

Extracting Resonances in Eta-Pi Diffractive Production

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Jefferson Lab
Joint Physics Analysis Center

Cake Seminar
JLab, May 2018

Introduction

Joint Physics Analysis Center (JPAC)

Exotic mesons

Reaction $\pi^- p \rightarrow \eta \pi^- p$

Past:

Extracted $a_2(1320)$ and $a_2(1700)$
pole position

Jackura et al (JPAC), PLB774, arXiv:1707.02848

Present:

Extraction of exotic meson pole position

A. Rodas, A. Pilloni et al (JPAC) in preparation

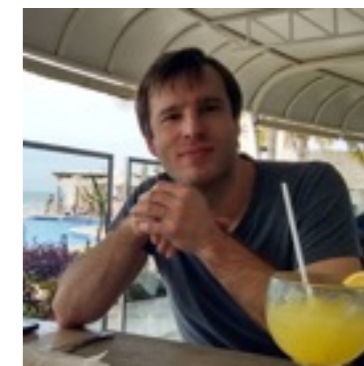
Future:

Implementation of DR constraining model
Transposition to GlueX/CLAS12 data

VM et al (JPAC), work in progress

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Cracow P.U.



Arkaitz Rodas Bilbao
Madrid U.



Viktor Mokeev
JLab



Miguel Albaladejo
Murcia U.

0^{--} 0^{-+} 0^{+-} 0^{++}

J^{PC} $q\bar{q}$ allowed

Ordinary

1^{--} 1^{-+} 1^{+-} 1^{++}



2^{--} 2^{-+} 2^{+-} 2^{++}

J^{PC} $q\bar{q}$ not allowed

Hybrid

3^{--} 3^{-+} 3^{+-} 3^{++}



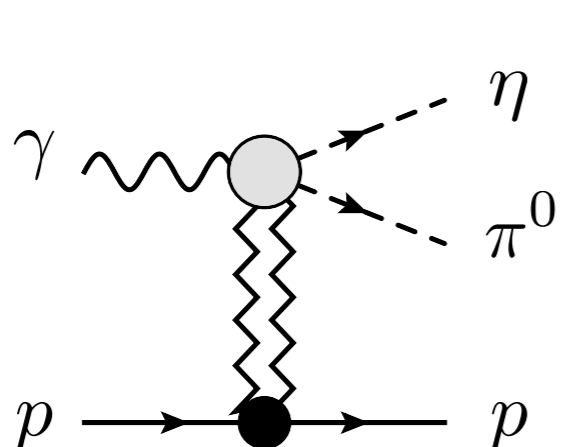
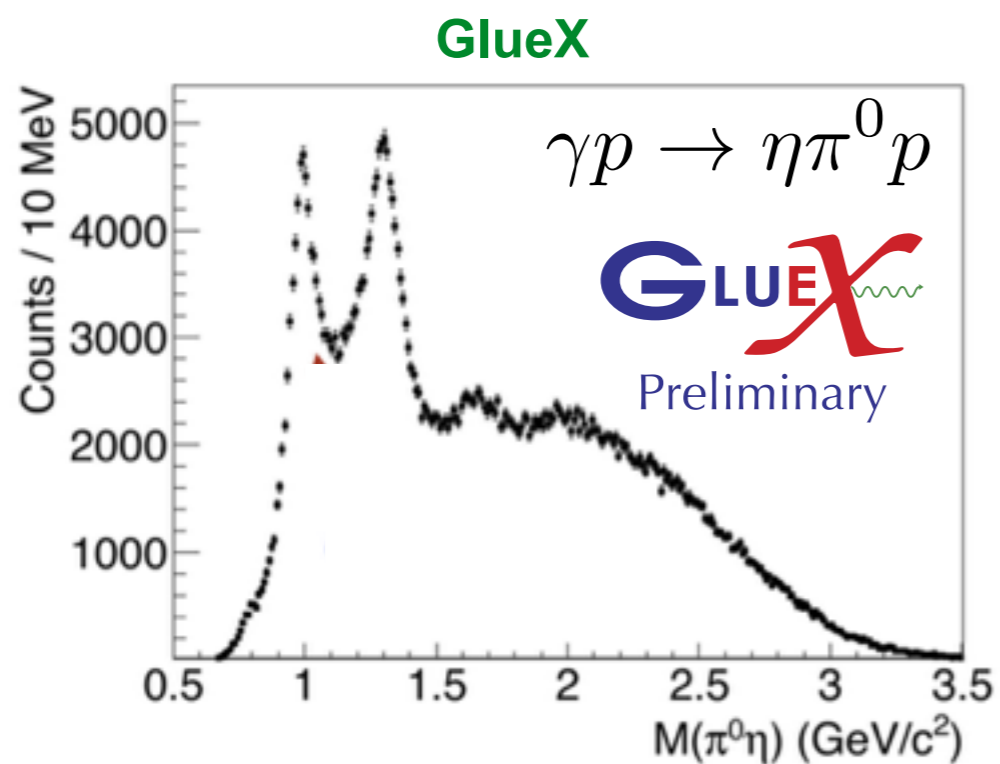
4^{--} 4^{-+} 4^{+-} 4^{++}

coupled to $\eta\pi$ and $\eta'\pi$
(isospin 1)

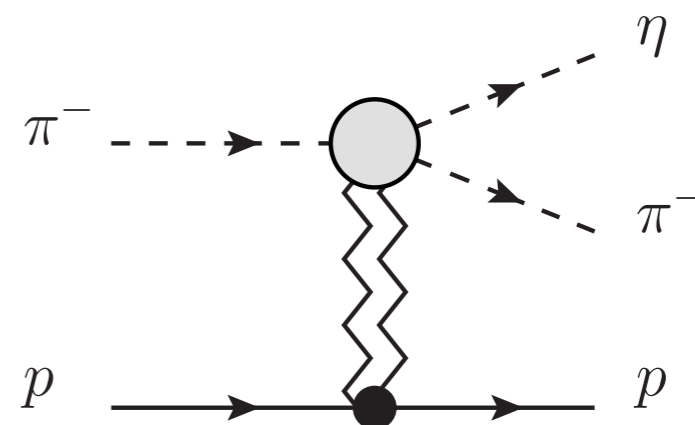
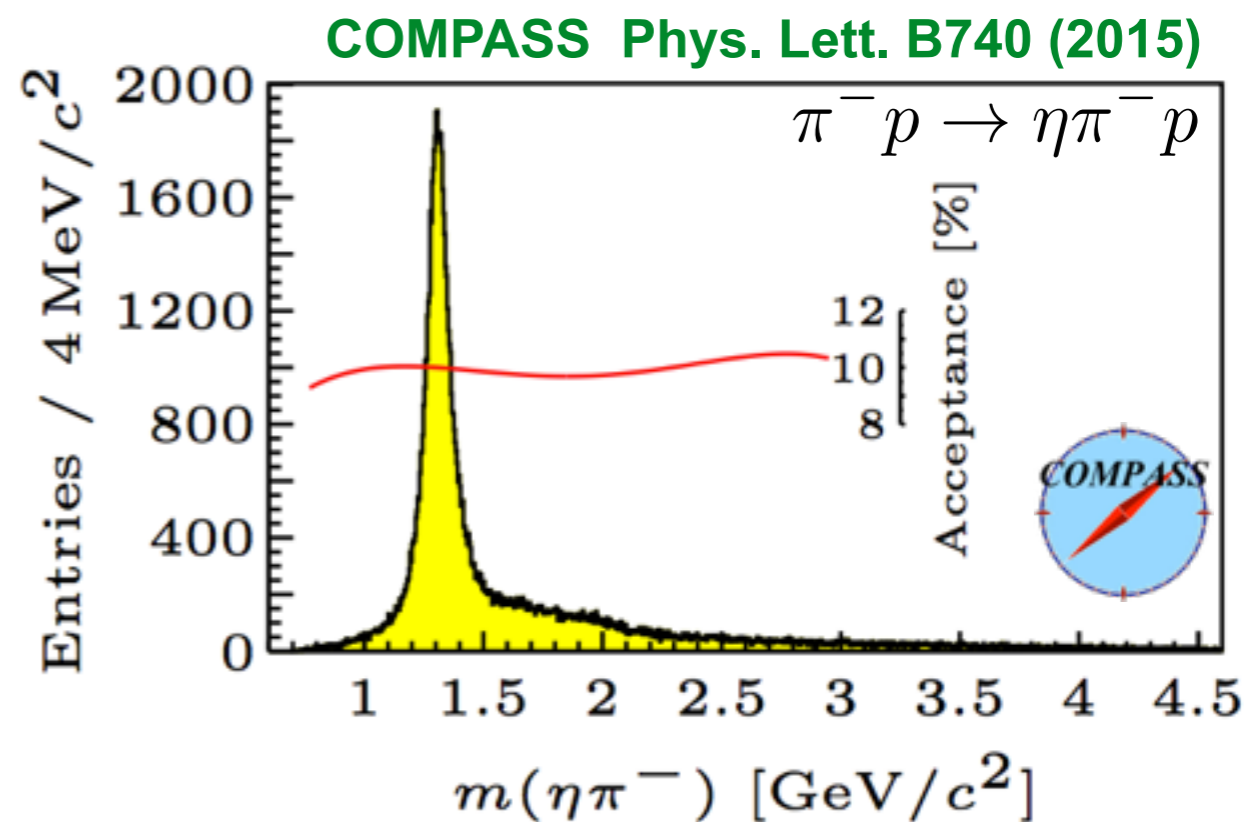
Quantum numbers filter ordinary mesons

Easier identification of **hybrid mesons with exotic quantum numbers**

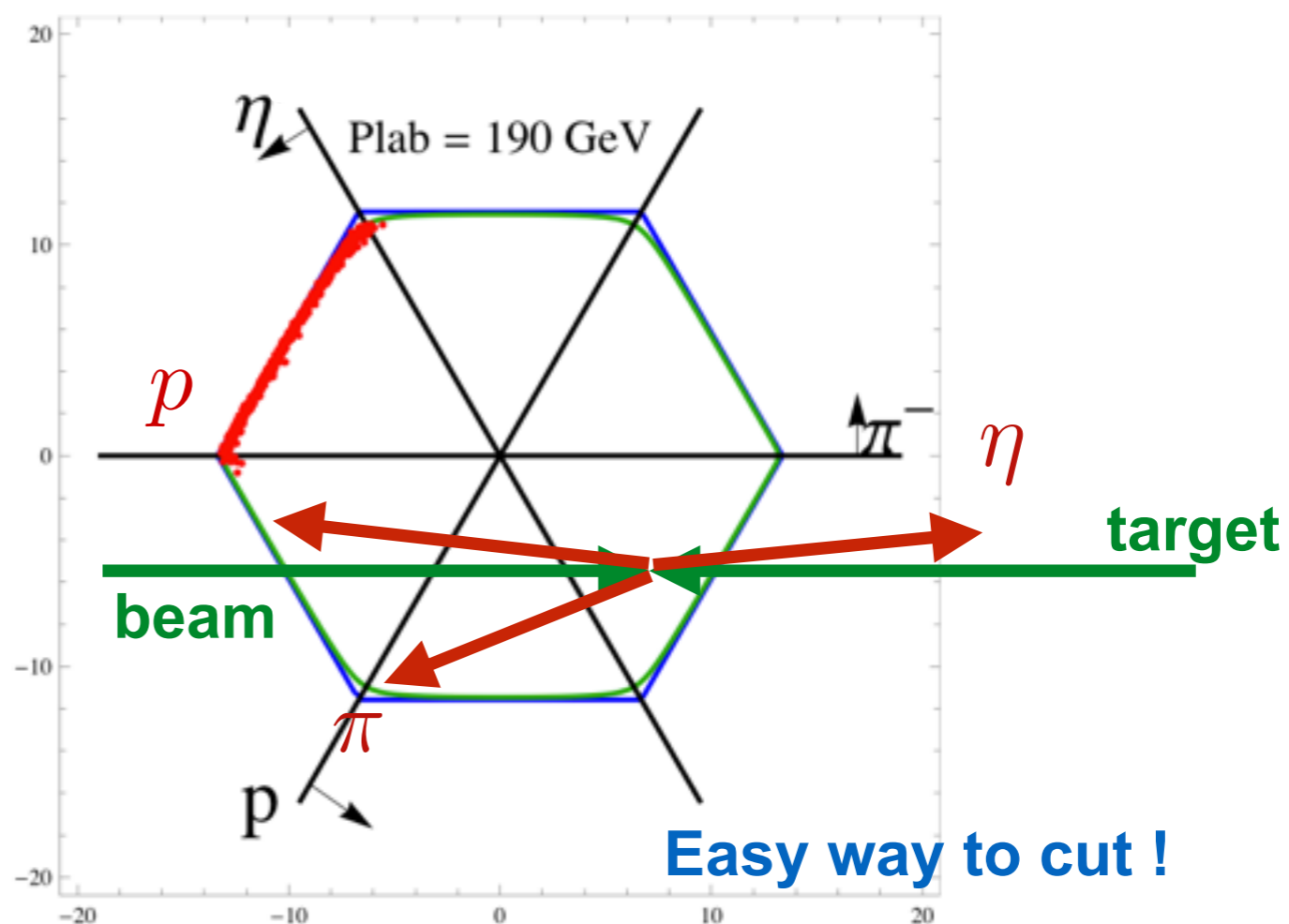
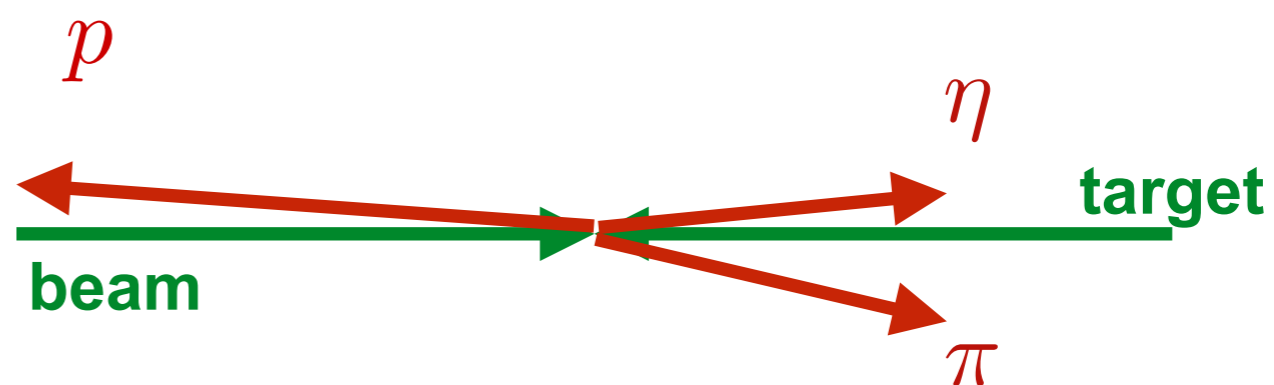
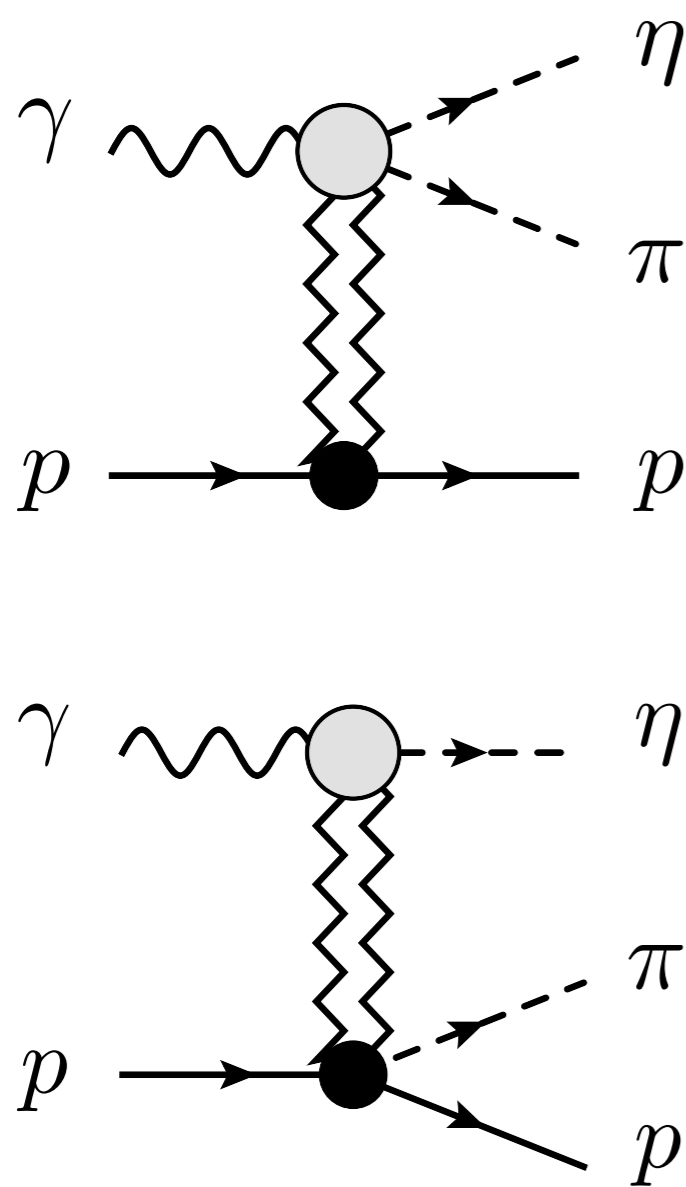
$$E_{\text{beam}} = 9 \text{ GeV}$$



$$E_{\text{beam}} = 190 \text{ GeV}$$

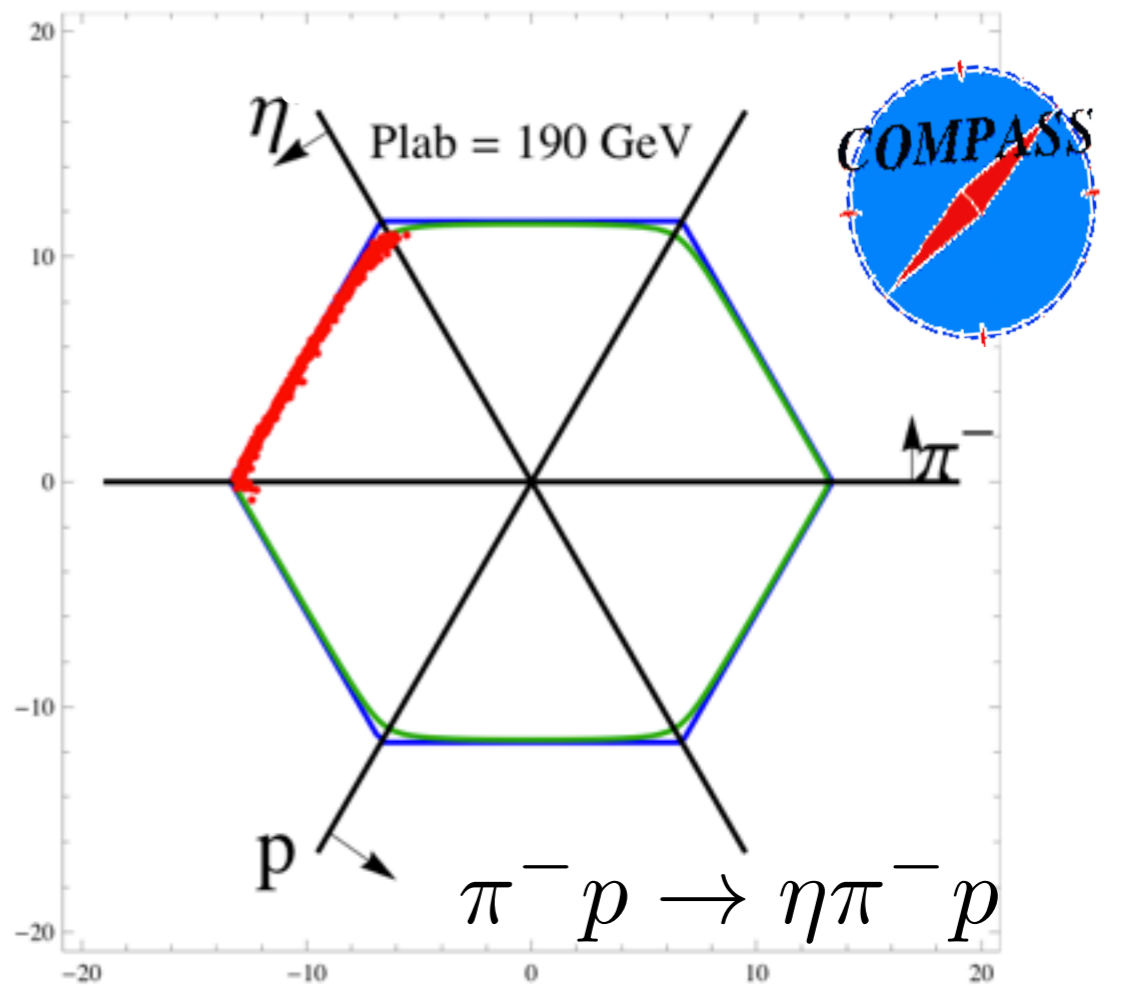
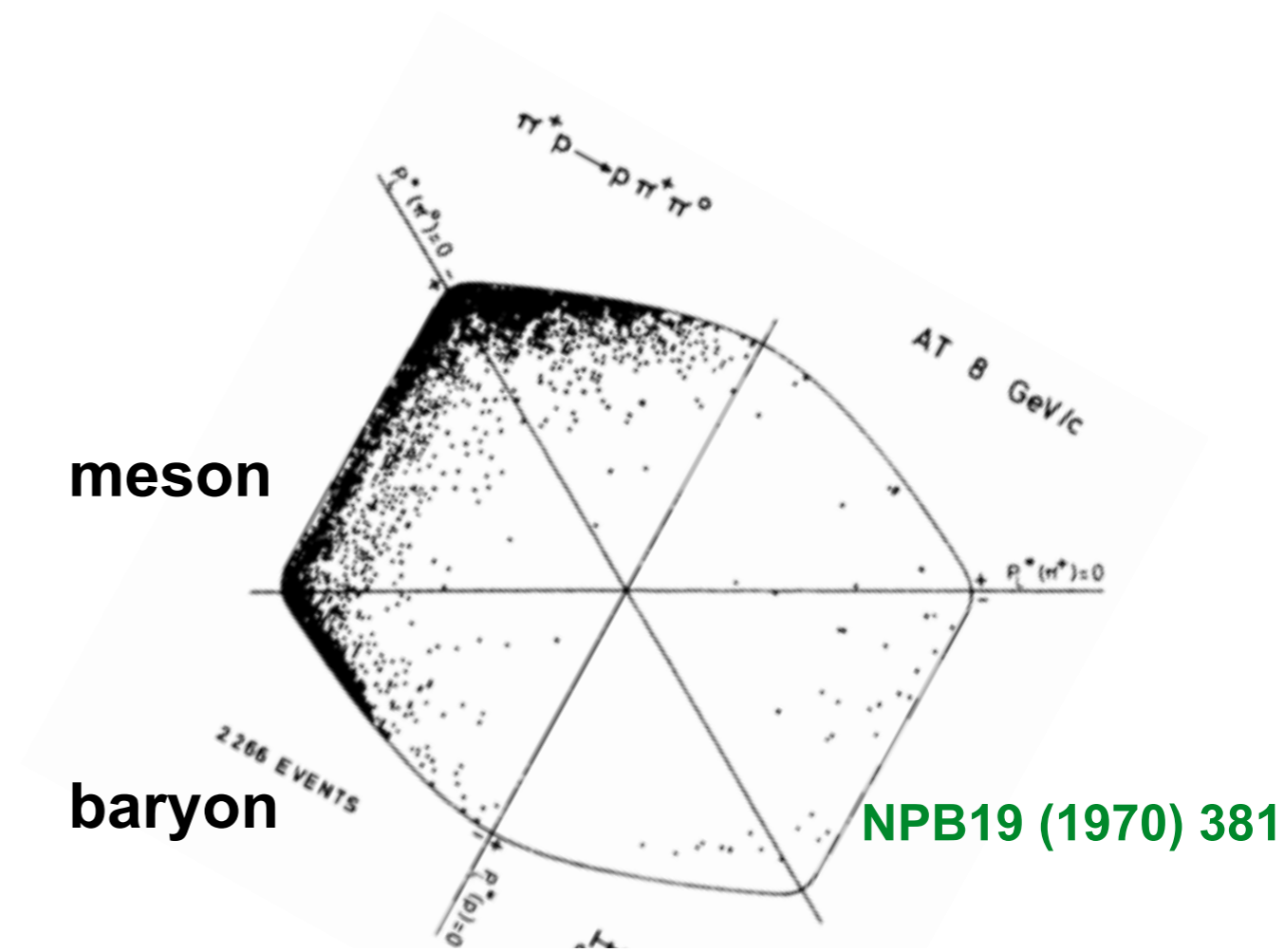
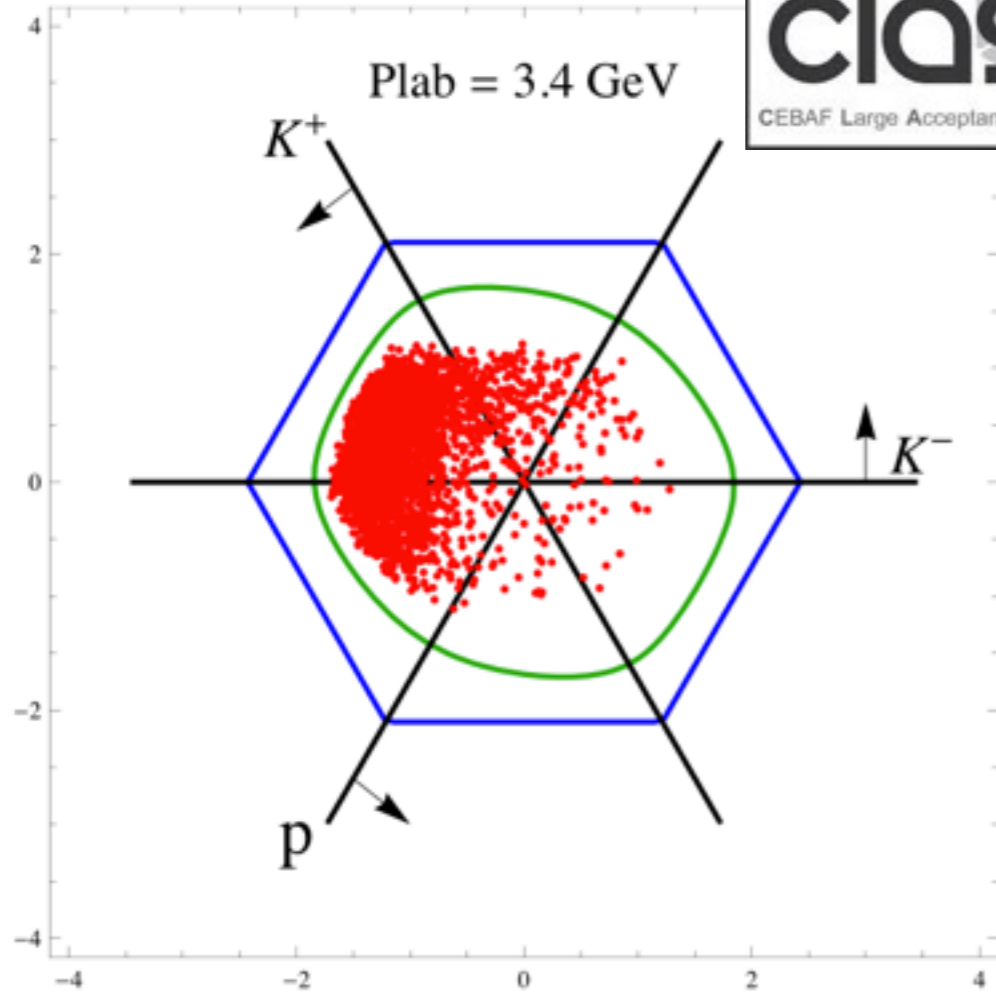


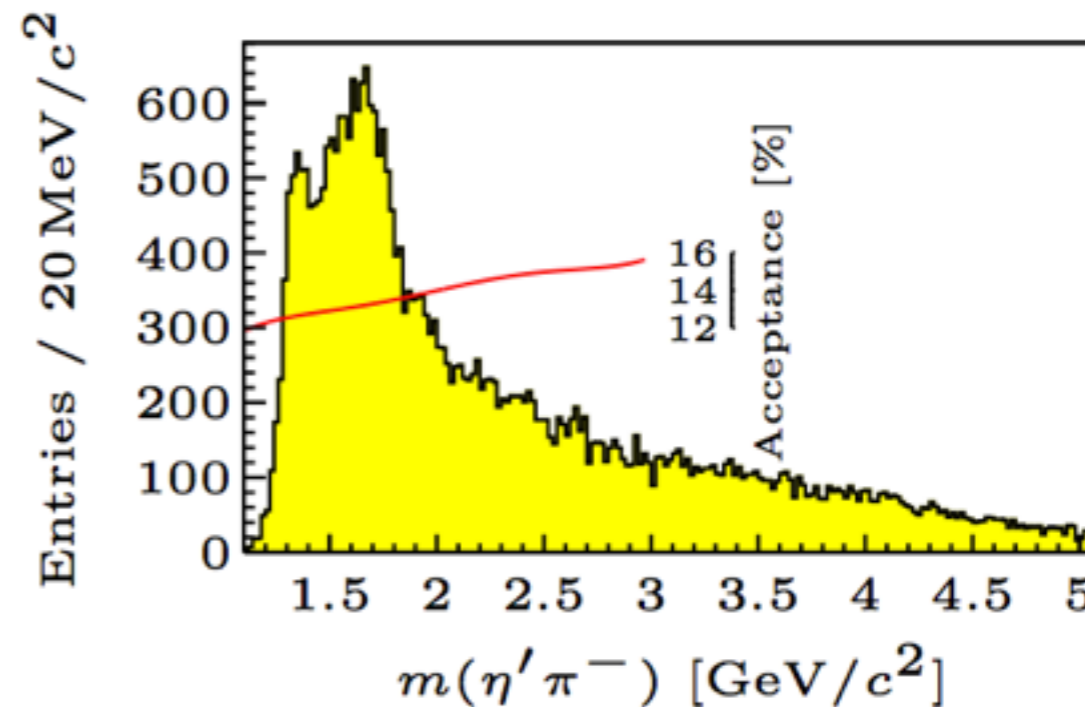
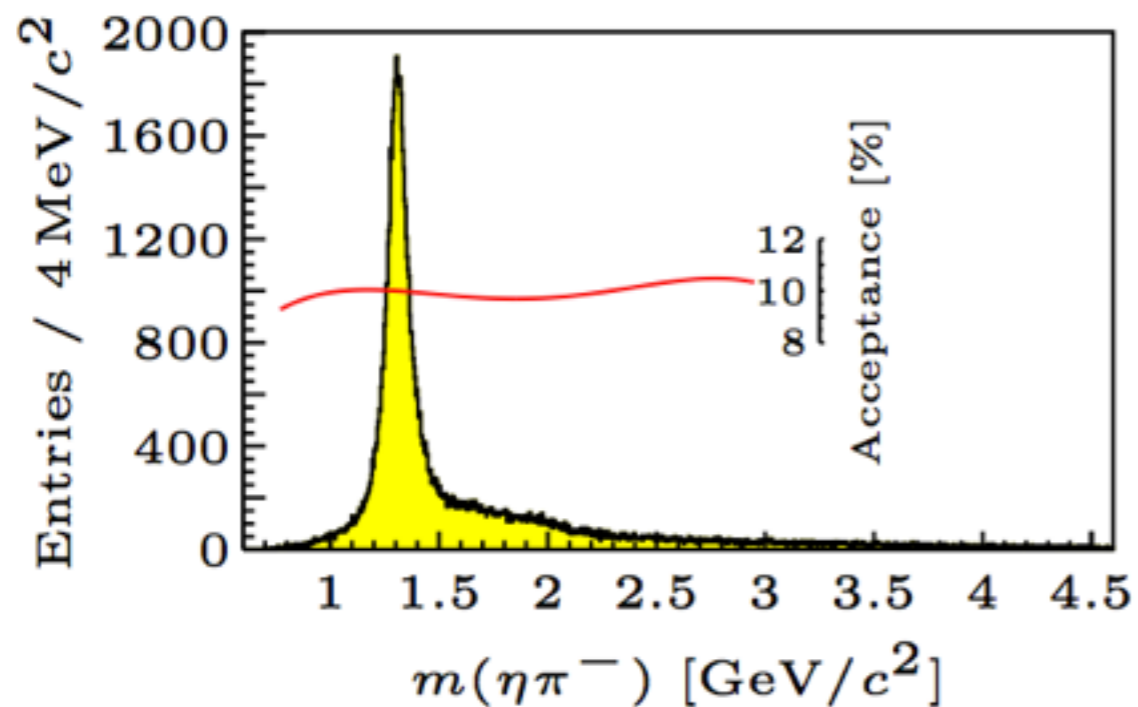
How do we select beam fragmentation ? \longrightarrow Boost in the rest frame



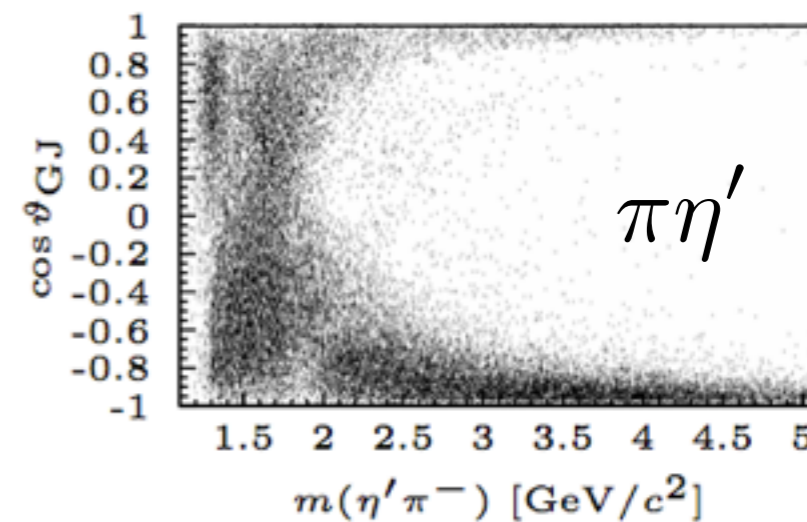
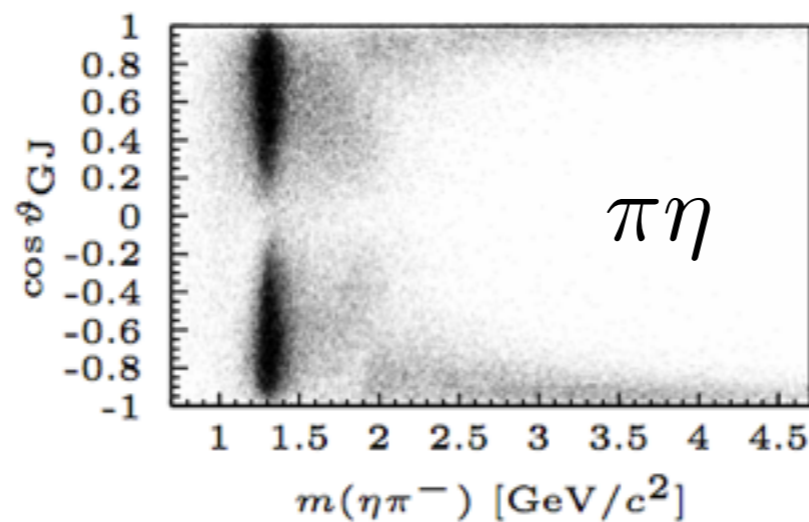
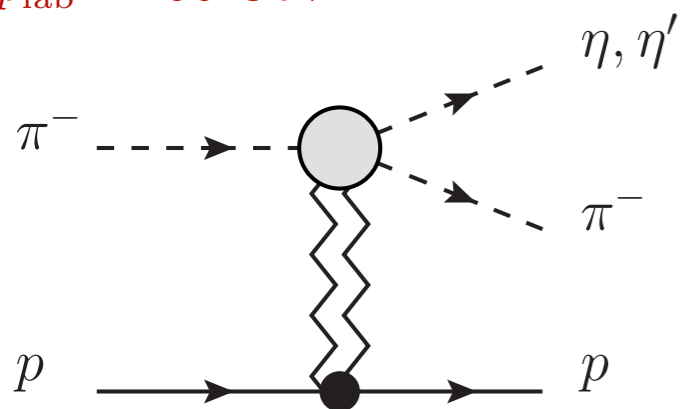
Van Hove NPB9 (1969) 331

M. Shi et al (JPAC) PRD91 (2015) 034007





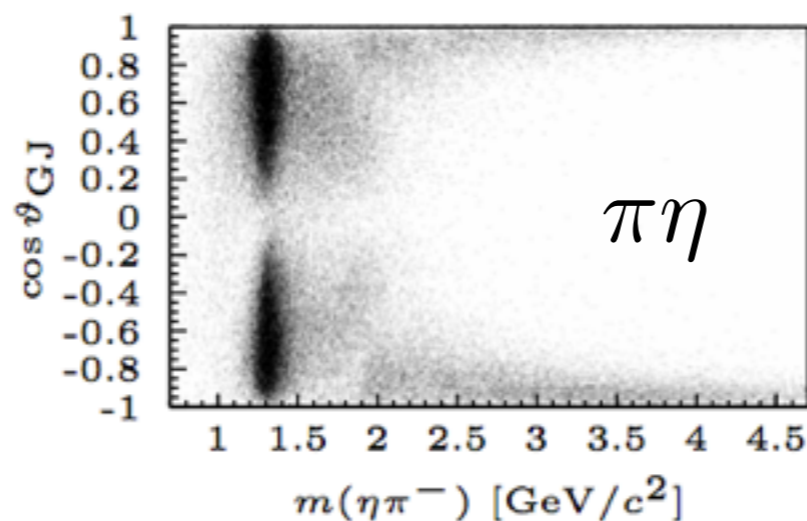
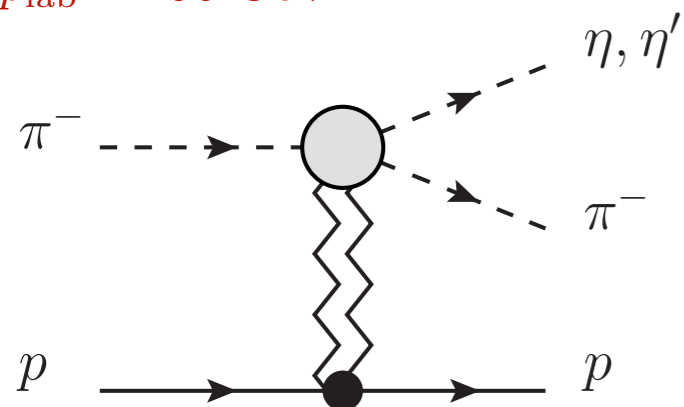
$p_{\text{lab}} = 190 \text{ GeV}$



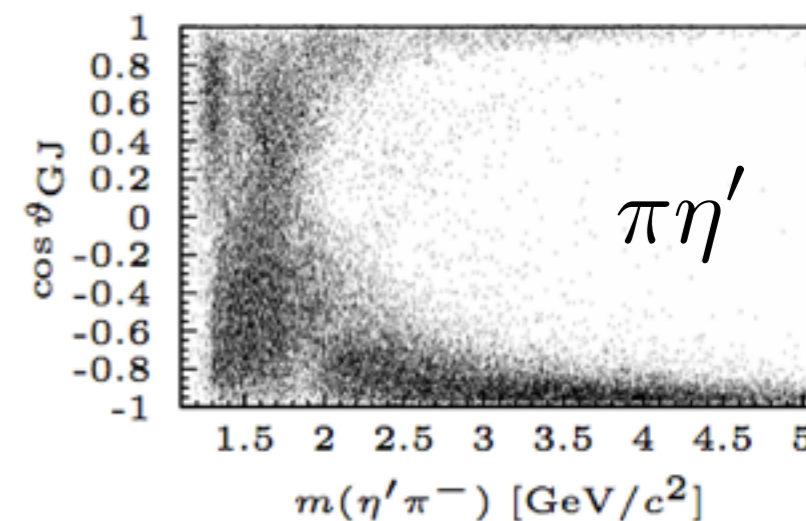
$(s, t_{\mathbb{P}}, s_{\eta\pi}, \theta, \phi)$

Gottfried-Jackson frame

$p_{\text{lab}} = 190 \text{ GeV}$



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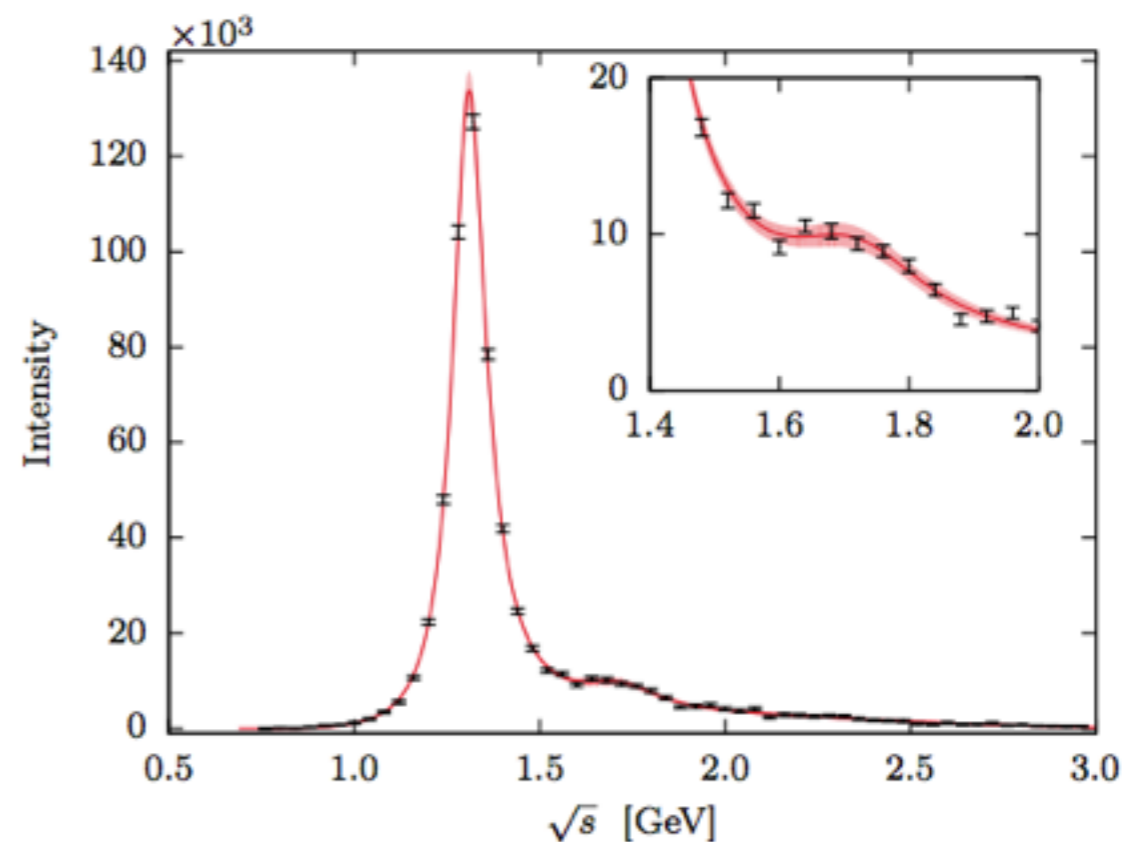


Quantum numbers determined by angular momentum

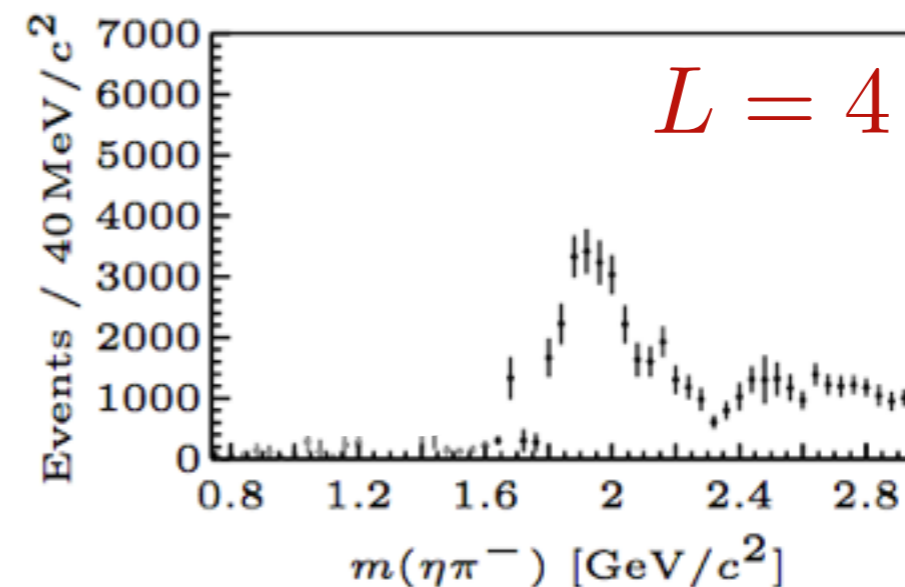
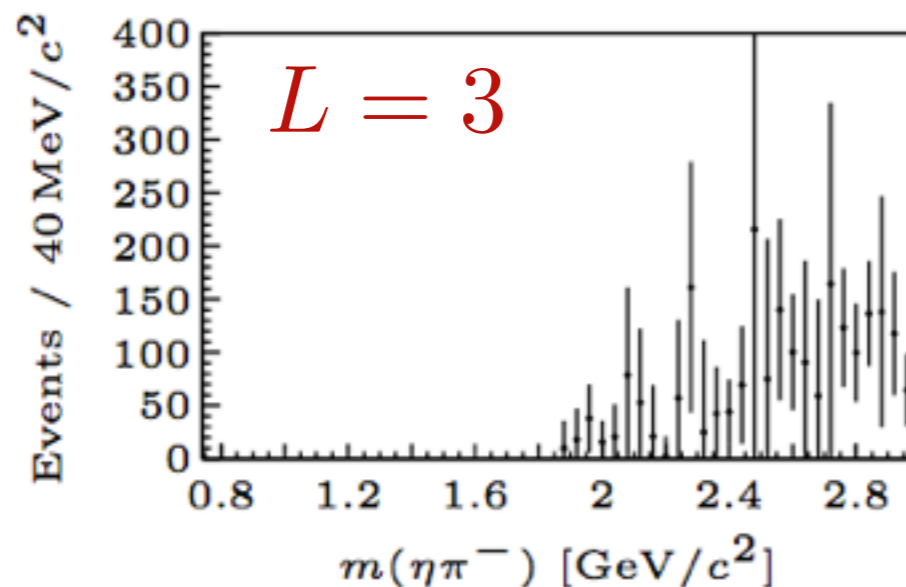
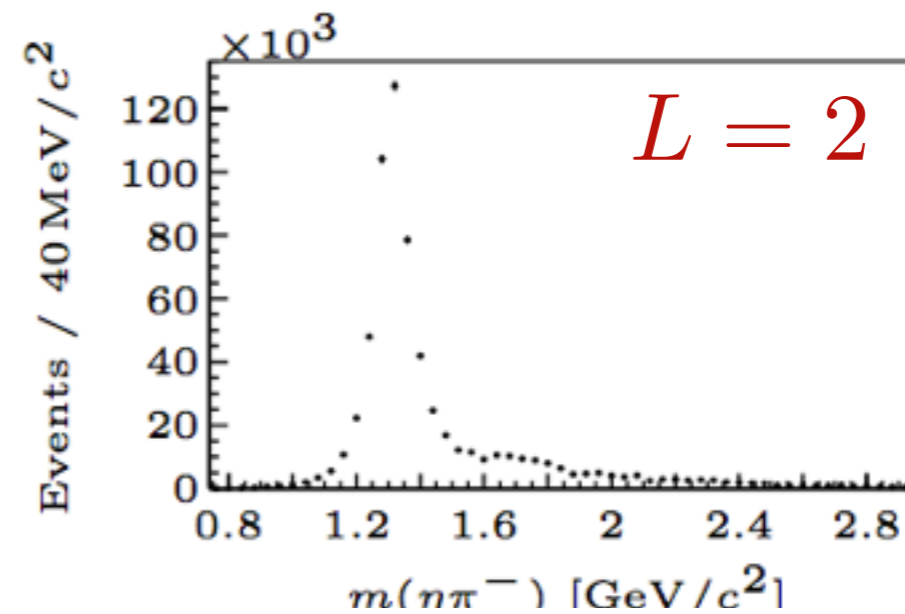
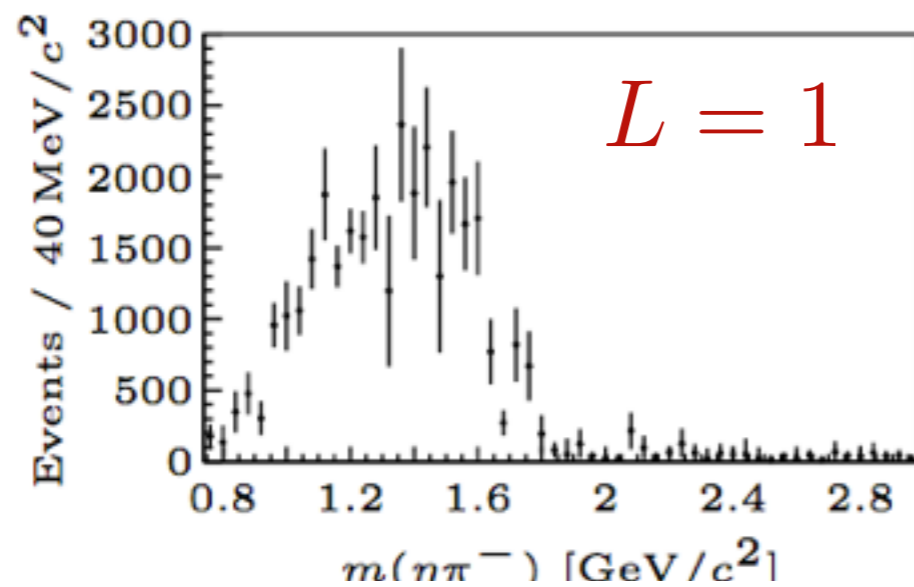
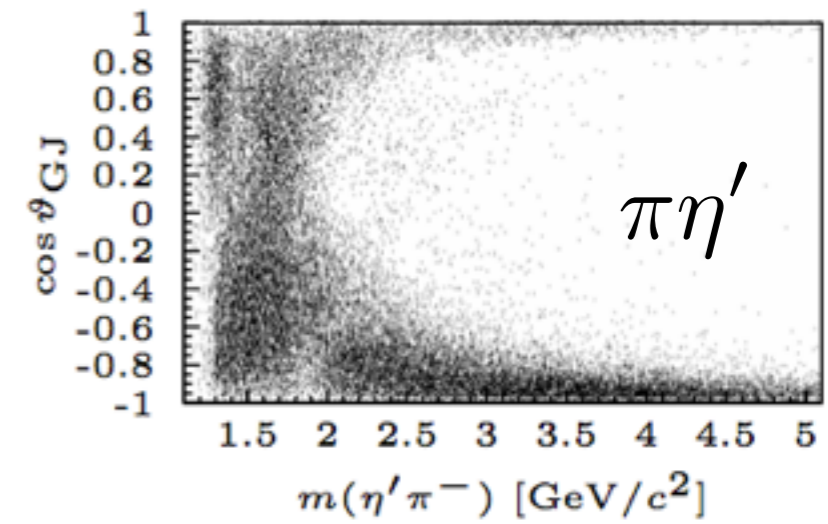
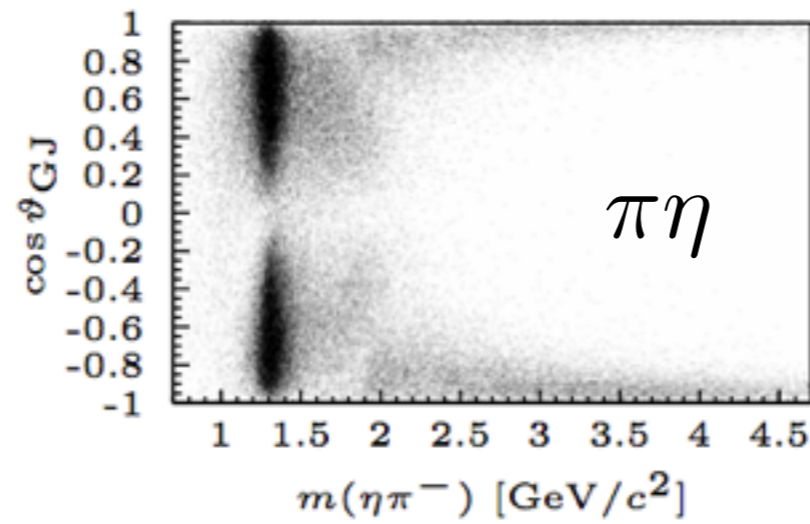
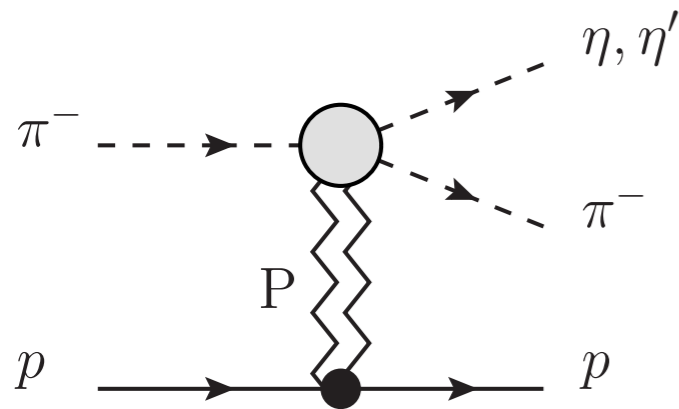
Which partial wave can yield the '8' in $\eta\pi$?

$$a_2(1320) : I^G J^{PC} = 1^- 2^{++}$$

$$d_{1,0}^2(\theta) \propto Y_2^1(\theta, 0) \propto \sin \theta \cos \theta$$

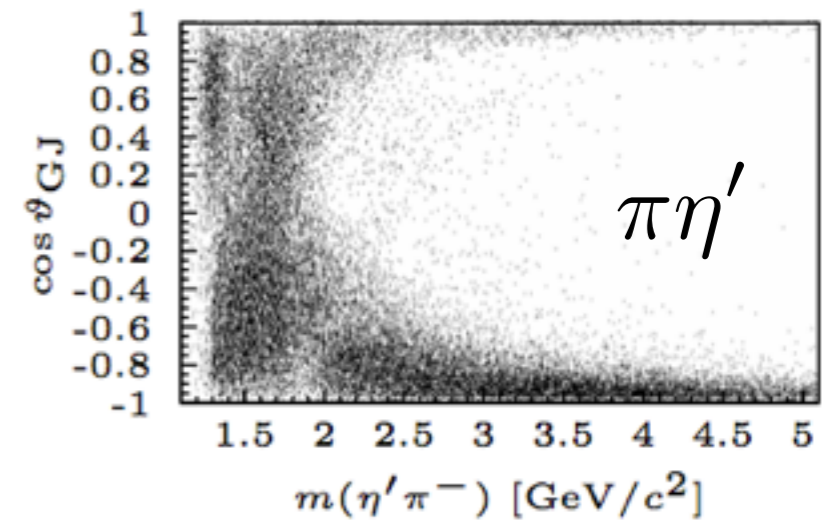
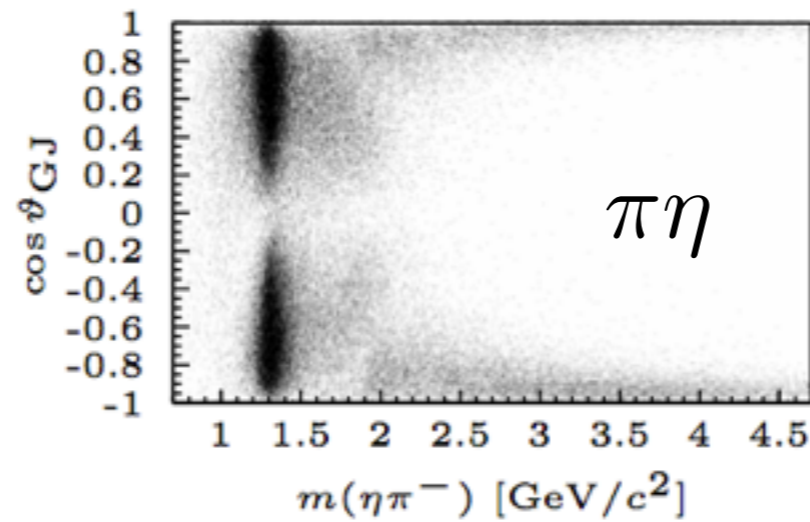
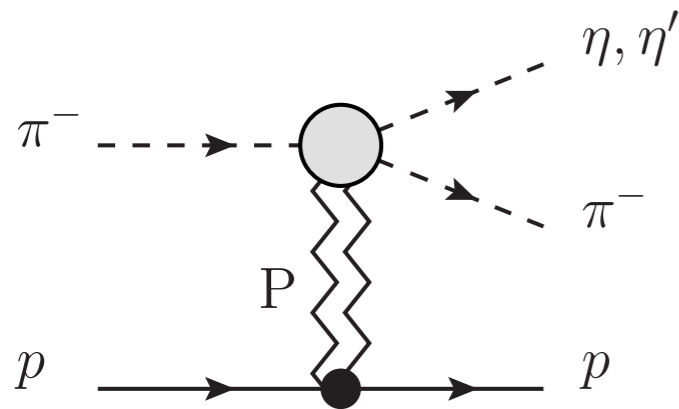


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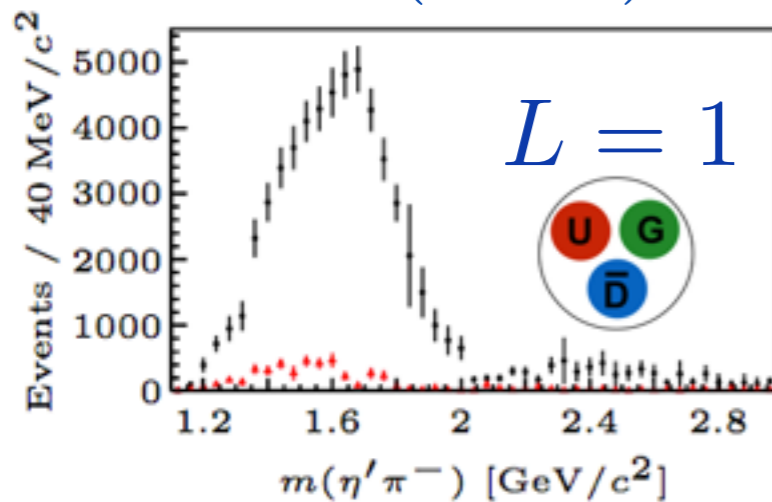


Partial Waves

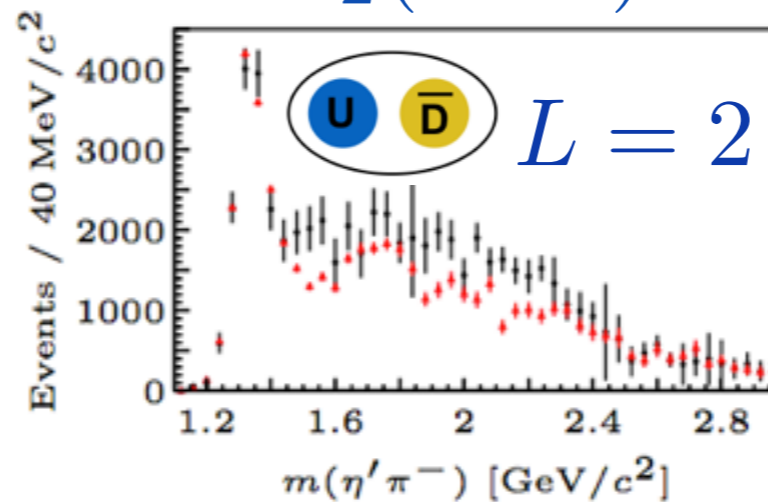
COMPASS Phys. Lett. B740 (2015)



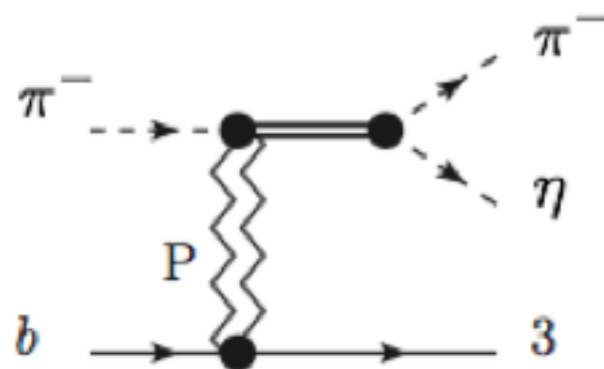
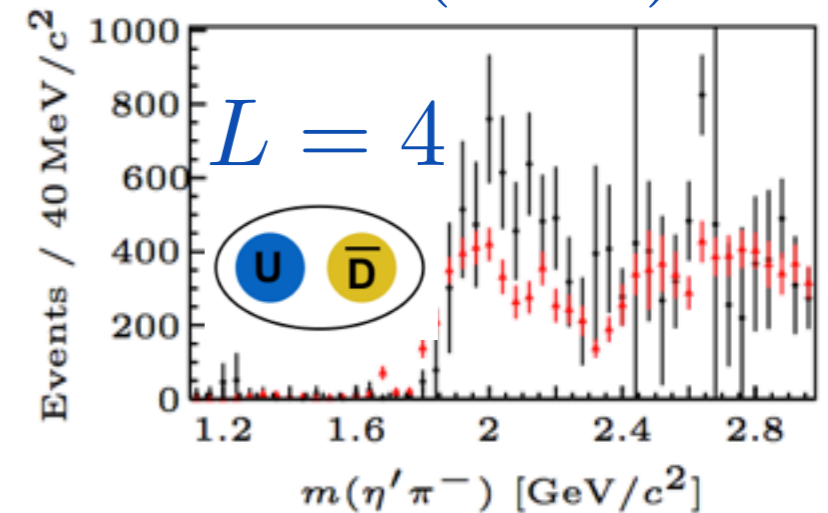
$\pi_1(1600)?$



$a_2(1320)$



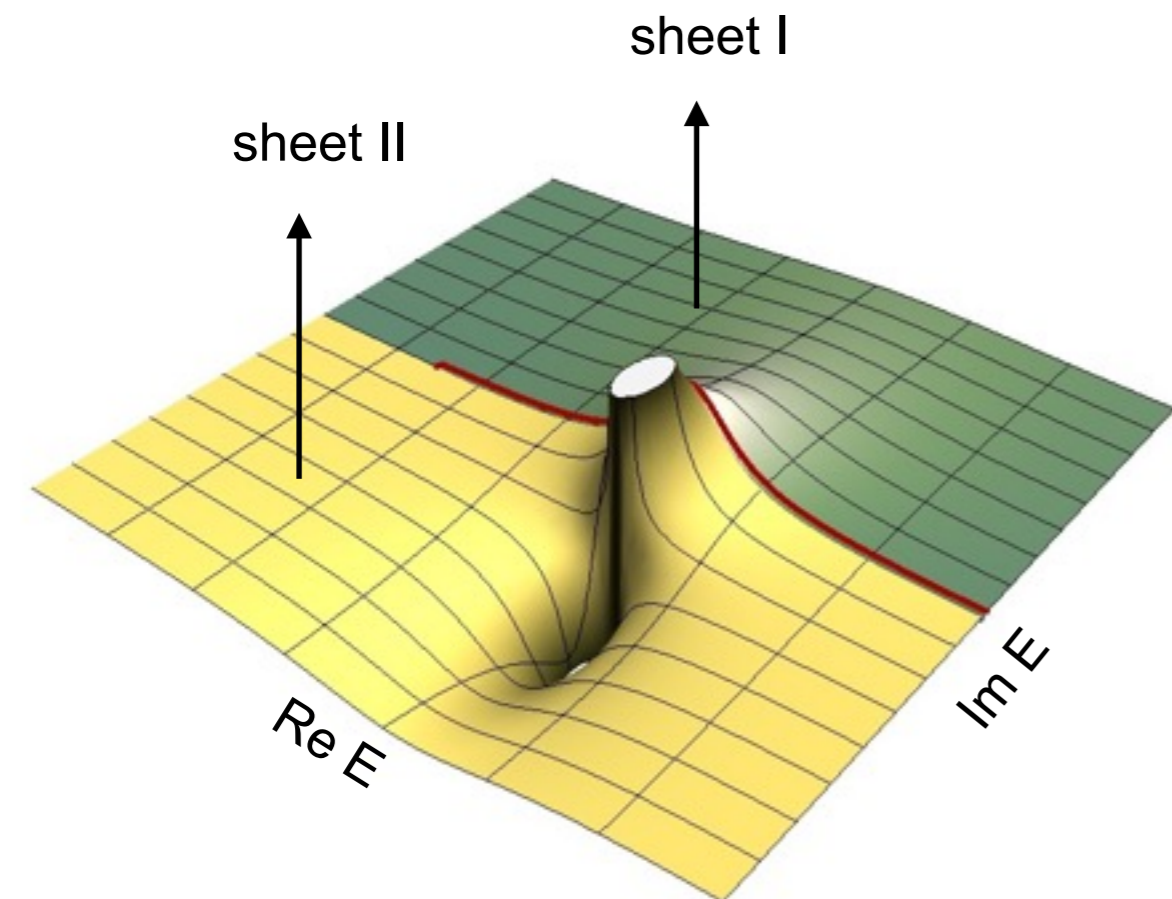
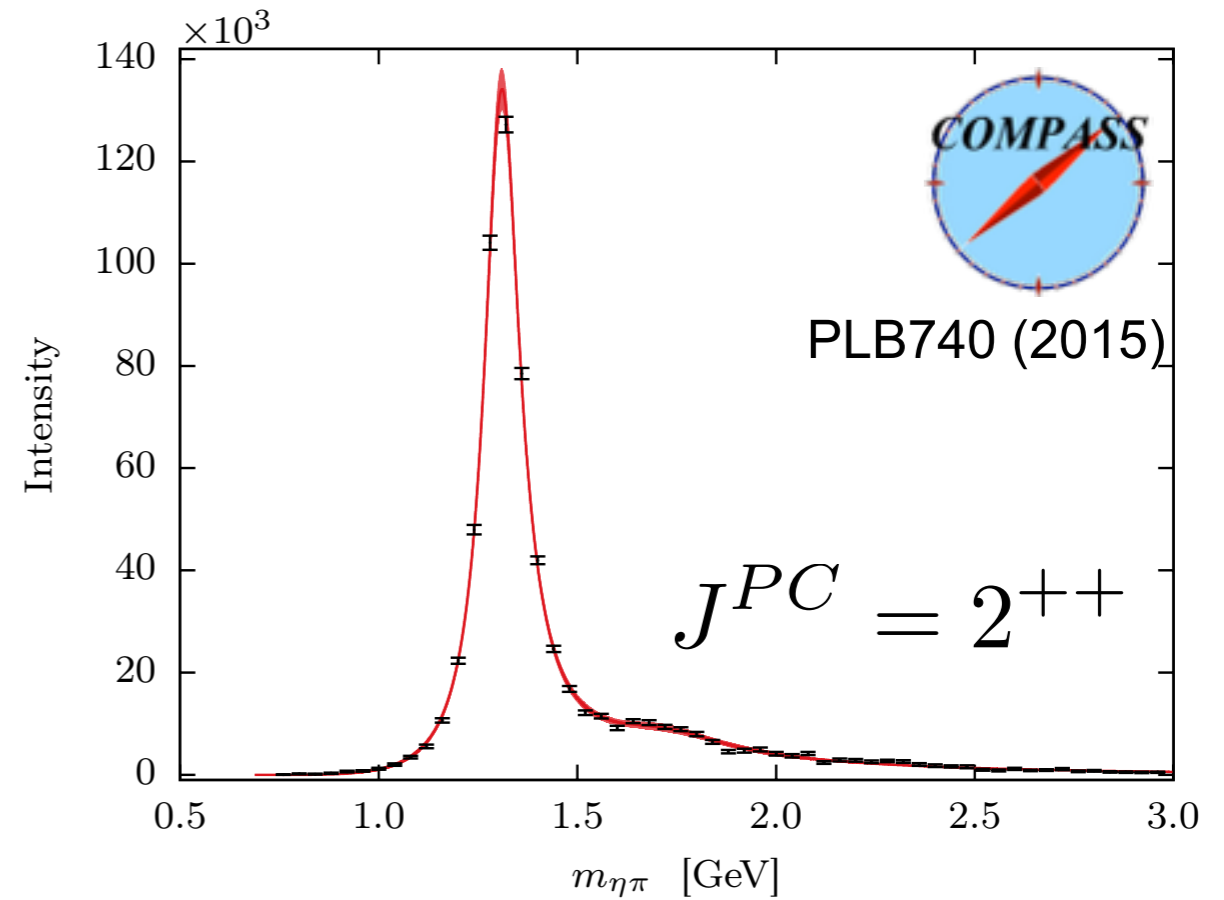
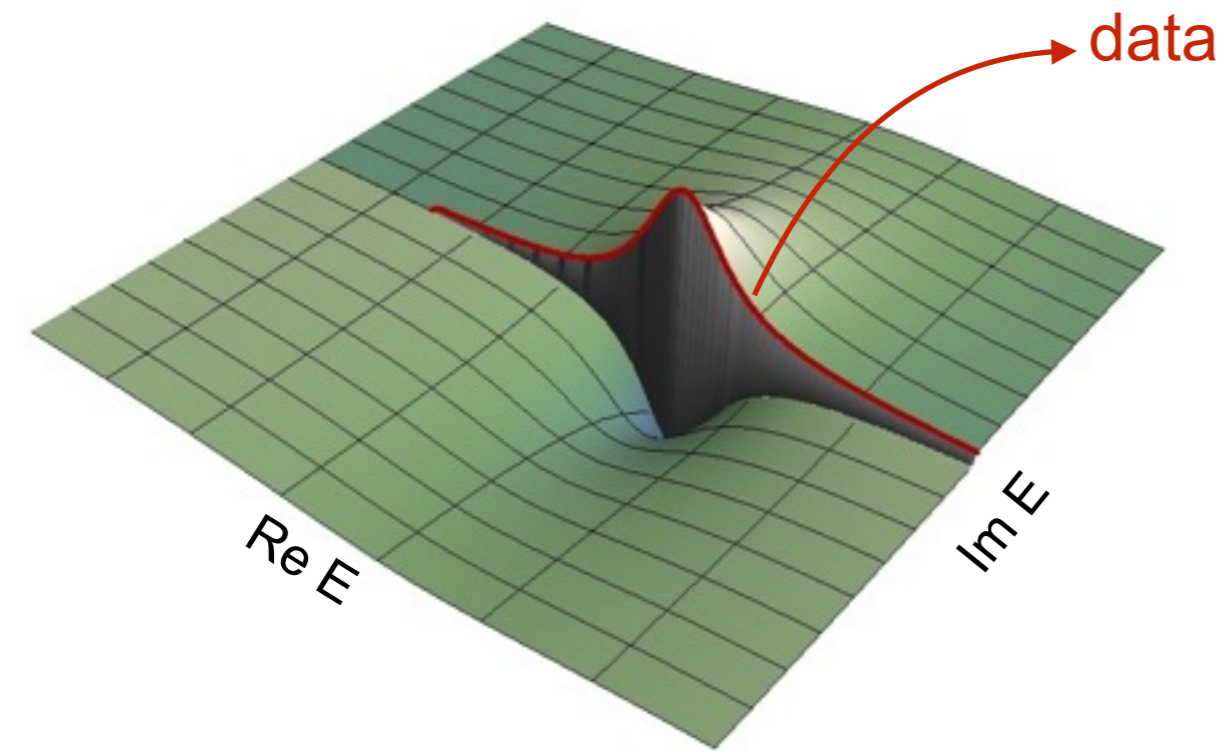
$a_4(2040)$



black: $\pi\eta'$

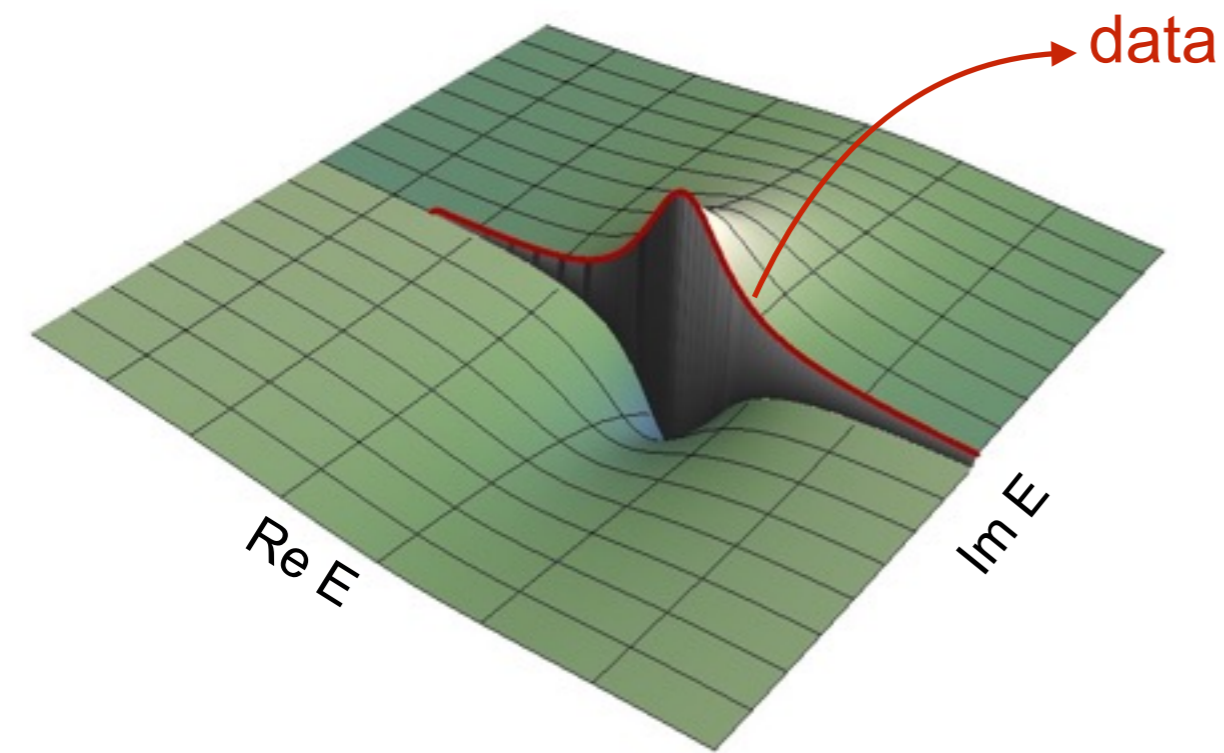
red: $\pi\eta$ (scaled)

Resonance in angular mom. $L = 1$?



Poles in the complex energy plane:
Real part \sim mass
Imaginary part \sim width
Residue \sim coupling

Poles or resonances are the universal building blocks of reactions



partial wave \leftarrow $\text{Im } t_\ell^{-1}(s) = -\rho(s)$ \rightarrow phase space

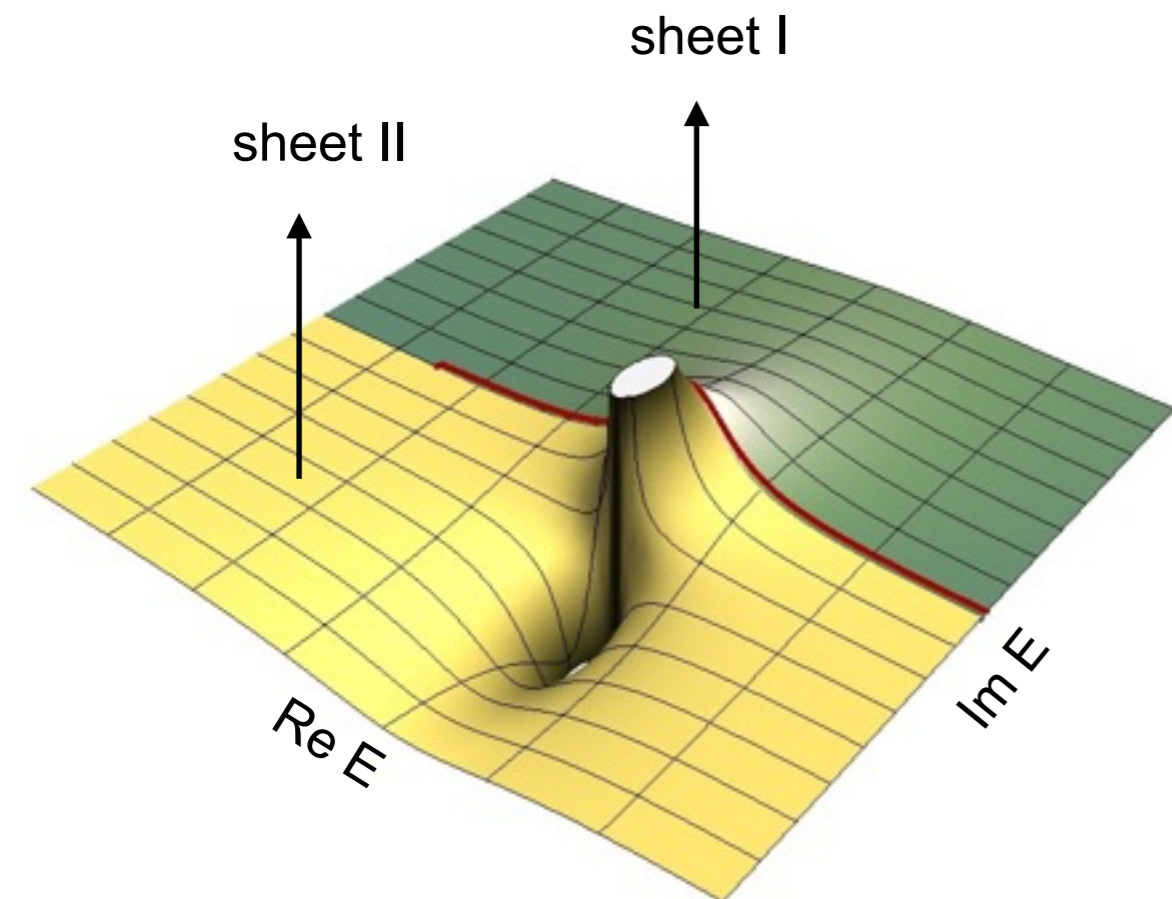
$$t_\ell(s \pm i\epsilon) = \frac{1}{K(s) \mp i\rho(s)}$$

satisfies causality
(regular outside the real axis)

define function on sheet II
on the lower half plane

$$t_\ell^{II}(s) = \frac{1}{K(s) - i\rho(s)}$$

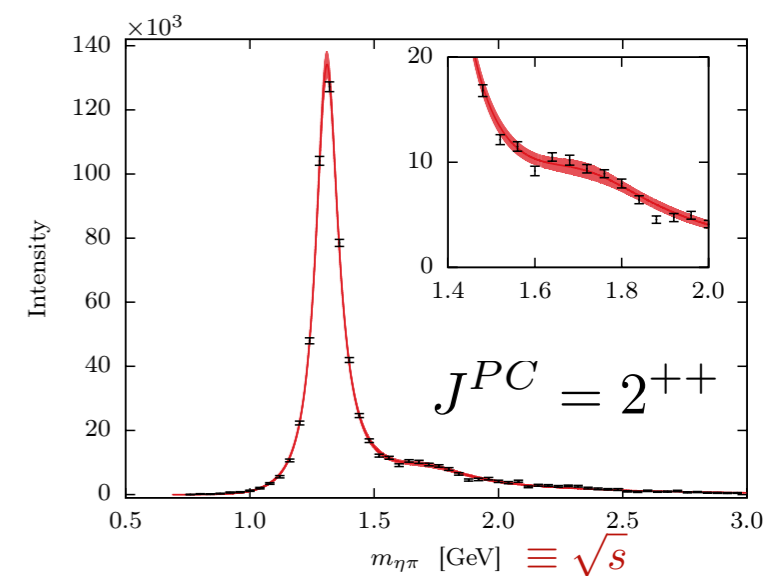
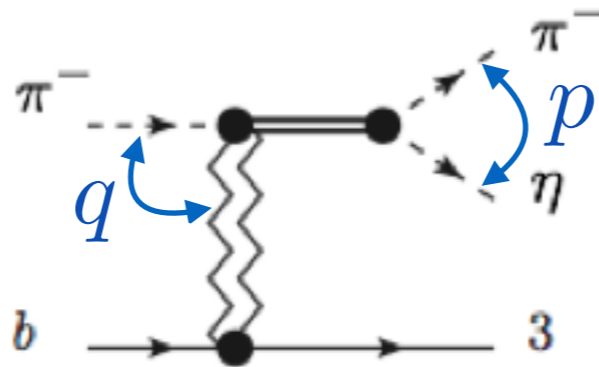
example $= \frac{m\Gamma}{m^2 - s - i\rho(s)m\Gamma}$



Eta-Pi@COMPASS

$$\frac{d\sigma}{d\sqrt{s}} = N p |a(s)|^2$$

normalization



$$a(s) = p^2 \frac{n(s)}{D(s)}$$

production

$$n(s) = \frac{q}{c_3 - s} \sum_n a_n T_n(\omega(s))$$

Chebyshev polynomials

$$\omega(s) = \frac{s}{s + \Lambda}$$

dynamics (poles)

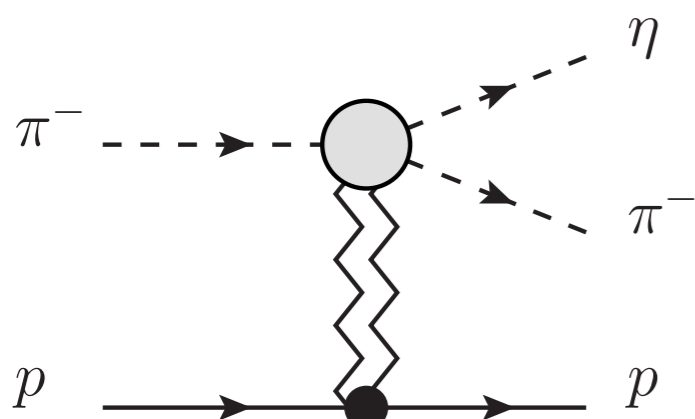
$$D(s) = D_0(s) - \frac{s}{\pi} \int_{s_{\text{th}}}^{\infty} \frac{ds'}{s'} \frac{\rho(s') N(s')}{s' - s}$$

real (masses)

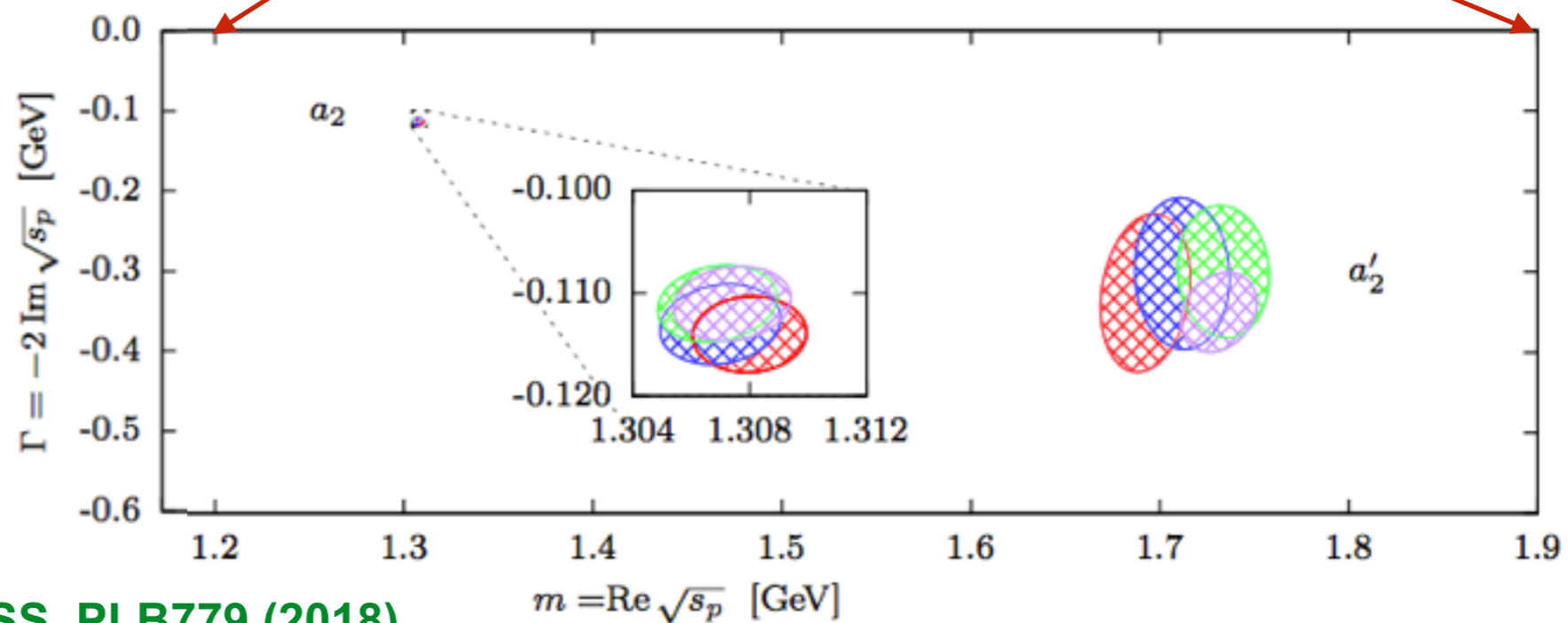
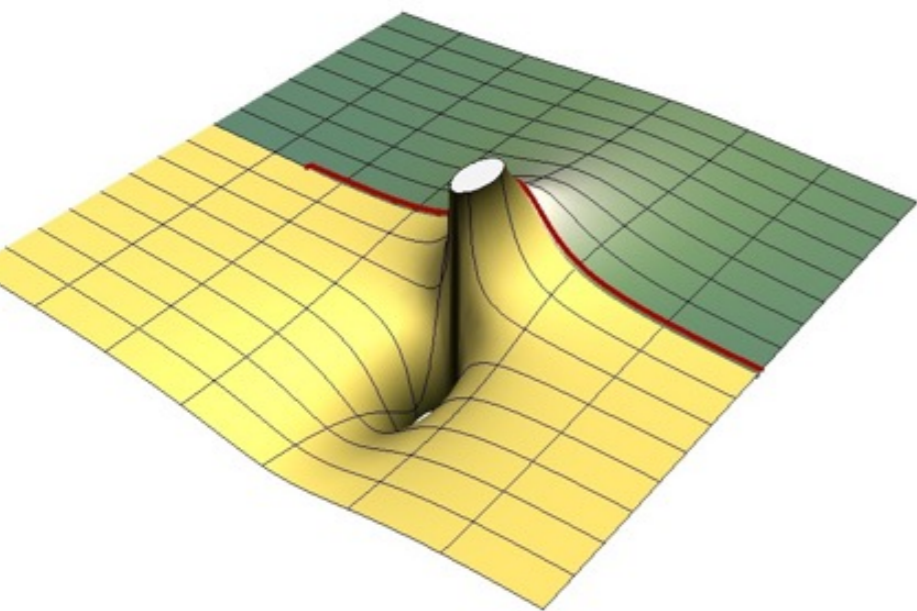
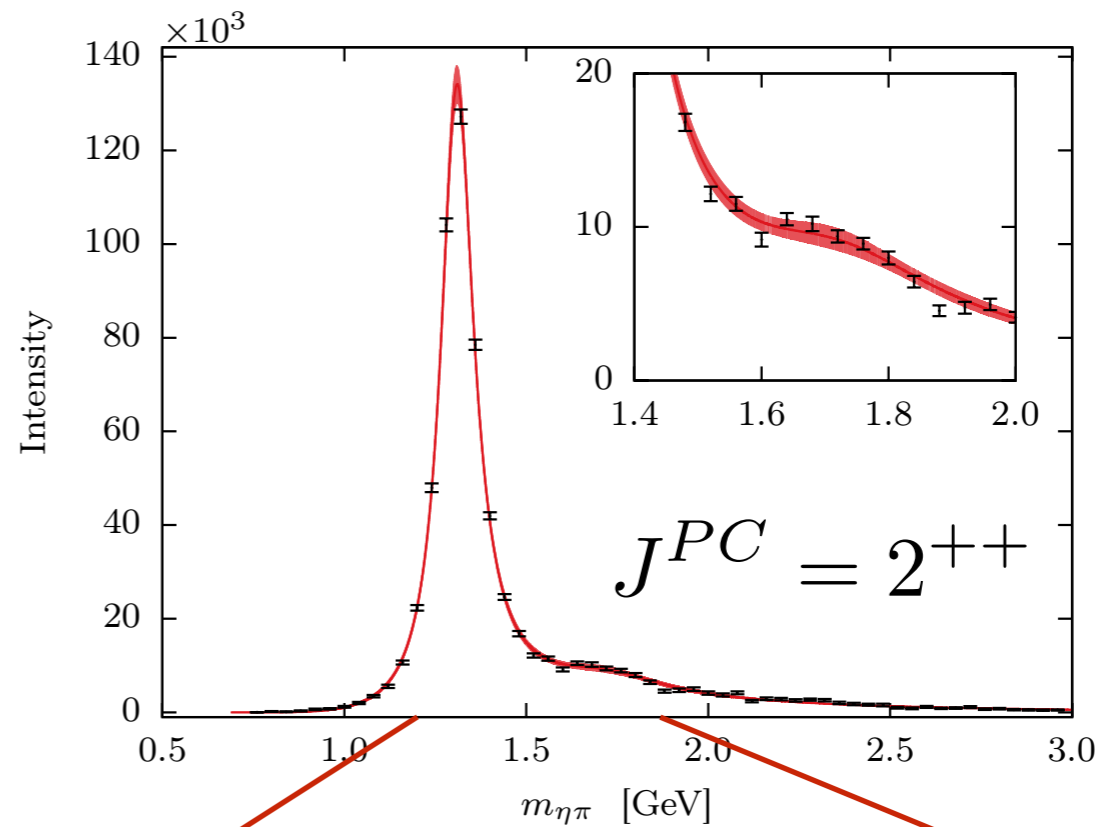
$$D_0(s) = c_0 - c_1 s$$

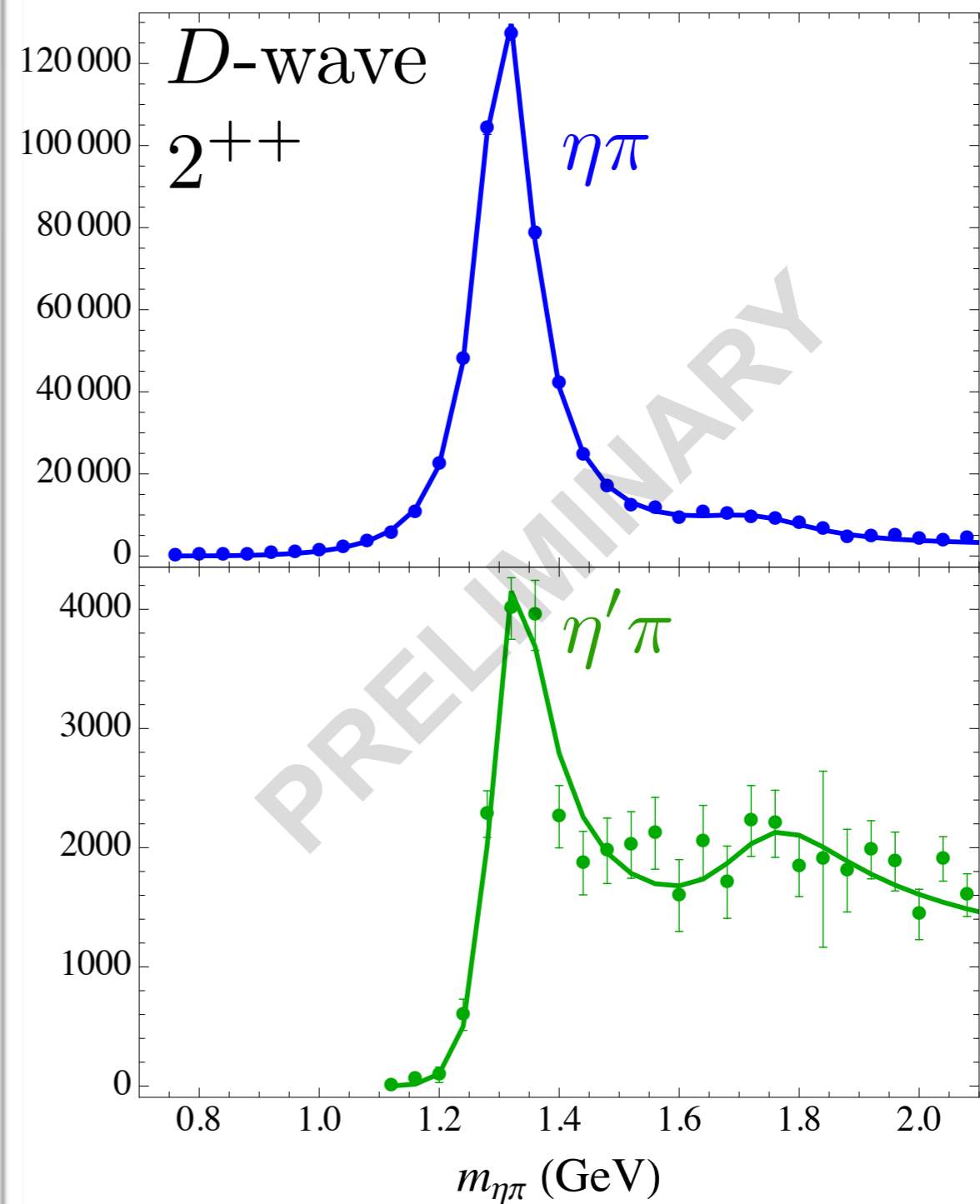
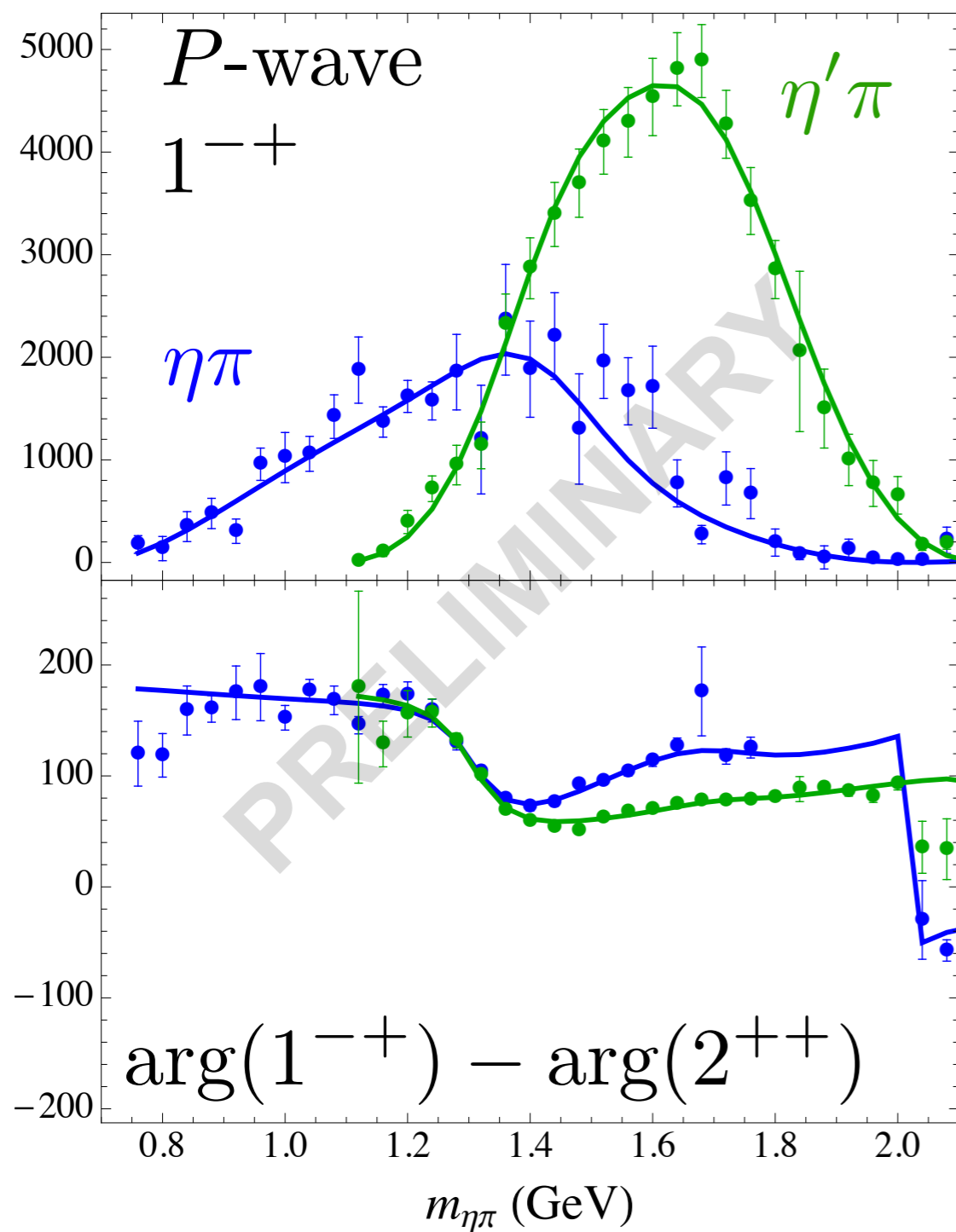
imaginary (widths)

$$\rho(s) N(s) = g \frac{\lambda^{J+\frac{1}{2}}(s, m_\eta^2, m_\pi^2)}{(s + s_R)^{2J+3}}$$

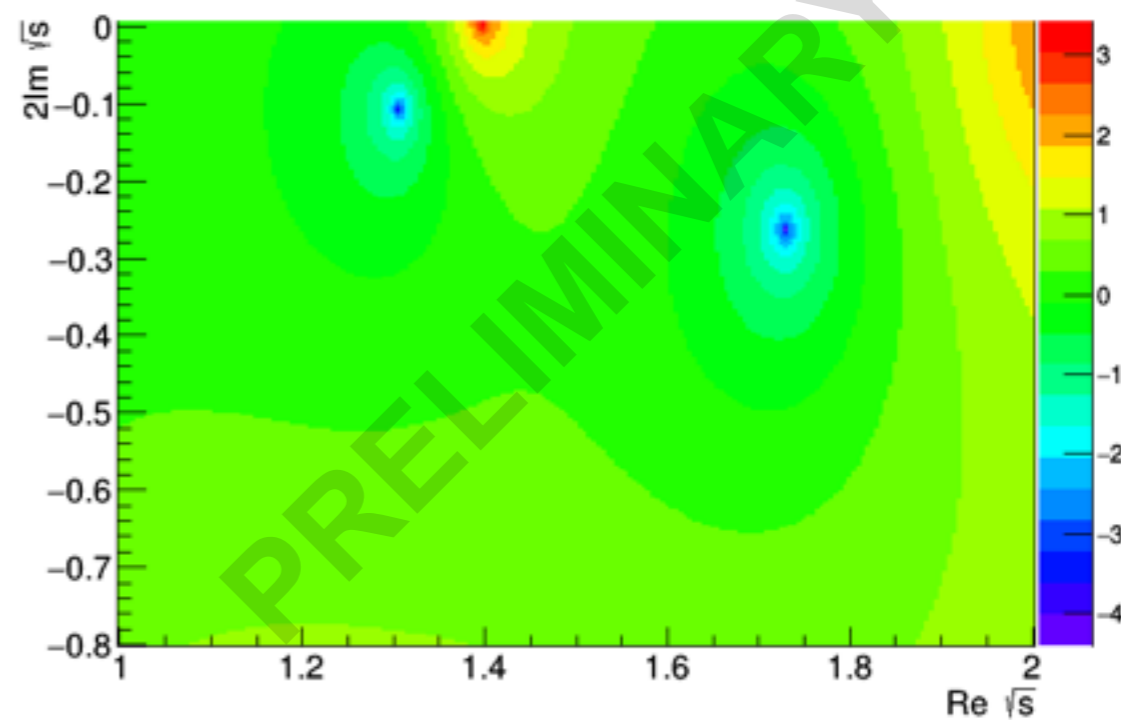
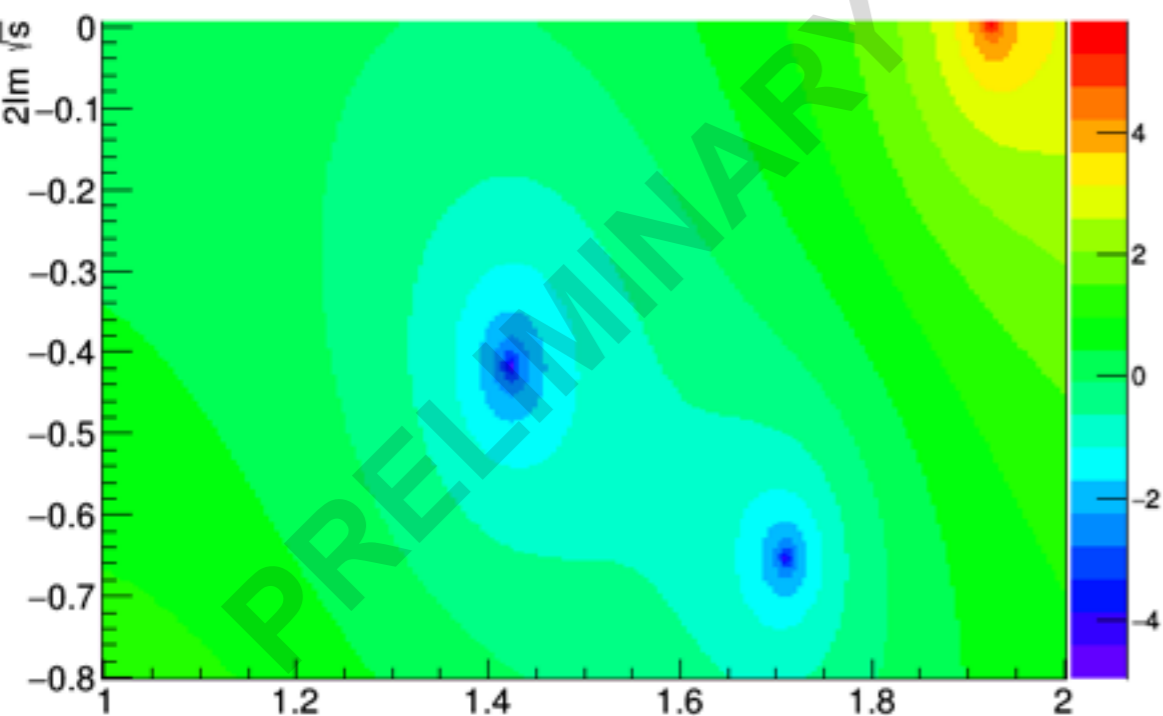
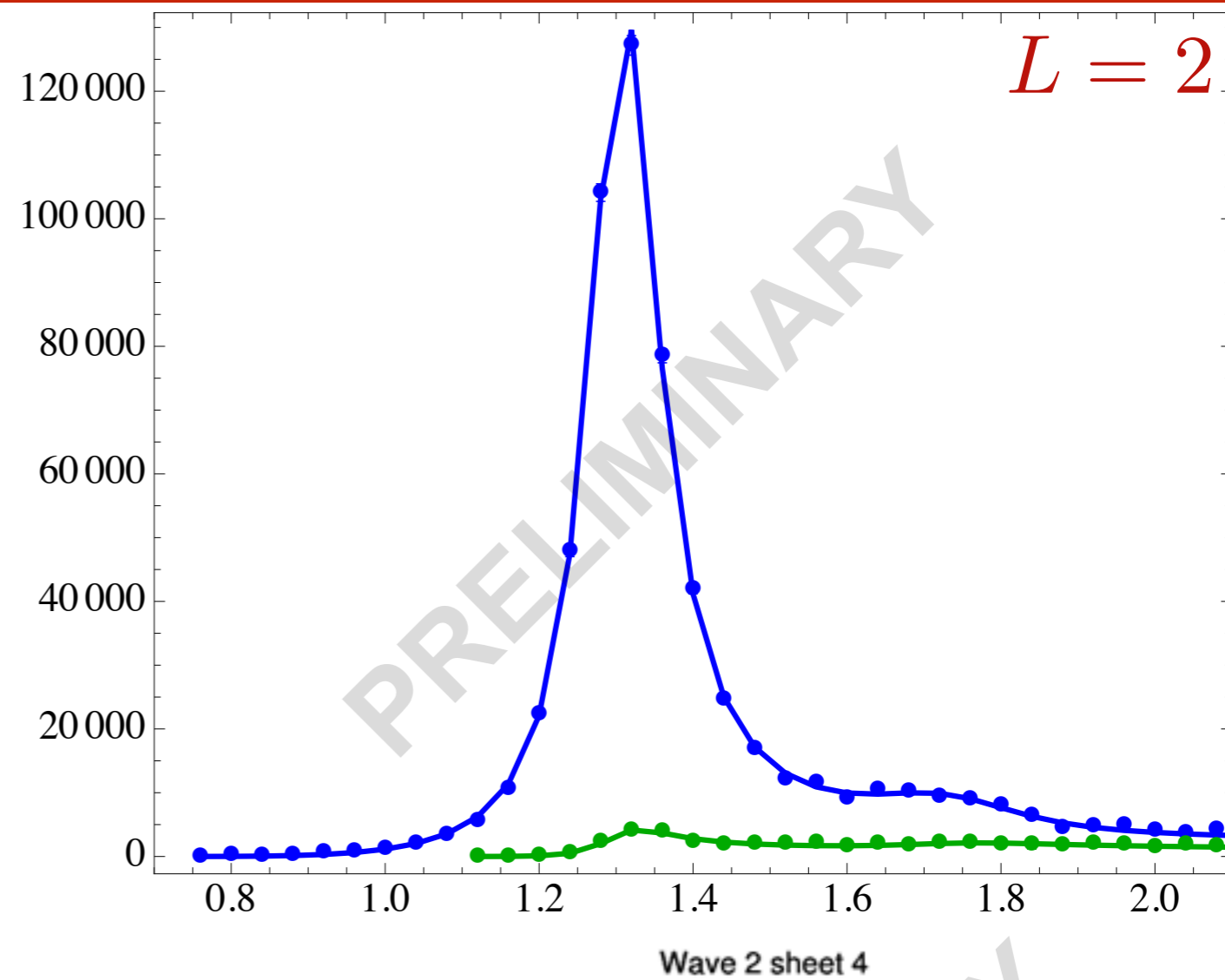
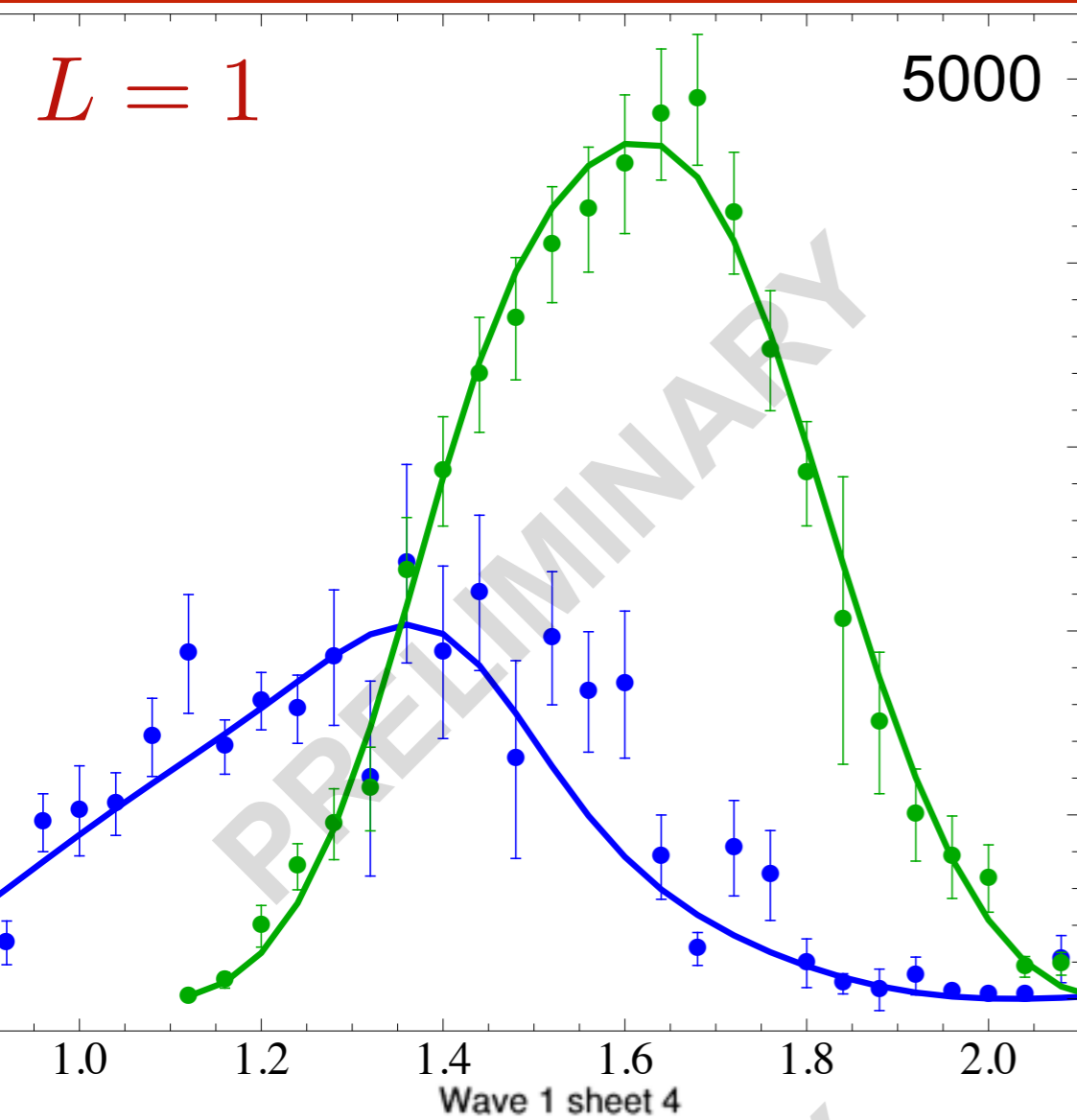


first precise determination of $a_2(1700)$ pole location

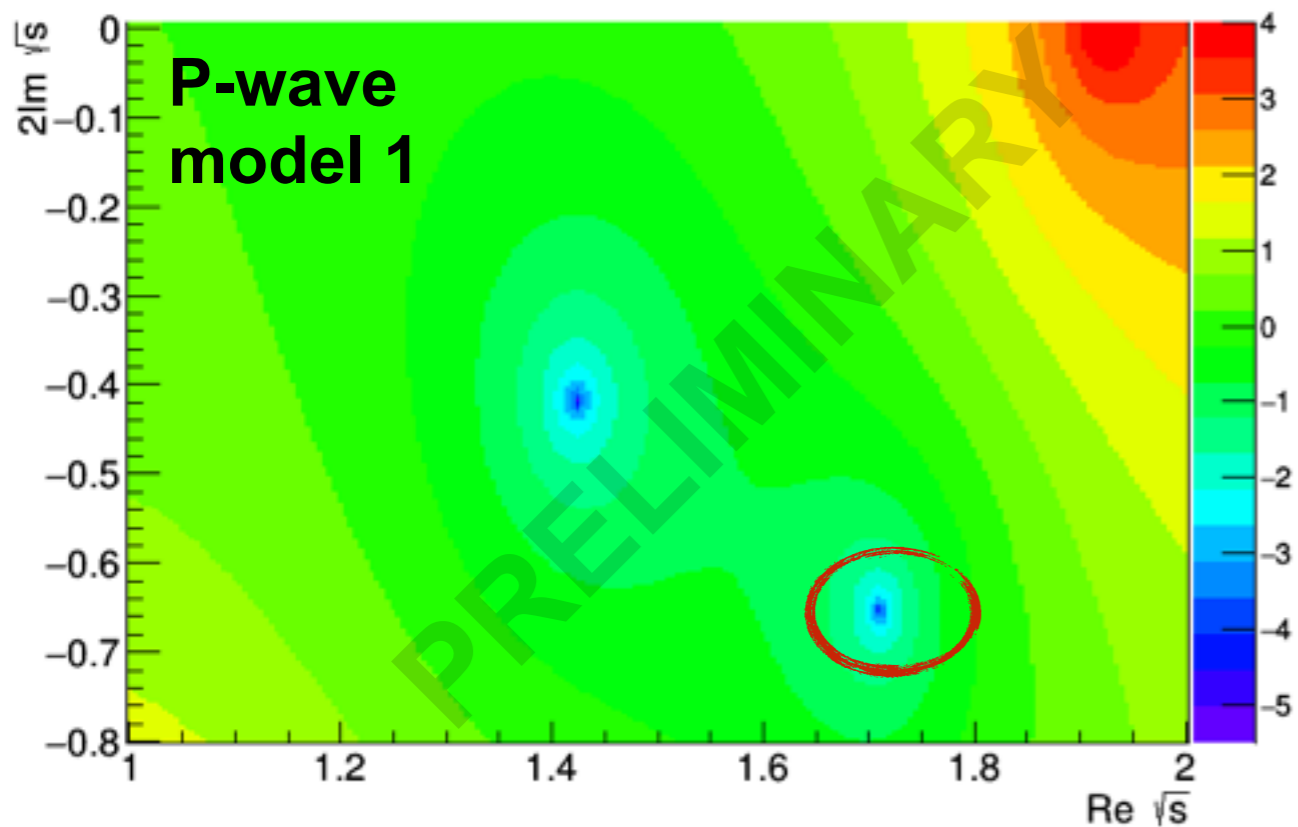




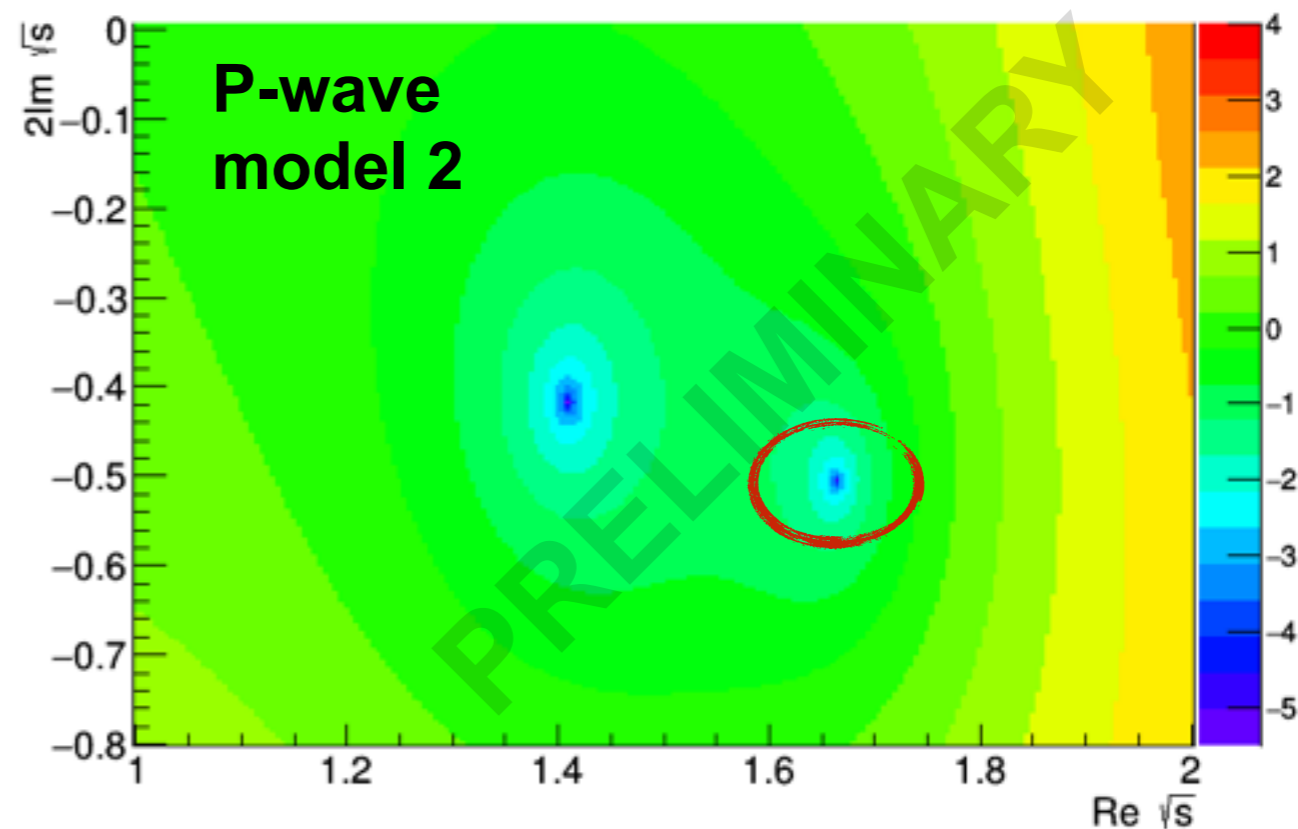
**On-going analysis (Arkaitz and Alessandro):
 Systematic studies and exploration of the complex plane**



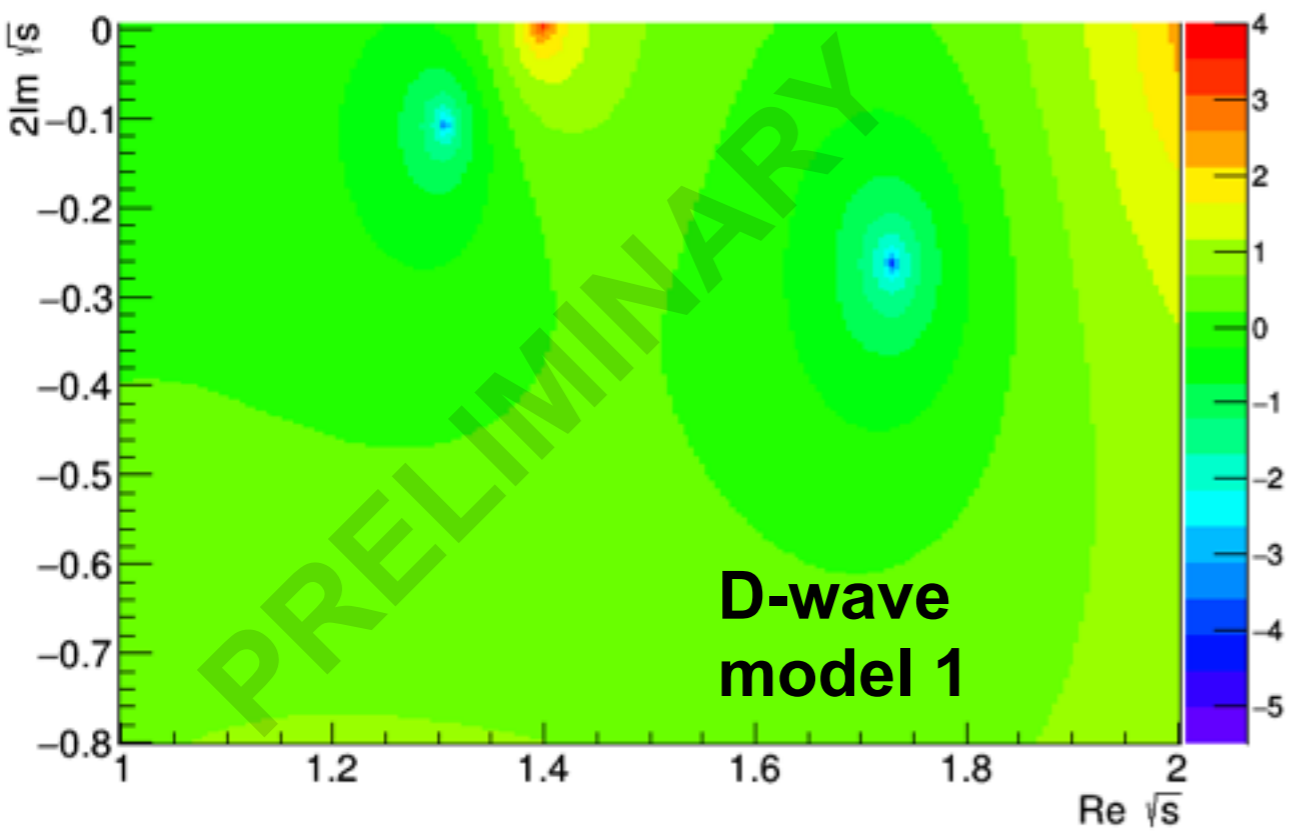
Wave 1 sheet 4



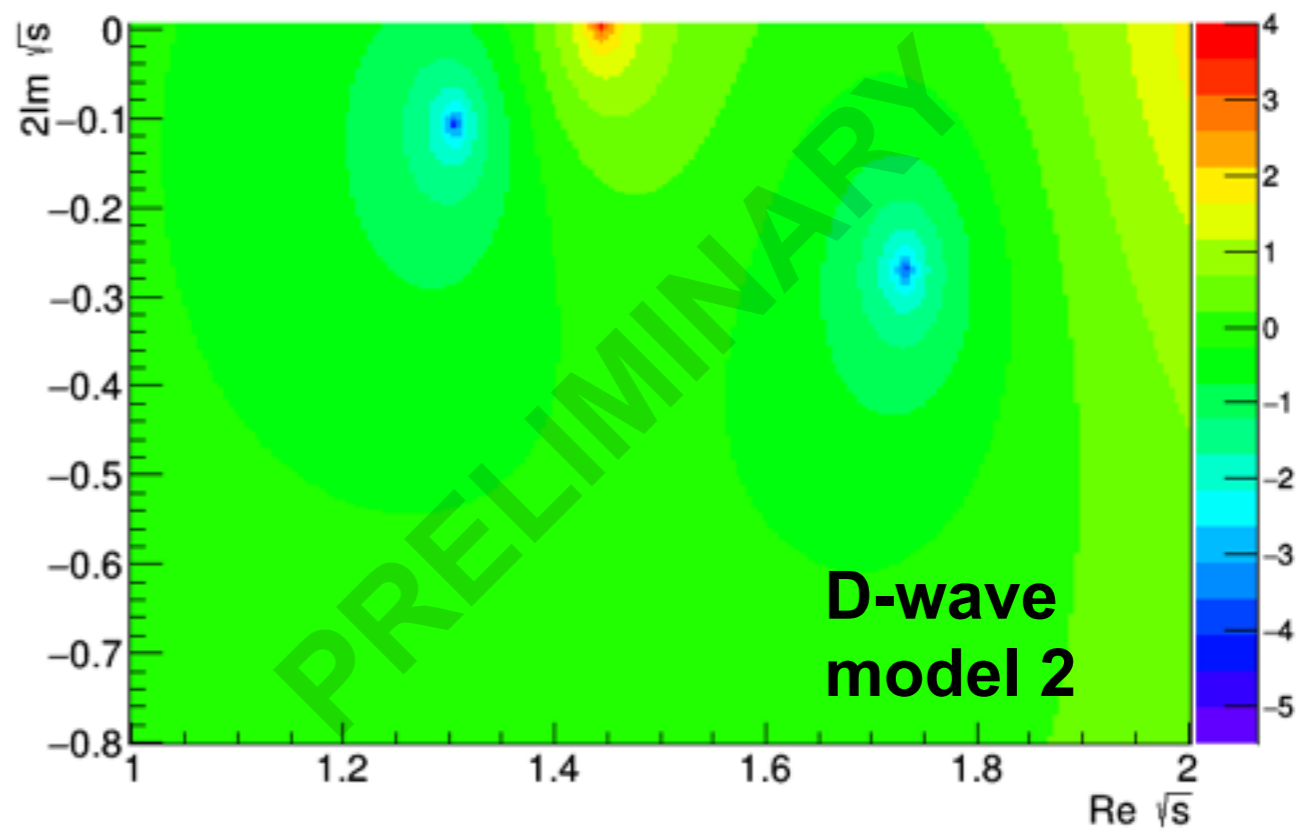
Wave 1 sheet 4

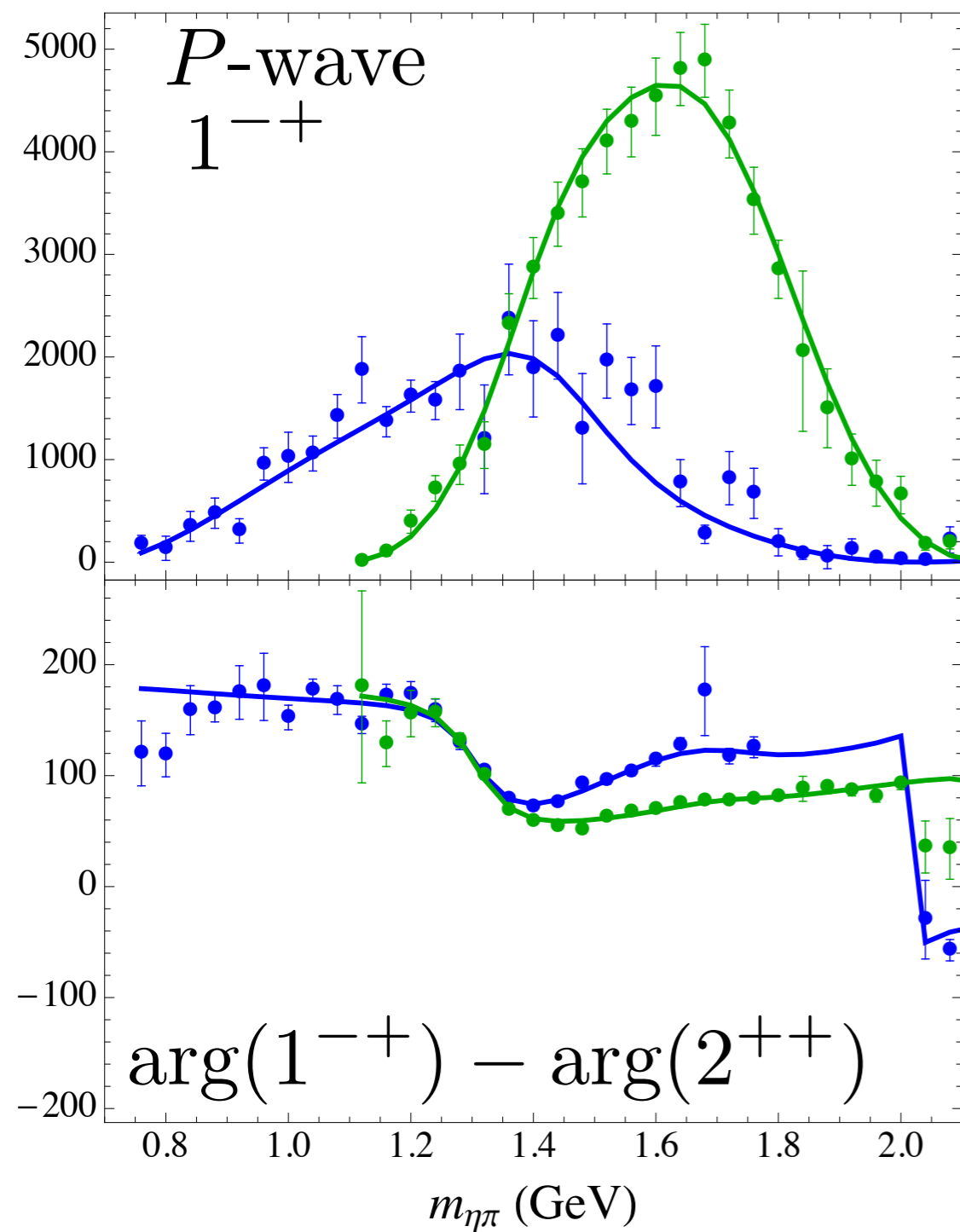


Wave 2 sheet 4



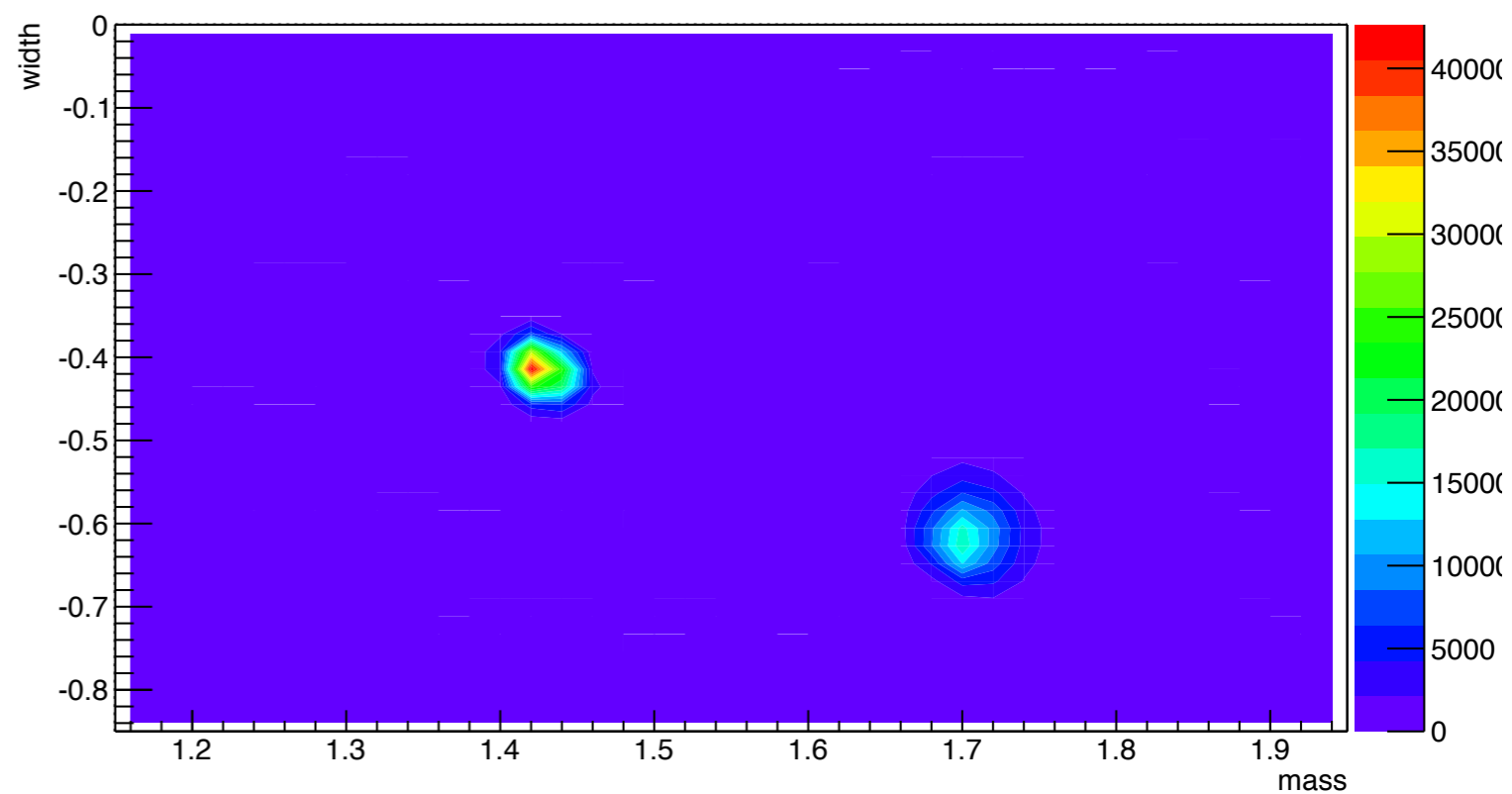
Wave 2 sheet 4



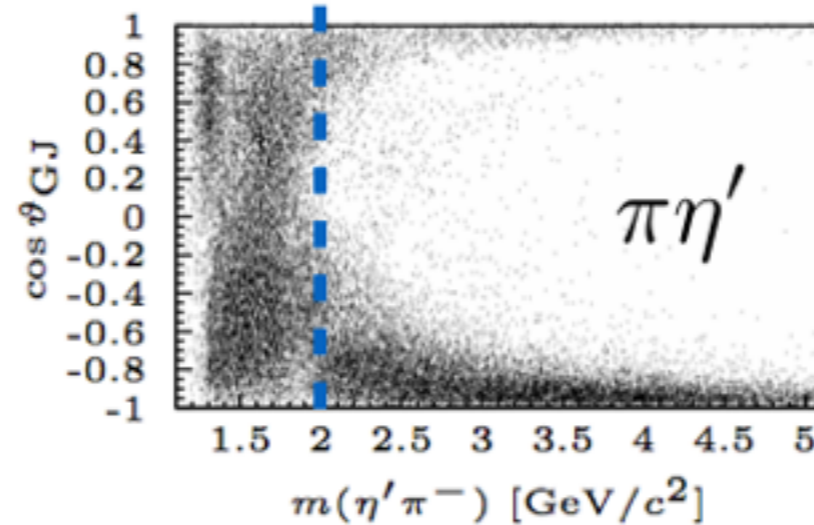
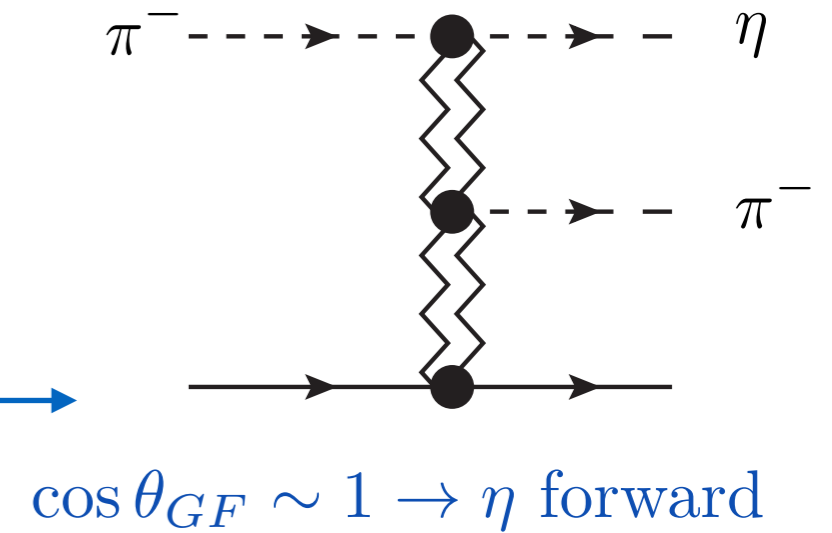
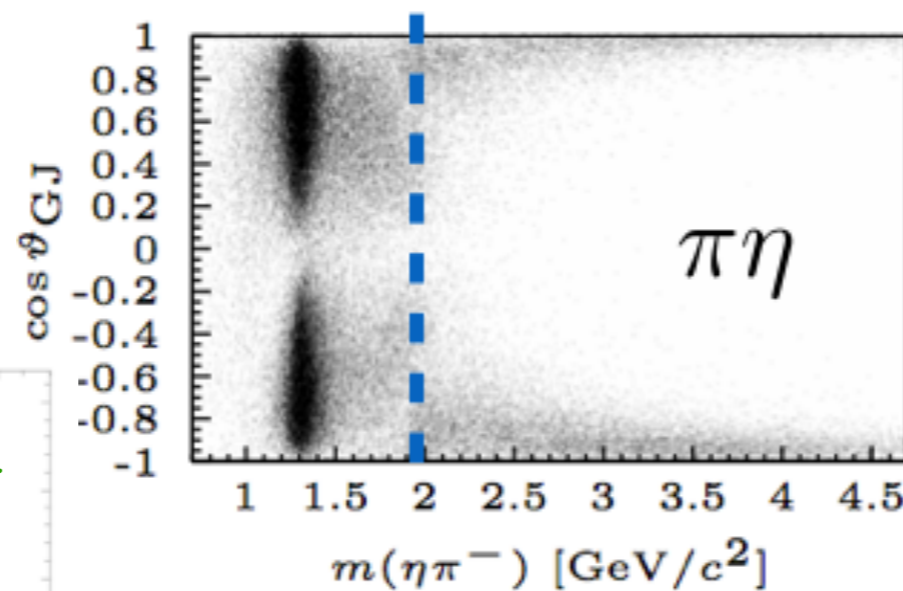
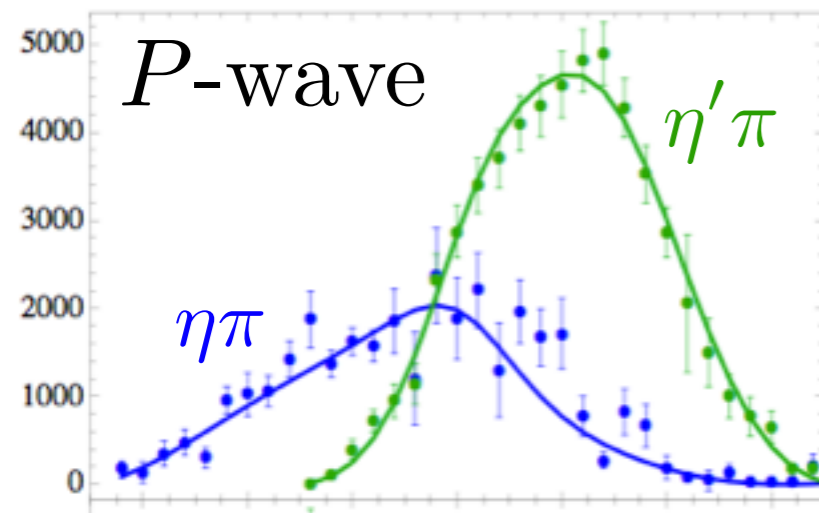


Generate 250k 'new data'
Fit with previous best fit as starting point

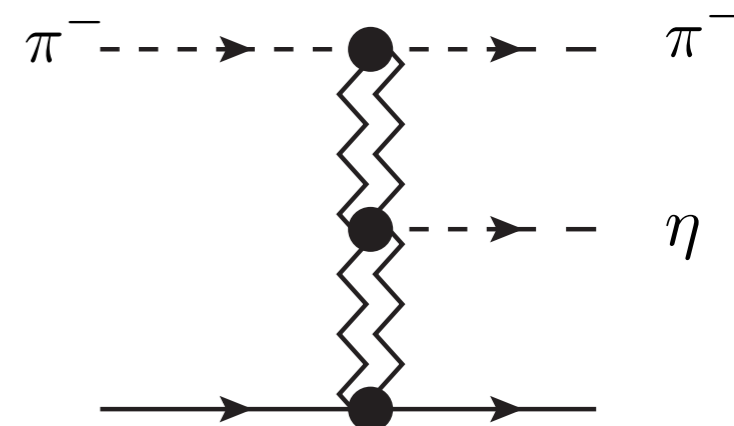
Display the 250k x2 P-wave poles:



Dispersion relation relates the high (exchanges) and the low (resonances) regions



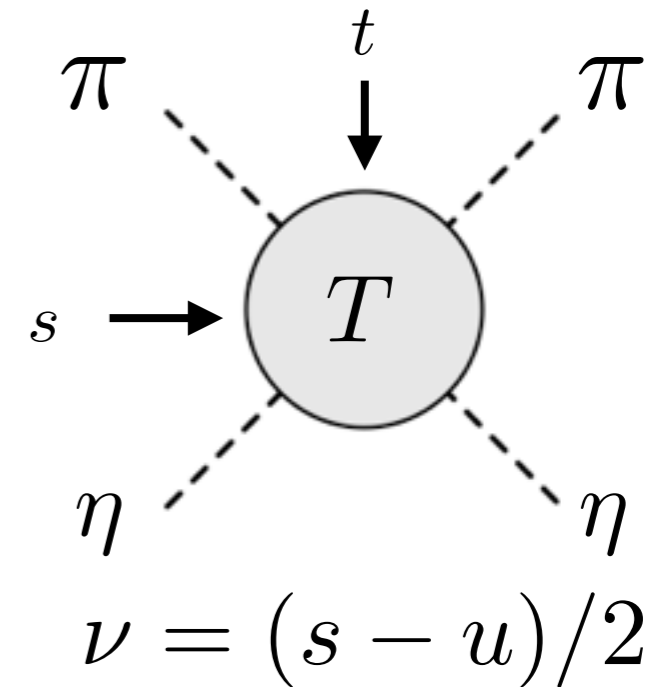
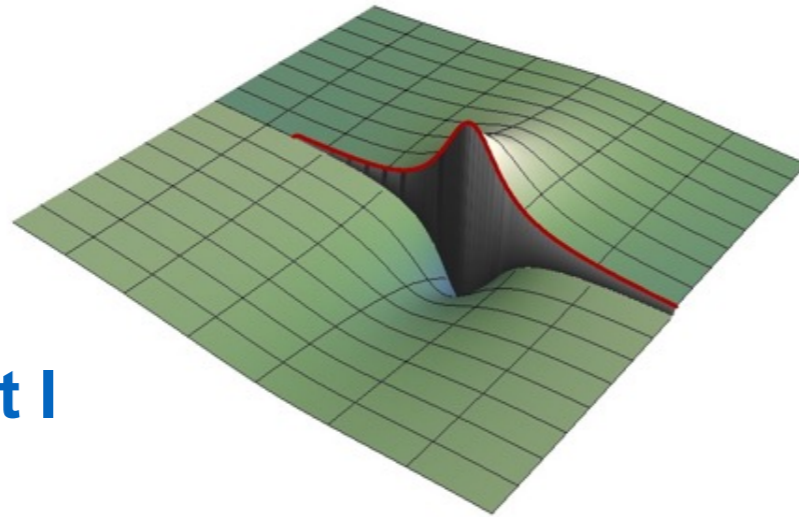
$\cos\theta_{GF} \sim -1 \rightarrow \eta$ backward



$$T(s, t)$$

has a right- and a left-hand cuts

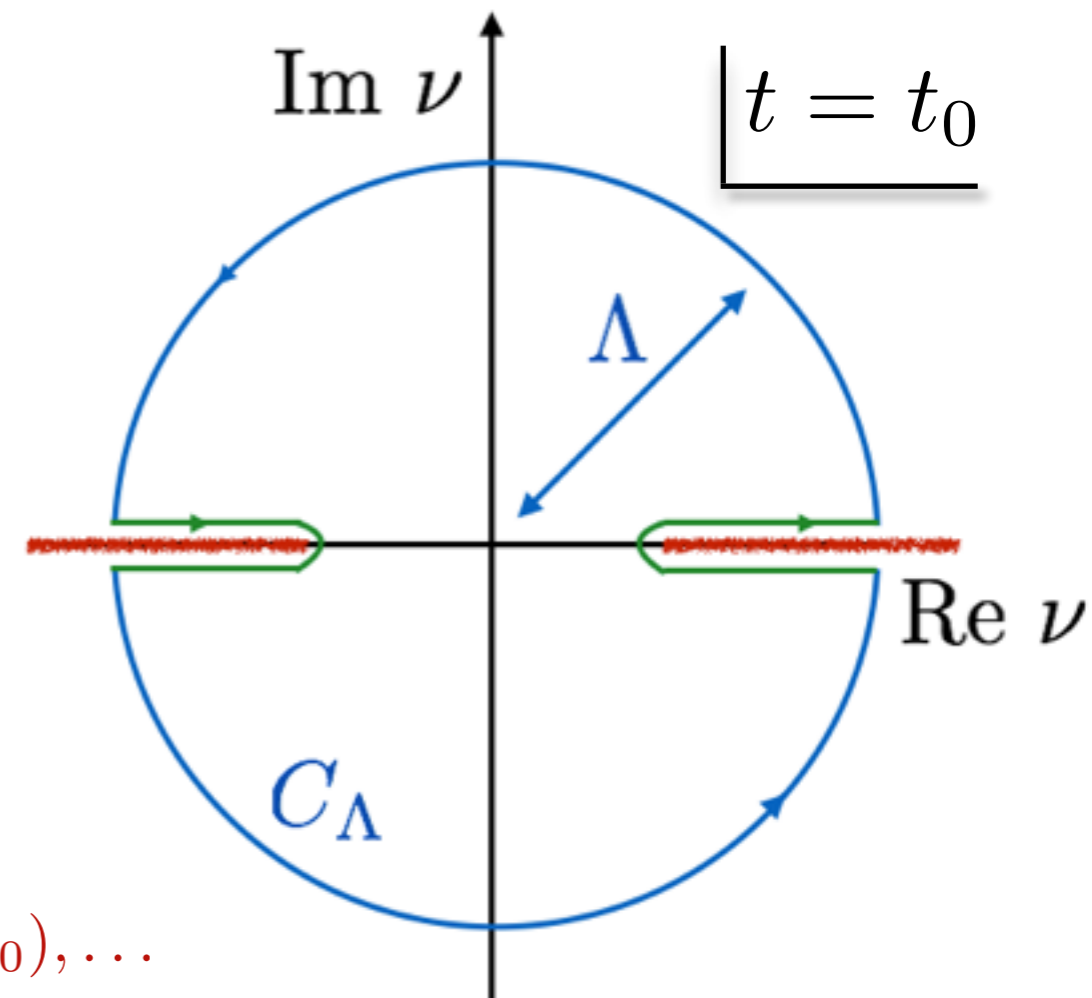
and has no other singularity on sheet I



$$f(\nu) = \int_{\nu_0}^{\Lambda} \left(\frac{\text{Im } f(\nu')}{\nu' - \nu} + \frac{\text{Im } f(-\nu')}{\nu' + \nu} \right) \frac{d\nu'}{\pi} + \oint_{C_{\Lambda}} \frac{f(\nu')}{\nu' - \nu} \frac{d\nu'}{2i\pi} \quad (+\text{sub.}) \quad (I)$$

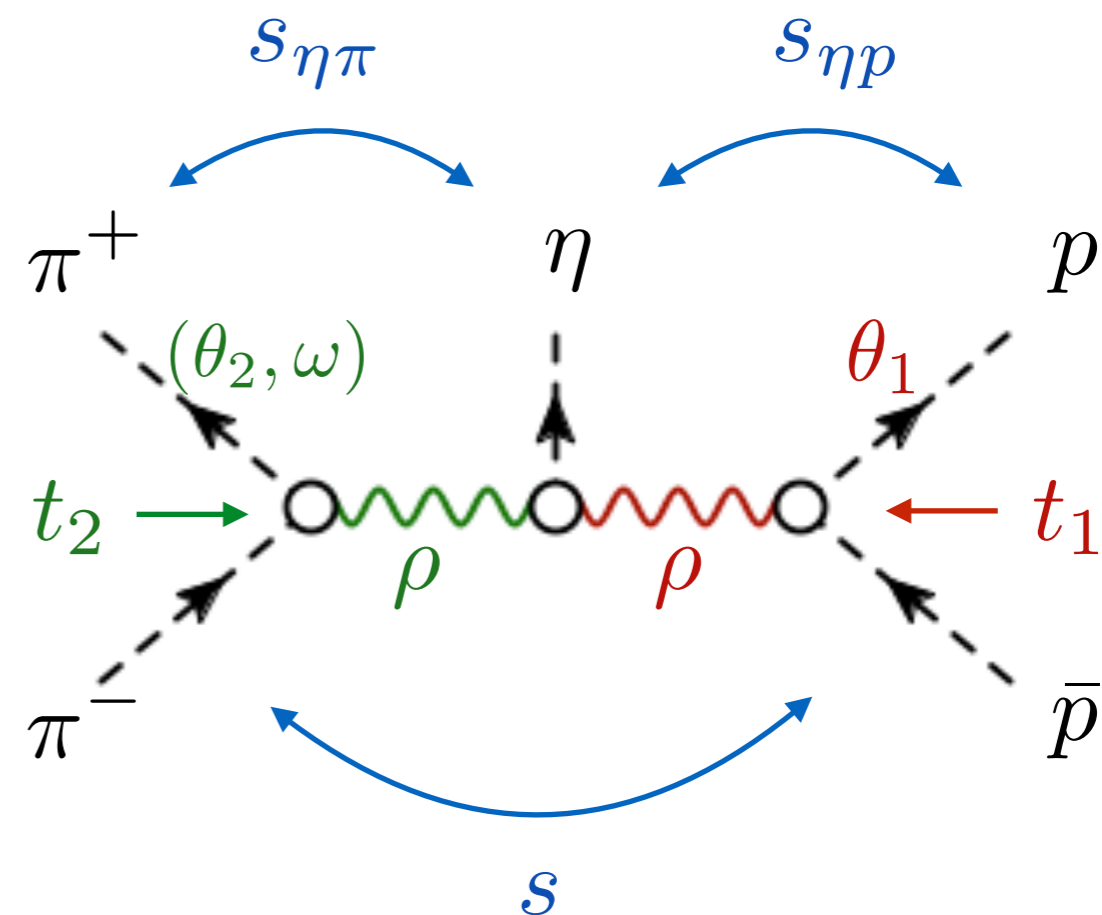
$$\int_{\nu_0}^{\Lambda} \text{Im } f(\nu') + \text{Im } f(-\nu') d\nu' = \oint_{C_{\Lambda}} f(\nu') \frac{d\nu'}{2i\pi} \quad (II)$$

$$f(\nu) = \dots, \frac{T(\nu, t_0)}{\nu^2}, \frac{T(\nu, t_0)}{\nu}, T(\nu, t_0), \nu T(\nu, t_0), \nu^2 T(\nu, t_0), \dots$$



case with 'scalar protons'

work at fixed t_1 and t_2



Partial wave expansion

$$A(\theta_1, \theta_2, \omega, t_1, t_2) = \sum_{J_1=0}^{\infty} \sum_{J_2=0}^{\infty} \sum_{\lambda=-M}^M d_{0\lambda}^{J_1}(\theta_1) e^{i\lambda\omega} d_{\lambda 0}^{J_2}(\theta_2) a_{\lambda}^{J_1 J_2}(t_1, t_2)$$

Example

$$A^{\rho\rho}(\theta_1, \theta_2, \omega, t, t_2) \propto \sin \theta_1 \sin \theta_2 \sin \omega$$

Mellin representation

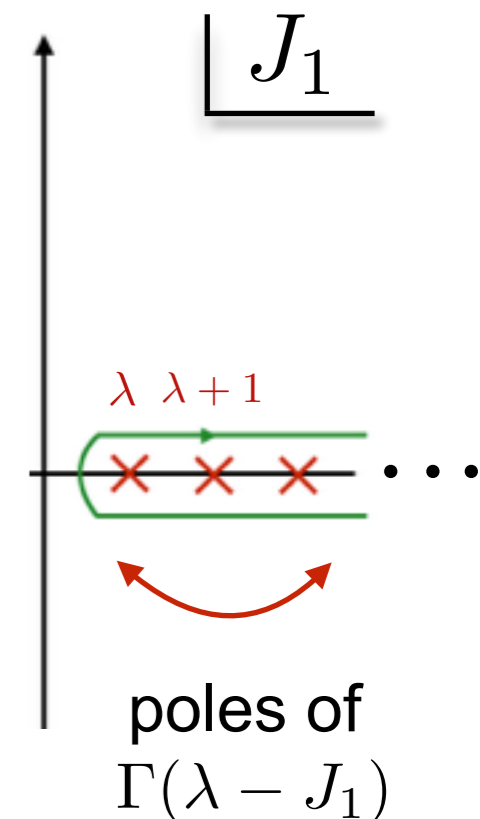
$$A(s, s_{\eta\pi}, s_{\eta p}, t_1, t_2) = \oint \frac{d\lambda}{2\pi i} \oint \frac{dJ_1}{2\pi i} \oint \frac{dJ_2}{2\pi i} \Gamma(1-\lambda) \Gamma(\lambda-J_1) \Gamma(\lambda-J_2) (-s)^\lambda (-s_{\eta p})^{J_1-\lambda} (-s_{\eta\pi})^{J_2-\lambda} a(J_1, J_2, \lambda, t_1, t_2)$$

Mellin representation:

$$A(s, s_{\eta\pi}, s_{\eta p}, t_1, t_2) = \oint \frac{d\lambda}{2\pi i} \oint \frac{dJ_1}{2\pi i} \oint \frac{dJ_2}{2\pi i} \Gamma(1 - \lambda) \Gamma(\lambda - J_1) \Gamma(\lambda - J_2) (-s)^\lambda (-s_{\eta p})^{J_1 - \lambda} (-s_{\eta\pi})^{J_2 - \lambda} a(J_1, J_2, \lambda, t_1, t_2)$$

Do the J_1 integration:

$$A = (-s_{\eta p})^{\alpha_1} \oint \frac{d\lambda}{2\pi i} \oint \frac{dJ_2}{2\pi i} \Gamma(1 - \lambda) \Gamma(\lambda - \alpha_1) \Gamma(\lambda - J_2) (s/s_{\eta p})^\lambda (-s_{\eta\pi})^{J_2 - \lambda} R(J_2, \lambda, t_1, t_2)$$



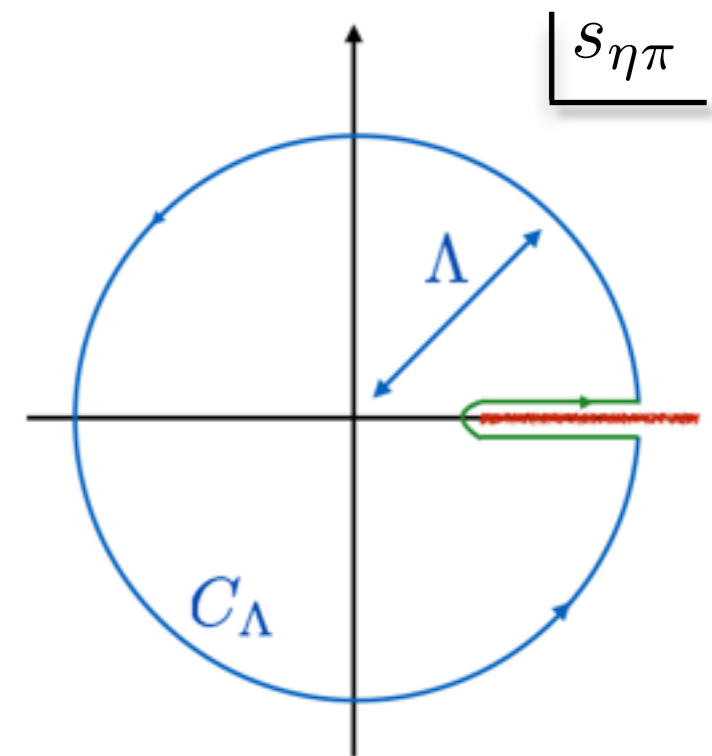
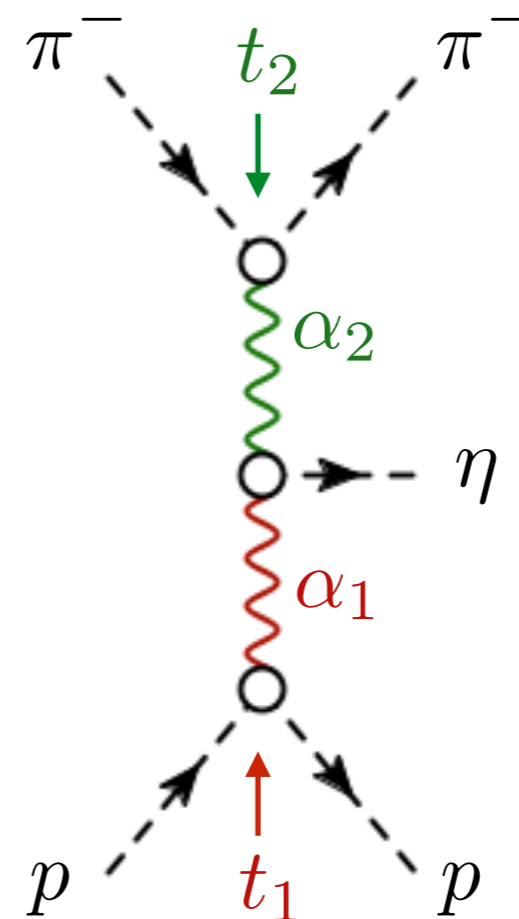
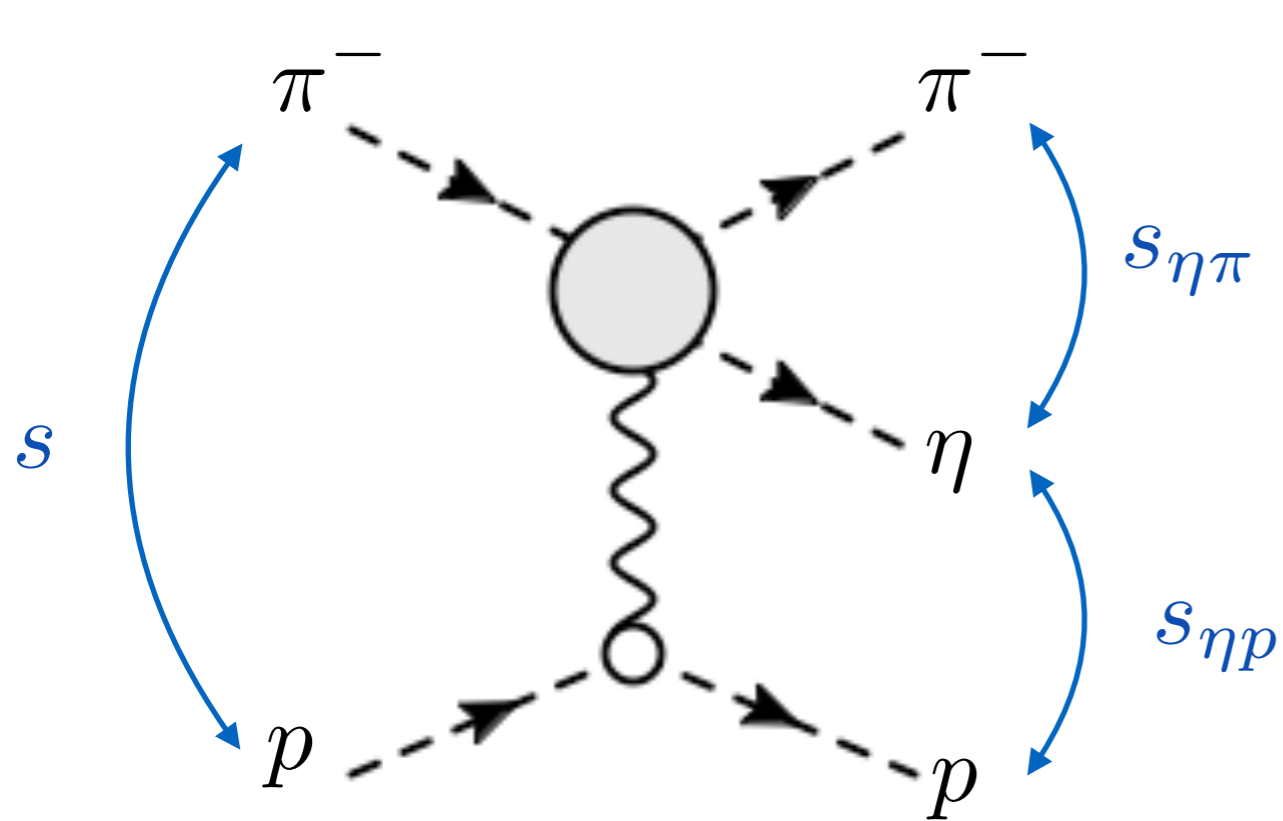
Do the λ integration:

$$A = (-s)^{\alpha_1} \left\{ \sum_{i=0}^{\infty} \frac{\Gamma(1 + i - \alpha_1)}{\Gamma(i + 1)} \left(-\frac{s_{\eta p}}{s} \right)^i \left[\oint \frac{dJ_2}{2\pi i} \Gamma(\alpha_1 - J_2 - i) (-s_{\eta\pi})^{J_2 - \alpha_1 + i} R(J_2, \alpha_1 - i, t_1, t_2) \right] \right. \\ \left. + \left(\frac{s_{\eta p}}{s} \right)^{\alpha_1 - \alpha_2} \sum_{i=0}^{\infty} \frac{\Gamma(1 + i - \alpha_1)}{\Gamma(i + 1)} \Gamma(\alpha_2 - \alpha_1 - i) \beta(\alpha_2 - i, t_1, t_2) \left(\frac{s_{\eta\pi} s_{\eta p}}{s} \right)^i \right\}$$

$$A = (-s)^{\alpha_1} \left\{ \sum_{i=0}^{\infty} \frac{\Gamma(1+i-\alpha_1)}{\Gamma(i+1)} \left(-\frac{s_{\eta p}}{s}\right)^i \left[\oint \frac{dJ_2}{2\pi i} \Gamma(\alpha_1 - J_2 - i) (-s_{\eta\pi})^{J_2 - \alpha_1 + i} R(J_2, \alpha_1 - i, t_1, t_2) \right] \right. \\ \left. + \left(\frac{s_{\eta p}}{s}\right)^{\alpha_1 - \alpha_2} \sum_{i=0}^{\infty} \frac{\Gamma(1+i-\alpha_1)}{\Gamma(i+1)} \Gamma(\alpha_2 - \alpha_1 - i) \beta(\alpha_2 - i, t_1, t_2) \left(\frac{s_{\eta\pi} s_{\eta p}}{s}\right)^i \right\}$$

infinite number of subtractions

Reggeon-particle amplitude



Consider $t_1, t_2, s_{\eta p}/s$ fixed

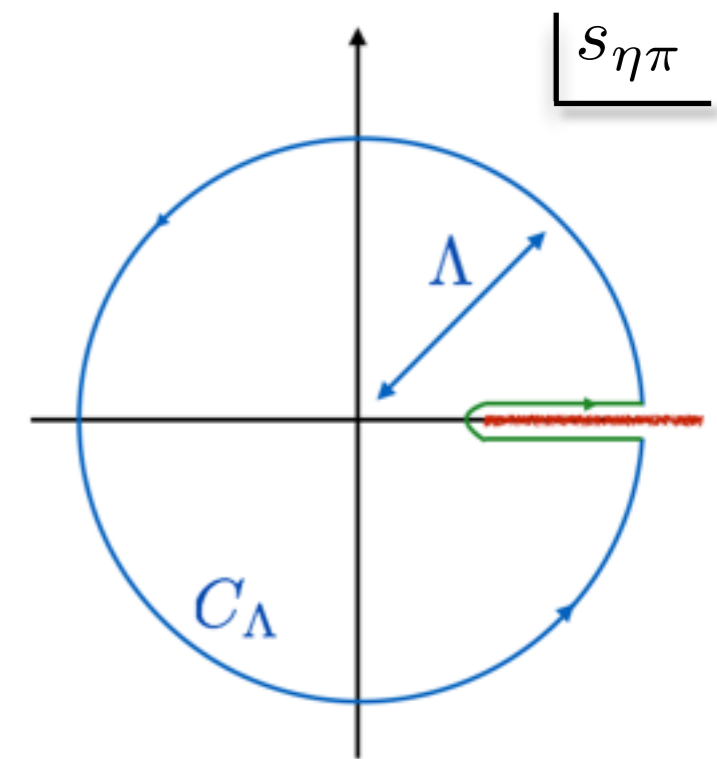
$$A = (-s)^{\alpha_1} \left\{ \sum_{i=0}^{\infty} \frac{\Gamma(1+i-\alpha_1)}{\Gamma(i+1)} \left(-\frac{s_{\eta p}}{s}\right)^i \left[\oint \frac{dJ_2}{2\pi i} \Gamma(\alpha_1 - J_2 - i) \underline{(-s_{\eta\pi})}^{J_2 - \alpha_1 + i} R(J_2, \alpha_1 - i, t_1, t_2) \right] \right. \\ \left. + \left(\frac{s_{\eta p}}{s}\right)^{\alpha_1 - \alpha_2} \sum_{i=0}^{\infty} \frac{\Gamma(1+i-\alpha_1)}{\Gamma(i+1)} \Gamma(\alpha_2 - \alpha_1 - i) \beta(\alpha_2 - i, t_1, t_2) \left(\frac{s_{\eta\pi} s_{\eta p}}{s}\right)^i \right\}$$

Consider $t_1, t_2, s_{\eta p}/s$ fixed

DR with infinite number of subtractions

Close the contour at finite Λ

and obtain sum rules:



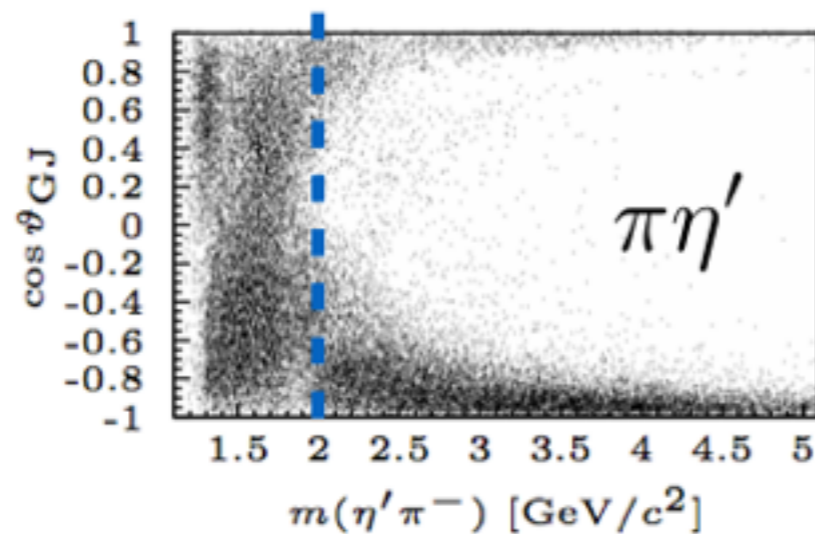
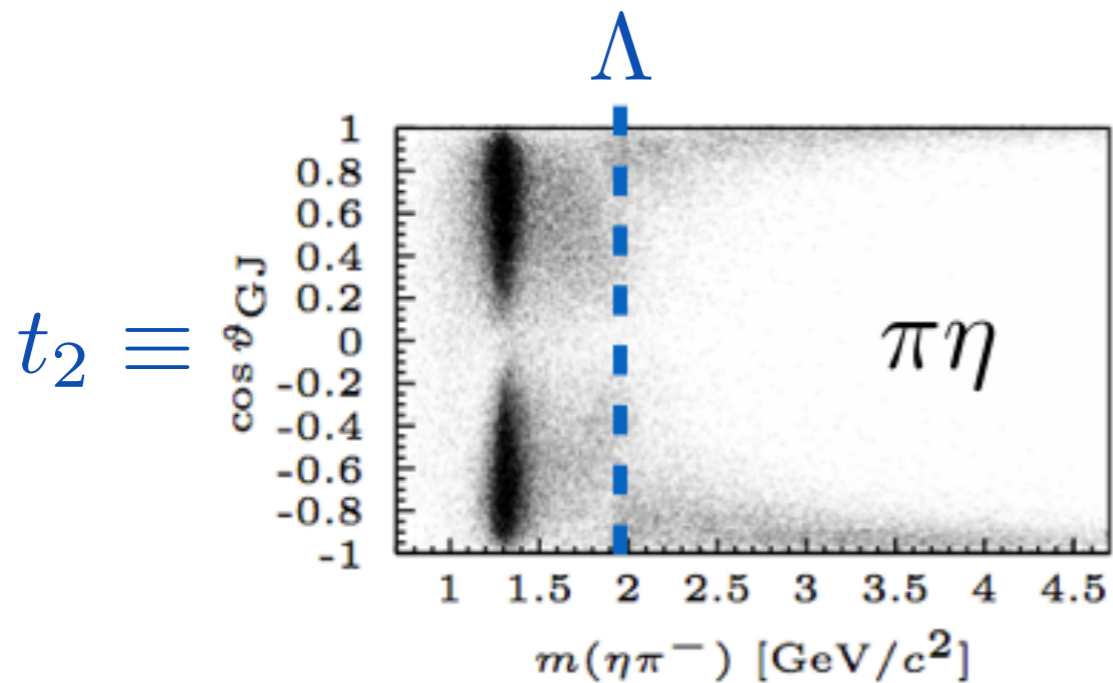
$$\int_{s_{\text{th}}}^{\Lambda} \text{Im} A(s, s_{\eta p}, s'_{\eta\pi}, t_1, t_2) ds'_{\eta\pi} = (-s)^{\alpha_1} \Lambda^{\alpha_2 - \alpha_1 + 1} \sum_{i=0}^{\infty} \frac{\Gamma(1+i-\alpha_1)}{\Gamma(i+1)} \frac{\beta(\alpha_1 - i, t_1, t_2)}{\Gamma(2+i+\alpha_2 - \alpha_1)} \left(\frac{\Lambda s_{\eta p}}{s}\right)^i$$

Sum rules at $t_1, t_2, s_{\eta p}/s$ fixed

$$\int_{s_{\text{th}}}^{\Lambda} \text{Im } A(s, s_{\eta p}, s'_{\eta\pi}, t_1, t_2) ds'_{\eta\pi} = (-s)^{\alpha_1} \Lambda^{\alpha_2 - \alpha_1 + 1} \sum_{i=0}^{\infty} \frac{\Gamma(1 + i - \alpha_1)}{\Gamma(i + 1)} \frac{\beta(\alpha_1 - i, t_1, t_2)}{\Gamma(2 + i + \alpha_2 - \alpha_1)} \left(\frac{\Lambda s_{\eta p}}{s} \right)^i$$

constraint on low mass model

fixed ; determined by high mass data



Analyticity and unitarity constrain amplitude construction

Past:

Extracted $a_2(1320)$ and $a_2(1700)$ pole position

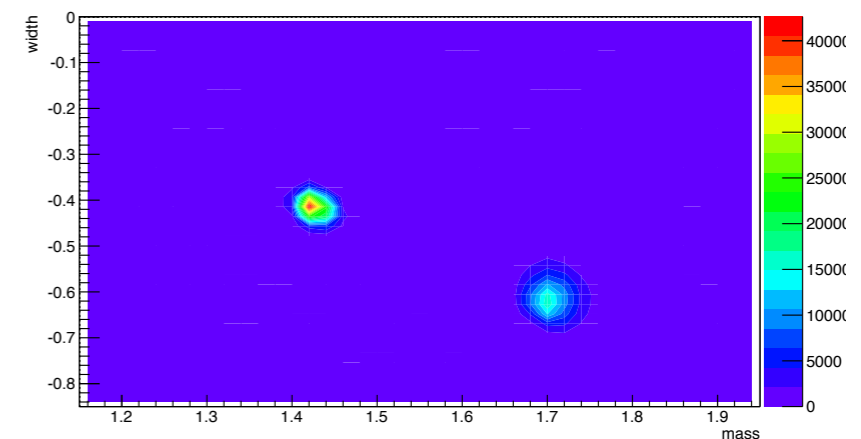
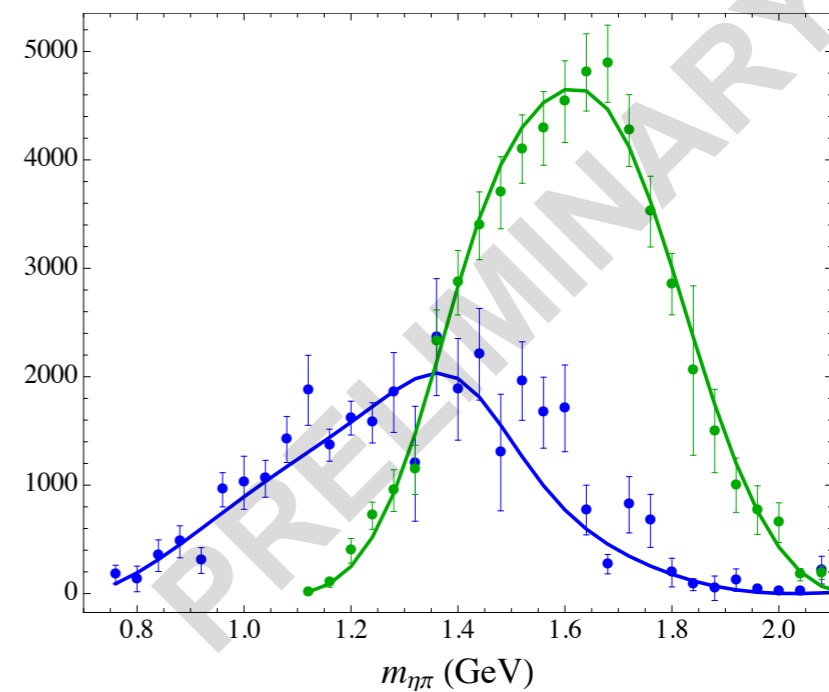
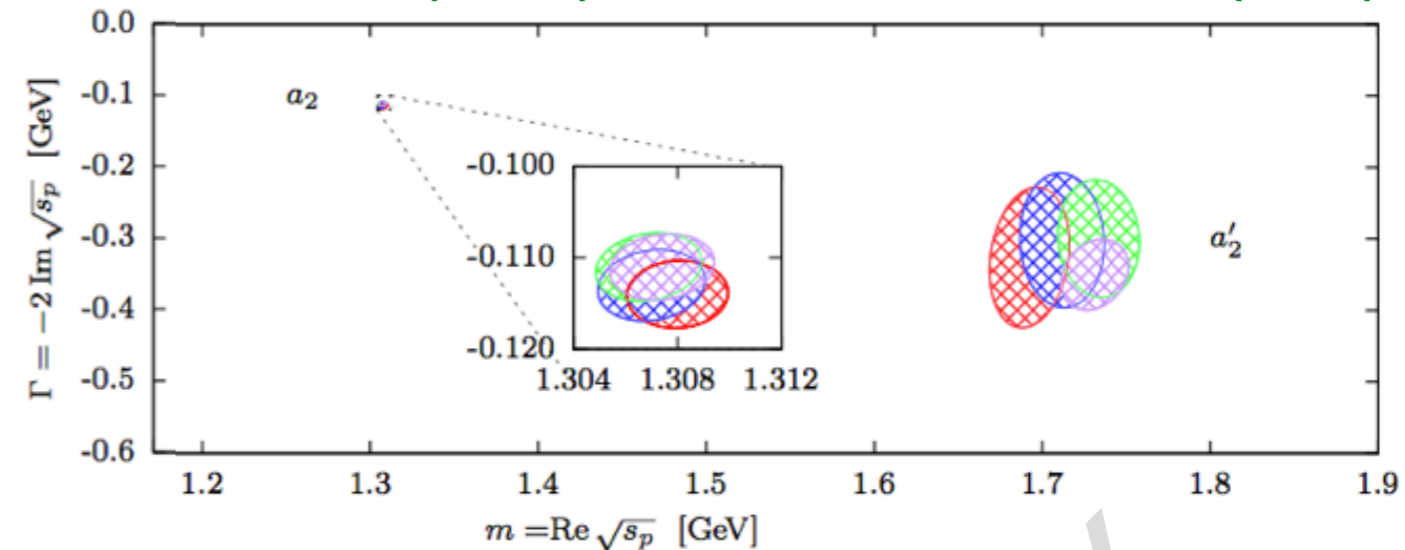
Present:

Extraction of exotic meson pole position

Future:

Implementation of DR constraining model
Transposition to GlueX/CLAS12 data

A. Jackura et al (JPAC) and COMPASS, PLB779 (2018)



Interactive webpage: <http://www.indiana.edu/~jpac/>



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Resources

- **Publication:** [Mat15a]
- **Fortran:** Fortran file, Input file, Output file
- **C/C++:** AmpTools class, C/C++ file, AmpTools class header
- **Mathematica:** notebook , converted in text
- **Data:** Anderson, All data
- **Contact person:** Vincent Mathieu
- **Last update:** November 2015

Description of the Fortran code: [show/hide]

Description of the C/C++ code: [show/hide]

Run the code

Choose the beam energy in the lab frame E_γ , the other variable (t or $\cos \theta$) and its minimal, maximal, and increment values. If you choose t (\cos) only the min, max and step values of t ($\cos \theta$) are read.

E_γ in GeV

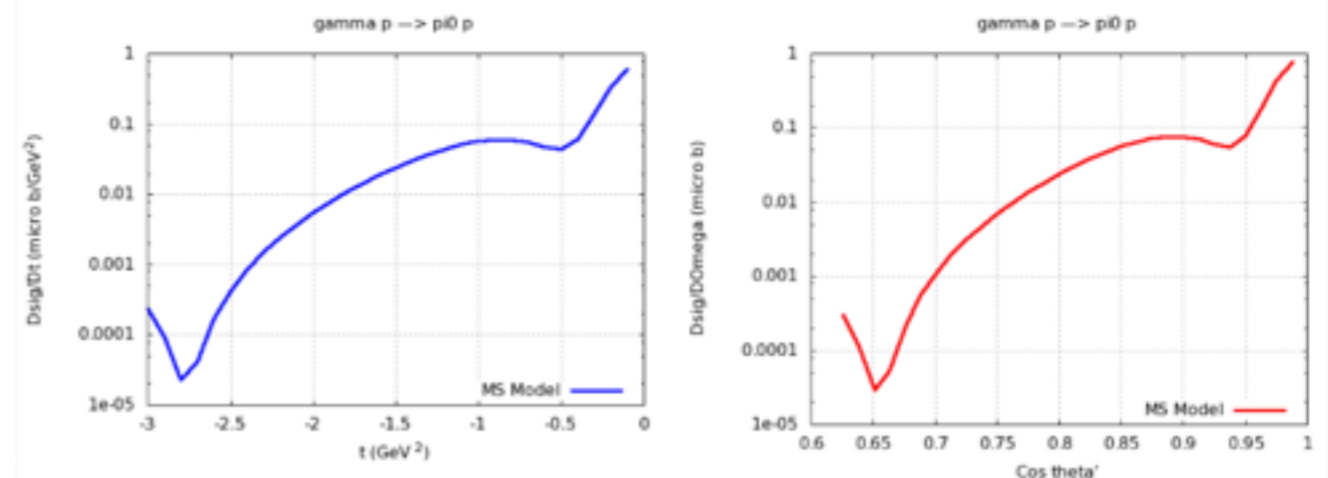
t \cos

t in GeV² (min max step)

$\cos \theta$ (min max step)

beam energy: 9 GeV
Observable: differential cross section
X variable: t with interval -3:0.1:-0.1

Download the output file, the plot with $O_x=t$, the plot with $O_x=\cos$.
In the file, the columns are: t (GeV²), \cos , $D\text{sig}/Dt$ (micro barn/GeV²), $D\text{sig}/D\Omega$ (micro barn)



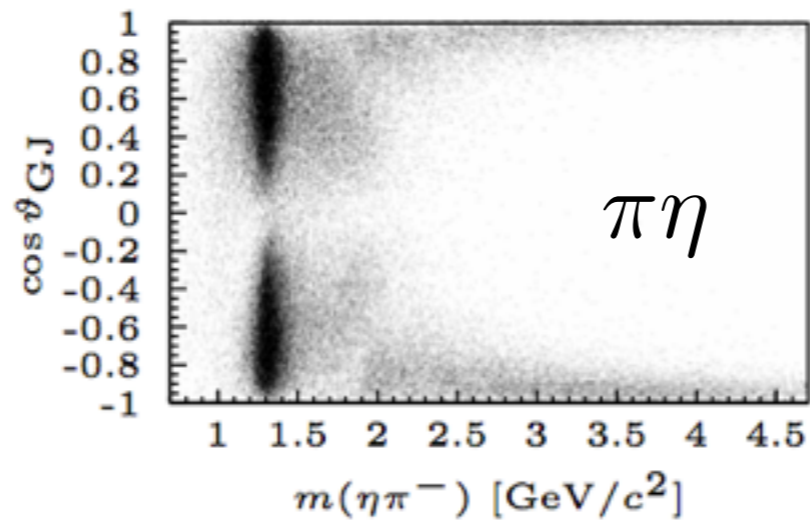
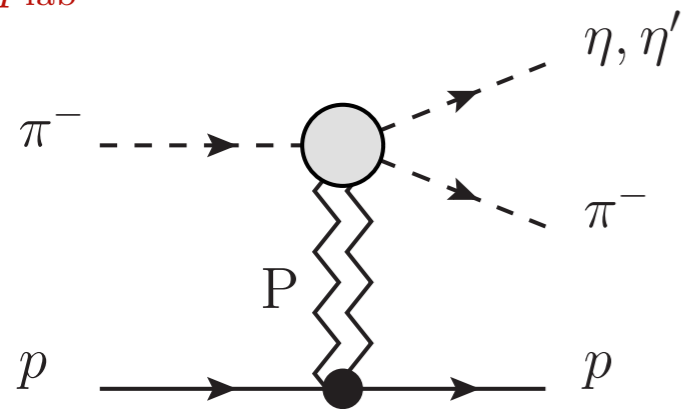
Interactive webpage: <http://www.indiana.edu/~jpac/>



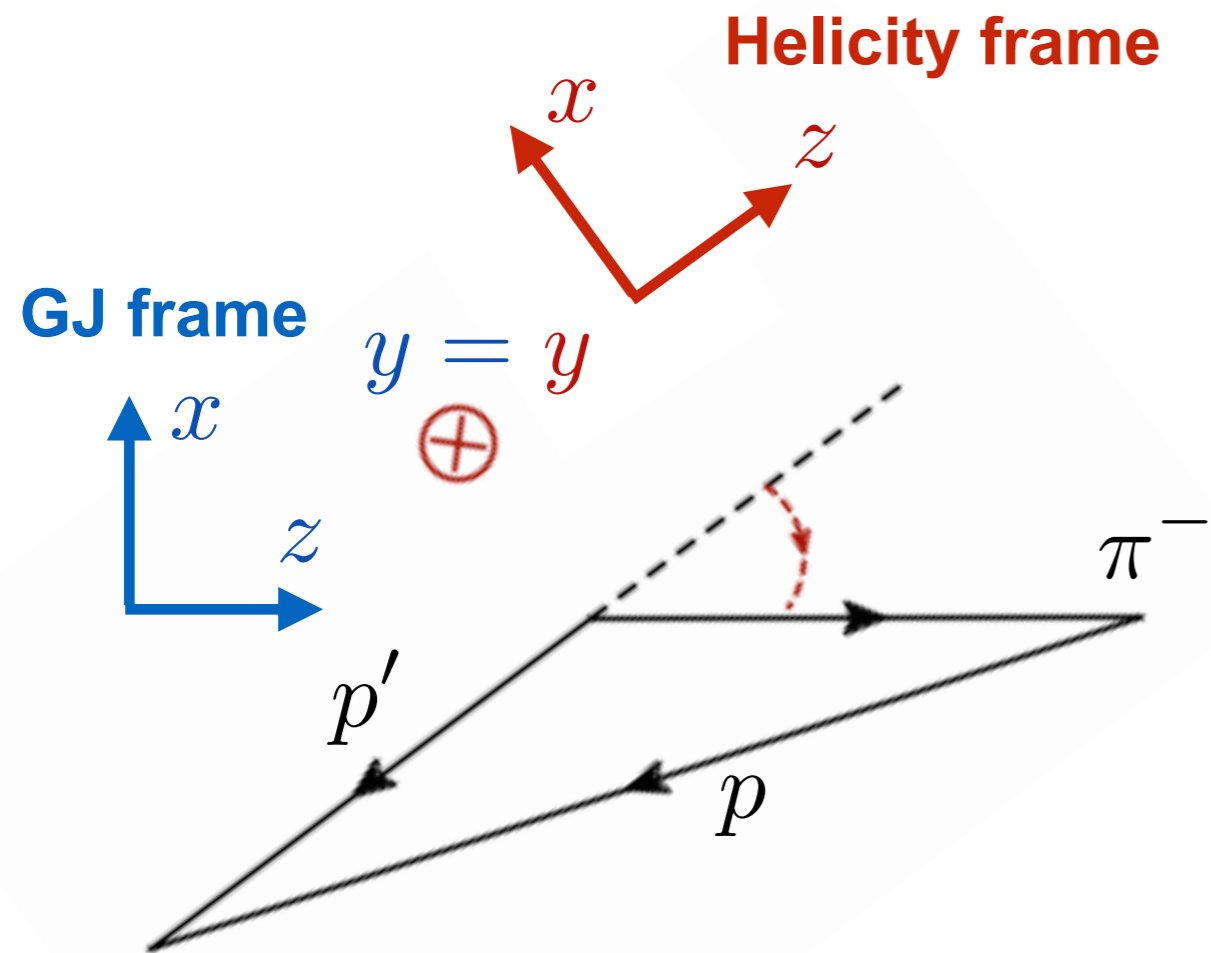
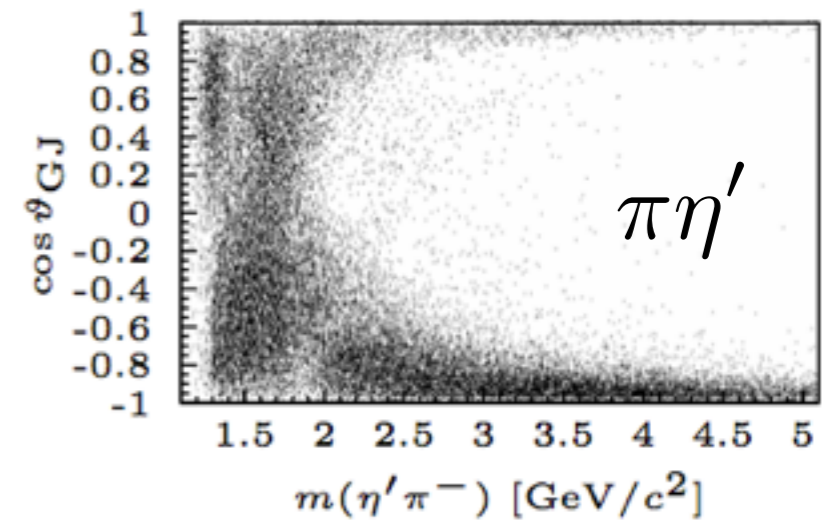
1.	United States	493 (32.18%)	14.	Switzerland	25 (1.63%)
2.	Germany	169 (11.03%)	15.	Russia	20 (1.31%)
3.	United Kingdom	150 (9.79%)	16.	Saudi Arabia	20 (1.31%)
4.	Mexico	103 (6.72%)	17.	Iraq	17 (1.11%)
5.	Italy	57 (3.72%)	18.	India	16 (1.04%)
6.	Spain	52 (3.39%)	19.	Canada	15 (0.98%)
7.	Belgium	38 (2.48%)	20.	Netherlands	15 (0.98%)
8.	China	37 (2.42%)	21.	France	12 (0.78%)
9.	(not set)	37 (2.42%)	22.	Bosnia & Herzegovina	10 (0.65%)
10.	Brazil	35 (2.28%)	23.	Greece	10 (0.65%)
11.	Poland	31 (2.02%)	24.	Portugal	9 (0.59%)
12.	Austria	30 (1.96%)	25.	South Korea	7 (0.46%)
13.	Japan	26 (1.70%)			

Backup Slides

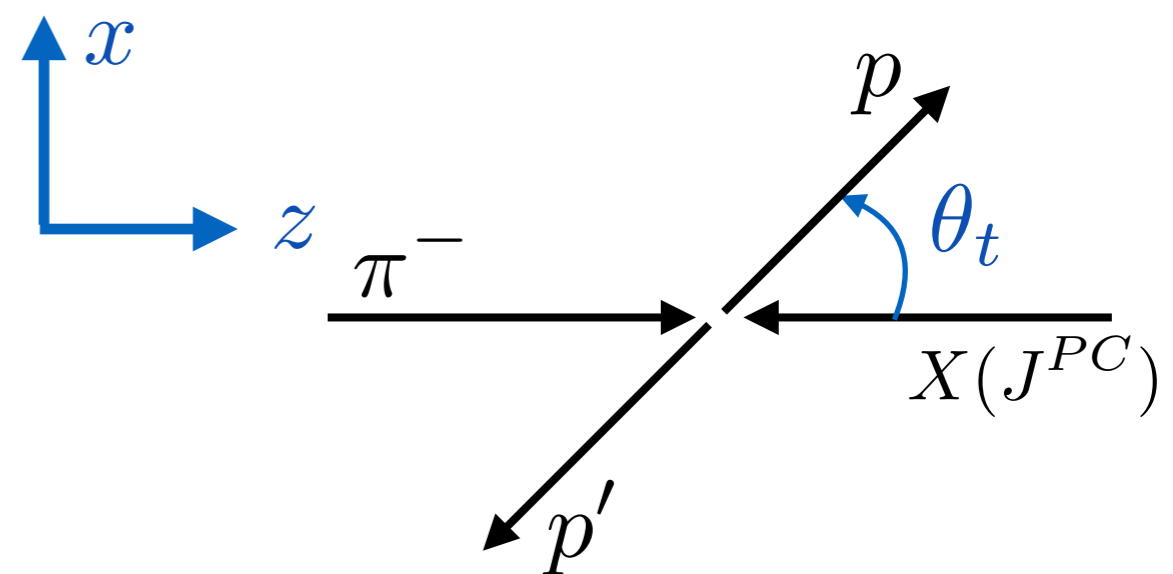
$p_{\text{lab}} = 190 \text{ GeV}$



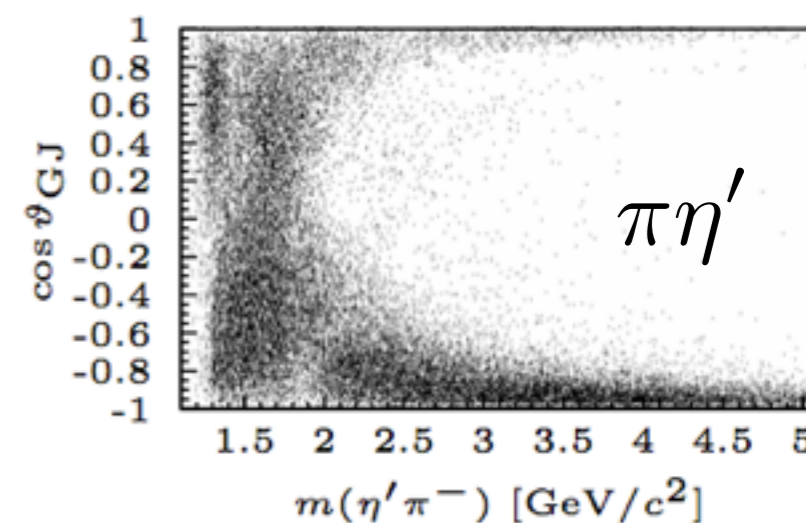
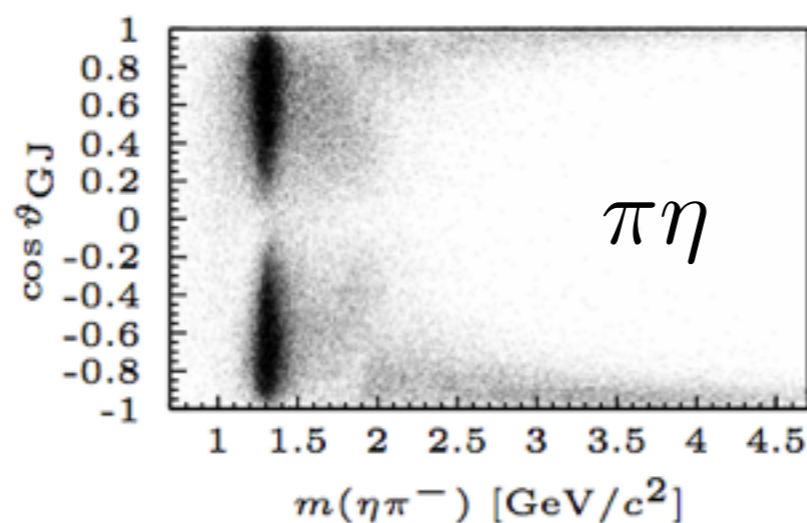
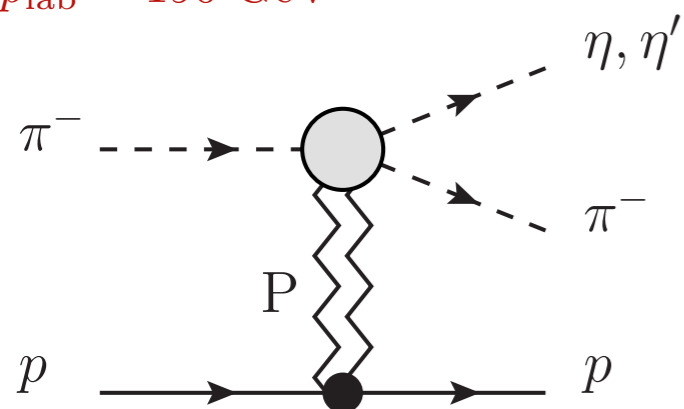
COMPASS Phys. Lett. B740 (2015)



t-channel frame

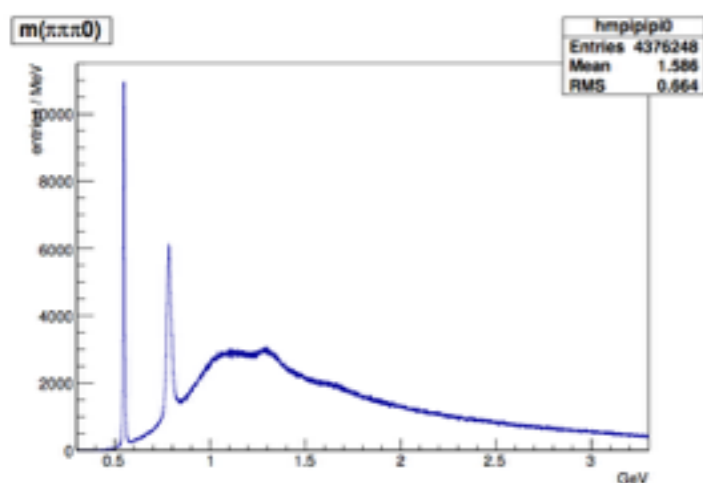


COMPASS Phys. Lett. B740 (2015)

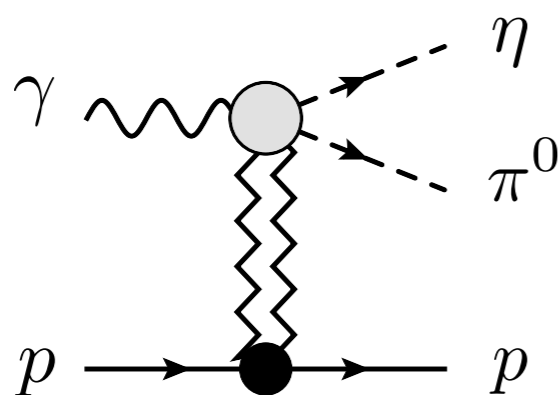
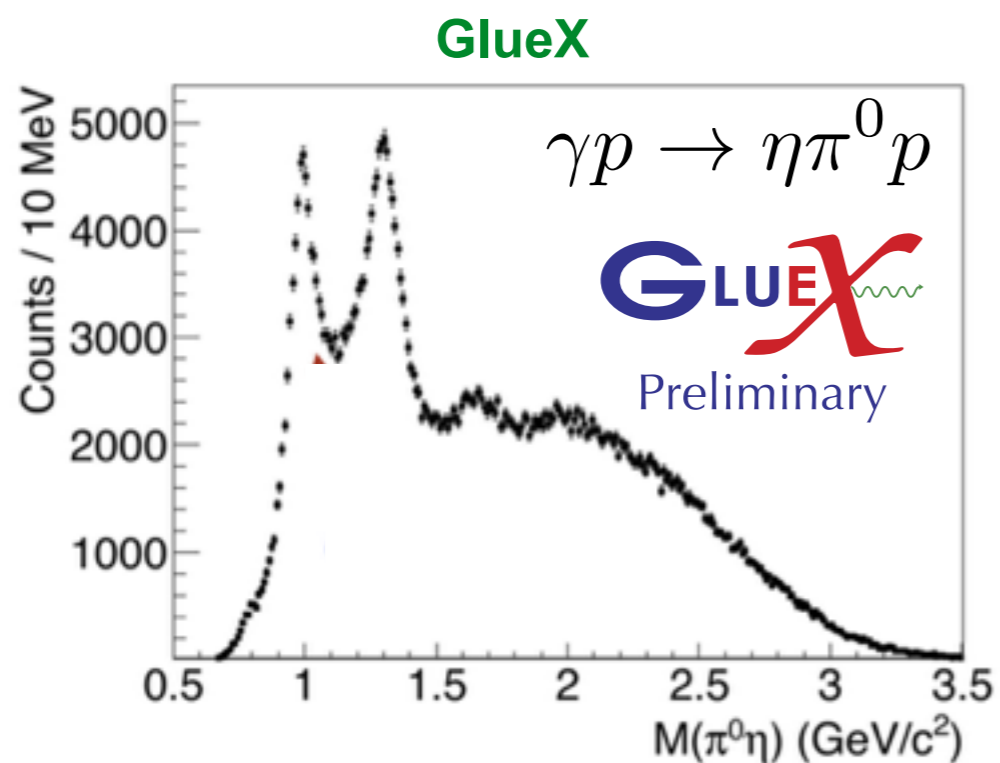
 $p_{\text{lab}} = 190 \text{ GeV}$


$$I(\tau) = \sum_{\epsilon} \left| \sum_{L,M} A_{LM}^{\epsilon} \psi_{LM}^{\epsilon}(\tau) \right|^2 + \text{non-}\eta^{(\prime)} \text{ background}$$

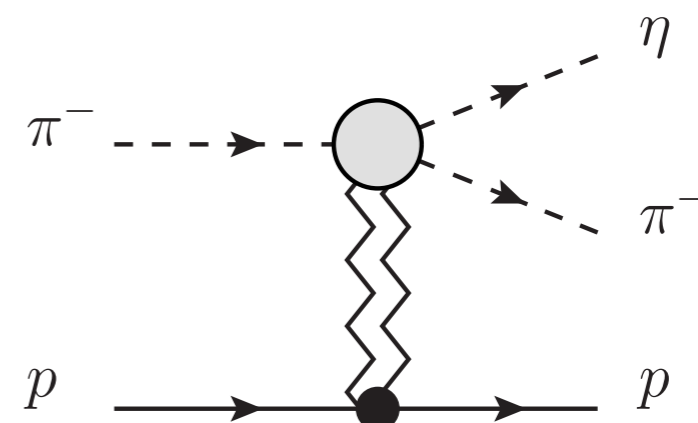
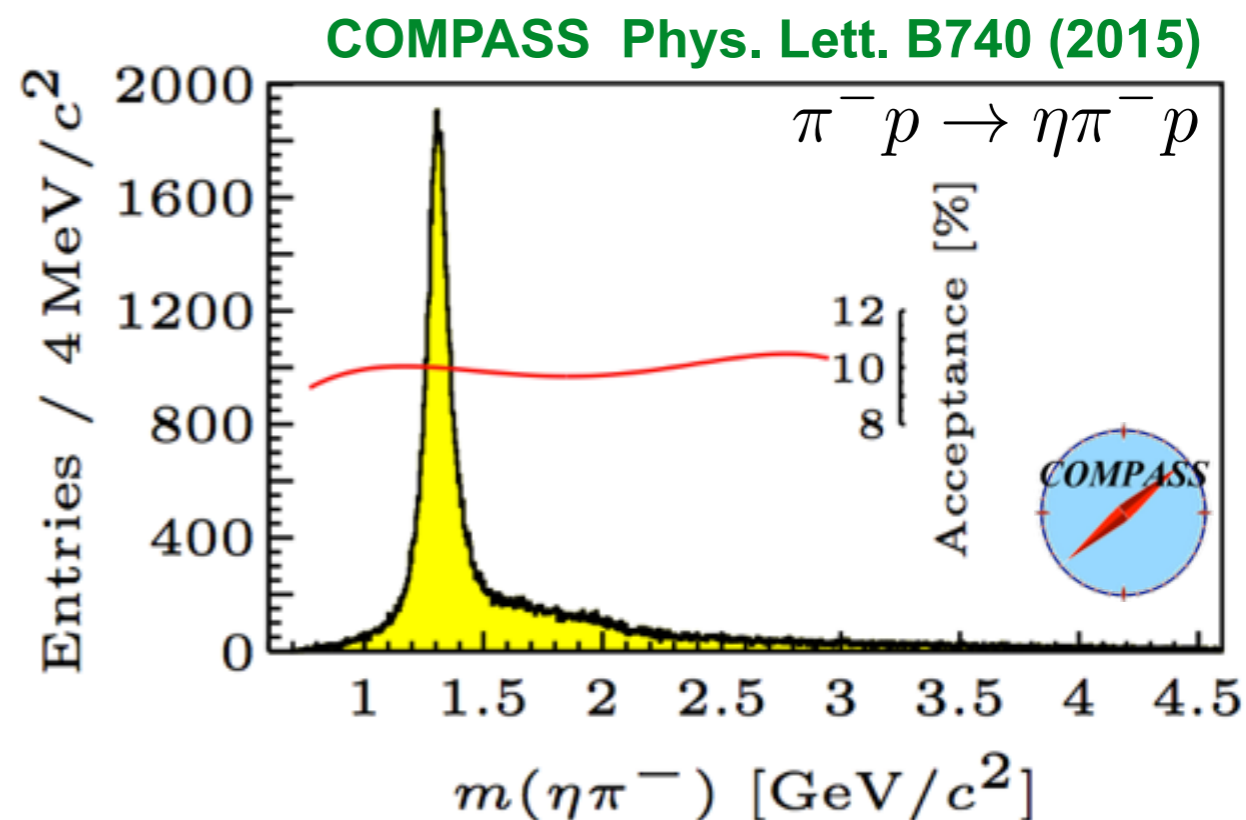
$$\psi_{LM}^{\epsilon}(\tau) = f_{\eta}(p_{\pi^{-}}, p_{\pi^{+}}, p_{\pi^0}) \times Y_L^M(\vartheta_{GJ}, 0) \times \begin{cases} \sin M\varphi_{GJ} & \text{for } \epsilon = +1 \\ \cos M\varphi_{GJ} & \text{for } \epsilon = -1 \end{cases}$$



$$E_{\text{beam}} = 9 \text{ GeV}$$



$$E_{\text{beam}} = 190 \text{ GeV}$$



$a_0(980)$ is produced by unnatural exchange $P(-1)^J = -1$
suppressed at high energy