DIS on a polarized spin-1 target with spectator tagging

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JLab Theory Seminar

in collaboration with Ch. Weiss, M. Sargsian JLab LDRD project on spectator tagging



Light ions at EIC: physics objectives



Neutron structure

- flavor decomposition of quark PDFs/GPDs/TMDs
- flavor structure of the nucleon sea
- singlet vs non-singlet QCD evolution, leading/higher-twist effects



Bound nucleons in QCD

- medium modification of quark/gluon structure
- QCD origin of short-range nuclear force
- Imaging nuclear bound states



Coherence and saturation

 interaction of high-energy probe with coherent quark-gluon fields

Dedicated workshop at Ghent last month

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JLab seminar

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Tagged spectator DIS process with deuteron



- DIS off a nuclear target with a slow (relative to nucleus c.m.) nucleon detected in the final state
- Control nuclear configuration
- Advantages for the deuteron
 - ► simple NN system, non-nucleonic (△△) dof suppressed
 - active nucleon identified
 - recoil momentum selects nuclear configuration (medium modifications)
 - limited possibilities for nuclear FSI, calculable
- Wealth of possibilities to study (nuclear) QCD dynamics
- Will be possible in a wide kinematic range @ EIC (**polarized**)
- Suited for colliders: no target material, forward detection, transverse pol.

fixed target CLAS BONuS limited to recoil momenta $\sim 70~\text{MeV}$

Pole extrapolation for on-shell nucleon structure



 Allows to extract free neutron structure in a model independent way

- ► Recoil momentum p_R controls off-shellness of neutron $t' \equiv t m_N^2$
- Free neutron at pole $t m_N^2 \rightarrow 0$: "on-shell extrapolation"
- Small deuteron binding energy results in small extrapolation length
- Eliminates nuclear binding and FSI effects [Sargsian,Strikman PLB '05]
- D-wave suppressed at on-shell point \rightarrow neutron \sim 100% polarized
- Precise measurements of neutron (spin) structure at an EIC

Outline

Theoretical Formalism

- ► General expression of SIDIS for a polarized spin 1 target
 - ► Tagged spectator DIS is SIDIS in the target fragmentation region

$$\vec{e} + \vec{T} \rightarrow e' + X + h$$

- Dynamical model to express structure functions of the reaction
 - First step: impulse approximation (IA) model
 - FSI corrections (unpolarized)
- Light-front structure of the deuteron
 - Natural for high-energy reactions as off-shellness of nucleons in LF quantization remains finite
- Neutron structure with pole extrapolation for EIC
- Experimental apparatus

Polarized spin 1 particle

Spin state described by a 3*3 density matrix in a basis of spin 1 states polarized along the collinear virtual photon-target axis

$$W_D^{\mu\nu} = Tr[\rho_{\lambda\lambda'}W^{\mu\nu}(\lambda'\lambda)]$$

Characterized by 3 vector and 5 tensor parameters

$$\mathcal{S}^{\mu} = \langle \hat{W}^{\mu}
angle$$
, $T^{\mu
u} = rac{1}{2} \sqrt{rac{2}{3}} \langle \hat{W}^{\mu} \hat{W}^{
u} + \hat{W}^{
u} \hat{W}^{\mu} + rac{4}{3} \left(\mathcal{g}^{\mu
u} - rac{\hat{P}^{\mu} \hat{P}^{
u}}{M^2}
ight)
angle$

Split in longitudinal and transverse components

$$\rho_{\lambda\lambda'} = \frac{1}{3} \begin{bmatrix} 1 + \frac{3}{2}S_L + \sqrt{\frac{3}{2}}T_{LL} & \frac{3}{2\sqrt{2}}S_T e^{-i(\phi_h - \phi_S)} & \sqrt{\frac{3}{2}}T_{TT} e^{-i(2\phi_h - 2\phi_{T_T})} \\ & -\sqrt{3}T_{LT} e^{-i(\phi_h - \phi_{T_L})} & \\ \frac{3}{2\sqrt{2}}S_T e^{i(\phi_h - \phi_S)} & 1 - \sqrt{6}T_{LL} & \frac{3}{2\sqrt{2}}S_T e^{-i(\phi_h - \phi_S)} \\ & -\sqrt{3}T_{LT} e^{i(\phi_h - \phi_{T_L})} & & +\sqrt{3}T_{LT} e^{-i(\phi_h - \phi_{T_L})} \\ \sqrt{\frac{3}{2}}T_{TT} e^{i(2\phi_h - 2\phi_{T_T})} & \frac{3}{2\sqrt{2}}S_T e^{i(\phi_h - \phi_S)} & 1 - \frac{3}{2}S_L + \sqrt{\frac{3}{2}}T_{LL} \\ & +\sqrt{3}T_{LT} e^{i(\phi_h - \phi_{T_L})} & \end{bmatrix}$$

Spin 1 SIDIS: General structure of cross section

To obtain structure functions, enumerate all possible tensor structures that obey hermiticity and transversality condition (qW = Wq = 0)
 Cross section has 41 structure functions.

$$\frac{d\sigma}{dx dQ^2 d\phi_{l'}} = \frac{y^2 \alpha^2}{Q^4 (1-\epsilon)} \left(F_U + F_S + F_T\right) d\Gamma_{P_h} \,,$$

 \blacktriangleright U + S part identical to spin 1/2 case [Bacchetta et al.]HEP ('07)]

$$F_{U} = F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} + \epsilon \cos 2\phi_h F_{UU}^{\cos2\phi_h} + \frac{h}{\sqrt{2\epsilon(1-\epsilon)}} \sin\phi_h F_{LU}^{\sin\phi_h}$$

$$\begin{split} F_{S} &= S_{L} \left[\sqrt{2\epsilon(1+\epsilon)} \sin \phi_{h} F_{US_{L}}^{\sin \phi_{h}} + \epsilon \sin 2\phi_{h} F_{US_{L}}^{\sin 2\phi_{h}} \right] \\ &+ S_{L} h \left[\sqrt{1-\epsilon^{2}} F_{LS_{L}} + \sqrt{2\epsilon(1-\epsilon)} \cos \phi_{h} F_{LS_{L}}^{\cos \phi_{h}} \right] \\ &+ S_{\perp} \left[\sin(\phi_{h} - \phi_{S}) \left(F_{US_{T},T}^{\sin(\phi_{h} - \phi_{S})} + \epsilon F_{US_{T},L}^{\sin(\phi_{h} - \phi_{S})} \right) + \epsilon \sin(\phi_{h} + \phi_{S}) F_{US_{T}}^{\sin(\phi_{h} + \phi_{S})} \\ &+ \epsilon \sin(3\phi_{h} - \phi_{S}) F_{US_{T}}^{\sin(3\phi_{h} - \phi_{S})} + \sqrt{2\epsilon(1+\epsilon)} \left(\sin \phi_{S} F_{US_{T}}^{\sin \phi_{S}} + \sin(2\phi_{h} - \phi_{S}) F_{US_{T}}^{\sin(2\phi_{h} - \phi_{S})} \right) \right] \\ &+ S_{\perp} h \left[\sqrt{1-\epsilon^{2}} \cos(\phi_{h} - \phi_{S}) F_{LS_{T}}^{\cos(\phi_{h} - \phi_{S})} + \\ & \sqrt{2\epsilon(1-\epsilon)} \left(\cos \phi_{S} F_{LS_{T}}^{\cos \phi_{S}} + \cos(2\phi_{h} - \phi_{S}) F_{LS_{T}}^{\cos(2\phi_{h} - \phi_{S})} \right) \right] , \end{split}$$

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 Cross section has 41 structure functions,

$$rac{d\sigma}{dx dQ^2 d\phi_{l'}} = rac{y^2 lpha^2}{Q^4 (1-\epsilon)} \left(F_U + F_S + F_T
ight) d\Gamma_{P_h}$$
 ,

> 23 SF unique to the spin 1 case (tensor pol.), 4 survive in inclusive (b_{1-4}) [Hoodbhoy, Jaffe, Manohar PLB'88]

$$\begin{split} F_{T} &= T_{LL} \left[F_{UT_{LL},T} + \epsilon F_{UT_{LL},L} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_{h} F_{UT_{LL}}^{\cos \phi_{h}} + \epsilon \cos 2\phi_{h} F_{UT_{LL}}^{\cos 2\phi_{h}} \right] \\ &+ T_{LL} h \sqrt{2\epsilon(1-\epsilon)} \sin \phi_{h} F_{LT_{LL}}^{\sin \phi_{h}} \\ &+ T_{L\perp} [\cdots] + T_{L\perp} h [\cdots] \\ &+ T_{L\perp} \left[\cos(2\phi_{h} - 2\phi_{T_{\perp}}) \left(F_{UT_{TT},T}^{\cos(2\phi_{h} - 2\phi_{T_{\perp}})} + \epsilon F_{UT_{TT},L}^{\cos(2\phi_{h} - 2\phi_{T_{\perp}})} \right) \right. \\ &+ \epsilon \cos 2\phi_{T_{\perp}} F_{UT_{TT}}^{\cos 2\phi_{T_{\perp}}} + \epsilon \cos(4\phi_{h} - 2\phi_{T_{\perp}}) F_{UT_{TT}}^{\cos(4\phi_{h} - 2\phi_{T_{\perp}})} \\ &+ \sqrt{2\epsilon(1+\epsilon)} \left(\cos(\phi_{h} - 2\phi_{T_{\perp}}) F_{UT_{TT}}^{\cos(\phi_{h} - 2\phi_{T_{\perp}})} + \cos(3\phi_{h} - 2\phi_{T_{\perp}}) F_{UT_{TT}}^{\cos(3\phi_{h} - 2\phi_{T_{\perp}})} \right) \right] \\ &+ T_{\perp \perp} h [\cdots] \end{split}$$

Tagged DIS with deuteron: model for the IA



 Hadronic tensor can be written as a product of nucleon hadronic tensor with deuteron light-front densities

$$W_D^{\mu\nu}(\lambda',\lambda) = 4(2\pi)^3 \frac{\alpha_R}{2-\alpha_R} \sum_{i=U,z,x,y} W_{N,i}^{\mu\nu} \rho_D^i(\lambda',\lambda) ,$$

 $\begin{aligned} & \text{All SF can be written as} \\ F_{ij}^k = \{ \text{kin. factors} \} \times \{ F_{1,2}(\tilde{x}, Q^2) \text{or } g_{1,2}(\tilde{x}, Q^2) \} \times \{ \text{bilinear forms} \\ & \text{in deuteron radial wave function } U(k), W(k) \} \end{aligned}$

• In the IA the following structure functions are $\mathbf{zero} \rightarrow \mathbf{sensitive}$ to FSI

- beam spin asymmetry $[F_{LU}^{\sin \phi_h}]$
- target vector polarized single-spin asymmetry [8 SFs]
- target tensor polarized double-spin asymmetry [7 SFs]

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Deuteron light-front wave function



- Up to momenta of a few 100 MeV dominated by NN component
- Can be evaluated in LFQM [Coester,Keister,Polyzou et al.] or covariant Feynman diagrammatic way [Frankfurt,Sargsian,Strikman]
- One obtains a Schrödinger (non-rel) like eq. for the wave function components, rotational invariance recovered
- Light-front WF obeys baryon and momentum sum rule

$$\Psi_{\lambda}^{D}(\boldsymbol{k}_{f},\lambda_{1},\lambda_{2}) = \sqrt{E_{k_{f}}} \sum_{\lambda_{1}^{\prime}\lambda_{2}^{\prime}} \mathcal{D}_{\lambda_{1}\lambda_{1}^{\prime}}^{\frac{1}{2}} [R_{fc}(k_{1_{f}}^{\mu}/m_{N})] \mathcal{D}_{\lambda_{2}\lambda_{2}^{\prime}}^{\frac{1}{2}} [R_{fc}(k_{2_{f}}^{\mu}/m_{N})] \Phi_{\lambda}^{D}(\boldsymbol{k}_{f},\lambda_{1}^{\prime},\lambda_{2}^{\prime})$$

- Differences with non-rel wave function:
 - appearance of the Melosh rotations to account for light-front quantized nucleon states
 - ▶ k_f is the relative 3-momentum of the nucleons in the light-front boosted rest frame of the free 2-nucleon state (so not a "true" kinematical variable)

Unpolarized structure function



- Extrapolation for $(m_N^2 t) \rightarrow 0$ corresponds to on-shell neutron $F_{2N}(x, Q^2)$, here equivalent to imaginary p_s
- Clear effect of deuteron D-wave, largest in the region dominated by the tensor part of the *NN*-interaction
- D-wave drops out at the on-shell point

Tagging: free neutron structure



C. Hyde talk

*F*_{2n} extracted with percent-level accuracy at x < 0.1

- Uncertainty mainly systematic due to intrinsic momentum spread in beam (JLab LDRD project: detailed estimates)
- In combination with proton data non-singlet $F_{2p} F_{2n}$, sea quark flavor asymmetry $\bar{d} \bar{u}$

On-shell extrapolation of double spin asymm. $A_{||} = \frac{\sigma(++) - \sigma(-+) - \sigma(+-) + \sigma(--)}{\sigma(++) + \sigma(-+) + \sigma(--)} [\phi_h \text{avg}] = \frac{F_{LS_L}}{F_T + \epsilon F_L} = D \frac{g_{1n}}{F_{1n}} + \cdots$



- Again clear contribution from D-wave at finite recoil momenta
- Relativistic nuclear effects through Melosh rotations, grow with recoil momenta
- Both effects drop out near the on-shell extrapolation point

Tagging: polarized neutron structure

On-shell extrapolation of double spin asymm.

$$A_{||} = \frac{\sigma(++) - \sigma(-+) - \sigma(+-) + \sigma(--)}{\sigma(++) + \sigma(-+) + \sigma(--)} [\phi_h \text{avg}] = \frac{F_{LS_L}}{F_T + \epsilon F_L} = D \frac{g_{1n}}{F_{1n}} + \cdots$$



JLab LDRD arXiv:1407.3236, arXiv:1409.5768

 Systematic uncertainties cancel in ratio (momentum smearing, resolution effects)

Statistics requirements

- Physical asymmetries
 ~ 0.05 0.1
- Effective polarization
 P_eP_D ~ 0.5
- Luminosity required $\sim 10^{34} \text{ cm}^{-2} \text{s}^{-1}$
- Precise measurement of neutron spin structure

Tagging: polarized neutron structure II



Precise measurement of neutron spin structure

- separate leading- /higher-twist
- non-singlet/singlet QCD evolution
- ▶ pdf flavor separation Δu , Δd . ΔG through singlet evolution
- non-singlet $g_{1p} g_{1n}$ and Bjorken sum rule

Tagging: EMC effect



- Medium modification of nucleon structure embedded in nucleus (EMC effect)
 - dynamical origin?
 - caused by which momenta/distances in nuclear WF
 - spin-isospin dependence?

tagged EMC effect

- recoil momentum as extra handle on medium modification (off-shellness, size of nuclear configuration) away from the on-shell pole
- EIC: Q² evolution, gluons, spin dependence!
- Interplay with final-state interactions!
 - use $\tilde{x} = 0.2$ to constrain FSI
 - constrain medium modification at higher x

Final-state interactions: three regimes



DIS regime, intermediate x



- Shadowing in inclusive DIS $x \ll 10^{-1}$
 - Diffractive DIS on single nucleon (leading twist, HERA)
 - Interference of DIS on nucleon 1 and 2
 - Calculable in terms of nucleon diffractive structure functions [Gribov 70s, Frankfurt, Guzey, Strikman '02+]
- FSI between slow hadrons from the DIS products and spectator nucleon, fast hadrons hadronize after leaving the nucleus.
 - Data show slow hadrons in the target fragmentation region are mainly nucleons.
 - ▶ Input needed from nucleon target fragmentation data \rightarrow also possible at EIC
 - M. Strikman, Ch. Weiss arXiv: 1706.02244
- rescattering of resonance-like structure with spectator nucleon in eikonal approximation [Deeps,BONuS].
 - WC,M. Sargsian arXiv:1704.06117

FSI 1: Shadowing at small *x* in tagged DIS





- Explore shadowing through recoil momentum dependence
- Shadowing enhanced in tagged DIS compared to inclusive
 - enhancement factor from AGK rules
 - shadowing term drops slower with *p_R* than IA
- FSI contributions between slow p and n in diffractive events
- Large FSI effects in diffractive amplitudes (~ 40%), also at zero spectator momenta due to orthogonality of *np* state to deuteron
- Effects smaller in tagged as diffractive are ~ 10% of total events
- Possibilities to study diffractive events by double tagging

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FSI 2: intermediate x model





Strikman, Weiss, arXiv:1706.02244, PRC accepted

- Features of the FSI of slow hadrons with spectator nucleon are similar to what is seen in quasi-elastic deuteron breakup.
- Inclusion FSI diagram adds two contributions: FSI term (~ absorption, negative) and FSI² term (~ refraction, postive)
- At low momenta (p_r < 200 MeV) FSI term dominates, at larger momenta FSI² dominates.
- Both contributions vanish at the pole → pole extrapolation still feasible
- Calculation with realistic deuteron wf (AV18)

JLEIC full-acceptance detector



Forward detector integrated in interaction region & beam optics

- Good acceptance for elastic recoil Rigidity same as beam. Large dispersion generated *after* IP Longitudinal momentum up to 99.5% of beam, angles down to 2 mrad (10 σ)
- Good acceptance for spectators and ion fragments
 Rigidity different from beam. Large magnet apertures, small gradients
- Good momentum and angular resolution Longitudinal $dp/p \sim 4 \times 10^{-4}$, angular $\delta \theta \sim 0.2$ mrad $p_{TR} \sim 15 \text{ MeV}/c$ resolution for tagged 50 GeV/A deuterium beam

JLEIC: Momentum spread in beam





[[]Ch. Hyde, K. Park et al.]

- Intrinsic beam spread in ion beam "smears" recoil momentum
 - ► transverse momentum spread of $\sigma \approx 20$ MeV ($\delta \sigma / \sigma \sim 10\%$)
 - p_R (measured) $\neq p_R$ (vertex)
 - Systematic correlated uncertainty, x,Q² independent
- Dominant syst. uncertainty at JLEIC, detector resolution much higher than beam momentum spread (diff for eRHIC)
- On-shell extrapolation feasible!!

Tagging: developments and extensions

Final-state interactions

- in tagged $\vec{e} + \vec{D}$
- ► maximized/minimized by choice of kinematics. Constrain FSI models.
- azimuthal and spin observables non-zero through FSI
- Tagging with complex nuclei A > 2
 - ▶ isospin dependence, universality of bound nucleon structure
 - ► A − 1 ground state recoil

Resolved final states: SIDIS on neutron, hard exclusive channels

R&D project at JLAB

- Develop simulation tools (physics models, event generators, analysis tools) for DIS on light ions with spectator tagging at MEIC and study physics impact.
- ran FY14-15

D. Higinbotham, W. Melnitchouk, P. Nadel-Turonski, K. Park, C. Weiss (JLab), Ch. Hyde (ODU), M. Sargsian (FIU), V. Guzey (PNPI), with collaborators W. Cosyn (Ghent), S. Kuhn (ODU), M. Strikman (PSU), Zh. Zhao (JLab)

- Tools, documentation, results publicly available. Open for collaboration!
- More info:

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https://www.jlab.org/theory/tag/
arXiv:1407.3236, arXiv:1409.5768v1, arXiv:1601.066665,
arXiv:1609.01970
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Conclusions

- General form of SIDIS with a spin 1 target, 23 tensor polarized structure functions unique to spin 1
- Results for the impulse approximation using deuteron light-front structure, relativistic nuclear spin effects contribute.
- FSI/shadowing effects calculable
- Spectator tagging in *eD* scattering with EIC enables next-generation measurements with maximal control and unprecedented accuracy
 - Neutron structure functions, including spin
 - Nuclear modifications of quark/gluon structure
- Extensions:
 - ► Tagging with A > 2: isospin dependence, universality of bound nucleon structure ; A 1 recoil
 - Coherent processes: nuclear GPDs
 - ▶ Resolved final states: SIDIS on neutron, hard exclusive channels