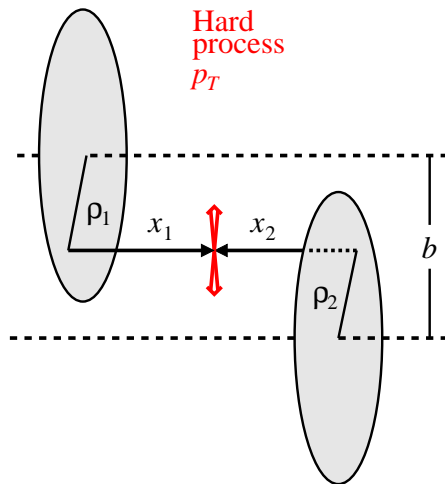
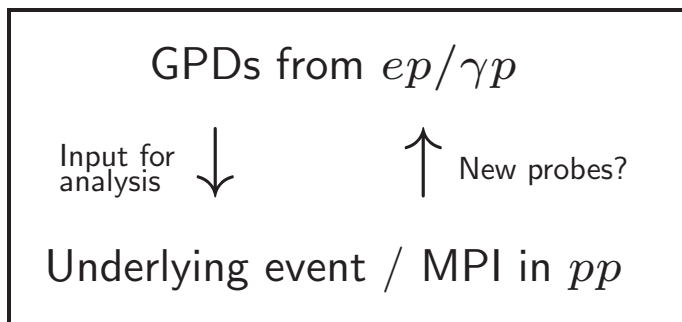


GPDs and transverse geometry in high-energy pp collisions

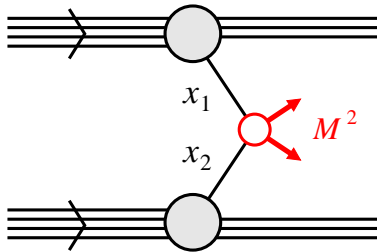
C. Weiss (JLab), INT Program INT-18-3, U. Washington, 04-Oct-2018



- Hard processes in pp
 - Inclusive cross section, factorization
 - Event observables, hard-soft correlations
- Transverse distribution of partons
 - GPDs and exclusive processes in $ep/\gamma p$
- Transverse geometry in pp collisions
 - Hard-soft correlations
 - Multiparton interactions
- Exclusive diffraction $pp \rightarrow p + H + p$
 - Rapidity gap survival probability



Based on work with L. Frankfurt, M. Strikman 2004–

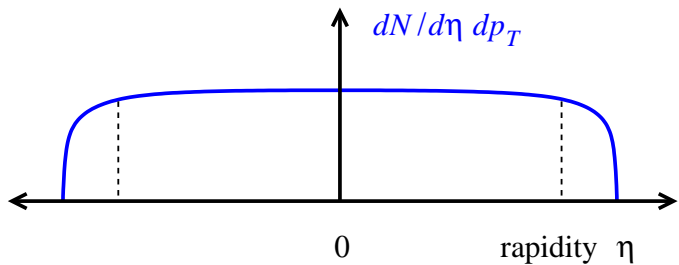


- Inclusive cross section

Separate $k_T^2 \sim \mu^2(\text{soft}) \longleftrightarrow M^2$

$$\sigma = f_1(x_1, \mu^2) f_2(x_2, \mu^2) \times \sigma_{\text{hard}}(\mu^2, M^2)$$

PDF same as in ep



- Underlying event characteristics

Many observables: Particle number distributions, average p_T , energy flow, . . .

Soft interactions mostly

- Hard-soft correlations

How does underlying event depend on hard process and its variables?

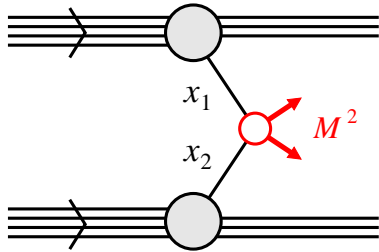
UE with hard processes \leftrightarrow UE with min bias trigger? Change with x_1, x_2, M^2 ?

Interest: Understand strong interaction dynamics in collision, facilitate search for rare hard processes through cuts

No strict factorization, additional assumptions needed

Here: Approach based on transverse geometry. Limited information, but model-independent. Very useful!

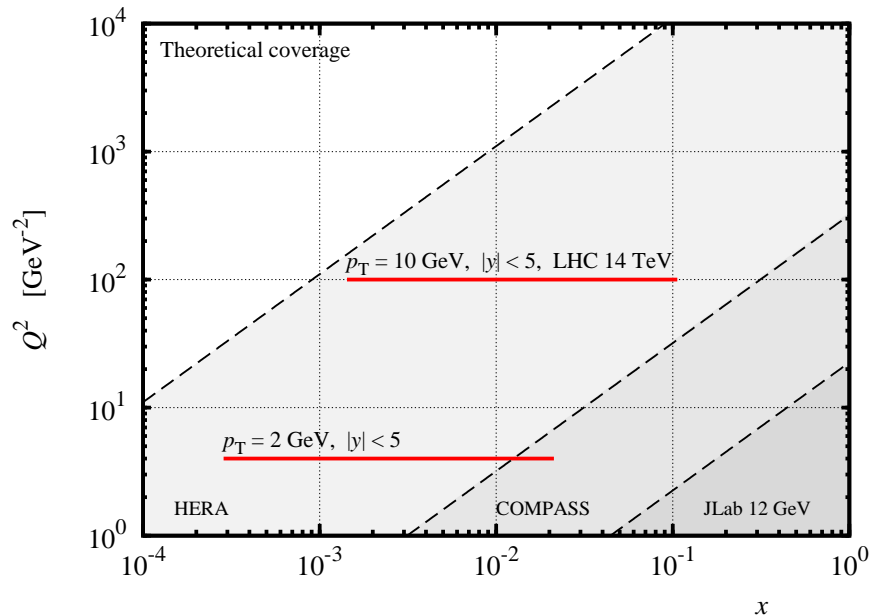
- Intuitive picture



Partonic wave function: Superposition of configurations with different particle number, spatial size, etc.

$$\text{PDF}(x) = \sum_{\text{configs}} [\text{parton with } x] \quad \text{one-body density}$$

Underlying event \leftrightarrow interactions of spectators

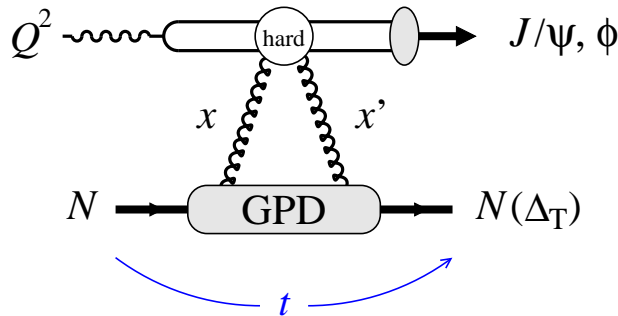


- Hard process kinematics

$$x_1 x_2 = \frac{M^2}{s}$$

$$M^2 = 4p_T^2 \quad \text{back-to-back dijet}$$

Typical x -values $\sim 10^{-1} - 10^{-3}$
HERA/COMPASS/EIC region



- Hard exclusive meson production

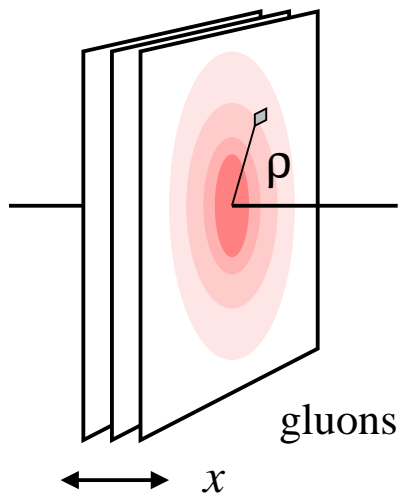
Meson produced in small-size $q\bar{q}$ configuration

QCD factorization theorem $Q^2 \gg \mu_{\text{had}}^2 \sim |t|$
 Collins, Frankfurt, Strikman 96

GPDs: Partonic form factor of nucleon, universal, process-independent

Ji 96, Radyushkin 96

Operator definition $\langle N' | \text{twist-2} | N \rangle$, renormalization, non-pert. methods



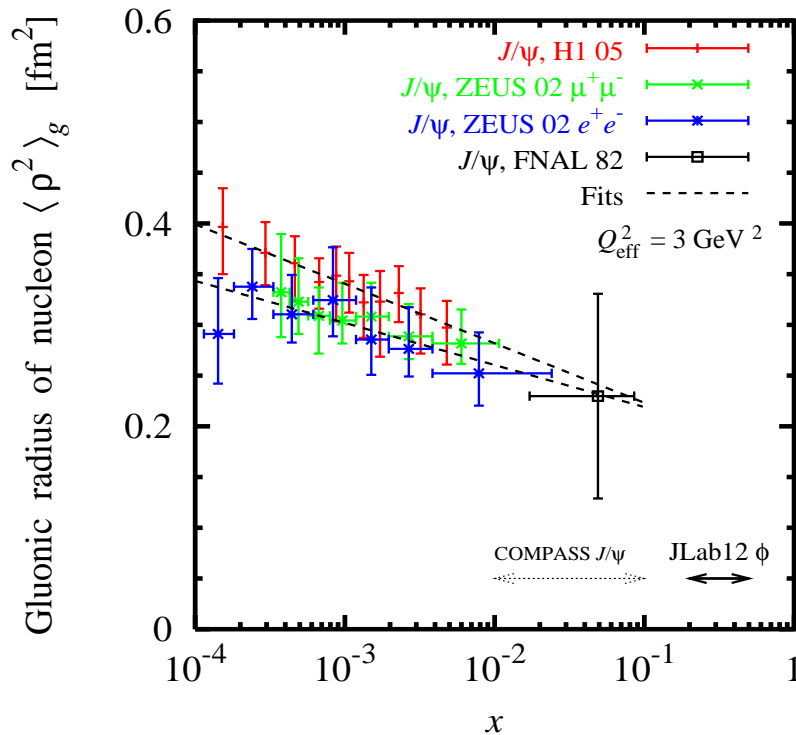
- Transverse spatial distribution of partons $x' = x$

$$f(x, \boldsymbol{\rho}) = \int \frac{d^2 \Delta_T}{(2\pi)^2} e^{-i\Delta_T \boldsymbol{\rho}} \text{GPD}(x, t) \quad \text{2D Fourier}$$

Tomographic image of nucleon at fixed x , changes with x and Q^2

- Large x : Quark GPDs, polarization, $x' \neq x$

JLab12: DVCS, meson production



- Transverse distribution of gluons

Exclusive J/ψ at HERA, also ϕ, ρ
 Large x : FNAL, COMPASS, JLab12 ϕ

Transverse profile from relative t -dep.

Gluonic radius from slope $\langle \rho^2 \rangle_g = 2B_{\text{excl}}$

- Important observations

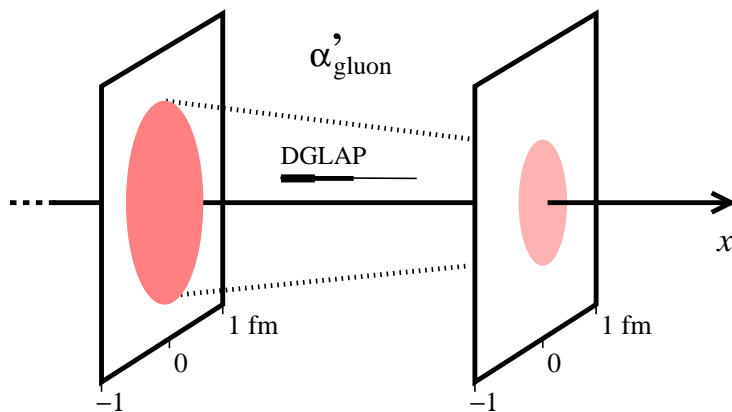
Gluonic radius $\langle \rho^2 \rangle_g$ much smaller than soft nucleon radius $\sim 1 \text{ fm}^2$

Grows with effective Regge slope
 $\alpha'_g \approx 0.14 \text{ GeV}^{-2} < \alpha'_{\text{soft}}$

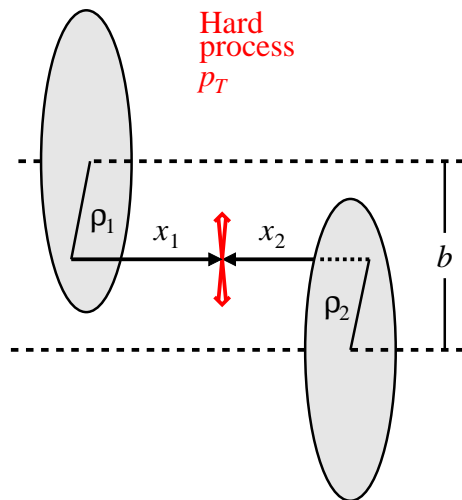
- Q^2 dep. from DGLAP evolution

Partons decay locally in transverse space

Size changes because initial partons at $x_0 > x$ sit at smaller transv. distances.
 Small effect at $Q^2 > \text{few GeV}^2$



Transverse geometry in pp: Hard processes

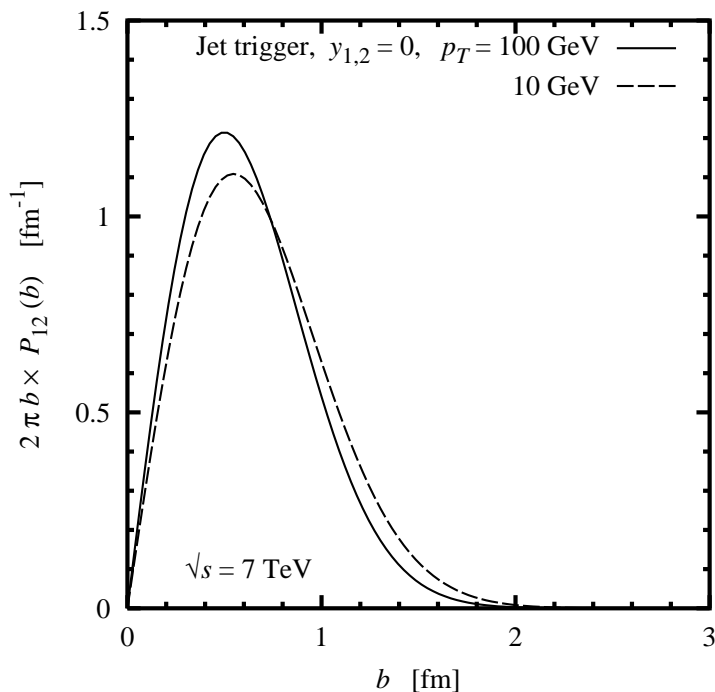


- Hard process from parton-parton collision
 Local in transverse space $p_T^2 \gg (\text{transv. size})^{-2}$
- Cross section as function of pp impact par

$$\sigma_{12}(b) = \int d^2\rho_1 d^2\rho_2 \delta(\mathbf{b} - \boldsymbol{\rho}_1 + \boldsymbol{\rho}_2) \times G(x_1, \rho_1) G(x_2, \rho_2) \sigma_{\text{parton}}$$

Calculable from known transverse distributions
 Integral $\int d^2b$ reproduces inclusive formula

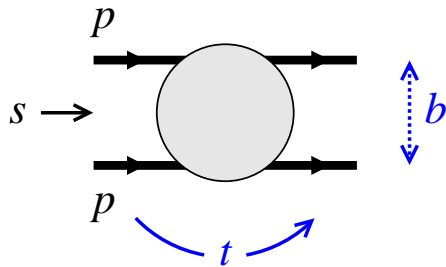
Normalized distribn $P_{12}(b) = \sigma_{12}(b) / [\int \sigma_{12}]$



- New information available
 - Model spectator interactions depending on b
 Underlying event
 - Predict probability of multiple hard processes
 Dynamical correlations? FSW04
 - Diffraction: Gap survival probability
 Determined largely by transverse geometry FHSW 07

Transverse geometry in pp: Soft interactions

- pp elastic scattering amplitude



$$A(s, t) = \frac{is}{4\pi} \int d^2b e^{-i\Delta_T b} \Gamma(s, b)$$

impact parameter representation ($t = -\Delta_T^2$)

Data ISR, Tevatron, LHC; empirical parametrizations

- Cross sections

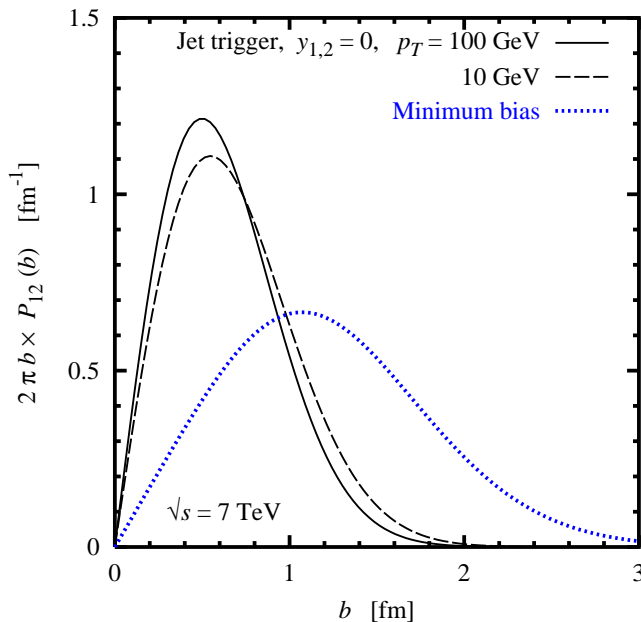
$$\sigma_{\text{el}}(s) \sim |A|^2 = \int d^2b |\Gamma(s, b)|^2 \quad \text{elastic}$$

$$\sigma_{\text{tot}}(s) \sim \text{Im } A = \int d^2b 2 \text{Re } \Gamma(s, b) \quad \text{total}$$

↓

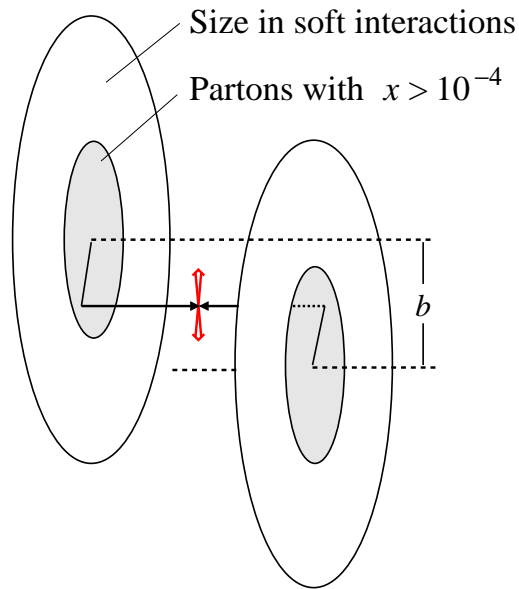
$$\begin{aligned} \sigma_{\text{in}}(s) &= \int d^2b [2 \text{Re } \Gamma - |\Gamma|^2] \\ &= \dots [1 - |1 - \Gamma|^2] \quad \text{inelastic} \end{aligned}$$

Normalized b -distributions



- Impact parameter distn in min bias events

Transverse geometry in pp: Hard vs soft



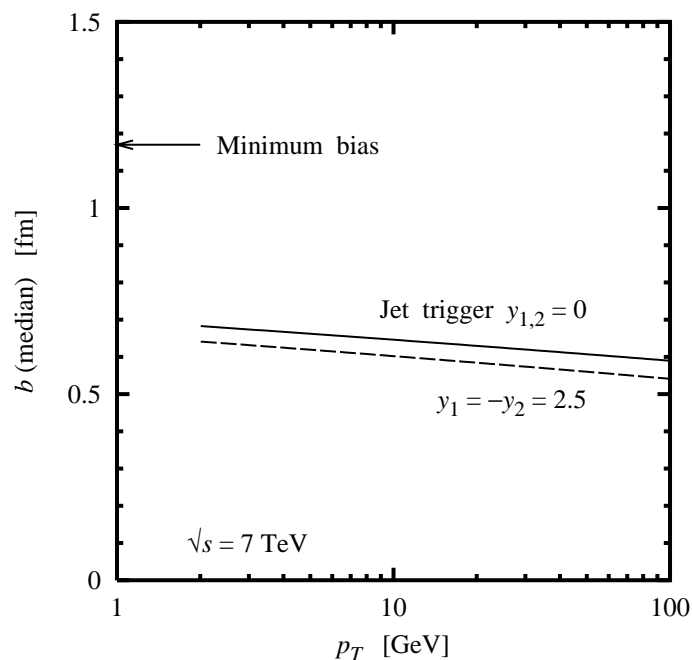
- Transverse proton size in soft interactions much larger than in hard processes

$$R^2(\text{soft}) \gg \langle \rho^2 \rangle_g(x > 10^{-4}) \quad \text{two scales!}$$

- Two classes of pp collisions

Peripheral: Most of inelastic cross section

Central: High probability for hard process



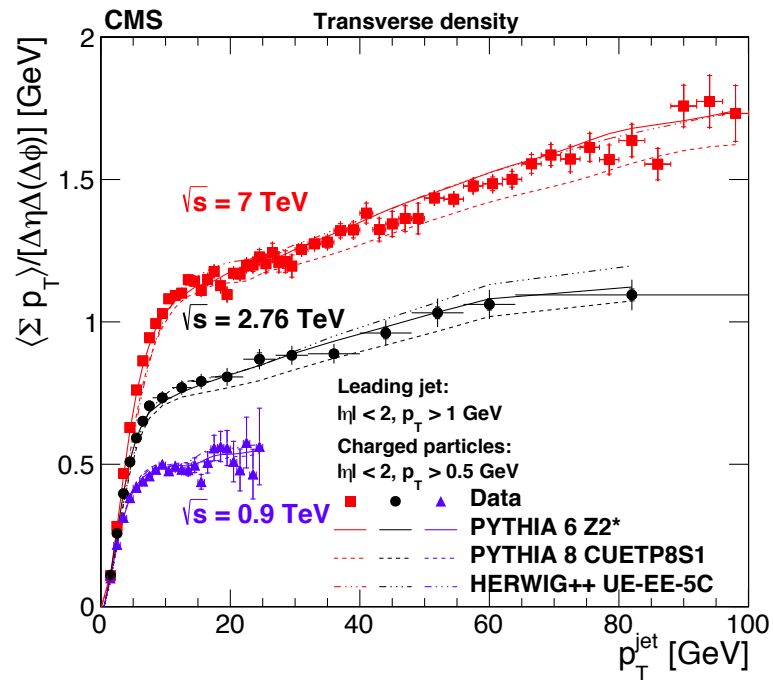
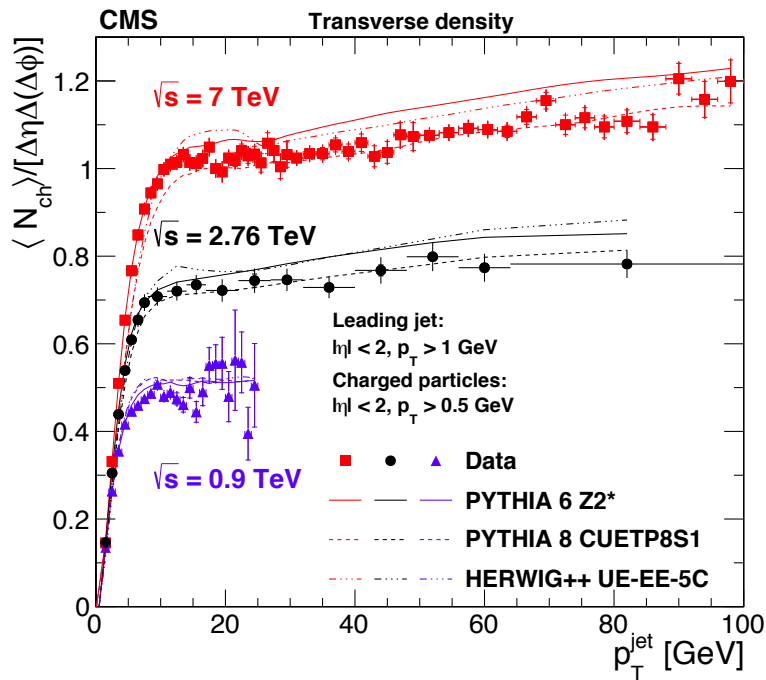
- Hard processes select central collisions

Underlying event in hard processes very different from min. bias collisions

Geometric correlations:
Hard process \leftrightarrow centrality \leftrightarrow event chars

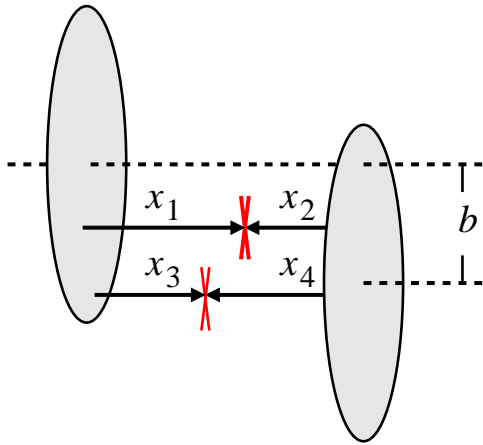
New tests of dynamical mechanisms in particle production

Transverse geometry in pp: Hard-soft correlations 9



CMS underlying event analysis, JHEP 1509 (2015) 137

- Underlying event activity as function of trigger p_T^{jet}
- $p_T^{\text{jet}} \sim \text{few GeV}$: No hard process, collisions mostly peripheral, low activity
- $p_T^{\text{jet}} \gtrsim 10$ GeV: Hard process, collisions central, high activity.
Little changes with further increase of p_T^{jet} because collision already central
- Geometric correlations — impact parameter as “hidden variable”



- Double collision rate parametrized by σ_{eff}^{-1}
- Mean field approximation

Calculable from transverse distributions

$$\sigma_{\text{eff}}^{-1} (\text{mean field}) = \int d^2b P_{12}(b) P_{34}(b)$$

Reference prediction

$$\langle \rho^2 \rangle_g (x \sim 0.1) \text{ gives } \sigma_{\text{eff}} \sim 34 \text{ mb}$$

- Enhancement observed

CDF/D0 3jet + γ rate about $2 \times$ larger than mean field

LHC MPI results forthcoming

Dynamical explanation? Correlations beyond MF

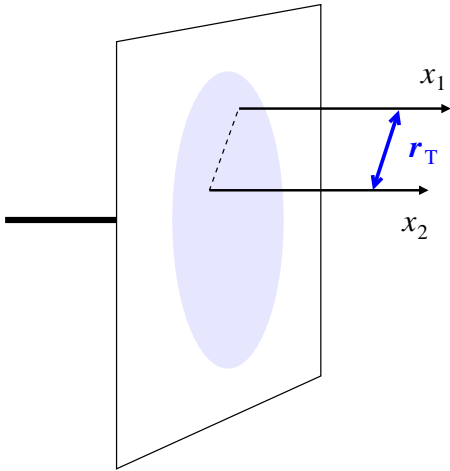
$$\frac{\sigma(12; 34)}{\sigma(12)\sigma(34)} = \frac{1}{\sigma_{\text{eff}}}$$

$$\times \frac{f(x_1, x_3)f(x_2, x_4)}{f(x_1)f(x_2)f(x_3)f(x_4)}$$

- Transverse distributions of partons determine mean field expectation for MPI

Frankfurt, Strikman, CW, PRD 69, 114010 (2004) [INSPIRE]

Multiparton interactions: Dynamical correlations 11



- Parton correlations in nucleon

How is the probability to find a parton influenced by having other parton nearby?

Fundamental property of many-body system: Condensed matter, nuclei

Multiparton distributions

Blok, Dokshitzer, Frankfurt, Strikman 10; Diehl, Ostermeier, Schafer 11

$$\langle N | O_{\text{tw}2}(x_1, \mathbf{r}_{1T}) O_{\text{tw}2}(x_2, \mathbf{r}_{2T}) | N \rangle_{\mathbf{r}_{1T} - \mathbf{r}_{2T} = \mathbf{r}_T}$$

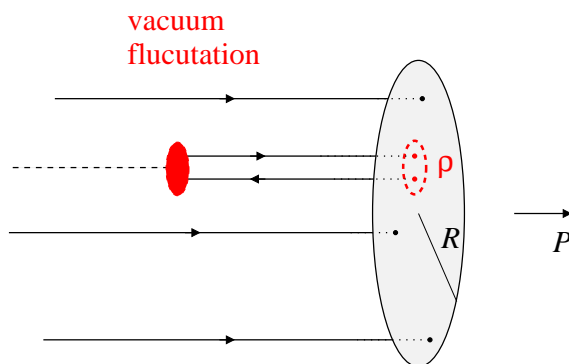
Subtleties: UV divergences, renormalization, mixing

- Perturbative and non-perturbative correlations

DGLAP evolution: Active parton from perturbative splitting, partner within range $r_T \sim \mu^{-1}$

Chiral symmetry breaking: Nonperturbative $q\bar{q}$ pairs with transverse size $\ll 1$ fm

Schweitzer, Strikman, CW 12. Cf. Shuryak 82; Diakonov, Petrov 84



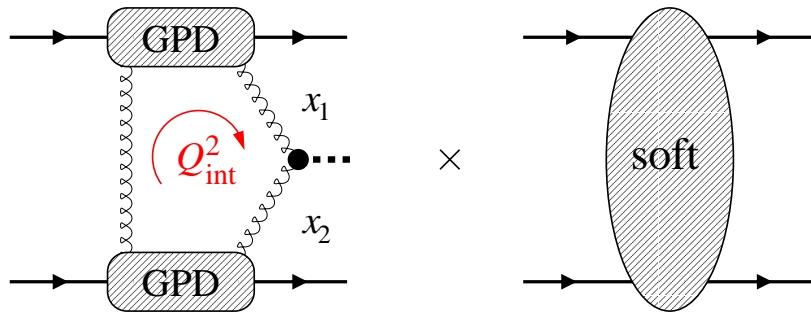
- Effect on MPI

Perturbative correlations can explain observed enhancement beyond mean field

Review Blok, Strikman 17

Exclusive diffraction: $pp \rightarrow p + H + p$

12



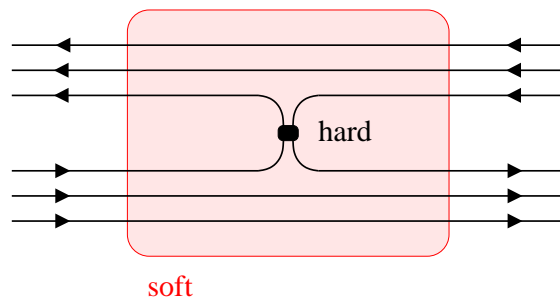
- H produced in hard process

$$\mu_{\text{soft}}^2 \ll Q_{\text{int}}^2 \ll M^2 \quad \text{Khoze et al. 97+}$$

$$x_{1,2} \sim \frac{M}{\sqrt{s}} \sim 10^{-2} \quad \text{Higgs at LHC}$$

- Soft spectator interactions must not produce particles

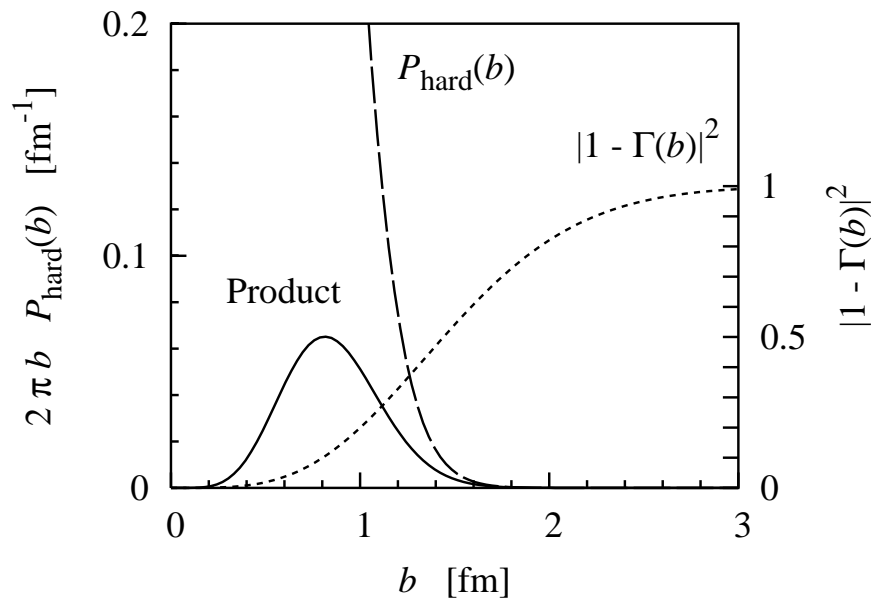
$$S^2 \equiv \frac{\sigma_{\text{diff}}(\text{full})}{\sigma_{\text{diff}}(\text{no soft})} \quad \text{Gap survival probability}$$



- Mean-field approximation:
 $[V_{\text{hard}}, H_{\text{soft}}] = 0$ independent,
 closure of partonic states

- Amplitude calculable in terms of
 - Gluon GPD, unintegrated
 - pp elastic S -matrix

Different time/distance scales



- Gap survival probability

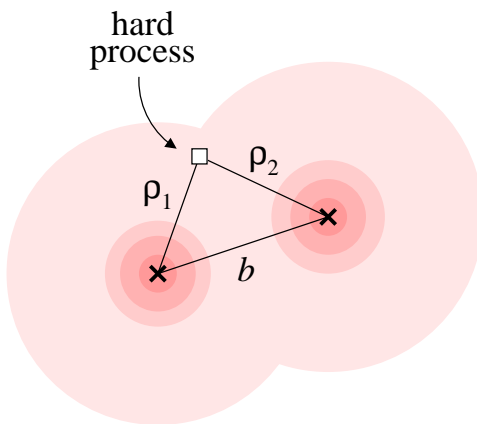
$$S^2 = \int d^2b P_{\text{hard}}(b) |1 - \Gamma(b)|^2$$

Probability for
two-gluon collision

Probability for
“no inelast. interaction”

favors small b

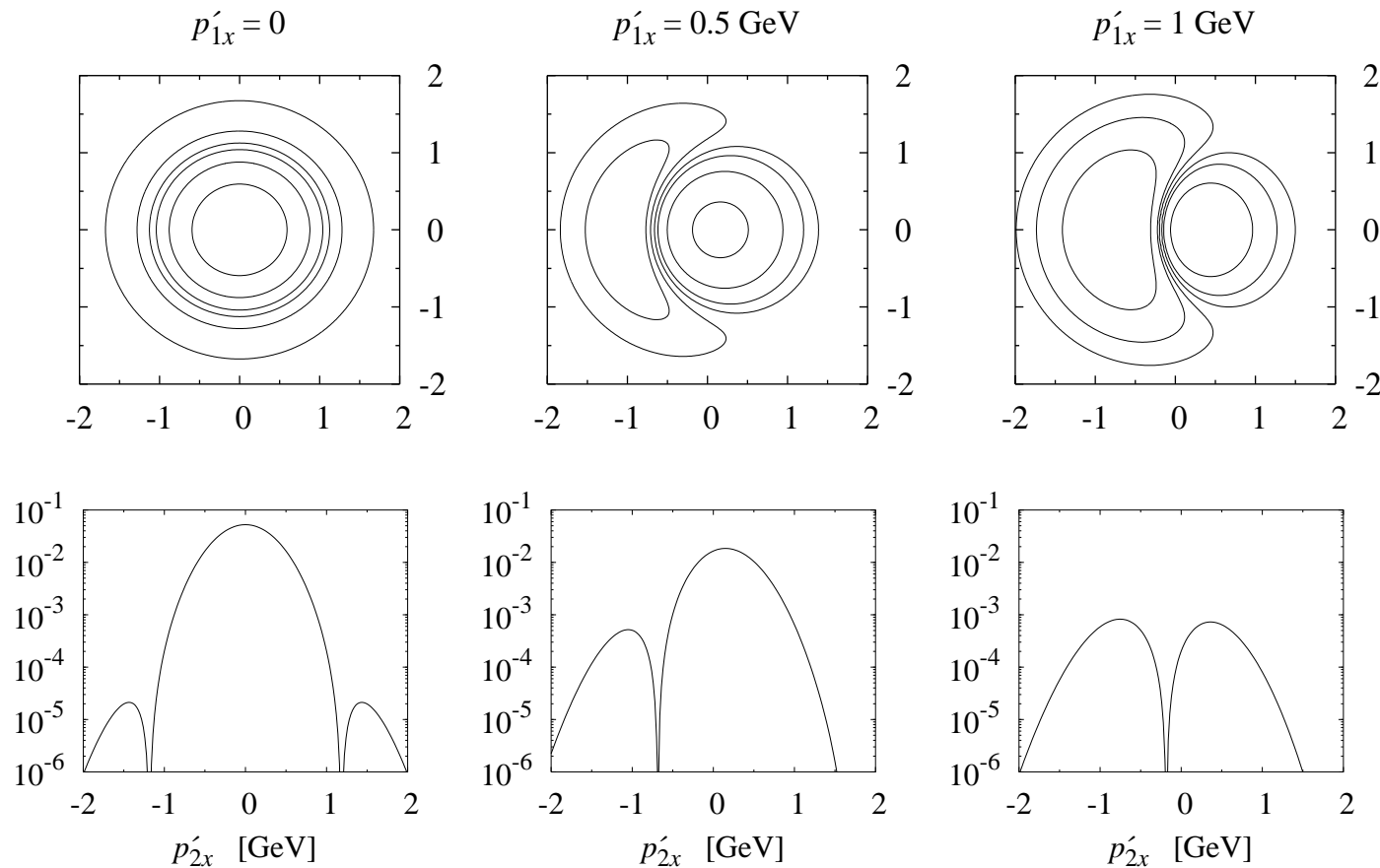
favors large b



$P_{\text{hard}}(b)$: Overlap of normalized
transverse gluon densities (squared)

- “Blackness” of pp amplitude $\Gamma(b) \sim 1$
suppresses diffraction at small b
- Numerical results in mean-field approx.
 $S^2 \sim 0.03 - 0.04$ Higgs at LHC
- Parton correlations further reduce
RGS probability

Exclusive diffraction: p_T dependence



- Gap survival probability depends on final proton transverse momenta \mathbf{p}_{1T} and \mathbf{p}_{2T}

Figure: Dependence on \mathbf{p}_{2T} for fixed \mathbf{p}_{1T} in x -direction

Observable diffraction pattern attests to interplay of hard and soft interactions

Calculable from GPDs and pp elastic amplitude

Frankfurt, Hyde, Strikman, CW, PRD 75, 054009 (2007) [INSPIRE]

- Centrality trigger for pp collisions

Hard process at $x_{1,2} \sim 10^{-1}-10^{-2}$ selects central pp collisions

Small- x gluon density in central pp collision comparable to heavy nucleus:
Black-disk regime in leading parton interactions, observables in forward particle production

- Quantum fluctuations of gluon density

Defined/quantified in context of collinear factorization

Measured in diffractive VM production $ep \rightarrow e + V + X(\text{low-mass})$

HERA, EIC

Influences MPI, rapidity gap survival in pp

- Parton correlations in ep scattering

Nonperturbative correlations between sea quarks due to chiral symmetry breaking

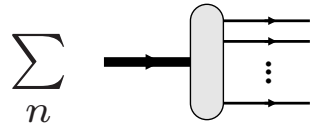
Observable in hadron correlations between current and target fragmentation regions
of DIS at intermediate energies ($W^2 \sim 30 \text{ GeV}^2$, $Q^2 \sim 2-3 \text{ GeV}^2$)

EIC

- pp collisions characterized by interplay between hard and soft interactions
 - Correlations between hard process and event observables
 - Proton structure beyond 1-body parton densities
- Transverse geometry essential aspect of pp collisions
 - Transverse distribution of partons from $ep/\gamma p$ (GPDs)
 - Hard processes select central pp collisions
 - Geometric correlations explain underlying-event characteristics
 - Baseline estimate of MPI rates from geometry and mean-field approximation
 - Baseline estimate of gap survival probability in central exclusive diffraction
- Synergies with EIC physics program
 - Conceptual connections
 - ep input for next-level analysis of pp collisions: Fluctuations, correlations
 - Highest available energies in pp at LHC

Supplementary material

- Nucleon quantum many-body system



Partonic wave function has components with different particle number, transverse size, etc.

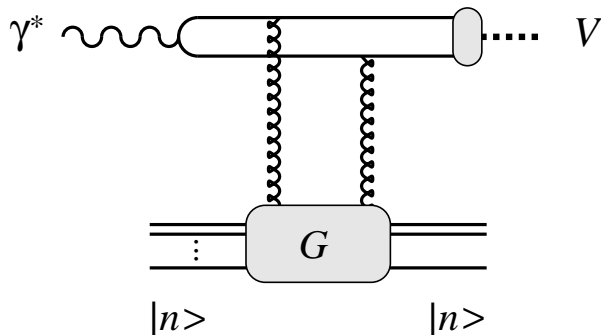
High-energy process intercepts instantaneous configurations, interactions "frozen"

Inclusive DIS measures average parton density

Fluctuations of parton density and transverse size?

Fundamental property of many-body system

Frankfurt, Strikman, Treleani, CW, PRL **101**:202003, 2008



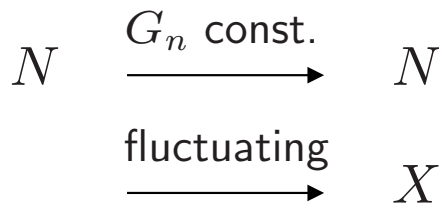
- Fluctuations of gluon density

Hard diffractive processes at small x

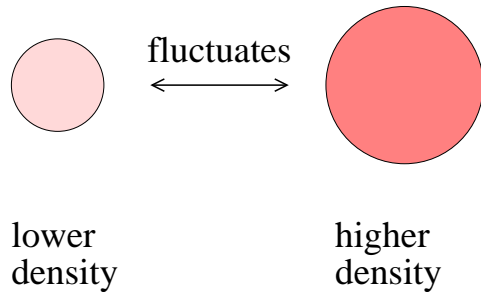
Amplitude diagonal in partonic states $|n\rangle$, proportional to configurations's gluon density G_n

Fluctuations of G_n lead to dissociation

Cf. soft diffraction: Good, Walker 60, Miettinen, Pumplin 78



$$\omega_g \equiv \frac{\langle G^2 \rangle - \langle G \rangle^2}{\langle G \rangle^2} = \frac{d\sigma/dt (\gamma^* N \rightarrow V X)}{d\sigma/dt (\gamma^* N \rightarrow V N)} \Big|_{t=0}$$



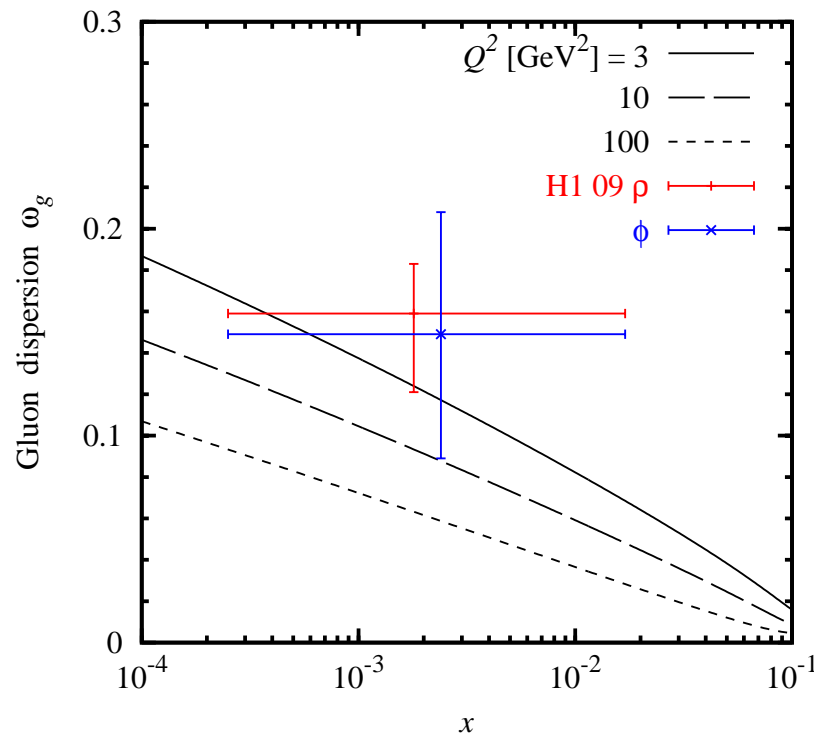
- Scaling model Close et al. 83: EMC effect

Fluctuations of size change effective scale of non-pert gluon density $\mu^2(\text{gluon}) \propto R^{-2}$

Size distribution from soft cross section fluctuations $\omega_\sigma \sim 0.25$ at $\sqrt{s} = 20 \text{ GeV}$

Gluon density fluctuations change with x, Q^2 through DGLAP evolution

Roughly consistent with HERA data

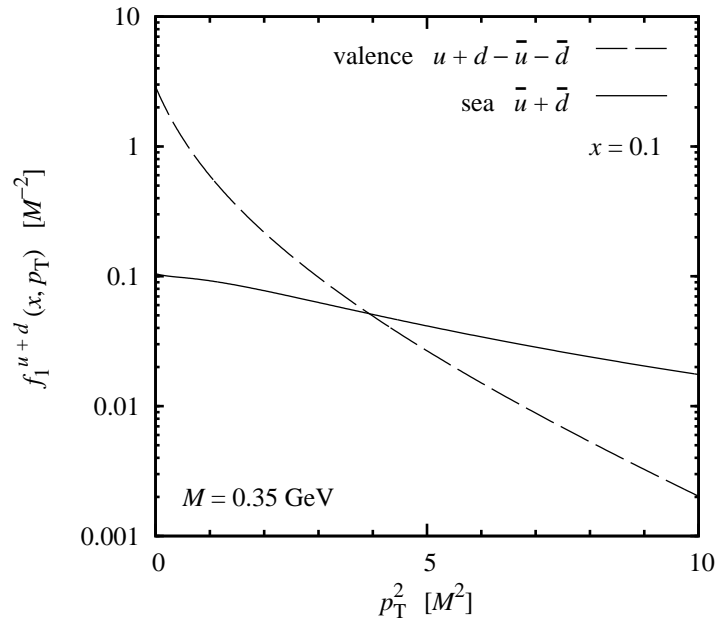


- Fluctuation effect on MPI

Small effect of gluon density fluctuations
 $\omega_g < 0.1$ at Tevatron

Moderate enhancement from size fluctuations
 $\sigma_{\text{eff}}(\text{fluct}) \approx (1 - \omega_\sigma/2) \sigma_{\text{eff}}(\text{mean field})$
 $\sim 10\text{-}15\%$ at Tevatron

Fluctuation effect on MPI small,
cannot explain experimental rates



- Model of nonperturbative correlations
Schweitzer, Strikman, CW 12

Chiral quark-soliton model: Dynamical quark mass, semiclassical approximation in large- N_c limit
Diakonov, Petrov, Polyitsa 88

Sea quark transverse momenta up to $p_T \sim \mu_{\chi SB}$
Different from valence quarks $p_T \sim R^{-1}$

Correlated $q\bar{q}$ pairs in nucleon wave function:
Spin/ flavor structure, σ/π quantum numbers

- Signals in deep-inelastic lepton scattering?

P_T distributions in semi-inclusive DIS
incl. spin asymmetries, particle correlations. JLab12, COMPASS

Particle correlations between current and target fragmentation regions
 $W \sim \text{few GeV}$ to avoid DGLAP radiation. COMPASS, EIC

Exclusive meson production at large x
Knockout of correlated $q\bar{q}$ pair. JLab12

