

Nucleon Spin Structure

Kees de Jager



Nucleon Spin Structure

- Introduction
- Q^2 -evolution of GDH integral
- Nucleon Spin Structure at large x
- Quark-Hadron (Spin) Duality
- $R = \sigma_L / \sigma_T$ in Resonance Region
- Real Compton Scattering up to Large t -values
- Summary

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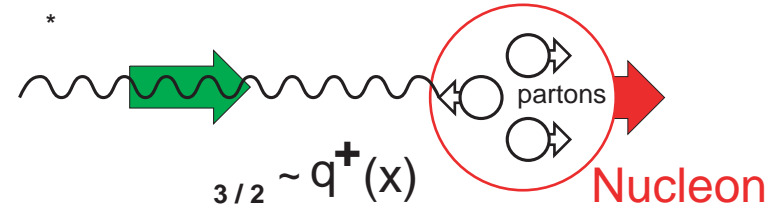
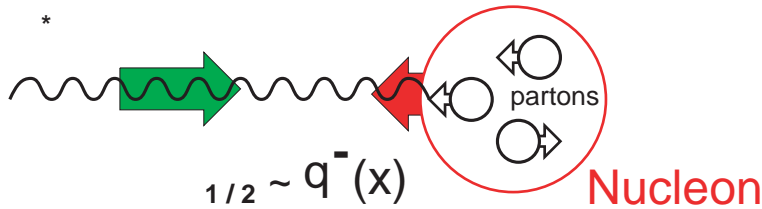
Spin Structure in Deep Inelastic Scattering

$$\langle S_z^N \rangle = \frac{1}{2} = J_q + J_g = \frac{1}{2} \Delta\Sigma + L_q + \Delta G + L_g$$

- Study of nucleon spin structure started with EMC (~'88):
 $\Delta\Sigma = \Delta d + \Delta u + \Delta s = 0.12 \pm 0.17$
- World-wide effort (SLAC, DESY, CERN) established that $\Delta\Sigma \sim 0.2 \dots 0.4$
- Focus shifted to other contributions to $\langle S_z \rangle$:
 - strange sea polarization semi-inclusive DSA
 - gluon polarization open charm, high p_T hadron pairs
 - orbital angular momentum Generalized Parton Distributions



Spin Structure in Deep Inelastic Scattering



Partonic Interpretation

x = fraction of nucleon momentum
carried by struck quark
 $q^{+/-}$ quark helicity parallel/antiparallel
to photon helicity
NO simple partonic picture of g_2

$$x = \frac{Q^2}{2M\nu}$$

$$q(x) = q^+(x) + q^-(x)$$

$$F_1(x) = \frac{1}{2} \sum_{flavor} e_f^2 q_f(x)$$

$$F_2(x) = xF_1(x)$$

$$\Delta q(x) = q^+(x) - q^-(x)$$

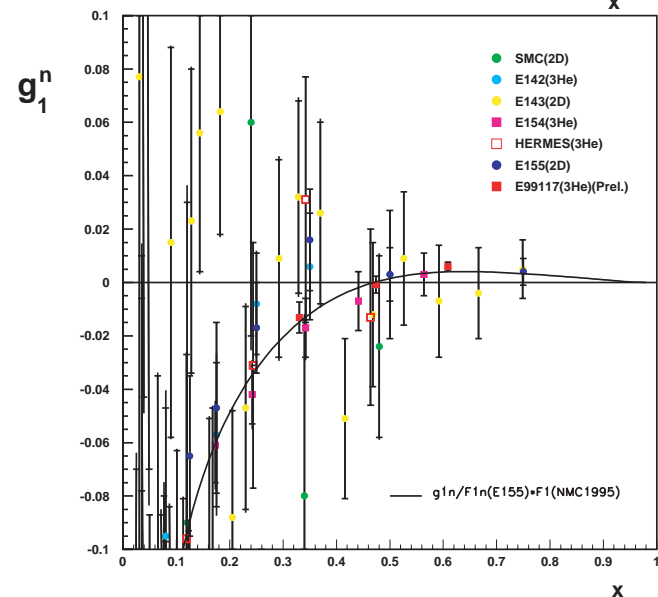
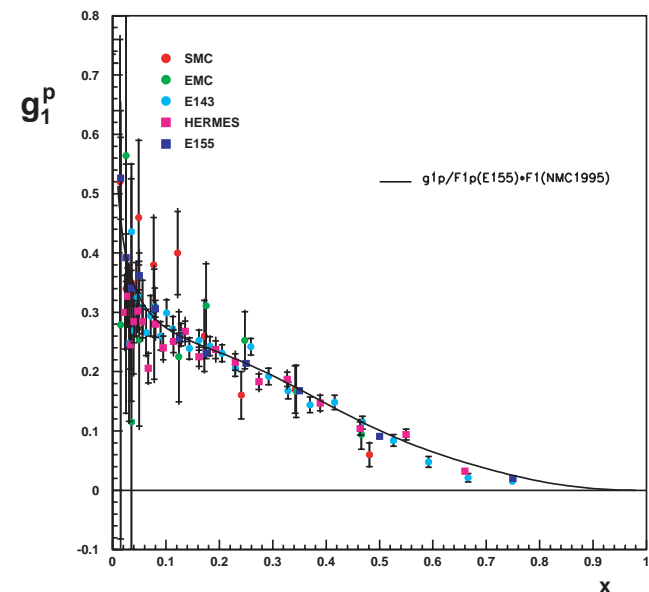
$$A_1(x) = \frac{g_1(x)}{F_1(x)}$$

$$g_1(x) = \frac{1}{2} \sum_{flavor} e_f^2 \Delta q_f(x)$$

$$g_2(x) = 0$$

Nucleon Spin Structure

- Fairly extensive data set on $g_1(x)$, but only for $x < 0.3$
- Low- x region dominated by **sea quarks**, predictions difficult
- High- x region (a single quark carries most of the nucleon momentum) dominated by **valence quarks**, predictions feasible
- Accurate data will allow selection of models
- JLab unique combination of energy and luminosity



Bjorken Sum Rule

$$\Gamma_1^p(Q^2) - \Gamma_1^n(Q^2) = \int \{g_1^p(x, Q^2) - g_1^n(x, Q^2)\} dx = \frac{1}{6} g_A C_{NS}$$

$$g_A = 1.2601 \pm 0.0025$$

C_{NS}

neutron β -decay coupling constant
 Q^2 -dependent QCD correction

Basic assumptions

- Isospin symmetry
- Current Algebra or Operator Product Expansion within QCD

Present status (at $Q^2 = 5 \text{ (GeV/c)}^2$)

Experiment $0.176 \pm 0.003 \pm 0.007$

Theory 0.182 ± 0.005

- Combined world data are consistent with the Bjorken Sum Rule at 5 % level



Gerasimov-Drell-Hearn Sum Rule

$$\int_{\nu_{in}}^{\infty} (\sigma_{1/2}(\nu) - \sigma_{3/2}(\nu)) \frac{d\nu}{\nu} = -\frac{2\pi^2 \alpha_{EM}}{M^2} K^2$$

GDH Sum Rule

- The Gerasimov-Drell-Hearn Sum Rule (at $Q^2 = 0$) is a fundamental test of the relation between the nucleon resonance excitation and its anomalous magnetic moment
- Rests on basic physics principles (**Lorentz invariance, gauge invariance, unitarity**) and on dispersion relation applied to forward Compton amplitude
- Technical developments have only recently allowed first measurement of GDH integral for the proton up to 800 MeV (Mainz) and up to 3 GeV (Bonn)
- Results agree with sum rule with assumptions for contributions at higher energies
- Many facilities (GRAAL, SPring-8, LEGS, HIGS, JLab) geared for extensive studies
- Will include data for the neutron from polarized deuterium



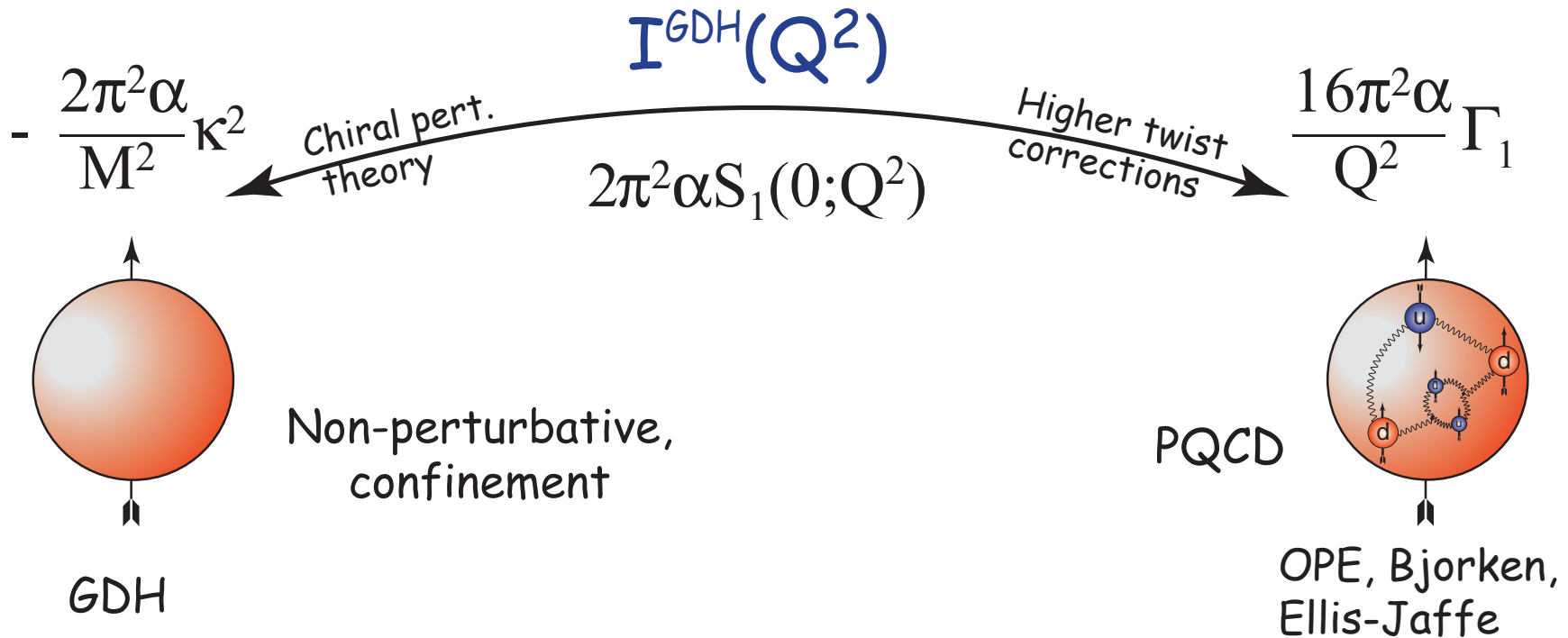
Transition from Strong to Perturbative QCD

- Dispersion relations on Compton Scattering Amplitudes lead to extension of GDH sum rule valid at all Q^2 (Ji and Osborne)
- Q^2 -evolution of Gerasimov-Drell-Hearn Sum Rule provides quantitative measure of transition from resonance (strong QCD) to DIS (pQCD) regimes
- Transition from Bjørken sum rule down to $\sim 1 \text{ GeV}^2$ can be predicted using Operator Product Expansion of higher twist contributions
- Transition from GDH sum rule up to $\sim 0.1 \text{ GeV}^2$ can be predicted using Chiral Perturbation Theory
- For intermediate region one awaits Lattice QCD calculations



Why is $I^{\text{GDH}}(Q^2)$ interesting?

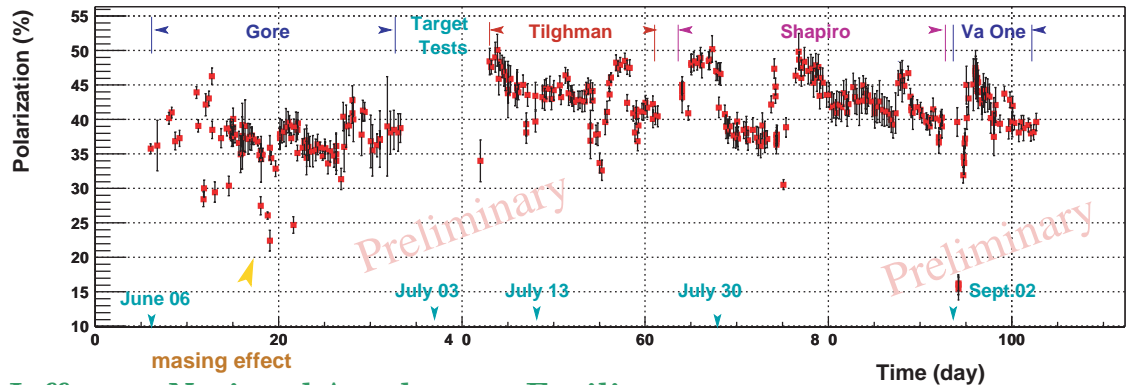
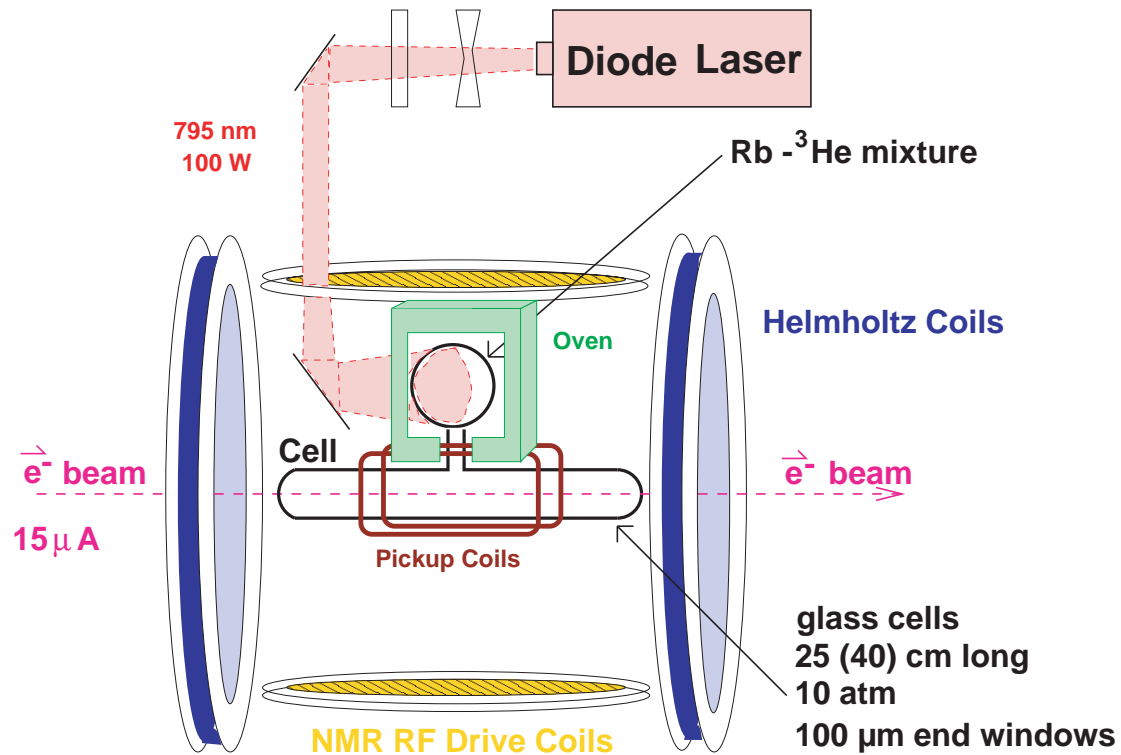
One of the few opportunities to "zoom out" from tiny length scales (DIS) to large length scales



Many of the underlying assumptions are the same as those being tested in high-energy spin-structure tests

Polarized ^3He Target (Hall A)

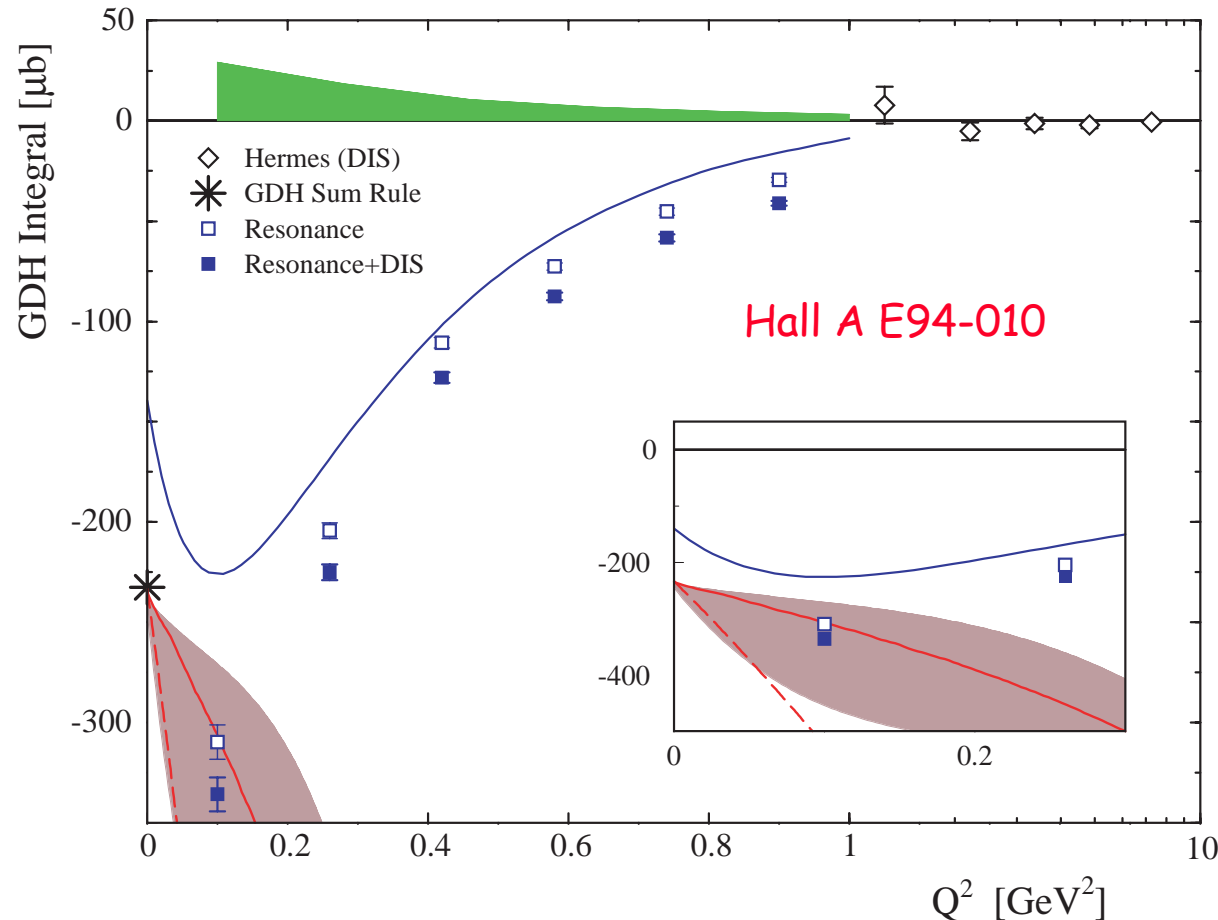
- Polarized ^3He is best approximation of polarized neutron $P_n=87\%$ and $P_p=2.7\%$
- Requires corrections for nuclear medium, investigated by many theorists
- Basic principle:
Optical pumping of Rb, followed by polarization transfer to ^3He through spin-exchange collisions
- Target polarization measured by EPR/NMR



Q²-Evolution of the Gerasimov-Drell-Hearn Integral

$$I_{GDH} = \int_{\nu_{in}}^{\infty} (\sigma_{\nu_2}(\nu, Q^2) - \sigma_{\nu_1}(\nu, Q^2)) \frac{d\nu}{\nu} = -2 \int_{\nu_{in}}^{\infty} \sigma_{TT} \frac{d\nu}{\nu} = -\frac{8\pi^2\alpha}{MK} \int_{\nu_{in}}^{\infty} \left(g_1 - \frac{Q^2}{\nu^2} g_2 \right) \frac{d\nu}{\nu}$$

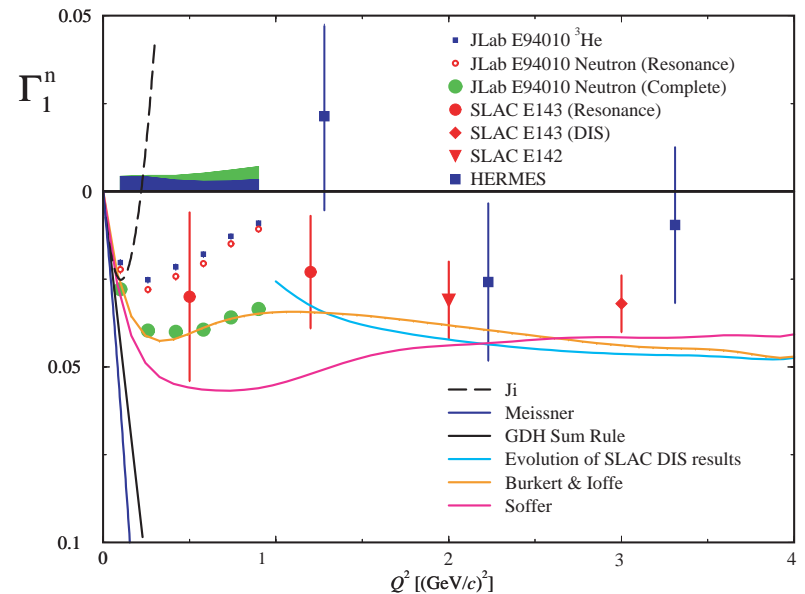
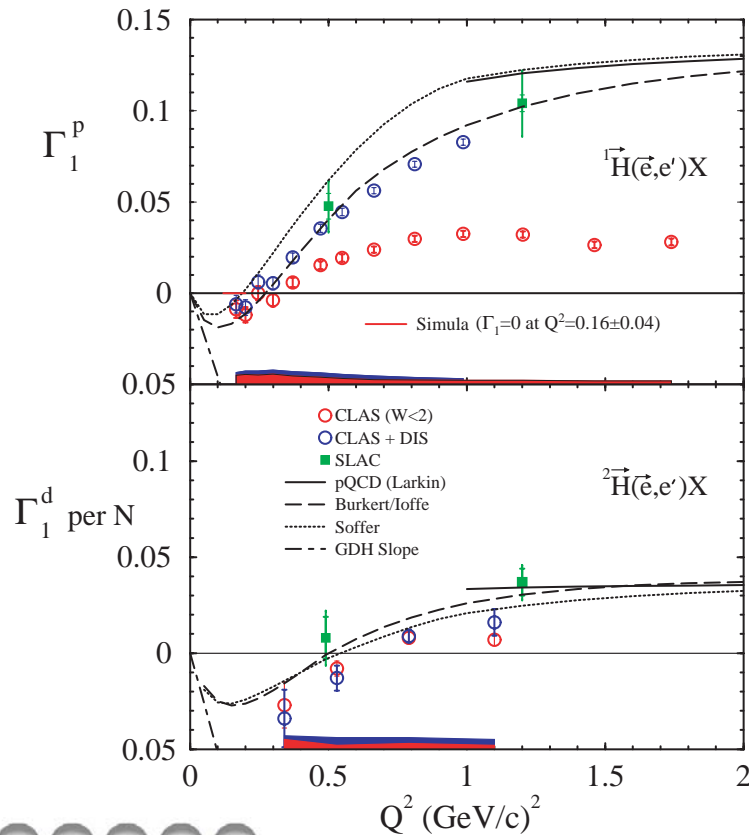
- Longitudinal and transverse target polarization allows separation of g_1 and g_2
- Kinematic coverage sufficient to integrate to $W \approx 2 \text{ GeV}$
- Nuclear medium corrections from Ciofi degli Atti and Scopetta
- Compared to calculations by Drechsel et al. which neglect contributions from DIS and by Ji and Bernard based on Chiral Perturbation Theory (band shows uncertainty in contribution from Δ -resonance)



Q²-Evolution of the Gerasimov-Drell-Hearn Integral (cont.)

$$\Gamma_1^{p/n} = \int_0^1 g_1^{p/n}(x_{Bj}, Q^2) dx_{Bj} \propto \int_0^1 (A_{//} + \alpha A_{\perp}) F_1(x_{Bj}, Q^2) dx_{Bj}$$

Hall B E91-023



The neutron A_1^n spin structure function

- Naïve SU(6) predictions:

$$A_1^p = 5/9, A_1^n = 0$$

- Broken SU(6):

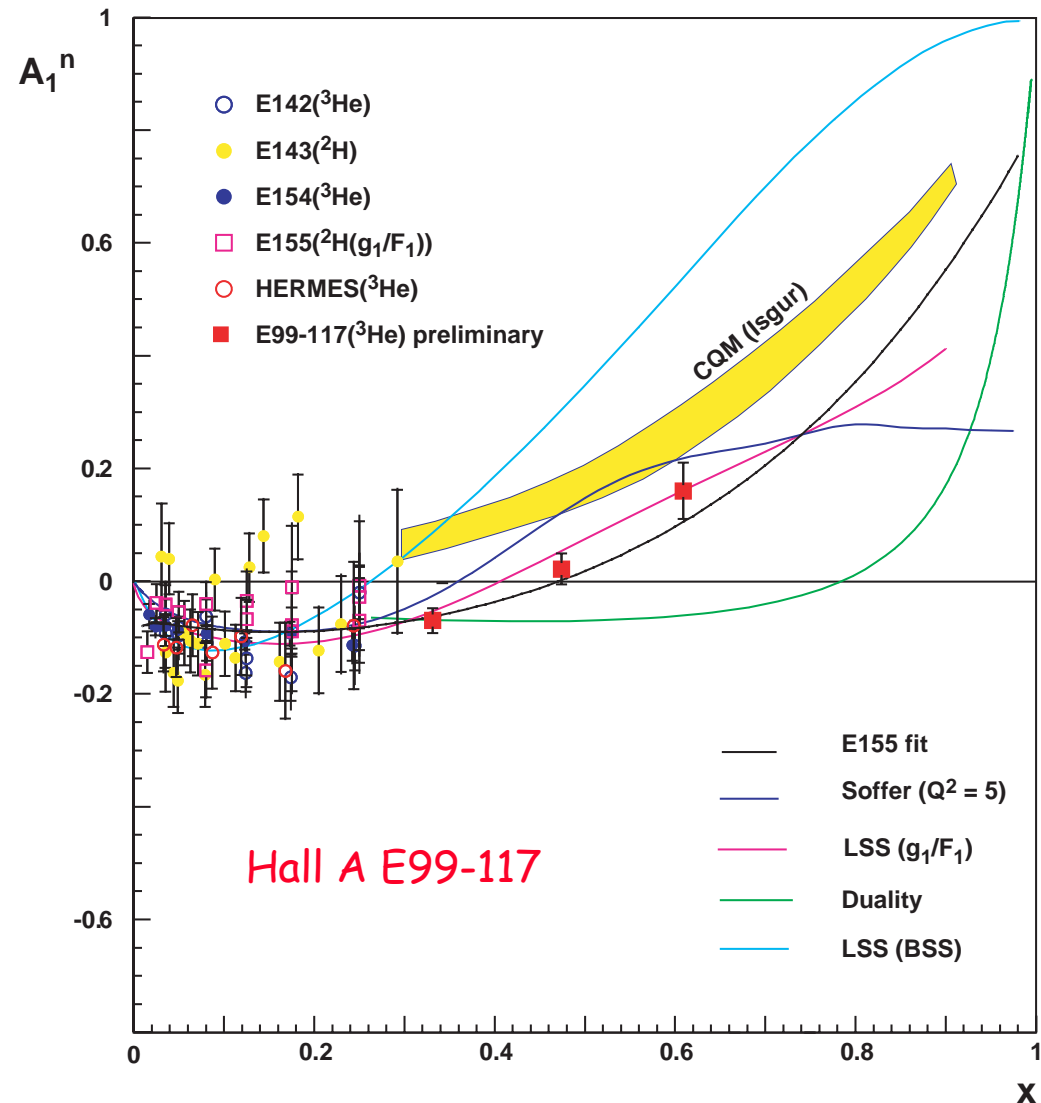
$$A_1^n \rightarrow 1 \text{ as } x \rightarrow 1$$

CQM: hyperfine
perturbed with simple
model for d/u

LSS: NLO polarized
parton densities

Soffer: global NLO
analysis of
(un)polarized DIS

Duality: local QHD using
available G_E, G_M data

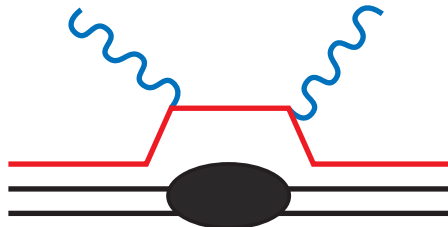


Quark-Gluon Correlations

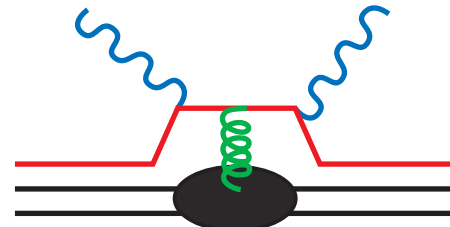
$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 g_1(y, Q^2) \frac{dy}{y}$$

- In simple partonic picture $g_2(x)=0$
- Wandzura and Wilczek have shown that g_2 can be written in two parts:
 - one given by g_1 in twist-2 contributions
 - the other originating solely from quark-gluon correlations (twist-3)



Twist-2



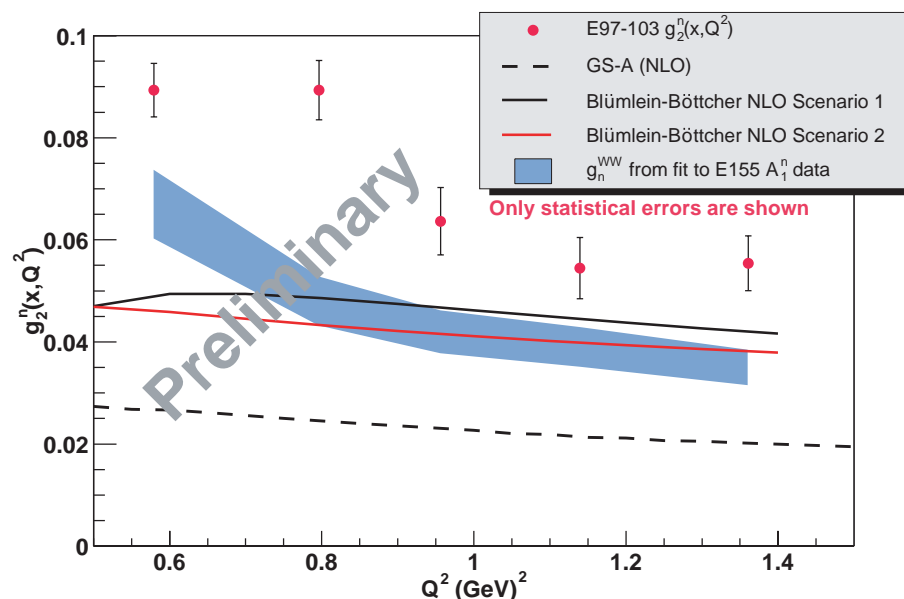
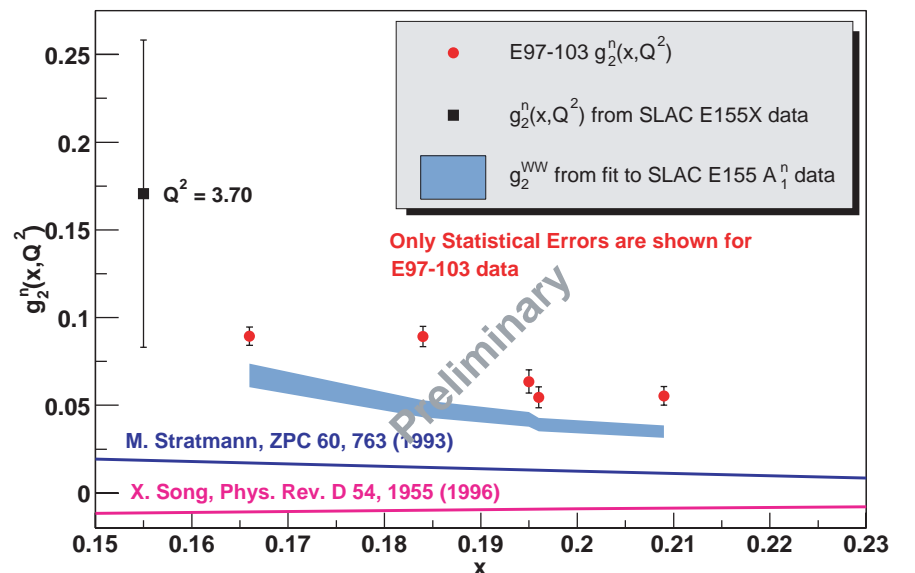
Twist-3

$$d_2^n(Q^2) = \int_0^1 x^2 [2g_1^n(x, Q^2) + 3g_2^n(x, Q^2)] dx \quad d_2 = (2\chi_B + \chi_E) / 3$$

The neutron g_2 structure function

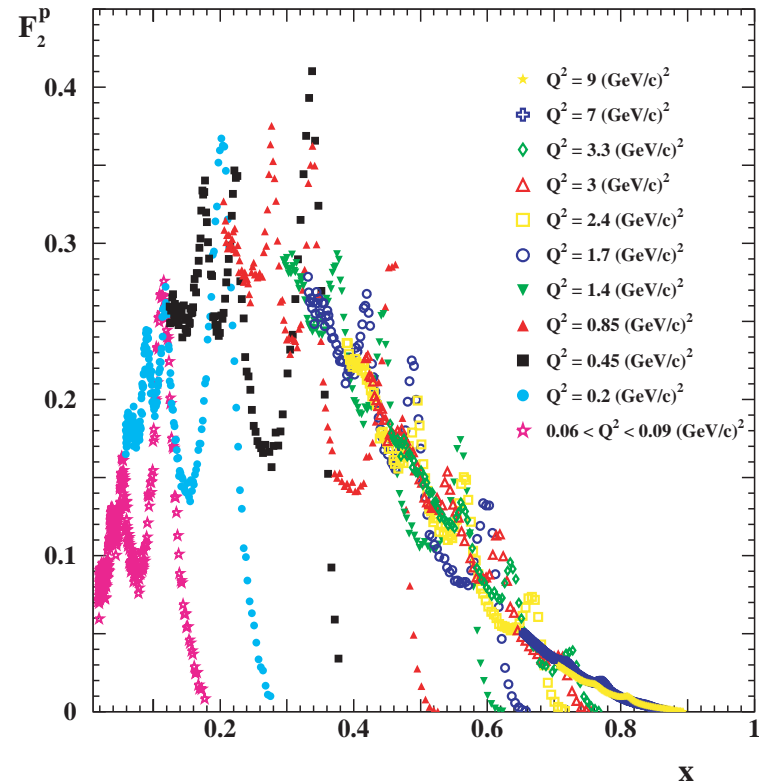
- First measurements of g_2 in Hall A order of magnitude improvement in accuracy over SLAC E155X
- Preliminary data indicate significant excess over simple prediction
- First quantitative information(?) of twist-3 effects

Hall A E97-103



Quark-Hadron Duality

- Quark-Hadron Duality implies that properly averaged hadronic observables can be described by perturbative QCD in a certain kinematic regime
- QHD must hold in the scaling region
- QHD must break down at very low Q^2
- Extensive data set from **Hall C** shows that QHD works well down to $Q^2 \sim 0.5 \text{ GeV}^2$
- Once QHD has been verified, it provides a relation between the resonance region and the DIS region

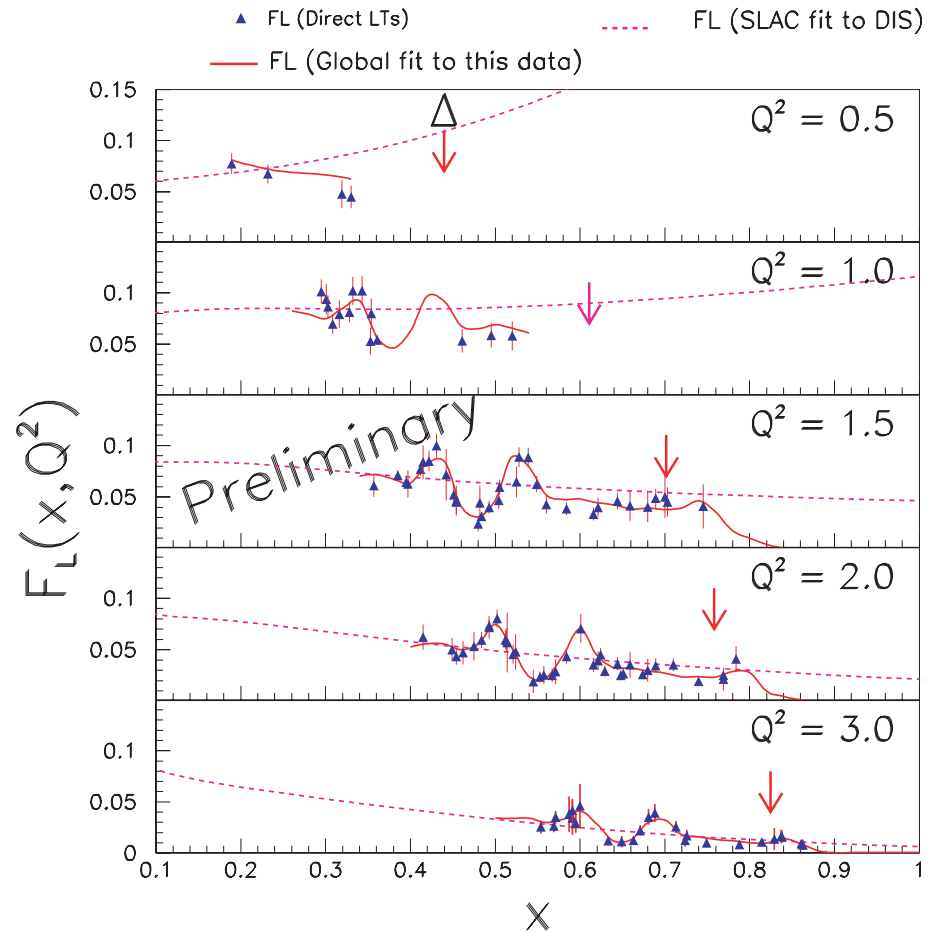


(c) Ioana Niculescu 02/02/00

$R = \sigma_L / \sigma_T$ in Resonance Region

- First measurements of R (yielding the longitudinal structure function F_L) in the resonance region
- Surprisingly strong resonance structure evident in F_L
- Allows test of QHD in F_L
- Moments of F_L can be directly compared to Lattice Gauge Theory calculations

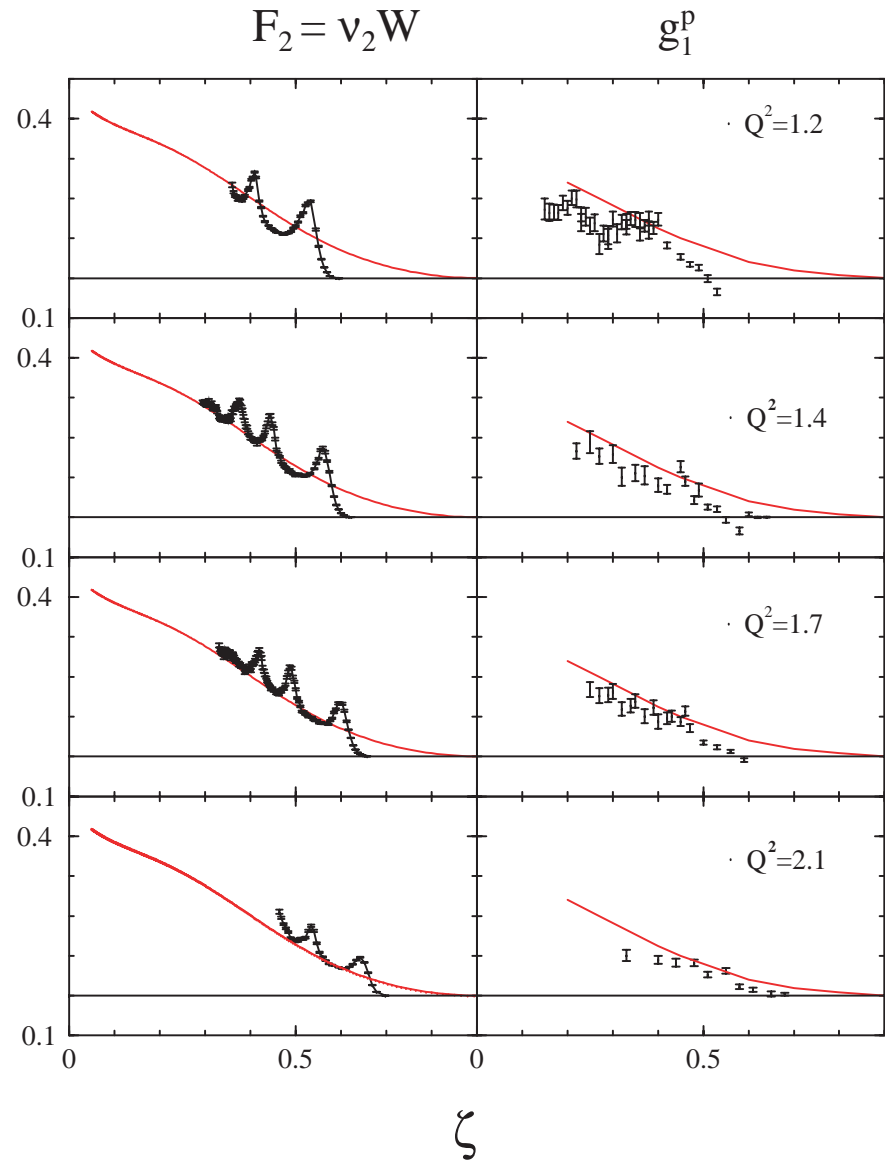
Hall C E94-110



Spin Duality

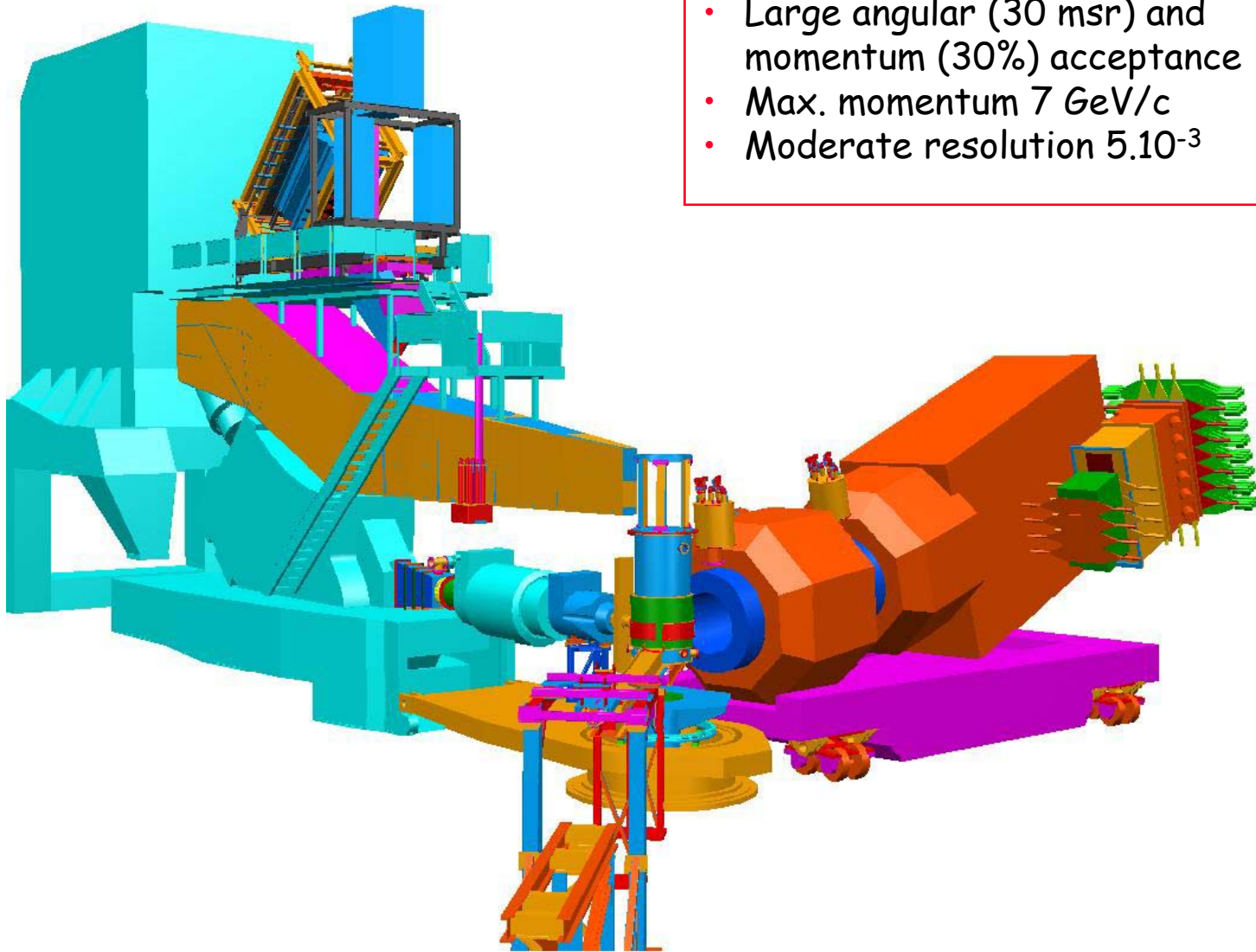
- First preliminary results of measurements of g_1^p in the resonance region
- Spin duality appears to set in at $Q^2 > 1.5 \text{ GeV}^2$
- Opens possibility to extend measurements of spin structure functions to smaller values of W (larger values of x , shown is the Nachtmann variable ξ , which is x with a target mass correction)

Hall B E91-023

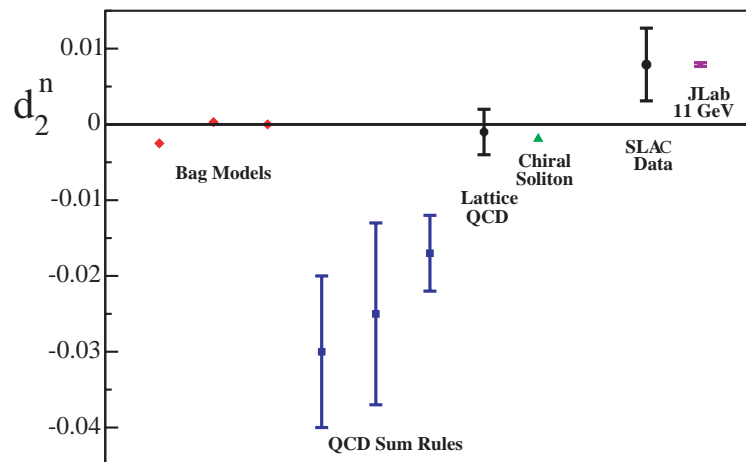
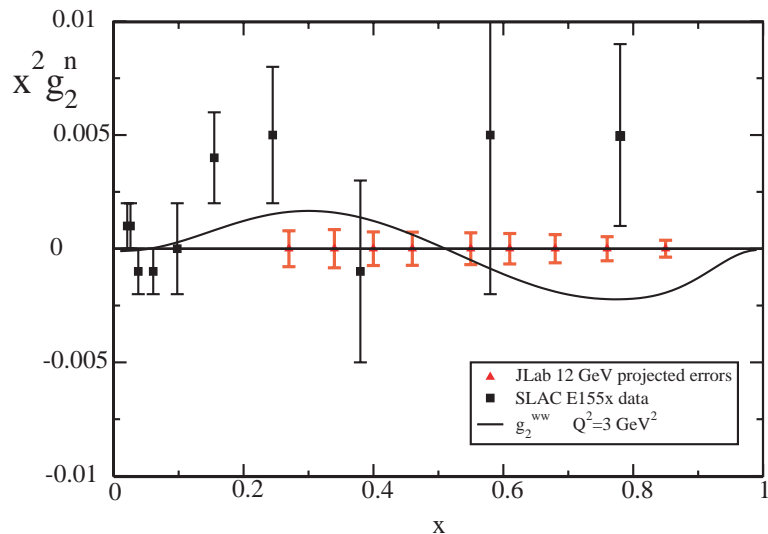
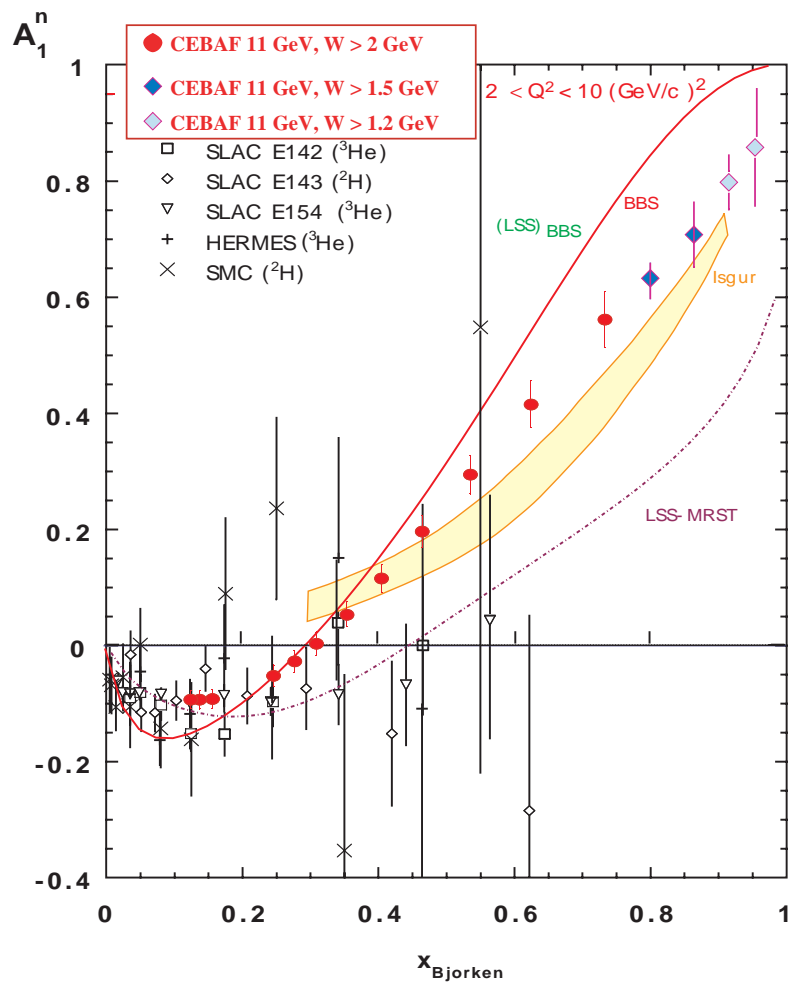


MAD Spectrometer in Hall A

- Large angular (30 msr) and momentum (30%) acceptance
- Max. momentum 7 GeV/c
- Moderate resolution $5 \cdot 10^{-3}$

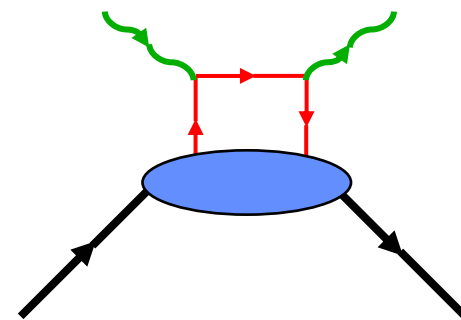
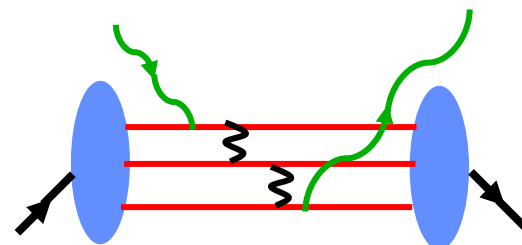


Expected results with 12 GeV upgrade

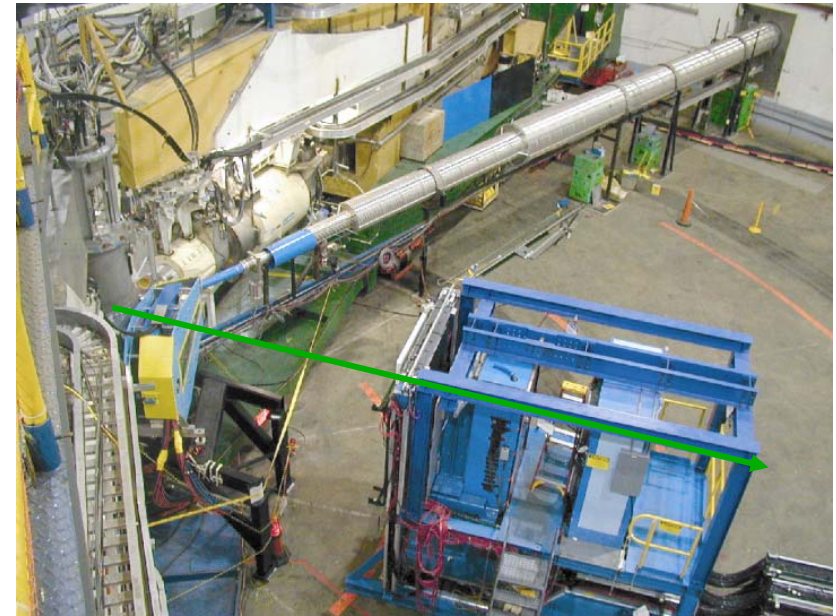
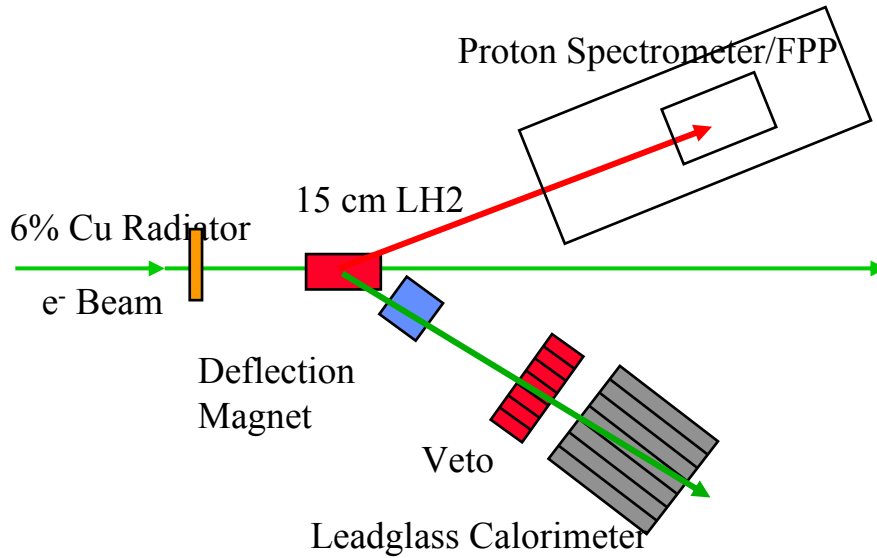


Real Compton Scattering

- Wide-Angle Compton Scattering (WACS) provides information on the partonic structure of the nucleon through the moments of the Generalized Parton Distributions
- First, the dominant mechanism of Real Compton Scattering at large values of s and t ($\sim 10 \text{ GeV}^2$) has to be established:
 - pQCD
 - momentum shared by hard gluon exchange
 - 3 active quarks
 - valence configuration dominates
 - scaling: $d\sigma/dt = f(\theta_{CM})/s^6$
 - Handbag diagram
 - hard scattering from single quark
 - momentum shared by soft overlap
 - 1-body form factor
 - soft gluon exchange neglected

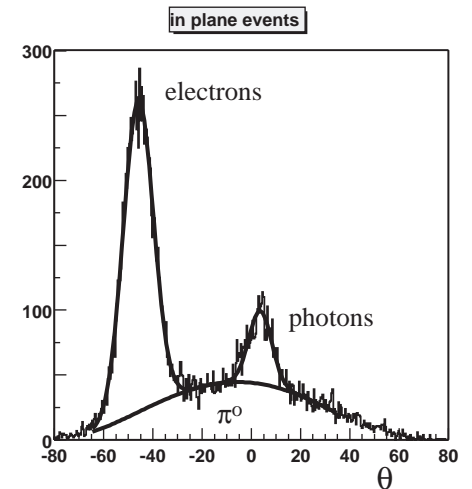
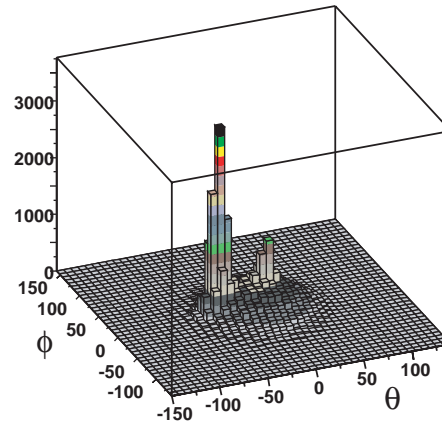


Real Compton Scattering (cont.)



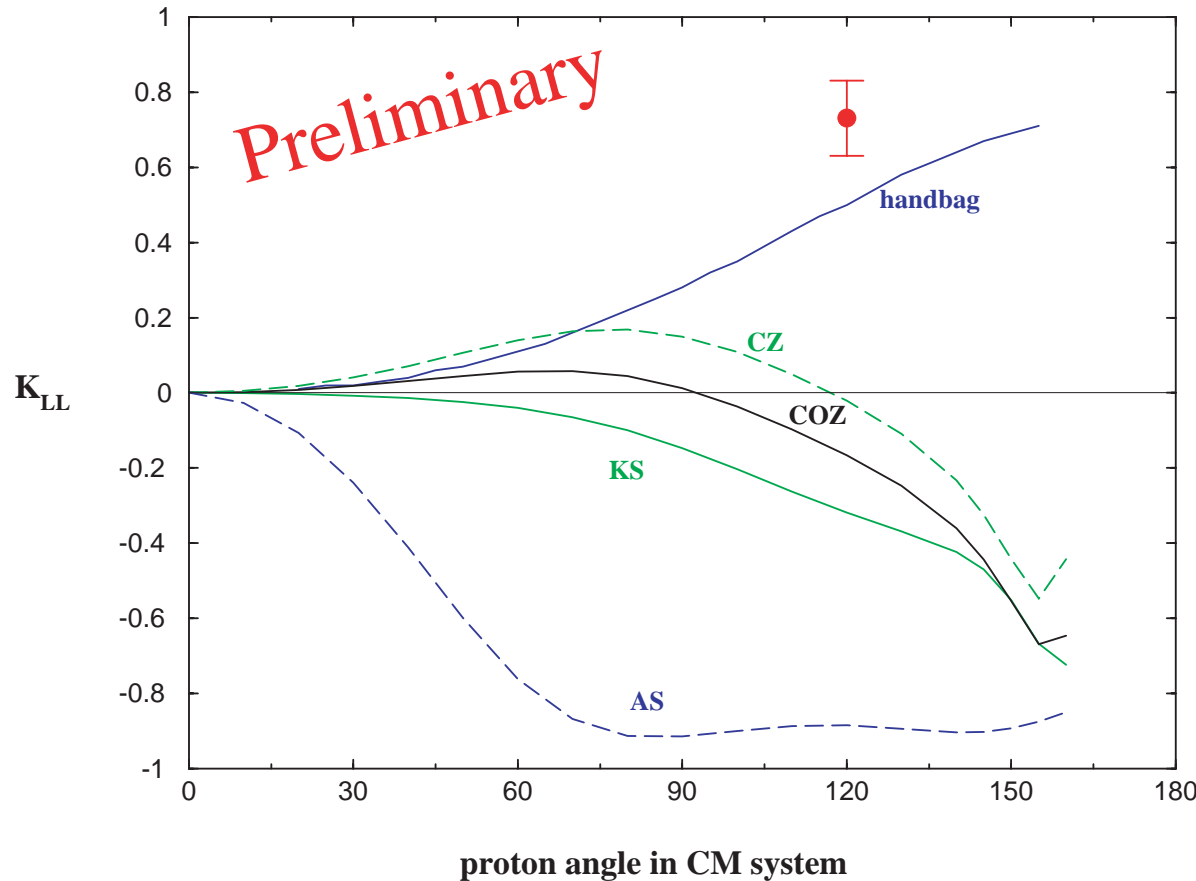
"Sweep" magnet essential for electron/photon separation

Hall A E99-114



Real Compton Scattering (cont.)

- On-line analysis of ~60 % of data
- Demonstrates feasibility of WACS at high luminosity (three orders higher than at Cornell)
- Proves dominance of soft-overlap mechanism (handbag)
- AS, KS, COZ, CZ: variety of pQCD calculations



Summary

- Study of Nucleon Spin Structure provides fascinating insight into the partonic structure of the nucleon
- Recent results from Jefferson Lab have contributed significantly to this field in a wide variety of aspects
 - Sensitive measurements of the Q^2 -evolution of the GDH integral
 - First accurate measurements of A_1^n at large x and of g_2^n
 - First L/T separation in the resonance region
 - First high-luminosity measurement of Wide Angle Compton Scattering
 -
- These studies will continue and be expanded strongly with the 12 GeV upgrade



Proton



Thomas Jefferson National Accelerator Facility