

Meson electro/photo- production from QCD

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OLD DOMINION
UNIVERSITY

 **Jefferson Lab**

February 2016

Questions to answer:

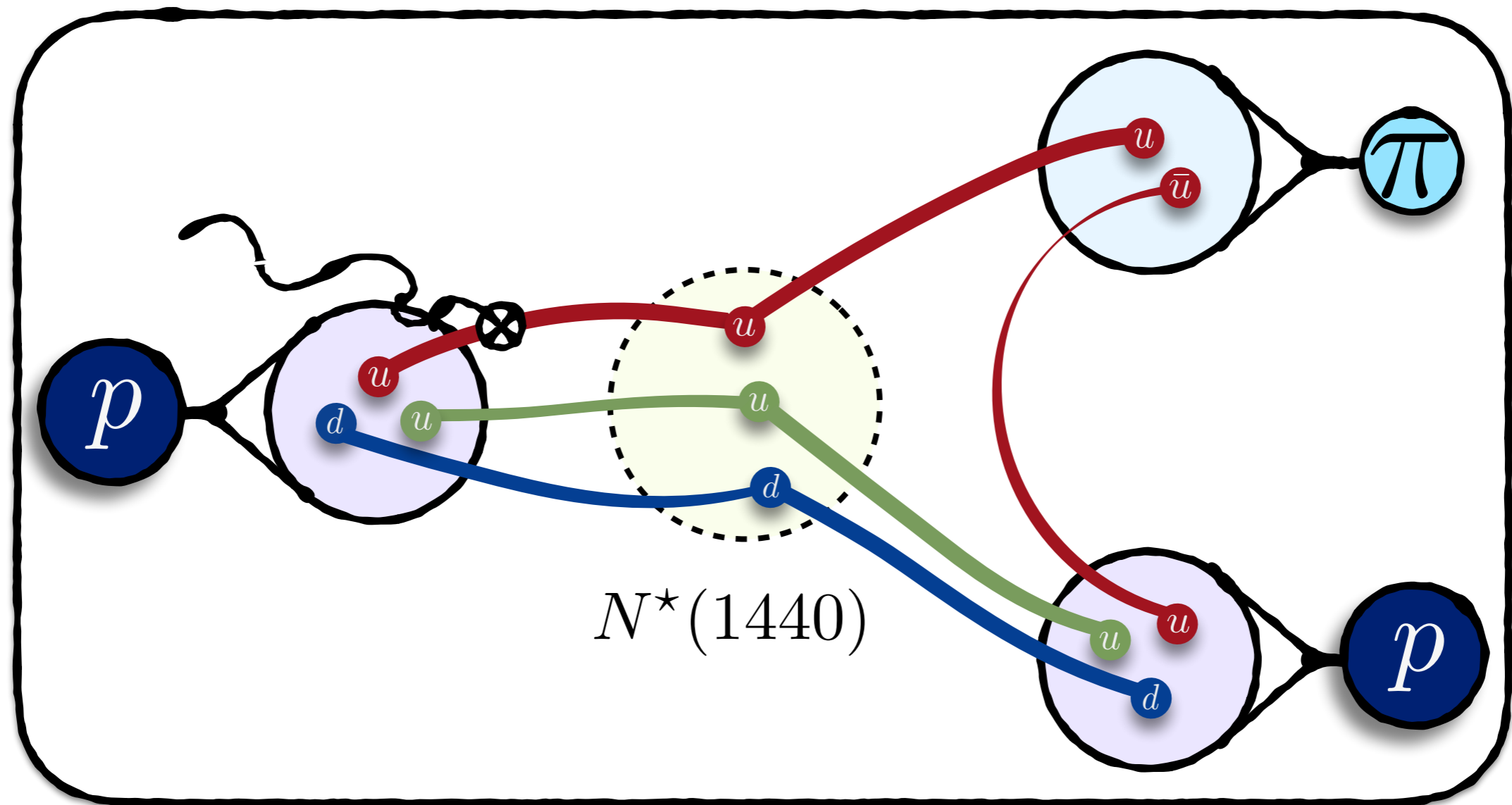
• What is possible from lattice QCD regarding electro/photo-production?

• How can lattice and experiment compare to each other?

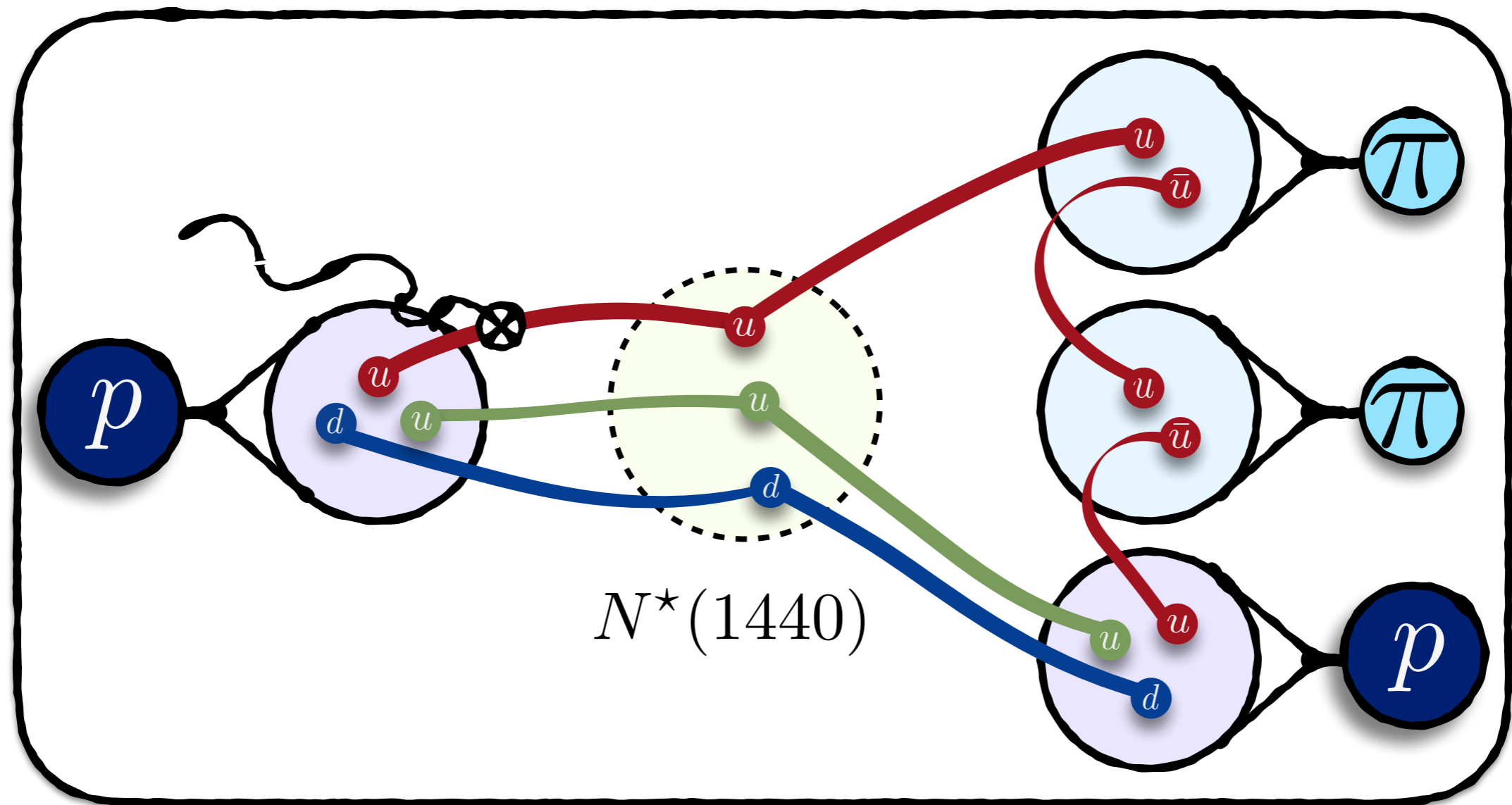
• What is the possibility/timescale for extending this to baryons?

• What does lattice “*need*” from experiment?

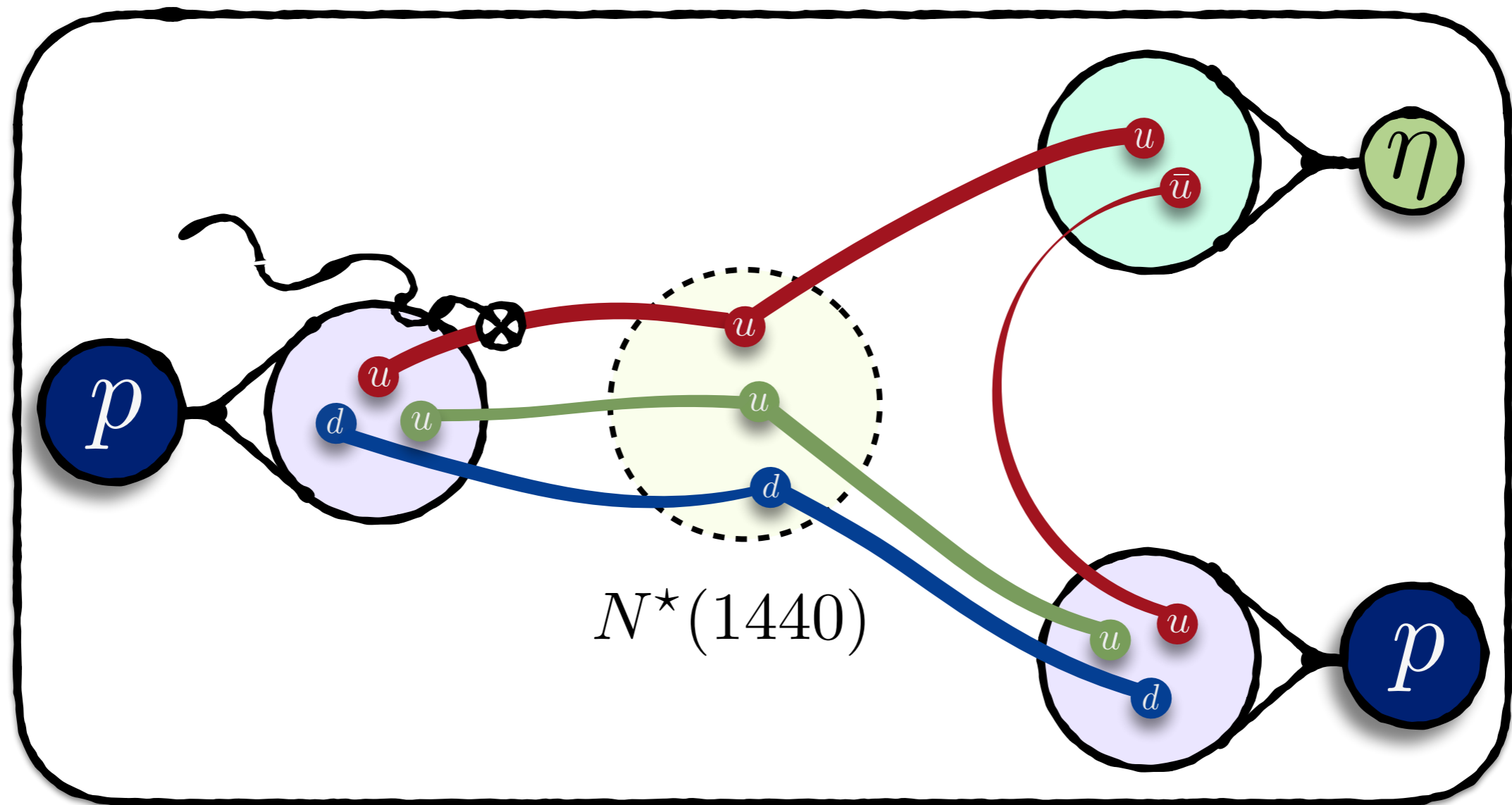
Working definition



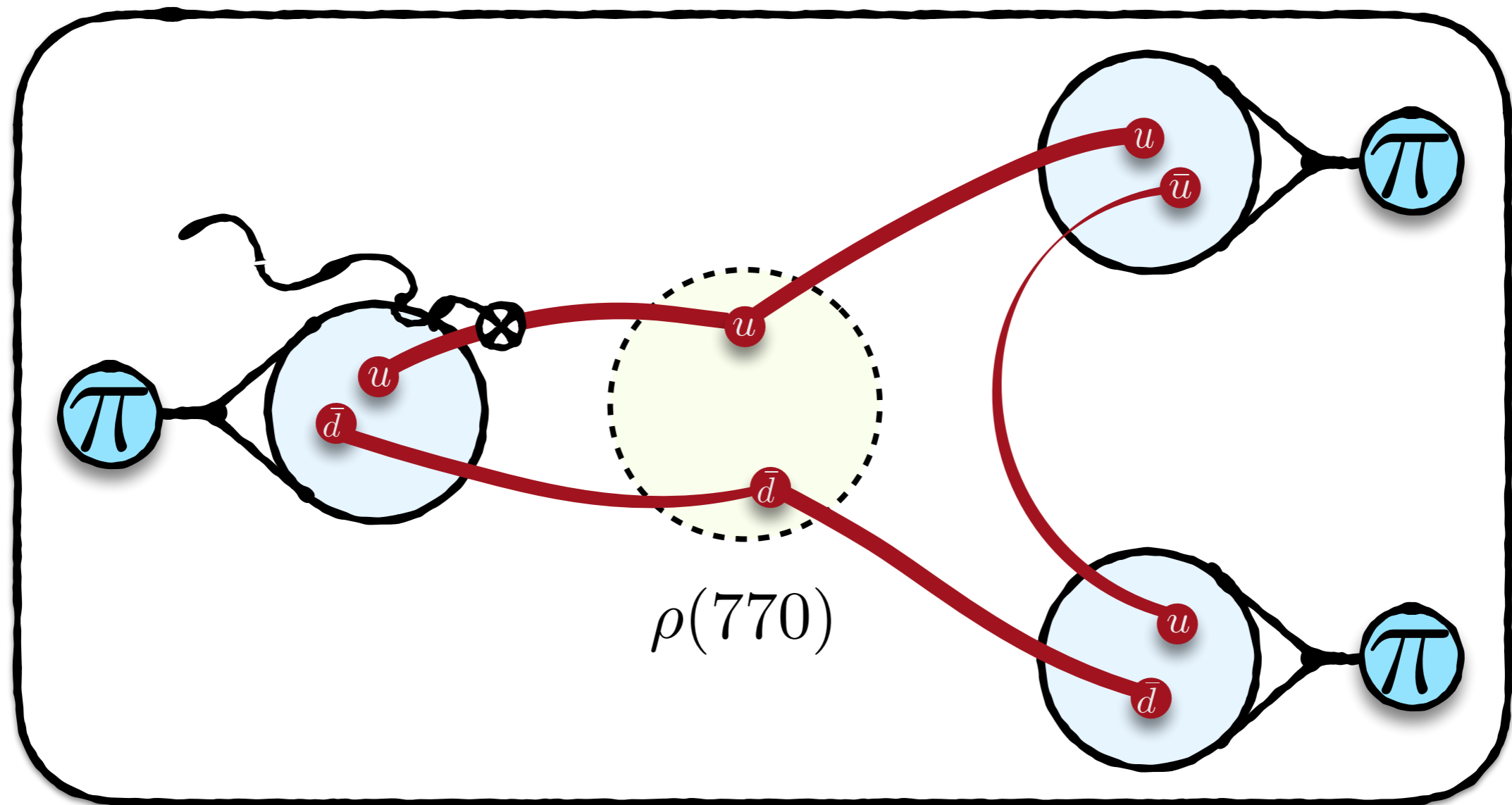
Working definition



Working definition



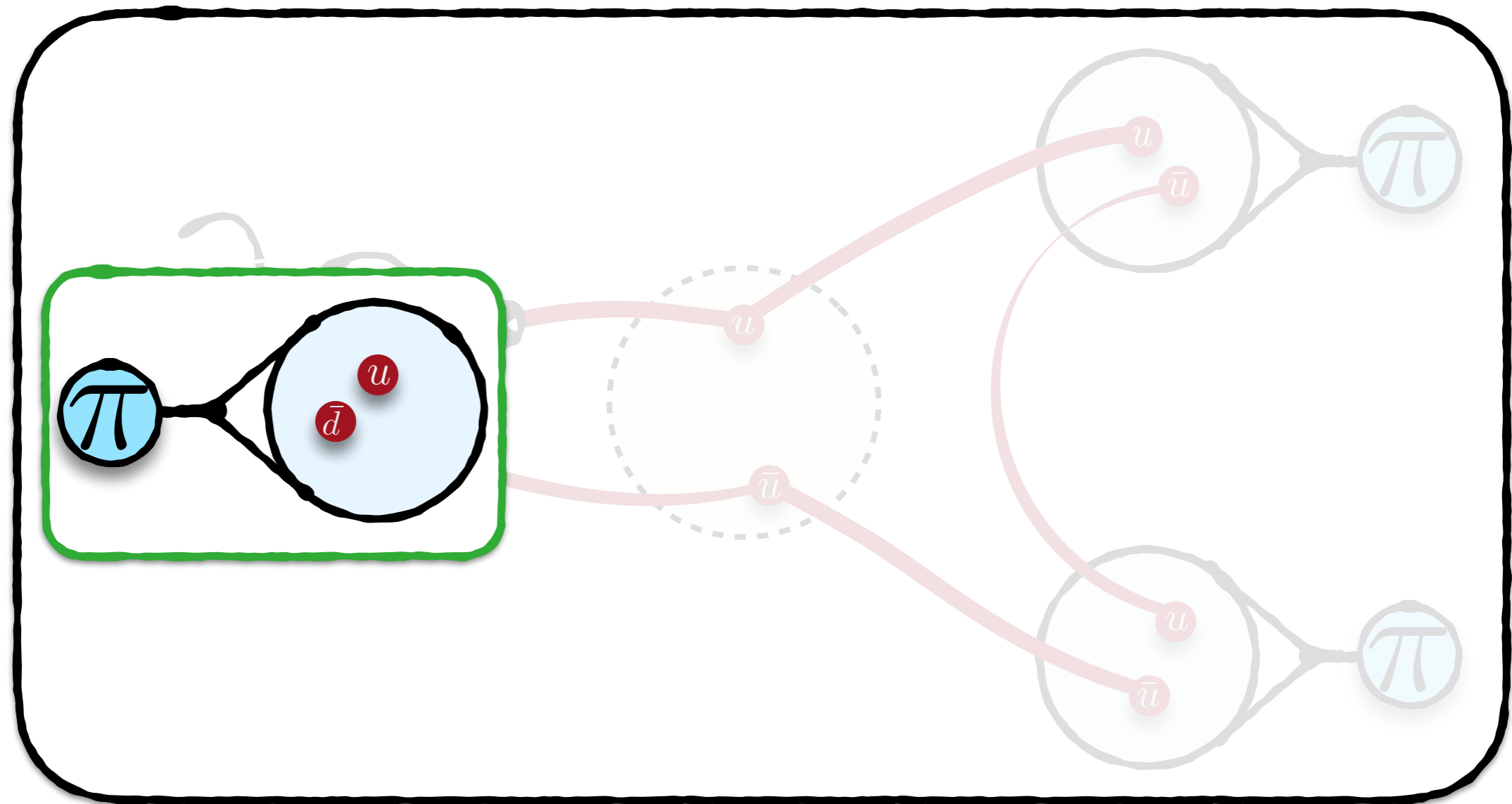
Working definition



Importance

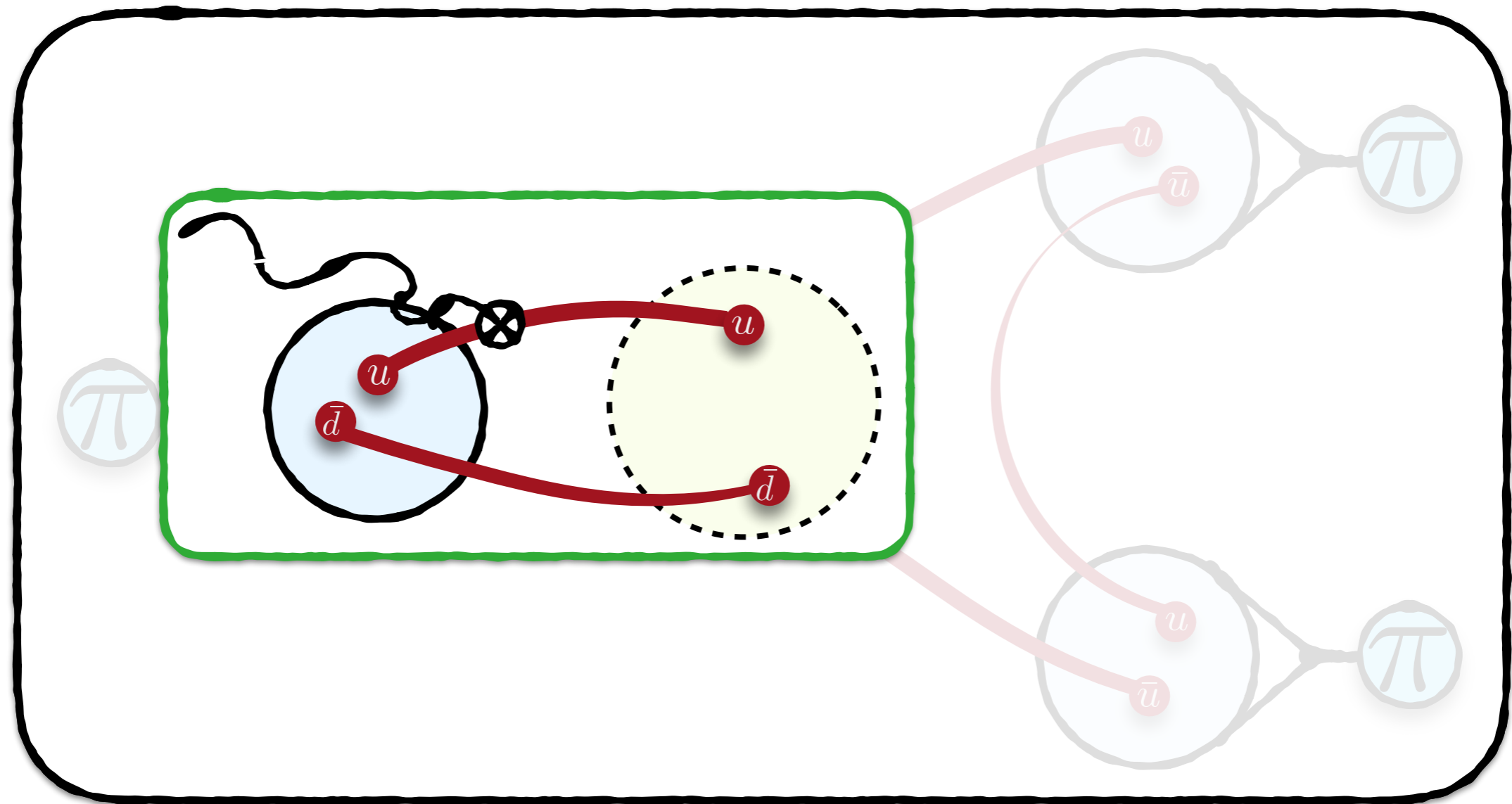
- Access the excited spectrum of QCD
- Test our understanding of QCD
- Probe the inner structure and shape of hadrons
- Test the limits of the standard model
- ...

Why lattice?



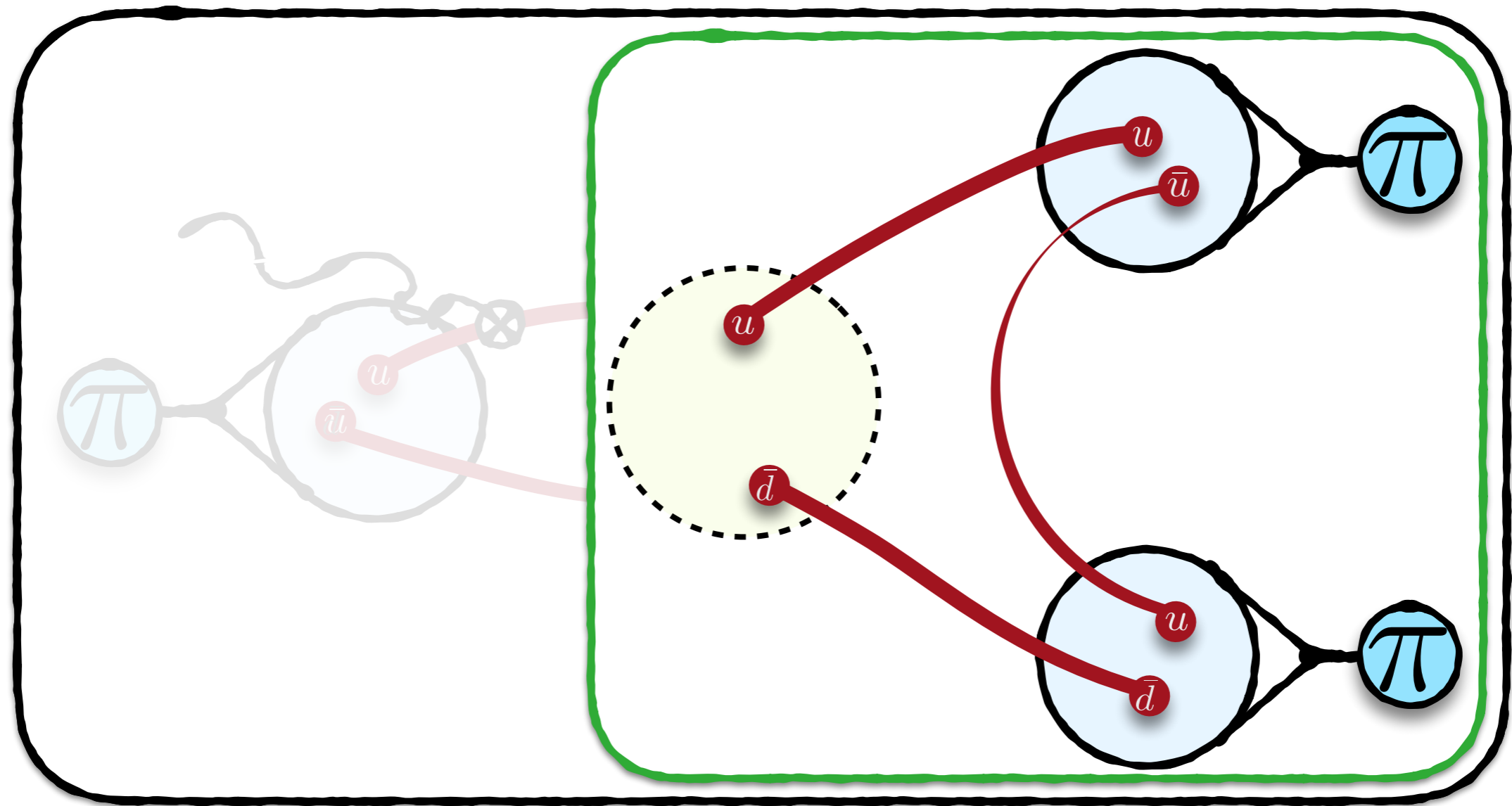
• QCD-stable states are generated exactly

Why lattice?



- QCD-stable states are generated exactly
- QED / weak sector can be treated perturbatively or non-perturbatively

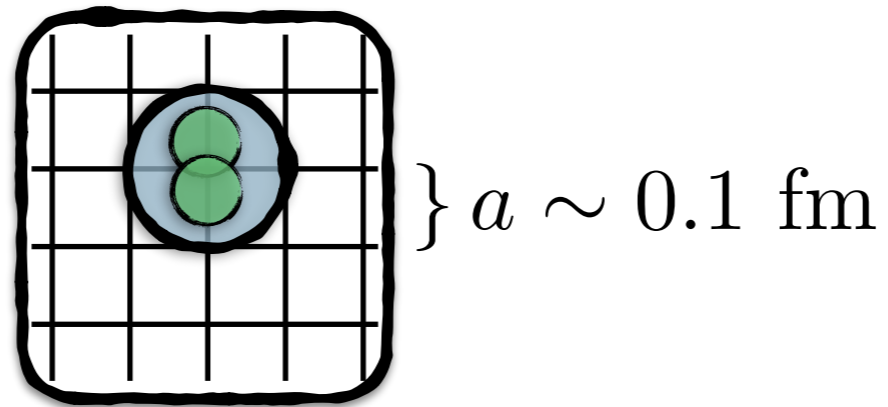
Why lattice?



- QCD-stable states are generated exactly
- QED / weak sector can be treated perturbatively or non-perturbatively
- Resonance are generated and decay

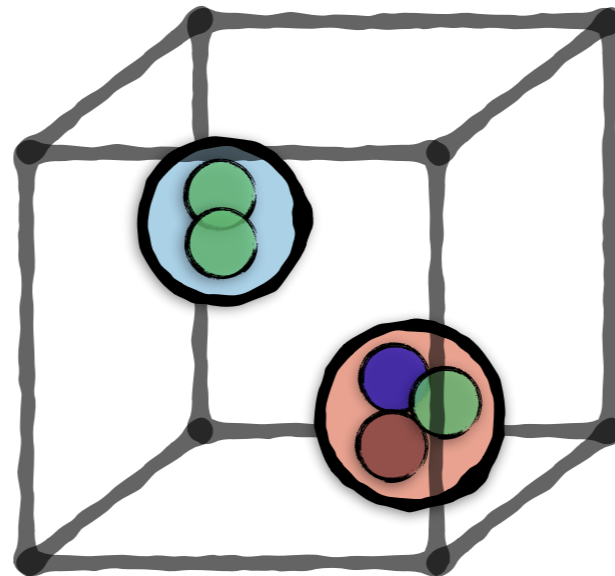
Lattice QCD

• Lattice spacing:



• Wick rotation [Euclidean spacetime]: $t_M \rightarrow it_E$

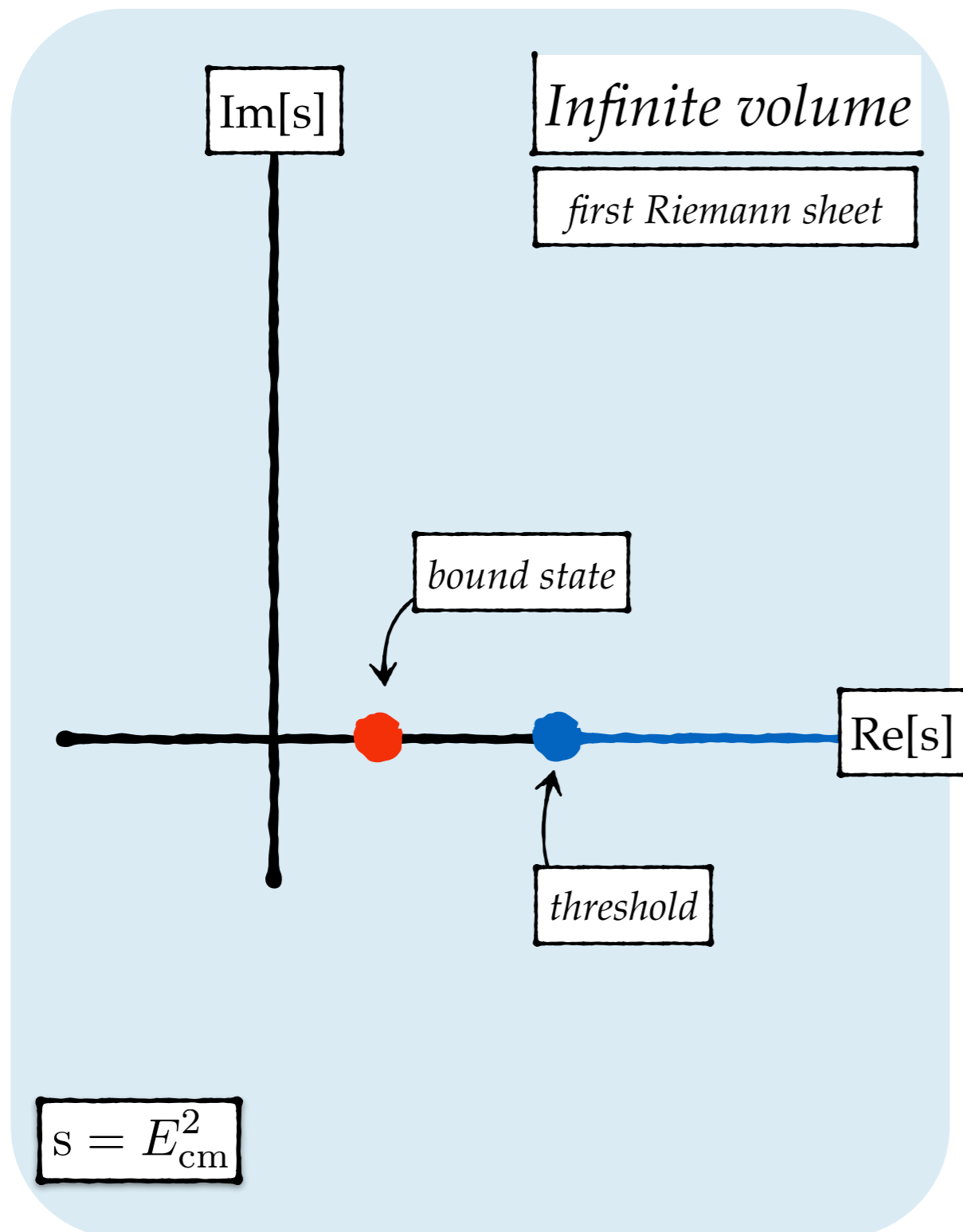
• Finite volume:



• Quark masses: $m_q \rightarrow m_q^{\text{phys.}}$

Have we 'mangled' QCD too much?

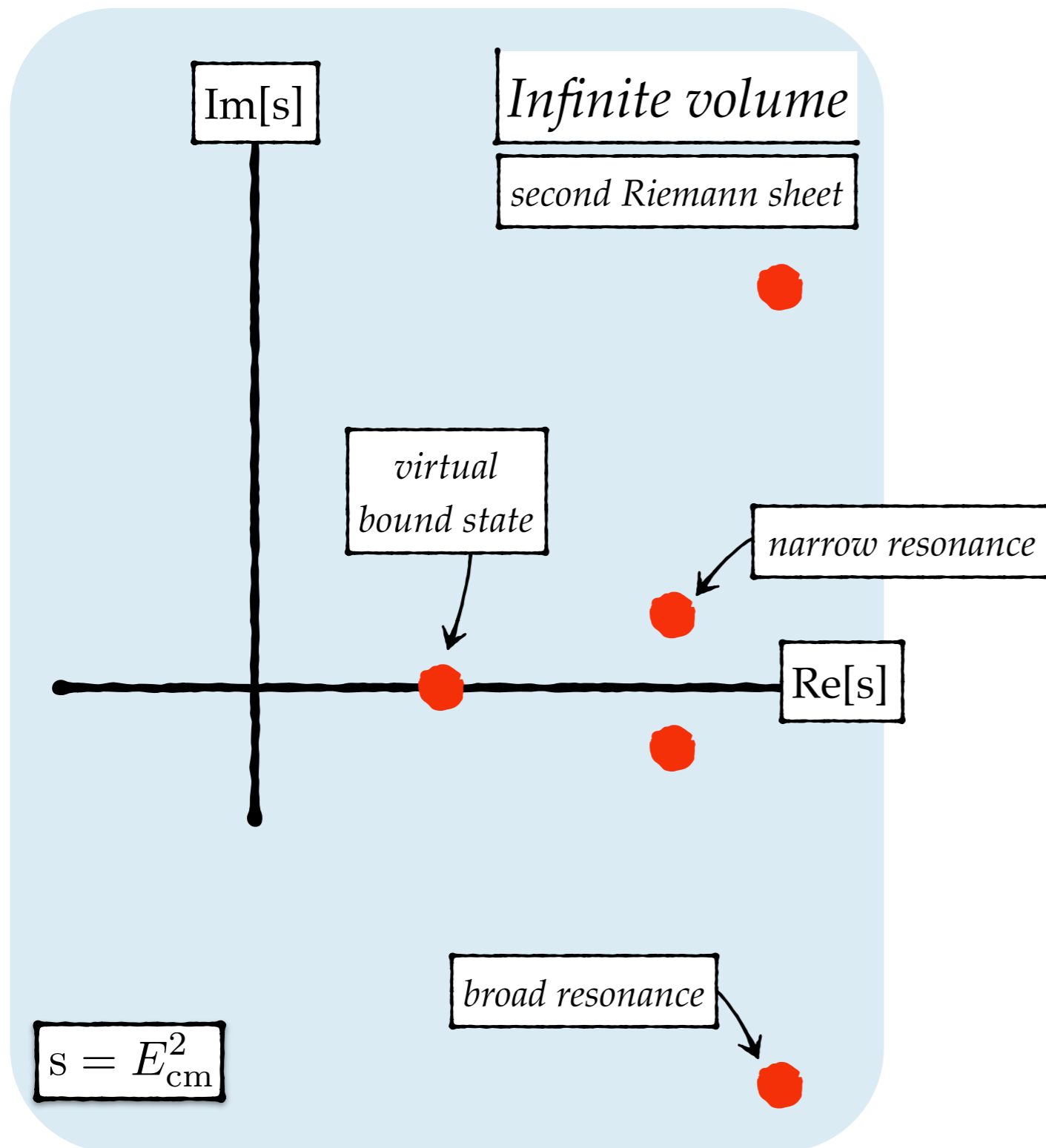
Finite vs. infinite volume spectrum



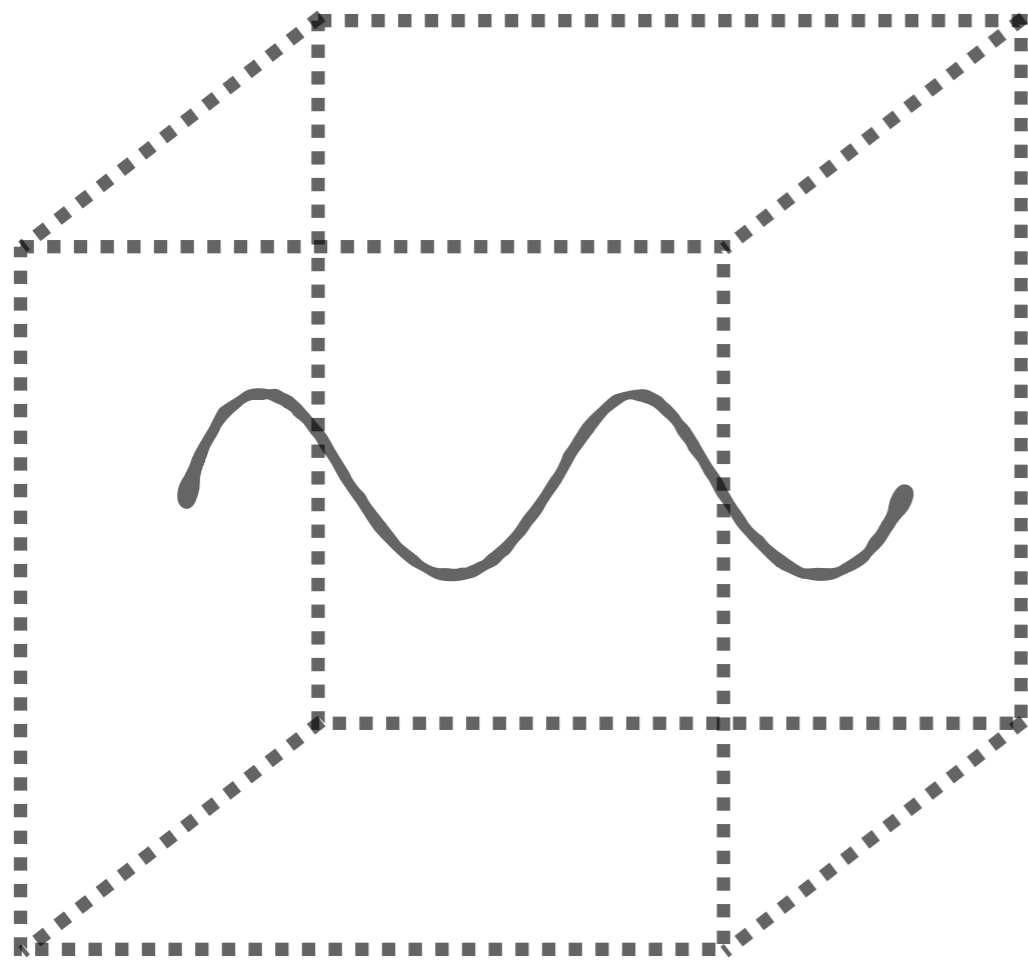
Consider any QCD channel with definite quantum numbers:

- angular momentum - J
- parity - P
- isospin - I
- ...

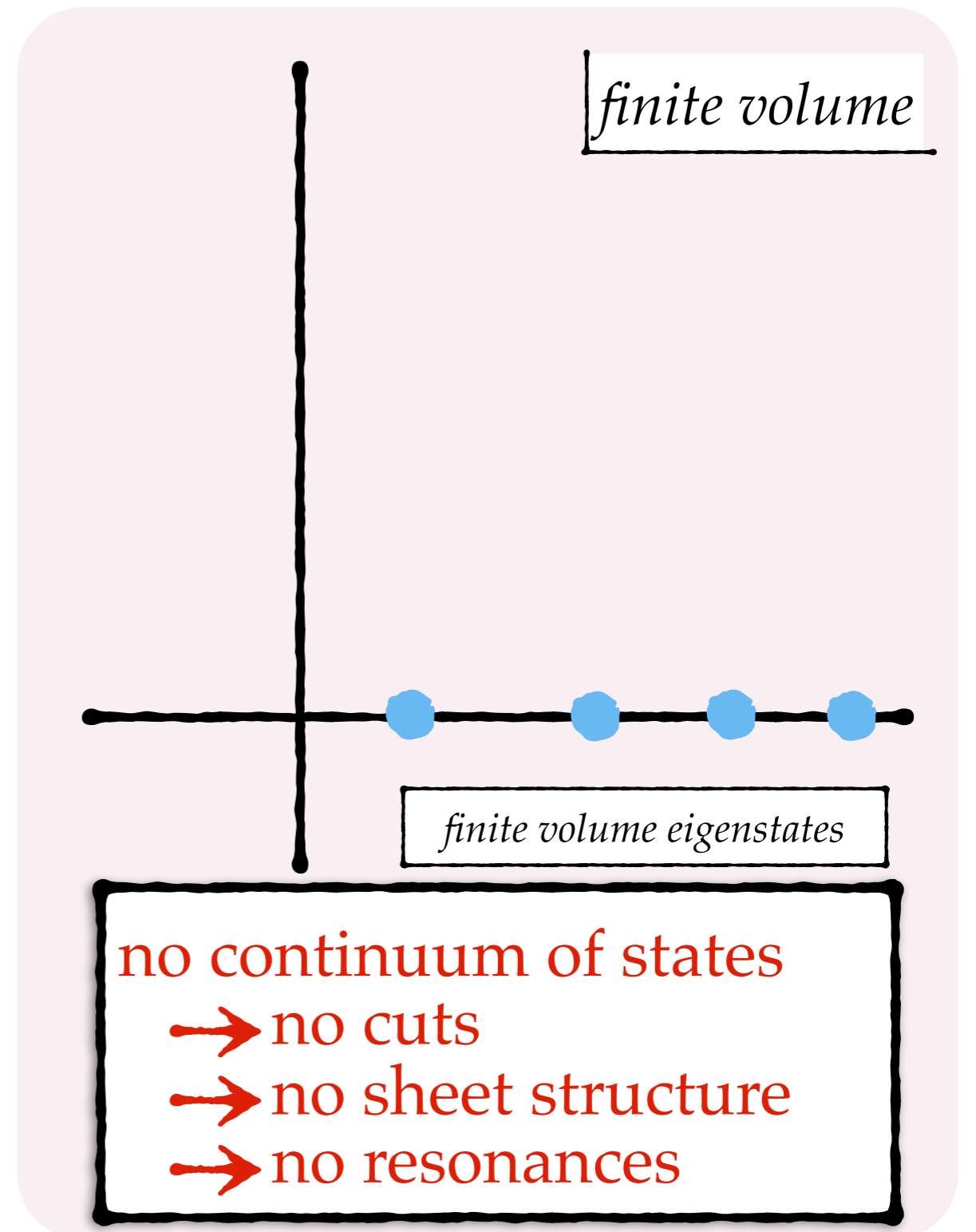
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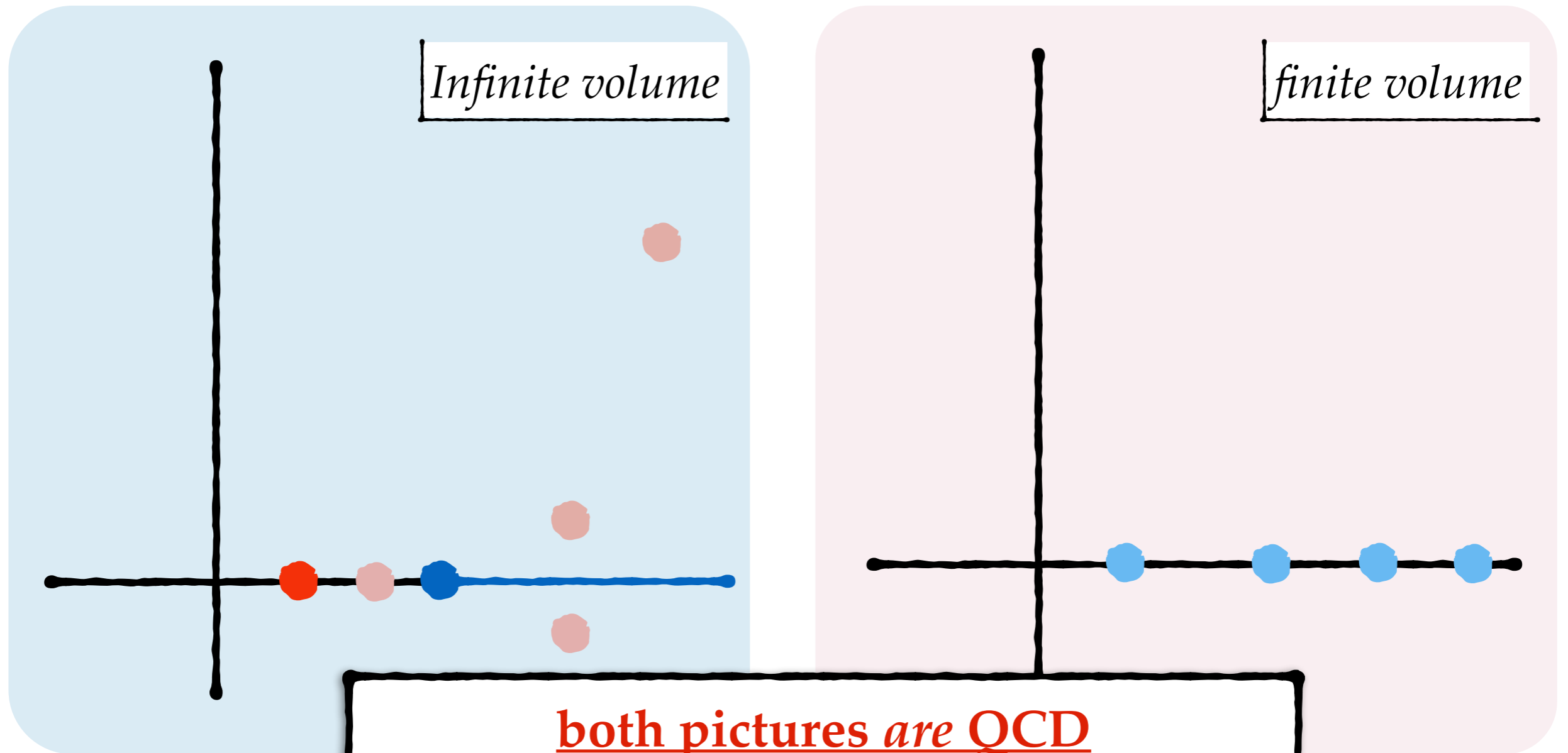
Finite vs. infinite volume spectrum



“only a finite number of modes can exist in a finite volume”



Finite vs. infinite volume spectrum



both pictures are QCD

the connection is perhaps not obvious since we have historically been "confined" to thinking about infinite volume physics

Questions to answer:

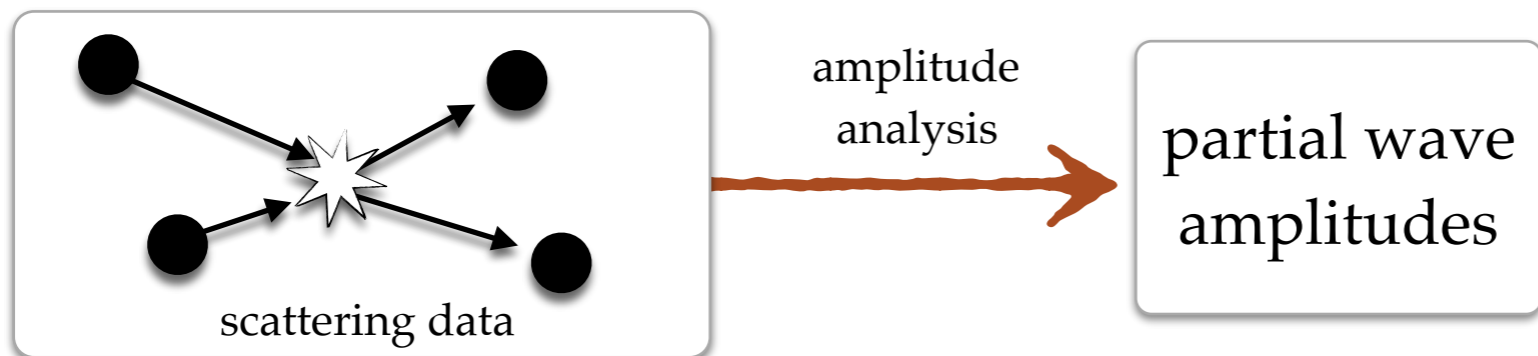
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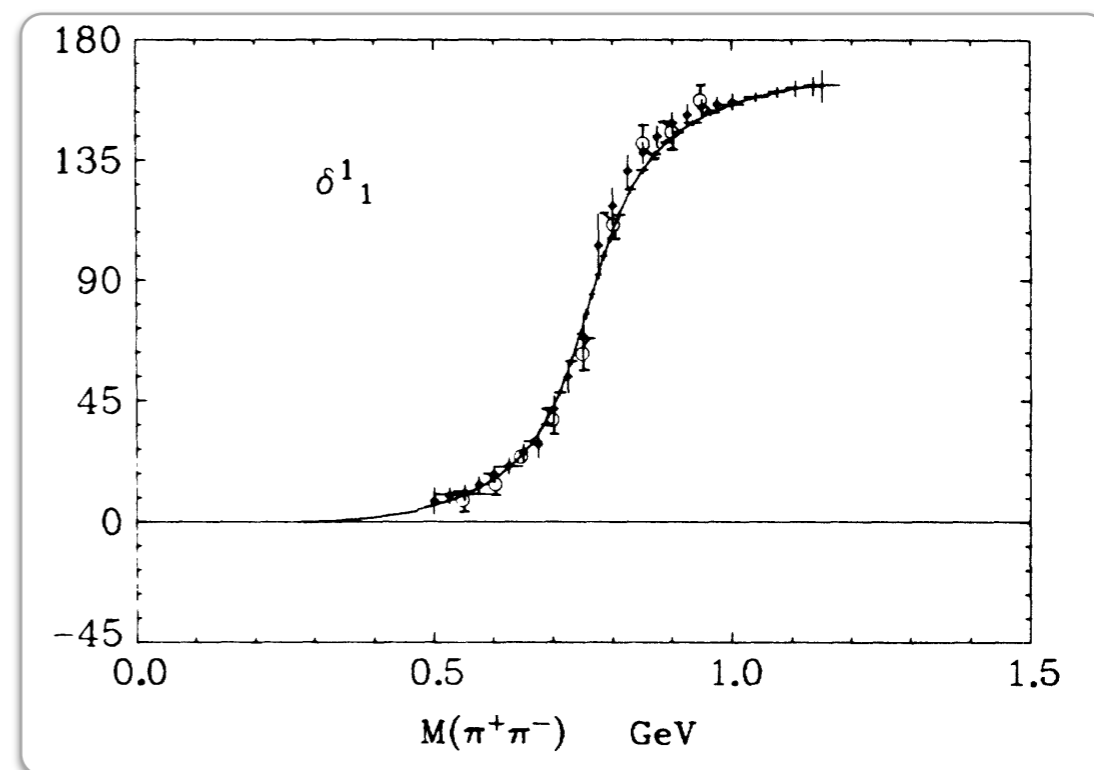
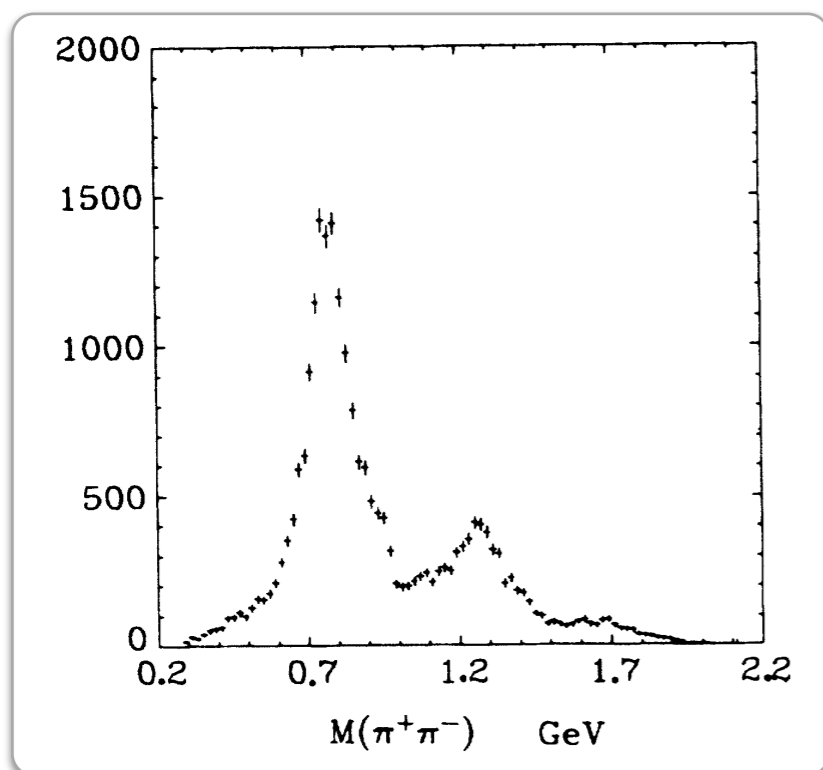
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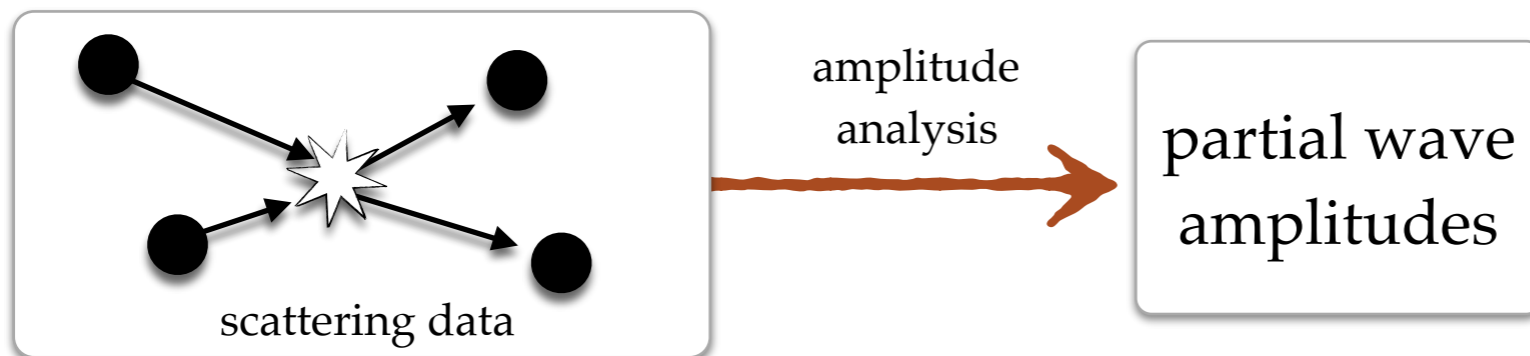
Experiment



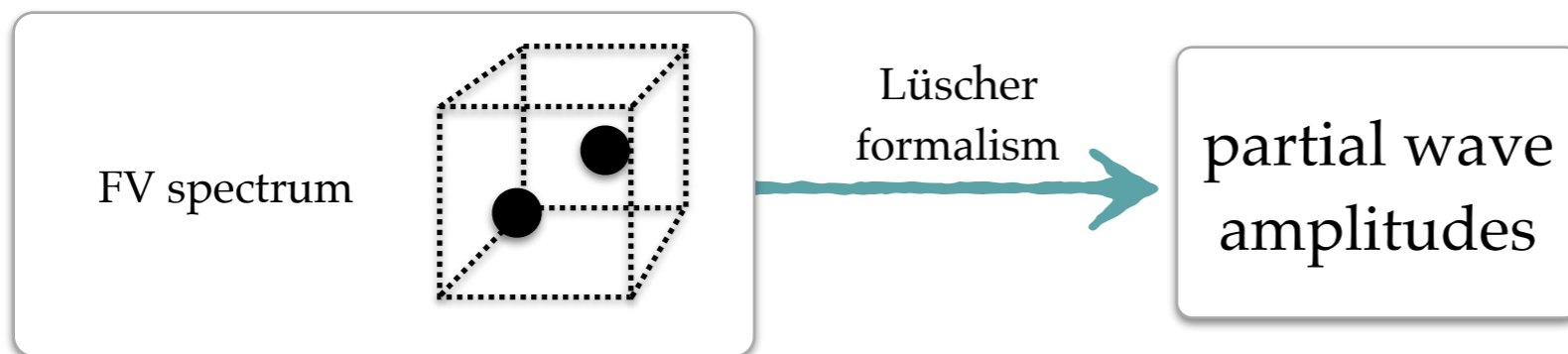
Protopopescu et al. (1972)



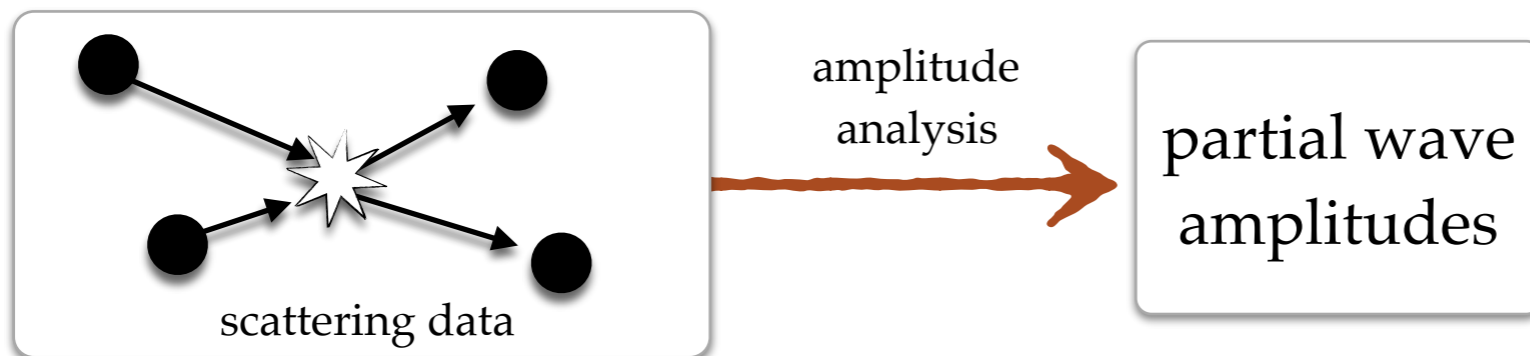
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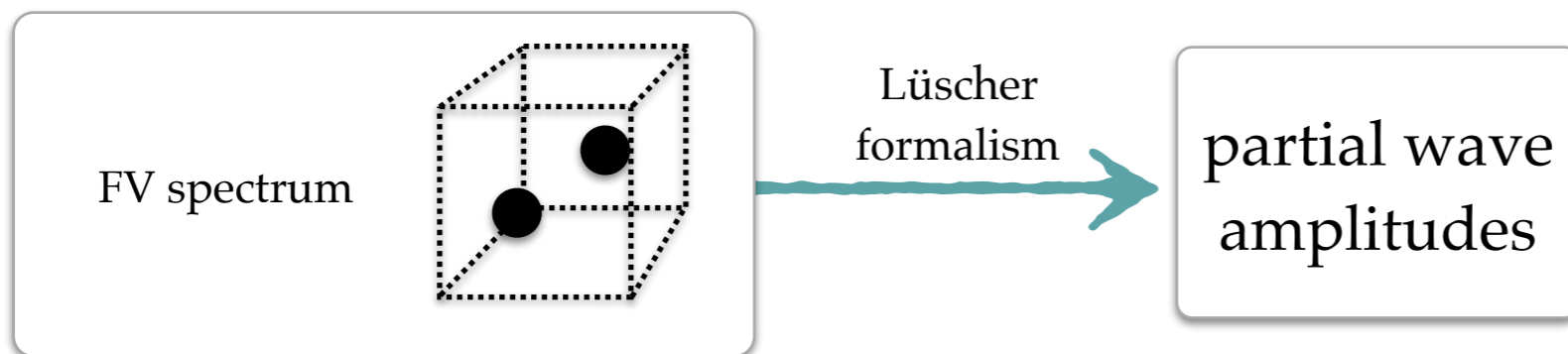
Lattice QCD



Experiment

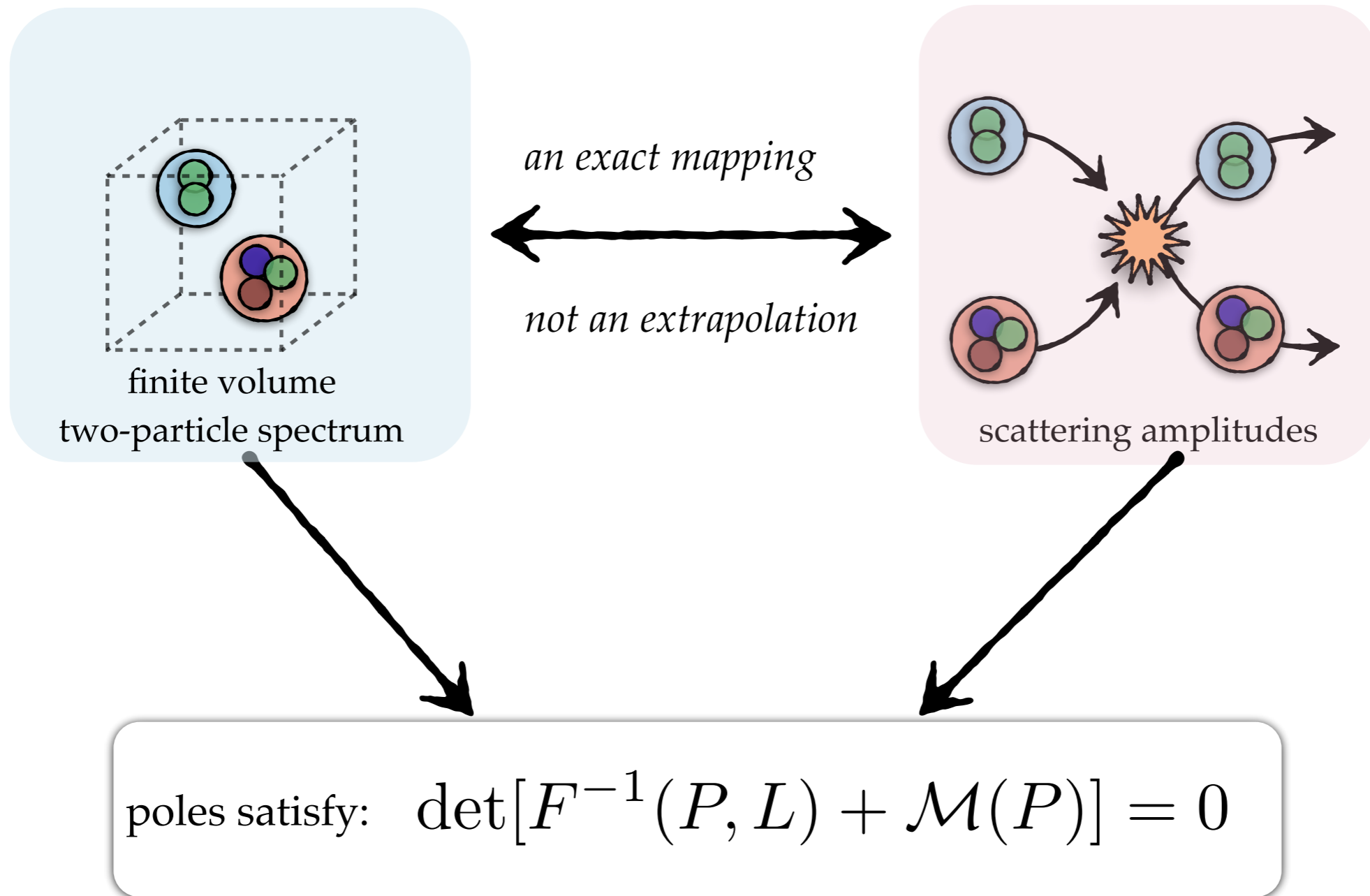


Lattice QCD



determining the spectrum is a science of its own, but lets assume it's done properly

Lüscher formalism

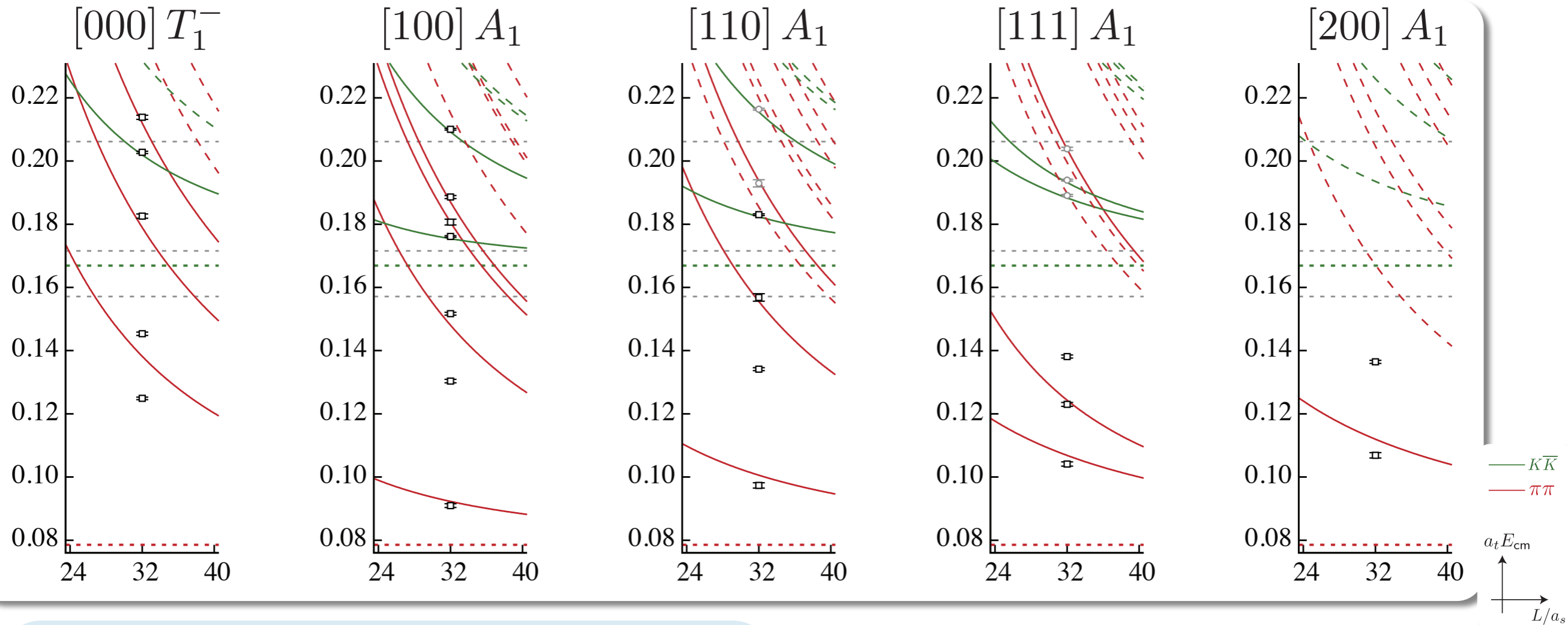


the most general two-particle quantization condition in a finite volume **[RB (2014)]**

$\pi\pi$ scattering

(I=1 channel)

A subset of the spectrum:



Wilson



Thomas



Dudek



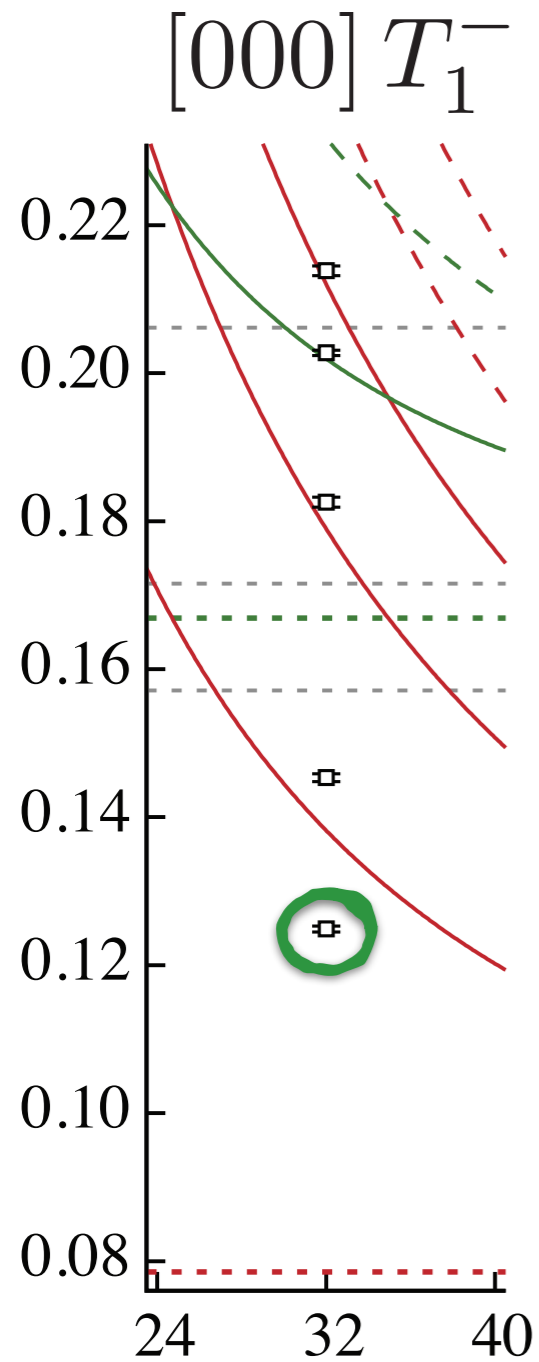
Edwards

**Hadron Spectrum
Collaboration**

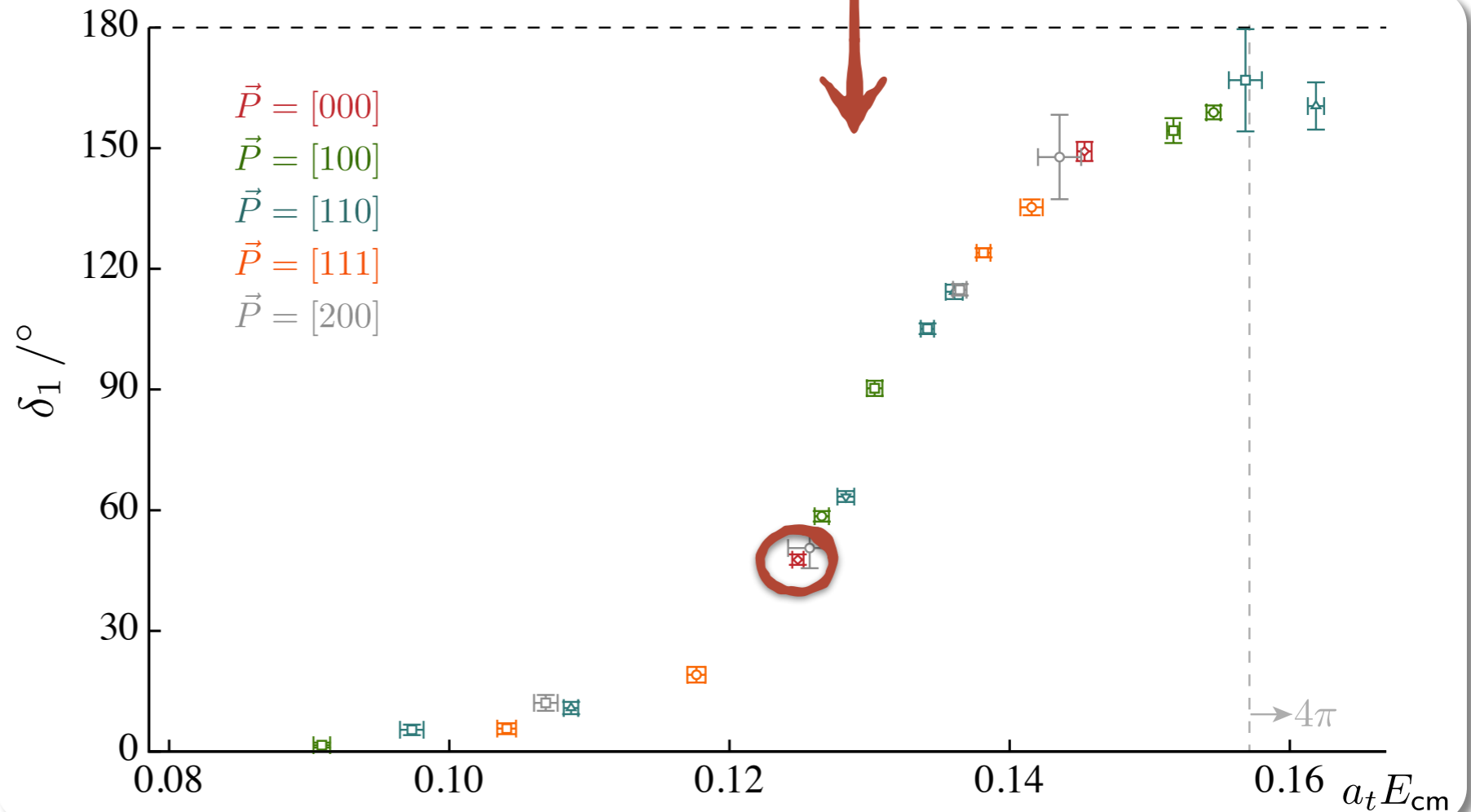
Wilson, RB, Dudek, Edwards & Thomas (2015)

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$$\det[\underbrace{F^{-1}(P, L)}_{\text{green}} + \underbrace{\mathcal{M}(P)}_{\text{red}}] = 0$$

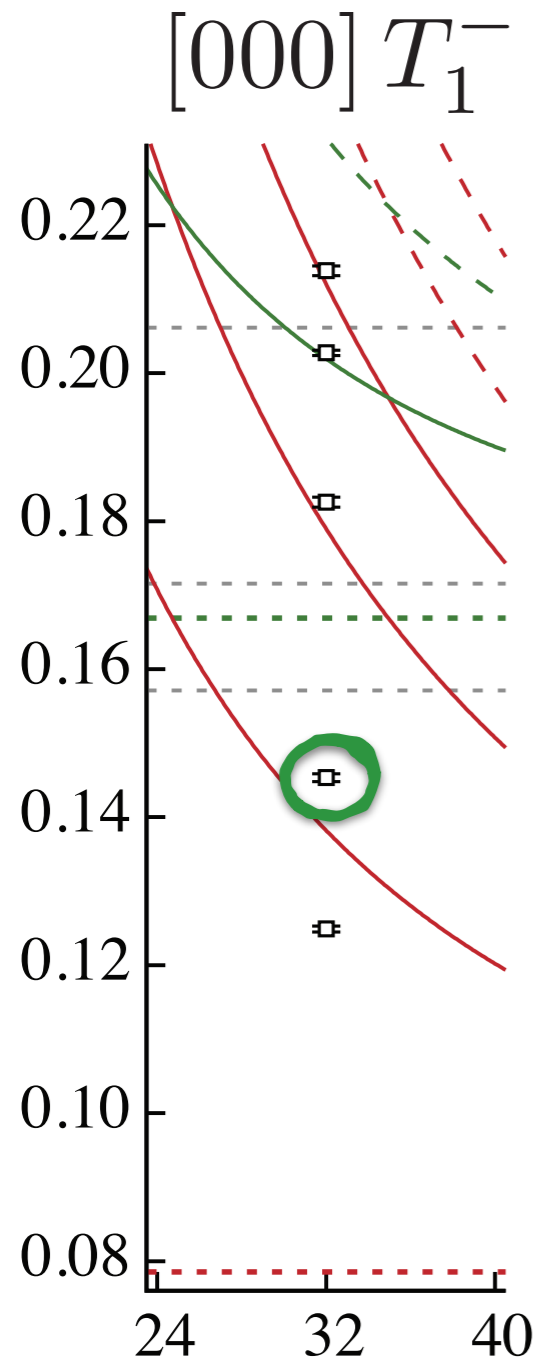


$$\mathcal{M} \propto \frac{1}{\cot \delta_1 - i}$$

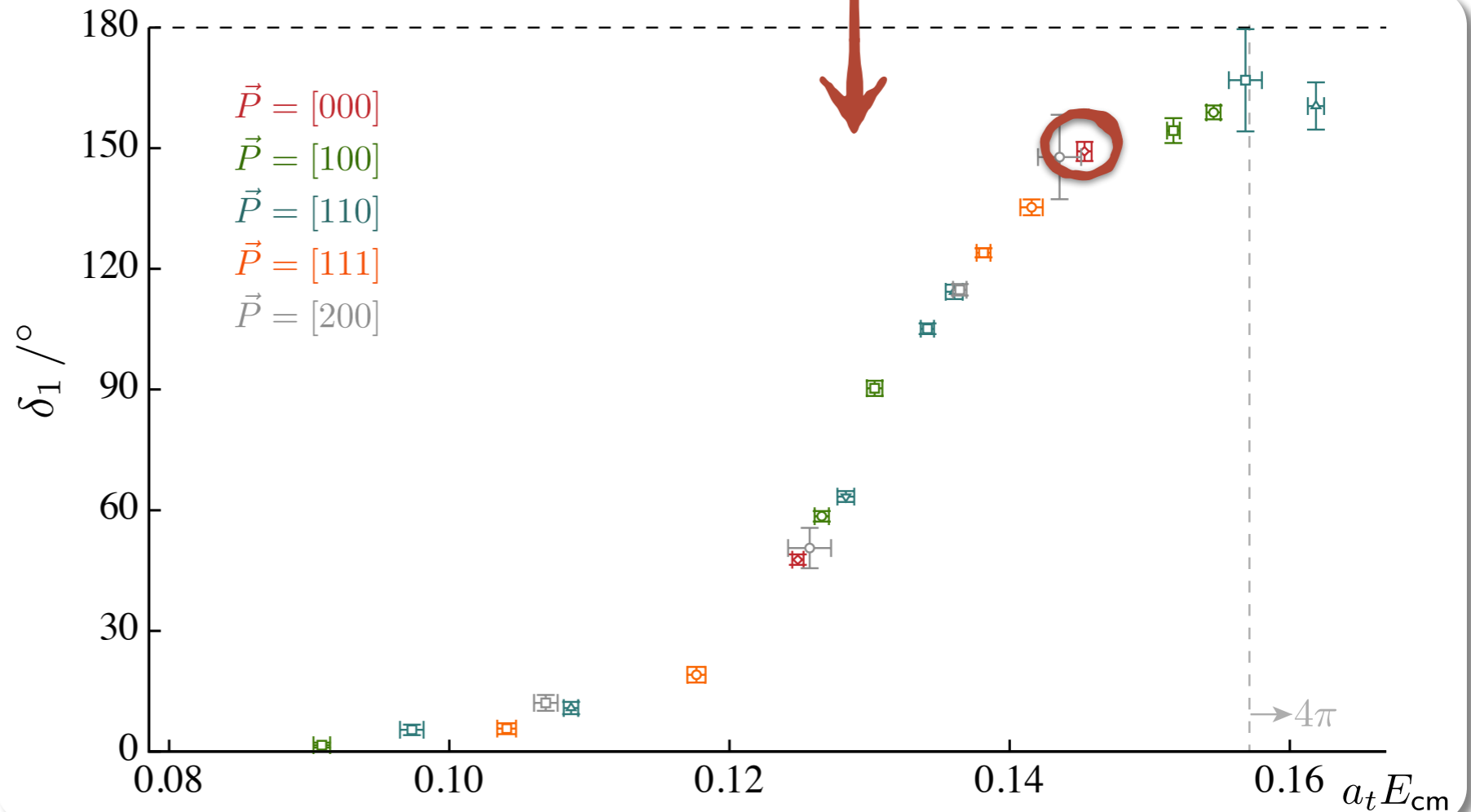
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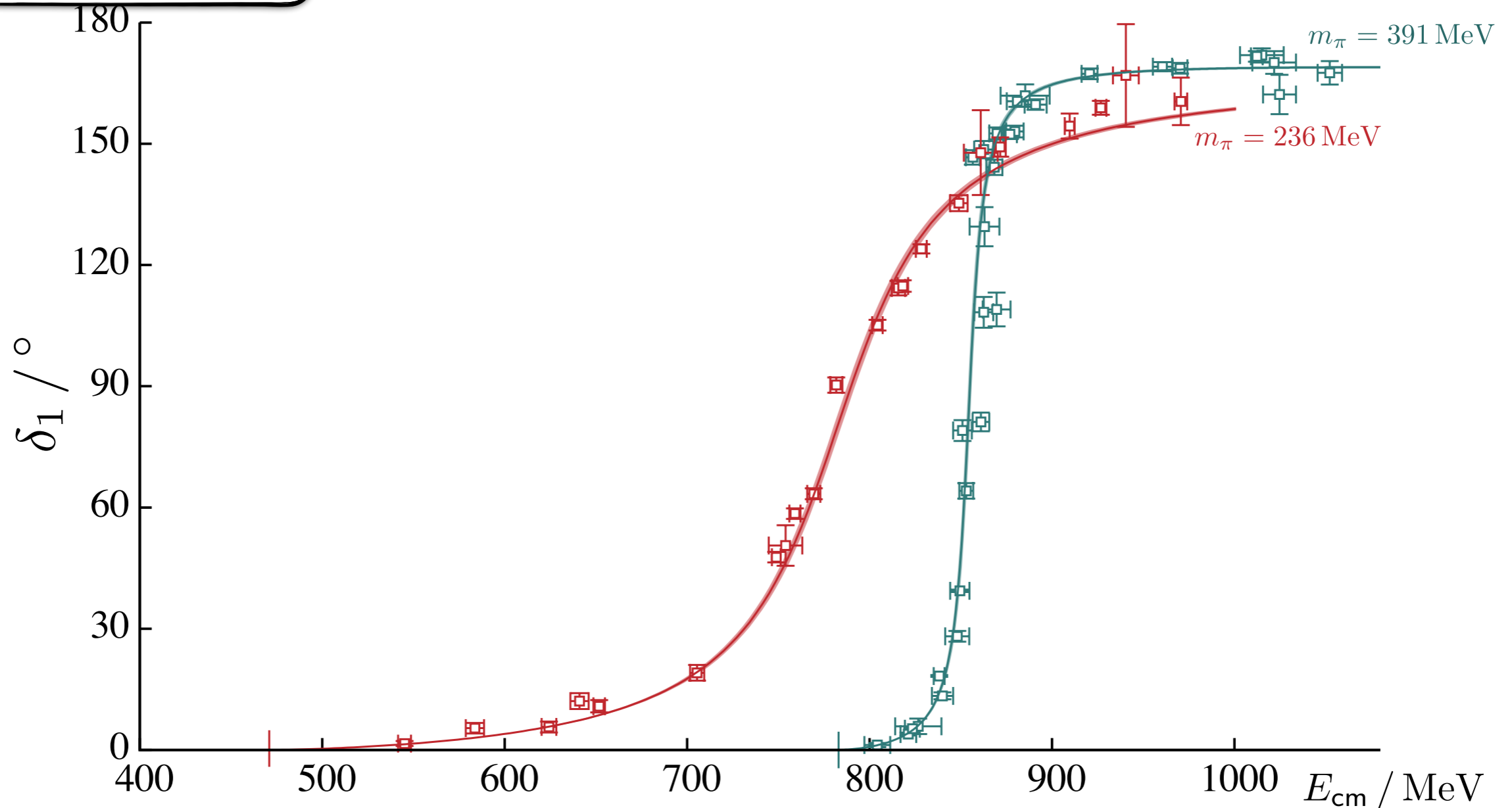


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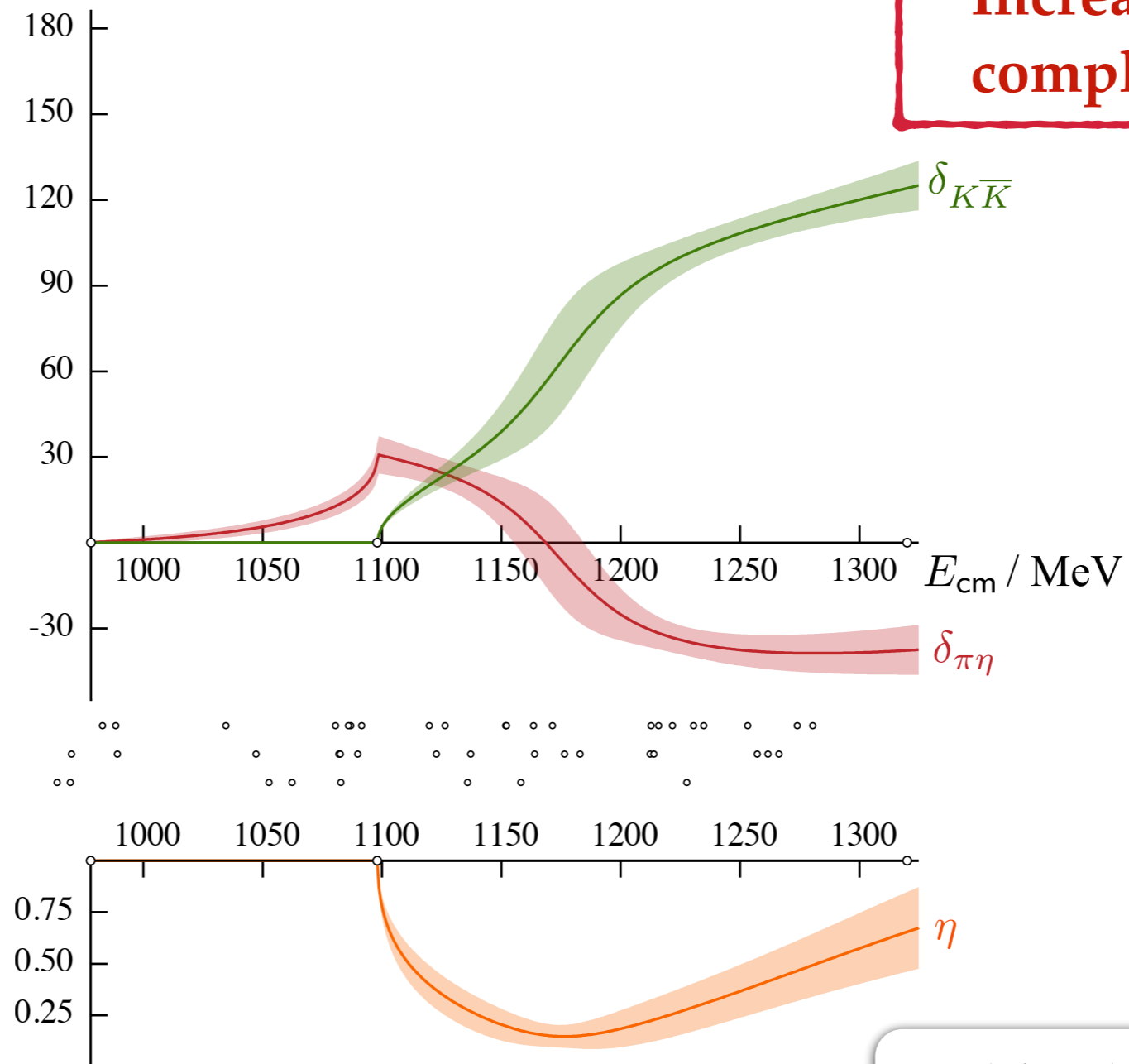
Dudek, Edwards & Thomas (2012)

Wilson, RB, Dudek, Edwards & Thomas (2015)

$\pi\eta-K\bar{K}$ scattering

(I=1 channel)

Increasingly
complex systems

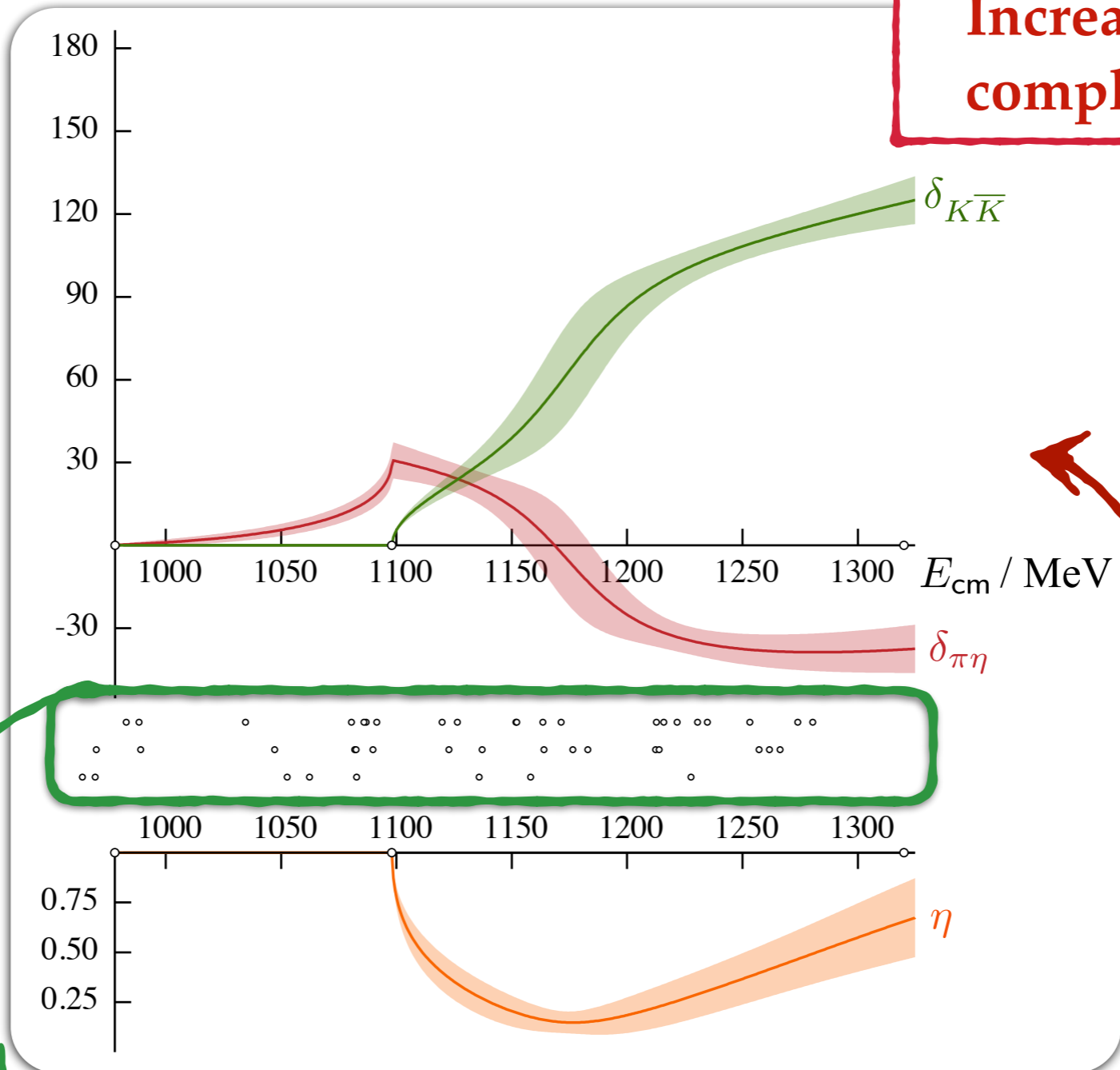


Dudek, Edwards & Wilson (2016)

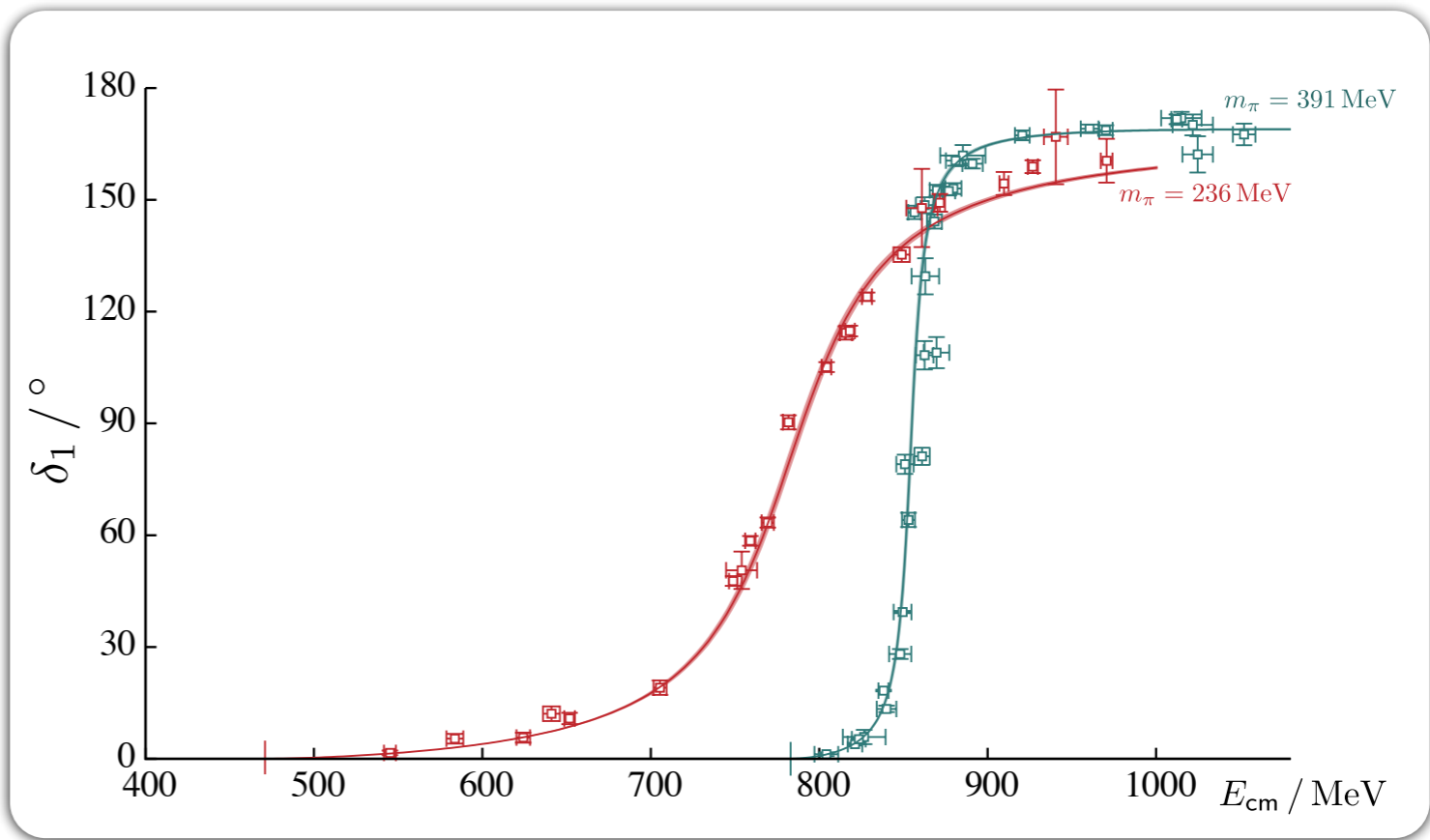
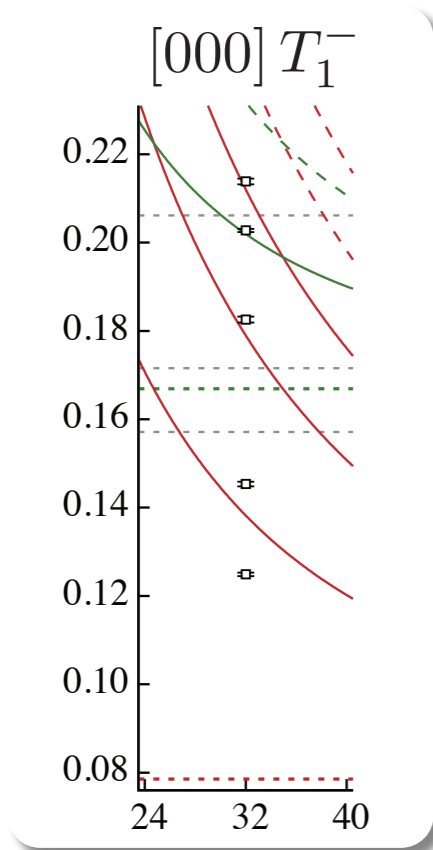
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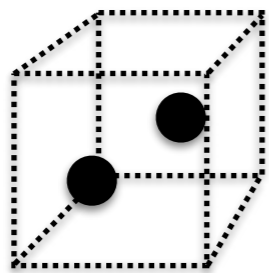


$$\det \left[\underbrace{\begin{pmatrix} F_{\pi\pi} & \\ & F_{K\bar{K}} \end{pmatrix}}^{-1} + \underbrace{\begin{pmatrix} \mathcal{M}_{\pi\pi,\pi\pi} & \mathcal{M}_{\pi\pi,K\bar{K}} \\ \mathcal{M}_{\pi\pi,K\bar{K}} & \mathcal{M}_{K\bar{K},K\bar{K}} \end{pmatrix}} \right] = 0$$



Lattice QCD

FV spectrum



Lüscher
formalism

partial wave
amplitudes

Some comments

- Model independent, **universal**, parametrization-independent, and exact

$$\det[F^{-1}(P, L) + \mathcal{M}(P)] = 0$$

RB (Jan 2014)

- Other on-going efforts:

- Model dependent
- process-dependent
- only suitable for few partial waves
- parametrization dependent
- low-energy approximations
- ignore spin
- ignore three-body

⋮

you get the idea

From Extraction of Nucleon Resonances to LQCD

T.-S. H. Lee,¹ Jia-Jun Wu,² and Hiroyuki Kamano³

¹Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA

²Special Research Center for the Subatomic Structure of Matter (CSSM),
School of Chemistry and Physics, University of Adelaide Adelaide 5005, Australia

³Research Center for Nuclear Physics,

Osaka University, Ibaraki, Osaka 567-0047, Japan

Feb 2016

Finite-volume Hamiltonian method for coupled channel interactions in lattice QCD

Jia-Jun Wu,¹ T.-S. H. Lee,¹ A. W. Thomas,^{2,3} and R. D. Young^{2,3}

¹Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA

²Special Research Center for the Subatomic Structure of Matter (CSSM),
School of Chemistry and Physics, University of Adelaide Adelaide 5005, Australia

³ARC Center of Excellence for Particle Physics at Terascale,
School of Chemistry and Physics, University of Adelaide Adelaide 5005,

Nov 2014

Finite-volume matrix Hamiltonian model for a $\Delta \rightarrow N\pi$ system

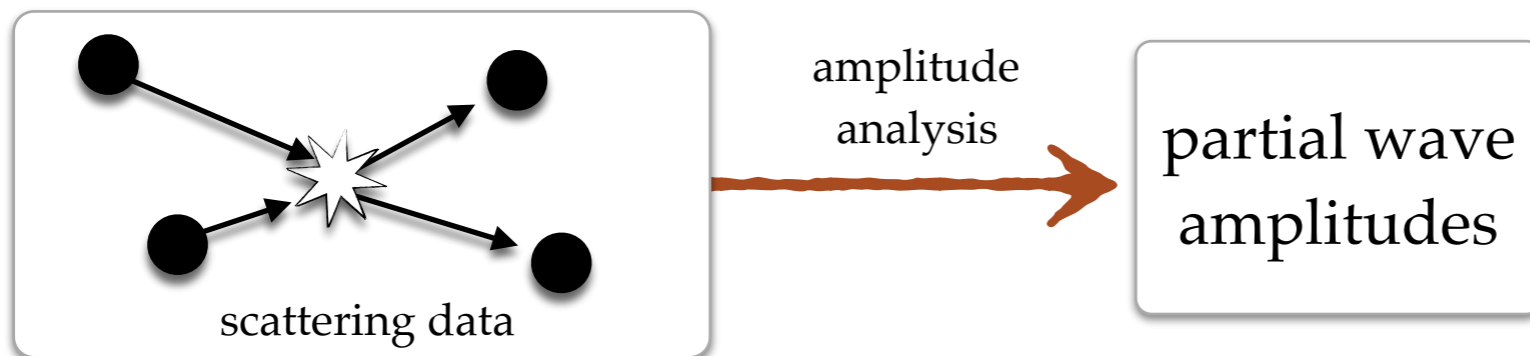
J. M. M. Hall,¹ A. C.-P. Hsu,¹ D. B. Leinweber,¹ A. W. Thomas,^{1,2} and R. D. Young^{1,2}

¹Special Research Centre for the Subatomic Structure of Matter (CSSM),
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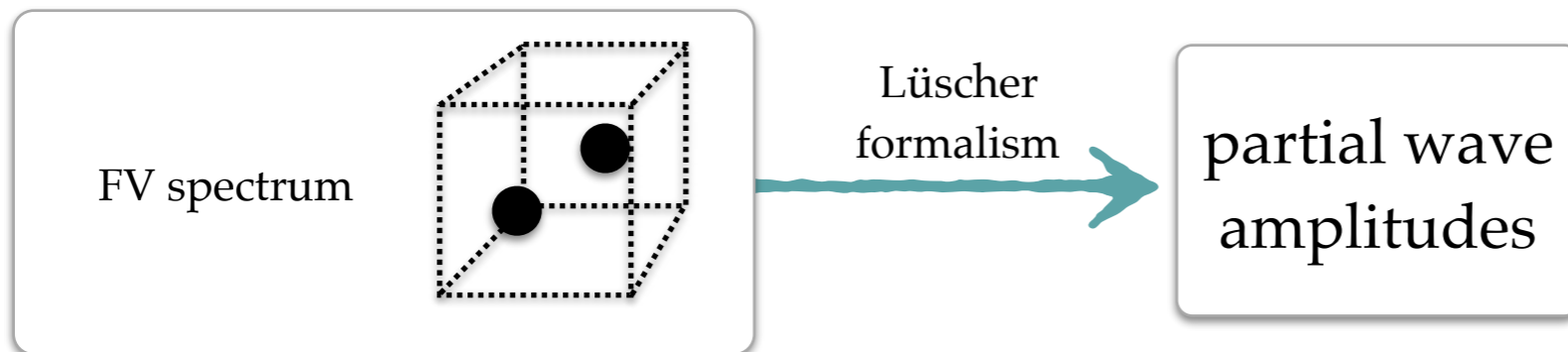
²ARC Centre of Excellence for Particle Physics at the Terascale,
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May 2013

Experiment



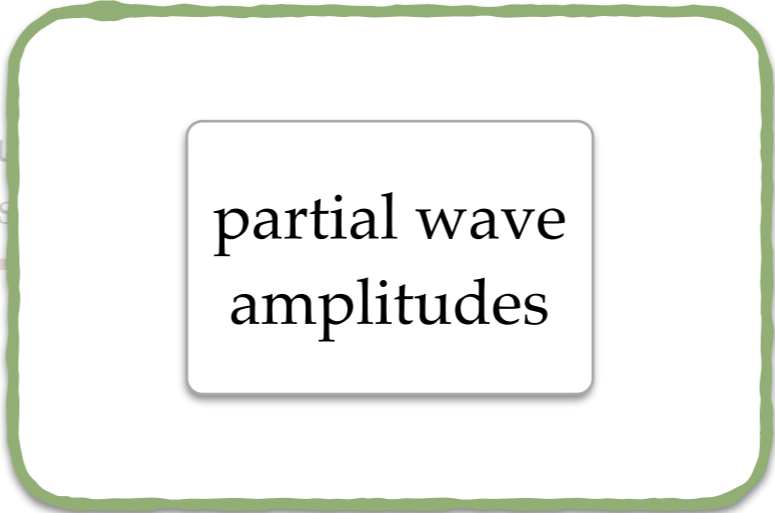
Lattice QCD



Experiment



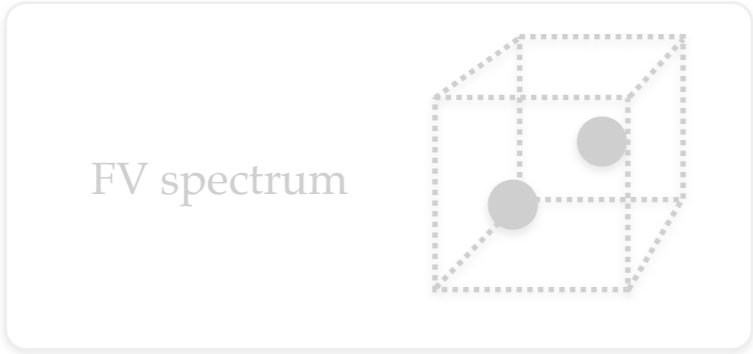
amplitude
analysis



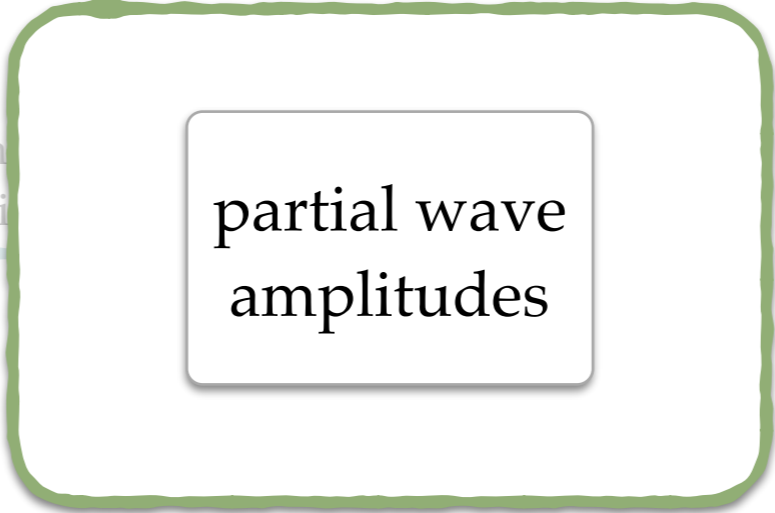
these can then be compared



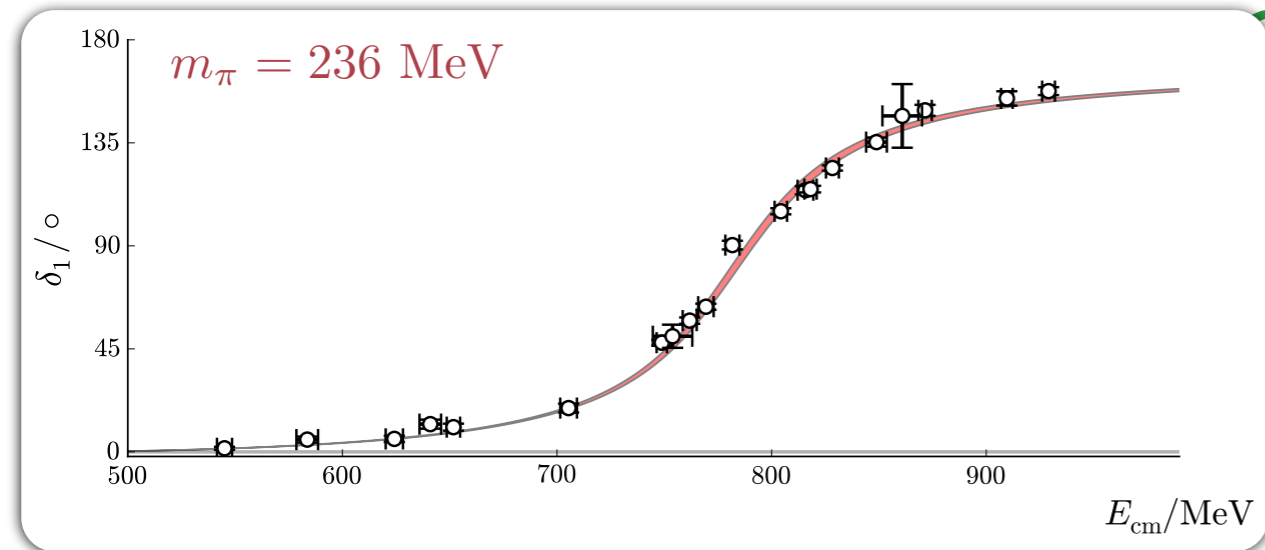
Lattice QCD



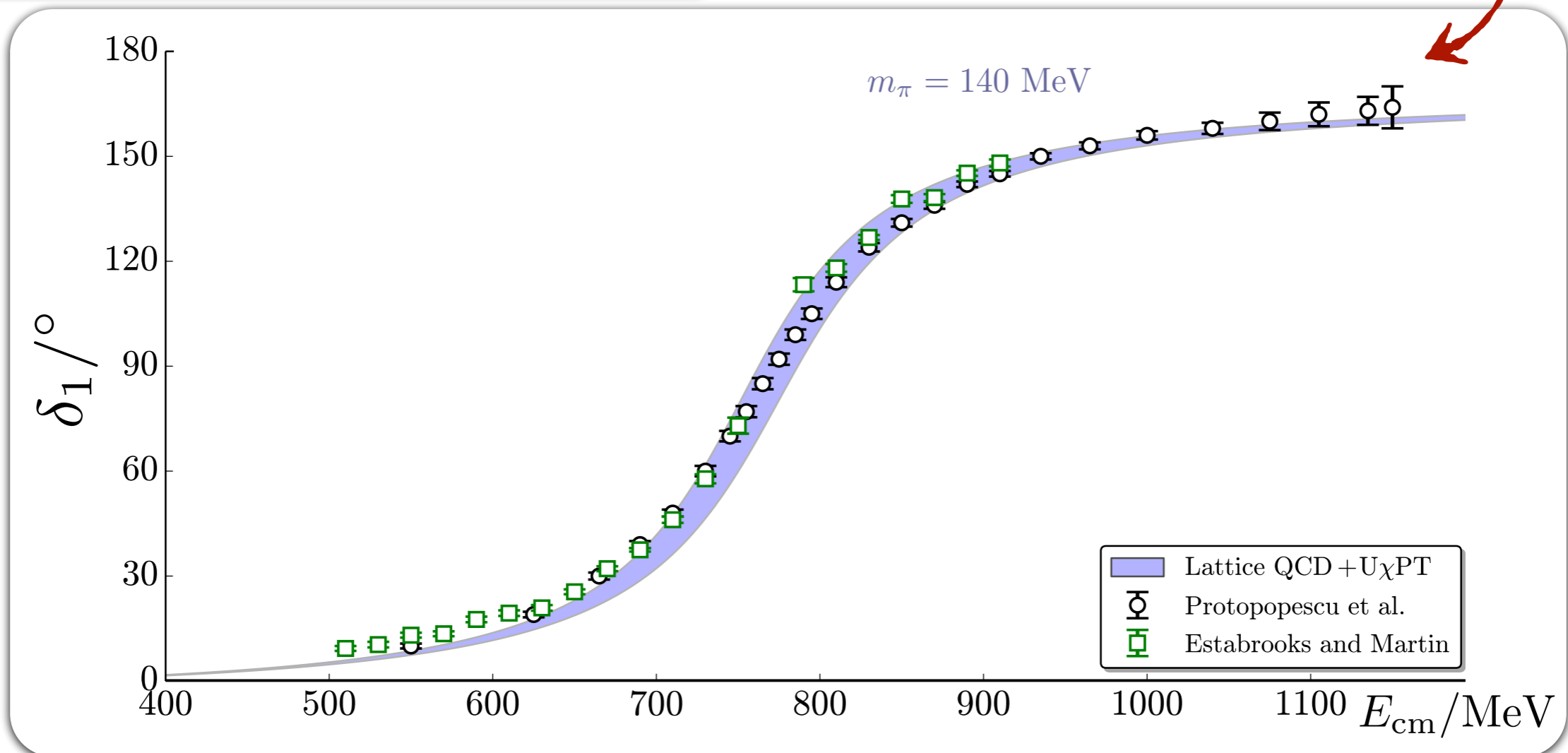
Lüscher
formalism



Comparing with experiment

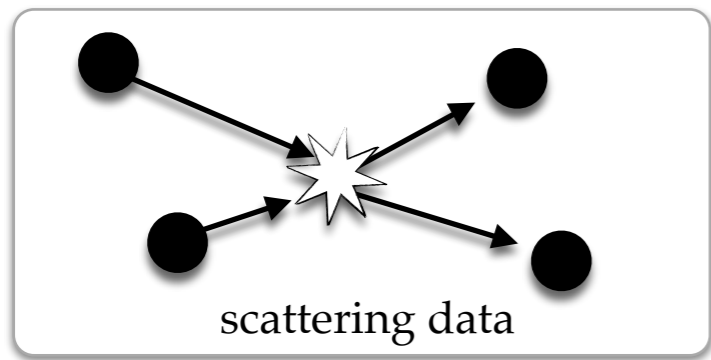


$$\det \left[\underline{F^{-1}} + \mathcal{M}(m_\pi, \{\alpha\}) \right] = 0$$



First chiral extrapolation of a resonant amplitude

Bolton, RB & Wilson (2015)



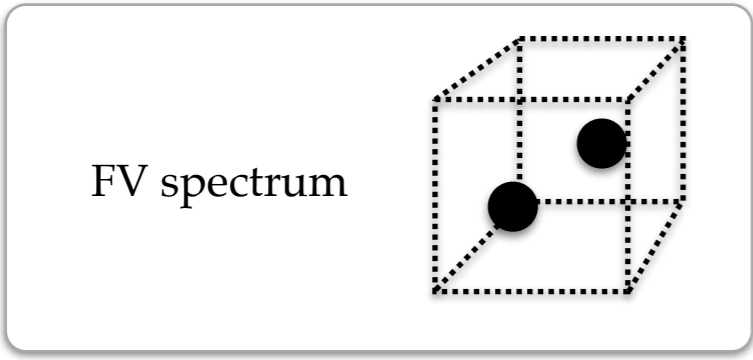
amplitude
analysis

partial wave
amplitudes

analytic
continuation

resonance
poles

Experiment



Lüscher
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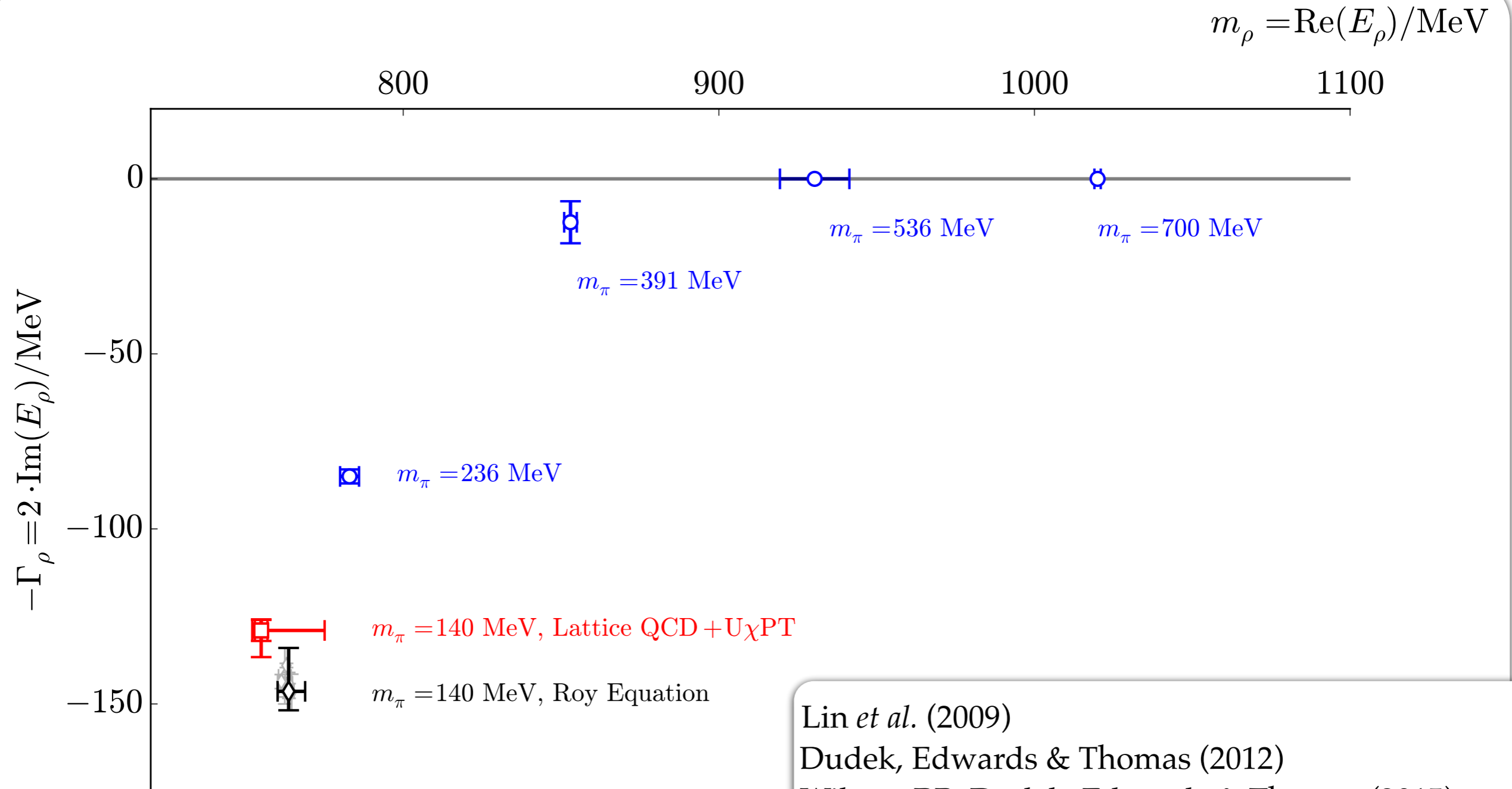
partial wave
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Lattice QCD

Quark-mass dependence of poles



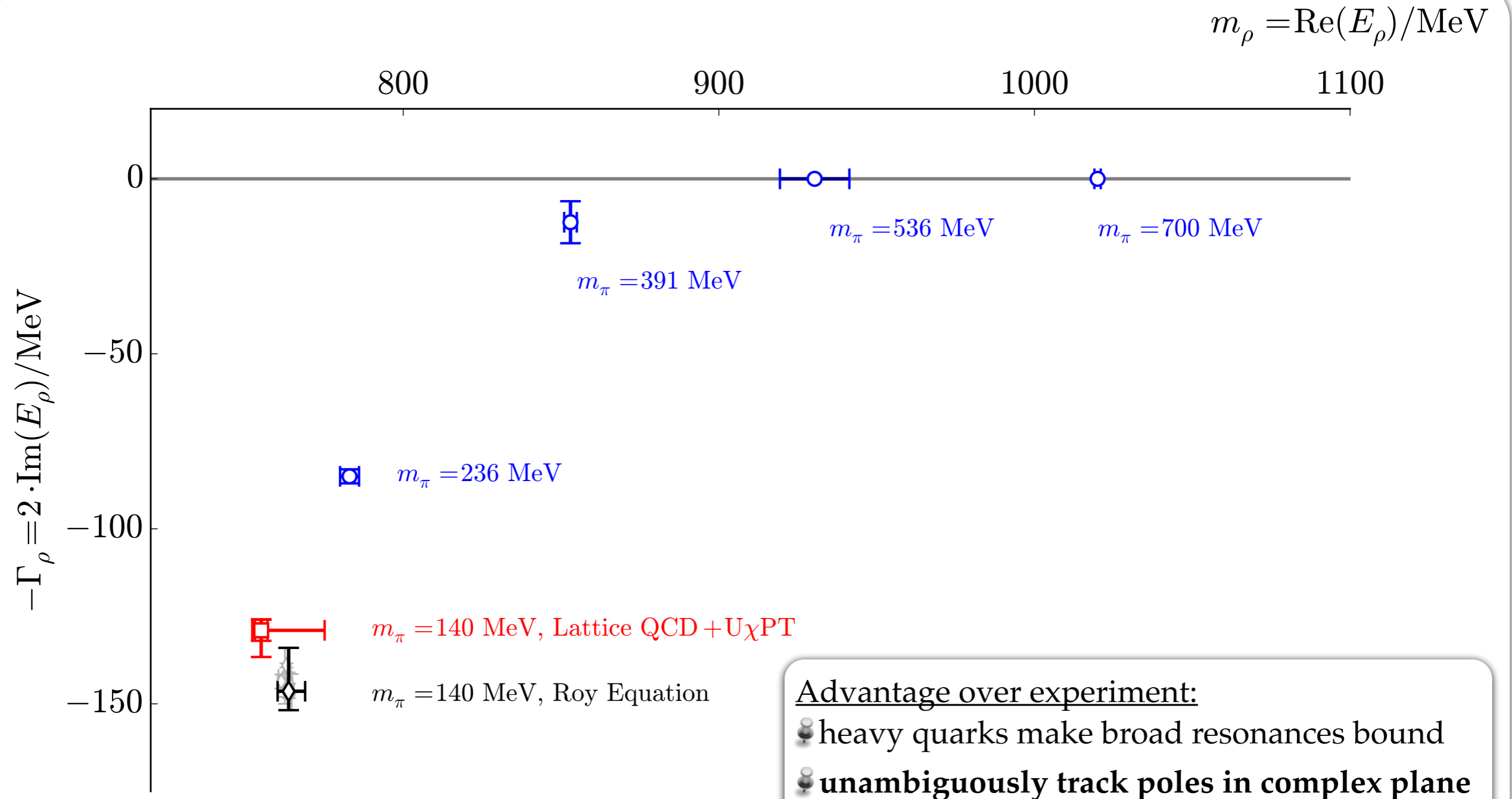
Lin *et al.* (2009)

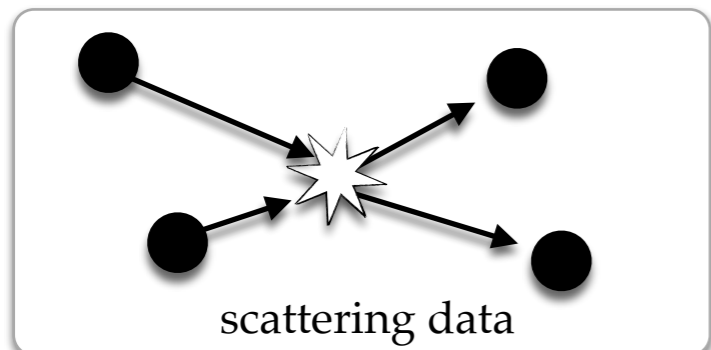
Dudek, Edwards & Thomas (2012)

Wilson, RB, Dudek, Edwards & Thomas (2015)

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Quark-mass dependence of poles





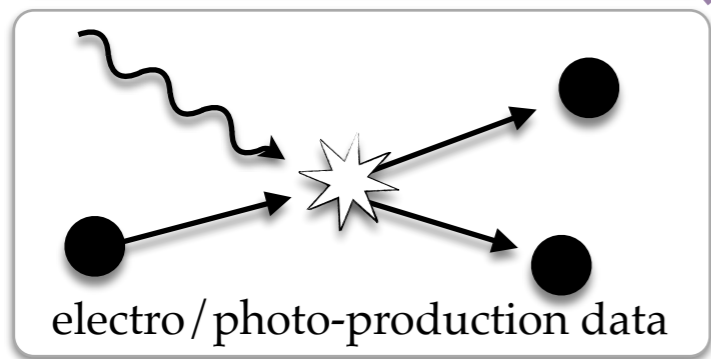
amplitude analysis

partial wave amplitudes

analytic continuation

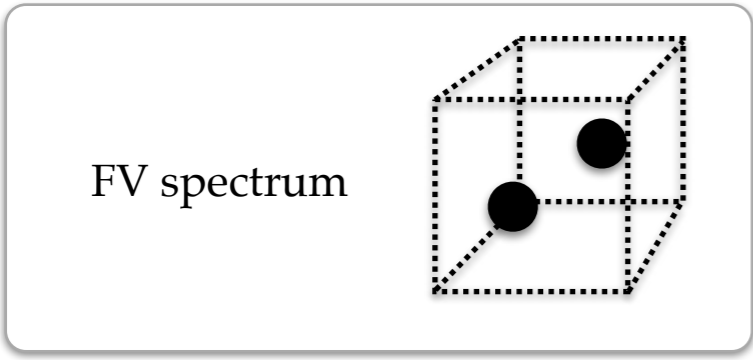
resonance poles

Experiment



amplitude analysis

transition amplitudes



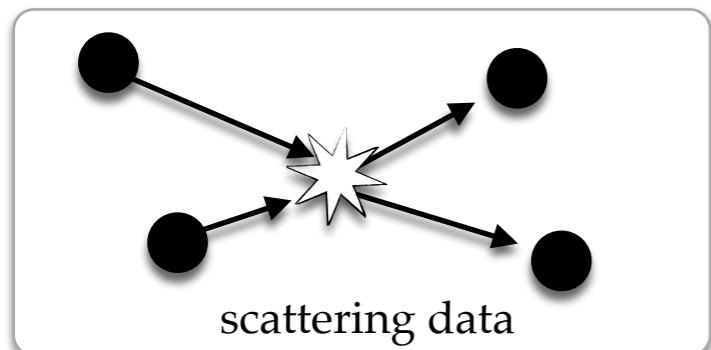
Lüscher formalism

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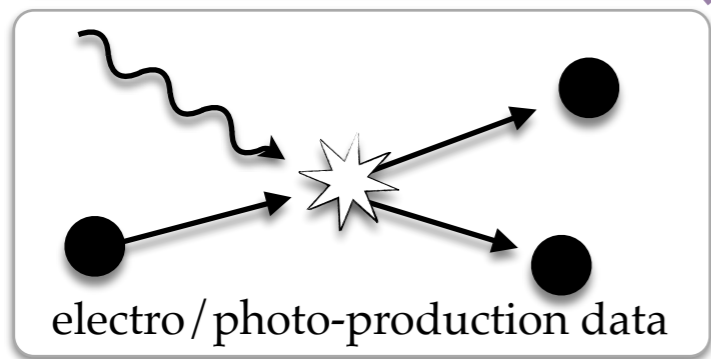


amplitude analysis

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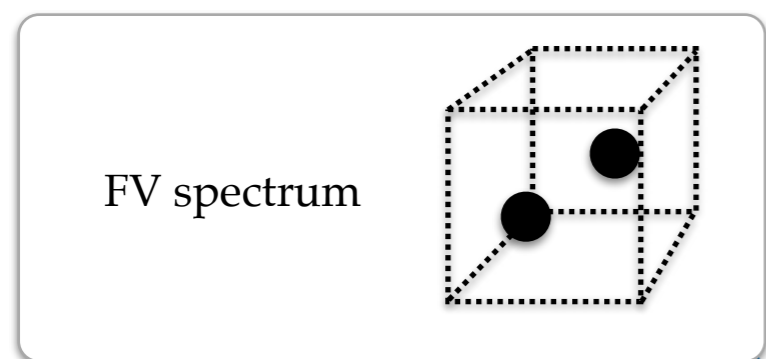
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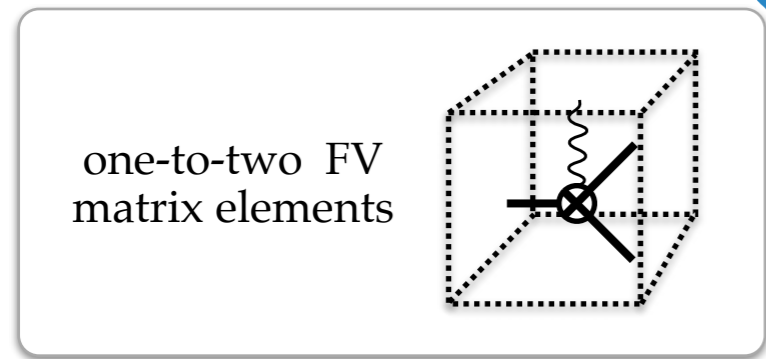
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Lellouch-Lüscher formalism

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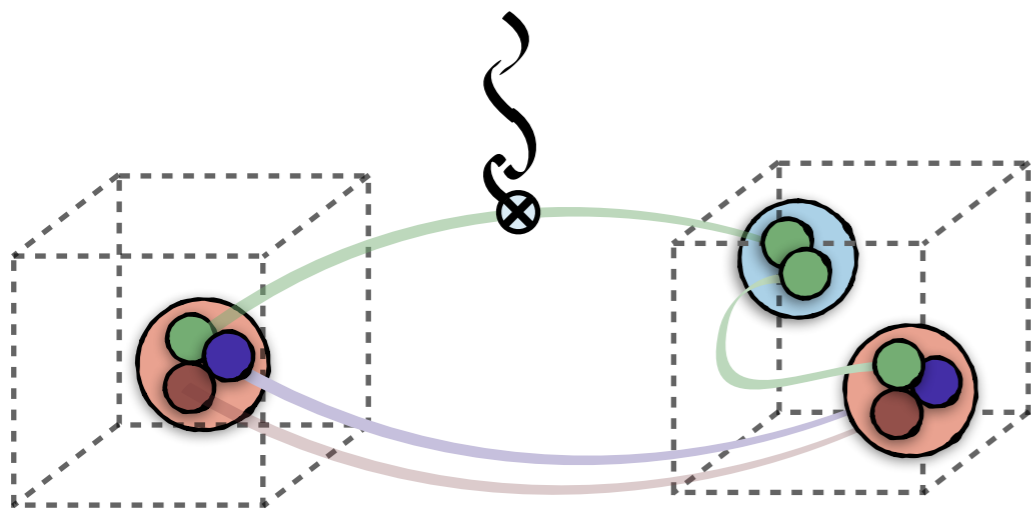
Matrix elements

1) Access matrix elements:

$$C_{\mathbf{2} \rightarrow \mathbf{1} \mathcal{J}}^{3pt.} = \langle \mathcal{O}_1(\delta t) \mathcal{J}(t) \mathcal{O}_2^\dagger(0) \rangle \longrightarrow \langle \mathbf{1} | \mathcal{J} | \mathbf{2} \rangle_L Z_1 Z_2^* e^{-(\delta t - t)E_1} e^{-tE_2} + \dots$$

2) Interpret matrix elements:

$$|\langle \mathbf{2} | \mathcal{J} | \mathbf{1} \rangle_L|^2 = \mathcal{H} \mathcal{R} \mathcal{H}$$



RB, Hansen & Walker-Loud (2014)

RB & Hansen (2015)

RB & Hansen (2015)

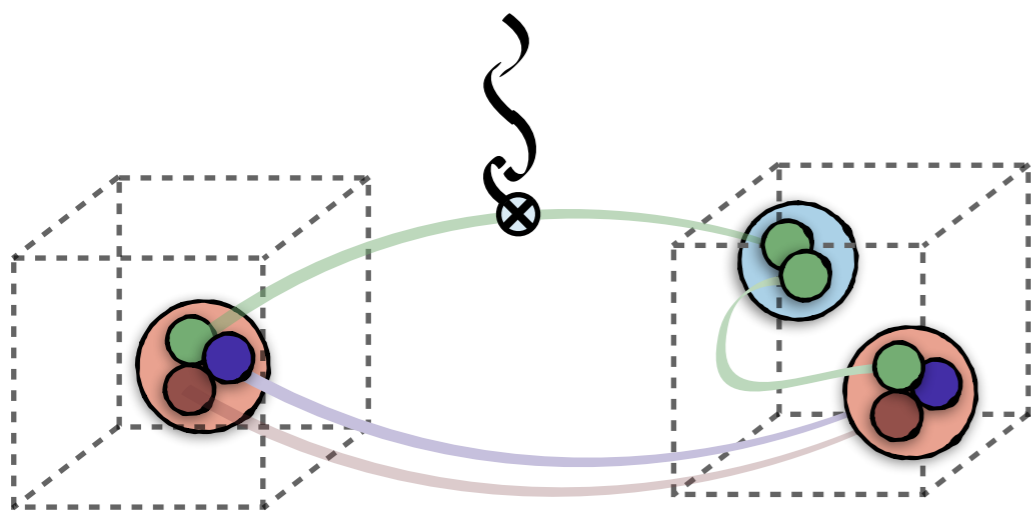
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known finite volume function

$$\mathcal{R} \left(E_2, L, \delta, \frac{\partial \delta}{\partial E_2} \right)$$

RB, Hansen & Walker-Loud (2014)

RB & Hansen (2015)

RB & Hansen (2015)

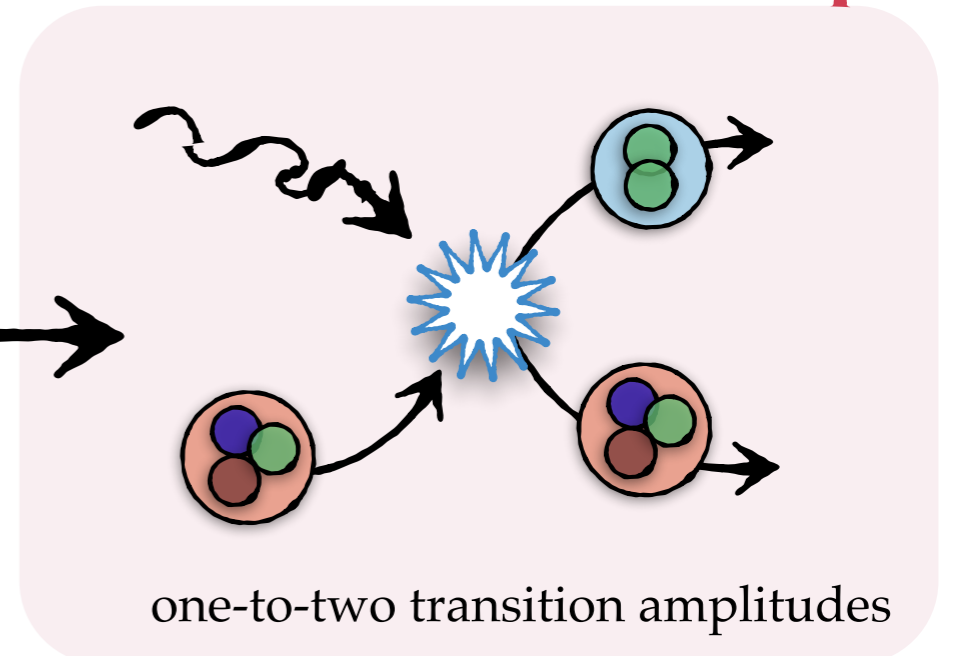
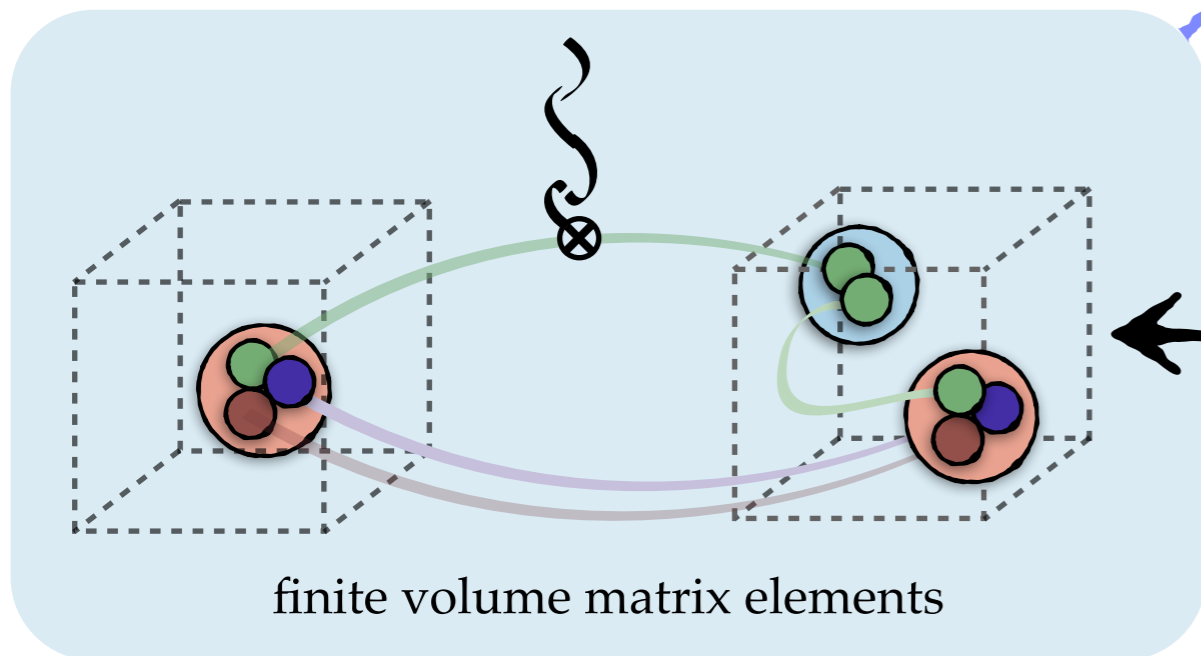
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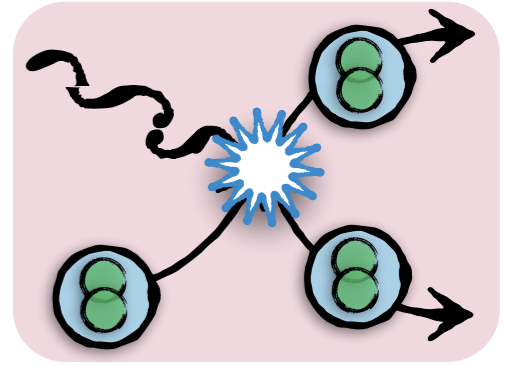
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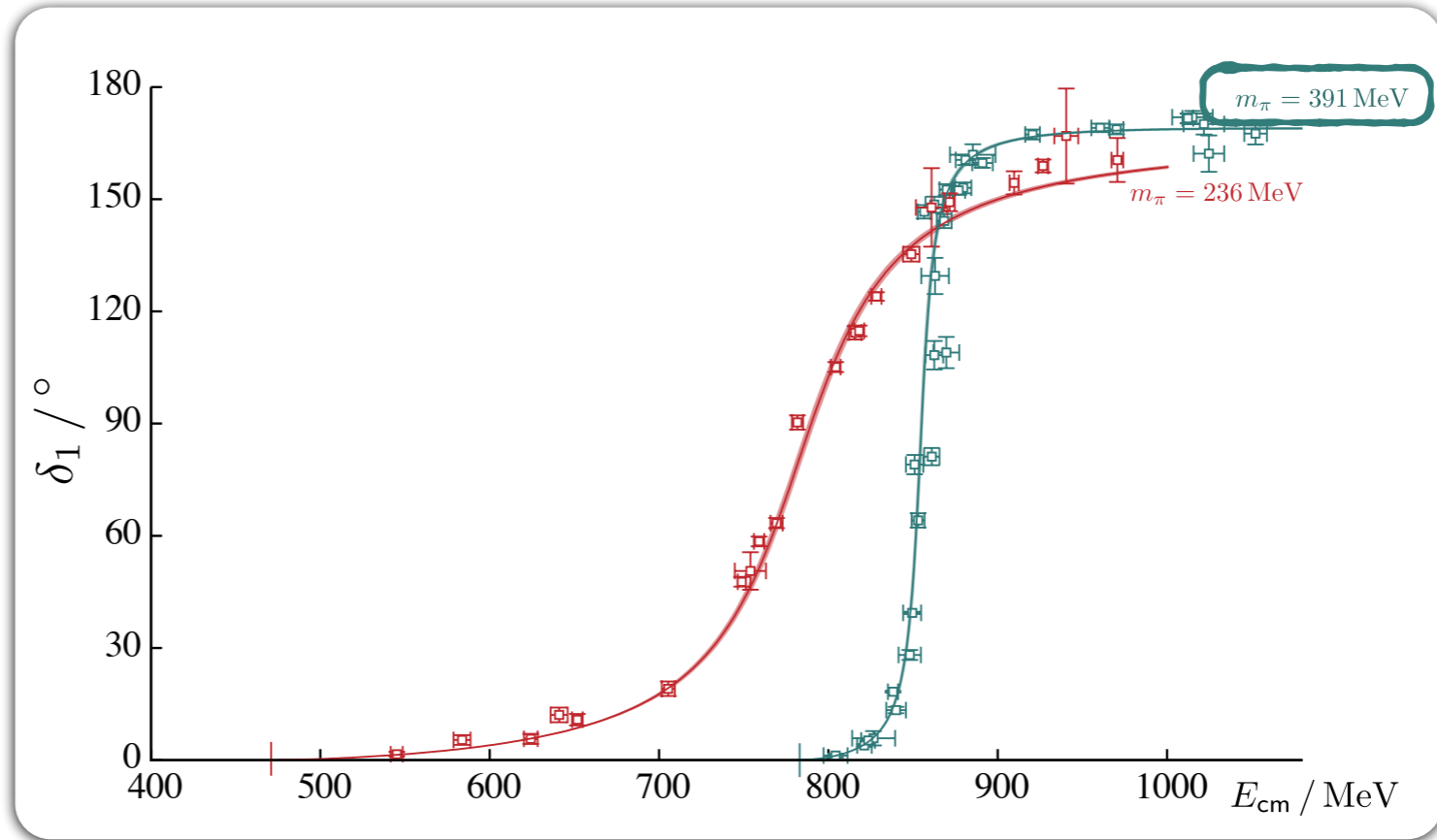


$\pi\gamma^*$ -to- $\pi\pi$



Exploratory $\pi\gamma^*$ -to- $\pi\pi$ / $\pi\gamma^*$ -to- ρ calculation:

$m_\pi = 391$ MeV



Matrix element determined in 42 kinematic point: $(E_{\pi\pi}, Q^2)$

Lorentz decomposition:

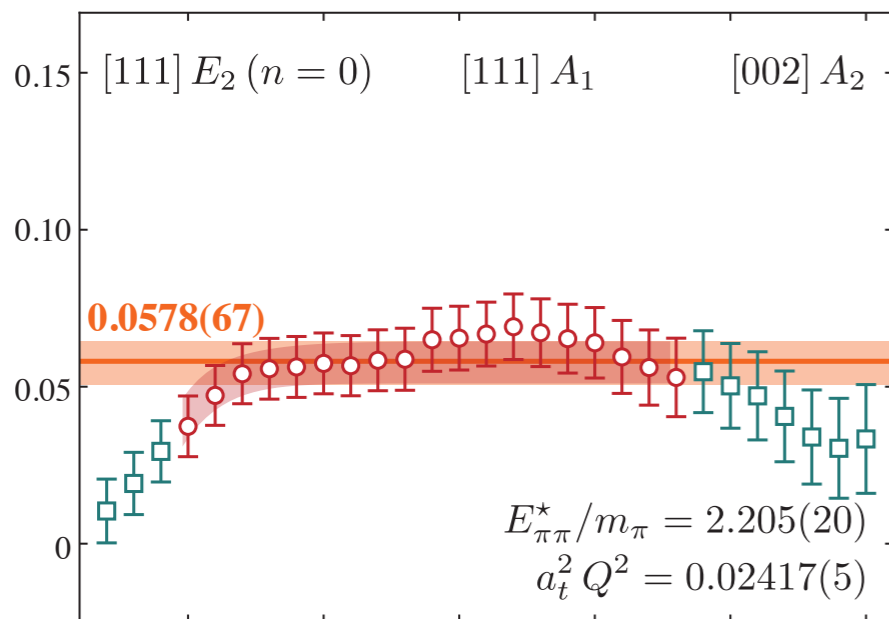
$$\mathcal{H}_{\pi\pi, \pi\gamma^*}^\mu = \epsilon^{\mu\nu\alpha\beta} P_{\pi, \nu} P_{\pi\pi, \alpha} \epsilon_\beta(\lambda_{\pi\pi}, \mathbf{P}_{\pi\pi}) \frac{2}{m_\pi} \mathcal{A}_{\pi\pi, \pi\gamma^*}$$

$\pi\pi/\rho$ polarization

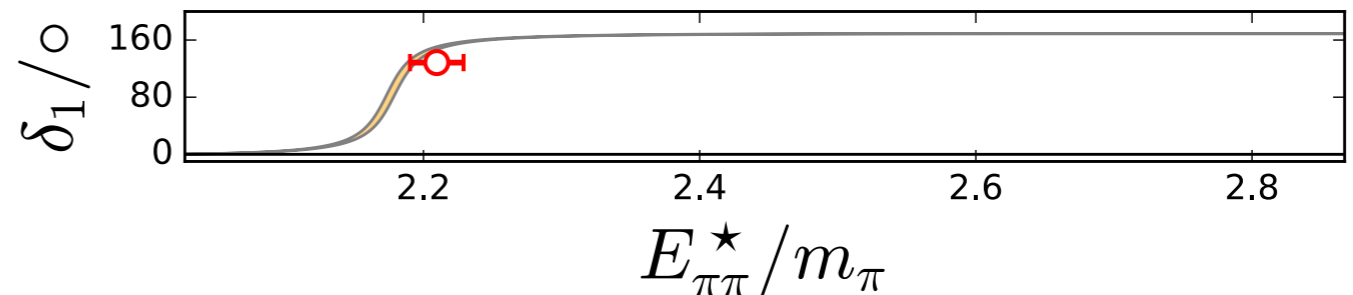
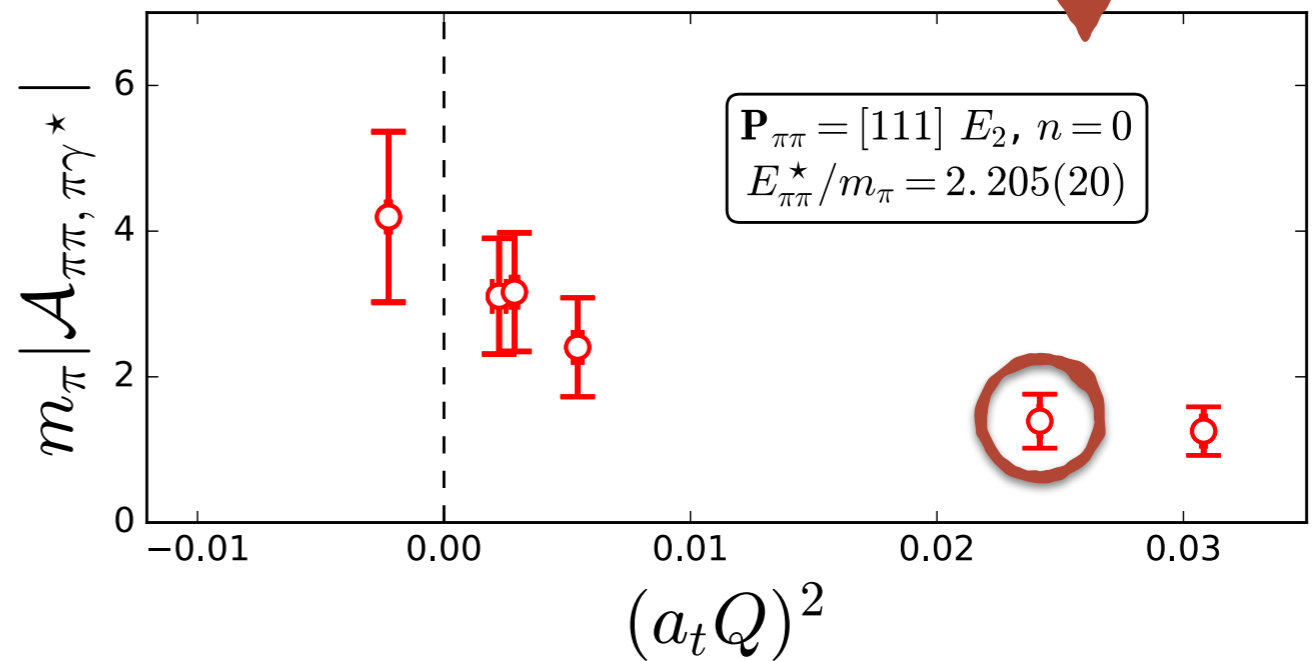
$\pi\pi/\rho$ helicity

Lorentz scalar

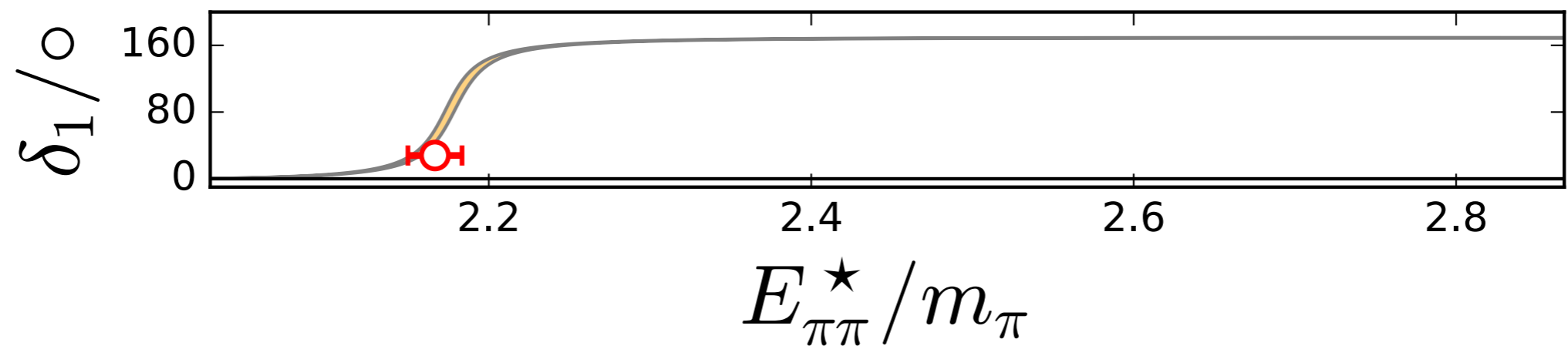
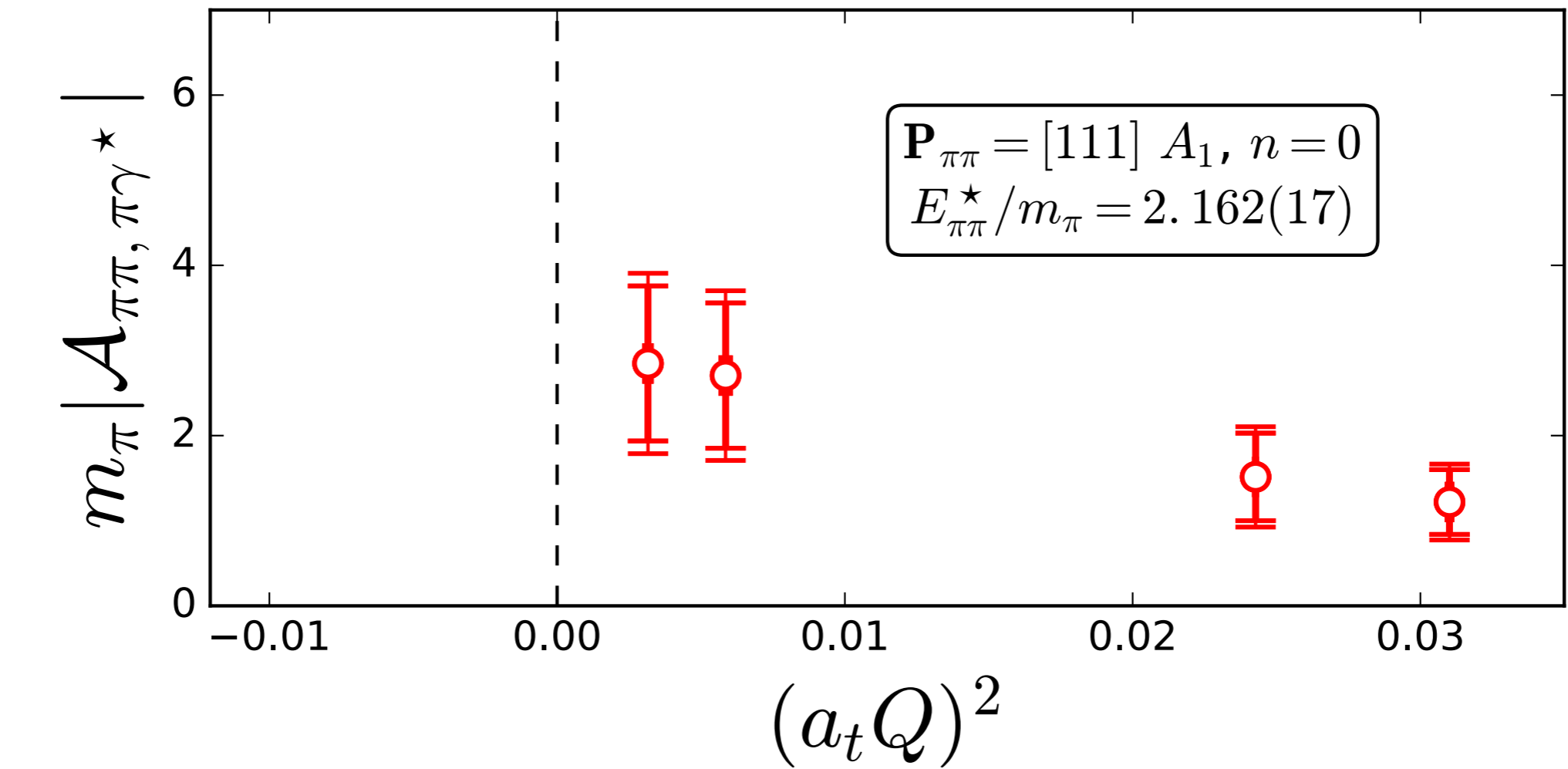
$\pi\gamma^*$ -to- $\pi\pi$ amplitude



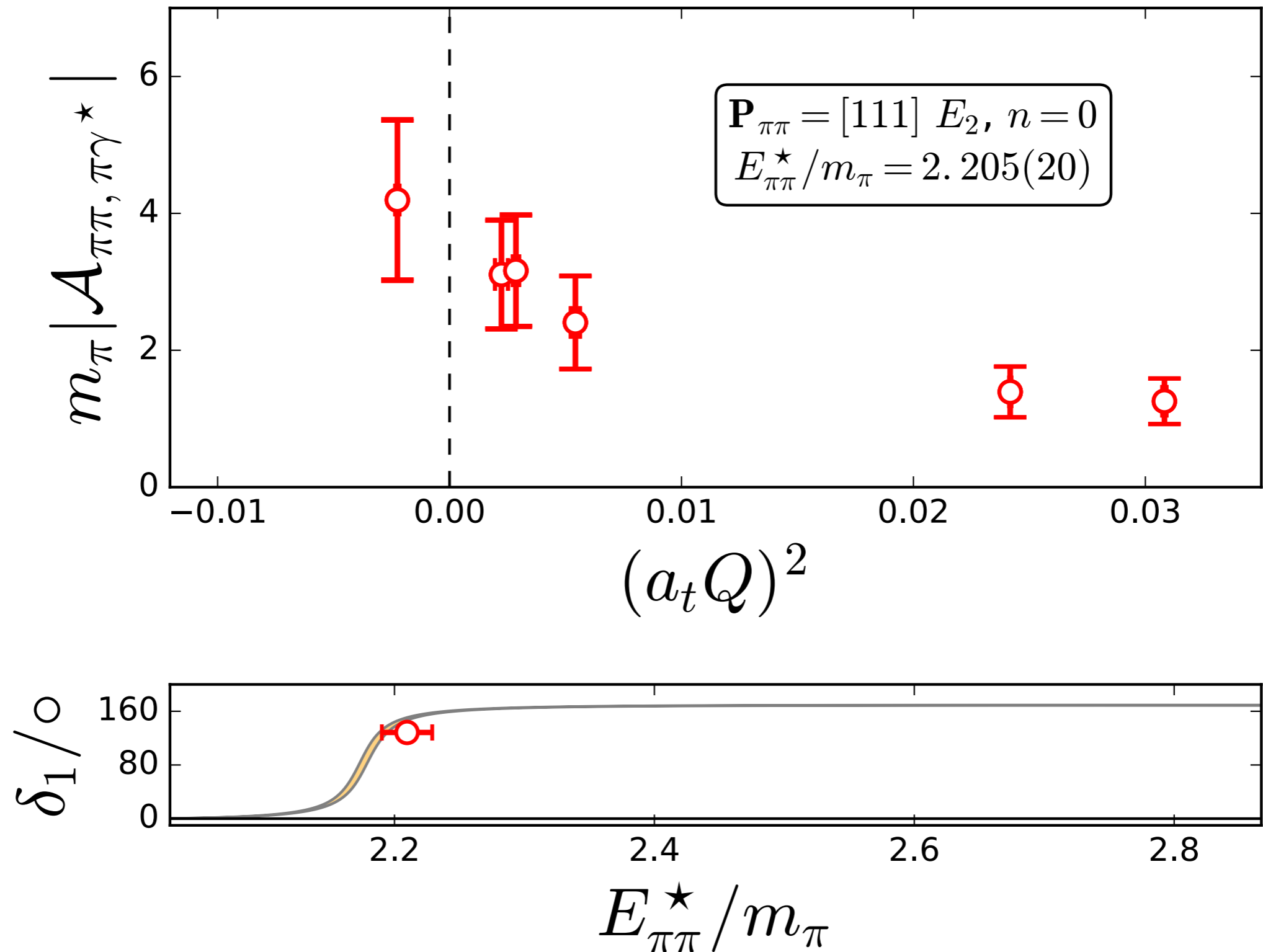
$$|\langle \mathbf{2} | \mathcal{J} | \mathbf{1} \rangle_L|^2 = \mathcal{H} \mathcal{R} \mathcal{H}$$



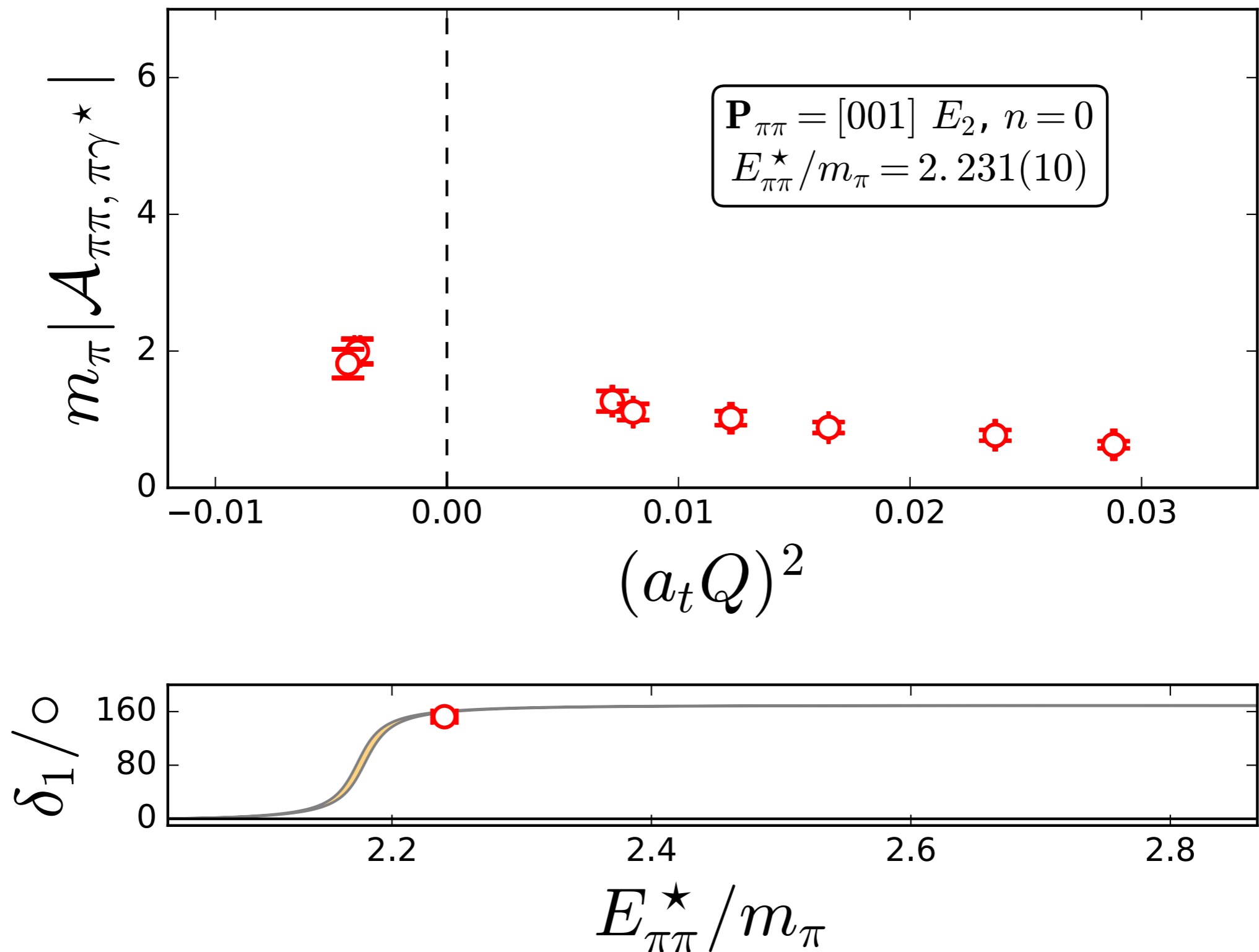
$\pi\gamma^*$ -to- $\pi\pi$ amplitude



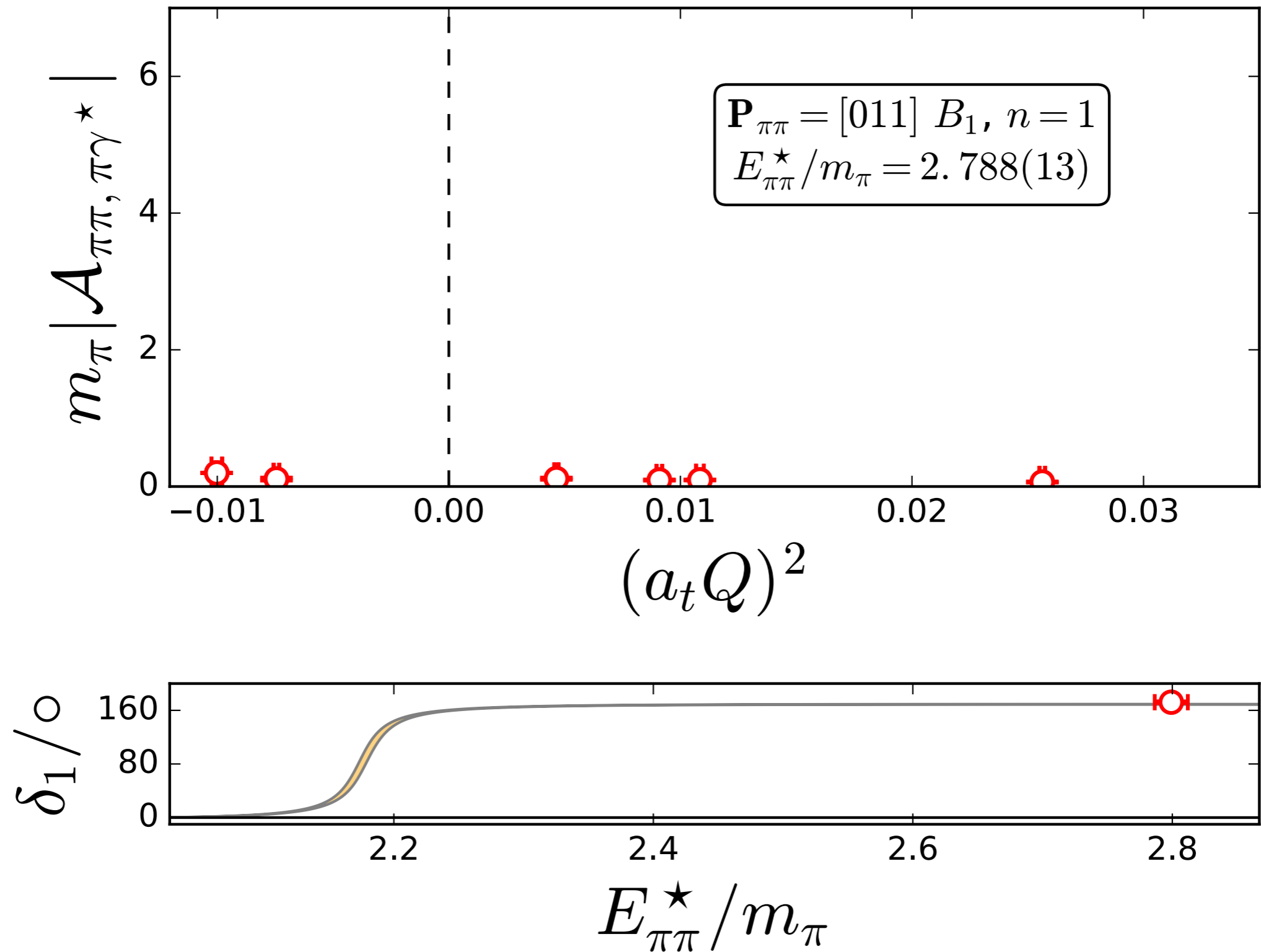
$\pi\gamma^*$ -to- $\pi\pi$ amplitude



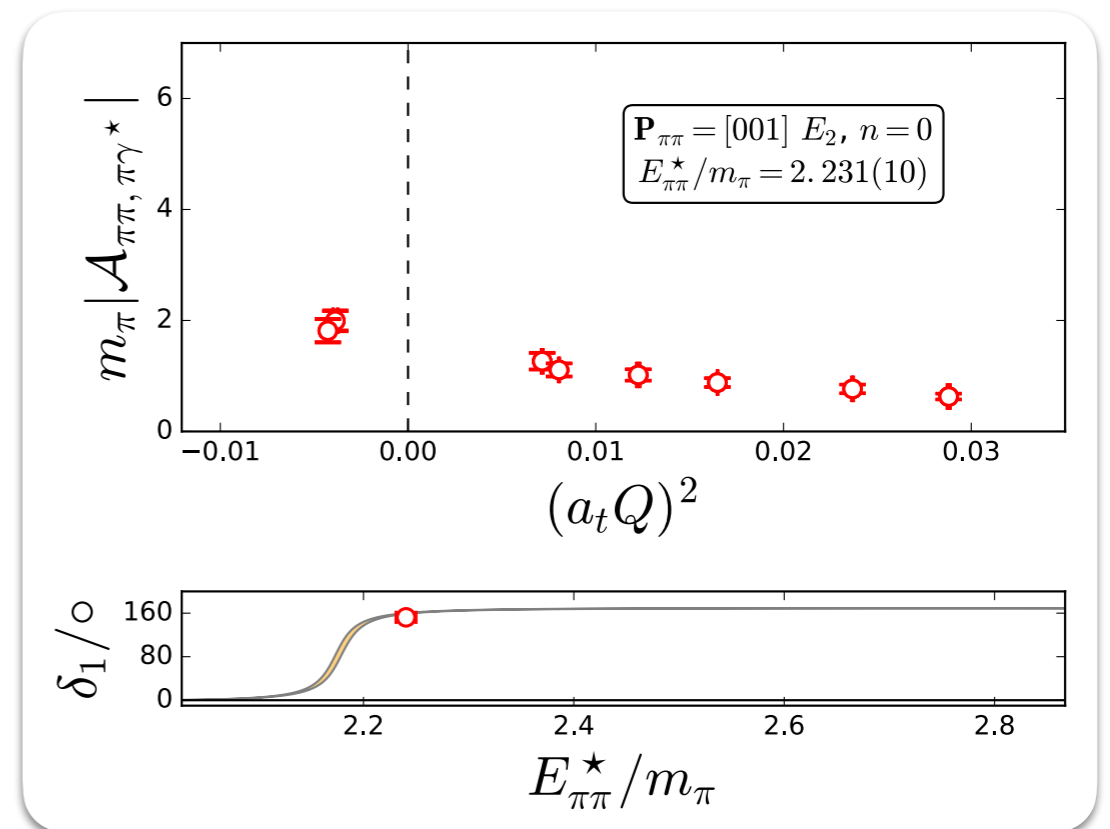
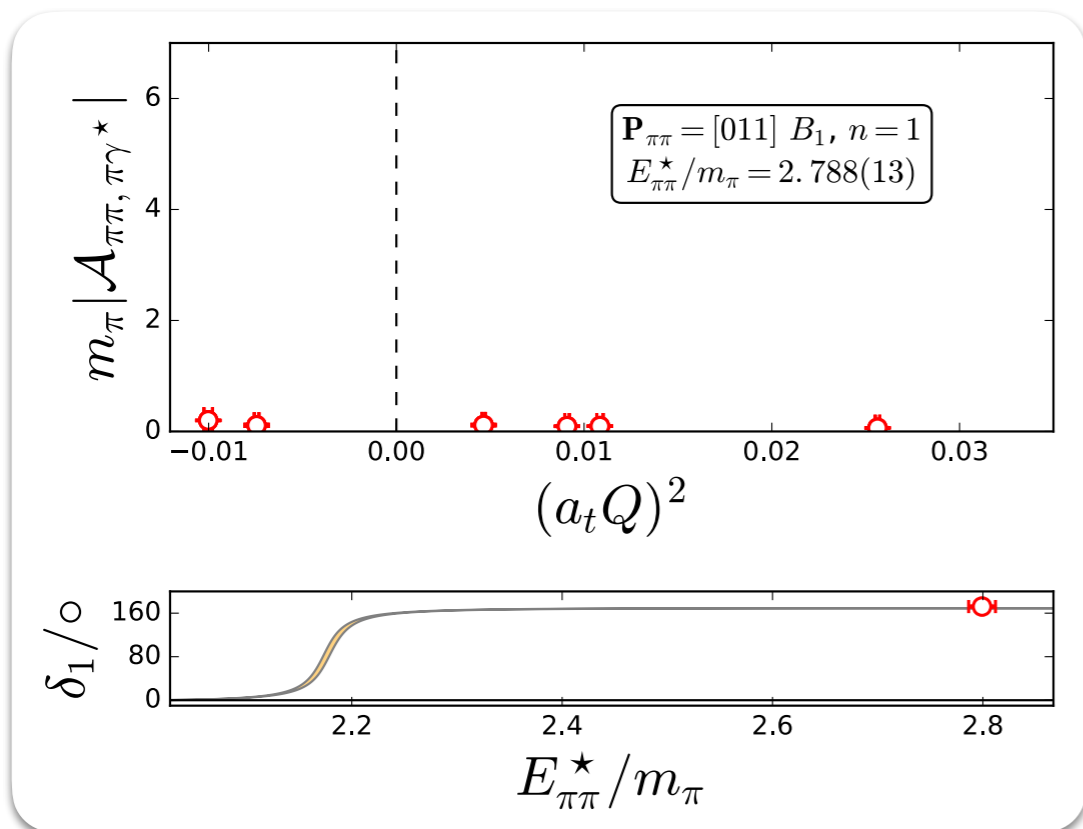
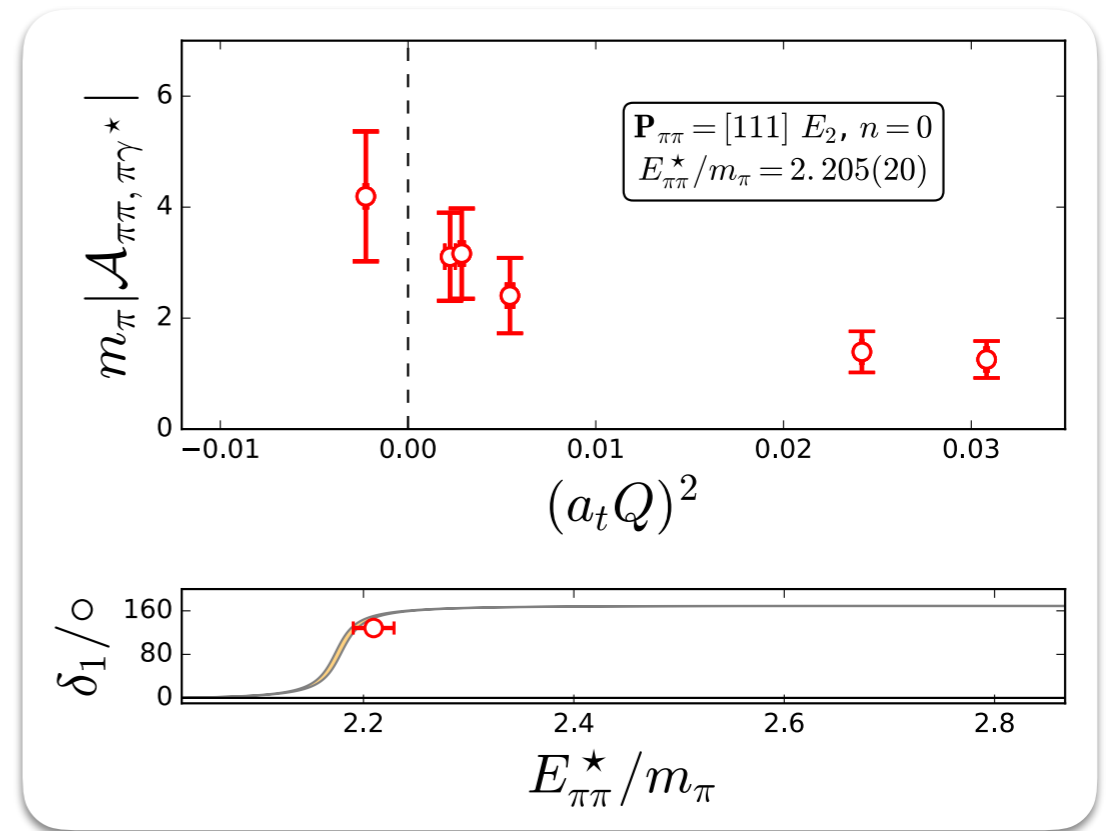
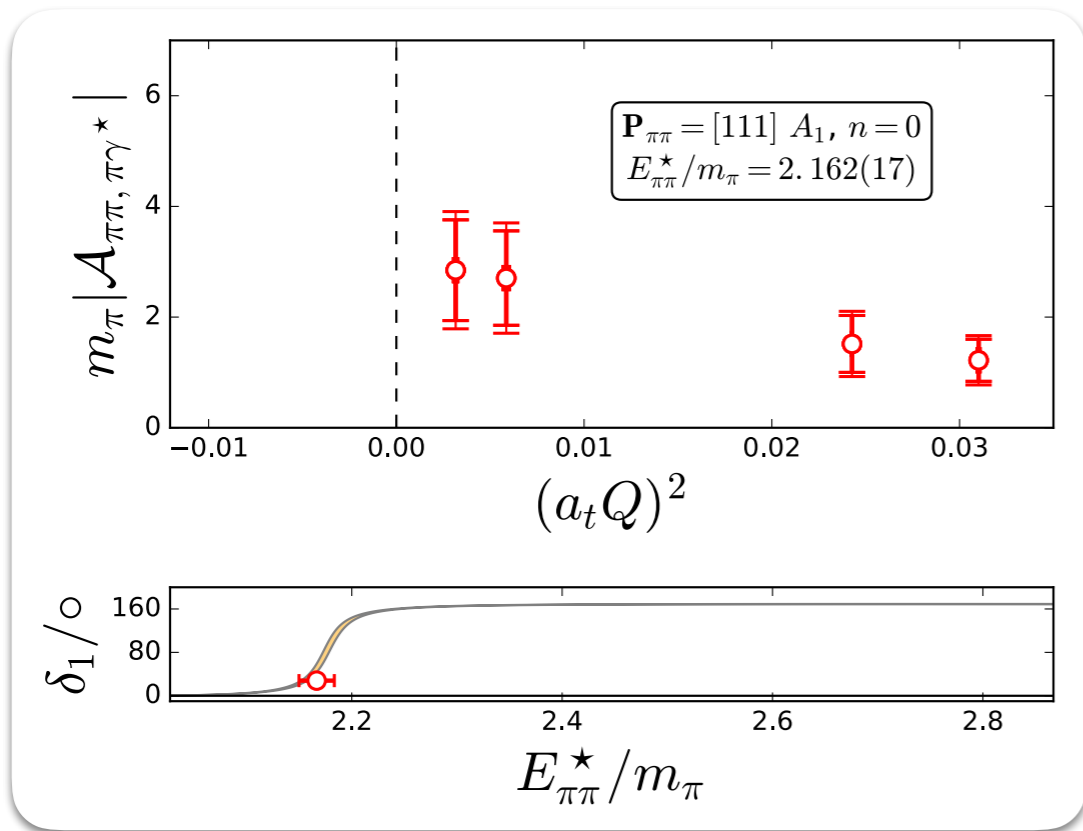
$\pi\gamma^*$ -to- $\pi\pi$ amplitude



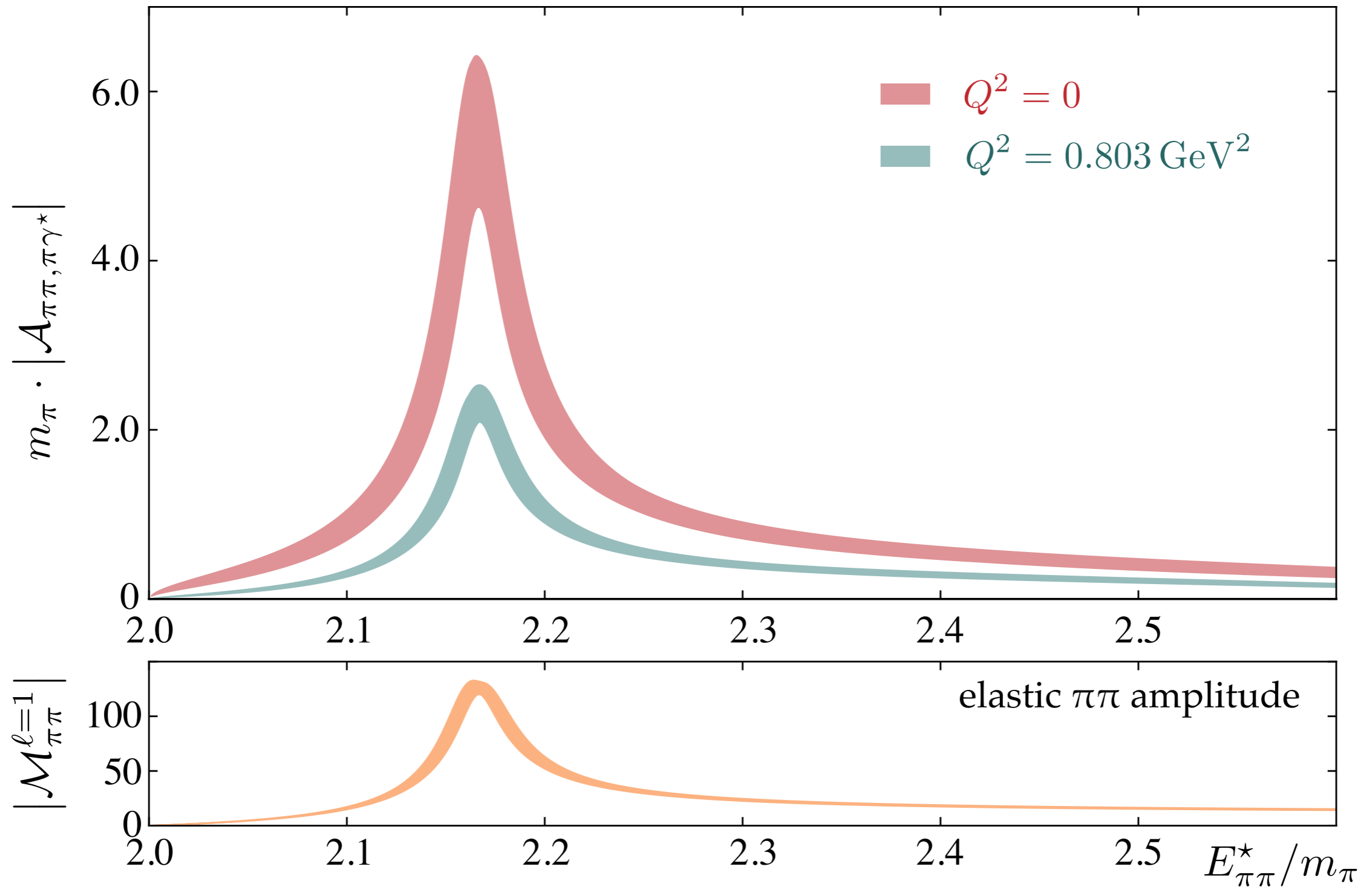
$\pi\gamma^*$ -to- $\pi\pi$ amplitude



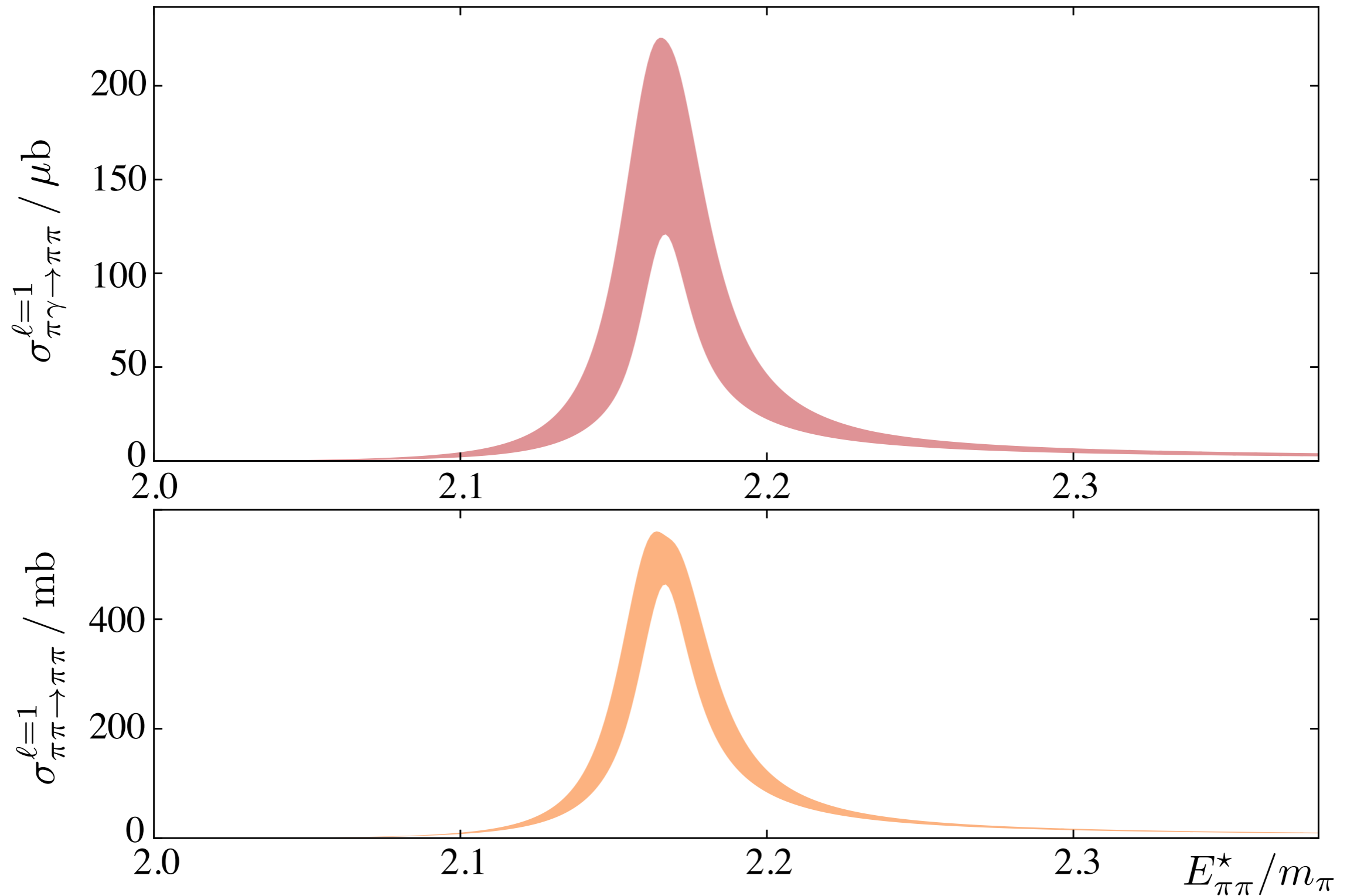
$\pi\gamma^*$ -to- $\pi\pi$ amplitude

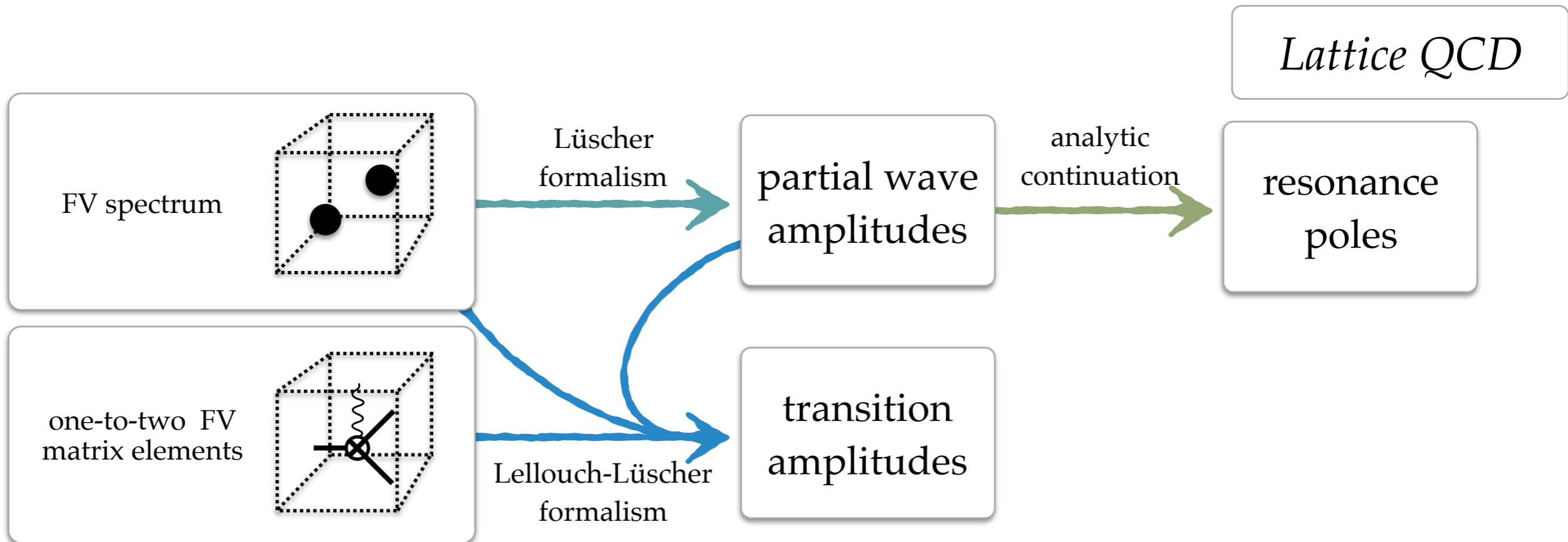
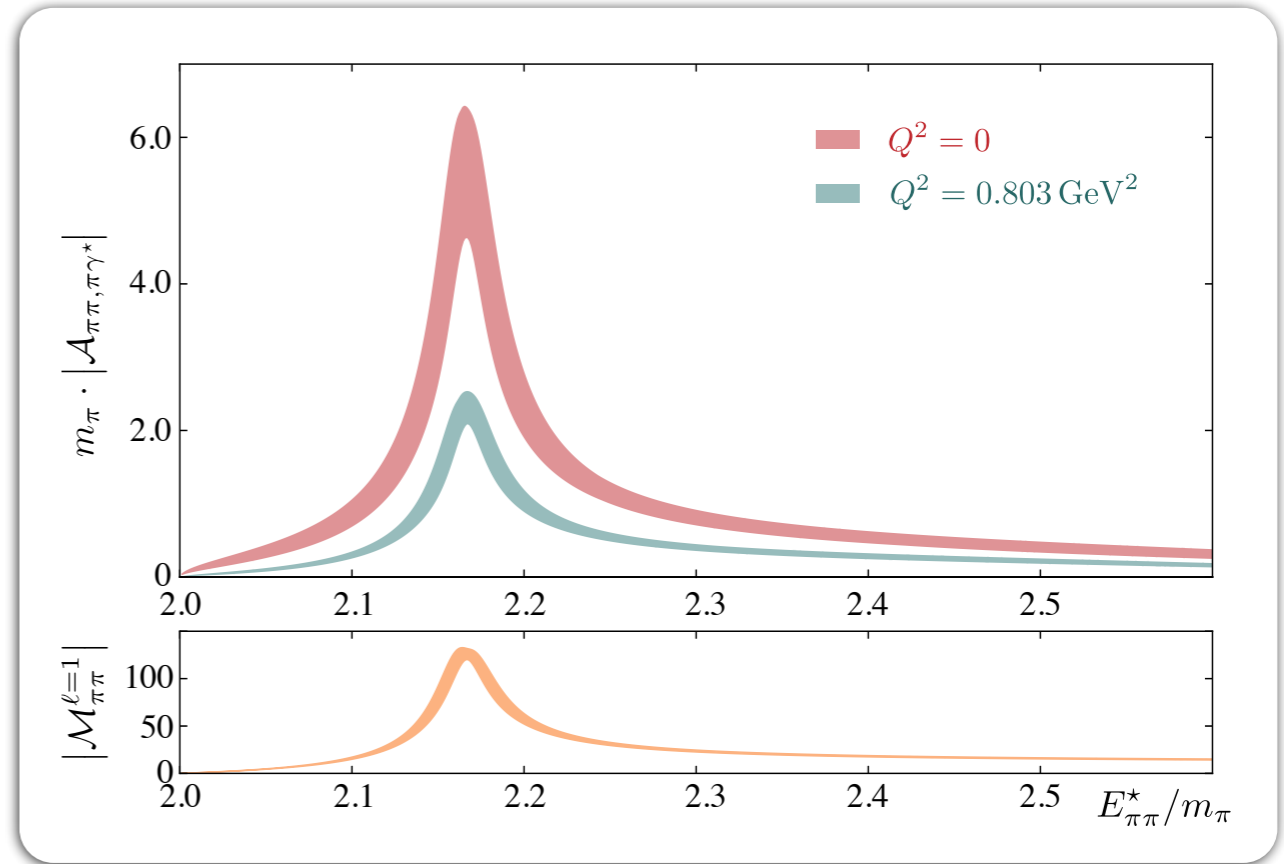
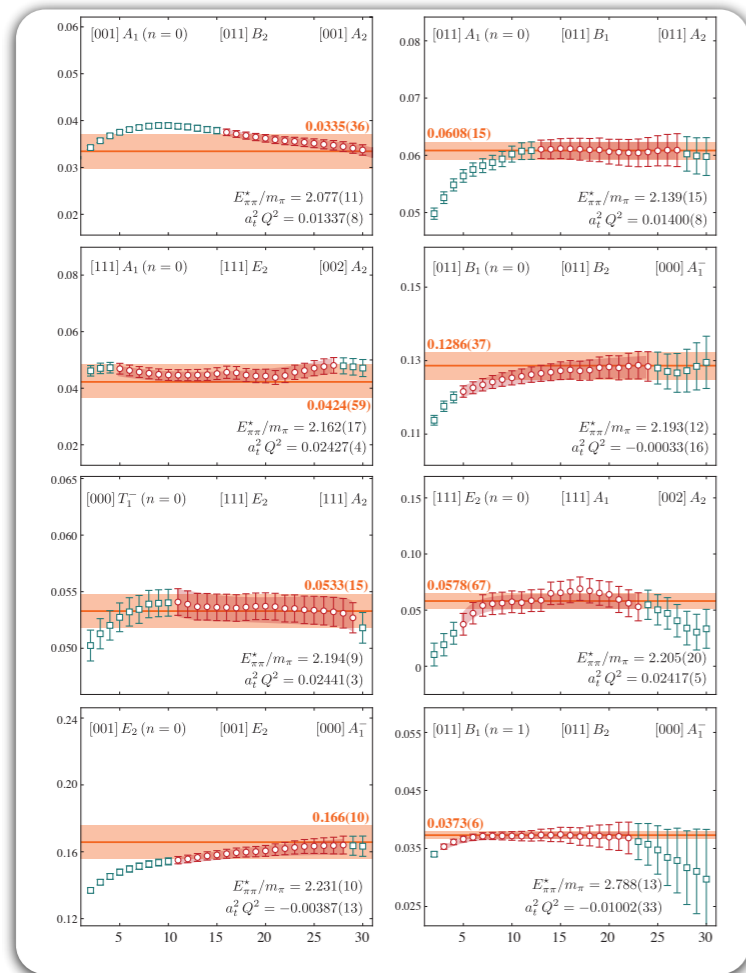


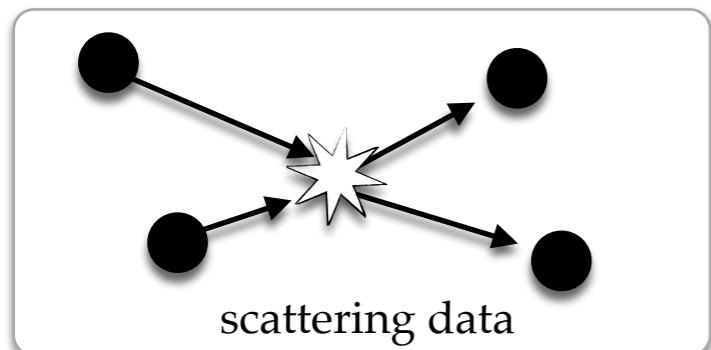
$\pi\gamma^*$ -to- $\pi\pi$ amplitude



$\pi\gamma^*$ -to- $\pi\pi$ cross section





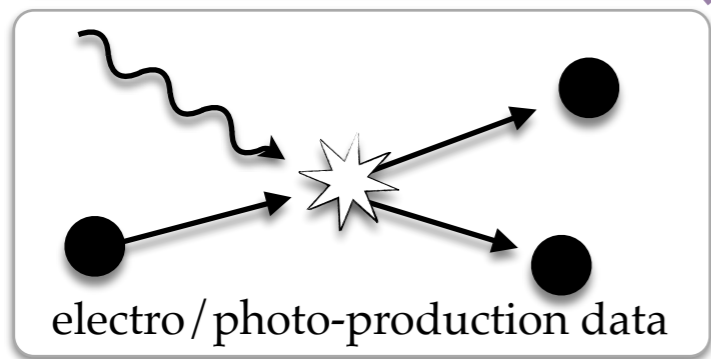


amplitude analysis

partial wave amplitudes

analytic continuation

resonance poles



amplitude analysis

transition amplitudes

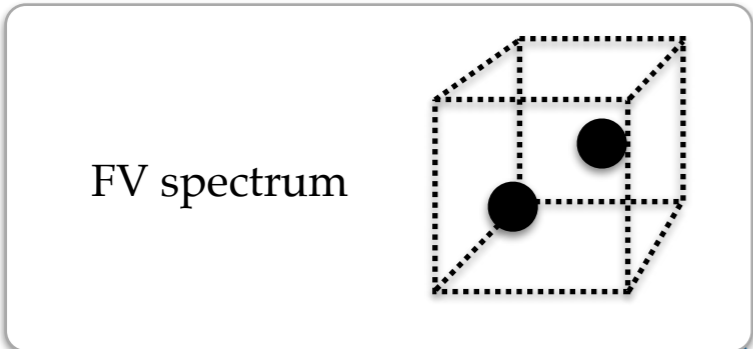
analytic continuation

form factors

Experiment



Lattice QCD

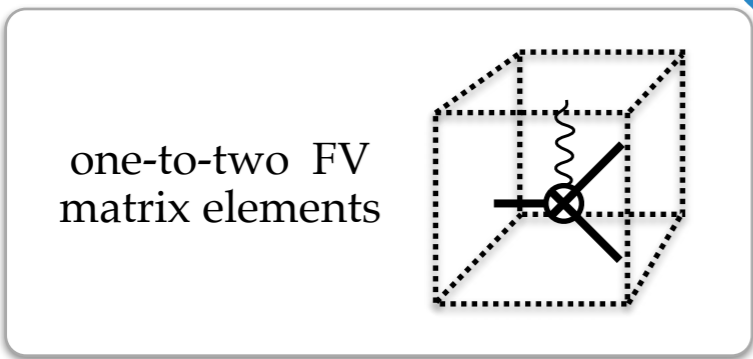


Lüscher formalism

partial wave amplitudes

analytic continuation

resonance poles



Lellouch-Lüscher formalism

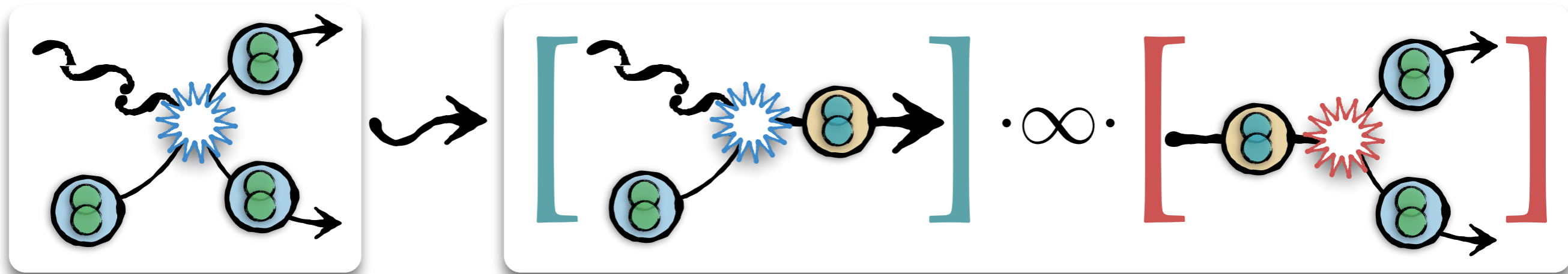
transition amplitudes

analytic continuation

form factors

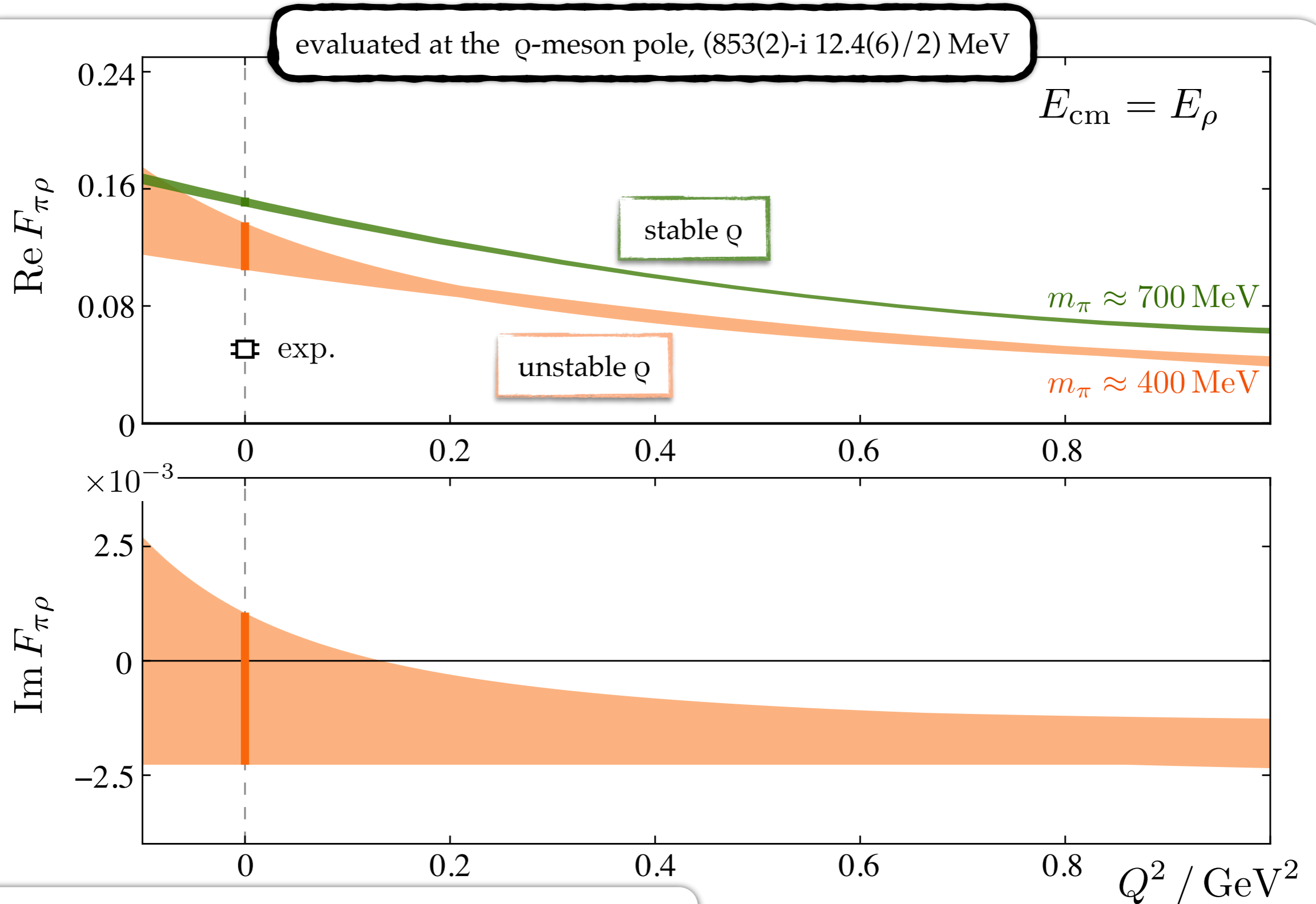
Form factor at ρ pole

- Near the ρ -pole, the $\pi\gamma^*$ -to- $\pi\pi$ diverges
- The residue encodes the $\pi\gamma^*$ -to- ρ form factor



$$\mathcal{A}_{\pi\pi, \pi\gamma^*}(E_{\pi\pi}, Q^2) = \underbrace{F(E_{\pi\pi}, Q^2)} \times \left[\frac{1}{\cot \delta_1(E_{\pi\pi}) - i} \right] \times \sqrt{\frac{16\pi}{\underbrace{q_{\pi\pi} \Gamma(E_{\pi\pi})}}}$$

Form factor at ρ pole

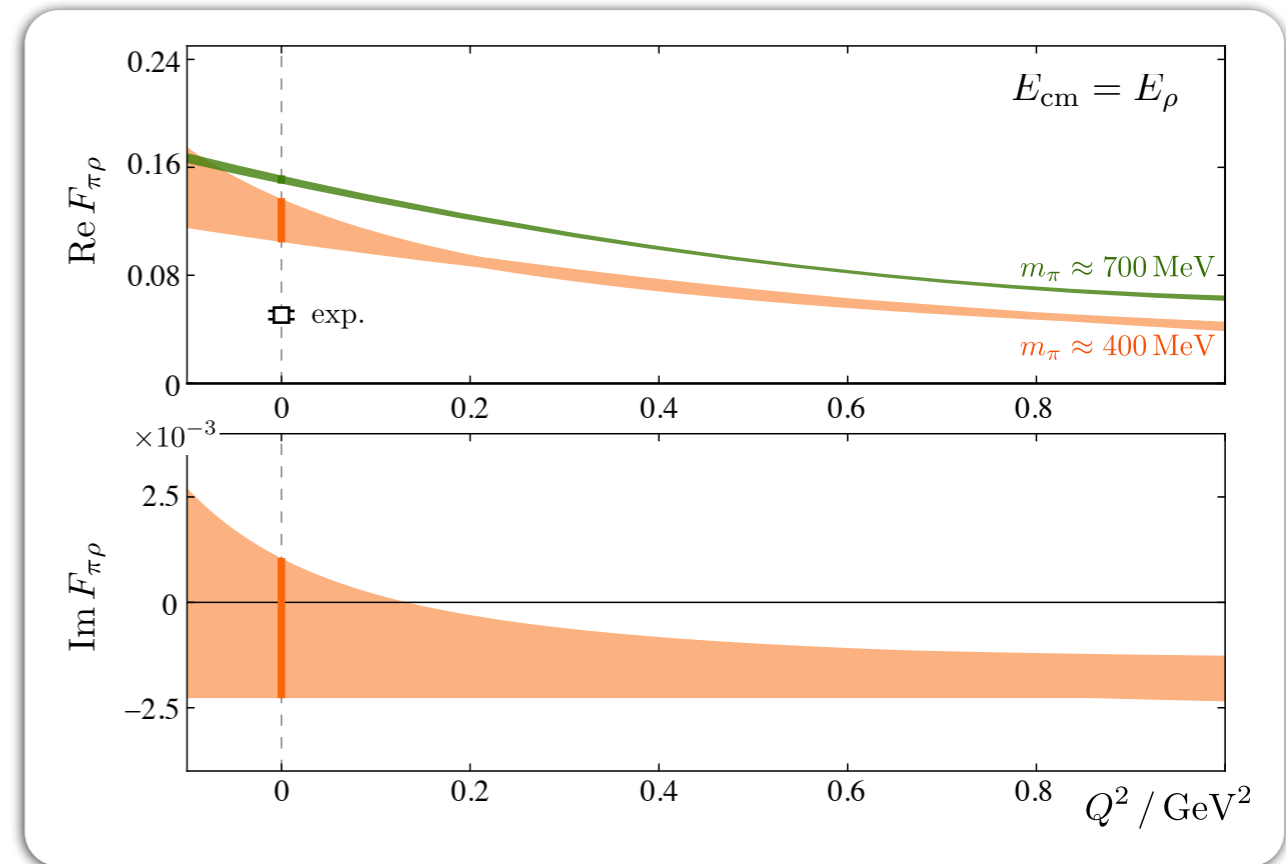


Shultz, Dudek, & Edwards (2014)

RB, Dudek, Edwards, Shultz, Thomas & Wilson (2015)

Some comments

- only lattice calculation of a resonance form factor
- all previous studies ignore resonant nature - it was not known how to
- these include...



PRL 114, 132002 (2015)

PHYSICAL REVIEW LETTERS

week ending
3 APRIL 2015

Lattice QCD Evidence that the $\Lambda(1405)$ Resonance is an Antikaon-Nucleon Molecule

Jonathan M. M. Hall,¹ Waseem Kamleh,¹ Derek B. Leinweber,^{1,*} Benjamin J. Menadue,^{1,2}
 Benjamin J. Owen,¹ Anthony W. Thomas,^{1,3} and Ross D. Young^{1,3}

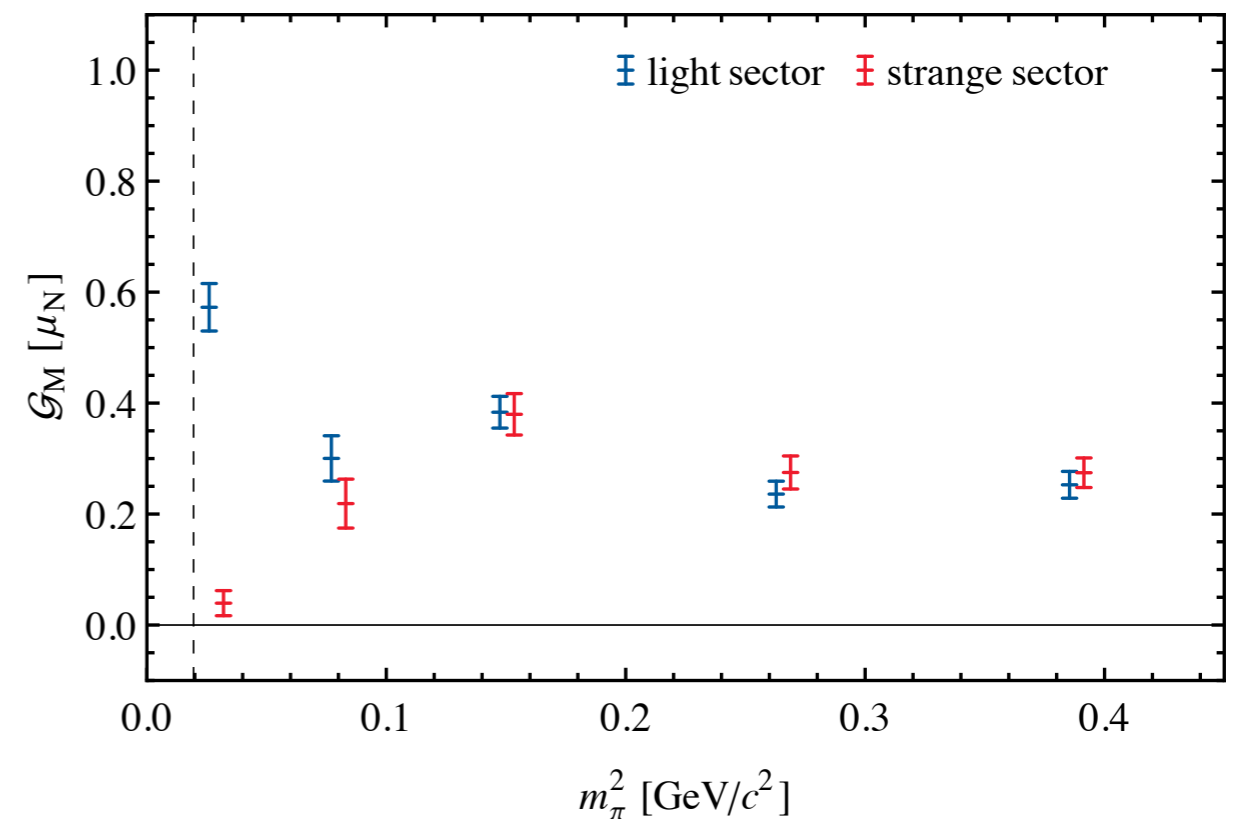
¹Special Research Centre for the Subatomic Structure of Matter (CSSM), Department of Physics,
 University of Adelaide, Adelaide, South Australia 5005, Australia

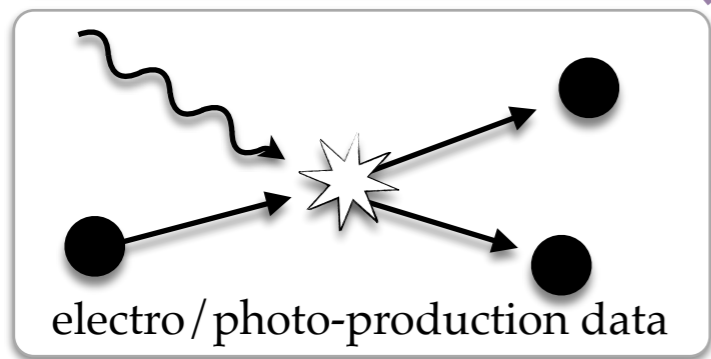
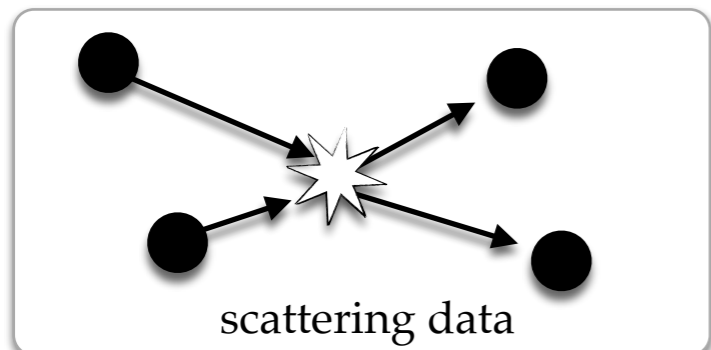
²National Computational Infrastructure (NCI), Australian National University, Canberra, Australian Capital Territory 0200, Australia

³ARC Centre of Excellence for Particle Physics at the Terascale (CoEPP), Department of Physics,
 University of Adelaide, Adelaide, South Australia 5005, Australia

(Received 12 November 2014; revised manuscript received 10 February 2015; published 1 April 2015)

No, there is no evidence *yet* from lattice QCD that this is in fact true!





amplitude analysis

partial wave amplitudes

analytic continuation

resonance poles

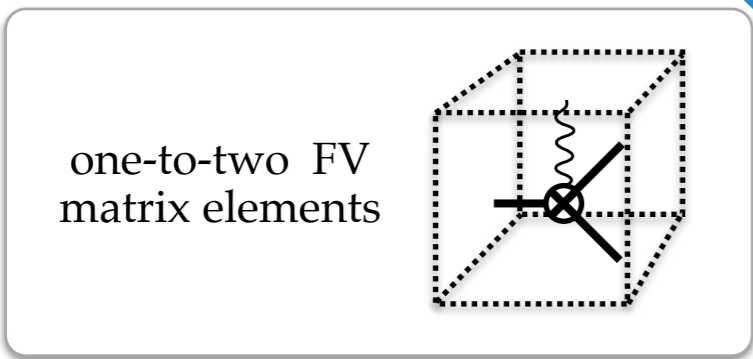
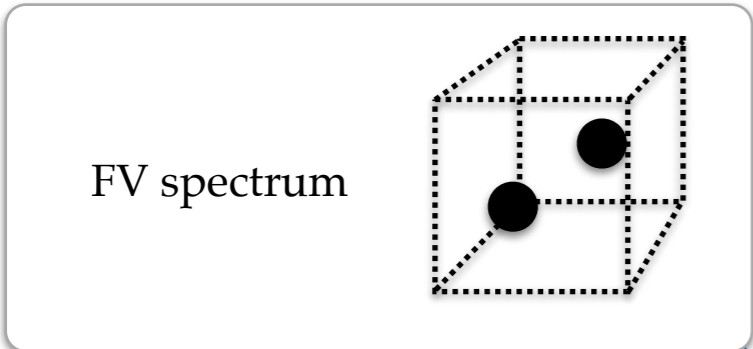
amplitude analysis

transition amplitudes

analytic continuation

form factors

Experiment



Lüscher formalism

partial wave amplitudes

analytic continuation

resonance poles

Lellouch-Lüscher formalism

transition amplitudes

analytic continuation

form factors

Lattice QCD

Questions to answer:

📌 What is possible from lattice QCD regarding electro/photo-production?

📌 How can lattice and experiment compare to each other?

📌 What is the possibility/timescale for extending this to baryons?

📌 What does lattice “*need*” from experiment?

Questions to answer:

What is possible from lattice QCD regarding electro/photo-production?

How can lattice and experiment compare to each other?

What is the possibility / timescale for extending this to baryons?



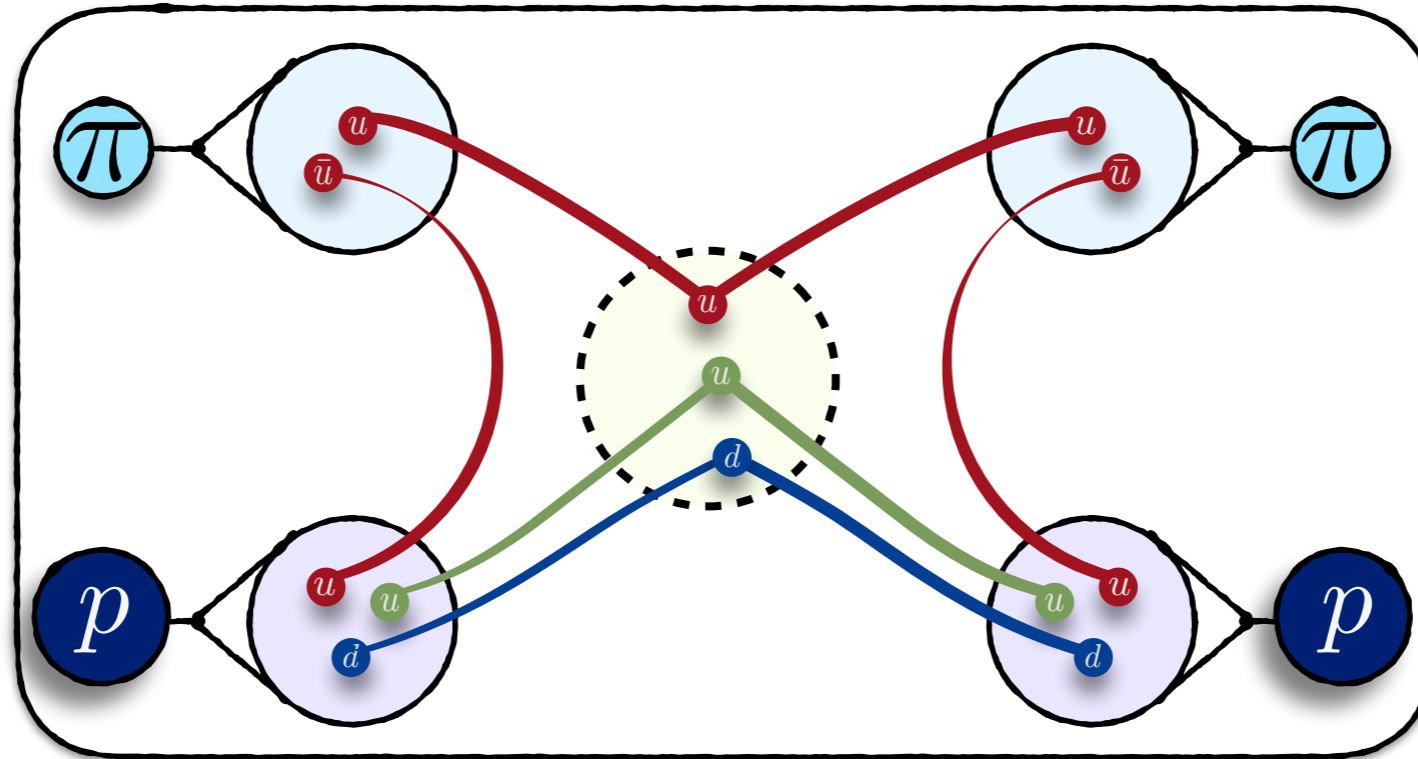
Why are baryons harder to study via lattice QCD?

What are we actively doing towards this goal?

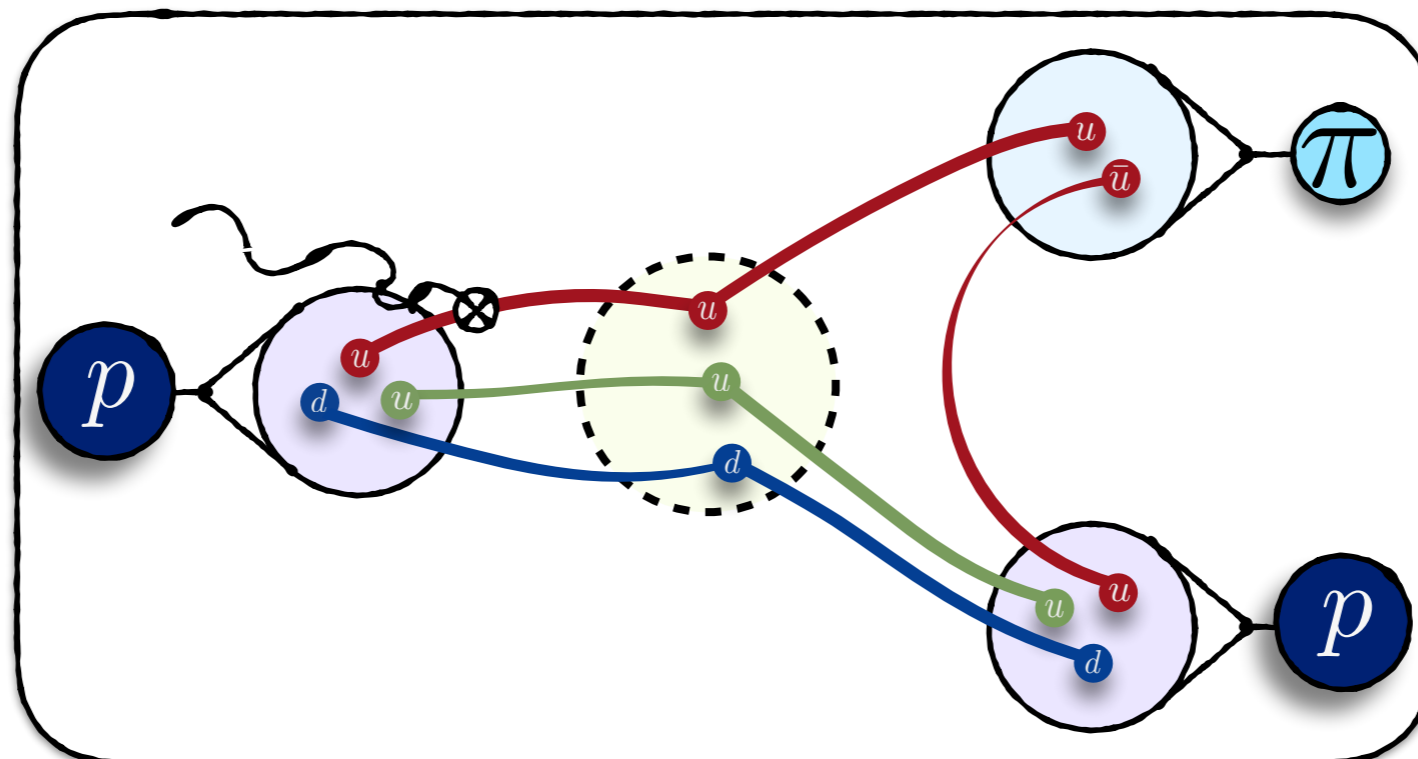
First things first

[no baryonic resonances from QCD, yet!]

First, we must calculate:

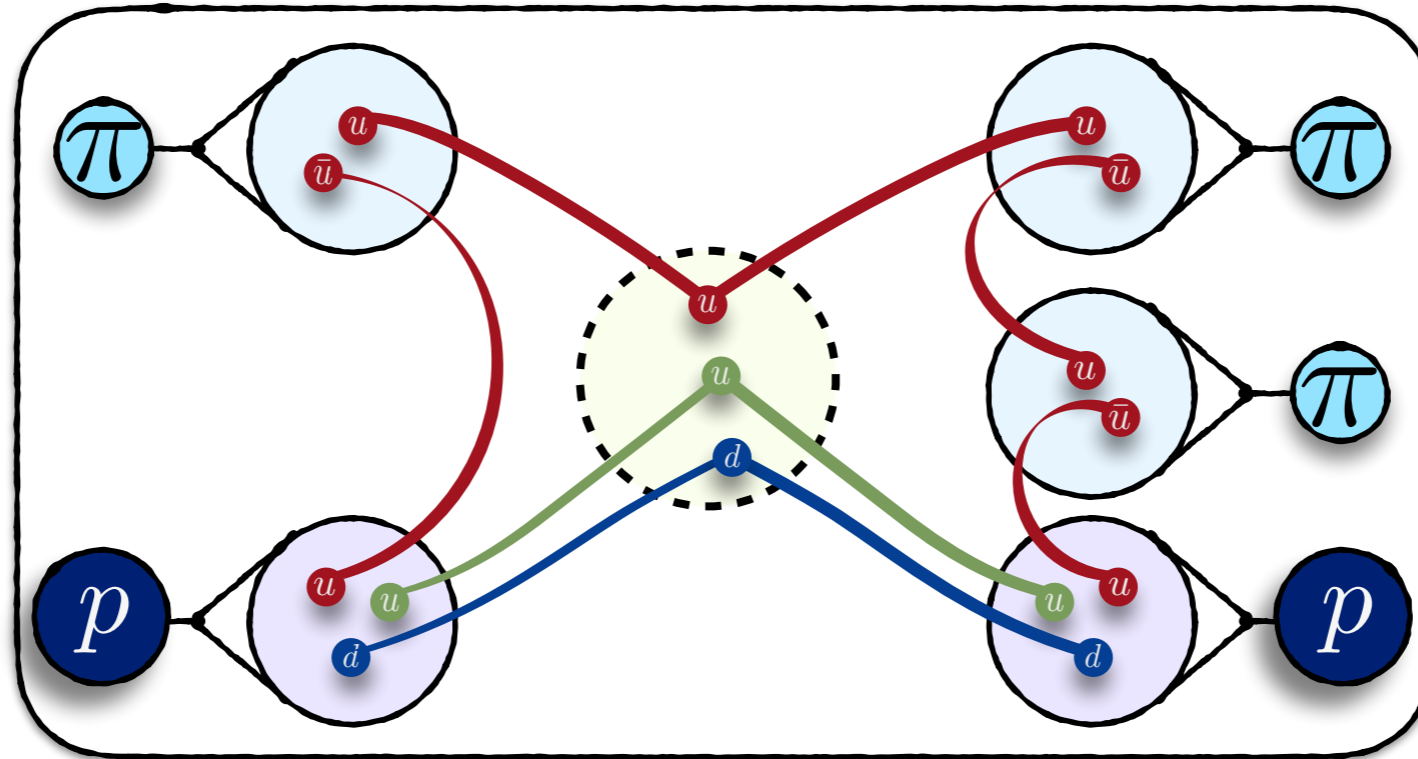


Before we study:



Challenges - three-body

Cannot ignore:



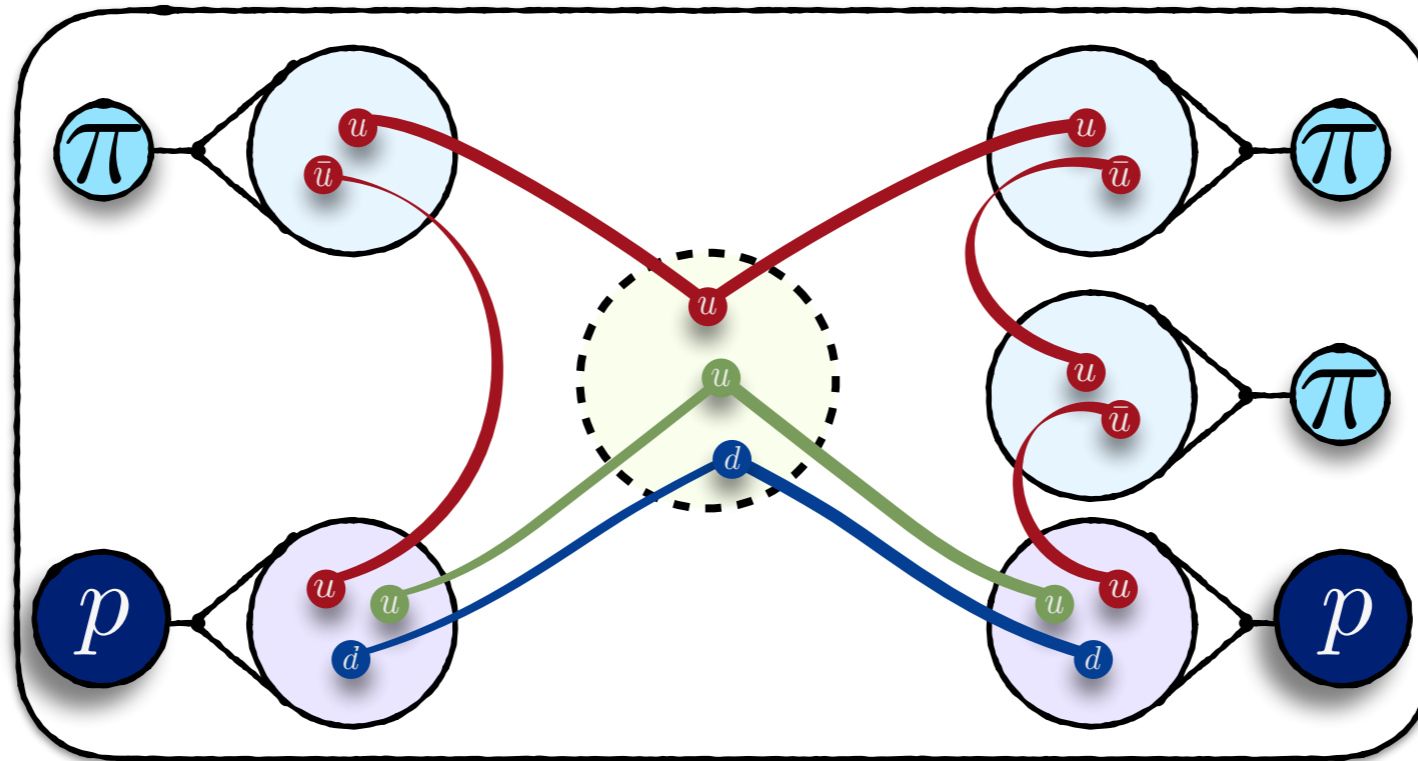
$\sim 30-40\%$

[More contractions,
more channels, etc.]

[Formal open question,
Harder to analyze]

Challenges - three-body

Cannot ignore:



~30-40%

[More contractions,
more channels, etc.]

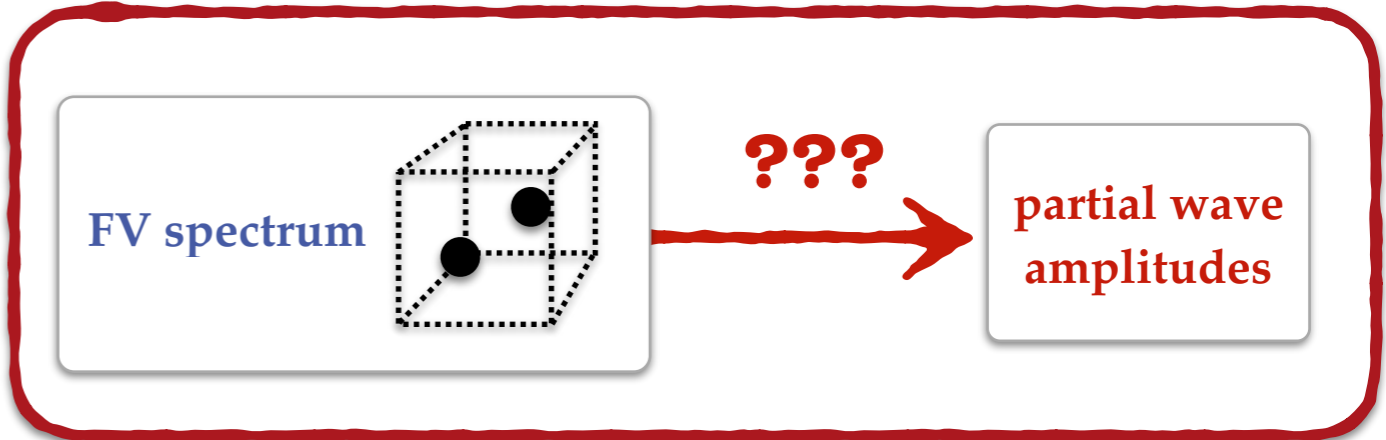
[Formal open question,
Harder to analyze]

Sharpe, Hansen, RB

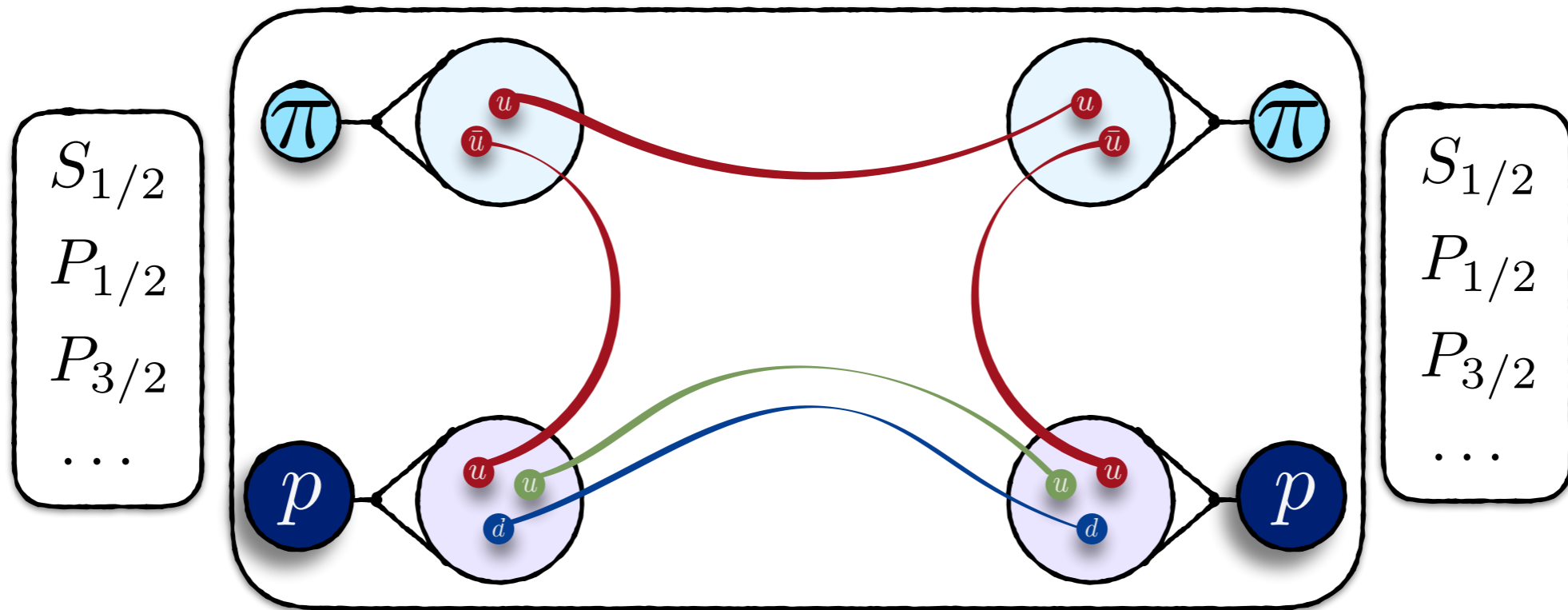
*had we had the spectrum,
nobody knows what it means!*

*obtaining FV spectrum
is harder, but doable*

Robert



Challenges - more partial waves and mixing



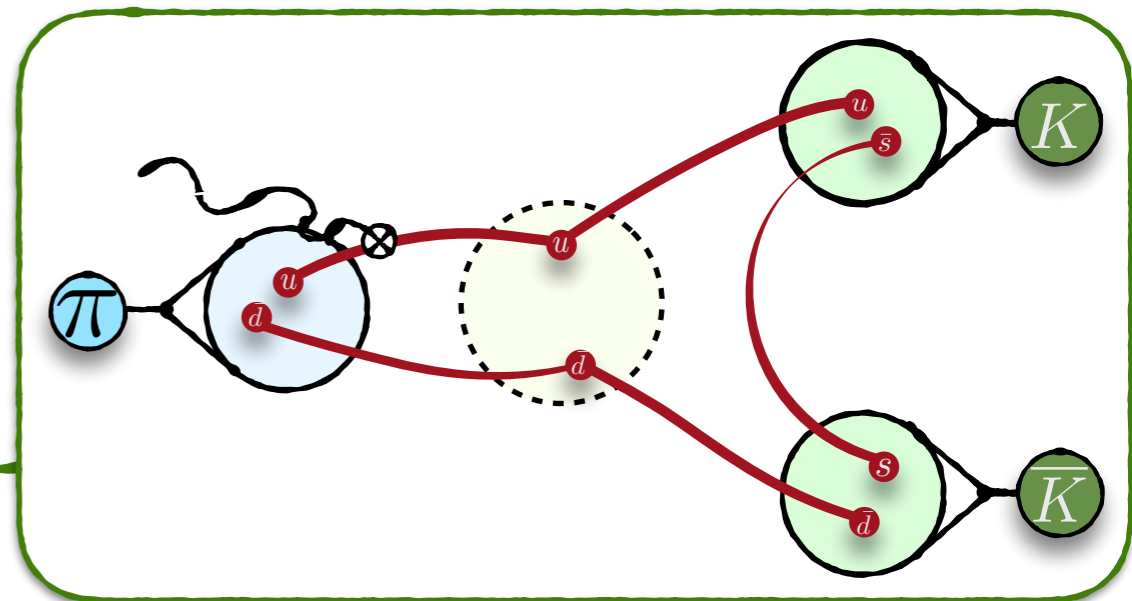
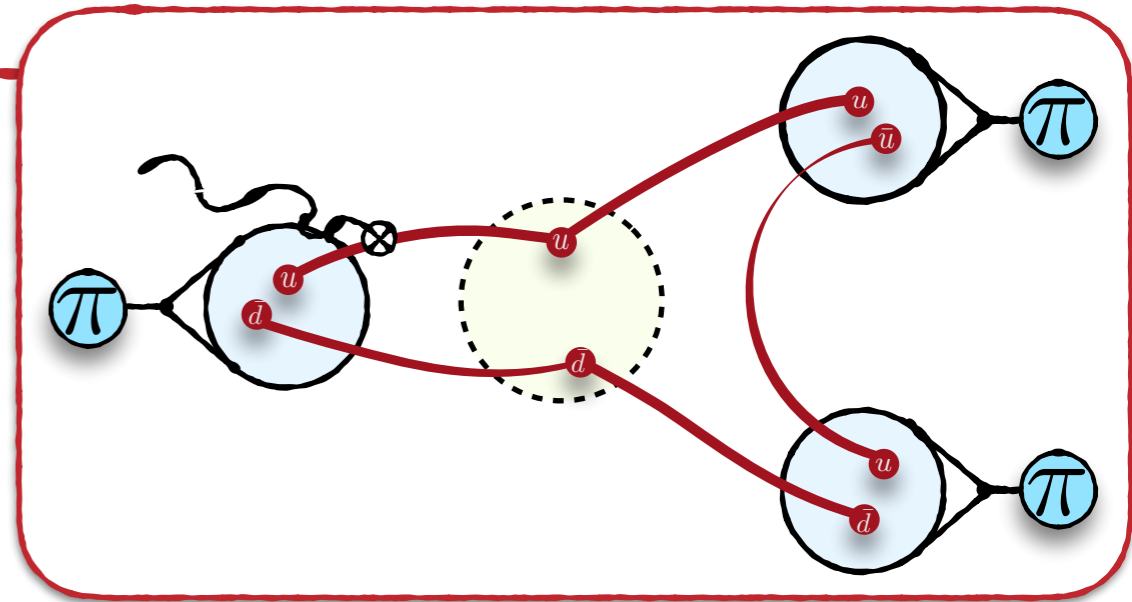
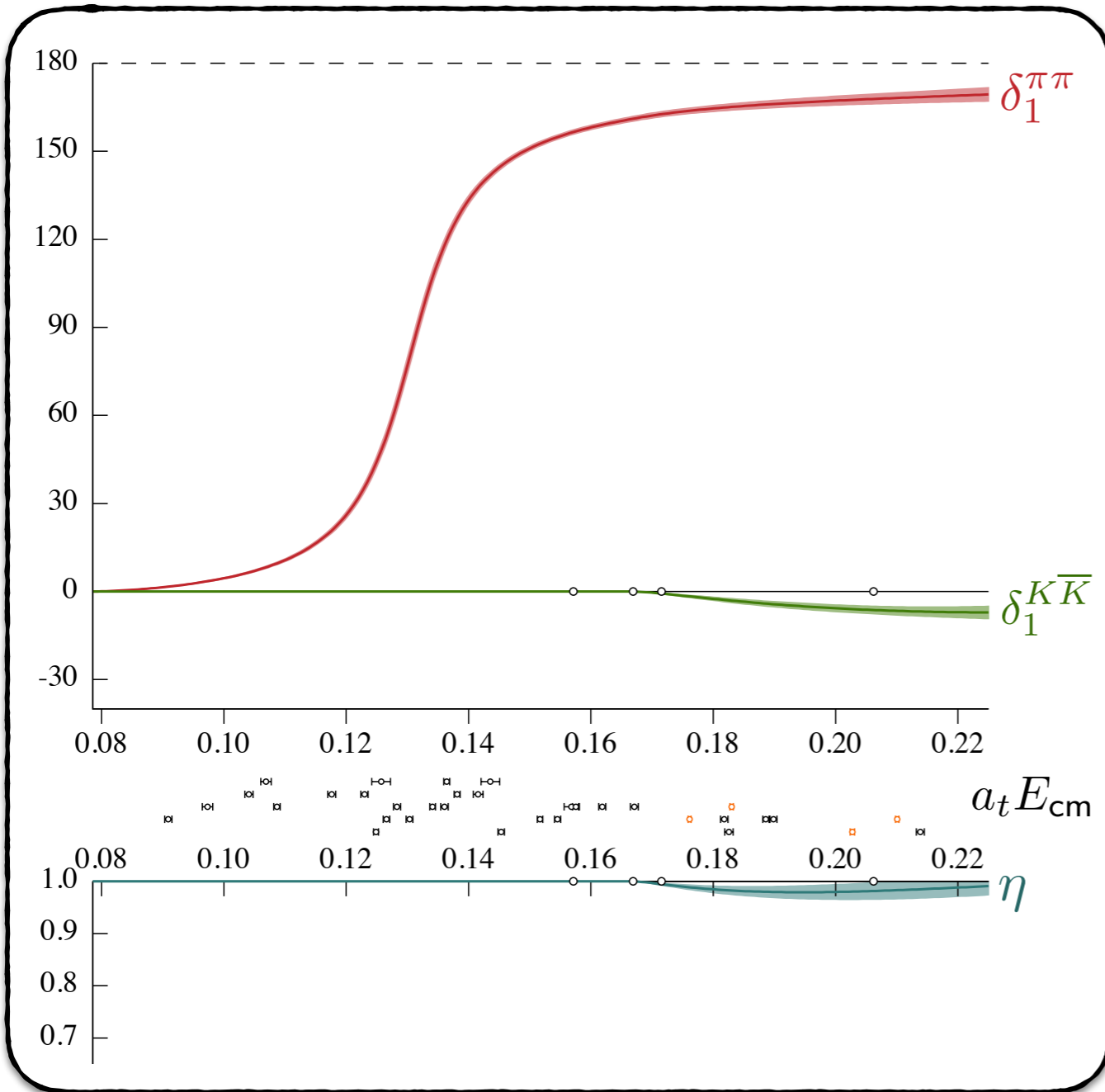
$$\det \left[\begin{pmatrix} F_{S_{1/2},S_{1/2}} & F_{S_{1/2},P_{1/2}} & F_{S_{1/2},P_{3/2}} & \dots \\ F_{P_{1/2},S_{1/2}} & F_{P_{1/2},P_{1/2}} & F_{P_{1/2},P_{3/2}} & \dots \\ F_{P_{3/2},S_{1/2}} & F_{P_{3/2},P_{1/2}} & F_{P_{3/2},P_{3/2}} & \dots \\ \dots & \dots & \dots & \dots \end{pmatrix}^{-1} + \begin{pmatrix} \mathcal{M}_{S_{1/2}} & & & \\ & \mathcal{M}_{P_{1/2}} & & \\ & & \mathcal{M}_{P_{3/2}} & \\ & & & \dots \end{pmatrix} \right] = 0$$

formally addressed - RB (2014)

currently being implemented - Jo & Dave Wilson

Expectations for the future

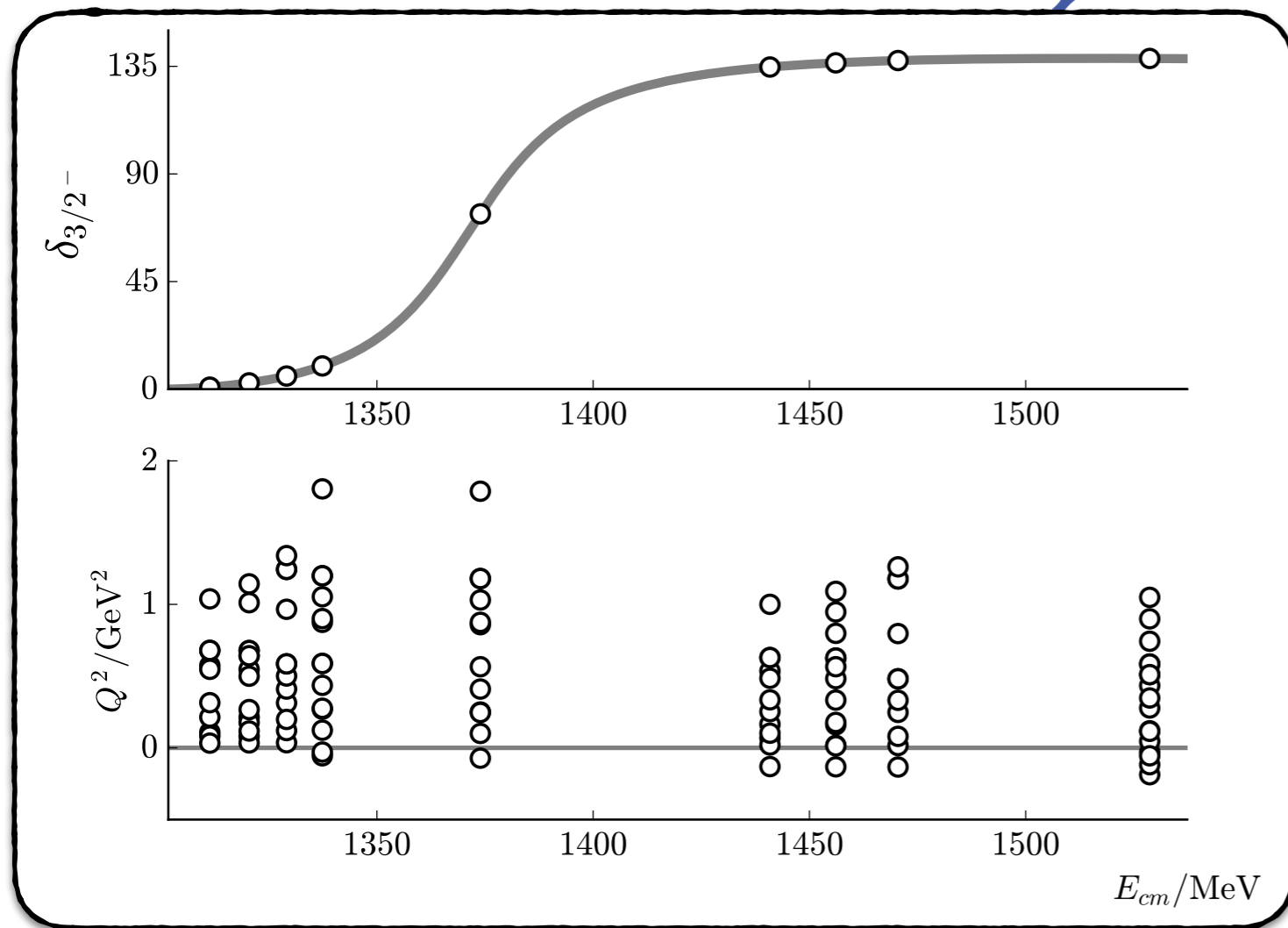
More mesons: reproducing calculation using $m_\pi \sim 236$ MeV, including coupled channels



Expectations for the future

Kinematic expectations for $m_\pi \sim 236$ MeV, single volume $L/a_s=32$

Second volume, $L/a_s=40$, is underway

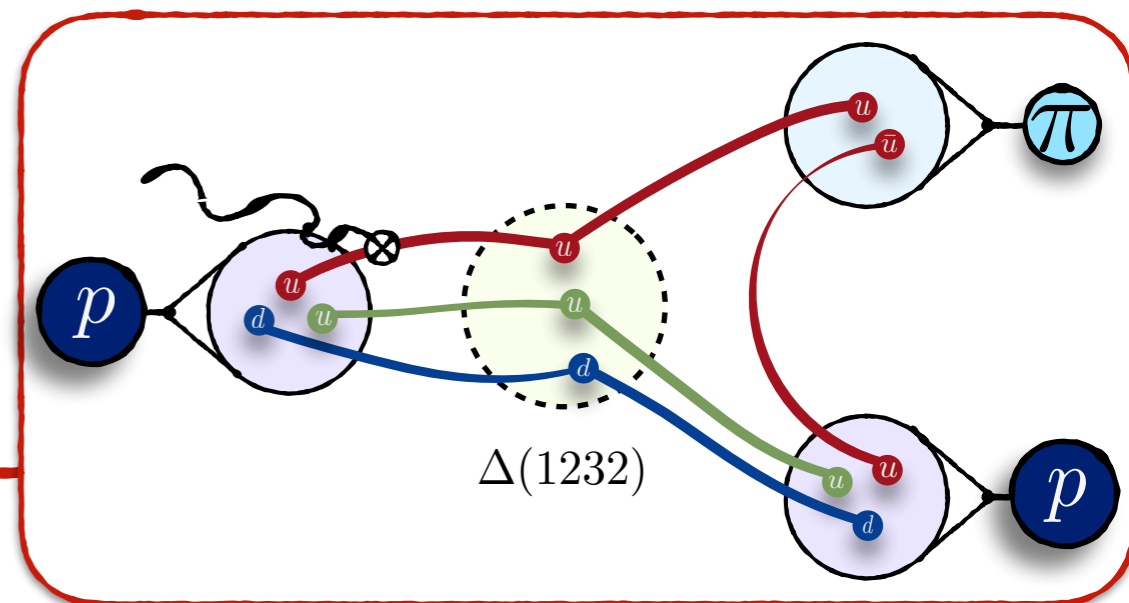
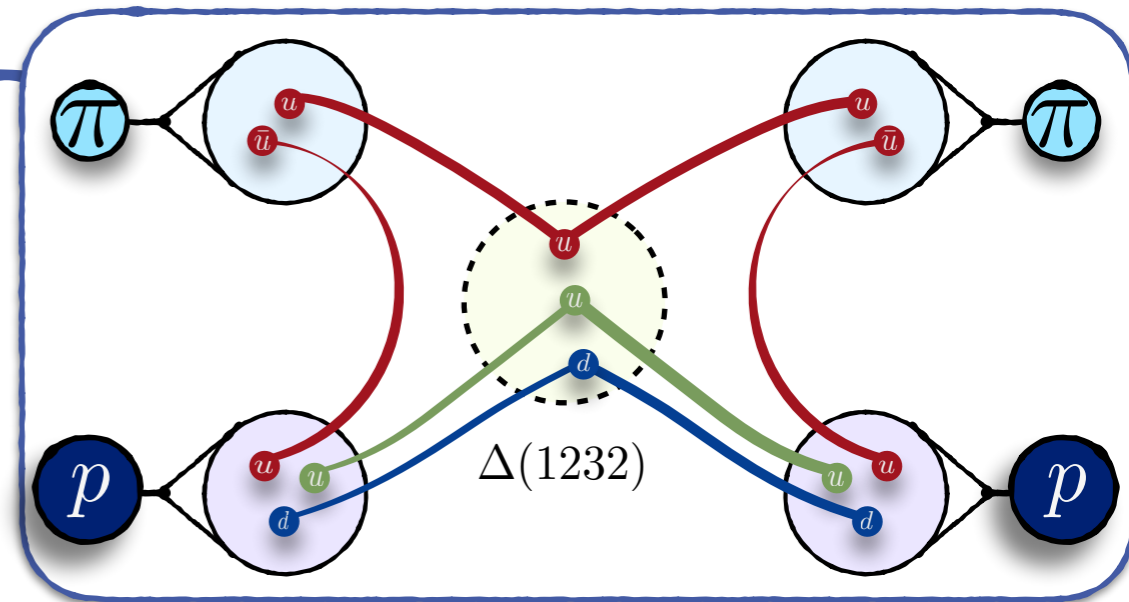


three amplitude/form factors:

G_{M1} - magnetic dipole

G_{E2} - electric quadrupole

G_C - Coulomb quadrupole



Questions to answer:

📍 What is possible from lattice QCD regarding electro/photo-production?

📍 How can lattice and experiment compare to each other?

📍 What is the possibility/timescale for extending this to baryons?

📍 What does lattice “need” from experiment?

What does lattice “*need*” from experiment?

- Absolutely nothing:

- Lattice QCD is fully predictive

- No inputs or approximations needed - [except quark masses]

What does lattice “*need*” from experiment?

- Absolutely nothing:

 - Lattice QCD is fully predictive

 - No inputs or approximations needed - [except quark masses]

- Support:

 - Shortage of **people power!**

 - Need more people:

 - developing: formalism, code, amplitude analysis, etc.

 - join effort with JPAC for amplitude analysis...

 - performing calculations, analysis, etc.

 - optimizing code

 - Will need more computer capabilities

Questions to answer:

📍 What is possible from lattice QCD regarding electro/photo-production?

📍 How can lattice and experiment compare to each other?

📍 What is the possibility/timescale for extending this to baryons?

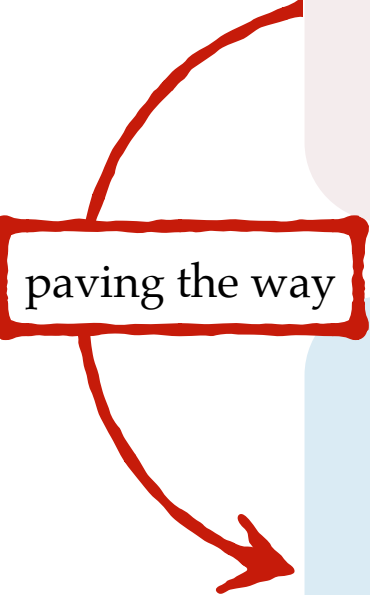
📍 What does lattice “*need*” from experiment?

Collective achievement / on-going efforts

Formal achievements:

Coupled-channels	2012, 2014
Meson electro / photo-production	2014, 2015
Elastic resonant form-factors	2015
Three-body - Sharpe, Hansen, RB	under way

paving the way



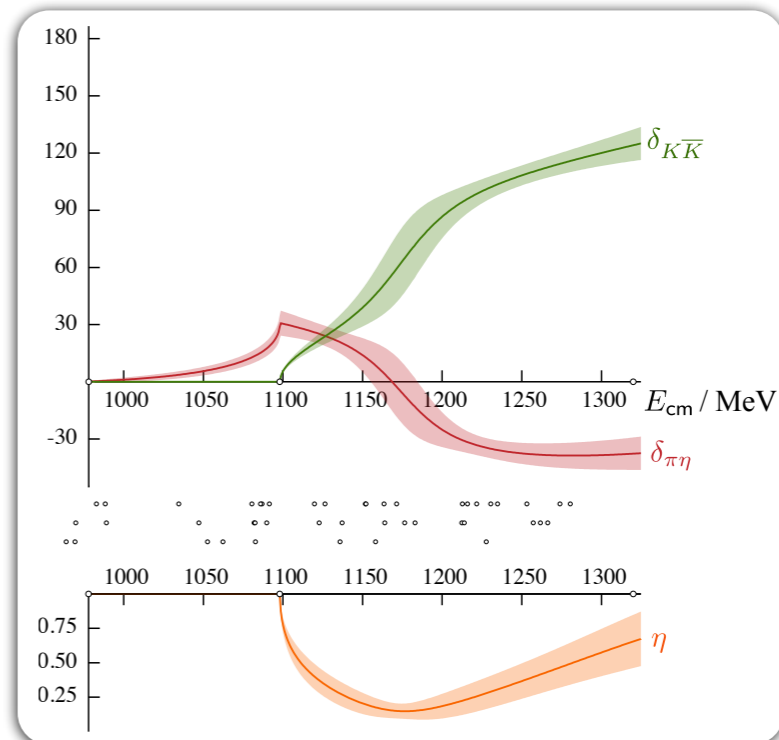
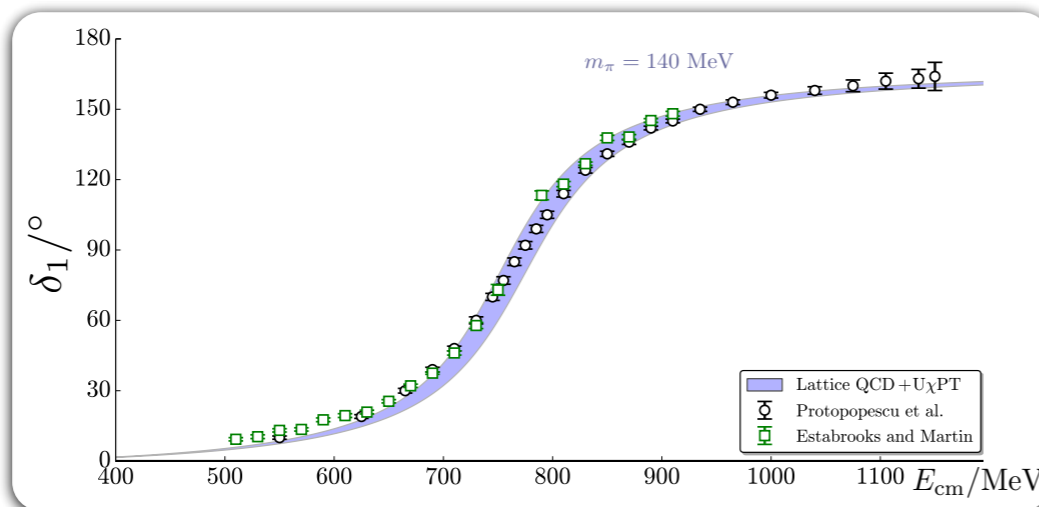
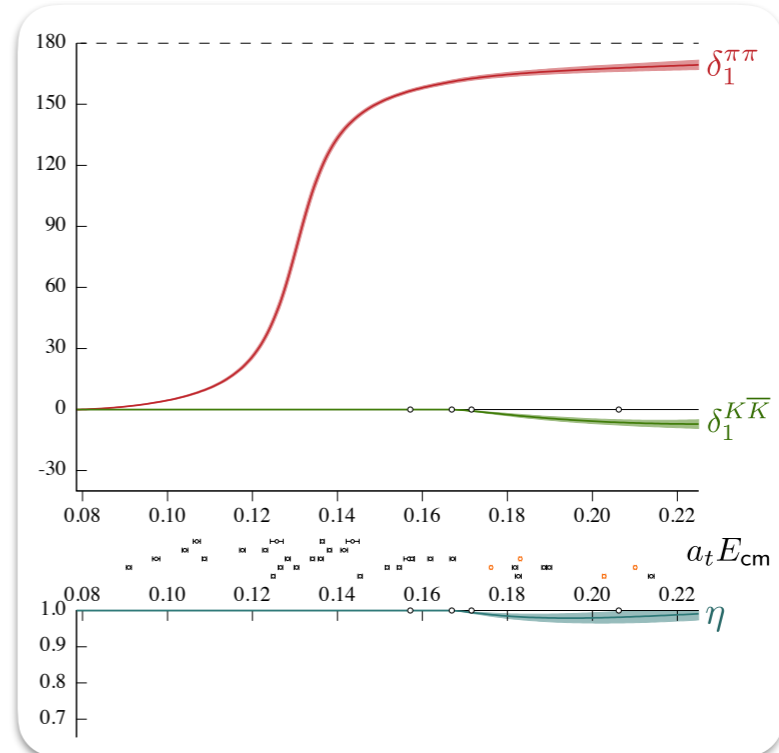
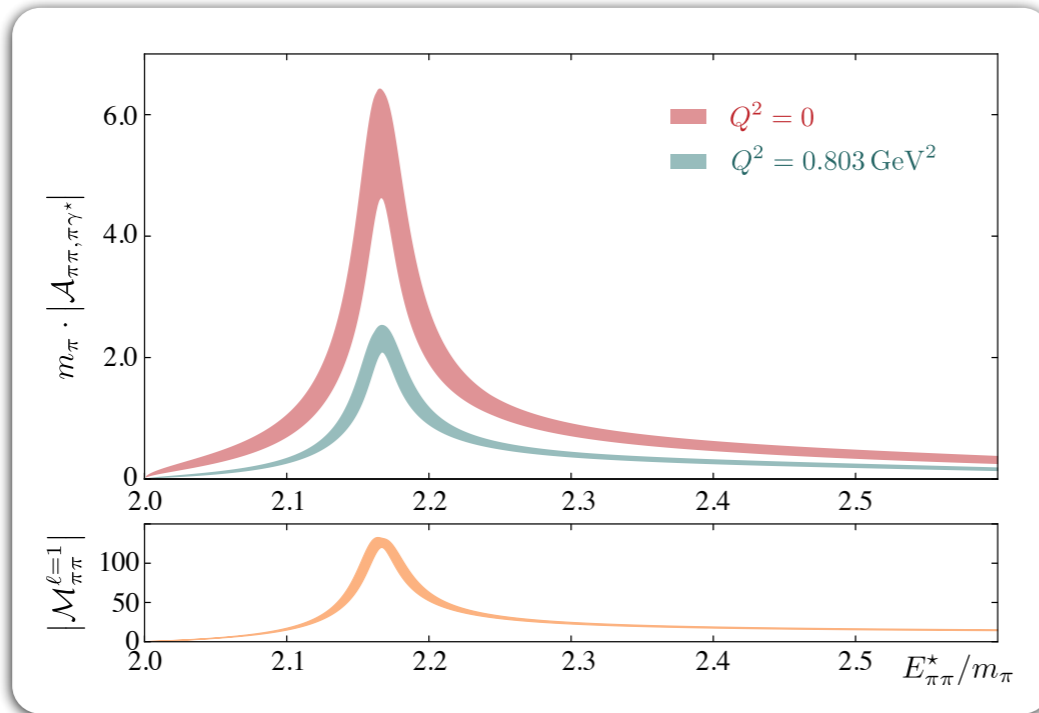
Numerical achievements:

First weakly coupled, two-channels	2014	RB
First resonant form-factor	2015	
First chiral extrapolation of resonant amplitude	2015	
First strongly coupled, two / three-channels	2016	RB
TOP SECRET	under way	

Code development - Robert

Analysis development [e.g., spin particles]- Jo & Dave Wilson

The big picture!



each one of these was 'impossible' 3 years ago...

Collaborators

formalism



Hansen

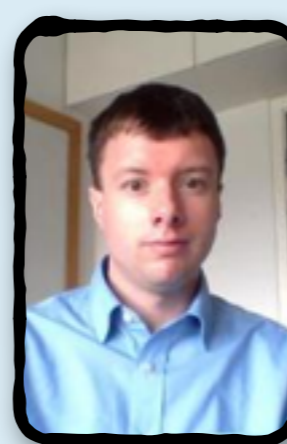


Sharpe

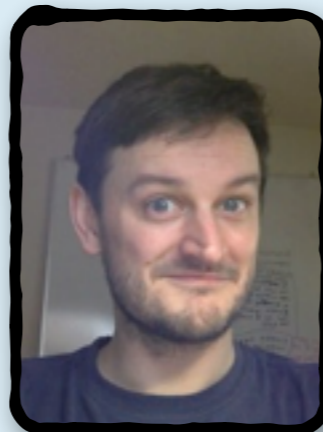
numerics



Wilson



Thomas



Dudek



Edwards

software



Joo



Winter

**HadSpec
Collaboration**

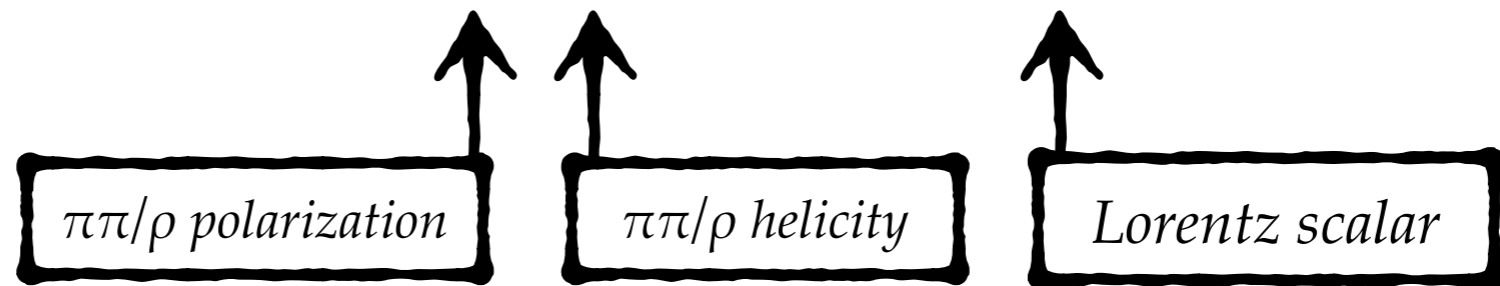
Back-up slides

$\pi\gamma^*$ -to- $\pi\pi$

(more details)

Lorentz decomposition:

$$\mathcal{H}_{\pi\pi,\pi\gamma^*}^\mu = \epsilon^{\mu\nu\alpha\beta} P_{\pi,\nu} P_{\pi\pi,\alpha} \epsilon_\beta(\lambda_{\pi\pi}, \mathbf{P}_{\pi\pi}) \frac{2}{m_\pi} \mathcal{A}_{\pi\pi,\pi\gamma^*}$$

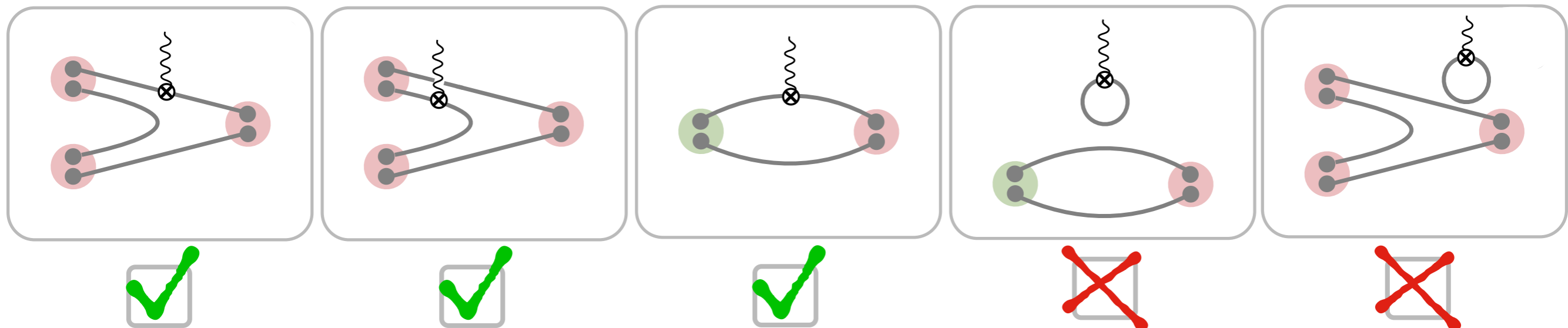


Approximations:

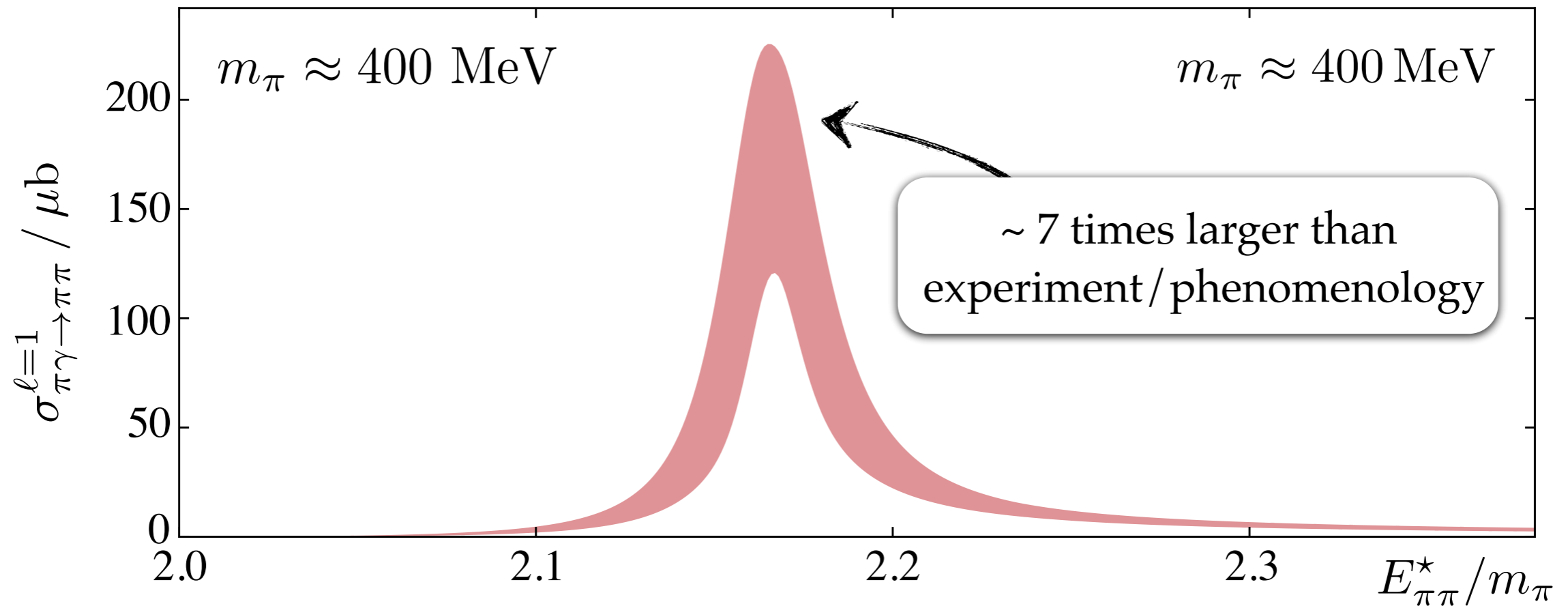
F-wave $\pi\gamma^*$ -to- $\pi\pi$ is ignored

kinematically and dynamically suppressed

contractions:

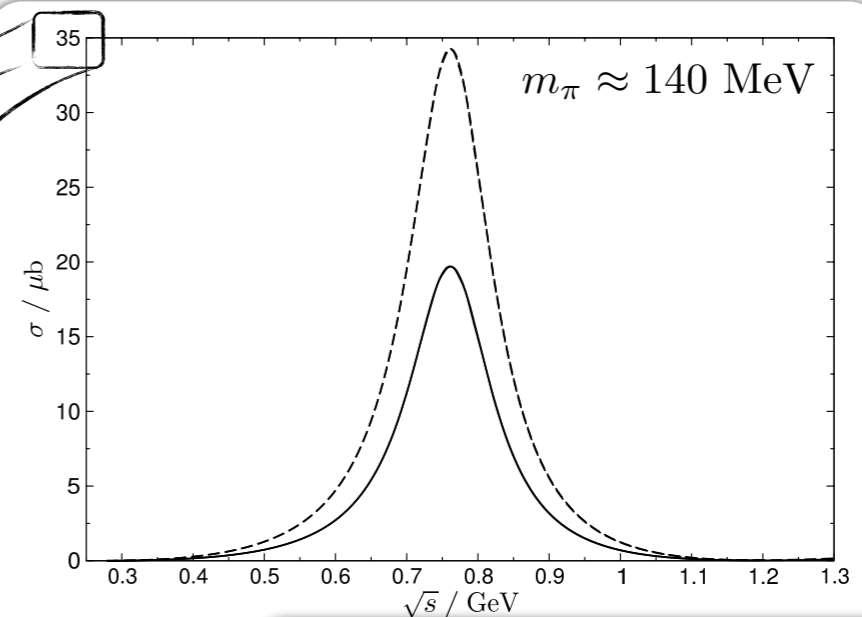


$\pi\gamma$ -to- $\pi\pi$ cross section

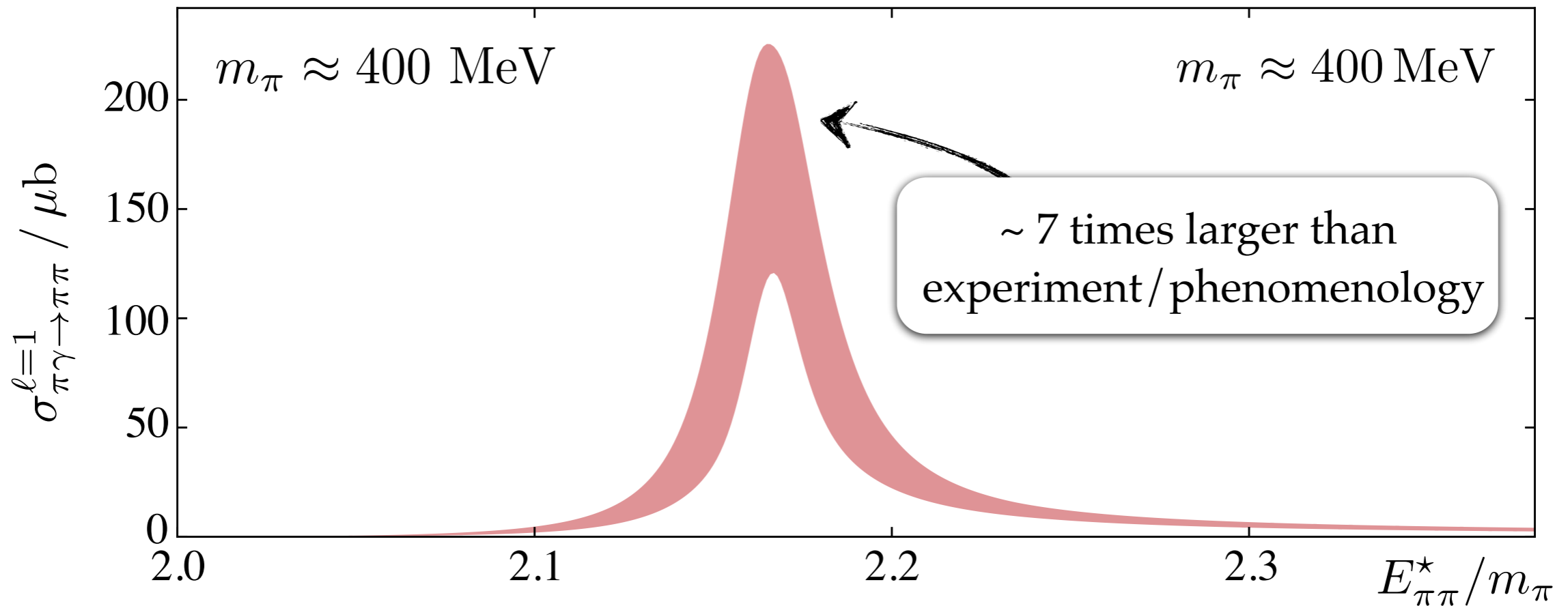


non trivial quark-mass dependence!

35



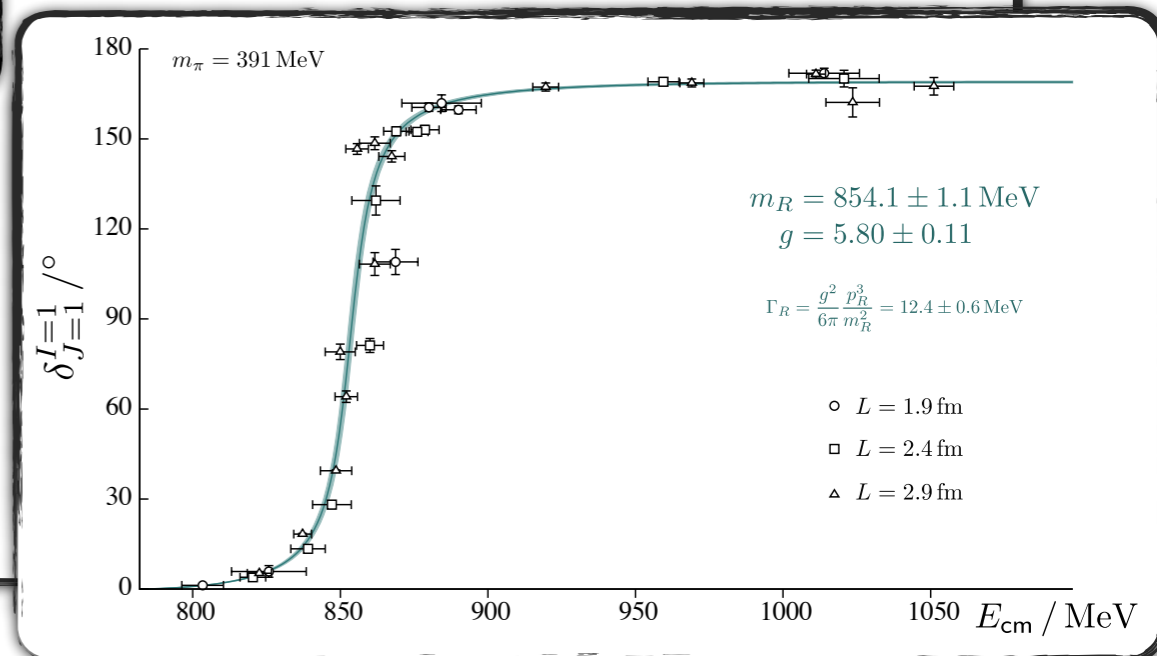
$\pi\gamma$ -to- $\pi\pi$ cross section



$$\lim_{E_{\pi\pi}^* \rightarrow m_\rho} \sigma(\pi^+ \gamma \rightarrow \pi^+ \pi^0) \propto \frac{q_{\pi\gamma}^* F_{\pi\rho}^2(m_\rho, 0)}{m_\pi^2} \times \frac{1}{\Gamma_1(m_\rho)}$$

0.60 x (physical)

12 x (physical)

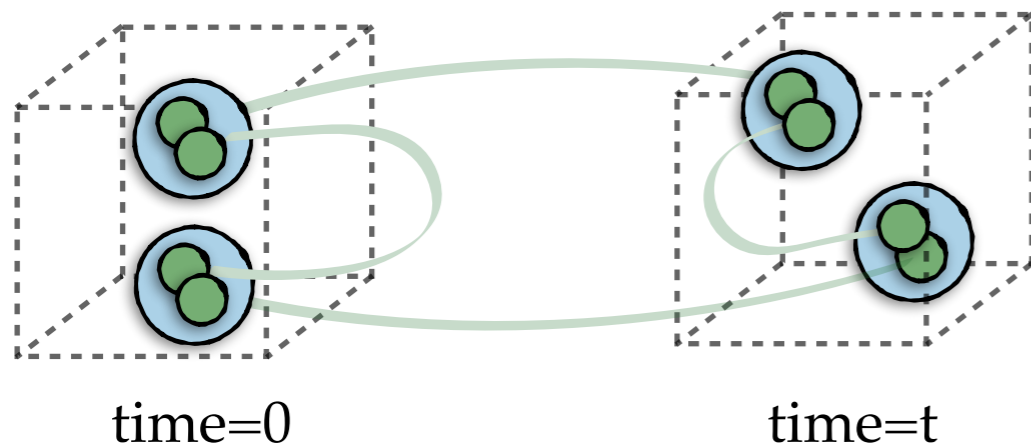
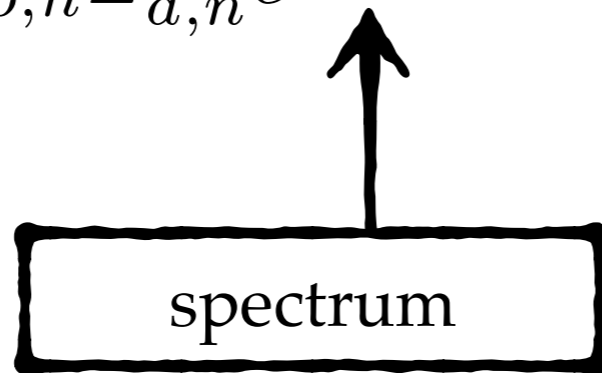


On determining correlation function
using small basis of operators

Extracting the spectrum

Two-point correlation functions:

$$\begin{aligned} C_{ab}^{2pt.}(t, \mathbf{P}) &\equiv \langle 0 | \mathcal{O}_b(t, \mathbf{P}) \mathcal{O}_a^\dagger(0, -\mathbf{P}) | 0 \rangle \\ &= \sum_n \langle 0 | \mathcal{O}_b(t, \mathbf{P}) | n, L \rangle \langle n, L | \mathcal{O}_a^\dagger(0, -\mathbf{P}) | 0 \rangle \\ &= \sum_n \langle 0 | e^{t\hat{H}_{QCD}} \mathcal{O}_b(0, \mathbf{P}) e^{-t\hat{H}_{QCD}} | n, L \rangle \langle n, L | \mathcal{O}_a^\dagger(0, -\mathbf{P}) | 0 \rangle \\ &= \sum_n Z_{b,n} Z_{a,n}^\dagger e^{-E_n t} \end{aligned}$$

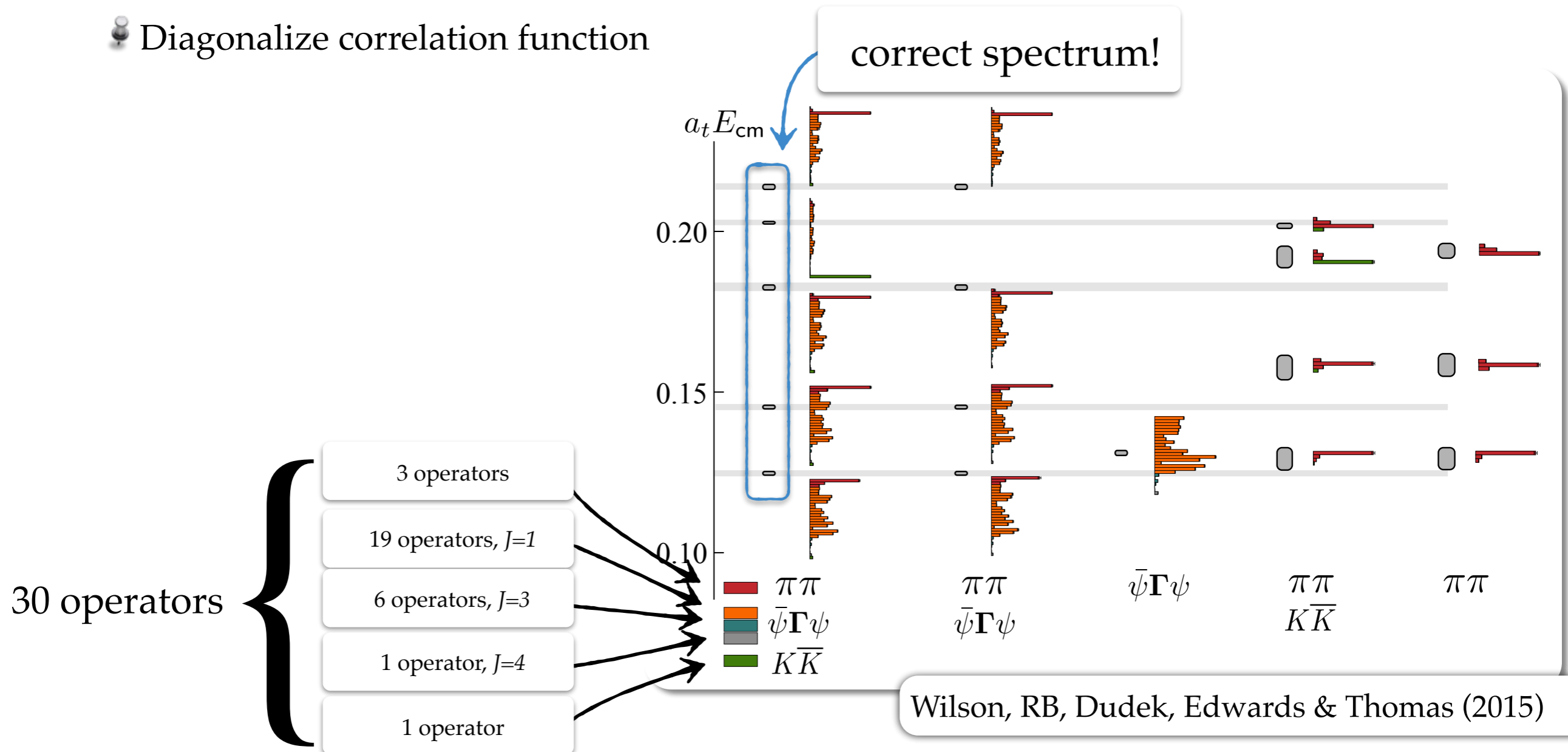


Extracting the spectrum

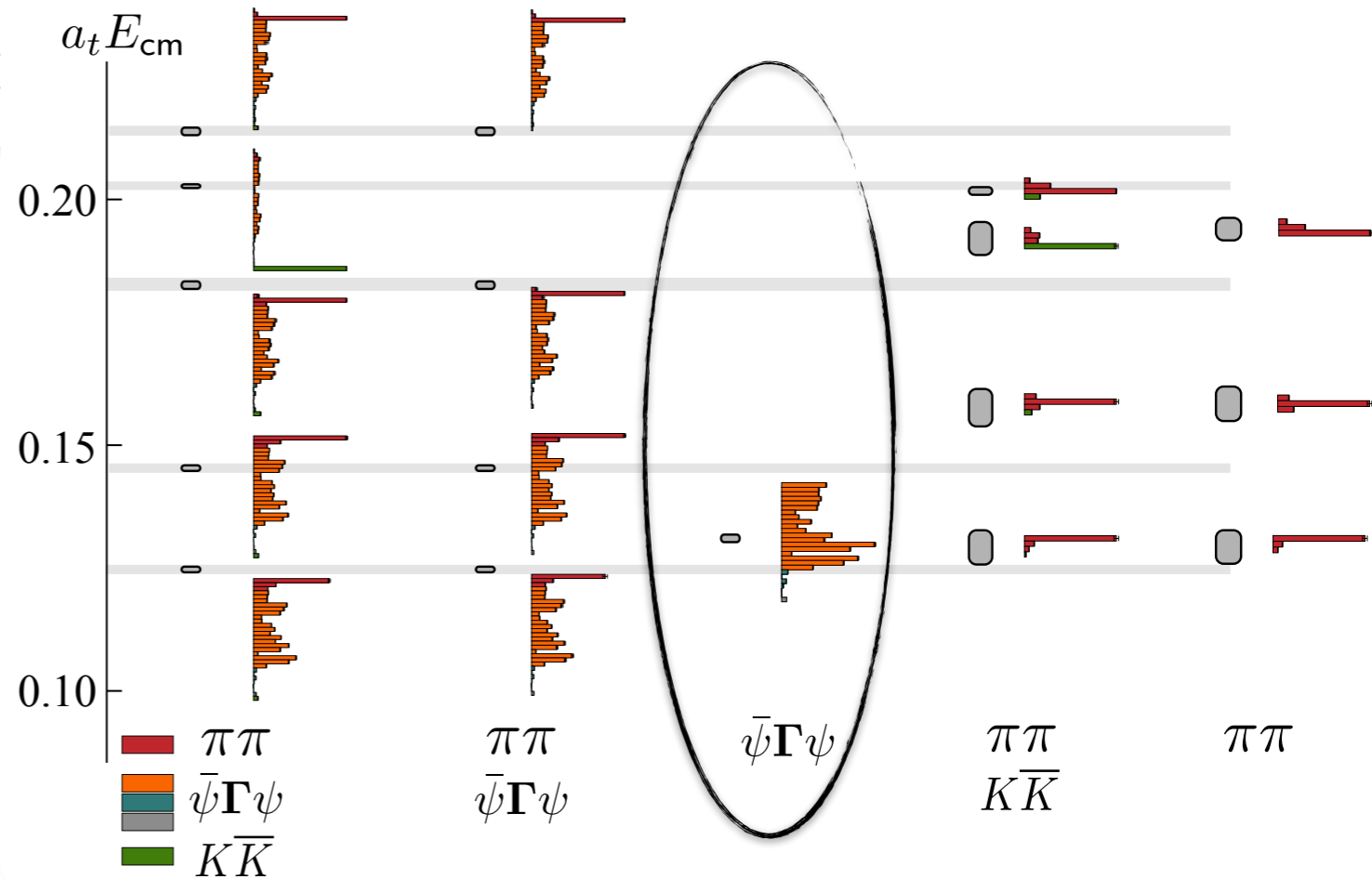
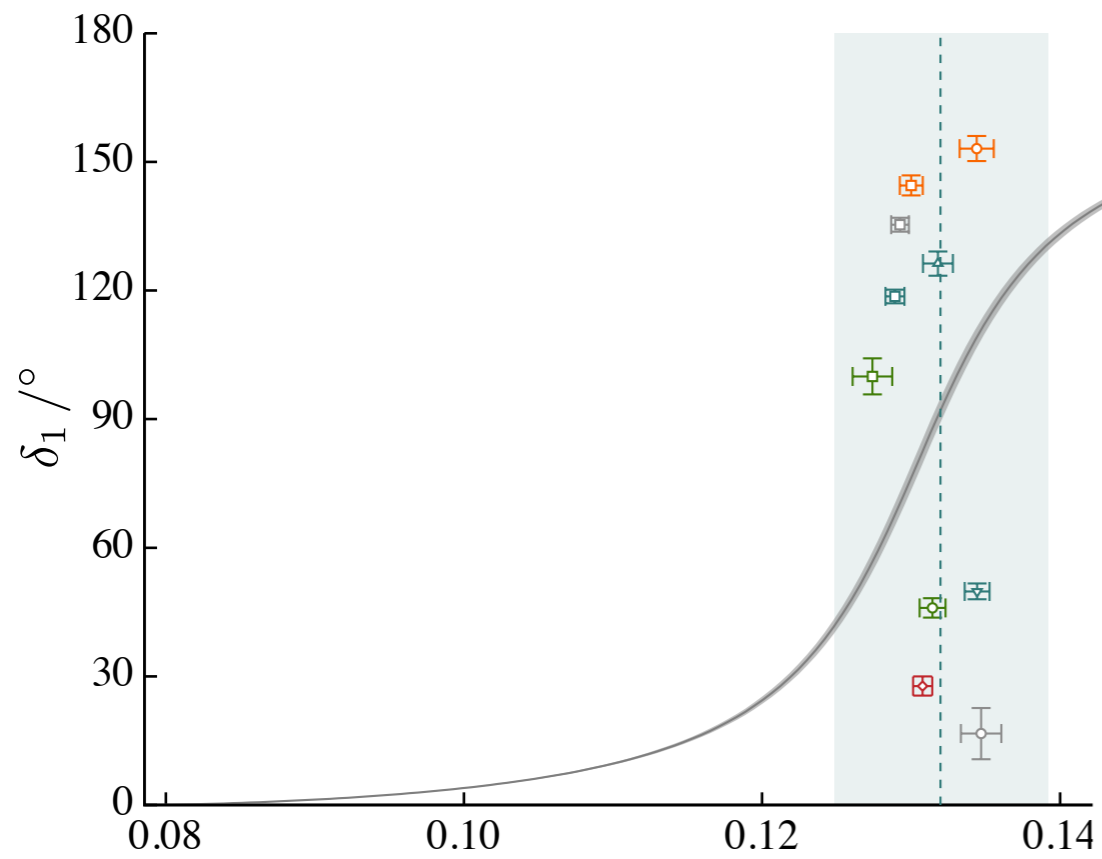
Two-point correlation functions:

$$C_{ab}^{2pt.}(t, \mathbf{P}) \equiv \langle 0 | \mathcal{O}_b(t, \mathbf{P}) \mathcal{O}_a^\dagger(0, -\mathbf{P}) | 0 \rangle = \sum_n Z_{b,n} Z_{a,n}^\dagger e^{-E_n t}$$

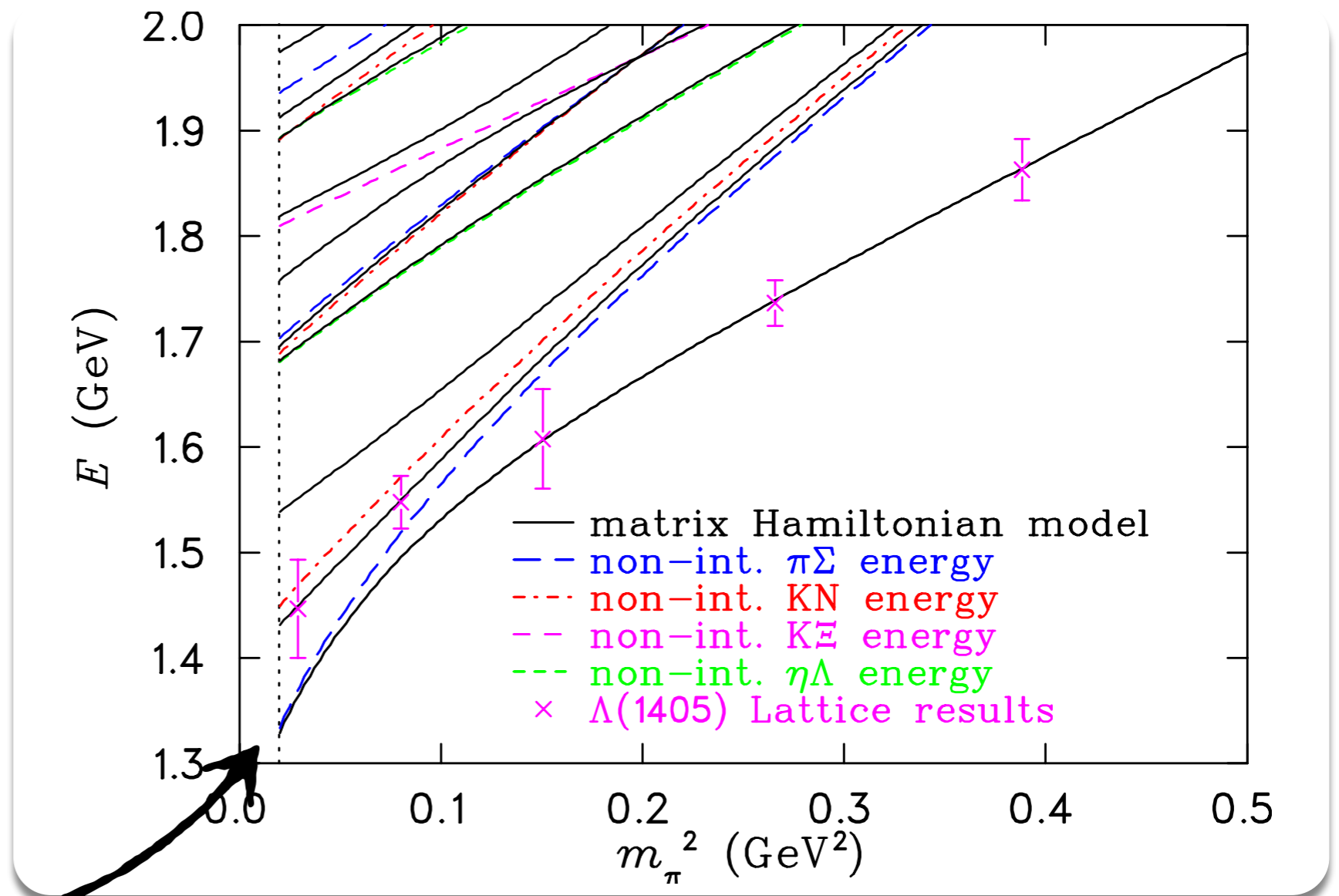
- Use a large basis of operators with the same quantum numbers
- Diagonalize correlation function



The incorrect answer

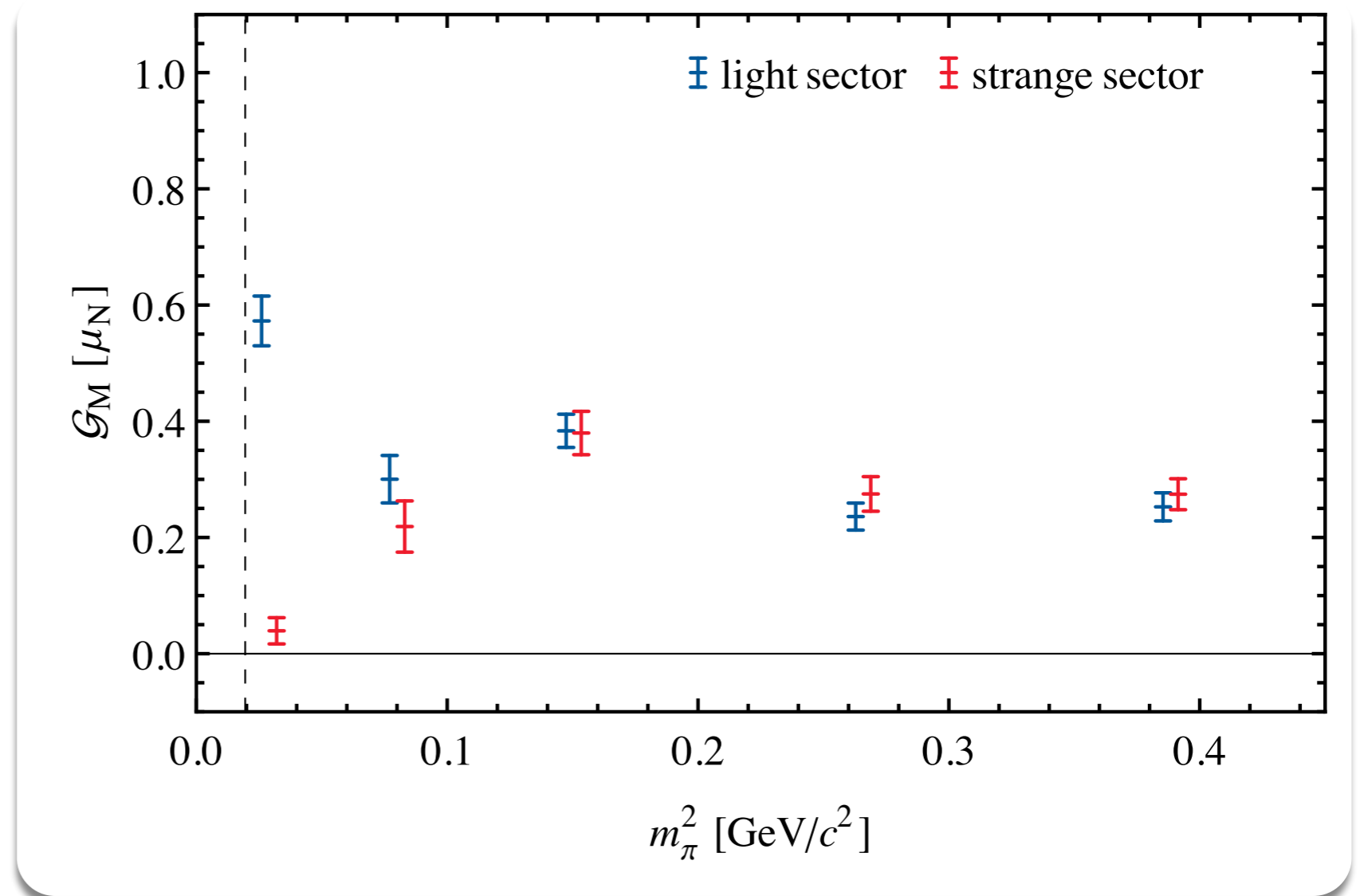


On the $\Lambda(1405)$ spectrum



- Noisy spectrum
- missing ground state
- found only one pole?
- claim resonance is a KN molecule, but did not use KN operators?
- where's the width of the $\Lambda(1405)$?

On the $\Lambda(1405)$ form factor



- claim resonance is a KN molecule, but did not use KN operators?
- claim: finite volume matrix element of QED current equal infinite volume form factor
- only true for stable states
- Ignore coupling with $\Sigma\pi$ in analysis of form factor
- for unstable states, the matrix element is **not proportional** to infinite volume form factor

Two-body maitre elements

$$|\langle \mathbf{2} | \mathcal{J} | \mathbf{2} \rangle_L| = \frac{1}{\sqrt{L^3}} \sqrt{\text{Tr} [\mathcal{R} \mathcal{W}_{L,\text{df}} \mathcal{R} \mathcal{W}_{L,\text{df}}]}$$

$\langle \mathbf{2} | \mathcal{J} | \mathbf{2} \rangle_L =$ FV matrix element

$$\mathcal{R} = \mathcal{R}(P, L, \mathcal{M}, \frac{\partial \mathcal{M}}{\partial E_2})$$

$$\mathcal{W}_{L,\text{df}} = \mathcal{W}_{\text{df}} + \mathcal{M} [G(L) \cdot w] \mathcal{M}$$

$G(L) =$ FV function

$w =$ single/stable particle form factor, e.g., N, π, \bar{K}, \dots

$\mathcal{W} =$ infinite volume $2 + \mathcal{J} \rightarrow 2$ amplitude

this was reported,
everything else was not

$$i\mathcal{W} = i\mathcal{W}_{\text{df}} + \left[\text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} \right]$$

form factors are defined
inside the residues of
this amplitude

Hamiltonian Effective Field Theory Study of the $N^*(1535)$ Resonance in Lattice QCD

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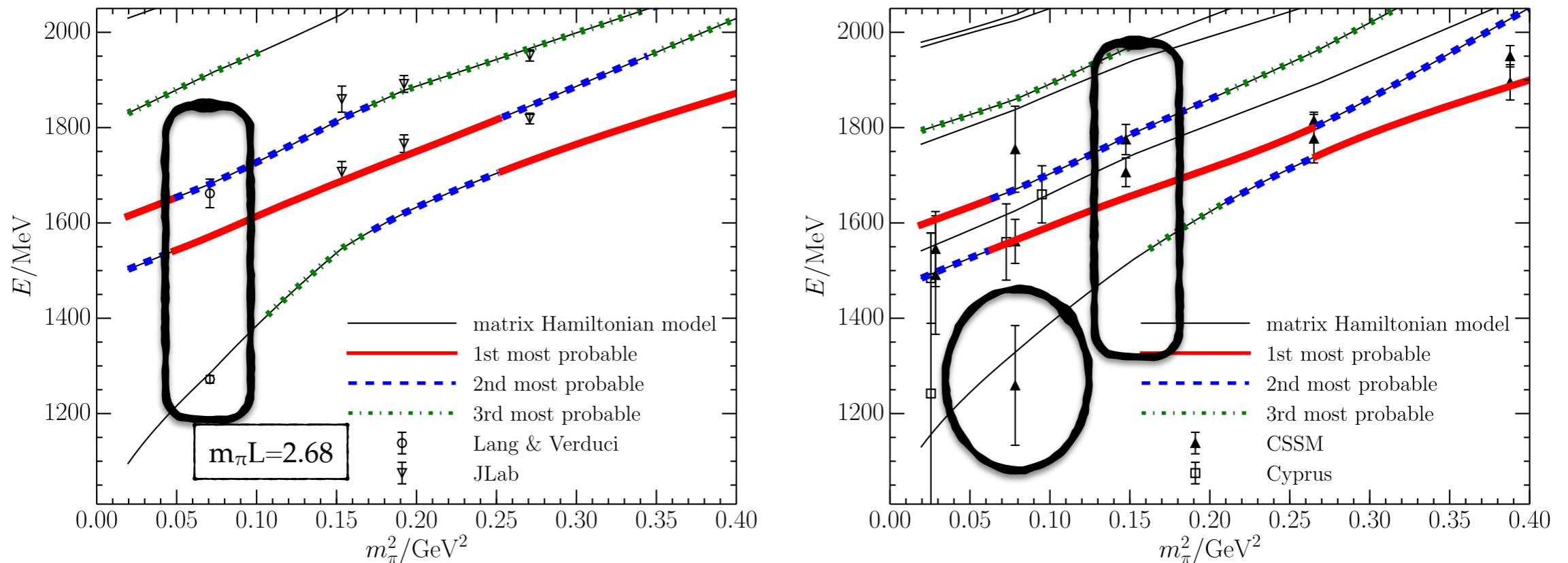


FIG. 2. The pion mass dependence of the $L \approx 1.98$ fm (left) and $L \approx 2.90$ fm (right) finite-volume energy eigenstates. The different line types and colors indicate the strength of the bare basis state in the Hamiltonian model eigenvector.


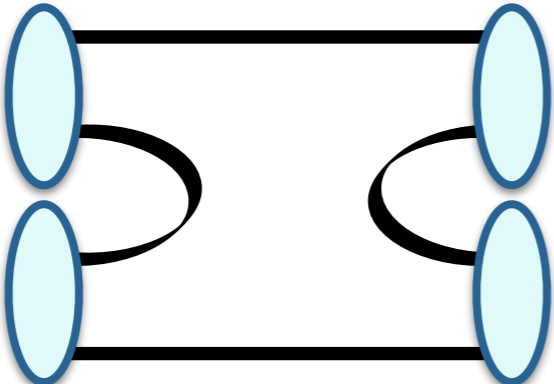
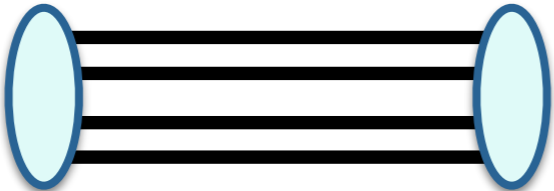
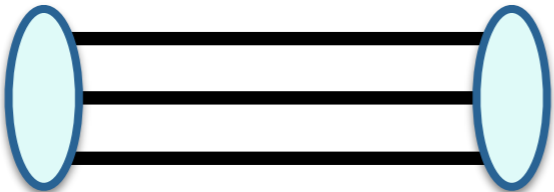
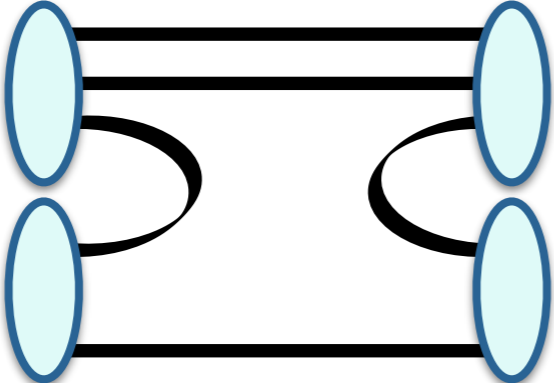
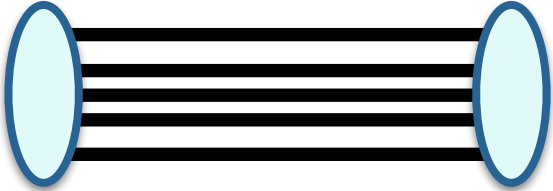
- Noisy spectrum
- missing states
- all, except Lang & Verduci, ignore $N\pi$, $N\pi\pi$, $N\eta$ operators
- requires experimental input to go from $\chi^2/N_{\text{dof}}=4.6$ to 1.7

Contraction cost

Estimates by Robert

Cost \propto Volume

assuming dense tensors

meson:		$\text{tr}(MPMP) \sim N^3$	"easy"
meson-meson:		$\sim 100 \times N^3$	affordable
tetraquark:		$\sim N^5$	sparse in spin, but not really affordable
baryon:		$\sim N^4$	affordable
baryon-meson:		$\sim 100 \times N^4$	to do a decent job, not really affordable
pentaquark:		$\sim N^6$	prohibitive