Monte Carlo Analysis of TMD Factorization in SIDIS

Ross Dempsey Wally Melnitchouk, Nobuo Sato



July 27, 2018

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Semi-Inclusive Deep Inelastic Scattering

- ► Deep inelastic scattering (DIS): $e^- + p \rightarrow e^- + X$
- In the parton model, e⁻ interacts with a single quark in the proton



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Semi-Inclusive Deep Inelastic Scattering

- ► Deep inelastic scattering (DIS): $e^- + p \rightarrow e^- + X$
- In the parton model, e⁻ interacts with a single quark in the proton
- Semi-inclusive deep inelastic scattering (SIDIS): $e^- + p \rightarrow e^- + h + X$



Figure: In DIS, an electron interacts with one of the quarks in a proton via a virtual photon.

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Semi-Inclusive Deep Inelastic Scattering

- ► Deep inelastic scattering (DIS): $e^- + p \rightarrow e^- + X$
- In the *parton* model, e⁻ interacts with a single quark in the proton
- Semi-inclusive deep inelastic scattering (SIDIS): $e^- + p \rightarrow e^- + h + X$
- Typical case: $h = \pi^+$

Figure: In DIS, an electron interacts with one of the quarks in a proton via a virtual photon.



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SIDIS Variables

 DIS is described by variables related to the electron-quark scattering:

$$Q^2 = -(k-k')^2 \equiv -q^2 =$$
 photon virtuality
 $x = rac{Q^2}{2p \cdot q} =$ quark momentum fraction

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 SIDIS requires additional variables to describe the final product:

> $P_T = \text{transverse momentum of } h$ $\eta = \tanh^{-1} \frac{P_z(h)}{E(h)} = \text{rapidity of } h$

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$$P_T = \text{transverse momentum of } h$$

 $\eta = \tanh^{-1} \frac{P_z(h)}{E(h)} = \text{rapidity of } h$

► The rapidity η is related to the more familiar z = P_h·q / p·q, but is more useful for delineating kinematic regions (Boglione et al. 2017)

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Parton Distributions

- Hadrons are built from partons (quarks and gluons)
- Parton structure of hadrons characterized through:



Figure: The CJ15 parton distribution functions for the proton.

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 - Longitudinal momentum. The parton distribution function (PDF) records the fraction of longitudinal momentum that is found in a parton species



Figure: The CJ15 parton distribution functions for the proton.

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Parton Distributions

- Hadrons are built from partons (quarks and gluons)
- Parton structure of hadrons characterized through:
 - Longitudinal momentum. The parton distribution function (PDF) records the fraction of longitudinal momentum that is found in a parton species
 - Transverse momentum. Various TMD-PDFs determine how transverse momentum is distributed among the partons in the hadron



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SIDIS Regions

 TMD-PDFs can be extracted from the current fragmentation region, with small p_T and large negative η (Boglione et al. 2017)



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- How do the SIDIS regions shift with energy?
- Can we distinguish the current fragmentation region from other regions at JLab energies?



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Pythia

- Pythia 8 is a Monte Carlo event generator for high-energy collisions of particles
- Using a combination of Standard Model theory and phenomenological models constrained by data, produces realistic collision events at sufficiently high energies



Figure: Events generated with $Q^2 \ge 1 \text{ GeV}^2$, $W^2 \ge 5 \text{ GeV}^2$, and $E_{cm} = 18 \text{ GeV}$ (COMPASS energy).

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Rapidity Plateau I

 \blacktriangleright The overall distribution of events in rapidity shows a plateau, with a spike at $\eta=0$ for soft and collinear events



Figure: Rapidity distribution at $\sqrt{s} = 100 \text{ GeV}$.

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- At low W, dominated by central and target fragmentation; at high W, current fragmentation becomes dominant (Boglione et al. 2017)
- This contributes to a widening of the rapidity distribution at sufficiently high W

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Evolution of Rapidity Distribution

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Evolution of Rapidity- p_T Distribution

 The splitting of the regions can be seen in rapidity-p_T space

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Region Splitting

- At sufficiently high W, we can see two peaks in the rapidity distribution, and measure the separation between them
- This separation has been used experimentally as a proxy for the isolation of the current region (Berger criterion)



Figure: Rapidity spread is erratic at low W.

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Region Splitting

- At sufficiently high W, we can see two peaks in the rapidity distribution, and measure the separation between them
- This separation has been used experimentally as a proxy for the isolation of the current region (Berger criterion)



Figure: Rapidity spread eventually grows linearly with W.

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SIDIS Experiments

 Different experiments vary widely in the amount of region separation they can expect



Figure: Higher energy SIDIS experiments will see clearer separation of regions.

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 The unpolarized SIDIS cross section takes the form Anselmino et al. 2014

$$\frac{d\sigma^{\text{SIDIS}}}{dx\,dQ^2\,dz\,dP_T^2} = \frac{2\pi^2\alpha^2}{(xs)^2}\frac{1+(1-y)^2}{y^2}\sum_q e_q^2A_q,$$

where

$$A_q = \int d^2 p_T \, d^2 k_T \, \delta(P_T - zk_T - p_T) f_{q/p}(x, k_T) D_{h/q}(z, p_T)$$

► This expresses the final-state P_T as a combination of the quark transverse momentum k_T and the hadron momentum p_T Monte Carlo Analysis of SIDIS

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Factorization II

We assume, following Anselmino et al. 2014, that the PDFs and fragmentation functions factorize:

$$f_{q/p}(x,k_T) = f_{q/p}(x) \frac{e^{-k_T^2/\langle k_T^2 \rangle}}{\pi \langle k_T^2 \rangle},$$
$$D_{h/q}(z,p_T) = D_{h/q}(z) \frac{e^{-k_T^2/\langle p_T^2 \rangle}}{\pi \langle p_T^2 \rangle}$$

By the convolution theorem,

$$A_q = f_{q/p}(x)D_{h/q}(z)\frac{e^{-P_T^2/\langle P_T^2\rangle}}{\pi\langle P_T^2\rangle}$$

where $\langle P_T^2 \rangle = \langle p_T^2 \rangle + z^2 \langle k_T^2 \rangle$

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At fixed x, Q², and z, this simple factorization model predicts

$$\frac{d\sigma^{\rm SIDIS}}{dx\,dQ^2\,dz\,dP_T^2} \propto \exp\left(-\frac{P_T^2}{\langle P_T^2\rangle}\right)$$

► Fitting this distribution at various values of z, we should see a quadratic increase in the parameter (P²_T)

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Model Fit II

• Anselmino et al. find this increase in COMPASS data with $\langle p_T^2 \rangle = 0.12 \,\text{GeV}^2$ and $\langle k_T^2 \rangle = 0.68 \,\text{GeV}^2$.



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Figure: The parameter $\langle P_T^2\rangle$ is given by $\langle p_T^2\rangle+z^2\langle k_T^2\rangle$ with $\langle p_T^2\rangle=0.11\,{\rm GeV}^2$ and $\langle k_T^2\rangle=0.68\,{\rm GeV}^2.$

- Pythia is an effective tool for exploring SIDIS kinematic regions
- SIDIS regions are sensitive to energy, and isolating the current fragmentation region may require high energy
- Simple factorization schemes, e.g. Anselmino et al. 2014, show promising agreement with Monte Carlo data

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- SIDIS regions are sensitive to energy, and isolating the current fragmentation region may require high energy
- Simple factorization schemes, e.g. Anselmino et al.
 2014, show promising agreement with Monte Carlo data
- Ongoing work:
 - Evaluate the validity of factorization across rapidity-p_T space
 - Suggest improvements to the Berger criterion for isolating the current fragmentation region

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