Does QCD predict light hybrid mesons?

Jozef Dudek Old Dominion University & Jefferson Lab

Robert Edwards (JLab) Balint Joo (JLab) David Richards (JLab) Christopher Thomas (JLab) Mike Peardon (Trinity Coll., Dublin)

for the Hadron Spectrum Collaboration

JLab Users' Group Workshop

(historical) empirical motivation

1₄(2040)





















but what if excited gluonic fields play a rôle - a hybrid meson, $q\bar{q}G$?

possibly exotic *J^{PC}* & extra 'non-exotic' states

must be 'heavier' or 'harder to produce' ?



tubes, bags & heavy glue















two-point correlator

$$C_{ij}(t) = \left\langle 0 \left| \mathcal{O}_i(t) \mathcal{O}_j^{\dagger}(0) \right| 0 \right\rangle$$

e.g.
$$\mathcal{O}(t) = \sum_{\vec{x}} \left(\bar{\psi} \gamma^5 \psi \right)_{\vec{x},t}$$





two-point correlator matrix

 $C_{11}(t) \quad C_{12}(t) \quad C_{13}(t)$ $\begin{array}{ccc} C_{21}(t) & C_{22}(t) & C_{23}(t) \\ C_{31}(t) & C_{32}(t) & C_{33}(t) \end{array}$





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optimal operator : $\Omega^{(\mathfrak{n})} = \sum_i v_i^{(\mathfrak{n})} \mathcal{O}_i$

solve by using a 'Rayleigh-Ritz'-style variational approach - "diagonalize the matrix"

each state comes from an orthogonal combination of \mathcal{O}_i



fermion bilinears with up to three covariant derivatives - project into good **J**^{PC}





this is by far the largest operator set ever used in a calculation like this

isovector spectrum



isovector spectrum



isovector spectrum







exotic J^{PC} - before



exotic J^{PC} - after



exotic J^{PC} - after



systematics of a $q\bar{q}$ pair



systematics of a $q\bar{q}$ pair



try (model-dependent) analysis of matrix elements

$$\left[Z_i^{\mathfrak{n}} \equiv \left\langle \mathfrak{n} \middle| \mathcal{O}_i \middle| 0 \right\rangle \right]$$





$$\left(\pi \times D_{J=1}^{[2]} \right)^{J=1} \qquad \begin{array}{c} D_{J=1}^{[2]} \equiv \langle 1, m_1; 1, m_2 | 1, m \rangle \overleftarrow{D}_{m_1} \overleftarrow{D}_{m_2} \\ \sim [D, D] \sim F \end{array}$$

 $ho \qquad \gamma_i$



$$\left[Z_i^{\mathfrak{n}} \equiv \left\langle \mathfrak{n} \middle| \mathcal{O}_i \middle| 0 \right\rangle\right]$$













$$\left[Z_i^{\mathfrak{n}} \equiv \left\langle \mathfrak{n} \middle| \mathcal{O}_i \middle| 0 \right\rangle \right]$$

$$egin{aligned} & \rho & {}^3S_1 \ & \left(
ho imes D_{J=2}^{[2]}
ight)^{J=1} \;\; {}^3D_1 \ & \left(\pi imes D_{J=1}^{[2]}
ight)^{J=1} \;\; ext{hybrid}? \end{aligned}$$



$$\left[Z_i^{\mathfrak{n}} \equiv \left\langle \mathfrak{n} \middle| \mathcal{O}_i \middle| 0 \right\rangle \right]$$

$$\begin{aligned} \rho & {}^{3}S_{1} \\ \left(\rho \times D^{[2]}_{J=2}\right)^{J=1} & {}^{3}D_{1} \\ \left(\pi \times D^{[2]}_{J=1}\right)^{J=1} & \text{hybrid?} \end{aligned}$$





$$\left[Z_i^{\mathfrak{n}} \equiv \left\langle \mathfrak{n} \middle| \mathcal{O}_i \middle| 0 \right\rangle \right]$$

$$\begin{pmatrix} \rho & {}^{3}S_{1} \\ \left(\rho \times D_{J=2}^{[2]} \right)^{J=1} & {}^{3}D_{1} \\ \left(\pi \times D_{J=1}^{[2]} \right)^{J=1} & \text{hybrid}?$$



$$Z_i^{\mathfrak{n}} \equiv \left\langle \mathfrak{n} \Big| \mathcal{O}_i \Big| 0 \right\rangle$$

the lightest hybrid supermultiplet ?



Does QCD predict light hybrid mesons?



Does QCD predict light hybrid mesons?



* Light hybrid mesons may be observed only for heavy quark masses, risk of large hadronic widths has not been determined. The Hadron Spectrum Collaboration offers no guarantee that they can be produced in photoproduction.



exotic

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calculate at lighter quark masses

increased computational cost

just a matter of time ...

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observing states as decaying resonances

scattering of composite objects in non-perturbative field theory !

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observing states as decaying resonances

Isospin=2 $\pi\pi$ scattering $\underline{}_{0.8} \quad k^2 \,/\, \mathrm{GeV}^2$ 0.4 0.5 0.6 0.7 0.1 0.2 0.3 ° 🛓 -10 ⊥_ IJ Î -20 Ĩ₽₽₽ Į^ĮĮ^II ∎ t ⁴ I4 ₹ ₹ -30 o / -40 δ_{0} -50 -60 $524\,\mathrm{MeV}$ $444\,\mathrm{MeV}$ $396\,\mathrm{MeV}$ Hoogland Losty Cohen Durusoy 16³ ----- $16^3 - ---$ axpt 16^{3} – 20³ — ------70 20^{3} 20^{3} -0- 24^3 - $_{0.8}$ $k^2 / \,{
m GeV}^2$ 0.1 0.2 0.3 0.4 0.5 0.6 0.7 <mark>₫</mark>Î 4 $\delta_2 / ^{\circ}$ -10 $524\,\mathrm{MeV}$ $444\,\mathrm{MeV}$ $396\,\mathrm{MeV}$ Hoogland Losty Cohen $16^3 - -$ expt. $16^3 - ----$ 20³ — - --15 20³ — - - 20^3 — — Durusoy 24^3 ____

scattering of composite objects in non-perturbative field theory !



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determining coupling to photons is relatively straighforward



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"Isoscalar meson spectroscopy from lattice QCD" - arXiv:1102.4299 [hep-lat] (2011) (PRD in press) "The phase-shift of isospin-2 ππ scattering from lattice QCD" - PRD.83.071504 (2011) "Toward the excited meson spectrum of dynamical QCD" - PRD.82.034508 (2010) "Highly excited and exotic meson spectrum from dynamical lattice QCD" - PRL.103.262001 (2009) "A novel quark-field creation operator construction for hadronic physics in lattice QCD" - PRD.80.054506 (2009)







isoscalar (with light & strange) :



diagonalising gives the $\,\ellar{\ell},sar{s}\,$ mixing

very few results due to the difficulty of calculation

C. Michael et al (UKQCD, 2001) [heavy quarks, 2-flavour theory] found f_1/a_1 , b_1/h_1 , ρ/ω splittings consistent with zero

ETMC (2009) [2-flavour, extrap. to phys. quark mass] ρ/ω splitting of 27(10) MeV







isoscalar spectrum



isoscalar spectrum



isoscalar spectrum







but we can't be satisfied with this ...



excited states should be **resonances**

enhancements in the mesonmeson scattering continuum



in finite volume only **discrete meson-meson** states

but we aren't seeing them !



0.6

 16^{3}

 20^{3}

 $\frac{16^3}{\text{increasing}}$

volume

 20^{3}



parity doubling & chiral symmetry restoration ?



 ρ_3

 a_3

 h_3



 a_4

 η_4









 ho_4



cubic complications ...

integer spin not a good quantum number

restricted rotational symmetry of a cube

0, 4... A_1 $T_1 \mid 1, 3, 4...$ $T_2 \mid 2, 3, 4...$ $E \mid 2, 4 \dots$ A_2 3...

cubic complications ...

integer spin not a good quantum number

restricted rotational symmetry of a cube





integer spin not a good quantum number

restricted rotational symmetry of a cube

 1^{++}

 0^{++}



 2^{++}



integer spin not a good quantum number

restricted rotational symmetry of a cube

0, 4... A_1 T_1 1, 3, 4... $T_2 \mid 2, 3, 4...$ E| 2, 4... A_2 3...





operators respect cubic symmetry, but are 'preconditioned' to be *J*-diagonal

... but does it work in practice ?







clear dominance of a single spin for each state