

# TMD FFs from SIDIS

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Fragmentation Functions  
2018

Stresa, Feb. 21<sup>st</sup> 2018



# Outline

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1) TMD FFs

2) extraction from SIDIS data

3) what's next

# TMD FFs

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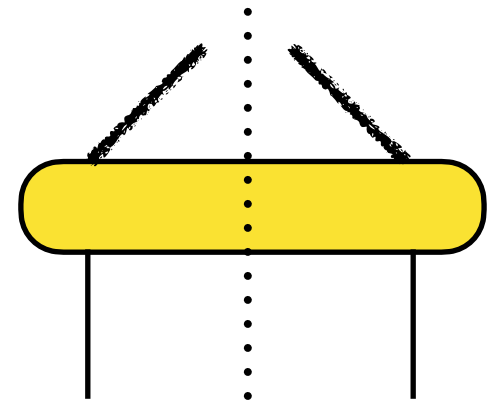
# Definition

[Parton Fragmentation Functions \(Metz-Vossen\) - DOI: 10.1016/j.ppnp.2016.08.003](https://doi.org/10.1016/j.ppnp.2016.08.003)

$$\Delta_{ij}(z, P_{h\perp}) = \frac{1}{2z} \sum_{\mathcal{X}} \int \frac{d\xi^+ \xi_T^2}{(2\pi^3)} e^{ik\xi} \langle 0 | \psi_i(\xi) | h\mathcal{X} \rangle \langle \mathcal{X} h | \bar{\psi}_j(0) | 0 \rangle$$

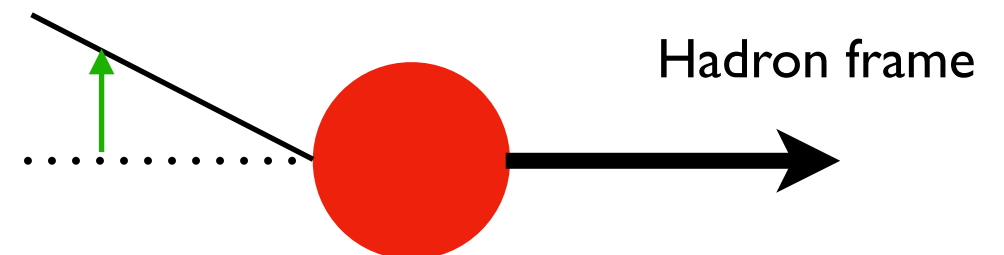
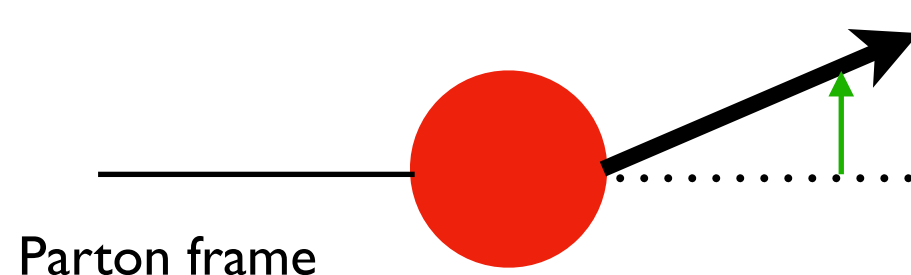
Hadronic variables  
(probabilistic interpretation)

Single hadron FF



J. Collins - Foundations of pQCD (2011)

Although the parton frame is a natural one for defining fragmentation functions as number densities, it is inconvenient for derivations of factorization. The problem is that, in a physical process, there is an integral over parton momentum, and so the parton-frame axes are not fixed. Neither parton momenta nor the resulting parton-frame axes can be determined from experimentally measured quantities. Therefore we will express the definitions of fragmentation functions in hadron-frame coordinates. In the derivation of factorization, we will use a hadron frame defined in terms of measured quantities.



Different frames for different purposes



# TMD FFs

		quark pol.		
		U	L	T
hadron pol.	U	$D_1$		$H_1^\perp$
	L		$G_{1L}$	$H_{1L}^\perp$
	T	$D_{1T}^\perp$	$G_{1T}$	$H_1, H_{1T}^\perp$

Twist-2 table

Diagonal: also collinear

Red: T-odd (but universal, unlike TMD PDFs)

Blue: T-even

A similar table exists for gluon TMD FFs

Correlator for spin 1/2 hadron:

Dirac matrix parametrized by  
**quark TMD FFs**

$$D_1^{a \rightarrow h}(z, P_\perp^2; \mu, \zeta)$$

Evolution equations with respect to  
two scales:

- UV renormalization
- rapidity renormalization

# TMDs and their evolution

FT of TMDs :

$$\tilde{F}_i(x, b_T; Q, Q^2) = \tilde{F}_i(x, b_T, \mu_{\hat{b}}, \mu_{\hat{b}}^2) \times \exp \left\{ \int_{\mu_{\hat{b}}}^Q \frac{d\mu}{\mu} \gamma_F[\alpha_s(\mu), Q^2/\mu^2] \right\} \left( \frac{Q^2}{\mu_{\hat{b}}^2} \right)^{-K(\hat{b}_T; \mu_{\hat{b}}) - g_K(\bar{b}_T; \{\lambda\})}$$

Sudakov form factor : perturbative and **nonperturbative** contributions

(input) TMD distribution : Wilson coefficients and **intrinsic part** Collinear distribution!

$$\tilde{F}_i(x, b_T; \mu_{\hat{b}}, \mu_{\hat{b}}^2) = \sum_{j=q, \bar{q}, g} C_{i/j}(x, \hat{b}_T; \mu_{\hat{b}}, \mu_{\hat{b}}^2) \otimes f_j(x; \mu_{\hat{b}}) \tilde{F}_{i, NP}(x, \bar{b}_T; \{\lambda\})$$

Nonperturbative parts : **power corrections to perturbative calculations**

# TMDs and their evolution

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Distribution for intrinsic transverse momentum  
(and its FT):

$$\tilde{F}_{i,NP}(x, \bar{b}_T; \{\lambda\})$$

Which form ?

Soft gluon emission

$$g_K(\bar{b}_T; \{\lambda\})$$

# TMDs and their evolution

Distribution for intrinsic transverse momentum  
(and its FT):

$$\tilde{F}_{i,NP}(x, \bar{b}_T; \{\lambda\})$$

Which form ?

Soft gluon emission

$$g_K(\bar{b}_T; \{\lambda\})$$

Separation of  $b_T$  regions

$$\hat{b}_T(b_T; b_{\min}, b_{\max}) \begin{cases} \nearrow b_{\max}, & b_T \rightarrow +\infty \\ \sim b_T, & b_{\min} \ll b_T \ll b_{\max} \\ \searrow b_{\min}, & b_T \rightarrow 0 \end{cases}$$

High  $b_T$  limit : avoid Landau pole

Low  $b_T$  limit : recover fixed order expression

# Extraction from SIDIS

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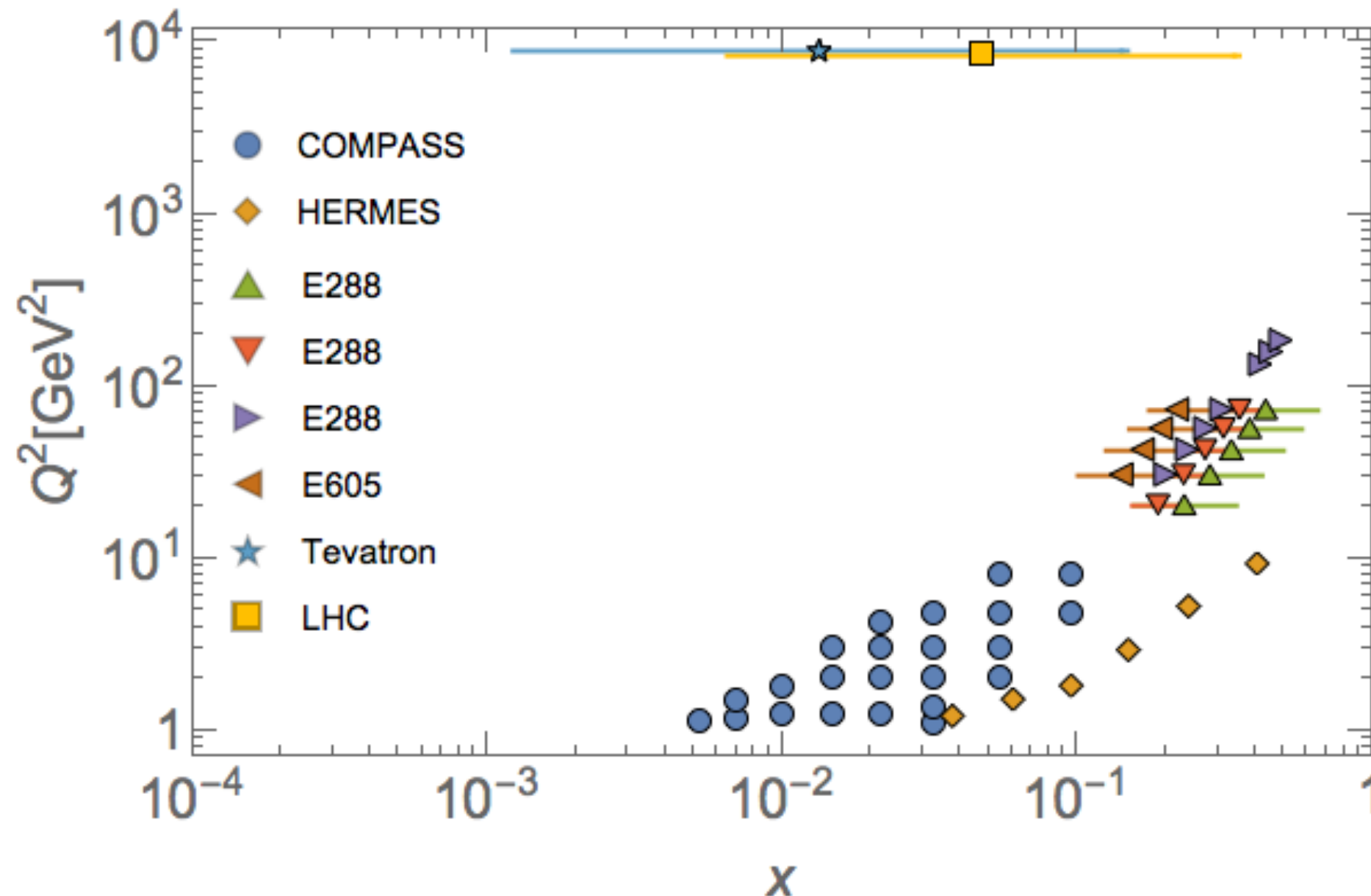
# What do we know ?

(only a selection of results!)

	Framework	HERMES	COMPASS	DY	Z production	N of points
KN 2006 <a href="#">hep-ph/0506225</a>	LO-NLL	✗	✗	✓	✓	98
Pavia 2013 (+Amsterdam, Bilbao) <a href="#">arXiv:1309.3507</a>	No evo (QPM)	✓	✗	✗	✗	1538
Torino 2014 (+JLab) <a href="#">arXiv:1312.6261</a>	No evo (QPM)	✓ (separately)	✓ (separately)	✗	✗	576 (H) 6284 (C)
DEMS 2014 <a href="#">arXiv:1407.3311</a>	NLO-NNLL	✗	✗	✓	✓	223
EIKV 2014 <a href="#">arXiv:1401.5078</a>	LO-NLL	1 (x,Q <sup>2</sup> ) bin	1 (x,Q <sup>2</sup> ) bin	✓	✓	500 (?)
Pavia 2017 <a href="#">arXiv:1703.10157</a>	LO-NLL	✓	✓	✓	✓	8059
SV 2017 <a href="#">arXiv:1706.01473</a>	NNLO- NNLL	✗	✗	✓	✓	309

[ courtesy A. Bacchetta ]

# Data sets and kinematic coverage

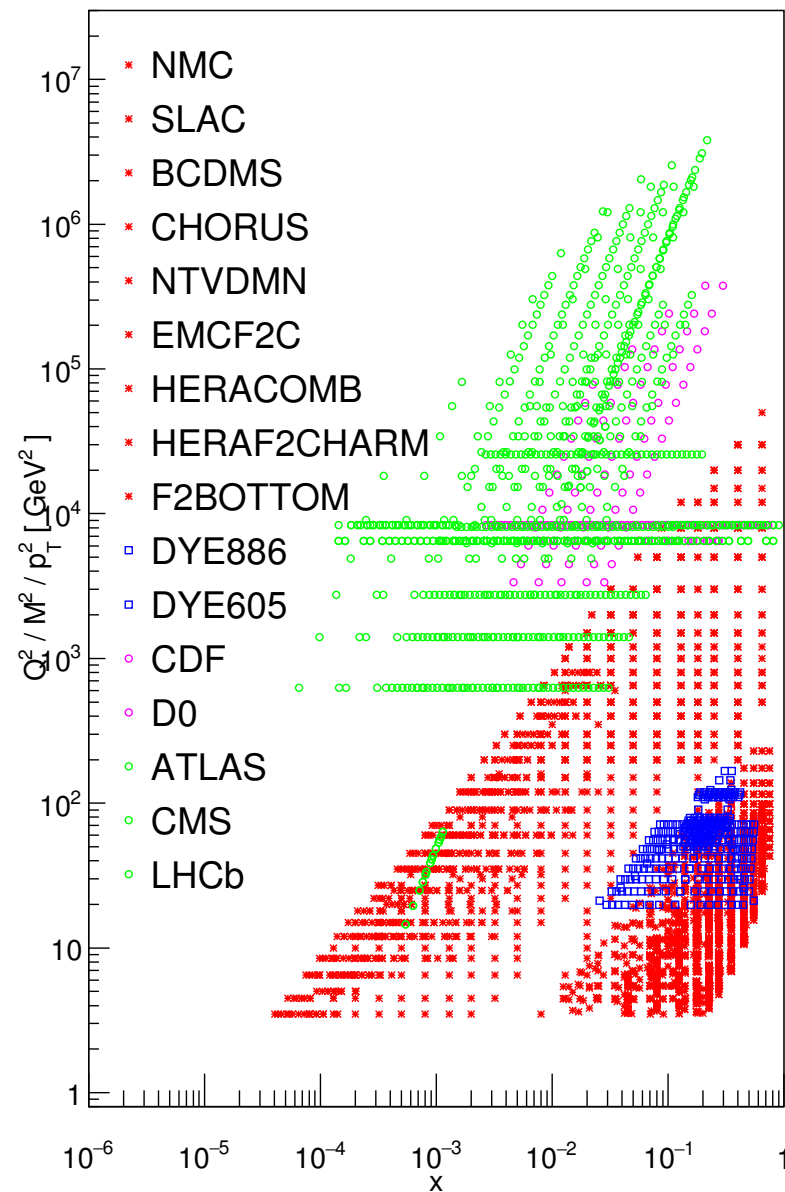


**Electron-positron** annihilation data are still **missing**  
(only some azimuthal asymmetries are available)

**crucial for analyses**  
**of TMD FFs !!**

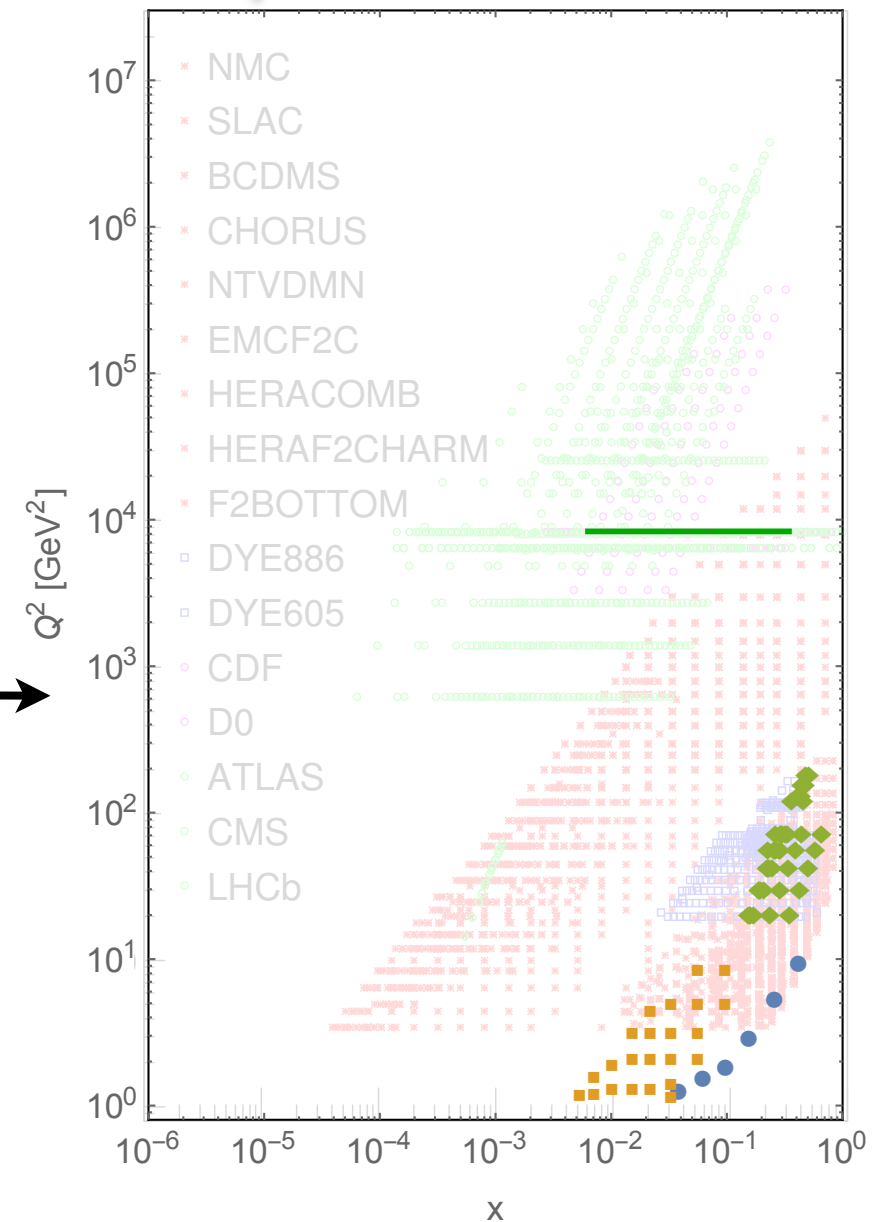
# Comparison with collinear PDF fits

see talk by E. Nocera at POETIC2016



data sets available:

← collinear PDFs  
vs  
TMD PDFs →

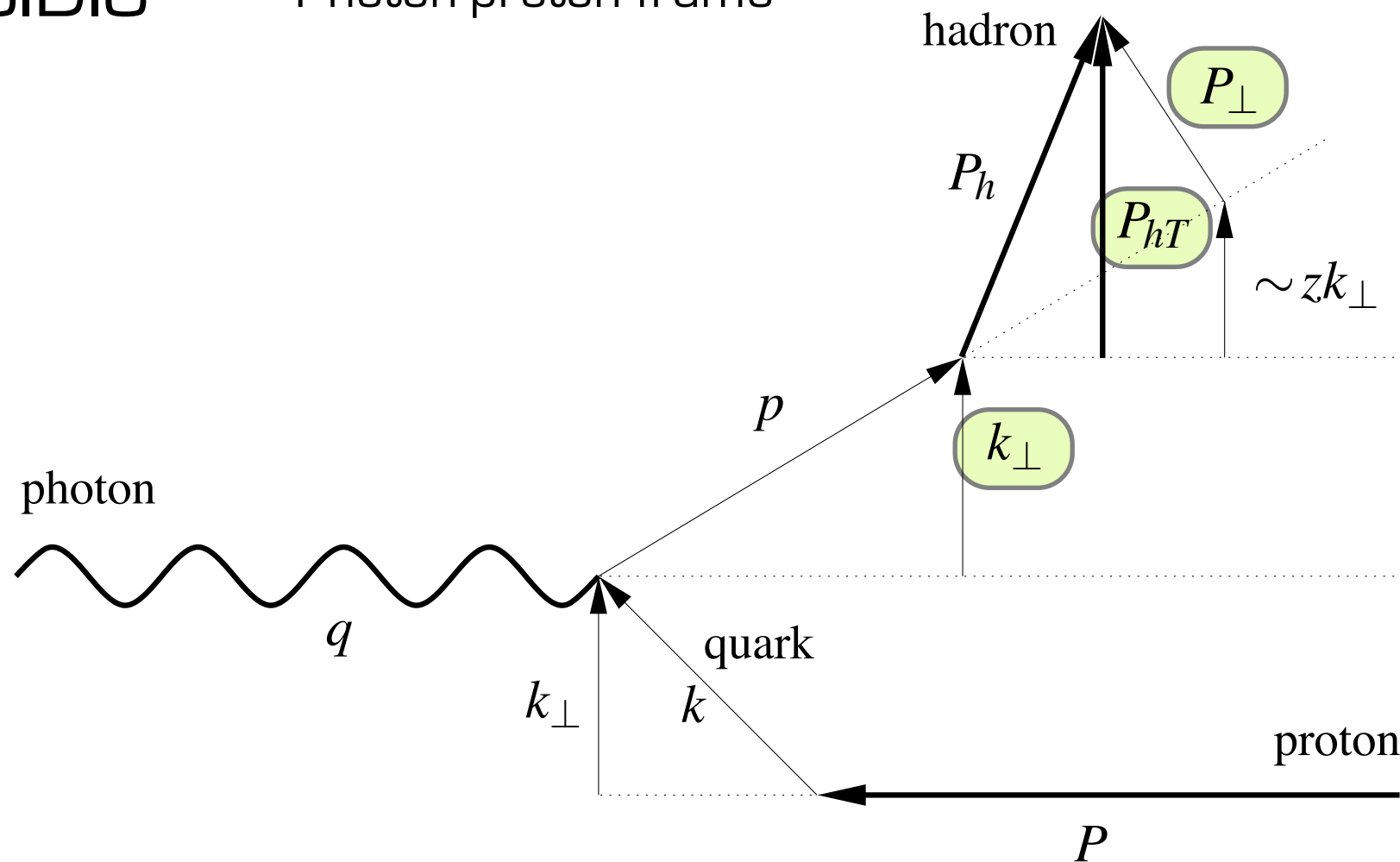




# Transverse momenta in SIDIS

SIDIS

Photon-proton frame



TMD FF

TMD PDF

Longitudinal scaling variable

$$z = \frac{P \cdot P_h}{P \cdot q}$$

Observed transverse momentum

$$P_{hT} \approx z k_\perp + P_\perp$$

# Data sets and selections - SIDIS

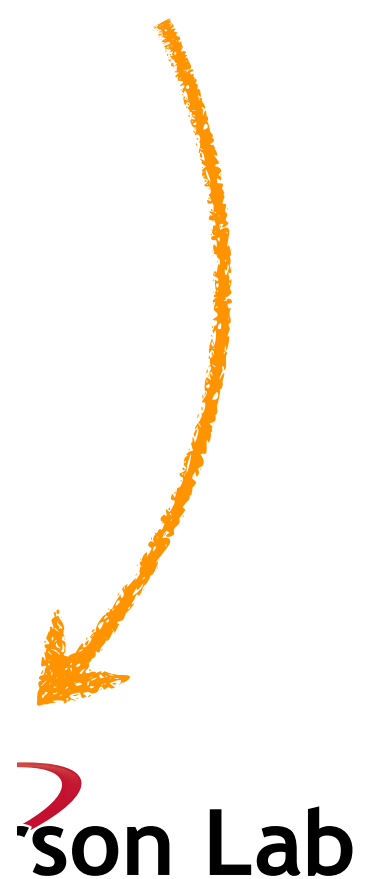
	HERMES $p \rightarrow \pi^+$	HERMES $p \rightarrow \pi^-$	HERMES $p \rightarrow K^+$	HERMES $p \rightarrow K^-$
Reference	[61]			
Cuts	$Q^2 > 1.4 \text{ GeV}^2$ $0.2 < z < 0.7$ $P_{hT} < \text{Min}[0.2 Q, 0.7 Qz] + 0.5 \text{ GeV}$			
Points	190	190	189	187
Max. $Q^2$	9.2 $\text{GeV}^2$			
$x$ range	$0.06 < x < 0.4$			

TMD factorization  
 $(P_{hT}^2/z^2 \ll Q^2)$

avoid target fragmentation [?]  
 (low  $z$ )  
 and exclusive contributions [?]  
 (high  $z$ )

Problem with normalization  
 in the previous release

	HERMES $D \rightarrow \pi^+$	HERMES $D \rightarrow \pi^-$	HERMES $D \rightarrow K^+$	HERMES $D \rightarrow K^-$	COMPASS $D \rightarrow h^+$	COMPASS $D \rightarrow h^-$
Reference	<a href="#">[74]</a>				<a href="#">[75]</a>	
Cuts	$Q^2 > 1.4 \text{ GeV}^2$ $0.20 < z < 0.74$ $P_{hT} < \text{Min}[0.2 Q, 0.7 Qz] + 0.5 \text{ GeV}$					
Points	190	190	189	189	3125	3127
Max. $Q^2$	9.2 GeV <sup>2</sup>				10 GeV <sup>2</sup>	
$x$ range	$0.04 < x < 0.4$				$0.005 < x < 0.12$	
Notes					Observable: $m_{\text{norm}}(x, z, \mathbf{P}_{hT}^2, Q^2)$ , eq. (3.1)	



# Features

	Framework	HERMES	COMPASS	DY	Z production	N of points
Pavia 2017 arXiv:1703.10157	LO-NLL	✓	✓	✓	✓	8059

## PROs

almost a **global fit** of  
quark unpolarized TMDs

includes **TMD evolution**

**replica (bootstrap)**

fitting methodology

**kinematic dependence**

in intrinsic part of TMDs

intrinsic momentum: **beyond  
the Gaussian** assumption

## CONs

no “pure” info on TMD FFs

accuracy of TMD evolution :  
not the state of the art

only “low” transverse momentum  
(no fixed order and Y-term)

flavor separation in  
the transverse  
plane : problematic

# Intrinsic transverse momentum

$$f_{1NP}^a(x, k_{\perp}^2) = \frac{1}{\pi} \frac{(1 + \lambda k_{\perp}^2)}{g_{1a} + \lambda g_{1a}^2} e^{-\frac{k_{\perp}^2}{g_{1a}}}$$

$$\hat{x} = 0.1$$

$$g_1(x) = N_1 \frac{(1-x)^{\alpha} x^{\sigma}}{(1-\hat{x})^{\alpha} \hat{x}^{\sigma}}$$

**weighted sum of two Gaussian distributions:**  
**same widths** for TMD PDFs  
**different widths** for TMD FFs

$$\hat{z} = 0.5$$

$$g_{3,4}(z) = N_{3,4} \frac{(z^{\beta} + \delta)(1-z)^{\gamma}}{(\hat{z}^{\beta} + \delta)(1-\hat{z})^{\gamma}}$$

$$D_{1NP}^{a \rightarrow h}(z, P_{\perp}^2) = \frac{1}{\pi} \frac{1}{g_{3a \rightarrow h} + (\lambda_F / z^2) g_{4a \rightarrow h}^2} \left( e^{-\frac{P_{\perp}^2}{g_{3a \rightarrow h}}} + \lambda_F \frac{P_{\perp}^2}{z^2} e^{-\frac{P_{\perp}^2}{g_{4a \rightarrow h}}} \right)$$

Inspired by **model calculations**:

Matevosyan et al.

Phys. Rev. D85, 014021 (2012), 1111.1740

Bacchetta et al.

Phys. Lett. B659, 234 (2008), 0707.3372

Bacchetta et al.

Phys. Rev. D65, 094021 (2002),

hep-ph/0201091

There are **11 free parameters** in a flavor independent scenario (one for evolution)

# Models - evolution and $b_T$ regions

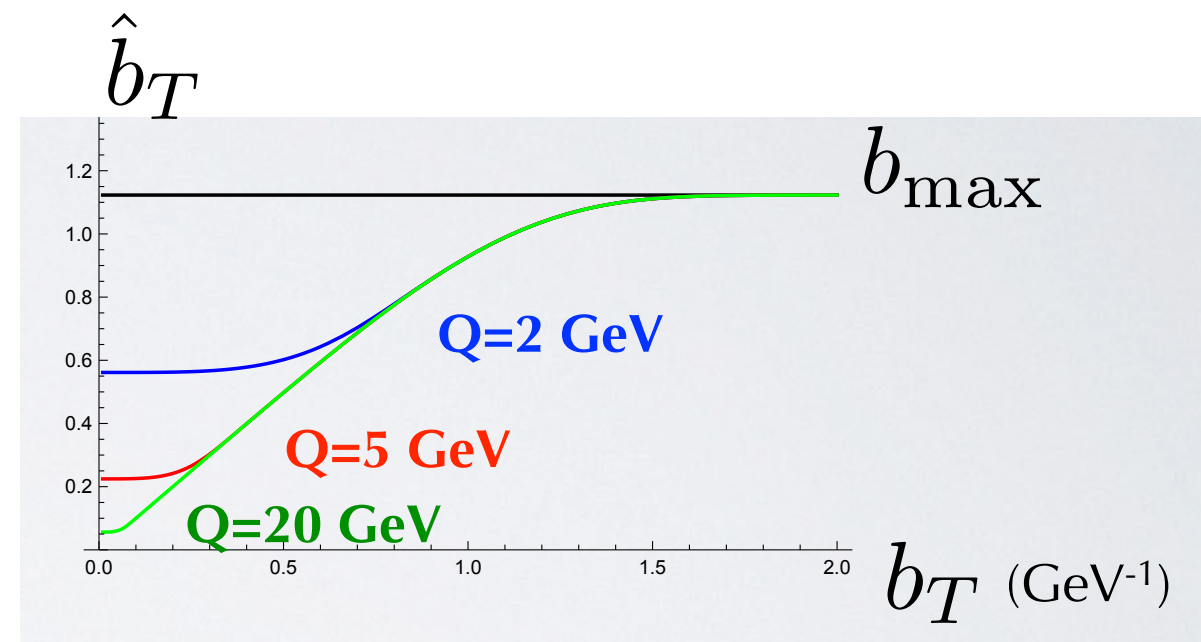
$$g_K(b_T; g_2) = -g_2 \frac{b_T^2}{2}$$

$$\hat{b}(b_T; b_{\min}, b_{\max}) = b_{\max} \left( \frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}} \right) \begin{matrix} \nearrow b_{\max}, & b_T \rightarrow +\infty \\ \searrow b_{\min}, & b_T \rightarrow 0 \end{matrix}$$

$$b_{\max} = 2e^{-\gamma_E}$$

$$b_{\min} = 2e^{-\gamma_E}/Q$$

These choices guarantee that for  $Q=1$  GeV the TMD coincides with the NP model



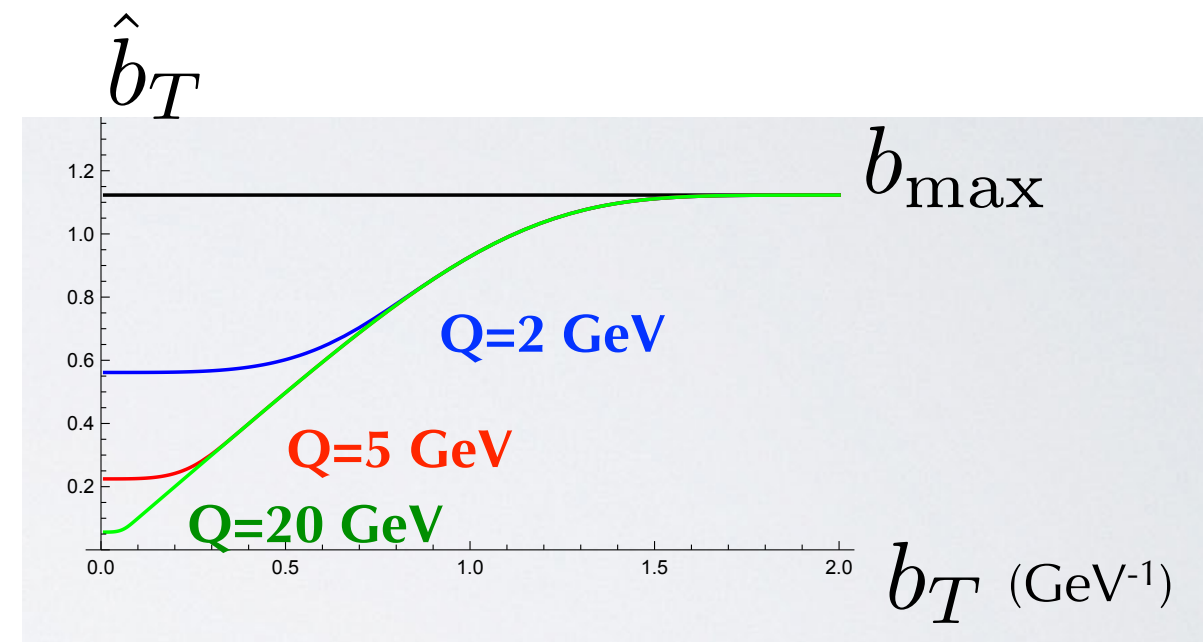
# Models - evolution and $b_T$ regions

$$g_K(b_T; g_2) = -g_2 \frac{b_T^2}{2}$$

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$$b_{\min} \sim 1/Q, \quad \mu_{\hat{b}} < Q$$

The phenomenological importance of  $b_{\min}$  is a signal that -especially in SIDIS data at **low  $Q$** - we are exiting the TMD region, entering the collinear factorization region



# Agreement data-theory

Flavor independent scenario

Flavor independent configuration | 11 parameters

Points	Parameters	$\chi^2$	$\chi^2/\text{d.o.f.}$
8059	11	$12629 \pm 363$	$1.55 \pm 0.05$

	HERMES $p \rightarrow \pi^+$	HERMES $p \rightarrow \pi^-$	HERMES $p \rightarrow K^+$	HERMES $p \rightarrow K^-$
Points	190	190	189	187
$\chi^2/\text{points}$	4.83	2.47	0.91	0.82

**Hermes** P/D into  $\pi^+$ :  
problems at low z

	HERMES $D \rightarrow \pi^+$	HERMES $D \rightarrow \pi^-$	HERMES $D \rightarrow K^+$	HERMES $D \rightarrow K^-$	COMPASS $D \rightarrow h^+$	COMPASS $D \rightarrow h^-$
Points	190	190	189	189	3125	3127
$\chi^2/\text{points}$	3.46	2.00	1.31	2.54	1.11	1.61

	E288 [200]	E288 [300]	E288 [400]	E605
Points	45	45	78	35
$\chi^2/\text{points}$	0.99	0.84	0.32	1.12

**Hermes** kaons better than pions:  
larger uncertainties from FFs

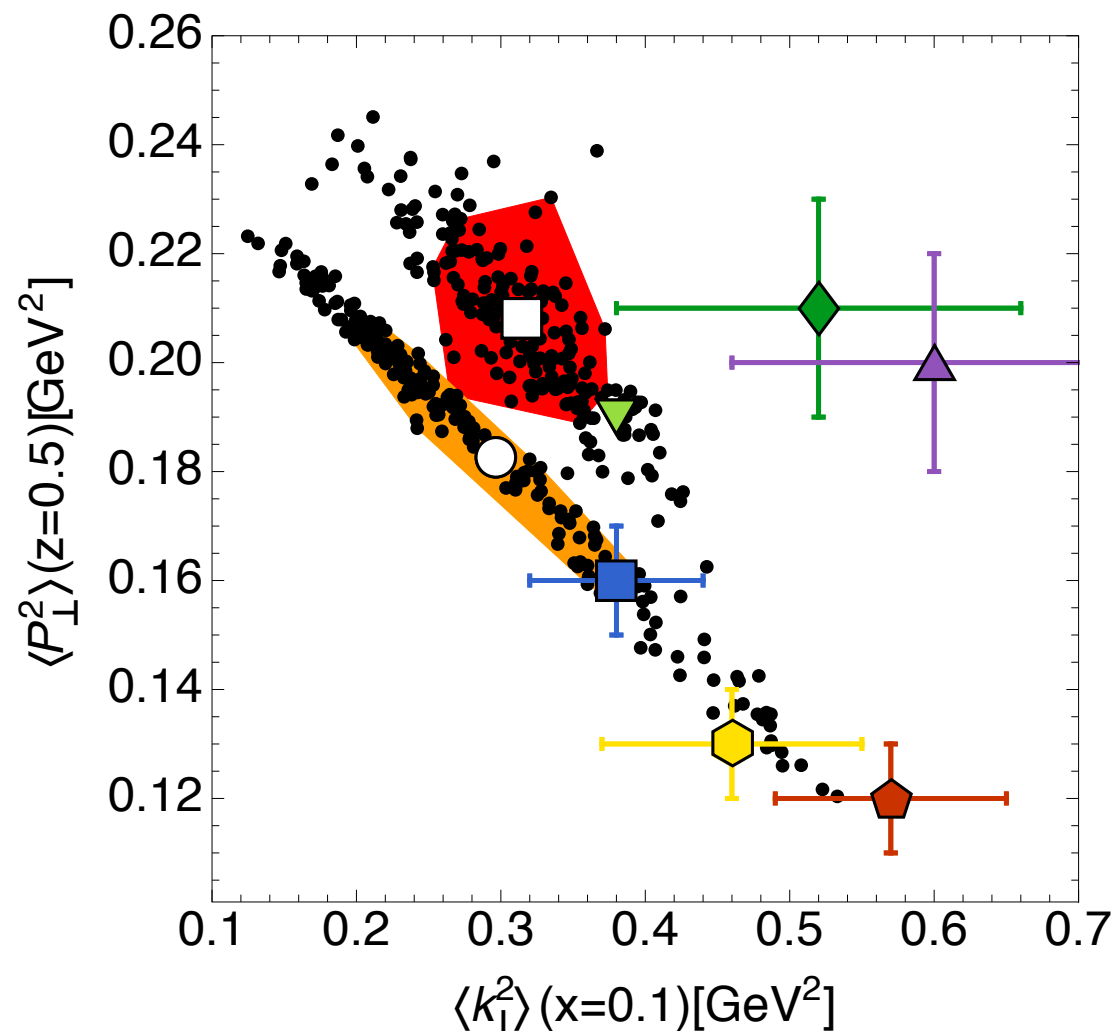
**Compass** : better agreement due to  
#points and normalization

**Let's see what**  
**happens with the new data**

	CDF Run I	D0 Run I	CDF Run II	D0 Run II
Points	31	14	37	8
$\chi^2/\text{points}$	1.36	1.11	2.00	1.73

# Average transverse momenta

Flavor ind. scenario



- Bacchetta, Delcarro, Pisano, Radici, Signori (JHEP 2017)
- Signori, Bacchetta, Radici, Schnell arXiv:1309.3507
- Schweitzer, Teckentrup, Metz, arXiv:1003.2190
- ⬡ Anselmino et al. arXiv:1312.6261 [HERMES]
- ⬡ Anselmino et al. arXiv:1312.6261 [HERMES, high z]
- ◆ Anselmino et al. arXiv:1312.6261 [COMPASS, norm.]
- ▲ Anselmino et al. arXiv:1312.6261 [COMPASS, high z, norm.]
- ▼ Echevarria, Idilbi, Kang, Vitev arXiv:1401.5078 ( $Q = 1.5 \text{ GeV}$ )

Red/orange regions : **68% CL** from replica method

Inclusion of **Compass** increases the  $\langle P_{\perp}^2 \rangle$   
and reduces its spread

Inclusion of **DY/Z** diminishes the correlation

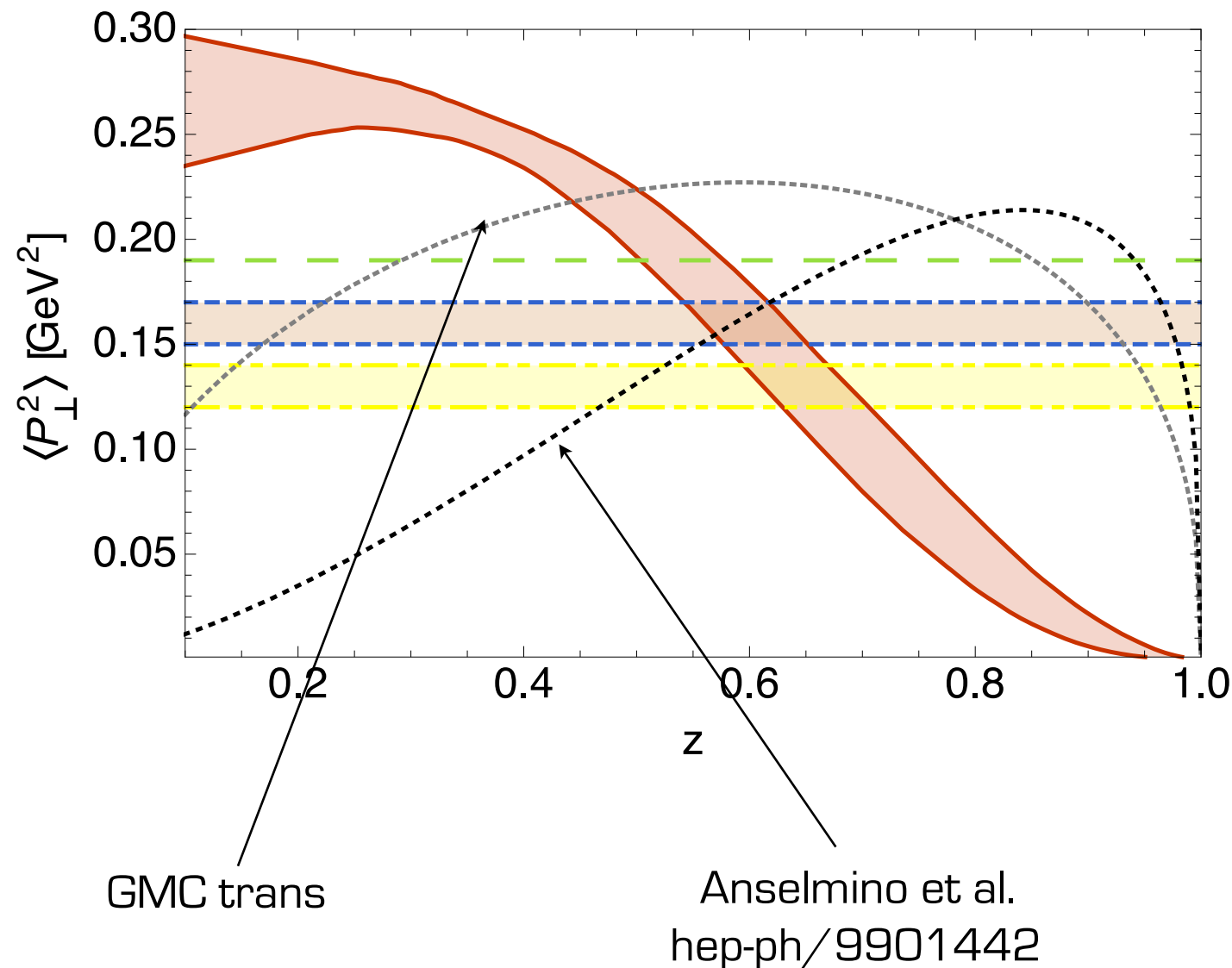
**e+e-** data would further reduce the correlation



# Kinematic dependence

$$\langle P_{\perp}^2 \rangle(z) = \frac{\int d^2 P_{\perp} P_{\perp}^2 D_1^{a \rightarrow h}(z, P_{\perp}^2, Q = 1 \text{ GeV})}{\int d^2 P_{\perp} D_1^{a \rightarrow h}(z, P_{\perp}^2, Q = 1 \text{ GeV})}$$

Average square  
transverse momentum in TMD FF



Color code : same as previous slide

Flavor-independent scenario:  
no differences in quark/hadron  
flavor

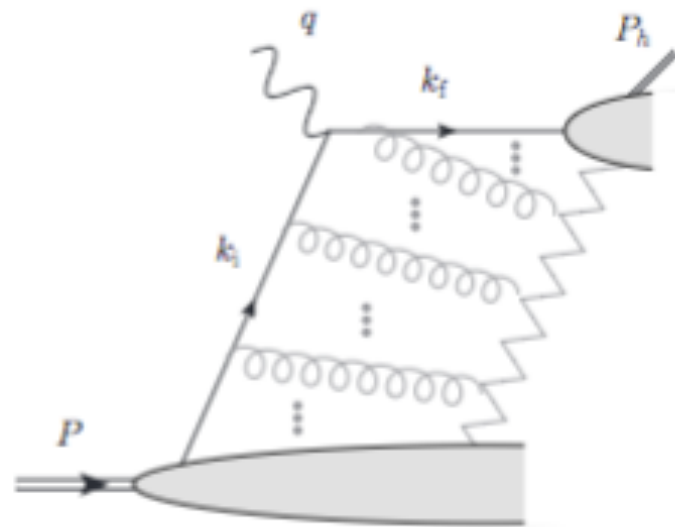
$z$ -dependence :  
important to fit the data

# What's next

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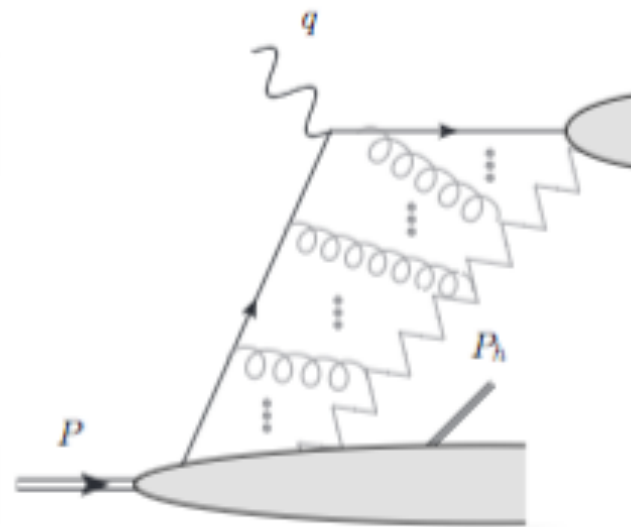
Just a selection of topics to feed the discussion

# Target **vs** current **vs** central regions



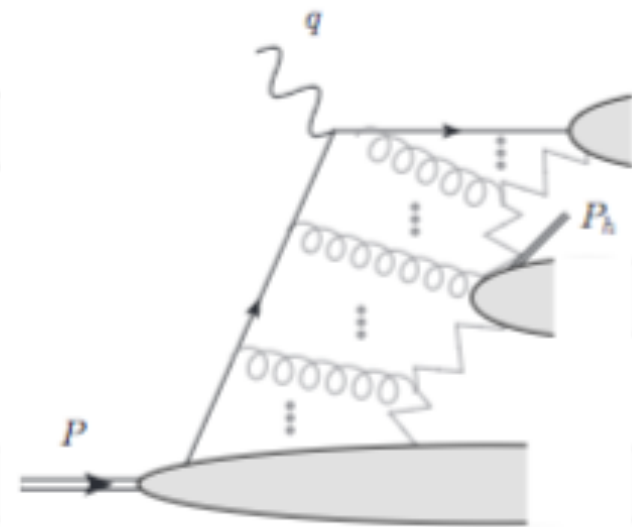
(a)

current fragmentation



(b)

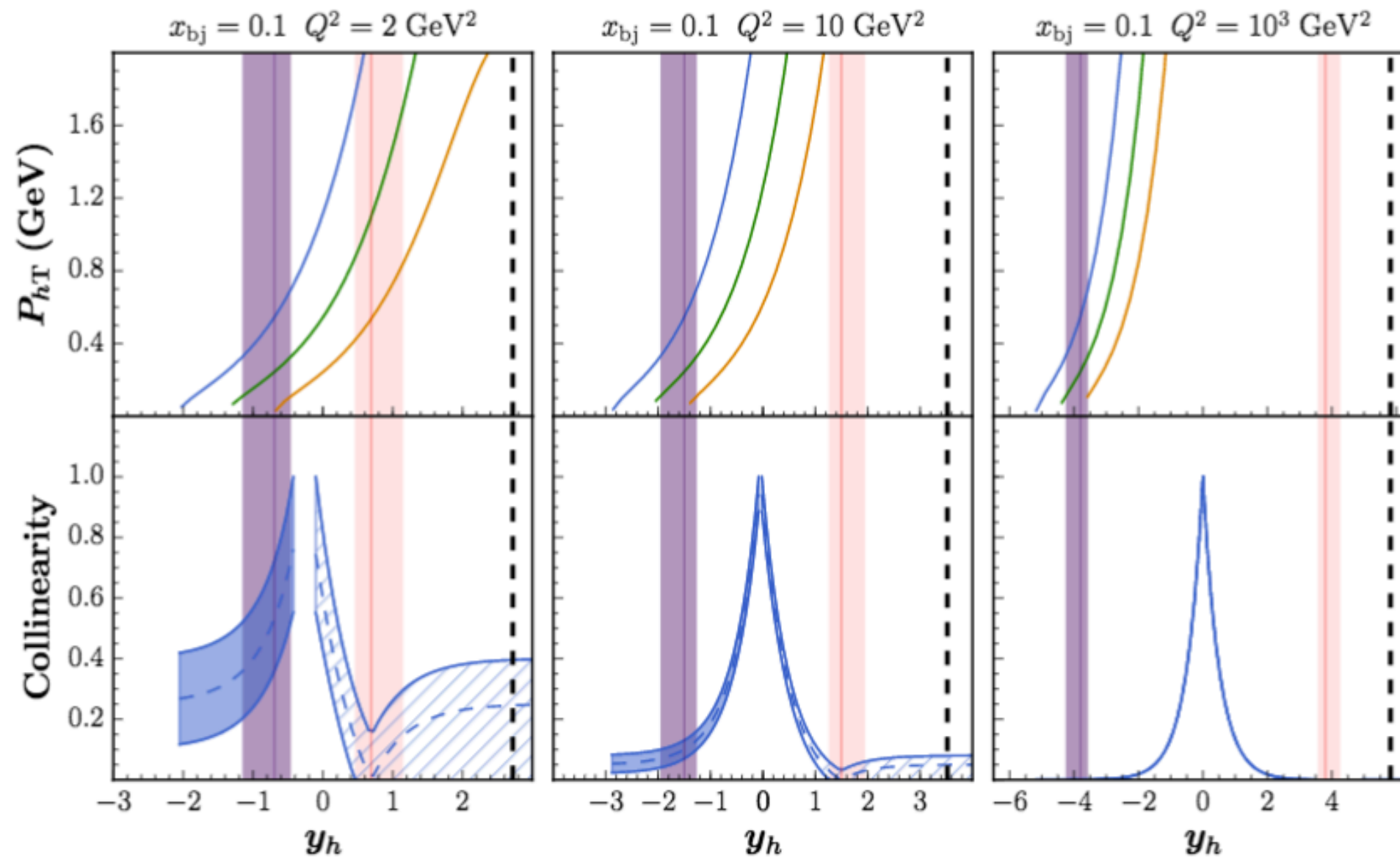
target fragmentation



(c)

“soft” fragmentation

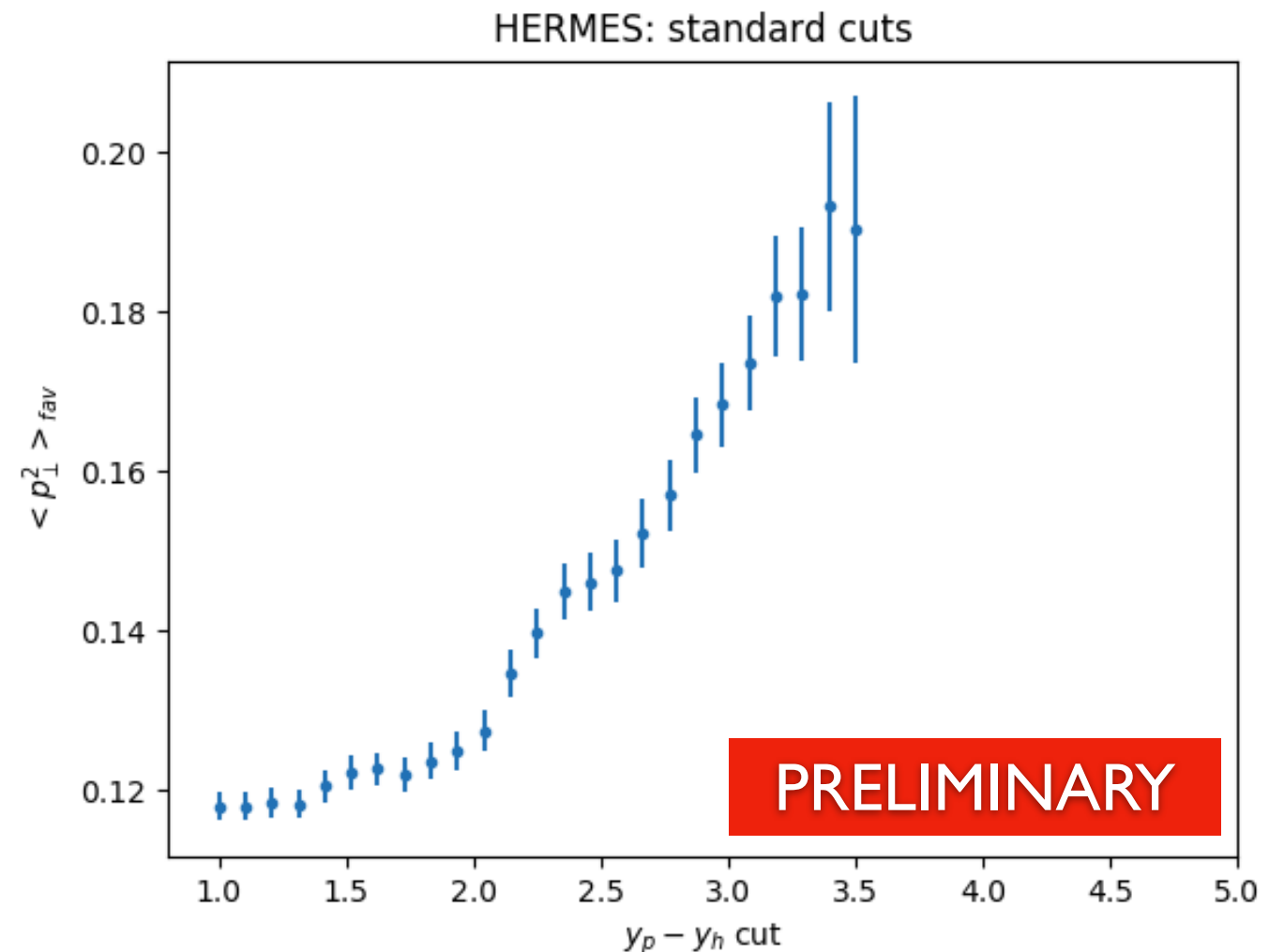
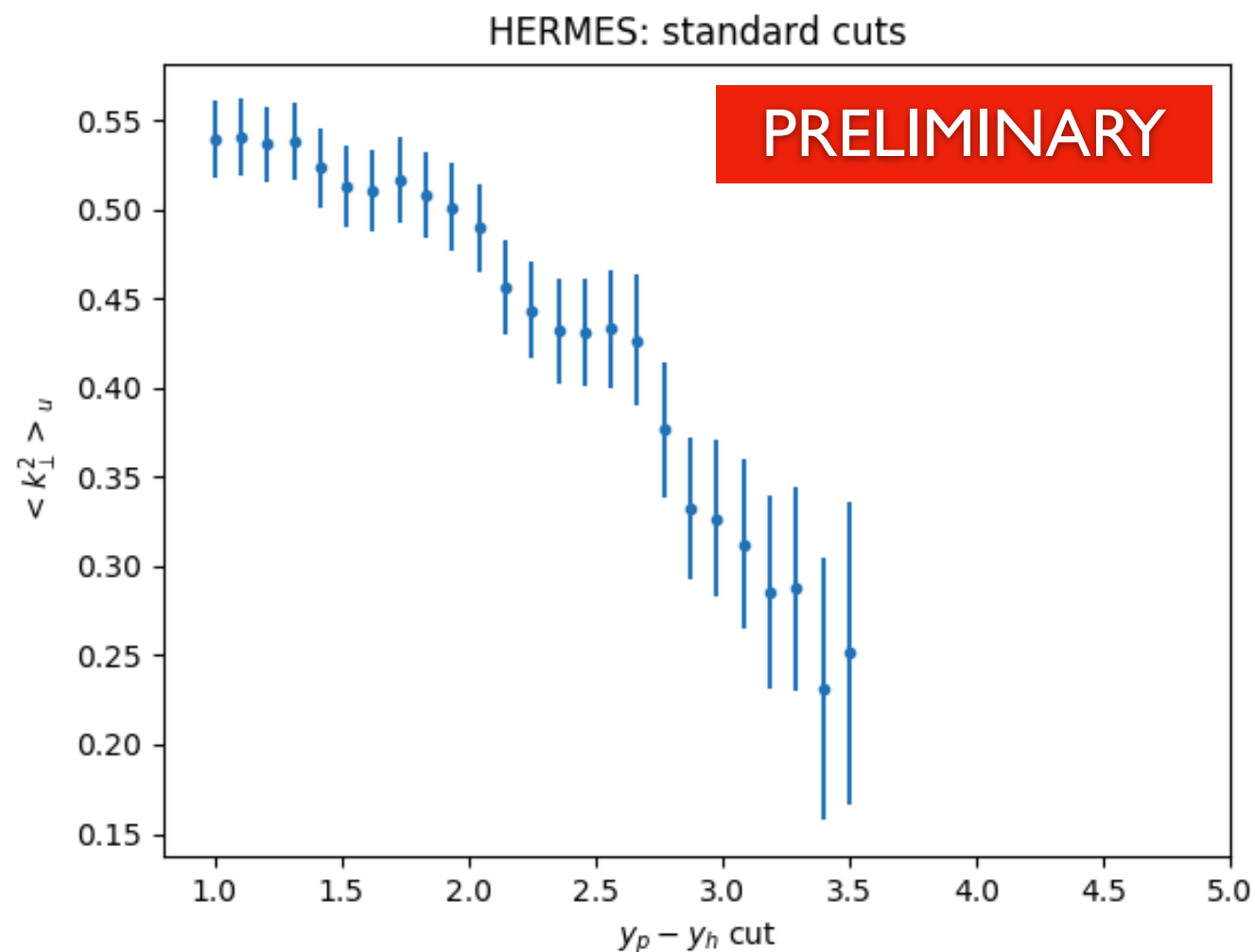
# Target *vs* current *vs* central regions



See O. Gonzalez's talk

“collinearity” criterion

# Target *vs* current *vs* central regions



Description of Hermes data within the quark parton model

Widths described as a function of the *rapidity difference*  
*between the incoming proton and outgoing hadron*

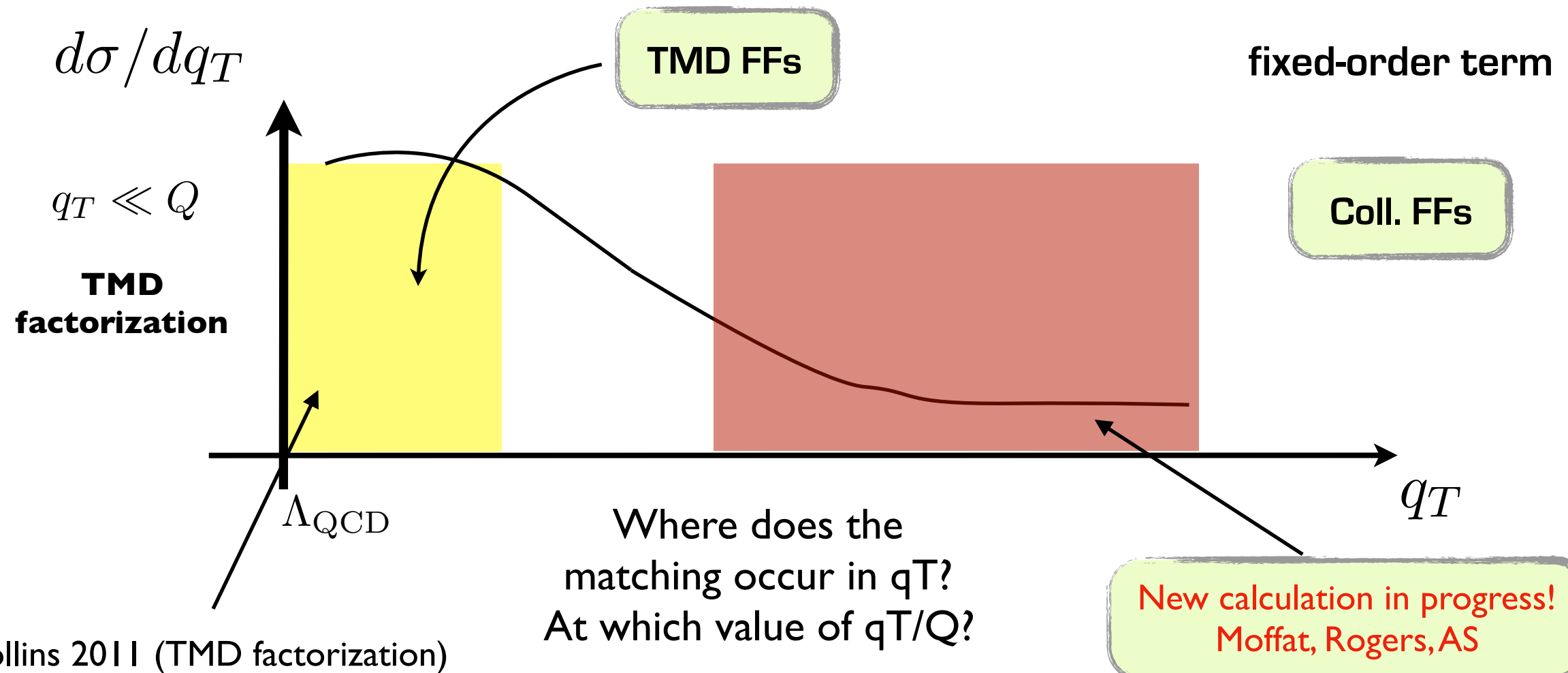
“rapidity difference”  
criterion

# TMD FFs from $e^+e^-$

Completing the formalism to study TMD FFs

fixed  $Q$ , variable  $q_T$

$$e^+e^- \rightarrow h_1 h_2 X$$



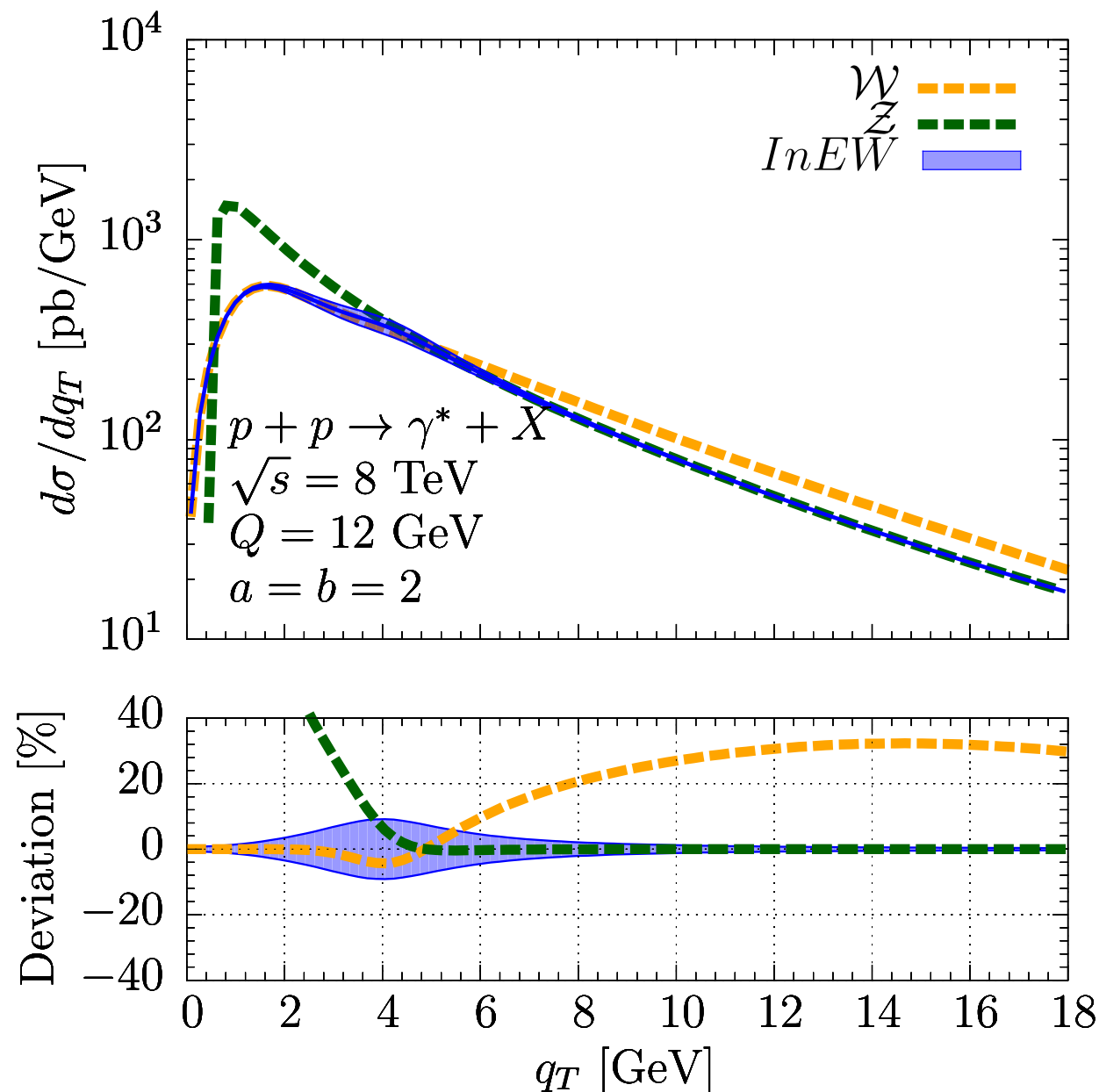
Several phenomenological works, e.g.:

\* Bacchetta, Echevarria, Mulders, Radici, AS - JHEP 2015

# Matching

Echevarria, Kasemets, Lansberg, AS, Pisano  
1801.01480

development of a new scheme (**InEW** - inverse error weighting)  
and comparison to improved CSS subtraction



$$d\sigma(q_T, Q) = \omega_1 \mathcal{W}(q_T, Q) + \omega_2 \mathcal{Z}(q_T, Q)$$

$$\omega_1 \sim \Delta\mathcal{W}^{-2}, \quad \Delta\mathcal{W} \sim \mathcal{O}\left(\frac{q_T}{Q}\right)^a + \mathcal{O}\left(\frac{m}{Q}\right)^{a'}$$

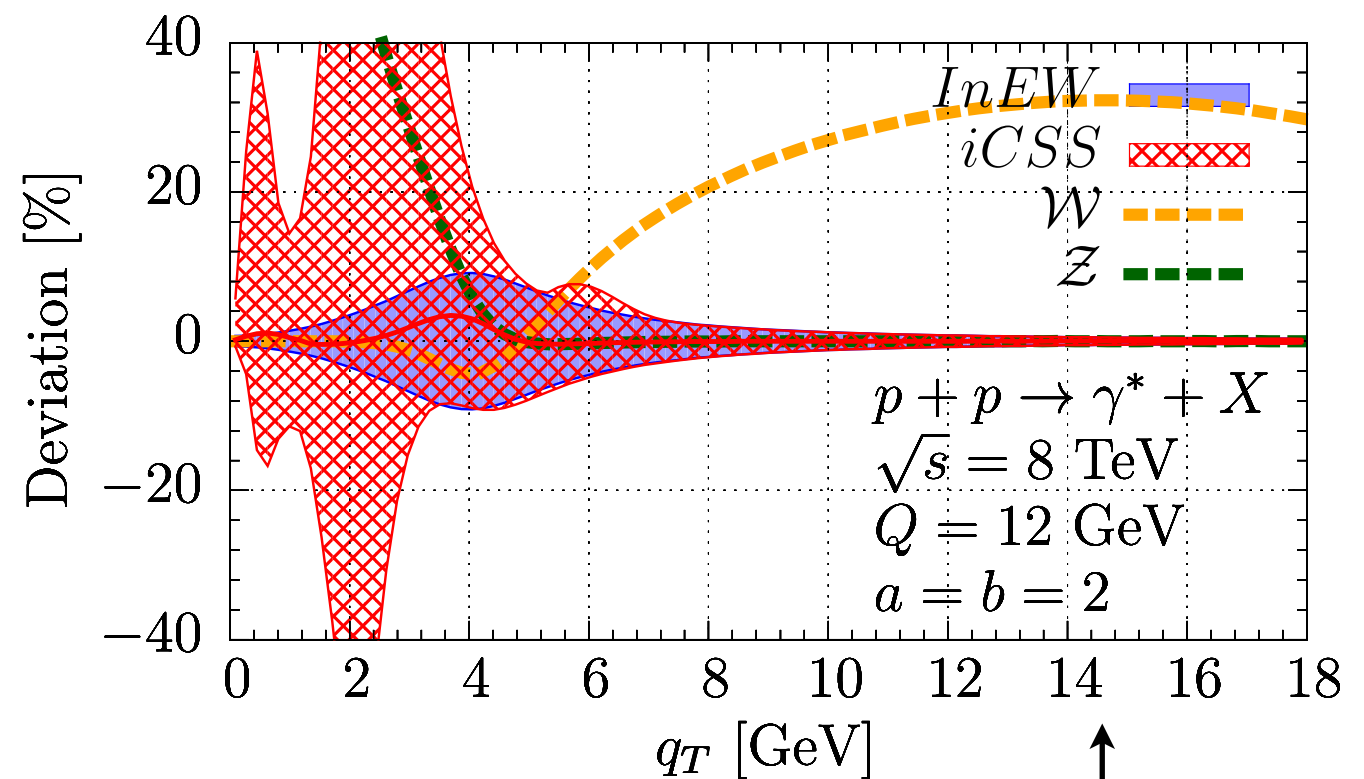
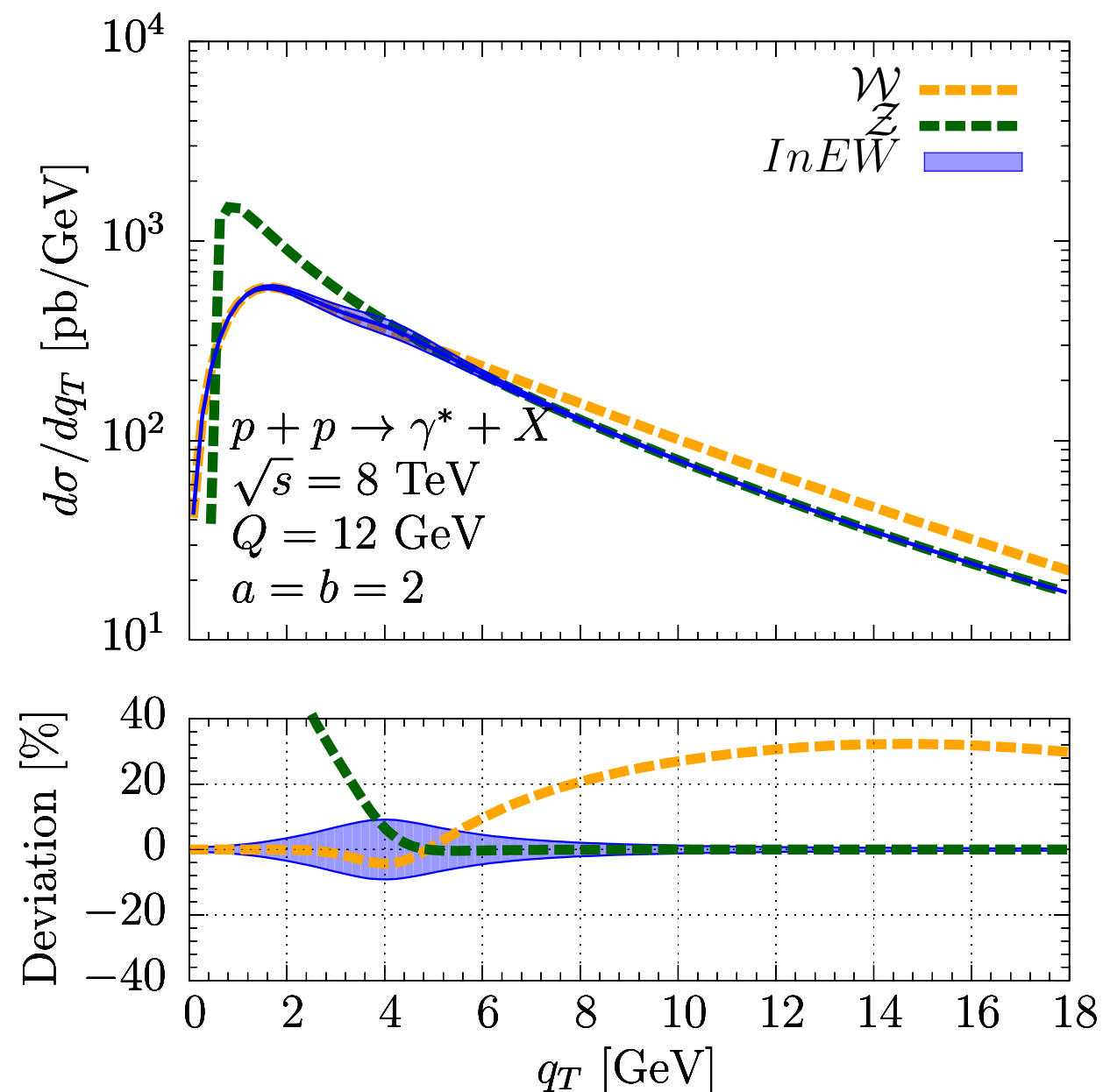
$$\omega_2 \sim \Delta\mathcal{Z}^{-2}, \quad \Delta\mathcal{Z} \sim \mathcal{O}\left(\frac{m}{q_T}\right)^b$$

Drell-Yan at LHC  
 $[Q=12 \text{ GeV}]$   
 [data available]

# Matching

Echevarria, Kasemets, Lansberg, AS, Pisano  
1801.01480

development of a new scheme (**InEW** - inverse error weighting)  
and comparison to improved CSS subtraction



Drell-Yan at LHC  
( $Q=12$  GeV)  
(data available)

InEW vs iCSS



# Inputs for discussion

---

## FORMALISM:

- \* definition of the **fragmentation regions**
- \* which **variables** shall we use to describe the momentum fractions? (see Gunar's talk)
- \* **matching** schemes: how to estimate **uncertainties** associated to the matching prescription
- \* **jet** fragmentation functions

## PERTURBATIVE ASPECTS:

- \* implementation of **evolution** (transverse momentum, threshold resummation, zeta prescription, ...?)
- \* **fixed-order** calculations in SIDIS : is it sufficient to describe data at higher  $q_T$  ? Or do we need power corrections/higher twist, contributions from soft region, etc ?

## NONPERTURBATIVE ASPECTS:

- \* functional form at **low transverse momentum**
- \* its **kinematic** dependence
- \* its **flavor** dependence
- \* nonperturbative contribution to TMD **evolution**

## DATA :

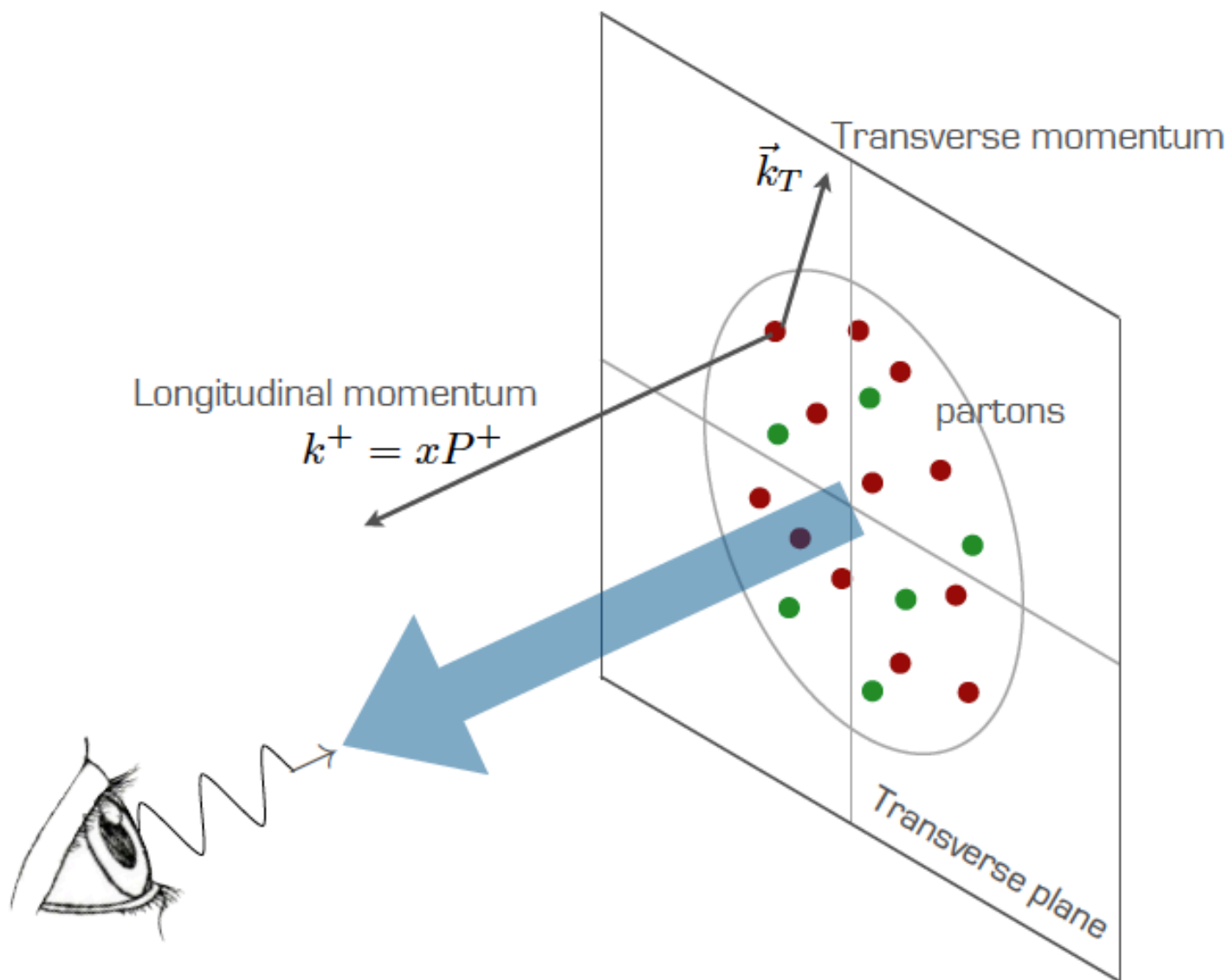
- \* impact of the new release of **Compass** data
- \* A **Fixed Target** Experiment at the **LHC** ?
- \* what can be done with the forthcoming  **$e^+e^-$**  data concerning TMD FFs (also including matching to high  $q_T$ )
- \* how well does the fixed order describes data at large transverse momentum
- \* ...

# Backup

---

# quark TMD PDFs

$$\Phi_{ij}(k, P; S_{\perp}) \sim \text{F.T.} \langle PS | \bar{\psi}_j(0) U_{[0,\xi]} \psi_i(\xi) | PS \rangle |_{LF}$$



extraction of a **quark**  
**not** collinear with the proton



# Beware of different notations...

---

Amsterdam

Seattle ([arXiv:1108.1713](#))

$p$

$k$  momentum of parton in distribution function

$p_T$

$k_\perp$  parton transverse momentum in distribution function

$k$

$p$  momentum of fragmenting parton

$k_T$

$p_\perp$  trans. momentum of fragmenting parton w.r.t. final hadron

$K_T$

$P_\perp$  trans. momentum of final hadron w.r.t. fragmenting parton

$P_{h\perp}$

$P_{hT}$  transverse momentum of final hadron w.r.t. virtual photon

Let's agree on the notation!

# Collinear and TMD factorization

Let's consider a process with  
**three separate scales:**

(SIDIS, Drell-Yan,  $e^+e^-$  to hadrons,  
pp to quarkonium, ... )

hadronic  
mass scale

$$\Lambda_{\text{QCD}} \ll q_T \ll Q$$

hard scale

(related to the)  
transverse momentum of the observed particle

The ratios

$$\Lambda_{\text{QCD}}/Q$$

$$\Lambda_{\text{QCD}}/q_T$$

$$q_T/Q$$

**select the factorization theorem** that we rely on.

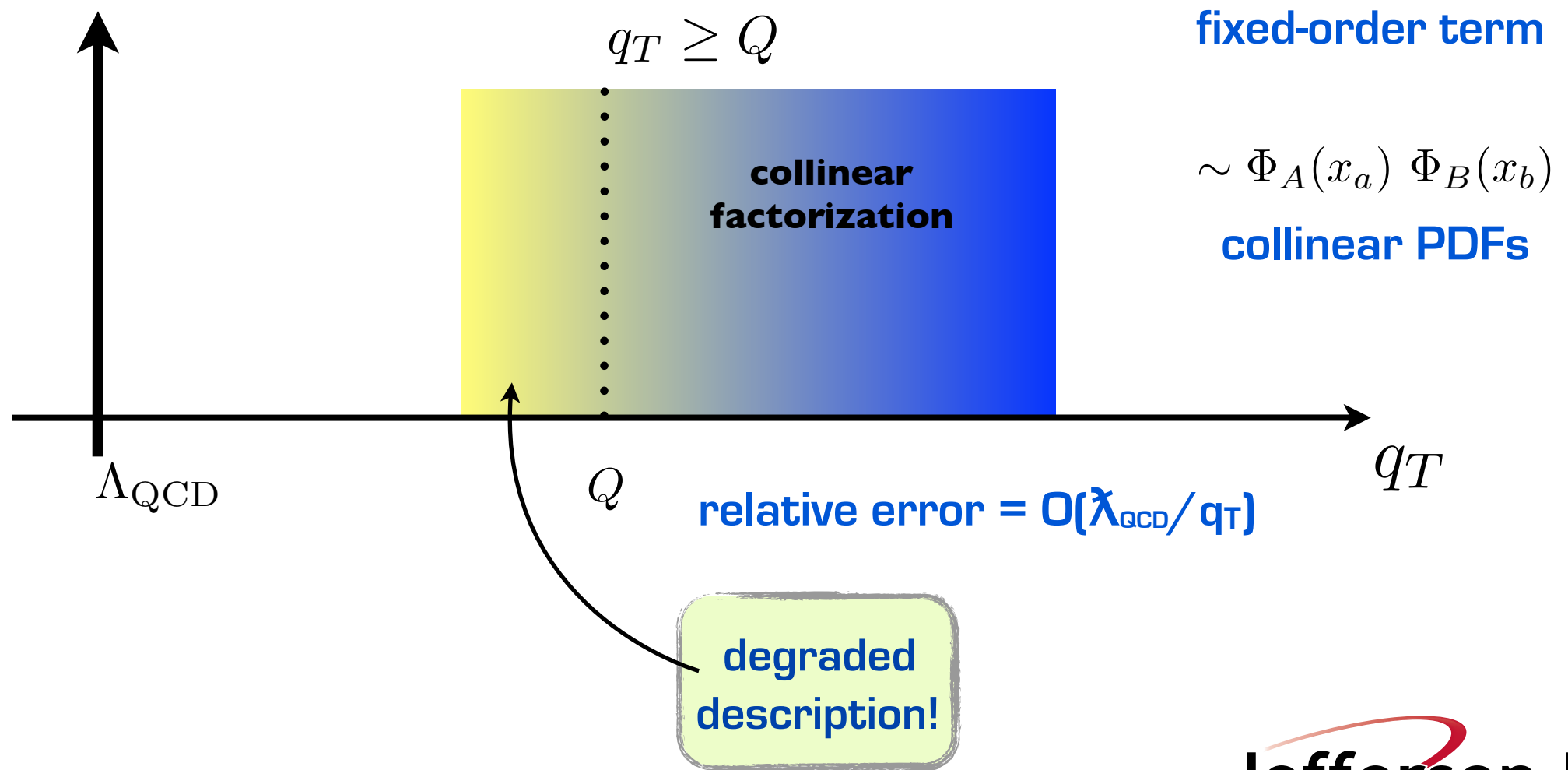
According to their **values** we can access **different**  
“**projections**” of hadron structure

# Collinear and TMD factorization

The key of phenomenology : emergence of TMD and collinear distributions from **factorization theorems**

fixed  $Q$ , variable  $q_T$

$d\sigma/dq_T$

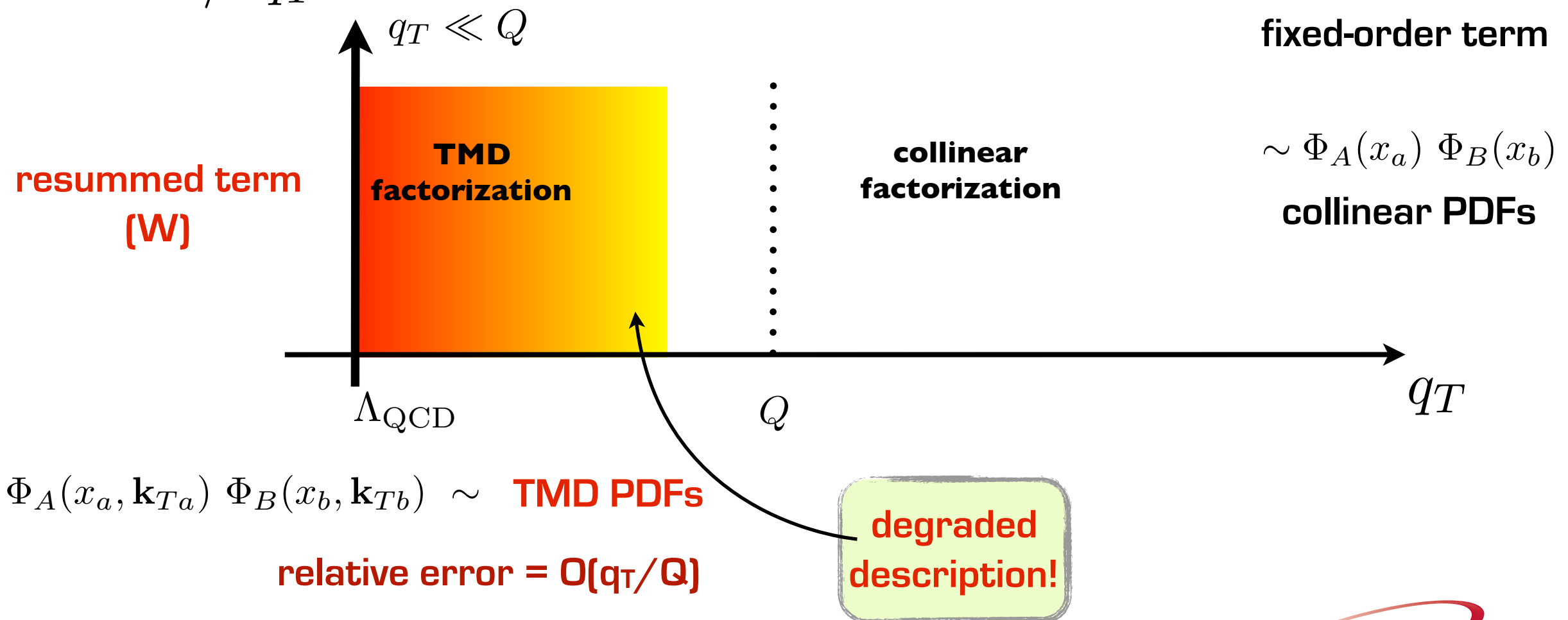


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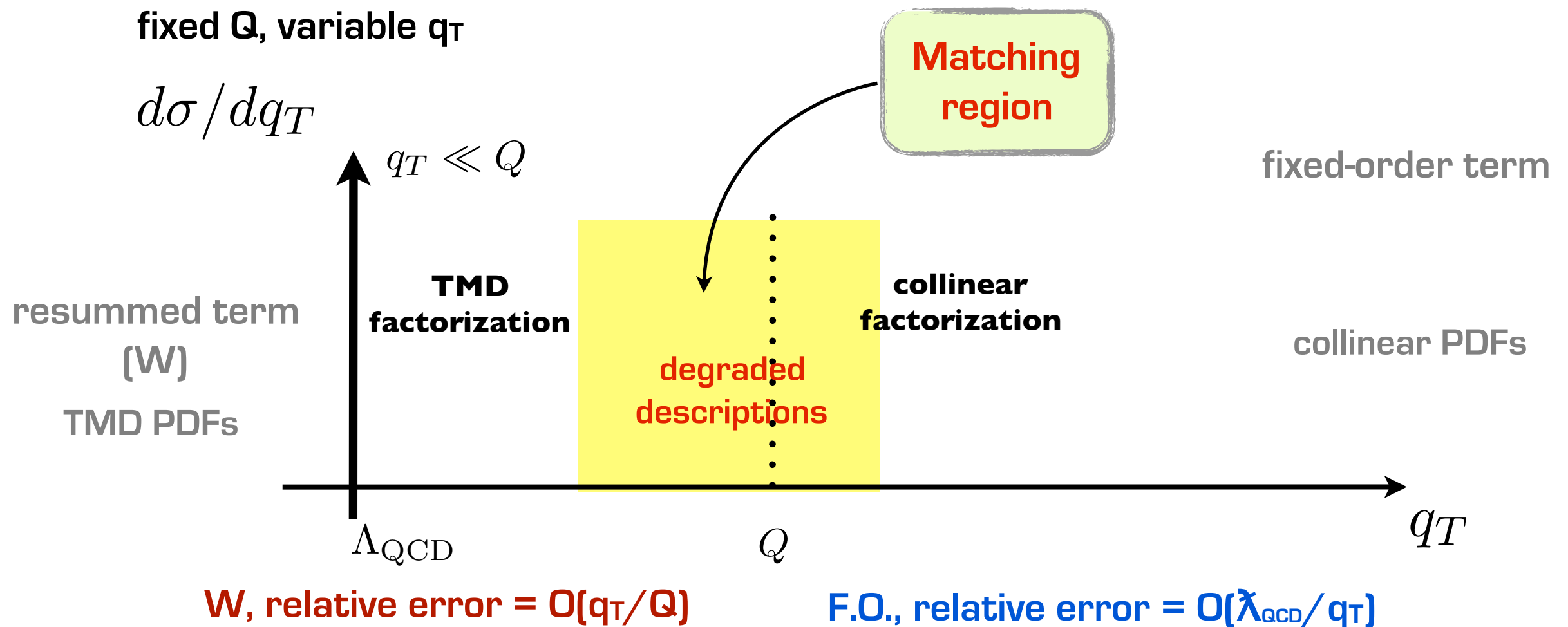
$d\sigma/dq_T$





# Collinear and TMD factorization

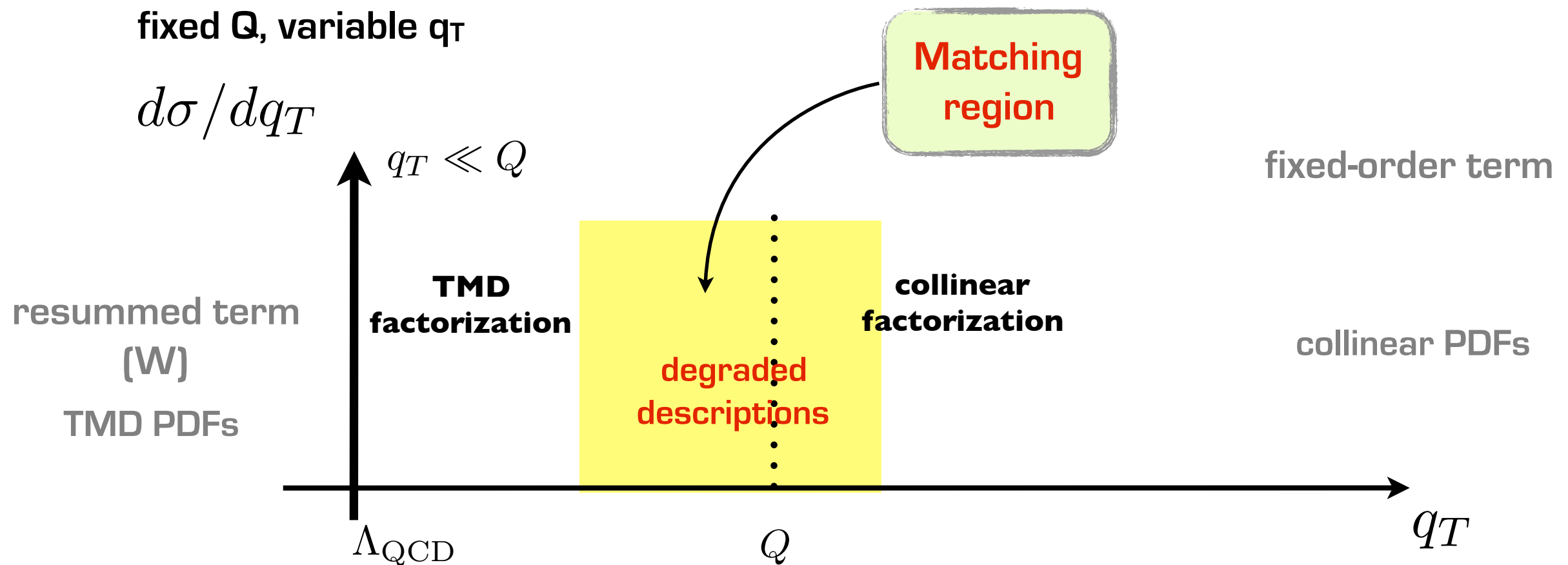
The key of phenomenology : emergence of TMD and collinear distributions from **factorization theorems**



We need a prescription to deal with the region where both descriptions are not good

# Collinear and TMD factorization

The key of phenomenology : emergence of TMD and collinear distributions from **factorization theorems**



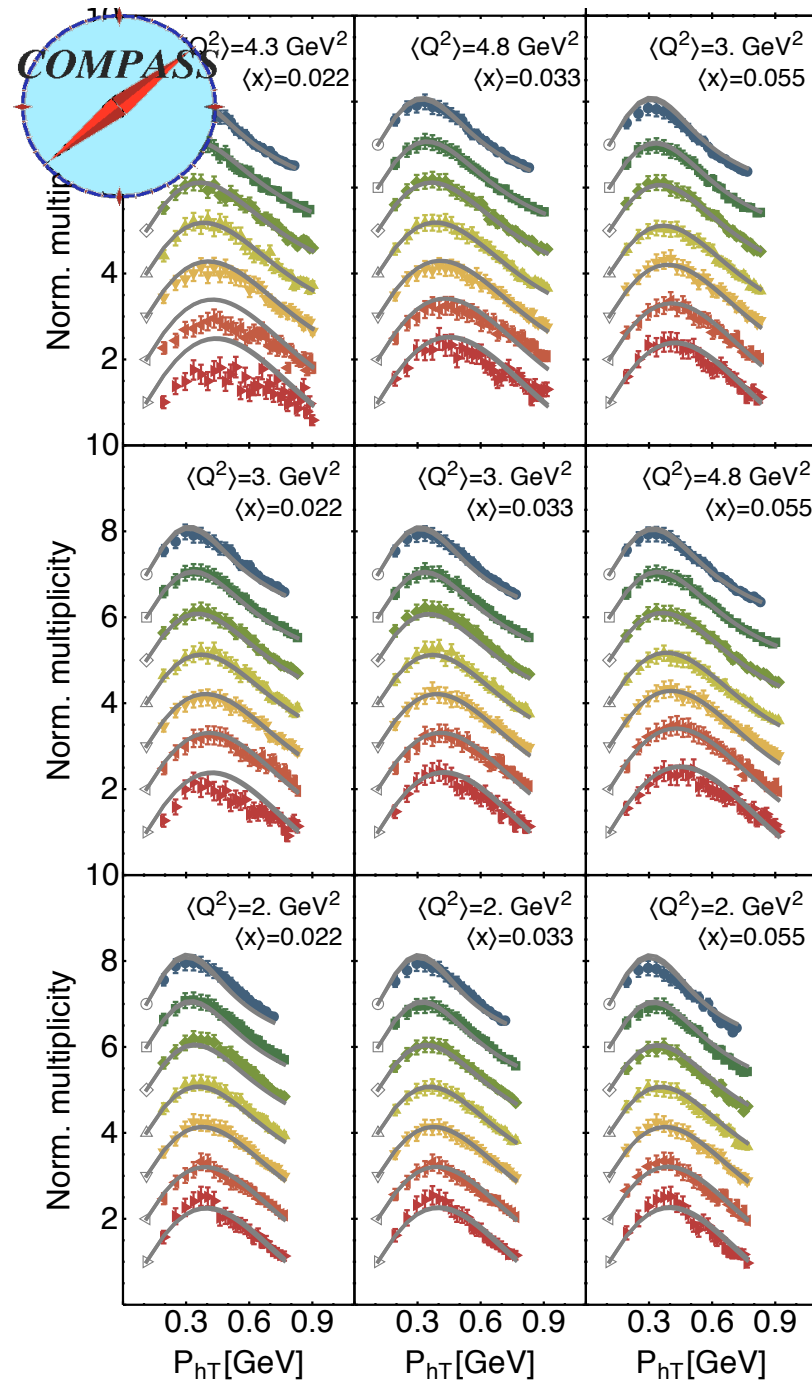
The extraction of the **nonperturbative part of TMDs** is affected by the description of the whole  $q_T$  range

Crucial, especially at **low  $Q$**  (e.g. JLab kinematics), where the **regions shrink**

polarization ?

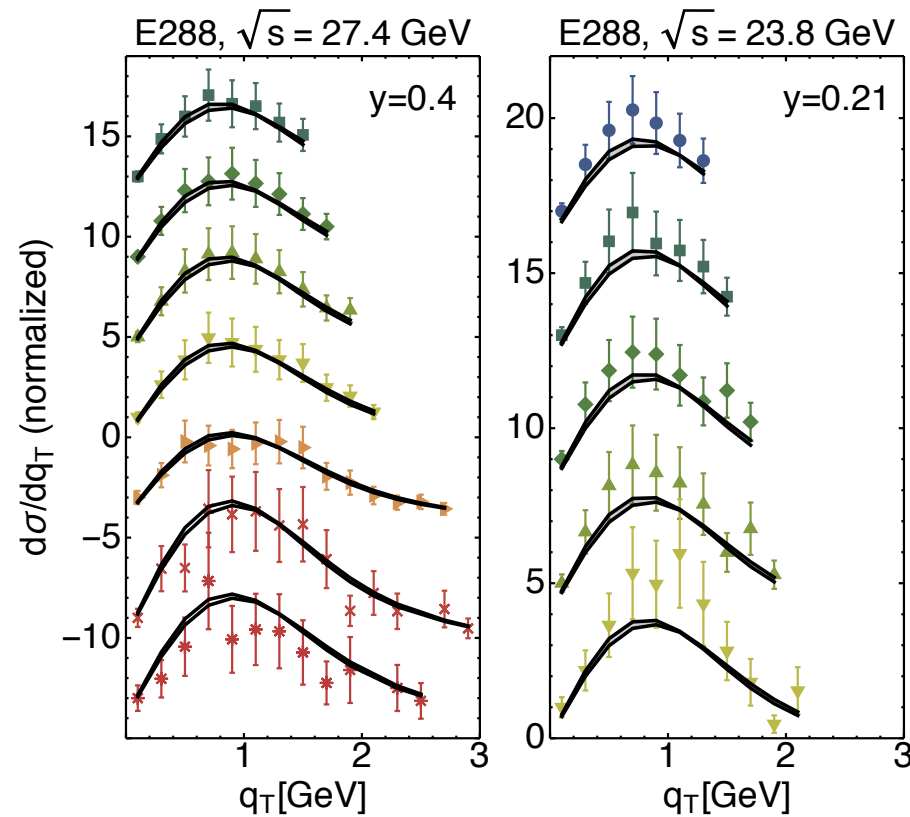
# Global fit

## SIDIS

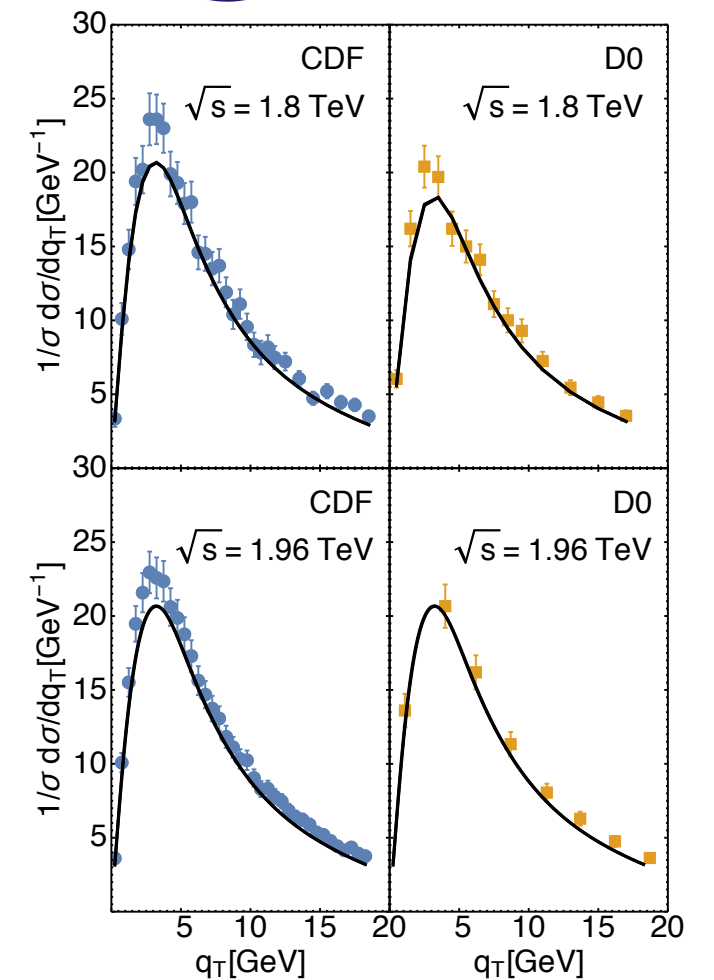
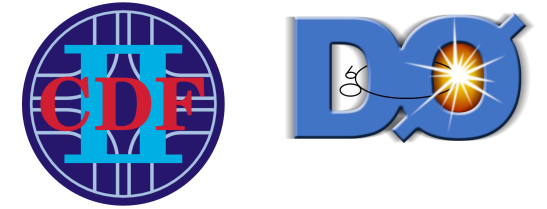


## Drell-Yan

 **Fermilab**



## Z production



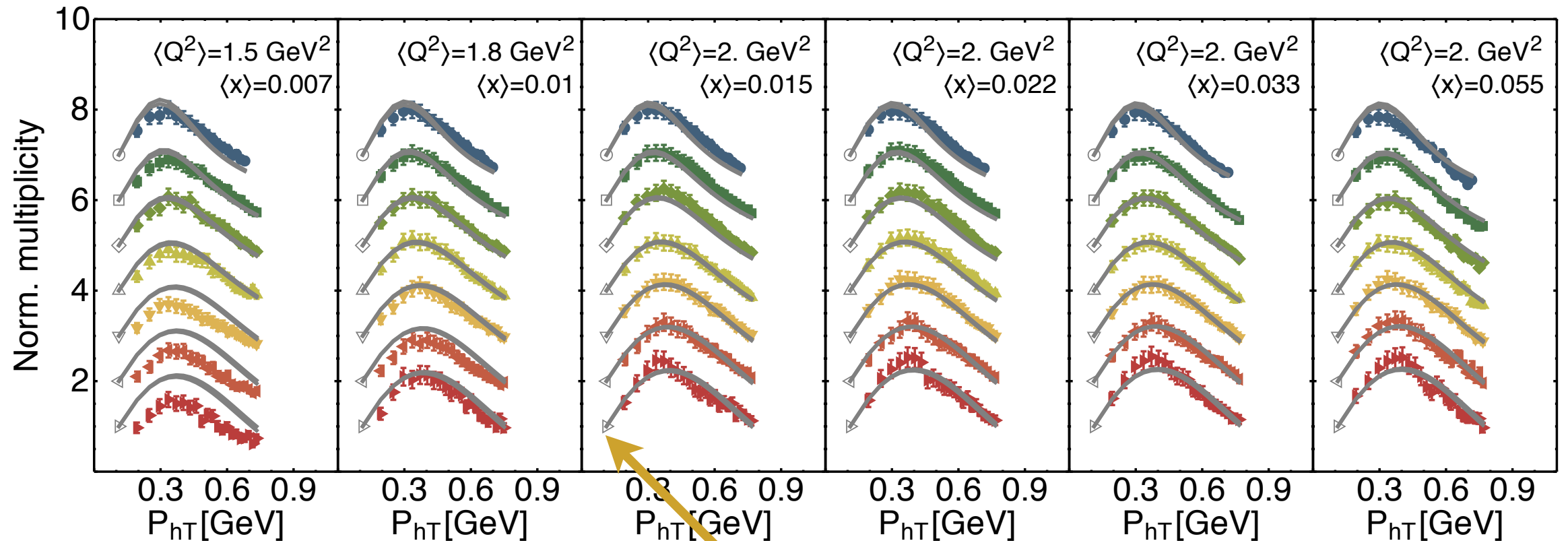
Bacchetta et al. **JHEP 1706 (2017) 081**

 **Jefferson Lab**

# COMPASS, selected bins



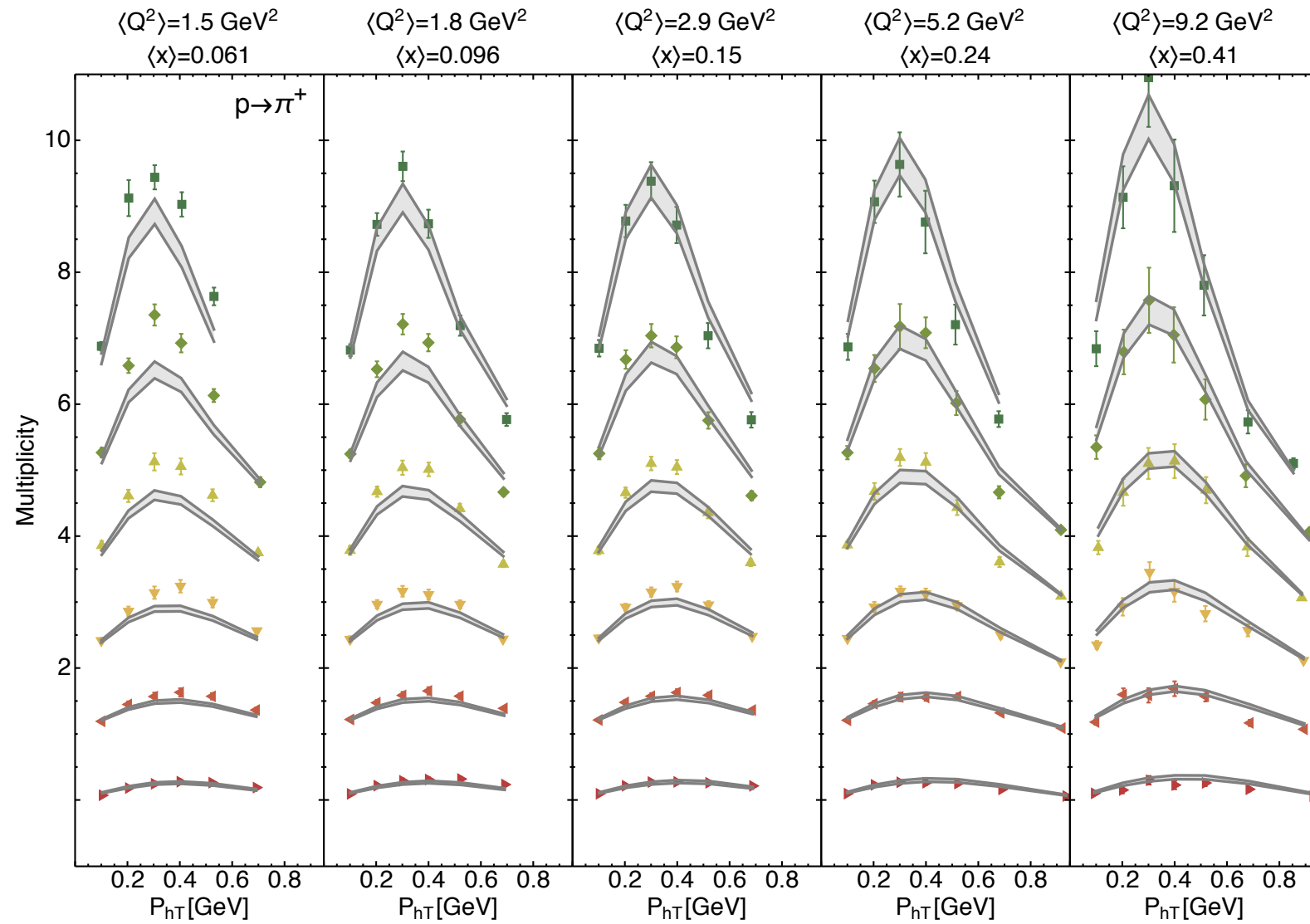
- $\langle z \rangle = 0.23$  (offset=6)
- $\langle z \rangle = 0.28$  (offset=5)
- ◆  $\langle z \rangle = 0.33$  (offset=4)
- ▲  $\langle z \rangle = 0.38$  (offset=3)
- ▼  $\langle z \rangle = 0.45$  (offset=2)
- ▲  $\langle z \rangle = 0.55$  (offset=1)
- ▼  $\langle z \rangle = 0.65$  (offset=0)



First points are not fitted, but used as normalization

Deuteron  $h^-$   $\chi^2/\text{dof} = 1.58$

# HERMES, selected bins



$$\chi^2/\text{dof} = 4.80$$

The worst of all channels...

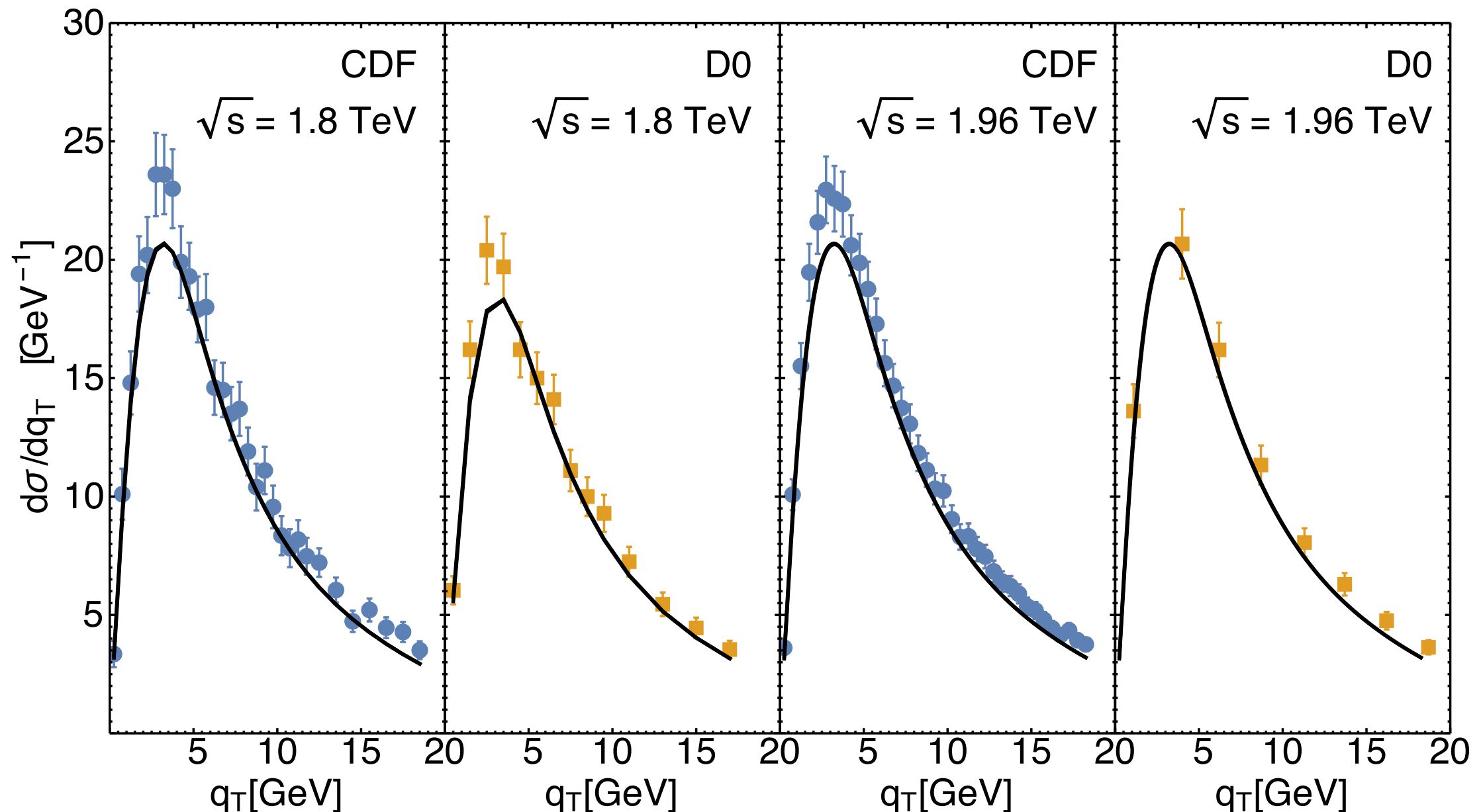
However **normalizing** the theory curves to the first bin, without changing the parameters of the fit,  $\chi^2/\text{dof}$  becomes good

Contributions to chi2 mainly from **normalization**, not shape  
(also in Z-boson production)

# Z-boson @ Fermilab

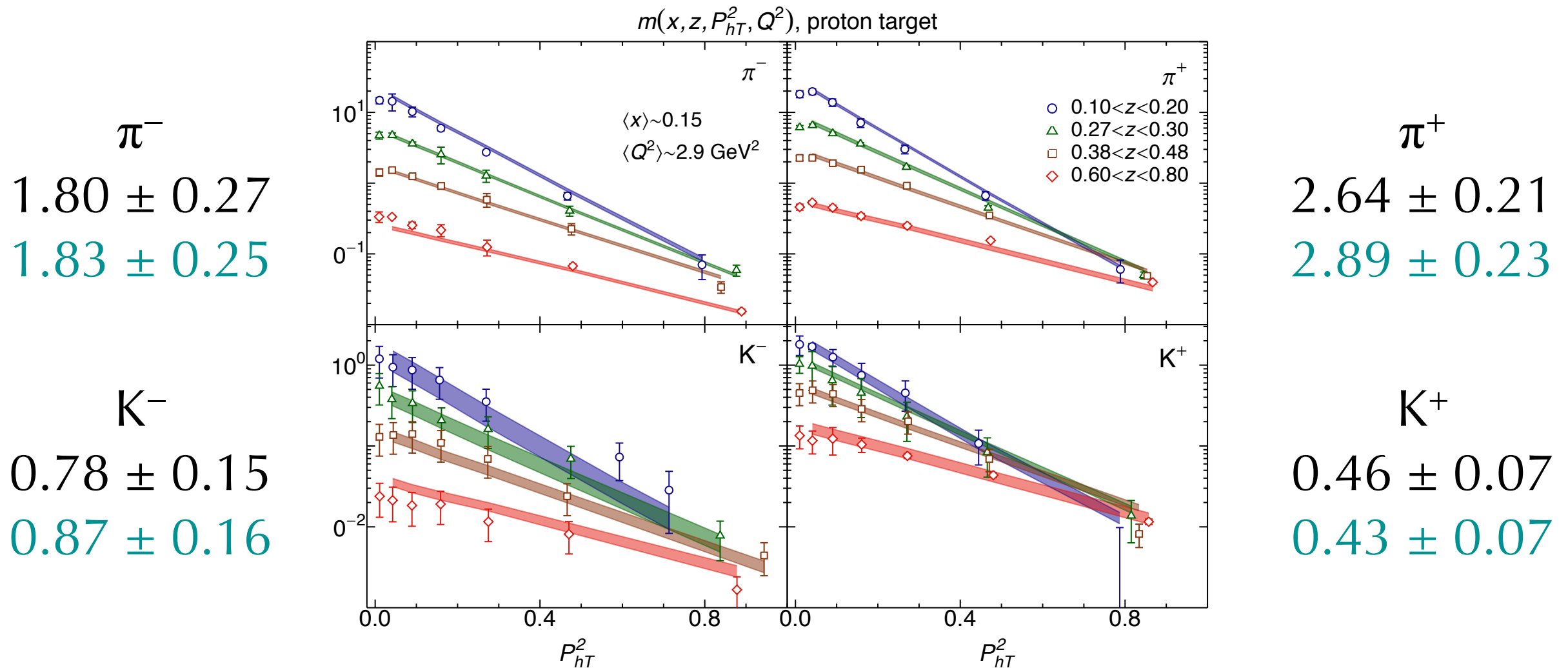
**Narrow bands**, driven mainly by  
 **$g_2$  values** (reduced sensitivity to  
intrinsic  $k_T$ )

Contributions to  $\chi^2$   
mainly from **normalization**,  
not shape



# Pavia / Amsterdam / Bilbao 2013

proton target	global $\chi^2 / \text{d.o.f.} = 1.63 \pm 0.12$
no flavor dep.	1.72 $\pm$ 0.11





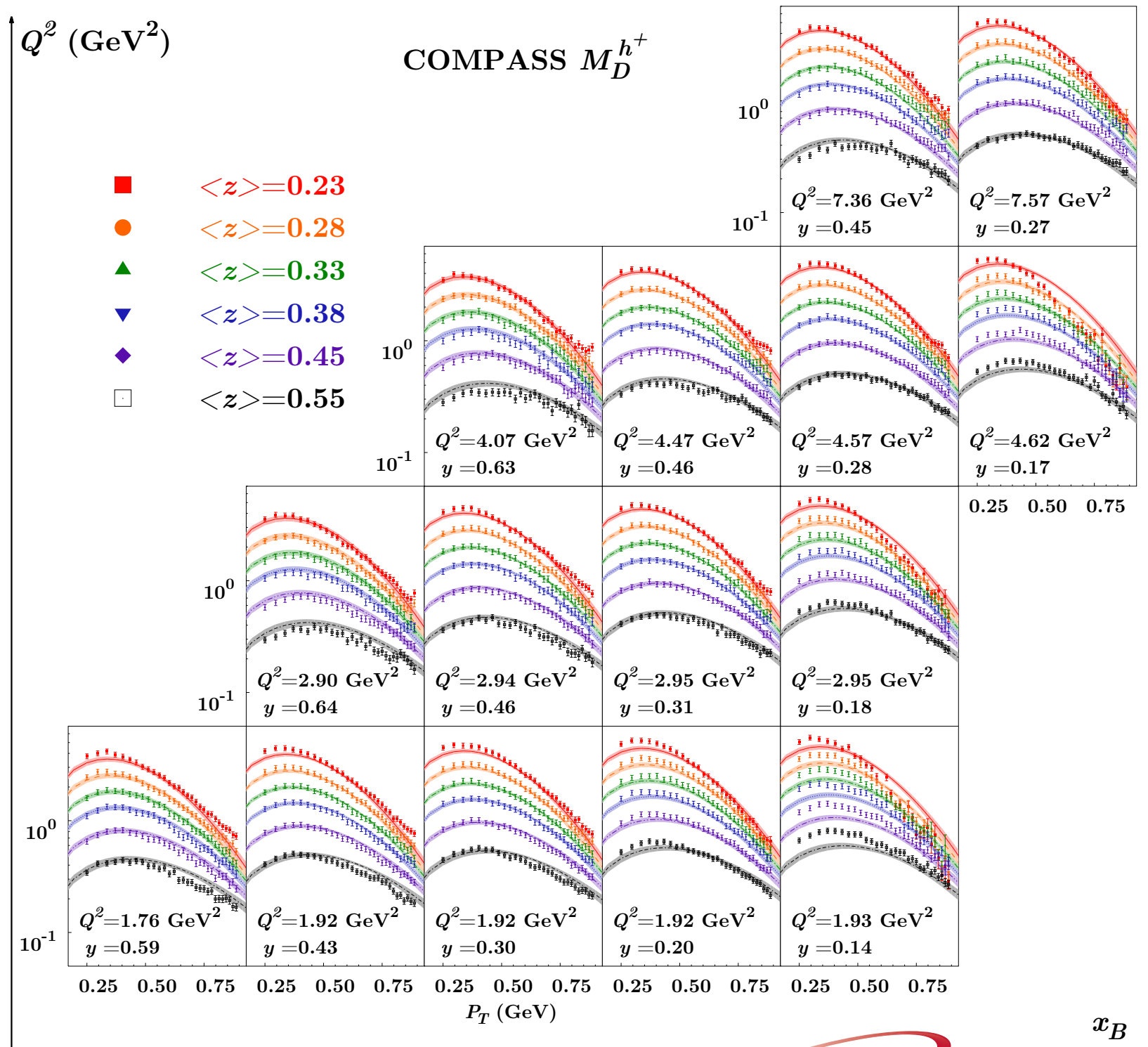
# Torino / JLab 2014

COMPASS  $M_D^{h^+}$

simple Gaussian ansatz

$\chi^2/\text{dof} = 3.79$   
with ad-hoc  
normalization

see Compass coll.  
Erratum

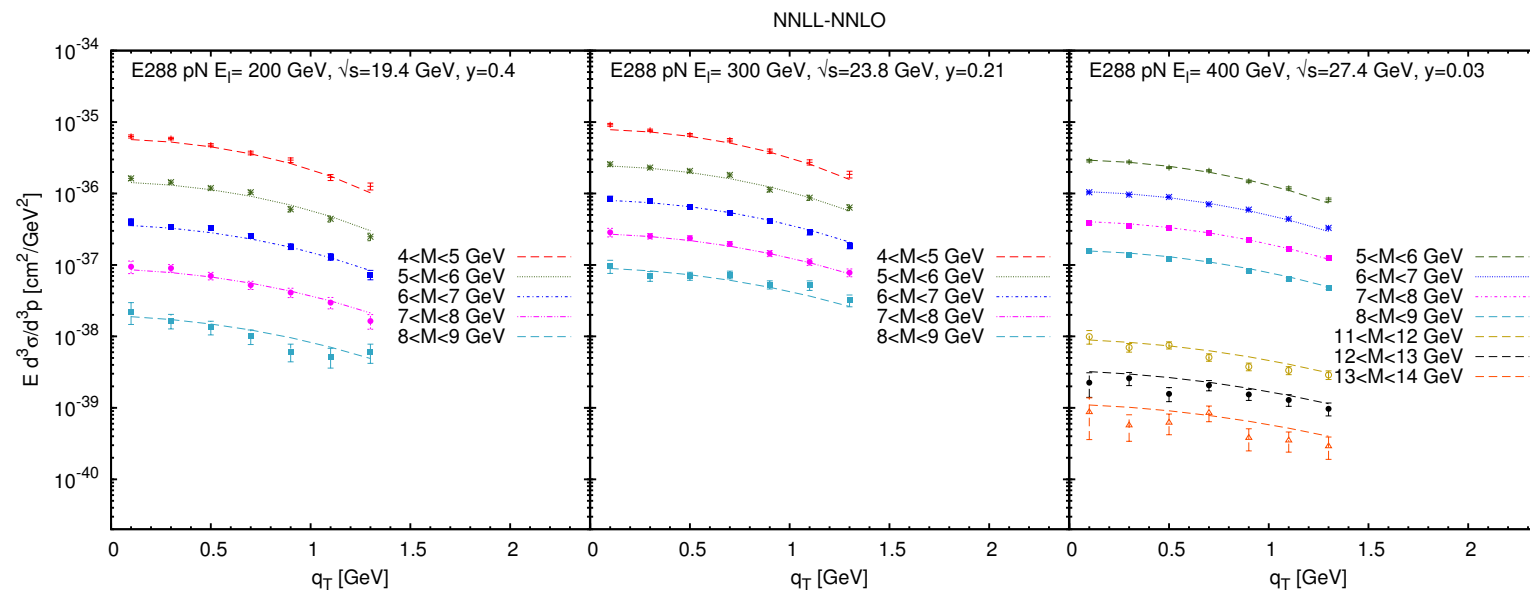
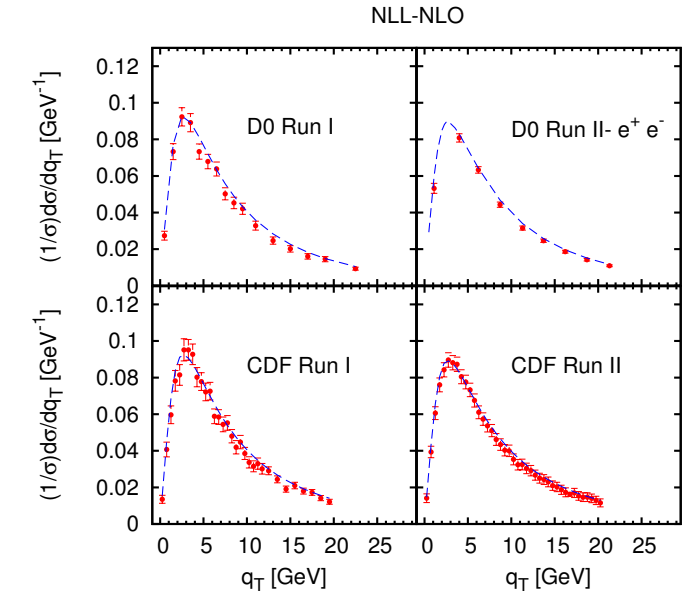
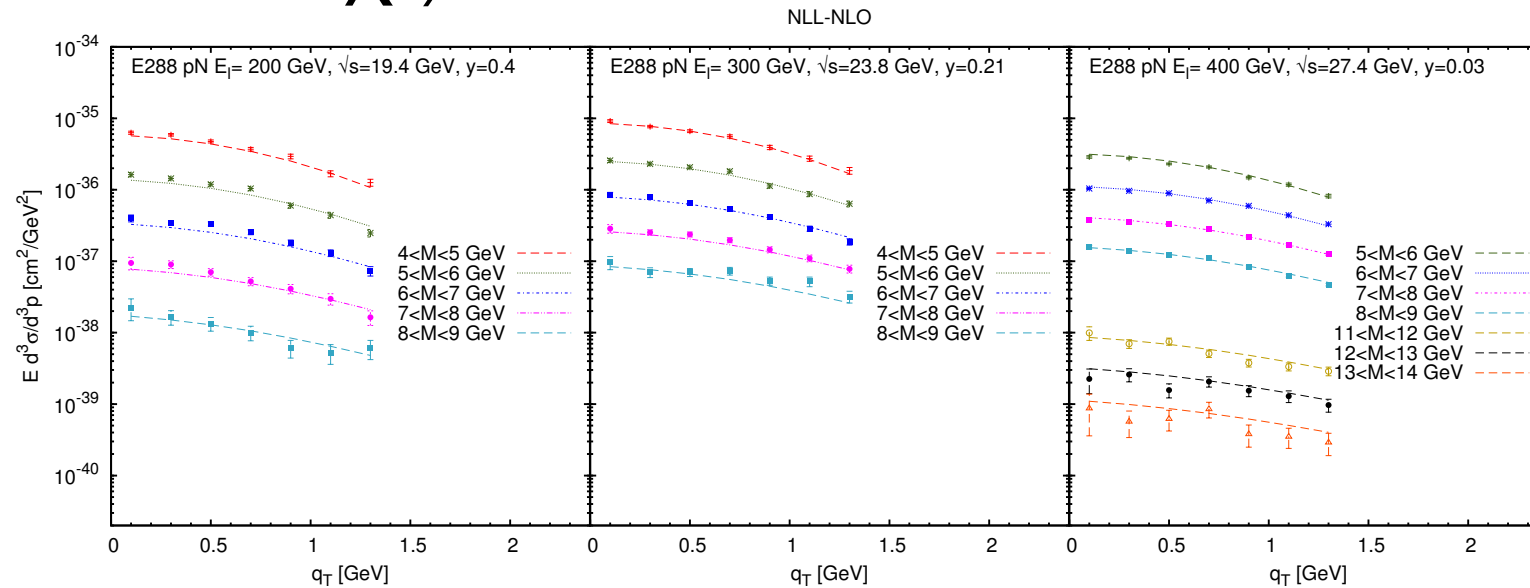




# DEMS 2014

$$\chi^2/\text{dof} = 0.81$$

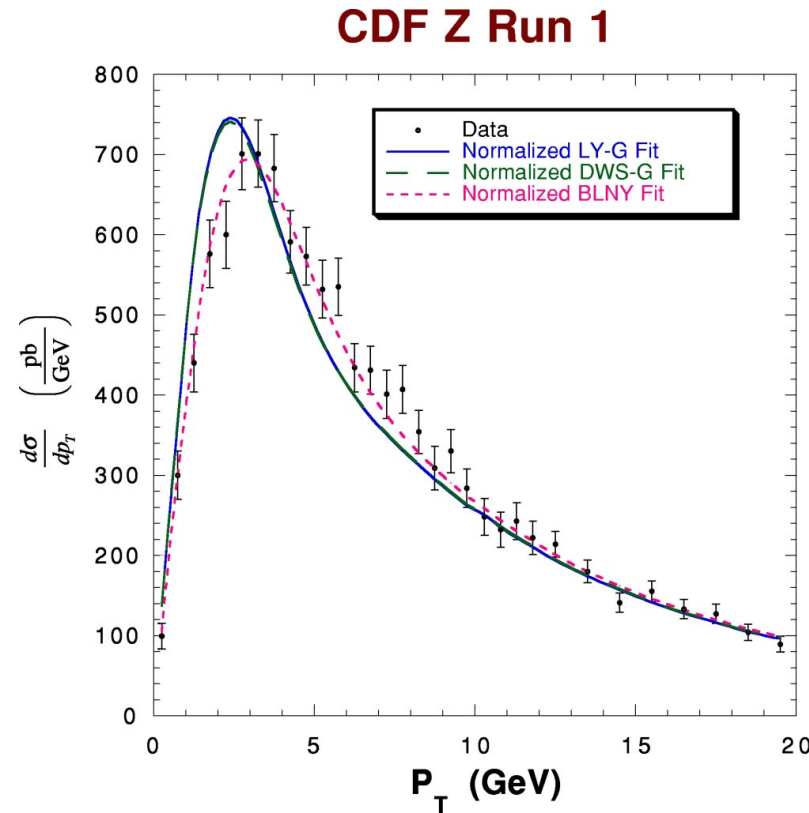
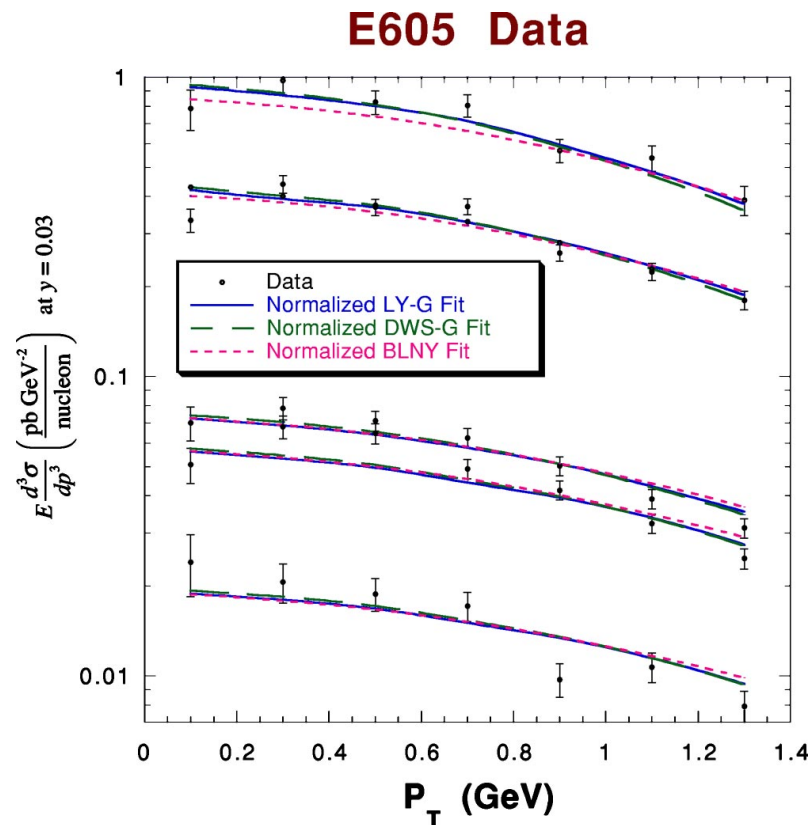
D'Alesio, Echevarria, Melis, Scimemi, JHEP 1411 [14]



NLO-NNLL analysis  
with evaluation of  
theoretical uncertainties

very good

# KN 2006



$\approx 100$  data points  
 $Q^2 > 4 \text{ GeV}^2$

$$Q_0 = 3.2 \text{ GeV}$$

$$b_{\text{max}} = 0.5 \text{ GeV}^{-1}$$

$$\frac{1}{\langle b_T^2 \rangle} = \frac{1}{2} \left( 0.21 + 0.68 \log \left( \frac{Q}{2Q_0} \right) - 0.25 \log(10x) \right)$$

Brock, Landry, Nadolsky, Yuan, PRD67 [03]

$$\frac{1}{\langle b_T^2 \rangle} = \frac{1}{2} \left( 0.20 + 0.184 \log \left( \frac{Q}{2Q_0} \right) - 0.026 \log(10x) \right)$$

$$b_{\text{max}} = 1.5 \text{ GeV}^{-1}$$

KN 2006

# EIKV 2014

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Parametrizations for intrinsic momenta  
and soft gluon emission :

$$F_{NP}(b_T, Q)^{\text{pdf}} = \exp \left[ -b_T^2 \left( g_1^{\text{pdf}} + \frac{g_2}{2} \ln(Q/Q_0) \right) \right]$$

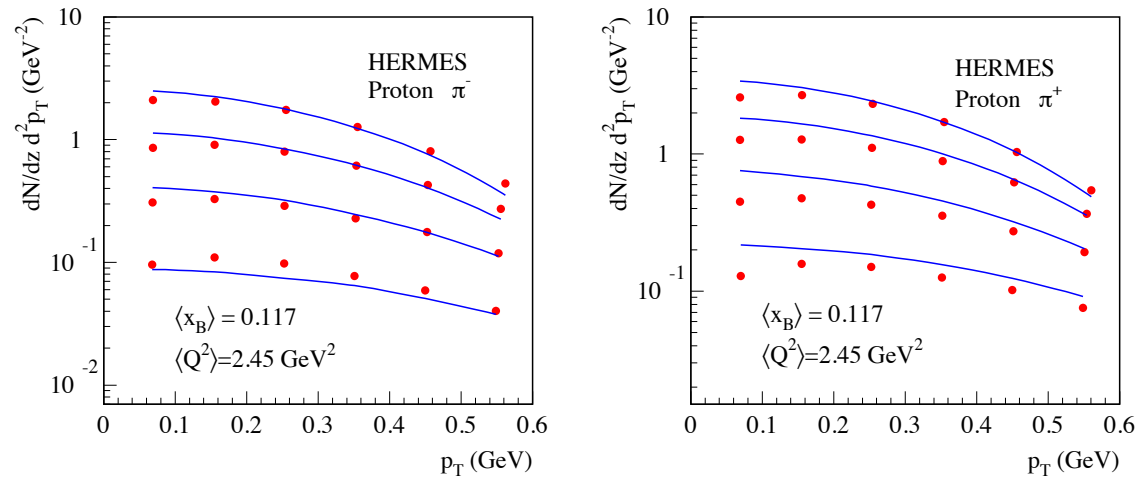
$$F_{NP}(b_T, Q)^{\text{ff}} = \exp \left[ -b_T^2 \left( g_1^{\text{ff}} + \frac{g_2}{2} \ln(Q/Q_0) \right) \right]$$

## Pros and Cons :

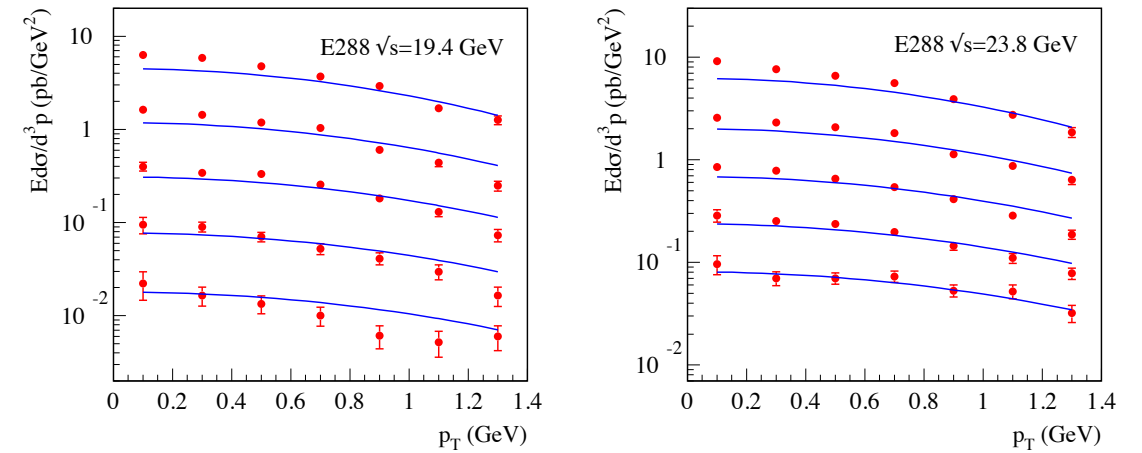
- 1) a global analysis of SIDIS and DY/Z/W data
- 2) TMD evolution at LO-NLL
- 3) multidimensionality not exploited
- 4) chi-square not provided
- 5) can't be considered as a "complete" fit**

# EIKV 2014

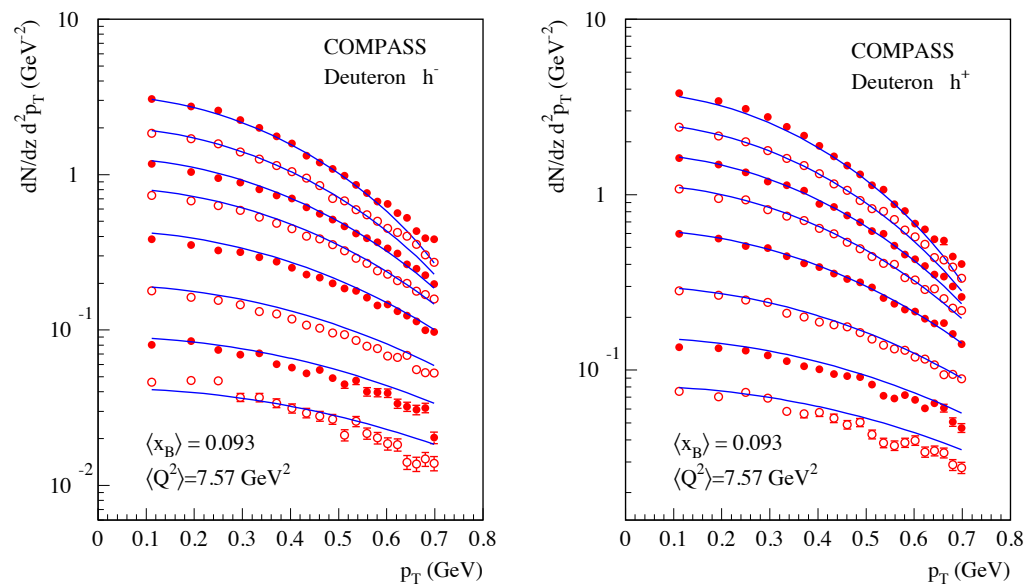
## SIDIS



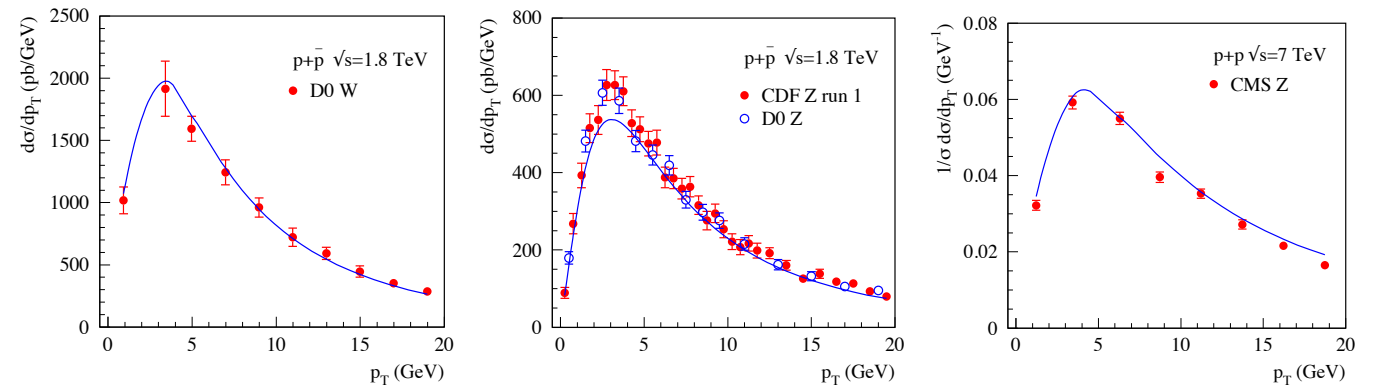
## DRELL-YAN



## SIDIS



## W AND Z PRODUCTION



$$b_{\text{max}} = 1.5 \text{ GeV}^{-1}$$

$$g_2 = 0.16$$

Echevarria et al. [arXiv:1401.5078](https://arxiv.org/abs/1401.5078)



# Other studies

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CSS formalism on DY/Z/W data:

- 1) Davies-Webber-Stirling [DOI: [10.1016/0550-3213\(85\)90402-X](https://doi.org/10.1016/0550-3213(85)90402-X)]
- 2) Ladinsky-Yuan [DOI: [10.1103/PhysRevD.50.R4239](https://doi.org/10.1103/PhysRevD.50.R4239)]
- 3) BLNY [DOI: [10.1103/PhysRevD.63.013004](https://doi.org/10.1103/PhysRevD.63.013004)]
- 4) Hirai, Kawamura, Tanaka [DOI: [10.3204/DESY-PROC-2012-02/136](https://doi.org/10.3204/DESY-PROC-2012-02/136)] - complex-b prescription

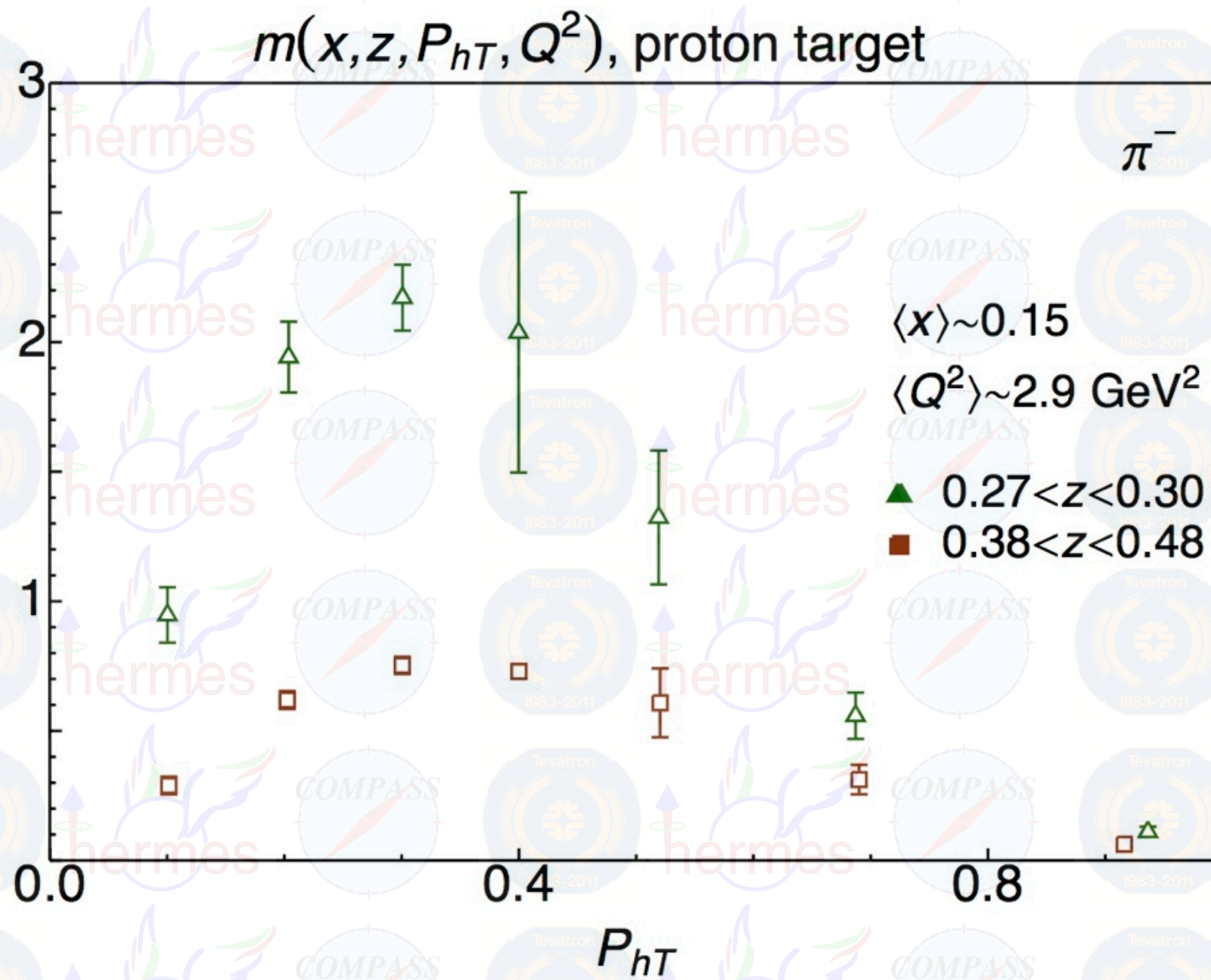
...

combined SIDIS/DY/W/Z :

- 5) Sun, Yuan [[arXiv:1308.5003](https://arxiv.org/abs/1308.5003)]
- 6) Isaacson, Sun, Yuan, Yuan [[arXiv:1406.3073](https://arxiv.org/abs/1406.3073)]

...

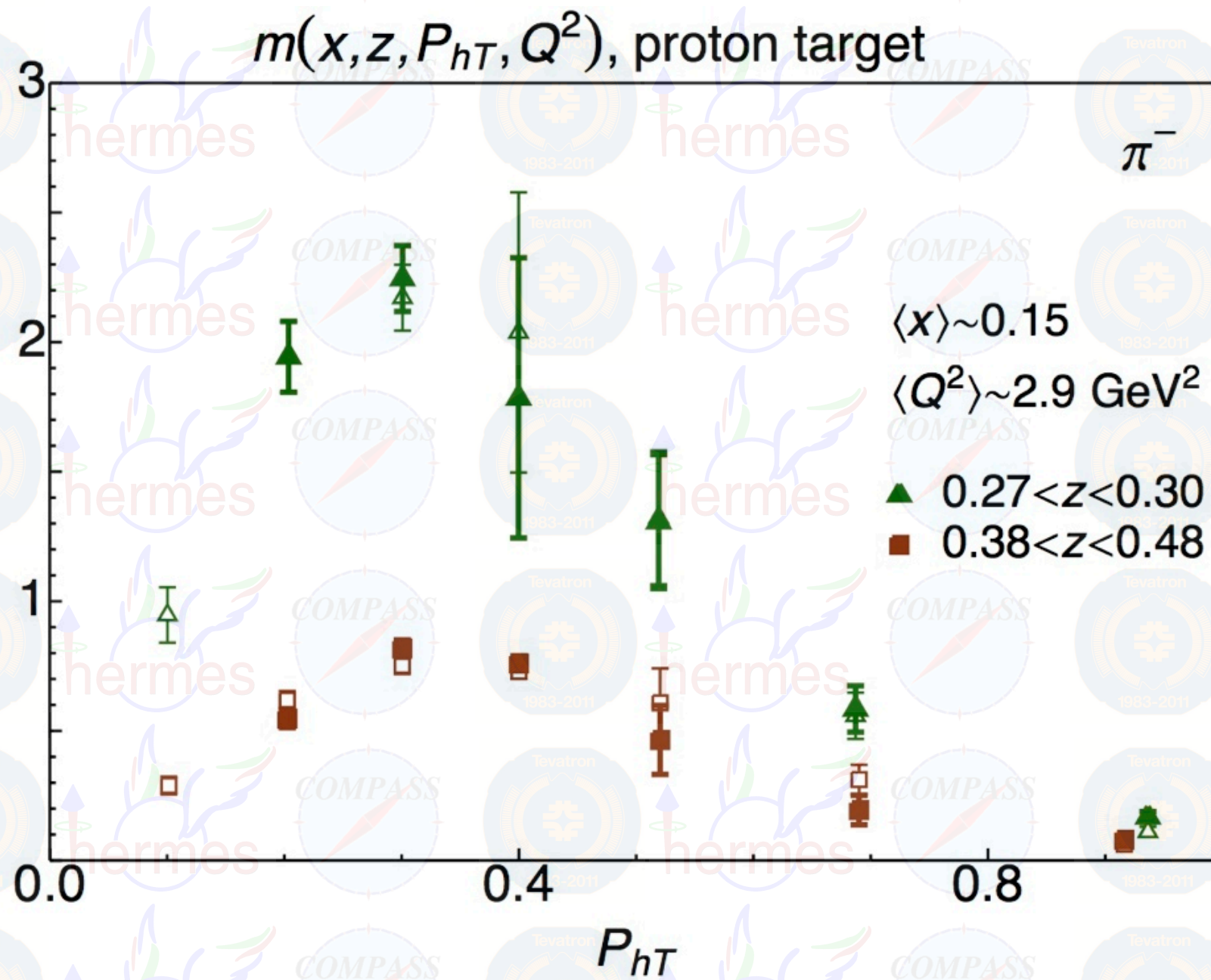
# The replica method



Sample of original data

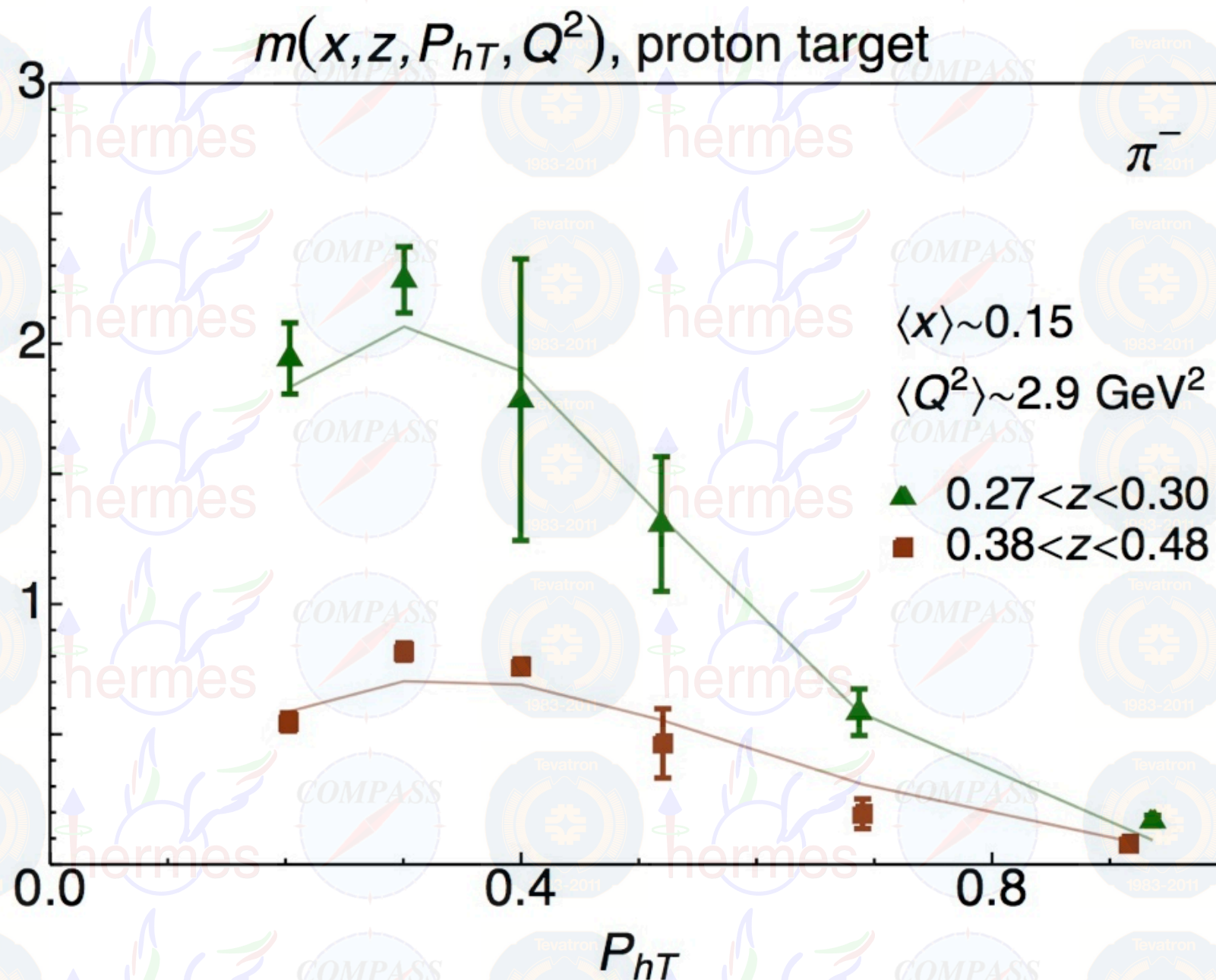


# The replica method



Replica of the original data with Gaussian noise

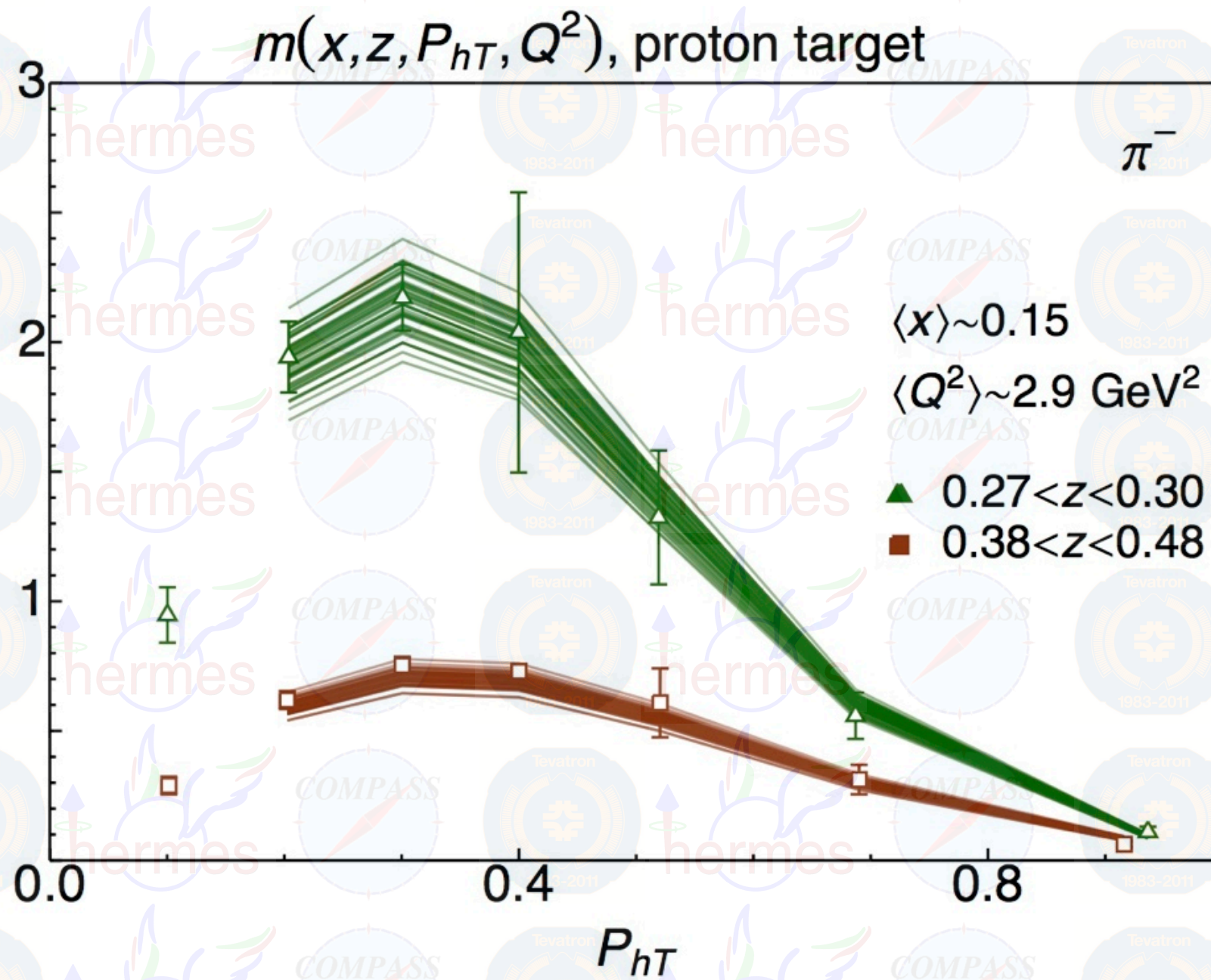
# The replica method



Fit of the replicated data

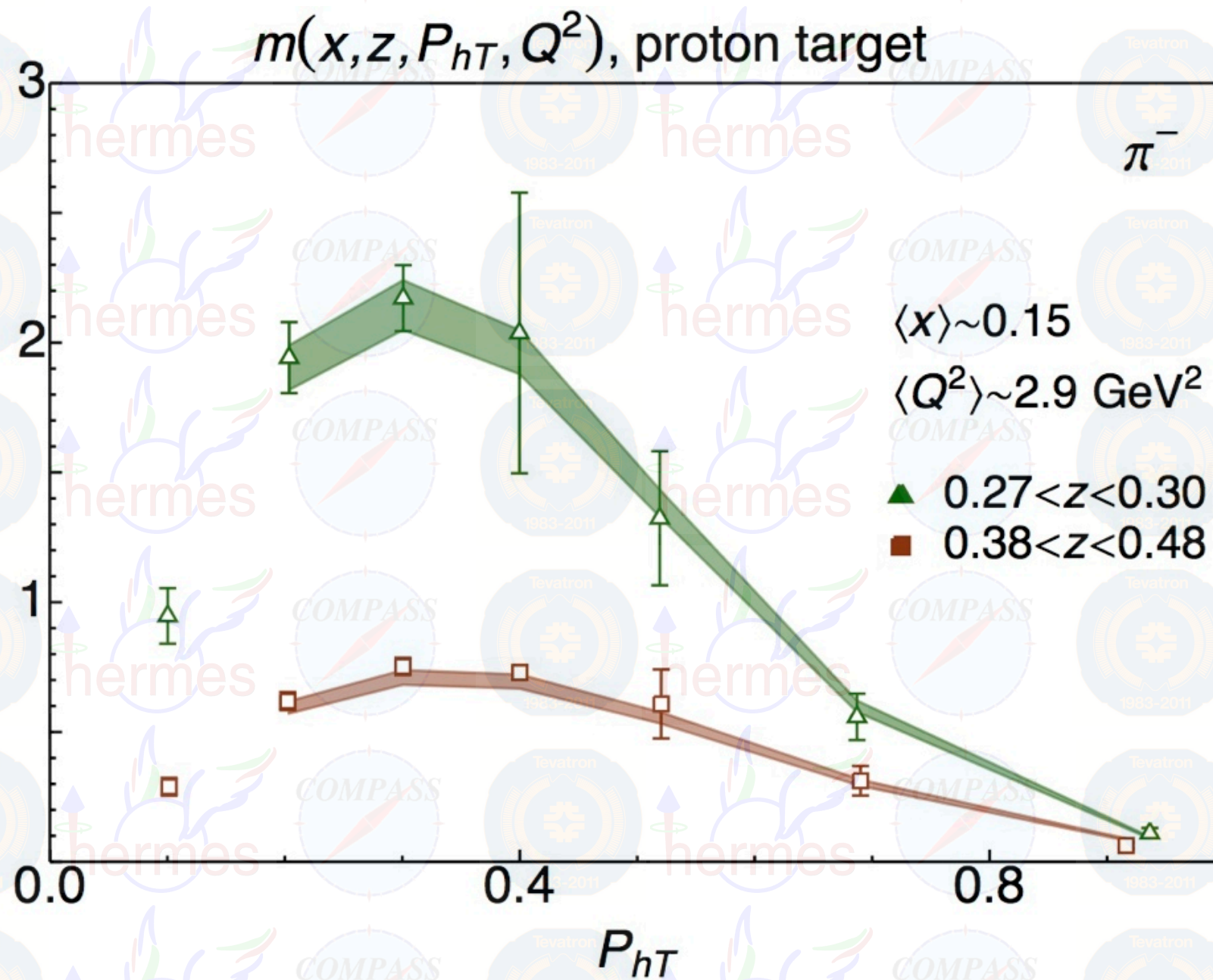


# The replica method



Repeat the generation and the fit N times

# The replica method



Obtain **distributions of best values** -  
calculate **68% CL bands**



