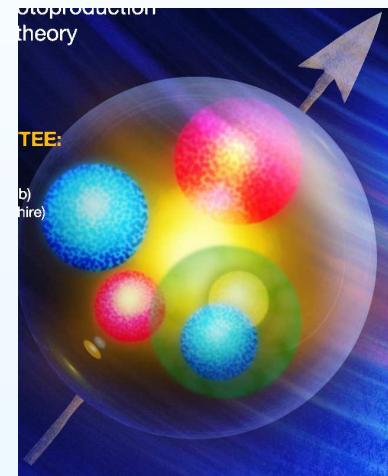


Neutron (^3He) Spin Structure Functions at Low Q^2



Vincent Sulkosky
Jefferson Laboratory

Spin Structure at Long Distance
March 12th 2009

Introduction

- Experiment E97-110:
 - Precise measurement of generalized GDH integral at low Q^2 , 0.02 to 0.3 GeV² for the neutron and ${}^3\text{He}$.
 - Cover an unmeasured region of kinematics to test rigorous theoretical calculations (Chiral Perturbation Theory).

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 - Precise measurement of **generalized GDH integral at low Q^2** , 0.02 to 0.3 GeV^2 for the **neutron** and **${}^3\text{He}$** .
 - Cover an **unmeasured region of kinematics** to **test rigorous theoretical calculations** (Chiral Perturbation Theory).
 - Data from **experiment E94-010** covered the transition region (0.1 to 0.9 GeV^2) from non-perturbative (mesons and baryons) to perturbative QCD (quarks and gluons).

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 - Precise measurement of generalized GDH integral at low Q^2 , 0.02 to 0.3 GeV^2 for the neutron and ${}^3\text{He}$.
 - Cover an unmeasured region of kinematics to test rigorous theoretical calculations (Chiral Perturbation Theory).
 - Data from experiment E94-010 covered the transition region (0.1 to 0.9 GeV^2) from non-perturbative (mesons and baryons) to perturbative QCD (quarks and gluons).
 - Preliminary results are now available and will be finalized in a few months.

Inclusive Cross Sections

- structure functions:

g_1 and g_2 (quark polarizations)

or σ_{TT} and σ_{LT}

$$\sigma_{TT} = \sigma_{1/2}(x, Q^2) - \sigma_{3/2}(x, Q^2)$$

- Polarized cross sections

$$\Delta\sigma_{||} = \frac{d^2\sigma_{\downarrow\uparrow}}{dE'd\Omega} - \frac{d^2\sigma_{\uparrow\uparrow}}{dE'd\Omega} = K \left[(E + E' \cos \theta) g_1(x, Q^2) - \left(\frac{Q^2}{\nu} \right) g_2(x, Q^2) \right]$$

$$\begin{aligned} \Delta\sigma_{\perp} &= \frac{d^2\sigma_{\downarrow\Rightarrow}}{dE'd\Omega} - \frac{d^2\sigma_{\uparrow\Rightarrow}}{dE'd\Omega} = KE' \sin \theta [g_1(x, Q^2) + \frac{2E}{\nu} g_2(x, Q^2)] \\ K &= \frac{4\alpha^2}{M\nu Q^2} \frac{E'}{E} \end{aligned}$$

$\downarrow \uparrow$ is for electron spin, $\uparrow \Rightarrow$ is for target spin direction

Gerasimov-Drell-Hearn (GDH) Sum Rule ($Q^2 = 0$)

$$I_{\text{GDH}} = \int_{\nu_{\text{th}}}^{\infty} \frac{\sigma_{\frac{1}{2}}(\nu) - \sigma_{\frac{3}{2}}(\nu)}{\nu} d\nu = -2\pi^2 \alpha \left(\frac{\kappa}{M} \right)^2$$

- Circularly **polarized photons** incident on a longitudinally polarized spin- $\frac{1}{2}$ target.
- $\sigma_{\frac{1}{2}}$ ($\sigma_{\frac{3}{2}}$) **photoabsorption cross section** with photon helicity parallel (anti-parallel) to the target spin.
- The sum rule is related to the **target's mass M** and **anomalous part of the magnetic moment κ** .
- Sum rules are solid theoretical predictions based on general principles.

GDH Measurements

The sum rule is **valid for any target** with definite spin- S .

	$M[\text{GeV}]$	Spin	κ	$I_{\text{GDH}}[\mu \text{ b}]$
Proton	0.938	$\frac{1}{2}$	1.79	-204.8
Neutron	0.940	$\frac{1}{2}$	-1.91	-233.2
Deuteron	1.876	1	-0.14	-0.65
Helium-3	2.809	$\frac{1}{2}$	-8.38	-498.0

- Proton sum rule was verified, Mainz, Bonn and LEGS.
- Measurements for the **neutron** are in progress.

See A. Sandorfi talk.

Generalized GDH Integral ($Q^2 > 0$)

$$I(Q^2) = \int_{\nu_{\text{th}}}^{\infty} \left[\sigma_{\frac{1}{2}}(\nu, Q^2) - \sigma_{\frac{3}{2}}(\nu, Q^2) \right] \frac{d\nu}{\nu}$$

$$\sigma_{1/2} - \sigma_{3/2} = \frac{8\pi^2\alpha}{MK} \left[g_1(\nu, Q^2) - \left(\frac{Q^2}{\nu^2} \right) g_2(\nu, Q^2) \right]$$

- Replace **photoproduction cross sections** with the corresponding **electroproduction cross sections**.
- The integral is related to the Compton scattering amplitude: $S_1(Q^2)$.

$$S_1(Q^2) = \frac{8}{Q^2} \int_0^1 g_1(x, Q^2) dx = \frac{8}{Q^2} \Gamma_1(Q^2)$$

X.-D. Ji and J. Osborne, J. Phys. **G27**, 127 (2001)

At $Q^2 = 0$, the **GDH sum rule is recovered**.

First moment of g_1 and g_2

$$\Gamma_1 = \int_0^1 g_1(x, Q^2) dx$$

$$\Gamma_2 = \int_0^1 g_2(x, Q^2) dx$$

- Γ_1 is closely related to generalized GDH integral as $Q^2 \rightarrow 0$.
- g_2 is suppressed at very low Q^2 .

Bjorken Sum Rule ($Q^2 \rightarrow \infty$)

- g_A is the nucleon axial charge.
- The sum rule has been confirmed to 10%.

$$\Gamma_1^p - \Gamma_1^n = \frac{g_A}{6}$$

J.D. Bjorken, Phys. Rev. 148, 1467 (1966)

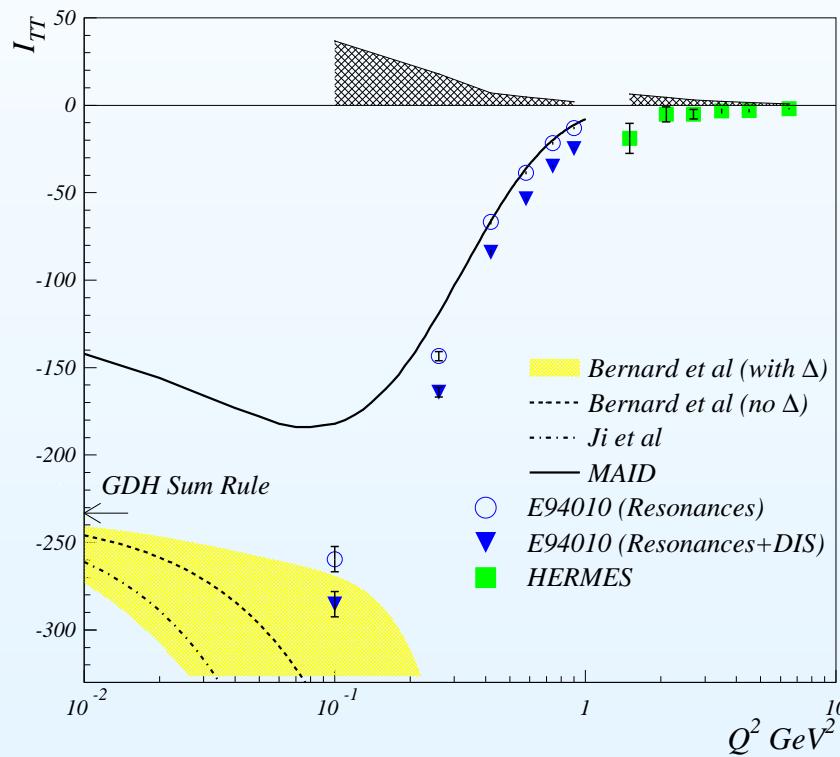
Importance of the Generalized GDH Sum Rule



- Constrained at the two ends of the Q^2 spectrum by known sum rules: GDH ($Q^2 = 0$) and Bjorken ($Q^2 \rightarrow \infty$).
- Generalized GDH Integral is **calculable at any Q^2** .
- Compare theoretical predictions to experimental measurements over the **entire Q^2 range**.
- Tool to **study non-perturbative QCD**, while starting on known theoretical grounds (pQCD).

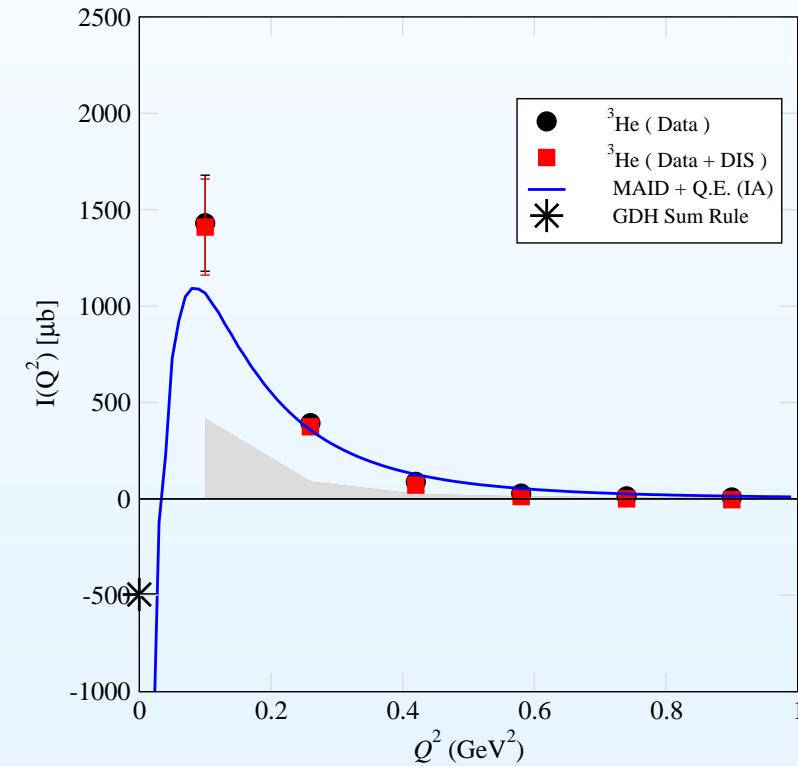
Hall A Neutron GDH Published Results

Neutron



M. Amarian *et al.*, PRL 89, 242301 (2002)

Helium-3

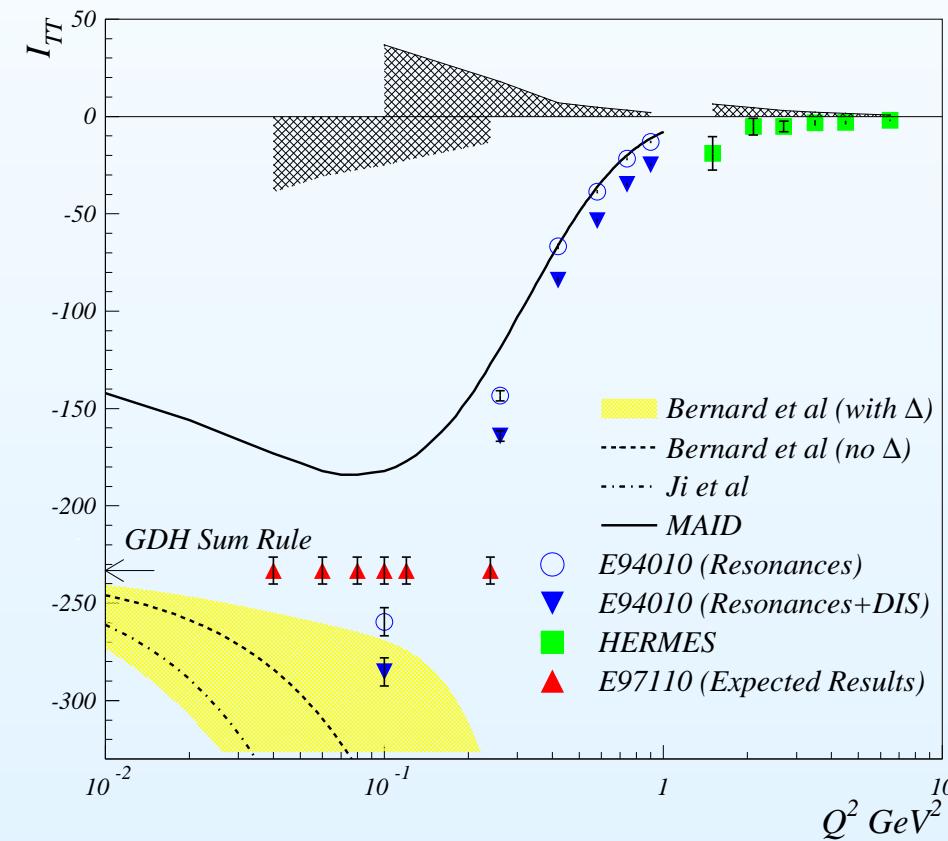


K. Slifer *et al.*, PRL 101, 022303 (2008).

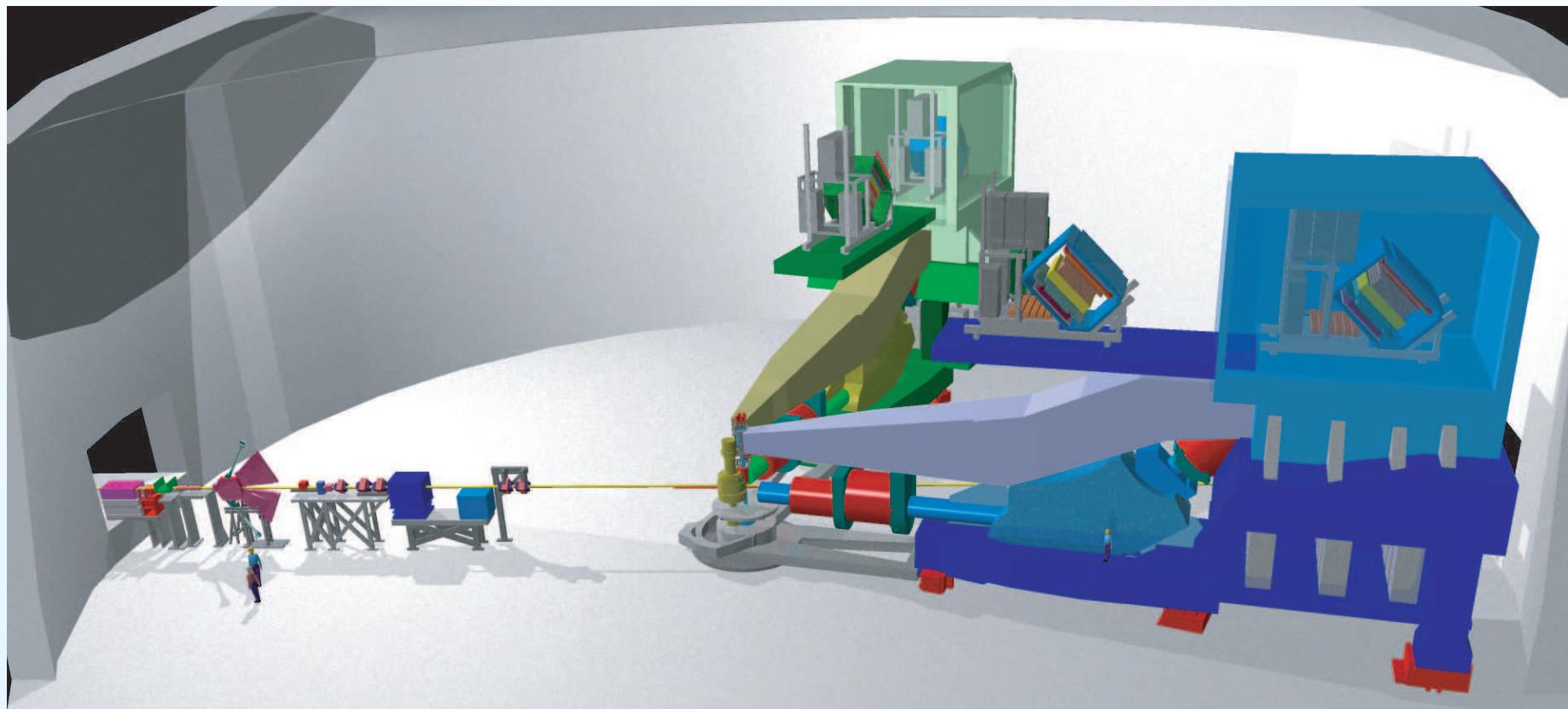
Experiment E97-110

Precise measurement of generalized GDH integral at low Q^2 , 0.02 to 0.3 GeV^2

- Ran in spring and summer 2003
- Inclusive experiment: ${}^3\text{He}(\vec{e}, e')X$
 - ⇒ Scattering angles of 6° and 9°
 - ⇒ Polarized electron beam:
 $\langle P_{\text{beam}} \rangle = 75\%$
 - ⇒ Pol. ${}^3\text{He}$ target (para & perp):
 $\langle P_{\text{targ}} \rangle = 40\%$
- Measured polarized cross-section differences

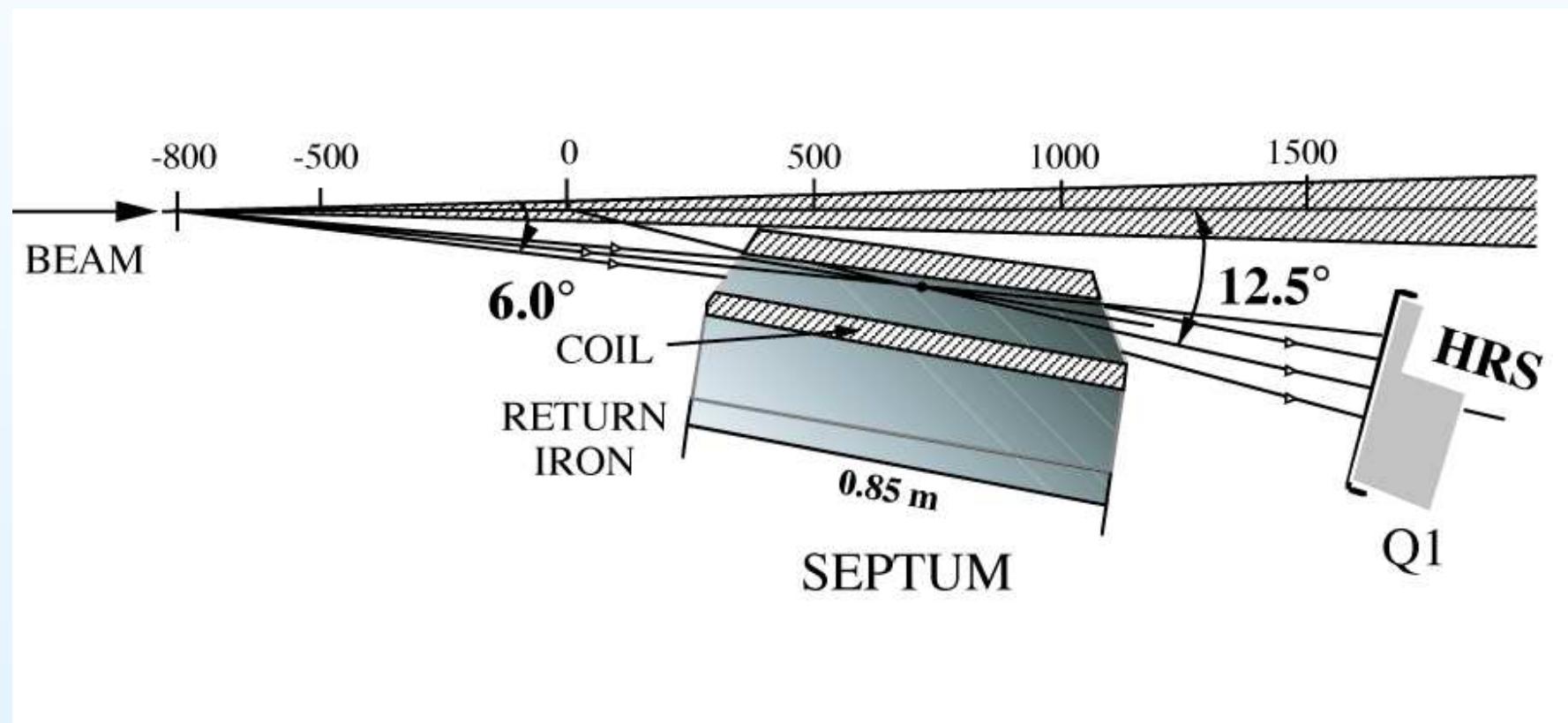


Experimental Setup

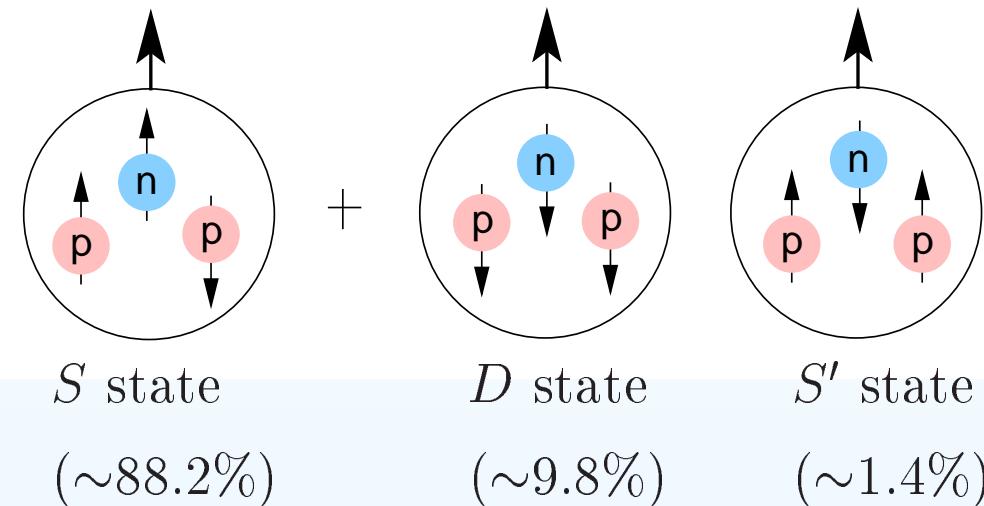


New Bending Magnet

- Low Q^2 requires forward angles.
- Minimum spectrometer angle is 12.5° .



^3He as an Effective Polarized Neutron Target



$$P_n = 86\% \text{ and } P_p = -2.8\%$$

J.L. Friar *et al.*, PRC **42**, (1990) 2310

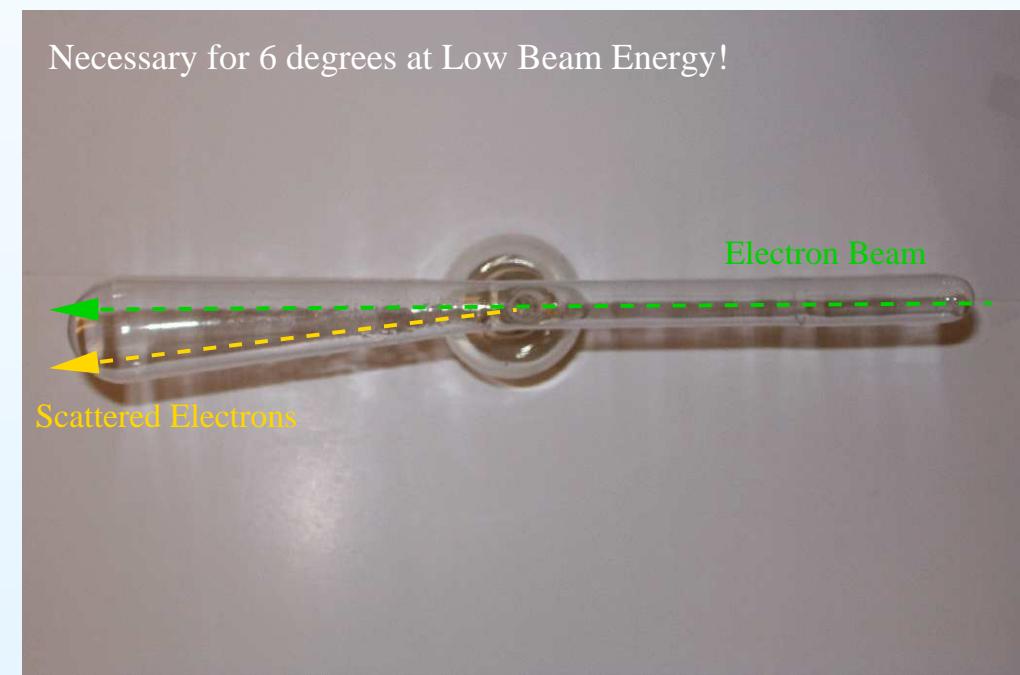
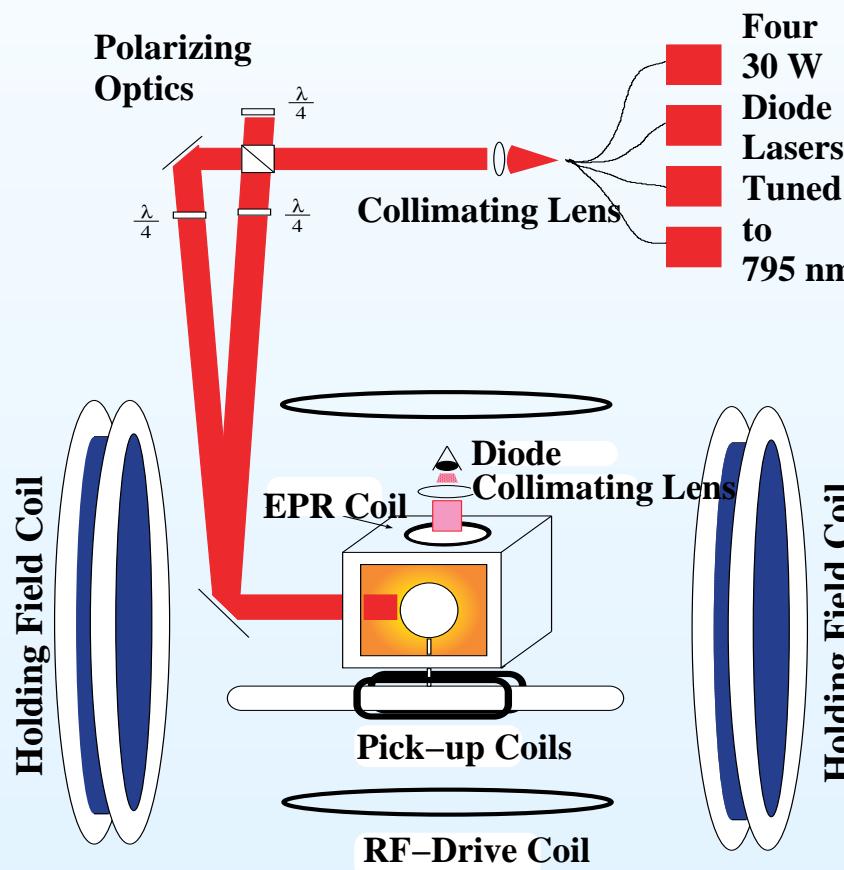
Extraction of Neutron Results

$$\Gamma_1^n(Q^2) = \frac{1}{P_n} [\Gamma_1^{^3\text{He}}(Q^2) - 2P_p\Gamma_1^p(Q^2)]$$

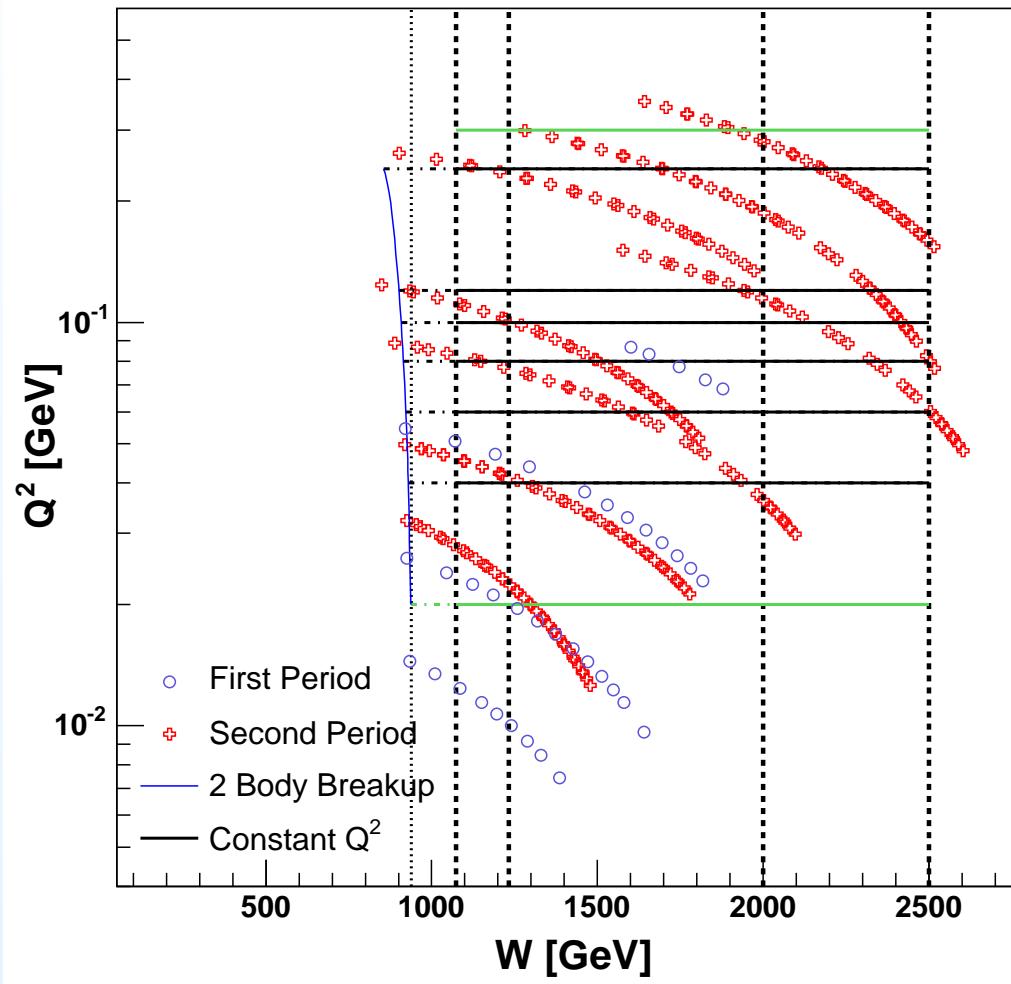
C. Ciofi degli Atti & S. Scopetta, PLB **404**, (1997) 223

Polarized ^3He System

- Both **longitudinal** and **transverse** configurations.
- Two independent polarimетries: **NMR** and **EPR**.

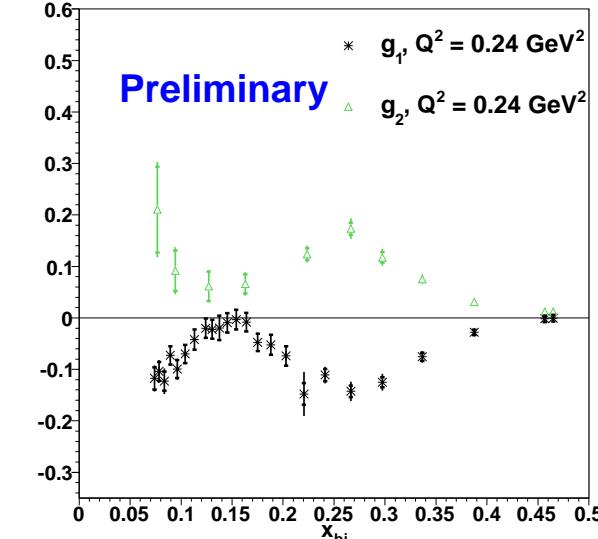
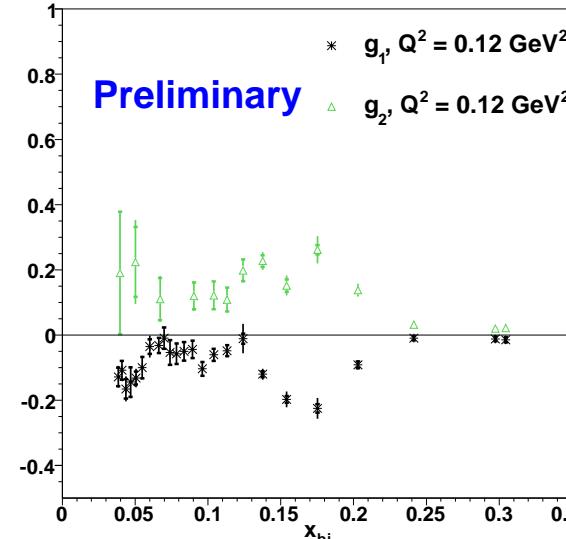
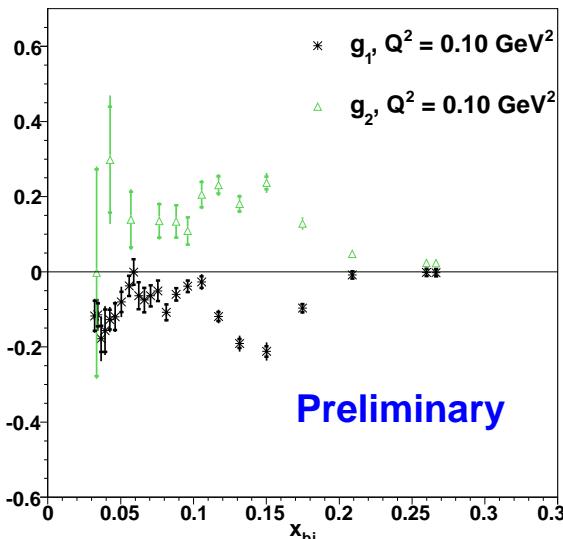
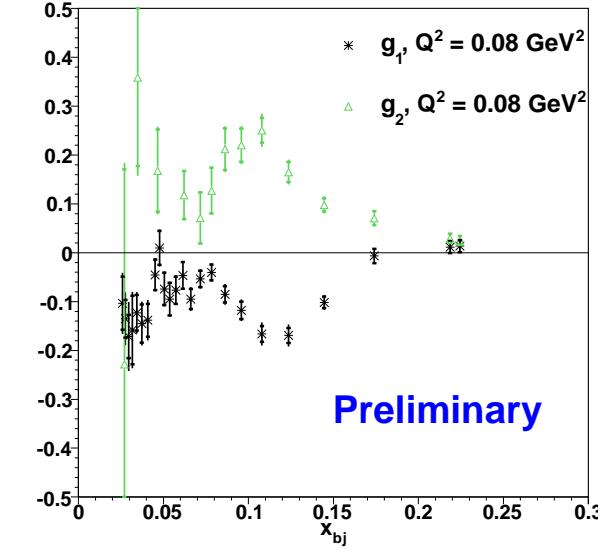
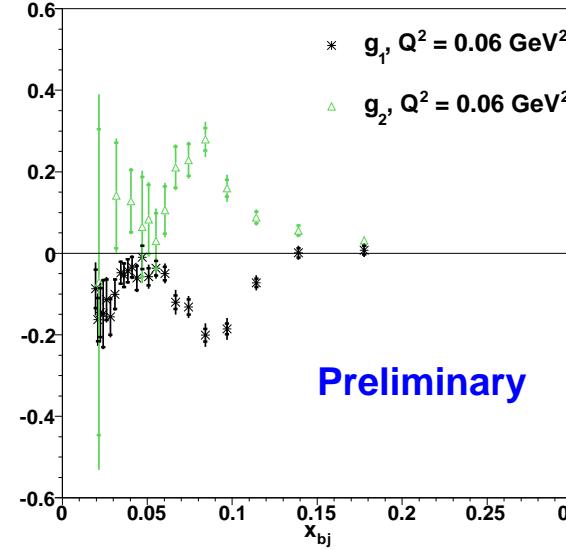
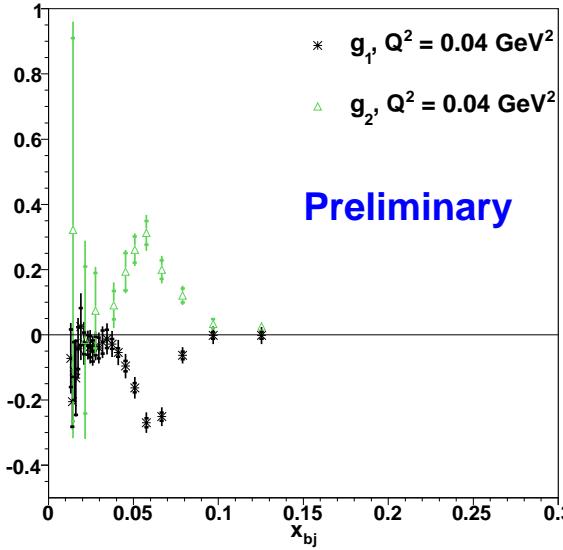


Kinematic Coverage and Interpolation

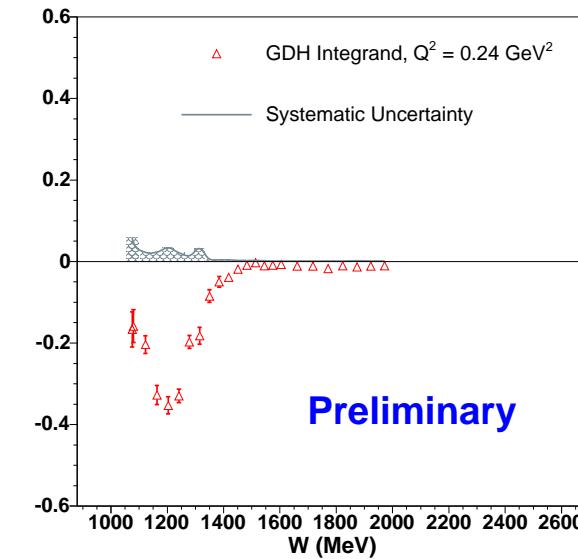
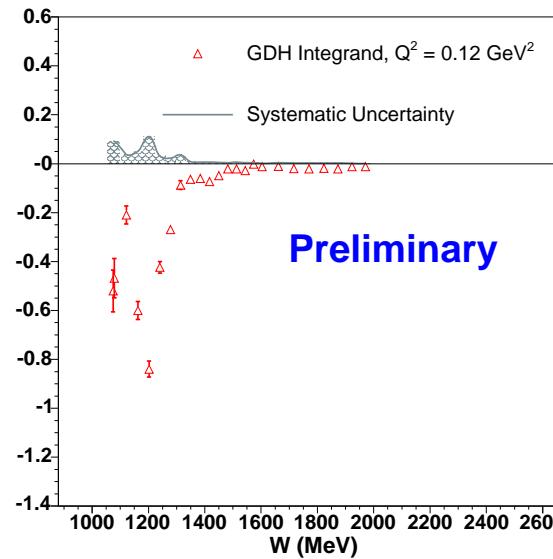
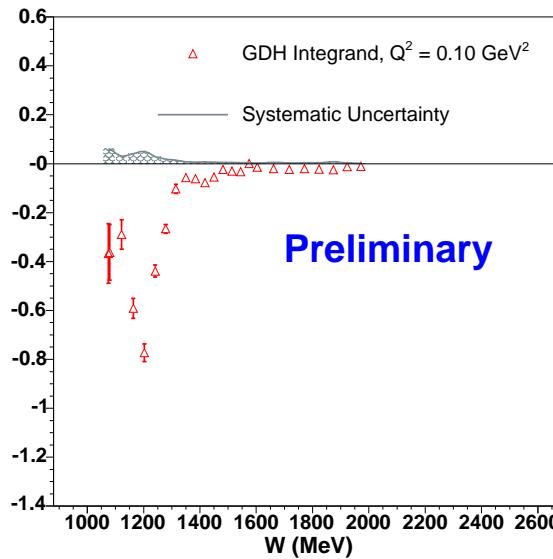
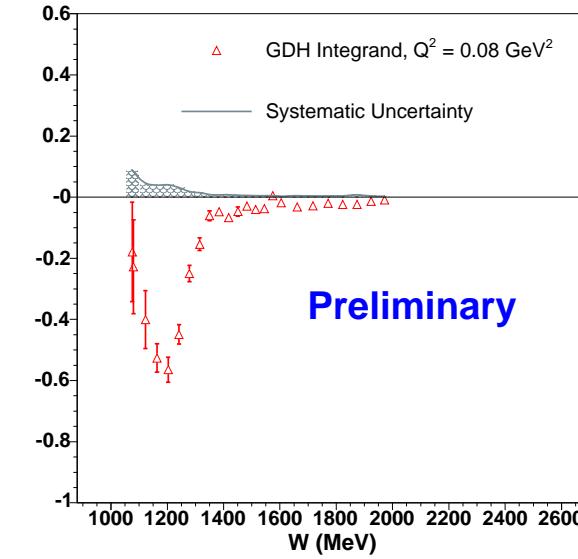
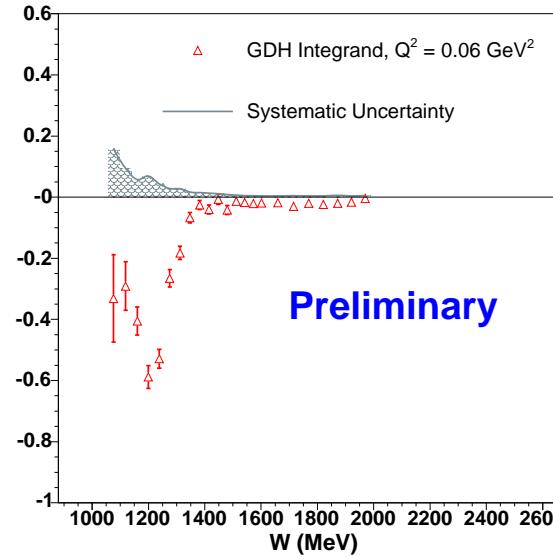
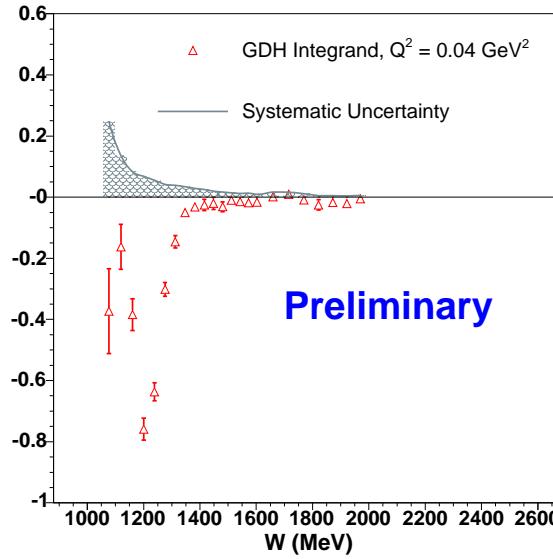


Six constant Q^2 points: 0.04, 0.06, 0.08, 0.1, 0.12 and 0.24 GeV^2 .

${}^3\text{He}$ - g_1, g_2 versus x at constant Q^2

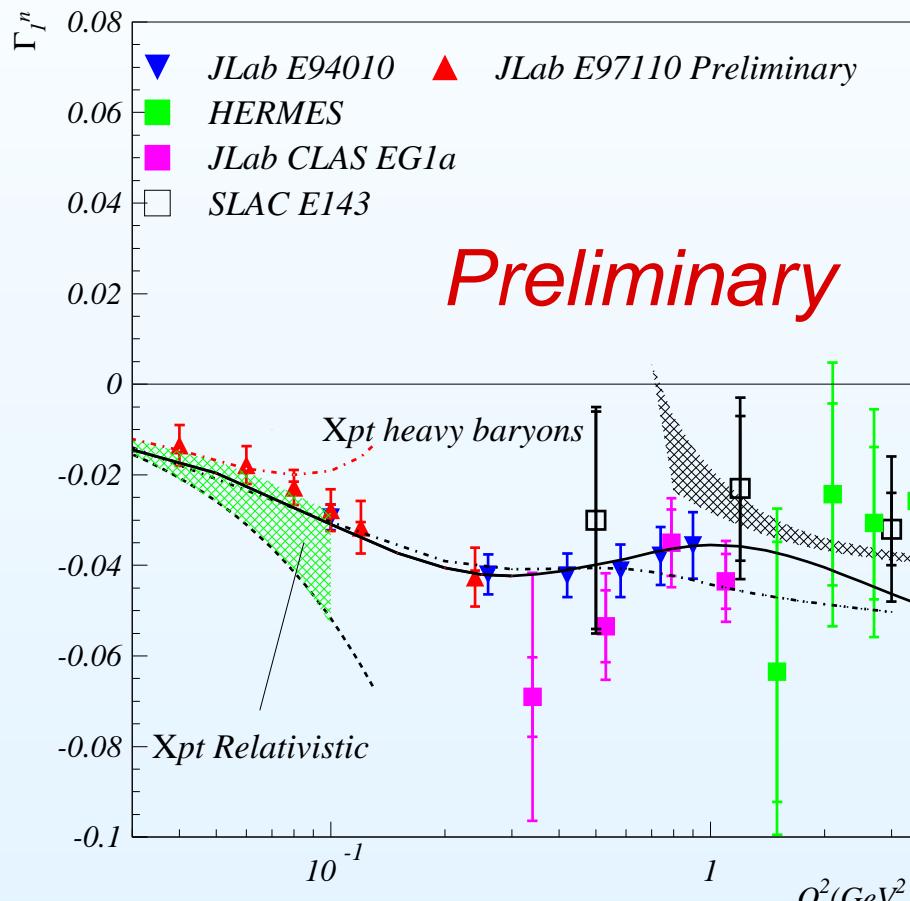


${}^3\text{He} - \frac{\sigma_{TT}}{\nu}$ versus W at constant Q^2



Γ_1^n : First Moment of g_1

$$\Gamma_1 = \int_0^1 g_1(x, Q^2) dx$$

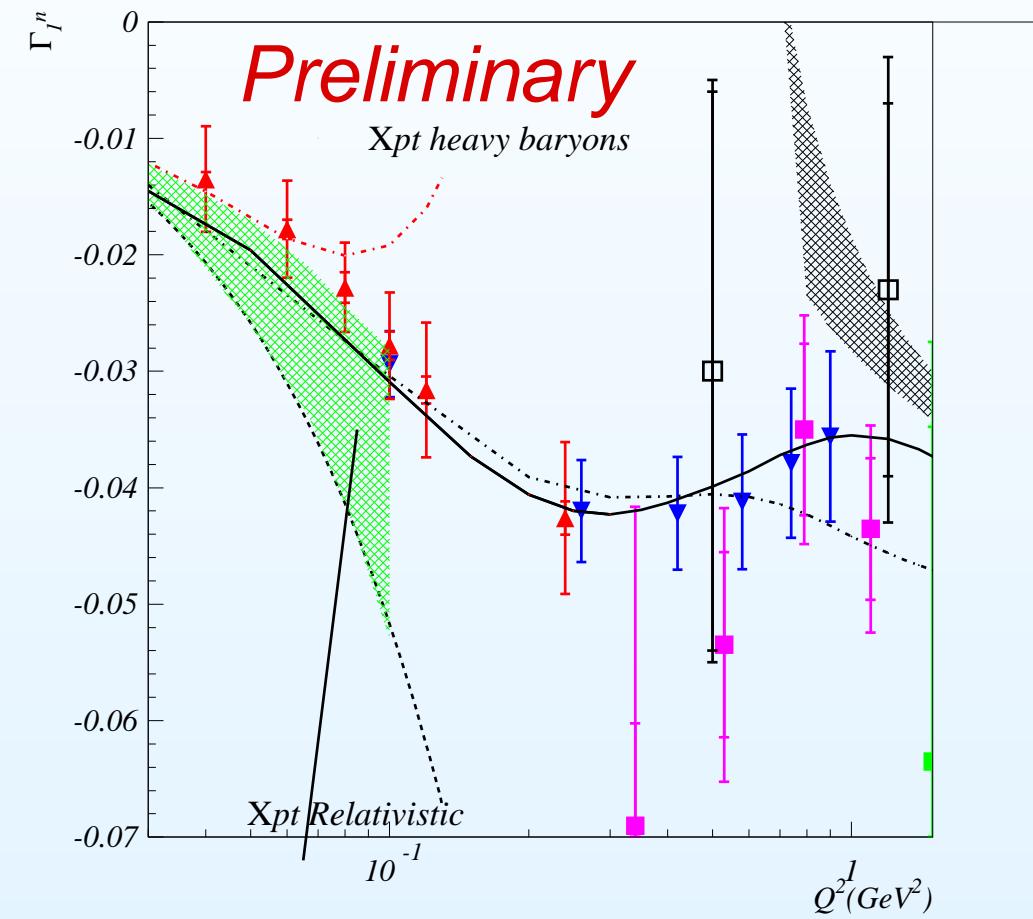
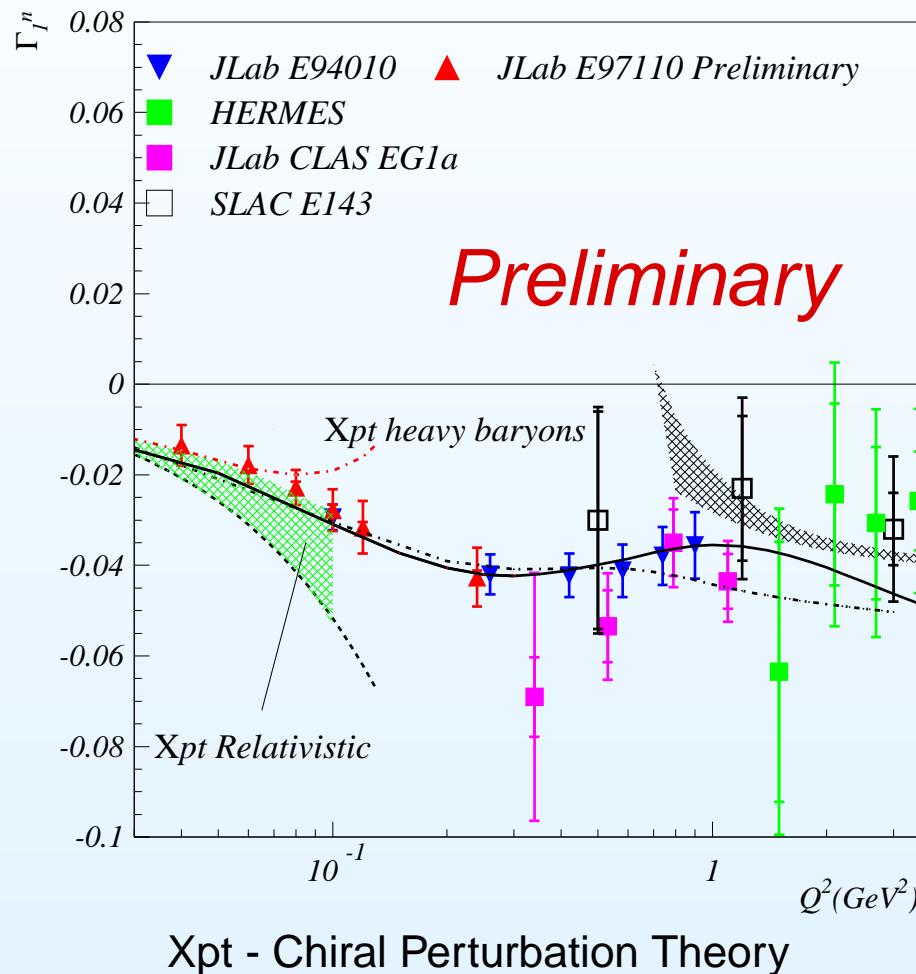


Xpt - Chiral Perturbation Theory

Preliminary

Γ_1^n : First Moment of g_1

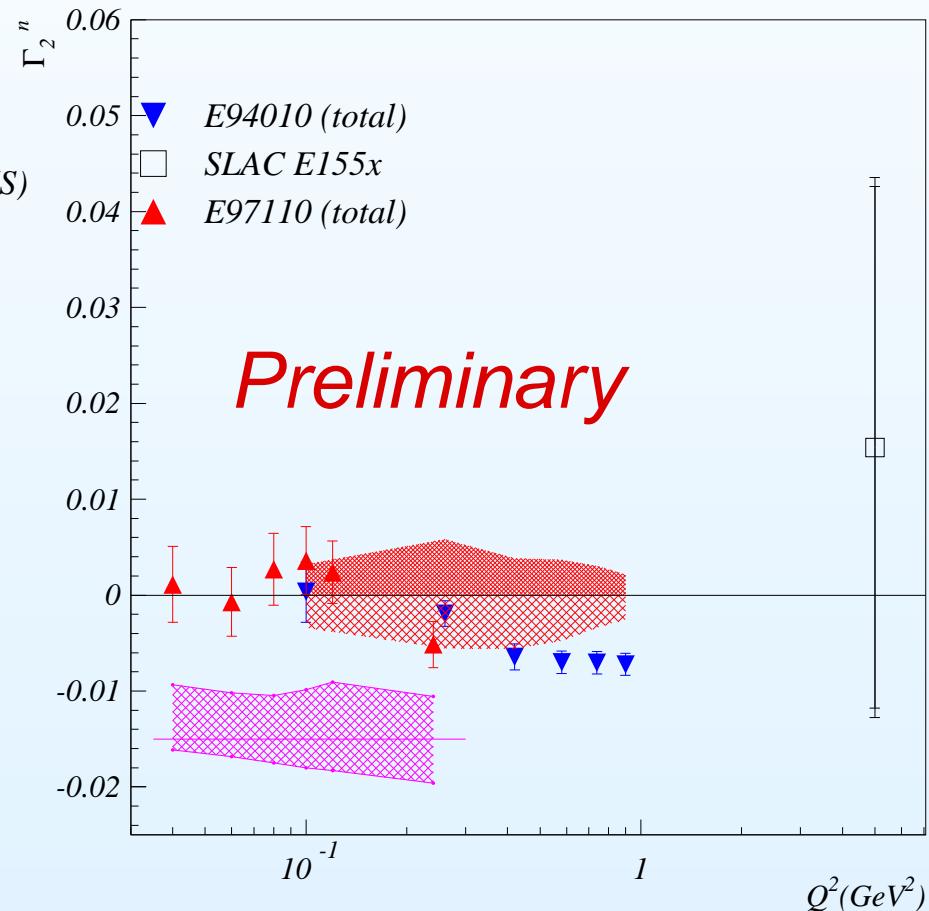
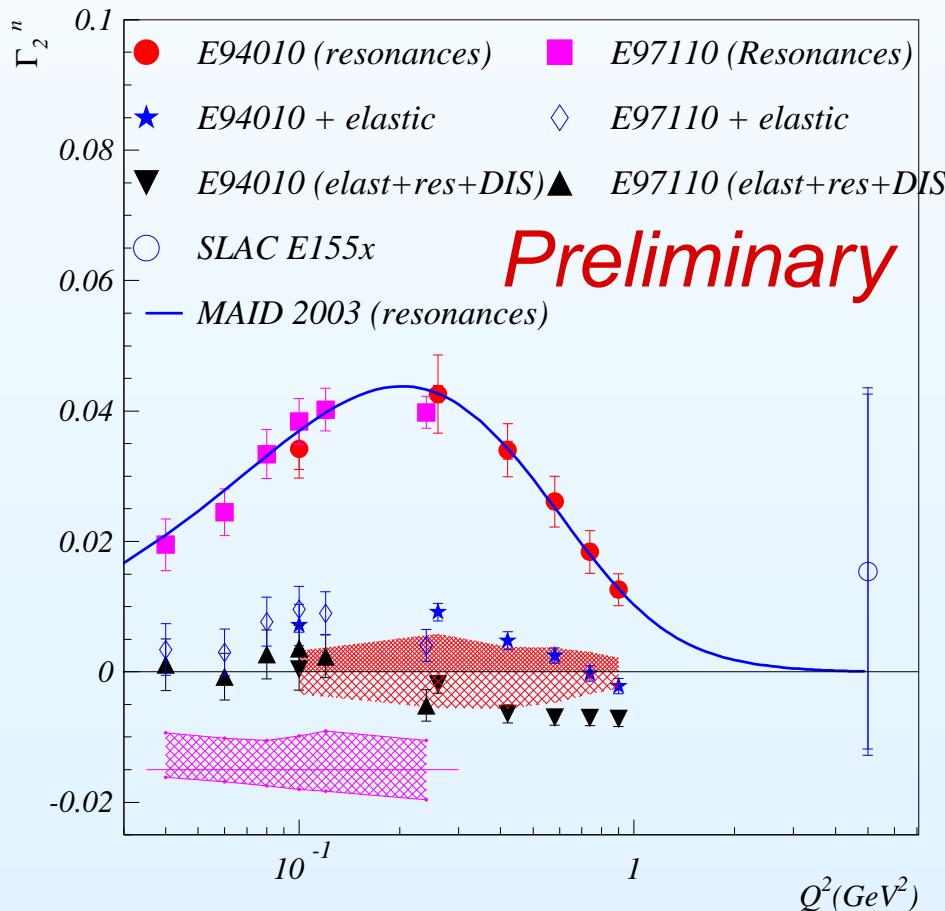
$$\Gamma_1 = \int_0^1 g_1(x, Q^2) dx$$



Γ_2^n : First Moment of g_2

$$\Gamma_2^n(Q^2) = \int_0^1 g_2(x, Q^2) dx = 0$$

Burkhardt-Cottingham Sum Rule



Summary and Conclusion

- The GDH integral is an important tool that can be used to study nucleon spin structure over the full Q^2 range:
 - in particular, the transition from **perturbative QCD** to **nonperturbative QCD**.
- Experiment E97-110 provides precision data for **moments of spin structure functions at low Q^2** : 0.02 to 0.3 [GeV/c]²
- Preliminary results of the **the neutron moments are available** and work is in progress to finalize the systematic effects.
- These data provide a **precision test of Chiral Perturbation Theory calculations** at a Q^2 where they are expected to be valid.
- Expect **final neutron results soon**.

Systematic Uncertainties

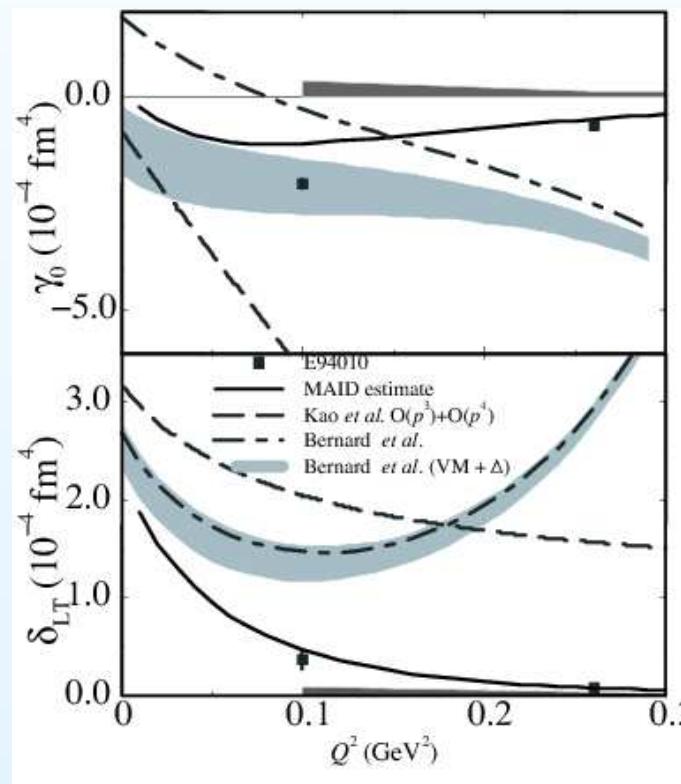
Source	Systematic Uncertainty		
Angle	6°	9°	3.775 GeV, 9°
Target density		2.0%	
Acceptance/Effects	5.0%	5.0%	15.0%
VDC efficiency	3.0%	2.5%	2.5%
Charge		1.0%	
PID Detector and Cut effs.		< 1.0%	
$\delta\sigma_{\text{raw}}$	6.4%	6.2%	15.5%
Nitrogen dilution		0.2–0.5%	
$\delta\sigma_{\text{exp}}$	6.5%	6.3%	15.5%
Beam Polarization		3.5%	
Target Polarization		7.5%	
Radiative Corrections*	20% (40% for $Q^2 \leq 0.08$)		
Total on $\Delta\sigma$	10.5%	10.4%	17.6%

* Radiative correction $\approx 5\text{--}10\%$ in delta region

Future Prospects: Spin Polarizabilities

$$\gamma_0 = \frac{16\alpha M^2}{Q^6} \int_0^{x_0} x^2 \left(g_1 - \frac{4M^2}{Q^2} x^2 g_2 \right) dx$$

$$\delta_{LT} = \frac{16\alpha M^2}{Q^6} \int_0^{x_0} x^2 (g_1 + g_2) dx$$



M. Amarian *et al.*, PRL 93, 152301 (2004)

The E97-110 Collaboration

S. Abrahamyan, K. Aniol, D. Armstrong, T. Averett, S. Bailey,
P. Bertin, W. Boeglin, F. Butaru, A. Camsonne, G.D. Cates,
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A. Vacheret, E. Voutier, K. Wang, L. Wan, B. Wojtsekowski,
S. Woo, H. Yao, **J. Yuan**, X. Zheng, L. Zhu

and the Jefferson Lab Hall A Collaboration

Extra Slides

Constant Q^2 Interpolation and Integral Extraction

Procedure:

- First interpolate to constant W for each energy.
- Second interpolation with respect to Q^2 .
- Integrals formed from $W = 1073$ GeV to 2000 GeV.
- We could **use our own data above $W = 2000$ GeV**.
- DIS contribution included up to $W = \sqrt{1000}$ using **Thomas and Bianchi parameterization**.
- Neutron extraction performed using calculation from Scopetta and Ciofi degli Atti paper for $Q^2 \geq 0.1$ GeV 2 .
- $Q^2 < 0.1$ GeV 2 use **effective polarization technique** (difference $\sim 5\text{--}10\%$).

Inclusive Electron Scattering

Energy transfer:

$$\nu = E - E'$$

4-momentum transfer squared:

$$\vec{q} = \vec{k} - \vec{k}'$$

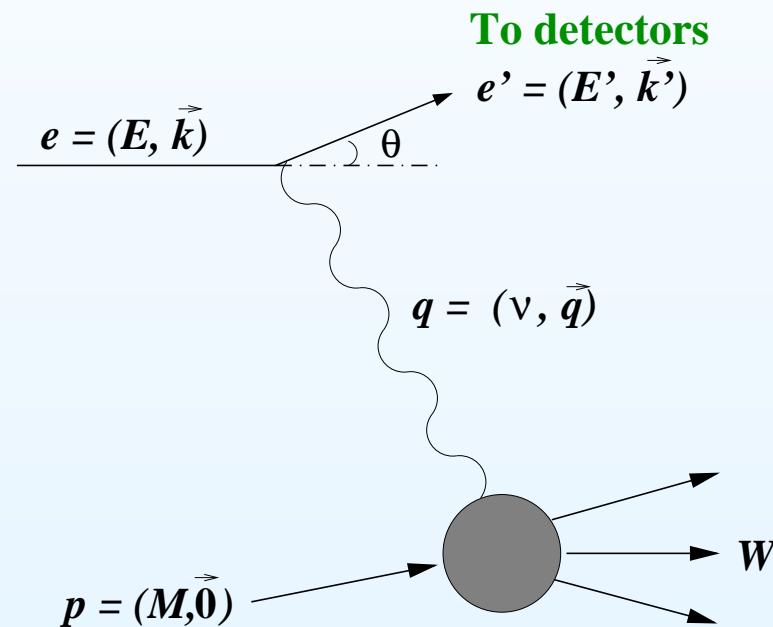
$$Q^2 = -q^2 = 4EE' \sin^2 \frac{\theta}{2}$$

Invariant Mass:

$$W^2 = M^2 + 2M\nu - Q^2$$

Bjorken variable:

$$x = \frac{Q^2}{2M\nu}$$



GDH Derivation for Real Photons

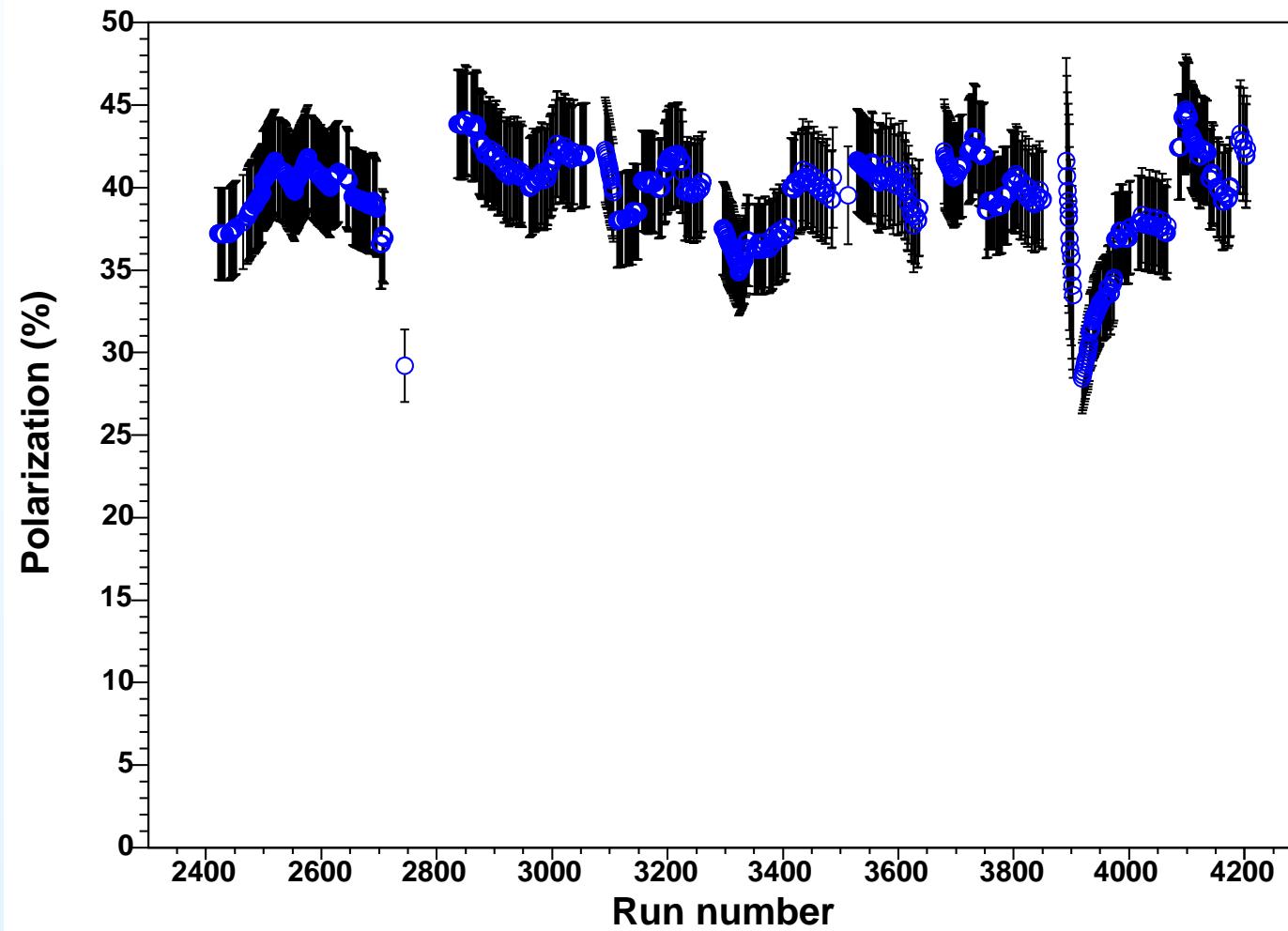
- Begin with the spin dependent part of the forward Compton amplitude, S_1
- Use the following dispersion relation and three assumptions:

$$\text{Re } S_1(\nu) = \frac{2\nu}{\pi} \int_{\nu_{\text{th}}}^{\infty} d\nu' \frac{\text{Im } S_1(\nu')}{\nu'^2 - \nu^2}$$

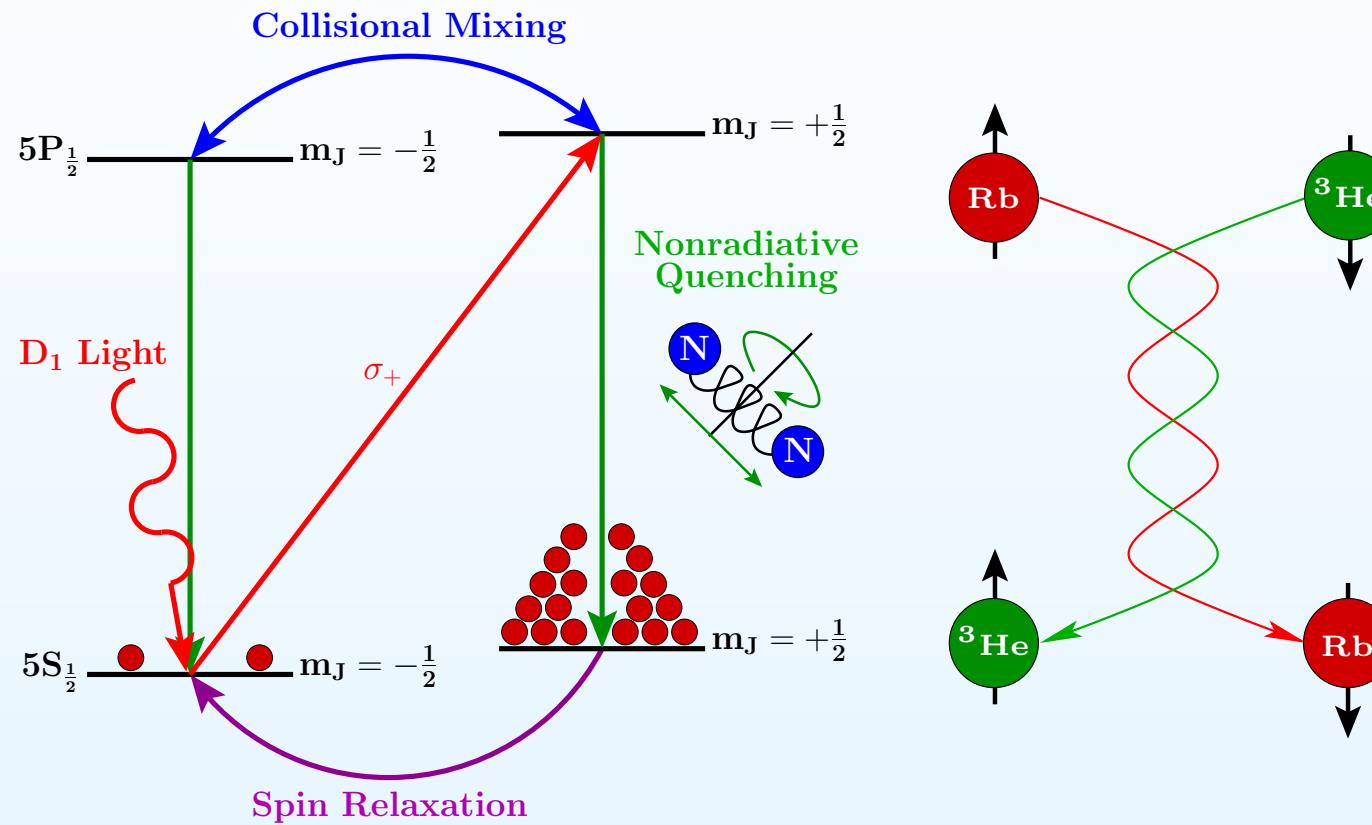
- Optical Theorem: $\text{Im } S_1(\nu) = \frac{\nu}{8\pi} \sigma_{TT}(\nu)$
- Low Energy Theorem: $\text{Re } S_1(\nu) = -\frac{e^2 \kappa^2}{8\pi M^2} \nu$
- Unsubtracted Dispersion Relation: assumption is convergence of the dispersion integral.

$$I_{\text{GDH}} = \int_{\nu_{\text{th}}}^{\infty} \frac{\sigma_{\frac{1}{2}}(\nu) - \sigma_{\frac{3}{2}}(\nu)}{\nu} d\nu = -2\pi^2 \alpha \left(\frac{\kappa}{M}\right)^2$$

Preliminary Target Polarization



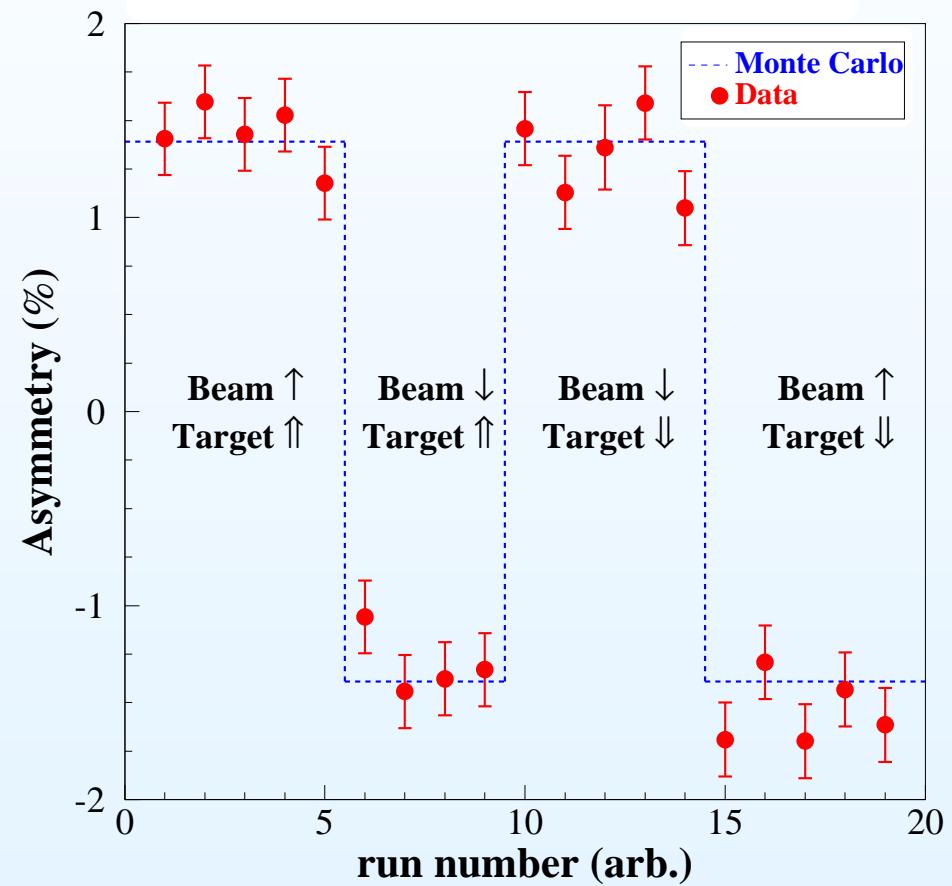
Spin Exchange Optical Pumping



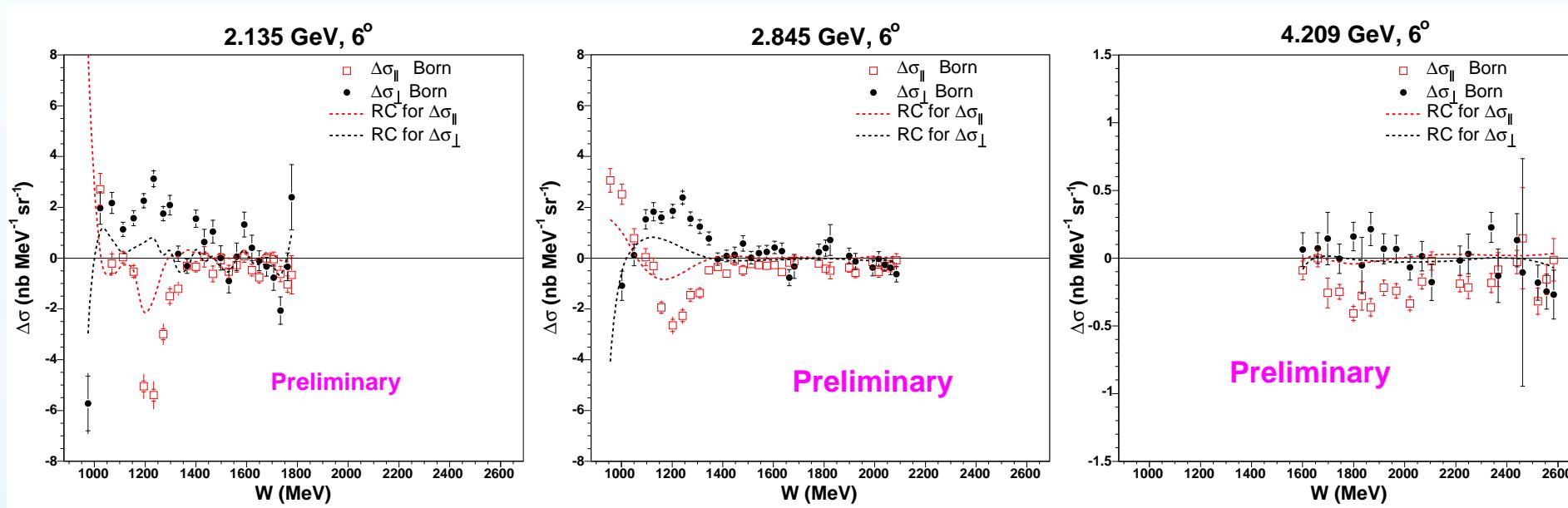
3He nucleus is polarized via **spin-exchange** with optically pumped Rb atoms.

^3He Elastic Asymmetry

- Monte Carlo prediction: 1.390%
- Preliminary data analysis:
 $(1.403 \pm 0.044)\%$ (stat. only)
 $\chi^2/N_{\text{dof}} = 1.08.$
- Four target and beam configurations
- For seven out of the twelve beam energies, elastic data were acquired.



Cross Section Differences



Cross Section Differences

