

Deeply Virtual Compton Scattering at Jefferson Lab

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Outline

- Introduction – physics motivation
- Experimental setup
- High Resolution Spectrometer optics optimization
- Outlook

Physics motivation

- Form factors \rightarrow partons position
- Parton distribution functions \rightarrow partons momentum

But no correlations between the two!

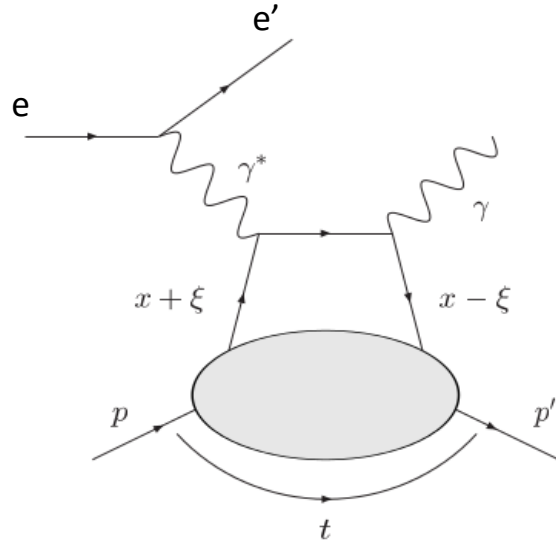
- \rightarrow Generalized Parton Distributions (GPDs) correlate longitudinal momentum and transverse position of partons.
- \rightarrow GPDs also give insights on nucleon spin structure.

GPDs accessible through Deeply Virtual Compton Scattering (DVCS) & interference with Bethe-Heitler.

My PhD subject: DVCS experiment at Jlab, Hall A!

- \rightarrow Extract DVCS cross-sections, as a function of Q^2 , for fixed values of x_B

The DVCS + Bethe-Heitler interactions $ep \rightarrow e'p'\gamma$



$Q^2 = -(e' - e)^2$: virtuality of γ^*

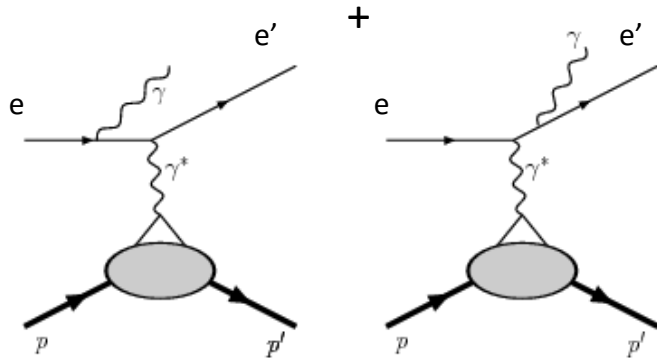
$v = E - E'$, energies of the electron before and after scattering

$$x_B = \frac{Q^2}{2Mv} \quad (\text{NB: } x_B \neq x)$$

$$\xi = \frac{x_B}{2 - x_B}$$

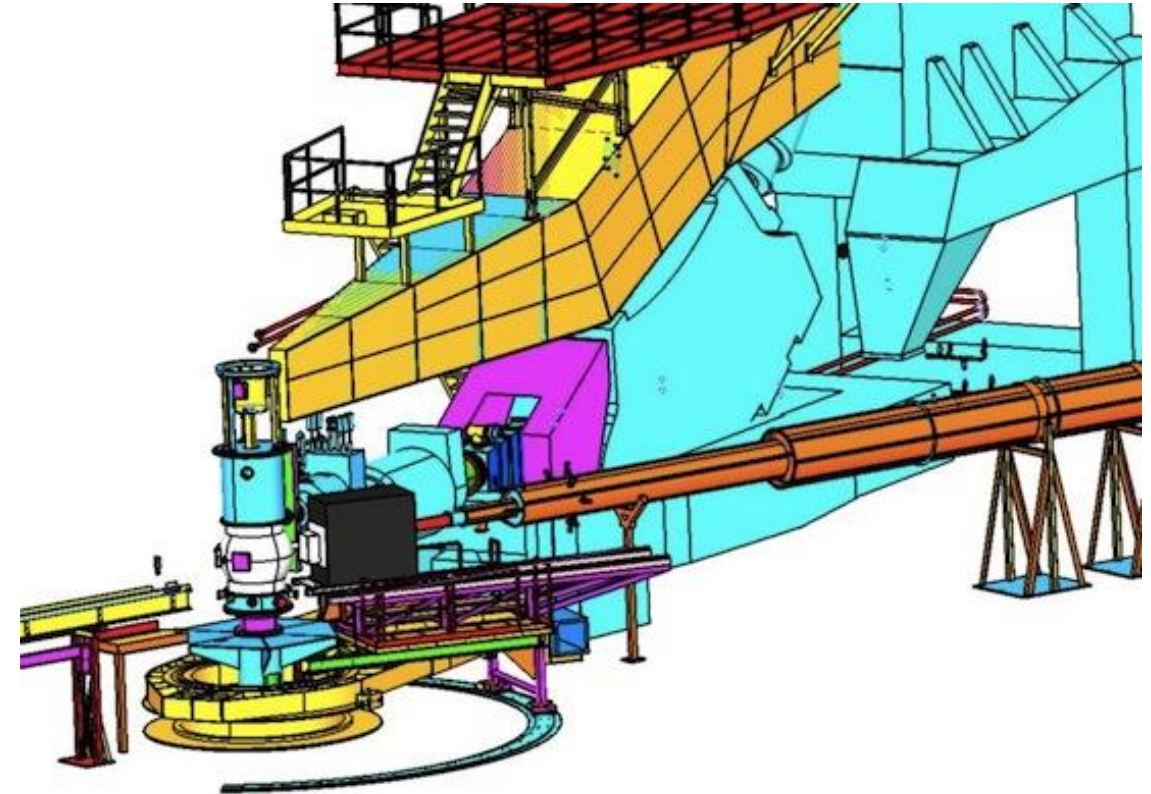
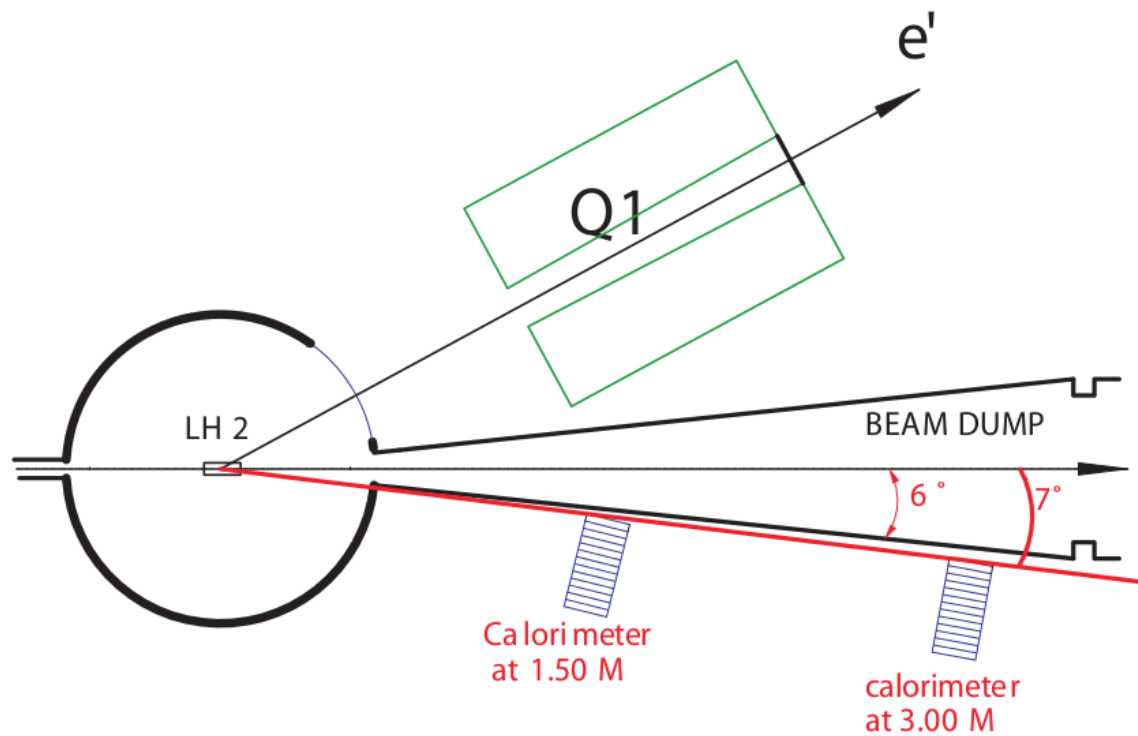
-2ξ : longitudinal momentum transfer to the struck quark.

$t = (p - p')^2$: squared momentum transfer to the proton

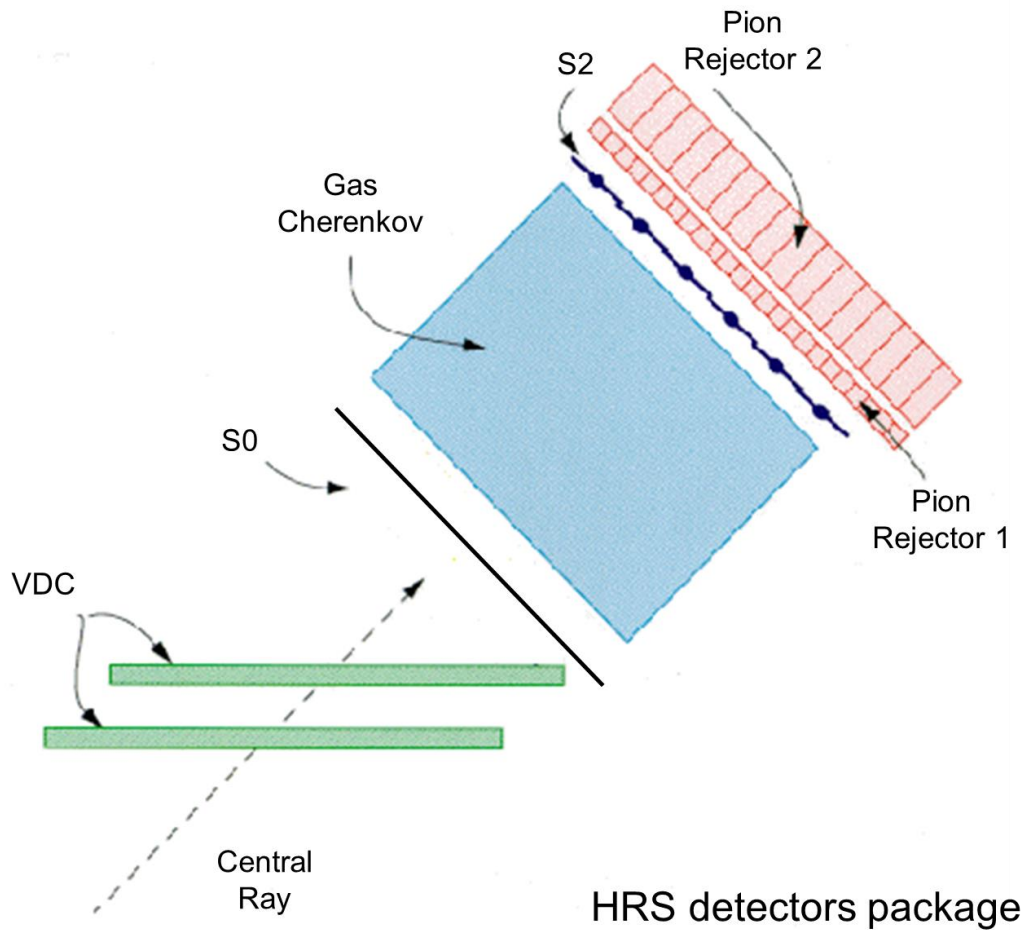


In the limit $Q^2 \rightarrow \infty$ and $v \rightarrow \infty$ but fixed x_B (Bjorken limit), the virtual photon γ^* interacts with a single quark in the proton.

The experimental setup

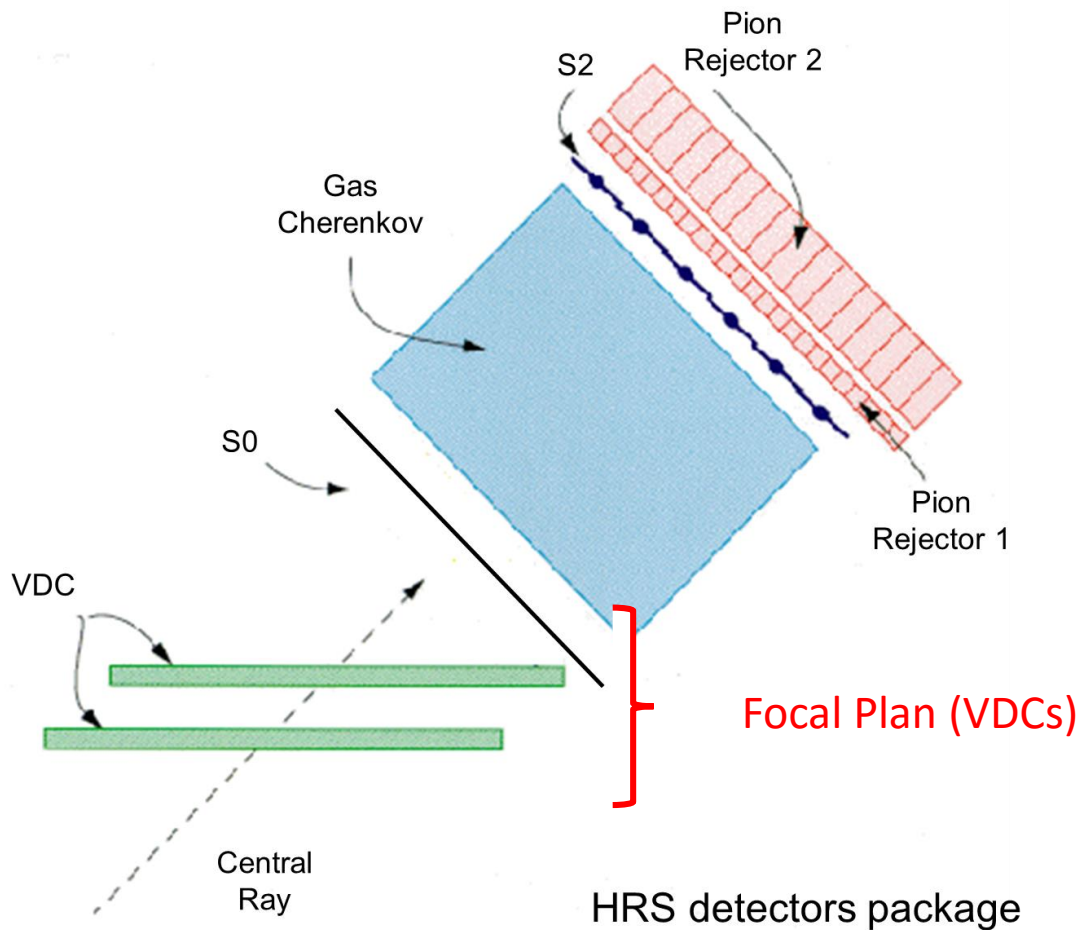


The LHRS detector package



- Vertical Drift Chambers : tracking
- Scintillators + Cherenkov : triggering
- Cherenkov + “Pion Rejector” (electromagnetic calorimeter) : particle identification (electrons VS pions)

The focal plan



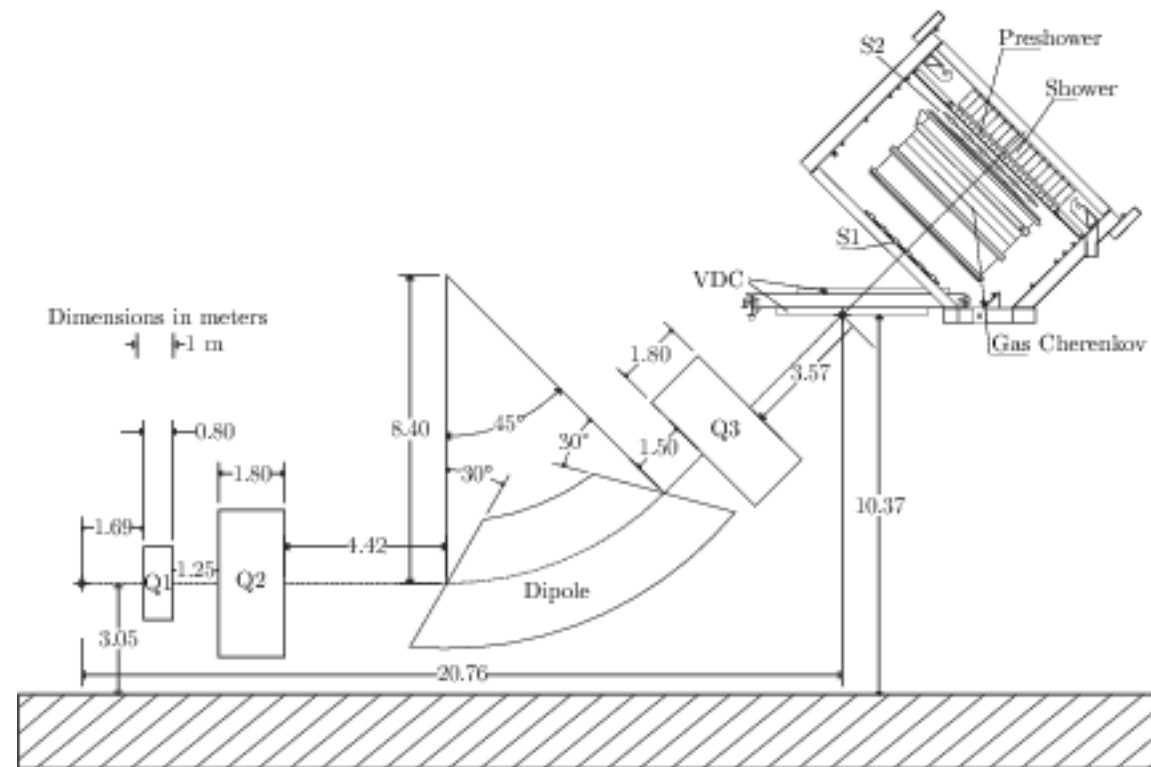
- LHRs focal plan ~ camera focal plan
- Electron detected at focal plan:
 - Position x_{fp}, y_{fp}
 - Angles θ_{fp}, ϕ_{fp} (electron trajectory VS central ray axis)
- From $x_{fp}, y_{fp}, \theta_{fp}$, and ϕ_{fp} : reconstruction at the target of :
 - Event vertex: y_{tg}
 - Angles θ_{tg}, ϕ_{tg} (electron trajectory VS central ray axis)
 - Momentum δ_{tg} ($p_{electron} - p_{LHRs}$)

Reconstructing vertex, trajectory & momentum

- How ?

LHRS : made of 4 superconducting magnets + detector package

Magnets : QQDQ configuration



The optics matrix

- Knowledge of how the electron trajectory is bent in the magnets allows the reconstruction.
 - Electric currents/Magnetic fields of each magnet “tuned” for specific bending.
- Optics Matrix

“On the paper” (1st order approximation):

$$\begin{bmatrix} \delta \\ \theta \\ y \\ \phi \end{bmatrix}_{tg} = \begin{bmatrix} \langle \delta | x \rangle & \langle \delta | \theta \rangle & 0 & 0 \\ \langle \theta | x \rangle & \langle \theta | \theta \rangle & 0 & 0 \\ 0 & 0 & \langle y | y \rangle & \langle y | \phi \rangle \\ 0 & 0 & \langle \phi | y \rangle & \langle \phi | \phi \rangle \end{bmatrix} \begin{bmatrix} x \\ \theta \\ y \\ \phi \end{bmatrix}_{fp}$$

The optics matrix coefficients

In real life:

$$y_{tg} = \sum_{j,k,l} Y_{jkl} \theta_{fp}^j y_{fp}^k \phi_{fp}^l,$$

$$Y_{jkl} = \sum_{i=1}^m C_i^{Y_{jkl}} x_{fp}^i$$

$$\theta_{tg} = \sum_{j,k,l} T_{jkl} \theta_{fp}^j y_{fp}^k \phi_{fp}^l,$$

$$\phi_{tg} = \sum_{j,k,l} P_{jkl} \theta_{fp}^j y_{fp}^k \phi_{fp}^l,$$

$$\delta = \sum_{j,k,l} D_{jkl} \theta_{fp}^j y_{fp}^k \phi_{fp}^l,$$

$$y_{tg} = \sum_{j,k,l} \sum_{i=1}^m C_i^{Y_{jkl}} x_{fp}^i \theta_{fp}^j y_{fp}^k \phi_{fp}^l$$

$C_i^{Y_{jkl}}$ “Optics matrix coefficients”

LHRS optics optimization

- What is the problem ?

Superconducting magnet: maximum current threshold above which loss of superconductivity.

(you do not want that to happen)

But Q1 is old! → threshold lower than in his youth!

→ Threshold lower than experimental requirement

→ Required current cannot be reached, magnetic field too low

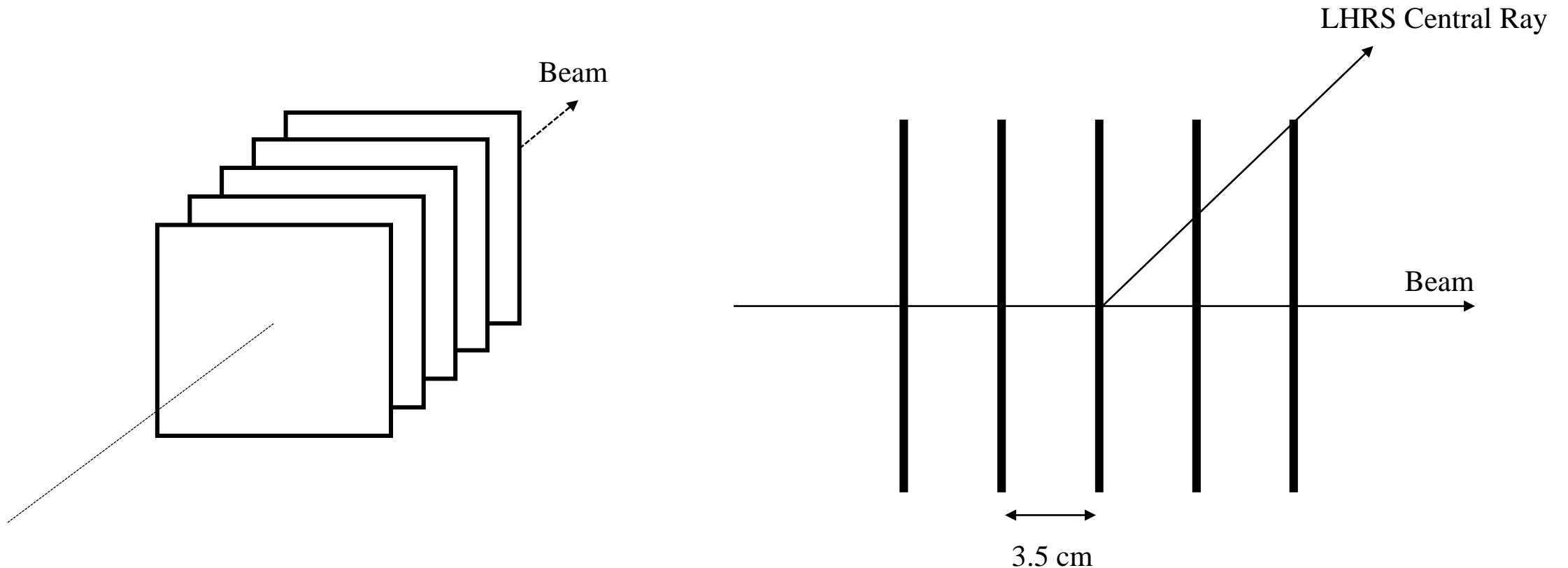
→ Different electron trajectory bending in magnet Q1

→ Q1 is “detuned”

→ Need to compute new optics matrix coefficients $C_i^{Y_{jkl}}$

Step 1 – optimization of the vertex reconstruction

- Data taken on a 5 thin carbon foils target (1mm thick)
→ Expected vertex values y_{tg}^0



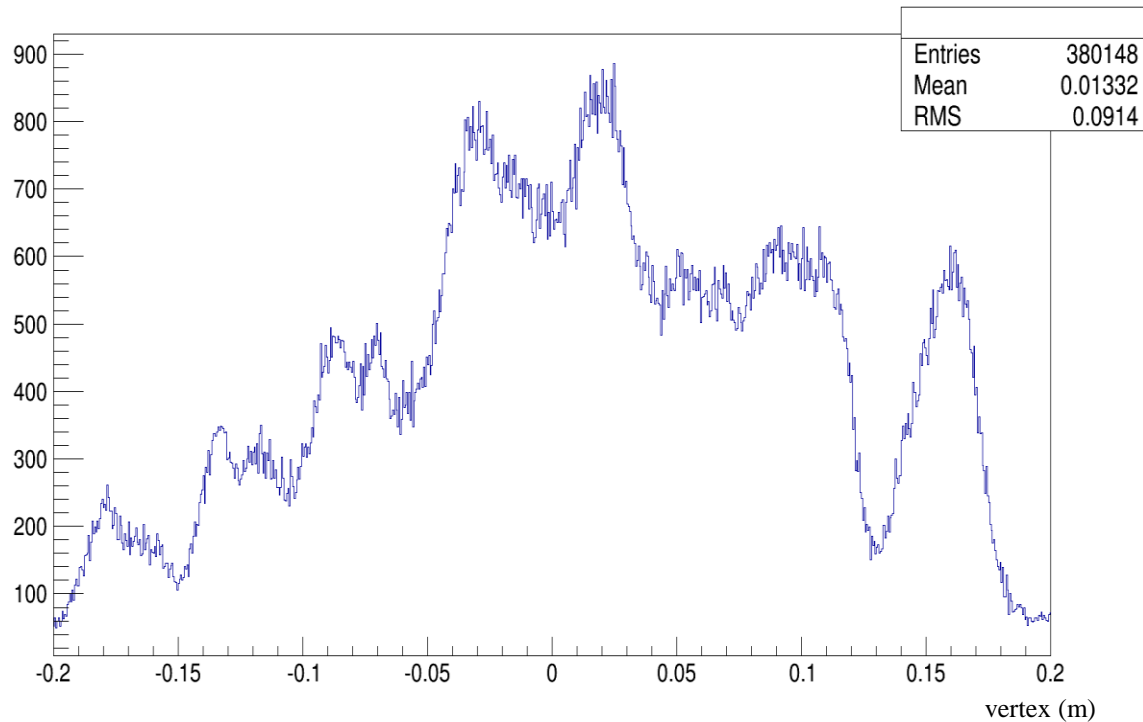
Step 1 – optimization of the vertex reconstruction

- Data taken on a 5 thin carbon foils target (1mm thick)
 - Expected vertex values y_{tg}^0
 - Computation of the new optics matrix coefficients $C_i^{Y_{jkl}}$ by minimizing:

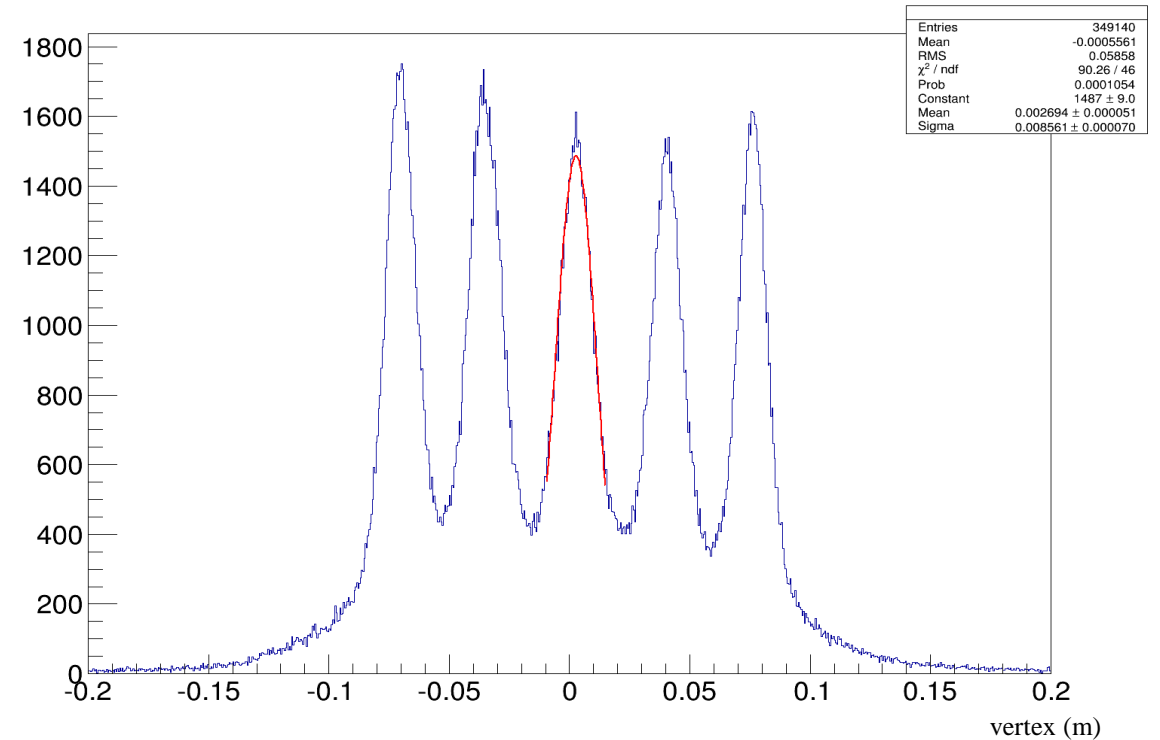
$$\Delta(y) = \sum_s \left[\frac{\sum_{j,k,l} Y_{jkl} \theta_{fp}^j y_{fp}^k \phi_{fp}^l - y_{tg}^0}{\sigma_y^s} \right]^2 \quad (\text{aberration function})$$

Step 1 – optimization of the vertex reconstruction

Vertex reconstruction



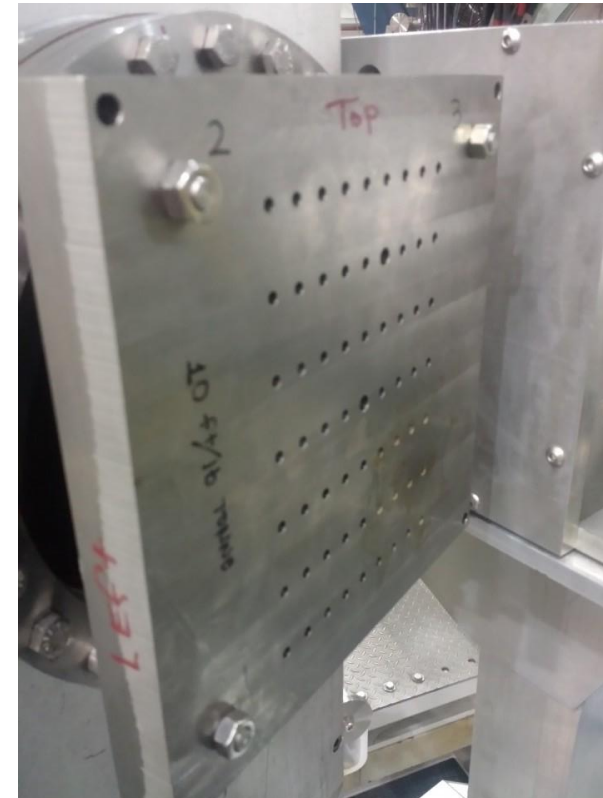
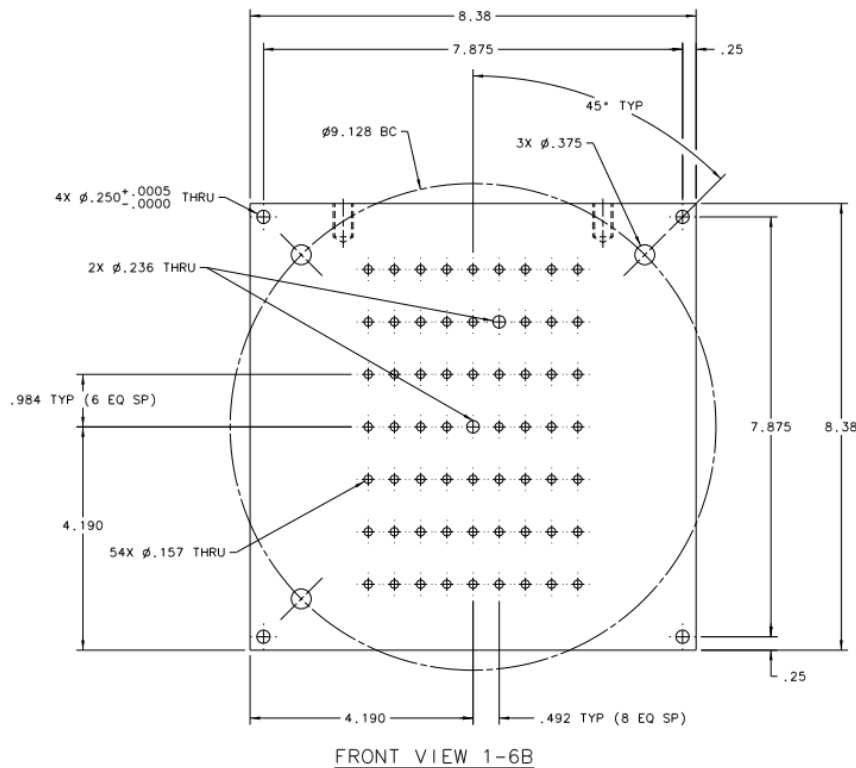
Before optimization



After optimization

Step 2 – optimization of angles reconstruction

- Data taken on a 5 thin carbon foils target (1mm thick)
- Thick metal plate with holes inserted in front of the LHRS entrance (Sieve)



Step 2 – optimization of angles reconstruction

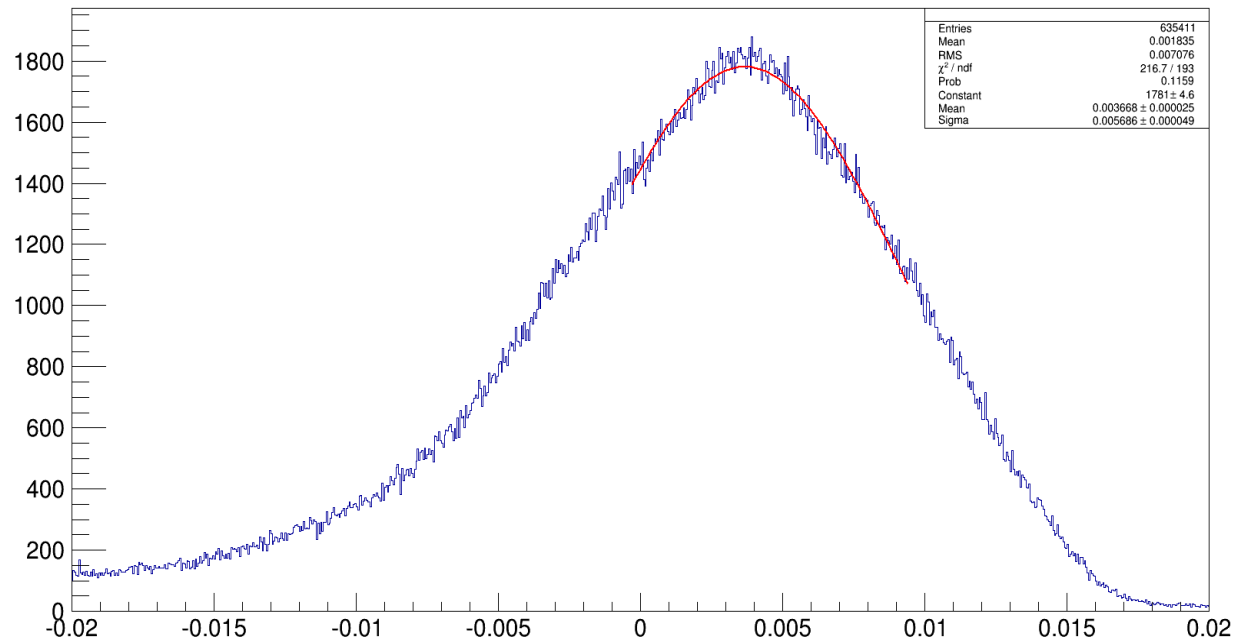
- Data taken on a 5 thin carbon foils target (1mm thick)
- Thick metal plate with holes inserted in front of the LHRS entrance (Sieve)
 - Holes = expected values for angles θ_{tg} and ϕ_{tg}
 - Computation of new optics matrix coefficients by minimization of aberration function.

Step 3 – optimization of the momentum reconstruction

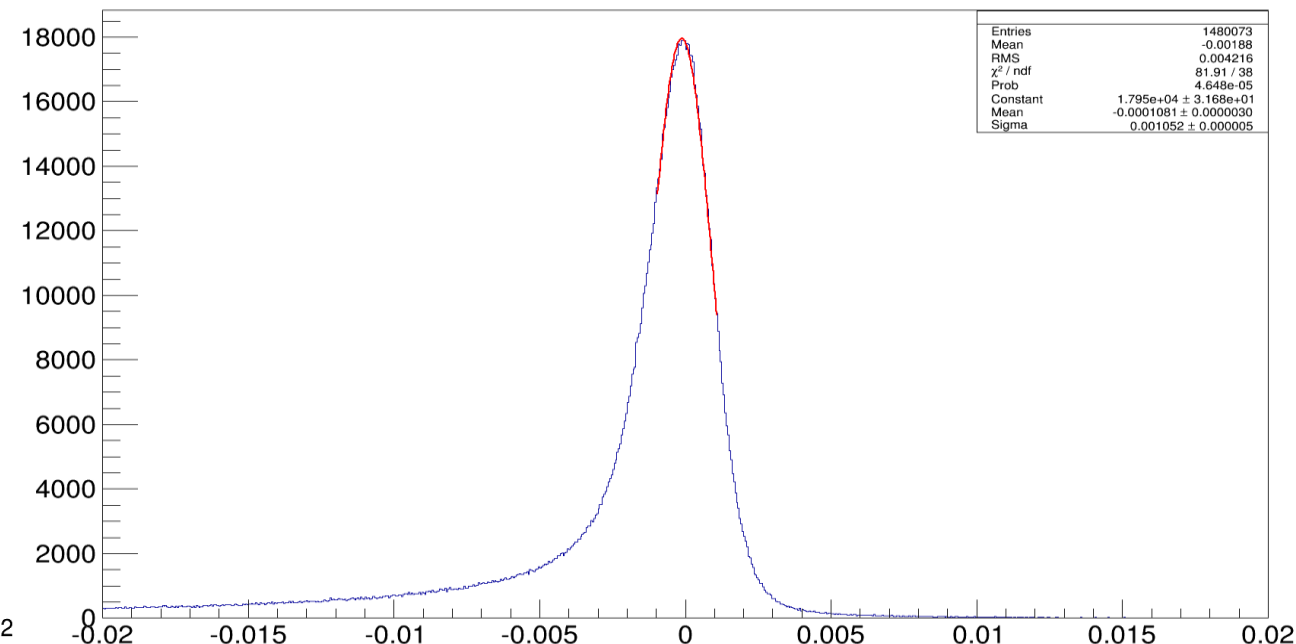
- Data taken on an LH₂ target, elastic scattering ep → ep setting
 - Constrained system: known scattering angle = known scattering momentum
 - Sieve out (not useful, would only decrease statistics)
 - “Delta Scan”
 - LHRS angle fixed
 - 5 runs varying central momentum setting (central momentum, ± 2%, ± 4%)
 - Central momentum selection made by varying magnetic fields
 - Because of elastic momentum-scattering angle correlation, need to “move” central ray to illuminate whole focal plane
- Expected values for momentum δ_{tg}
- Computation of new optics matrix coefficients by minimization of aberration function.

Step 3 – optimization of the momentum reconstruction

Relative momentum dp/p reconstruction



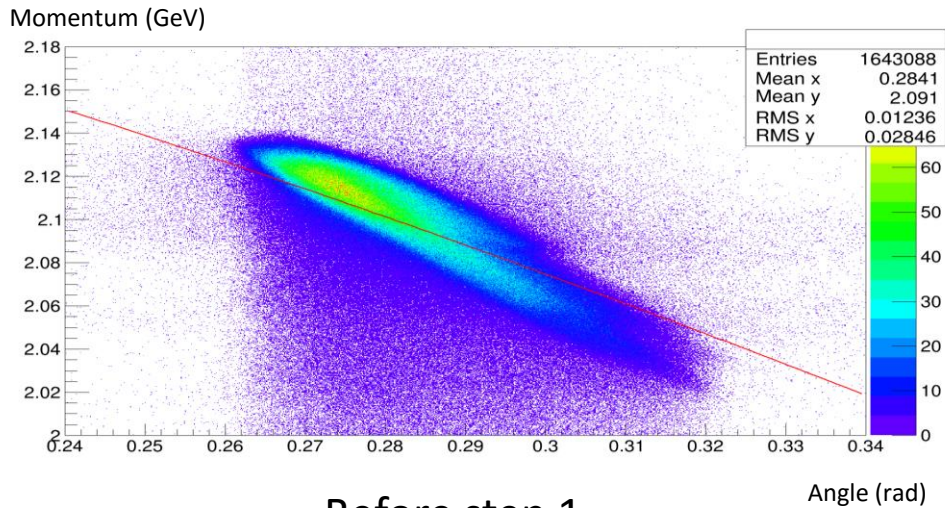
Before optimization



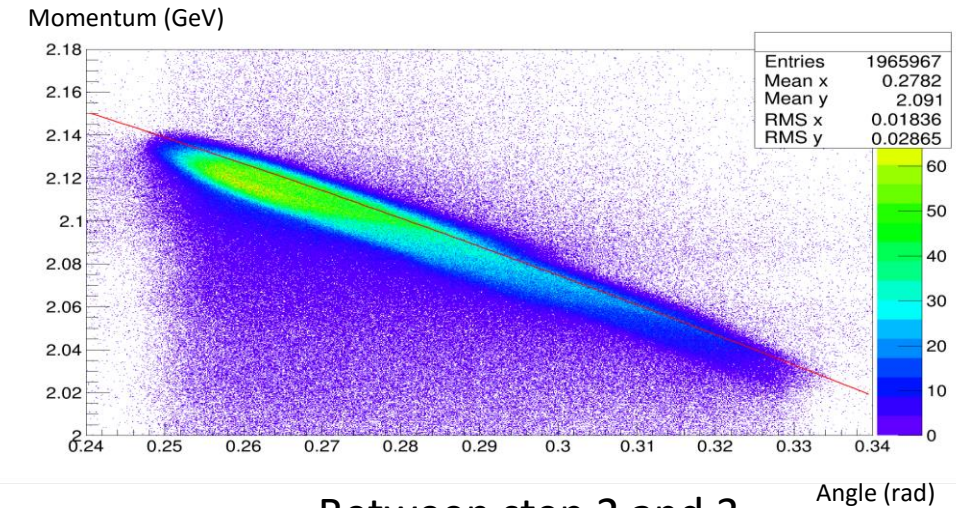
After optimization

Step 1, 2 and 3 – optimization of all 4 variables reconstruction

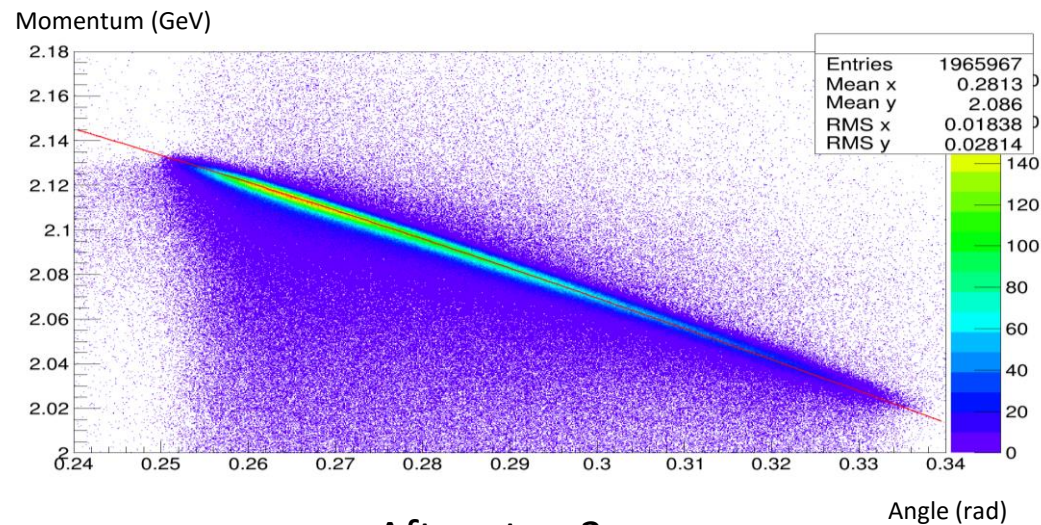
Elastic line: momentum – scattering angle



Before step 1



Between step 2 and 3



After step 3

Conclusion and outlook

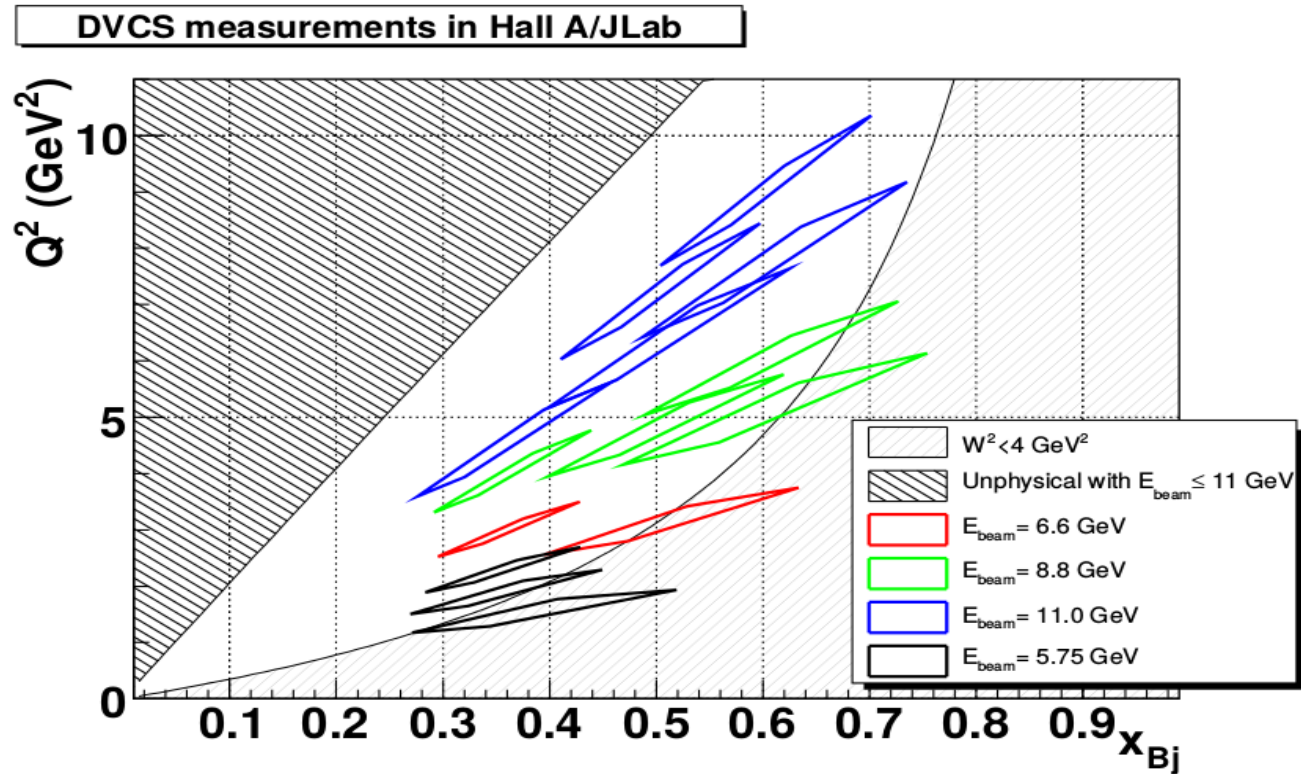
- DVCS experiment was running Winter/Spring 2016. Will come back Fall 2016.
 - Spectrometer optics calibration almost complete (small corrections).
 - However, still needs a thorough acceptance study.
 - And a lot more, done by the DVCS collaboration...
-
- Long term: DVCS cross-sections extraction

Thank you for your attention!

Questions ?

Annexes

DVCS experiment (Q^2, x_B) phase space



$W^2 = (p + q)^2$ invariant mass of the hadronic final state.

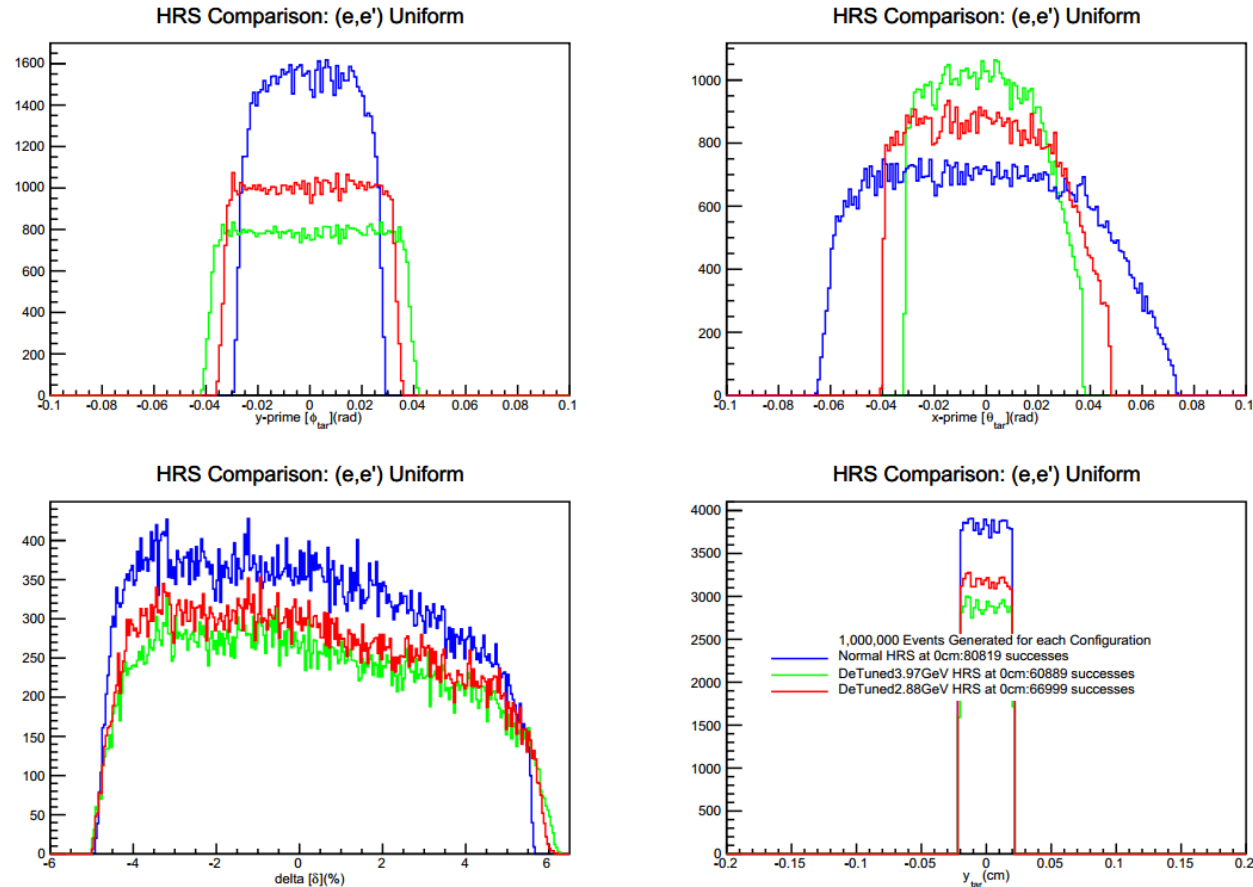
From DVCS cross-section to GPDs

DVCS cross-section is parametrized by GPDs, but integrated over x . We cannot directly extract GPDs from DVCS cross-section.

Extracted cross-section is proportional to $|BH|^2 + |DVCS|^2 + \text{Interference}$

- BH term: QED, well known and computed
- DVCS term: parametrized by bilinear combination of Compton Form Factors
 - Compton Form Factors are defined from GPDs combinations, integrated over x . Dependence on ξ and t .
 - Gives access to Compton Form Factors modulus
 - Needs several values of Q^2 and x_B to extract them
- Interference term: parametrized by real and imaginary parts of Compton Form Factors.
- DVCS + Interference terms allow to extract Compton Form Factors, from which GPDs are then extracted.

LHRS optics - acceptance study



Simulation shows that acceptance is affected. Will need further study to fully understand the new acceptance (correlations between parameters).