

# From quarks and gluons to the structure of hadrons

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# Outline

#### Part I:

- Introduction
- Bethe-Salpeter & Faddeev equations
- Form factors
- Tetraquarks
- Compton scattering

Part II: Ongoing and future directions

- Meson & baryon spectroscopy
- Transition form factors
- Compton scattering, PDFs, GPDs
- From quarks and gluons to nuclei

Ambitious experimental program underway with JLab @ 12 GeV:



- Search for exotic mesons
- · Nucleon form factors, flavor separation, proton radius
- · Nucleon resonances and transition form factors
- · Valence quark distributions and flavor structure
- Spatial and momentum tomography of nucleons (GPDs, TMDs)
- Short-range structure of nuclei, quarks and gluons in the nucleus, EMC effect

**QCD Lagrangian:**  $\mathcal{L} = \bar{\psi} \left( \partial + ig A + m \right) \psi + \frac{1}{4} F^{a}_{\mu\nu} F^{\mu\nu}_{a}$ 

- contains propagators and interactions:
- fully dressed n-point Green functions contain all quantum effects:

#### But!

•  $\alpha(Q^2)$  becomes large at low momenta  $\Rightarrow$ need **nonperturbative methods**!

Origin of mass generation and confinement?

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 Quarks and gluons are confined: we don't measure quarks & gluons, but hadrons



 $\Rightarrow$  need to understand spectrum and interactions!

Requires combined efforts of experiment, theory and phenomenology:

- · Amplitude analyses
- · Hadronic reaction models
- · Chiral effective field theory
- · Microscopic approaches and models
- Lattice QCD
  - $\langle 0 | T \psi(x) \overline{\psi}(y) \dots | 0 \rangle = \int \mathcal{D}A \mathcal{D}\psi \mathcal{D}\overline{\psi} e^{-S} \psi(x) \overline{\psi}(y) \dots$
- · Dyson-Schwinger, Bethe-Salpeter, Faddeev equations
- DSEs: quantum eqs. of motion
- · nonperturbative, covariant
- all momentum scales, light and heavy quarks
- · chiral symmetry
- truncations: model / neglect higher n-point functions to obtain closed system









# **QCD's Green functions**

Quark propagator



Dynamical chiral symmetry breaking generates 'constituentouark masses' Gluon propagator



Three-gluon vertex

 $\begin{array}{c} F_1 \left[ \ \delta^{\mu\nu} (p_1 - p_2)^{\rho} + \delta^{\nu\rho} (p_2 - p_3)^{\mu} \\ + \ \delta^{\rho\mu} (p_3 - p_1)^{\nu} \right] + \dots \end{array}$ 

see also Huber & von Smekal, JHEP 1304 (2013)

Agreement between lattice, DSE & FRG within reach

Quark-gluon vertex



Hadron properties are encoded in higher n-point functions.
 For example, quark four-point function contains all possible meson poles:





• Hadron properties are encoded in higher n-point functions. For example, quark four-point function contains all possible meson poles:



• Bethe-Salpeter approach: use scattering equation to obtain G in the first place:  $G = G_0 + G_0 K G$ 

Homogeneous BSE for **BS wave function:** 



 Kernel is connected to quark Dyson-Schwinger equation via chiral symmetry (can be derived from nPI effective action):



- $\rightarrow$  no constant quark mass unless NJL contact interaction
- → no crossed-ladder unless consistent quark-gluon vertex
- $\rightarrow$  cannot add **confinement** potential, drop **spin-orbit terms**, etc.
- In turn: em. gauge invariance, chiral symmetry, massless pion in chiral limit ... for free

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Homogeneous BSE for **BS wave function:** 



• Kernel is connected to **quark Dyson-Schwinger equation** via **chiral symmetry** (can be derived from nPI effective action):



Maris, Roberts, Tandy, PRC 56 (1997), PRC 60 (1999)

Rainbow-ladder:

gluon exchange with effective interaction

$$\alpha(k^2) = \alpha_{\text{IR}}\left(\frac{k^2}{\Lambda^2}, \eta\right) + \alpha_{\text{UV}}(k^2)$$

adjust scale  $\Lambda$  to observable, keep width  $\eta$  as parameter

see also: Qin, Chang et al., PRC 84 (1011), Binosi et al., PLB 742 (2015)

• In turn: em. gauge invariance, chiral symmetry, massless pion in chiral limit ... for free

# Baryons

• Covariant Faddeev equation for **baryons:** keep 2-body interactions & rainbow-ladder, but no further approximations:  $M_N = 0.94 \text{ GeV}$ GE, Alkofer, Krassnigg, Nicmorus, PRL 104 (2010), GE, PRD 84 (2011)

#### Relativistic bound states:

64 / 128 tensor structures for nucleon /  $\varDelta$ 

 Octet & decuplet baryons, pion cloud effects, beyond rainbow-ladder

Sanchis-Alepuz, Fischer, PRD 90 (2014), Sanchis-Alepuz, Fischer, Kubrak, PLB 733 (2014), Sanchis-Alepuz, Williams PLB 749 (2015)

Baryon form factors:

nucleon and  $\varDelta$  FFs,  $N \rightarrow \varDelta \gamma$  transition

GE, PRD 84 (2011), Sanchis-Alepuz, Williams, Alkofer, PRD 87 (2013), Alkofer, GE, Sanchis-Alepuz, Williams, Hyp. Int. 234 (2015)



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# **Pion form factor**



A. Krassnigg (Schladming 2010), Maris & Tandy, Nucl. Phys. Proc. Suppl. 161 (2006)

· Form factor from



• Timelike vector meson poles automatically generated by quark-photon vertex BSE!



- $\Rightarrow \Gamma^{\mu} = \begin{array}{l} {\rm Ball-Chiu} \\ ({\rm em.\ gauge\ invariance}) \end{array}$ 
  - + Transverse part (vm. poles & dominance)
- Form factor at large  $Q^2$ Chang, Cloet, Roberts, Schmidt, Tandy, PRL 111 (2013)
- Include pion cloud effects: GE, Fischer, Kubrak, Williams, in preparation

### Nucleon em. form factors



Three-body results:

all ingredients calculated, model dependence shown by bands GE, PRD 84 (2011)

- electric proton form factor: consistent with data, possible zero crossing
- magnetic form factors: missing pion effects at low Q<sup>2</sup>
- Similar for axial & ps. FFs,  $\Delta$  elastic and  $N \rightarrow \Delta \gamma$  transition GE, Fischer, EPJ A 48 (2012), Sanchis-Alepuz et al., PRD 87 (2013), Alkofer et al., Hyp. Int. 234 (2015)
  - ⇒ "quark core without pion-cloud effects"

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# Nucleon em. form factors

#### Nucleon charge radii:

isovector (p-n) Dirac (F1) radius



 Pion-cloud effects missing (⇒ divergence!), agreement with lattice at larger quark masses.



#### Nucleon magnetic moments:

isovector (p-n), isoscalar (p+n)



• But: pion-cloud cancels in  $\kappa^s \Leftrightarrow$  quark core Exp:  $\kappa^s = -0.12$ Calc:  $\kappa^s = -0.12(1)$ GE, PRD 84 (2011)

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# Nucleon- $\Delta$ - $\gamma$ transition



-10

-12

0.0

OOPS (Sparveris '05) MAMI (Stave '08)

CLAS (Aznaurvan '09

 $Q^2 [GeV^2]$ 

0.2 0.4 0.6 0.8 1.0 1.2

- small & negative, encode deformation. Reproduced without pion cloud: **OAM** from **p waves!** GE, Nicmorus, PRD 85 (2012)
- First three-body results similar Alkofer, GE, Sanchis-Alepuz, Williams, Hyp. Int. 234 (2015)

#### **Meson spectrum**



# **Tetraquarks?**

Light scalar (0<sup>++</sup>) mesons don't fit into the conventional meson spectrum:







- Why are *a*<sub>0</sub>, *f*<sub>0</sub> mass-degenerate?
- Why are their decay widths so different?

 $\Gamma(\sigma, \kappa) \approx 550 \text{ MeV}$  $\Gamma(a_0, f_0) \approx 50-100 \text{ MeV}$ 

 Why are they so light? Scalar mesons ~ p-waves, should have masses similar to axialvector & tensor mesons ~ 1.3 GeV

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# **Tetraquarks?**

What if they were tetraquarks (diquark-antidiquark)? Jaffe 1977, Close, Tornqvist 2002, Maiani, Polosa, Riquer 2004



Pelaez 2004, Weinberg 2013, Cohen, Llanes-Estrada, Pelaez, Ruiz de Elvira 2014, Londergan, Nebreda, Pelaez, Szczepaniak 2014, Parganlija, Giacosa, Rischke 2010, ...

### **Tetraquarks**

#### Four-quark bound-state equation:



Keep **two-body interactions** with **rainbow-ladder kernel:** well motivated by many other studies, tetraquark is **s-wave** 

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# Structure of the amplitude

General structure of **Bethe-Salpeter amplitude**  $\Gamma(p,q,k,P)$  complicated:



#### Arrange Lorentz invariants into multiplets of permutation group S4:

GE, Fischer, Heupel, PRD 92 (2015)

 $\Rightarrow f_i(\mathcal{S}_0, \nabla, \mathbf{O})$ 

- Singlet:  $S_0 = \frac{1}{4} (p^2 + q^2 + k^2)$
- **Doublet:**  $D_0 = \frac{1}{4S_0} \begin{bmatrix} \sqrt{3}(q^2 p^2) \\ p^2 + q^2 2k^2 \end{bmatrix}$
- 2 Triplets: ( ), ( )



Keep **s waves** only: Fierz-complete, **16** tensors:

e.g. 
$$\left\{ \begin{array}{c} C^T \gamma_5 \otimes \gamma_5 C \\ C^T \gamma^\mu \otimes \gamma^\mu C \\ \dots \end{array} \right\}$$
 in (12)(34)

automatically includes also  $\gamma_5 \otimes \gamma_5$  in (23)(14), (31)(24)

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### Four quarks $\Rightarrow$ meson molecule

 Four-quark equation dynamically generates pion poles outside the integration domain



- Dynamical formation of a 'meson molecule', diquarks almost irrelevant!
- Pion poles drive tetraquark mass from 1.4 GeV to ~350 MeV
- Poles enter integration domain above threshold  $M > 2m_{\pi}$ : the tetraquark becomes a **resonance**
- Four-body equation generates **bound state** together with its **decay channels**!



GE, Fischer, Heupel, PLB 753 (2016)



#### Four quarks $\Rightarrow$ meson molecule

Same pattern for multiplet partners: GE, Fischer, Heupel, PLB 753 (2016)

σ, κ,  $a_0 / f_0 \sim$  **350**, **750**, **1080 MeV** 

- Light scalar 'mesons' are light because they 'feel' **Goldstone nature of** *π*, *η*, *K*
- $\sigma$  is resonance close to  $\pi\pi$  threshold, becomes bound state in charm region
- Similar results from meson-meson / diquark-antidiquark two-body equations: analogous to quark-diquark model for baryons



Oettel, Hellstern, Alkofer, Reinhardt, PRC 58 (1998) GE, Cloet, Alkofer, Krassnigg, Roberts, PRC 79 (2009)



Same mechanism: **baryons** dominated by **diquarks**, **tetraquarks** by **pseudoscalar mesons**. Resolves problem with dq- $\overline{dq}$  interpretation: ' 2 x 800 MeV - binding energy ' ~ 500 MeV?!

# **Compton scattering**



Four independent variables:

$$\eta_{+} = \frac{Q^{2} + Q'^{2}}{2m^{2}}, \quad \eta_{-} = \frac{Q \cdot Q'}{m^{2}}$$
  
 $Q^{2} - Q'^{2}, \quad \eta_{-} = \frac{Q \cdot Q'}{m^{2}}$ 

$$\omega = \frac{Q^2 - Q'^2}{2m^2}, \qquad \lambda = \frac{p \cdot \Sigma}{m^2}$$





• VCS: generalized polarizabilities

- DVCS: handbag dominance, GPDs
- Forward limit: structure functions in DIS
- Timelike region: pp annhihilation at PANDA
- Spacelike region: two-photon corrections to nucleon form factors, proton radius puzzle?



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# **Compton scattering**

• Two-photon corrections to form factors: can explain difference between Rosenbluth and polarization transfer measurements Guichon, Vanderhaeghen, PRL 91 (2003)





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Arrington, Blunden, Melnitchouk Prog. Part. Nucl. Phys. 66 (2011)

#### · Proton radius puzzle:

can  $2\gamma$  corrections explain difference between electron and muon measurements?

So far: probably not, but . . .

Carlson, Vanderhaeghen, 2011 Birse, McGovern, EPJ A 48 (2012)

### Compton scattering ...

Compton amplitude = sum of Born terms + 1PI structure part:



# ... at the quark level



But only **sum** is **gauge invariant**, not individual diagrams  $\Rightarrow$  problem! Solved by projecting onto full tensor basis (transverse + gauge)

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# **Proton polarizabilities**



#### First results: GE, 1601.04154

- bands = results inside cone (70% of radius)
- compared to GPs from dispersion relation
   Pasquini et al., EPJ A11 (2001),
   Downie & Fonvieille, EPJ ST 198 (2011)



- ⇒ cf. meson electroproduction: "QCD background"!
- Large Δ contribution to β<sub>M</sub>, expect large pion effects!

Pascalutsa, Phillips, PRC 68 (2003),...

 $n_{-}$ 

↓ η<sub>+</sub>

vcs

RCS

#### In total: polarizabilities ~

Impulse app. (handbag + t-channel poles)

- + nucleon resonances (mostly Δ)
- + pion cloud (at low  $\eta_+$ )?

# **Summary Part I**

#### Nucleon and Delta masses & form factors

- · microscopic description works reasonably well, improvements underway
- · need to include meson cloud effects
- nucleon resonances?  $\rightarrow$  see Part II

#### Light scalar mesons as tetraquarks

- · transition from four quarks to "meson molecule"
- resonances!
- many future applications → see Part II

#### **Compton scattering**

- hadronic vs. quark-level decomposition (general!)
- polarizabilities
- meson electroproduction from quark level? → see Part II

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# Outline

Part I:

- Introduction
- Bethe-Salpeter & Faddeev equations
- Form factors
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#### Part II: Ongoing and future directions

- Meson & baryon spectroscopy
- Transition form factors
- Compton scattering, PDFs, GPDs
- · From quarks and gluons to nuclei

- Beyond rainbow-ladder is becoming state-of-the-art for mesons Fischer, Williams, PRL 103 (2009); Chang, Roberts, PRL 103 (2009); Williams, Fischer, Heupel, PRD 93 (2016)
- Light meson spectrum beyond rainbow-ladder: • Sanchis-Alepuz, Williams, PLB 749 (2015)



- · Gluon propagator, qg and ggg vertex solved in the process, no need for model interaction!
- Badial excitations and exotics now in right ballpark
- Becent results: scalar mesons above 1 GeV Williams, Fischer, Heupel, PRD 93 (2016)

• Tetraquarks in charmonium & bottomonium spectrum: X(3872), Y(4260), charged Z states?



 Four-body BSE dynamically determines strengths of these components: four quarks rearrange themselves into dq-dq, molecule, hadroquarkonium



- Glueball calculations underway
  Meyers, Swanson, PRD 87 (2013), Sanchis-Alepuz, Fischer, Kellermann, von Smekal, PRD 92 (2015)
- Hybrid mesons with DSEs and BSEs: Collab. with C. Fischer, J. Segovia, R. Williams



Study both within **three-body** and "**quark-diquark**" inspired approach:  $q\bar{q}$  octet repulsive, qg triplet attractive

Application to hybrid baryons?

• Exotic quantum numbers: can be obtained with BSEs, but not on lattice?



#### **Baryon spectroscopy**



• "Missing resonances": three-quark vs. quark-diquark composition?



- "Quark core" vs. meson-baryon coupled channel effects?
- Hybrid baryons?
- Connection to quark-gluon substructure in QCD?



### **Baryon spectroscopy**

Roper: 1st radial excitation of nucleon

#### Nucleon resonances:

efforts with quark-diquark and three-quark approaches underway



 Same trend for N(1535), Δ(1620), Δ(1700), Δ(1910) (dominated by negative-parity diquarks) GE, 1602.03462

### **Baryon spectroscopy**

#### Meson-baryon interactions?

- Calculate "quark core", assume that chiral interactions will provide the rest. (But what is the core? Where do we stop?)
- Implement quark-gluon topologies that produce pion cloud effects... difficult! 4PI?
- Implement effective pion cloud at quark level
   Fischer, Nickel, Wambach, PRD 76 (2007), Cloet, Bentz, Thomas, PRC 90 (2014),
   Sanchis-Alepuz, Fischer, Kubrak, PLB 733 (2014)
- New avenue? Solve five-quark equation ⇒ system will dynamically rearrange itself into Nπ,...if this is dominant







• LHCb pentaquark? Aaij et al., PRL 115 (2015)

Technical challenge: resonances! Luescher method, finite volume?

# **Transition form factors**



High-precision data on  $N \rightarrow \gamma N^*$  transition form factors available. Theory predictions?

- Roper: in ballpark of data, missing pion cloud Segovia et al., PRL 115 (2015)
- N(1535), etc.?





# **Transition form factors**



High-precision data on  $N \rightarrow \gamma N^*$  transition form factors available. Theory predictions?

- Roper: in ballpark of data, missing pion cloud Segovia et al., PRL 115 (2015)
- N(1535), etc.? Example: N(1535) elastic FF (preliminary)

 $Q^2$  range limited by quark + diquark "threshold":

 $Q^2 < 4 \left[ (m_q + m_{\rm ps})^2 - M^2 \right]$ 

Trick: calculate FFs with lower M, approach calculated mass from below. "Neutron" charge  $\rightarrow$  0: gauge invariance ok!

 Implement pion cloud effects: start with pion electromagnetic FF GE, Fischer, Kubrak, Williams, in preparation

# **Pion electroproduction**

Extraction of nucleon resonances from electroproduction amplitudes depends on knowledge of non-resonant "QCD background"

 decomposition analogous to Compton scattering (there: large contribution to α + β!)



- same phase space, but relevant kinematic region difficult to access
  - ⇒ effective scaling behavior, similar to Compton scattering?
  - ⇒ need amplitude decomposition with correct implementation of gauge invariance & analyticity

# **Compton scattering**

Formalism for Compton scattering at quark level established:

- · Amplitude decomposition with correct implementation of gauge invariance & analyticity
- t-channel meson poles reproduced GE, Fischer, PRD 87 (2013)
- general form of  $N \rightarrow \gamma N^*$  transition form factors w/o kinematic constraints and offshell extension ( $\rightarrow$  for hadronic reaction models) GE, Ramalho, in preparation
- scalar polarizabilities GE, 1601.04154

Next steps:

- · spin polarizabilities
- two-photon corrections to nucleon electromagnetic FFs
- proton radius puzzle?
- RCS/WACS: handbag vs. pQCD? Fanelli et al., PRL 115 (2015)
- DVCS, generalized parton distributions
- · Forward limit and nucleon structure functions
- Timelike Compton scattering?



### **Nucleon structure**

 Non-local four-point correlator: Wigner distribution, "mother" of TMDs, GPDs, PDFs, FFs

 $W(p,\Delta,z) = \langle N \,|\, \psi(\frac{z}{2}) \,\bar{\psi}(-\frac{z}{2}) \,|\, N' \,\rangle$ 





- With gauge link: sum of diagrams with arbitrarily many gluons coming out of G
- At same light-cone time  $z^+ = 0$ : Lorcé, Pasquini, Vanderhaeghen, JHEP 1105 (2011)



- Same kinematics as in **Compton amplitude:** 4 Lorentz invariants, 32 "form factors"
- Establish basis free of kinematic constraints
   ⇒ only singularities in "FFs" are dynamical
   ⇒ facilitates extrapolating DSE results & modeling
- DSE studies for pion GPDs underway: Mellin moments → double distributions → GPDs Mezrag, Moutarde, Rodriguez-Quintero, 1602.07722

# From quarks and gluons to nuclei

Transition from **quark-gluon to nuclear degrees of freedom**  $\Rightarrow$  generalize tetraquark studies to pentaquarks and hexaquarks



- · only input are quarks and gluons
- · dynamical generation of hadron poles, system is dominated by lowest-lying poles
- resonances!

Numerous open questions:

- Six ground states, one of them deuteron Dyson, Xuong, PRL 13 (1964)
- Dibaryon vs. hidden-color configurations? Bashkanov, Brodsky, Clement, PLB 727 (2013)
- d\*(2380), H-Dibaryon?
- Deuteron form factors from quark level?
- Microscopic origins of nuclear binding?

⇒ complementary to ongoing lattice calculations!

NPLQCD collab.: Beane et al., PRD 87 (2013), ...

# From quarks and gluons to nuclei

#### NN interaction?

- NN potential: long-range pion exchange vs. short-range repulsive core
- Nuclear ab-initio calculations
  Bedaque, van Kolck 2002, Bogner, Furnstahl, Schwenk 2010,
  Machleidt & Entem 2011, Epelbaum, Meissner 2012,
  Carlson et al. 2015, Vary et al. 2015
- Lattice (HAL-QCD): calculate scattering matrix, retroactively extract NN potential Aoki et al. 2012
- **Microscopic decomposition** analogous to FFs and other scattering amplitudes:





- only input are quarks and gluons
- quark interchange and pion exchange automatically included
- dibaryon exchanges

# From quarks and gluons to nuclei

#### Nucleons in a nuclear environment:

- **Baryon effects** important at high nuclear densities: phase structure of QCD, critical endpoint?
- Nucleon form factors and structure functions: medium effects vs. short-range structure? Underlying QCD mechanism?
- Finite T and  $\mu$ : DSE, FRG, model studies Qin et al., PRL 106 (2011), Fischer, Luecker, PLB 718 (2013), Wang et al., PRD 87 (2013), Cloet, Bentz, Thomas, PLB 642 (2006), ...

Baryon back-reaction on phase diagram GE, Fischer, Welzbacher, PRD 93 (2016)

Need to solve **meson BSE** and **baryon Faddeev equation** at T and  $\mu$ , analyze phase structure, calculate form factors and structure functions





# **Summary Part II: near future**

#### Meson spectroscopy:

- · beyond rainbow-ladder
- · tetraquarks, glueballs, hybrids

#### Baryon spectroscopy:

- · nucleon resonances & transition form factors
- · microscopic background in electroproduction

#### Hadron structure:

- · Compton scattering and its (many) applications
- Longitudinal, transverse and spin structure: GPDs, TMDs, PDFs

#### QCD and nuclei:

- · Dibaryons and NN interaction from quark level
- · Medium modifications of form factors and structure functions

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# **Backup slides**

### **Bethe-Salpeter equations**

• Extract hadron properties from **poles** in  $q\bar{q}$ , qqq,  $qq\bar{q}\bar{q}$  scattering matrices:



· defines onshell Bethe-Salpeter amplitude. Simplest example: pion

 $\psi(q, P) = \gamma_5 \left( f_1 + f_2 \not P + f_3 \not q + f_4 \left[ \not q, \not P \right] \right) \otimes \text{Color} \otimes \text{Flavor}$ 

most general Dirac-Lorentz structure, Lorentz-invariant dressing functions:

$$f_i = f_i(q^2, q \cdot P, P^2 = -m^2)$$
  $\Rightarrow$  (relative momentum ~ orbital angular momentum)

nion is made of a wayoe and n wayoel

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Same in lattice QCD: construct gauge-invariant current correlators

$$\langle (\bar{\psi} \, \mathcal{O} \, \psi)_x \, (\bar{\psi} \, \mathcal{O} \, \psi)_y \rangle = \int \mathcal{D}[\psi, \bar{\psi}, A] \, e^{-iS} \, (\bar{\psi} \, \mathcal{O} \, \psi)_x \, (\bar{\psi} \, \mathcal{O} \, \psi)_y \quad \longrightarrow \quad e^{-mt}$$

• Bethe-Salpeter approach: use scattering equation to obtain G in the first place:  $G = G_0 + G_0 K G$ 

Homogeneous BSE for **BS wave function:** 



• BS wave function only makes sense **onshell**, but homogeneous BSE = **eigenvalue equation**, can be solved for offshell momenta:

$$K \psi_i = \lambda_i(P^2) \psi_i$$
,  $\lambda_i \xrightarrow{P^2 \longrightarrow -m_i^2} 1$ 

Largest eigenvalue  $\Leftrightarrow$  ground state, smaller ones  $\Leftrightarrow$  excitations



#### Mesons

 Dynamical chiral symmetry breaking generates "constituent- quark masses"

$$S_0(p) = \frac{-i\not\!\!p + m}{p^2 + m^2} \ \longrightarrow \ S(p) = \frac{1}{A(p^2)} \frac{-i\not\!\!p + M(p^2)}{p^2 + M^2(p^2)}$$

• Pion is Goldstone boson:  $m_{\pi}^{2} \sim m_{q}$ 



Also heavy mesons

Fischer, Kubrak, Williams, EPJ A 51 (2015), Hilger et al., PRD 91 (2015)



# **Axial form factors**



- looks like magnetic form factors: missing structure at low  $Q^2 \Rightarrow g_A$  too small
- Timelike meson poles:  $a_1$  in  $G_A$ ,  $\pi \& \pi(1300)$  in  $G_P$ ,  $G_{\pi NN}$
- Goldberger-Treiman relation reproduced for all quark masses:

$$G_A(0) = \frac{f_{\pi}}{M_N} G_{\pi NN}(0)$$
 GE & Fischer, EPJ A 48 (2012)



## Quark-diquark model

 Assumption: separable qq scattering matrix ⇒ Faddeev equation simplifies to quark-diquark BSE



Oettel, Hellstern, Alkofer, Reinhardt, PRC 58 (1998), Cloet, GE, El-Bennich, Klahn, Roberts, FBS 46 (2009) Segovia, Cloet, Roberts, Schmidt, FBS 55 (2014)

- Quark exchange between quark & diquark binds nucleon
- Calculate all quark and diquark ingredients from quark level ⇒ direct link to quark-gluon interaction! Rainbow-ladder: scalar diquark ~ 800 MeV, axialvector diquark ~ 1 GeV
- N and ∆ masses & form factors very similar: quark-diquark model is good approximation for three-body equation
- Nucleon and Δ electromagnetic FFs, N→Δγ and N→Δπ transition GE, Cloet, Alkofer, Krassnigg, Roberts, PRC 79 (2009), Nicmorus, GE, Alkofer, PRD 82 (2010), Mader, GE, Blank, Krassnigg, PRD 84 (2011), GE, Nicmorus, PRD 85 (2012)



Oettel, Pichowsky, von Smekal, EPJ A 8 (2000)

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#### GE, Nicmorus, PRD 85 (2012)



# $\varDelta$ electromagnetic FFs

Almost no experimental information since  $\Delta$  unstable:  $\Delta \rightarrow N\pi$ 

**Magnetic moment**  $\mu_{\Delta} \sim 3.5$  with large errors ( $\Delta^+$ ). But  $\Omega^-$  (spin 3/2, sss) is stable w.r.t strong interaction, magnetic moment  $|\mu_{\Omega}| = 3.6(1)$ . Accidental?



$$J^{\mu,\rho\sigma}(P,Q) = i \mathbb{P}^{\rho\alpha}(P_f) \left[ \left( F_1^{\star} \gamma^{\mu} - F_2^{\star} \frac{\sigma^{\mu\nu}Q^{\nu}}{2M_{\Delta}} \right) \delta^{\alpha\beta} - \left( F_3^{\star} \gamma^{\mu} - F_4^{\star} \frac{\sigma^{\mu\nu}Q^{\nu}}{2M_{\Delta}} \right) \frac{Q^{\alpha}Q^{\beta}}{4M_{\Delta}^2} \right] \mathbb{P}^{\beta\sigma}(P_i)$$

Form factors at  $Q^2=0$ :

$G_{E_0}(0) = e_\Delta$ $G_{E_2}(0) = \mathcal{Q}$	charge electric quadrupole moment	$G_{M1}(0)$	4 -	
$G_{M_1}(0) = \mu_\Delta$ $G_{M_3}(0) = \mathcal{O}$	magnetic dipole moment magnetic octupole moment		2 -	+

almost quark-mass independent, match  $\Omega^-$  magnetic moment Nicmorus, GE, Alkofer, PRD 82 (2010)

Three-body results similar (except  $G_{M_3}$ ) Sanchis-Alepuz, Alkofer, Williams, PRD 87 (2013)



# **Pion cloud effects**

#### • Hadron level:

 $N\pi$  contributions to nucleon self-energy; charge radii diverge in chiral limit,  $\Delta \rightarrow N\pi$  decay cusps, etc.



- Baryons: pion effects reduce N, Δ masses but also f<sub>π</sub> (sets the scale) by similar amount: net effect small Sanchis-Alepuz, Fischer, Kubrak, PLB 733 (2014)
- Pion form factor: photon also couples to pion (necessary for gauge invariance), π exchange in guark-photon vertex



#### • Quark level:

 $\pi$  contributions to quark self-energy, effective  $\pi$  exchange between quarks; pion not elementary field!





GE, Fischer, Kubrak, Williams, in preparation

# DSE / Faddeev landscape $N ightarrow N^* \gamma$

	Quark-diquark			Three-quark		
	Contact interaction	QCD-based model	DSE (RL)	RL	bRL	bRL + 3q
$N, \Delta$ masses $N, \Delta$ em. FFs $N \rightarrow \Delta \gamma$	$\sqrt[]{}$	$\checkmark$ $\checkmark$	$\sqrt[]{}$	$\sqrt[n]{}$	$\checkmark$	
Roper $N \to N^* \gamma$						
$N^*(1535), \ldots$ $N \to N^*\gamma$						

# N\*(1535)?



Form factors: no kinematic constraints CLAS data & toy parametrization with "ho bump"

...vs. helicity amplitudes

in  $[10^{-3} GeV^{-1/2}]$ 

kinematic zeros at  $Q^2 = -(m_R \pm m)^2$ 

see also Ramalho & Tsushima, PRD 84 (2011)

# Negative parity resonances?



# N\*(1535): the recipe

• Calculate quark DSE and (pseudoscalar, vector) diquark BSEs & propagators in complex plane





pseudoscalar diquark  $\sim$  1 GeV vector diquark  $\sim$  1.1 GeV

 Solve Faddeev equation, obtain N\*(1535) mass and wave function



Calculate quark-photon and (pseudoscalar, vector scalar, axialvector) diquark-photon vertices



• Insert everything here and calculate transition form factor:



# N\*(1535): expectations vs. results

- Mesons and (opposite-parity) diquarks intrinsically linked in rainbow-ladder
- Ps & v mesons ok, but sc & av mesons too light, repulsive shifts beyond RL Fischer, Williams, PRL 103 (2009). Chang, Roberts, PRL 103 (2009)  $\Rightarrow$  also **ps & v diquarks** should be "too light"  $\Rightarrow$  same for  $N^*(1535)$ , etc. Chen et al, FBS 53 (2012)
- Three-body Faddeev equation:  $M(1535) \sim 1.1 \, {\rm GeV}$  Sanchis-Alepuz, Williams, PLB 749 (2015)
- But quark-diquark BSE with ps & v diquarks only:  $M(1535) \sim 1.65 \dots 1.70 \text{ GeV}$  !
- Including sc & av diquarks for  $N^*(1535)$ ,  $N^*(1520)$ : not yet fully consistent (nor stable), but important for  $N^*(1520)$ . Near mass degeneracy?
- Transition form factor (with ps & v diquarks only) extremely small ⇒ ps & v diquarks not enough?



# Nucleon resonances: $J^P = \frac{1}{2}^{\pm}$

 Offshell nucleon-photon vertex depends on 12 tensor structures (4 gauge + 8 transverse)

$\gamma^{\mu}$	$t^{\mu u}_{QQ}\gamma^{ u}$	$t^{\mu u}_{QQ}ik^{ u}$	
$ik^{\mu}$	$t^{\mu\nu}_{QQ}k\!\cdot\!Q\tfrac{i}{2}[\gamma^\nu,k\!\!\!/]$	$t^{\mu u}_{QQ}k^{ u}k$	
$k^{\mu}k$	$\frac{i}{2} \left[ \gamma^{\mu}, Q \right]$	$t^{\mu\nu}_{Qk}k\!\cdot\!Q\gamma^\nu$	
$k\cdot Q \stackrel{i}{\underline{\imath}} [\gamma^\mu, k]$	$\frac{1}{6} \left[ \gamma^{\mu}, k, Q \right]$	$t_{Qk}^{\mu\nu}\tfrac{i}{2}[\gamma^\nu,k]$	
gauge	transverse		
		/	



 $J^{\mu} = \bar{u}_f \left[ F_1 \gamma^{\mu} + \frac{F_2}{2m} \frac{i}{2} [\gamma^{\mu}, \mathcal{Q}] \right] u_i$ 

• No gauge part for  $N \to N^*(\frac{1}{2}^{\pm})$  transition form factors (no conserved charge), hence

$$J^{\mu} = \bar{u}_{f} \begin{bmatrix} 1\\ i\gamma_{5} \end{bmatrix} \begin{bmatrix} \frac{F_{1}}{m^{2}} \underbrace{t^{\mu\nu}_{QQ} \gamma^{\nu}}_{Q^{2}} + \frac{F_{2}}{2m} \frac{i}{2} [\gamma^{\mu}, Q] \end{bmatrix} u_{i}$$

 $\Rightarrow$   $F_1$  and  $F_2$  free of kinematic constraints



+ crossed term

$$t^{\mu\nu}_{AB} = A \cdot B \,\delta^{\mu\nu} - B^{\mu}A^{\nu}$$

• What about offshell form factors in nucleon Born term?

 $\gamma^{\mu}$  clashes with gauge invariance, only  $t_{QQ}^{\mu\nu} \gamma^{\nu}$ ,  $\frac{i}{2} [\gamma^{\mu}, Q]$  allowed: must use **Dirac current** 

Gauge invariance restored by 1PI part GE, Fischer, PRD 87 (2013)

⇒ careful with offshell form factors!

# Nucleon resonances: $J^P = \frac{3}{2}^{\pm}$

 Offshell N → N\*(<sup>3±</sup>) transition vertices: Rarita-Schwinger rep. contains spin-3/2 & 1/2; must get rid of offshell spin-1/2 background



Shklyar, Lenske, PRC 80 (2009) GE, Ramalho, in preparation + crossed term

Vertices must satisfy electromagnetic & spin-3/2 gauge invariance and invariance under point transformations: Pascalutsa. Timmermanns. PRC 60 (1999)

$$Q^{\mu}\,\Gamma^{\alpha\mu}_{\rm R}=0,\qquad k^{\alpha}\,\Gamma^{\alpha\mu}_{\rm R}=0,\qquad \gamma^{\alpha}\,\Gamma^{\alpha\mu}_{\rm R}=0$$

• General form of offshell  $N \rightarrow \frac{3}{2}^{\pm}$  transition currents: 12 form factors



· onshell everything collapses into 3 form factors:

$$J^{\mu} = \frac{1}{m^2} \, \bar{u}_f^{\alpha} \begin{bmatrix} i \gamma_5 \\ 1 \end{bmatrix} \begin{bmatrix} F_1 \, \varepsilon_{kQ}^{\alpha \mu} - F_2 \, t_{kQ}^{\alpha \mu} - \frac{F_3}{m} \, i t_{k\gamma}^{\alpha \beta} \, t_{QQ}^{\beta \mu} \end{bmatrix} u_i$$

 $\Rightarrow$  no kinematic constraints.

Jones-Scadron decomposition not good offshell!

### **Tetraquarks**



$$\begin{array}{c} -p_{1} \\ -p_{2} \\ p_{2} \\ p_{1} \\ p_{1} \\ + \end{array} = \begin{array}{c} 0 \\ p_{2} \\ p_{1} \\ p_{1} \\ + \end{array} = \begin{array}{c} 0 \\ p_{2} \\ p_{1} \\ p_{1} \\ + \end{array} = \begin{array}{c} 0 \\ p_{2} \\ p_{1} \\ p_{2} \\ p_{1} \\ p_{1} \\ + \end{array} = \begin{array}{c} 0 \\ p_{1} \\ p_{2} \\ p_{1} \\ p_{2} \\ p_{1} \\ p_{1} \\ p_{2} \\ p_{2} \\ p_{1} \\ p_{2} \\$$

 BSE dynamically generates pion poles in wave function, drive σ mass from 1.5 GeV to ~350 MeV

diquark



Four quarks rearrange to "meson molecule", diquarks irrelevant

Tetraquark is at the same time dynamically generated resonance!



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# Muon g-2

• Muon anomalous magnetic moment: total SM prediction deviates from exp. by ~3 $\sigma$ 

$$\int_{p}^{q} = ie \, \bar{u}(p') \left[ F_1(q^2) \, \gamma^{\mu} - F_2(q^2) \, \frac{\sigma^{\mu\nu}q_{\nu}}{2m} \right] u(p)$$

• Theory uncertainty dominated by **QCD:** Is QCD contribution under control?



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Hadronic light-by-light scattering

$a_{\mu} [10^{-10}]$	Jegerlehner, Nyffeler, Phys. Rept. 477 (2009)			
Exp:	11	659 208.9	(6.3)	_
QED:	11	658 471.9	(0.0)	
EW:		15.3	(0.2)	
Hadronic:				
• VP (LO+H	O)	685.1	(4.3)	
• LBL		10.5	(2.6)	?
SM:	11	659 182.8	(4.9)	
Diff:		26.1	(8.0)	

LbL amplitude: ENJL & MD model results

Bijnens 1995, Hakayawa 1995, Knecht 2002, Melnikov 2004, Prades 2009, Jegerlehner 2009, Pauk 2014



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SM: 11	659 182.8	(4.9)	-
Diff:	26.1	(8.0)	

• LbL amplitude at quark level, derived from gauge invariance: GE, Fischer, PRD 85 (2012), Goecke, Fischer, Williams, PRD 87 (2013)



- no double-counting, gauge invariant!
- need to understand structure of amplitude GE, Fischer, Heupel, PRD 92 (2015)