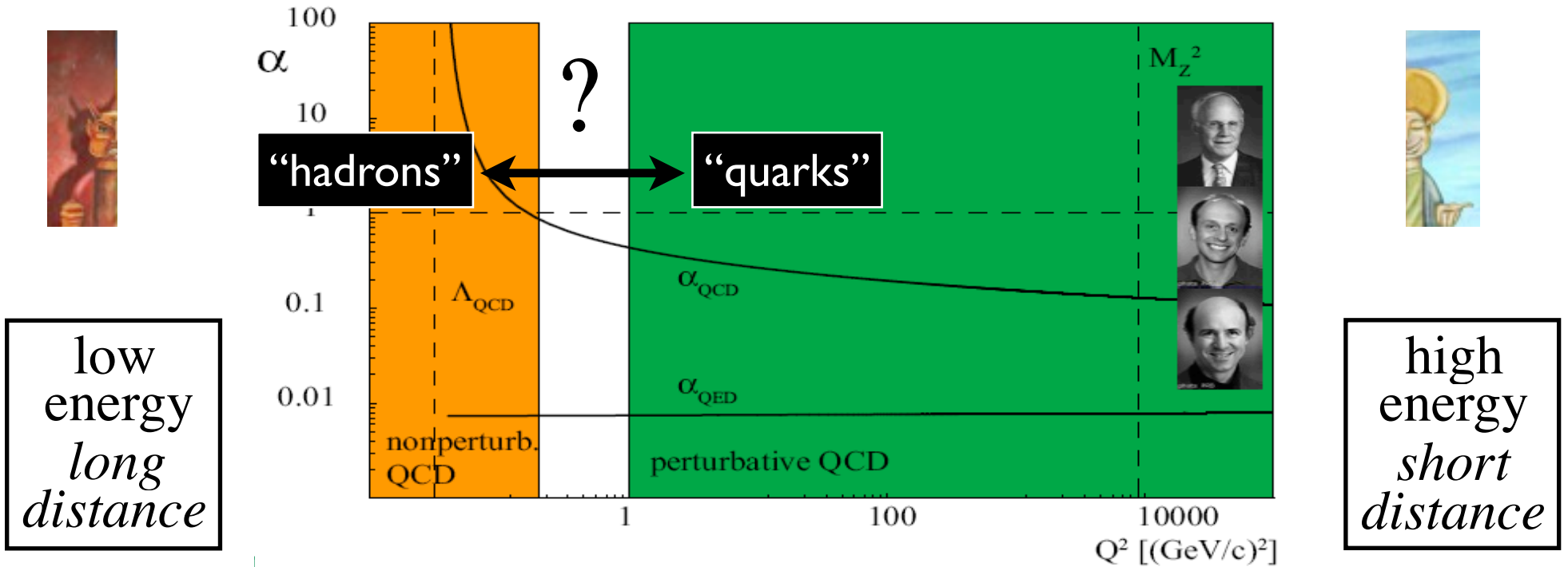


N 2013 Conference
Peñíscola, Spain, May 27, 2013*

N^* Resonances and Duality in Deep-Inelastic Scattering

Wally Melnitchouk

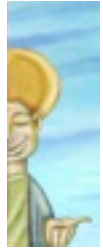
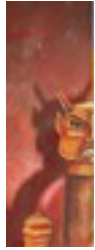




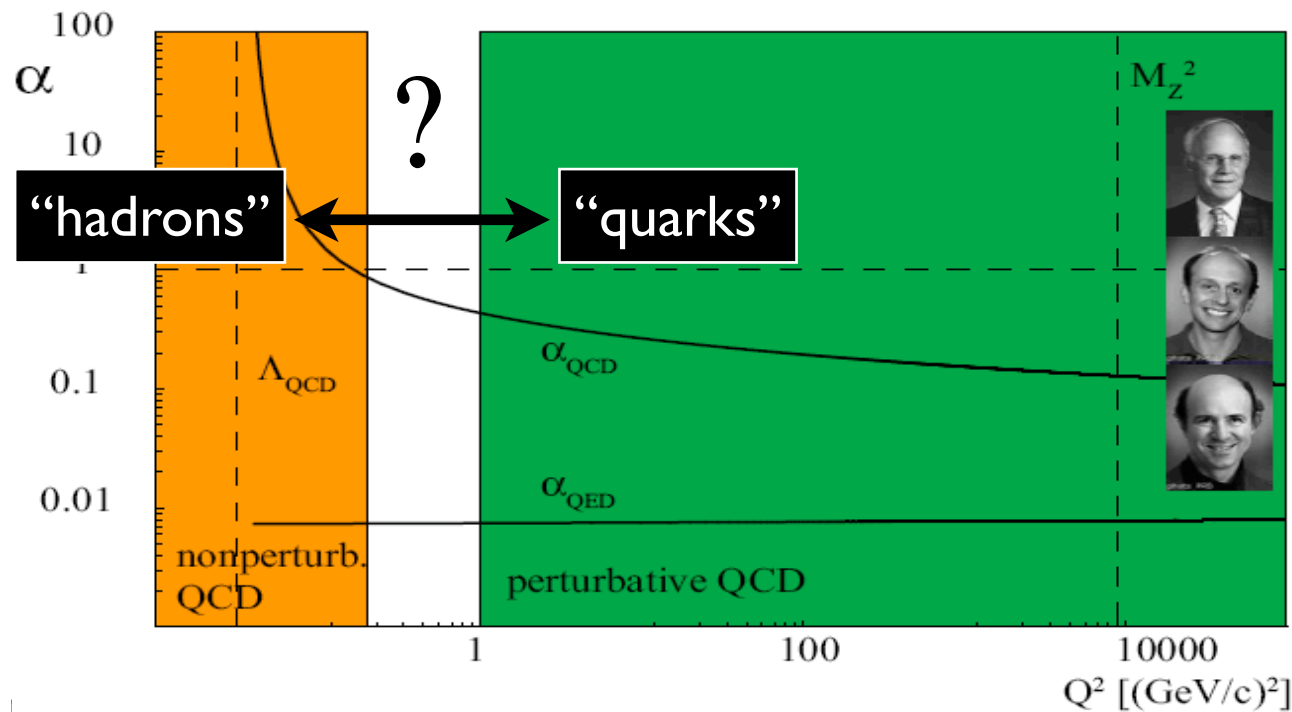
- Duality hypothesis: complementarity between *quark* and *hadron* descriptions of observables

$$\sum_{hadrons} = \sum_{quarks}$$

→ can use either set of *complete* basis states to describe physical phenomena



low
energy
long
distance



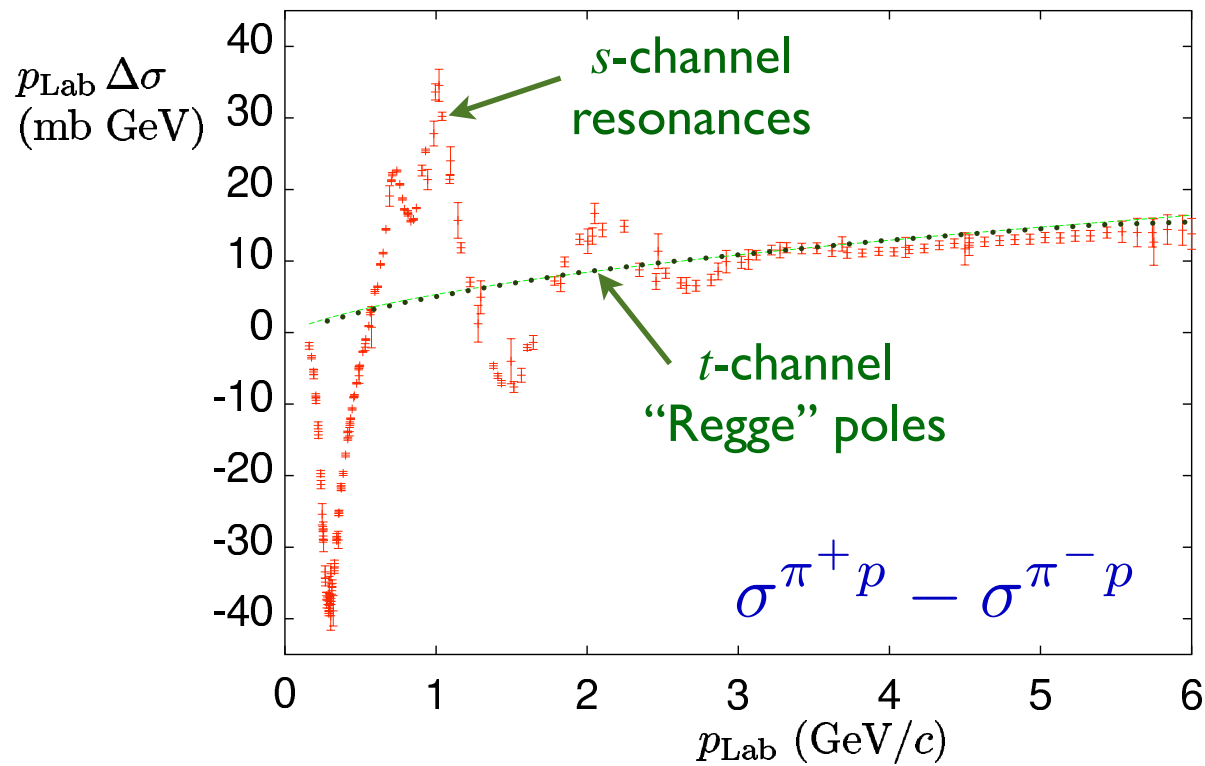
high
energy
short
distance

- In practice, at finite energy typically have access only to *limited* set of basis states

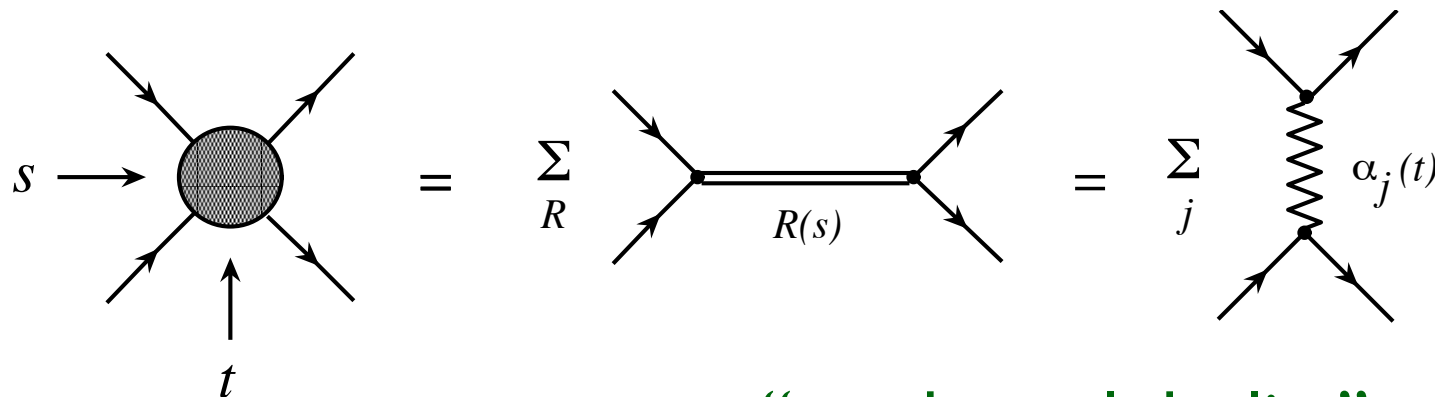


- In practice, at finite energy typically have access only to *limited* set of basis states
- Question is not *why* duality exists, but *how* it arises where it exists, and how we can *make use of it*

Duality in hadron-hadron scattering



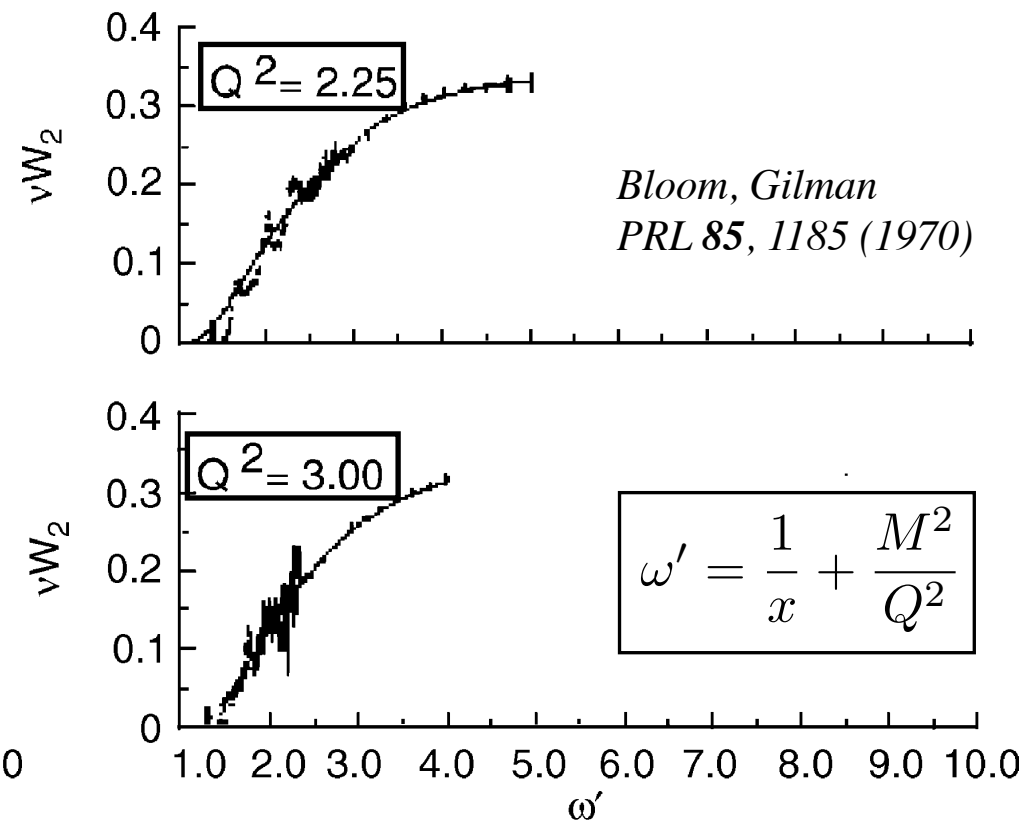
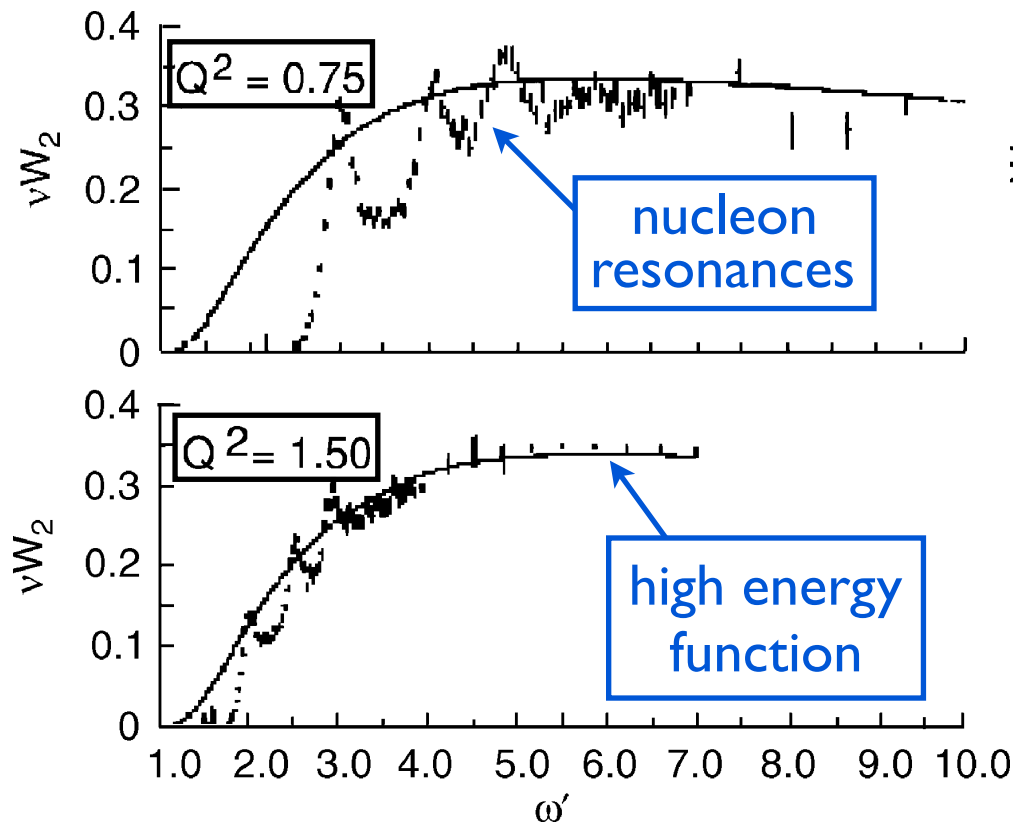
Igi (1962)
Dolen, Horn, Schmidt (1968)



" s - t channel duality"

Duality in electron-nucleon scattering

“Bloom-Gilman duality”



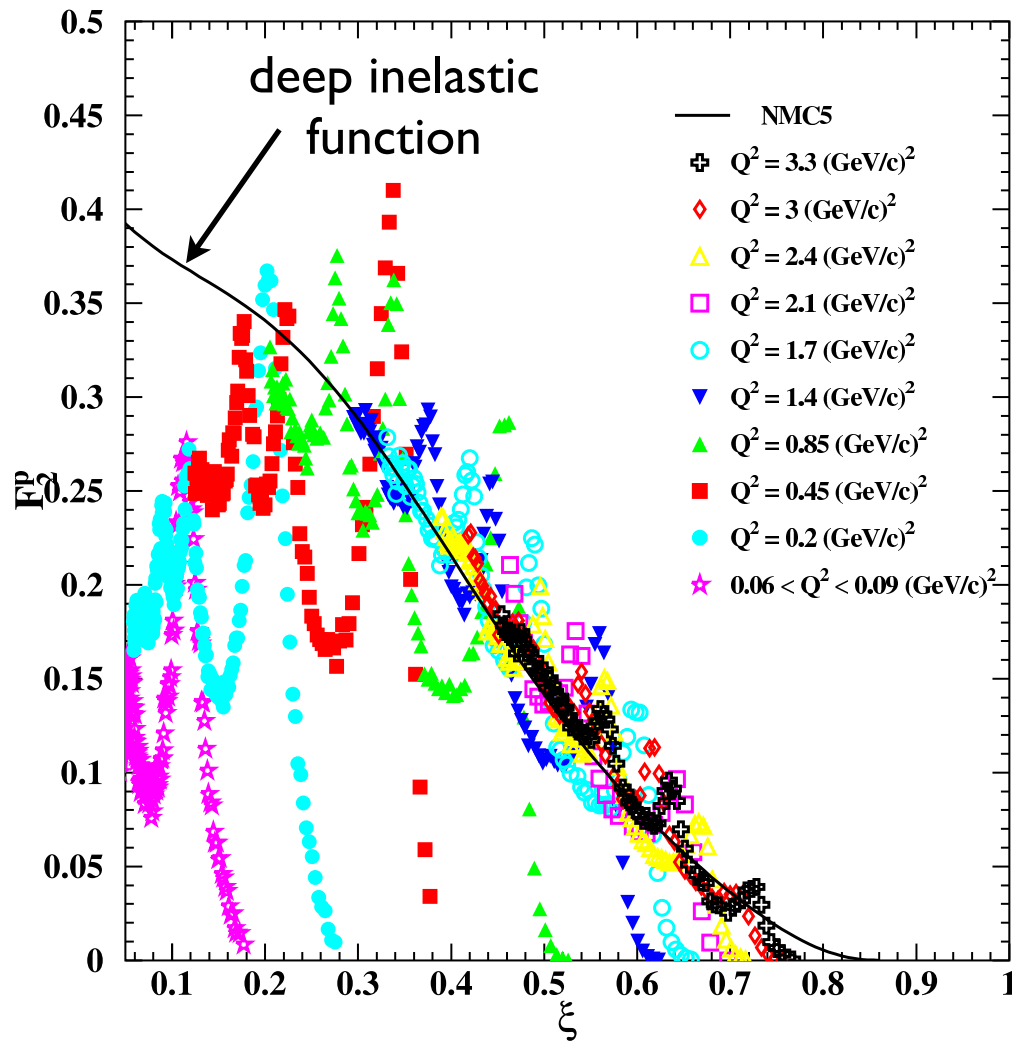
$$\frac{2M}{Q^2} \int_0^{\nu_m} d\nu \nu W_2(\nu, Q^2) = \int_1^{\omega'_m} d\omega' \nu W_2(\omega')$$

“hadrons”

“quarks”

finite-energy
sum rules

Duality in electron-nucleon scattering



Niculescu et al., PRL 85, 1182 (2000)

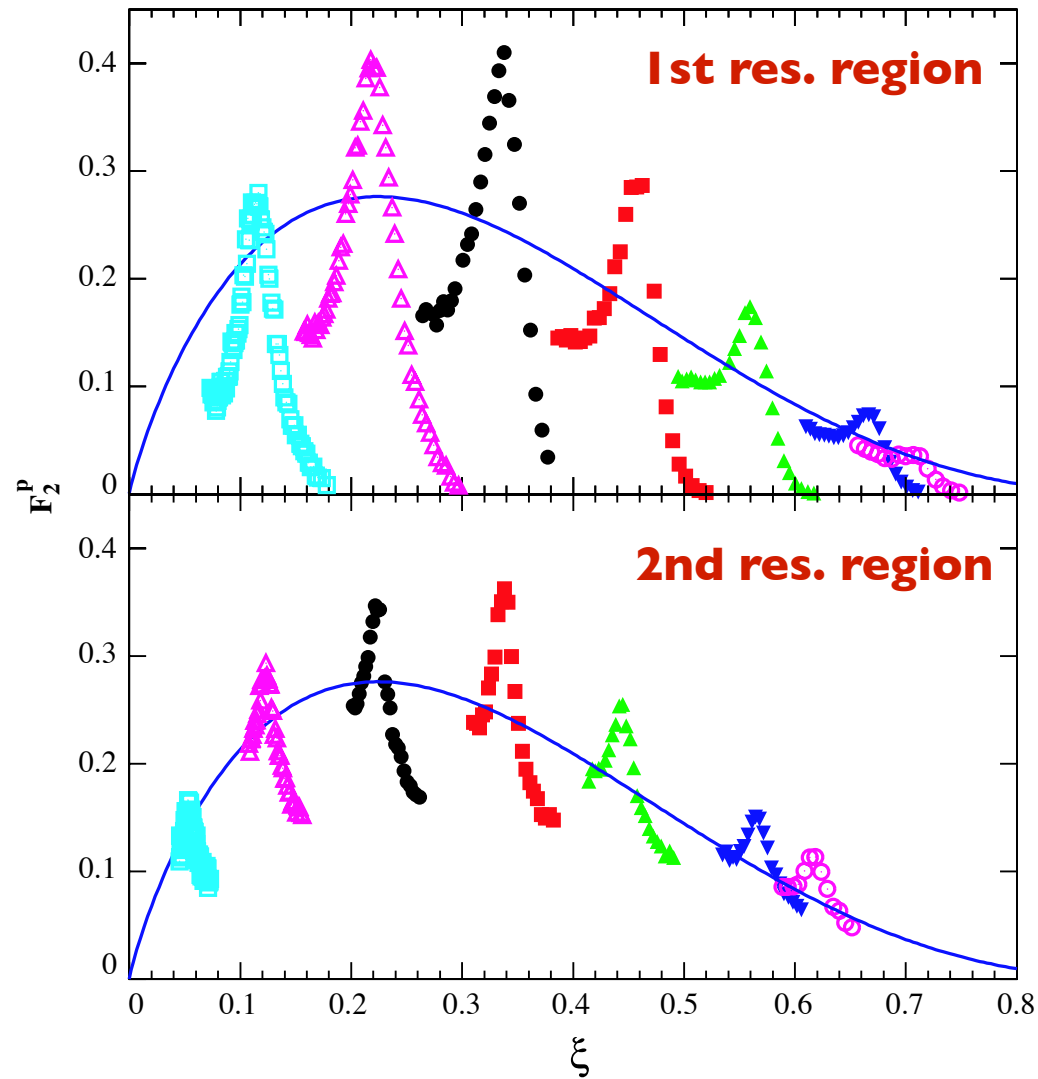
WM, Ent, Keppel, PRep. 406, 127 (2005)

average over
(strongly Q^2 dependent)
resonances
 $\approx Q^2$ independent
scaling function

“Nachtmann” scaling variable

$$\xi = \frac{2x}{1 + \sqrt{1 + 4M^2 x^2 / Q^2}}$$

Duality in electron-nucleon scattering

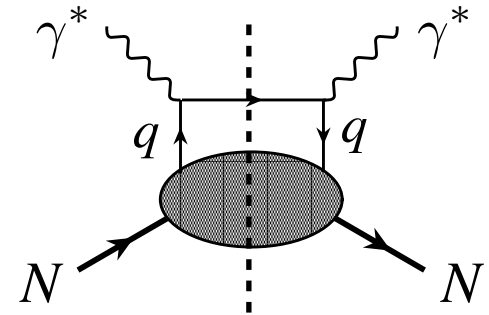


→ also exists *locally* in individual resonance regions

Duality in electron-nucleon scattering

- In *deep-inelastic* region ($W \gtrsim 2 \text{ GeV}$, $Q^2 \gtrsim 1 \text{ GeV}^2$) structure function given by parton distributions

$$F_2(x, Q^2) = x \sum_q e_q^2 q(x, Q^2)$$



- In *resonance* region ($W \lesssim 2 \text{ GeV}$), or at low Q^2 ($Q^2 \lesssim 1 \text{ GeV}^2$) can no longer resolve individual quark structure
- Resonance and DIS regions intimately connected
 - resonances an *integral* part of scaling structure function
 - e.g.* in large- N_c limit, spectrum of zero-width resonances is “maximally dual” to quark-level (smooth) structure function

How to build up a scaling structure function from $\gamma^* NN^*$ transitions?

■ Earliest attempts predate QCD

→ e.g. harmonic oscillator spectrum $M_n^2 = (n + 1)\Lambda^2$
including states with spin = 1/2, ..., $n+1/2$

(n even: $I = 1/2$, n odd: $I = 3/2$)

Domokos et al., PRD 3, 1184 (1971)

→ at large Q^2 magnetic coupling dominates

$$G_n(Q^2) = \frac{\mu_n}{(1 + Q^2 r^2 / M_n^2)^2} \quad r^2 \approx 1.41$$

→ in Bjorken limit, $\sum_n \longrightarrow \int dz$, $z \equiv M_n^2 / Q^2$

$$F_2 \sim (\omega' - 1)^{1/2} (\mu_{1/2}^2 + \mu_{3/2}^2) \int_0^\infty dz \frac{z^{3/2} (1 + r^2/z)^{-4}}{z + 1 - \omega' + \Gamma_0^2 z^2}$$

→ scaling function of $\omega' = \omega + M^2 / Q^2$ ($\omega = 1/x$)

How to build up a scaling structure function from $\gamma^* NN^*$ transitions?

■ Earliest attempts predate QCD

→ e.g. harmonic oscillator spectrum $M_n^2 = (n + 1)\Lambda^2$
including states with spin = $1/2, \dots, n+1/2$

(n even: $I = 1/2$, n odd: $I = 3/2$)

Domokos et al., PRD 3, 1184 (1971)

→ in $\Gamma_n \rightarrow 0$ limit

$$F_2 \sim (\mu_{1/2}^2 + \mu_{3/2}^2) \frac{(\omega' - 1)^3}{(\omega' - 1 + r^2)^4}$$

cf. Drell-Yan-West relation

$$G(Q^2) \sim \left(\frac{1}{Q^2} \right)^m \iff F_2(x) \sim (1 - x)^{2m-1}$$

→ similar behavior found in many models

Einhorn, PRD 14, 3451 (1976) ('t Hooft model)

Greenberg, PRD 47, 331 (1993) (NR scalar quarks in HO potential)

Pace, Salme, Lev, PRC 57, 2655 (1995) (relativistic HO with spin)

Isgur et al., PRD 64, 054005 (2001) (transition to scaling)

....

How to build up a scaling structure function from $\gamma^* NN^*$ transitions?

■ More recent phenomenological analyses at finite Q^2

→ additional constraints from threshold behavior at $q \rightarrow 0$
and asymptotic behavior at $Q^2 \rightarrow \infty$

*Davidovsky, Struminsky,
Phys. Atom. Nucl. **66**, 1328 (2003)*

$$\left(1 + \frac{\nu^2}{Q^2}\right) F_2^R = M\nu \left[|G_+^R|^2 + 2|G_0^R|^2 + |G_-^R|^2 \right] \delta(W^2 - M_R^2)$$

→ 21 isospin-1/2 & 3/2 resonances (with mass < 2 GeV)

$$|G_{\pm}^R(Q^2)|^2 = |G_{\pm}^R(0)|^2 \left(\frac{|\vec{q}|}{|\vec{q}|_0} \frac{\Lambda'^2}{Q^2 + \Lambda'^2} \right)^{\gamma_1} \left(\frac{\Lambda^2}{Q^2 + \Lambda^2} \right)^{m_{\pm}} \quad m_{+,0,-} = 3, 4, 5$$

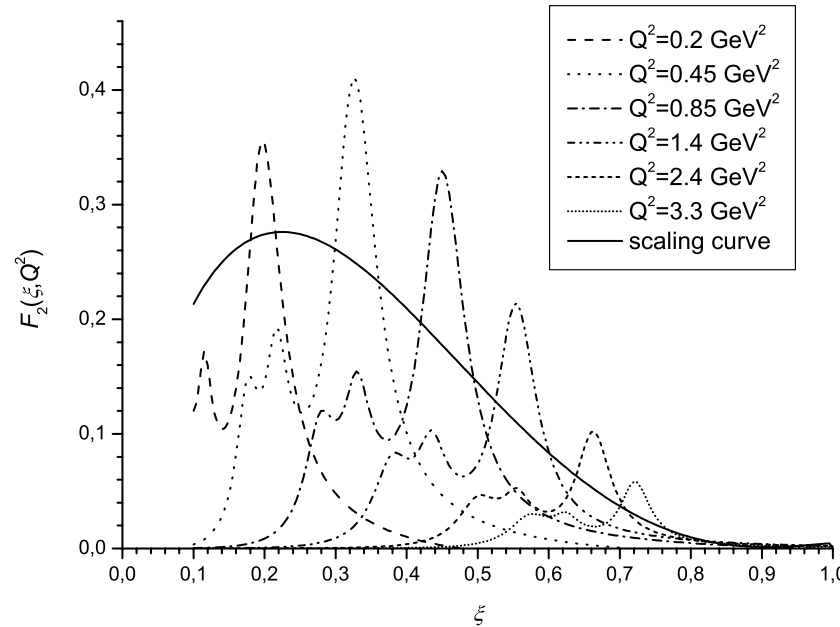
$$|G_0^R(Q^2)|^2 = C^2 \left(\frac{Q^2}{Q^2 + \Lambda''^2} \right)^{2a} \frac{q_0^2}{|\vec{q}|^2} \left(\frac{|\vec{q}|}{|\vec{q}|_0} \frac{\Lambda'^2}{Q^2 + \Lambda'^2} \right)^{\gamma_2} \left(\frac{\Lambda^2}{Q^2 + \Lambda^2} \right)^{m_0}$$

→ in $x \rightarrow 1$ limit,

$$F_2(x) \sim (1 - x)^{m_+}$$

How to build up a scaling structure function from $\gamma^* NN^*$ transitions?

■ More recent phenomenological analyses at finite Q^2



Davidovsky, Struminsky,
Phys. Atom. Nucl. **66**, 1328 (2003)

→ valence-like structure of dual function suggests
“two-component duality”:

- valence (Reggeon exchange) dual to resonances $F_2^{(\text{val})} \sim x^{0.5}$
- sea (Pomeron exchange) dual to background $F_2^{(\text{sea})} \sim x^{-0.08}$

Freund, *PRL* **20**, 235 (1968)
Harari, *PRL* **20**, 1395 (1969)

Duality and QCD

■ Operator product expansion

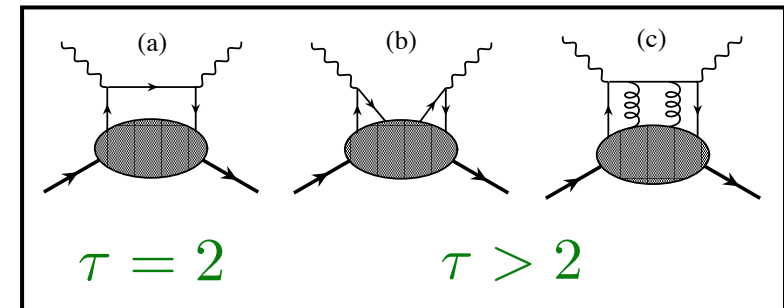
→ expand *moments* of structure functions
in powers of $1/Q^2$

$$x \rightarrow 1 \iff W \rightarrow M$$

$$M_n(Q^2) = \int_0^1 dx \, x^{n-2} F_2(x, Q^2)$$
$$= A_n^{(2)} + \frac{A_n^{(4)}}{Q^2} + \frac{A_n^{(6)}}{Q^4} + \dots$$

matrix elements of operators with
specific “twist” τ

$\tau = \text{dimension} - \text{spin}$



Duality and QCD

■ Operator product expansion

→ expand *moments* of structure functions
in powers of $1/Q^2$

$$\begin{aligned} M_n(Q^2) &= \int_0^1 dx \, x^{n-2} F_2(x, Q^2) \\ &= A_n^{(2)} + \frac{A_n^{(4)}}{Q^2} + \frac{A_n^{(6)}}{Q^4} + \dots \end{aligned}$$

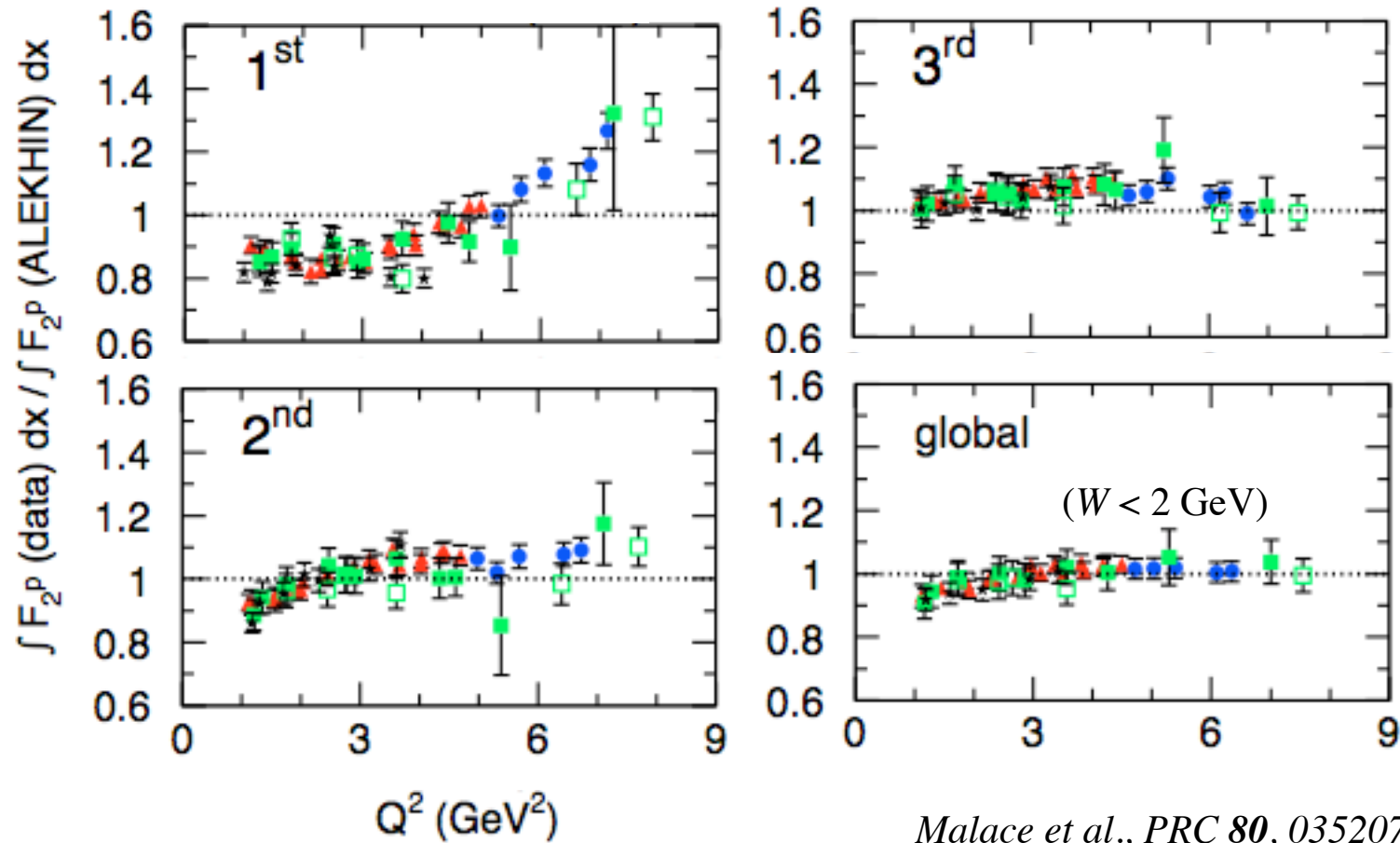
de Rujula, Georgi, Politzer
Ann. Phys. **103**, 315 (1975)

■ If moment \approx independent of Q^2

→ “higher twist” terms $A_n^{(\tau>2)}$ small

■ Duality \longleftrightarrow suppression of higher twists

■ Truncated moments of F_2^p in resonance region



Malace et al., PRC 80, 035207 (2009)
[JLab Hall C]

→ higher twists < 10–15% for $Q^2 > 1$ GeV²

Resonances & twists

- Total “higher twist” is *small* at scales $Q^2 \sim \mathcal{O}(1 \text{ GeV}^2)$
 - On average, nonperturbative interactions between quarks and gluons not dominant (at these scales)
 - nontrivial interference between resonances
-

- Can we understand this dynamically, at quark level?
 - is duality an accident?
- Can we use resonance region data to learn about *leading twist* structure functions (and *vice versa*)?
 - expanded data set has potentially significant implications for global quark distribution studies

- Consider simple quark model with spin-flavor symmetric wave function

low energy

→ *coherent* scattering from quarks $d\sigma \sim \left(\sum_i e_i \right)^2$

high energy

→ *incoherent* scattering from quarks $d\sigma \sim \sum_i e_i^2$

- For duality to work, these must be equal

→ how can square of a sum become sum of squares?

■ Dynamical cancellations

→ *e.g.* for toy model of two quarks bound in a harmonic oscillator potential, structure function given by

$$F(\nu, \mathbf{q}^2) \sim \sum_n |G_{0,n}(\mathbf{q}^2)|^2 \delta(E_n - E_0 - \nu)$$

→ charge operator $\sum_i e_i \exp(i\mathbf{q} \cdot \mathbf{r}_i)$ excites
even partial waves with strength $\propto (e_1 + e_2)^2$
odd partial waves with strength $\propto (e_1 - e_2)^2$

→ resulting structure function

$$F(\nu, \mathbf{q}^2) \sim \sum_n \{ (e_1 + e_2)^2 G_{0,2n}^2 + (e_1 - e_2)^2 G_{0,2n+1}^2 \}$$

→ if states degenerate, *cross terms* ($\sim e_1 e_2$) *cancel* when averaged over nearby *even and odd parity* states

■ Dynamical cancellations

→ duality is realized by summing over at least one complete set of even and odd parity resonances

Close, Isgur, PLB 509, 81 (2001)

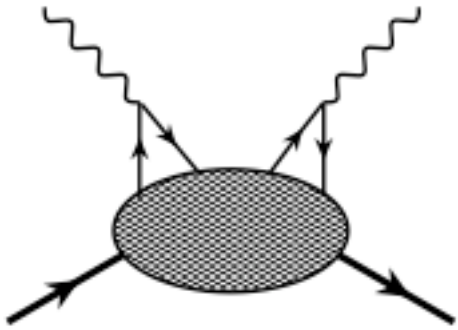
→ in NR Quark Model, even & odd parity states generalize to **56** ($L=0$) and **70** ($L=1$) multiplets of spin-flavor SU(6)

representation	${}^2\mathbf{8}[\mathbf{56}^+]$	${}^4\mathbf{10}[\mathbf{56}^+]$	${}^2\mathbf{8}[\mathbf{70}^-]$	${}^4\mathbf{8}[\mathbf{70}^-]$	${}^2\mathbf{10}[\mathbf{70}^-]$	Total
F_1^p	$9\rho^2$	$8\lambda^2$	$9\rho^2$	0	λ^2	$18\rho^2 + 9\lambda^2$
F_1^n	$(3\rho + \lambda)^2/4$	$8\lambda^2$	$(3\rho - \lambda)^2/4$	$4\lambda^2$	λ^2	$(9\rho^2 + 27\lambda^2)/2$

$\lambda(\rho) =$ (anti) symmetric component of ground state wave function

Close, WM, PRC 68, 035210 (2003)
PRC 79, 055202 (2009)

■ Accidental cancellations of charges?



cat's ears diagram (4-fermion higher twist $\sim 1/Q^2$)

$$\propto \sum_{i \neq j} e_i e_j \sim \left(\sum_i e_i \right)^2 - \sum_i e_i^2$$

\uparrow coherent \uparrow incoherent

proton HT $\sim 1 - \left(2 \times \frac{4}{9} + \frac{1}{9} \right) = 0 !$

neutron HT $\sim 0 - \left(\frac{4}{9} + 2 \times \frac{1}{9} \right) \neq 0$

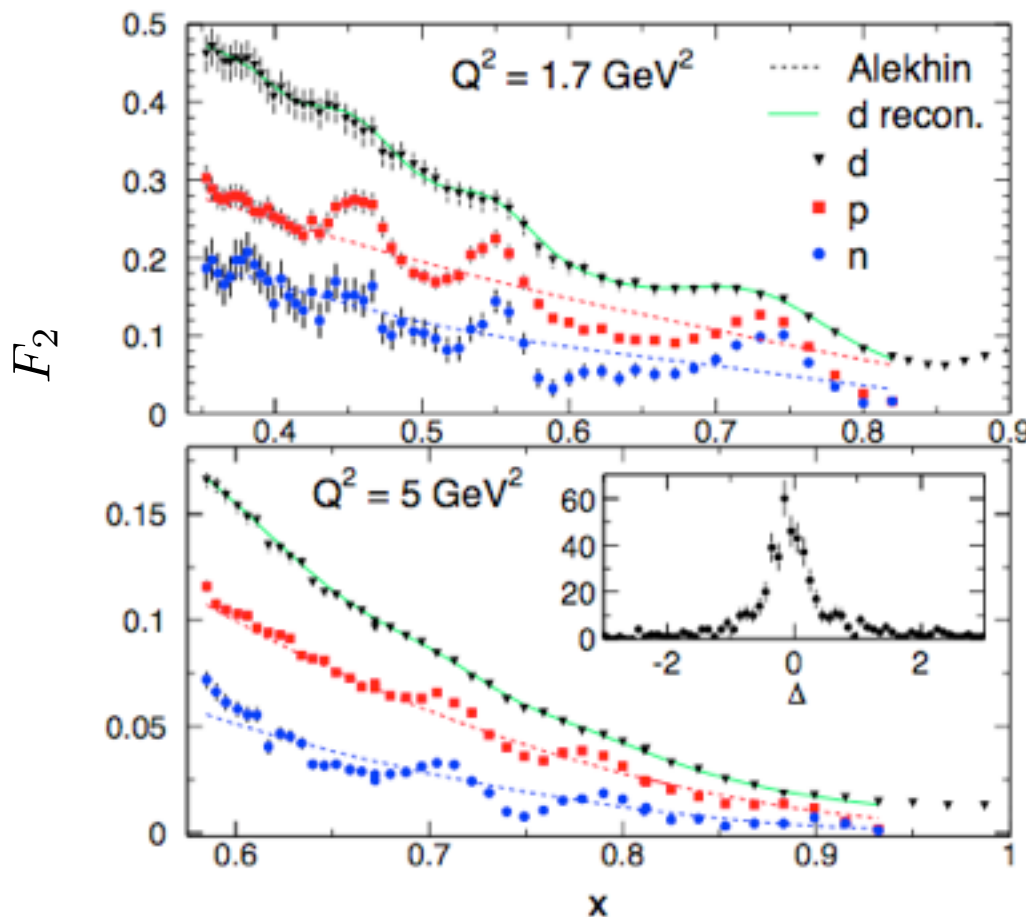
S. Brodsky (2000)

→ duality in proton a *coincidence!*

→ should not hold for neutron !!

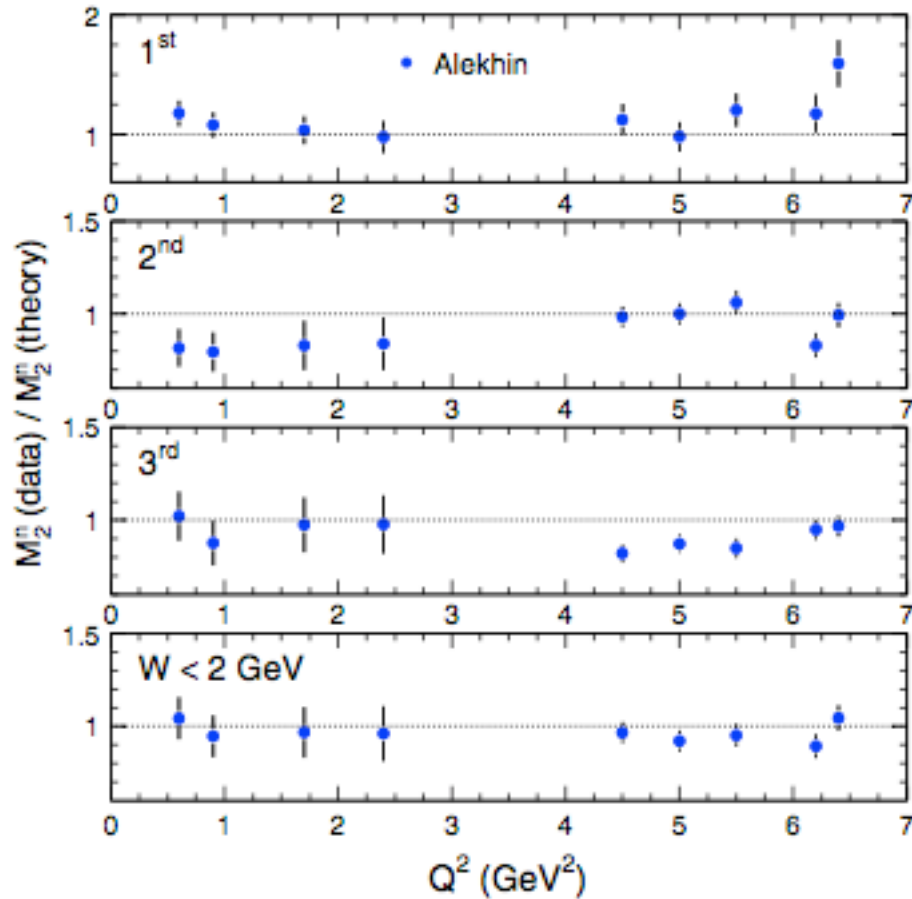
Neutron: the smoking gun

- Duality in *neutron* more difficult to test because of absence of free neutron targets
- New extraction method (using iterative procedure for solving integral convolution equations) has allowed first determination of F_2^n in resonance region & test of neutron duality



Malace, Kahn, WM, Keppel
PRL **104**, 102001 (2010)

Neutron: the smoking gun



→ “theory”: fit to $W > 2 \text{ GeV}$ data

Alekhin et al., 0908.2762 [hep-ph]

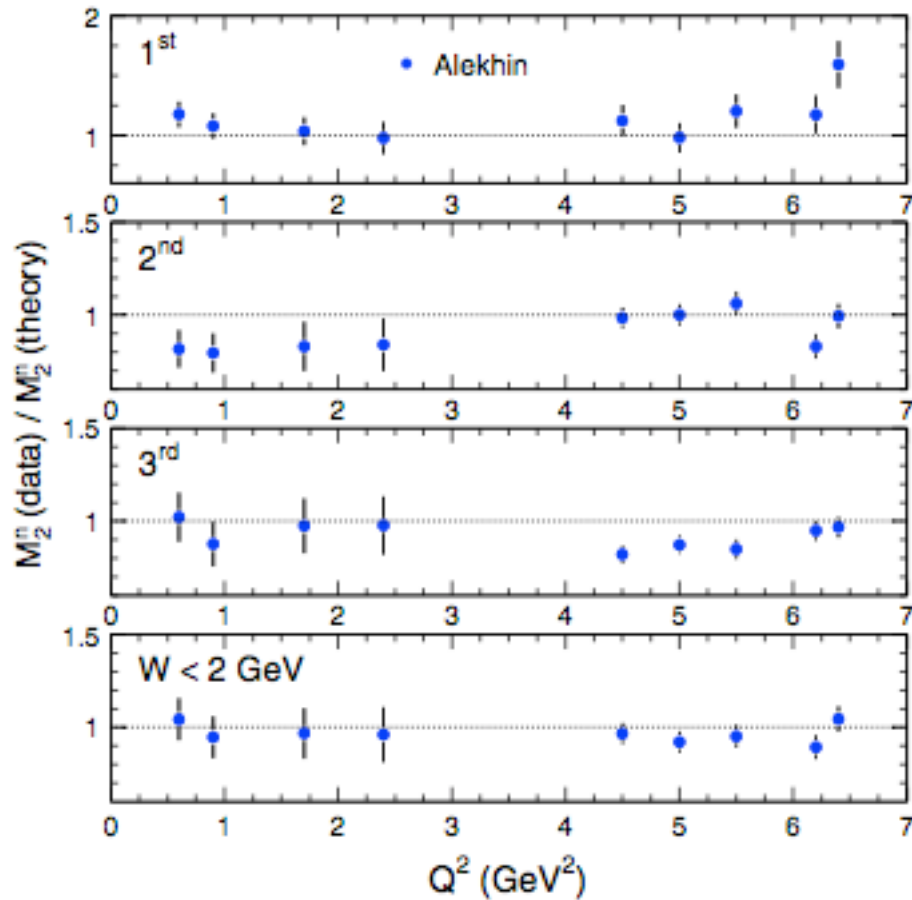
→ *locally*, violations of duality in resonance regions $< 15\text{--}20\%$ (largest in Δ region)

→ *globally*, violations $< 10\%$

Malace, Kahn, WM, Keppel
PRL 104, 102001 (2010)

➡ duality is not accidental, but a general feature of resonance–scaling transition!

Neutron: the smoking gun



→ “theory”: fit to $W > 2$ GeV data

Alekhin et al., 0908.2762 [hep-ph]

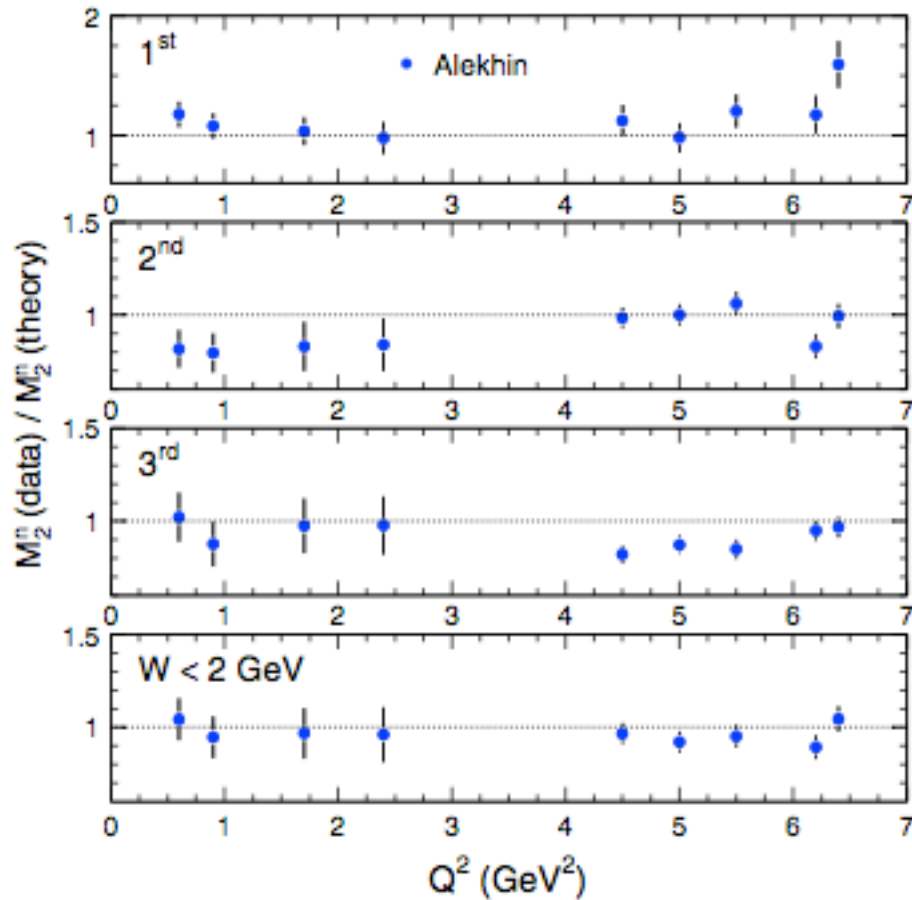
→ *locally*, violations of duality in
resonance regions $< 15\text{--}20\%$
(largest in Δ region)

→ *globally*, violations $< 10\%$

Malace, Kahn, WM, Keppel
PRL 104, 102001 (2010)

→ analysis using recent (model-independent)
BoNuS data in progress

Neutron: the smoking gun



→ “theory”: fit to $W > 2$ GeV data

Alekhin et al., 0908.2762 [hep-ph]

→ *locally*, violations of duality in resonance regions < 15–20% (largest in Δ region)

→ *globally*, violations < 10%

Malace, Kahn, WM, Keppel
PRL 104, 102001 (2010)

→ use resonance region data to learn about *leading twist* structure functions?

CTEQ-JLab (CJ) global PDF analysis *

- New global NLO analysis of expanded set of p and d data (DIS, pp , pd) including large- x , low- Q^2 region

*Owens, Accardi WM
arXiv:1212.1702 (PRD, in print)*

- Systematically study effects of Q^2 & W cuts
→ down to $Q \sim m_c$ and $W \sim 1.7$ GeV

cut0: $Q^2 > 4 \text{ GeV}^2$, $W^2 > 12.25 \text{ GeV}^2$

cut1: $Q^2 > 3 \text{ GeV}^2$, $W^2 > 8 \text{ GeV}^2$

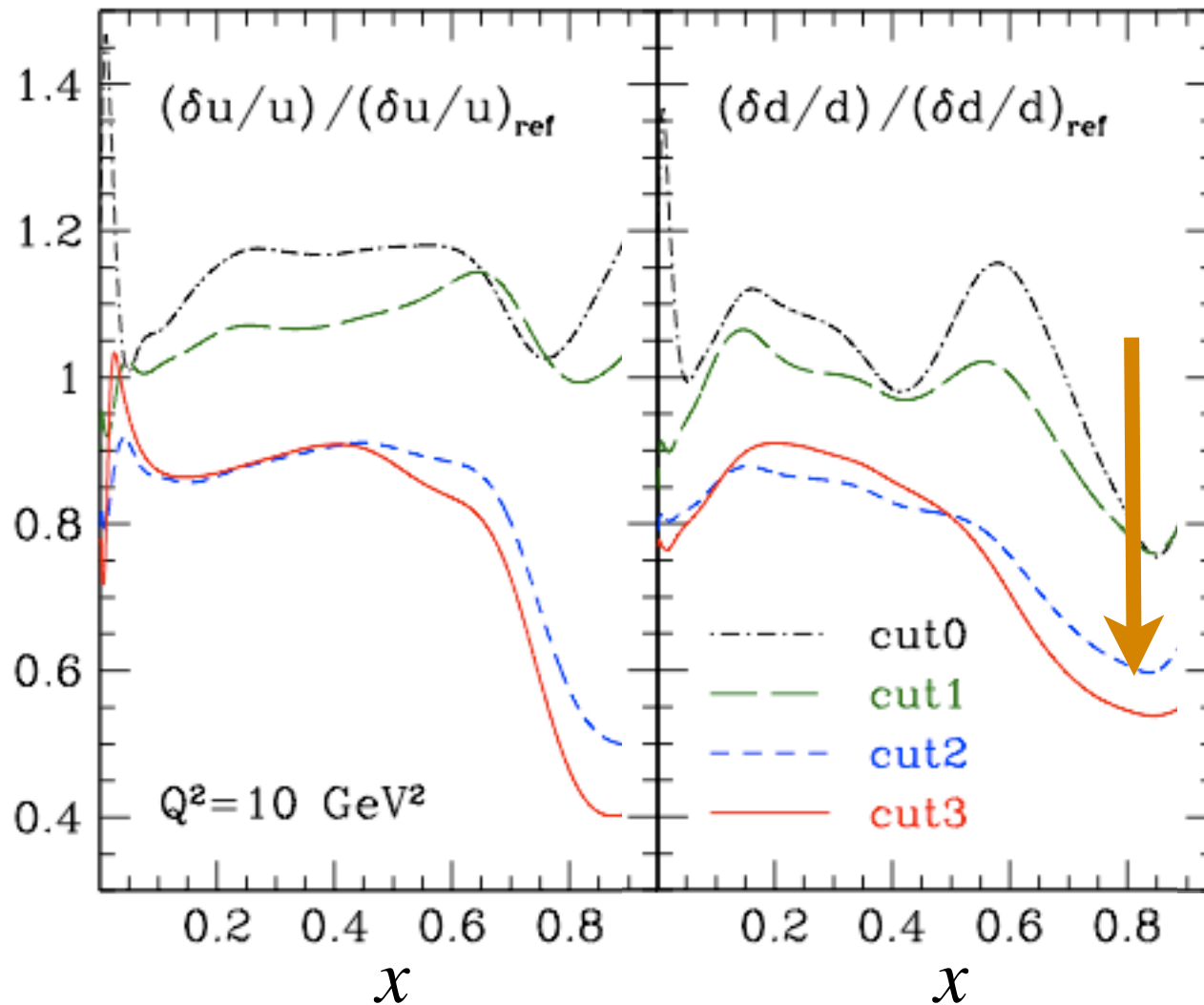
cut2: $Q^2 > 2 \text{ GeV}^2$, $W^2 > 4 \text{ GeV}^2$

cut3: $Q^2 > m_c^2$, $W^2 > 3 \text{ GeV}^2$

factor 2 increase
in DIS data from
cut0 → cut3

* CJ collaboration: <http://www.jlab.org/CJ>

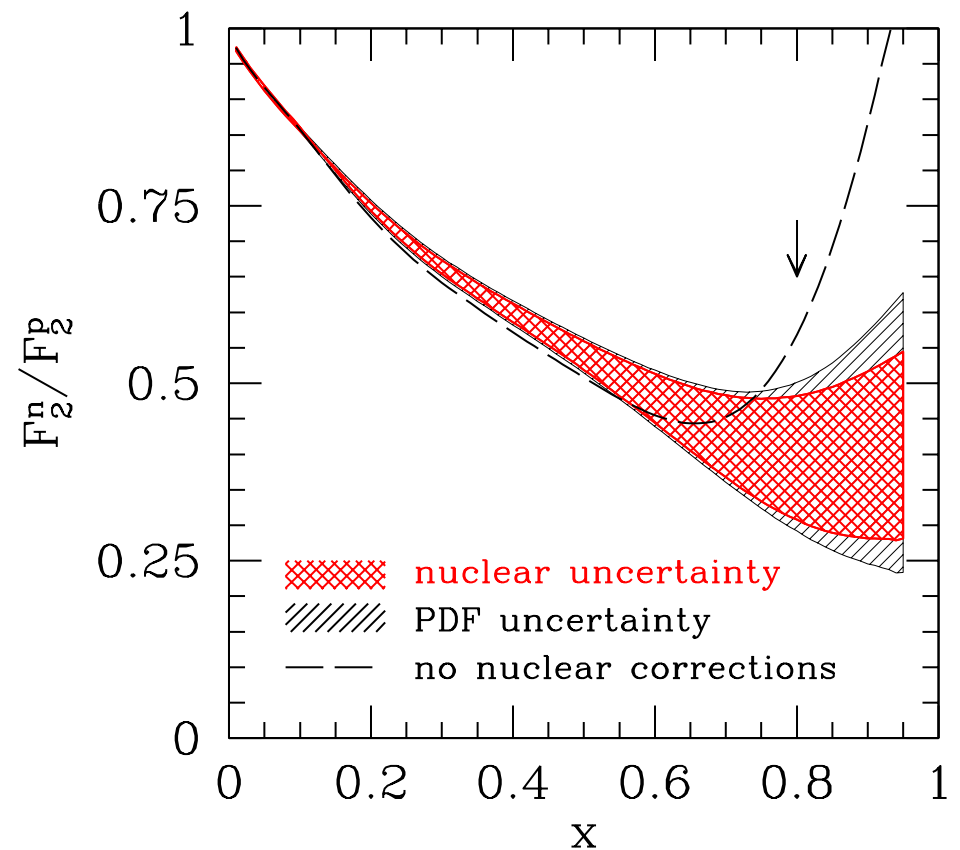
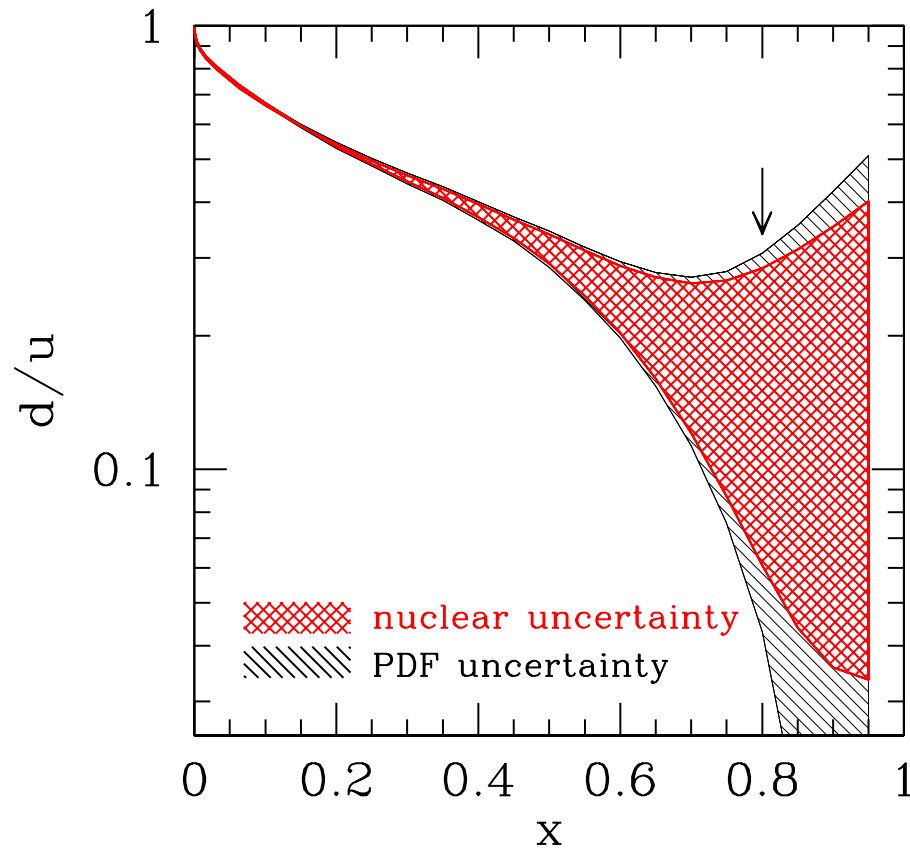
- Larger database with weaker cuts leads to significantly reduced errors, especially at large x



Accardi et al.
PRD 81, 034016 (2010)

→ up to 40–60% error reduction when cuts extended into resonance region

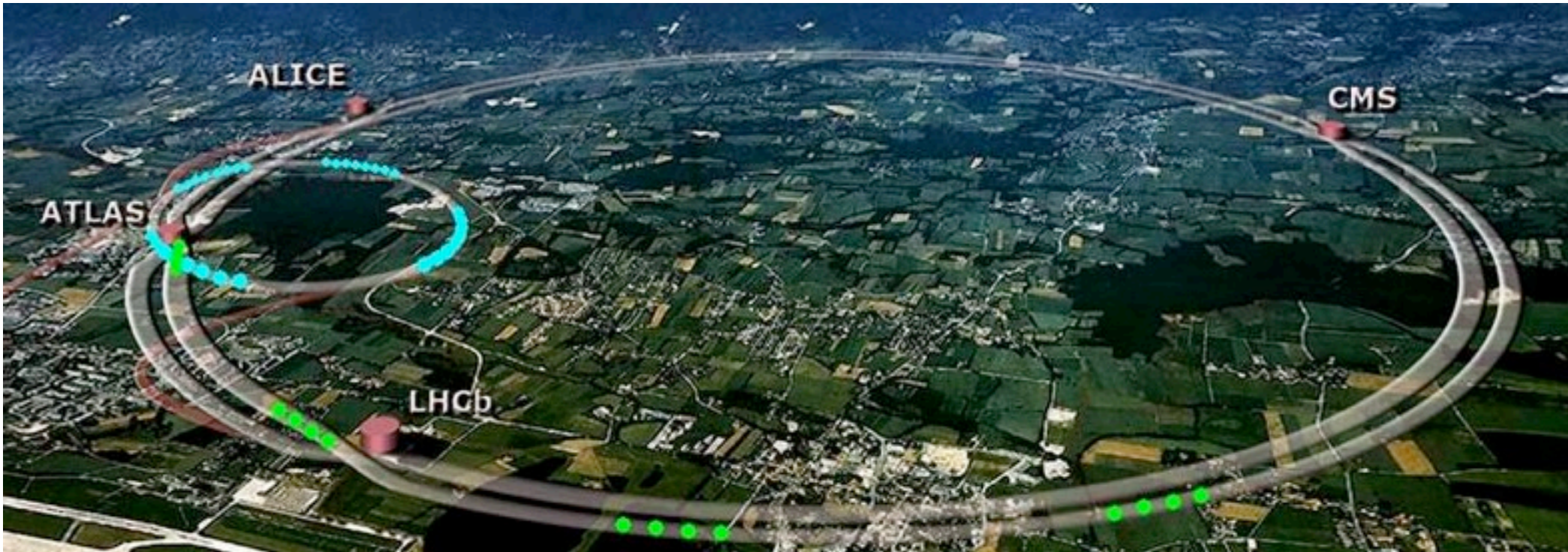
- Vital for large- x analysis, which currently suffers from large uncertainties (mostly due to nuclear corrections)



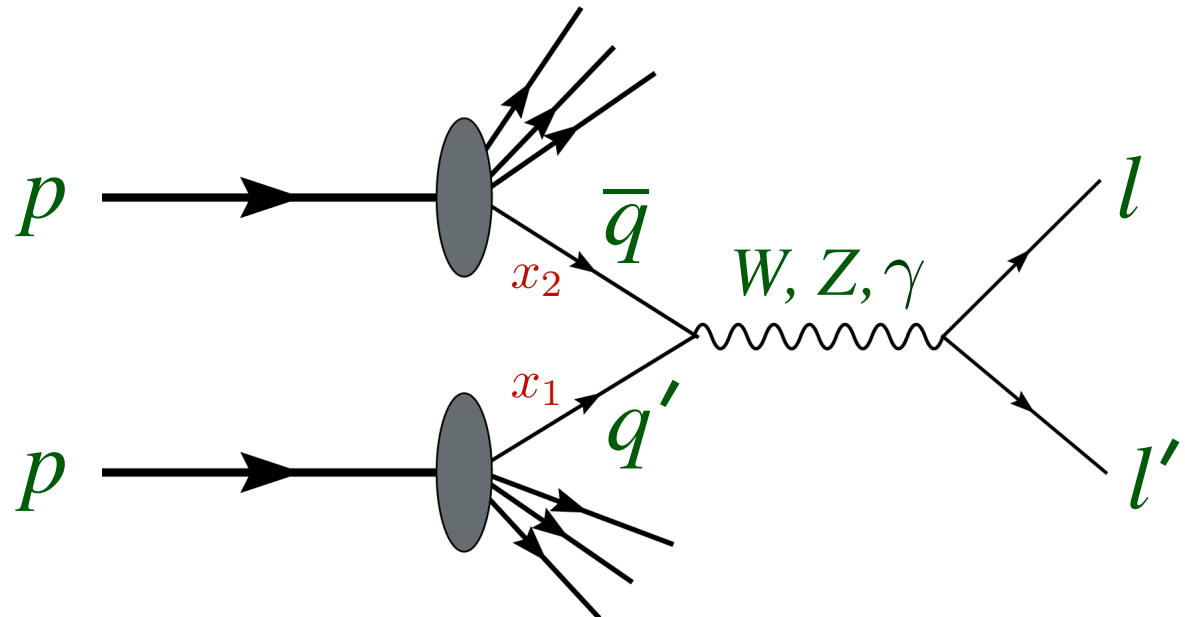
Accardi et al., PRD 84, 014008 (2011)

→ uncertainty in d feeds into larger uncertainty in g at high x (important for LHC physics!)

Large Hadron Collider (CERN)



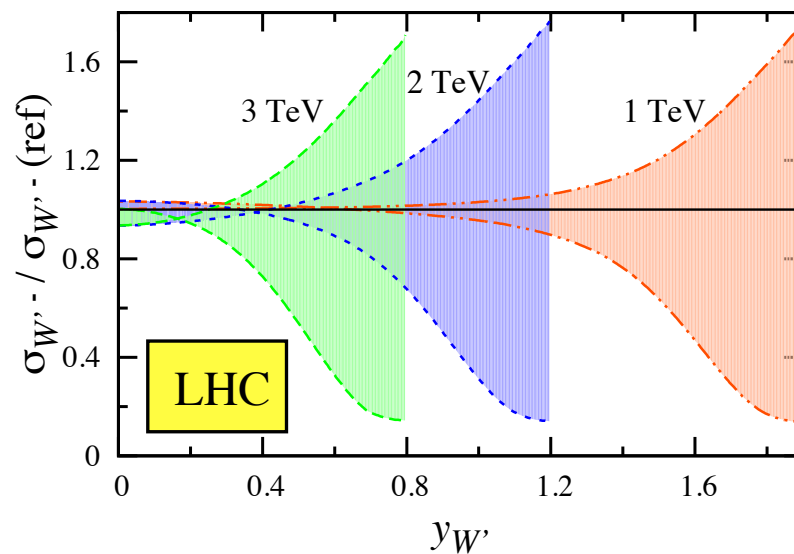
pp collisions
at $\sqrt{s} = 7$ TeV



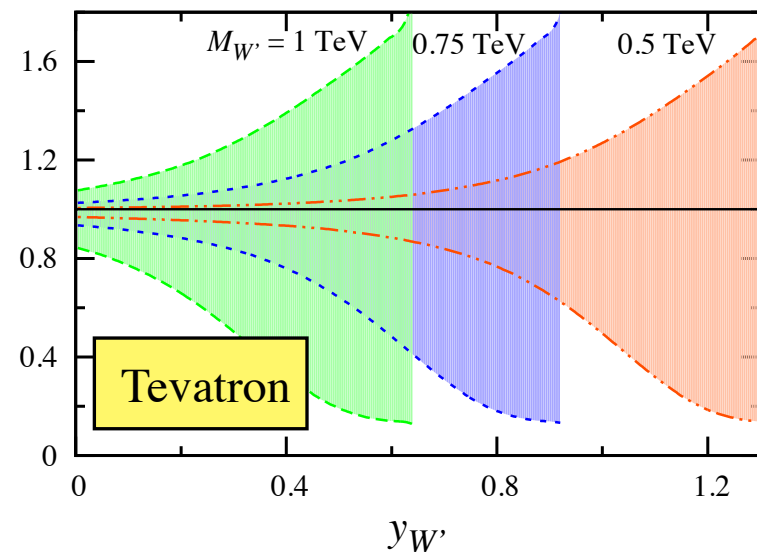
Heavy Z' , W' boson production

- Observation of new physics signals requires accurate determination of QCD backgrounds — depend on PDFs!
(since $x_{1,2} \sim M_{Z',W'}$, large- x uncertainties scale with mass!)

- for W'^- production



→ dominated by $d * \bar{u}$



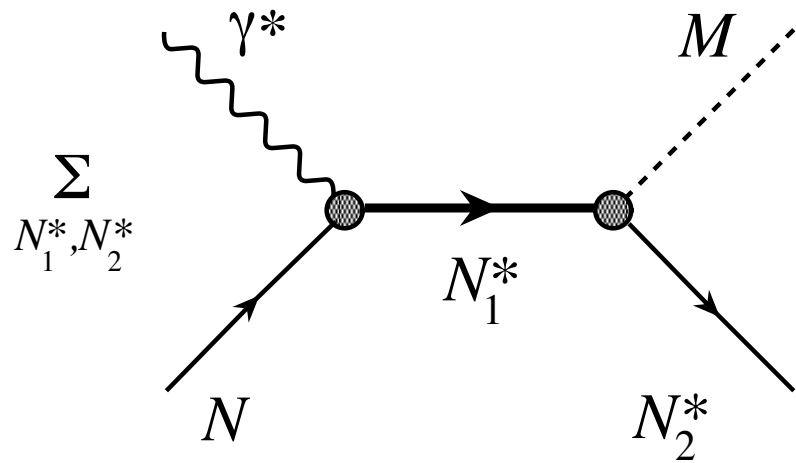
→ dominated by $d * u + u * d$

> 100% uncertainties at large y !

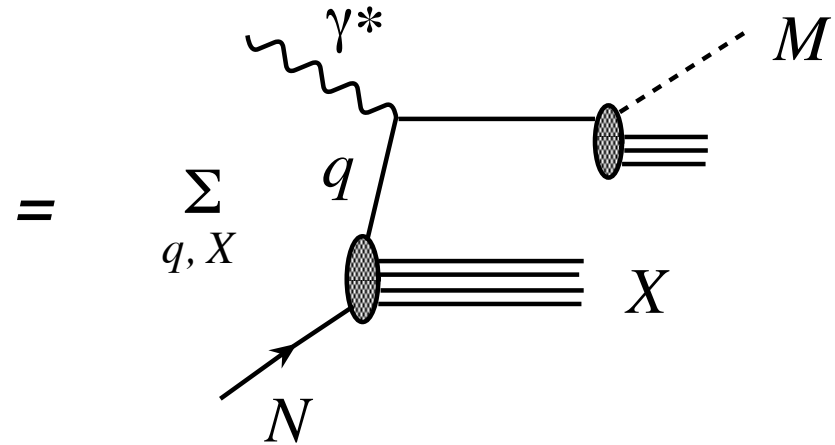
Brady et al., JHEP 1206, 019 (2012)

Duality in (semi-inclusive) meson production

- Extend duality to less inclusive processes, such as meson electroproduction



s -channel resonance
excitation and decay



parton level scattering
and fragmentation

$$\sum_{N_2^*} \left| \sum_{N_1^*} F_{\gamma N \rightarrow N_1^*}(Q^2, M_1^*) \mathcal{D}_{N_1^* \rightarrow N_2^* M}(M_1^*, M_2^*) \right|^2 = \sum_q e_q^2 q(x, Q^2) D_q^M(z, Q^2)$$

Duality in exclusive reactions

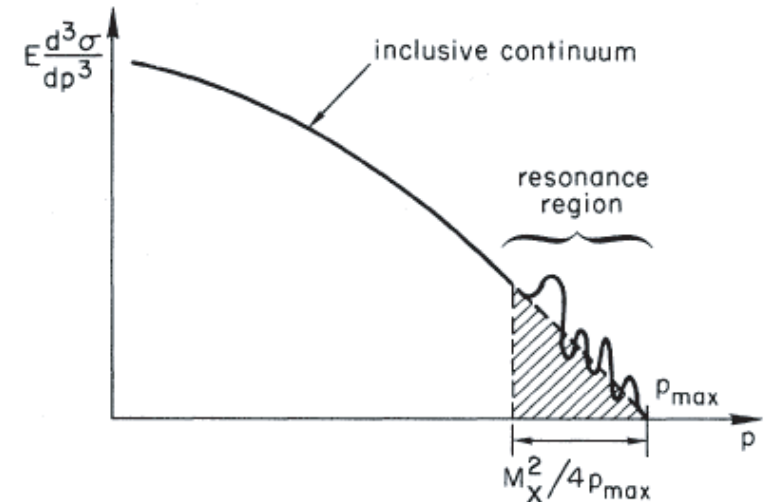
■ Exclusive–inclusive correspondence principle:

→ continuity of dynamics from one (known) region to another (poorly known)

$$\int_{p_{\max} - M_X^2/4p_{\max}}^{p_{\max}} dp \left. E \frac{d^3\sigma}{dp^3} \right|_{\text{incl}} \sim \sum_{\text{res}} \left. E \frac{d\sigma}{dp_T^2} \right|_{\text{excl}}$$

\nearrow
 $\gamma^* N \rightarrow M X$

\nearrow
 $\gamma^* N \rightarrow M N^*$



→ resonance contribution to $d\sigma$ should be comparable to the continuum contribution extrapolated from high energy

$$\frac{E}{\sigma} \frac{d^3\sigma}{dp^3} \equiv f(x, p_T^2, sQ^2) \longrightarrow f(x, p_T^2, sQ^2) \xrightarrow{s \rightarrow \infty} f(x, p_T^2)$$

Bjorken, Kogut, PRD 8, 1341 (1973)

Conclusion

- Confirmation of duality (experimentally & theoretically) suggests origin in dynamical cancelations between resonances
 - explore more realistic descriptions based on phenomenological $\gamma^* NN^*$ form factors
 - incorporate *nonresonant* background in same framework
- Practical application of duality
 - use resonance region data to constrain PDFs at high x
(better knowledge of resonances could be relevant for LHC !)
- Extend quark-hadron duality concept to electroproduction
 - application to semi-inclusive DIS, DVCS / GPDs, ...

Gracias!