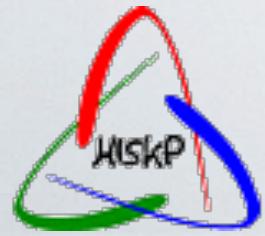


# Chiral effective field theory for hadron and nuclear physics

Jose Manuel Alarcón

Helmholtz-Institut für Strahlen- und Kernphysik  
University of Bonn



*Biographical presentation*

# *Biographical presentation*

- Born in 1983 in Cartagena in the Region of Murcia (Spain).
- 2001 - I started my studies in Physics at the University of Murcia.
- 2005 - I received the 1st prize in the physics contest “Celebrando la Física” organized by the University of Murcia.
- 2006 - I completed my studies with the best grades, receiving a special prize for it (Premio extraordinario fin de carrera).
- 2007 - I finished the “Master of Advanced Physics”, with specialization in theoretical physics, at the University of Valencia (maximum grade in Master Thesis).
- 2007 - I started to work on my thesis under the supervision of Prof. Jose Antonio Oller, at the University of Murcia.
- Topic: Relativistic formulations of chiral EFT with baryons and application to  $\pi N$  scattering.

# Biographical presentation

- 2012 - I defended my thesis, entitled “Baryon Chiral Perturbation Theory in its manifestly covariant forms and the study of the  $\pi N$  dynamics & On the  $Y(2175)$  resonance”. (Sobresaliente Cum Laude).
- 2014 - Thesis awarded with the “Premio extraordinario de doctorado”, given to the best thesis in physics defended in the period 2012-2013 at the University of Murcia.
- 2015 - Thesis awarded by the Nuclear Physics Division of the European Physical Society with the Dissertation Award (2012-2014)
- July 2012 - I started to work in the group of Prof. Marc Vanderhaeghen at the Johannes Gutenberg University, Mainz.
  - Working with V. Pascalutsa: Nucleon Polarizabilities and  $\mu H$  Lamb shift.
- September 2014 - I started to work in Bonn.
  - Application of EFT to *ab initio* many-body nuclear calculations.

# *Research topics*

$\pi N$  scattering

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.
  - Fundamental hadronic reaction involving one baryon.

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.
  - Fundamental hadronic reaction involving one baryon.
  - Long range part of the NN forces.

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.
  - Fundamental hadronic reaction involving one baryon.
  - Long range part of the NN forces.
  - Not completely understood in the context of chiral dynamics

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.
  - Fundamental hadronic reaction involving one baryon.
  - Long range part of the NN forces.
  - Not completely understood in the context of chiral dynamics   
Disagreement with dispersive approaches [Becher and Leutwyler, JHEP (2001)].

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.
  - Fundamental hadronic reaction involving one baryon.
  - Long range part of the NN forces.
  - Not completely understood in the context of chiral dynamics   
Disagreement with dispersive approaches [Becher and Leutwyler, JHEP (2001)].
  - Information about the internal scalar structure of the nucleon

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.
  - Fundamental hadronic reaction involving one baryon.
  - Long range part of the NN forces.
  - Not completely understood in the context of chiral dynamics → Disagreement with dispersive approaches [Becher and Leutwyler, JHEP (2001)].
  - Information about the internal scalar structure of the nucleon →  $\sigma_{\pi N}$

# $\pi N$ scattering

- Application of chiral EFT with baryons to fundamental hadronic reactions:
  - Understanding of the chiral dynamics of the hadronic processes at low energies.
  - Insight into the internal structure of the nucleon.
- $\pi N$  scattering.
  - Fundamental hadronic reaction involving one baryon.
  - Long range part of the NN forces.
  - Not completely understood in the context of chiral dynamics → Disagreement with dispersive approaches [Becher and Leutwyler, JHEP (2001)].
  - Information about the internal scalar structure of the nucleon →  $\sigma_{\pi N}$ 
    - Important hadronic uncertainty in direct detection of DM [Bottino, Donato, Fornengo and Scopel, Astropart. Phys. 13, (2000); Astropart. Phys. 18, (2002)] [Ellis, Olive and Savage PRD 77, (2008)].
    - Formation of elements needed for life [Berengut et. al., PRD 87, (2013)].

# $\pi N$ scattering

- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).

# $\pi N$ scattering

- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [*Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)*]:

# $\pi N$ scattering

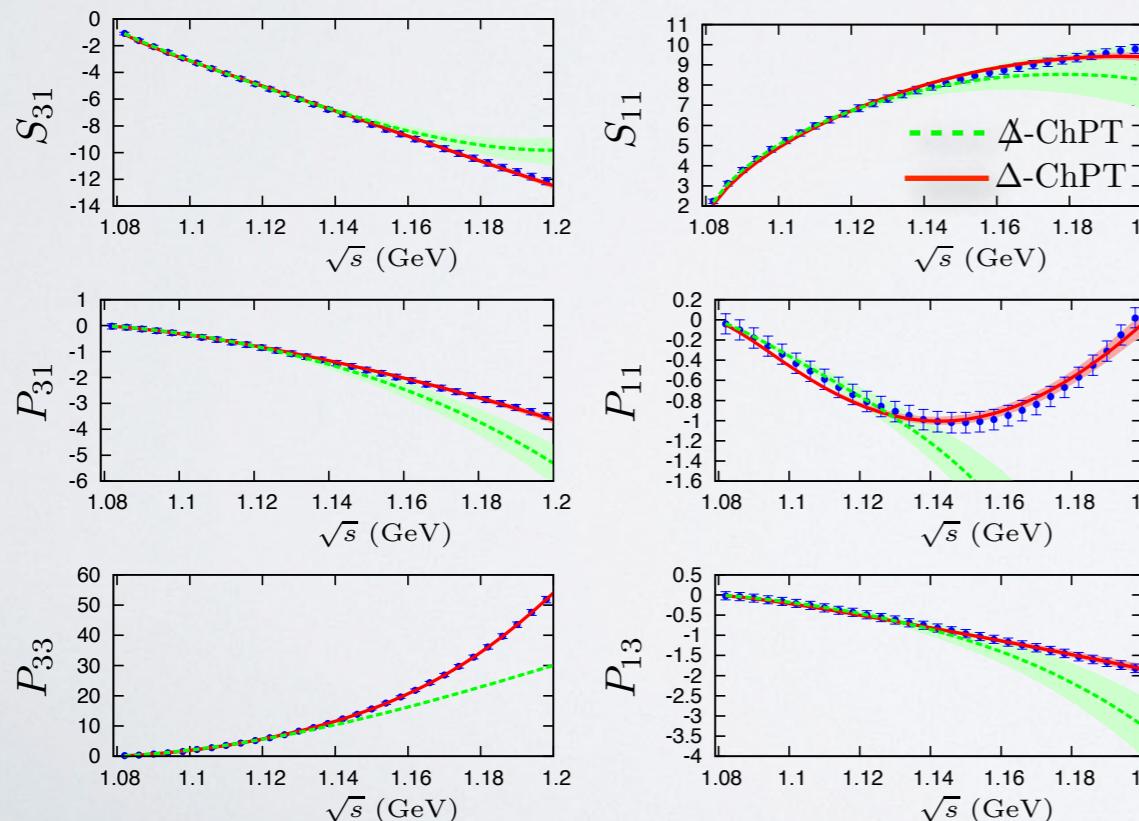
- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [*Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)*]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.

# $\pi N$ scattering

- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [*Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)*]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.
  - The inclusion of the  $\Delta(1232)$  makes a difference in the prediction of the phenomenology at low energies

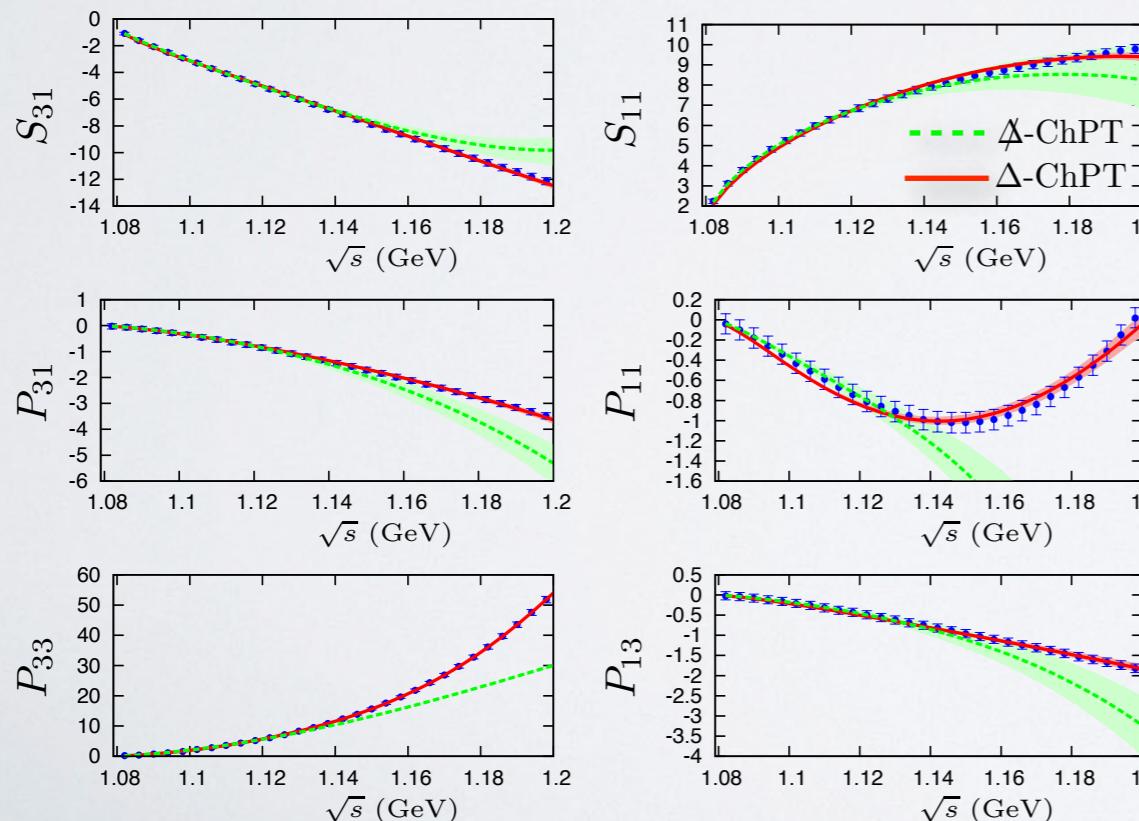
# $\pi N$ scattering

- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.
  - The inclusion of the  $\Delta(1232)$  makes a difference in the prediction of the phenomenology at low energies → EOMS is clearly the best.



# $\pi N$ scattering

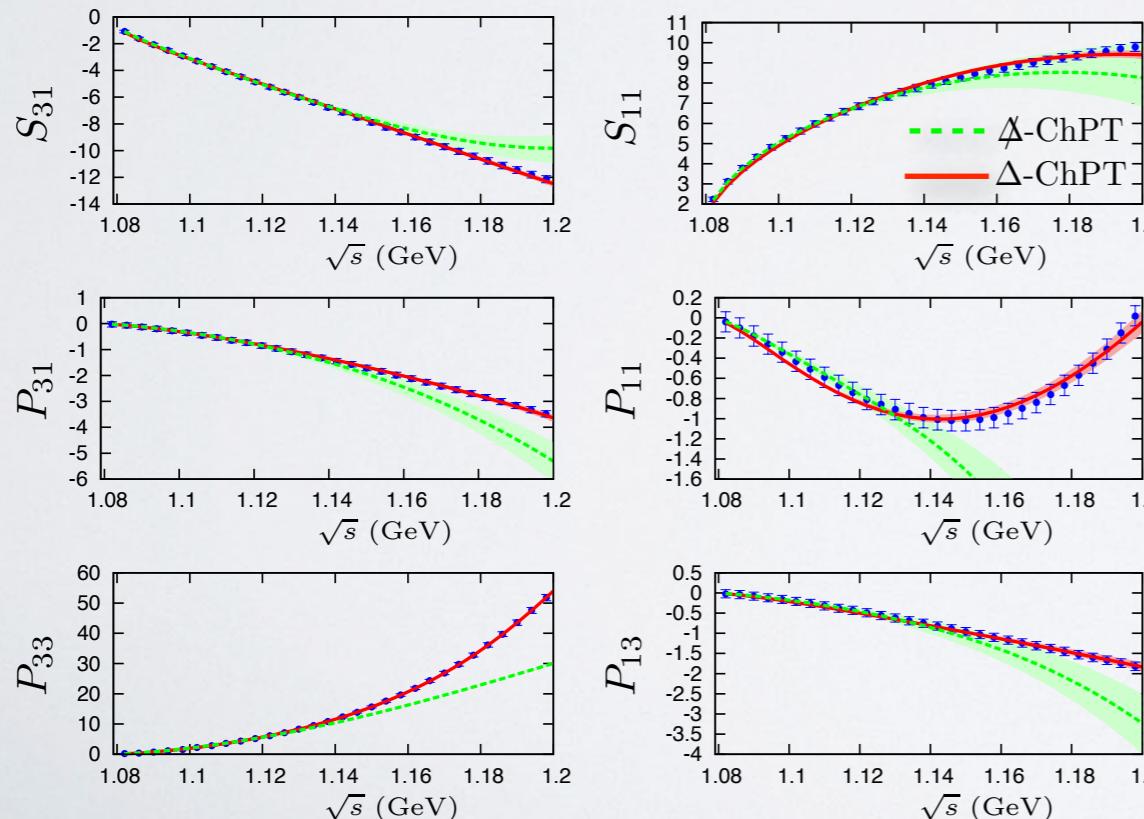
- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.
  - The inclusion of the  $\Delta(1232)$  makes a difference in the prediction of the phenomenology at low energies → EOMS is clearly the best.



- EOMS +  $\Delta(1232)$ :

# $\pi N$ scattering

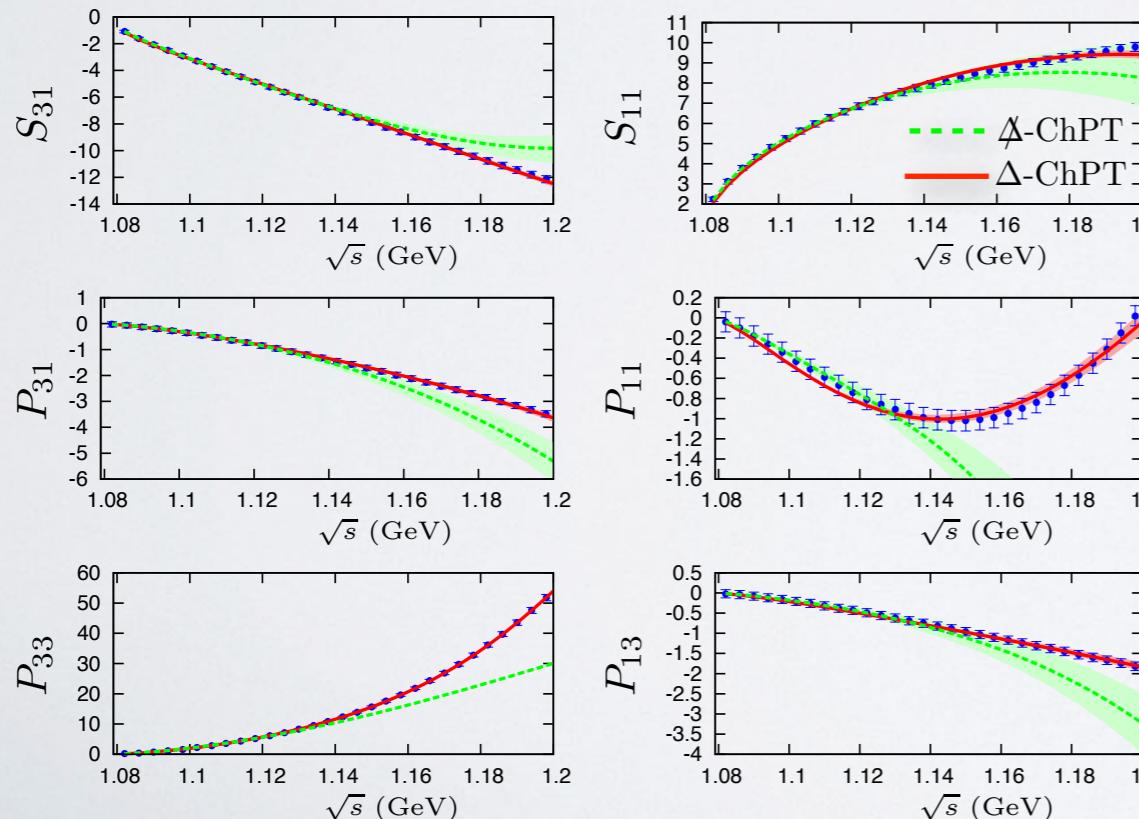
- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.
  - The inclusion of the  $\Delta(1232)$  makes a difference in the prediction of the phenomenology at low energies → EOMS is clearly the best.



- EOMS +  $\Delta(1232)$ :
- Agreement with the dispersive approaches

# $\pi N$ scattering

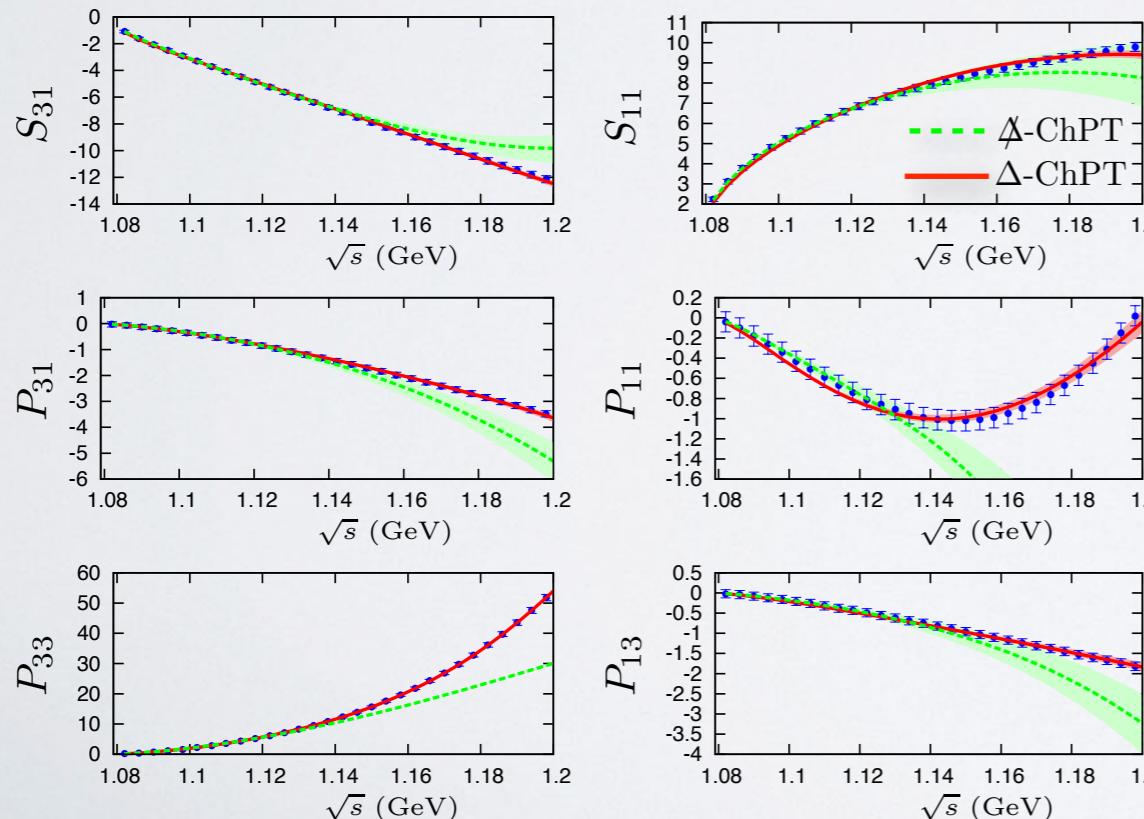
- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.
  - The inclusion of the  $\Delta(1232)$  makes a difference in the prediction of the phenomenology at low energies → EOMS is clearly the best.



- EOMS +  $\Delta(1232)$ :
- Agreement with the dispersive approaches  
→ BChPT is reliable!

# $\pi N$ scattering

- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.
  - The inclusion of the  $\Delta(1232)$  makes a difference in the prediction of the phenomenology at low energies → EOMS is clearly the best.



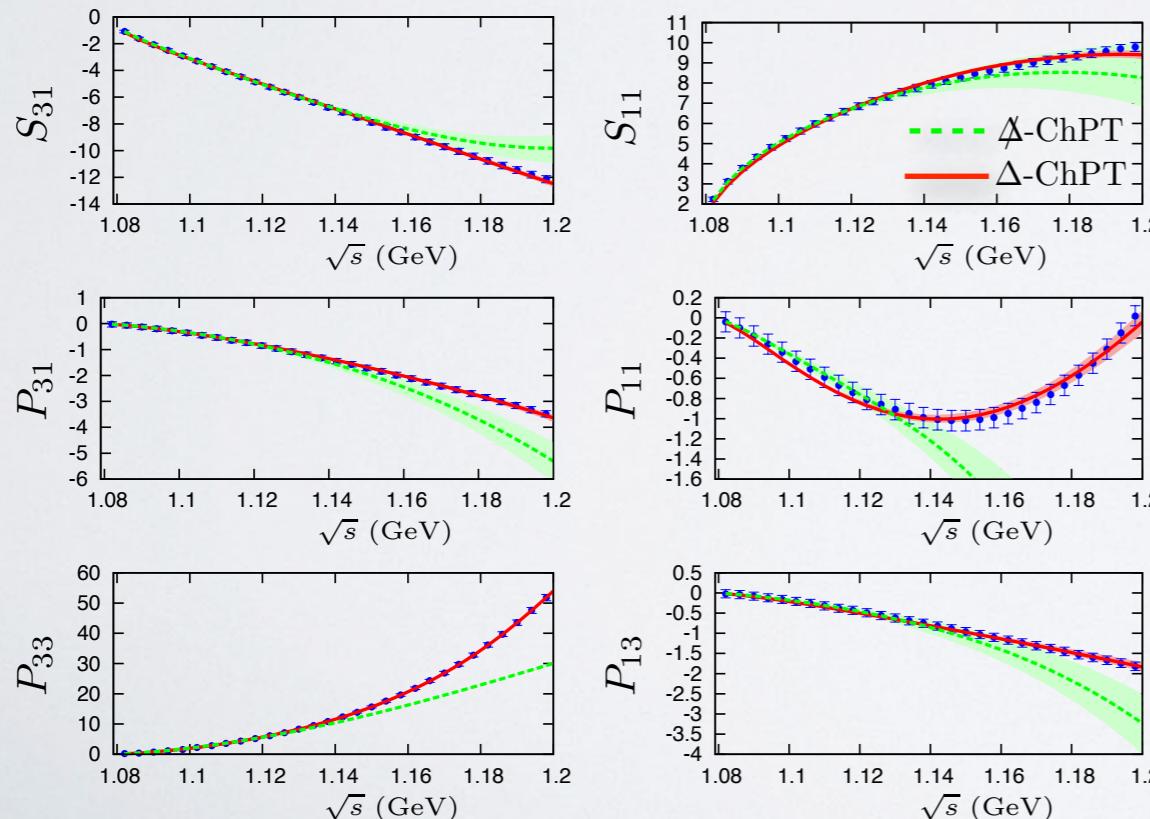
- EOMS +  $\Delta(1232)$ :
  - Agreement with the dispersive approaches  
→ BChPT is reliable!
  - Determination of  $\sigma_{\pi N}$  from modern phenomenological information.

$$\sigma_{\pi N} = 59(7) \text{ MeV}$$

[Alarcón, Martín Camalich and Oller, PRD 85 (2012)]

# $\pi N$ scattering

- The two manifestly Lorentz invariant schemes were used: Infrared Regularization (IR) and Extended-On-Mass-Shell (EOMS).
- Conclusions [Alarcón, Martín Camalich and Oller, Ann. of Phys. 336 (2013)]:
  - At low energies above threshold IR, EOMS as well as Heavy Baryon ChPT give very similar results.
  - The inclusion of the  $\Delta(1232)$  makes a difference in the prediction of the phenomenology at low energies → EOMS is clearly the best.



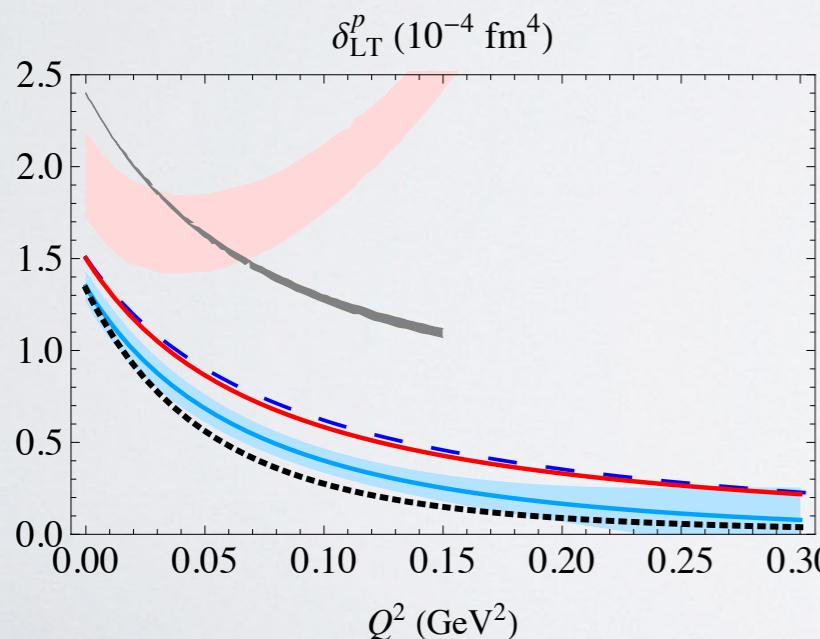
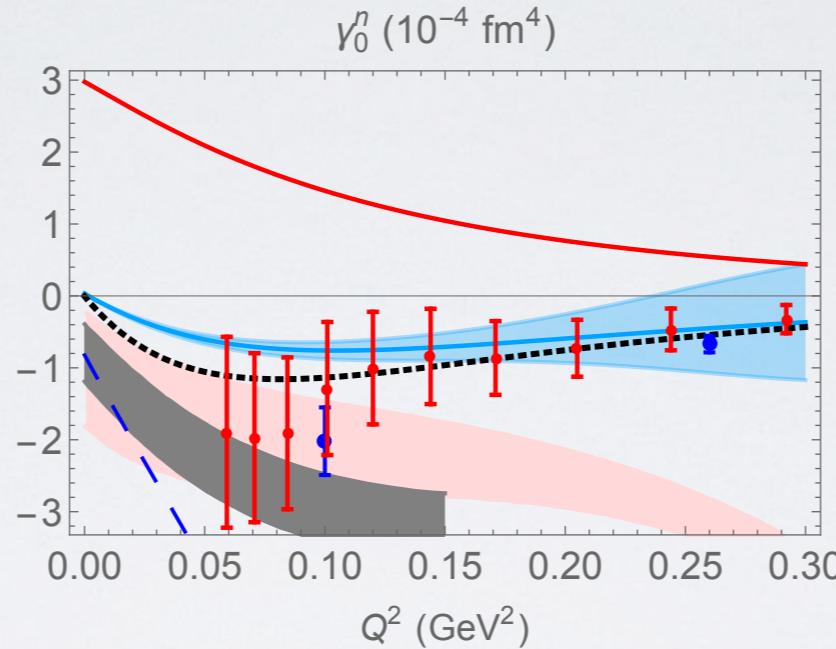
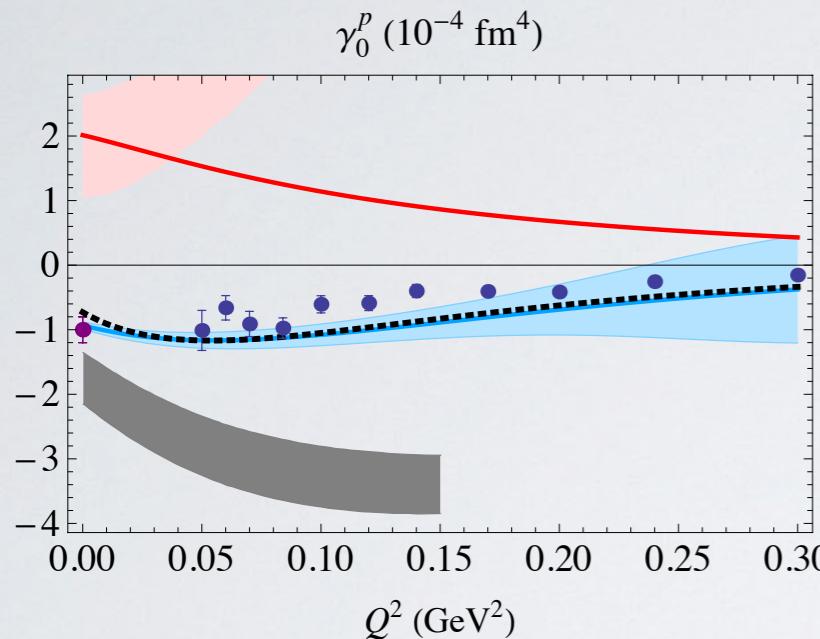
- EOMS +  $\Delta(1232)$ :
  - Agreement with the dispersive approaches  
→ BCChPT is reliable!
  - Determination of  $\sigma_{\pi N}$  from modern phenomenological information.
- $\sigma_{\pi N} = 59(7)$  MeV [Alarcón, Martín Camalich and Oller, PRD 85 (2012)]
- Analysis confirmed point by point by the Roy-Steiner analysis of [Hoferichter et al., PRL 115 (2015)]

# *Nucleon Polarizabilities and Lamb shift*

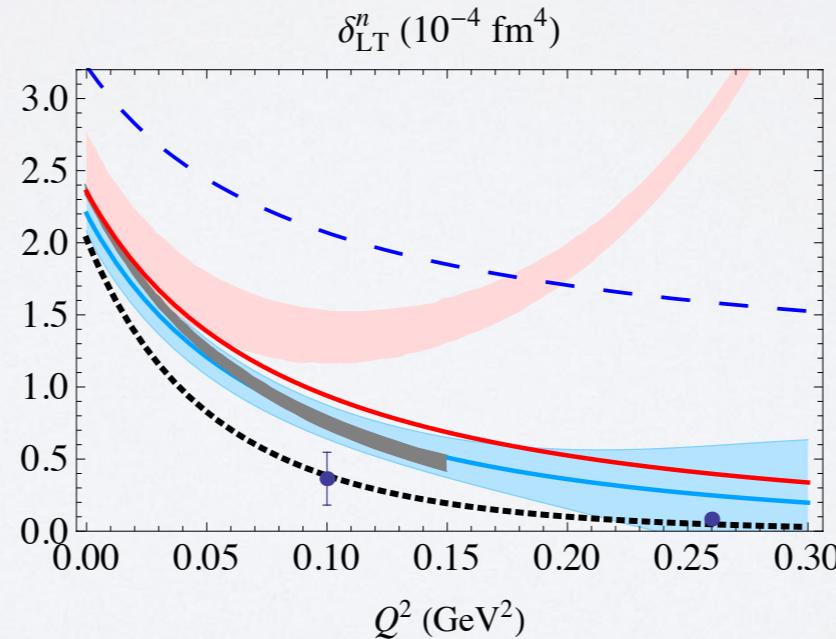
# Nucleon Polarizabilities

- Relativistic Chiral EFT+ $\Delta$  calculation of forward VVCS up to  $\mathcal{O}\left(\frac{p^4}{\Delta}\right)$

- Spin Polarizabilities:



[Lensky, Alarcón and Pascalutsa, PRC 90 (2014)]

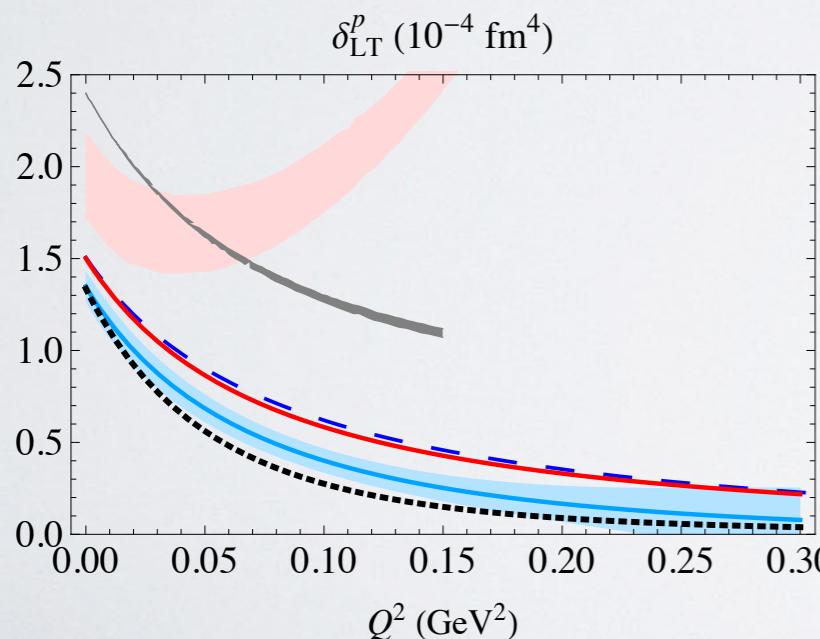
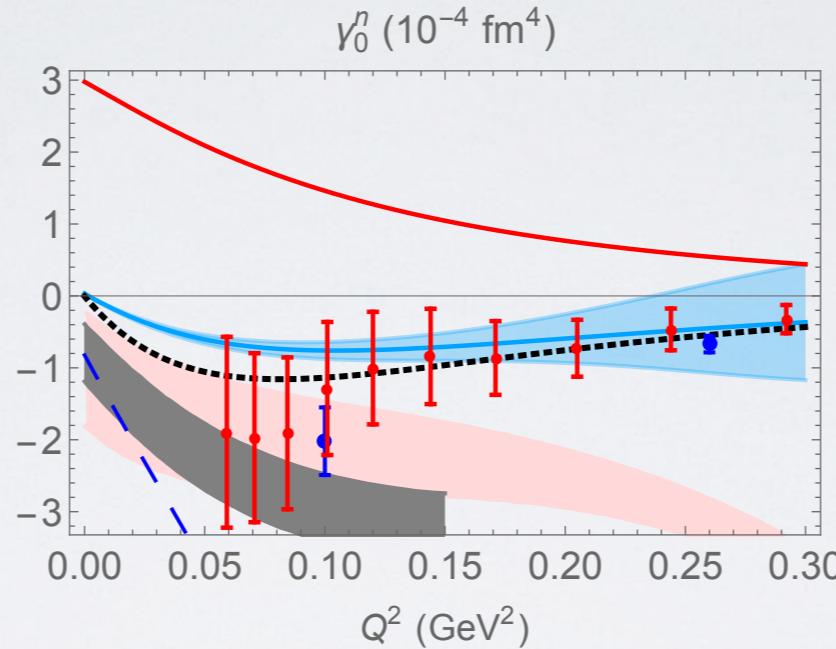
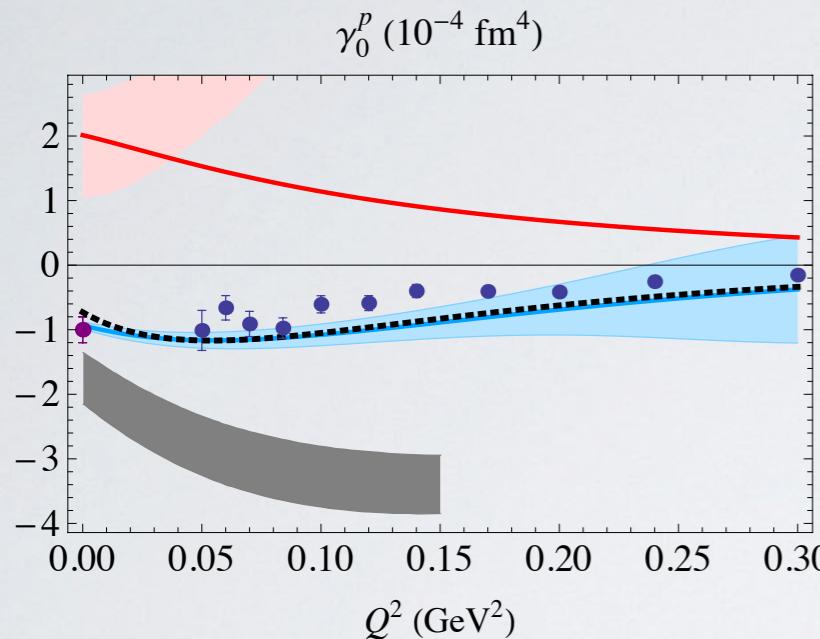


- BChPT+ $\Delta$
- LO BChPT
- LO HB
- [Kao et al., PRD 67 (2003)]
- - - MAID
- BChPT+ $\Delta^*$
- [Bernard et al., PRD 87 (2013)]
- IR+ $\Delta$
- [Bernard et al., PRD 67 (2003)]
- (Proton) Prok et al., PLB 672 (2009)
- (Neutron) Amarian et al. PRL 93 (2004)
- Dutz, et al., PRL 91(2003)
- Guler et al., PRC 92 (2015)

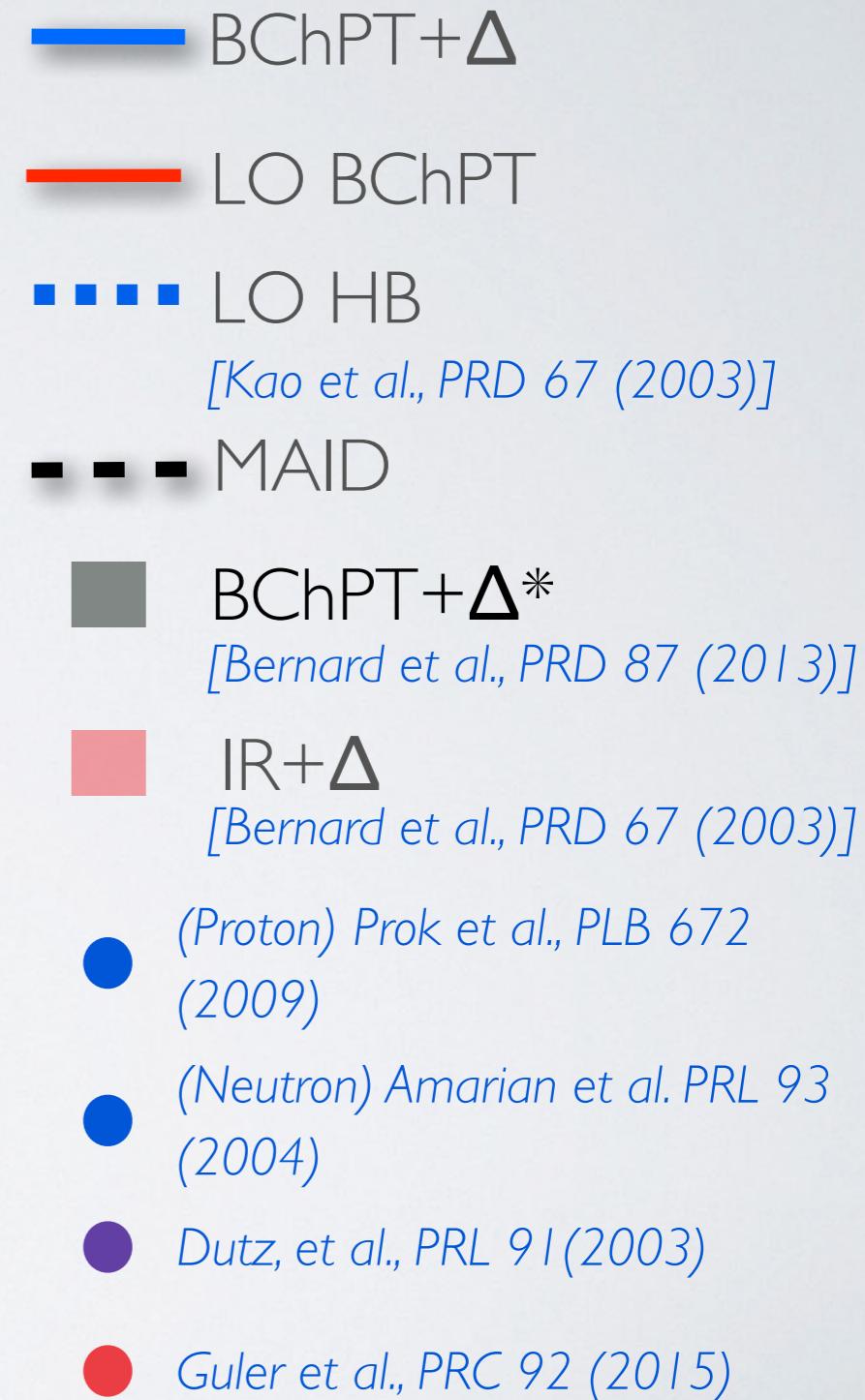
# Nucleon Polarizabilities

- Relativistic Chiral EFT+ $\Delta$  calculation of forward VVCS up to  $\mathcal{O}\left(\frac{p^4}{\Delta}\right)$

- Spin Polarizabilities:



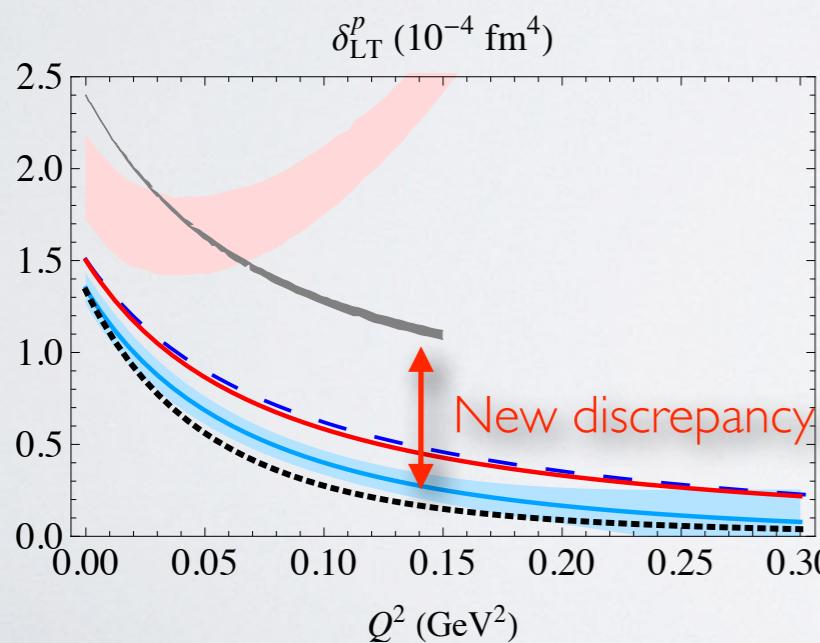
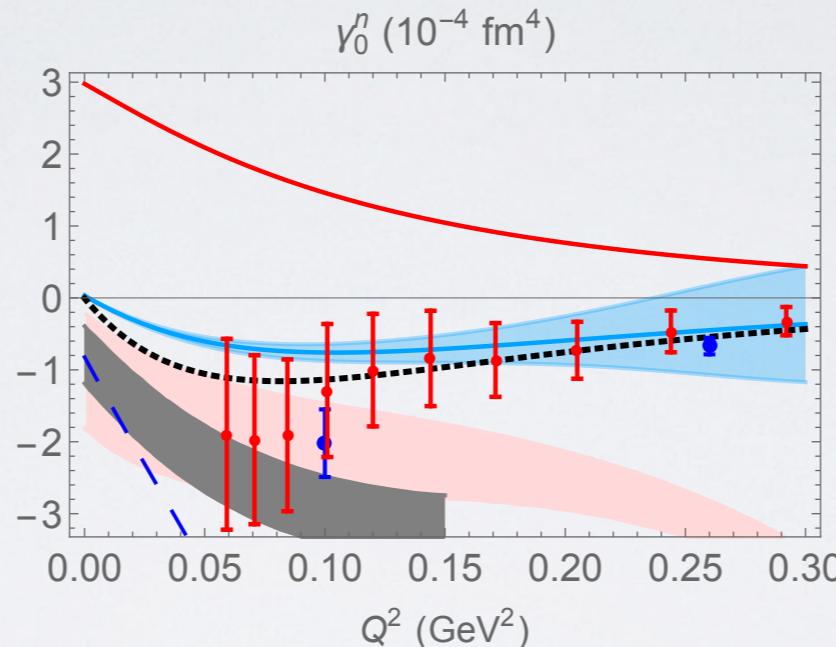
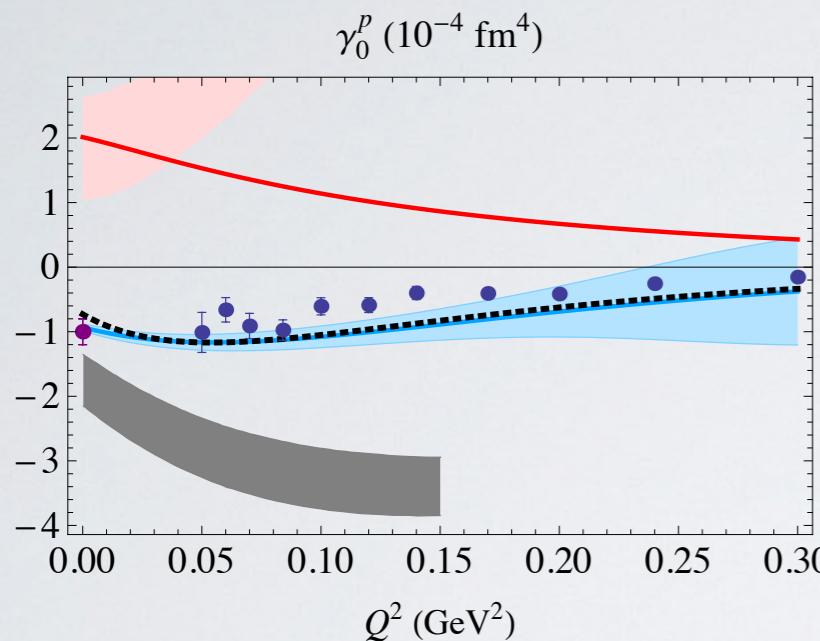
[Lensky, Alarcón and Pascalutsa, PRC 90 (2014)]



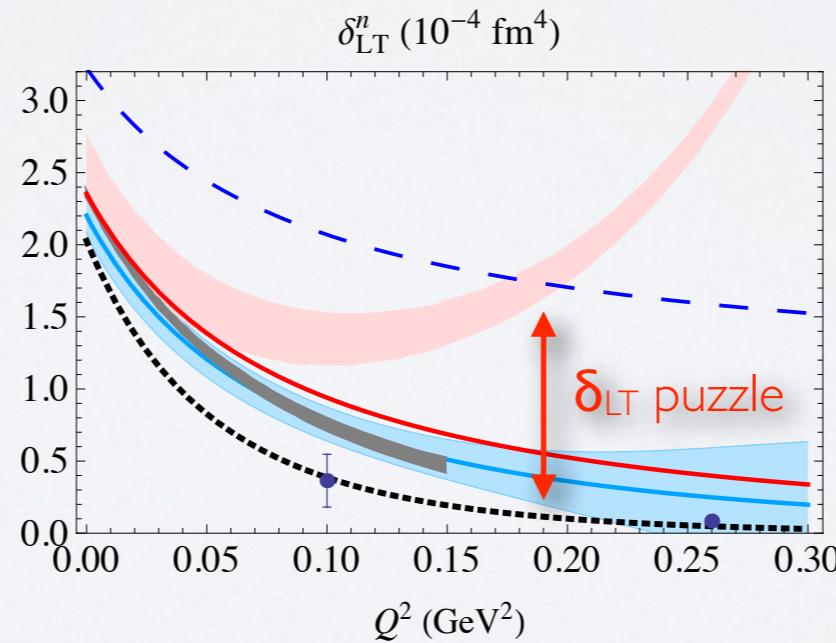
# Nucleon Polarizabilities

- Relativistic Chiral EFT+ $\Delta$  calculation of forward VVCS up to  $\mathcal{O}\left(\frac{p^4}{\Delta}\right)$

- Spin Polarizabilities:



[Lensky, Alarcón and Pascalutsa, PRC 90 (2014)]



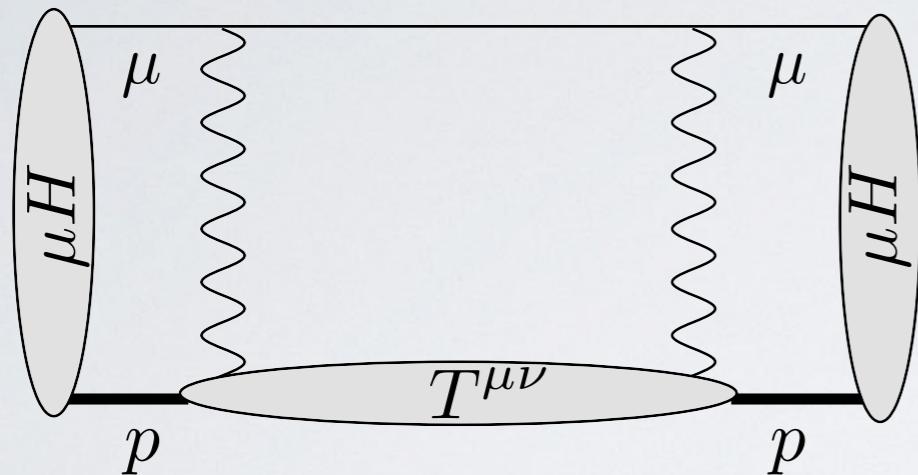
- BChPT+ $\Delta$
- LO BChPT
- LO HB
- [Kao et al., PRD 67 (2003)]
- - - MAID
- BChPT+ $\Delta^*$
- [Bernard et al., PRD 87 (2013)]
- IR+ $\Delta$
- [Bernard et al., PRD 67 (2003)]
- (Proton) Prok et al., PLB 672 (2009)
- (Neutron) Amarian et al. PRL 93 (2004)
- Dutz, et al., PRL 91(2003)
- Guler et al., PRC 92 (2015)

# Lamb shift

- The relativistic structure is important to agree with phenomenological determinations of  $\Delta E_{2S}^{(\text{pol})}$  [Alarcón, Lensky, Pascalutsa, EPJ C 74 (2014)].

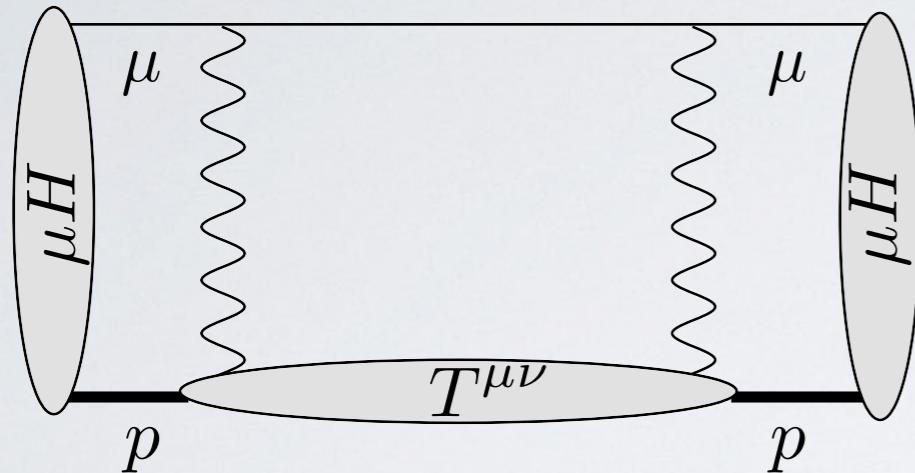
# Lamb shift

- The relativistic structure is important to agree with phenomenological determinations of  $\Delta E_{2S}^{(\text{pol})}$  [Alarcón, Lensky, Pascalutsa, EPJ C 74 (2014)].



# Lamb shift

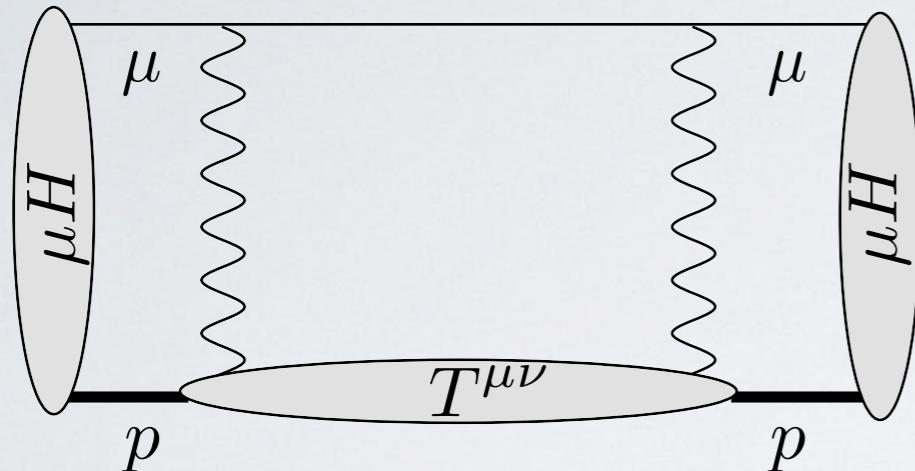
- The relativistic structure is important to agree with phenomenological determinations of  $\Delta E_{2S}^{(\text{pol})}$  [Alarcón, Lensky, Pascalutsa, EPJ C 74 (2014)].



$$\begin{aligned} T^{\mu\nu}(P, q) = & - \left( g^{\mu\nu} + \frac{q^\mu q_\nu}{q^2} \right) T_1(\nu, Q^2) \\ & + \frac{1}{m_N^2} \left( P^\mu - \frac{P \cdot q}{q^2} q^\mu \right) \left( P^\nu - \frac{P \cdot q}{q^2} q^\nu \right) T_2(\nu, Q^2) \end{aligned}$$

# Lamb shift

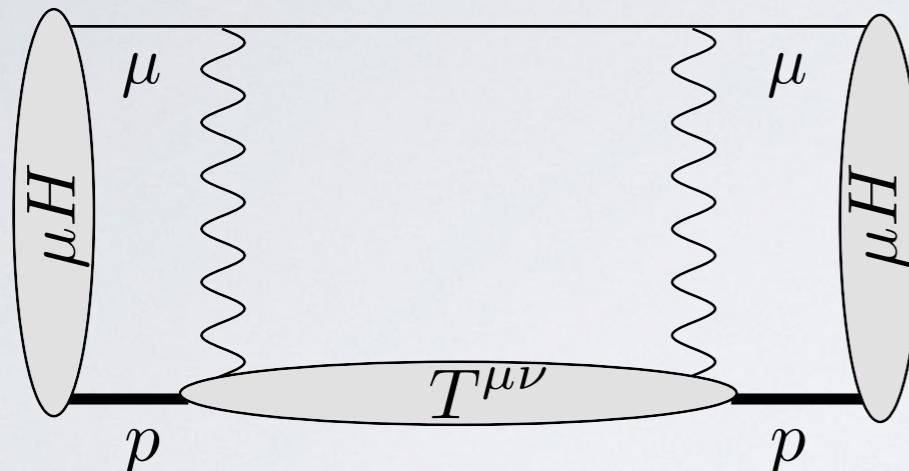
- The relativistic structure is important to agree with phenomenological determinations of  $\Delta E_{2S}^{(pol)}$  [Alarcón, Lensky, Pascalutsa, EPJ C 74 (2014)].



$$\begin{aligned} T^{\mu\nu}(P, q) = & - \left( g^{\mu\nu} + \frac{q^\mu q_\nu}{q^2} \right) T_1(\nu, Q^2) \\ & + \frac{1}{m_N^2} \left( P^\mu - \frac{P \cdot q}{q^2} q^\mu \right) \left( P^\nu - \frac{P \cdot q}{q^2} q^\nu \right) T_2(\nu, Q^2) \\ \Delta E_{2S}^{(pol)} \approx & \frac{\alpha_{em}}{\pi} \phi_{n=2}^2 \int_0^\infty \frac{dQ}{Q^2} w(\tau_\ell) \left[ T_1^{(NB)}(0, Q^2) - T_2^{(NB)}(0, Q^2) \right] \end{aligned}$$

# Lamb shift

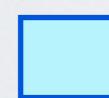
- The relativistic structure is important to agree with phenomenological determinations of  $\Delta E_{2S}^{(pol)}$  [Alarcón, Lensky, Pascalutsa, EPJ C 74 (2014)].



$$T^{\mu\nu}(P, q) = -\left(g^{\mu\nu} + \frac{q^\mu q_\nu}{q^2}\right)T_1(\nu, Q^2) + \frac{1}{m_N^2} \left(P^\mu - \frac{P \cdot q}{q^2} q^\mu\right) \left(P^\nu - \frac{P \cdot q}{q^2} q^\nu\right) T_2(\nu, Q^2)$$

$$\Delta E_{2S}^{(pol)} \approx \frac{\alpha_{em}}{\pi} \phi_{n=2}^2 \int_0^\infty \frac{dQ}{Q^2} w(\tau_\ell) \left[ T_1^{(NB)}(0, Q^2) - T_2^{(NB)}(0, Q^2) \right]$$

(μeV)	[1]	[2]	[3]	[4]	[5]	[6]	[7]	[8]
$\Delta E_{2S}^{(pol)}$	-12(2)	-11.5	-18.5	-7.4(2.4)	-8.5(1.1)	-15.3(5.6)	$-8.2^{+2.0}_{-2.5}$	-26.5



Chiral EFT calculations



Phenomenological determinations (dispersion relations+data)

[1] K. Pachucki, Phys. Rev. A 60 (1999).

[2] A. P. Martynenko, Phys. Atom. Nucl. 69 (2006).

[3] D. Nevado and A. Pineda, Phys. Rev. C 77 (2008).

[4] C. E. Carlson and M. Vanderhaeghen, Phys. Rev. A 84, (2011).

[5] M. C. Birse and J. A. McGovern, Eur. Phys. J. A 48, (2012).

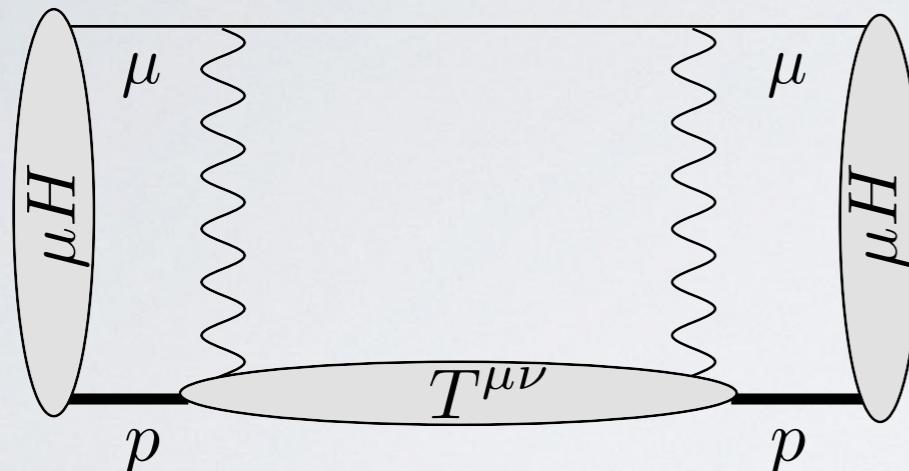
[6] M. Gorchtein, F. J. LLanes-Estrada and A. P. Szczepaniak, Phys. Rev. A 87 (2013).

[7] J. M. Alarcón, V. Lensky, V. Pascalutsa, Eur. Phys. J. C 74 (2014).

[8] C. Peset and A. Pineda Eur. Phys. J. A 51 (2015).

# Lamb shift

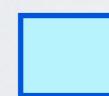
- The relativistic structure is important to agree with phenomenological determinations of  $\Delta E_{2S}^{(pol)}$  [Alarcón, Lensky, Pascalutsa, EPJ C 74 (2014)].



$$T^{\mu\nu}(P, q) = -\left(g^{\mu\nu} + \frac{q^\mu q_\nu}{q^2}\right)T_1(\nu, Q^2) + \frac{1}{m_N^2} \left(P^\mu - \frac{P \cdot q}{q^2} q^\mu\right) \left(P^\nu - \frac{P \cdot q}{q^2} q^\nu\right) T_2(\nu, Q^2)$$

$$\Delta E_{2S}^{(pol)} \approx \frac{\alpha_{em}}{\pi} \phi_{n=2}^2 \int_0^\infty \frac{dQ}{Q^2} w(\tau_\ell) \left[ T_1^{(NB)}(0, Q^2) - T_2^{(NB)}(0, Q^2) \right]$$

(μeV)	[1]	[2]	[3]	[4]	[5]	[6]	[7]	[8]
$\Delta E_{2S}^{(pol)}$	-12(2)	-11.5	-18.5	-7.4(2.4)	-8.5(1.1)	-15.3(5.6)	-8.2 <sup>+2.0</sup> <sub>-2.5</sub>	-26.5



Chiral EFT calculations



Phenomenological determinations (dispersion relations+data)

[1] K. Pachucki, Phys. Rev. A 60 (1999).

[2] A. P. Martynenko, Phys. Atom. Nucl. 69 (2006).

[3] D. Nevado and A. Pineda, Phys. Rev. C 77 (2008).

[4] C. E. Carlson and M. Vanderhaeghen, Phys. Rev. A 84, (2011).

[5] M. C. Birse and J. A. McGovern, Eur. Phys. J. A 48, (2012).

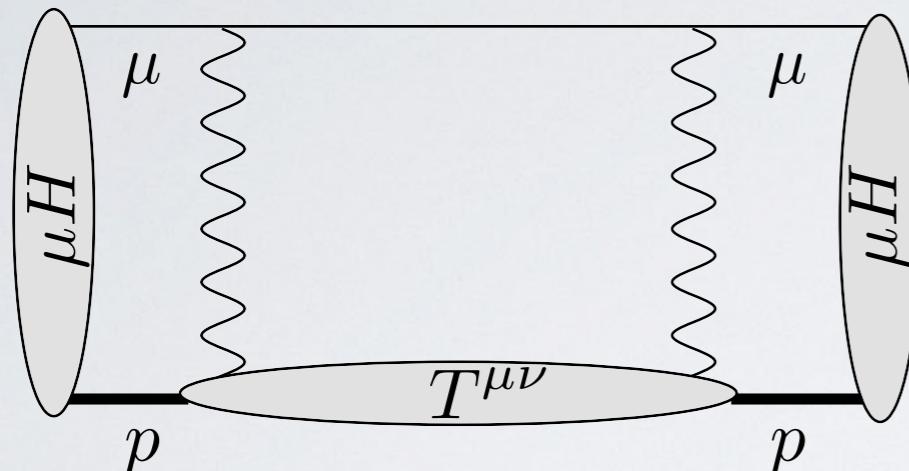
[6] M. Gorchtein, F. J. LLanes-Estrada and A. P. Szczepaniak, Phys. Rev. A 87 (2013).

[7] J. M. Alarcón, V. Lensky, V. Pascalutsa, Eur. Phys. J. C 74 (2014).

[8] C. Peset and A. Pineda Eur. Phys. J. A 51 (2015).

# Lamb shift

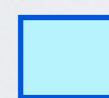
- The relativistic structure is important to agree with phenomenological determinations of  $\Delta E_{2S}^{(pol)}$  [Alarcón, Lensky, Pascalutsa, EPJ C 74 (2014)].



$$T^{\mu\nu}(P, q) = -\left(g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2}\right)T_1(\nu, Q^2) + \frac{1}{m_N^2} \left(P^\mu - \frac{P \cdot q}{q^2} q^\mu\right) \left(P^\nu - \frac{P \cdot q}{q^2} q^\nu\right) T_2(\nu, Q^2)$$

$$\Delta E_{2S}^{(pol)} \approx \frac{\alpha_{em}}{\pi} \phi_{n=2}^2 \int_0^\infty \frac{dQ}{Q^2} w(\tau_\ell) \left[ T_1^{(NB)}(0, Q^2) - T_2^{(NB)}(0, Q^2) \right]$$

(μeV)	[1]	[2]	[3]	[4]	[5]	[6]	[7]	[8]
$\Delta E_{2S}^{(pol)}$	-12(2)	-11.5	-18.5	-7.4(2.4)	-8.5(1.1)	-15.3(5.6)	-8.2 <sup>+2.0</sup> <sub>-2.5</sub>	-26.5



Chiral EFT calculations



Phenomenological determinations (dispersion relations+data)

[1] K. Pachucki, Phys. Rev. A 60 (1999).

[2] A. P. Martynenko, Phys. Atom. Nucl. 69 (2006).

[3] D. Nevado and A. Pineda, Phys. Rev. C 77 (2008).

[4] C. E. Carlson and M. Vanderhaeghen, Phys. Rev. A 84, (2011).

[5] M. C. Birse and J. A. McGovern, Eur. Phys. J. A 48, (2012).

[6] M. Gorchtein, F. J. LLanes-Estrada and A. P. Szczepaniak, Phys. Rev. A 87 (2013).

[7] J. M. Alarcón, V. Lensky, V. Pascalutsa, Eur. Phys. J. C 74 (2014).

[8] C. Peset and A. Pineda Eur. Phys. J. A 51 (2015).

# *Nuclear Lattice Effective Field Theory*

# *NLEFT*

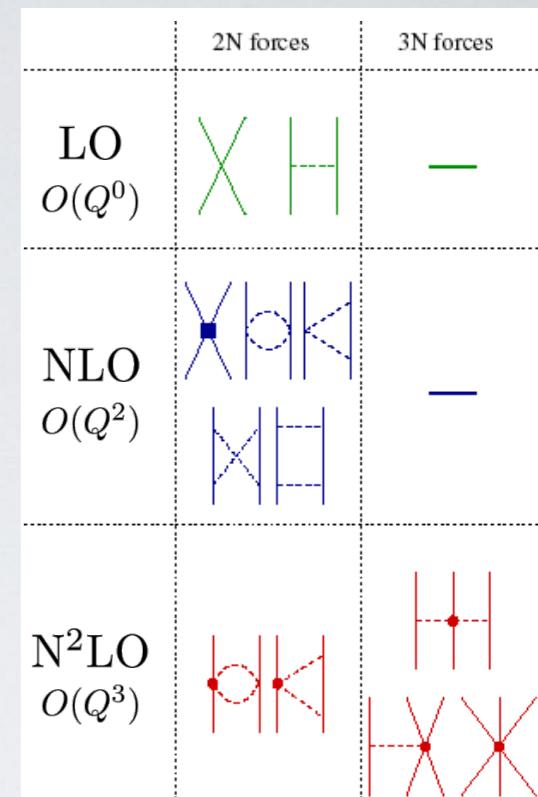
- Lattice techniques to study nuclear many-body problems with chiral EFT.

# *NLEFT*

- Lattice techniques to study nuclear many-body problems with chiral EFT.
- $2N$  forces is crucial in studies of many-body nuclear interactions with chiral EFT:

# NLEFT

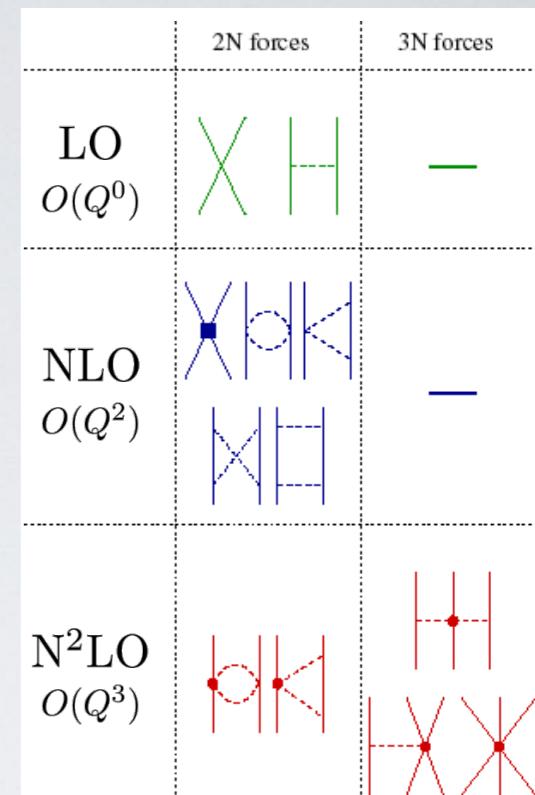
- Lattice techniques to study nuclear many-body problems with chiral EFT.
- 2N forces is crucial in studies of many-body nuclear interactions with chiral EFT:



[Courtesy of Dean Lee]

# NLEFT

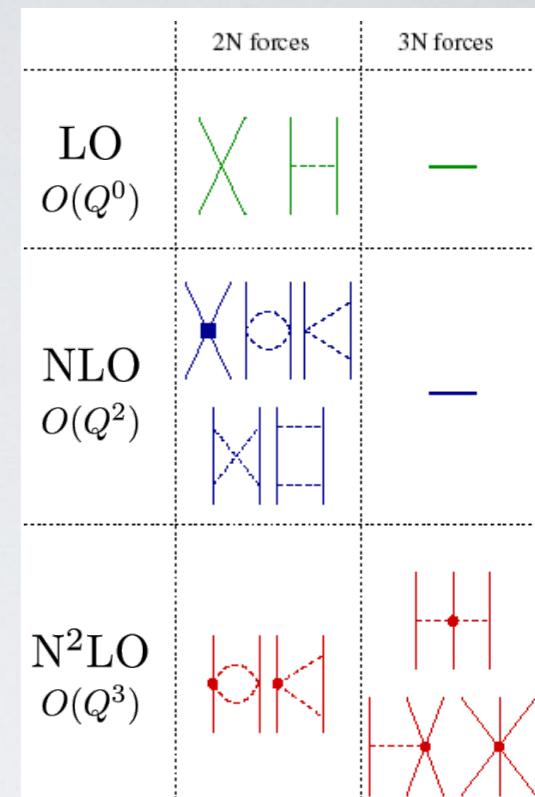
- Lattice techniques to study nuclear many-body problems with chiral EFT.
- 2N forces is crucial in studies of many-body nuclear interactions with chiral EFT:
  - Determination of NN LECs.



[Courtesy of Dean Lee]

# NLEFT

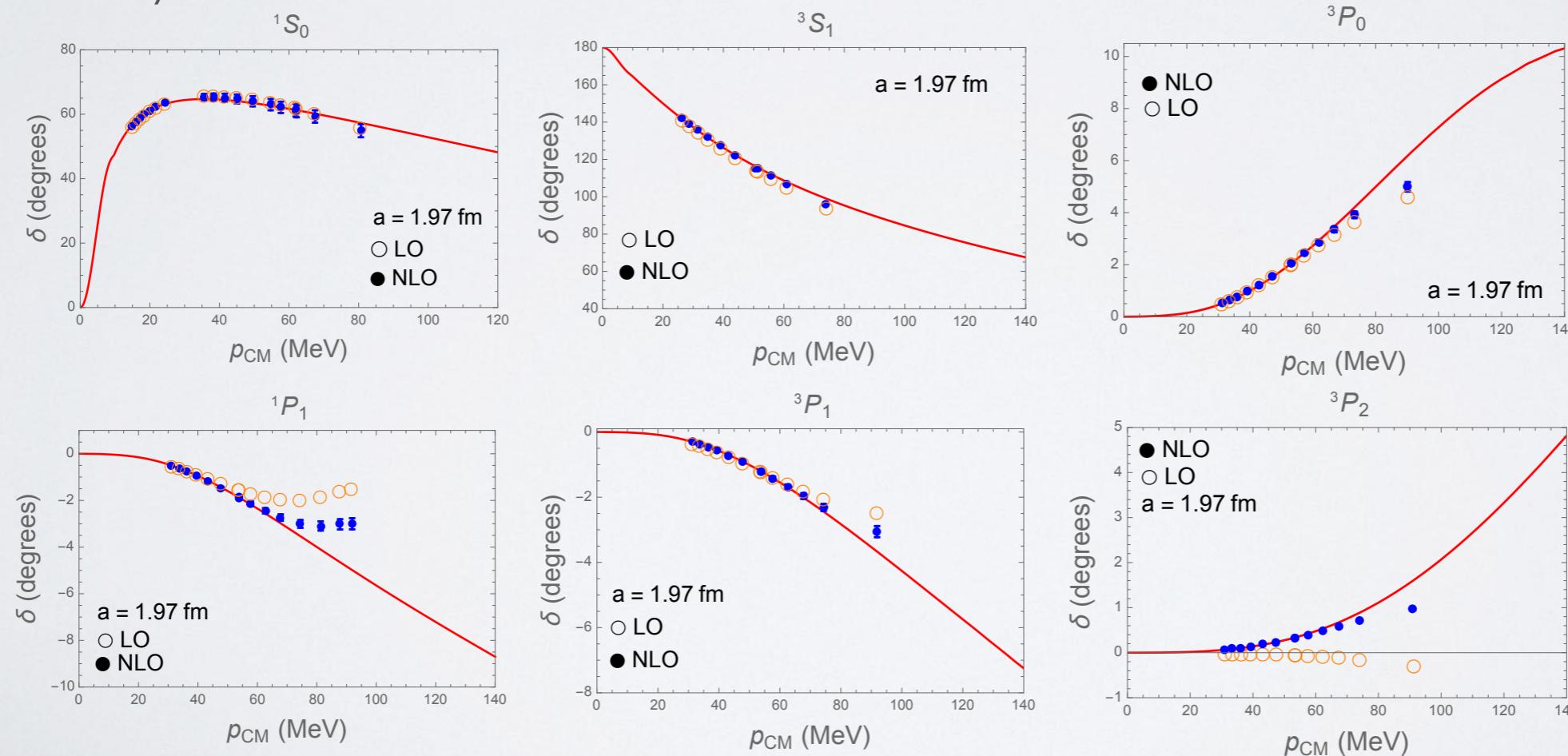
- Lattice techniques to study nuclear many-body problems with chiral EFT.
- 2N forces is crucial in studies of many-body nuclear interactions with chiral EFT:
  - Determination of NN LECs.
  - Study of the uncertainties.



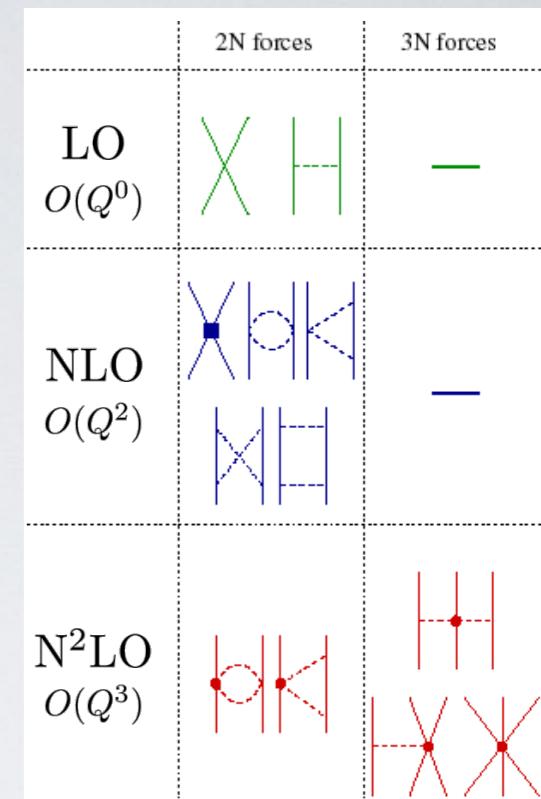
[Courtesy of Dean Lee]

# NLEFT

- Lattice techniques to study nuclear many-body problems with chiral EFT.
- 2N forces is crucial in studies of many-body nuclear interactions with chiral EFT:
  - Determination of NN LECs.
  - Study of the uncertainties.



[J. M. Alarcon, Chiral Dynamics Workshop 2015]



## *Future Projects*

# *Chiral predictions of transverse structure of baryons*

- Provide spatial picture of the charge and magnetic density of the baryon octet and densities of the decuplet.

# *Chiral predictions of transverse structure of baryons*

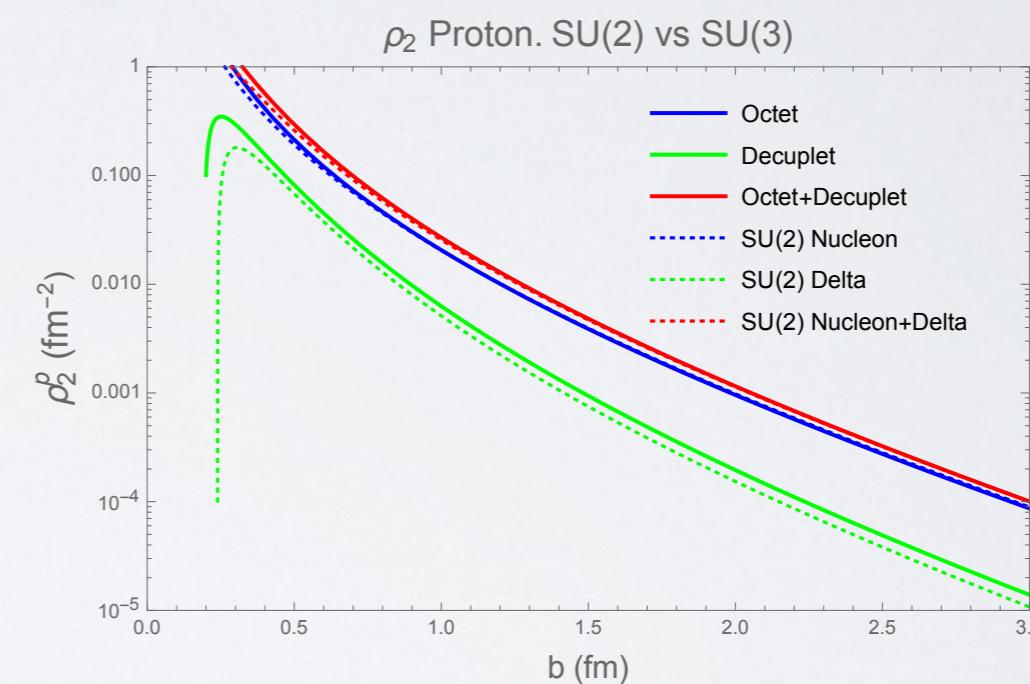
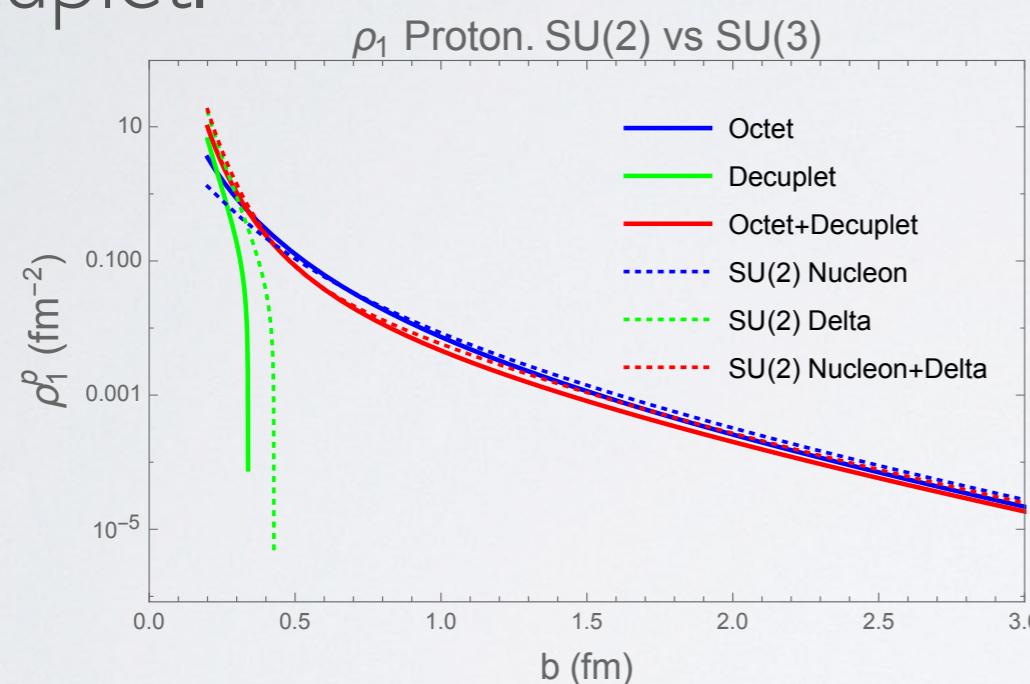
- Provide spatial picture of the charge and magnetic density of the baryon octet and densities of the decuplet.
- Chiral EFT provides predictions for the peripheral region  $b \sim l/M_\pi$ .

# *Chiral predictions of transverse structure of baryons*

- Provide spatial picture of the charge and magnetic density of the baryon octet and densities of the decuplet.
- Chiral EFT provides predictions for the peripheral region  $b \sim l/M_\pi$ .
- Extend the work of [Granados & Weiss, *JHEP* 1401 (2014)] to the octet and decuplet.

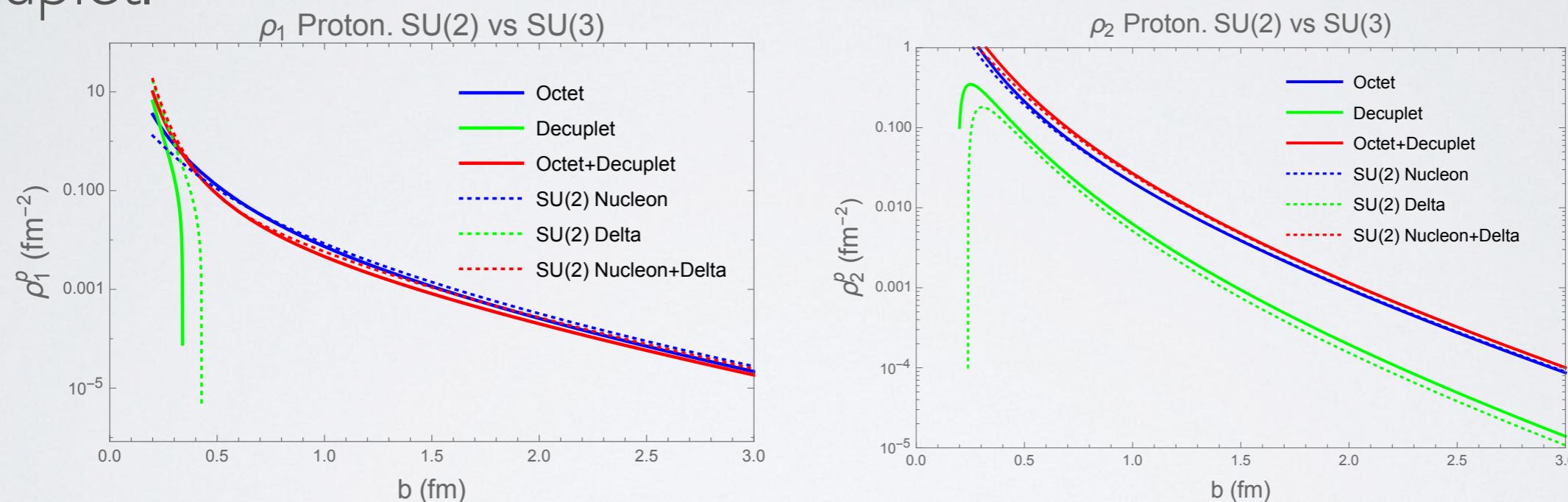
# Chiral predictions of transverse structure of baryons

- Provide spatial picture of the charge and magnetic density of the baryon octet and densities of the decuplet.
- Chiral EFT provides predictions for the peripheral region  $b \sim l/M_\pi$ .
- Extend the work of [Granados & Weiss, JHEP 1401 (2014)] to the octet and decuplet.



# Chiral predictions of transverse structure of baryons

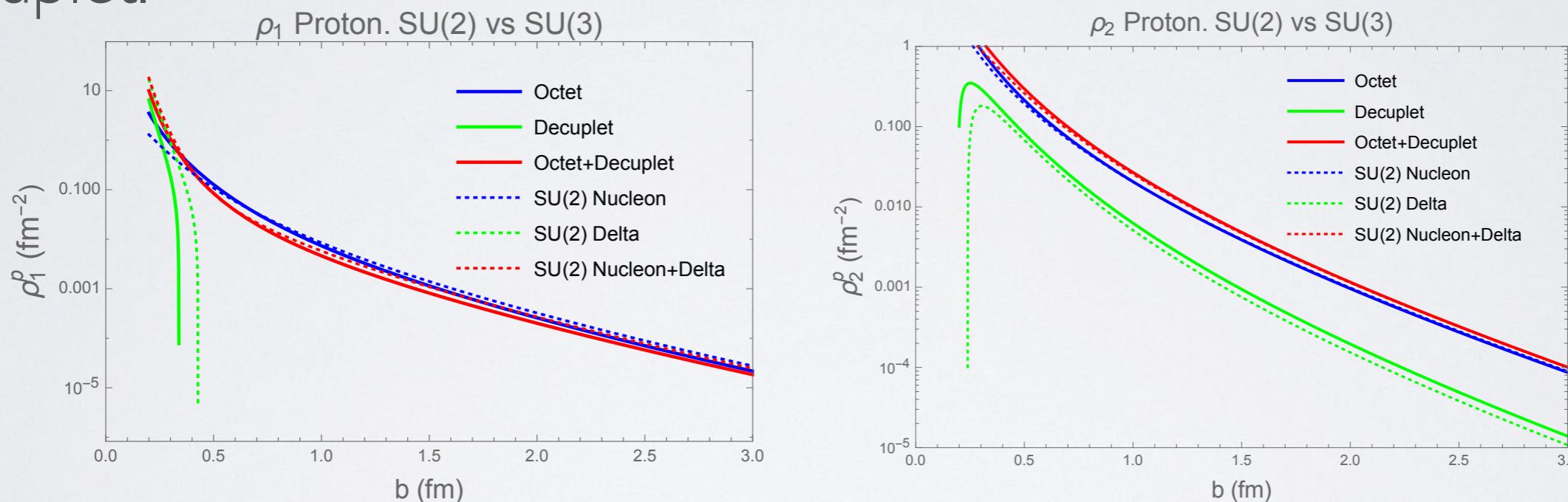
- Provide spatial picture of the charge and magnetic density of the baryon octet and densities of the decuplet.
- Chiral EFT provides predictions for the peripheral region  $b \sim l/M_\pi$ .
- Extend the work of [Granados & Weiss, JHEP 1401 (2014)] to the octet and decuplet.



- Check the large- $N_c$  relations in chiral EFT.

# Chiral predictions of transverse structure of baryons

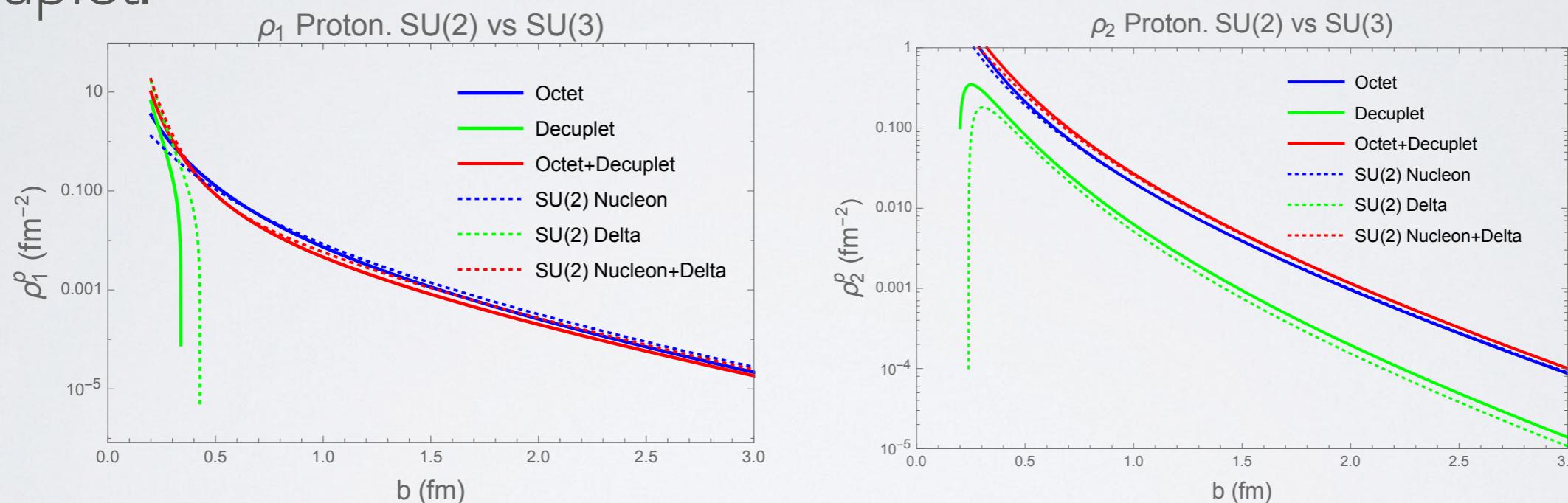
- Provide spatial picture of the charge and magnetic density of the baryon octet and densities of the decuplet.
- Chiral EFT provides predictions for the peripheral region  $b \sim l/M_\pi$ .
- Extend the work of [Granados & Weiss, JHEP 1401 (2014)] to the octet and decuplet.



- Check the large- $N_c$  relations in chiral EFT.
  - We saw that different formulations with  $\Delta$  are subdominant in  $N_c$ .

# Chiral predictions of transverse structure of baryons

- Provide spatial picture of the charge and magnetic density of the baryon octet and densities of the decuplet.
- Chiral EFT provides predictions for the peripheral region  $b \sim l/M_\pi$ .
- Extend the work of [Granados & Weiss, JHEP 1401 (2014)] to the octet and decuplet.



- Check the large- $N_c$  relations in chiral EFT.
  - We saw that different formulations with  $\Delta$  are subdominant in  $N_c$ .  
→ They do not modify the correct  $N_c$  scaling.

## *Summary and Conclusions*

# Summary and Conclusions

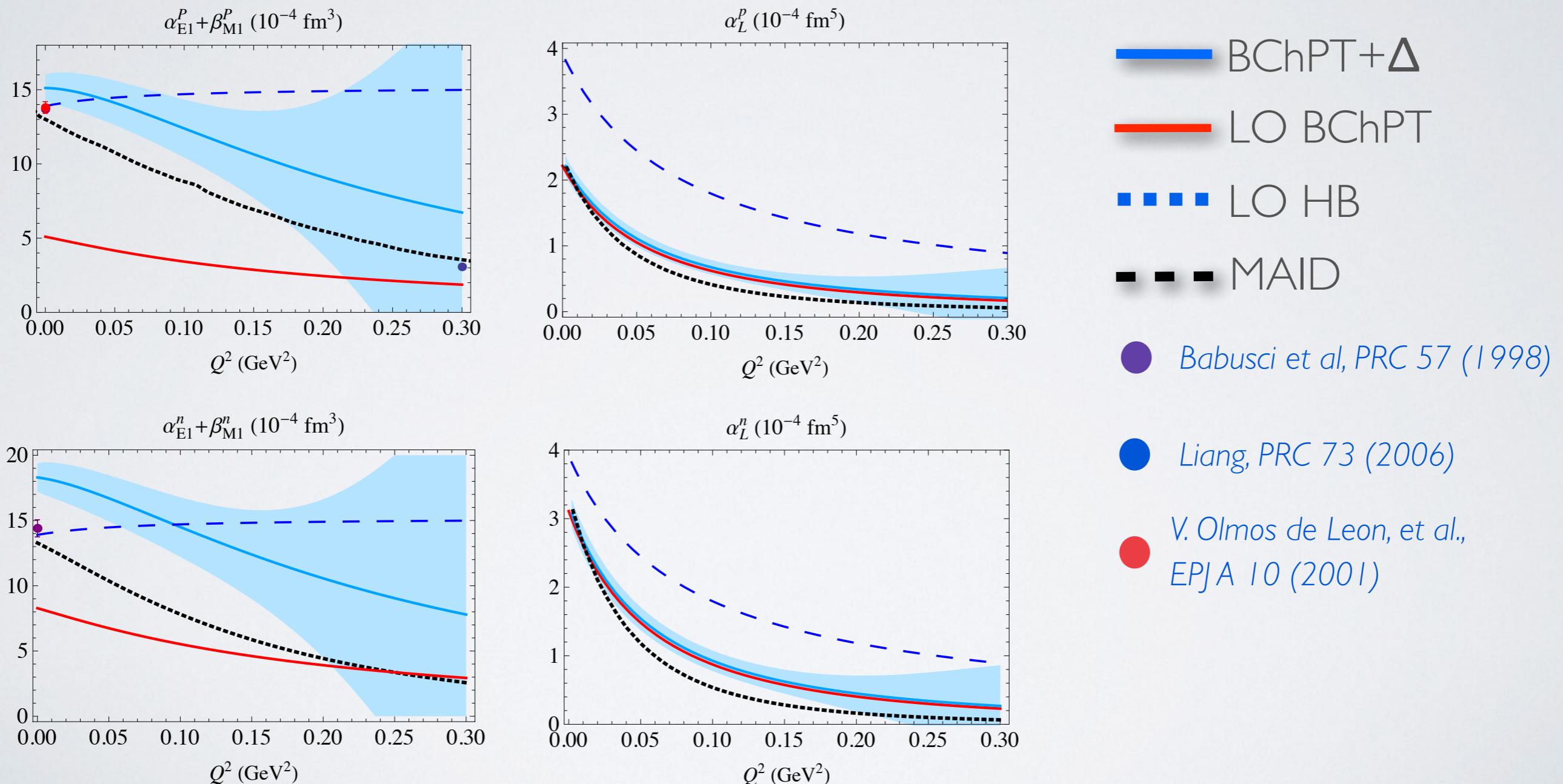
- Chiral EFT with baryons is an excellent tool to investigate fundamental hadronic interactions involving nucleons on QCD grounds.
- $\pi N$ 
  - Good description of modern scattering data below the  $\Delta$  peak.
  - Agreement with dispersive extractions.
  - Extraction of important quantities from phenomenology (  $\sigma_{\pi N}$  )
- Forward doubly virtual Compton scattering
  - Prediction of polarizabilities in agreement with MAID model and experimental data → Improves previous ChPT predictions.
- Nuclear Lattice EFT
  - Fundamental piece in *ab initio* many-body nuclear calculations.
  - Further improvements possible including spin-flavor symmetry.

**FIN**

*Spares*

# Polarizabilities

- Relativistic baryon chiral EFT with electromagnetic probes:
  - Scalar and spin VVCS Polarizabilities.
- Scalar Polarizabilities:

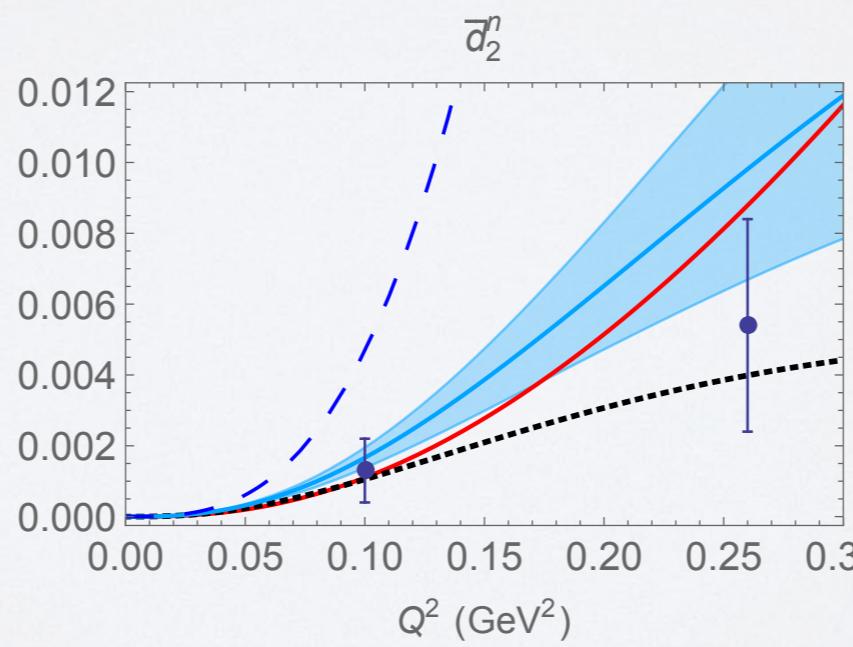
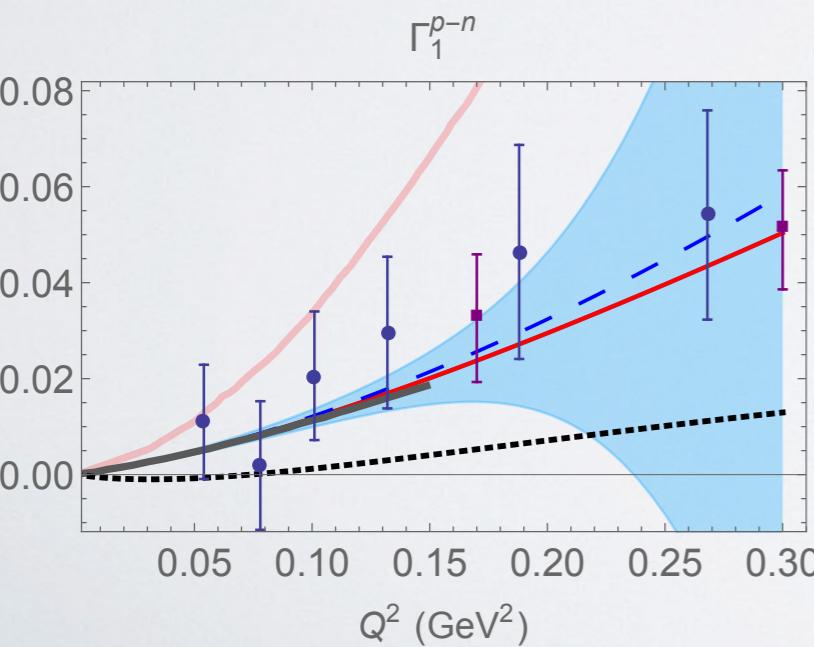
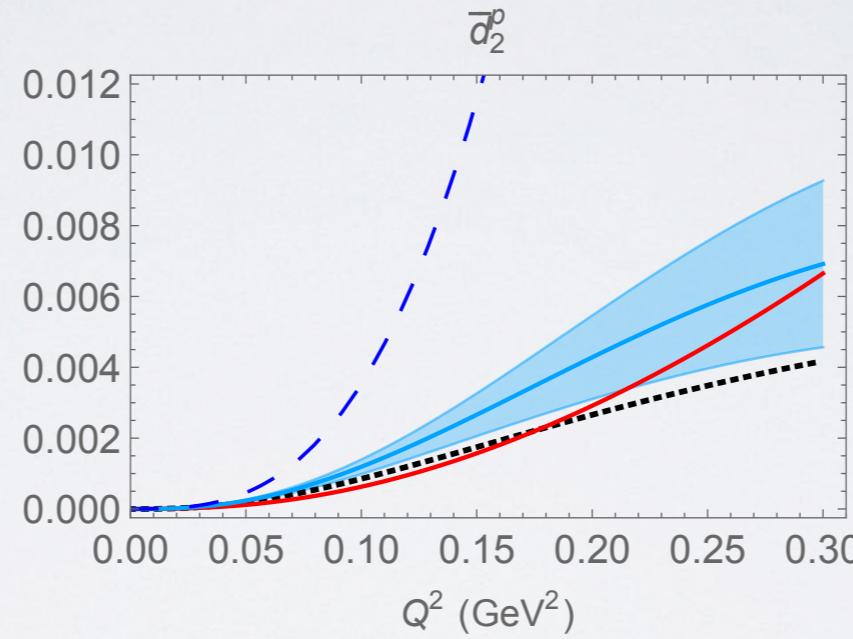
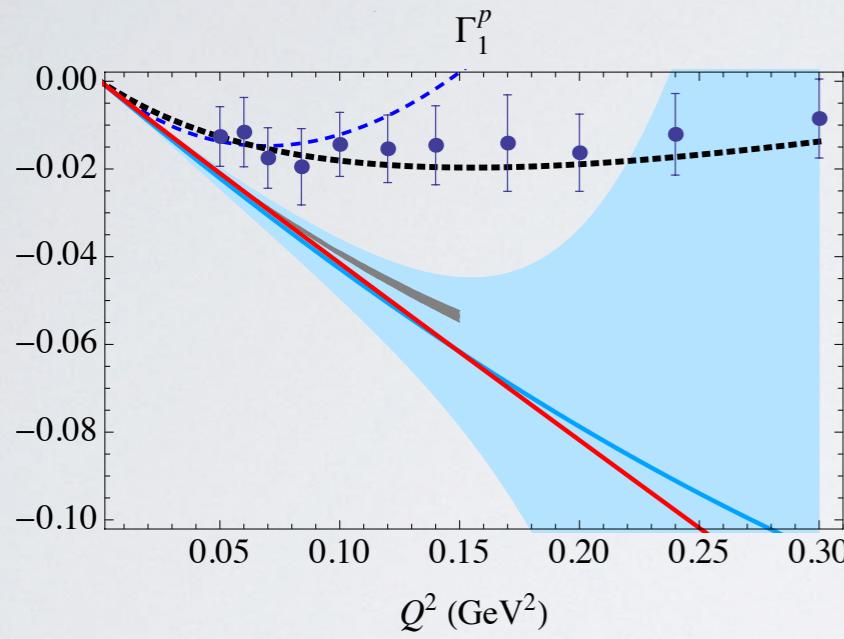


# Polarizabilities

- Some interesting moments:

$$\Gamma_1(Q^2) = \int_0^{x_0} dx g_1(x, Q^2)$$

$$\bar{d}_2(Q^2) = \int_0^{x_0} dx x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)]$$



- BChPT+ $\Delta$
- LO BChPT
- LO HB  
[Kao et al., PRD 67 (2003)]
- - - MAID
- BChPT+ $\Delta^*$   
[Bernard et al., PRD 87 (2013)]
- IR+ $\Delta$   
[Bernard et al., PRD 67 (2003)]
- $(\Gamma_1^p)$  Prok et al., PLB 672 (2009)
- $(\Gamma_1^{p-n})$  Deur et al., PRD 78 (2008)
- $(\Gamma_1^{p-n})$  Deur et al., PRL 93 (2004)
- $(d_2^n)$  Amarian et al., PRL 92 (2004)

