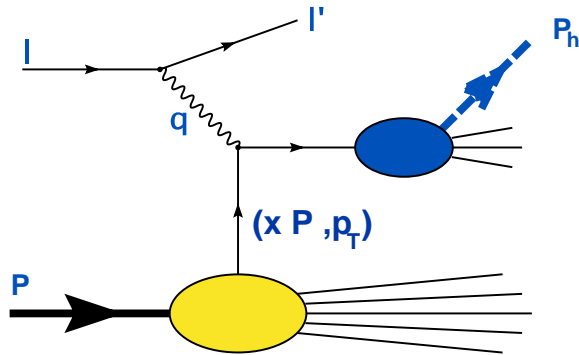


# **“Time-reversal odd” (T-odd) effects in semi-inclusive deep-inelastic lepton-nucleon scattering**

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# Partonic picture of semi-inclusive deep-inelastic scattering (SIDIS)



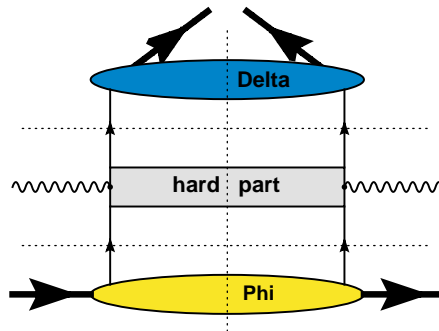
kinematic variables:

$$P^+ \text{ large} \quad ; \quad -q^2 = Q^2 \text{ large}$$

$$x_B = \frac{Q^2}{2P \cdot q}; \quad z_h = \frac{P \cdot P_h}{P \cdot q}; \quad |\vec{P}_{h\perp}| \ll Q$$

## The SIDIS cross section

quark-quark correlators:



$$\Phi_{ij}(x, \vec{p}_T) = \int \frac{d\xi^- d^2\xi_T}{(2\pi)^3} e^{ixP^+\xi^- - i\vec{p}_T \cdot \vec{\xi}_T} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle \Big|_{\xi^+=0}$$

$$\Delta_{ij}(z, \vec{k}_T) = \mathcal{FT} \sum_X \langle 0 | \psi_i(\xi) | P_h; X \rangle \langle P_h; X | \bar{\psi}_j(0) | 0 \rangle \Big|_{\xi^-=0}$$

$$\sigma_{\text{SIDIS}} = \Phi \otimes (\text{hard}) \otimes \Delta$$

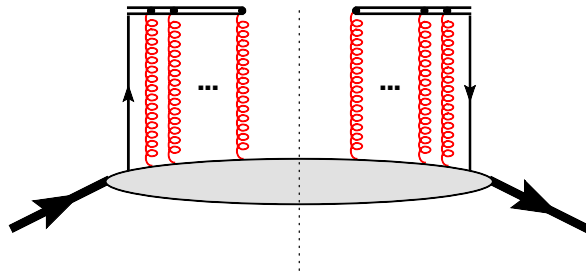
- Intrinsic dependence on the transv. parton momentum  $\vec{p}_T$ .
- $\Phi \longrightarrow$   $p_T$ -dependent parton distributions.
- $\Delta \longrightarrow$   $k_T$ -dependent fragmentation functions.

## Color gauge invariance

mathematical definition of the gauge link:

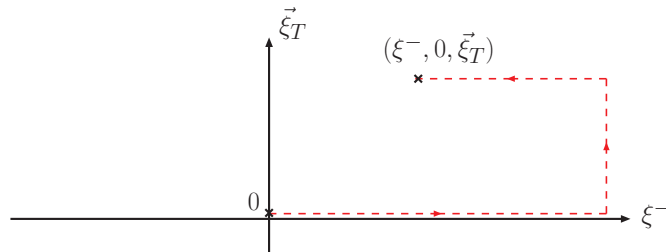
$$\mathcal{W}(a, b | \text{path}) = \mathcal{P} \exp \left( -ig \int_a^b ds^\mu A_\mu(s) \right)$$

physical picture:



$$\Phi(x, \vec{p}_T) = \mathcal{FT} \left[ \langle P | \bar{\psi}_j(0) \mathcal{W}(0, \xi) \psi_i(\xi) | P \rangle \Big|_{\xi^+ = 0} \right]$$

Wilson line in SIDIS:



$$\mathcal{W}(0, \xi | n_-) = \left[ 0; \infty, 0, \vec{0}_T \right] \times \left[ \infty, 0, \vec{0}_T; \infty, 0, \vec{\xi}_T \right] \times \left[ \infty, 0, \vec{\xi}_T; \xi^-, 0, \vec{\xi}_T \right]$$

- Consequences of the Gauge link: T-odd parton distributions, contributions to single-spin asymmetries.

## Parameterization of $qq$ -correlators

1. )  $\Phi_{ij}, \Delta_{ij}$  matrices in Dirac-space  $\longrightarrow$  decomposition into  $\mathbb{1}, \gamma_5, \gamma^\mu, \gamma^\mu \gamma_5, i\sigma^{\mu\nu} \gamma_5$
2. ) **Coefficients**  $\implies$  parameterization of traces in PDFs and FFs, e.g.

$$\frac{1}{2} \text{Tr} \left[ \Phi(x, \vec{p}_T) \gamma^+ \right] = f_1(x, \vec{p}_T^2) - \frac{\vec{p}_T \times \vec{S}_T}{M} f_{1T}^\perp(x, \vec{p}_T^2)$$

3. ) Parameterization restricted by constraints due to **parity (P)** and **time-reversal (T)**-transformations.
4. ) **T-constraint** only valid if the Wilson line is neglected  $\implies$  Wilson line enables **T-odd PDFs**.
5. ) **Choice of Wilson line**  $\implies$  dependence of  $\Phi$  and  $\Delta$  on lightcone vectors  $n_-$  and  $n_+$ ,

$$\Phi = \Phi_{ij}(x, \vec{p}_T; S | n_-) \quad ; \quad \Delta = \Delta_{ij}(z, \vec{k}_T; S_h | n_+)$$

$\implies$  induces additional **T-odd twist-3 PDFs and FFs**  $g^\perp, e_T^\perp, f_T', G^\perp$ .

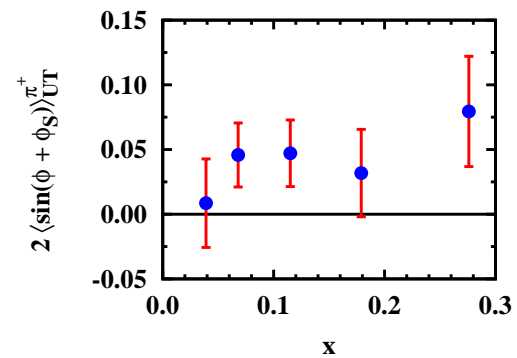
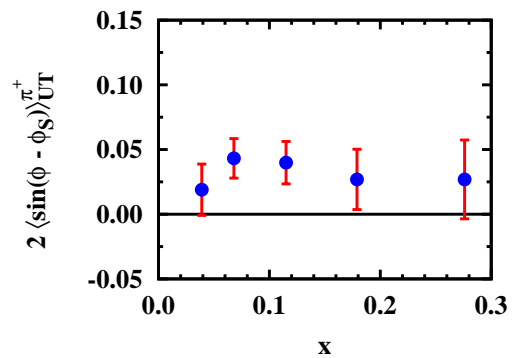
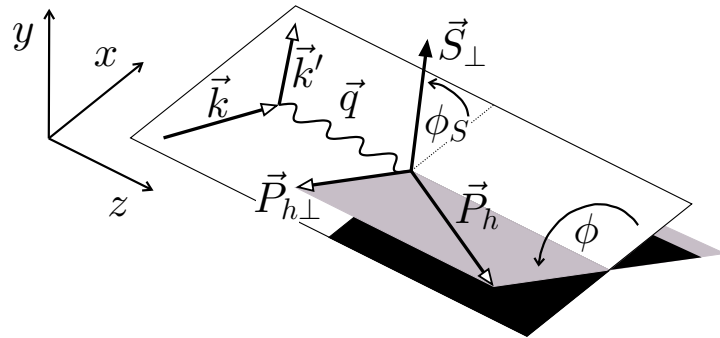
6. ) **Complete parameterization.**

# Single-Spin asymmetries $A = \frac{\sigma(\uparrow) - \sigma(\downarrow)}{\sigma(\uparrow) + \sigma(\downarrow)}$ in SIDIS

Twist-2 SSA for transversely pol. target:

$A_{UT}^{\text{Sivers}}$	$\propto$	$\sin(\phi_h - \phi_s) f_{1T}^\perp D_1$
$A_{UT}^{\text{Collins}}$	$\propto$	$\sin(\phi_h + \phi_s) h_1 H_1^\perp,$

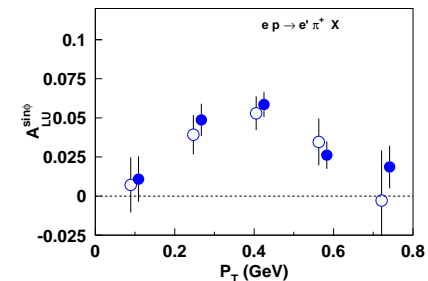
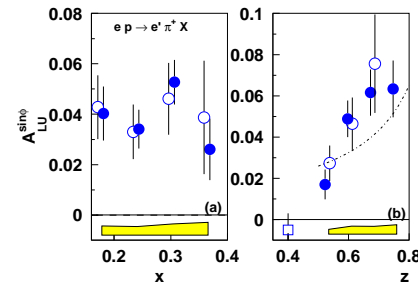
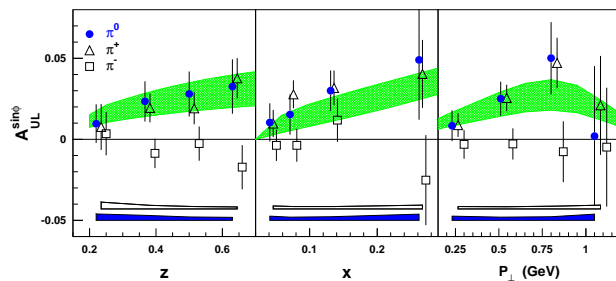
Experiments by HERMES (also by COMPASS):



Twist-3 SSAs for longitudinally pol. target and beam:

$A_{UL}^{\sin \phi}$  measured by HERMES (2001):

$A_{LU}^{\sin \phi}$  measured by CLAS (2003):



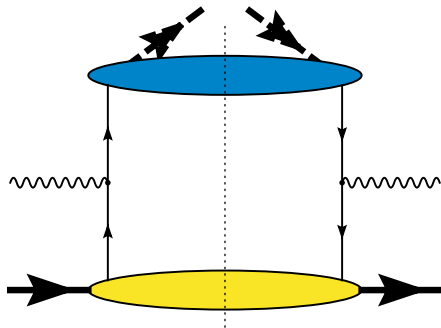
$\implies$  Twist-3 observables theoretically more complicated!

# “Tree-level” formalism of twist-3 observables

(Mulders, Tangerman, 1996; Boer, Mulders, Pijlman, 2003)

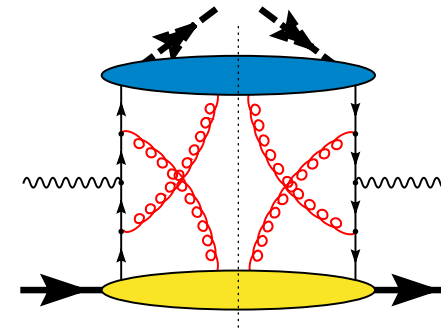
## Contributions to twist-3 observables:

### 1. Twist-3 correlation functions:



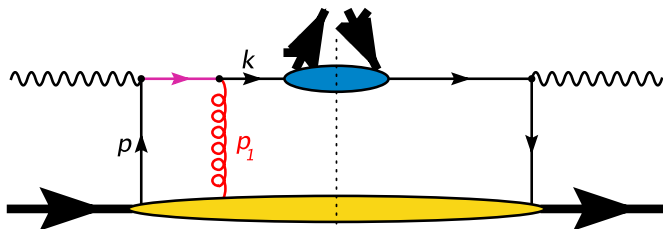
Contribution contains **twist-3 PDFs/FFs**

### 2. Genuine twist-3:



Quark-Gluon-Quark correlators

## Treatment of $qgq$ -Correlators:



$$\propto \int dp_1 \mathcal{F}_T \left( \langle P, S | \bar{\psi}(0) \gamma^\mu \Delta \gamma^\rho \frac{\not{k} - \not{p}_1 + m}{(k - p_1)^2 - m^2 + i0} g A^\rho(\eta) \gamma^\nu \psi(\xi) | P, S \rangle \right)$$

### 1/Q-expansion of the propagator:

- $A^+$  ( $A_T^i(x^- = \infty)$ )-gluons  $\rightarrow$  Gauge Link
- Transverse  $A_T^i(x)$ -gluons  $\rightarrow$  Twist-3  $\rightarrow$  color g.i.  $qgq$ -corr.  $\Phi_A^i$

- Parameterization of  $qqq$ -correlator  $\Phi_A^i \propto \langle P | \bar{\psi} (\int \mathcal{W} F^{+i} \mathcal{W}) \psi | P \rangle$  in terms of *tilde*-functions, e.g.

$$\text{Tr} \left[ \Phi_{A,i} i\sigma^{i+} \gamma_5 \right] = 2Mx \left( (\tilde{h}_L + i\tilde{e}_L) + \frac{\vec{p}_T \cdot \vec{S}_T}{M} (\tilde{h}_T + i\tilde{e}_T) \right), \dots$$

- Connection between *tilde*-functions and PDFs/FFs via QCD-equation of motion, e.g.

$$\tilde{h}_L = h_L + \frac{\vec{p}_T^2}{2M^2} \frac{h_{1L}^\perp}{x} - \frac{m}{M} \frac{g_{1L}}{x} \quad ; \quad \tilde{g}_T = g_T - \frac{\vec{p}_T^2}{2M^2} \frac{g_{1T}}{x} - \frac{m}{M} \frac{h_1}{x}, \dots$$

$$\tilde{e}_T = e_T \quad ; \quad \tilde{e}_T^\perp = e_T^\perp - \frac{m}{M} \frac{f_{1T}^\perp}{x}, \dots$$

- Decomposition of SIDIS-cross section into structure functions  $F(x_B, z_h, Q^2, \vec{P}_{h\perp}^2)$ :

$$\text{Twist-2: } F_{UU}, F_{UU}^{\cos 2\phi_h}, F_{UL}^{\sin 2\phi_h}, F_{LL}, F_{UT}^{\sin(\phi_h - \phi_s)}, F_{UT}^{\sin(\phi_h + \phi_s)}, F_{UT}^{\sin(3\phi_h - \phi_s)}, F_{LT}^{\cos(\phi_h - \phi_s)}$$

$$\text{Twist-3: } F_{UU}^{\cos \phi_h}, F_{LU}^{\sin \phi_h}, F_{UL}^{\sin \phi_h}, F_{LL}^{\cos \phi_h}, F_{UT}^{\sin \phi_s}, F_{UT}^{\sin(2\phi_h - \phi_s)}, F_{LT}^{\cos \phi_h}, F_{LT}^{\cos(2\phi_h - \phi_s)}$$

- Structure functions in terms of PDFs and FFs, e.g.

(Mulders, Tangerman, 1996; Boer, Mulders, 1998; Bacchetta, Mulders, Pijlman, 2004; Bacchetta, Diehl, Goeke, Metz, Mulders, M.S., 2006)

$$F_{LT}^{\cos(\phi_h - \phi_s)} \propto I [\alpha g_{1T} D_1],$$

$$F_{LT}^{\cos \phi_s} \propto \frac{M}{Q} I \left[ \alpha_1 g_T D_1 + \alpha_2 h_1 \tilde{E} + \alpha_3 e_T H_1^\perp + \alpha_4 g_{1T} \tilde{D}^\perp + \alpha_5 e_T^\perp H_1^\perp + \alpha_6 f_{1T}^\perp \tilde{G}^\perp \right],$$

$$F_{LU}^{\sin \phi_h} \propto \frac{M}{Q} I \left[ \beta_1 e H_1^\perp + \beta_2 f_1 \tilde{G}^\perp + \beta_3 g^\perp D_1 + \beta_4 h_1^\perp \tilde{E} \right], \dots$$

- Simplifications:

1.) Wandzura-Wilczek-approximation: neglect *tilde*-functions

$$F_{LT}^{\cos \phi_s} \longrightarrow \frac{M}{Q} I [\beta g_{1T} D_1]$$

2.) Semi-inclusive jet production:  $D_1 \longrightarrow \delta(1 - z) \delta^{(2)}(\vec{k}_T)$ ,  $H_1^\perp = \dots = 0$

jet asymmetries

$$\boxed{A_{LU, \text{jet}}^{\sin \phi_h} \propto \frac{M}{Q} \frac{g_1^\perp(1)(x)}{f_1(x)}} \quad ; \quad \boxed{A_{UL, \text{jet}}^{\sin \phi_h} \propto \frac{M}{Q} \frac{f_L^\perp(1)(x)}{f_1(x)}}$$

- Results valid when applying “tree level” formalism for twist-3 observables.

Check of tree-level formalism?



Model calculations

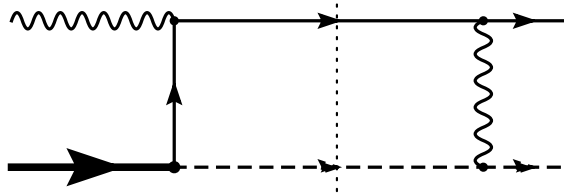


## Longitudinal jet-SSAs in the diquark-spectator-model:

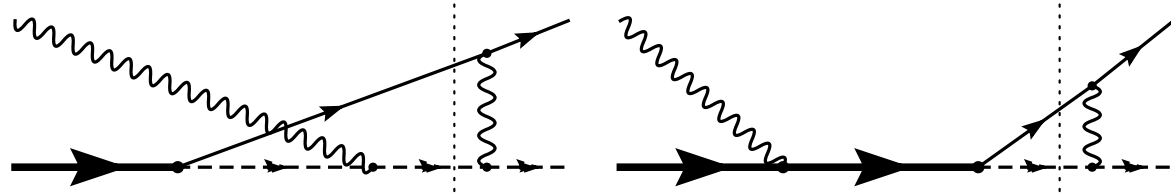
(Afanasev, Carlson, 2003, 2006; Metz, M.S., 2004)

Diagrams with imaginary parts necessary for SSAs:

Rescattering effect (RE):



non-compatible (NC) diagrams:



Non-compatible diagrams needed to restore electromagnetic gauge invariance ( $M = \varepsilon_\mu(q) J^\mu$ ):

$$q_\mu (\text{Im} (J_{RE}^\mu) + \text{Im} (J_{NC,D}^\mu) + \text{Im} (J_{NC,N}^\mu)) = 0$$

Difference to twist-2 observables:

- 1)  $A_{UT}$ : NC-diagrams ( $J_{NC}^\perp$ ) are suppressed by  $\mathcal{O}(1/Q^2)$ .
- 2)  $A_{UL}, A_{LU}$ : NC-diagrams ( $J_{NC}^\pm$ ) are of the same order as RE-diagram ( $J_{RE}^\pm$ ).

$\Rightarrow$  Result: Non-vanishing asymmetries  $A_{UL} \neq 0$ ,  $A_{LU} \neq 0$ .

## Model calculation of $g^\perp$

(Gamberg, Hwang, Metz, M.S., 2006)

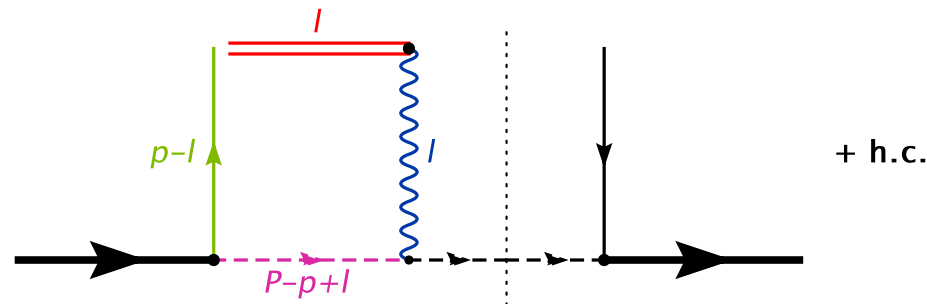
Beam-spin asymmetry:  $A_{LU, \text{jet}}^{\sin \phi} \propto g^{\perp(1)}(x)$   $\longrightarrow$   $g^\perp$  in the **diquark-spectator-model**

- Starting “from scratch”:

$$\Phi_{ij}(x, \vec{p}_T; P, S|n) = \mathcal{FT} \left( \langle P, S | \bar{\psi}_j(0) \mathcal{W}^{[-]}(0; \infty^- | n) |dq\rangle \langle dq| W^{[-]}((\infty^-, \vec{\xi}_T); \xi | n) \psi_i(\xi) | P, S \rangle \right)$$

$$g^\perp \propto \text{Tr} \left[ \Phi(x, \vec{p}_T) \gamma^i \gamma_5 \right] \Big|_{\text{unpol}}$$

- Leading non-trivial order:



$$g^\perp \propto \sum_{\pm} \int \frac{d^4 l}{(2\pi)^4} \frac{v \cdot (2P - 2p + l) \left[ \epsilon_T^{ij} p_T^j \left( P^+ l^- - P^- l^+ \right) + \epsilon_T^{ij} l_T^j \left( P^- p^+ - P^+ p^- \right) \right]}{[(l \cdot n) \pm i0] [l^2 \mp i0] [(P - p + l)^2 - m_s^2 \mp i0] [(p - l)^2 - m_q^2 \mp i0]}$$

For  $n = [1^-, 0^+, \vec{0}_T]$  on the light-cone  $\longrightarrow$  **Divergence!**

- Regularization: “non lightlike” Wilson lines:  $n \longrightarrow v = [v^-, v^+, \vec{0}_T], \left| \frac{v^+}{v^-} \right| \ll 1$

- Result:

$$g^\perp(x, \vec{p}_T^2, v) \propto \frac{1-x}{x} \ln \left( \frac{v^2 \tilde{m}^2}{2(v \cdot P)^2} \right) + \text{finite} + \mathcal{O} \left( \left| \frac{v^+}{v^-} \right| \right)$$

- Shows **LC-divergence** explicitly  $\longrightarrow$  same divergence for **all T-odd twist-3 PDFs**, also in quark-target-model.
- Finite Box-graph contributions for twist-2 PDFs, especially  $f_{1T}^\perp, h_1^\perp$ .
- Regularization procedure: “Tree-level” predictions  $A_{LU} \propto g^\perp$  must be modified.
- Factorization theorem for twist-2 observables: (Ji, Ma, Yuan, 2004)

$$\frac{d\sigma_{UU}}{dx_B dy dz_h d^2 P_{h\perp}} \propto \int d^2 p_T d^2 k_T \left( \int d^2 l_T S(\vec{l}_T) \delta^{(2)}(\vec{p}_T - \frac{\vec{P}_{h\perp}}{z_h} - \vec{k}_T + \vec{l}_T) \right) f_1(x_B, \vec{p}_T^2) D_1(z_h, \vec{k}_T^2) + \dots$$

Soft factor  $S(\vec{l}_T)$  due to soft gluon radiation  $\longrightarrow$  modifies  $\delta$ -function.

PDFs/FFs: “non light-like” Wilson lines  $\longrightarrow$  “non-light-likeness” parameter  $\zeta = \sqrt{\frac{2(P \cdot v)^2}{v^2}}$

- Generalization of “**all-order factorization**” for twist-3 observables possible?

## Summary & Conclusions

### × Semi-inclusive DIS:

Transverse parton momentum has to be taken into account

→  $p_T$ -dependent parton distributions and fragmentation functions

→ Gauge link ensures **color gauge invariance** and enables **T-odd parton distributions**.

→ **Single-spin asymmetries**.

### × Analysis of the SIDIS-cross section:

Quark-Quark Correlators  $\Phi_{ij}(x, \vec{p}_T)$  und  $\Delta_{ij}(z, \vec{k}_T)$ .

→ **complete determination** of the twist-3 contribution to spin-observables in SIDIS in the parton-model using “tree-level” formalism.

### × Testing “tree-level” approximation:

Calculation of longitudinal asymmetries  $A_{UL}$  and  $A_{LU}$  in a diquark-spectator-model.

→ **(IR and collinear) finite** results.

Calculation of T-odd twist-3 parton distributions

→ **Light-cone divergences**.

→ **Modification** of existing twist-3 predictions necessary.