



PDFs from Jefferson Lab to the LHC

Wally Melnitchouk



JLab Angular Momentum (JAM) collaboration: <http://www.jlab.org/JAM>



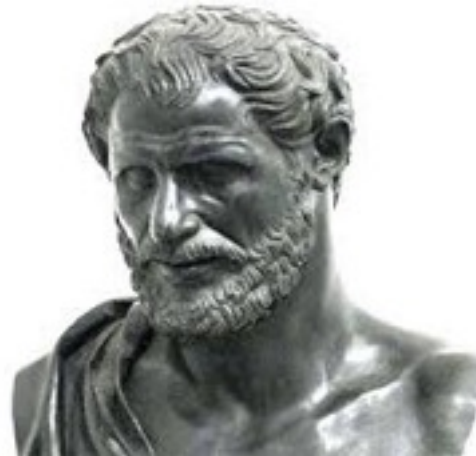
CTEQ-JLab (CJ) collaboration: <http://www.jlab.org/CJ>

☀ Quest to understand nature in terms of fundamental building blocks — “atoms” — is ancient!

→ circa 5th century BC, not far from here...

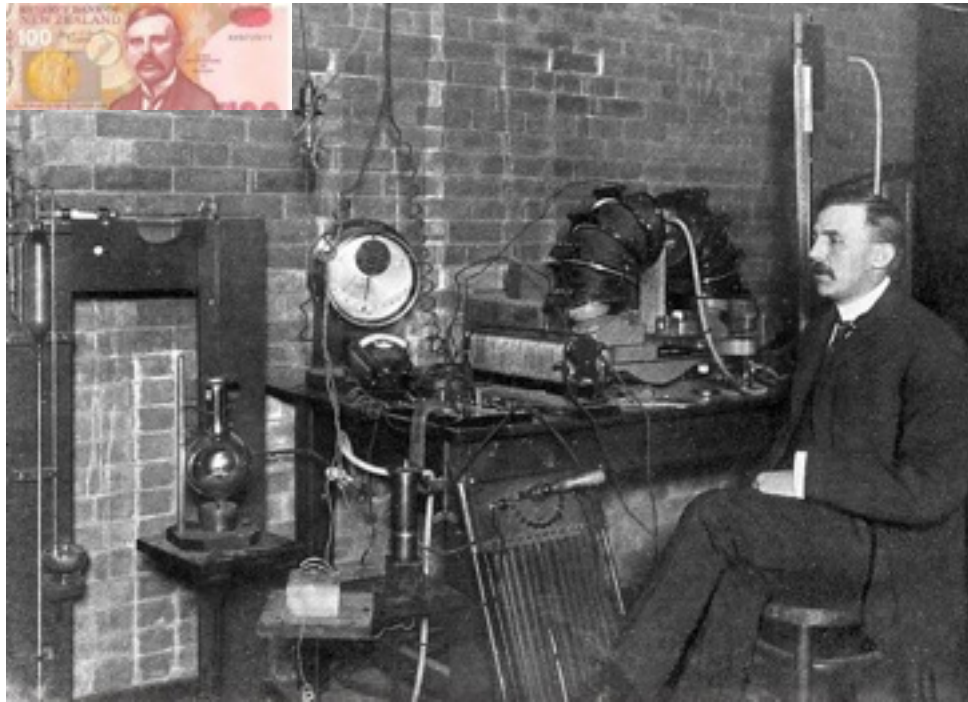


Leucippus



Democritus

☀ 23 centuries later, experimental discovery of atomic nucleus



Rutherford (1911)

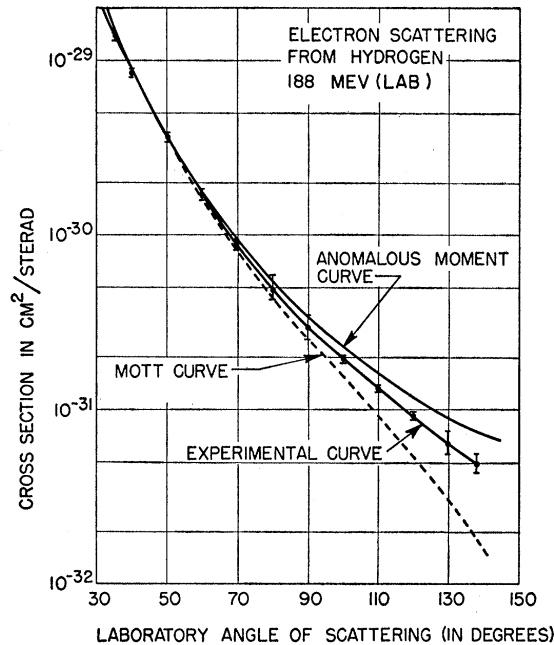
☀ ... followed by discoveries of proton (*Rutherford, 1917*)
and neutron (*Chadwick, 1932*)





Almost 100 years later, what do we know about the nucleon?

→ it has finite size



Hofstadter (1955)



elastic $ep \rightarrow ep$
scattering cross section

$$\frac{d\sigma}{d\Omega} = \left. \frac{d\sigma}{d\Omega} \right|_{\text{point}} \times F^2(Q^2)$$

form factor

→ from slope of form factor $F(Q^2)$ vs. momentum transfer squared Q^2 , obtained r.m.s. charge radius

$$\sim 0.75 \times 10^{-15} \text{ m}$$

(precise value currently under hot debate!)

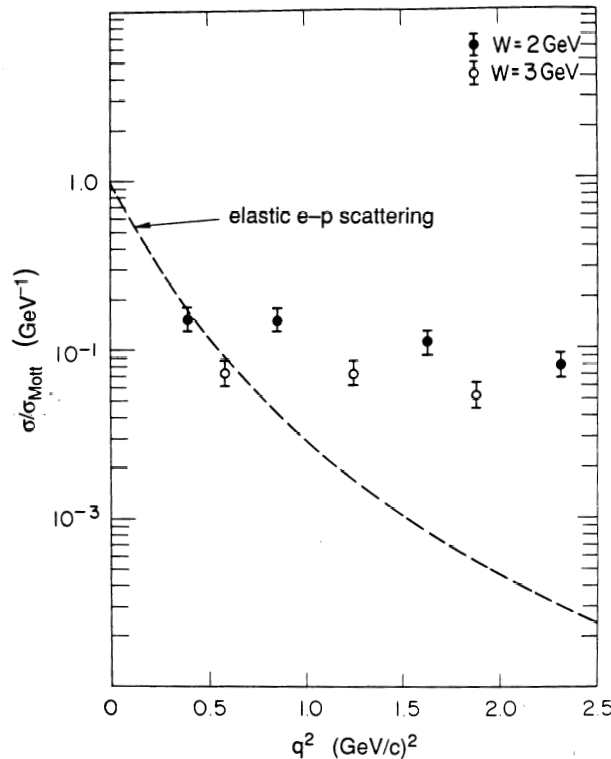


R. Hill talk
Tue. 9:30



Almost 100 years later, what do we know about the nucleon?

→ remarkably, at higher Q^2 inelastic cross section exhibits point-like behavior!



Friedman, Kendall, Taylor (1969)



deep-inelastic $ep \rightarrow eX$ scattering
(DIS) cross section

$$\frac{d^2\sigma}{d\Omega dE'} = \frac{d^2\sigma}{d\Omega dE'} \Big|_{\text{point}} \left(2F_1 \tan^2 \frac{\theta}{2} + F_2 \right)$$

structure functions

→ how can nucleon have finite-size and point-like properties at the same time?

☀ At high energies, scattering is point-like...
but from *constituents* of nucleon

→ “partons” (*Feynman, 1970*) = quarks (*Gell-Mann, Zweig, 1964*)
+ gluons (*Gell-Mann... 1972*)

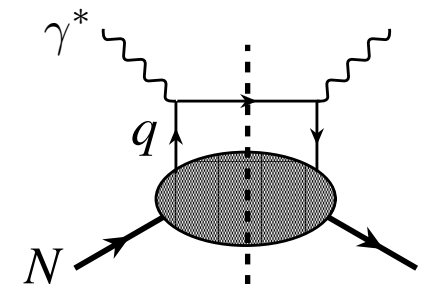


☀ Measurement of structure functions in DIS reveals
how nucleon is made up of quarks & gluons

→ in Feynman’s parton model, structure functions
given by parton distribution functions (PDFs)

$$F_2 = x \sum_q e_q^2 q(x)$$

$q(x)$ = probability distribution to find quark q
in nucleon with momentum fraction x



☀ In QCD, parton distributions are universal functions which are process-independent

→ established formally through factorization theorems
(*e.g.* collinear, TMD, ...)

Collins, Soper, Sterman (“CSS”), 1980s



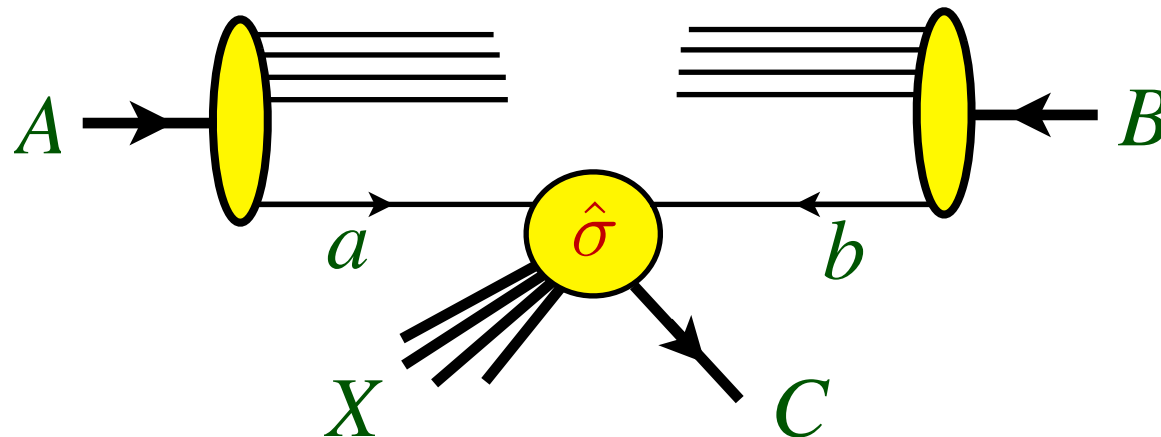
→ allows high-energy cross sections to be factorized into “hard scattering partonic cross sections” (calculated from QCD using perturbation theory), and “soft” matrix elements (parametrized via PDFs)



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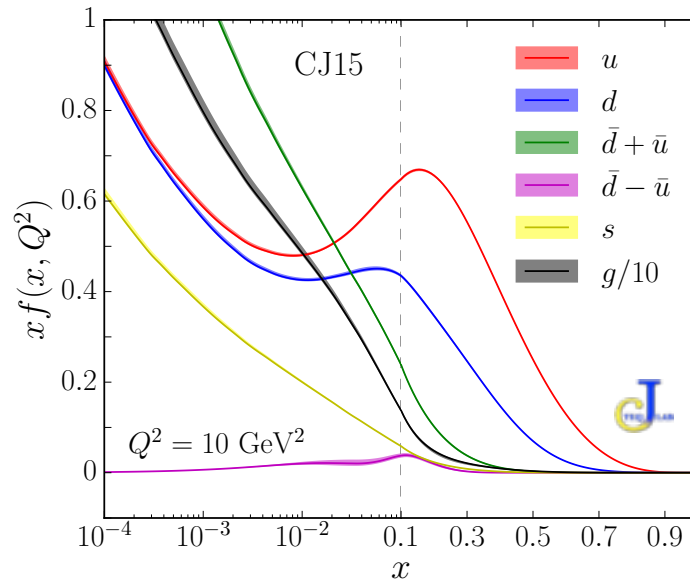
$$\sigma_{AB \rightarrow CX}(p_A, p_B) = \sum_{a,b} \int dx_a dx_b \boxed{f_{a/A}(x_a, \mu)} \boxed{f_{b/B}(x_b, \mu)} \\ \times \hat{\sigma}_{ab \rightarrow CX}(x_a p_A, x_b p_B, Q/\mu)$$



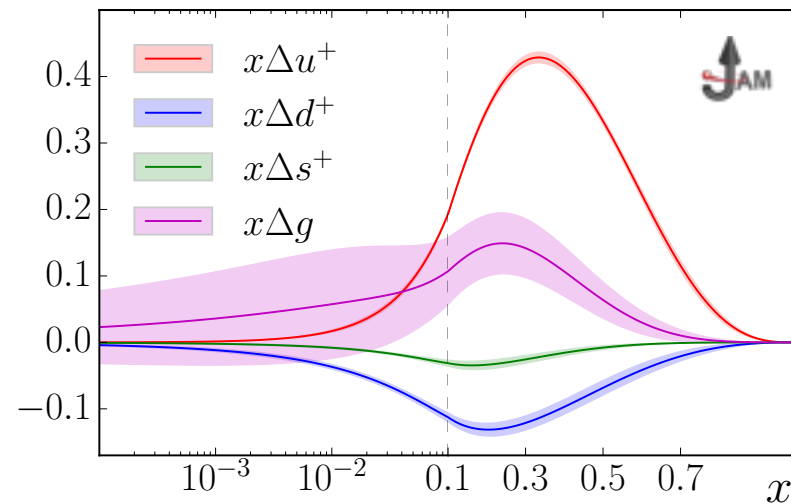
Universality of PDFs allows data from many different processes (DIS, SIDIS, weak boson/jet production in pp , Drell-Yan, ...) to be analyzed simultaneously

→ global QCD analyses of spin-averaged ($f = f^\uparrow + f^\downarrow$) and spin-dependent ($\Delta f = f^\uparrow - f^\downarrow$) PDFs

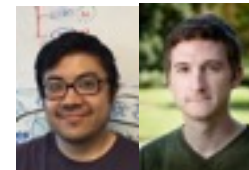
→ e.g. CTEQ-JLab (CJ), JLab Angular Momentum (JAM), ...



CJ (Owens, Accardi, Keppel, WM...)



JAM (Sato, Ethier, WM...)



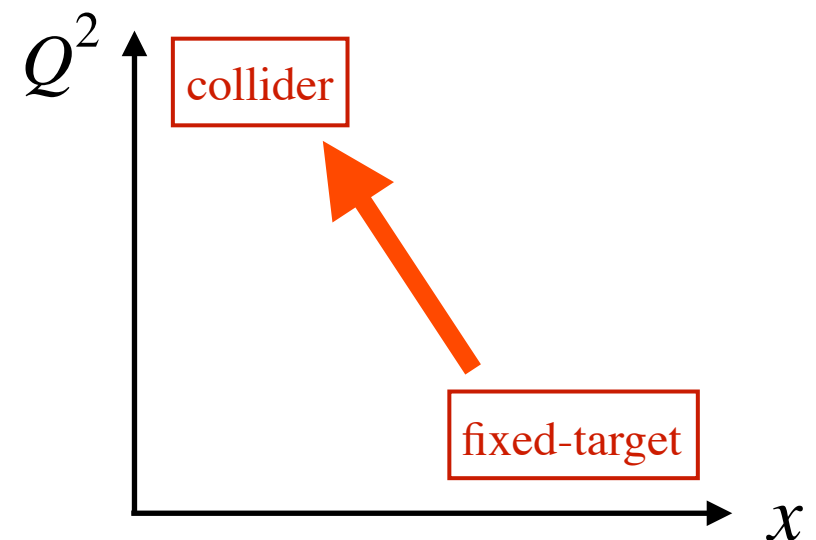
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→ global QCD analyses of spin-averaged ($f = f^\uparrow + f^\downarrow$) and spin-dependent ($\Delta f = f^\uparrow - f^\downarrow$) PDFs

☀ Precision PDFs needed to

- (1) understand basic structure of QCD bound states
- (2) compute backgrounds in searches for BSM physics

→ Q^2 evolution feeds
low x , high Q^2 (“LHC”)
from high x , low Q^2 (“JLab”)



Valence quarks & QCD models

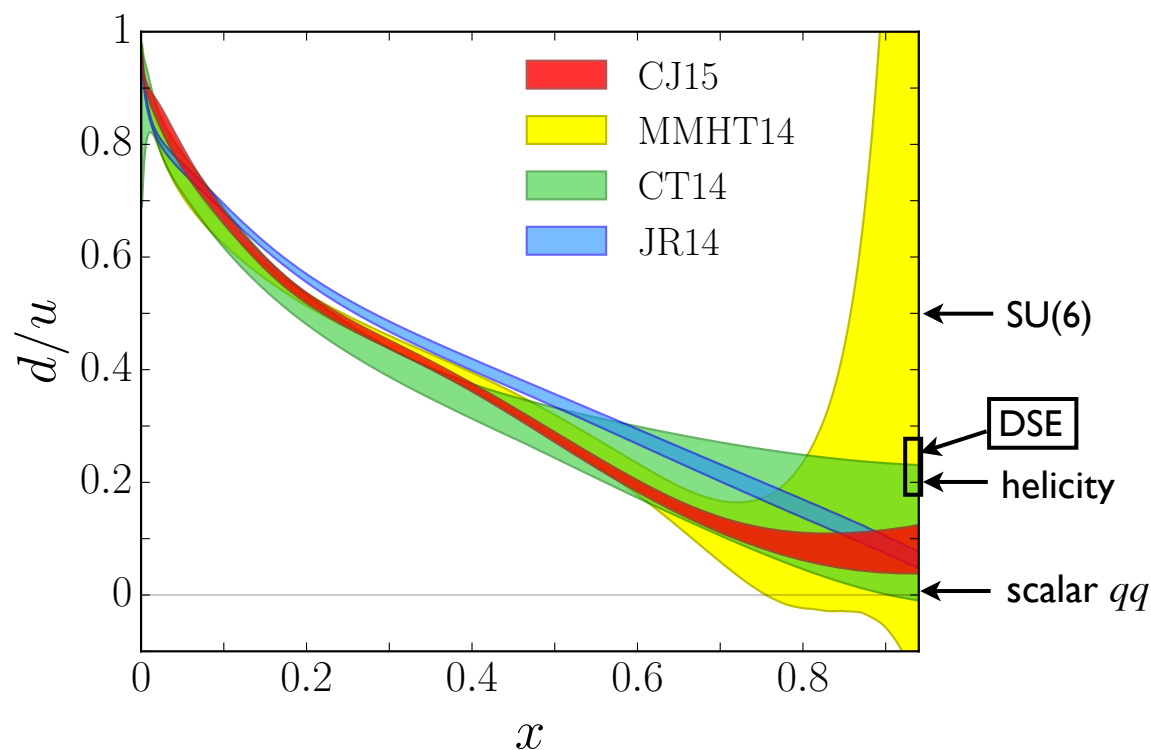
☀ Valence d/u ratio at high x of particular interest

→ testing ground for nucleon models in $x \rightarrow 1$ limit

- $d/u \rightarrow 1/2$
SU(6) symmetry
- $d/u \rightarrow 0$
 $S = 0$ qq dominance
(color-hyperfine interaction)
- $d/u \rightarrow 1/5$
 $S_z = 0$ qq dominance
(perturbative gluon exchange)
- $d/u \rightarrow 0.18 - 0.28$
DSE with qq correlations



Roberts

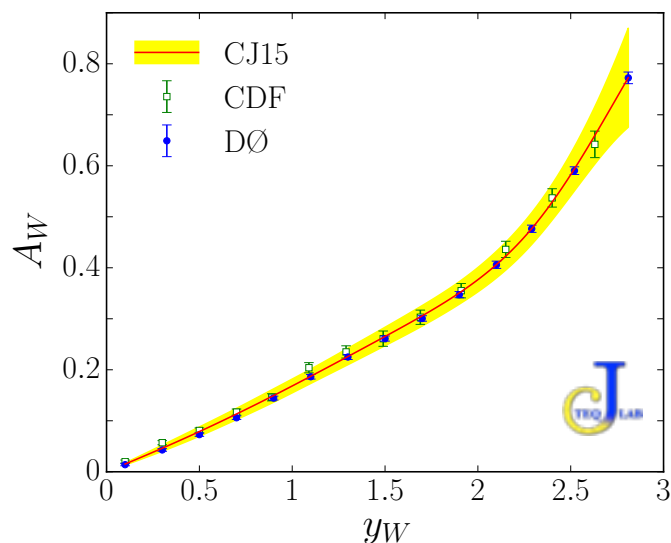


→ considerable uncertainty at high x from deuterium corrections (no free neutrons!)

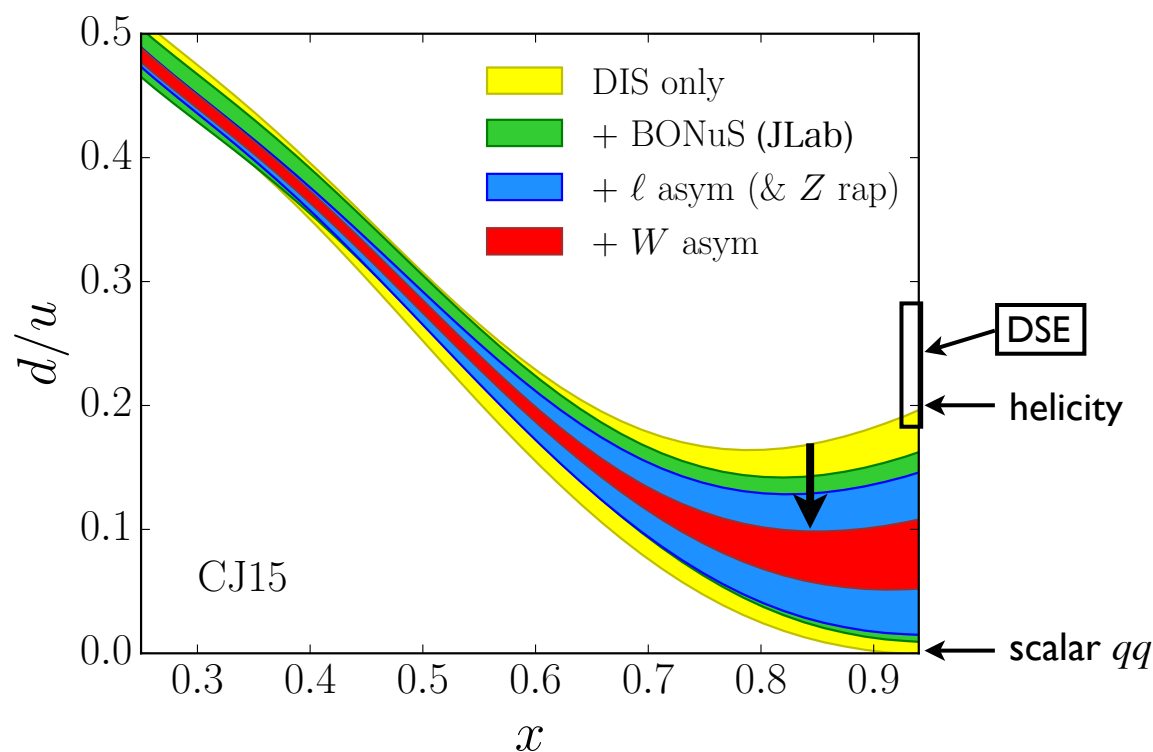
Valence quarks & QCD models

☀ Valence d/u ratio at high x of particular interest

→ significant reduction of PDF errors with new JLab tagged neutron & FNAL W -asymmetry data



Accardi et al. (2016)



→ extrapolated ratio at $x = 1$

$$d/u \rightarrow 0.09 \pm 0.03$$

does not match any model!

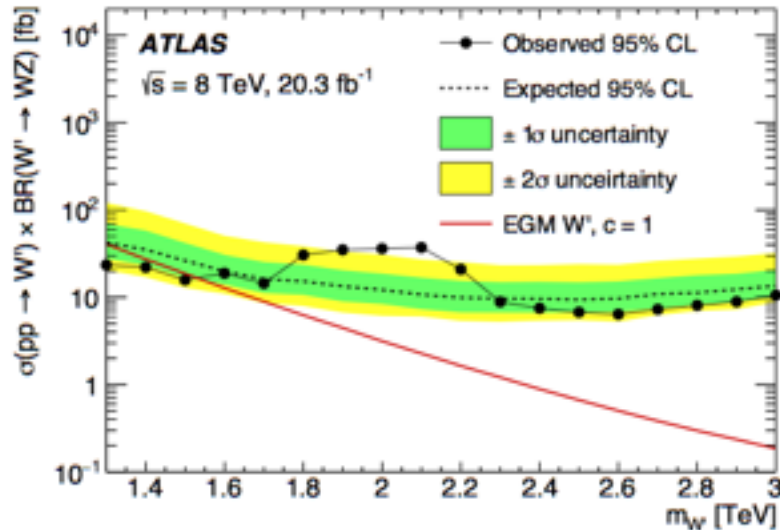
→ upcoming experiments at JLab (MARATHON, BONuS, SoLID) will determine d/u up to $x \sim 0.85$

Valence quarks & BSM searches



Observation of new physics signals requires accurate determination of QCD backgrounds, which depend on PDFs

→ e.g. Higgs boson, or heavy W' boson production at LHC



- 3.4 σ excess in WZ diboson channel at ~ 2 TeV
- extended gauge model $W' \rightarrow WZ$ with $M < 1.5$ TeV excluded at 95% c.l.

→ for W'^- production, parton luminosity is

$$\begin{aligned} \mathcal{L}_{W'^-} &\sim x_1 x_2 \left[\cos^2 \theta_C (d(x_1) \bar{u}(x_2) + s(x_1) \bar{c}(x_2)) \right. \\ &\quad \left. + \sin^2 \theta_C (s(x_1) \bar{u}(x_2) + d(x_1) \bar{c}(x_2)) \right] + (x_1 \leftrightarrow x_2) \\ &\sim d(x_1) \bar{u}(x_2) \quad \text{at large rapidity } y_{W'} \end{aligned}$$

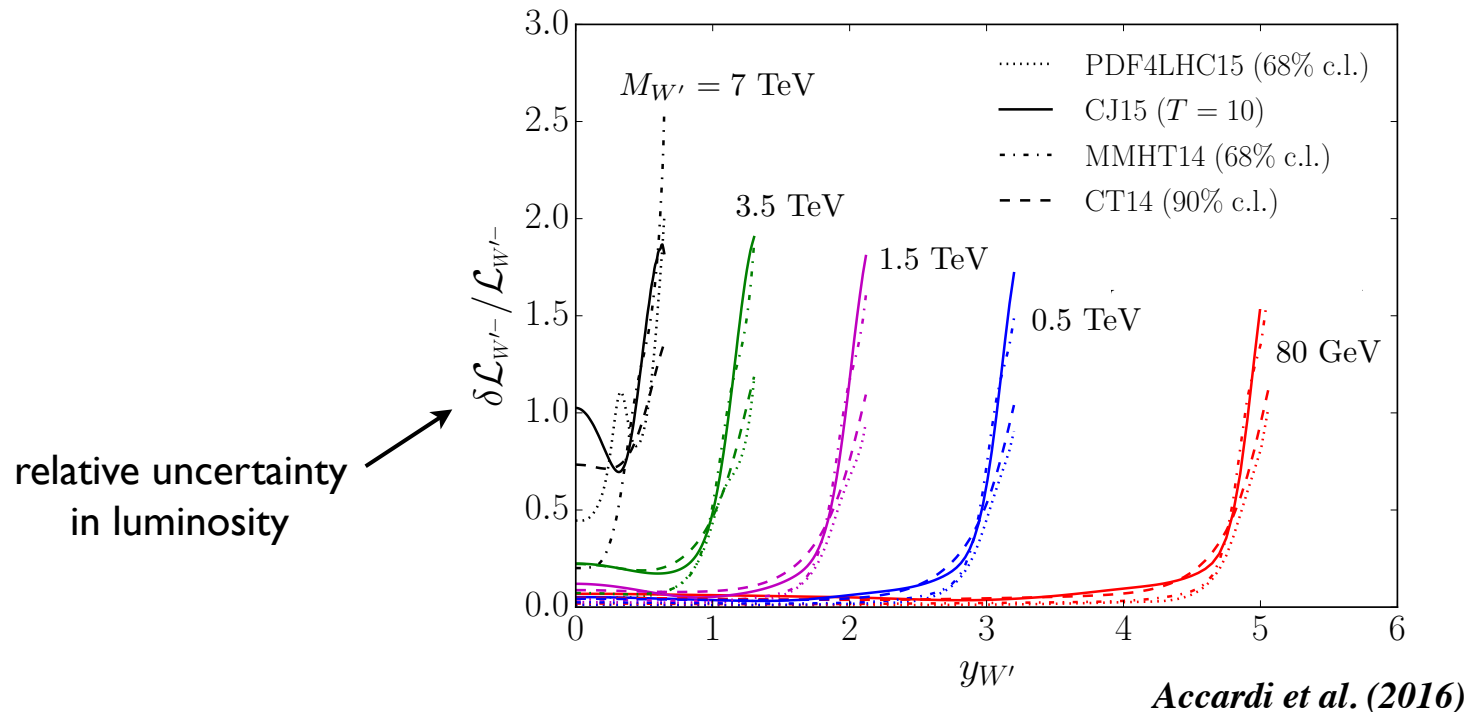
$$x_{1,2} = \frac{M_{W'}}{\sqrt{s}} e^{\pm y_{W'}}$$

Valence quarks & BSM searches



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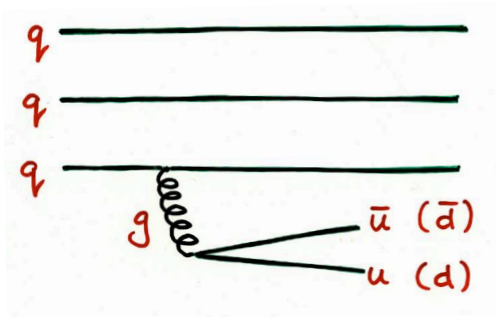
→ large- x uncertainties scale with mass



→ PDF uncertainty small at low $y_{W'}$, rises dramatically at large $y_{W'}$, for all $M_{W'}$

Light quark sea

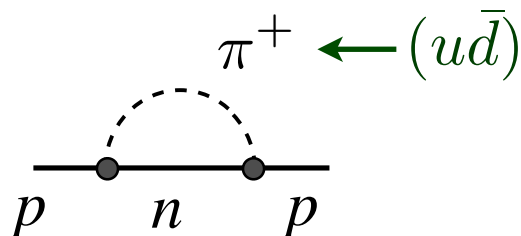
- From perturbative QCD expect symmetric $q\bar{q}$ sea generated by gluon radiation into $q\bar{q}$ pairs (if quark masses are the same)



→ since u and d quarks nearly degenerate, expect flavor-symmetric light-quark sea

$$\bar{d} \approx \bar{u}$$

- In 1984 Thomas made audacious suggestion that chiral symmetry of QCD (important at low energies) should have consequences for antiquark PDFs in the nucleon (at high energies)

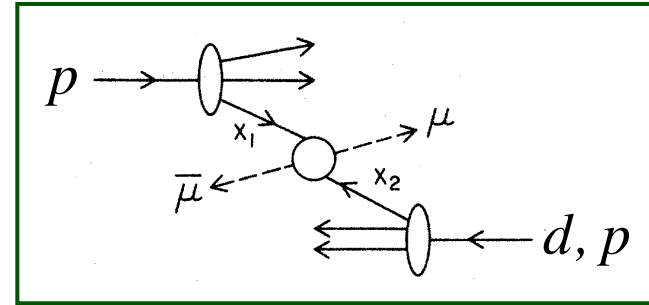
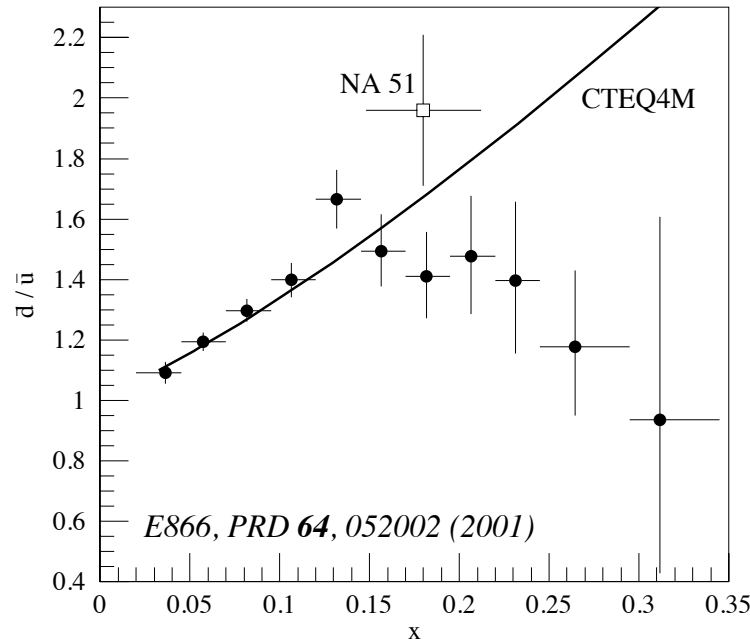


→ $\bar{d} > \bar{u}$

Light quark sea



Asymmetry spectacularly confirmed (more than a decade later)
in high-precision DIS and Drell-Yan experiments



$$\frac{d\sigma}{dx_1 dx_2} \sim \sum_q e_q^2 q(x_1) \bar{q}(x_2) + (x_1 \leftrightarrow x_2)$$

$$\frac{\sigma^{pd}}{\sigma^{pp}} \approx 1 + \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \text{ for } x_1 \gg x_2$$

→ firmly established role of chiral symmetry and pion cloud
as central to understanding of nucleon's quark structure

$$(\bar{d} - \bar{u})(x) = (f_\pi \otimes \bar{q}_\pi)(x)$$

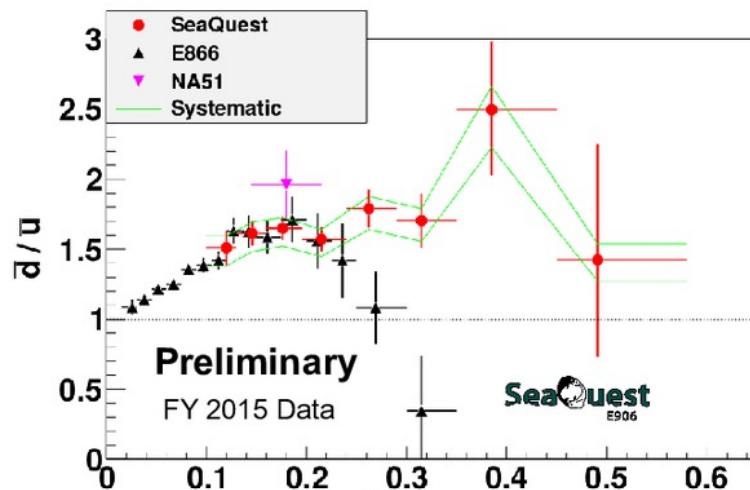
pion distribution
in nucleon

pion PDF

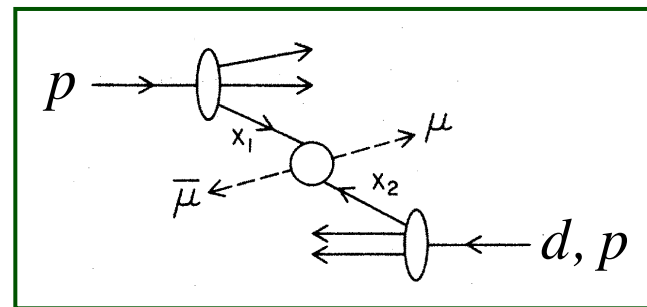
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B. Kerns (April 2016)



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pion distribution
in nucleon

pion PDF

Light quark sea



Early calculations used phenomenological models
— more recently rigorous connection with QCD
established via effective chiral field theory

$$\mathcal{L}_{\text{eff}} = \frac{g_A}{2f_\pi} \bar{\psi}_N \gamma^\mu \gamma_5 \vec{\tau} \cdot \partial_\mu \vec{\pi} \psi_N - \frac{1}{(2f_\pi)^2} \bar{\psi}_N \gamma^\mu \vec{\tau} \cdot (\vec{\pi} \times \partial_\mu \vec{\pi}) \psi_N$$



Weinberg (1967)

→ lowest order πN interaction includes
pion rainbow and tadpole contributions

→ matching quark- and hadron-level operators

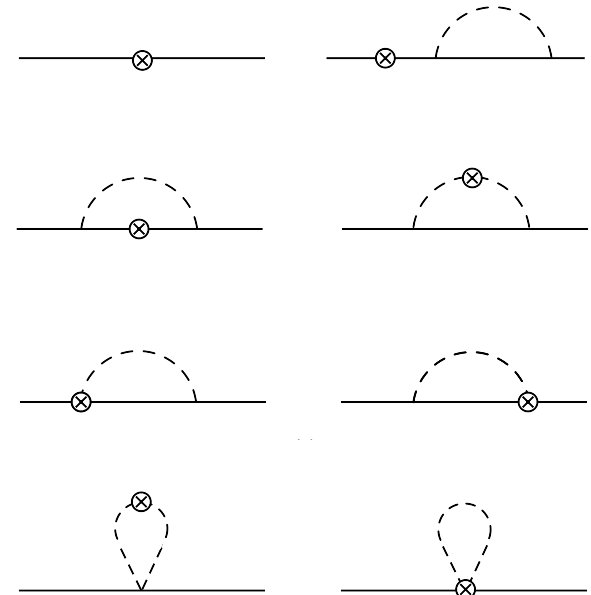
$$\mathcal{O}_q^{\mu_1 \dots \mu_n} = \sum_h c_{q/h}^{(n)} \mathcal{O}_h^{\mu_1 \dots \mu_n}$$

yields convolution representation

$$q(x) = \sum_h \int_x^1 \frac{dy}{y} f_h(y) q_v^h(x/y)$$



Xiangdong Ji (2001)



Light quark sea



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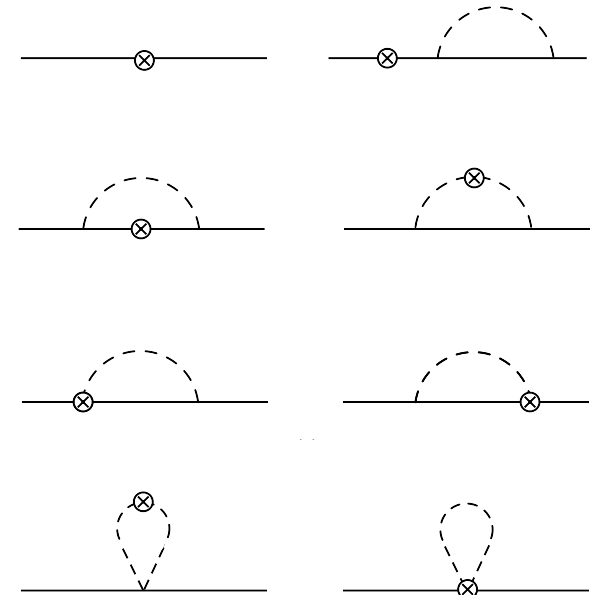
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Stelios Kazantzidis (?)



Light quark sea



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Weinberg (1967)

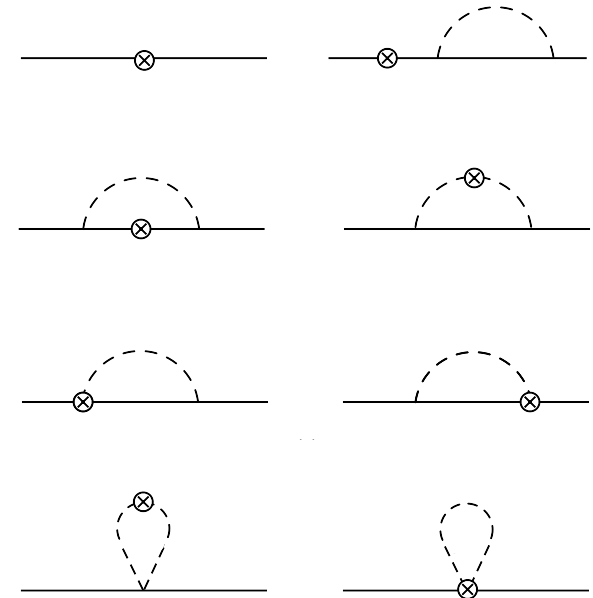
→ expanding PDF moments in powers of m_π ,
coefficients of leading nonanalytic (LNA)
terms are model-independent

Thomas, WM, Steffens (2000)

Chuang Ji, WM, Thomas (2012)



**Chuang Ji talk
Mon. 15:00**

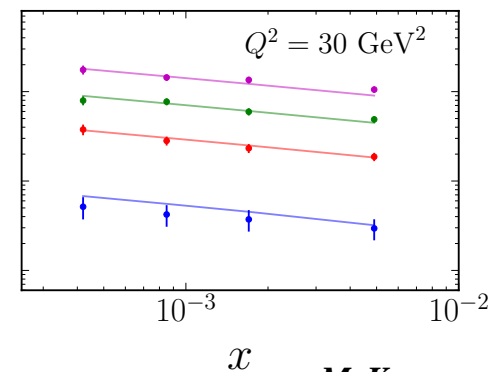
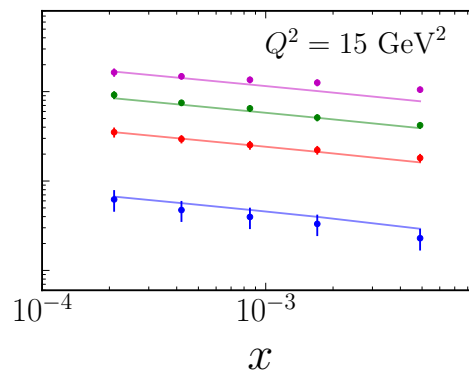
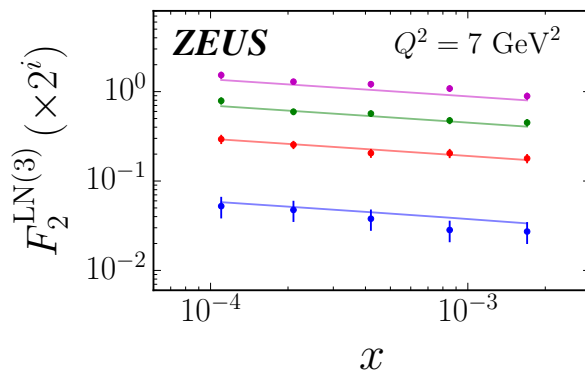
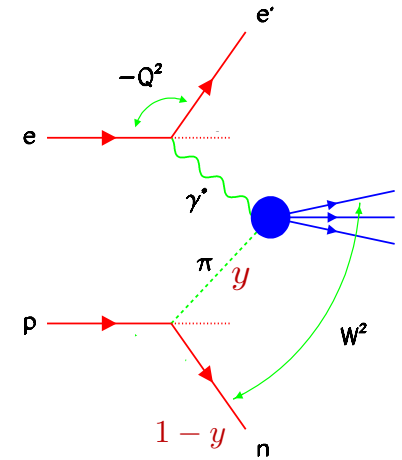


Light quark sea

☀ Drell-Yan $\bar{d} - \bar{u}$ data can be described using range of UV regulators (shapes of pion momentum distributions)

→ semi-inclusive production of “leading neutrons” at HERA can discriminate between different shapes

$$\frac{d^3\sigma^{\text{LN}}}{dx dQ^2 dy} \sim F_2^{\text{LN}(3)}(x, Q^2, y) = f_\pi(y) F_2^\pi(x/y, Q^2)$$



McKenney et al. (2016)

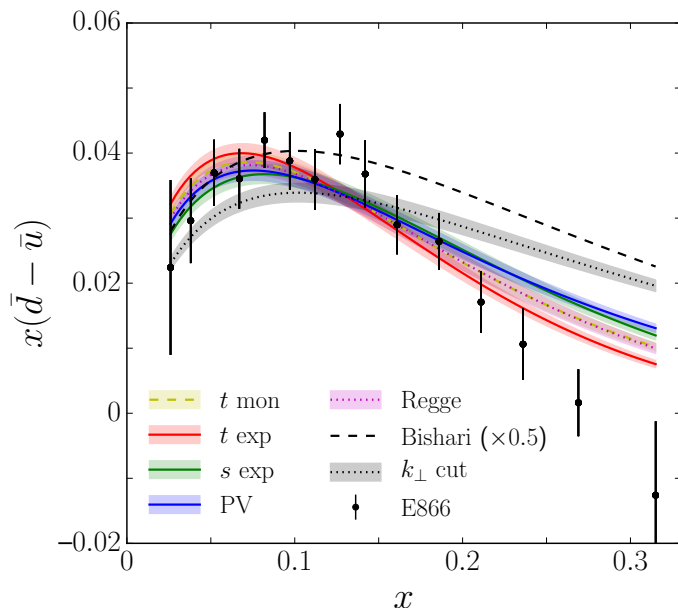
→ best fit to ZEUS & H1 data prefers t -dependent exponential regulator

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→ constrain shape of F_2^π at $10^{-4} \lesssim x_\pi \lesssim 0.03$ from combined HERA + Drell-Yan fit

McKenney, Sato, WM, Chueng Ji (2016)

→ global analysis under way of HERA LN, Drell-Yan πN + pd/pp (+ future JLab TDIS data) to determine pion PDFs at all x

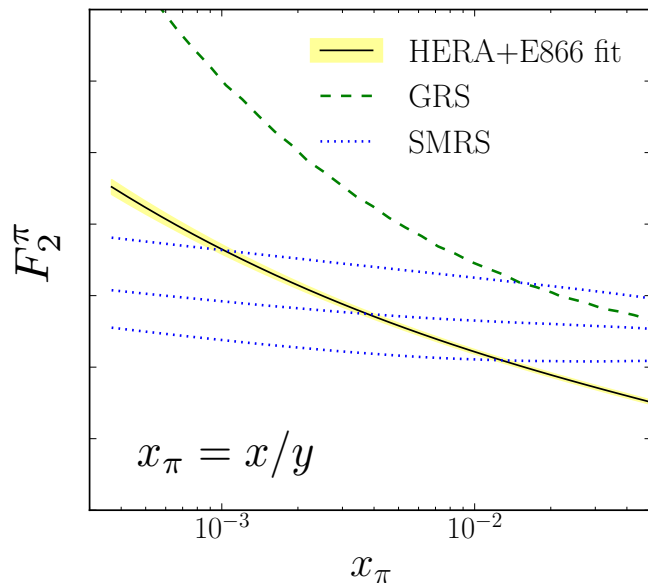
Barry, Chueng Ji, WM, Sato (2016)

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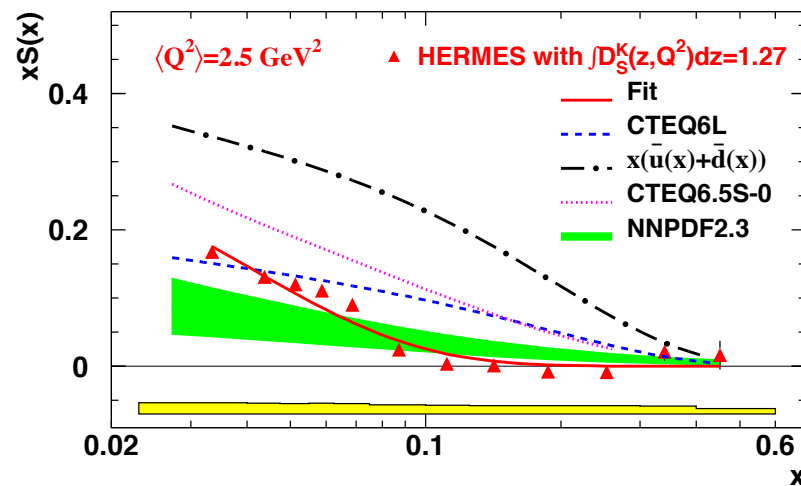
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Barry, Chueng Ji, WM, Sato (2016)

Strange quarks

☀ Strange quark PDFs most directly determined from dimuon production in (anti)neutrino-nucleus DIS ($W^+ s \rightarrow c$ / $W^- \bar{s} \rightarrow \bar{c}$)

- but significant uncertainty from nuclear corrections, semileptonic branching ratio uncertainty
- tension with HERMES semi-inclusive K -production data



historically, strange to nonstrange ratio

$$\kappa = \frac{s + \bar{s}}{\bar{u} + \bar{d}} \sim 0.2 - 0.5$$

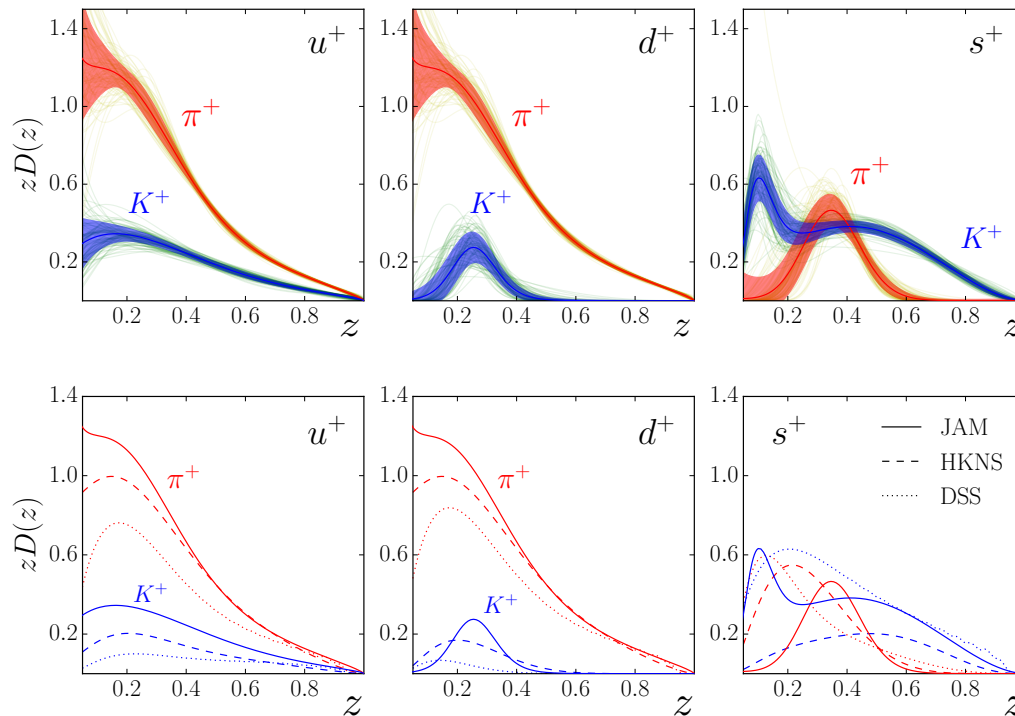
... but uncertainty from K fragmentation functions

Strange quarks



Fragmentation functions (FFs) determined from single-inclusive meson production in e^+e^- annihilation

→ new “iterative Monte Carlo” (IMC) global analysis (including new Belle & BaBar data) suggests differences with previous extractions



Nobuo Sato et al. (2016)
(arXiv this week)



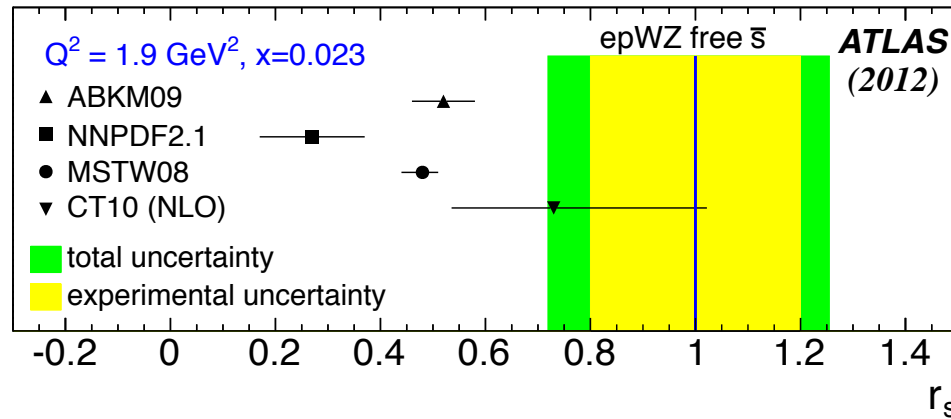
→ SIDIS data also constrain fragmentation functions,
but require simultaneous PDF + FF fit (currently in progress)

Strange quarks



Alternatively, probe strange PDF in W/Z production at LHC

$pp \rightarrow W(Z) + X$ free of nuclear effects

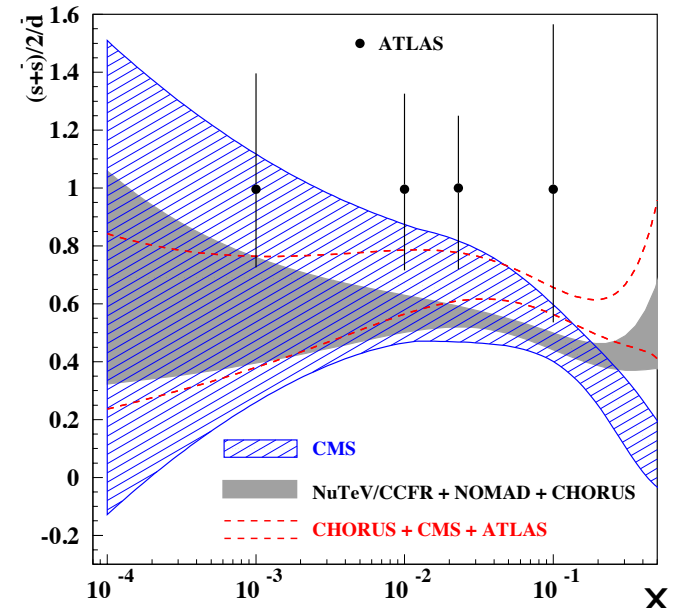


→ surprisingly large strangeness

$$r_s = (s + \bar{s})/2\bar{d}$$

$$= 1.00^{+0.25}_{-0.28}$$

→ more recent reanalysis of neutrino DIS (CHORUS, NOMAD) and ATLAS data does not support enhanced strange PDF

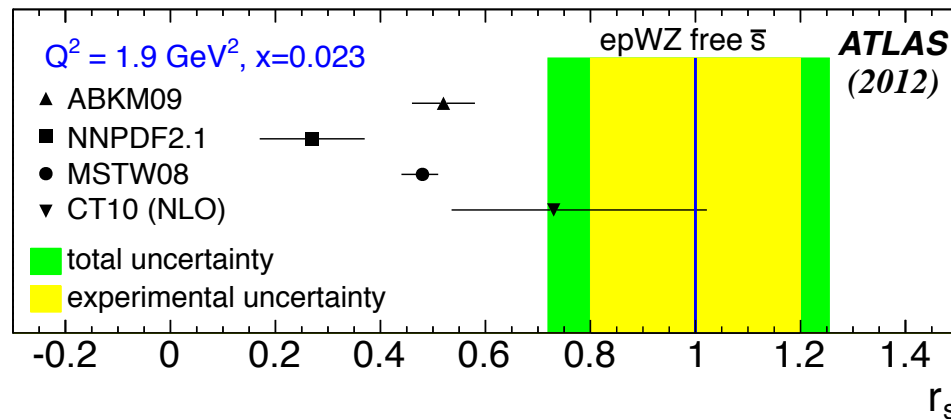


Alekhin et al. (2015)

Strange quarks



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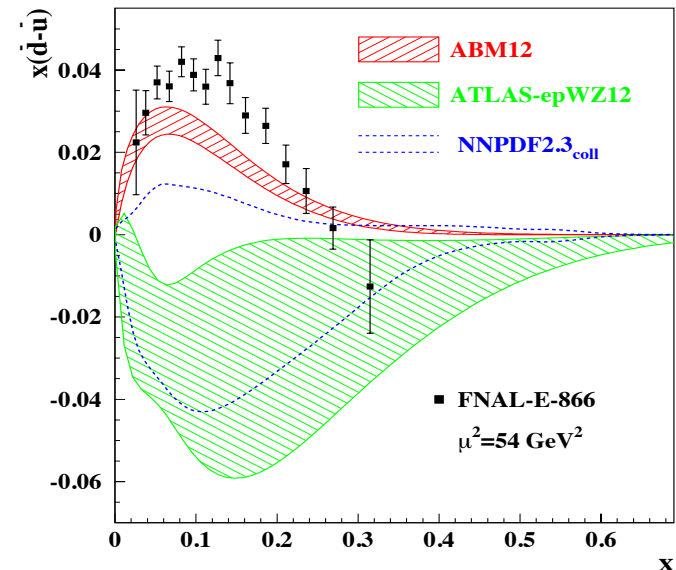


→ surprisingly large strangeness

$$r_s = (s + \bar{s})/2\bar{d}$$

$$= 1.00^{+0.25}_{-0.28}$$

→ suggests effect related to underestimated \bar{d} PDF
 — collider data alone cannot (yet) disentangle flavors

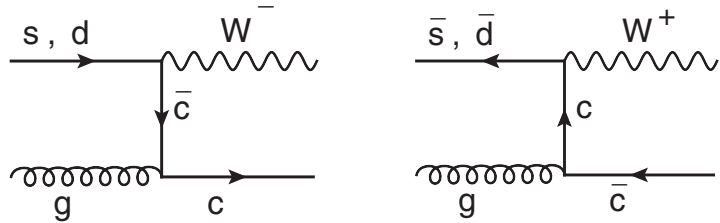


Alekhin et al. (2015)

Strange quarks



More reliable determination may come from associated production of W + charm jet events



→ discriminate between s and \bar{s} in the proton!

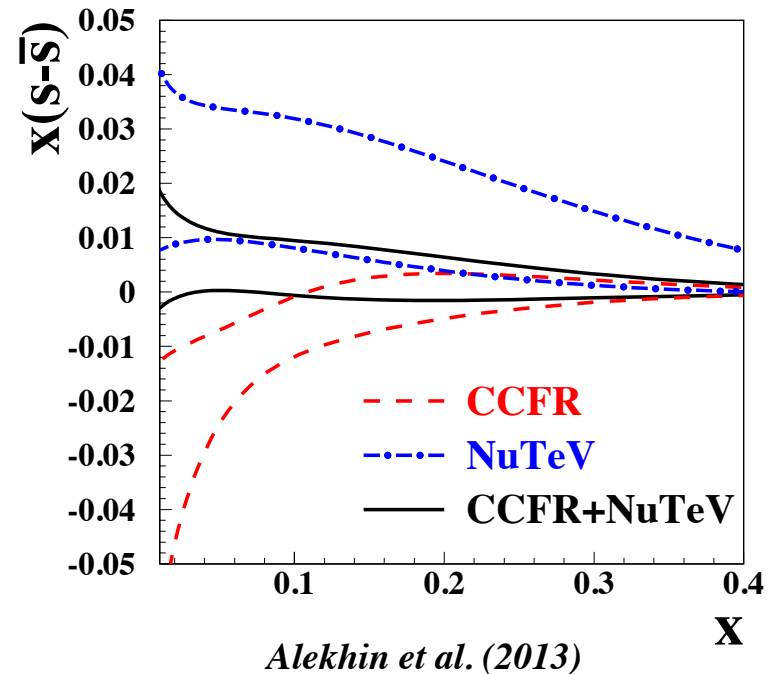
→ some indication from (anti)neutrino DIS from NuTeV experiment

$$S^- = \int_0^1 dx x(s - \bar{s})$$

$$= (2.0 \pm 1.4) \times 10^{-3}$$

NuTeV (2007)

... but global analysis including other data is inconclusive

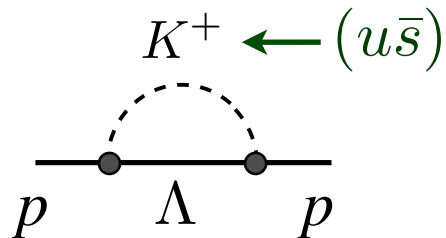


Strange quarks

☀ Strange–antistrange asymmetry important indicator of nonperturbative physics

→ predicted from chiral $SU(3)$ symmetry breaking via “kaon cloud”

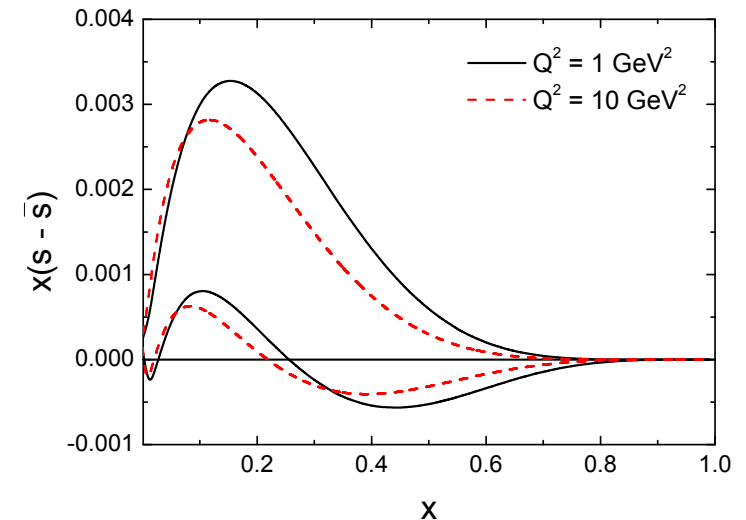
Signal, Thomas (1987)



... but shape more difficult to constrain

☀ Recent chiral effective theory analysis (including rainbow, tadpole & Kroll-Ruderman) favors small positive moment

$$S^- \lesssim 1 \times 10^{-3}$$



Xuangong Wang et al. (2016)

Strange quarks



Parity-violating DIS measurements (*e.g.* SoLID at JLab12) will allow strange contribution to be isolated, when combined with e.m. p and n DIS data at low/intermediate x

→ at leading order

$$F_2^{\gamma p} = \frac{4}{9}x(u + \bar{u}) + \frac{1}{9}x(d + \bar{d} + s + \bar{s}) + \dots$$

$$F_2^{\gamma n} = \frac{4}{9}x(d + \bar{d}) + \frac{1}{9}x(u + \bar{u} + s + \bar{s}) + \dots$$

$$F_2^{\gamma Z, p} = \left(\frac{1}{3} - \frac{8}{9} \sin^2 \theta_W \right) x(u + \bar{u}) + \left(\frac{1}{6} - \frac{2}{9} \sin^2 \theta_W \right) (d + \bar{d} + s + \bar{s}) + \dots$$

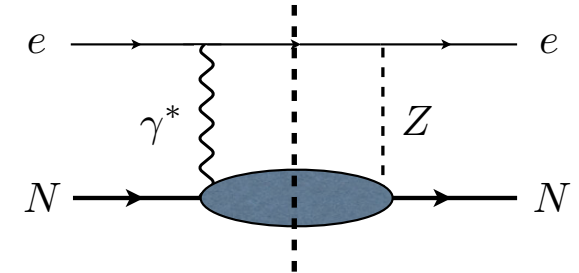
$$\approx \frac{1}{9}x(u + \bar{u} + d + \bar{d} + s + \bar{s}) + \dots \quad \text{for } \sin^2 \theta_W \approx 1/4$$

→ 3 equations with 3 unknowns;

order of magnitude greater sensitivity of γZ to strange PDF

→ $V \times A$ term also sensitive to $s - \bar{s}$

$$F_3^{\gamma Z, p} = \frac{2}{3}(u - \bar{u}) + \frac{1}{3}(d - \bar{d} + s - \bar{s}) + \dots$$

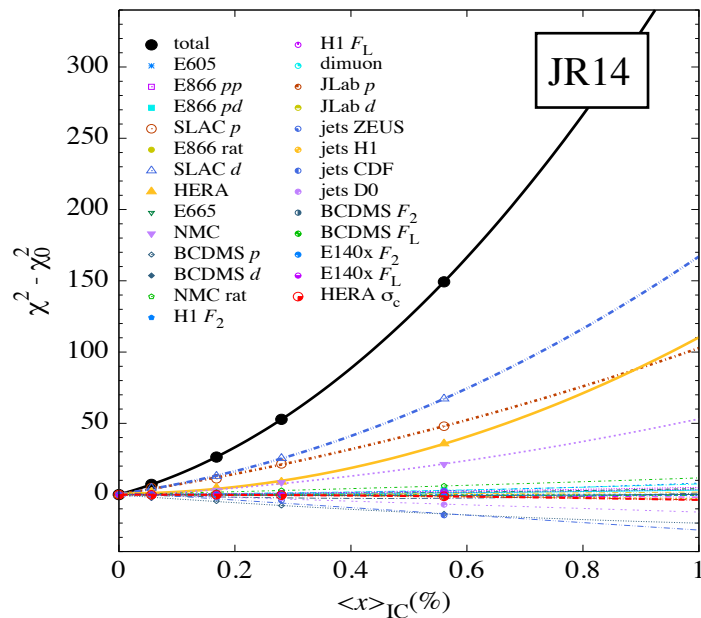


Charm in the nucleon

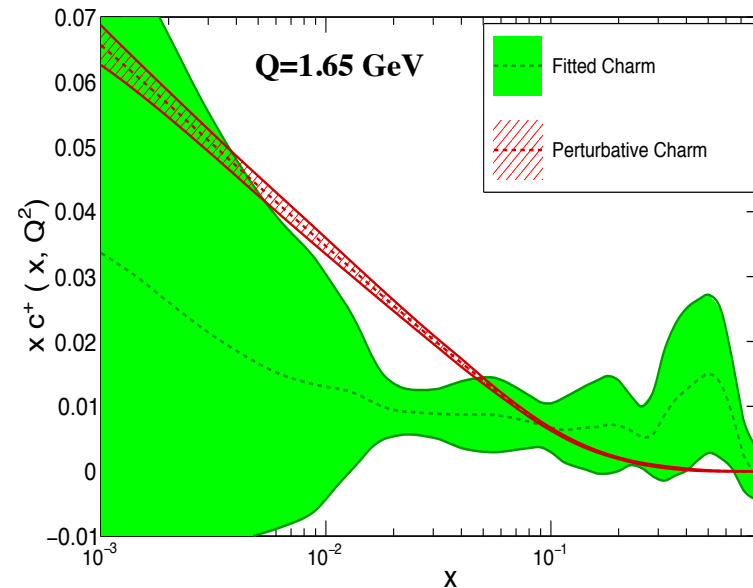


Global PDF analyses find no evidence for large “intrinsic charm” (IC)

→ different upper limits obtained depend on definition of tolerance & treatment of nonperturbative effects (thresholds, masses, ...)



Jimenez-Delgado et al. (2015)



NNPDF (2016)

- with standard fitting technology
momentum carried by “IC”
 $\langle x \rangle_{IC} < 0.1\%$ at 5σ CL

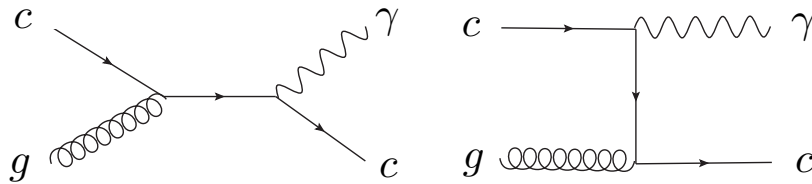
- with “FONNL-B” HQ scheme,
 $\langle x \rangle_{IC} < 1\%$ at 1σ CL, but can go
 < 0 at low x to fit EMC F_2^c data

Charm in the nucleon



Associated prompt photon + charm production

$pp \rightarrow \gamma + c + X$ may reveal “intrinsic” charm component



Bednyakov et al. (2014)

→ “smoking gun” would be observation of asymmetric distributions $c(x) \neq \bar{c}(x)$



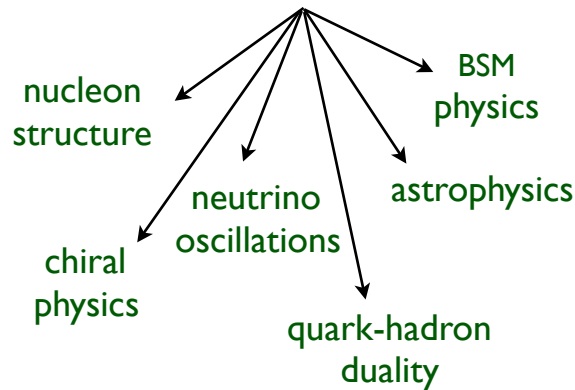
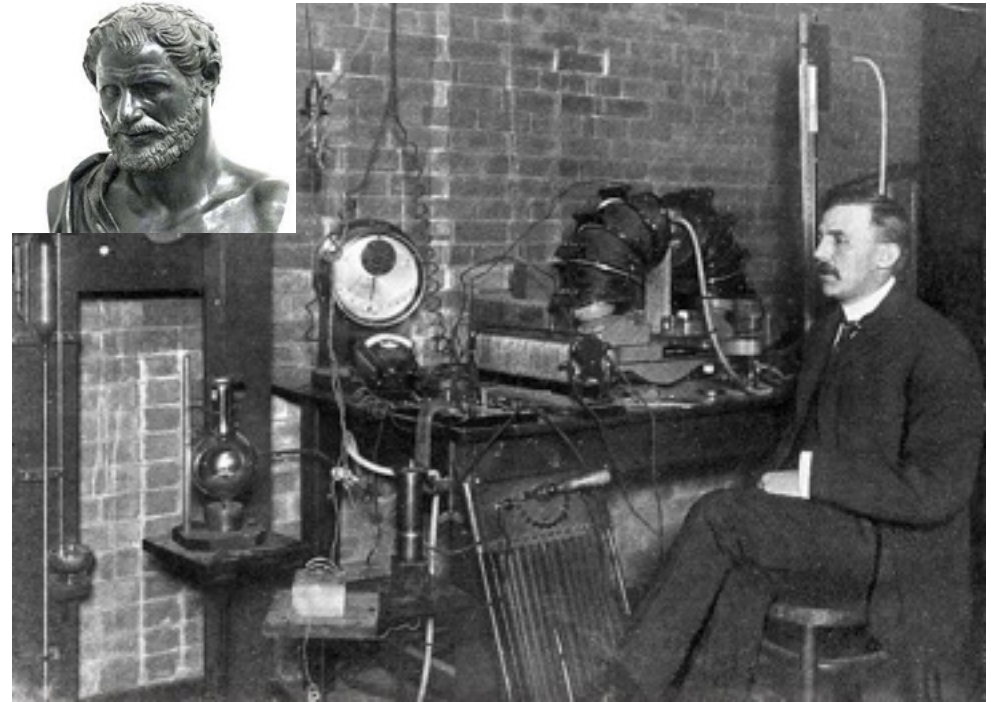
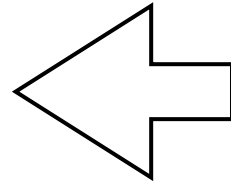
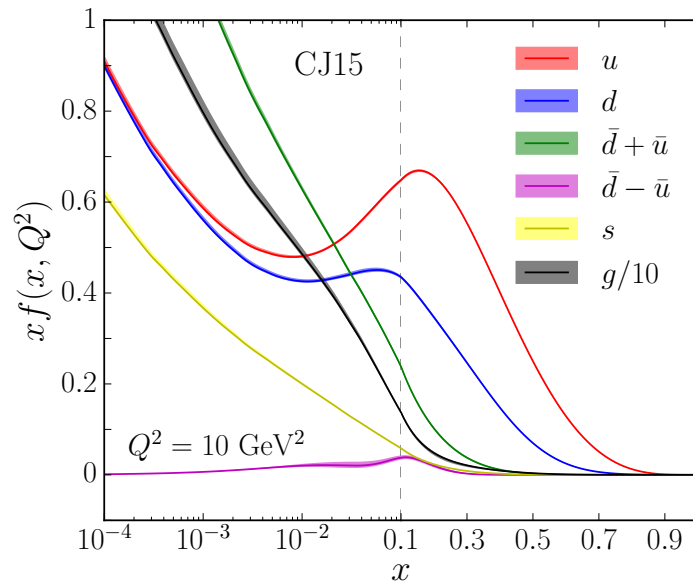
***Tim Hobbs talk
Tue. 18:30***



Outlook and new directions



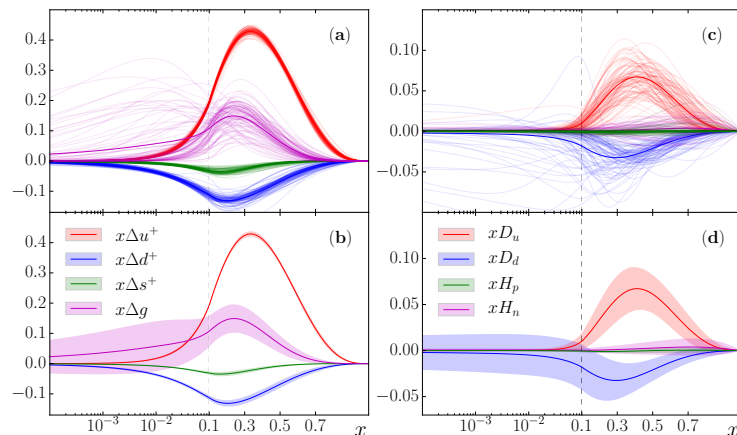
Study of PDFs has brought together essential elements of nuclear and high-energy physics



Outlook and new directions

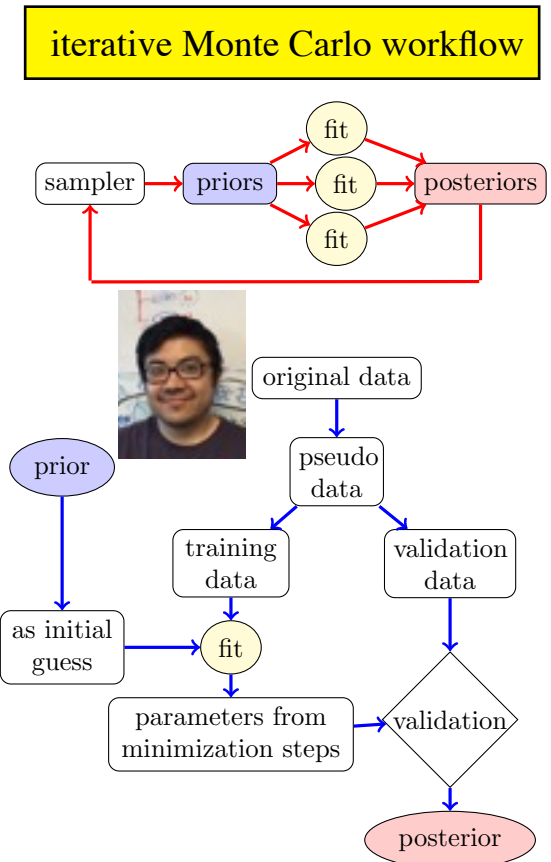
☀ New approach to global QCD analysis — “IMC”

- minimize bias from choice of initial parameters
- statistically rigorous PDF uncertainties, without assumptions about “tolerance” or Gaussianity



Sato et al. (2016)

- ultimate goal of simultaneous fit of unpolarized and polarized PDFs and fragmentation functions in sight (eventually generalize to TMDs)



Outlook and new directions



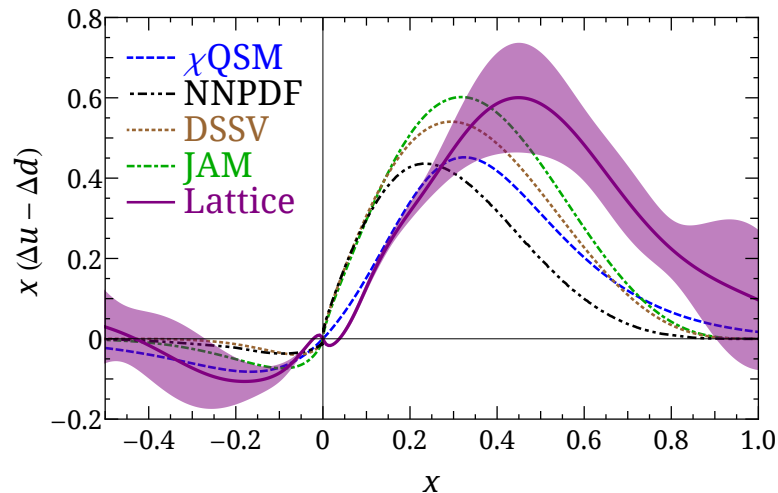
PDFs on the lattice

- compute light-cone correlations of partons by boosting matrix elements of spatial correlations to large momentum
 - large momentum effective field theory (LaMET)

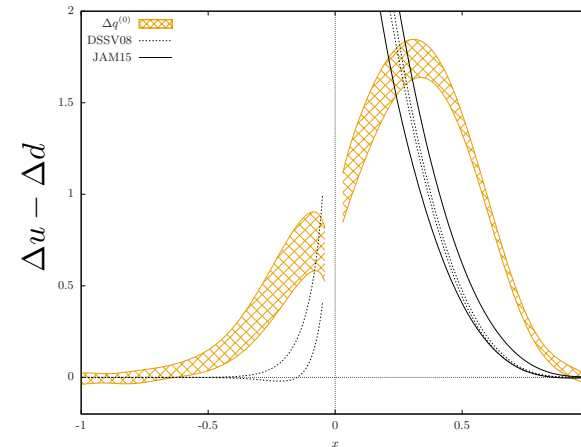


Xiangdong Ji (2013)

- challenges to overcome, but first computations promising



H.-W. Lin, J.-W. Chen et al. (2015-16)



Alexandrou et al. (2016)

Το τέλος.

Ευχαριστώ!