

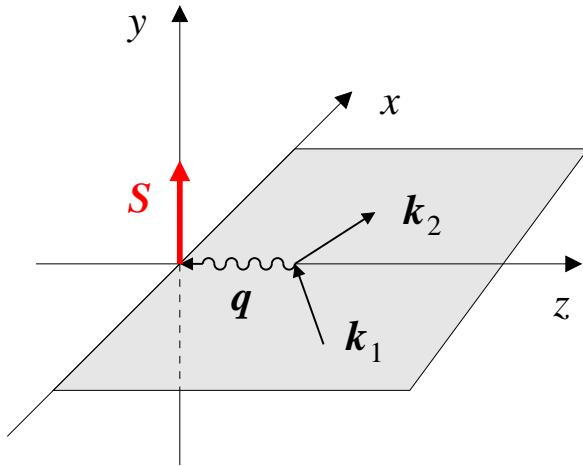
Transversity and spin-orbit correlations in two-photon DIS

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Transverse target single-spin asymmetry
in inclusive $eN(\uparrow) \longrightarrow e'X$

- “Pure” two-photon exchange effect!
- Probes helicity-flip amplitudes at quark level (h, g_T)
- Approved JLab Hall A experiment PR-07-013 [X. Jiang et al.]
Sensitivity $\sim 10^{-4}$ cf. SLAC 1970 $\sim 10^{-2}$

Transverse target spin dependence in $eN \rightarrow e'X$



- Target spin dependence of cross section

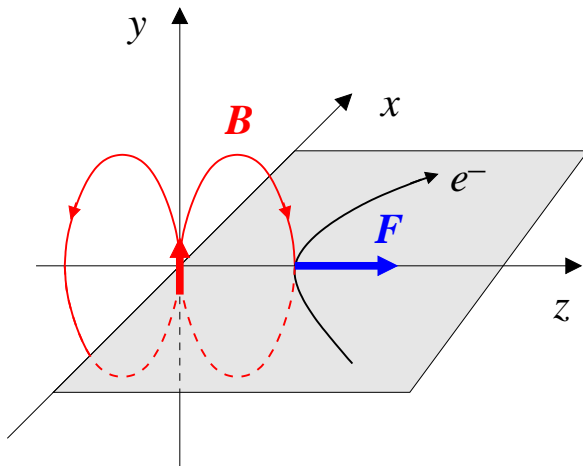
$$\sim \mathbf{S} \cdot (\mathbf{k}_1 \times \mathbf{k}_2) \quad \text{“normal spin”}$$

- Relative asymmetry

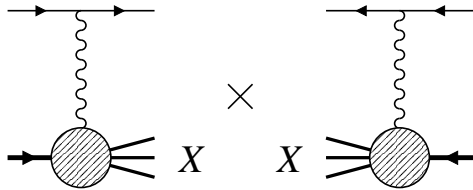
$$A_y = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$$

- Classical analog: Scattering from magnetic dipole (Lorentz force)

→ Sign, p/n ratio



Spin dependence with two-photon exchange



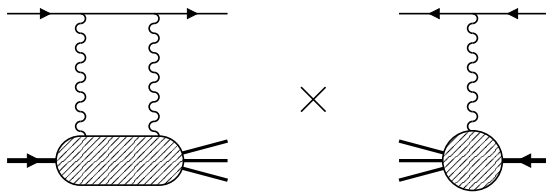
- Transverse spin dependence zero in one-photon exchange [Christ, Lee 66]

$$L_{\mu\nu} = L_{\nu\mu}$$

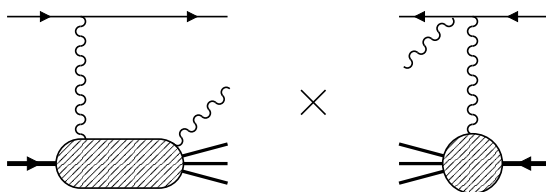
leptonic tensor symmetric (unpol. beam)

$$W_{\mu\nu}(\mathbf{S}) = -W_{\mu\nu}(-\mathbf{S})$$

hadronic tensor antisymmetric (P, T inv.)



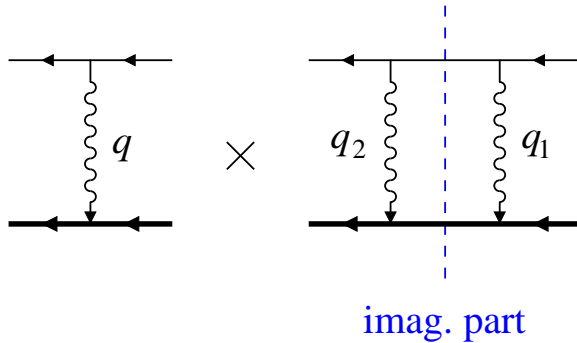
- Nonzero at $O(\alpha^3)$: Two-photon exchange and real emission



- Contributions **individually IR-finite**
 - No cancellations (cf. elastic FF)
 - Clean two-photon exchange effect

Example: Pointlike target

[Barut, Fronsda 60; . . .]



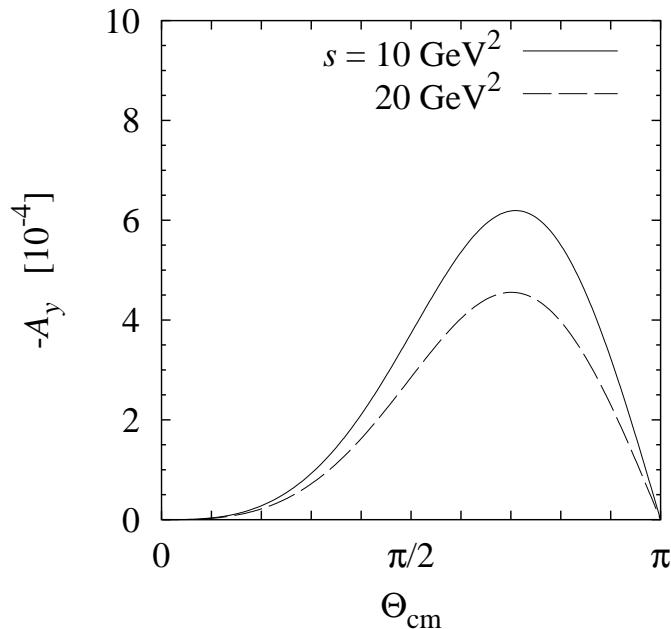
$$A_y = \frac{\alpha M}{\sqrt{s}} f(\Theta_{\text{cm}}) \quad (\Theta_{\text{cm}} \leftrightarrow Q^2/s)$$

- Only **imaginary part** of two-photon exchange enters; no IR divergences

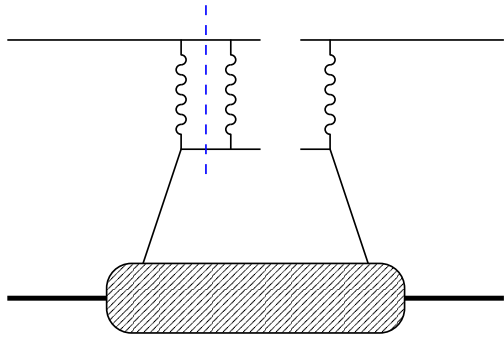
- Photon virtualities $-q_1^2, q_2^2 \sim Q^2$

- Include strong interactions:
No QED collinear divergences thanks to gauge invariance

[Afanasev, Strikman, CW 07]



Transverse spin dependence in QCD

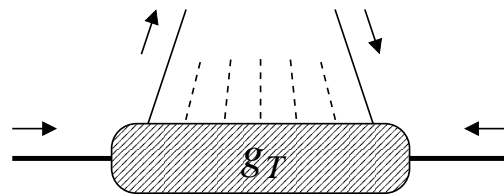


- Dominance of scattering from same quark (no “anomalous” IR/collinear enhancement)

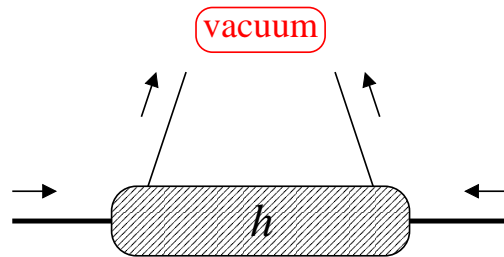
- Two contributions

- I) Quark helicity non-flip and interactions w. spectators

[Goeke, Metz, Schlegel 06 . . . gauge invariance!]

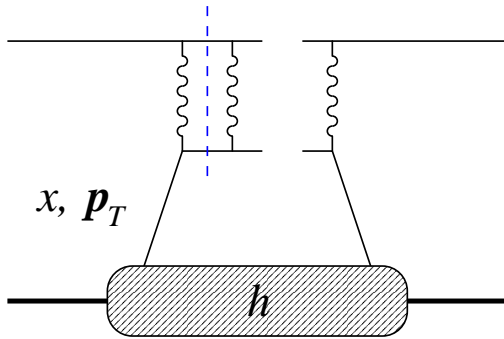


- II) Quark helicity flip by interaction with vacuum fields:
Chiral symmetry breaking



No Sudakov suppression if
IR cutoff $\sim \mu^2(\text{chiral}) \gg \Lambda_{\text{QCD}}^2$

Composite nucleon approximation [ASW 06/07]



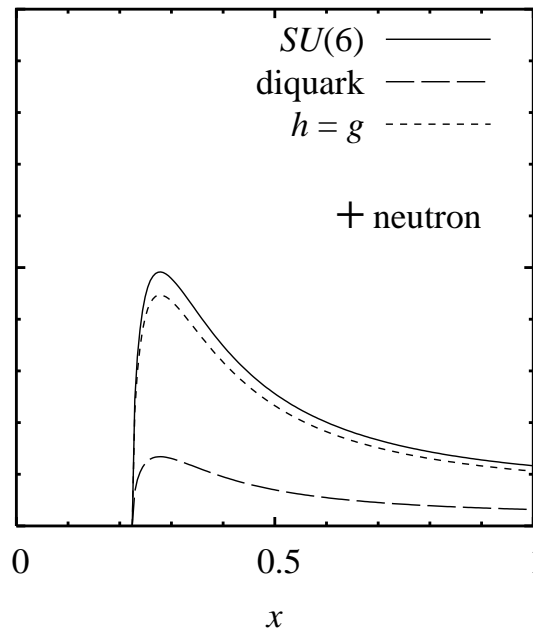
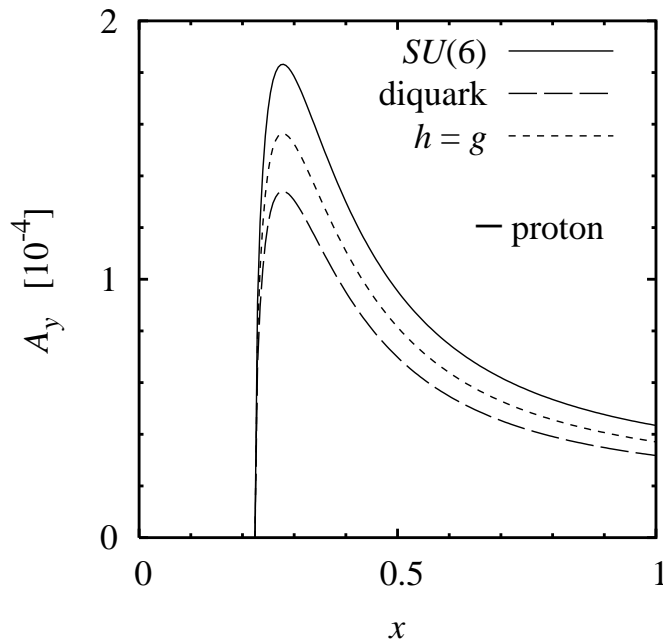
- Assume “composite” nucleon

$$R_N^{-2} \sim \langle p_T^2 \rangle \ll M_q^2$$

→ Quark helicity flip dominates!

→ Light-front constituent quark model

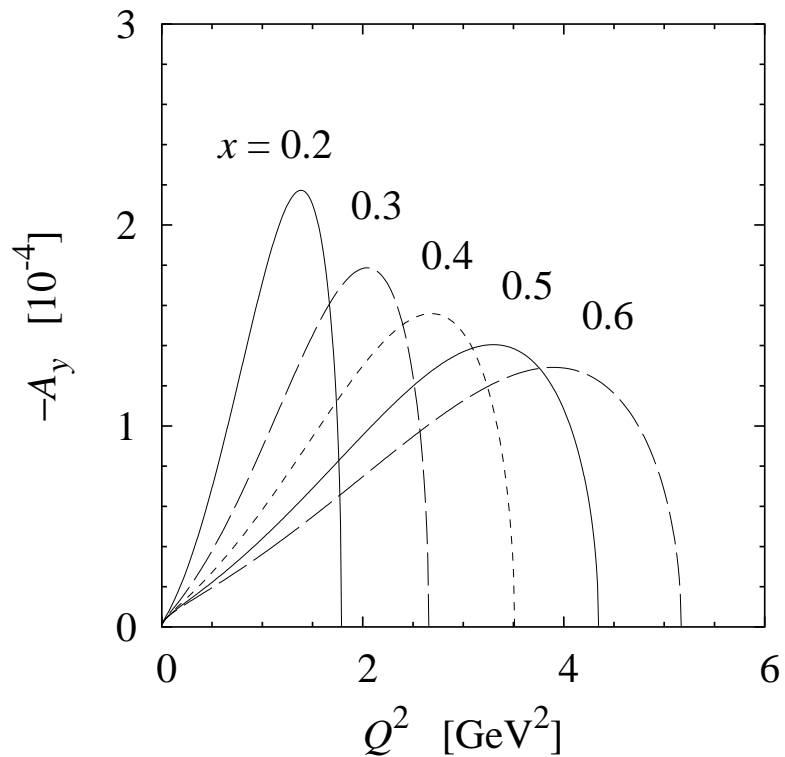
[cf. Miller 02]



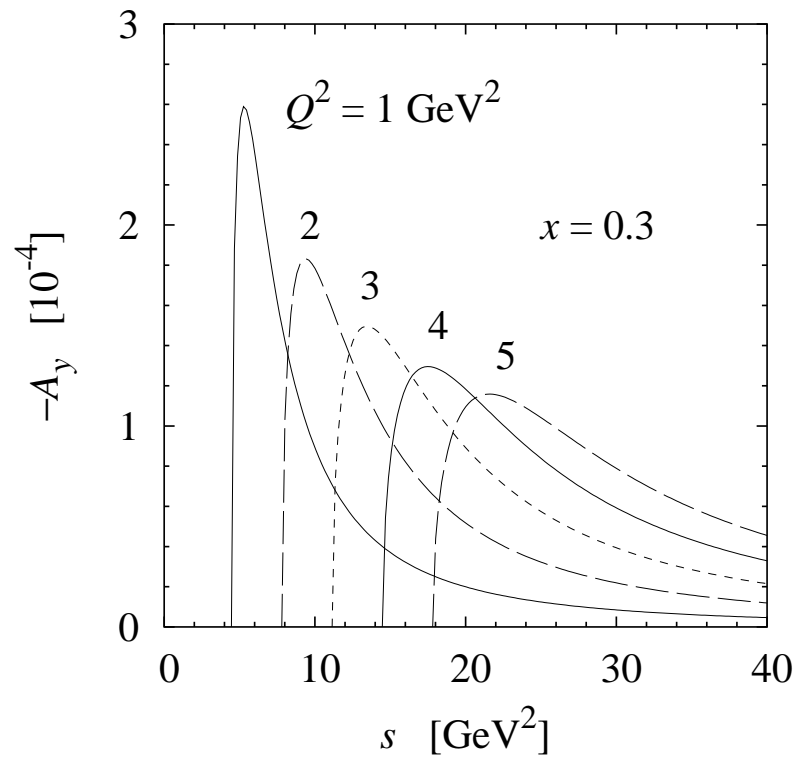
$$A_y = \frac{\sum e_q^3 h_q(x)}{\sum e_q^2 f_q(x)} \times A_y(\text{quark})$$

$$[\propto M_q \approx 300 \text{ MeV}]$$

Predictions for kinematic dependences



$s = 10 \text{ GeV}^2$ [JLab 6 GeV]



- Asymmetry vanishes in high-energy limit $A_y \sim s^{-2}$ ($s \gg Q^2$)
cf. photon polarizations in 2γ box [Gribov, Lipatov, Frolov 70]

Summary

- Very interesting/challenging problem!
 - Higher-order QED corrections
 - QCD factorization
 - Vacuum structure
- “Cleanest” two-photon exchange observable
 - IR finite — no IR cancellations with real emission
 - How large are finite contributions from real emission?
- Probes helicity-flip amplitudes at quark level
 - Composite Nucleon Approximation $\rightarrow h(x)$ transversity
 - How large are helicity-conserving contributions?