

The quest to unveil partonic degrees of freedom in hadrons using high energy reactions

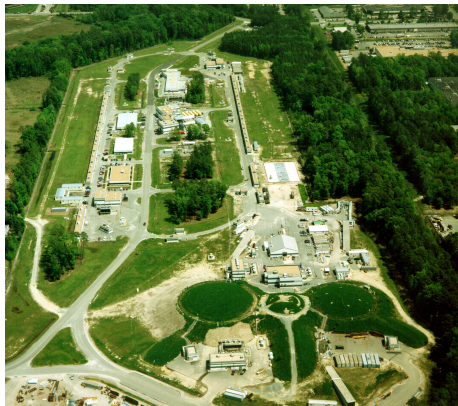
Nobuo Sato

University of Connecticut/JLab

JLab seminar, Mar 1 2017

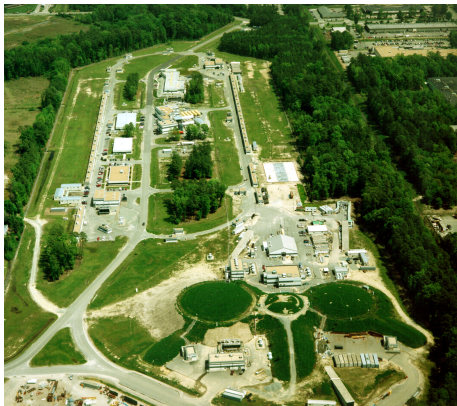
Physics questions for JLab12

- How does QCD work at energy scales of a few GeV?



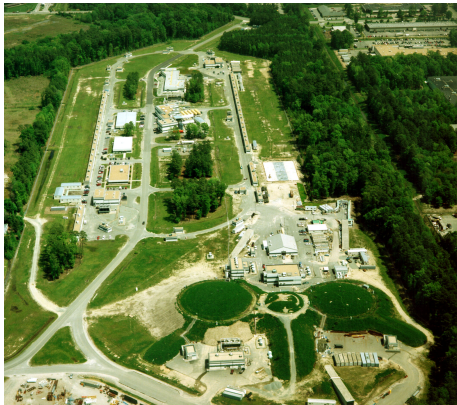
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- Can we use small coupling techniques?



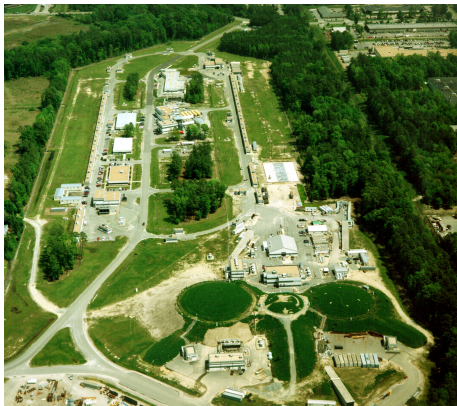
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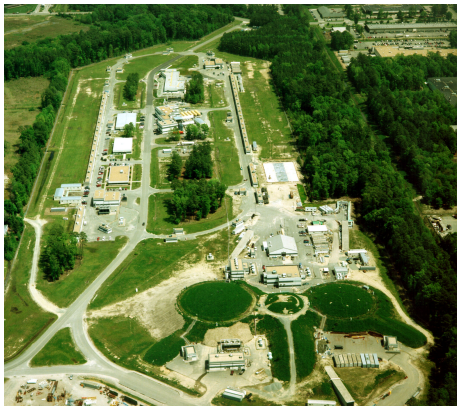
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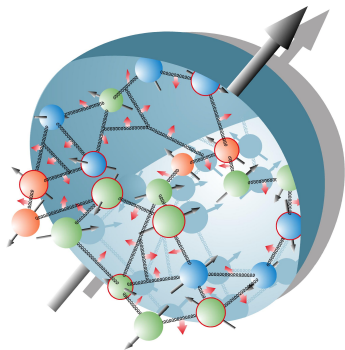
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- Are the factorization theorems valid at such scales?
- How can we model the transition from nonperturbative physics to perturbative physics?
- How do final state partons combine to form hadrons?



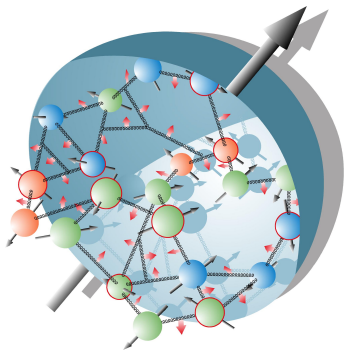
Physics questions for JLab12

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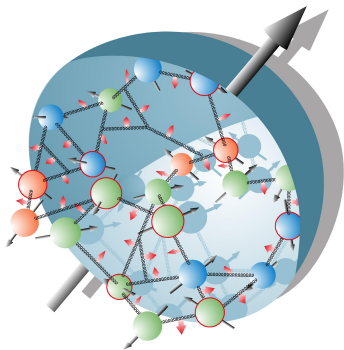
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- What is the detailed partonic structure of the nucleons?



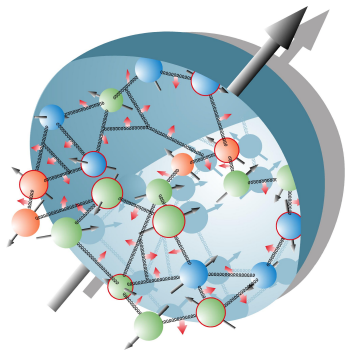
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- How does confinement work?



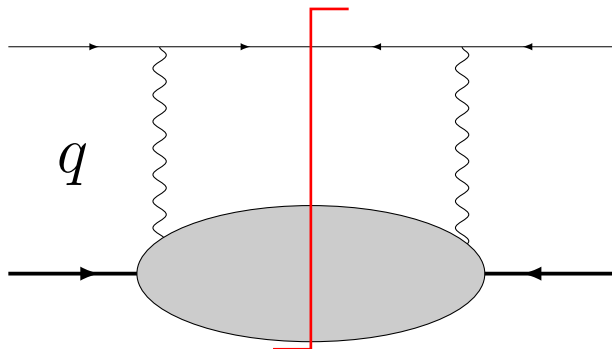
Physics questions for JLab12

- How do the quarks and gluons carry the spin of the proton?
- What is the detailed partonic structure of the nucleons?
- How does confinement work?
- Can we interpret single spin asymmetries?



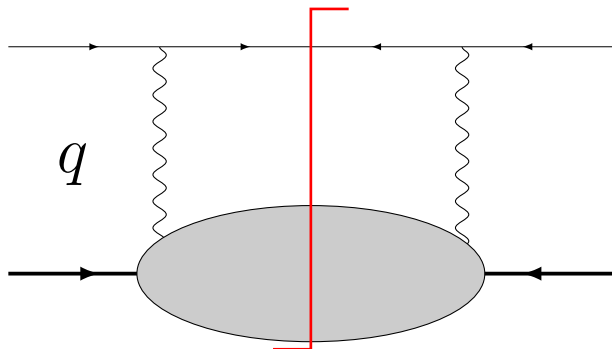
Quick review of QCD framework in DIS

DIS theoretical framework



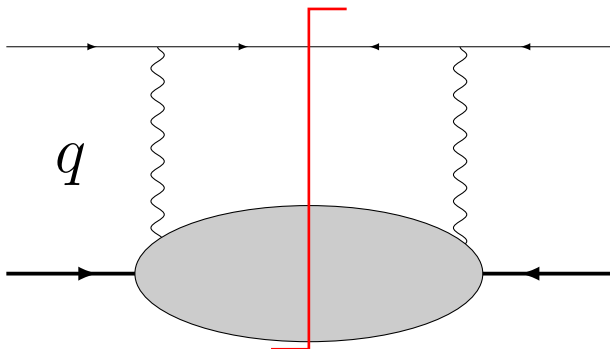
- The upper part of the diagram can be calculated using QED
- The lower part is more complicated → **Factorization**

Factorization in inclusive DIS



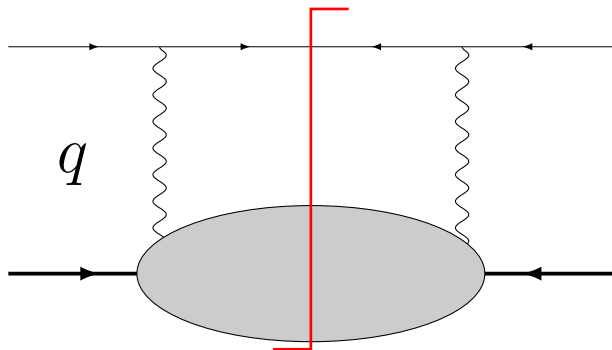
- The lower blob represents a collection of Feynman diagrams

Factorization in inclusive DIS



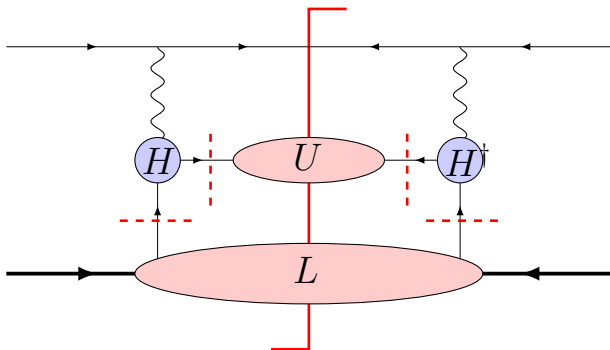
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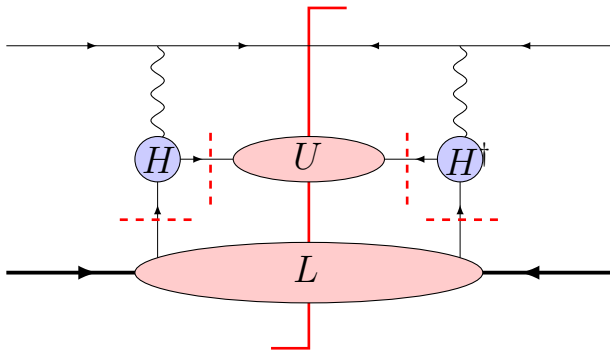
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- Using kinematic approximations one can disentangle the various regions assuming $Q^2 = -q^2$ is large

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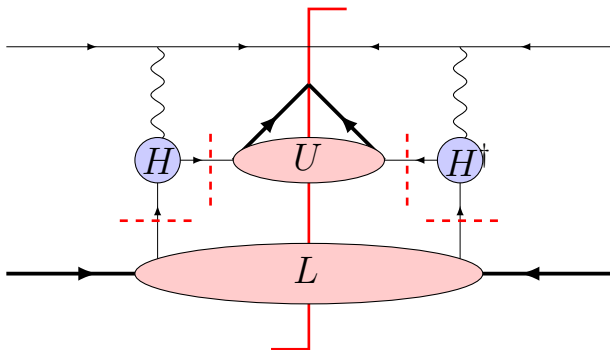
Factorization in inclusive DIS



Schematically

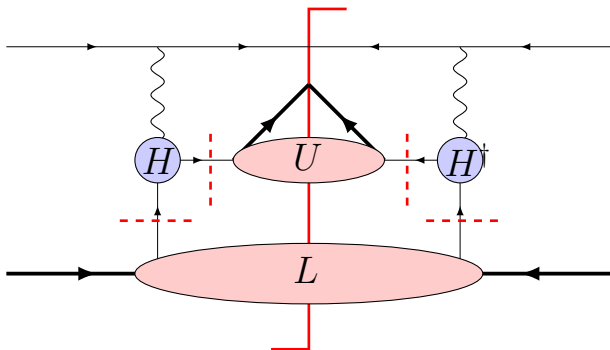
$$d\sigma \sim \underbrace{|H|^2}_{\text{Hard}} \underbrace{\int dk^- dk^\perp L(xP^+, k^-, k^\perp)}_{\text{Parton densities}} \underbrace{\int dl^+ U(l^+) d\Omega}_{\text{Jet function}}$$

Factorization in SIDIS



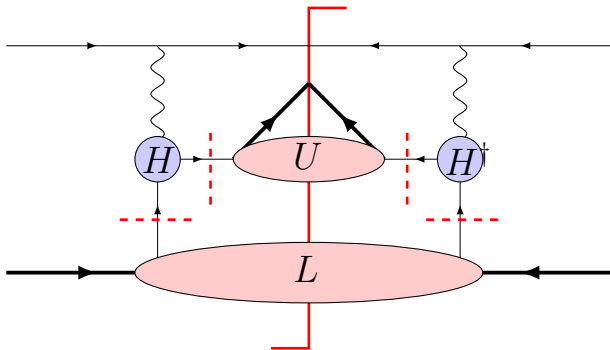
- Additional information about the lower blob can be accessed by measuring hadrons in the final state (i.e., flavor separation, intrinsic transverse momentum)

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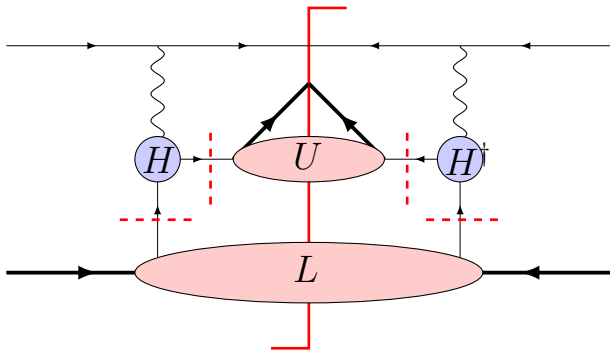
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- The upper blob along with tagged hadrons characterizes formation of hadrons from partons

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- The upper blob along with tagged hadrons characterizes formation of hadrons from partons
- Factorization of these regions assumes large $Q^2 = -q^2$

Factorization in SIDIS



Schematically

$$d\sigma \sim \underbrace{|H|^2}_{\text{Hard}} \underbrace{\int dk^- dk^\perp L(xP^+, k^-, k^\perp)}_{\text{Parton densities}} \underbrace{\int dk_h^+ dk_h^\perp U(k_h^+, k_h^\perp) d\Omega}_{\text{Fragmentation function}}$$

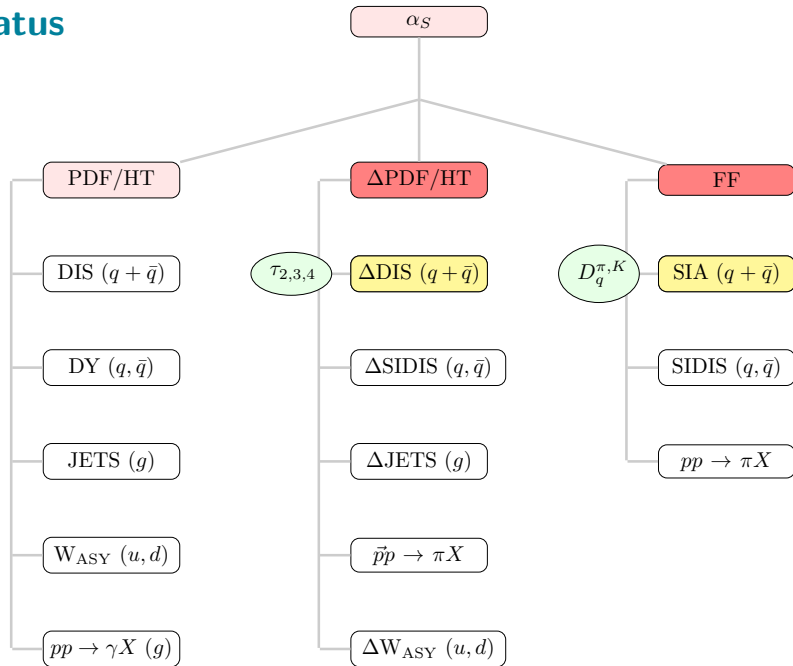
Longitudinal distributions

The collinear master plan



- Extract reliable collinear polarized and unpolarized parton distribution functions (PDFs) and fragmentation functions (FFs)
- Improvements in fitting methodology:
Iterative Monte Carlo
- Reliable description of q_T integrated SIDIS data.

Status



Milestones

- Global analysis of all inclusive polarized DIS data

Iterative Monte Carlo analysis of spin-dependent parton distributions

Nobuo Sato,¹ W. Melnitchouk,¹ S. E. Kuhn,² J. J. Ethier,³ and A. Accardi^{4,1}
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■ Global analysis of all SIA data

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■ Invited talk in FCCee workshop

Parton Radiation and Fragmentation from LHC to FCC-ee

Workshop Proceedings, CERN, Geneva, Nov. 22nd–23th, 2016

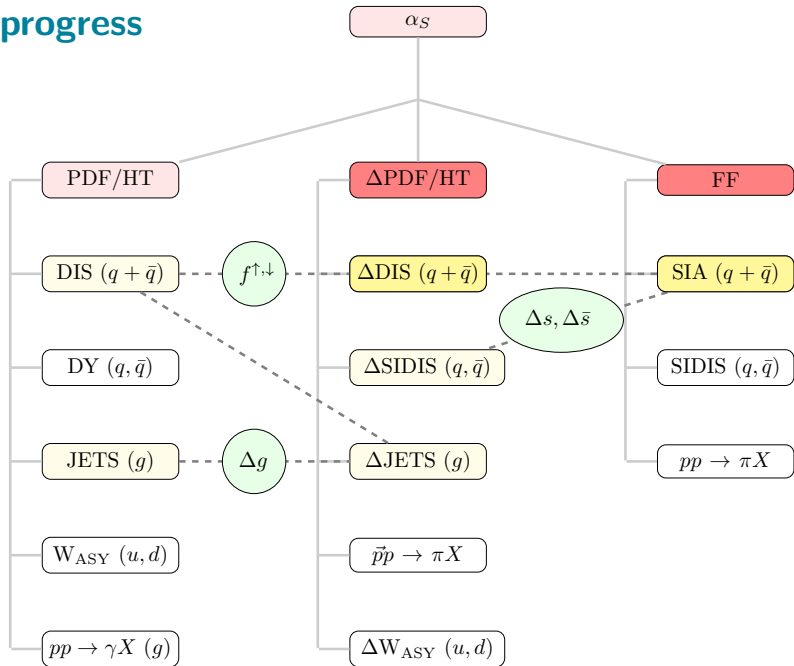
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David d'Enterria (CERN), Peter Z. Skands (Monash)

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S. Jadach (IFJ-PAN Krakow), T. Kaufmann (U. Tübingen), S. Kluth (T.U. München).

In progress



Innovations in fitting methodology

Theory of fitting

The goal is to estimate:

$$E[\mathcal{O}] = \int d^n a \mathcal{P}(\mathbf{a}|data) \mathcal{O}(\mathbf{a})$$

$$V[\mathcal{O}] = \int d^n a \mathcal{P}(\mathbf{a}|data) [\mathcal{O}(\mathbf{a}) - E[\mathcal{O}]]^2$$

Monte Carlo methods

- $\mathcal{P}(\mathbf{a}|data) \rightarrow \{\mathbf{a}_k\}$
- $E[\mathcal{O}] \approx \frac{1}{N} \sum_k \mathcal{O}(\mathbf{a}_k)$
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Maximum Likelihood

- Maximize $\mathcal{P}(\mathbf{a}|data) \rightarrow \mathbf{a}_0$
- $E[\mathcal{O}] \approx \mathcal{O}(\mathbf{a}_0)$
- $V[\mathcal{O}] \approx \text{hessian}, \Delta\chi^2 \text{ envelope}, \dots$

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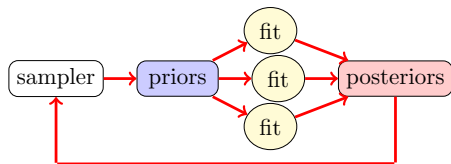
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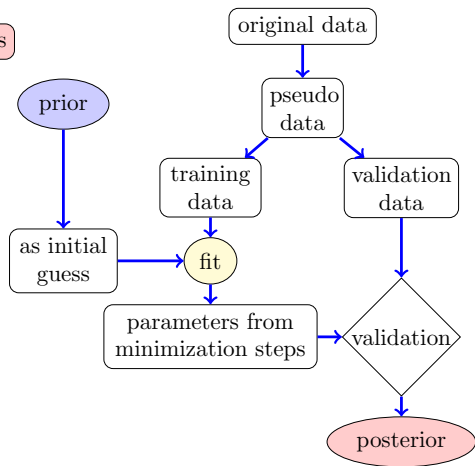
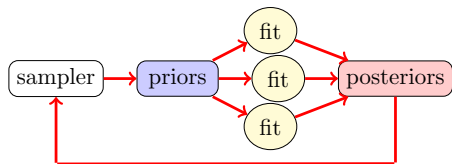
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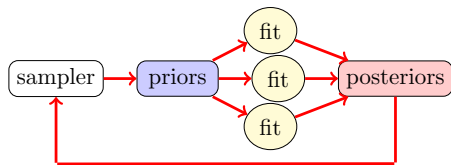
Iterative Monte Carlo analysis (IMC)



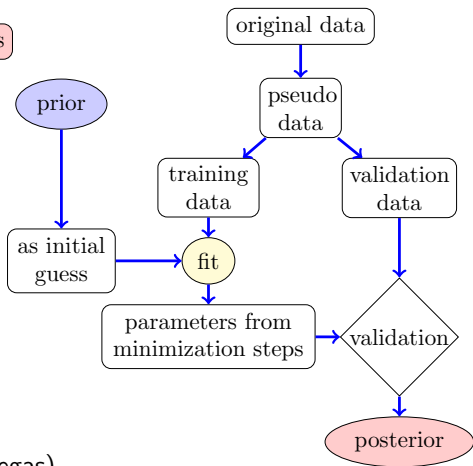
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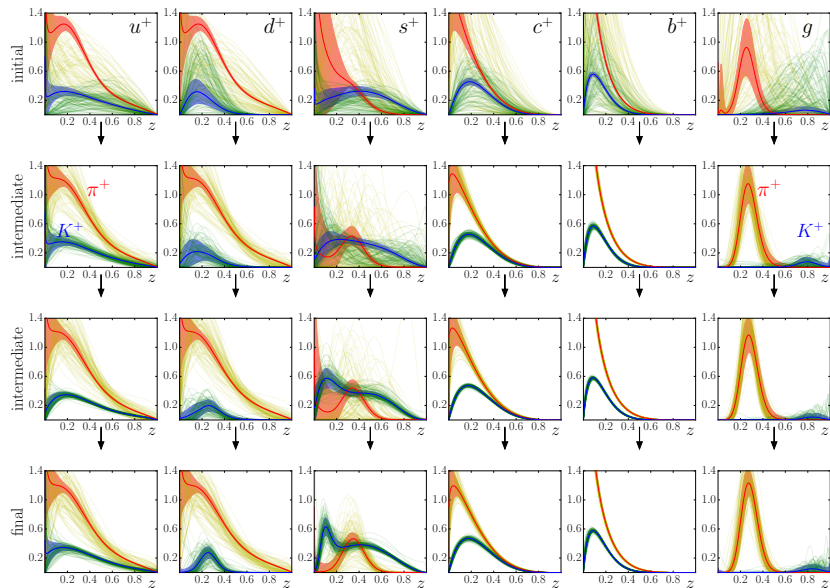
Iterative Monte Carlo analysis (IMC)



- Use traditional ansatz
$$xf(x) = Nx^a(1-x)^b(1+c\sqrt{x}+dx)$$
- Keep all the parameters free.
No assumptions on the exponents
- Avoid over-fitting by Cross-Validation
- **Iterative procedure**
→ Adaptive MC integration (like in Vegas)
- Robust estimation of uncertainties

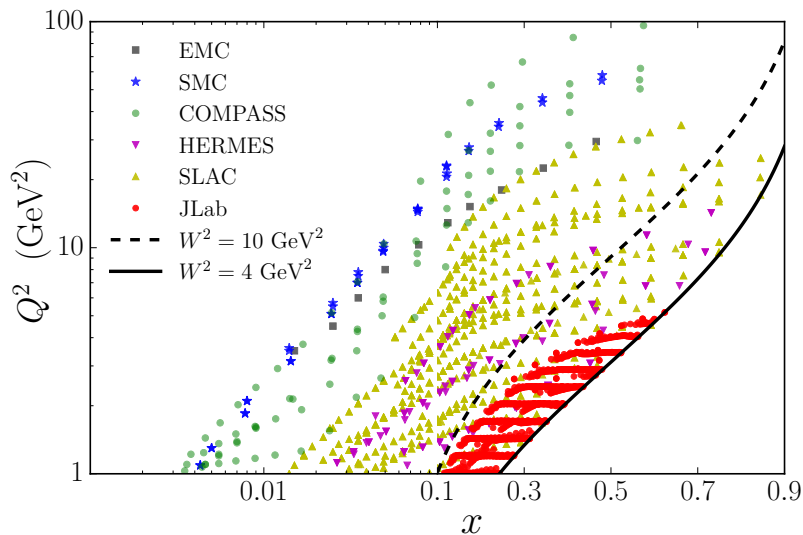


IMC in action (e.g., FFs analysis)

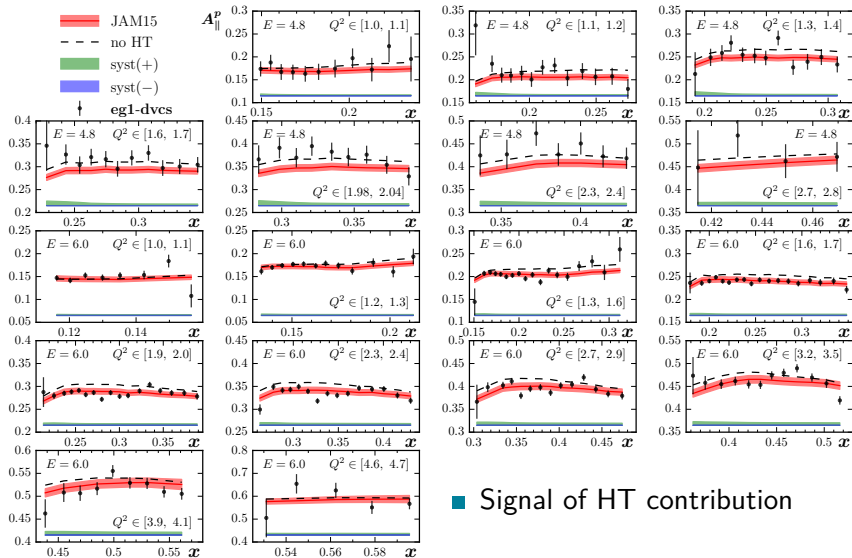


Spin PDFs from polarized DIS

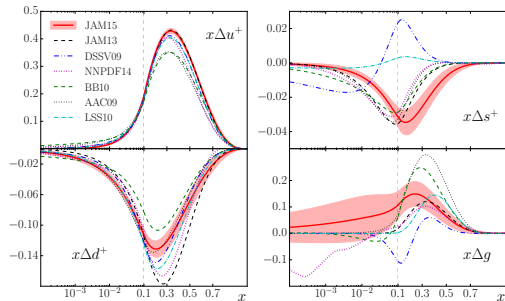
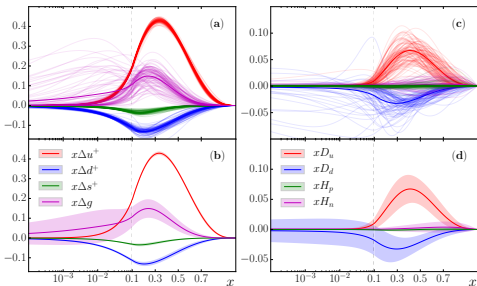
Global polarized DIS data



Data vs theory: proton JLab eg1-dvcs



Results

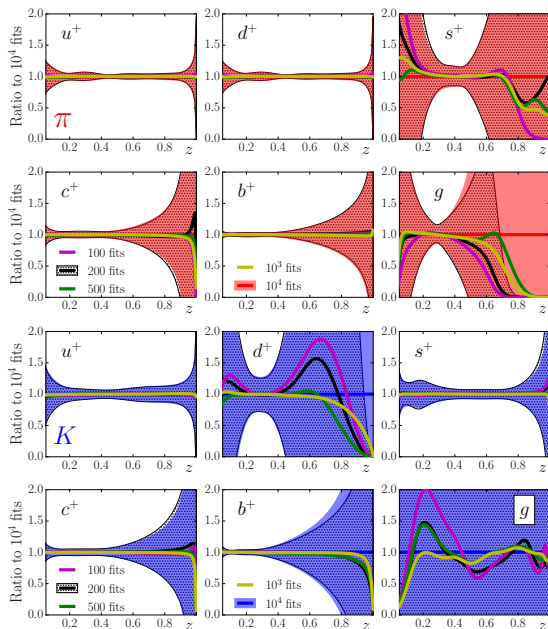


moment	truncated	full
Δu^+	0.82 ± 0.01	0.83 ± 0.01
Δd^+	-0.42 ± 0.01	-0.44 ± 0.01
Δs^+	-0.10 ± 0.01	-0.10 ± 0.01
$\Delta \Sigma$	0.31 ± 0.03	0.28 ± 0.04
ΔG	0.5 ± 0.4	1 ± 15
d_2^p	0.011 ± 0.004	0.011 ± 0.004
d_2^n	-0.002 ± 0.002	-0.002 ± 0.002
h_p	-0.000 ± 0.001	0.000 ± 0.001
h_n	0.001 ± 0.002	0.001 ± 0.003

- $\chi^2/N_{npts} = 1.07$
- Sign of τ_3 distributions is the same as τ_2
- **Negative Δs^+**
- Δg compatible with the most recent DSSV fits
- Moment of Δg not constrained

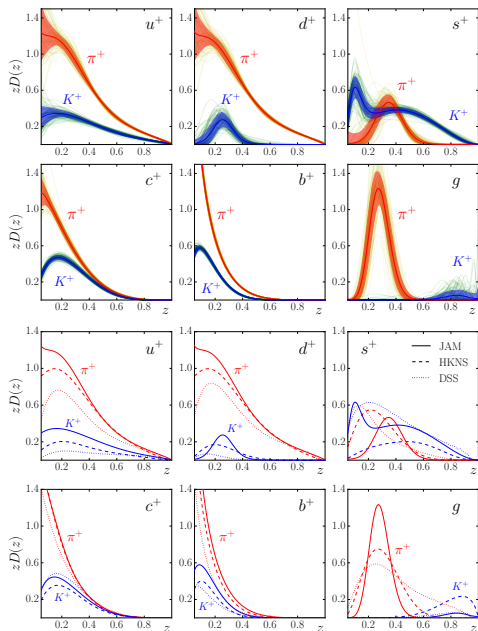
Fragmentation Functions from SIA data

Fragmentation functions



- Stronger constraints on favored FFs than unfavored FFs
- More sensitivity on D_g^π than D_g^K
- D_g^K is unknown
- Convergence attained with ~ 200 posteriors

Fragmentation functions

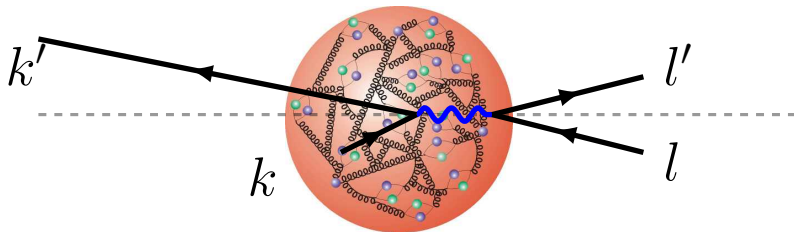


- D_g^π and D_g^K behave differently
- HQ distributions becomes comparable in size with favored distributions
- JAM $D_{s^+}^K$ is compatible with DSS. **Will it change the sign of Δ_{s^+} ?**

TMD physics

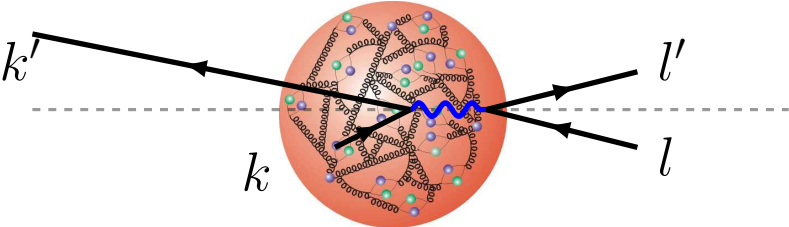
Intuitive picture

The Breit frame: $q = (0, 0, 0, -Q)$ and $P = (P^0, 0, 0, P_z)$



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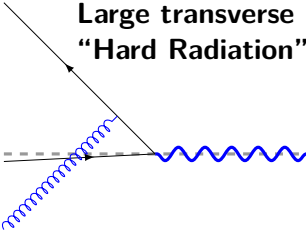
The Breit frame: $q = (0, 0, 0, -Q)$ and $P = (P^0, 0, 0, P_z)$



**Small transverse mom.
"Intrinsic"**



**Large transverse mom.
"Hard Radiation"**



Formal developments

- Factorization in TMD observables

$$\Gamma = d\sigma/dq_T$$
$$q_T = p^h/z$$

$$\begin{aligned}\Gamma &= \Gamma \\ &= \mathbf{T}_{\text{TMD}}\Gamma + [\Gamma - \mathbf{T}_{\text{TMD}}\Gamma] \\ &= \underbrace{\mathbf{T}_{\text{TMD}}\Gamma}_{\mathbf{W}} + \underbrace{\mathbf{T}_{\text{coll}}[\Gamma - \mathbf{T}_{\text{TMD}}\Gamma]}_{\mathbf{Y}} + \mathcal{O}(m^2/Q^2)\Gamma\end{aligned}$$

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- Region of $q_T \gtrsim Q$

- Collinear approx. dominates $\rightarrow \Gamma \approx \mathbf{T}_{\text{coll}}\Gamma$
- At large Q , $\mathbf{T}_{\text{TMD}}\Gamma$ is mostly perturbative

Formal developments

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$$\mathbf{W} = \mathbf{T}_{\text{TMD}}\Gamma$$

$$\mathbf{FO} = \mathbf{T}_{\text{coll}}\Gamma$$

$$\mathbf{ASY} = \mathbf{T}_{\text{coll}}\mathbf{T}_{\text{TMD}}\Gamma$$

$$\mathbf{Y} = \mathbf{FO} - \mathbf{ASY}$$

SIDIS (One of the main programs of JLab12)

■ Cross section and structure functions

$$\frac{d^5\sigma(S_\perp)}{dx_B dQ^2 dz_h d^2P_{h\perp}} = \sigma_0 \left[F_{UU} + \sin(\phi_h - \phi_s) F_{UT}^{\sin(\phi_h - \phi_s)} \right. \\ \left. + \sin(\phi_h + \phi_s) \frac{2(1-y)}{1+(1-y)^2} F_{UT}^{\sin(\phi_h + \phi_s)} + \dots \right]$$

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■ CSS formalism

$$F_{UU} = H_{\text{SIDIS}} \frac{1}{z_h^2} \int_0^\infty \frac{db b}{(2\pi)} J_0(q_{h\perp} b) \widetilde{W}_{UU}(b_*) + Y_{UU}$$
$$F_{UT}^{\sin(\phi_h - \phi_s)} = -H_{\text{SIDIS}} \frac{M_P}{z_h^2} \int_0^\infty \frac{db b^2}{(2\pi)} J_1(q_{h\perp} b) \widetilde{W}_{UT}^{\sin(\phi_h - \phi_s)}(b_*) + Y_{UT}^{\sin(\phi_h - \phi_s)}$$
$$F_{UT}^{\sin(\phi_h + \phi_s)} = H_{\text{SIDIS}} \frac{M_h}{z_h^2} \int_0^\infty \frac{db b^2}{(2\pi)} J_1(q_{h\perp} b) \widetilde{W}_{UT}^{\sin(\phi_h + \phi_s)}(b_*) + Y_{UT}^{\sin(\phi_h + \phi_s)}$$

SIDIS: small transverse momentum

■ W term formulation in b_T space

$$\widetilde{W}_{UU}(b_*) \equiv e^{-S_{pert}(Q,b_*) - S_{NP}^{f_1}(Q,b) - S_{NP}^{D_1}(Q,b)} \widetilde{F}_{UU}(b_*)$$

$$\widetilde{W}_{UT}^{\sin(\phi_h - \phi_s)}(b_*) \equiv e^{-S_{pert}(Q,b_*) - S_{NP}^{f_{1T}^\perp}(Q,b) - S_{NP}^{D_1}(Q,b)} \widetilde{F}_{UT}^{\sin(\phi_h - \phi_s)}(b_*)$$

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$$\widetilde{W}_{UT}^{\sin(\phi_h - \phi_s)}(b_*) \equiv e^{-S_{pert}(Q,b_*) - S_{NP}^{f_{1T}^\perp}(Q,b) - S_{NP}^{D_1}(Q,b)} \widetilde{F}_{UT}^{\sin(\phi_h - \phi_s)}(b_*)$$

$$\widetilde{W}_{UT}^{\sin(\phi_h + \phi_s)}(b_*) \equiv e^{-S_{pert}(Q,b_*) - S_{NP}^{h_1}(Q,b) - S_{NP}^{H_1^\perp}(Q,b)} \widetilde{F}_{UT}^{\sin(\phi_h + \phi_s)}(b_*)$$

■ Small b_T contribution

$$\widetilde{F}_{UU}(b_*) = \sum_q e_q^2 \left(C_{q \leftarrow i}^{f_1} \otimes f_1^i(x_B, \mu_b) \right) \left(\hat{C}_{j \leftarrow q}^{D_1} \otimes D_{h/j}(z_h, \mu_b) \right)$$

$$\widetilde{F}_{UT}^{\sin(\phi_h - \phi_s)}(b_*) = \sum_q e_q^2 \left(C_{q \leftarrow i}^{f_{1T}^\perp} \otimes f_{1T}^{\perp(1)i}(x_B, \mu_b) \right) \left(\hat{C}_{j \leftarrow q}^{D_1} \otimes D_{h/j}(z_h, \mu_b) \right)$$

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SIDIS: small transverse momentum

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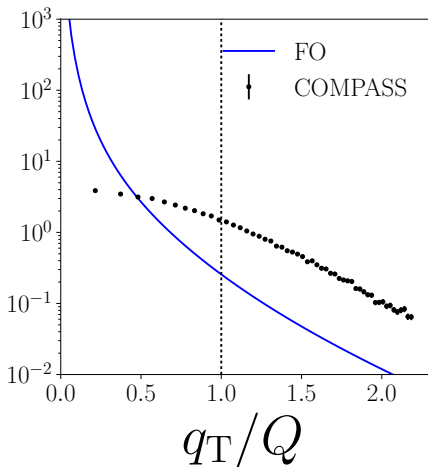
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■ Collinear distribution are important in TMDs

Does it work?

$$\frac{\frac{d\sigma}{dx dz dQ^2 dp_T^2}}{\frac{d\sigma}{dx dQ^2}}$$



■ Kinematics

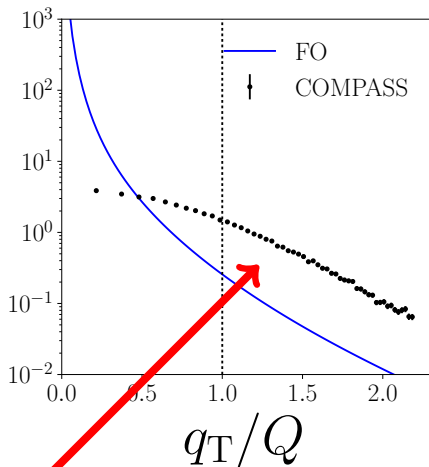
$$Q^2 = 1.92 \text{ GeV}^2$$

$$x = 0.0318$$

$$z = 0.375$$

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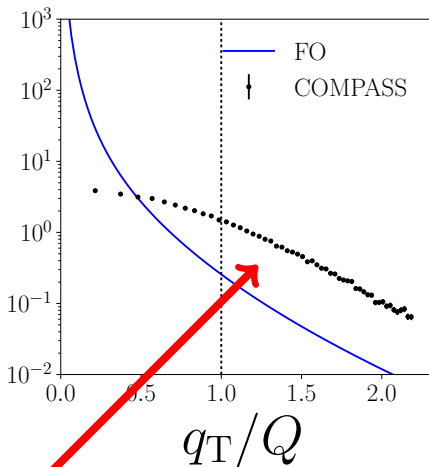
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- Need order α_S^2 or beyond?
- Soft gluon resummation?
- Subleading power corrections?

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The unpolarized SIDIS cross sections needs to be ready to interpret upcoming TMD data from JLab 12

Progress

- Improve the matching \rightarrow extensions to $W + Y$

PHYSICAL REVIEW D **94**, 034014 (2016)

Relating transverse-momentum-dependent and collinear factorization theorems in a generalized formalism

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- SIDIS Kinematics analysis



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Physics Letters B

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- Applicability of collinear factorization at low energies

What are the moderate- Q and large- x limits of collinear QCD factorization theorems?

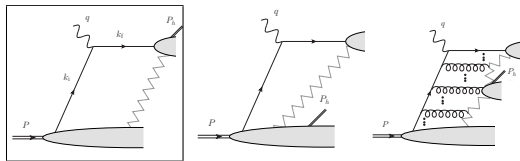
E. Moffat,^{1,*} W. Melnitchouk,^{2,†} T. C. Rogers,^{1,2,‡} and N. Sato^{2,§}

SIDIS kinematics analysis

- Can we apply factorization theorems in SIDIS measurements?

SIDIS kinematics analysis

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- Factorization demands that



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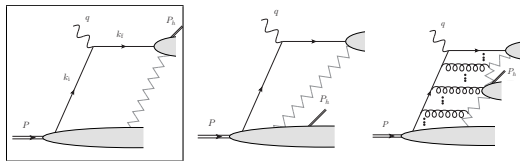
$$p_h \cdot k_i = \mathcal{O}(Q^2)$$

- Define a *collinearity* parameter

$$R = \frac{(p_h \cdot k_f)}{(p_h \cdot k_i)} = \mathcal{O}(m^2/Q^2)$$

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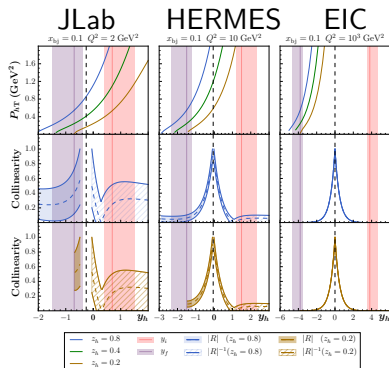


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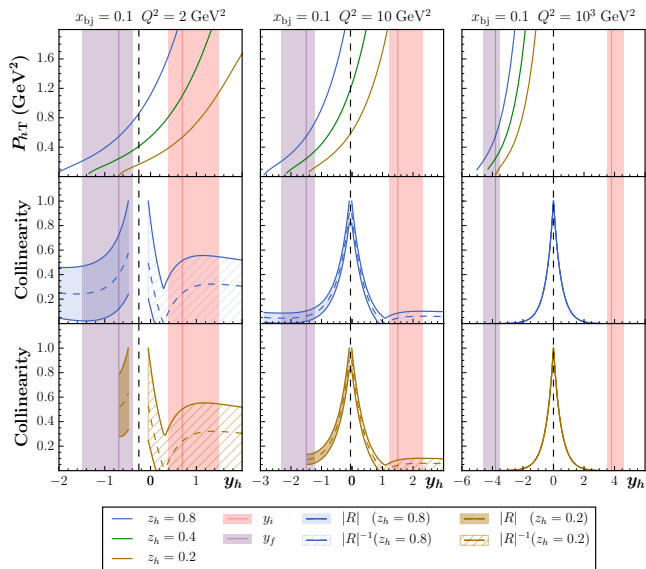
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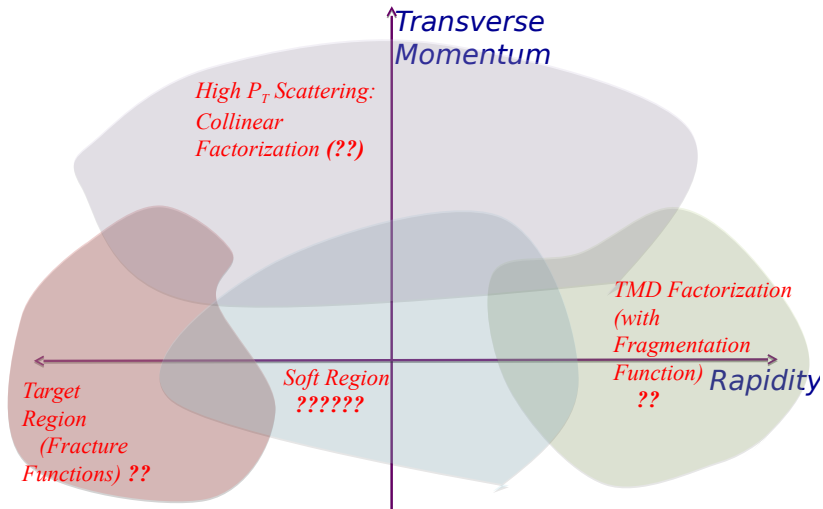
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SIDIS kinematics analysis



Map of Leading Twist SIDIS



$\mathcal{O}(\alpha_S^2)$ corrections for the Y term

- NLO QCD correction to q_T dependent SIDIS
- The steps:
 1. Computation of the Born contribution $d\sigma_{\text{Born}}$

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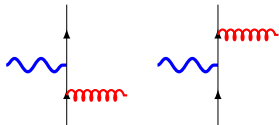
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5. $d\sigma = d\sigma_{\text{Born}} + d\sigma_{\text{Virtual}} + d\sigma_{\text{Real}} + d\sigma_{\text{CCS}} \rightarrow$ IR finite

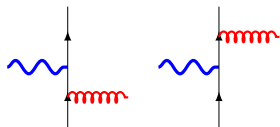
$\mathcal{O}(\alpha_s^2)$ corrections to $\gamma^* + q \rightarrow q + g$

Born contribution

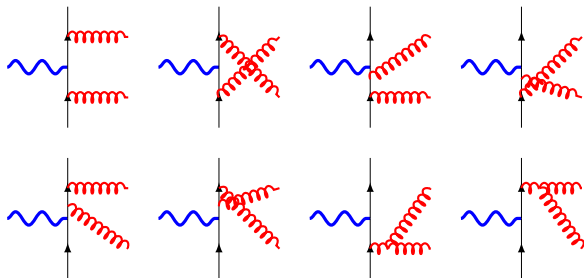


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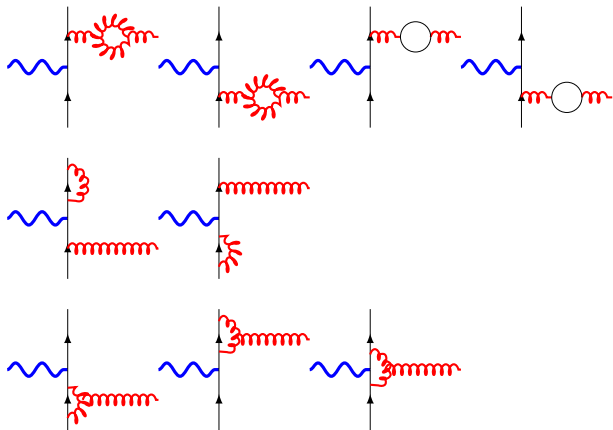
Real contribution



- integration of the extra gluon emission is the most non trivial part

$\mathcal{O}(\alpha_S^2)$ corrections to $\gamma^* + q \rightarrow q + g$

Virtual contribution



Framework for automatic Feynman amps calculation

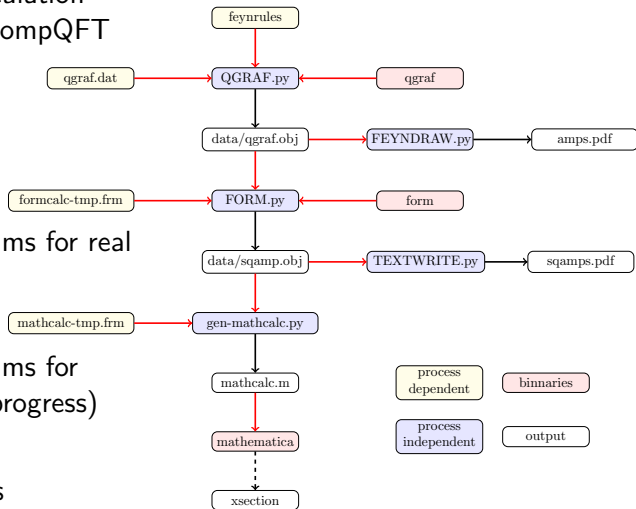
- We perform the calculation using the package CompQFT

- CompQFT is still in development

- ✓ Real emission diagrams for real photons

- Real emission diagrams for virtual photons (in progress)

- Virtual contributions (in progress)



MCEG for TMDs studies



Phenomenological Study of Hadronization in Nuclear and High-Energy Physics Experiments



LDRD Proposal
2017-LDRD-5

Jefferson Lab
Thomas Jefferson National Accelerator Facility

- MCEG for SIDIS with spin
- Understanding hadronization process: comparison between Lund string models and CSS
- Extension of CSS to describe hadronization from the soft region

Summary and outlook

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 - ✓ New MC fitting methodology
 - ✓ New IMC analysis for polarized PDFs and FFs
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- **MCEG for SIDIS:** Language dictionary
 - Pythia8 validation of Hermes multiplicities (in progress)
 - Extraction of FFs from Pythia 8: test of DGLAP and Pythia8's parton shower