

(E12-10-007 Update to Jefferson Lab PAC-50)

Precision Measurement of Parity-violation in Deep Inelastic Scattering Over a Broad Kinematic Range

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Executive Summary

3 We provide an update to the SoLID PVDIS experiment, E12-10-007, approved by PAC37 in 2010.
4 The main goal of SoLID PVDIS is to test the Standard Model (SM) and search for beyond-the-SM (BSM)
5 physics. Even though the SM is consistent with a large number of experimental tests, evidence that the SM
6 is incomplete and that it belongs to a larger theoretical framework remains, such as the long-lasting dark
7 matter and hierarchy problems. Testing of the SM and search for BSM physics will remain a long-term
8 effort pursued by communities from low energy atomic physics to the high energy program at the LHC,
9 and SoLID PVDIS belongs to this EW/BSM program yet remains unique given its unambiguous access
10 to a specific combination of the electron-quark effective couplings. On the hadronic physics side, PVDIS
11 on the proton provides access to the PDF ratio d/u at high x , free from nuclear effects that cloud even
12 the latest experiments that utilized innovative experimental techniques. In short, the physics motivation
13 of SoLID PVDIS remains strong.

14 Extensive progress has been made on the development of the SoLID spectrometer, from procure-
15 ment and refurbishment of the magnet, to detector pre-R&D activities, to progress in data acquisition and
16 simulation. Discussions with theory groups are also in progress on the interpretation of the data. The mea-
17 surement described in our original proposal remains feasible and the budget for systematic uncertainties
18 valid. In this document, we will review the physics motivation of PVDIS and updates to the landscape of
19 EW/BSM physics and PDF ratio d/u study, and provide progress on the apparatus. In 2010, we requested
20 338 PAC days of beam time and were approved for 169 days. We request 338 PAC days of beam time,
21 same as the original proposal.

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22 1 Introduction

23 The Standard Model (SM) is a theoretical framework that successfully explains nearly all existing phe-
 24 nomenon of particle physics. However, there is long-standing evidence of new physics, such as dark
 25 matter and neutrino mass, that suggest many possibilities to extend the SM to higher energy scales. Thus
 26 it is useful to carry out as many high-precision measurements as possible to test the SM and to shed light
 27 on where beyond-the-standard model (BSM) physics might occur.

28 The Solenoid Large Intensity Device (SoLID), to be used in the measurement proposed here, is a gen-
 29 eral purpose spectrometer that will be used to realize a large physics program. The large acceptance, high
 30 intensity nature of SoLID will provide a unique opportunity to measure the parity violation asymmetry in
 31 electron deep inelastic scattering (PVDIS). The first PVDIS experiment, SLAC E122 [1, 2], provided a
 32 pivotal role in establish the SM. During the 6 GeV era of JLab, the PVDIS asymmetry was measured and a
 33 value of the effective vector electron axial-vector quark neutral current coupling combination $2g_{VA}^{eu} - g_{VA}^{ed}$
 34 was extracted, showing for the first time this quantity is non-zero [3].

35 PVDIS is one of a number of PVES experiments performed in progress, see left panel of Fig. 1.
 36 Remarkable progress over time has been made in both how small an asymmetry can be measured and also
 37 how precise (relative uncertainty) the measurement can be made. It is in the context of this record that we
 38 are willing to tackle an experiment as challenging as the present one, where we are pushing the envelope
 39 in how precisely an asymmetry ($\delta A_{PV}/A_{PV} \sim 0.5\%$) can be measured, see right panel of Fig. 1.

40 The specific physics case for the SoLID PVDIS experiment is presented in PR-12-9-012 and PR-12-
 41 007, and was updated in Section 2.3 of the SoLID Pre-CDR [4]. The experimental details are also in the
 42 Pre-CDR. In this document, we will focus on new developments that have impact on the motivation for
 43 the experiment as well as new information on the design, much of which is contained in the recently-
 44 completed PRE-R&D reports.

45 2 Phenomenology of PVDIS

46 The fundamental parity-violating weak neutral current Lagrangian can be approximated as four fermion
 47 contact interactions at low energy scales, $Q^2 \ll M_Z^2$. Counting only parity-violating terms:

$$\mathcal{L}_{\text{NC}}^{eq} = \frac{1}{2v^2} \left(\bar{e}\gamma^\mu\gamma^5 e \sum_{q=u,d} g_{AV}^{eq} \bar{q}\gamma_\mu q + \bar{e}\gamma^\mu e \sum_{q=u,d} g_{VA}^{eq} \bar{q}\gamma_\mu\gamma^5 q \right), \quad (1)$$

48 where $v = (\sqrt{2}G_F)^{-1/2} = 246.22$ GeV is the Higgs vacuum expectation value setting the electroweak
 49 scale with G_F the Fermi constant. For the SM at the tree level, the electron-quark effective couplings are

$$g_{AV}^{eu} = C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W; \quad g_{AV}^{ed} = C_{1d} = \frac{1}{2} - \frac{2}{3} \sin^2 \theta_W$$

50

$$g_{VA}^{eu} = C_{2u} = -\frac{1}{2} + 2 \sin^2 \theta_W; \quad g_{VA}^{ed} = C_{2d} = \frac{1}{2} - 2 \sin^2 \theta_W.$$

51 In the previous proposal, the notations $C_{1q,2q}$ were used. The new notation g 's [5] differ from the more
 52 familiar couplings C_{iq} in the detailed radiative corrections.

53 In this notation, the asymmetry for PVDIS for the deuteron can be written approximately as

$$A_{RL,d}^{\text{DIS}} \approx \frac{3}{20\pi\alpha} \frac{Q^2}{v^2} \left[(2g_{AV}^{eu} - g_{AV}^{ed}) + (2g_{VA}^{eu} - g_{VA}^{ed}) \frac{1 - (1-y)^2}{1 + (1-y)^2} \right], \quad (2)$$

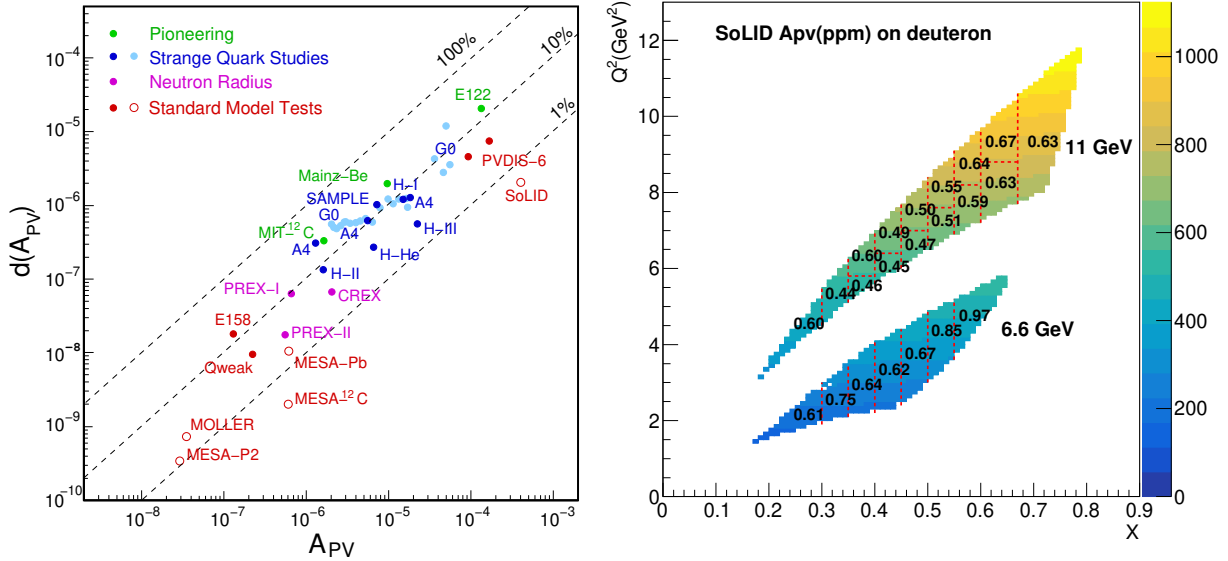


Figure 1: Left panel: History of PVES experiments in the size and precision of the asymmetry measured. Right panel: Anticipated kinematic coverage for SoLID in the PVDIS configuration. The top band is for a 11 GeV beam and the bottom band is for a 6.6 GeV beam. The number in each bin (with bin boundary shown by red dashed lines) represent the relative statistical precision in percent, while the color scale shows the size of A_{PV} for the deuteron in ppm.

54 where α is the fine structure constant and the PDF and many other possible corrections cancel to first order
 55 due to the isoscalar nature of the target. The SoLID spectrometer is designed to have large acceptance at
 56 large y , so that A_{PV} is very sensitive to the g_{VA}^{eq} couplings.

57 The main motivation for the experiment is that PVDIS is the only known experiment that can isolate
 58 one combination of the coupling constants, $2g_{VA}^{eu} - g_{VA}^{ed}$ with minimal theoretical errors. For processes at
 59 lower values of Q^2 , large, uncertain radiative corrections involving hadrons limit the possible precision.
 60 The fundamental motivation for PVDIS is that it is important to measure as many of the couplings as
 61 possible, each to as high a precision as possible.

62 From Eq. 2, we see that PVDIS is also sensitive to the g_{AV}^{eq} , but presently these are well known
 63 from Atomic parity-violation and the recently-published Qweak experiment [6]. In the future, the P2
 64 experiment [7] at MESA will further improve on the precision of these couplings. With the new data
 65 published by Qweak and anticipated from P2, the uncertainty on the g_{AV}^{eq} will essentially be completely
 66 negligible in extracting $2g_{VA}^{eu} - g_{VA}^{ed}$, an important development since the proposal was submitted. The
 67 projected uncertainties on the couplings from SoLID and P2 are shown in the left panel of Fig. 2. We note
 68 from Fig 1 that the the ^{12}C part of P2 proposes to measure an asymmetry at the 0.5% level.

69 3 Physics beyond the Standard Model

70 One possibility for BSM physics is new interactions that occur at the multi-TeV scale. At low energies,
 71 these interactions are manifest as contact interactions of the form of Eq. 1. In total, there are 86 possible

72 operators for SM particles. The theoretical framework to model these heavy-particle new physics is called
 73 Standard Model Effective Field Theory (SMEFT), and global analyses can be done to identify regions of
 74 parameter space where new physics is not excluded. In this context, the goal of SoLID is to reduce the
 75 volume of this multi-dimensional parameter space by an order of magnitude by measuring one specific
 76 combination of the coupling constants.

77 One experiment at the LHC, the production of Drell-Yan e^+e^- pairs, has relevance for SoLID. Al-
 78 though the observables in the experiment are parity conserving, there are large interference terms with
 79 the parity-violating SM that give sensitivity to parity-violating couplings. The experiment thus measures
 80 linear combinations of 14 couplings. The SoLID measurement will provide an additional constraint. An
 81 SMEFT analysis of the both LHC Drell-Yan and PVES data has been published [8]. The analysis was
 82 limited to scenarios where many of the couplings were arbitrarily set to zero. We are awaiting a full
 83 model-independent analysis that can show the full impact of SoLID.

84 Another possibility for BSM physics is the existence of “dark light” in which there is a light boson
 85 that also couples to dark matter [9]. The resulting modification to PVES experiments is that $\sin^2 \theta_W$ has
 86 an additional Q^2 -dependence beyond that predicted by the SM. Here, PVDIS is unique in that its $Q^2 \sim$
 87 $7(\text{GeV}/c)^2$. In this scenario, the analysis process is different from the determination of the couplings.
 88 Instead of extracting $2g_{VA}^{eu} - g_{VA}^{ed}$, the SM values of the couplings are used and the asymmetry becomes
 89 a function of $\sin^2 \theta_W$. The value of $\sin^2 \theta_W$ can thus be extracted from the measured asymmetry and
 90 compared with other high-precision measurements at different Q^2 values.

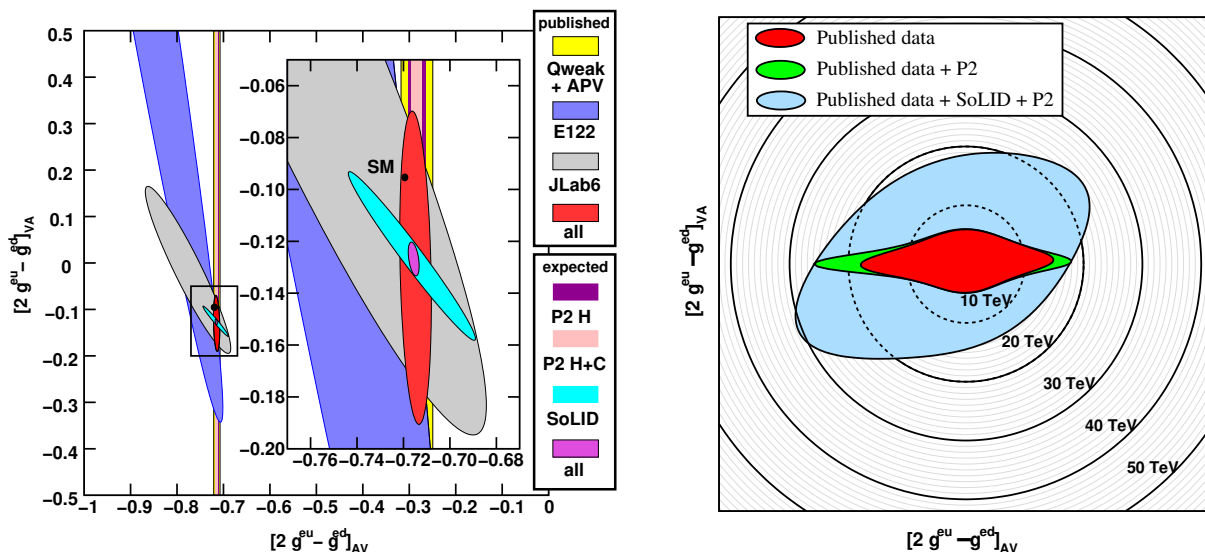


Figure 2: Left panel: Expected limits on the coupling constants from future experiments. Right panel: Projection limits on the scale of BSM physics for SoLID. Definitions are given in Ref.[3].

91 4 Measurement of the PDF d/u Independent of Nuclear Structure

92 A measurement of A_{RL}^{DIS} on a proton target is sensitive to the ratio of the d to u quark PDF. The standard
 93 determination of the d/u ratio relies on fully inclusive DIS on a proton target compared to a deuteron
 94 target. In the large x region, nuclear corrections in the deuteron target lead to large uncertainties in the
 95 d/u ratio. However, they can be completely eliminated if the d/u ratio is obtained from the proton target

96 alone, such as from A_{RL}^{DIS} for a proton target, given approximately by

$$A_{RL,p}^{\text{DIS}} \approx \frac{1}{4\pi\alpha} \frac{Q^2}{v^2} \left[\frac{12 g_{AV}^{eu} - 6 g_{AV}^{ed} d/u}{4 + d/u} \right].$$

97 Precision measurements of A_{RL}^{DIS} on a proton target can thus be a powerful and unambiguous probe of the
98 d/u ratio.

99 Since the proposal was submitted, there have been two important developments in the field on this
100 topic. First, W production data from Fermilab have greatly improved the measurement of d/u at high
101 Q^2 [10]. At low Q^2 , the most important development is the publication of recent data from the MARATHON
102 collaboration. The MARATHON data [11] have been interpreted in two different ways [11, 12], high-
103 lighting the importance of the proposed measurement of PVDIS asymmetry with a hydrogen target. It is
104 important that the nuclear physics be understood and that the determinations of d/u be understood both
105 from W production and light nuclei.

106 5 PVDIS at the Future EIC

107 The Electron Ion Collider (EIC) will be the first lepton-ion collider that polarizes both the electron and the
108 proton (ion) beams, and the first collider with the capacity to flip its electron beam helicity at the source
109 level. These unique design features will allow a direct measurement of A_{PV} at a collider setting for the
110 first time.

111 The impact of EIC on the determination of the weak mixing angle and NC couplings is marginal [13]
112 because of two reasons: First, while EIC is expected to reach a $10^{33} \text{ cm}^{-2} \text{ s}^{-1}$ luminosity, it is still
113 four orders of magnitude below that of SoLID. Second, the electron beam polarimetry at the EIC is ex-
114 pected to reach 1% uncertainty, limiting the precision of the measurement. We note the larger uncertainty
115 on $\sin^2 \theta_W$ expected for the annual integrated luminosity of EIC than our proposed measurement using
116 SoLID. In addition, as outlined in Section 2, only eD scattering will allow a direct access to the combi-
117 nation $2g_{VA}^{eu} - g_{VA}^{ed}$ with minimal ambiguity from PDFs, and it is unlikely that all beam time of EIC will
118 be spent on eD collision. This further limits the impact of EIC on EW NC study in terms of a precision
119 determination of the specific EW couplings $2g_{VA}^{eu} - g_{VA}^{ed}$. Finally, the EIC access, and thus tests the SM,
120 in a different energy range from SoLID.

121 6 Progress on the SoLID Apparatus

122 Since the approval in 2010 of five SoLID experiments with high rating by the JLab PAC, the collaboration
123 has developed a Pre-CDR [4] which gives the details for the SoLID apparatus as well as a cost estimate.
124 In 2015, SoLID together with the approved experiments received a strong endorsement from the Nuclear
125 Physics Long Range Plan. In September of 2019, the Pre-CDR successfully passed the second of two
126 JLab Director's Reviews. The latter review covered a detailed cost estimate and detailed analysis that
127 concluded that all critical items were low risk. The Pre-CDR was the basis of the SoLID MIE submitted
128 to the DOE in February of 2020. In 2020, the DOE funded a Pre-R&D plan, which has demonstrated that
129 there are no show-stoppers in the design of SoLID. In March of 2021, the DOE performed the Science Re-
130 view of SoLID. We give a brief overview of the features of SoLID apparatus and pre-R&D progress below.

131

132 In the Pre-CDR, we have demonstrate that the apparatus can achieve the following goals:

- 133 1. High luminosity ($10^{37} \text{ cm}^{-2}\text{s}^{-1}$ for SIDIS and J/ψ , $10^{39} \text{ cm}^{-2}\text{s}^{-1}$ for PVDIS with baffles)
- 134 2. Large acceptance: 2π in azimuthal angle ϕ ; In polar angle θ : $8^\circ - 24^\circ$ for SIDIS and J/ψ , $22^\circ - 35^\circ$
- 135 for PVDIS: Momentum range: $1 - 7 \text{ GeV}/c$
- 136 3. High rates (trigger rate limit 100 KHz for SIDIS and J/ψ , 600 KHz of the 30 sectors for PVDIS)
- 137 4. High background tolerance ($\sim 1 \text{ GHz}$, dominated by low energy photons/electrons)
- 138 5. High radiation environment tolerance (10^{2-3} krad)
- 139 6. Moderate resolutions ($1 - 2\%$ in momentum, $1 - 2 \text{ mrad}$ in θ , 6 mrad in ϕ)
- 140 7. Good electron PID; moderate pion PID (SIDIS), and demanding Kaon PID (SIDIS-Kaon, enhanced
- 141 configuration)
- 142 8. High precision with low systematic effects

143 Since the SoLID experiments were approved, the CLEO magnet, located at Cornell University, was
 144 chosen. JLAB requested the magnet in 2012 and it was moved to JLAB in 2016, along with most of the
 145 required return iron. JLAB is currently performing minor refurbishment of the magnet and preparing for
 146 a cold test to establish the magnet's operational condition. The cold test is scheduled to be completed
 147 before the end of 2022.

148 The SoLID apparatus in its PVDIS configuration is shown in Fig. 3. Significant progress has been
 149 made in all subsystems, and many components have been tested and shown to meet specifications. The
 150 UVa group has built large GEM chambers with sizes what SoLID needs. They also successfully built and
 151 operated GEM detectors for the PRad experiment and SBS experiment. During our recent execution of
 152 the Pre-R&D plan, we have tested new VMM3 GEM readout chips, which we plan to use instead of the
 153 obsolete APV25 chips listed in the original proposals. We have demonstrated that with the VMM3 chips
 154 operated in a mode with a 25 ns shaping time, the readout can handle the high rates required for SoLID.

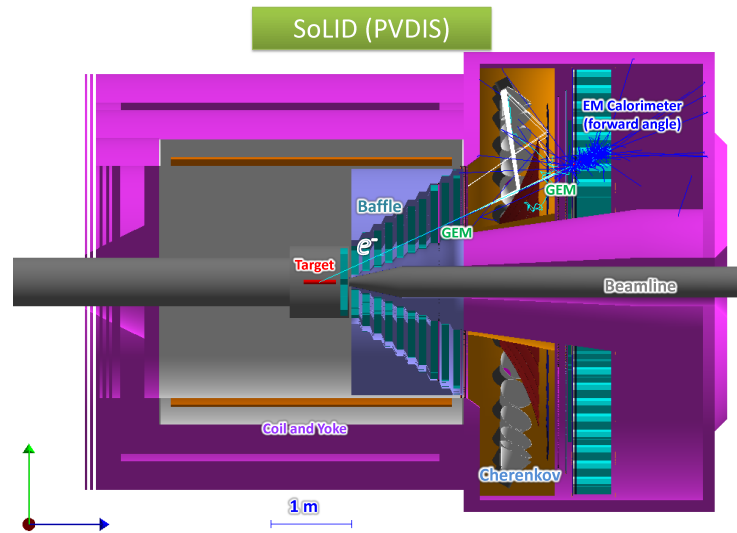


Figure 3: SoLID apparatus in its PVDIS configuration [4].

155 A prototype Cherenkov was tested with beam at JLab Hall C in 2020, both at low rates and also at high
 156 rates equal to those expected in the SoLID spectrometer. Photoelectrons were detected with MaMPT's as
 157 described in the proposal. In both the low and high rate configurations, clean signals could be identified

158 and the overall electronics performed very well. In addition, an Incom LAPPD (or large-area micro-
159 channel plate PMT), which can operate at higher magnetic fields, was tested. So far, this new technology
160 appears to be a promising alternative of the Cherenkov readout.

161 The ECal prototype modules were successfully tested at the Fermilab Test Beam Facility in January
162 of 2021. Data showed that the energy resolution of the 3-module setup reached the SoLID requirement of
163 $\delta E/E = \sqrt{10\%}/\sqrt{E}$ and the position resolution exceeded the required 1 cm specification.

164 The SoLID DAQ system is based on JLab 12 GeV FADC base pipelined electronics, which now have
165 been successfully used in Hall B and D. The system can meet the PVDIS deadtime requirement of 0.1%.
166 Special features of the JLab FADC required for SoLID, which were built into the JLab design but never
167 used in previous applications, have now been successfully tested for all systems as part of the SoLID
168 Pre-R&D plan.

169 Currently, a beam test of a full set of SoLID detector prototypes – GEM, LGC, ECal, DAQ and
170 associated electronics– is in preparation. The goal of the test is to fully characterize the functionality of
171 the detector system under a high-rate, high radiation environment that are similar to SoLID operation.

172 The SoLID simulation software, based on GEMC/GEANT4, is a single package that is used to model
173 all approved experiments and run-group experiments. It is being used to optimize the design and evaluate
174 the impact of the results from the past and ongoing tests. A Kalman Filter based track finding and fitting
175 algorithm was developed and tested with fully digitized GEM simulation data. Tracking resolutions with
176 good tracking efficiency were obtained with background taken into account. We are actively working on
177 assembling simulation, reconstruction, and analysis into one software framework.

178 Lastly, we anticipate that beam polarimetry will be one of the largest systematic errors. The recently
179 completed CREx experiment achieved an uncertainty in the Compton polarimetry of 0.52% for a 2.18 GeV
180 beam [14]. This is already close to the uncertainty budgeted in our original proposal, and is a promising
181 progress towards the polarimetry precision required by the SoLID PVDIS measurement.

182 **7 Update on Systematic Uncertainties**

183 **7.1 Contributions from PDFs**

184 Our current estimate of the systematic uncertainties is given in Table 1. Compared with the original pro-
185 posal, we have added uncertainty related to the use of PDFs in our analysis process following the recom-
186 mendation from the March 2021 Science Review. To estimate the effect of PDF uncertainty, we calculated
187 the value of A_{PV} using CT18NLO [15] (LHAPDF [16] ID 14400–14458) and MMHT2014nlo_68cl [17]
188 (ID 25100–25150) PDF sets. Considering data points at $x = 0.4$ and above (which are most relevant
189 for BSM and CSV search), the average relative uncertainty on the deuteron PVDIS asymmetry from
190 CT18NLO is 0.18% (at 90% C.L.) and the values calculated using CT18NLO and MMHT14NLO are
191 consistent within uncertainties. The uncertainty from MMHT14NLO is smaller. We thus use a conserva-
192 tive estimation of 0.2% as an overall uncertainty. The total systematic uncertainty increased from 0.53%
193 in our original proposal to 0.57% and the total uncertainty from 0.61% to 0.64%.

194 Since the March 2021 Science Review, we have conducted extensive studies of the PDF effect on the
195 extraction of electroweak parameters following the prescription described in Section IV of [13]. We found
196 the effect from PDF uncertainty to be consistent with Table 1.

Polarimetry	0.4
Q^2	0.2
Radiative Corrections	0.2
Event reconstruction	0.2
Statistics	0.3
PDF uncertainty	0.2
Total systematic	0.57
Total statistic and systematic	0.64

Table 1: Updated error budget for the SoLID deuterium PVDIS asymmetry measurement in percent (%).

7.2 Progress on Radiative Corrections

Since 2010, and in particular prompted by the March 2021 Science Review, extensive discussions with theory groups have been made on how a high-precision radiative correction for the measurement can be carried out. First, progress has been made on the uncertainty estimation of external radiative corrections using data that will be collected in the low W region simultaneously with the measurement. For internal radiative correction, we have adapted the event generator `Djangoh` [18], originally developed for HERA cross section analysis, to fixed-target experiments and to nuclear targets. We have made modification to `Djangoh` such that it can be used to calculate parity violation asymmetry to high precision, immune from the statistical limit of a Monte-Carlo program. While there are still detailed work to be done, we anticipate that the 0.2% uncertainty projected on the radiative corrections as given in Table 1 can be reached.

8 Summary

Although the Standard Model has been successful, many experiments at the LHC, JLab, Mainz, Fermilab, etc. are searching for new physics by making precision measurements. One important type of experiments, PVES, includes Qweak, MOLLER and SoLID PVDIS with a deuterium target at JLab and P2 at Mainz. These PVES experiments have made many advances over the years that give them impressive reach for new physics. PVDIS with a hydrogen target can also make a significant contribution to the issues surrounding the measurement the PDF ratio d/u at large x .

The design of the SoLID spectrometer is well underway. The solenoid magnet from the CLEO experiment has arrived at JLab and is being tested. An MIE describing the spectrometer has been submitted to the DOE. The physics case was reviewed by the DOE in March 2021. Extensive simulations have shown that the design is capable of completing the physics goals. A successful Pre-R&D plan has demonstrated that the proposed novel apparatus is on track to meeting the specifications.

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