TMDs and the W boson mass

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Effect of flavor-dependent partonic transverse momentum on the determination of the W boson mass in hadronic collisions

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Within the framework of transverse-momentum-dependent factorization, we investigate for the first time the impact of a flavor-dependent intrinsic transverse momentum of quarks on the production of W± bosons in proton-proton collisions at \( \sqrt{s} = 7 \) TeV. We estimate the shift in the extracted value of the W boson mass \( M_W \) induced by different choices of flavor-dependent parameters for the intrinsic quark transverse momentum by means of a template fit to the transverse-mass and the lepton transverse-momentum distributions of the W-decay products. We obtain \(-11 \leq \Delta M_W^+ \leq 4\) MeV and \(-6 \leq \Delta M_W^- \leq 2\) MeV with a statistical uncertainty of \(\pm 4\) MeV. Our findings call for more detailed investigations of flavor-dependent nonperturbative effects linked to the proton structure at hadron colliders.

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Introduction and motivation.

Nonperturbative effects in transverse-momentum-dependent (TMD) phenomena are a central topic in the hadronic physics community with potentially important applications to high-energy physics. The study of nonperturbative corrections originates from the work of Parisi and Petronzio [1] and Collins, Soper, and Sterman [2], which focused on the role of the hard scale of the process compared to the infrared scale of QCD. TMD factorization and evolution have been extensively studied in the literature [3–6], together with the matching to collinear factorization [2, 7–12]. Despite the limited amount of data available and the many open theoretical questions, in the past years we started gaining phenomenological information about TMD parton distribution functions (TMD PDFs) with increasing level of accuracy. Recently, the unpolarized quark TMD PDF was extracted for the first time from a global fit of data from deep-inelastic scattering (SIDIS) and \(e^+e^-\) production [13].

Experimental measurements and uncertainties.

The determination of the W boson mass, \( M_W \), from the global electroweak fit (\( M_W = 80.356 \pm 8 \) MeV) [20] features a very small uncertainty that sets a goal for the precision of the experimental measurements at hadron colliders.

Precise determinations of \( M_W \) have been extracted from \( pp \) collisions at D0 [21] and at CDF [22], and from \( pp \) collisions at ATLAS [23] with a total uncertainty of 23 MeV, 19 MeV and 19 MeV, respectively. The current world average, based on these measurements and the ones performed at LEP, is \( M_W = 80.379 \pm 12 \) MeV [24]. The experimental analyses are based on a template-fit procedure on the differential distributions of the decay products: in particular, the transverse momentum of the final lepton, \( p_{T\ell} \), the transverse momentum of the neutrino \( p_{T\nu} \) (only at the Tevatron), and the transverse mass \( m_T \) of the lepton pair (where \( m_{T} = \sqrt{2 \, p_T \cdot p_T} \), with \( \phi_{\ell,\nu} \) being the azimuthal angles of the lepton and the neutrino, respectively).

In a template-fit procedure, several histograms are generated, each corresponding to the results of the different calculations used to compute the templates (choice of scales, PDFs, perturbative order, resummation of logarithmically enhanced contributions, nonperturbative effects linked to the proton structure at hadron colliders).

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TMDs and determination of W mass

CDF collaboration, PRD77 (08)

TABLE XVI. Systematic uncertainties in units of MeV on the combination of the six fits in the electron and muon channels. Each uncertainty has been estimated by removing its covariance and repeating the sixfold combination.

<table>
<thead>
<tr>
<th>Source</th>
<th>Uncertainty (MeV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Lepton scale</td>
<td>23.1</td>
</tr>
<tr>
<td>Lepton resolution</td>
<td>4.4</td>
</tr>
<tr>
<td>Lepton efficiency</td>
<td>1.7</td>
</tr>
<tr>
<td>Lepton tower removal</td>
<td>6.3</td>
</tr>
<tr>
<td>Recoil energy scale</td>
<td>8.3</td>
</tr>
<tr>
<td>Recoil energy resolution</td>
<td>9.6</td>
</tr>
<tr>
<td>Backgrounds</td>
<td>6.4</td>
</tr>
<tr>
<td>PDFs</td>
<td>12.6</td>
</tr>
<tr>
<td>W boson $p_T$</td>
<td>3.9</td>
</tr>
<tr>
<td>Photon radiation</td>
<td>11.6</td>
</tr>
</tbody>
</table>

$m_W = 80.398 \pm 0.025$ GeV. \hfill (53)

...eight years later we managed to published something about it...
Our findings

The fact that quark intrinsic transverse momentum can be flavor-dependent leads to an additional uncertainty on $M_W$, not considered so far.

$-11 \leq M_{W^+} \leq 4 \text{ MeV}$

$-6 \leq M_{W^-} \leq 2 \text{ MeV}$
The state of the art

\[ m_W = 80370 \pm 7 \text{ (stat.)} \pm 11 \text{ (exp. syst.)} \pm 14 \text{ (mod. syst.)} \text{ MeV} \]
\[ = 80370 \pm 19 \text{ MeV}, \]
\[ m_{W^+} - m_{W^-} = -29 \pm 28 \text{ MeV}. \]
How is the W mass determined?
Lepton transverse momentum

The four-momentum of the neutrino is difficult to determine. Other variables have to be measured, e.g., the lepton $p_T$

Additional variables are missing $p_T$ and transverse mass $m_T$
Inclusion of corrections

If the $W$ were exactly collinear ($p_{TW}=0$, no TMD effects), the distribution of events would look like this

Detector effects cause further changes

If TMDs are taken into consideration, the distribution gets modified like this
Effect of changing $W$ mass

A change of 10 MeV in the $W$ mass ($M_W$) induces distortions at the per mille level only: challenging!

see, e.g., Bozzi, Rojo, Vicini, arXiv:1104.2056
Template fit technique

- Using Monte Carlo generators that include all known corrections, several high-statistics “templates” are produced with different $M_W$.

- The template that fits data best determines the value of $M_W$. 
How are sys. uncertainties estimated?

• The Monte Carlo generator is used to produce pseudodata with known $M_W$, but with some other differences (e.g., changing the PDF set).

• The template fit is applied to the pseudodata and the difference between the extracted $M_W$ and the input one is used to determine $\delta M_W$
### Uncertainties

<table>
<thead>
<tr>
<th>Kinematic distribution</th>
<th>( W^+ )</th>
<th>( W^- )</th>
<th>Combined</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \delta m_W ) [MeV]</td>
<td>( p_T^\ell )</td>
<td>( p_T^\ell )</td>
<td>( p_T^\ell )</td>
</tr>
<tr>
<td>Fixed-order PDF uncertainty</td>
<td>13.1</td>
<td>12.0</td>
<td>8.0</td>
</tr>
<tr>
<td>AZ tune</td>
<td>3.0</td>
<td>3.0</td>
<td>3.0</td>
</tr>
<tr>
<td>Charm-quark mass</td>
<td>1.2</td>
<td>1.2</td>
<td>1.2</td>
</tr>
<tr>
<td>Parton shower ( \mu_F ) with heavy-flavour decorrelation</td>
<td>5.0</td>
<td>5.0</td>
<td>5.0</td>
</tr>
<tr>
<td>Parton shower PDF uncertainty</td>
<td>3.6</td>
<td>2.6</td>
<td>1.0</td>
</tr>
<tr>
<td>Angular coefficients</td>
<td>5.8</td>
<td>5.8</td>
<td>5.8</td>
</tr>
<tr>
<td>Total</td>
<td>15.9</td>
<td>14.8</td>
<td>11.6</td>
</tr>
</tbody>
</table>

This contribution contains also intrinsic transverse momentum of partons. The MC has been tuned to describe Z-boson data.

ATLAS Collab. arXiv:1701.07240
Z and W production involve different flavor combinations.
Our work
Flavor contributions: $Z$ boson

$u \bar{u}$ and $d \bar{d}$ are the most important channels
Flavor contributions: $W^+$ boson

$u\,d\,\bar{u}$ is the most important channel
Monte Carlo generation

- DYRes code of Catani, de Florian, Ferrera, Grazzini (2015)
- We assume the conditions of LHC 7 TeV and ATLAS acceptance cuts
- The cross section involves Transverse Momentum Distributions (TMDs)

\[ f_1^a(x, k_\perp; \mu^2) = \frac{1}{2\pi} \int d^2b_\perp e^{-ib_\perp \cdot k_\perp} \tilde{f}_1^a(x, b_\perp; \mu^2) \]

\[ \tilde{f}_1^a(x, b_T; \mu^2) = \sum_i (\tilde{C}_{a/i} \otimes f_1^i)(x, b_*; \mu_b) e^{\tilde{S}(b_*; \mu_b, \mu)} f_1^{aNP}(x, b_T) \]

Perturbative parts at order \( \alpha_s \) — NLL
Nonperturbative parts

\[ f_1^{\alpha NP}(b_T^2) \propto e^{-g_{NP}^\alpha b_T^2} \]

\[ \exp(-g_{NP}^\alpha b_T^2) \longrightarrow \exp[-[g_{evo} \ln(Q^2/Q_0^2) + g_\alpha] b_T^2] \]

this component is flavor-independent (gluon radiation)

this component can be flavor-dependent

see, e.g., Bacchetta, Delcarro, Pisano, Radici, Signori, arXiv:1703.10157
Possible flavor dependence

Signori, Bacchetta, Radici, Schnell, arXiv:1409.3507

Our studies of SIDIS data indicate that there is a lot of room for flavor dependence. More flavor-sensitive data are needed!
Choice of nonperturbative parameters

We considered initially:

- **50 flavour-dependent sets** \( \{ g_{np}^{u}, g_{np}^{d}, g_{np}^{u}, g_{np}^{d}, g_{np}^{s} \} \) with \( g_{np}^{a} \in [0.2, 0.6] \) GeV

- **1 flavour-independent set** with \( g_{np}^{a} = 0.4 \) GeV

We selected the sets that give a description of Z boson data equivalent to the flavor-independent set ("Z-equivalent")

We then chose a few sets with interesting characteristics
Inclusion of flavor dependence

<table>
<thead>
<tr>
<th>Set</th>
<th>$u_v$</th>
<th>$d_v$</th>
<th>$u_s$</th>
<th>$d_s$</th>
<th>$s$</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>0.34</td>
<td>0.26</td>
<td>0.46</td>
<td>0.59</td>
<td>0.32</td>
</tr>
<tr>
<td>2</td>
<td>0.34</td>
<td>0.46</td>
<td>0.56</td>
<td>0.32</td>
<td>0.51</td>
</tr>
<tr>
<td>3</td>
<td>0.55</td>
<td>0.34</td>
<td>0.33</td>
<td>0.55</td>
<td>0.30</td>
</tr>
<tr>
<td>4</td>
<td>0.53</td>
<td>0.49</td>
<td>0.37</td>
<td>0.22</td>
<td>0.52</td>
</tr>
<tr>
<td>5</td>
<td>0.42</td>
<td>0.38</td>
<td>0.29</td>
<td>0.57</td>
<td>0.27</td>
</tr>
</tbody>
</table>

narrow, medium, large
narrow, large, narrow
large, narrow, large
large, medium, narrow
medium, narrow, large
Templates and pseudo data

**TEMPLATES**

- high statistics (750M events)
- different values of $M_W$
  $\Delta M_W = -15$ MeV to $+15$ MeV
- no flavor-dependent intrinsic transverse momentum

**PSEUDODATA**

- “low” statistics (75M events)
- central value $M_W = 80.385$ GeV
- flavor-dependent intrinsic transverse momentum
We compute the chi$^2$ between templates and pseudo data, find which template gives the best description and determine $\Delta M_W$
Resulting shifts

<table>
<thead>
<tr>
<th>Set</th>
<th>$\Delta M_{W^+}$</th>
<th>$\Delta M_{W^-}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>2</td>
<td>-2</td>
</tr>
<tr>
<td>2</td>
<td>1</td>
<td>-1</td>
</tr>
<tr>
<td>3</td>
<td>0</td>
<td>-3</td>
</tr>
<tr>
<td>4</td>
<td>1</td>
<td>-2</td>
</tr>
<tr>
<td>5</td>
<td>2</td>
<td>0</td>
</tr>
</tbody>
</table>
Conclusions

No flavor-blind analysis of $M_W$ should be allowed.

Data sensitive to the flavor dependence of TMDs are needed to reduce this uncertainty.