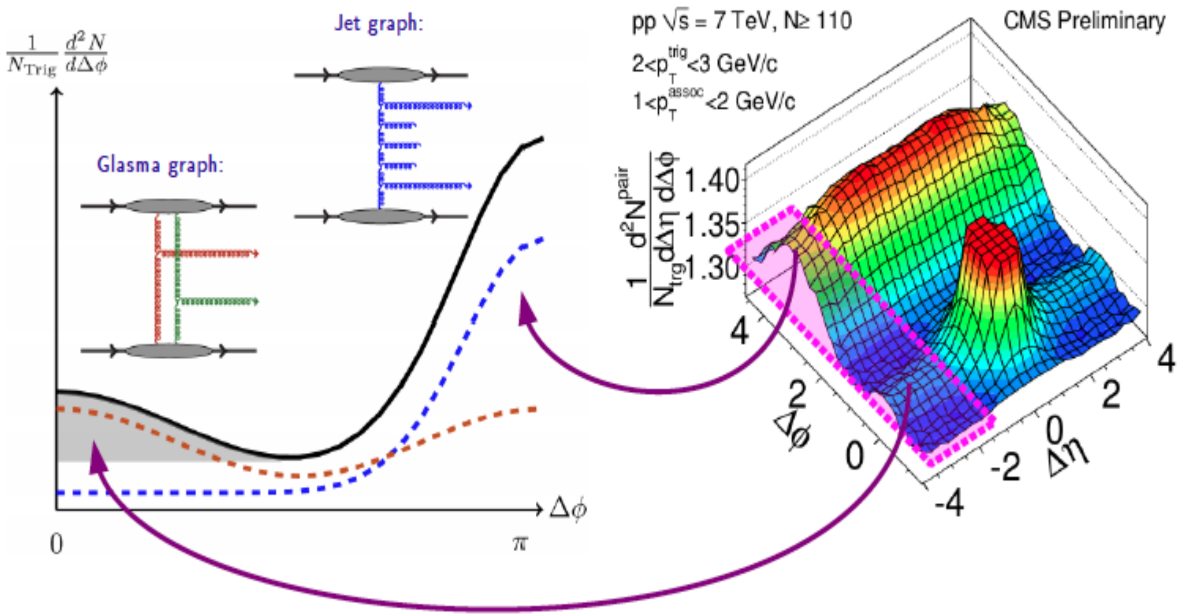


Gluons with $k_T \sim Q_S$ resolve
$n \sim 1/g^2$ color sources
Effective coupling: $g*n \sim 1/g$

Huge enhancement of Glasma graphs for high parton densities…

Dumitru,Dusling,Gelis,Jalilian-Marian,
Lappi,Venugopalan, PLB697 (2011)21
Dusling,Venugopalan,PRL108 (2012)262001

Anatomy of long range collimations



RG evolution of Glasma graphs:

$$C(\mathbf{p},\mathbf{q}) \propto \frac{g^4}{\mathbf{p}_\perp^2 \mathbf{q}_\perp^2} \int \mathrm{d}^2\mathbf{k}_{1\perp}\, \Phi_{A_1}^2(y_p,\mathbf{k}_{1\perp}) \Phi_{A_2}(y_p,\mathbf{p}_\perp - \mathbf{k}_{1\perp}) \Phi_{A_2}(y_q,\mathbf{q}_\perp - \mathbf{k}_{1\perp})$$
$$+ \quad \text{permutations}$$

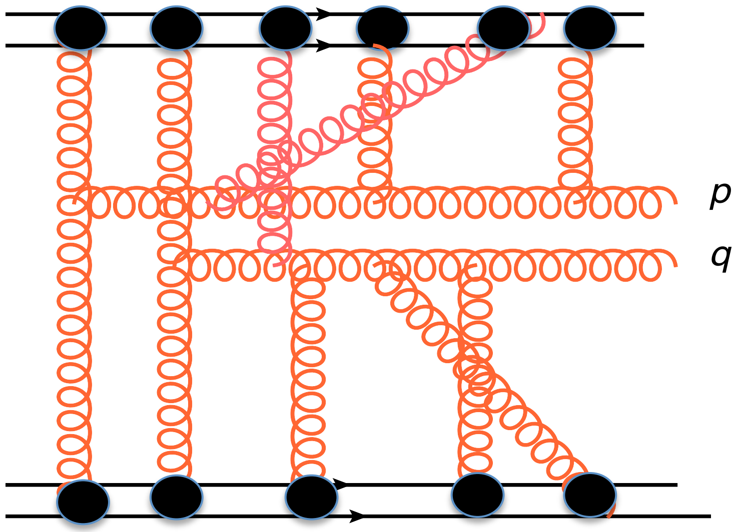
RG evolution of the mini-jets: $C_{\text{dijet}}(\mathbf{p},\mathbf{q}) \propto \Phi_A \otimes \Phi_B \otimes G_{\text{BFKL}}$

Good agreement with data for $p_T > Q_S$

However no odd harmonics v_3, v_5 for gluons because C(p,q) = C(p,-q)

Dusling,RV, PRD 87, 051502 (R) (2013); PRD87 (2013) 094034
Dusling,Tribedy,RV,PRD93 (2016) 014034

Beyond glasma graphs



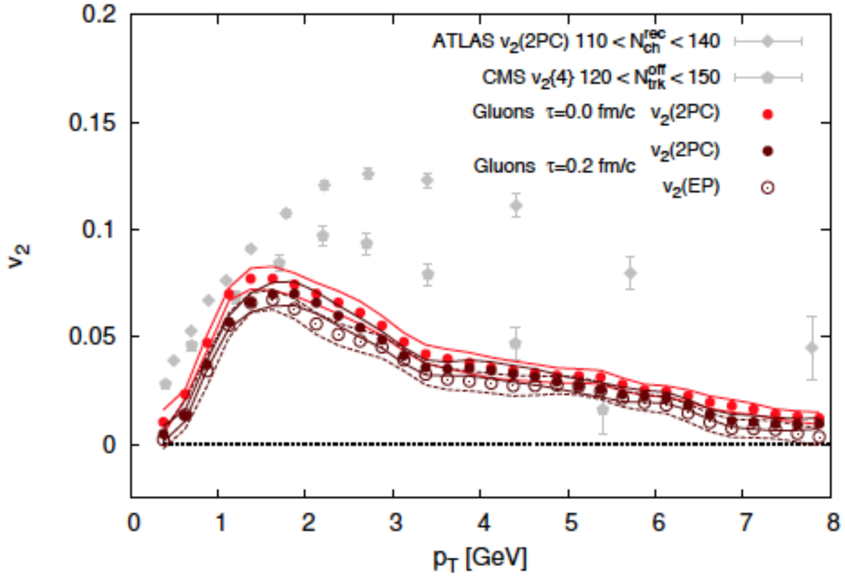
Coherent multiple scattering is of the same order in the coupling: power suppressed for $p_T >> Q_S$, important for $p_T < Q_S$

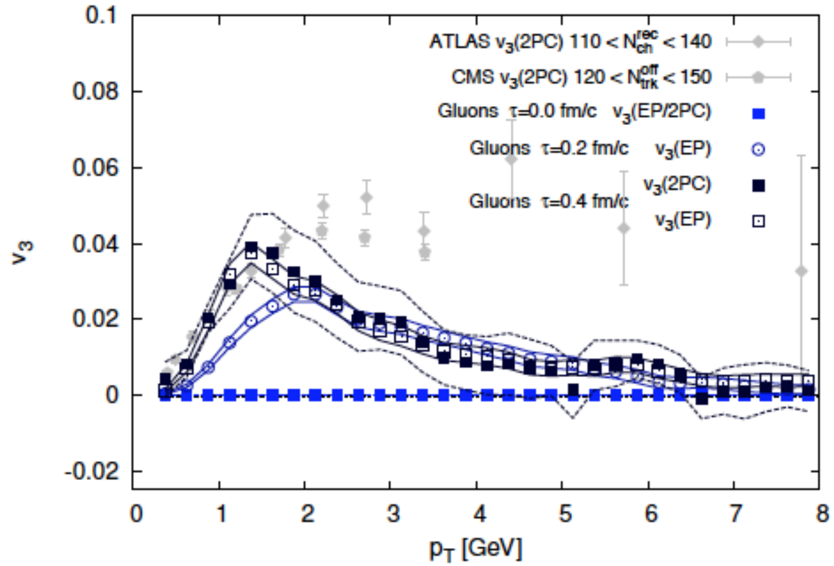
Compute (numerically) by solving Yang-Mills equations in presence of two light cone sources: IP-Glasma model Schenke,Tribedy,Venugopalan, arXiv:1202.6646, 1206.6805

$$\left\langle \frac{d^2 N}{d^2 p_T q_T} \right\rangle = \int D\rho_A \, D\rho_B \, e^{-\int d^2 x_T \rho_A^2 / Q_{s,A}^2} \, e^{-\int d^2 x_T \rho_B^2 / Q_{s,B}^2} \, \frac{dN}{d^2 p_T}[\rho_A, \rho_B] \, \frac{dN}{d^2 q_T}[\rho_A, \rho_B]$$

C(p,q) ≠ C(p,-q) -- all harmonics contribute Lappi,Srednyak,Venugopalan, arXiv:0911.2068

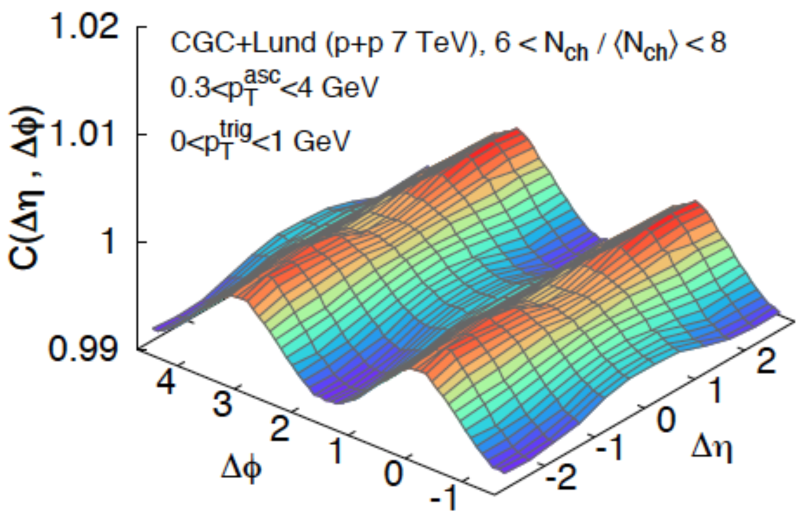
Azimuthal anistropy from Yang-Mills dynamics





Schenke,Schlichting,RV, PLB747(2015)76

Recent analytical work in dilute-dense approx:
Kovchegov,Wertepny,NPA906 (2013)50
McLerran, Skokov arXiv:1611.09870
Kovner,Lublinsky,Skokov, arXiv:1612.07790



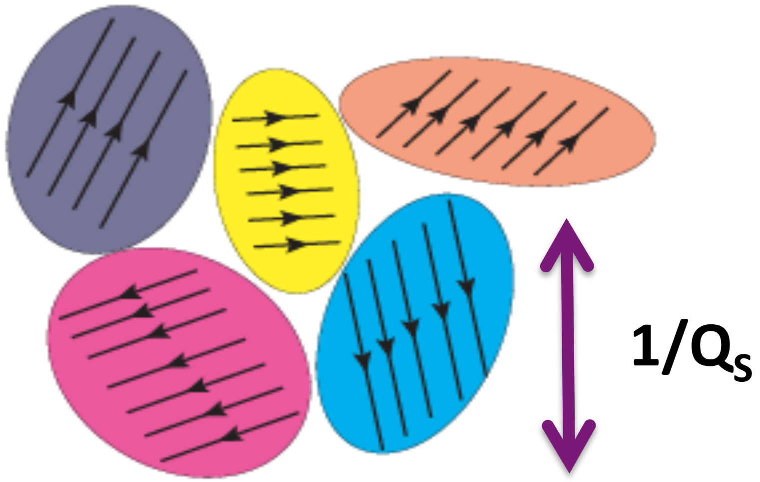
Combining gluon distributions from CGC with PYTHIA reproduces ridge distributions

However 4-particle collectivity is computationally challenging at present

Tracing azimuthal initial state correlations

Simple ab initio initial state model:
Multi-particle correlations from Eikonal scattering of partons off color domains in a nuclear target

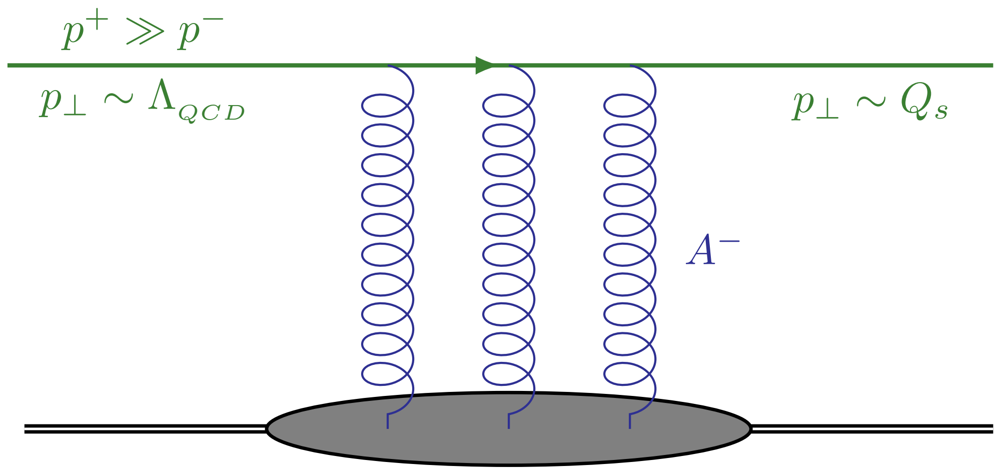
Lappi, arXiv:1501.05505
Lappi,Schenke,Schlichting,RV, arXiv:1509.03499
Dusling,Mace,RV, arXiv:1705.00745



Color domain model:

Kovner,Lublinsky,arXiv:1012.3398,1109.0347
Dumitru,Gianini, arXiv:1406.5781
Dumitru,Skokov,arXiv:1411.6630,
Dumitru,McLerran,Skokov,arXiv:1410.4844

Eikonal scattering: the parton model



Color rotation of parton in external field by a lightlike Wilson line

$$W[A](x) = \mathcal{P}\exp\left[ig\int dz^+ A_a^-(z^+, x)\right]$$

Parton distribution after coherent multiple scattering off nucleus:

$$\frac{dN_q}{d^2p} = \frac{1}{\pi^2}\int d^2b \int \frac{d^2k}{(2\pi)^2}\int d^2r \underbrace{e^{-b^2/B}\, e^{-k^2 B}}_{\text{Wigner function}} \left\langle \underbrace{D\left(b + r/2, b - r/2\right)}_{\text{Dipole correlator}} \right\rangle$$

$$D(x, \bar{x}) = \frac{1}{N_c}\mathrm{Tr}\left[W(x)W^\dagger(\bar{x})\right]$$

B is transverse area of proton

Bjorken, Kogut, Soper, Phys. Rev., D3:1382, (1971)
Dumitru, Jalilian-Marian, Phys. Rev. Lett., 89:022301, (2002)

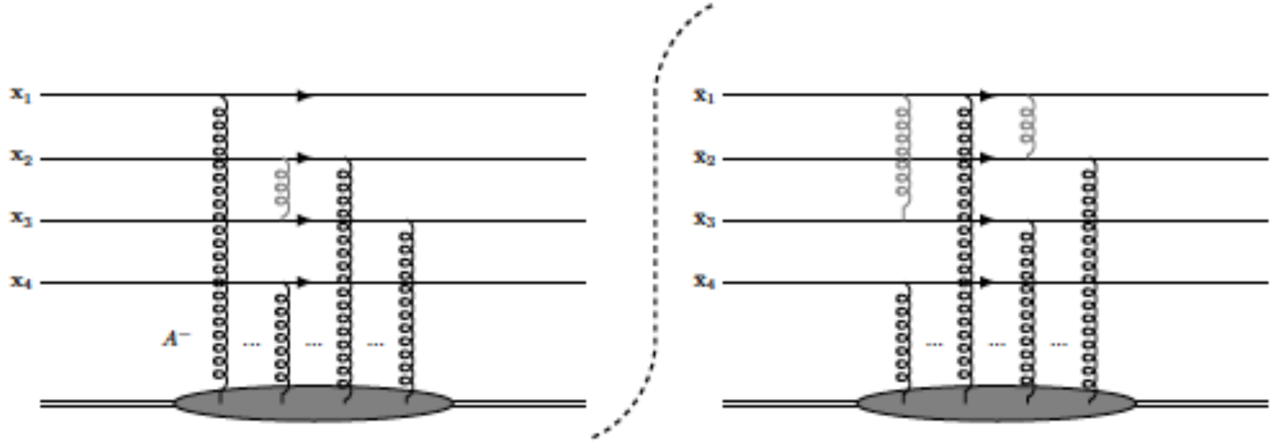
Multiparton Eikonal scattering

Two partons:

$$\frac{d^2N}{d^2p_1 d^2p_2} \simeq \frac{1}{(\pi B_p)^2} \int_{x\bar{x}y\bar{y}} e^{-(x^2+\bar{x}^2)/2B_p} \; e^{-(y^2+\bar{y}^2)/2B_p} e^{ip_1\cdot(x-\bar{x})} e^{ip_2\cdot(y-\bar{y})}$$

$$\times \left\langle \frac{1}{N_c} \mathrm{Tr}\left[W(x)W^\dagger(\bar{x})\right] \frac{1}{N_c}\mathrm{Tr}\left[W(y)W^\dagger(\bar{y})\right] \right\rangle \;\; \propto \langle D\, D \rangle$$

Four parton



$$d^4N \sim \int \langle \mathrm{Tr}\left[W(w)W^\dagger(\bar{w})\right] \mathrm{Tr}\left[W(x)W^\dagger(\bar{x})\right] \mathrm{Tr}\left[W(y)W^\dagger(\bar{y})\right] \mathrm{Tr}\left[W(z)W^\dagger(\bar{z})\right] \rangle$$

$\propto \langle D\, D\, D\, D \rangle$ **and so on** ...

Oversimplification: n-particle Wigner distributions factorize

$$W_{q^n}(\mathbf{b_1}, \mathbf{k_1}, ..., \mathbf{b_n}, \mathbf{k_n}) = W_q(\mathbf{b_1}, \mathbf{k_1}) \cdot ... \cdot W_q(\mathbf{b_n}, \mathbf{k_n})$$

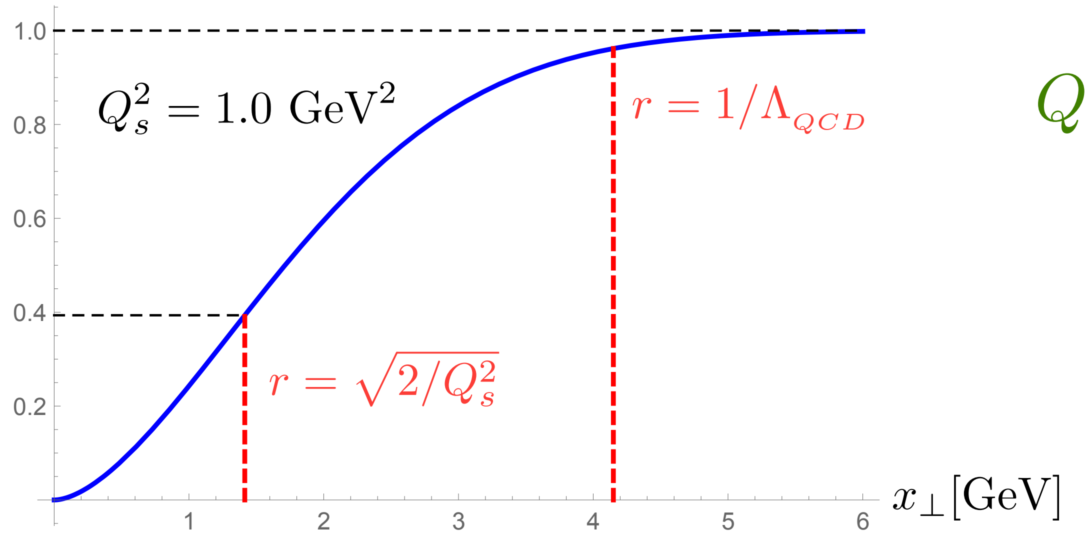
Averaging over color representations in the target

Dipole correlator is evaluated in the MV model where color correlations in the target are Gaussian (random walk in color)

$$g^2 \left\langle A_a^-(x) A_b^-(y) \right\rangle = \delta^{ab} L_{xy} \qquad L_{xy} = -\frac{g^4 \mu^2}{16\pi} |x-y|^2 \ln \frac{1}{\Lambda |x-y|}$$

gives $D(x,\bar{x}) = \frac{1}{N_c} \mathrm{Tr}\left[W(x) W^\dagger(\bar{x})\right] = \exp(C_F L_{x\bar{x}})$

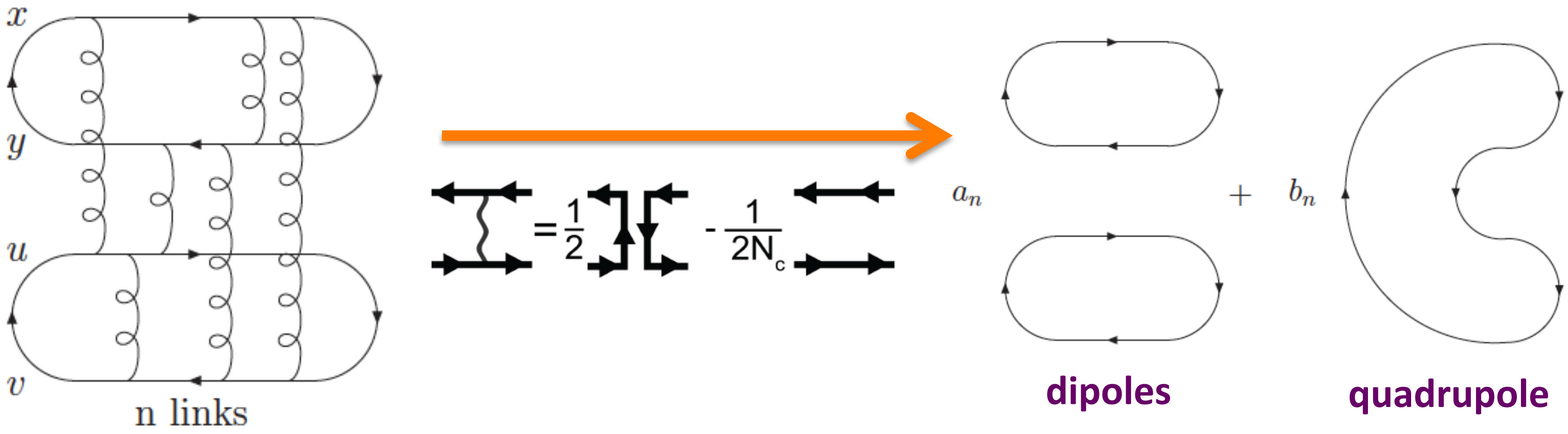
$N(x_\perp) = 1 - D(x_\perp)$



$Q_s \propto g^2 \mu$

Averaging over multi-point dipole correlators

To compute the 2-dipole correlator:



Use $W(x) \equiv \mathcal{P} \exp\left[ig \int dz^+ A_a^-(z^+, x)\right] \simeq V(x)\left[1 + igA_a^-(\xi, x)T^a + \cdots\right]$

This gives $\langle D_{x\bar{x}} D_{y\bar{y}} \rangle_W \simeq \alpha_{x\bar{x}y\bar{y}} \langle D_{x\bar{x}} D_{y\bar{y}} \rangle_V + \beta_{xy\bar{x}\bar{y}} \langle Q_{x\bar{y}y\bar{x}} \rangle_V$

Equivalently,

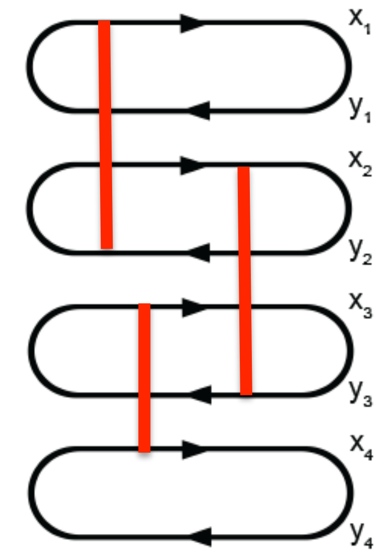
$$\begin{pmatrix} \langle D_{x\bar{x}} D_{y\bar{y}} \rangle \\ \langle Q_{x\bar{y}y\bar{x}} \rangle \end{pmatrix}_W = \begin{pmatrix} \alpha_{x\bar{x}y\bar{y}} & \beta_{xy\bar{x}\bar{y}} \\ \beta_{xy\bar{y}\bar{x}} & \alpha_{x\bar{y}y\bar{x}} \end{pmatrix} \begin{pmatrix} \langle D_{x\bar{x}} D_{y\bar{y}} \rangle \\ \langle Q_{x\bar{y}y\bar{x}} \rangle \end{pmatrix}_V$$

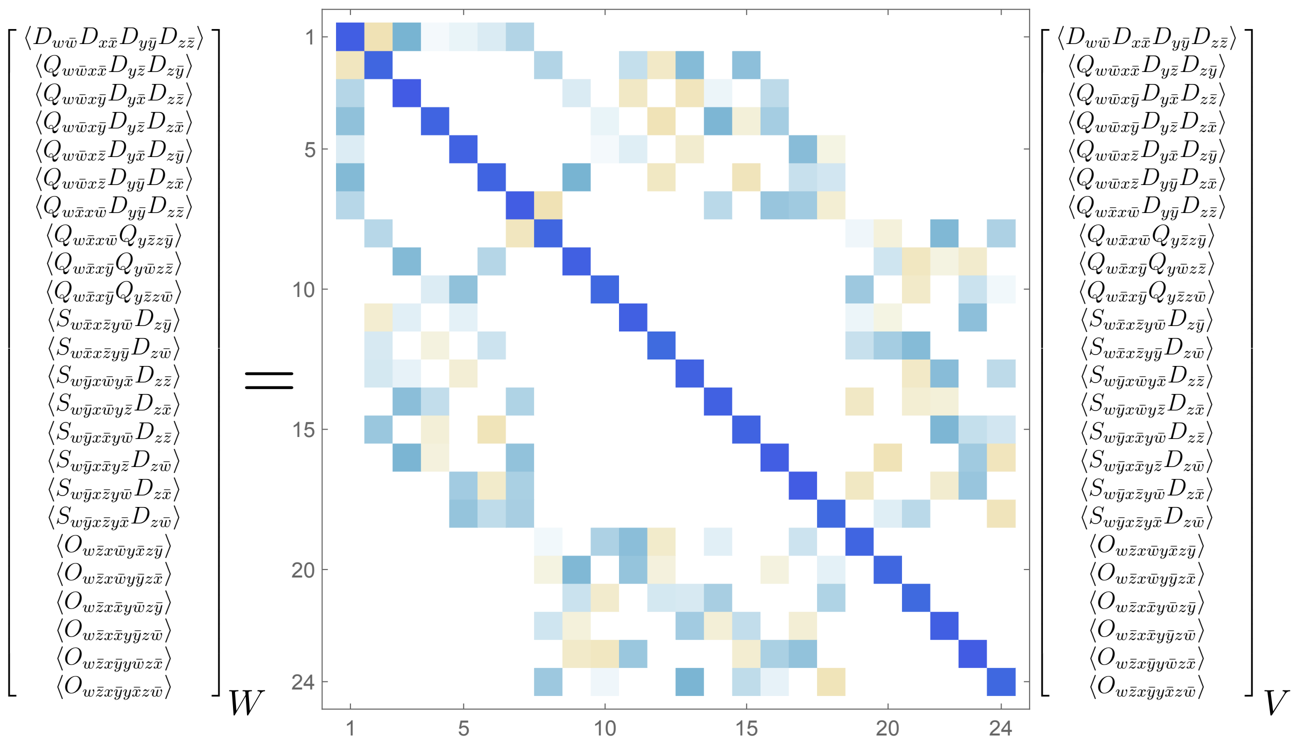
Kovner, Wiedemann, Phys. Rev., D64:114002, (2001)
Fujii, Nucl. Phys., A709:236 (2002).
Blaizot, Gelis, Venugopalan. Nucl. Phys., A743:57, (2004)
Dominguez, Marquet, Wu, Nucl. Phys., A823:99, (2009)

Iterate, diagonalize,exponentiate, to compute correlator

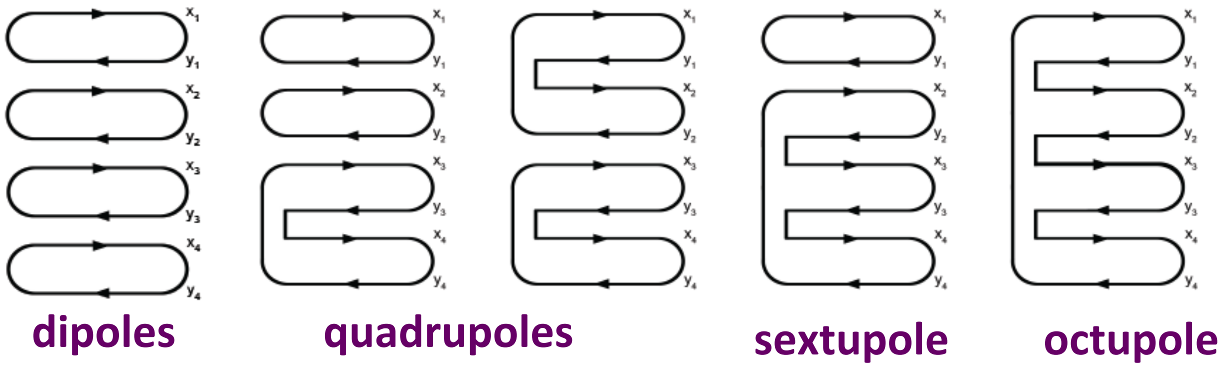
Averaging over multi-point dipole correlators

4-dipole correlator:





To compute n-gluon exchange, diagonalize 24×24 matrix and exponentiate



Dusling,Mace,Venugopalan,arXiv:1705.00745

Results for azimuthal anisotropies from multiparton eikonal scattering

Dusling,Mace,Venugopalan,arXiv:1705.00745
Dusling,Mace,Venugopalan, to appear next week

Objects to be computed

N-particle distributions:

$$\frac{d^n N}{d^2p \cdots} \sim \int \overbrace{e^{-(x^2+\bar{x}^2)/2B_p \cdots}}^{B_p = 4\ \mathrm{GeV}^{-2}} \left\langle \overbrace{\frac{1}{N_c}\mathrm{Tr}\left[W(x)W^\dagger(\bar{x})\right] \cdots}^{Q_s^2 \sim 1-3\ \mathrm{GeV}^2} \right\rangle e^{ip\cdot(x-\bar{x})\cdots}$$

Two-particle cumulants:

$$c_n\{2\} = \frac{\kappa_n\{2\}}{\kappa_0\{2\}}, \quad v_n\{2\} = \sqrt{c_n\{2\}}$$

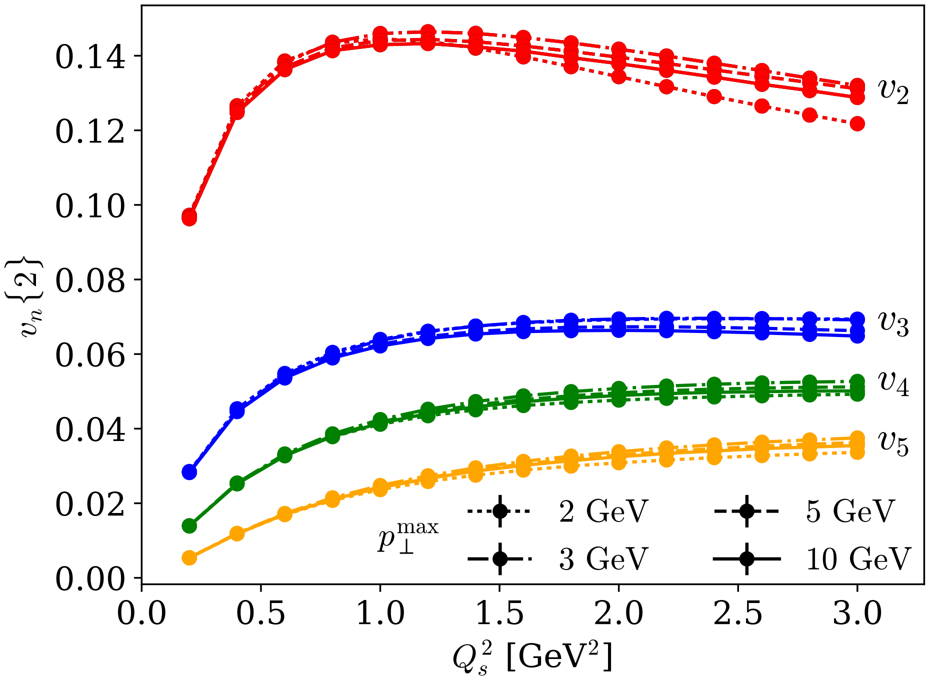
$$\kappa_n\{2\} = \int d^2p_1 d^2p_2 \cos[n(\phi_{p1} - \phi_{p2})] \frac{d^2N}{d^2p_1 d^2p_2}$$

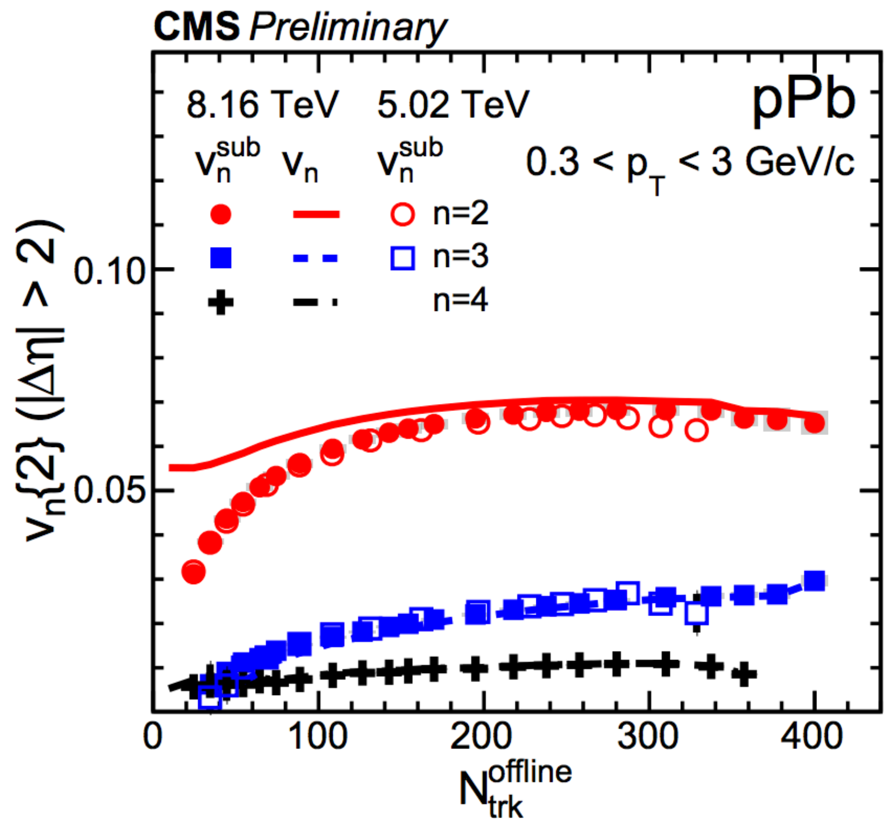
Four-particle cumulants:

$$c_n\{4\} = \frac{\kappa_n\{4\}}{\kappa_0\{4\}} - 2\left(\frac{\kappa_n\{2\}}{\kappa_0\{0\}}\right)^2, \qquad v_n\{4\} = (-c_n\{4\})^{1/4}$$

$$\kappa_n\{4\} = \int d^2p_1 d^2p_2 d^2p_3 d^2p_4 \cos[n(\phi_{p1} + \phi_{p2} - \phi_{p3} - \phi_{p4})] \frac{d^4N}{d^2p_1 d^2p_2 d^2p_3 d^2p_4}$$

Integrated anisotropy coefficients



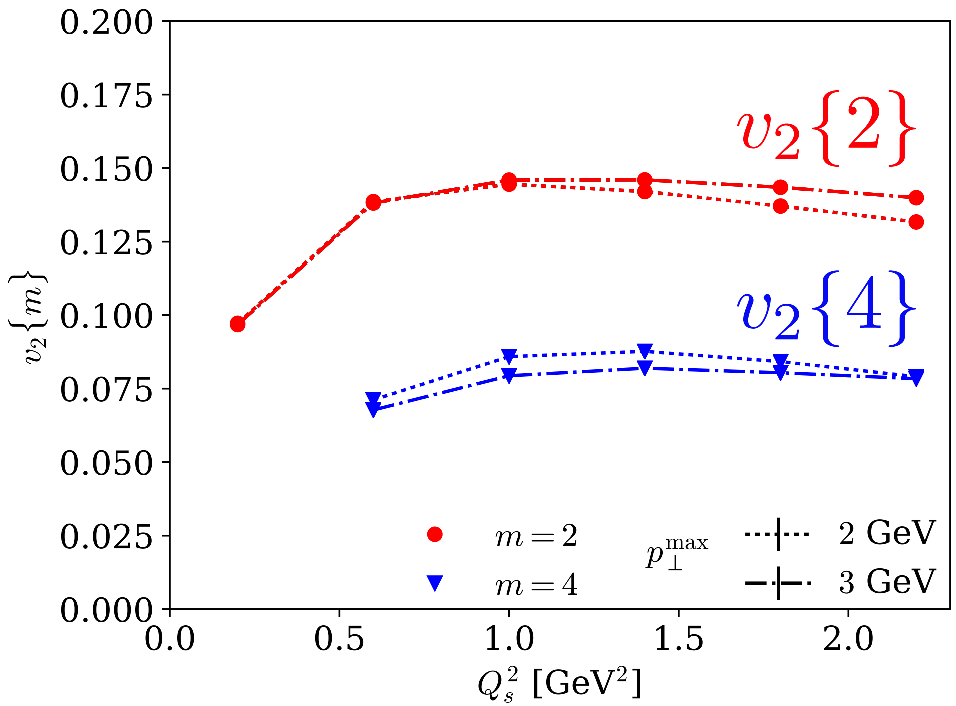


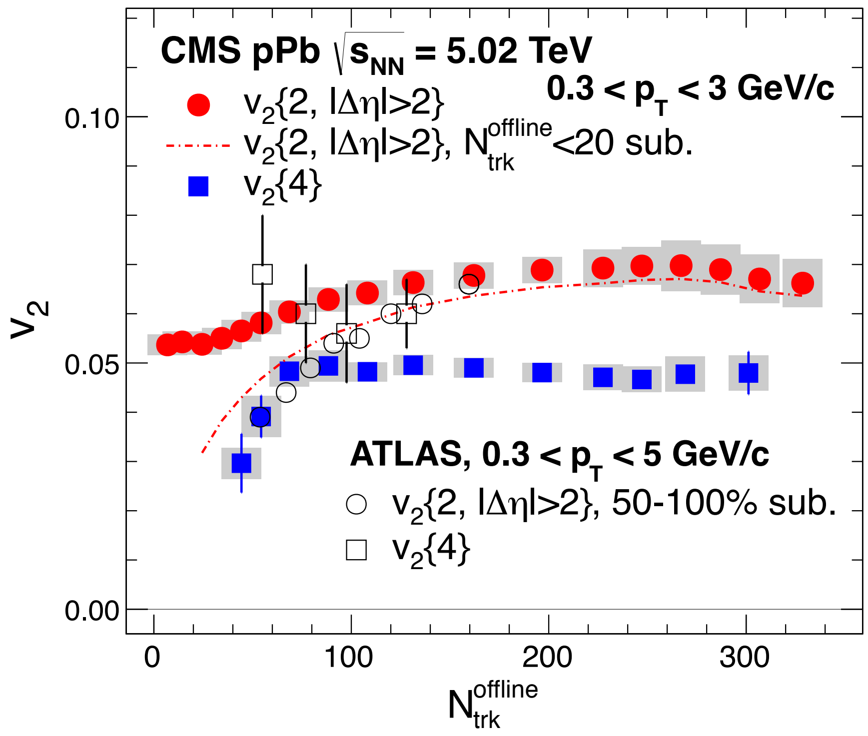
Similar ordering of "Flow" coefficients as seen in the data

Important caveat: No simple map between theory and experiment
Theory results are for quarks, Q_s^2 is the saturation scale in the target

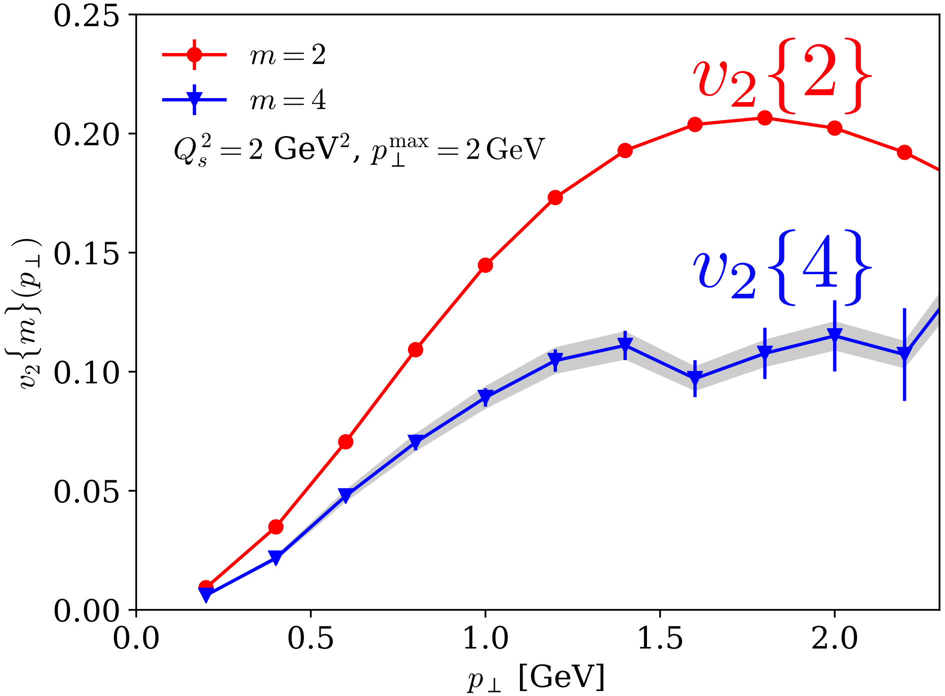
Lappi, Phys. Lett., B744:315, (2015)
Lappi, Schenke, Schlichting, Venugopalan. JHEP, 01:061, (2016)

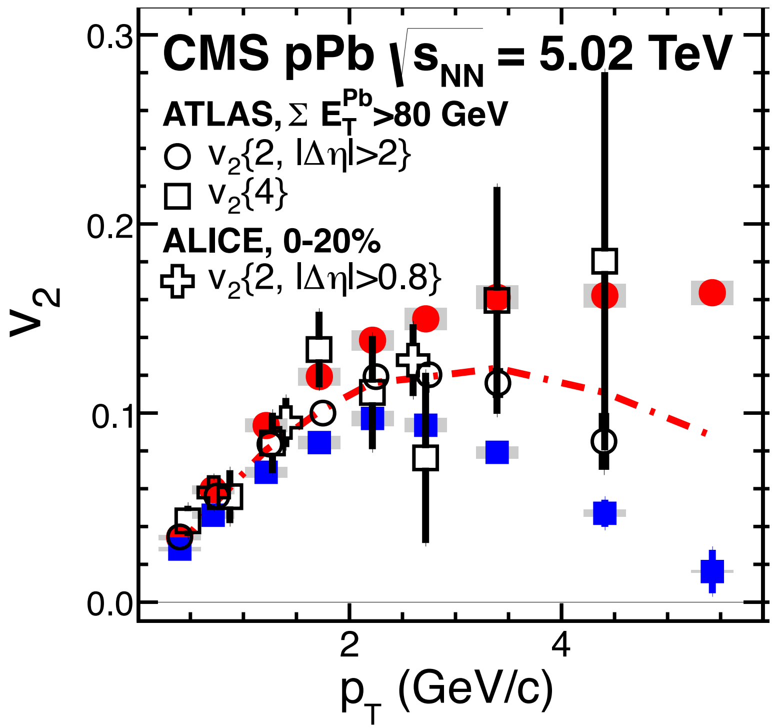
Integrated anisotropy coefficients



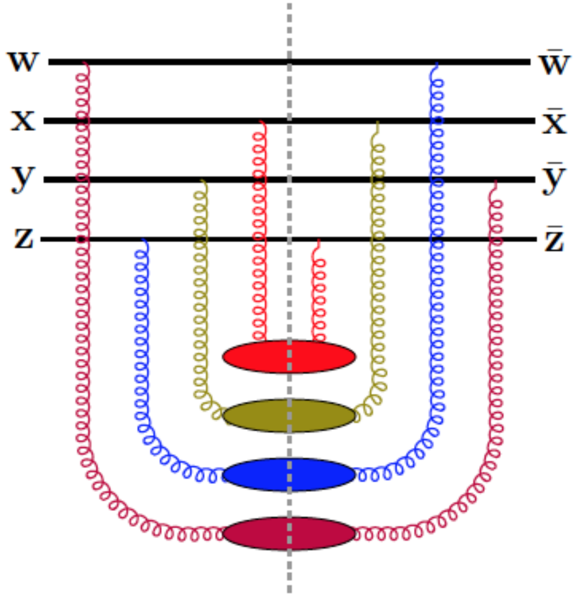


p_T dependence of anisotropy coefficients



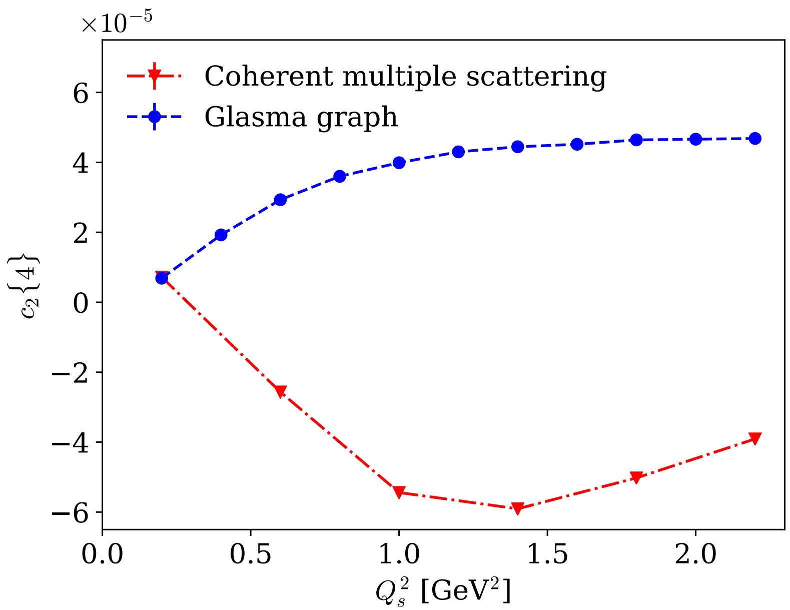


Back to glasma graphs



Glasma graph (single scattering) correlations are very strong – the n-particle distribution is close to a *Bose distribution* – *as in a laser*

Gelis,Lappi,McLerran, arXiv:0905.3234

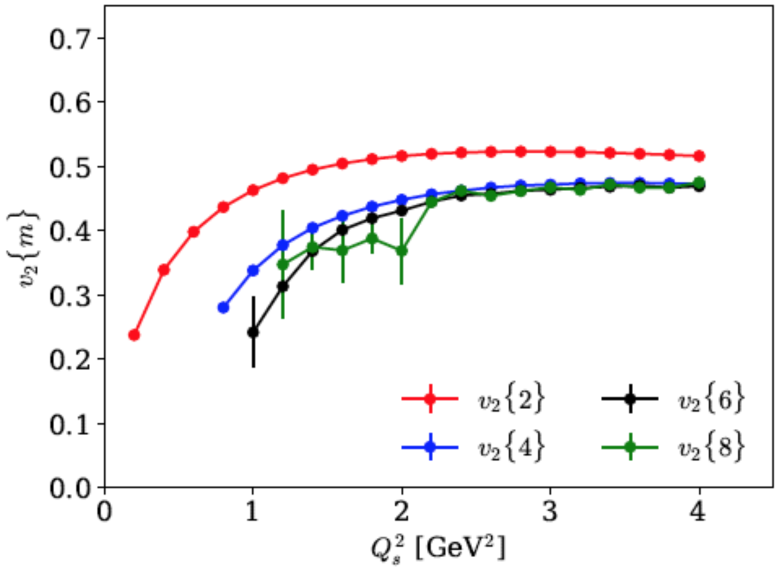


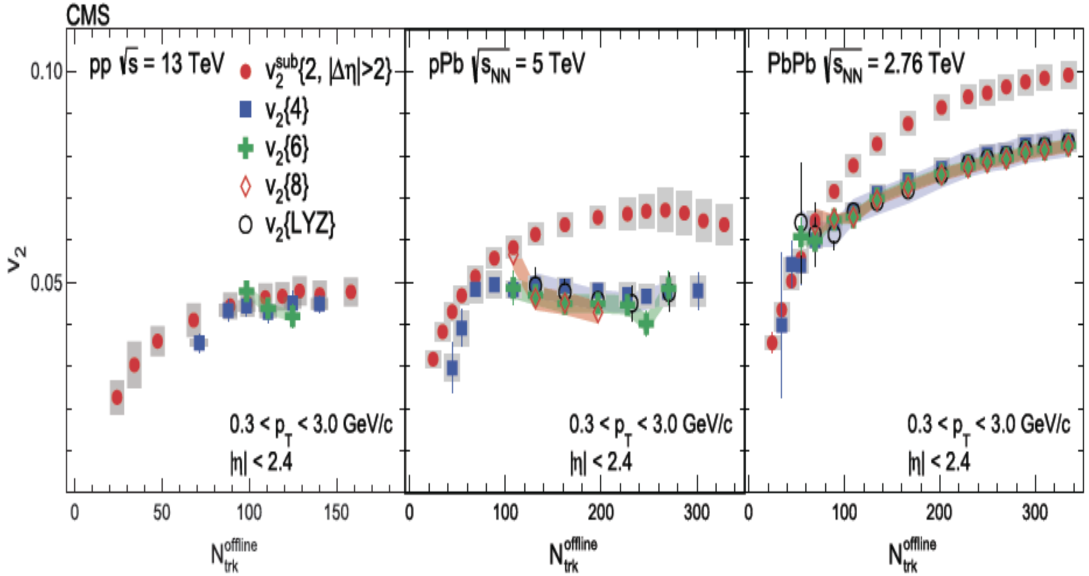
But $v_2\{4\}$ is imaginary…

For real $v_2\{4\}$, must have dominance of first two moments of distribution
-- achieved by coherent multiple scattering…

Dusling,Mace,Venugopalan,to appear

Higher cumulants from scattering off coherent Abelian fields

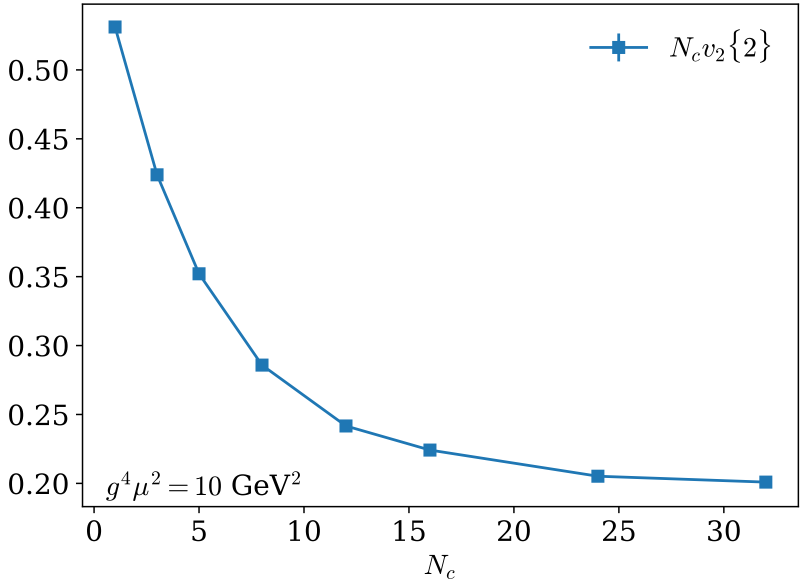


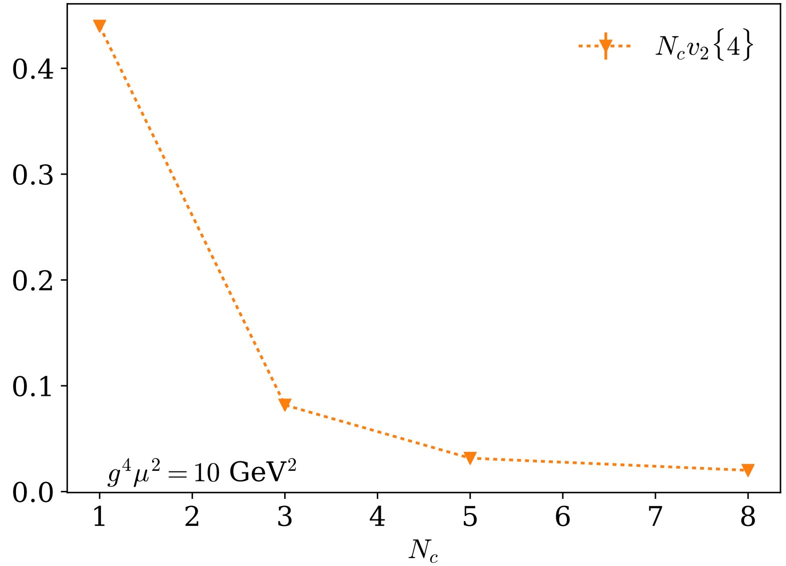


Replace $N_C \times N_C$ trace with simple path ordered exponentials (N_C=1)

2m-particle collectivity reproduced in this simple parton model…

Nc scaling of anisotropy coefficients





Both $v_2\{2\}$ and $v_2\{4\}$ scale as $1/N_C$ for large N_C

Summary-I

Hydrodynamic paradigm appears to describe multi-particle correlations even in the smallest systems

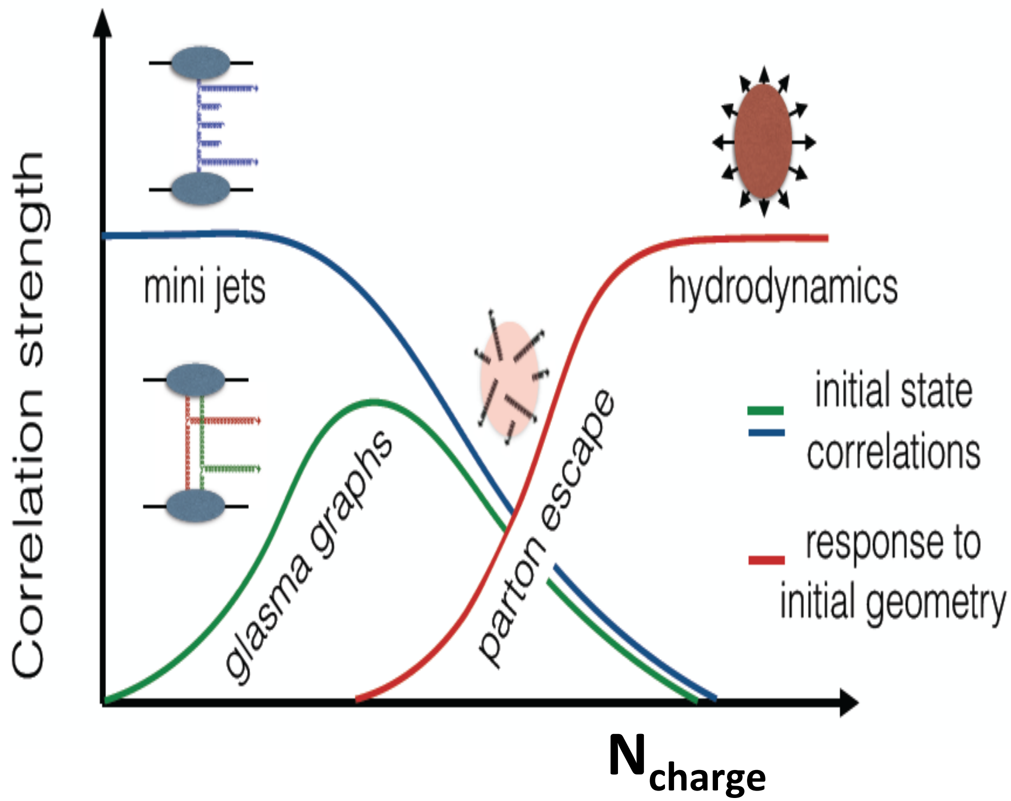
There are however puzzling features of the data,

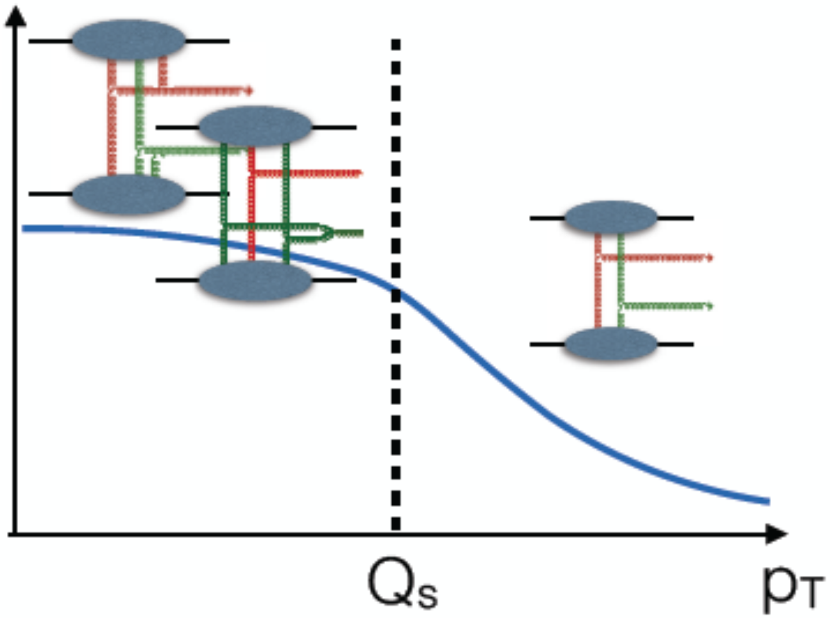
questions about the the validity of hydro, fine tuning of initial conditions (requiring implicitly strong initial state correlations), absence of jet quenching,··· and explanation of anisotropies for p_T > few GeV

Initial state QCD frameworks now also able to explain many features of the data but systematic treatments are still in their infancy

Despite much progress, no completely satisfactory explanation of the data -- the problem is still wide open

Summary-II





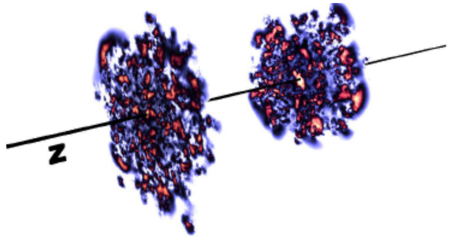
Event engineering across system sizes, energies, and varieties of probes:
Offers exciting possibility of exploring dynamical evolution of strongly correlated quark-gluon matter from high occupancy, out of equilibrium, dynamics…
to hydrodynamics

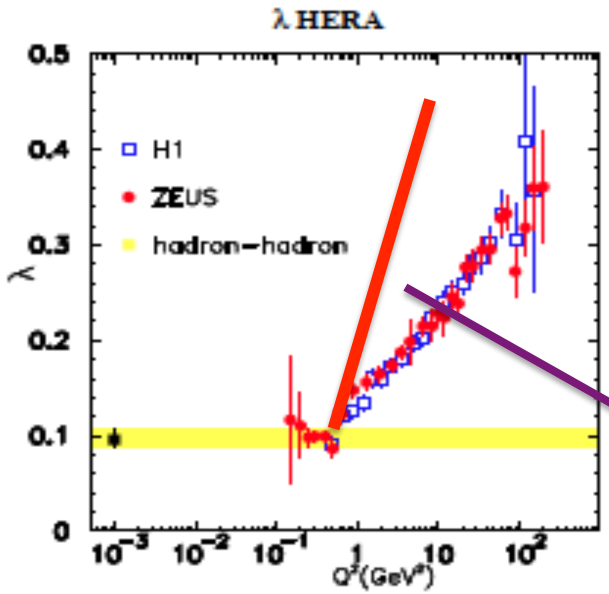
Figures: S. Schlichting at Quark Matter 2015

Thanks for listening!

High multiplicities + small systems = gluon saturation

What does it take to produce~ 150 hadrons per 5 units of rapidity in a single p+p event ?





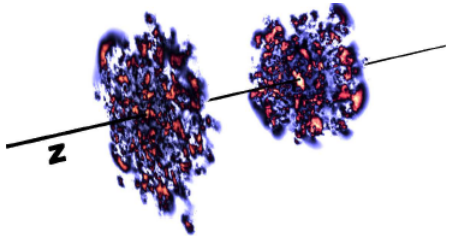
For λ=0.14, get about 13 gluons produced in 5 units ~ min.bias hadron multiplicity

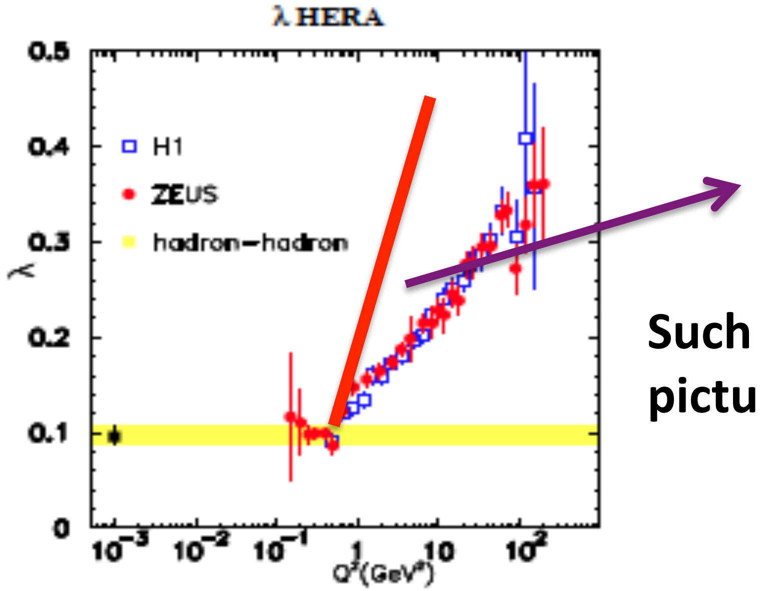
λ=0.3: ~45 gluons in 5 units,
λ=0.4: ~90 gluons in 5 units, in ball park…

Very rapid growth of gluon dist. in such events…

High multiplicities + small systems = gluon saturation

What does it take to produce~ 150 hadrons per 5 units of rapidity in a single p+p event ?



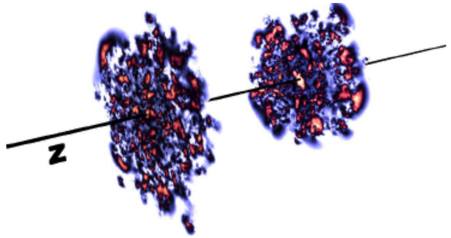


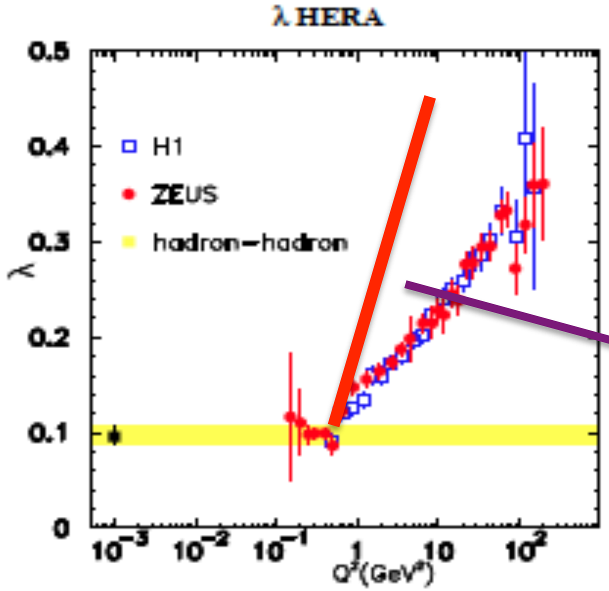
$$\frac{4\pi}{Q^2} * N_g(Q^2) = \pi R^2_{\text{glue}}$$

Such rapid growth, in an "independent parton" picture lead to very large gluon radii, R_g > 1 fm

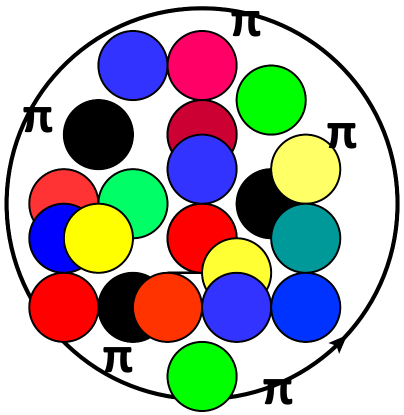
High multiplicities + small systems = gluon saturation

What does it take to produce~ 150 hadrons per 5 units of rapidity in a single p+p event ?





Saturation regulates this by adding increasingly "smaller" gluons of size 1/Q_s(x) with decreasing x (increasing energy)



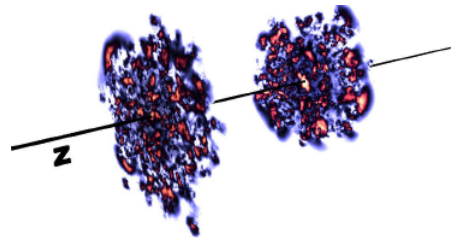
$$\frac{dN_g^{\text{prot.}}}{d\eta} \approx \frac{1.1 C_F}{2\pi^2} \frac{S_\perp Q^2_{\text{S,prot.}}}{\alpha_S}$$

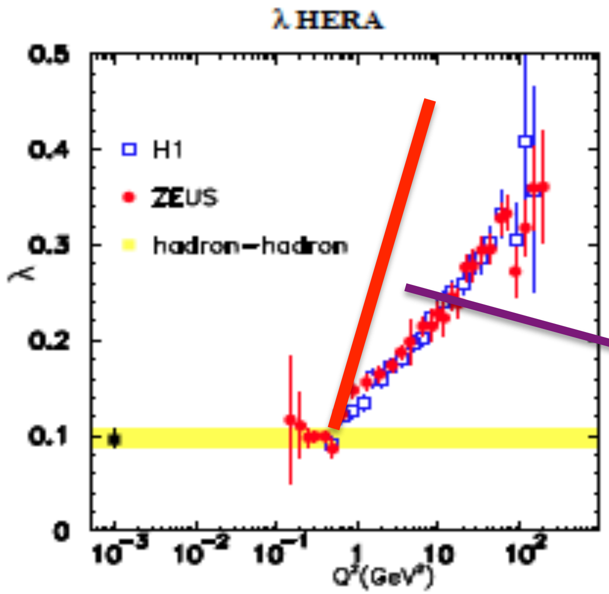
Lappi: arXiv 0711.3039

N_g ~ 100 in 5 units for Q_S^2 ~ 2 GeV^2 : a semi-hard scale !

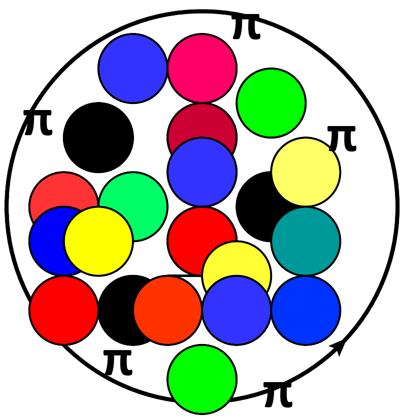
High multiplicities + small systems = gluon saturation

What does it take to produce~ 150 hadrons per 5 units of rapidity in a single p+p event ?



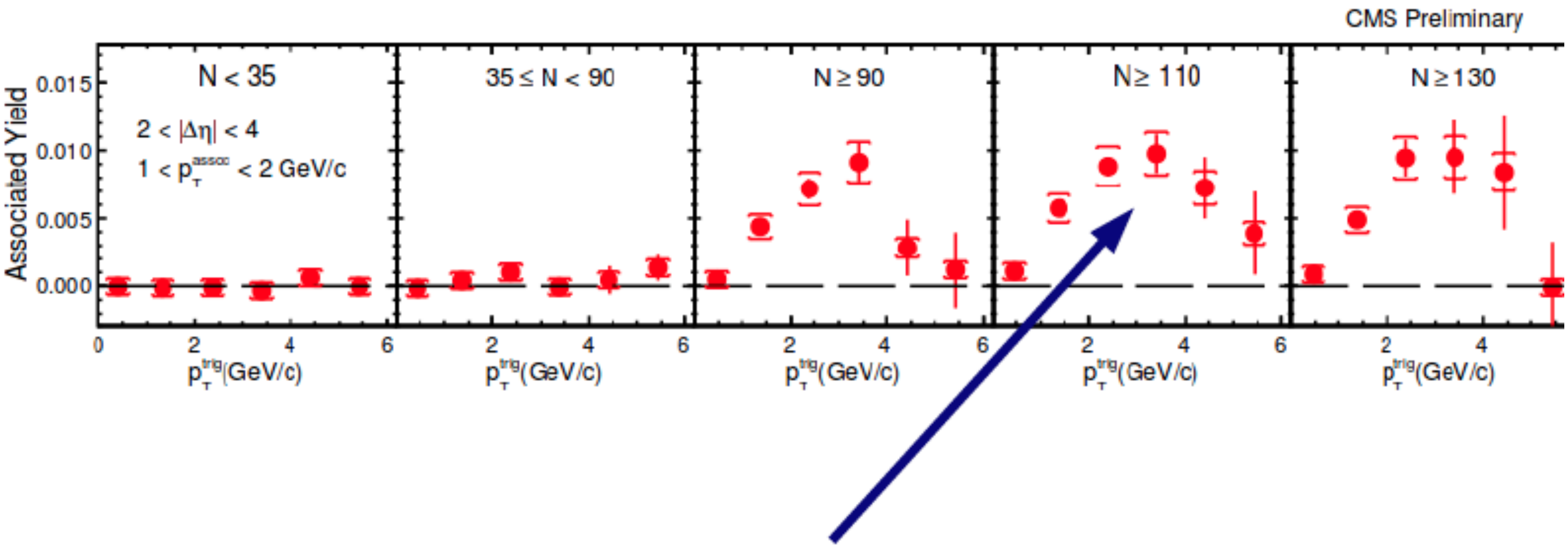


Saturation regulates this by adding increasingly "smaller" gluons of size $1/Q_s(x)$ with decreasing x (increasing energy)



Event generators (such as EPOS) that describe data in high multiplicity events, build in a saturation scale …

The ridge



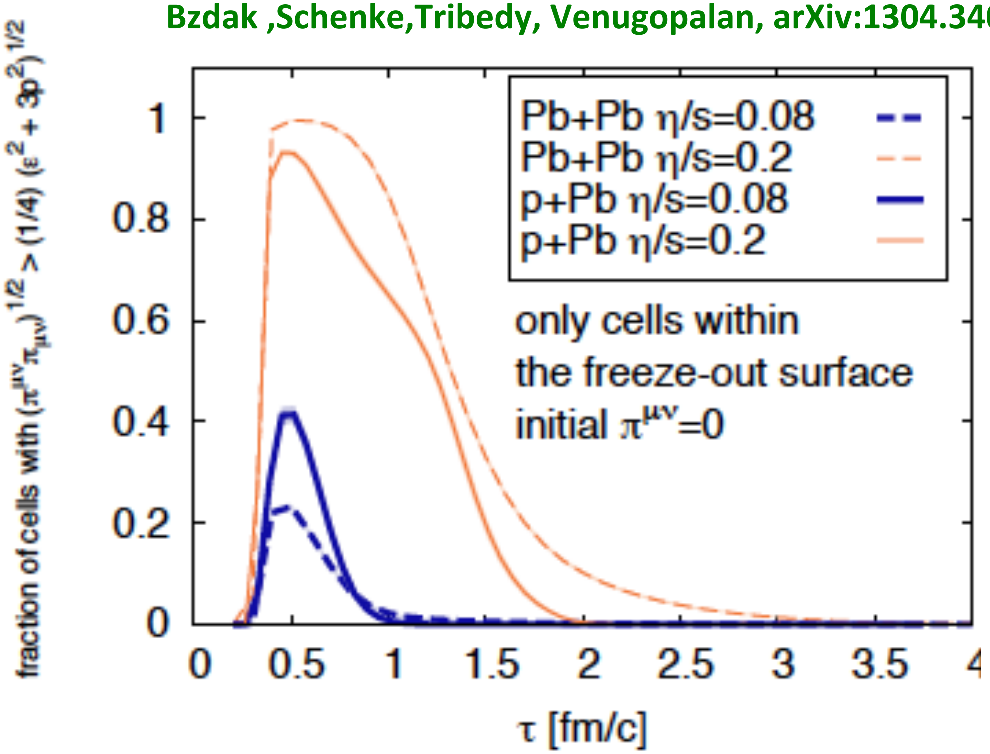
Evidence of a semi-hard scale in the data ?

Issues with the hydrodynamic paradigm: I

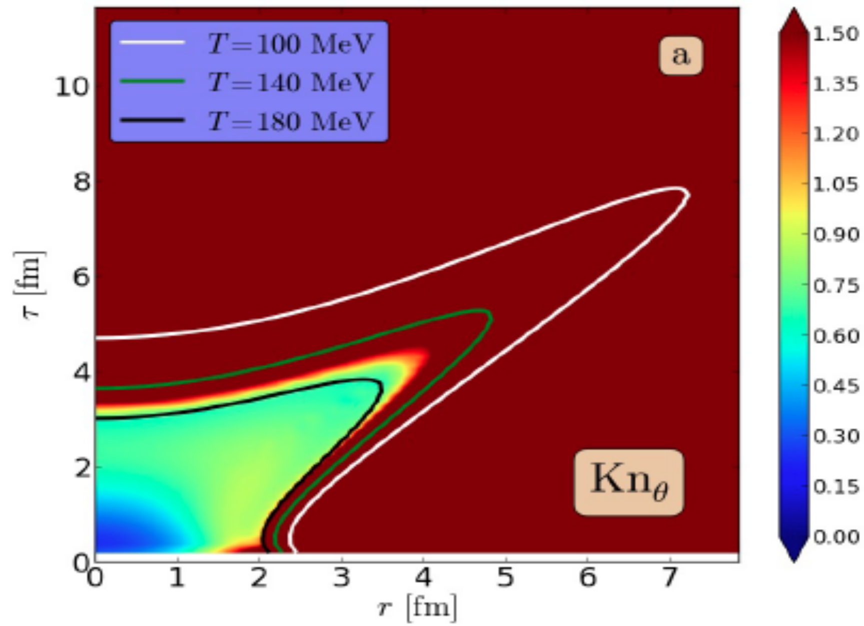
Two frequently used measures: Reynolds # and Knudsen #

$$R^{-1} \propto (\Pi^{\mu\nu}\Pi_{\mu\nu})^{1/2}/(\epsilon^2 + 3P^2)^{1/2} \qquad \mathrm{Kn} = \frac{\tau_\pi}{L} \;\; ; \;\; \tau_\pi \propto \frac{\eta}{sT}$$

Bzdak ,Schenke,Tribedy, Venugopalan, arXiv:1304.3403



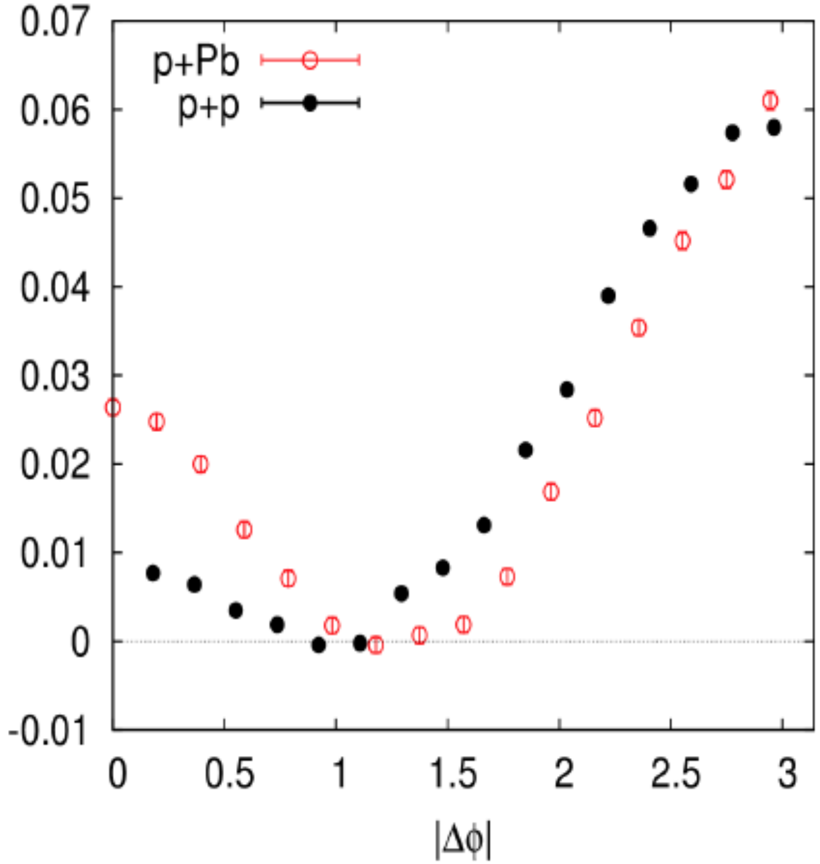
Dumitru,Molnar,Nara,arXiv:0706.2233
Denicol, Niemi, arXiv:1404.7327

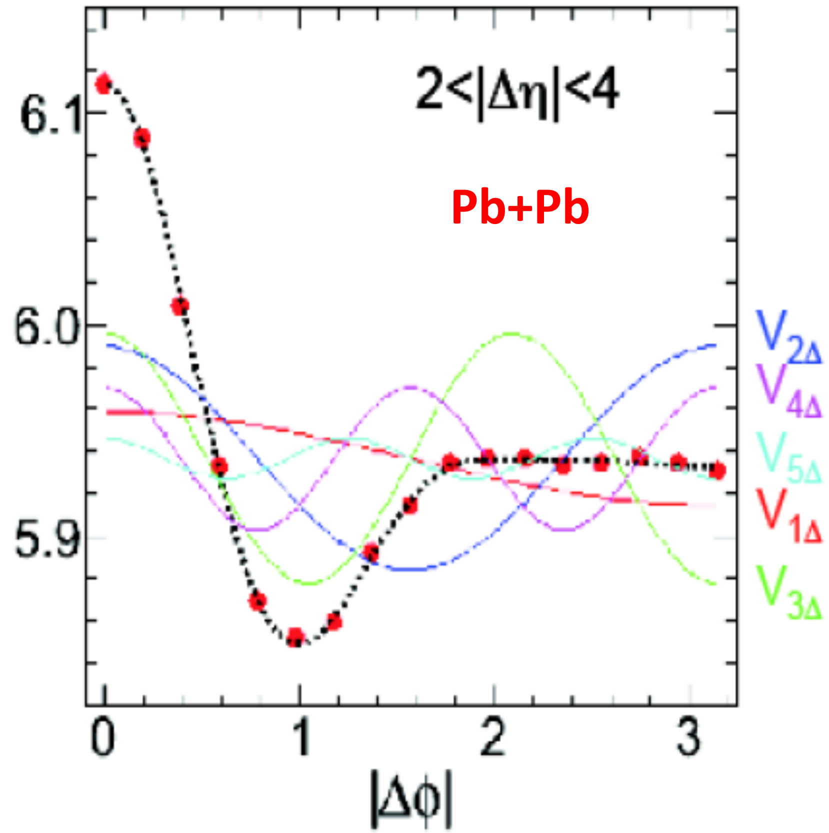


Hydro good for Kn < 0.5, marginal for K < 1 transient regime; K > 1 free streaming

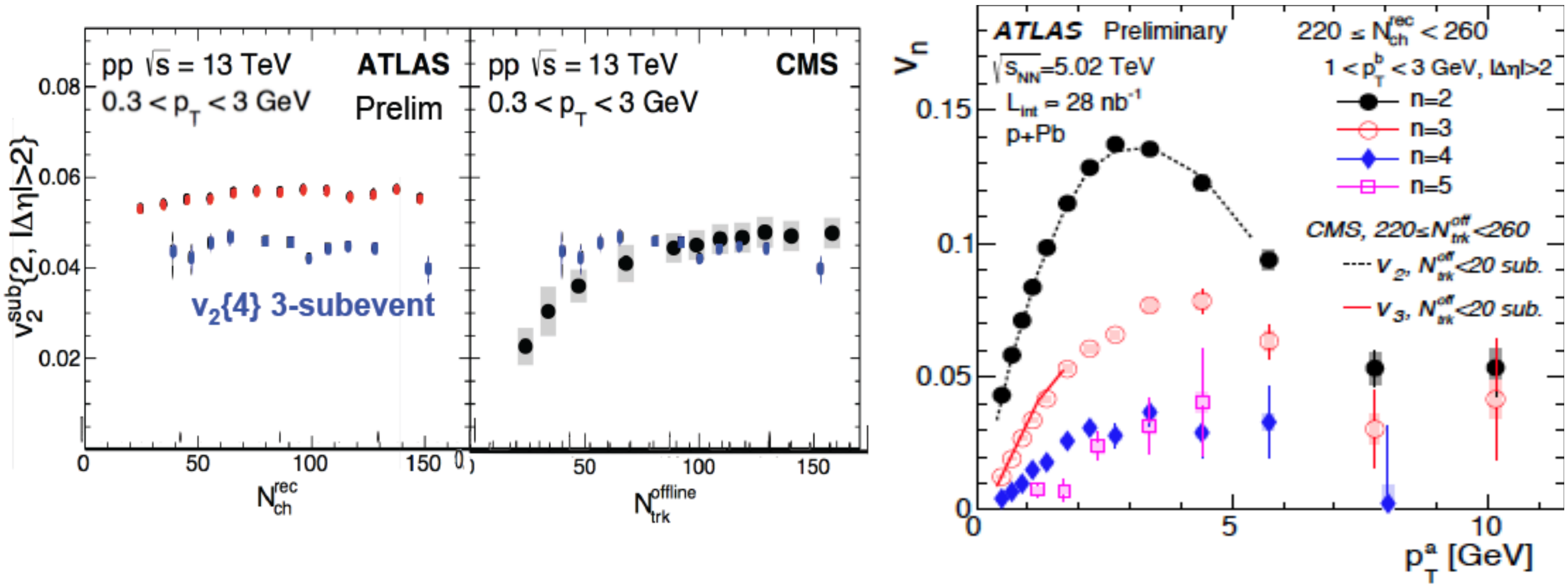
Issues with the hydrodynamic paradigm: II

No (mini-) jet quenching seen in the smaller systems





Issues with the hydrodynamic paradigm: III

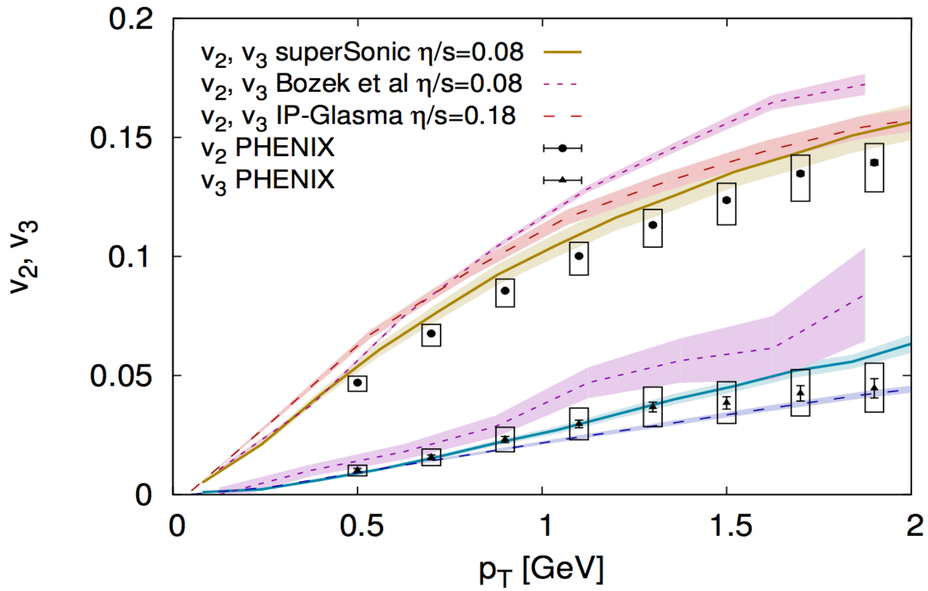


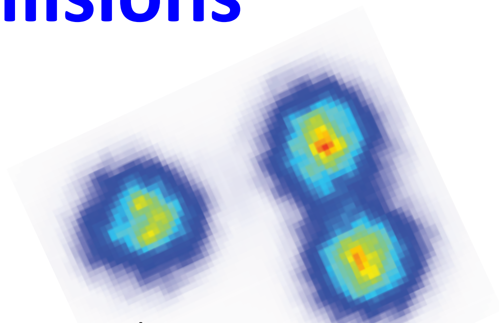
Large anistropies at larger p_T and smaller N_{ch} than one might reconcile with a hydrodynamic description

Four-particle collectivity seen in minimum bias events. . .

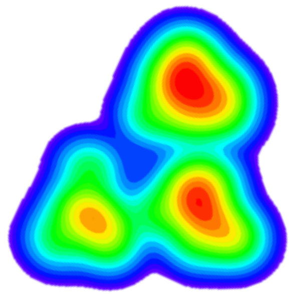
Collectivity in 3He+Au collisions

A. Adare et al. (PHENIX Collaboration)
Phys. Rev. Lett. 115, 142301 (2015)

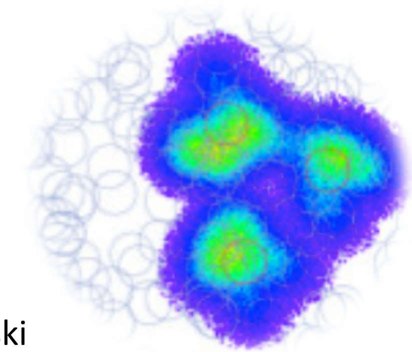




Schenke, Venugopalan
Nucl. Phys. A931 (2014) 1039-1044



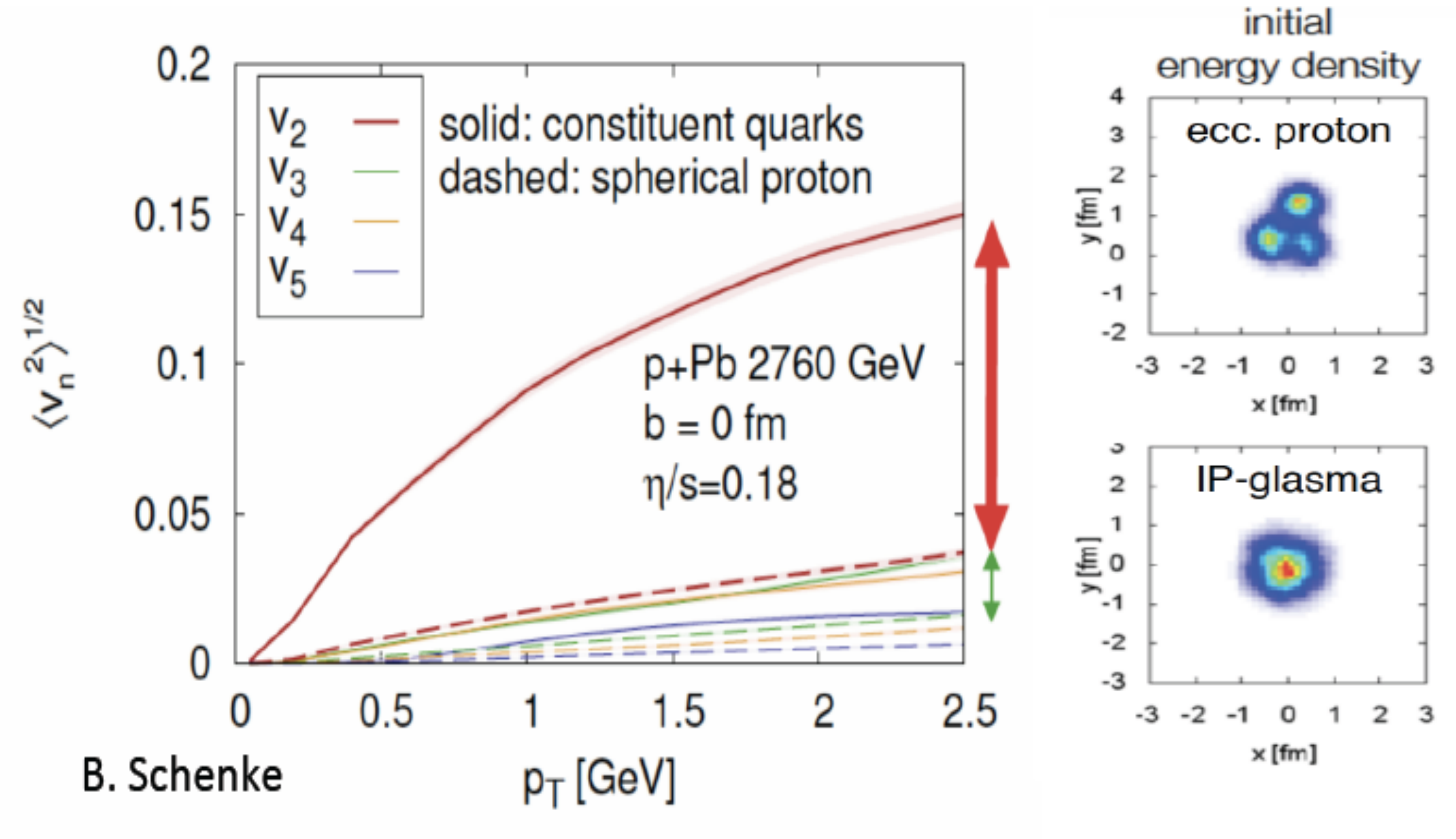
J.L. Nagle, et al.
Phys. Rev. Lett. 113, 112301



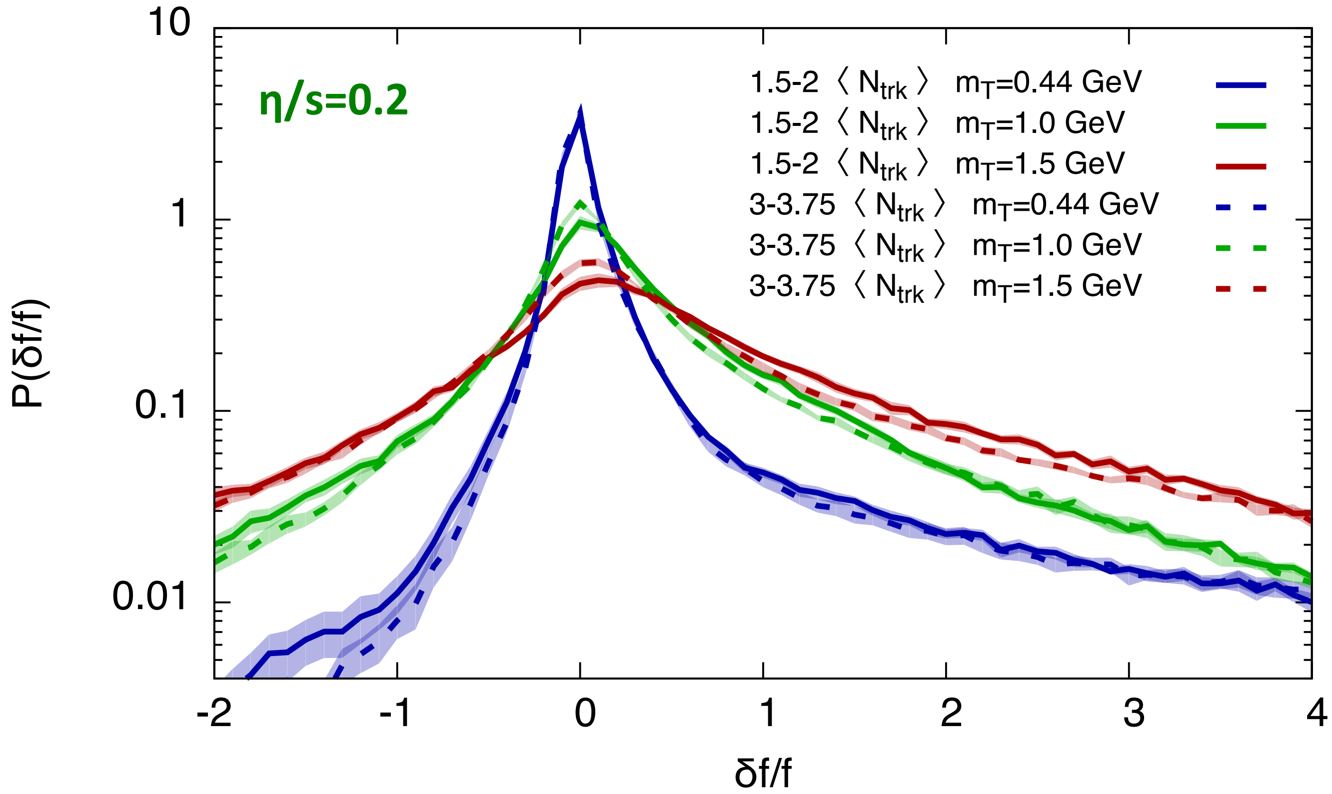
Bożek, Broniowski
Physics Letters B 747 (2015) 135–138

Shape matters ?

Schenke,Schlichting, 1407.8458



Freeze-out corrections in p+Pb as function of p_T

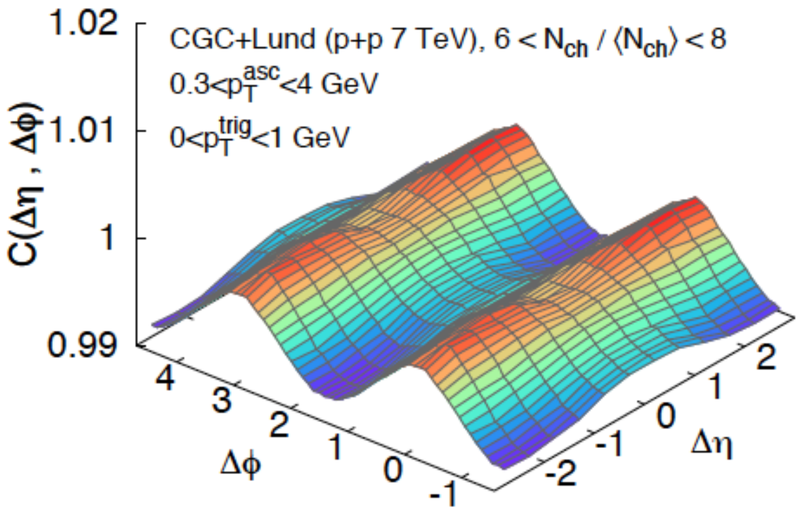


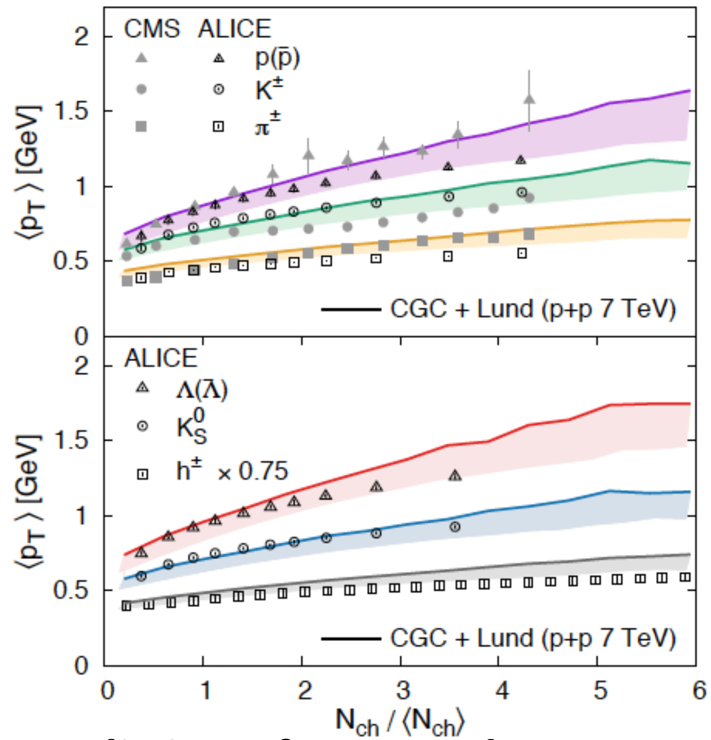
Plot by Bjoern Schenke

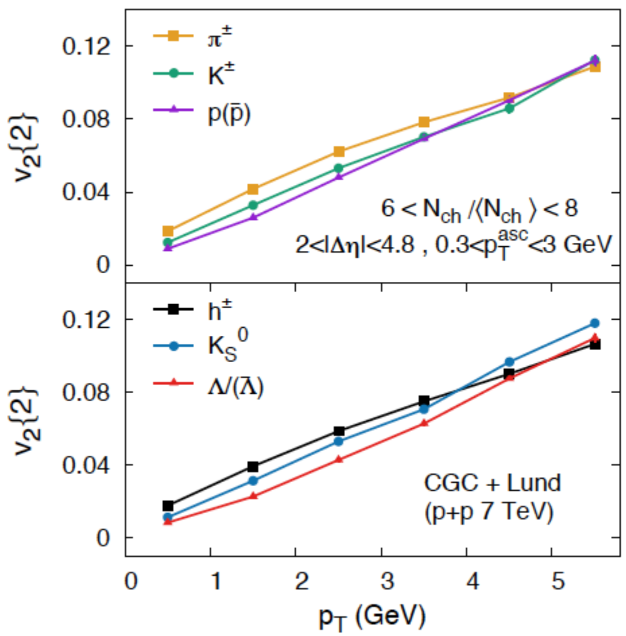
For m_T = 1 GeV, 26% of hydro cells have a 100% correction
For m_T = 1.5 GeV, 43% have a 100% correction

Mantysaari,Schenke,Shen,Tribedy, arXiv:1705.03177

IP-Glasma+Lund fragmentation





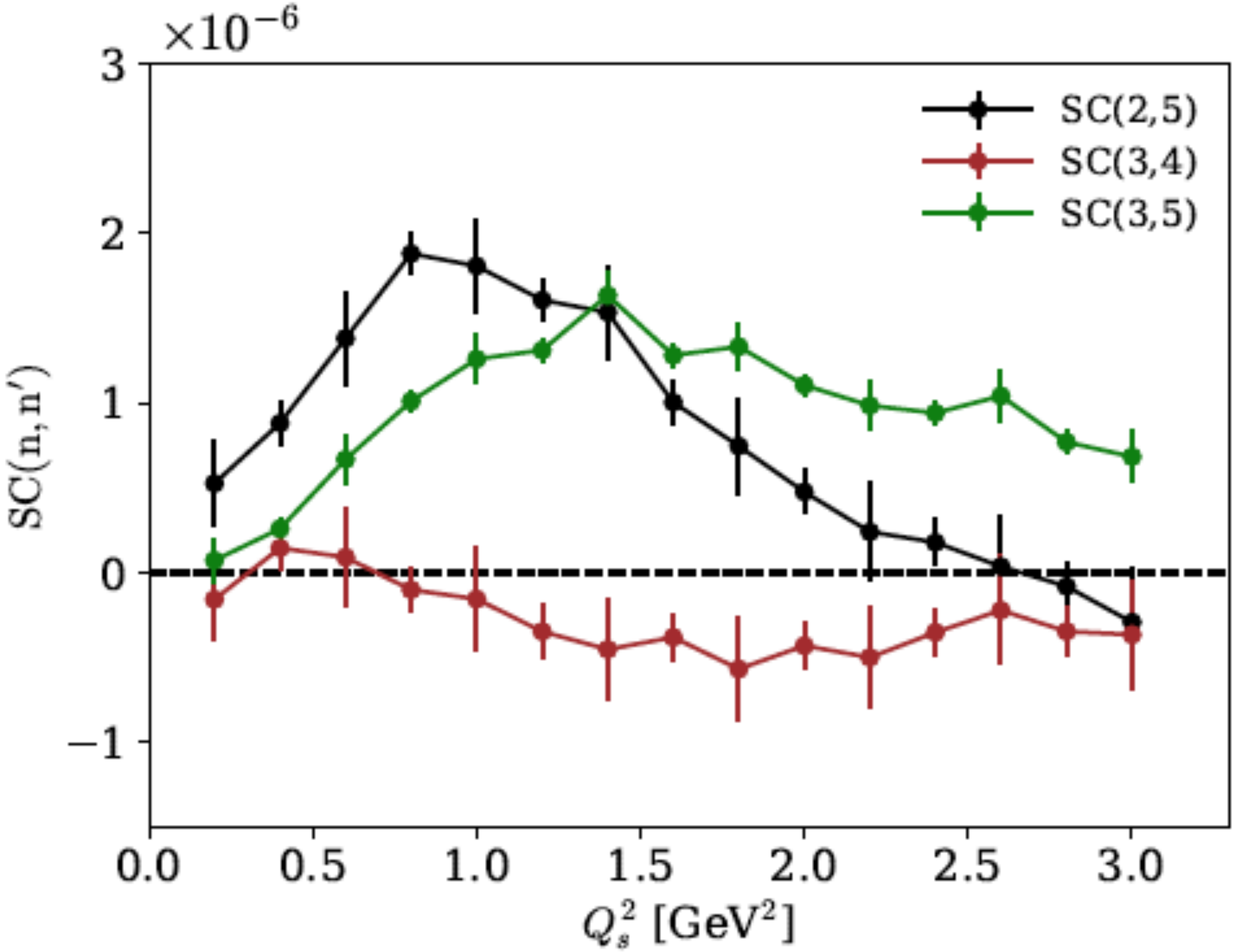


Pattern of mass splitting of $<p_T>$ and v_2 seen in high multiplicity events is reproduced Schenke,Schlichting,Tribedy,RV, PRL117(2016)162301

What about 4-particle collectivity?
Numerically very challenging-in progress

Schenke,Schlichting,Tribedy,RV

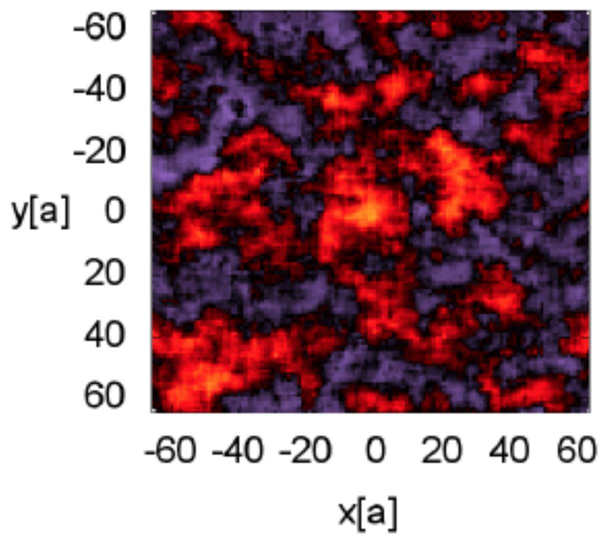
Predictions for p+A symmetric cumulants



Higher cumulants in the color domain model

Dumitru,McLerran,Skokov, 1410.4844

$\mathrm{Re}\,\mathrm{Tr}\left(V^\dagger(0,0)V(x,y)\right)/N_c$



Color domain model: express intrinsic higher point correlators as correlators of produced particles with a target field in a color domain, averaged over all orientations of the field.

$$c_2\{2\} = \frac{1}{N_D}\left(\mathcal{A}^2 + \frac{1}{4(N_c^2-1)}\right)$$

$$c_2\{4\} = -\frac{1}{N_D^3}\left(\mathcal{A}^4 - \frac{1}{4(N_c^2-1)^3}\right)$$

"A" term is the correlation induced between projectile particles due to color field orientation of target (more generically, non-Gaussian correlations)

The N_c term is the "connected Glasma graph" (Gaussian correlations)

N_D is # of color domains – few in p+A, several in A+A