



*Large- x ("non-small x ") Meeting
Jefferson Lab, October 4, 2016*

Pion structure from leading particle production

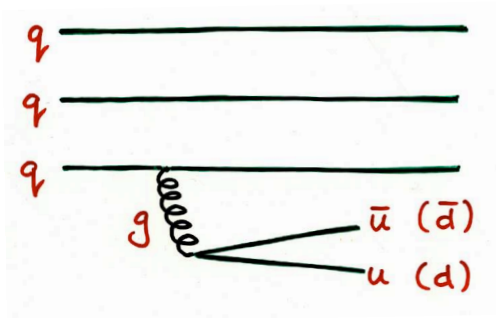
Wally Melnitchouk



*with Chueng Ji (NCSU), Josh McKinney (UNC),
Nobuo Sato (JLab), Tony Thomas (Adelaide)*

Light quark sea

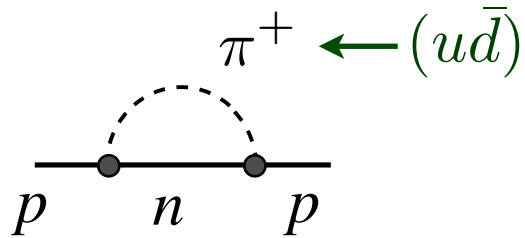
- From perturbative QCD expect symmetric $q\bar{q}$ sea generated by gluon radiation into $q\bar{q}$ pairs (if quark masses are the same)



→ since u and d quarks nearly degenerate, expect flavor-symmetric light-quark sea

$$\bar{d} \approx \bar{u}$$

- In 1980s Thomas argued that chiral symmetry of QCD (important at low energies) should have consequences for antiquark PDFs in the nucleon (at high energies)

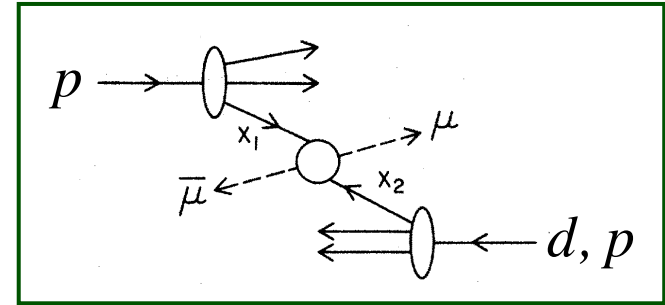
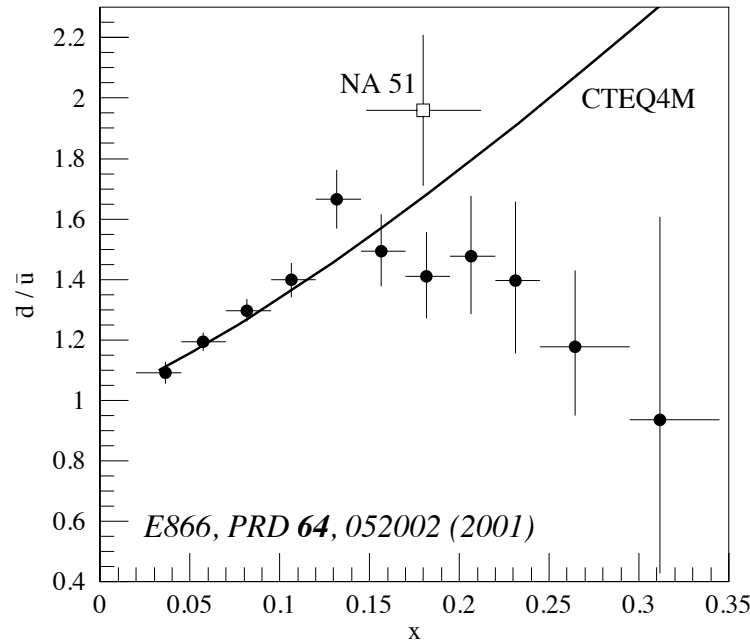


$$\bar{d} > \bar{u}$$

Light quark sea



Asymmetry spectacularly confirmed in high-precision DIS and Drell-Yan experiments



$$\frac{d\sigma}{dx_1 dx_2} \sim \sum_q e_q^2 q(x_1) \bar{q}(x_2) + (x_1 \leftrightarrow x_2)$$

$$\frac{\sigma^{pd}}{\sigma^{pp}} \approx 1 + \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \text{ for } x_1 \gg x_2$$

→ strongly suggested role of chiral symmetry and pion cloud as central to understanding of nucleon's quark structure

$$(\bar{d} - \bar{u})(x) = (f_\pi \otimes \bar{q}_\pi)(x)$$

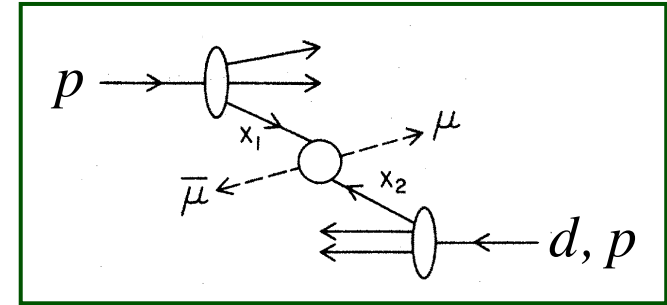
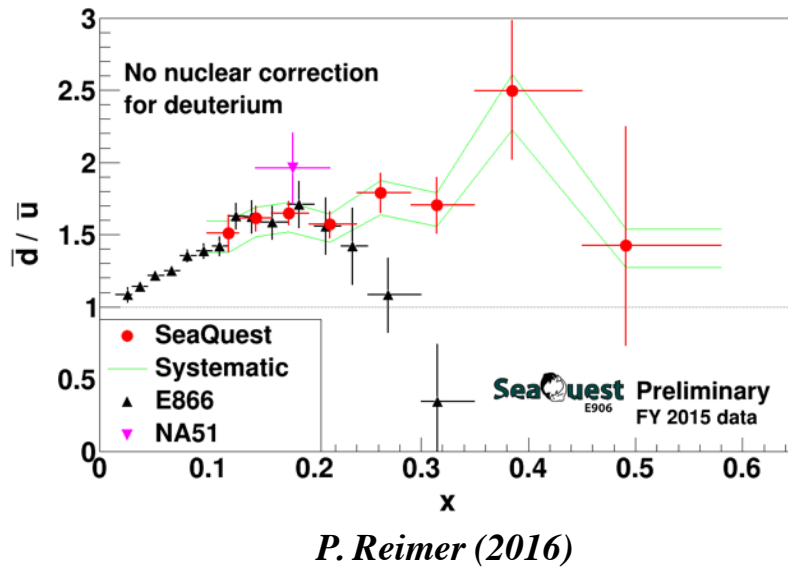
pion distribution
in nucleon

pion PDF

Light quark sea



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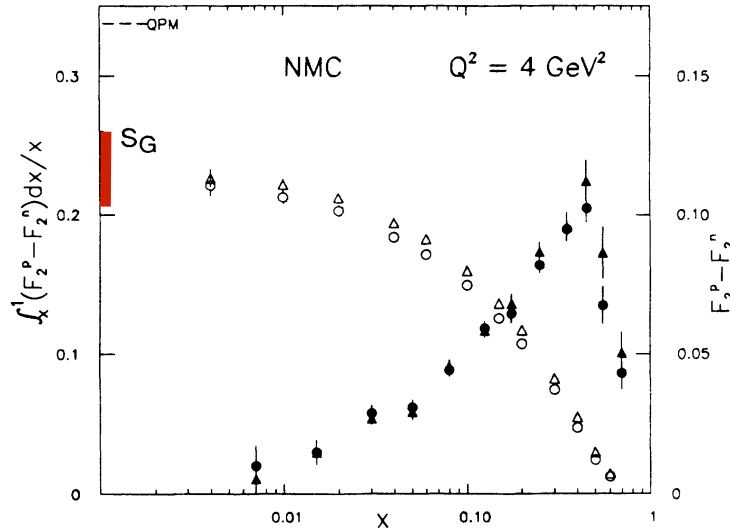
pion distribution
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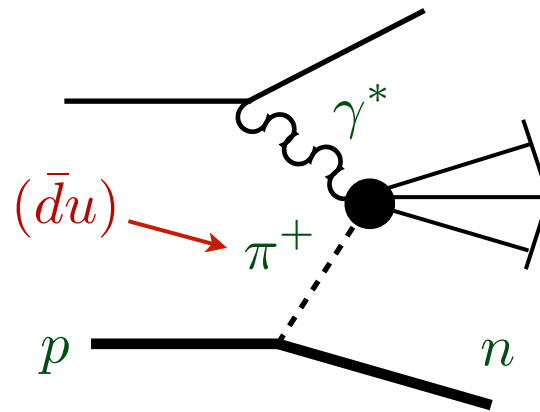


NMC (1994)

$$\int_0^1 \frac{dx}{x} (F_2^p - F_2^n) = \frac{1}{3} - \frac{2}{3} \int_0^1 dx (\bar{d} - \bar{u}) = 0.235(26)$$

violation of Gottfried sum rule!

→ Sullivan process —
DIS from pion cloud
of the nucleon



Sullivan (1972)

Chiral effective theory



Early calculations used phenomenological models
— more recently rigorous connection with QCD
established via effective chiral field theory

$$\mathcal{L}_{\text{eff}} = \frac{g_A}{2f_\pi} \bar{\psi}_N \gamma^\mu \gamma_5 \vec{\tau} \cdot \partial_\mu \vec{\pi} \psi_N - \frac{1}{(2f_\pi)^2} \bar{\psi}_N \gamma^\mu \vec{\tau} \cdot (\vec{\pi} \times \partial_\mu \vec{\pi}) \psi_N$$

Weinberg (1967)

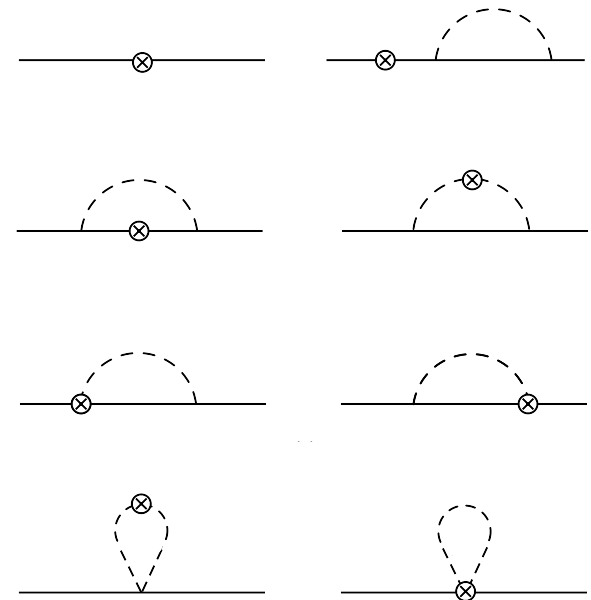
→ lowest order πN interaction includes
pion rainbow and tadpole contributions

→ matching quark- and hadron-level operators

$$\mathcal{O}_q^{\mu_1 \cdots \mu_n} = \sum_h c_{q/h}^{(n)} \mathcal{O}_h^{\mu_1 \cdots \mu_n}$$

yields convolution representation

$$q(x) = \sum_h \int_x^1 \frac{dy}{y} f_h(y) q_v^h(x/y)$$



Chiral effective theory

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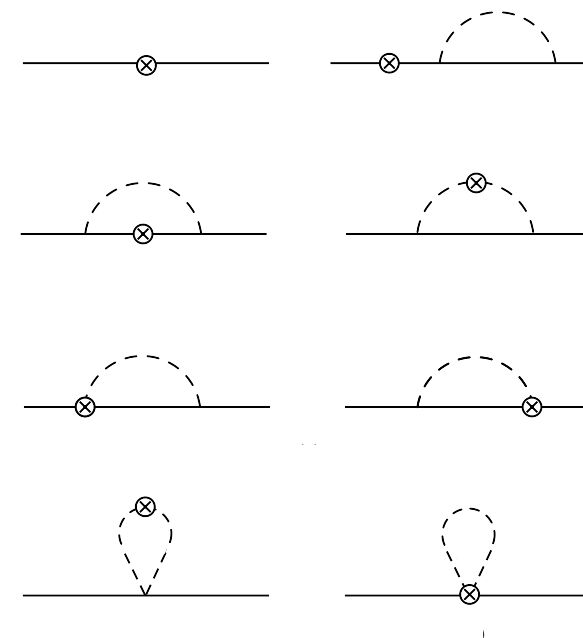
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Weinberg (1967)

- expanding PDF moments in powers of m_π , coefficients of leading nonanalytic (LNA) terms are model-independent!

Thomas, WM, Steffens (2000)

Chuang Ji, WM, Thomas (2012)



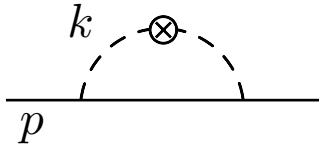
- nonanalytic behavior vital for chiral extrapolation of lattice data on PDF moments

$$\langle x \rangle_{u-d}^{\text{LNA}} \sim m_\pi^2 \log m_\pi^2$$

Detmold et al. (2001)

Pion splitting functions

- Spitting functions for various diagrams computed in chiral theory
e.g. pion rainbow diagram



$$f_{\pi}(y) = f^{(\text{on})}(y) + f^{(\delta)}(y)$$

has on-shell ($y > 0$)
and $\delta(y)$ contributions!

$$f^{(\text{on})}(y) = \frac{g_A^2 M^2}{(4\pi f_{\pi})^2} \int dk_{\perp}^2 \frac{y(k_{\perp}^2 + y^2 M^2)}{[k_{\perp}^2 + y^2 M^2 + (1-y)m_{\pi}^2]^2}$$

$$f^{(\delta)}(y) = \frac{g_A^2}{4(4\pi f_{\pi})^2} \int dk_{\perp}^2 \log \left(\frac{k_{\perp}^2 + m_{\pi}^2}{\mu^2} \right) \delta(y)$$

- For point-like nucleons and pions, integrals divergent

→ finite size of nucleon provides natural regularization scale
(but does not prescribe form of regularization)

$$\mathcal{F} = \Theta(\Lambda^2 - k_{\perp}^2) \quad k_{\perp} \text{ cutoff}$$

$$\mathcal{F} = \exp [(M^2 - s)/\Lambda^2] \quad s\text{-dep. exponential}$$

$$\mathcal{F} = \left(\frac{\Lambda^2 - m_{\pi}^2}{\Lambda^2 - t} \right) \quad t \text{ monopole}$$

$$\mathcal{F} = \left[1 - \frac{(t - m_{\pi}^2)^2}{(t - \Lambda^2)^2} \right]^{1/2} \quad \text{Pauli-Villars}$$

$$\mathcal{F} = \exp [(t - m_{\pi}^2)/\Lambda^2] \quad t \text{ exponential}$$

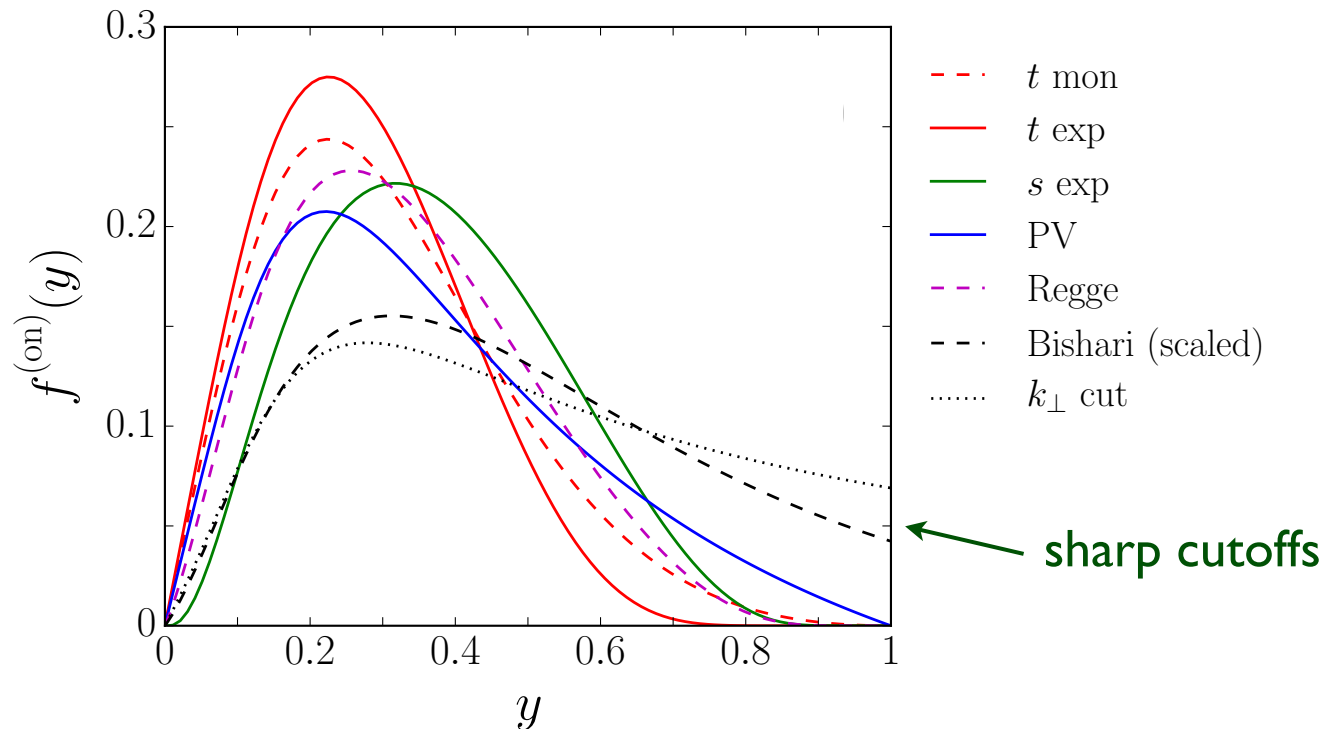
$$\mathcal{F} = y^{-\alpha_{\pi}(t)} \exp [(t - m_{\pi}^2)/\Lambda^2] \quad \text{Regge}$$

Pion splitting functions



Detailed shape of splitting function depends on regularization, but common general features

e.g. on-shell function



$$\mathcal{F} = \Theta(\Lambda^2 - k_{\perp}^2) \quad k_{\perp} \text{ cutoff}$$

$$\mathcal{F} = \left(\frac{\Lambda^2 - m_{\pi}^2}{\Lambda^2 - t} \right) \quad t \text{ monopole}$$

$$\mathcal{F} = \exp \left[(t - m_{\pi}^2) / \Lambda^2 \right] \quad t \text{ exponential}$$

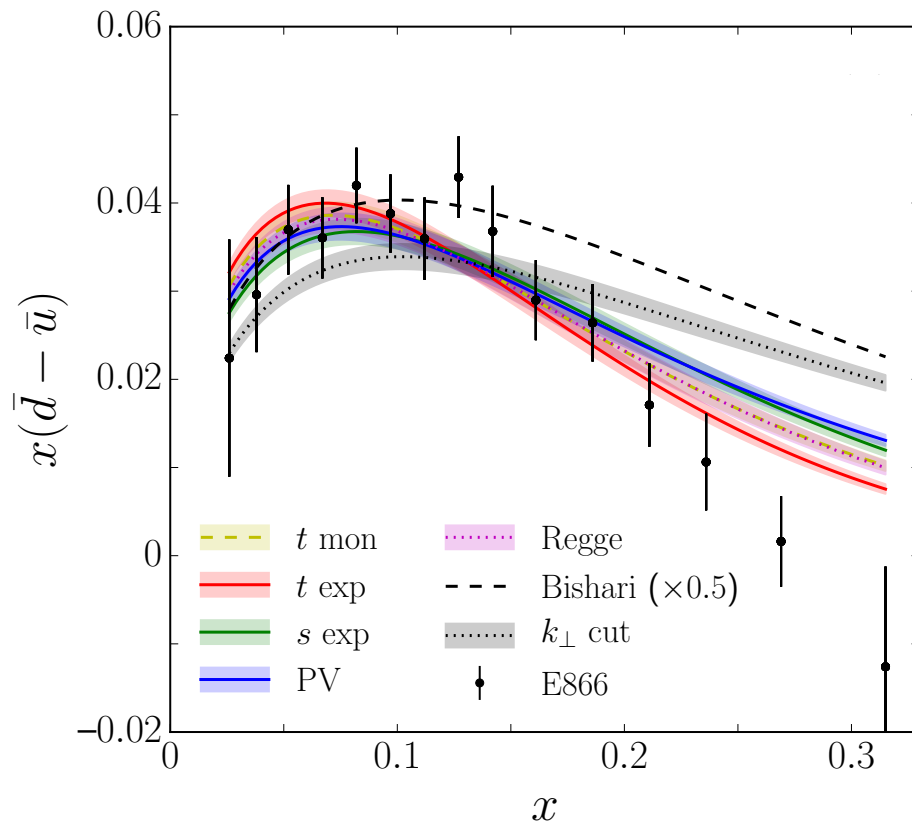
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Pion splitting functions

☀ E866 $\bar{d} - \bar{u}$ data can be fitted with range of regulators



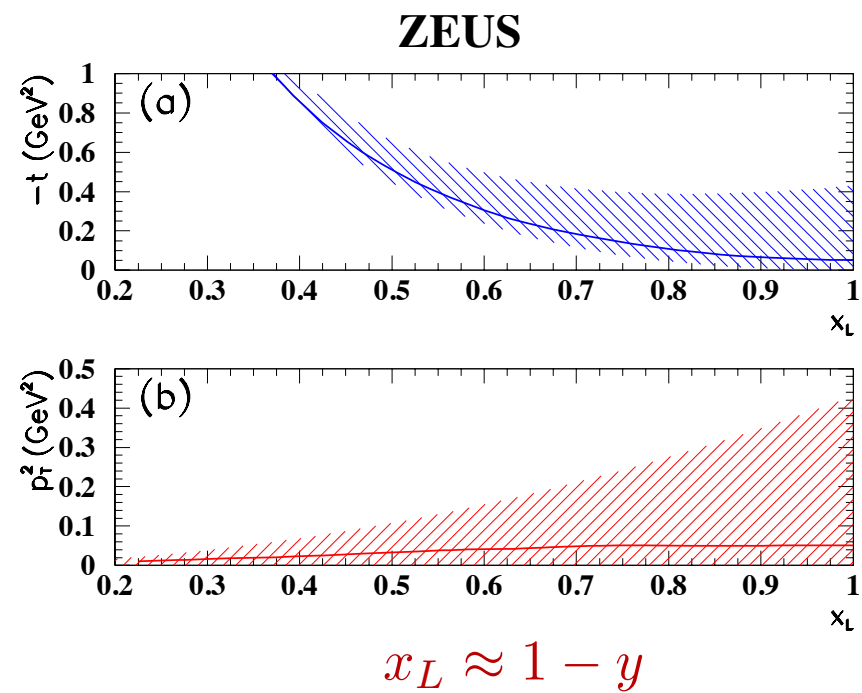
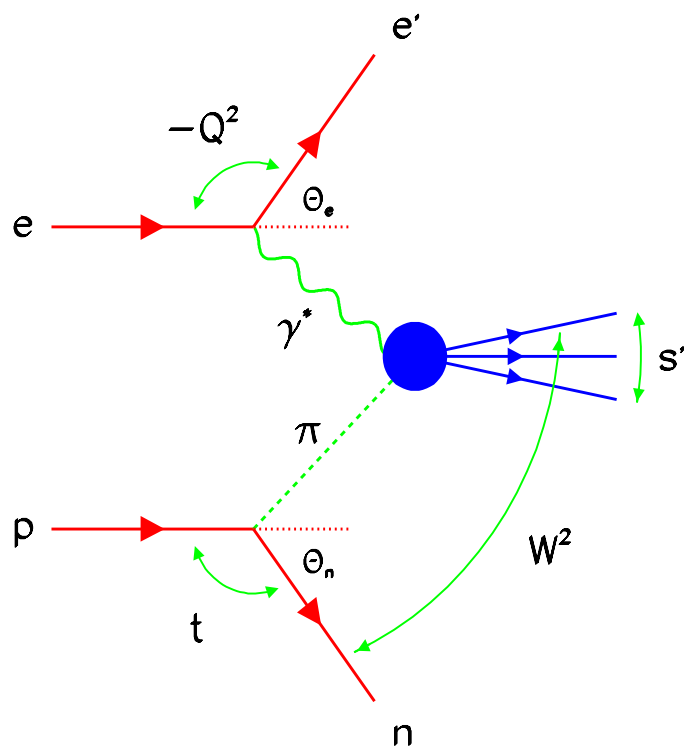
average pion “multiplicity”

$$\langle n \rangle_{\pi N} = 3 \int_0^1 dy f_N^{(\text{on})}(y)$$
$$\sim 0.25 - 0.3$$

- with exception of k_{\perp} cutoff and Bishari models, all others give reasonable fits, $\chi^2 \lesssim 1.5$
- are there other data that can be more discriminating?

Leading neutron production at HERA

- ZEUS & H1 collaborations measured spectra of neutrons produced at very forward angles, $\theta_n < 0.8$ mrad



- can data be described within same framework as E866 asymmetry?
- simultaneous fit never previously been performed!

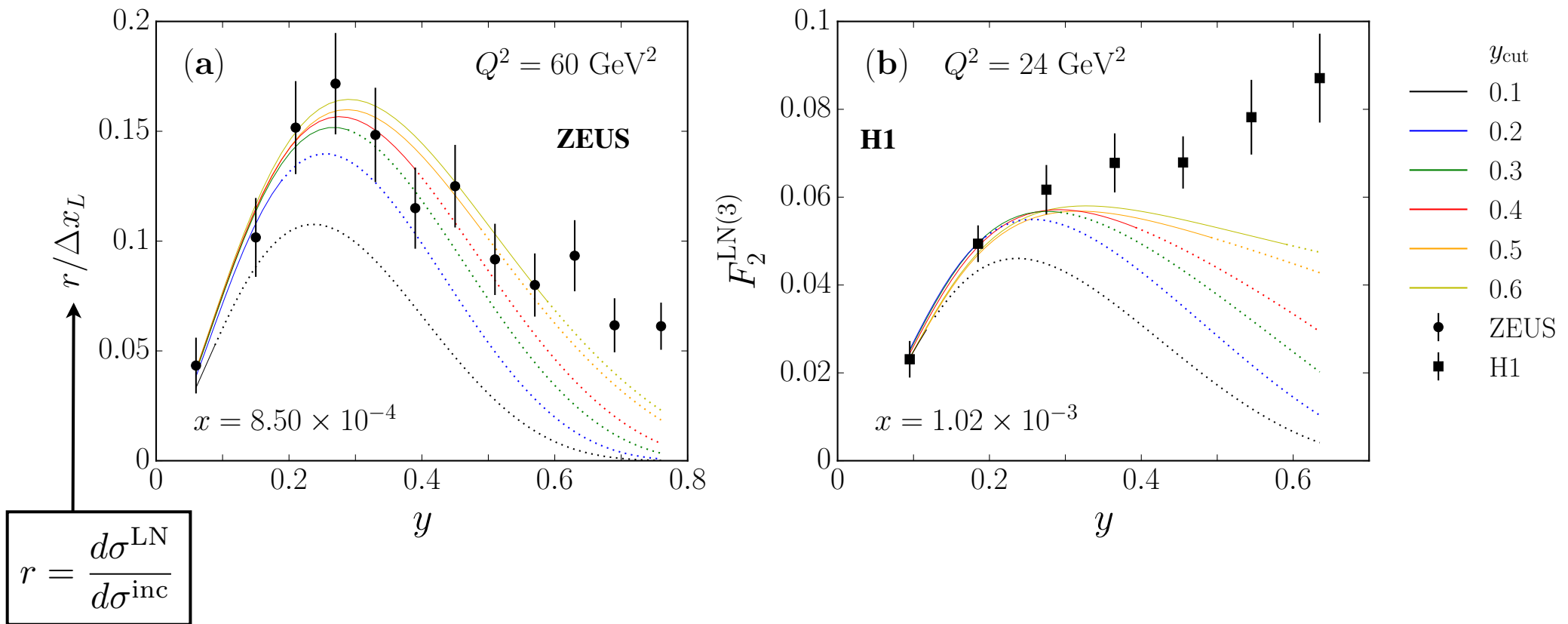
Leading neutron production at HERA

■ Measured LN differential cross section (integrated over p_\perp)

$$\frac{d^3\sigma^{\text{LN}}}{dx dQ^2 dy} \sim F_2^{\text{LN}(3)}(x, Q^2, y)$$

$$2f_N^{(\text{on})}(y) F_2^\pi(x/y, Q^2) \text{ for } \pi \text{ exchange}$$

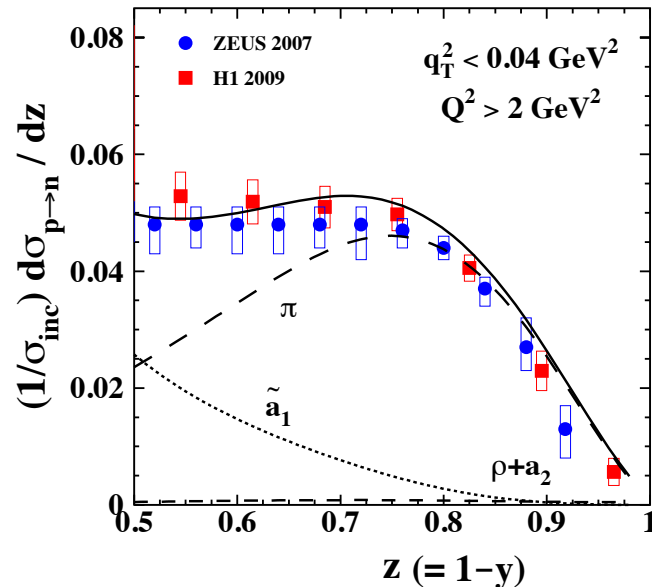
e.g.



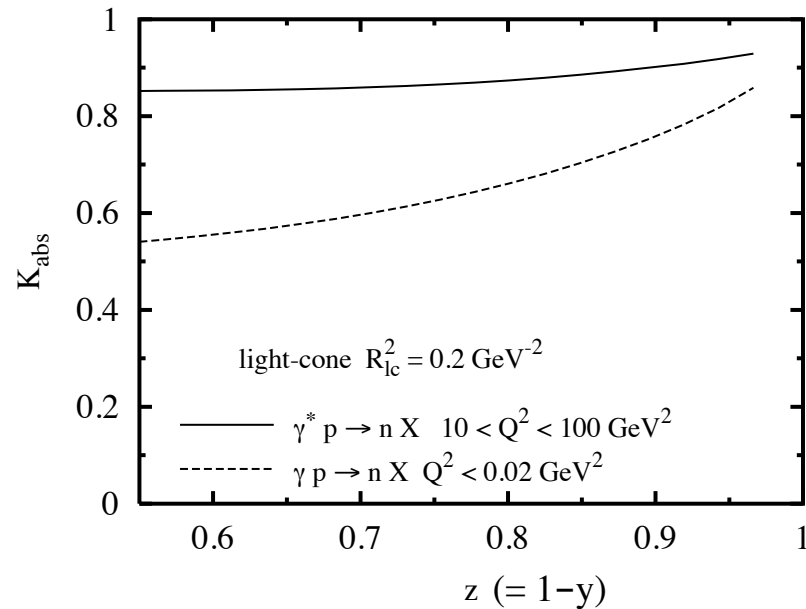
→ quality of fit depends on range of y fitted

Leading neutron production at HERA

- At large y non-pionic mechanisms contribute (*e.g.* heavier mesons, absorption)



Kopeliovich et al. (2012)

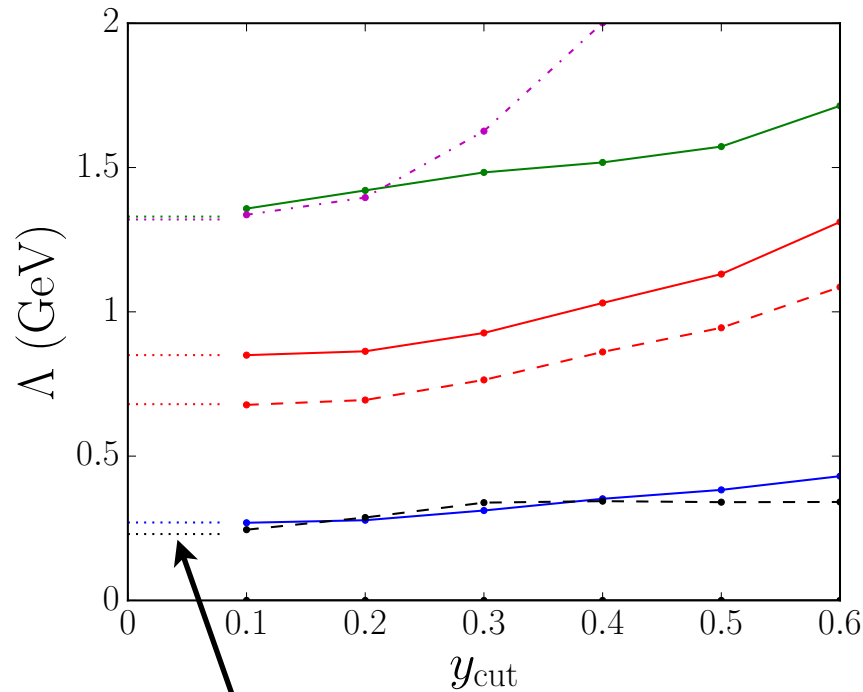


D'Alesio, Pirner (2000)

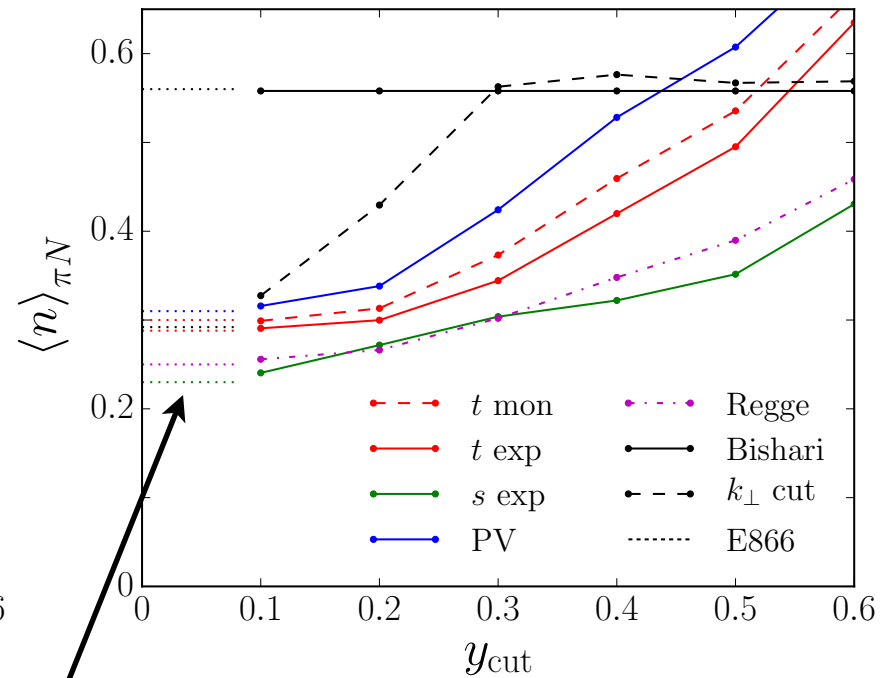
- To reduce model dependence, fit the value of y_{cut} up to which data can be described in terms of π exchange

Leading neutron production at HERA

- Fit requires higher momentum pions with increasing y_{cut}



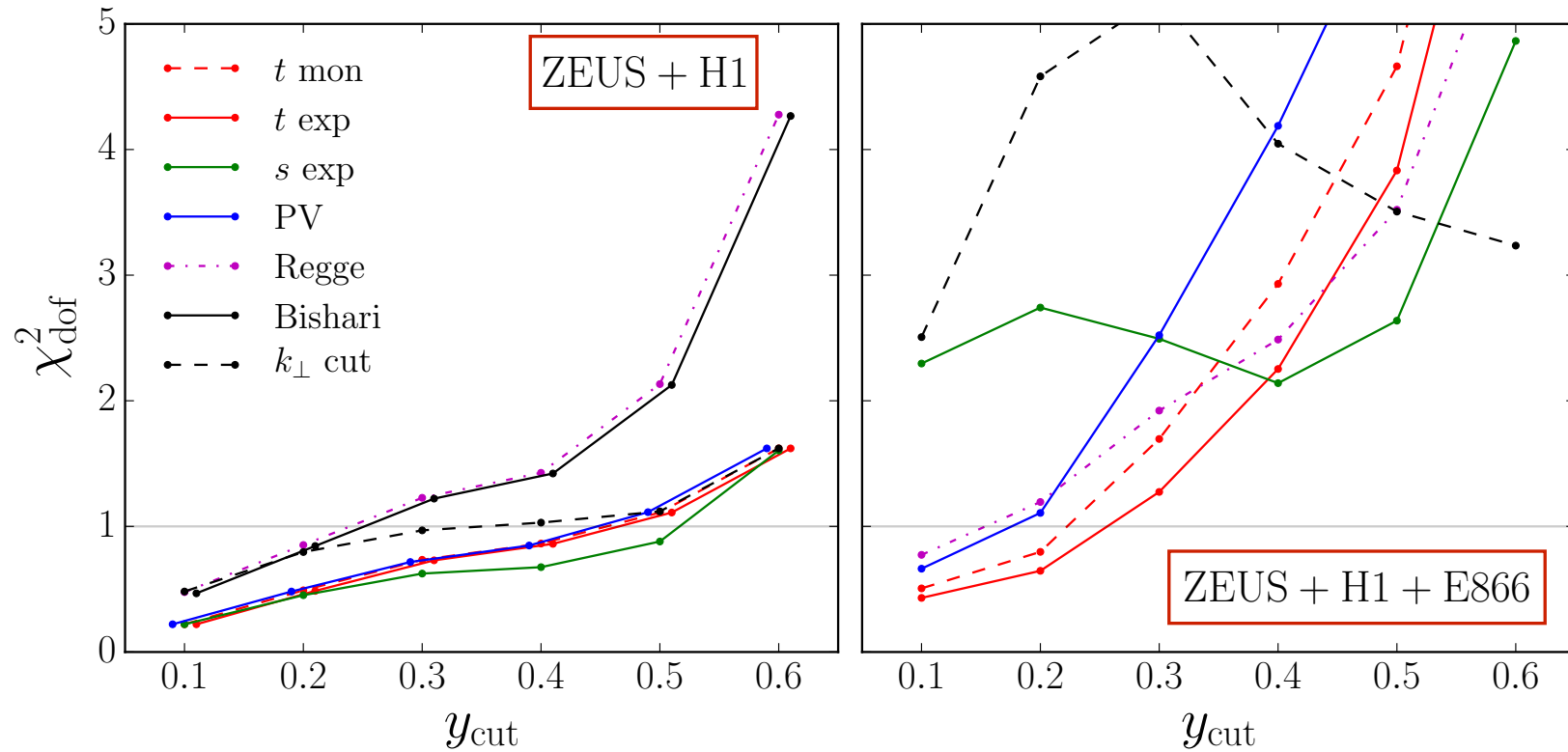
values from fit to E866 data only



→ larger values of y_{cut} more in conflict with E866 data

Leading neutron production at HERA

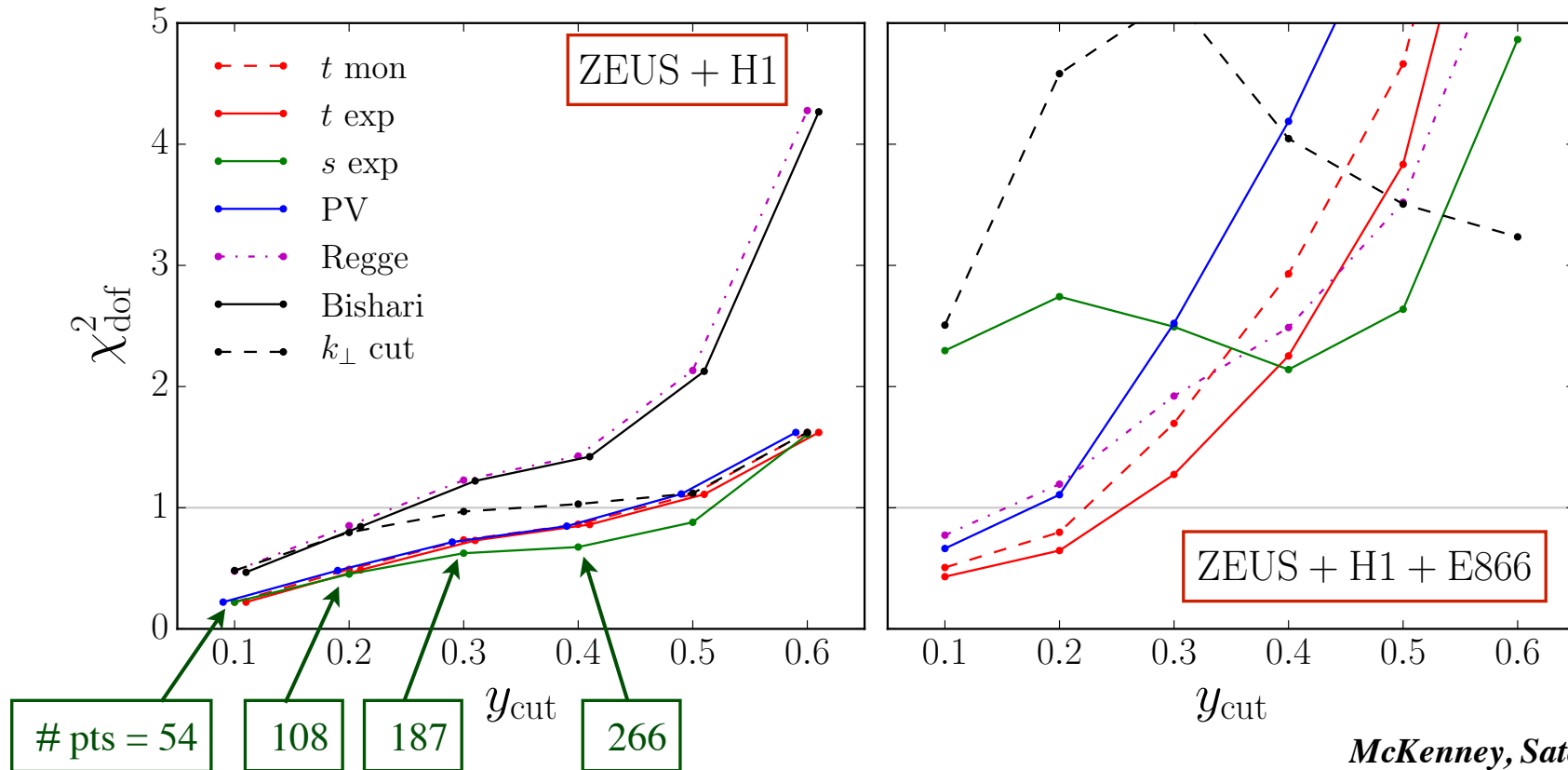
■ Combined fit to HERA LN and E866 Drell-Yan data



McKenney, Sato, WM, Ji (2016)

Leading neutron production at HERA

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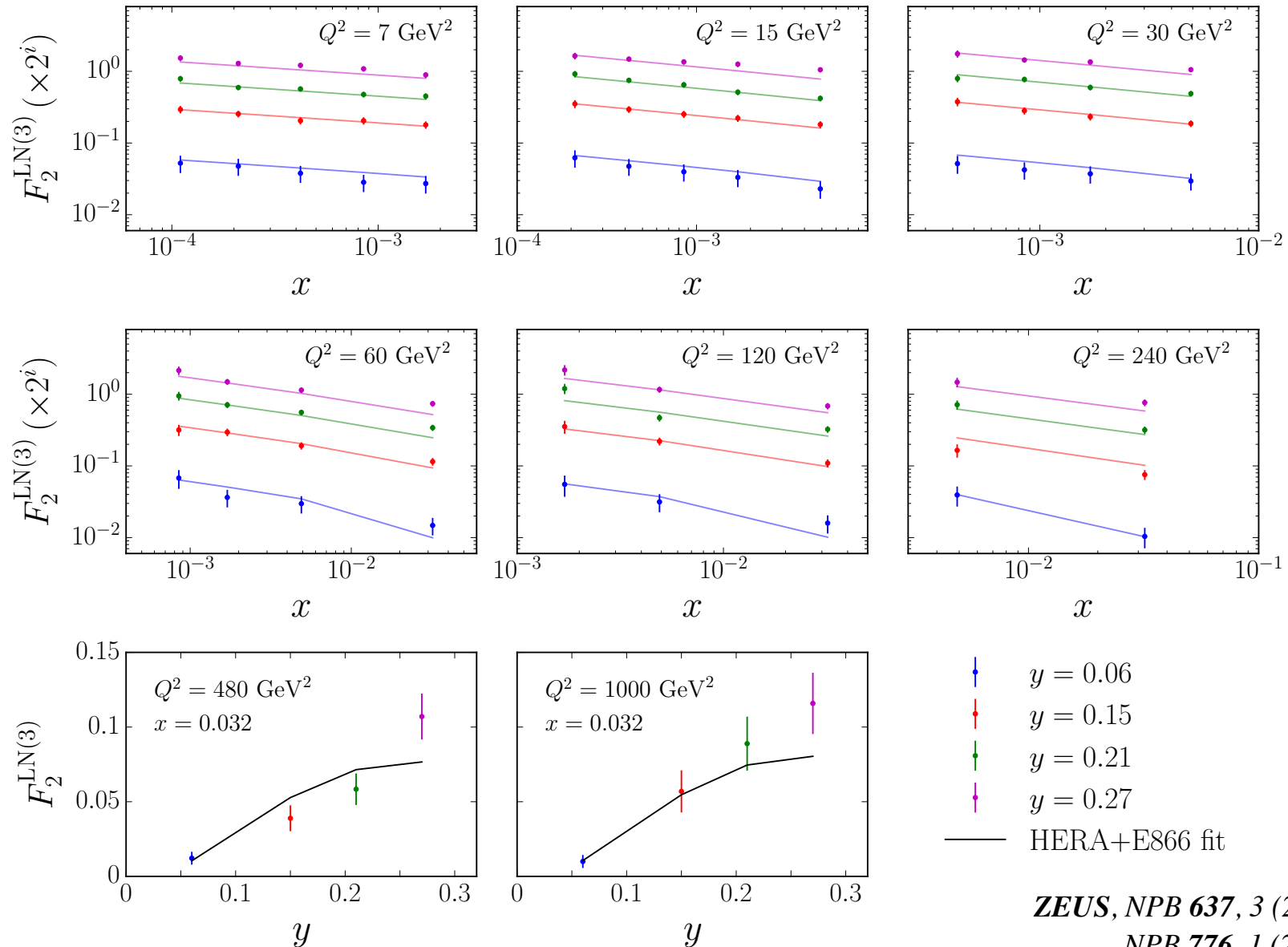


McKenney, Sato, WM, Ji (2016)

→ best fits for largest number of points afforded by t -dependent exponential (and t monopole) regulators

Leading neutron production at HERA

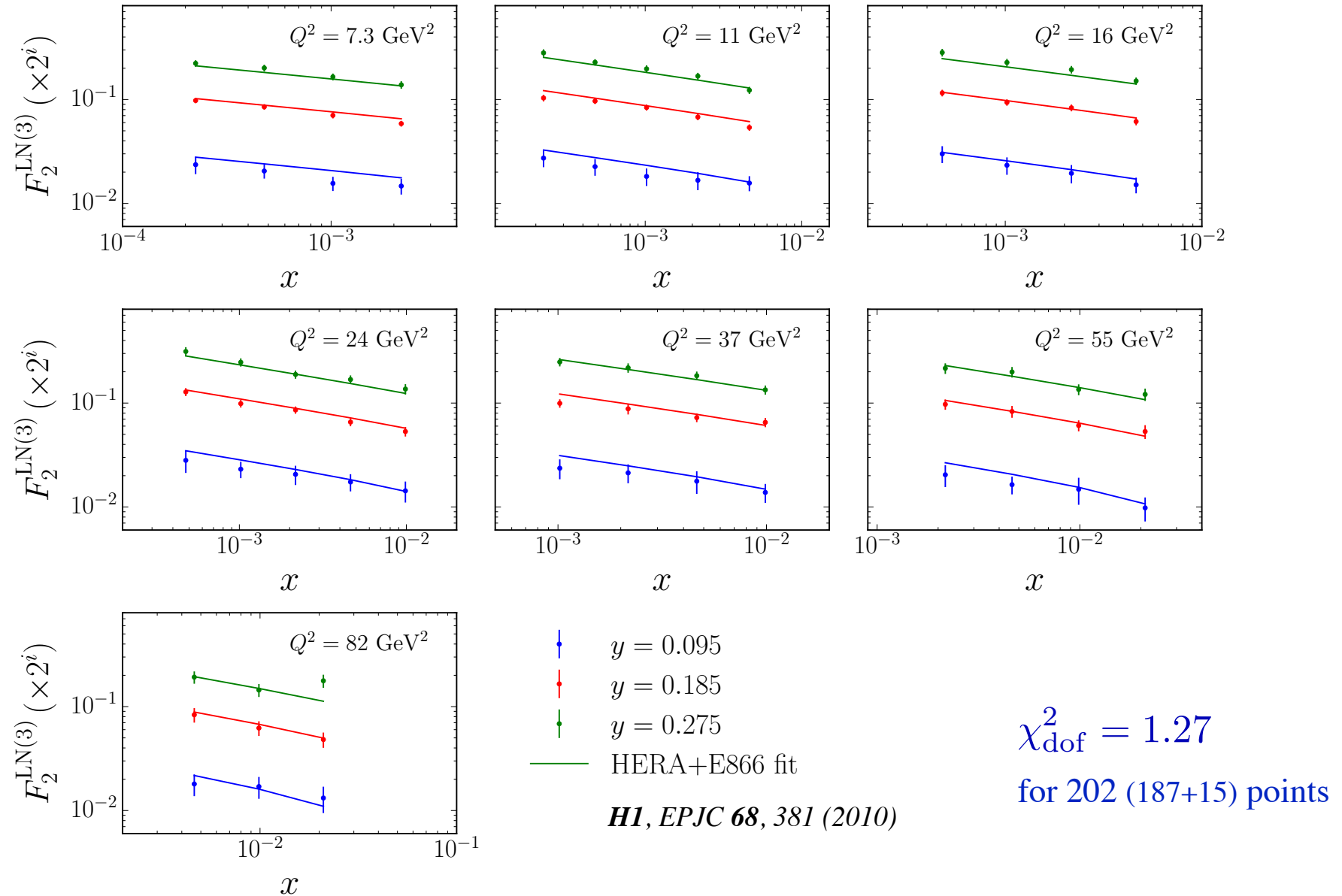
■ Fit to ZEUS LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



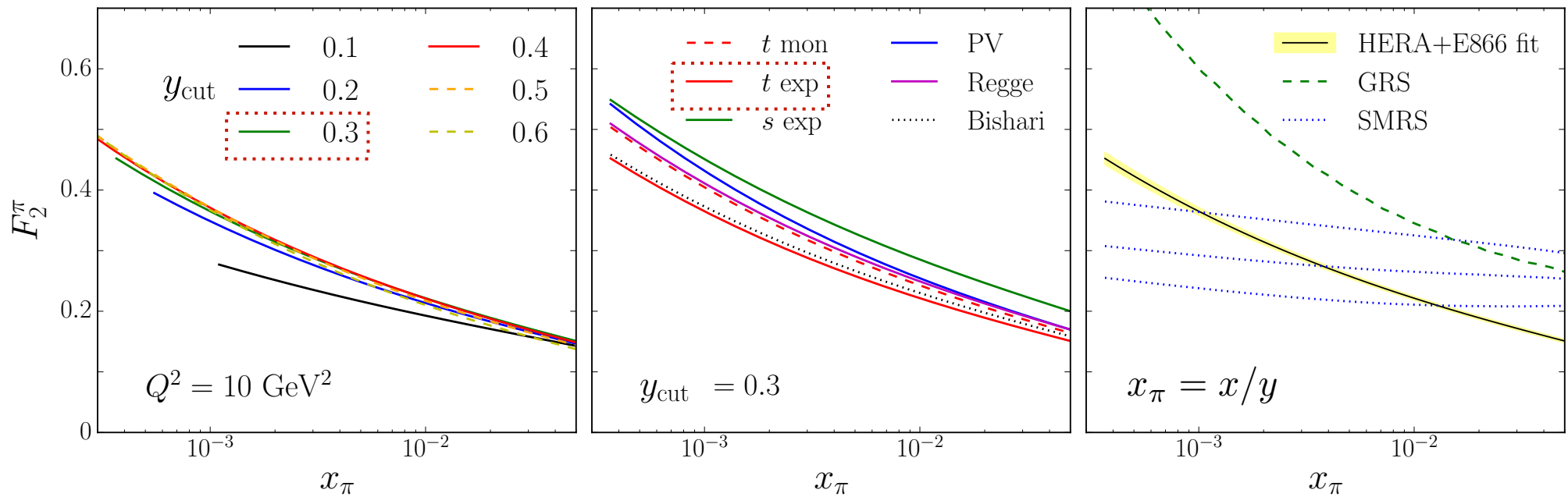
ZEUS, NPB 637, 3 (2002)
NPB 776, 1 (2007)

Leading neutron production at HERA

■ Fit to H1 LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



Extracted pion structure function



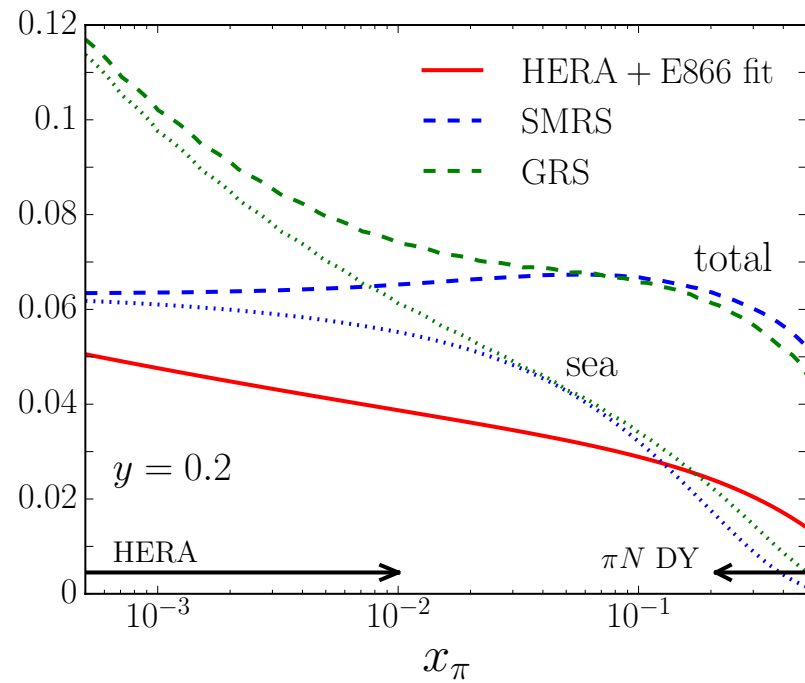
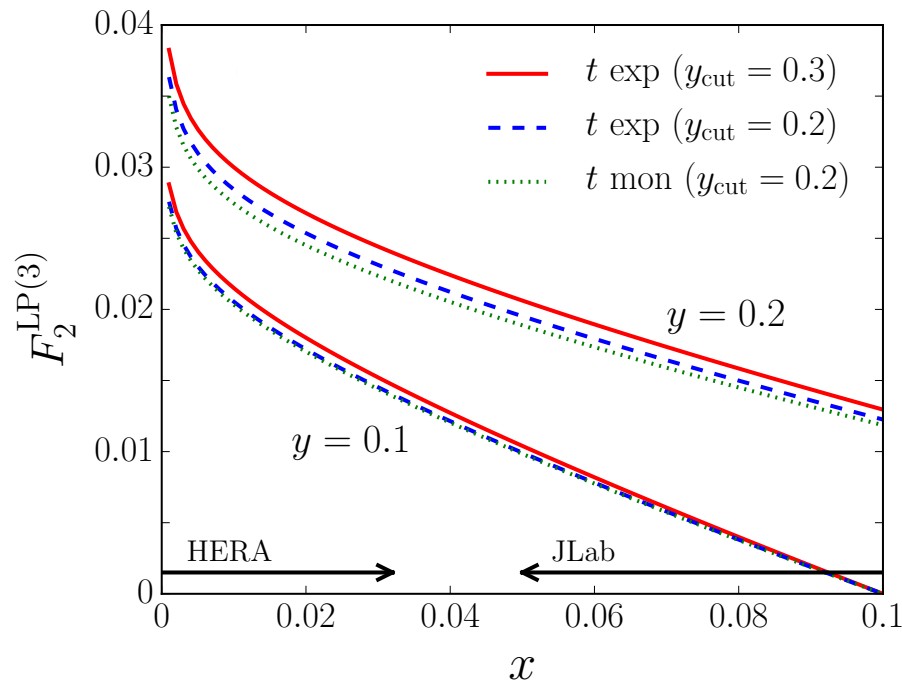
McKenney, Sato, WM, Ji (2016)

$$F_2^\pi = N x_\pi^a (1 - x_\pi)^b, \quad a = a_0 + a_1 \eta$$

$$\eta \sim \log(\log Q^2)$$

- stable values of F_2^π at $4 \times 10^{-4} \lesssim x_\pi \lesssim 0.03$ from combined fit
- shape similar to GRS fit to πN Drell-Yan data (for $x_\pi \gtrsim 0.2$), but smaller magnitude

Predictions at TDIS kinematics



McKenney, Sato, WM, Ji (2016)

→ JLab TDIS experiment can fill gap in x_π coverage between HERA and πN Drell-Yan kinematics

Outlook

- Combined analysis can be extended by including πN DY data
→ constrain large- x_π region ($x_\pi \gtrsim 0.2$)
- Generalize parametrization by fitting individual pion valence and sea quark PDFs, rather than F_2^π
- Longer-term goal is to use all data sensitive to pion structure (including TDIS, EIC) to constrain pion PDFs over full range $10^{-4} \lesssim x_\pi \lesssim 1$
→ global analysis under way of HERA LN, Drell-Yan $\pi N + pd/pp$ (+ future JLab TDIS data) to determine pion PDFs at all x

Patrick Barry, Chueng Ji (NCSU), Nobuo Sato, WM (2016)