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Exploring GPDs: energy-momentum tensor & applications*

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Outline

Introduction

hard-exclusive reactions: $GPDs \rightarrow ?$ crucial quantity: energy-momentum tensor

• Energy-momentum tensor

Ji sum rule & mechanical properties last unknown global nucleon property(!)

Applications

motivation & vision: mechanical stability practical use: from hard-exclusive reactions at JLab to charmonium pentaquark spectroscopy at LHCb

Outlook

cool and promising future!

* in collaboration with Irina Perevalova, Maxim Polyakov based on work supported by JLab, DOE, NSF

Introduction

hard-exclusive reactions
 factorization, access to GPDs
 Ji; Radyushkin; Collins, Frankfurt, Strikman



$$\int \frac{\mathrm{d}\lambda}{2\pi} e^{i\lambda x} \langle N'(p') | \overline{\psi}_q(-\frac{\lambda n}{2}) n_\mu \gamma^\mu [-\frac{\lambda n}{2}, \frac{\lambda n}{2}] \psi_q(\frac{\lambda n}{2}) | N(p) \rangle$$
$$= \overline{u}(p') \left[n_\mu \gamma^\mu H^q(x, \xi, t) + \overline{u}(p') \frac{i\sigma^{\mu\nu} n_\mu \Delta_\nu}{2M_N} E^q(x, \xi, t) \right] u(p)$$

• meaning: "microsurgery"



definitions for completeness: $\xi = (n \cdot \Delta)/(n \cdot P), \quad t = \Delta^2$ $P = \frac{1}{2}(p' + p), \quad \Delta = p' - p$ $n^2 = 0, \quad n \cdot P = 2, \quad k = xP$ renormalization scale μ analog gluon GPDs

• what do we learn?

We learn a lot, because GPDs ...

- generalize PDFs & em or axial form factors $\int dx \ H^q(x,\xi,t) = F_1^q(t)$ $\lim_{\Delta \to 0} H^q(x,\xi,t) = f_1^q(x)$
- explore impact parameter space
 & tomography (M. Burkardt, ...)

$$H^{q}(x,b_{\perp}) = \int \frac{\mathrm{d}^{2} \Delta_{\perp}}{(2\pi)^{2}} \left[\lim_{\xi \to 0} H^{q}(x,\xi,t) \right] e^{i \Delta_{T} b_{T}}$$

- allow to access (polynomiality) gravitational form factors $\int dx \ x \ H^q(x,\xi,t) = A^q(t) + \frac{4}{5}\xi^2 d_1^q(t)$ $\int dx \ x \ E^q(x,\xi,t) = B^q(t) - \frac{4}{5}\xi^2 d_1^q(t)$
- and gravity couples to
 energy momentum tensor
 probably most fundamental quantity



Energy-momentum tensor (EMT)

- instead of arguing how important EMT is, question: are you aware of introductory QFT text books which *do not* discuss EMT in first chapters?*
- if a theory can be solved: construct T_{μν} and generators of Poincaré group learn what is mass, spin, D-term (?) of the particles (in introductory QFT text books: free fields, so it can be done & is instructive)
- even if a theory cannot be solved, studies of EMT insightfull 3 examples (in QCD)

(in chronological order)

^{*} interestingly, advanced QFT books discuss EMT in later chapters: *trace anomaly* $\hat{T}_{\mu}^{\ \mu} \equiv \frac{\beta}{2g} F^{\mu\nu}F_{\mu\nu} + (1 + \gamma_m) \sum_q m_q \bar{\psi}_q \psi_q$ Adler, Collins, Duncan, PRD15 (1977) 1712; Nielsen, NPB 120, 212 (1977); Collins, Duncan, Joglekar, PRD 16, 438 (1977) • mass: celebrated Higgs mechanism explains $m_u \approx 2 \text{ MeV}$, $m_d \approx 5 \text{ MeV}$ $\hookrightarrow 2 \text{ MeV} + 2 \text{ MeV} + 5 \text{ MeV} \stackrel{!?}{=} 940 \text{ MeV}$ 2 orders of magnitude mismatch! Cool!!

trace anomaly \rightarrow glue! Can we do better? Ji, PRL 74, 1071 & PRD 52, 271 (1995) nucleon mass decomposition: how do E_{kin} , E_{pot} of quarks & gluons contribute?

• **spin** structure: **Ji** sum rule X.-D. Ji, Phys. Rev. Lett. 78 (1997) 610 prominent motivator since its appearance!

$$\int dx \ x \left(H^q(x,\xi,t) + E^q(x,\xi,t) \right) = A^q(t) + B^q(t) \xrightarrow{t \to 0} 2J^q(0)$$

 D-term, last unknown global property M. V. Polyakov, C. Weiss, PRD60 (1999) 114017 stress tensor & mechanical properties M. V. Polyakov, PLB555 (2003) 57

Energy momentum tensor $T^{\mu\nu}$

def

finition
$$\langle P' | \hat{T}^{q,g}_{\mu\nu} | P \rangle = \bar{u}(p') \left[A^{q,g}(t) \frac{\gamma_{\mu}P_{\nu} + \gamma_{\nu}P_{\mu}}{2} + B^{q,g}(t) \frac{i(P_{\mu}\sigma_{\nu\rho} + P_{\nu}\sigma_{\mu\rho})\Delta^{\rho}}{4M_{N}} + d_{1}^{q,g}(t) \frac{\Delta_{\mu}\Delta_{\nu} - g_{\mu\nu}\Delta^{2}}{5M_{N}} \pm \bar{c}(t)g_{\mu\nu} \right] u(p)$$

- $\hat{T}^{q}_{\mu\nu}$ and $\hat{T}^{g}_{\mu\nu}$ separately gauge-invariant (but not conserved)
- total EMT $\hat{T}_{\mu\nu} = \hat{T}^q_{\mu\nu} + \hat{T}^g_{\mu\nu}$ is conserved $\partial_{\mu}\hat{T}^{\mu\nu} = 0$
- constraints: mass $\Leftrightarrow A^q(0) + A^g(0) = 1$ (100% of nucleon momentum carried by quarks + gluons)

spin \Leftrightarrow $J^q(0) + J^g(0) = \frac{1}{2}$ (100% of nucleon spin carried by quarks + gluons)

• property: **D-term** \Leftrightarrow $d_1^q(0) + d_1^g(0) \equiv d_1 \rightarrow$ also conserved Noether charge!

but unconstrained! Unknown!

 \hookrightarrow last unknown global property(!?) of the nucleon

• How do we learn about nucleon?

 $|N\rangle = \text{strong}$ interaction particle. Use other forces to probe it!

em:
$$\partial_{\mu}J_{em}^{\mu} = 0$$
 $\langle N'|J_{em}^{\mu}|N \rangle \longrightarrow Q, \mu, ...$
weak: PCAC $\langle N'|J_{weak}^{\mu}|N \rangle \longrightarrow g_{A}, g_{p}, ...$
gravity: $\partial_{\mu}T_{grav}^{\mu\nu} = 0$ $\langle N'|T_{grav}^{\mu\nu}|N \rangle \longrightarrow M, J, d_{1}, ...$
1st global properties: $Q_{prot} = 1.602176487(40) \times 10^{-19}$ C
 $\mu_{prot} = 2.792847356(23)\mu_{N}$
 $g_{A} = 1.2694(28)$
 $g_{p} = 8.06(0.55)$
 $M = 938.272013(23) \text{ MeV}$
 $J = \frac{1}{2}$
 2^{nd} partonic structure: $\dots \dots \dots$
 $\dots \dots \dots$

• interpretation (M.V.Polyakov, PLB 555 (2003) 57)

Breit frame with $\Delta^{\mu} = (0, \vec{\Delta})$: static EMT $T_{\mu\nu}(\vec{r}, \vec{s}) = \int \frac{\mathrm{d}^{3}\vec{\Delta}}{2E(2\pi)^{3}} e^{i\vec{\Delta}\cdot\vec{r}} \langle P'|\hat{T}_{\mu\nu}|P\rangle$

interpretation okay for large- N_c : $M_N \sim N_c$, $t \sim N_c^0 \Rightarrow$ recoil corrections $t/M_N^2 \sim 1/N_c^2$ (formulae correct $\forall N_c$)

$$\int d^3r \ T_{00}(\vec{r}) = M_N \quad \text{known}$$

$$\int d^3r \ \varepsilon^{ijk} s_i r_j \ T_{0k}(\vec{r}) = \frac{1}{2} \quad \text{known}$$

$$- \frac{M_N}{2} \int d^3r \ \left(r^i r^j - \frac{r^2}{3} \delta^{ij}\right) \ T_{ij}(\vec{r}) \equiv d_1 \quad \text{new!}$$

with:
$$T_{ij}(\vec{r}) = s(r) \left(\frac{r_i r_j}{r^2} - \frac{1}{3} \delta_{ij} \right) + p(r) \delta_{ij}$$
 stress tensor

 $\left. \begin{array}{c} s(r) \ \text{related to distribution of shear forces} \\ p(r) \ \text{distribution of pressure inside hadron} \end{array} \right\} \longrightarrow$ "mechanical properties"

• relation to stability: EMT conservation $\Leftrightarrow \partial^{\mu} \hat{T}_{\mu\nu} = 0 \Leftrightarrow \nabla^{i} T_{ij}(\vec{r}) = 0$

 \hookrightarrow necessary condition for stability $\int dr r^2 p(r) = 0$ (von Laue, 1911)

$$d_{1} = -\frac{4\pi}{3} M_{N} \int_{0}^{\infty} dr \ r^{4}s(r) = 5\pi M_{N} \int_{0}^{0} dr \ r^{4} p(r)$$

 $\hookrightarrow d_1$ shows how internal forces balance



insights:

$$\begin{split} T_{00}(0) &= 1.70 \,\text{GeV/fm}^3 \approx 3 \times 10^{15} \,\rho(\text{H}_2\text{O}) \approx 13 \times (\text{nuclear density}) \\ p(0) &= 0.23 \,\text{GeV/fm}^3 \approx 4 \times 10^{34} \,\text{N/m}^2 \quad \gtrsim 100 \times (\text{pressure in center of neutron star}) \\ & \text{ in chiral quark soliton model} \quad (\text{Goeke et al, PRD75 (2007) 094021}) \end{split}$$

... and in QCD? wouldn't it be fascinating to know??

• investigating forces

prominent property of proton: life time $\tau_{\rm prot} > 2.1 \times 10^{29}$ years!

question: how do strong forces balance to produce stability?

- answer in model: strong cancellation of repulsive forces due to quark core, and attractive forces from pion cloud
- answer in **QCD:** we do not know

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nice pictures(!), attractive insights(!)
underexplored propaganda(?)
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(be aware: same picture for neutron,

\tau_{\text{neut}} = 14 \min 40 \sec \gg 10^{-23} \sec)
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• as mental support for GPD program: okay

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... but is there any practical use of that?
answer before: not really ...
answer today: <u>Yes!</u>
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in chiral quark soliton model chiral symmtry breaking \checkmark realization of QCD in large- $N_c \checkmark$ built on instanton vacuum calculus \checkmark not bad, but after all a model ... Goeke et al, PRD75 (2007)





from hard-exclusive reactions at JLab, COMPASS ...

... to spectroscopy of $\bar{c}c$ -pentaquarks at LHCb

not usual hadrons, not just any exotic hadron

only $\bar{c}c$ -baryon bound states \rightarrow rich enough!

• discovery of charmonium pentaquarks in Λ_b^0 decays at LHCb

Aaij et al. PRL 115, 072001 (2015)

$$\Lambda^0_b \longrightarrow J/\Psi\, p\, K^-$$
 seen

 $\begin{array}{ll} \Lambda^0_b & m = 5.6 \ {\rm GeV}, \ \ \tau = 1.5 \ {\rm ps} \\ J/\Psi & m = 3.1 \ {\rm GeV}, \ \ \Gamma = 93 \ {\rm keV}, \ \ \Gamma_{\mu^+\mu^-} = 6 \ \% \\ \Lambda^* & m = 1.4 \ {\rm GeV} \ {\rm or \ more}, \ \Lambda^* \to K^-p \ {\rm in} \ 10^{-23} {\rm s} \end{array}$



state	$m \; [MeV]$	Γ [MeV]	Γ _{rel}	mode	J^P
P_c^+(4380)	$4380\pm8\pm29$	$205\pm18\pm86$	$(4.1\pm0.5\pm1.1)\%$	J/\psip	$\frac{3}{2}^{\mp}$ or $\frac{5}{2}^{+}$
$P_{c}^{+}(4450)$	$4450\pm2\pm3$	$39\pm5\pm19$	$(8.4 \pm 0.7 \pm 4.2)$ %	J/\psip	$\frac{5}{2}^{\pm}$ or $\frac{3}{2}^{-}$

• how can GPDs help? M. I. Eides, V. Y. Petrov and M. V. Polyakov, PRD93, 054039 (2016) $R_{J/\psi} \ll R_N \Rightarrow$ non-relativistic multipole expansion Gottfried, PRL 40 (1978) 598 baryon-quarkonium interaction dominated by 2 virtual chromoelectric dipole gluons

 $V_{\text{eff}} = -\frac{1}{2} \alpha \vec{E}^2$ Voloshin, Sov. J. Nucl. Phys. **36**, 143 (1982) [Yad. Fiz. **36**, 247 (1982)]

 α = chromoelectric polarizability heavy quark mass limit & large- N_c limit \rightarrow "perturbative result" Peskin, NPB 156 (1979) 365

value for $2S \rightarrow 1S$ transition from phenomenological analysis of $\psi' \rightarrow J/\psi \pi \pi$ data Voloshin, Prog. Part. Nucl. Phys. 61 (2008) 455

(chromoelectric field)² =
$$\vec{E}^2 = \frac{1}{2}(\vec{E}^2 + \vec{H}^2) + \frac{1}{2}(\vec{E}^2 - \vec{H}^2) = g^2 \left(\frac{8\pi^2}{bg^2}T^{\mu}_{\ \mu} + T^G_{00}\right)$$

 $\alpha(1S) \approx 0.2 \,\mathrm{GeV^{-3}}$ (pert).

 $\alpha(2S) \approx 12 \,\mathrm{GeV}^{-3}$ (pert),

 $\alpha(2S \to 1S) \approx \begin{cases} -0.6 \,\text{GeV}^{-3} \,\text{(pert)}, \\ +2 \,\text{GeV}^{-3} \,\text{(pheno)}. \end{cases}$

 $b = \frac{11}{3} N_c - \frac{2}{3} N_F$ leading coeff. of β -function $g = \text{strong coupling at low (nucleon) scale} \lesssim 1 \text{ GeV}$ $g_s = \text{strong coupling at scale of heavy quark (e.g. <math>m_c \sim 1.5 \text{ GeV}, g_s \neq g$) $T_{00}^G = \xi T_{00}$ with $\xi = \text{fractional contributions of gluon to nucleon mass}$

 $T^{\mu}{}_{\mu} = T^{00} - T^{ii}$ with stress tensor $T^{ij} = \left(\frac{r^i}{r}\frac{r^j}{r} - \frac{1}{3}\delta^{ij}\right)s(r) + \delta^{ij}p(r)$ M. V. Polyakov, PLB **555**, 57 (2003)

$$V_{
m eff} = -rac{1}{2} lpha \; rac{8 \pi^2}{b} rac{g^2}{g_s^2} \Bigl[
u \, T_{00}(r) + 3 p(r) \Bigr] \,, \quad
u = 1 + \xi_s rac{b \, g_s^2}{8 \pi^2}$$

universal potential

 $\nu \approx 1.5$ Eides et al, op. cit. Novikov & Shifman, Z.Phys.C8, 43 (1981); X. D. Ji, Phys. Rev. Lett. **74**, 1071 (1995)

- in future GPDs can help: GPDs \Rightarrow EMT form factors \Rightarrow EMT densities \Rightarrow universal potential V_{eff} for quarkonium-baryon interaction!
- **Currently:** use e.g. chiral quark soliton model (done in Eides et al, op. cit.)
- compute quarkonium-nucleon bound state

solve
$$\left(-\frac{\vec{\nabla}^2}{2\mu} + V_{\rm eff}(r)\right)\psi = E_{\rm bind}\psi$$

(here μ = reduced quarkonium-baryon mass)

• results:



nucleon and J/ψ form no bound state

nucleon and $\psi(2S)$ form 2 bound states with nearly degenerate masses around 4450 MeV in L = 0 channel, $J^P = \frac{1}{2}^+$ and $\frac{3}{2}^+$ if $\alpha(2S) \approx 17$ GeV⁻³ (consistent with guideline from pert. calc.)

• decay

cannot decay directly to $\psi(2S)$ and nucleon, as $M_{\psi(2S)} + M_N > 4450 \,\text{MeV}$

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instead transition (2S) \rightarrow (1S) governed by the same V_{eff}
but with small \alpha(2S \rightarrow 1S) transition polarizability
\Rightarrow it "takes time"
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after transition "completed," prompt decay to J/\psi + nucleon (as observed)
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estimated width is tens of MeV \rightarrow compatible with data!

 based on χQSM (Eides et al) model-dependent? Need to confirm in different models! (Perevalova, Polyakov, PS, in progress)

- Skyrme model (Cebulla et al 2007) incorporates chiral symmetry & also soliton however, different model, different way to realize stability (Skyrme term) ideal to provide an independent cross-check
- result:

same conclusions!

Very robust predictions for mass and decay width of P_c^+ (4450) important confirmation!

• what about $P_c^+(4380)$? much broader more possibilities under investigation

Summary & Outlook

- GPDs important objects, but what do we learn?
- crucial quantity: form factors of energy momentum tensor mass decomposition, spin decomposition, and *D*-term!
- what's D-term? Last unknown global property, related to forces attractive and physically appealling \rightarrow "motivation"
- recent development: knowledge of internal forces and energy density \rightarrow quarkonium-baryon interaction $V_{\rm eff}$
- naturally explains properties of $P_c^+(4450)$ observed at LHCb rich potential, ongoing work



