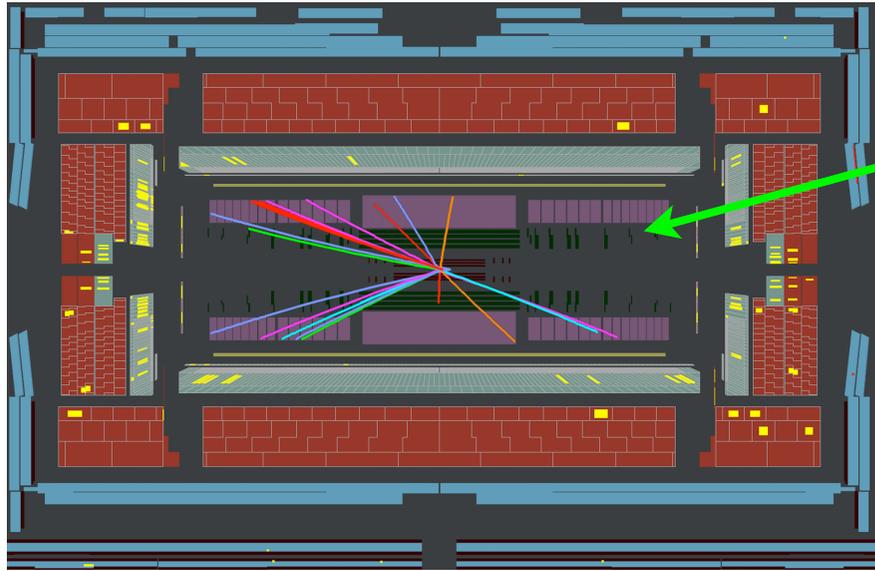


Precision light and heavy meson physics from Lattice QCD

Christine Davies
University of Glasgow
HPQCD collaboration

Jefferson Lab
Feb 2017

QCD is a key part of the Standard Model but quark confinement is a complication/interesting feature.



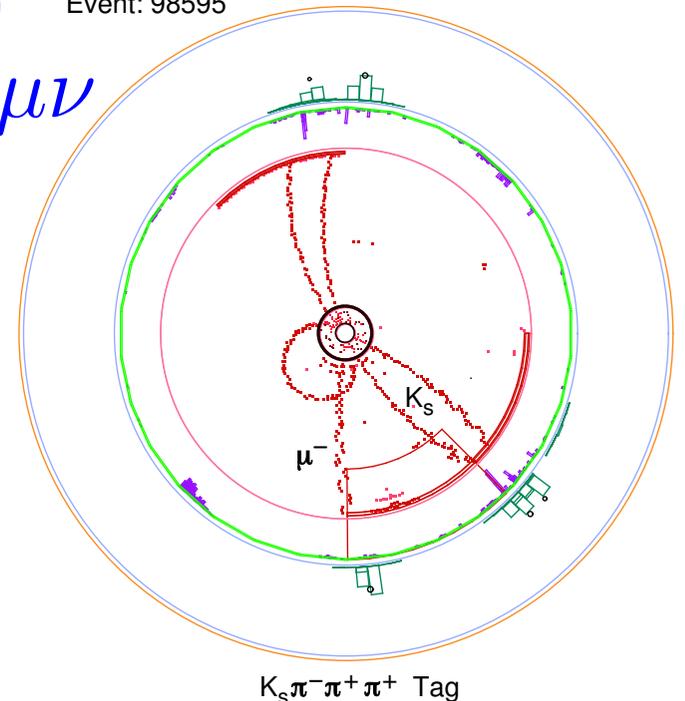
ATLAS@LHC

Connecting observed hadron properties to those of quarks requires full nonperturbative treatment of Quantum Chromodynamics.

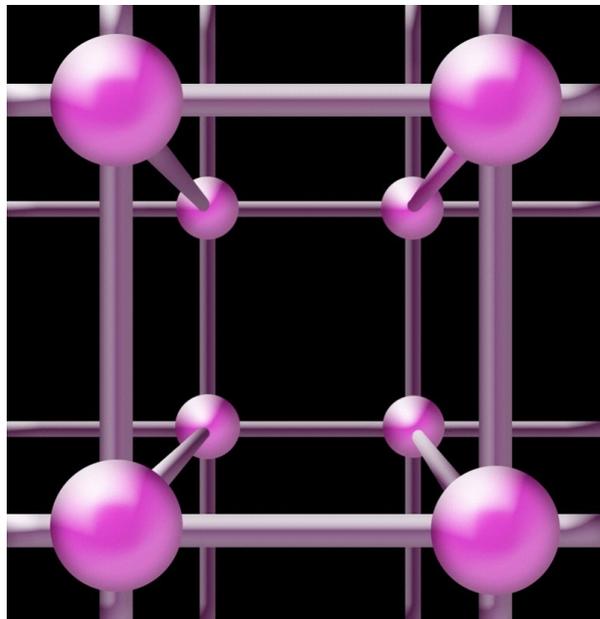
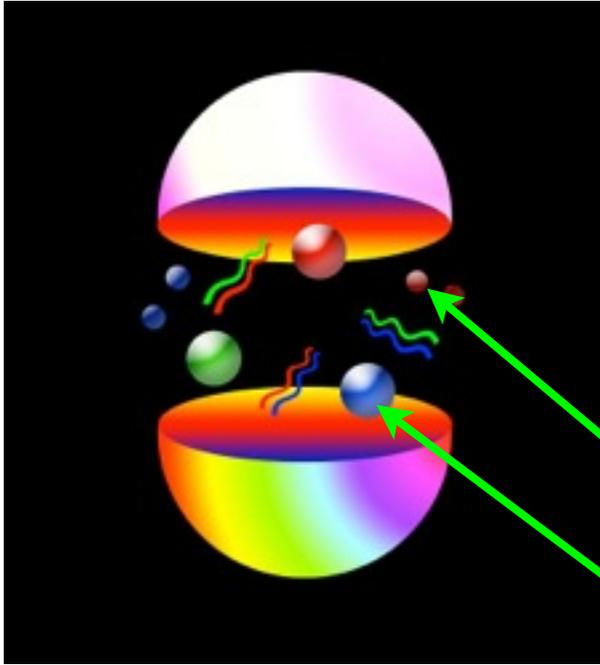
Some properties of hadrons can be very accurately measured and are calculable in lattice QCD - can test SM and determine parameters very accurately (1%).

CLEO
 $D \rightarrow \mu\nu$

Event: 98595



$K_s \pi^- \pi^+$ Tag



a

Lattice QCD: fields defined on 4-d discrete space-(Euclidean) time.

Lagrangian parameters: $\alpha_s, m_q a$

- 1) Generate sets of gluon fields for Monte Carlo integrn of Path Integral (inc effect of u, d, s, (c) sea quarks)
- 2) Calculate valence quark propagators and combine for “hadron correlators” . Fit for hadron masses and amplitudes

- Determine a to convert results in lattice units to physical units. Fix m_q from hadron mass

numerically extremely challenging

- cost increases as $a \rightarrow 0, m_u/d \rightarrow \text{phys}$ and with statistics, volume.

Darwin@Cambridge,
part of UK's £15m HPC facility
for theoretical particle physics
and astronomy



State-of-the-art commodity
cluster: 9600 Intel Sandybridge
cores, infiniband interconnect,
fast switch and 2 Pbytes storage

DiRAC

www.dirac.ac.uk

Allows us to calculate
quark propagators
rapidly and store them
for flexible re-use.



Quark formalisms

Many ways to discretise Dirac Lagrangian onto lattice.
All should give same answers.

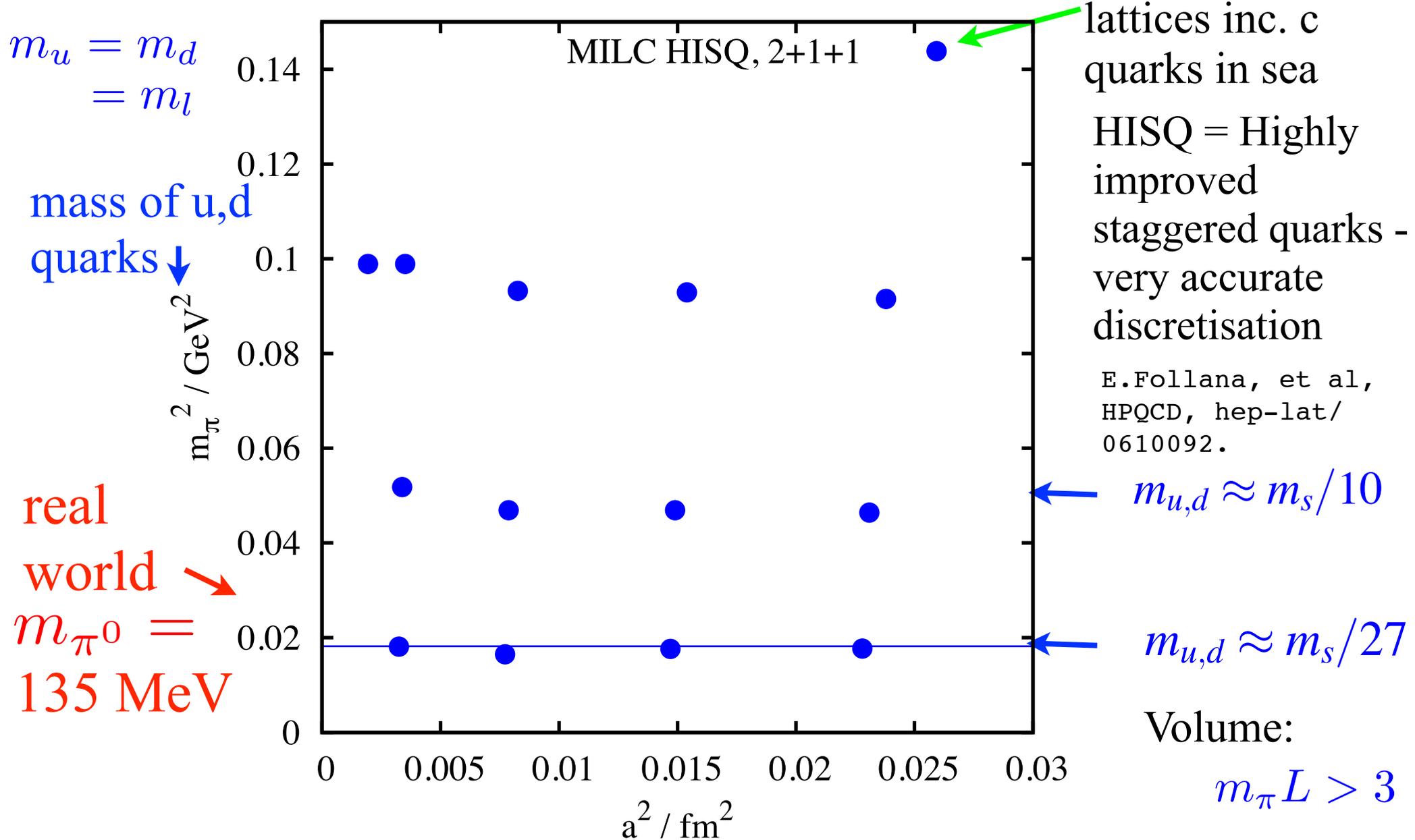
Issues are: Discretisation errors at power a^n
Numerical speed of matrix inversion
Chiral symmetry
Quark doubling

We use Highly Improved Staggered Quarks (HISQ) for u, d, s and c. Also (with extrapolation) for b.

Disc. errors $\alpha_s^2 a^2, a^4$. Numerically fast. Chiral symm.
Some complications from doublers ('tastes').

For b quarks, also use NRQCD. Fast for b quarks at moderate a, disc. errors now $\alpha_s^2 a^2$. Current normln perturbative, through α_s

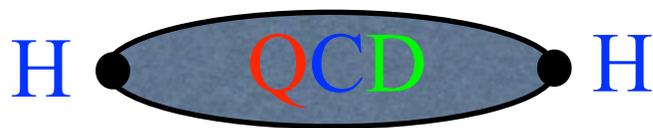
Example parameters for ‘2nd generation’ calculations now being done with staggered quarks.



other formalisms: clover (Fermilab), domain-wall, twisted-mass ...

Hadron correlation functions ('2point functions') give masses and decay constants.

$$\langle 0 | H^\dagger(T) H(0) | 0 \rangle = \sum_n A_n e^{-m_n T} \xrightarrow{T \text{ large}} A_0 e^{-m_0 T}$$

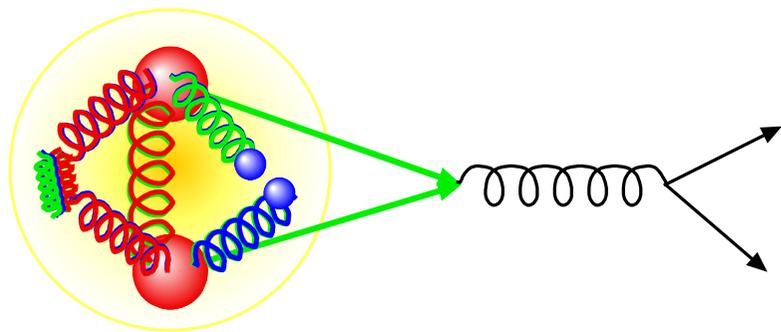


masses of all hadrons with quantum numbers of H

$$A_n = \frac{|\langle 0 | H | n \rangle|^2}{2m_n} = \frac{f_n^2 m_n}{2}$$

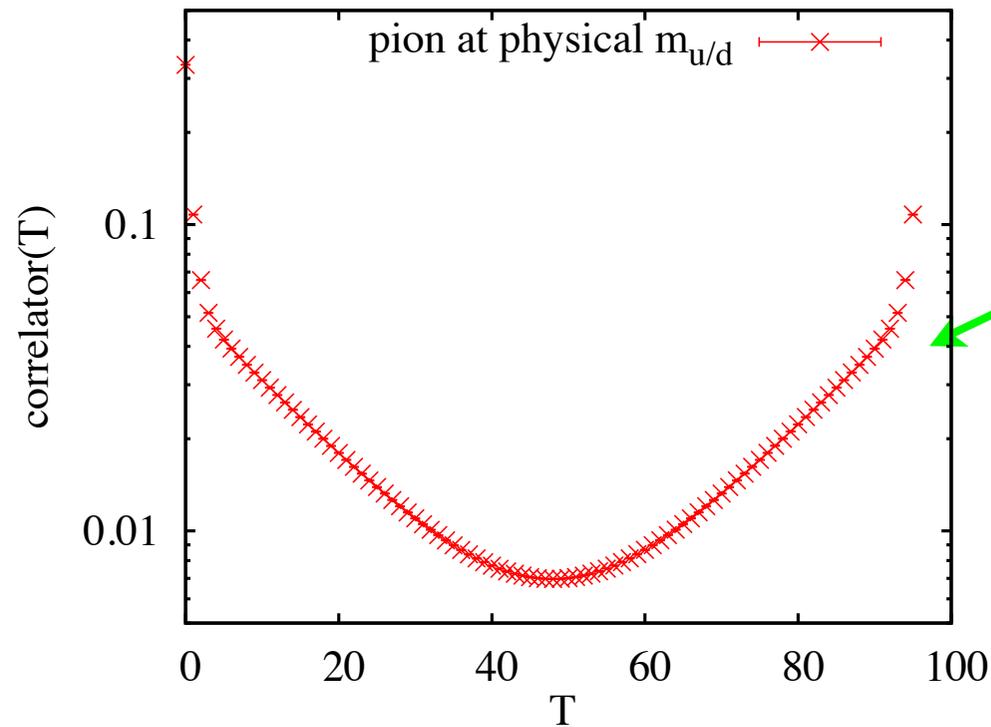
decay constant parameterises amplitude to annihilate - a property of the meson calculable in QCD. Can often relate to experimental decay rate.

1% accurate experimental info. for f and m for many mesons!
Need accurate determination from lattice QCD to match



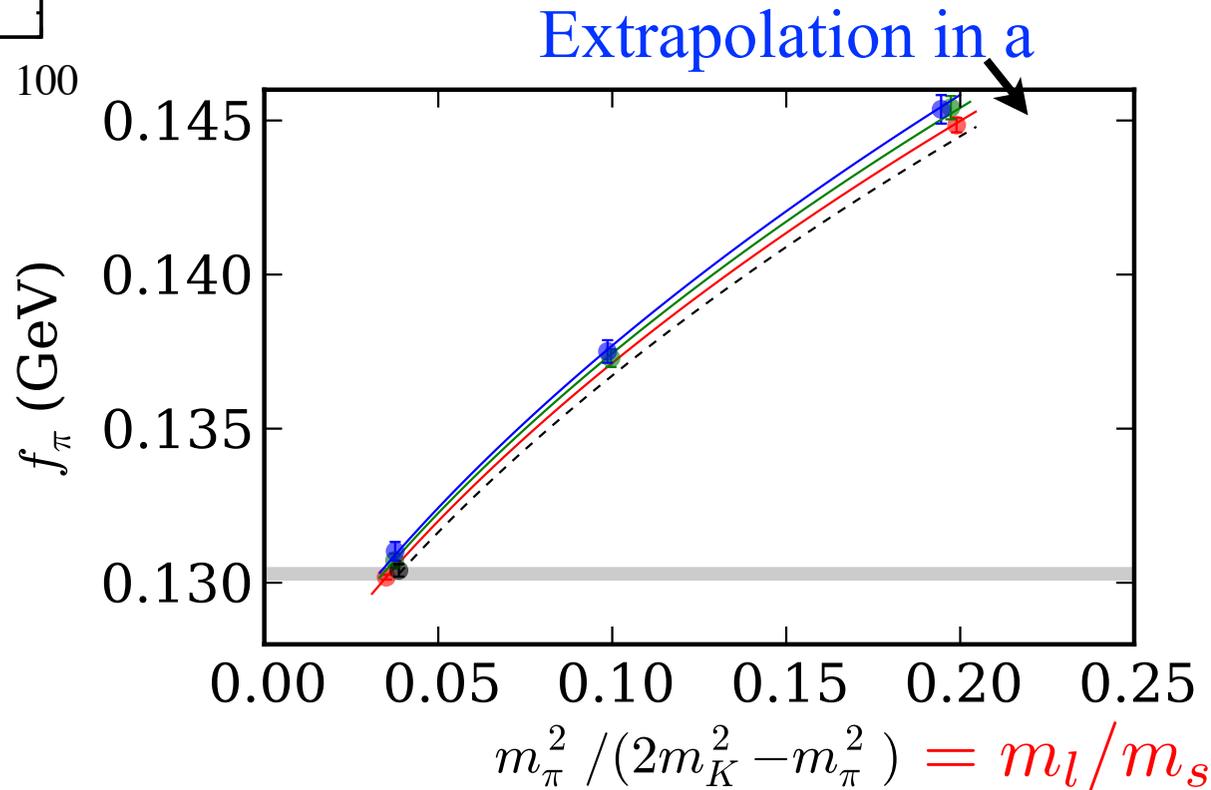
Example (state-of-the-art) calculation

R. Dowdall et al, HPQCD,
1303.1670.



Extract meson mass and amplitude=decay constant from correlator for multiple lattice spacings and $m_{u/d}$. Very high statistics (16,000 samples of the correlator)

Convert decay constant to GeV units using w_0 to fix relative lattice spacing. Very small discretisation errors. PCAC reln means no renormln factors needed.

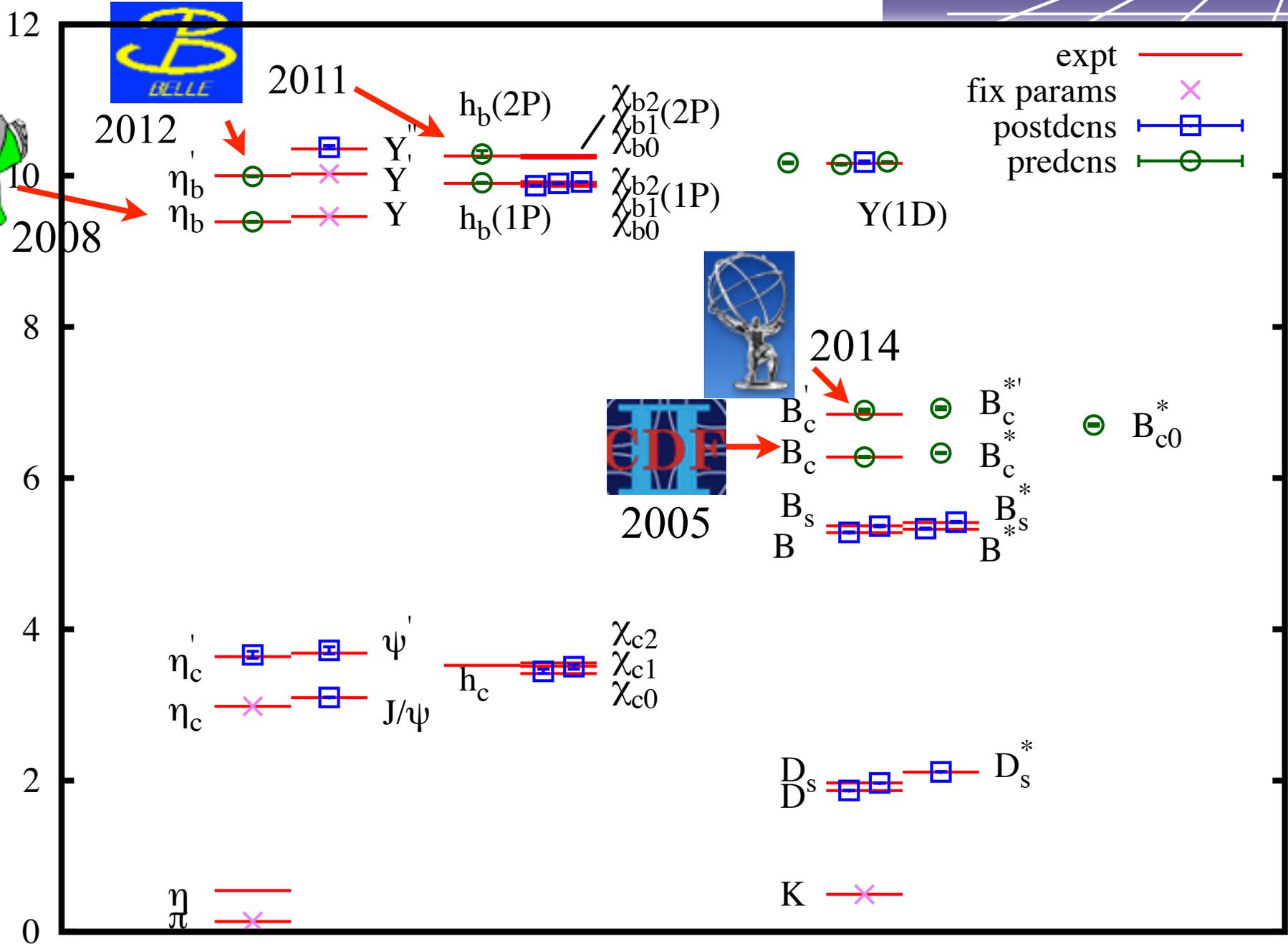


The gold-plated meson spectrum

HPQCD



MESON MASS (GeV/c^2)

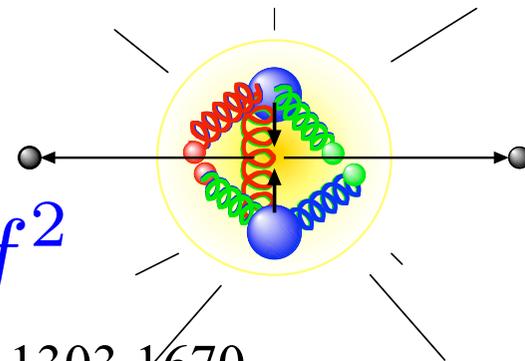


Uncerties at few MeV level : Future: inc QED and $m_u \neq m_d$

Meson decay constants

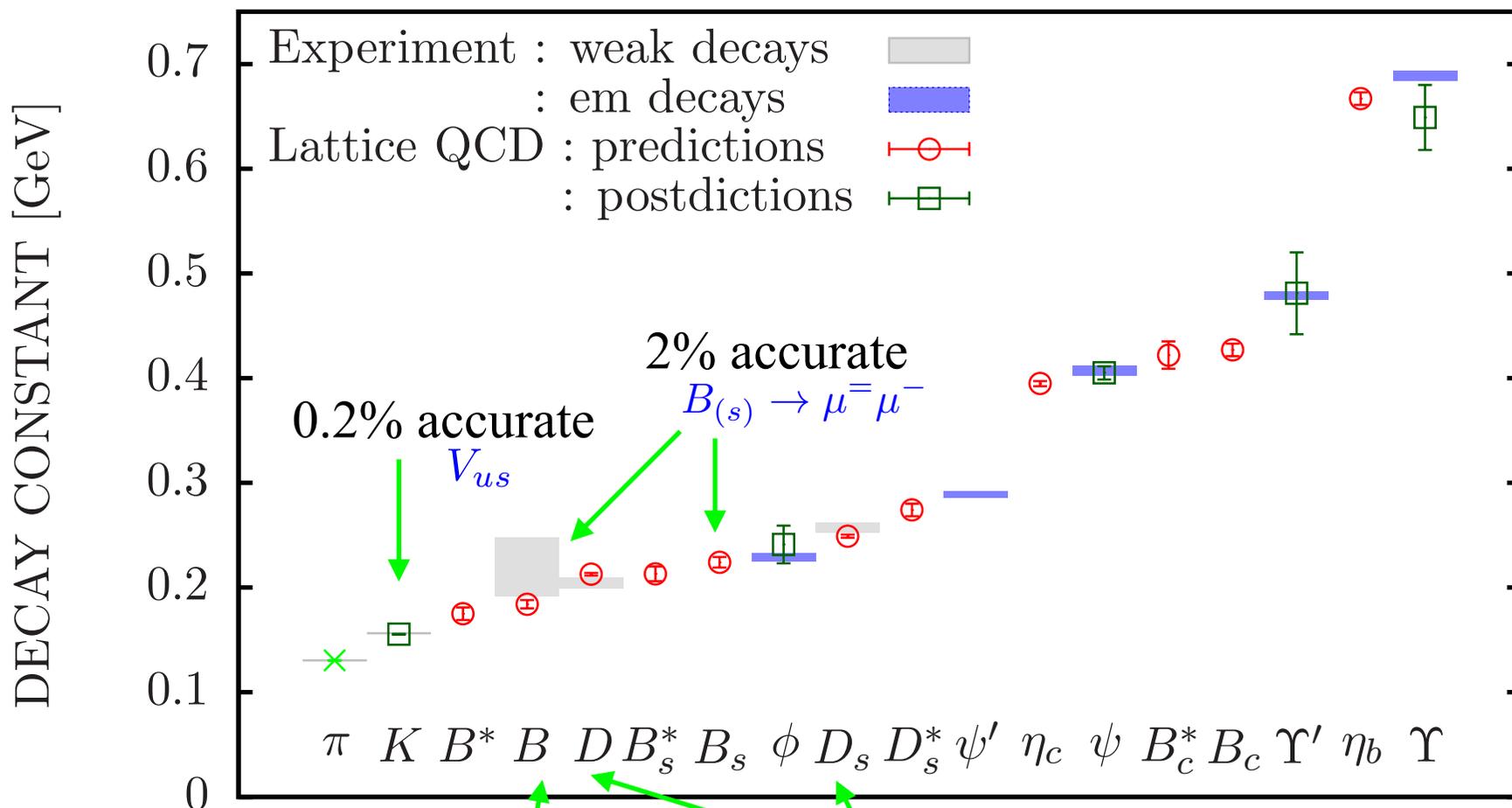
Parameterises hadronic information needed for annihilation rate to W or photon:

$$\Gamma \propto f^2$$



HPQCD

1503.05762, 1408.5768, 1302.2644, 1303.1670



decay constants of vector mesons now being pinned down



2012
 $B \rightarrow \tau \nu$
 V_{ub}

0.5% accuracy from lattice QCD
 now : **FNAL/MILC** 1407.3772
 BES will improve expt. V_{cd} V_{cs}

Vector meson decay constants HPQCD 1312.5264, 1503.05762

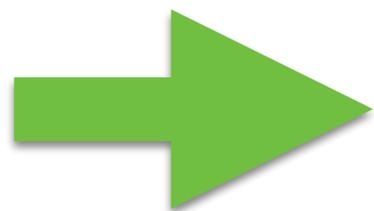
We calculate:

allows us to fix

$$f_{D_s^*} / f_{D_s} = 1.10(2) \quad \text{total width}$$

may be possible to see this?

$$\Gamma(D_s^* \rightarrow D_s \gamma) = 0.066(26) \text{keV}$$



$$\text{Br}(D_s^* \rightarrow \ell \nu) = 3.4(1.4) \times 10^{-5}$$

We also find:

$$f_{B_s^*} / f_{B_s} = 0.953(23)$$

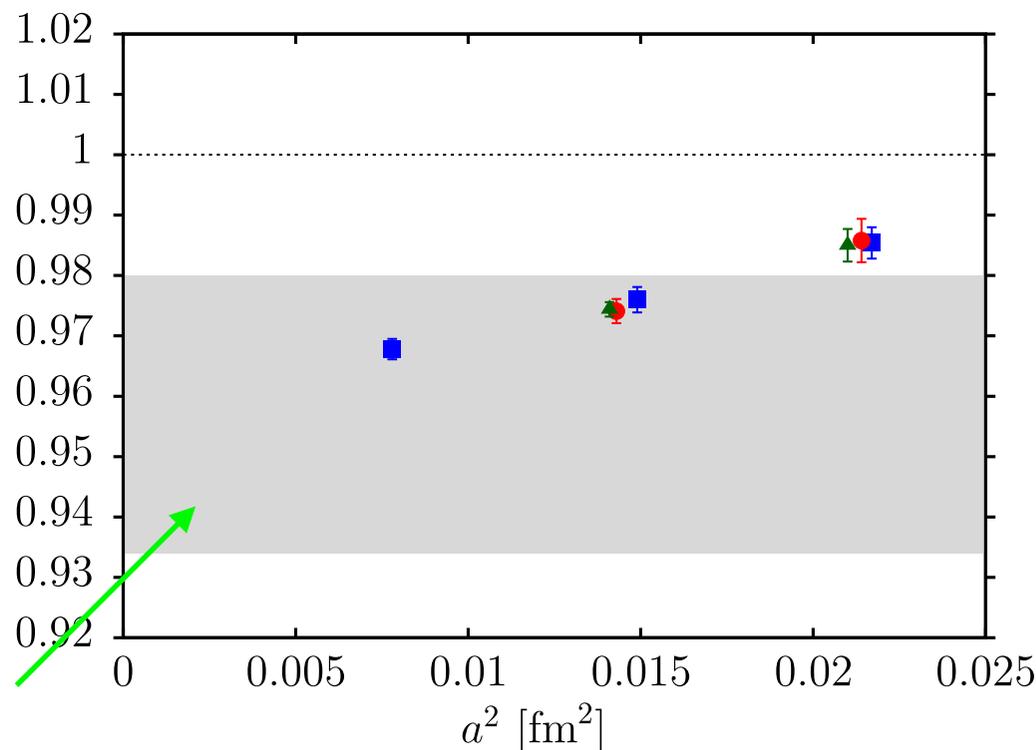
< 1

$$\frac{f_{B_s^*} \sqrt{M_{B_s^*}}}{f_{B_s} \sqrt{M_{B_s}}}$$

R_s

can test HQET

pert. matching
uncertainty dominates



Further probes of hadron structure: form factors as functions of q^2

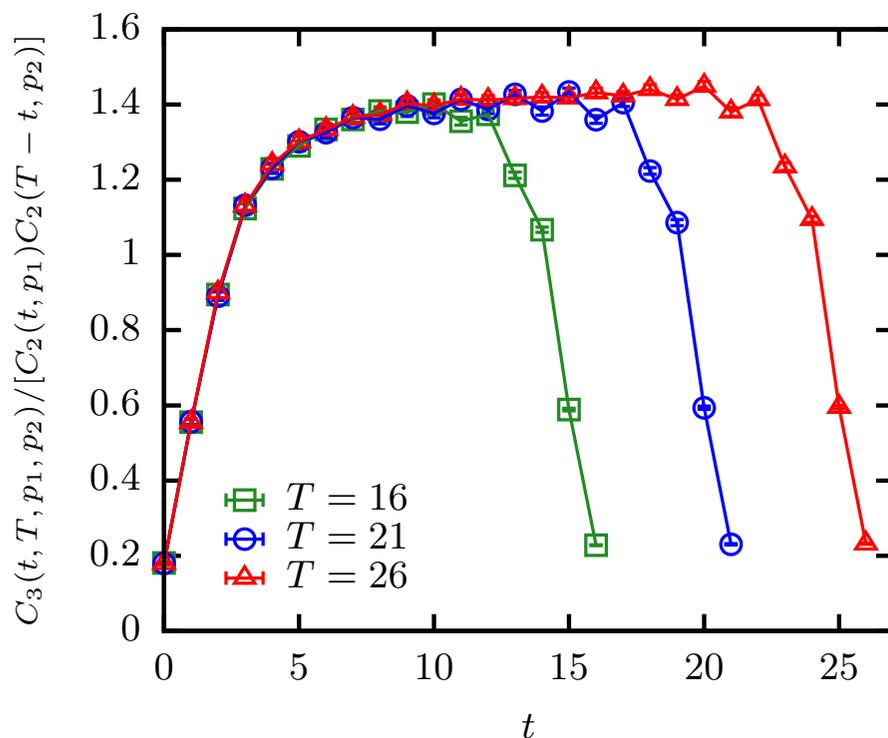
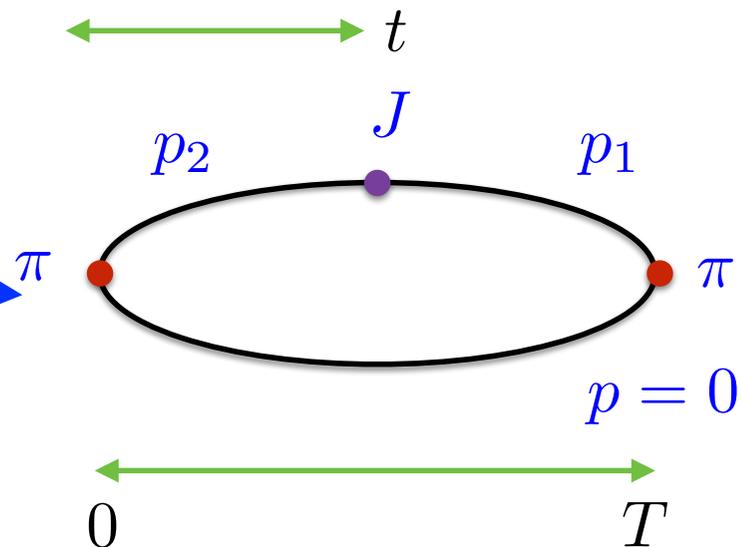
Needs '3-point' functions
e.g. for pion to pion transition
via vector current J

Need to calculate correlators
for multiple T values and
 $0 < t < T$ and fit as a function of
 t, T simultaneously with 2pt.

$$C_{3pt} = \sum_{i,j} b_i J_{ij} b_j e^{-E_i t} e^{-E_j (T-t)}$$

$\langle \pi | V_\mu | \pi \rangle / (2Z \sqrt{E_i E_j})$

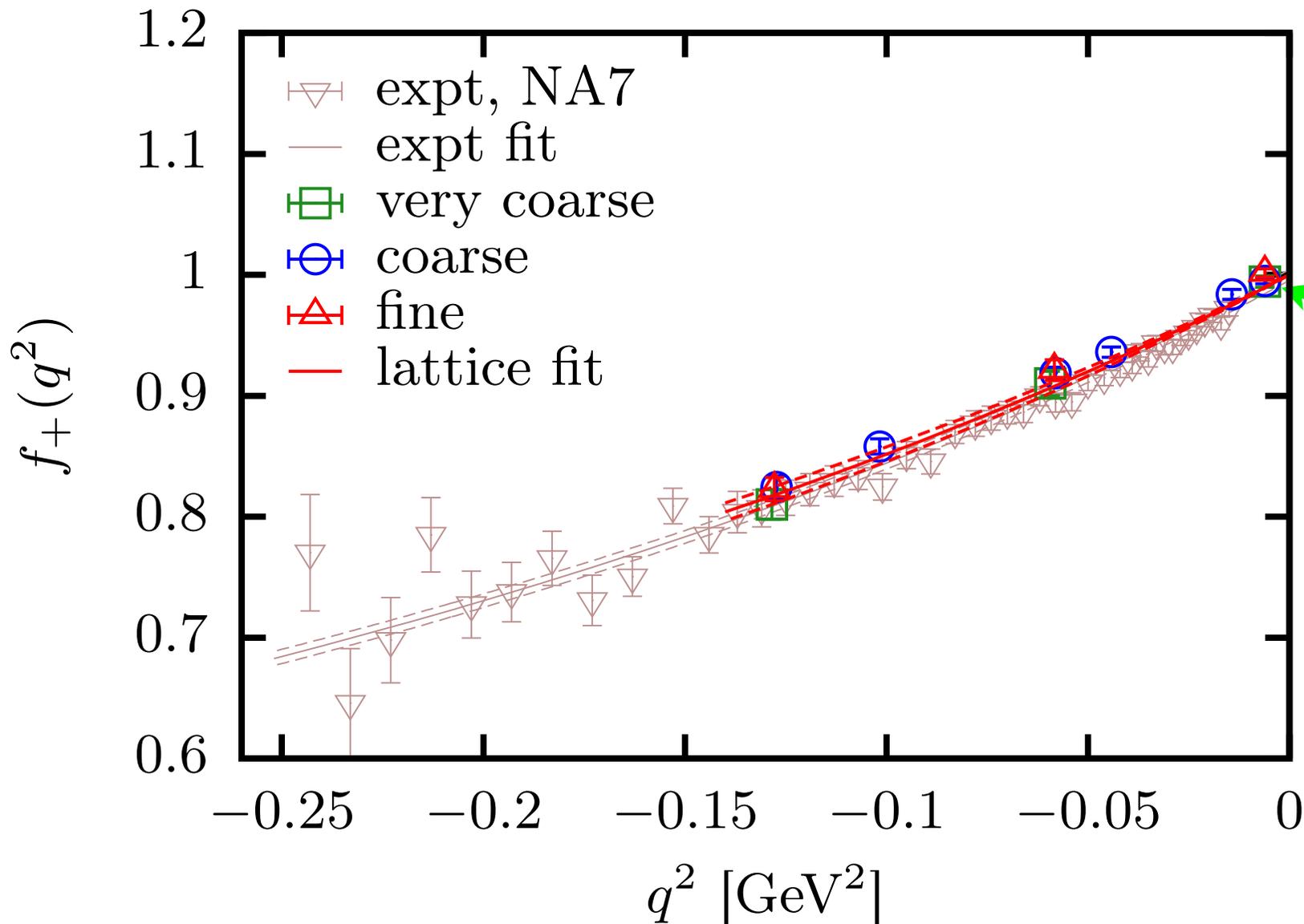
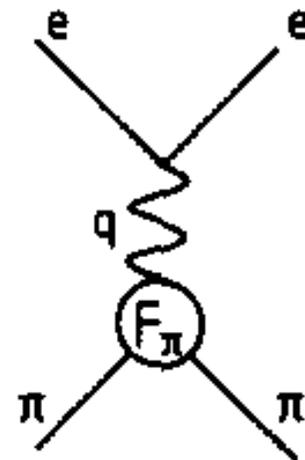
Normln of J must be fixed, e.g.
here $f_+(0) = 1$ from charge cons.



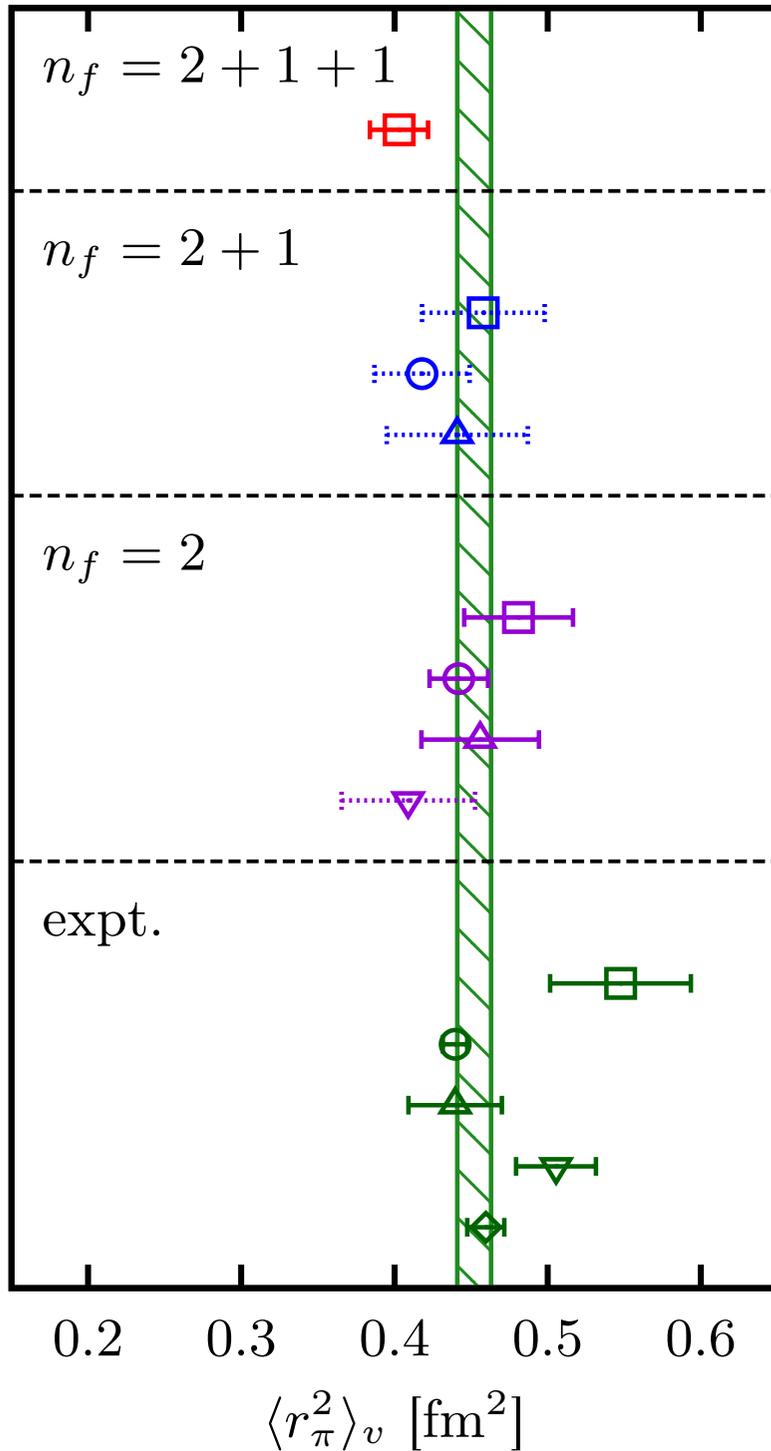
π electromagnetic form factor at small q^2

$\pi - e$ scattering probes π electric charge distn

$$\langle \pi(p_1) | V_\mu | \pi(p_2) \rangle = f_+(q^2) (p_1 + p_2)_\mu$$



Working at **physical** u/d quark masses on HISQ 2+1+1 configs, lattice QCD raw results on top of experiment



HPQCD

$$\langle r^2 \rangle_{V,\text{lat}} = 0.403(18)(6) \text{fm}^2$$

JLQCD

RBC/UKQCD

PACS-CS

At small q^2 can fit pole:

$$f_+(q^2) = \frac{1}{1 - q^2/M^2}$$

Mainz

QCDSF

ETMC

JLQCD/TWQCD

$$M^2 = 6/\langle r^2 \rangle \approx M_\rho^2$$

expt.

MAMI '99

NA7 '86

Fermilab F2 '82

Cornell '78

Orsay '78

$$\langle r^2 \rangle_{V,\text{NA7}} = 0.431(10) \text{fm}^2$$

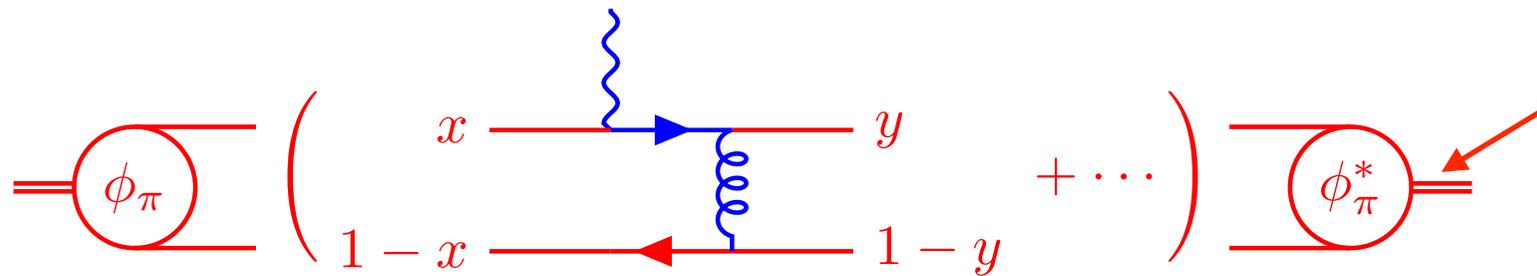
Meson form factors at high (space-like) q^2

To 'understand' form factor want to map it experimentally *and* theoretically up to high q^2 where perturbative QCD becomes valid. But where is this point?

EXPT: JLAB E12-06-101

Perturbative QCD: High q^2 photon must be accompanied by high momentum gluon exchange Lepage, Brodsky, PRD22(1980)2157

$$Q^2 = -q^2 \rightarrow \infty$$



coeffs of expansion of distn amplitude in Gegenbauer polyn., run to 0 with inc. Q^2

$$f_+(q^2) = F_\pi(Q^2) = \frac{8\pi f_\pi^2 \alpha_s(Q/2)}{Q^2} \left| 1 + \sum_{n=2}^{\infty} a_n^\pi(Q/2) \right|^2$$

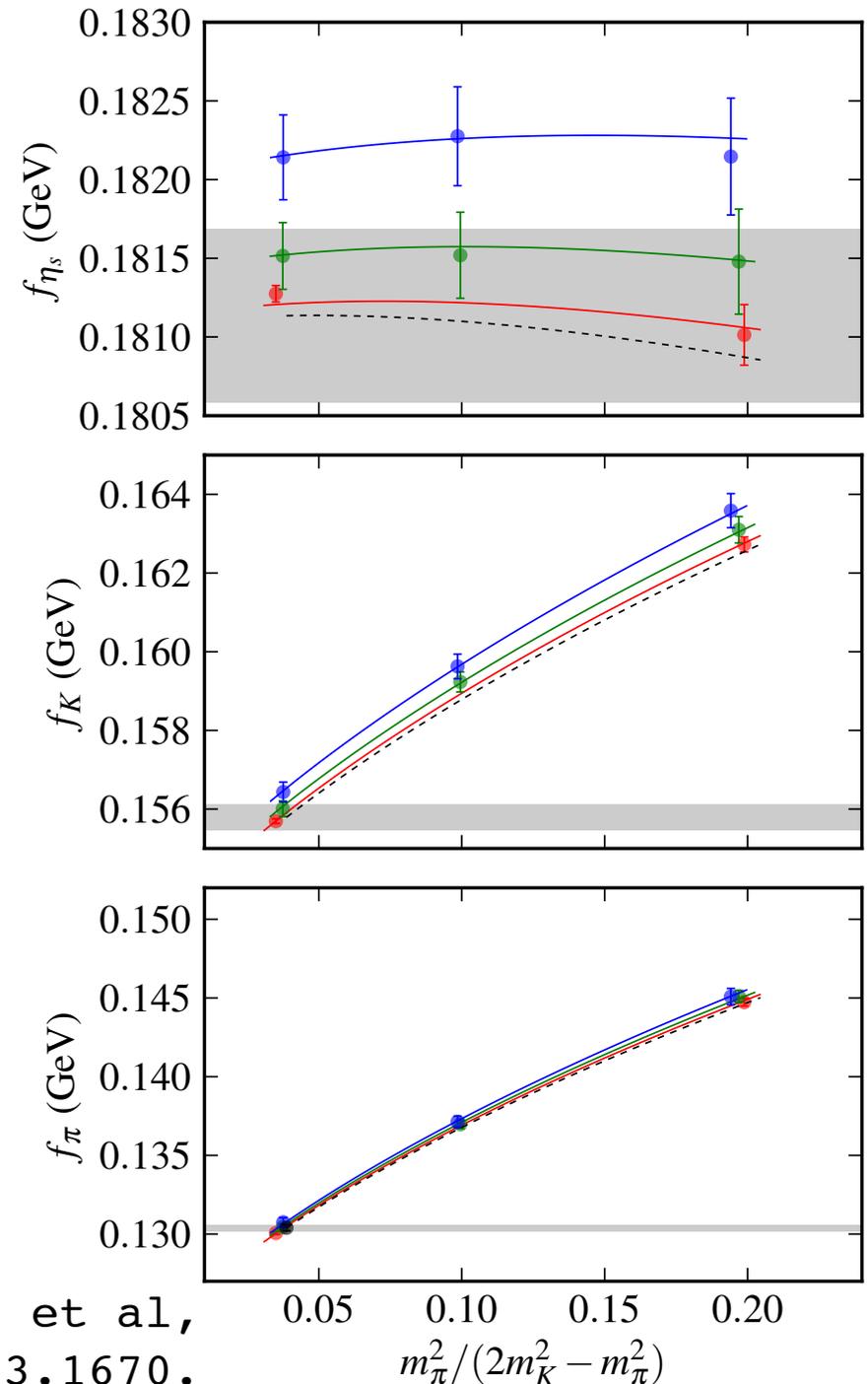
Enter lattice QCD ...

Need a formalism with small discretisation errors that is numerically fast. ***HISQ***

Instead of π use η_s pseudoscalar made of s quarks for added speed.

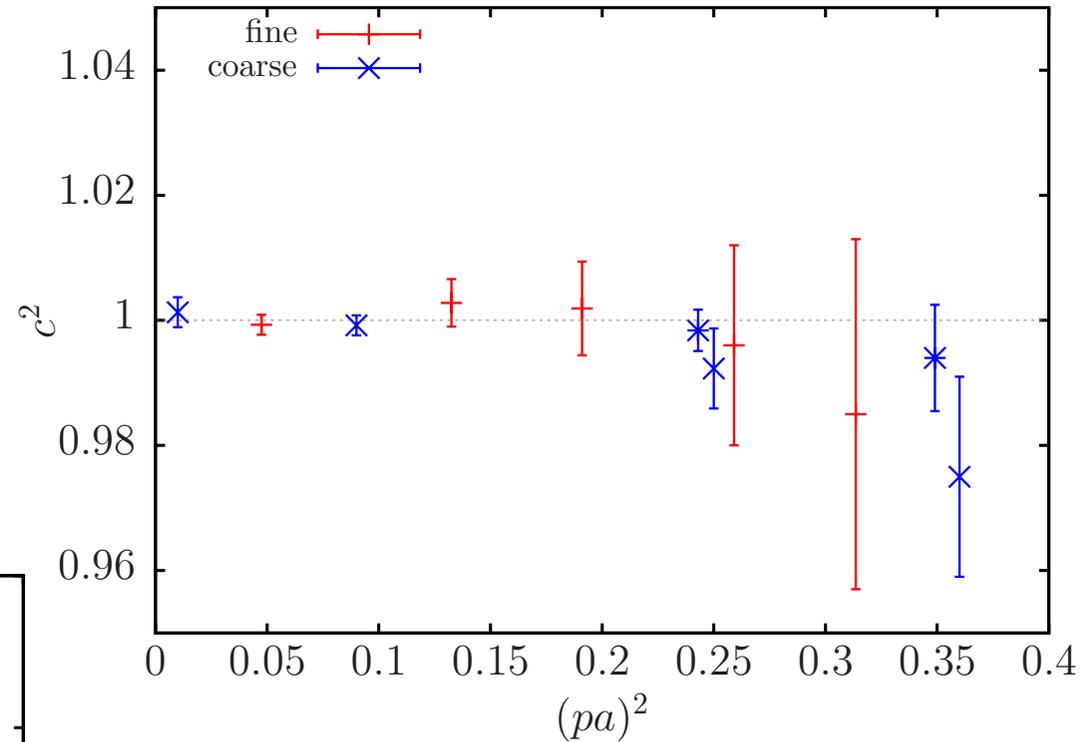
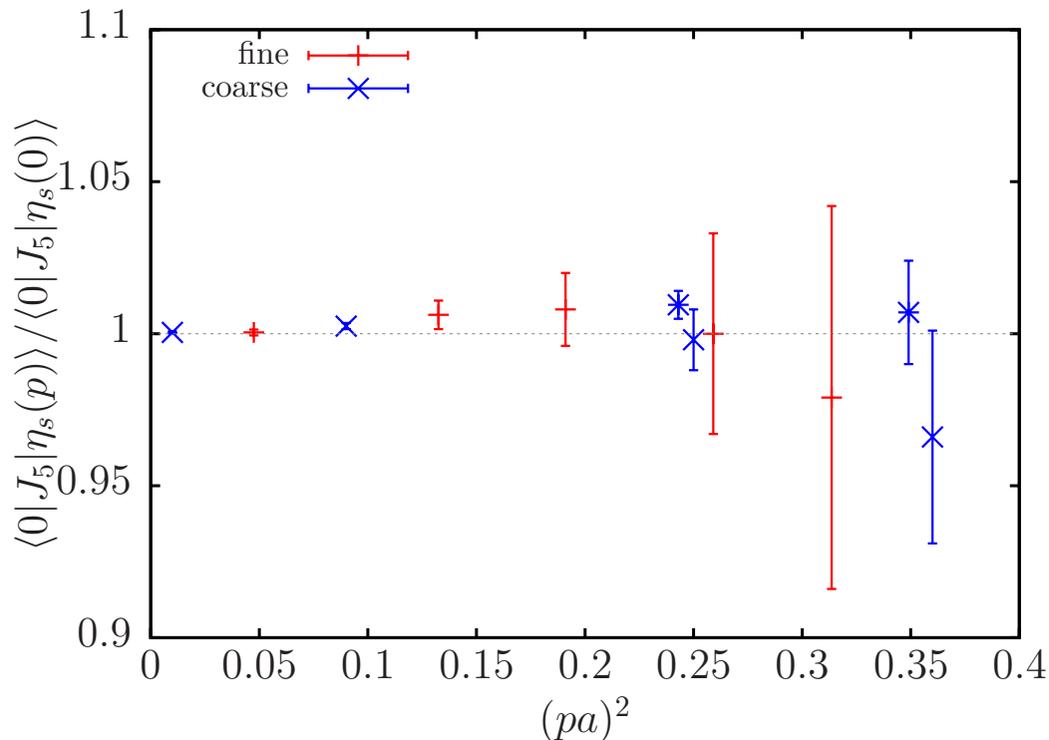
η_s prevented from mixing with light states on lattice so not physical, but properties can be mapped out. $M_{\eta_s} = 688.5(2.2) \text{ MeV}$

$$f_{\eta_s} = 181.14(55) \text{ MeV} \quad \text{R. Dowdall et al, HPQCD, 1303.1670.}$$



Work at two values of lattice spacing (0.12fm and 0.09fm)
and two values of light sea quark mass ($m_s/5$ and $m_s/10$)

Use ‘twisted boundary conditions’ to insert momentum and test discretisation errors as a function of (pa)



For 3pt functions use

Breit frame $\vec{p}_i = -\vec{p}_f$

Maximises Q^2 for a given (pa)

J. Koponen et al, HPQCD, 1701.04250.

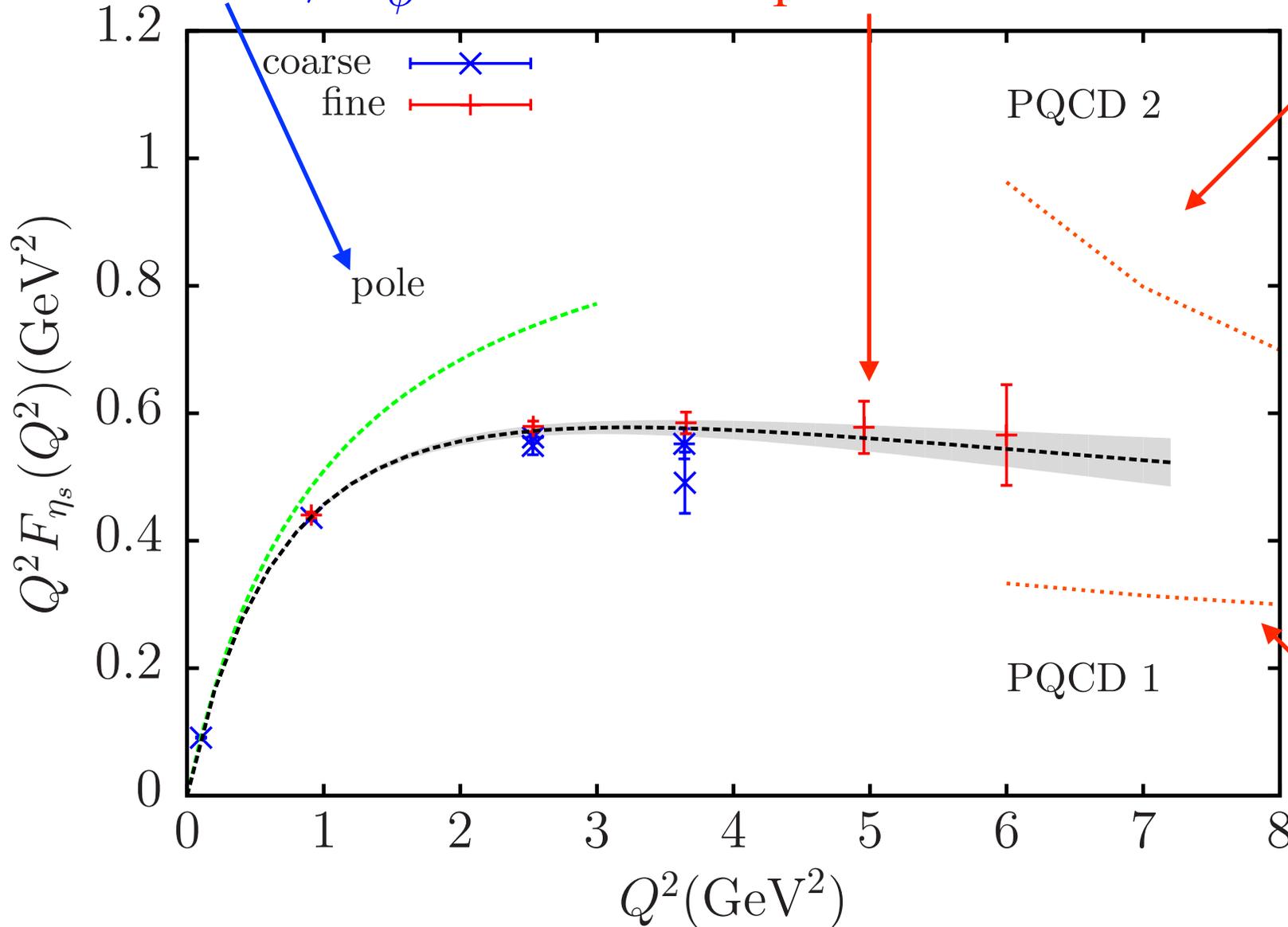
Results - Can reach Q^2 of 6 GeV^2 for $(pa) < 0.6$

Disc. and sea quark mass effects very small

$$F = \frac{1}{1 + Q^2/M_\phi^2}$$

physical point limit of
'z-expansion' fit

pert. QCD inc.



$$\phi_\pi(2\text{GeV}) = [x(1-x)]^{0.52}$$

Braun et al,
1503.03656

asymp. pert.
QCD

J. Koponen et
al, HPQCD,
1701.04250.

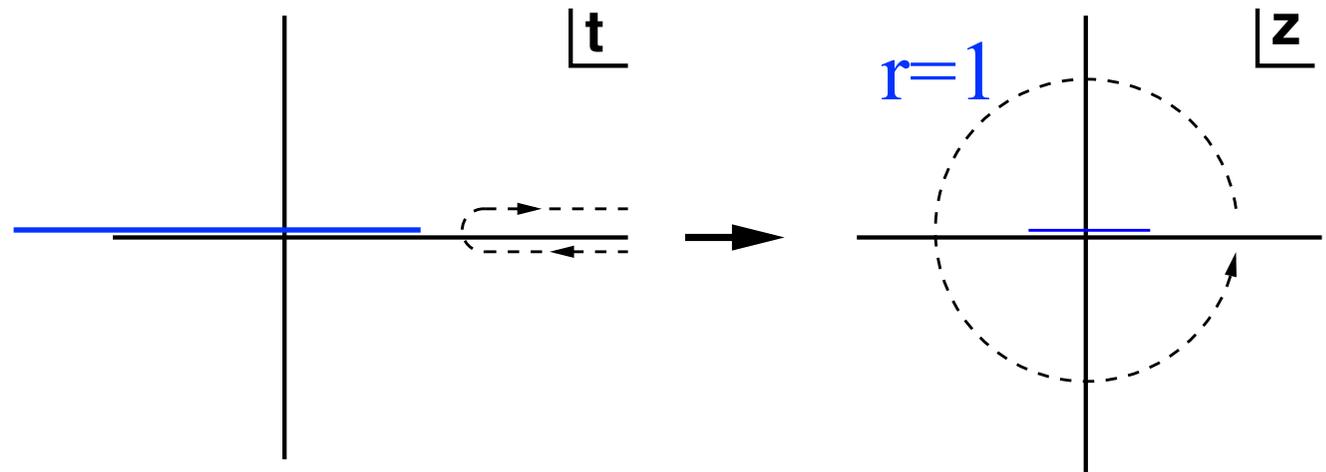
z-expansion

$$z(t, t_{cut}) = \frac{\sqrt{t_{cut} - t} - \sqrt{t_{cut}}}{\sqrt{t_{cut} - t} + \sqrt{t_{cut}}}$$

$$t = q^2$$

$$t_{cut} = 4M_K^2$$

Maps t region into
 $-1 < z < 1$

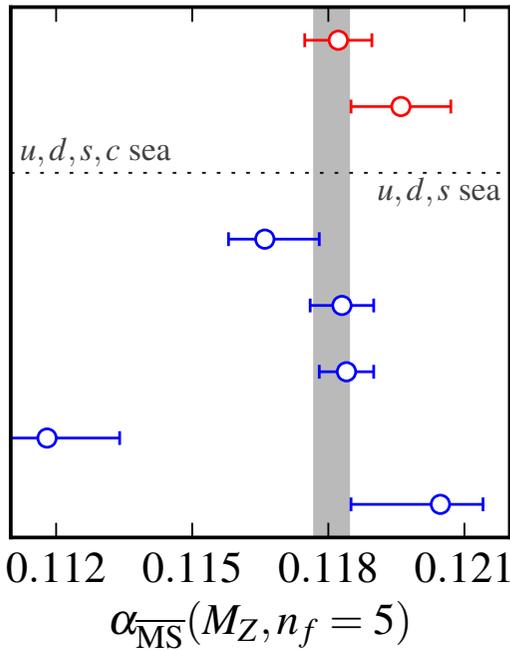


Fit:

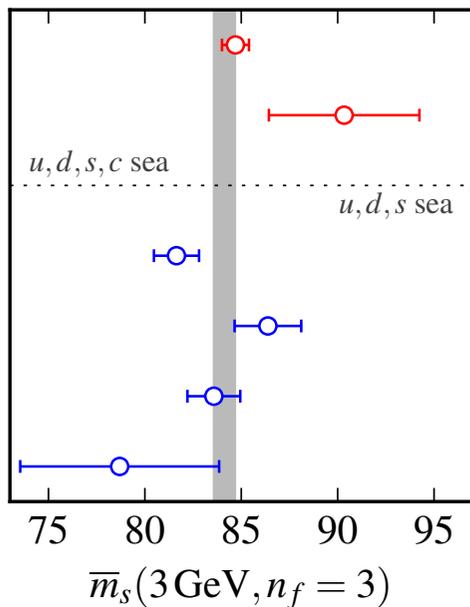
$$(1 + Q^2/M_\phi^2)F = 1 + \sum_i z^i A_i \left[1 + B_i (a\Lambda)^i + C_i (a\Lambda)^4 + D_i \frac{\delta m_{sea}}{10} \right]$$

Further tests of pert. QCD underway: scaling with decay constant as change mass and ff of helicity-nonconserving scalar current.

Quark masses and strong coupling constant

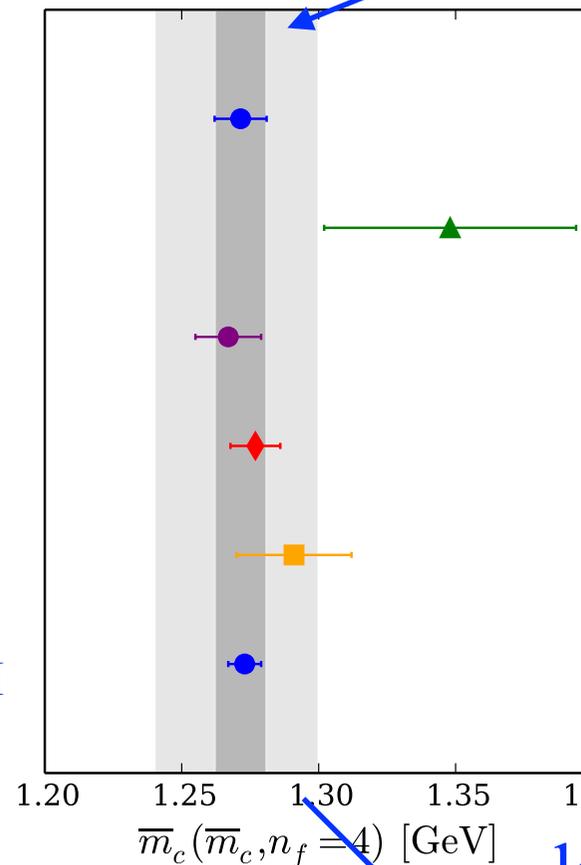


HPQCD- jj 1408.4169
 ETMC 1310.3763
 ghost-gluon
 static potential
 Basavov et al 1407.8437v2
 HPQCD- jj 1004.4285
 HPQCD- W_{nm} 1004.4285
 JLQCD 1002.0371 VPF
 PACS-CS 0906.3906
 SF method



m_c+m_c/m_s
 HPQCD 1408.4169
 ETMC 1403.4504 RI-MOM
 RI-SMOM
 RBC/UKQCD 1411.7017
 Durr et al 1011.2403 RI-MOM
 HPQCD 0910.3102
 HPQCD (pert) 0511160

Multiple lattice QCD methods now that agree to high accuracy



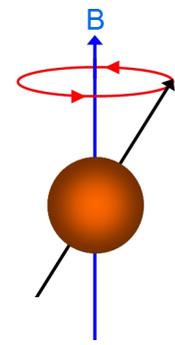
PDG

1408.4169
 HPQCD HISQ4 JJ
 1403.4504
 ETMC RI-mom
 1606.08798
 M+P HISQ3 JJ
 1511.09163
 JLQCD dw3 JJ
 1410.3343
 χ QCD ovp3 RI-m
 1004.4285
 HPQCD HISQ3 JJ

latt, $n_f=4$
 1.2715(95) GeV

Muon anomalous magnetic moment

$$\vec{\mu} = g \frac{e}{2m} \vec{S} \quad a_{\mu} = \frac{g - 2}{2}$$



Measure using polarised muons circulating in E and B fields. At a momentum where $\beta \times E$ terms cancel, difference between precession and cyclotron frequencies:

$$\omega_a = -\frac{e}{m} a_{\mu} B$$

BNL result:

$$a_{\mu}^{expt} = 11659208.9(6.3) \times 10^{-10}$$

E989 (FNAL) will
reduce exptl uncty to
1.6, starting 2017

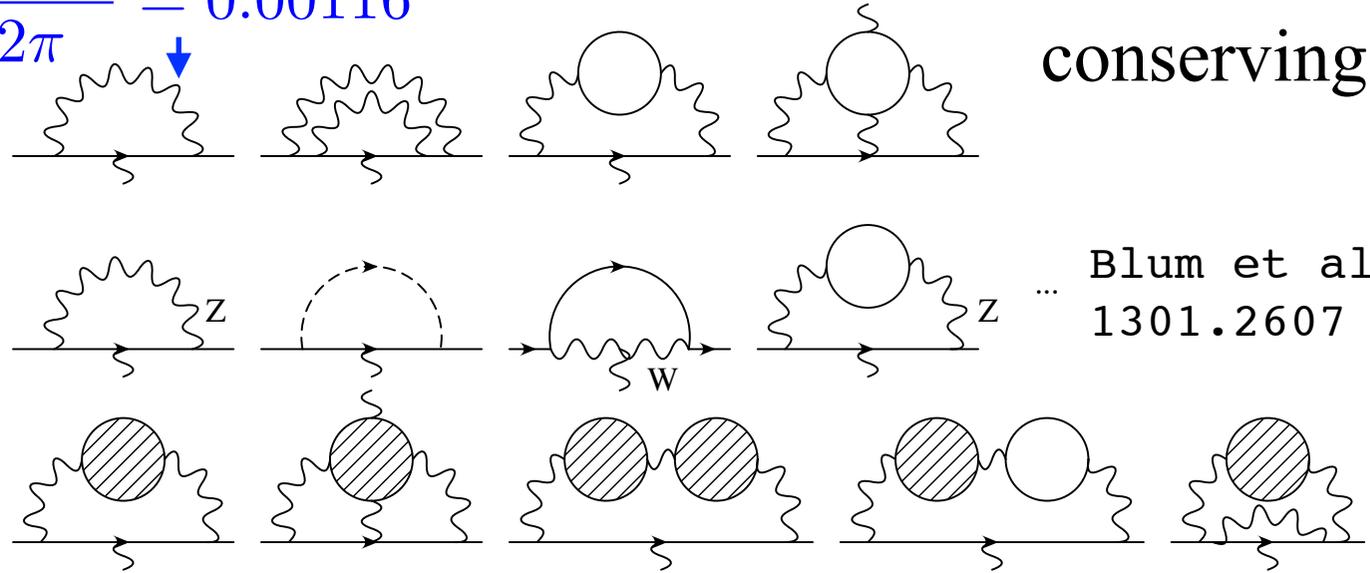


Standard Model theory expectations

Contributions from QED, EW and QCD interactions.

QED dominates.
QCD contriibs start at α_{QED}^2

$$\frac{\alpha_{QED}}{2\pi} = 0.00116$$



flavour and CP conserving

... Blum et al, 1301.2607

LO Hadronic vacuum polarisation (HVP) dominates uncertainty in SM result

$$a_{\mu}^{QED} = 11658471.885(4) \times 10^{-10}$$

$$a_{\mu}^{EW} = 15.4(2) \times 10^{-10}$$

$$a_{\mu}^{E821} = 11659208.9(6.3) \times 10^{-10}$$

Hadronic (and other) contributions = EXPT - QED - EW

$$a_{\mu}^{expt} - a_{\mu}^{QED} - a_{\mu}^{EW} = 721.7(6.3) \times 10^{-10}$$
$$= a_{\mu}^{HVP} + a_{\mu}^{HOHVP} + a_{\mu}^{HLbL} + a_{\mu}^{new\ physics}$$

Focus on lowest order hadronic vacuum polarisation,
so assume:

$$a_{\mu}^{HLbL} = 10.5(2.6) \times 10^{-10}$$

$$a_{\mu}^{HOHVP} = -8.85(9) \times 10^{-10} \leftarrow \text{NLO+NNLO}$$

Kurz et al,
1403.6400

$$a_{\mu}^{HVP, no\ new\ physics} = 719.8(6.8) \times 10^{-10}$$

Lattice calculation of HVP,LO

Analytically continue to Euclidean q^2 .

$$a_{\mu}^{HVP,i} = \frac{\alpha}{\pi} \int_0^{\infty} dq^2 f(q^2) (4\pi\alpha e_i^2) \hat{\Pi}_i(q^2)$$

connected contribution for flavour i

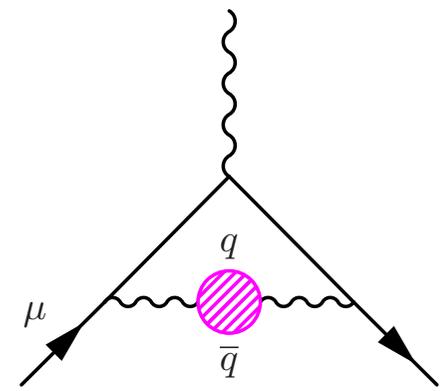
$f(q^2)$ divergent function with scale set by m_{μ}

$$\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$$

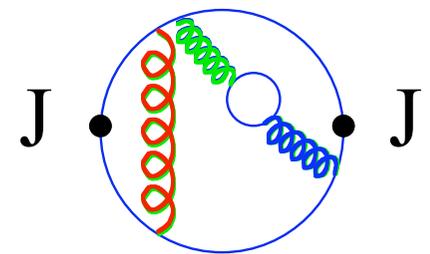
HPQCD method: time-moments of vector JJ correlators give expansion around $q^2=0$

$$G_n \equiv \sum_{t, \vec{x}} t^n Z_V^2 \langle J^j(\vec{x}, t) J^j(0) \rangle \quad \hat{\Pi}(q^2) = \sum_{k=1}^{\infty} q^k \Pi_k$$

$$\Pi_k = (-1)^{k+1} \frac{G_{2k+2}}{(2k+2)!}$$



Blum, hep-lat/
0212018



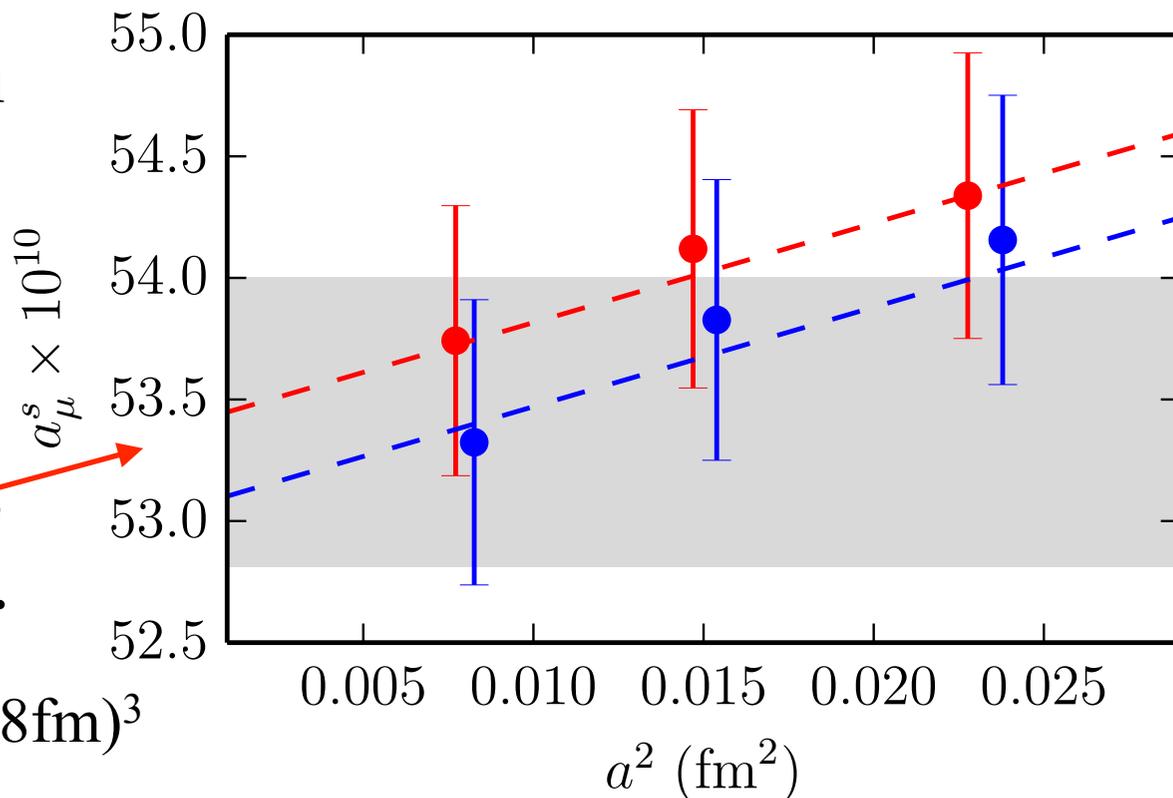
replace with
[2,2] Padé

Test on STRANGE contribution

HISQ valence quarks on
MILC 2+1+1 HISQ
configs. Local J_V -
nonpert. Z_V .

multiple a (fixed by w_0),
 m_l (inc. phys.), volumes.

Tune s from η_s ↑
up to $(5.8\text{fm})^3$



	a_μ^s
Uncertainty in lattice spacing (w_0, r_1):	1.0%
Uncertainty in Z_V :	0.4%
Monte Carlo statistics:	0.1%
$a^2 \rightarrow 0$ extrapolation:	0.1%
QED corrections:	0.1%
Quark mass tuning:	0.0%
Finite lattice volume:	< 0.1%
Padé approximants:	< 0.1%
Total:	1.1%

$$a_\mu^{HVP,s} = 53.41(59) \times 10^{-10}$$

Also

$$a_\mu^{HVP,c} = 14.4(4) \times 10^{-10}$$

$$a_\mu^{HVP,b} = 0.27(4) \times 10^{-10}$$

UP/DOWN contribution

$$m_u = m_d$$

Chakraborty et al, HPQCD, 1601.03071

Much noisier and sensitive to u/d mass. Use

$$G(t) = \begin{cases} G_{\text{data}}(t) & \text{for } t \leq t^* \\ G_{\text{fit}}(t) & \text{for } t > t^* \end{cases}$$

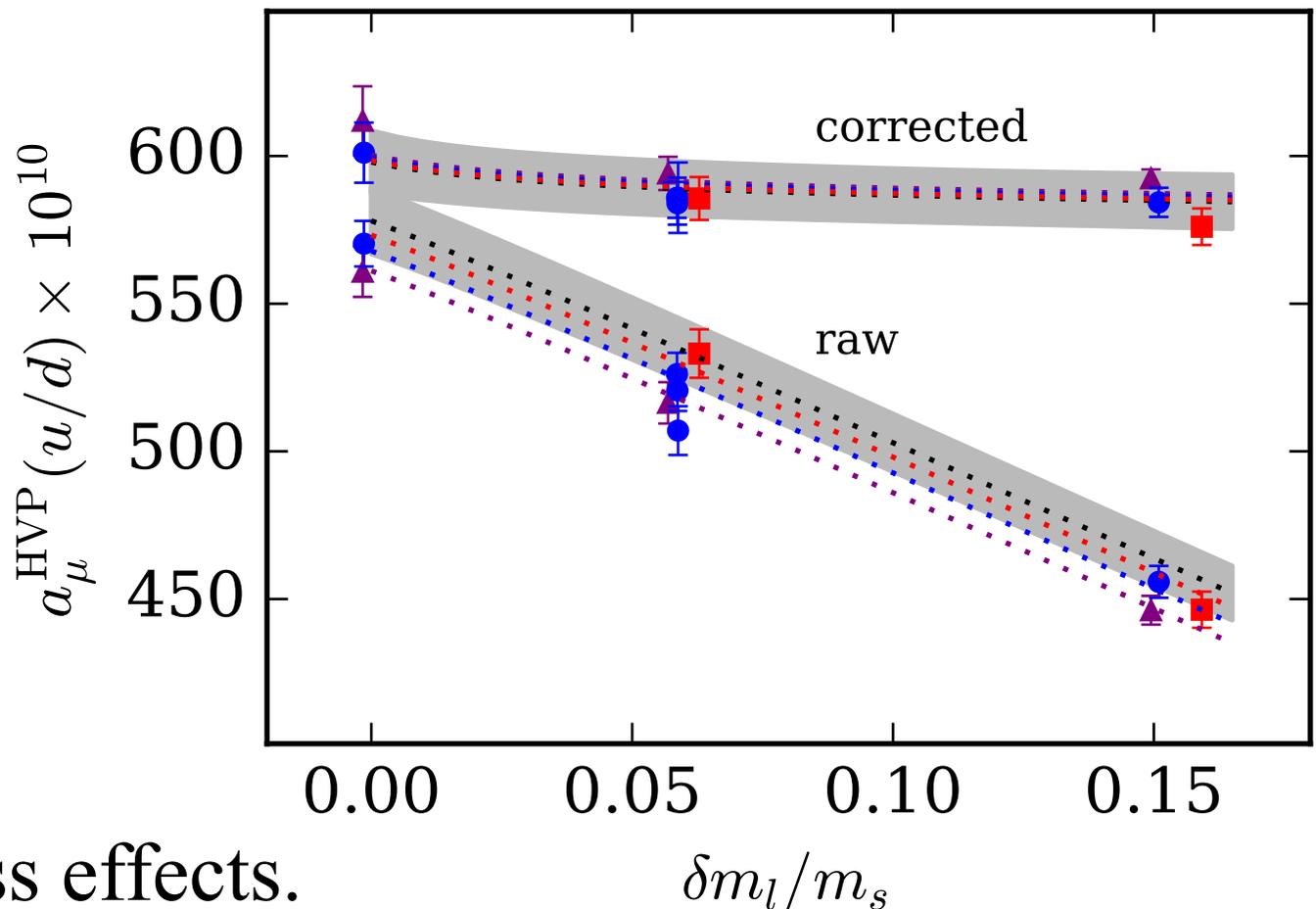
← from Monte Carlo ← from multi-exponential fit

$$t^* = 1.5\text{fm} = 6/m_\rho \quad \text{so 70\% of result from } G_{\text{data}}$$

Must correct for finite vol. effects in $\pi\pi$ contribn using scalar QED (7%)

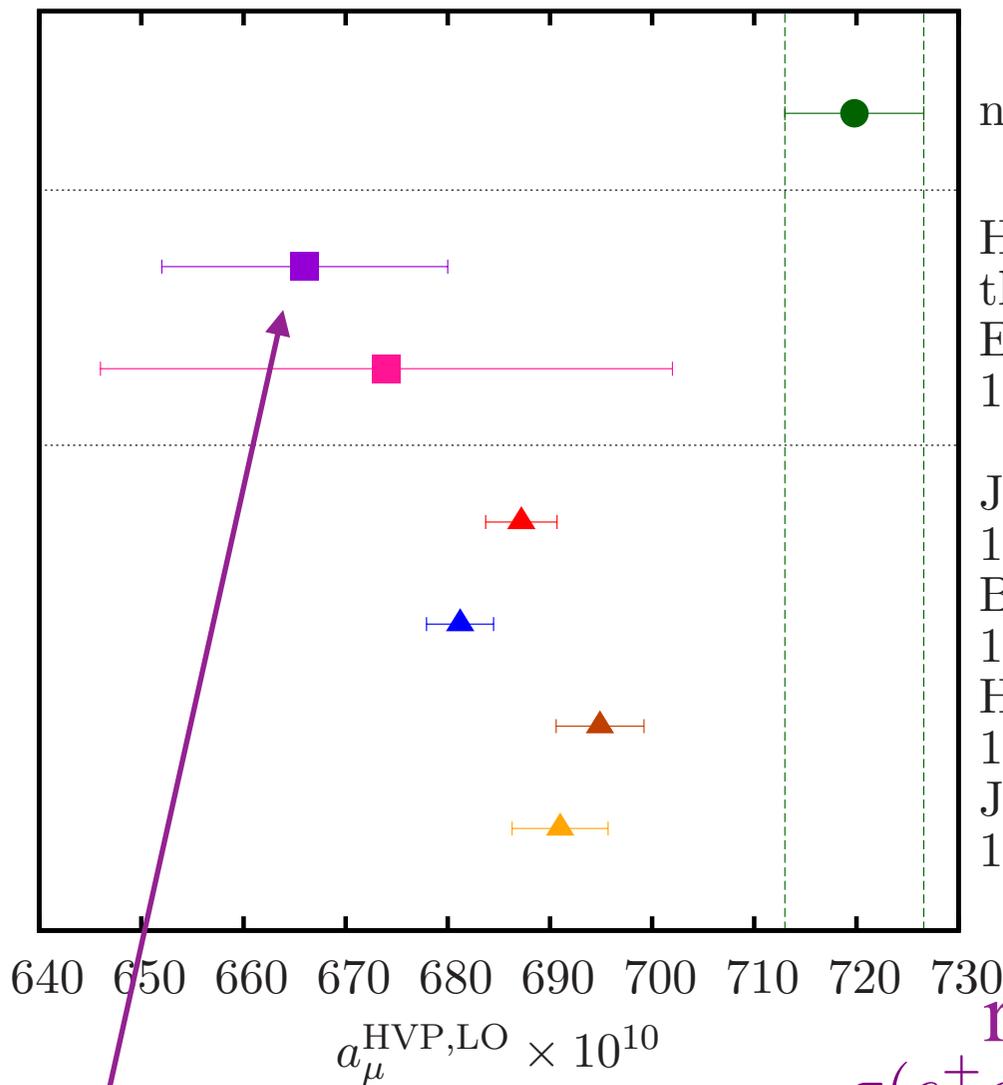
Rescale Π_j by $(m_\rho^{\text{latt}}/m_\rho^{\text{expt}})^{2j}$

to reduce u/d mass effects.



Combining numbers for a total

Chakraborty et al, HPQCD,
1601.03071



$a_\mu^{\text{HVP,LO}} \times 10^{-10}$

598(11) *u/d*

53.4(6) *s*

14.4(4) *c*

0.27(4) *b*

Total 666(6)(12)

add syst from
disc. diags
(1.5%) in quad

results using
 $\sigma(e^+e^- \rightarrow \text{hadrons})$

3.5 σ discrepancy with no new physics

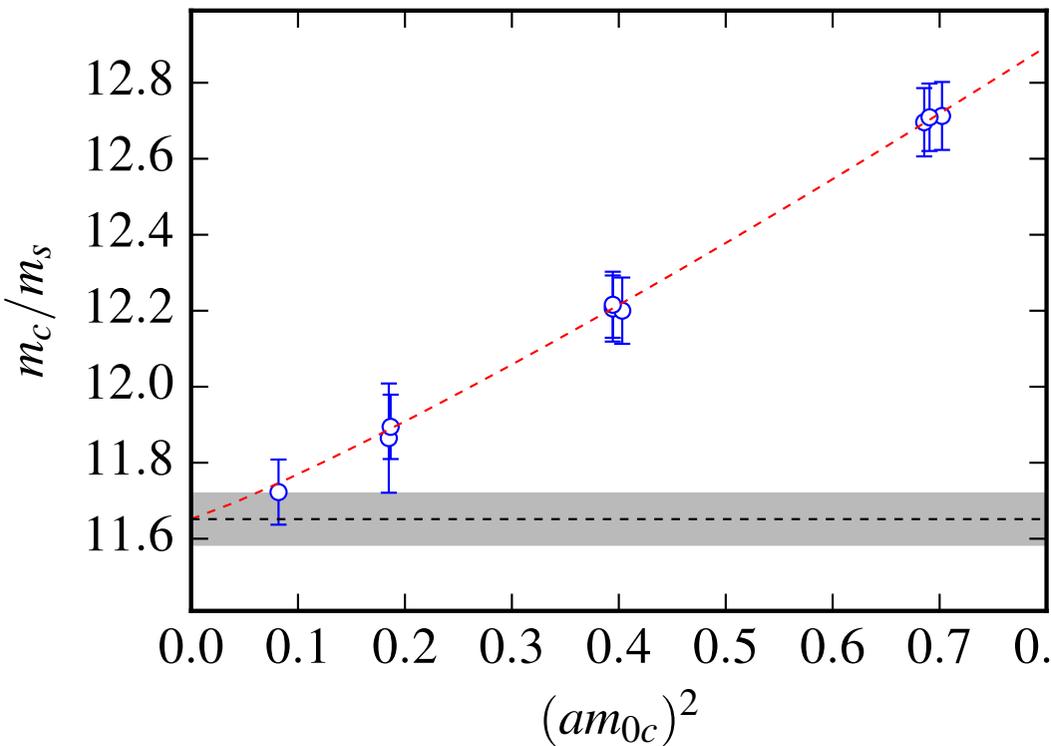
Conclusion

- Lattice QCD calculations now on ‘2nd generation’ gluon configs with charm in the sea and $m_{u,d}$ at physical value (so no extrapln).
- Gold-plated hadron masses and decay constants provide stringent tests of QCD/SM.
- NEW calculation of light meson electromagnetic form factor up to $Q^2 = 6 \text{ GeV}^2$ shows behaviour inconsistent with asymptotic pert. QCD. Further calc. underway ...
- Accurate quark mass determinations being tested with multiple methods
- sub-1% uncertainties on lattice QCD calculations for HVP contribution to a_μ are within sight.

Spares

$$m_c/m_s$$

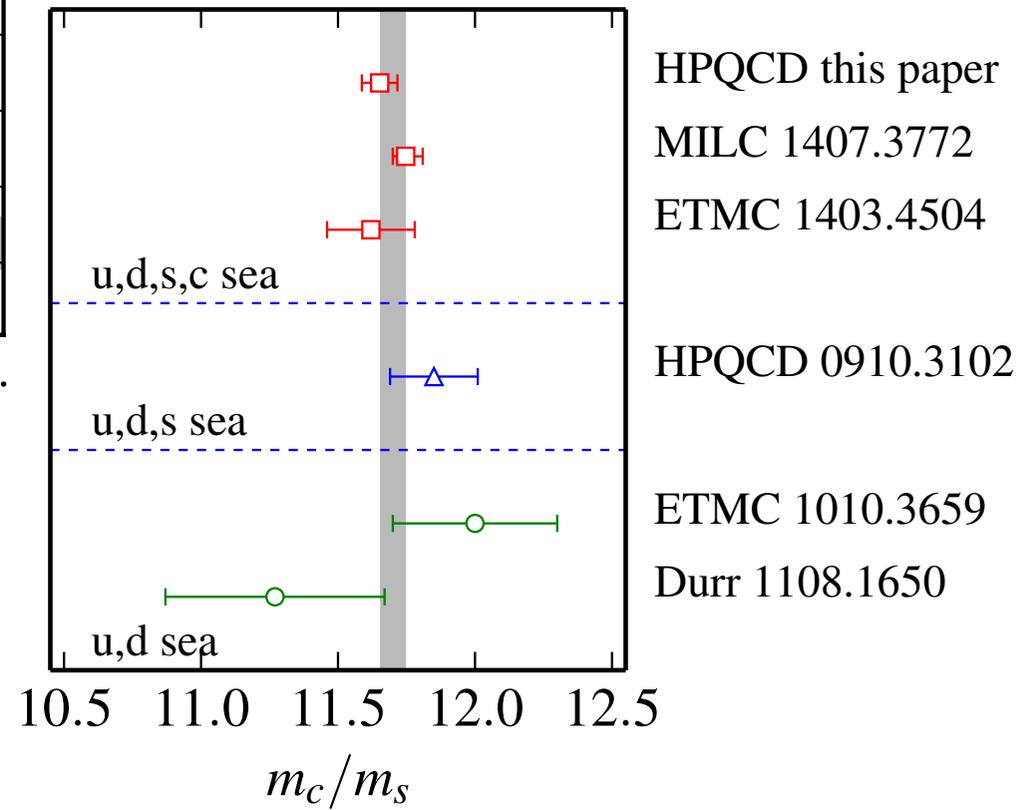
Mass ratio direct from lattice QCD using same formalism for both quarks (HISQ). Not possible to connect heavy and light masses without lattice.



$$\frac{m_c}{m_s} = 11.652(65)$$

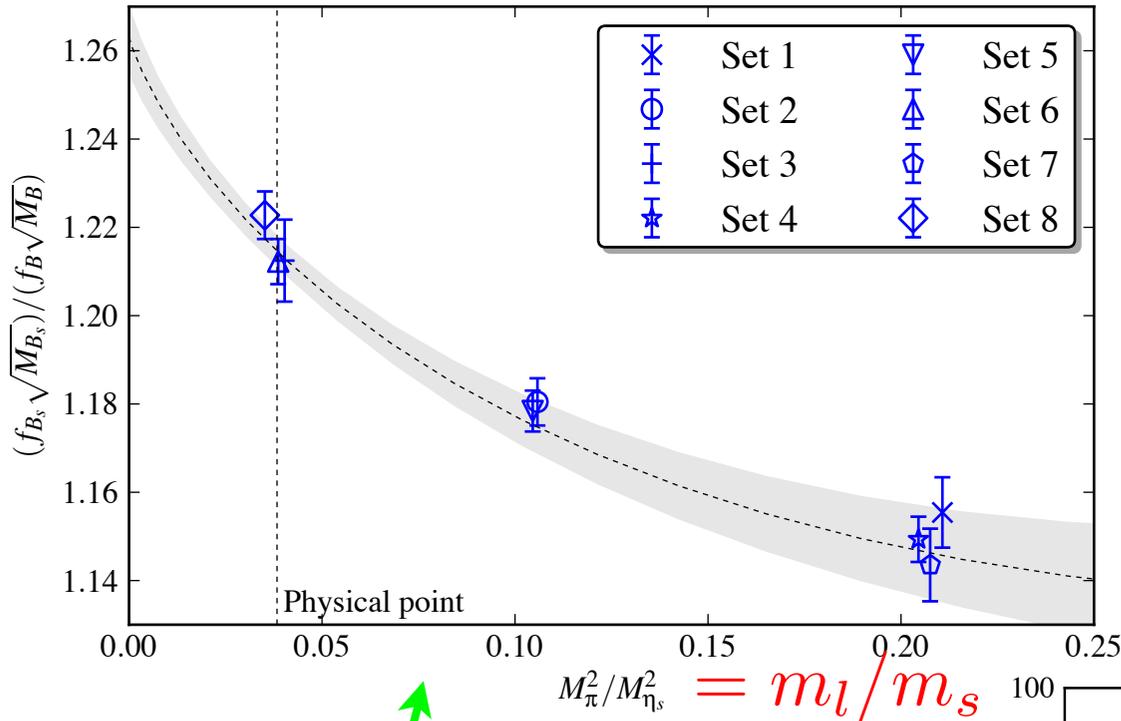
HPQCD, 1408.4169

Good consistency
between different lattice
actions



B meson decay constants: HPQCD results from NRQCD b and physical u/d quarks

HPQCD: R Dowdall et al, 1302.2644.

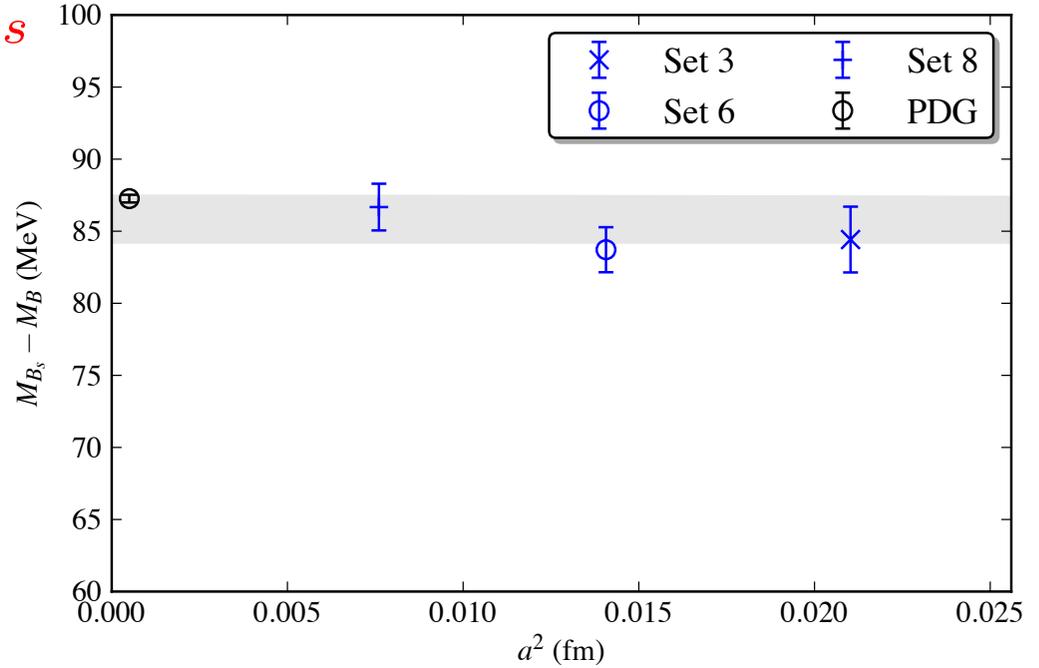


Use 2+1+1 configs with u/d down to physical values + improved NRQCD

meson mass difference correct to 2%



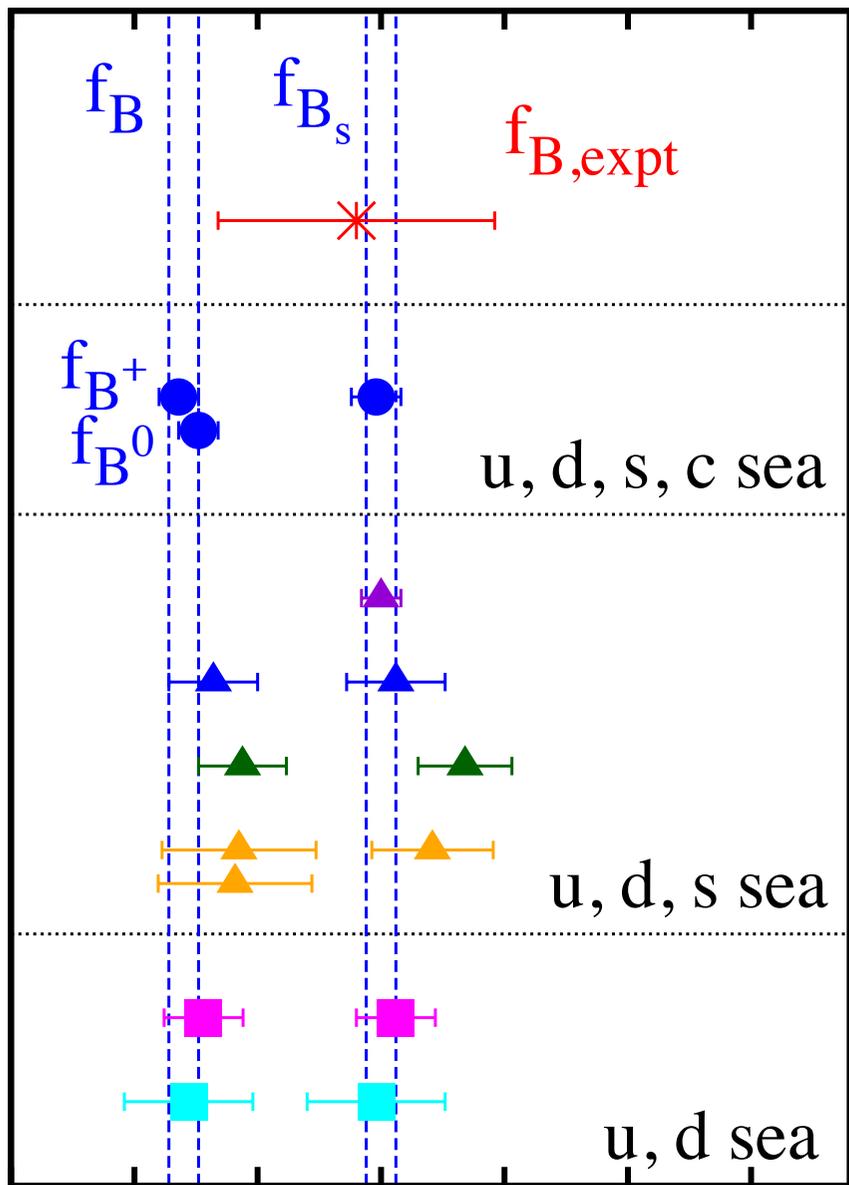
B_s to B decay constant ratio accurate to 0.6% - since Z factors cancel. Separate decay constants to 2%



185(3) 225(3)

averages

B, B_s decay constant
world averages



PDG av. branching fraction
+ unitarity V_{ub}

HPOCD NRQCD
1302.2644

HPQCD HISQ 1110.4510
HPQCD NRQCD
1202.4914

FNAL/MILC 1112.3051

RBC/UKQCD 1404.4670

ETMC 1308.1851

ALPHA 1404.3590

Different
approaches to
b quark
being used
and to
normalising
weak current

150 175 200 225 250 275 300

f_{B_x} / MeV

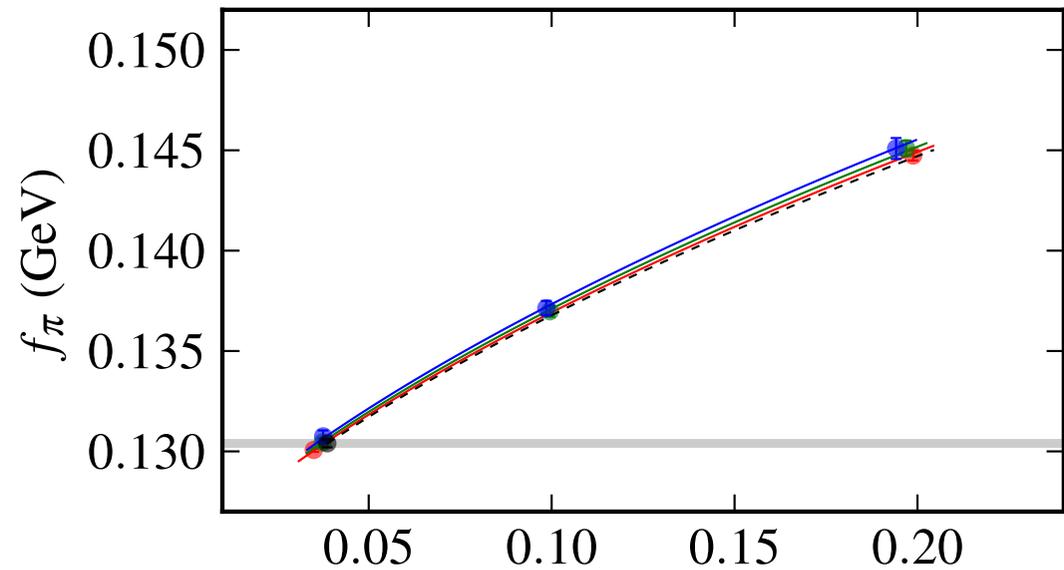
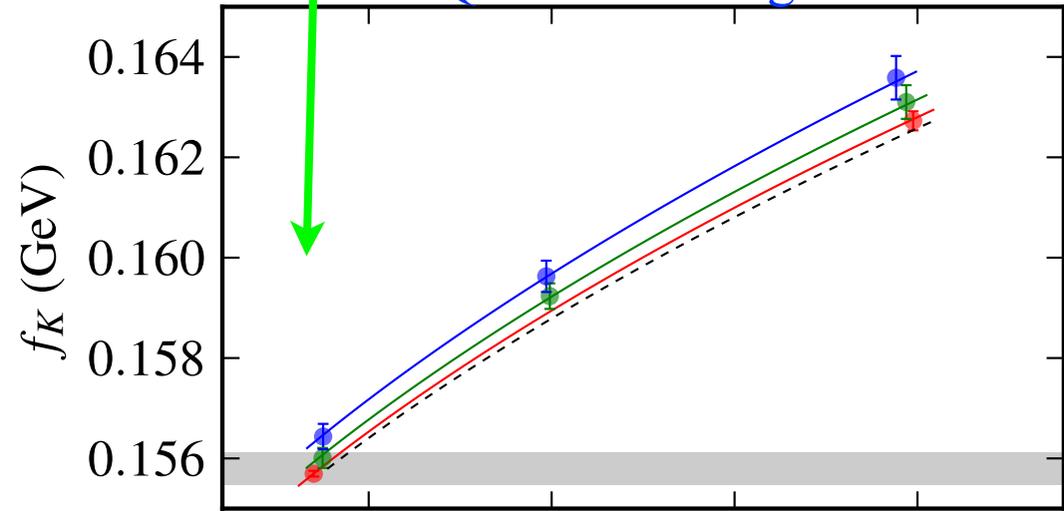
NOTE:

$f_{B_s} < f_{D_s}$ now quite clear

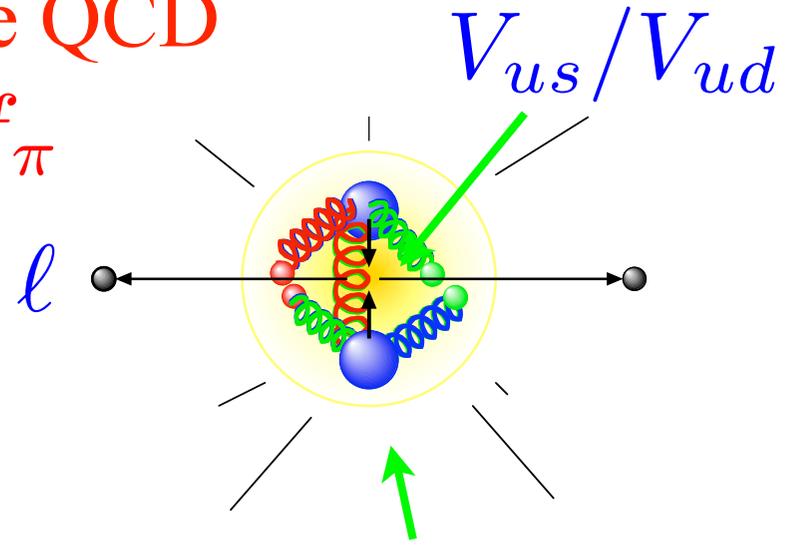
Constraining new physics with lattice QCD

* results at physical u/d quark masses*

HISQ 2+1+1 configs



f_K / f_π



Annihilation of K/π to W allows CKM element determination given decay constants from lattice QCD

expt for $\frac{\Gamma(K^+ \rightarrow l\nu)}{\Gamma(\pi^+ \rightarrow l\nu)}$

$$\frac{|V_{us}|f_{K^+}}{|V_{ud}|f_{\pi^+}} = 0.27598(35)_{\text{Br}(K^+)}(25)_{EM}$$

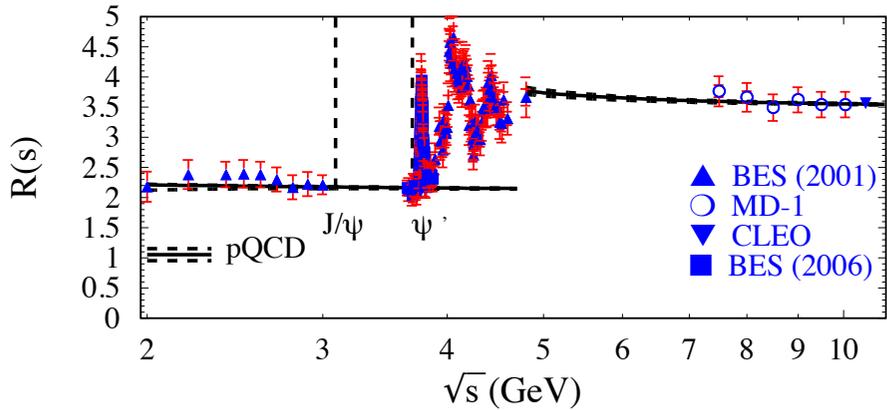
$\frac{f_{K^+}}{f_{\pi^+}}$ from lattice gives CKM

R. Dowdall et al, HPQCD: 1303.1670

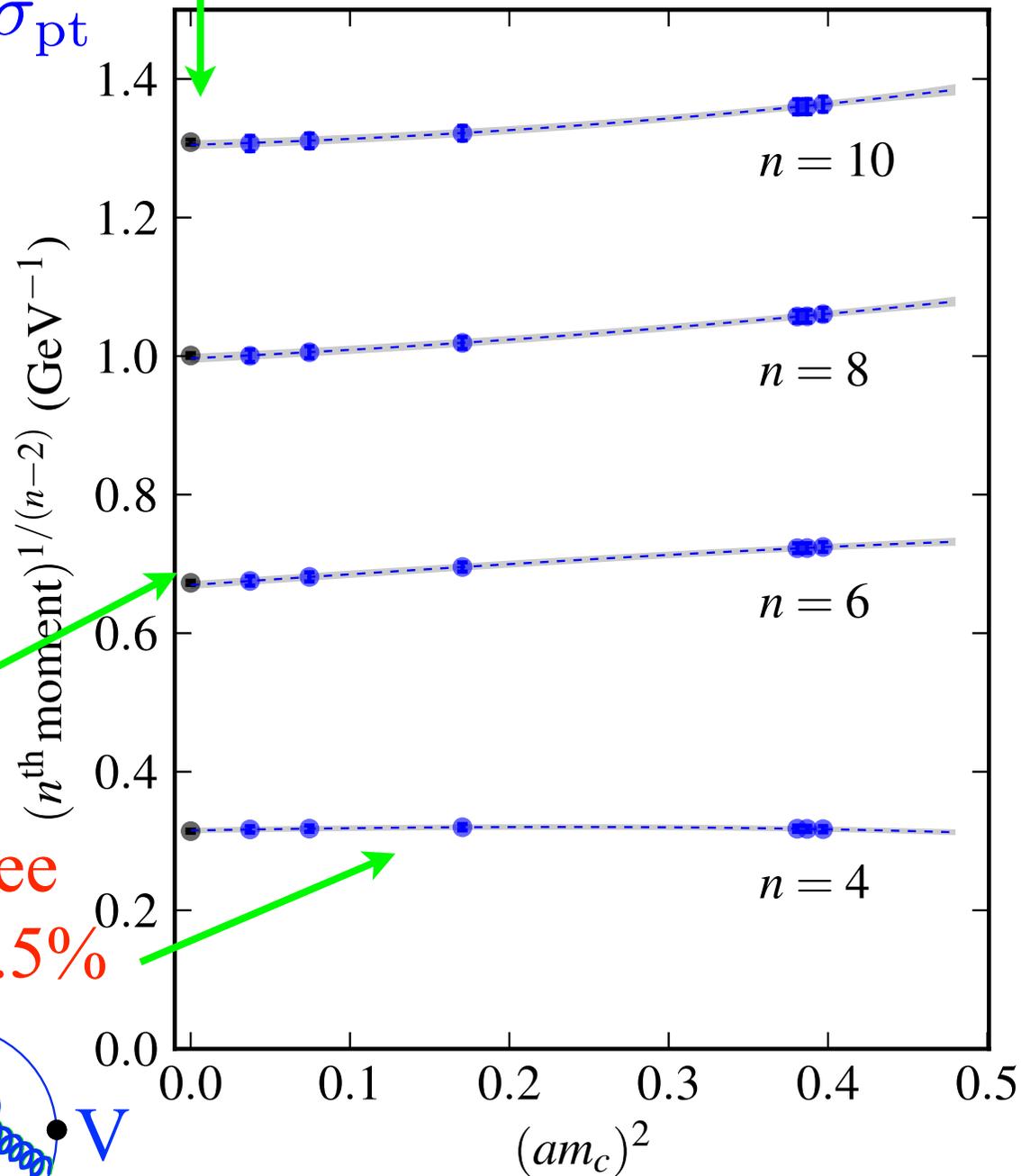
$$\frac{m_\pi^2}{(2m_K^2 - m_\pi^2)} = m_{u,d}/m_s$$

Charm contribution to

$$R = \sigma(e^+e^- \rightarrow \text{hadrons}) / \sigma_{\text{pt}}$$



'expt' ← from J. Kuhn et al, hep-ph/0702103



Subtract u,d,s with pert. th. to isolate charm contribution.
 Calculate inverse-s-moment:

$$\mathcal{M}_n \equiv \int \frac{ds}{s^{n+1}} R_c(s)$$

Lattice calcln:

$$M_n = \sum_t t^n G(t)$$

Agree to 1.5%

Error budget for HVP,LO calculation

TABLE III: Error budget for the connected contributions to the muon anomaly a_μ from vacuum polarization of u/d quarks.

	$a_\mu^{\text{HVP,LO}}(u/d)$
QED corrections:	1.0 %
Isospin breaking corrections:	1.0 %
Staggered pions, finite volume:	0.7 %
Noise reduction (t^*):	0.5 %
Valence m_ℓ extrapolation:	0.4 %
Monte Carlo statistics:	0.4 %
Padé approximants:	0.4 %
$a^2 \rightarrow 0$ extrapolation:	0.3 %
Z_V uncertainty:	0.4 %
Correlator fits:	0.2 %
Tuning sea-quark masses:	0.2 %
Lattice spacing uncertainty:	< 0.05 %
Total:	1.9 %