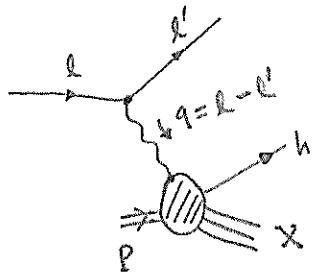
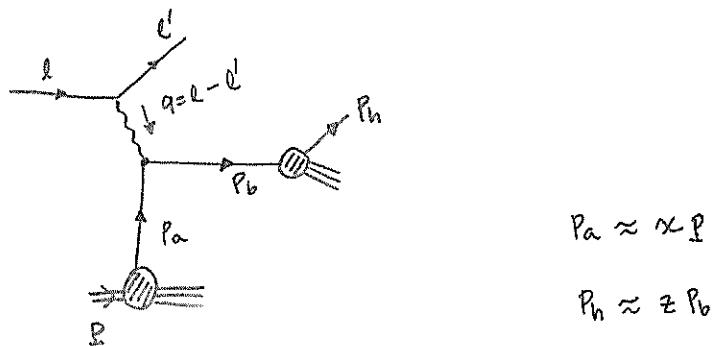


Semi-inclusive deep inelastic Scattering (SIDIS)

$$e(l) + P(p) \rightarrow e(l') + h(p_h) + X$$



partonic level



usual SIDIS kinematic variables

$$S_{\text{sep}} = (P + l)^2 = 2P \cdot l$$

$$Q^2 = -q^2 = -(l - l')^2$$

$$x_B = \frac{Q^2}{2P \cdot q} \quad z_h = \frac{P \cdot p_h}{P \cdot q} \quad y = \frac{P \cdot q}{P \cdot l} = \frac{Q^2}{x_B S_{\text{sep}}}$$

- Normalization of D2S-type process

$$d\sigma = \left| \begin{array}{c} l \\ l' \\ p \\ X \end{array} \right|^2$$

$$= \left| \begin{array}{c} l \\ l' \\ q \\ p \\ p_X \end{array} \right|^2 + \left| \begin{array}{c} l' \\ l \\ \mu \\ p \\ p \end{array} \right|^2$$

$$= \frac{1}{F} \overline{|M|^2} d\sigma$$

initial-flux $F = 4 \sqrt{(p_A \cdot p_B)^2 - m_A^2 m_B^2} = 4 p \cdot l = 2 S_{\text{sep}}$

$$\overline{|M|^2} d\sigma = \frac{1}{2} \text{Tr} [\gamma^\mu \gamma' \gamma^\nu \gamma'] e^2 * \frac{d^3 l'}{(2\pi)^3 2E'} (2\pi)^4 \delta^4(p+q-p_X) * \left(\frac{1}{Q^2}\right)^2$$

$$* e^2 \frac{1}{2} \sum_s \left(\frac{d^3 p_X}{(2\pi)^3 2E_X} \sum_{S_X} \right) \langle ps | J_\mu^+(0) | X \rangle \langle X | J_\nu(0) | ps \rangle$$

denote $\sum_X \equiv \left(\frac{d^3 p_X}{(2\pi)^3 2E_X} \right) \sum_{S_X}$

$$L^{\mu\nu} = \frac{1}{2} \text{Tr} [\gamma^\mu \gamma' \gamma^\nu \gamma']$$

$$= L^{\mu\nu} e^4 \left(\frac{1}{Q^2} \right) \frac{d^3 l'}{(2\pi)^3 2E'} \frac{1}{2} \sum_{S_X} \langle ps | J_\mu^+(0) | X \rangle \langle X | J_\nu(0) | ps \rangle * (2\pi)^4 \delta^4(p+q-p_X)$$

$$\Downarrow e^4 = (4\pi \alpha_{ew})^2$$

$$= \frac{\alpha_{ew}^2}{Q^2} \left(\frac{d^3 l'}{E'} \right) L^{\mu\nu} \left(\frac{1}{\pi} \right) * \frac{1}{2} \sum_{S_X} \langle ps | J_\mu^+(0) | X \rangle \langle X | J_\nu(0) | ps \rangle * (2\pi)^4 \delta^4(p+q-p_X)$$

define

$$W_{\mu\nu} = \frac{1}{4\pi} - \frac{1}{2} \sum_{S,V} \langle \bar{s} s | J_M^+(0) | x \rangle \times \langle \bar{v} v | J_V(0) | p_S \rangle \quad (2\pi)^4 \delta^4(p+q-p_S)$$

Then we have

$$d\sigma = \frac{1}{2S} \frac{\alpha_{em}^2}{Q^4} \frac{d^3 k'}{E} 4 W_{\mu\nu} L^{\mu\nu}$$

$$\boxed{E \frac{d\sigma}{d^3 k'} = \frac{2}{S} \frac{\alpha_{em}^2}{Q^4} L^{\mu\nu} W_{\mu\nu}}$$

further using $\delta^4(p+q-p_S) = \int \frac{d^4 y}{(2\pi)^4} e^{i(p+q-p_S) \cdot y}$

one might show

$$W_{\mu\nu} = \frac{1}{4\pi} \int d^4 y e^{i q \cdot y} - \frac{1}{2} \sum_S \langle \bar{s} s | J_M^+(y) J_V(0) | p_S \rangle$$

- lowest order of inclusive DIS

$$W^{\mu\nu} = \text{Feynman diagram} * \frac{1}{4\pi}$$

$$= \int dx f_{q/p}(x) * \text{Feynman diagram} * \frac{1}{4\pi}$$

$$= \int dx f_{q/p}(x) \frac{1}{4\pi} \operatorname{Tr} \left[\frac{\not{p}}{2} \not{v}(xp+q) \not{v}^* \right] e_q^2 + 2\pi \delta[(xp+q)^2]$$

$$= \int dx f_{q/p}(x) \left[p^\mu (xp+q)^\nu + p^\nu (xp+q)^\mu - g^{\mu\nu} (xp+q) \cdot p \right] e_q^2 \delta[(xp+q)^2]$$

Choose a frame $p^\mu = p^+ \bar{n}^\mu$

$$q^\mu = -x_B p^\mu + \frac{Q^2}{2x_B p^+} n^\mu \quad \bar{n}^\mu = [1, 0, 0]$$

$$n^\mu = [0, 1, 0]$$

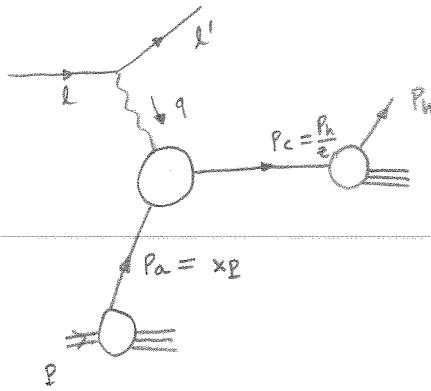
$$v^+ = \frac{1}{\sqrt{2}}(v^0 + v^z)$$

$$v^- = \frac{1}{\sqrt{2}}(v^0 - v^z)$$

$$W^\mu = \int dx f_{q/p}(x) e_q^2 \frac{1}{2} d^{\mu\nu} \delta(x - x_B)$$

$$= \frac{1}{2} d^{\mu\nu} \sum_i e_i^2 f_{q/p}(x_i)$$

$$\text{where } d^{\mu\nu} = -g^{\mu\nu} + \bar{n}^\mu n^\nu + \bar{n}^\nu n^\mu$$



$$S = (P + \ell)^2 \approx 2P \cdot \ell$$

$$Q^2 = -q^2$$

$$x_B = \frac{Q^2}{2P \cdot q} \quad z_h = \frac{P \cdot P_h}{P \cdot q} \quad y = \frac{P \cdot q}{P \cdot \ell} = \frac{Q^2}{x_B S}$$

$$\text{define } \hat{x} = \frac{x_B}{x} \quad \hat{z} = \frac{z_h}{z}$$

work in the so-called hadron frame

$$\bar{n}^\mu = [1^+, 0^-, 0_\perp] \quad n^\mu = [0^+, 1^-, 0_\perp]$$

$$p^\mu = p^+ \bar{n}^\mu$$

$$q^\mu = -x_B p^+ \bar{n}^\mu + \frac{Q^2}{2x_B p^+} n^\mu$$

$$\text{from } z_h = \frac{P \cdot P_h}{P \cdot q} = \frac{P^+ P_h^-}{Q^2/(2x_B)} \Rightarrow P_h^- = z_h \frac{Q^2}{2x_B p^+}$$

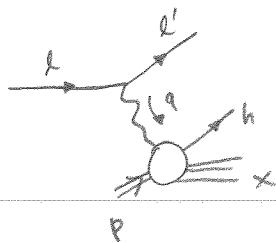
$$P_h^2 = 2P_h^+ P_h^- - \vec{P}_{h\perp}^2 \Rightarrow P_h^+ = \frac{\vec{P}_{h\perp}^2}{2P_h^-} = \frac{x_B \vec{P}_{h\perp}^2}{z_h Q^2} p^+$$

$$\text{thus } P_h^\mu = \frac{x_B \vec{P}_{h\perp}^2}{z_h Q^2} p^+ \bar{n}^\mu + \frac{z_h Q^2}{2x_B p^+} n^\mu + P_{hT}^\mu \quad (P_{hT}^\mu P_{hT\mu} = -\vec{P}_{h\perp}^2)$$

$$P_c^\mu = \frac{1}{z} P_h^\mu \quad (P_{c\perp} = \frac{P_{h\perp}}{z})$$

$$= \frac{x_B P_{c\perp}^2}{z Q^2} p^+ \bar{n}^\mu + \frac{\hat{z} Q^2}{2x_B p^+} n^\mu + P_{cT}^\mu$$

DIS normalization



from CTEQ handbook

$$E' \frac{d\sigma}{d^3 k'} = \left(\frac{2}{3}\right) \left(\frac{\alpha_{em}}{Q^2}\right)^2 L^{\mu\nu} W_{\mu\nu}$$

where $L^{\mu\nu} = \frac{1}{2} \text{Tr}[k \gamma^\mu k' \gamma^\nu]$

$$W_{\mu\nu} = \frac{1}{4\pi} \int dy e^{iq \cdot y} - \frac{1}{2} \sum_s \langle ps | J_\mu^+(y) J_\nu(0) | ps \rangle$$

Note $\frac{d^3 k'}{E'} = \frac{\pi Q^2}{x_B s} dx_B dQ^2$

↓ define $y = \frac{Q^2}{x_B s}$

$$= \pi s y dx_B dy$$

$$\frac{d\sigma}{dx_B dy} = \frac{2\pi \alpha_{em}^2 y}{(Q^2)^2} L^{\mu\nu} W_{\mu\nu}$$

↓ take $\frac{1}{4\pi}$ out from $W_{\mu\nu}$

$$= \frac{\alpha_{em} y}{2(Q^2)^2} L^{\mu\nu} W_{\mu\nu}$$

In a so-called hadron frame, one could write

$$\frac{2}{Q^2} L^{\mu\nu} = (1 + \cosh^2 \psi) (x^\mu x^\nu + y^\mu y^\nu) + 2 \sinh \psi T^\mu T^\nu$$

$$\cosh \psi = \frac{2}{y} - 1$$

$$X^{\mu}X^{\nu} + Y^{\mu}Y^{\nu} = -g^{\mu\nu} + T^{\mu}T^{\nu} - Z^{\mu}Z^{\nu}$$

$$\begin{aligned} T^{\mu} &= \frac{1}{\alpha} (q^{\mu} + 2x_B p^{\mu}) \\ Z^{\mu} &= -\frac{q^{\mu}}{\alpha} \end{aligned}$$

drop all q^{μ}, q^{ν} since $q^{\mu} W_{\mu\nu} = q^{\nu} W_{\mu\nu} = 0$

$$= -g^{\mu\nu} + \frac{4x_B^2}{\alpha^2} p^{\mu} p^{\nu}$$

thus

$$\begin{aligned} \frac{2}{\alpha^2} L^{\mu\nu} &\Rightarrow \frac{2}{y^2} \left[\underbrace{\left(-g^{\mu\nu} + \frac{4x_B^2}{\alpha^2} p^{\mu} p^{\nu} \right)}_{\text{called transverse projection}} \left(1 + (1-y)^2 \right) + 2 \underbrace{\frac{4x_B^2}{\alpha^2} p^{\mu} p^{\nu} (2(1-y))}_{\text{longitudinal projection}} \right] \\ &= \frac{2}{y^2} \left[(-g^{\mu\nu}) (1 + (1-y)^2) + \left(\frac{4x_B^2}{\alpha^2} p^{\mu} p^{\nu} \right) (1 + 4(1-y) + (1-y)^2) \right] \\ &\quad \begin{matrix} \uparrow & \curvearrowleft \\ \text{refer to "Metric" contribution} & \text{longitudinal contribution} \\ \text{See, e.g. hep-ph/0411212} & \end{matrix} \end{aligned}$$

If we're only interested in Metric contribution, then we'll have

$$\frac{2}{\alpha^2} L^{\mu\nu} \rightarrow \frac{2}{y^2} [1 + (1-y)^2] (-g^{\mu\nu})$$

Then

$$\begin{aligned} \frac{d\sigma}{dx_B dy} &= \frac{\alpha^2 y}{2(\alpha^2)^2} \cdot \frac{\alpha^2}{2} \cdot \frac{2}{y^2} [1 + (1-y)^2] (-g^{\mu\nu}) W_{\mu\nu} \\ &= \frac{\alpha^2}{\alpha^2} \cdot \frac{1 + (1-y)^2}{2y} (-g^{\mu\nu}) W_{\mu\nu} \end{aligned}$$

At the partonic level

$$\frac{d\sigma}{dx dy} = \frac{2em^2}{Q^2} \frac{1+(1-y)^2}{2y} \int \frac{dx}{x} dz f_{g/p}(x) D_{q \rightarrow h}(z) [-g^{\mu\nu} H_{\mu\nu}] dp_{S(n)}$$

for example, at leading order

$$\begin{aligned} -g^{\mu\nu} H_{\mu\nu} &= \text{Diagram showing a quark-gluon-gluon vertex with momenta } q, p, \bar{p}, \text{ and a gluon loop with momentum } Q. \\ &= \frac{1}{2} \nabla[(x_p) \gamma^\nu (x \bar{x} + \chi) \gamma^\mu] (-g_{\mu\nu}) \\ &= 4(1-\epsilon) x_p \cdot q \\ &= (1-\epsilon) 2 \frac{\chi}{x_p} Q^2 \end{aligned}$$

$$\begin{aligned} dP_S^{(n)} &= \frac{d^{n-1} p_c}{(2\pi)^{n-1} 2E_c} (2\pi)^n \delta^n(x_p + q - p_c) \\ &\Downarrow p_c = \frac{1}{2} p_h \\ &= \frac{1}{z^{n-2}} \frac{d^{n-1} p_h}{(2\pi)^{n-1} 2E_h} (2\pi)^n \delta^n(x_p + q - p_c) \\ &= \frac{1}{z^{n-2}} \frac{d^n p_h}{(2\pi)^n} 2\pi \delta(p_h^2) (2\pi)^n \delta^n(x_p + q - p_c) \\ &= \frac{1}{z^{n-2}} dP_h^+ dP_h^- d^{n-2} P_{hT} \underbrace{2\pi \delta(zp_h^+ + p_h^- - \vec{p}_{hT}^2)}_{\frac{1}{2p_h^-} \delta(P_h^+ - \frac{\vec{p}_{hT}^2}{2p_h^-})} \delta(x_p^+ + q^+) \delta(q^- - p_c^-) \delta^{n-2}(p_{c\perp}) \\ &\Downarrow \frac{dp_h^-}{p_h^-} = \frac{dz_h}{z_h} \quad \delta^{n-2}(P_{hT}) = \delta^{n-2}(P_{hT}/z) = z^{n-2} \delta^{n-2}(P_{hT}) \\ &= \frac{1}{z^{n-2}} \frac{d^{2n}}{2z_h} d^{n-2} P_{hT} z^{n-2} \delta^{n-2}(P_{hT}) \frac{1}{p_h^+} \delta(x - x_B) \frac{1}{q^-} \delta(1 - \hat{z}) * 2\pi \end{aligned}$$

$$d\psi^{(i)} = \frac{dz_h}{z z_h} - \frac{1}{p+q} + \frac{1}{x} \delta(1-\hat{x}) \delta(1-\hat{z}) + 2\pi$$

$$\begin{aligned} &\Downarrow 2p+q = 2p+q = \frac{Q^2}{x_B} \\ &= \frac{dz_h}{z_h} - \frac{x_B}{x} \frac{1}{Q^2} \delta(1-\hat{x}) \delta(1-\hat{z}) + 2\pi = dz_h * \frac{x_B}{z x Q^2} \delta(1-\hat{x}) \delta(1-\hat{z}) + 2\pi \\ &\Downarrow z_h = z \end{aligned}$$

$$(-g^{(\mu)} H_{\mu\nu} d\psi^{(i)}) = z(1-\epsilon) \cancel{\frac{x}{x_B Q^2}} * \frac{dz_h}{z} \cancel{\frac{x_B}{x} \frac{1}{Q^2}} \delta(1-\hat{x}) \delta(1-\hat{z}) * 2\pi$$

Thus

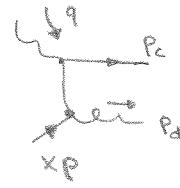
$$\begin{aligned} \frac{d\sigma}{dx dy dz_h} &= \frac{2\pi \cancel{J_{em}^2}}{Q^2} \frac{1+(1-y)^2}{2y} * z(1-\epsilon) \int \frac{dx}{x} \frac{dz}{z} f_{q/p}(x) D_{q \rightarrow h}(z) \\ &\quad * \delta(1-\hat{x}) \delta(1-\hat{z}) \end{aligned}$$

define $\sigma_0 = \frac{2\pi \cancel{J_{em}^2}}{Q^2} \frac{1+(1-y)^2}{y} (1-\epsilon)$, then

$$\frac{d\sigma}{dx dy dz_h} = \sigma_0 \int \frac{dx}{x} \frac{dz}{z} f_{q/p}(x) D_{q \rightarrow h}(z) * \delta(1-\hat{x}) \delta(1-\hat{z})$$

Study higher order

Now for real diagram, we have $dps^{(2)}$



$$dps^{(2)} = \frac{d^{n-1}p_c}{(2\pi)^{n-1} 2E_c} \frac{d^{n-1}p_d}{(2\pi)^{n-1} 2E_d} (2\pi)^n \delta^n(x_p + q - p_c - p_d)$$

$$= \frac{d^{n-1}p_h}{(2\pi)^{n-1} 2E_h} \frac{1}{z^{n-2}} \frac{d^n p_d}{(2\pi)^n} 2\pi \delta(p_d^2) * (2\pi)^n \delta^n(x_p + q - p_c - p_d)$$

$$= \frac{d^n p_h}{(2\pi)^n} 2\pi \delta(p_h^2) \frac{1}{z^{n-2}} 2\pi \delta(p_d^2)$$

$$= dP_h^+ dP_h^- d^{n-2} P_{hL} \underbrace{\delta(2P_h^+ P_h^- - \vec{P}_{hL}^2)}_{\frac{1}{2P_h^- \delta(P_h^+ - \frac{\vec{P}_{hL}^2}{2P_h^-})}} \frac{1}{z^{n-2} (2\pi)^{n-2}} \delta[(x_p + q - p_c)^2]$$

$$\downarrow \frac{dP_h^-}{P_h^-} = \frac{dz_h}{z_h}$$

$$= \frac{dz_h}{2z_h} d^{n-2} P_{hL} \frac{1}{(2\pi z)^{n-2}} \delta[(x_p + q - p_c)^2]$$

$$(x_p + q - p_c)^2 = (x_p + q)^2 - 2p_c \cdot (x_p + q)$$

$$= -Q^2 + x_2 p \cdot q - x_2 p_c \cdot p - 2p_c \cdot q$$

define $\hat{s} = (x_p + q)^2 = -Q^2 + x_2 p \cdot q = -Q^2 + x \frac{Q^2}{x_B} = \frac{Q^2(1-\hat{x})}{\hat{x}}$

$$\begin{aligned} \hat{t} &= (p_c - q)^2 = -Q^2 - 2p_c \cdot q = -Q^2 - [2p_c^+ q^- + 2p_c^- q^+] \\ &= -Q^2 - \left[2 \frac{x_B p_c^2}{\hat{x} Q^2} p^+ \frac{Q^2}{2x_B p^+} + 2 \frac{\hat{x} Q^2}{2x_B p^+} (-x_B p^+) \right] \\ &= -Q^2 - \left[\frac{p_c^2}{\hat{x}} - \frac{\hat{x} Q^2}{2} \right] \\ &= - \left[(1-\hat{x}) Q^2 + \frac{p_c^2}{\hat{x}} \right] \end{aligned}$$

$$\hat{u} = (x_p - p_c)^2 = x (-2p_0 p_c) = x (-2) p^+ \frac{\hat{z} Q^2}{2 x_B p^+} = -\frac{z^2}{x} Q^2$$

Thus from $0 = (x_p + q - p_c)^2$

$$\begin{aligned} \Rightarrow \delta[(x_p + q - p_c)^2] &= \delta[\hat{s} + \hat{t} + \hat{u} + Q^2] \\ &= \delta\left[\frac{Q^2(1-\hat{z})}{\hat{z}} - (1-\hat{z}) Q^2 - \frac{p_{cL}^2}{\hat{z}} - \frac{z^2}{x} Q^2 + Q^2\right] \\ &= \delta\left[\frac{p_{cL}^2}{\hat{z}} - \frac{Q^2}{\hat{z}} (1-\hat{z})(1-\hat{z})\right] \\ &= \hat{z} \delta\left[p_{cL}^2 - Q^2 \frac{\hat{z}(1-\hat{z})(1-\hat{z})}{\hat{z}}\right] \end{aligned}$$

$$\text{Thus } p_{cL}^2 = \frac{Q^2 \hat{z}(1-\hat{z})(1-\hat{z})}{\hat{z}}$$

$$\begin{aligned} \Rightarrow \text{Thus } \hat{t} &= -\left[(1-\hat{z}) Q^2 + \frac{p_{cL}^2}{\hat{z}}\right] = -\left[(1-\hat{z}) Q^2 + Q^2(1-\hat{z}) \frac{(1-\hat{z})}{\hat{z}}\right] \\ &= -\left[Q^2(1-\hat{z}) \frac{1}{\hat{z}}\right] \end{aligned}$$

$\hat{s} = \frac{1-\hat{z}}{\hat{z}} Q^2$
$\hat{t} = -\frac{1-\hat{z}}{\hat{z}} Q^2$
$\hat{u} = -\frac{z^2}{\hat{z}} Q^2$

$$\begin{aligned} \frac{\hat{t} \hat{u} \hat{s}}{(\hat{s} + Q^2)^2} &= \frac{1-\hat{z}}{\hat{z}} Q^2 \times \frac{\hat{z}}{\hat{z}} Q^2 \times \frac{1-\hat{z}}{\hat{z}} Q^2 \\ &= \frac{\hat{z}(1-\hat{z})(1-\hat{z})}{\hat{z}} Q^2 = p_{cL}^2 \end{aligned}$$

\Rightarrow

$p_{cL}^2 = \frac{\hat{t} \hat{u} \hat{s}}{(\hat{s} + Q^2)^2}$
--

$p_{hL}^2 = z^2 p_{cL}^2 = z^2 \frac{\hat{t} \hat{u} \hat{s}}{(\hat{s} + Q^2)^2}$

$$\delta[(x_p + q - p_c)^2] = \hat{z} \delta\left[\vec{p}_{cl}^2 - \frac{Q^2 \hat{z}(\hat{t}-\hat{z})(\hat{t}-\hat{z})}{\hat{z}}\right]$$

$$\Downarrow \quad p_{cl}^2 = \frac{p_{cl}^2}{\hat{z}^2}$$

$$= \hat{z} z^2 \delta\left[p_{cl}^2 - \frac{z^2 Q^2 \hat{z}(\hat{t}-\hat{z})(\hat{t}-\hat{z})}{\hat{z}}\right]$$

Thus

$$dp_s^{(2)} = \frac{dz_n}{2z_n} d^{n-2} p_{cl} \frac{1}{(2\pi z)^{n-2}} \hat{z} z^2 \delta\left[p_{cl}^2 - \frac{z^2 Q^2 \hat{z}(\hat{t}-\hat{z})(\hat{t}-\hat{z})}{\hat{z}}\right]$$

$$\begin{aligned} \text{Note } \int d^d p_{cl} &= \int p_{cl}^{d-1} d p_{cl} * \sqrt{\omega_d} \\ &= \frac{1}{2} (p_{cl}^2)^{\frac{d-2}{2}} d p_{cl}^2 * \frac{2\pi^{d/2}}{\Gamma(d/2)} \\ &= \frac{\pi^{d/2}}{\Gamma(d/2)} (p_{cl}^2)^{\frac{d-2}{2}} d p_{cl}^2 \\ &\Downarrow d = n-2 = 2-\epsilon \\ &= \frac{\pi^{-\epsilon}}{\Gamma(-\epsilon)} (p_{cl}^2)^{-\epsilon} d p_{cl}^2 \end{aligned}$$

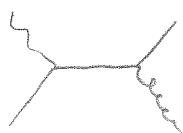
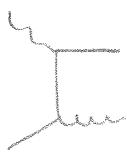
$$\begin{aligned} dp_s^{(2)} &= \frac{dz_n}{2z_n} * \frac{\pi^{-\epsilon}}{\Gamma(-\epsilon)} (p_{cl}^2)^{-\epsilon} d p_{cl}^2 \frac{1}{(2\pi z)^{2-2\epsilon}} * \hat{z} z^2 \\ &* \delta\left[p_{cl}^2 - \frac{z^2 Q^2 \hat{z}(\hat{t}-\hat{z})(\hat{t}-\hat{z})}{\hat{z}}\right] \end{aligned}$$

$$= \left(dz_n \frac{1}{\hat{z}}\right) * \frac{1}{8\pi} \left(\frac{4\pi}{Q^2}\right)^\epsilon \frac{1}{\Gamma(-\epsilon)} \left[\hat{z}(\hat{t}-\hat{z})\right]^{-\epsilon} \left[(\hat{t}-\hat{z})^{-\epsilon} \hat{z}^\epsilon\right]$$

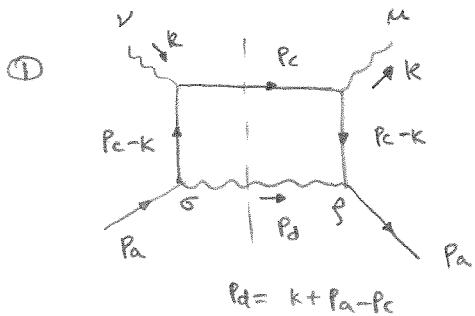
thus for spin-averaged one

$$\frac{d\sigma}{dx dy dz} = \frac{4\pi m^2}{Q^2} \frac{1+(t-4)^2}{2y} \int \frac{dx}{x} \frac{dz}{z} f_{q\bar{q}}(x) D_{q\rightarrow q}(z) [-g^{\mu\nu} H_{\mu\nu}]$$
$$* \frac{1}{8\pi} \left(\frac{4\pi}{Q^2}\right) \frac{\epsilon}{\Gamma(\ell+\epsilon)} z^{-\epsilon} (-z)^{-\epsilon} \bar{z}^\epsilon (-\bar{z})^{-\epsilon}$$

Let's study unpolarized cross-section first



$$k^2 = -Q^2$$



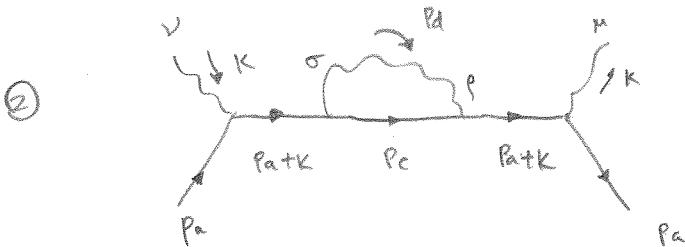
$$\text{define } \hat{s} = (p_a + k)^2 = -Q^2 + 2p_a \cdot k$$

$$\hat{t} = (p_c - k)^2 = -Q^2 - 2p_c \cdot k$$

$$\hat{u} = (p_a - p_c)^2 = -2p_a \cdot p_c$$

$$\text{Fig 1} = \frac{1}{2} \text{Tr} [\gamma_\mu \gamma^\nu (k - k) \gamma^\mu k_\nu \gamma^\nu (k - k) \gamma^\sigma] (-g_{\mu\nu}) d\phi_s(p_d)$$

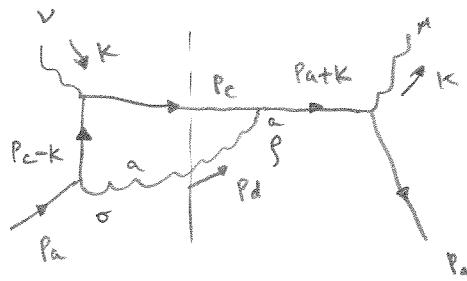
$$* \left[\frac{1}{(P_c - k)^2} \right]^2 * g_s^2$$



$$\text{Fig 2} = \frac{1}{2} \text{Tr} [\gamma_\mu \gamma^\nu (p_a + k) \gamma^\mu k_\nu \gamma^\sigma (p_a + k) \gamma^\nu] (-g_{\mu\nu}) d\phi_s(p_d)$$

$$* \left[\frac{1}{(P_a + k)^2} \right]^2 * g_s^2$$

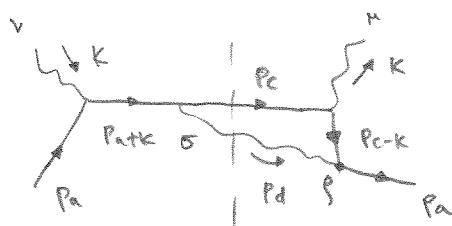
(3)



$$\omega_{\mu\nu} = \frac{1}{N} \text{Tr}[\tau_\alpha \tau^\alpha] = C_F$$

$$\text{Fig 3} = \frac{1}{2} \text{Tr} [\gamma_\alpha \gamma^\mu (\gamma_\mu + K) \gamma^\rho \gamma_c \gamma^\nu (\gamma_c - K) \gamma^\sigma] (-g_{\mu\nu}) d_{\rho\sigma}(p_a) \\ * \frac{1}{(p_c - k)^2} \frac{1}{(p_{a+K})^2}$$

(4)



$$\text{Fig 4} = \frac{1}{2} \text{Tr} [\gamma_\alpha \gamma^\rho (\gamma_c - K) \gamma^\mu \gamma_c \gamma^\sigma (\gamma_a + K) \gamma^\nu] (-g_{\mu\nu}) d_{\rho\sigma}(p_a) \\ * \frac{1}{(p_c - k)^2} \frac{1}{(p_{a+K})^2}$$

$$\text{Eq 2} + 2+3+4 = 4(1-\epsilon) \frac{1}{\hat{s}\hat{t}} \left[-(1-\epsilon)(\hat{s}^2 + \hat{t}^2) + 2\epsilon \hat{s}\hat{t} - 2Q^2 \underbrace{(\hat{Q}^2 + \hat{s} + \hat{t})}_{-\hat{u}} \right]$$

$$= 4(1-\epsilon) \left[(1-\epsilon) \left(\frac{\hat{s}}{-\hat{t}} + \frac{-\hat{t}}{\hat{s}} \right) + \frac{2Q^2 \hat{u}}{\hat{s}\hat{t}} + 2\epsilon \right]$$

Eventually we have

$$\begin{aligned} \frac{d\sigma}{dx dy dz_h} &= \frac{d\sigma_{em}^2}{Q^2} \frac{1+(1-y)^2}{2y} \int \frac{dx}{x} \frac{dz}{z} f_{q/p}(x) D_{q+h}(z) \\ &\quad * (g_S \mu^\epsilon)^2 * 4(1-\epsilon) \left[(1-\epsilon) \left(-\frac{\hat{s}}{\hat{t}} - \frac{\hat{t}}{\hat{s}} \right) + \frac{2Q^2 \hat{u}}{\hat{s}\hat{t}} + 2\epsilon \right] \\ &\quad * \frac{1}{8\pi} \left(\frac{4\pi}{Q^2} \right)^\epsilon \frac{1}{\Gamma(1-\epsilon)} \hat{z}^{-\epsilon} (\hat{-z})^{-\epsilon} \hat{x}^\epsilon (\hat{-x})^{-\epsilon} \\ &= \frac{\frac{2\pi d\sigma_{em}^2}{Q^2}}{y} \frac{1+(1-y)^2}{y} * \frac{ds}{2\pi} \int \frac{dx}{x} \frac{dz}{z} f_{q/p}(x) D_{q+h}(z) \\ &\quad * \left(\frac{4\pi \mu^2}{Q^2} \right)^\epsilon \frac{1}{\Gamma(1-\epsilon)} \hat{z}^{-\epsilon} (\hat{-z})^{-\epsilon} \hat{x}^\epsilon (\hat{-x})^{-\epsilon} \\ &\quad * (1-\epsilon) \left[(1-\epsilon) \left(-\frac{\hat{s}}{\hat{t}} - \frac{\hat{t}}{\hat{s}} \right) + \frac{2Q^2 \hat{u}}{\hat{s}\hat{t}} + 2\epsilon \right] \end{aligned}$$

define (like before) $\sigma_0 = \frac{2\pi d\sigma_{em}^2}{Q^2} \frac{1+(1-y)^2}{y} (1-\epsilon)$

$$\frac{dx}{x} = \frac{d\hat{x}}{\hat{x}} \quad \frac{dz}{z} = \frac{d\hat{z}}{\hat{z}}$$

Color = c_F

$$\frac{d\sigma}{dx dy dz_h} = \sigma_0 \frac{ds}{2\pi} \int \frac{dx}{x} \frac{dz}{z} f_{q/p}(x) D_{q+h}(z)$$

$$* \left(\frac{4\pi \mu^2}{Q^2} \right)^\epsilon \frac{1}{\Gamma(1-\epsilon)} \hat{z}^{-\epsilon} (\hat{-z})^{-\epsilon} \hat{x}^\epsilon (\hat{-x})^{-\epsilon}$$

$$* \left[(1-\epsilon) \left(-\frac{\hat{s}}{\hat{t}} - \frac{\hat{t}}{\hat{s}} \right) + \frac{2Q^2 \hat{u}}{\hat{s}\hat{t}} + 2\epsilon \right]$$

$$\hat{s} = \frac{1-\hat{x}}{\hat{x}} Q^2 \quad \hat{t} = -\frac{1-\hat{z}}{\hat{x}\hat{z}} Q^2 \quad \hat{u} = -\frac{\hat{z}}{\hat{x}} Q^2$$

$$[\dots] = \left\{ (-\epsilon) \left[\frac{1-\hat{x}}{1-\hat{z}} + \frac{1-\hat{z}}{1-\hat{x}} \right] + \frac{2\hat{x}}{1-\hat{x}} \frac{\hat{z}}{1-\hat{z}} + 2\epsilon \right\}$$

$$\frac{d\sigma}{dx dy dz} = \sigma_0 \frac{dc}{2\pi} \int \frac{dx}{x} \frac{dz}{z} f_{g_F}(x) D_{q \rightarrow h}(z)$$

$$* \left(\frac{4\pi \mu^2}{Q^2} \right) \epsilon \frac{1}{\Gamma(-\epsilon)} \hat{z}^{-\epsilon} (-\hat{z})^{-\epsilon} \hat{x} \in (-\hat{x})^{-\epsilon}$$

$$* \left[(-\epsilon) \left(\frac{1-\hat{x}}{1-\hat{z}} + \frac{1-\hat{z}}{1-\hat{x}} \right) + \frac{2\hat{x}}{1-\hat{x}} \frac{\hat{z}}{1-\hat{z}} + 2\epsilon \right]$$

$$\begin{aligned} \hat{z}^{-\epsilon} (-\hat{z})^{-\epsilon-1} &= -\frac{1}{\epsilon} \delta(-\hat{z}) + \frac{1}{(-\hat{z})_+} - \epsilon \left(\frac{\ln(-\hat{z})}{1-\hat{z}} \right)_+ - \epsilon \frac{\ln \hat{z}}{1-\hat{z}} + O(\epsilon^2) \\ \hat{x} \in (-\hat{x})^{-\epsilon} &= (-\hat{x}) \left[1 + \epsilon \ln \frac{\hat{x}}{1-\hat{x}} \right] \end{aligned}$$

$$\hat{z}^{-\epsilon} (-\hat{z})^{+\epsilon} = (-\hat{z}) \left[1 - \epsilon (\ln \hat{z} + \ln(-\hat{z})) \right]$$

$$\hat{x}^\epsilon (-\hat{x})^{-\epsilon-1} = -\frac{1}{\epsilon} \delta(-\hat{x}) + \frac{1}{(-\hat{x})_+} - \epsilon \left(\frac{\ln(-\hat{x})}{1-\hat{x}} \right)_+ - \epsilon \frac{\ln \hat{x}}{1-\hat{x}} + O(\epsilon^2)$$

$$\hat{z}^{+\epsilon} (-\hat{z})^{-\epsilon-1} = -\frac{1}{\epsilon} \delta(-\hat{z}) + \frac{\hat{z}}{(-\hat{z})_+} - \epsilon \hat{z} \left(\frac{\ln(-\hat{z})}{1-\hat{z}} \right)_+ - \epsilon \frac{\hat{z}}{1-\hat{z}} \ln \hat{z}$$

$$\hat{x}^{1+\epsilon} (-\hat{x})^{-\epsilon-1} = -\frac{1}{\epsilon} \delta(-\hat{x}) + \frac{\hat{x}}{(-\hat{x})_+} - \epsilon \hat{x} \left(\frac{\ln(-\hat{x})}{1-\hat{x}} \right)_+ + \epsilon \frac{\hat{x}}{1-\hat{x}} \ln \hat{x}$$

$$\hat{z}^{-\epsilon} (-\hat{z})^{-\epsilon} = 1 - \epsilon (\ln \hat{z} + \ln(-\hat{z}))$$

$$\hat{x}^\epsilon (-\hat{x})^{-\epsilon} = 1 + \epsilon \ln \frac{\hat{x}}{1-\hat{x}}$$

$$\frac{d\sigma}{dx dy dz} = \sigma_0 \frac{dz}{2\pi} \int \frac{dx}{x} \frac{dz}{z} f_{g \rightarrow h}(x) D_{g \rightarrow h}(z) \left(\frac{4\pi \mu^2}{Q^2} \right)^{\epsilon} \frac{1}{\Gamma(\epsilon)} \frac{1}{\Gamma(\epsilon)}$$

$$* \left\{ (\epsilon) \left[-\frac{1}{\epsilon} \delta(\epsilon z) + \frac{1}{(\epsilon z)_+} \right] \left[1 + \epsilon \ln \frac{\hat{x}}{\epsilon x} \right] (\epsilon x) \right.$$

$$+ (\epsilon) (\epsilon z) \left[1 + \epsilon \ln \hat{z} (\epsilon z) \right] \left[-\frac{1}{\epsilon} \delta(\epsilon - \hat{x}) + \frac{1}{(\epsilon - \hat{x})_+} \right]$$

$$+ 2 \left[-\frac{1}{\epsilon} \delta(\epsilon - \hat{x}) + \frac{\hat{x}}{(\epsilon - \hat{x})_+} - \epsilon \hat{x} \left(\frac{\ln(\epsilon - \hat{x})}{\epsilon - \hat{x}} \right)_+ + \epsilon \hat{x} \frac{\ln \hat{x}}{\epsilon - \hat{x}} \right]$$

$$* \left[-\frac{1}{\epsilon} \delta(\epsilon - \hat{z}) + \frac{\hat{z}}{(\epsilon - \hat{z})_+} - \epsilon \hat{z} \left(\frac{\ln(\epsilon - \hat{z})}{\epsilon - \hat{z}} \right)_+ - \epsilon \hat{z} \frac{\ln \hat{z}}{\epsilon - \hat{z}} \right]$$

$$+ 2 \epsilon \left\{ \right.$$

$$\{ \dots \} = (\epsilon x) \left[-\frac{1}{\epsilon} \underbrace{\delta(\epsilon z)}_{\sim} + \frac{1}{(\epsilon z)_+} + \left(1 - \ln \frac{z}{\epsilon x} \right) \delta(\epsilon z) \right]$$

$$+ (\epsilon z) \left[-\frac{1}{\epsilon} \underbrace{\delta(\epsilon x)}_{\sim} + \frac{1}{(\epsilon x)_+} + \left(1 + \ln \hat{z} (\epsilon z) \right) \delta(\epsilon x) \right]$$

$$+ 2 \left[-\frac{1}{\epsilon^2} \delta(\epsilon - \hat{x}) \delta(\epsilon - \hat{z}) - \frac{1}{\epsilon} \delta(\epsilon - \hat{x}) \frac{\hat{z}}{(\epsilon - \hat{z})_+} - \frac{1}{\epsilon} \delta(\epsilon - \hat{z}) \frac{\hat{x}}{(\epsilon - \hat{x})_+} \right.$$

$$\left. + \frac{\hat{x} \hat{z}}{(\epsilon - \hat{x})_+ (\epsilon - \hat{z})_+} + \delta(\epsilon - \hat{z}) \left(\hat{x} \left(\frac{\ln(\epsilon - \hat{z})}{\epsilon - \hat{z}} \right)_+ - \hat{x} \frac{\ln \hat{z}}{\epsilon - \hat{z}} \right) \right]$$

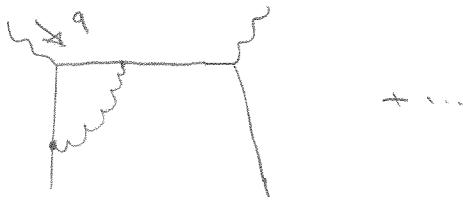
$$+ \delta(\epsilon - \hat{x}) \left(\hat{z} \left(\frac{\ln(\epsilon - \hat{x})}{\epsilon - \hat{x}} \right)_+ + \hat{z} \frac{\ln \hat{x}}{\epsilon - \hat{x}} \right) \right]$$

$$= \frac{2}{\epsilon^2} \delta(\epsilon x) \delta(\epsilon z) - \frac{1}{\epsilon} \delta(\epsilon x) \frac{1 + \hat{z}^2}{(\epsilon - \hat{z})_+} - \frac{1}{\epsilon} \delta(\epsilon - \hat{z}) \frac{1 + \hat{x}^2}{(\epsilon - \hat{x})_+}$$

$$+ \frac{1 + (\epsilon - \hat{x} - \hat{z})^2}{(\epsilon - \hat{x})_+ (\epsilon - \hat{z})_+} + \delta(\epsilon - \hat{z}) \left[(\epsilon - \hat{x}) \left(1 - \ln \frac{\hat{x}}{\epsilon x} \right) + 2 \hat{x} \left(\frac{\ln(\epsilon - \hat{x})}{\epsilon - \hat{x}} \right)_+ - 2 \hat{x} \frac{\ln \hat{x}}{\epsilon - \hat{x}} \right]$$

$$+ \delta(\epsilon - \hat{x}) \left[(\epsilon - \hat{z}) \left(1 + \ln \hat{z} (\epsilon z) \right) + 2 \hat{z} \left(\frac{\ln(\epsilon - \hat{z})}{\epsilon - \hat{z}} \right)_+ + 2 \hat{z} \frac{\ln \hat{z}}{\epsilon - \hat{z}} \right]$$

Now for virtual diagram



$$\Gamma^k(q) = g^k \left\{ 1 + \frac{ds}{4\pi} \ln \left(\frac{4\pi m^2}{-q^2} \right) \epsilon \frac{\Gamma(1+\epsilon) \Gamma^2(1-\epsilon)}{\Gamma(1-2\epsilon)} \left(-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} - 8 \right) \right\}$$

$2 \operatorname{Re} (\text{Virtual} * \text{lowest order})$

$$\Rightarrow \frac{ds}{2\pi} \ln \left(\frac{4\pi m^2}{\omega^2} \right) \epsilon \frac{1}{\Gamma(1-\epsilon)}$$

$$* \frac{\Gamma(1+\epsilon) \Gamma^2(1-\epsilon)}{\Gamma(1-2\epsilon)} \left(-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} - 8 \right)$$

$$\Downarrow = \left(-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} - 8 \right)$$

$$= \frac{ds}{2\pi} \ln \left(\frac{4\pi m^2}{\omega^2} \right) \epsilon \frac{1}{\Gamma(1-\epsilon)} \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} - 8 \right]$$

Note

$$2\hat{x} \left(\frac{\ln(1-\hat{x})}{1-\hat{x}} \right)_+ = [1+x^2 - (1-x)^2] \left(\frac{\ln(1-x)}{1-x} \right)_+$$

$$= (1+x^2) \left(\frac{\ln(1-x)}{1-x} \right)_+ - (1-x) \ln(1-x)$$

likewise for \hat{z} , we thus have (Real + virtual)

$$\frac{d\sigma}{dx dy dz_h} = \sigma_0 \frac{ds}{2\pi} \int \frac{dx}{x} \frac{dz}{z} f_{q/p}(x) D_{q \rightarrow L}(z) \left(\frac{4\pi\mu^2}{Q^2} \right) \epsilon \frac{1}{F(1-\epsilon)} \times \left[\begin{array}{l} \left\{ -\frac{1}{\epsilon} \delta(1-\hat{x}) \left[\frac{1+\hat{x}^2}{(1-\hat{x})_+} + \frac{1}{2} \delta(1-\hat{x}) \right] \right. \\ \left. -\frac{1}{\epsilon} \delta(1-\hat{z}) \left[\frac{1+\hat{z}^2}{(1-\hat{z})_+} + \frac{1}{2} \delta(1-\hat{z}) \right] \right\} \\ + \left\{ \frac{1+(1-x-\hat{x})^2}{(1-x)+(1-\hat{x})_+} \right. \\ \left. + \delta(1-\hat{x}) \left[(1+\hat{x}^2) \left(\frac{\ln(1-\hat{x})}{1-\hat{x}} \right)_+ - \frac{1+\hat{x}^2}{1-\hat{x}} \ln \hat{x} + (1-\hat{x}) \right] \right. \\ \left. + \delta(1-\hat{z}) \left[(1+\hat{z}^2) \left(\frac{\ln(1-\hat{z})}{1-\hat{z}} \right)_+ + \frac{1+\hat{z}^2}{1-\hat{z}} \ln \hat{z} + (1-\hat{z}) \right] \right. \\ \left. - \delta(1-\hat{x}) \delta(1-\hat{z}) \right\} \end{array} \right]$$

This result is consistent with NPB 160 (1979) 301

Altarelli-Ellis-Martinelli-pi

(after convert D2S scheme to MS scheme)

Expansion

$$\hat{z}^{-\epsilon} (1-\hat{z})^{-\epsilon-1} = -\frac{1}{\epsilon} \delta(1-\hat{z}) + \frac{1}{(1-\hat{z})_+} - \epsilon \left(\frac{\ln(1-\hat{z})}{1-\hat{z}} \right)_+ - \epsilon \frac{\ln \hat{z}}{1-\hat{z}}$$

$$\hat{x}^\epsilon (1-\hat{x})^{-\epsilon-1} = -\frac{1}{\epsilon} \delta(1-\hat{x}) + \frac{1}{(1-\hat{x})_+} - \epsilon \left(\frac{\ln(1-\hat{x})}{1-\hat{x}} \right)_+ + \epsilon \frac{\ln \hat{x}}{1-\hat{x}}$$

$$\hat{x}^\epsilon (1-\hat{x})^{-\epsilon} = 1 + \epsilon \ln \frac{\hat{x}}{1-\hat{x}}$$

$$\hat{z}^{-\epsilon} (1-\hat{z})^{-\epsilon} = 1 - \epsilon \ln \hat{z} - \epsilon \ln(1-\hat{z})$$

$$I_0 = z^{-\epsilon} (1-z)^{-\epsilon-1}$$

$$I = \int_0^1 dz z^{-\epsilon} (1-z)^{-\epsilon-1} f(z)$$

$$= \int_0^1 dz z^{-\epsilon} (1-z)^{-\epsilon-1} [f(z) - f(1) + f(1)]$$

$$= \int_0^1 dz z^{-\epsilon} (1-z)^{-\epsilon-1} f(1) + \int_0^1 dz z^{-\epsilon} (1-z)^{-\epsilon-1} [f(z) - f(1)]$$

$$= \frac{\Gamma(1-\epsilon) \Gamma(-\epsilon)}{\Gamma(1-2\epsilon)} f(1) + \int_0^1 dz [f(z) - f(1)] \left[\frac{1}{1-z} - \epsilon \frac{\ln(1-z)}{1-z} - \epsilon \frac{\ln z}{1-z} + O(\epsilon^2) \right]$$

$$= \left[-\frac{1}{\epsilon} + \frac{\pi^2}{6} \epsilon + O(\epsilon^2) \right] f(1) + \int_0^1 dz \frac{1}{1-z} (f(z) - f(1))$$

$$- \epsilon \int_0^1 dz \frac{\ln(1-z)}{1-z} (f(z) - f(1))$$

$$- \epsilon \int_0^1 dz \frac{\ln z}{1-z} (f(z) - f(1))$$

$$+ O(\epsilon^2)$$

Note $\frac{\ln z}{1-z} \rightarrow -1$ when $z \rightarrow 1$ thus finite

$$\int_0^1 dz \frac{\ln z}{1-z} = -\frac{\pi^2}{6}$$

also $\int_0^1 dz [w(z)]_+ f(z) = \int_0^1 dz w(z) (f(z) - f(1))$

$$= \left[-\frac{1}{\epsilon} + \frac{\pi^2}{6} \epsilon \right] f(1) + \int_0^1 dz \frac{1}{(1-z)_+} f(z) - \epsilon \int_0^1 dz \left(\frac{\ln(1-z)}{1-z} \right)_+ f(z)$$

$$- \epsilon \int_0^1 dz \frac{\ln z}{1-z} f(z) - \frac{\pi^2}{6} f(1) \epsilon + O(\epsilon^2)$$

$$= \int_0^1 dz f(z) \left[-\frac{1}{\epsilon} \delta(1-z) + \frac{1}{(1-z)_+} - \epsilon \left(\frac{\ln(1-z)}{1-z} \right)_+ - \epsilon \frac{\ln z}{1-z} + O(\epsilon^2) \right]$$

Thus $\boxed{z^{-\epsilon} (1-z)^{-\epsilon-1} = -\frac{1}{\epsilon} \delta(1-z) + \frac{1}{(1-z)_+} - \epsilon \left(\frac{\ln(1-z)}{1-z} \right)_+ - \epsilon \frac{\ln z}{1-z} + O(\epsilon^2)}$

Make sense the result

$$\left(\frac{4\pi \mu^2}{Q^2}\right)^{\epsilon} \frac{1}{\Gamma(1+\epsilon)} (-\frac{1}{\epsilon}) = \underbrace{-\frac{1}{\epsilon}}_{-\frac{1}{\epsilon} + \gamma_E - \ln(4\pi) + \ln(\frac{Q^2}{\mu^2}) + O(\epsilon)} (\overline{MS} \text{ scheme})$$

$$= -\frac{1}{\epsilon} + \ln \frac{\mu_F^2}{\mu^2} + \ln \left(\frac{Q^2}{\mu_F^2}\right) + O(\epsilon)$$

then $f_{qF}(x_B, \mu_F^2) = f_{qF}^{(0)}(x_B) + \frac{\alpha_S}{2\pi} \left(-\frac{1}{\epsilon} + \ln \frac{\mu_F^2}{\mu^2}\right) \int_{x_B}^1 \frac{dx}{x} f_{qF}(x_B) P_{qq}(\hat{x})$

where $P_{qq}(\hat{x}) = C_F \left[\frac{1+\hat{x}^2}{(1-\hat{x})_+} + \frac{3}{2} \delta(1-\hat{x}) \right]$

Thus $\frac{\partial}{\partial \ln \mu_F^2} f_{qF}(x_B, \mu_F^2) = \frac{\alpha_S}{2\pi} \int_{x_B}^1 \frac{dx}{x} f_{qF}(x_B, \mu_F^2) P_{qq}(\hat{x})$

In other words, we "reabsorb" the divergence (collinear)

into the redefinition of parton distribution function

(similar) for fragmentation function

to become "renormalized" PDFs and FFS

finally we have NLO result

$$\begin{aligned}
 \frac{d\sigma}{dx_b dy dz_h} &= \sigma_0 \frac{ds}{2\pi} \sum_q e_q^2 \int \frac{dx}{x} \frac{dz}{z} f_{q\bar{q}}(x, \mu_F^2) D_{q\bar{q}h}(z, \mu_F^2) \\
 &\times \left\{ \ln \frac{Q^2}{\mu_F^2} \left[P_{qq}(\hat{x}) \delta(1-\hat{z}) + P_{q\bar{q}}(\hat{z}) \delta(1-\hat{x}) \right] \right. \\
 &+ C_F \left[\frac{1+(1-\hat{x}-\hat{z})^2}{(1-\hat{x})+(1-\hat{z})_+} - g \delta(1-\hat{x}) \delta(1-\hat{z}) \right] \\
 &+ \delta(1-\hat{z}) C_F \left[(1+\hat{z}^2) \left(\frac{\ln(1-\hat{z})}{1-\hat{z}} \right)_+ - \frac{(1+\hat{z}^2)}{1-\hat{z}} \ln \hat{z} + (1-\hat{z}) \right] \\
 &+ \delta(1-\hat{x}) C_F \left[(1+\hat{z}^2) \left(\frac{\ln(1-\hat{x})}{1-\hat{x}} \right)_+ + \frac{1+\hat{x}^2}{1-\hat{x}} \ln \hat{x} + (1-\hat{x}) \right] \left. \right\}
 \end{aligned}$$