

# Linearly Polarized Gluons in $J/\psi$ and $\Upsilon$ Production

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# Outline

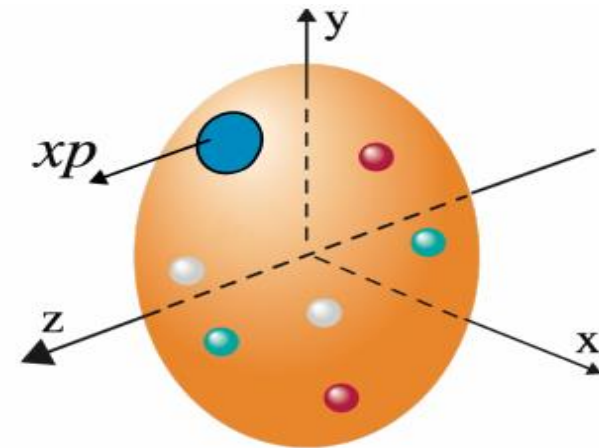
Gluon TMDs

Quarkonium Models

$J/\psi$  and  $\Upsilon$  production

Conclusion

# Transverse Momentum Dependent (TMD) Distributions

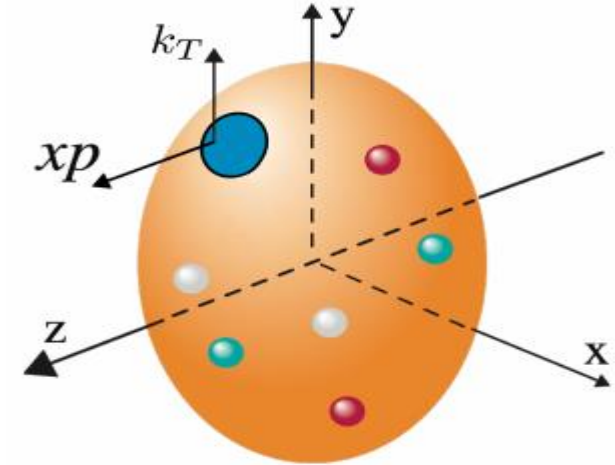


1D

Collinear pdf  $f(x)$

Universality

$$lp \rightarrow lX$$



3D

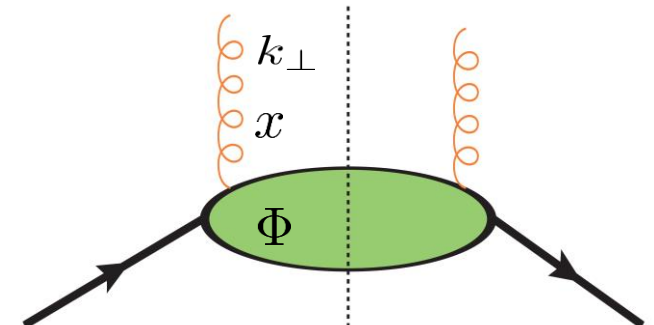
TMD pdf  $f(x, k_{\perp})$

Non-trivial Universality

$$lp \rightarrow lhX$$

$$pp \rightarrow hX$$

# Gluon Correlator



$$\Phi_g^{\mu\nu}(x, \mathbf{k}_\perp) = \frac{n_\rho n_\sigma}{(k \cdot n)^2} \int \frac{d(\lambda \cdot P) d^2 \lambda_T}{(2\pi)^3} e^{ik \cdot \lambda} \langle P | \text{Tr} [F^{\mu\rho}(\lambda) W(\lambda, 0) F^{\nu\sigma}(0)] | P \rangle |_{LF}$$

P. J. Mulders and J. Rodrigues PRD 63 094021(2001)

Parameterization of gluon correlator at “Leading Twist” is

$$= -\frac{1}{2x} \left\{ g_T^{\mu\nu} f_1^g(x, \mathbf{k}_\perp^2) - \left( \frac{k_\perp^\mu k_\perp^\nu}{M_h^2} + g_T^{\mu\nu} \frac{\mathbf{k}_\perp^2}{2M_h^2} \right) h_1^{\perp g}(x, \mathbf{k}_\perp^2) \right\}$$

Gluons	Unpolarized	Circularly	Linearly
Target			
Unpolarized	$f_1^g$		$h_1^{\perp g}$
Longitudinal		$g_{1L}^g$	
Transverse	$f_{1T}^{\perp g}$	$g_{1T}^g$	$h_{1T}^g, h_{1T}^{\perp g}$

D. Boer et al.  
arXive  
1507.05267

# Linearly Polarized Gluons

No experimental investigation has been carried out to extract the  $h_1^{\perp g}$  until now .

- ★ Theoretical upper bound of  $h_1^{\perp g}$  is only known.

D. Boer et al. PRL 106 132001 (2011)

## In proton-proton collision

- ★  $pp \rightarrow \gamma\gamma X$  at RHIC Qiu, Schlegel, Vogelsang, PRL 107, 062001 (2011)

- ★  $pp \rightarrow \Upsilon\gamma X$  at LHC Dunnen, Lansberg, Pisano, Schlegel, PRL 112, 212001 (2014)

- ★  $pp \rightarrow H X$  D. Boer et al. PRL 108, 032002 (2012)

- ★  $pp \rightarrow \eta_{c,b}$  or  $\chi_{c,b} + X$  at LHCb and AFTER D. Boer, C. Pisano, PRD 86, 094007 (2012)

## In electron-proton scattering

- ★  $ep \rightarrow eQ\bar{Q}X$  or  $e + jet + jet + X$  C. Pisano et al. JHEP 10 (2013) 024

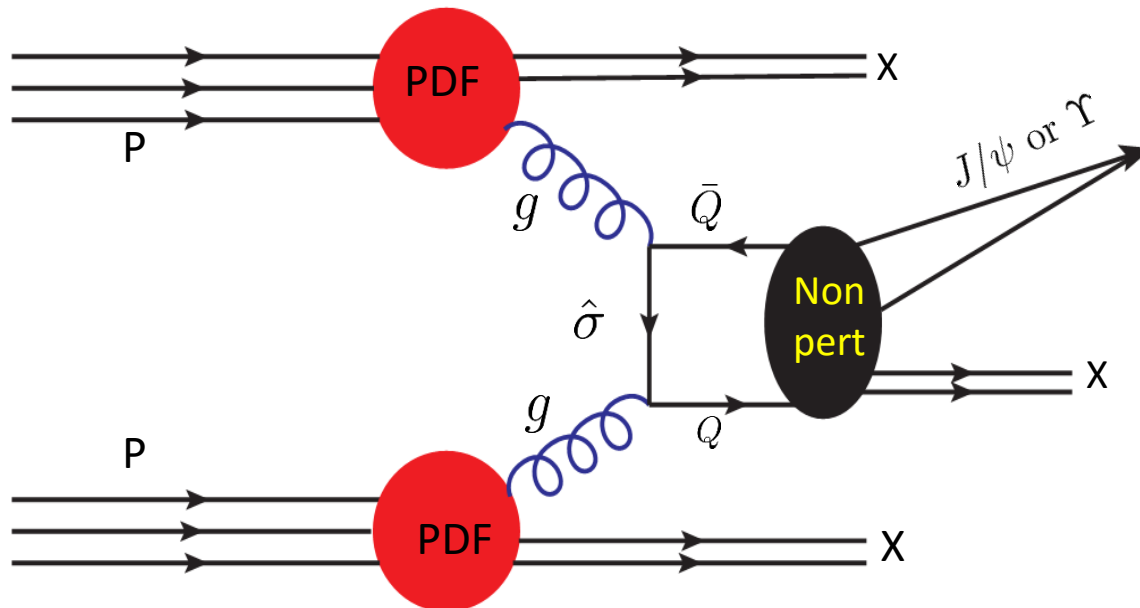
# Quarkonium Models

Color Singlet Model (CSM)

Color Octet Model (COM)

Color Evaporation Model (CEM)

$$\sigma_{J/\psi, \Upsilon} = \hat{\sigma} \times \text{Nonperturbative term}$$

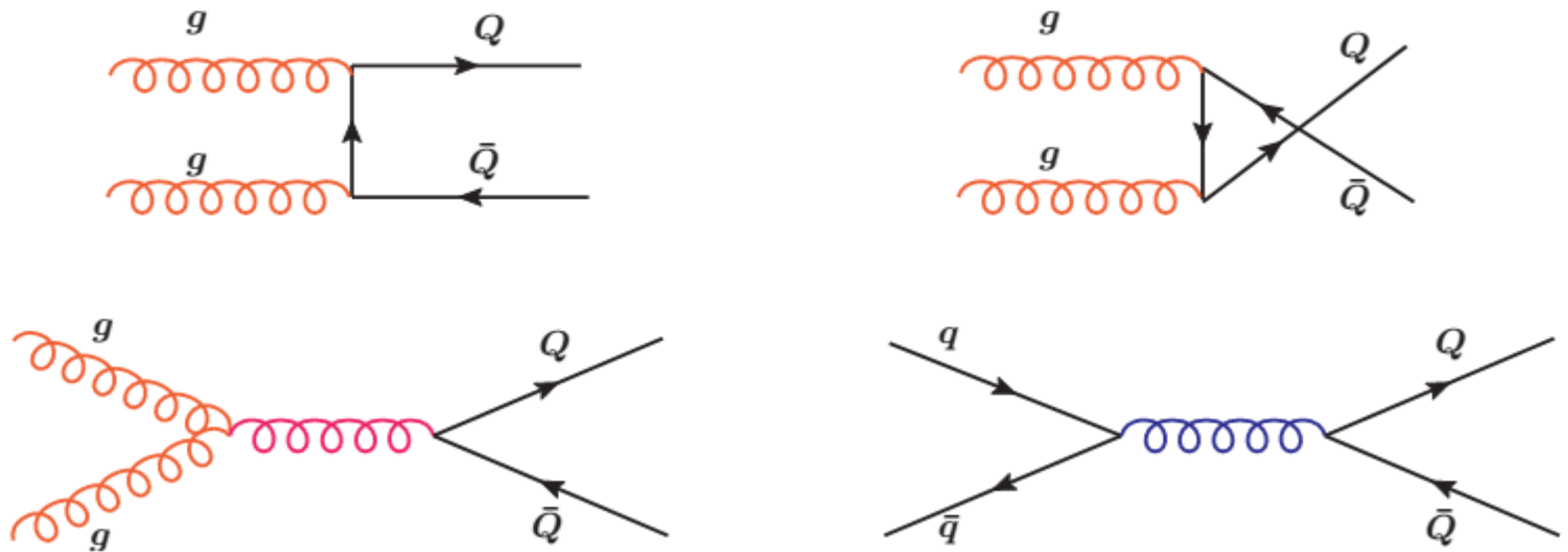


# $pp \rightarrow J/\psi \text{ or } \Upsilon + X$

The cross section for Quarkonium production in CEM is

$$\sigma = \frac{\rho}{9} \int_{2m_Q}^{2m_{Q\bar{q}}} dM \frac{d\hat{\sigma}_{Q\bar{Q}}}{dM}$$

Where  $m_Q = m_c(m_b)$  and  $m_{Q\bar{q}} = m_D(m_B)$  for charmonium (bottomonium)



# $pp \rightarrow J/\psi \text{ or } \Upsilon + X$

Using QCD factorization theorem

$$d\sigma = \frac{\rho}{9} \int dx_a dx_b d^2\mathbf{k}_{\perp a} d^2\mathbf{k}_{\perp b} \left\{ \Phi_g^{\mu\nu}(x_a, \mathbf{k}_{\perp a}) \Phi_{g\mu\nu}(x_b, \mathbf{k}_{\perp b}) d\hat{\sigma}^{gg \rightarrow Q\bar{Q}} + \left[ \Phi^q(x_a, \mathbf{k}_{\perp a}^2) \Phi^{\bar{q}}(x_b, \mathbf{k}_{\perp b}^2) + \Phi^{\bar{q}}(x_a, \mathbf{k}_{\perp a}^2) \Phi^q(x_b, \mathbf{k}_{\perp b}^2) \right] d\hat{\sigma}^{q\bar{q} \rightarrow Q\bar{Q}} \right\}$$

## TMDs Parameterization

TMDs exhibit Gaussian distribution

$$f_1^g(x, \mathbf{k}_{\perp}^2, Q^2) = f_1^g(x, Q^2) \frac{1}{\pi \langle k_{\perp}^2 \rangle} e^{-\mathbf{k}_{\perp}^2 / \langle k_{\perp}^2 \rangle}$$

Evolution in collinear PDFs

No Evolution in  $k_{\perp}$

DGLAP Evolution



$$pp \rightarrow J/\psi \text{ or } \Upsilon + X$$

Model independent theoretical upper bound

$$\frac{\mathbf{k}_\perp^2}{2M_h^2} |h_1^{\perp g}(x, \mathbf{k}_\perp^2)| \leq f_1^g(x, \mathbf{k}_\perp^2)$$

D. Boer et al. PRL 106 132001 (2011)

Assuming linearly polarized gluons also exhibit Gaussian form

$$h_1^{\perp g}(x, \mathbf{k}_\perp^2) = \frac{M_h^2 f_1^g(x, Q^2)}{\pi \langle k_\perp^2 \rangle^2} \frac{2(1-r)}{r} e^{1 - \mathbf{k}_\perp^2 \frac{1}{r \langle k_\perp^2 \rangle}}$$

$$0 < r < 1$$

D. Boer, C. Pisano, PRD 86, 094007 (2012)

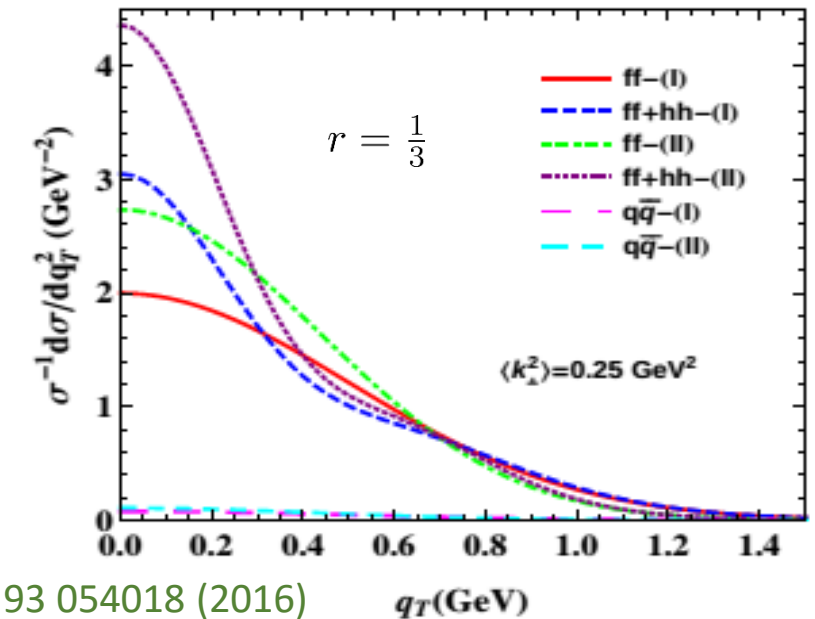
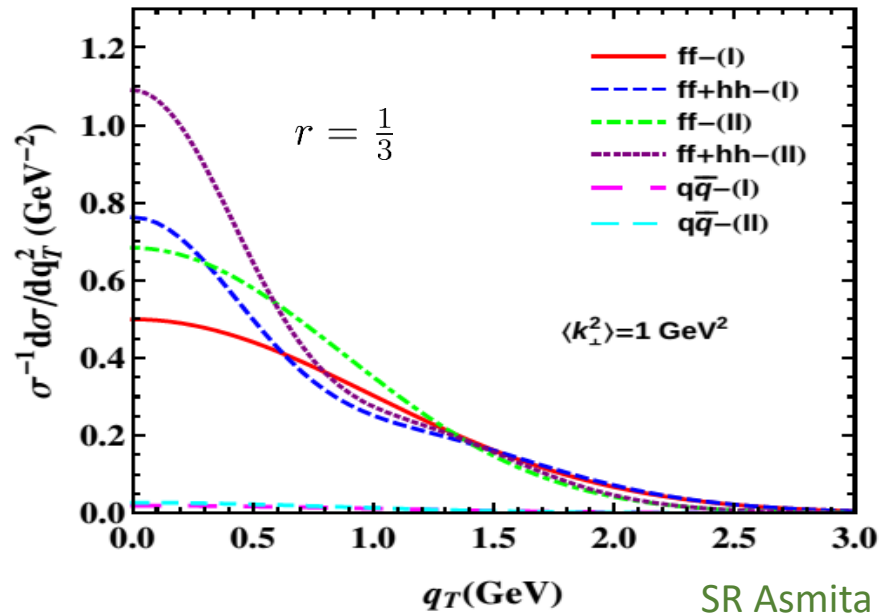
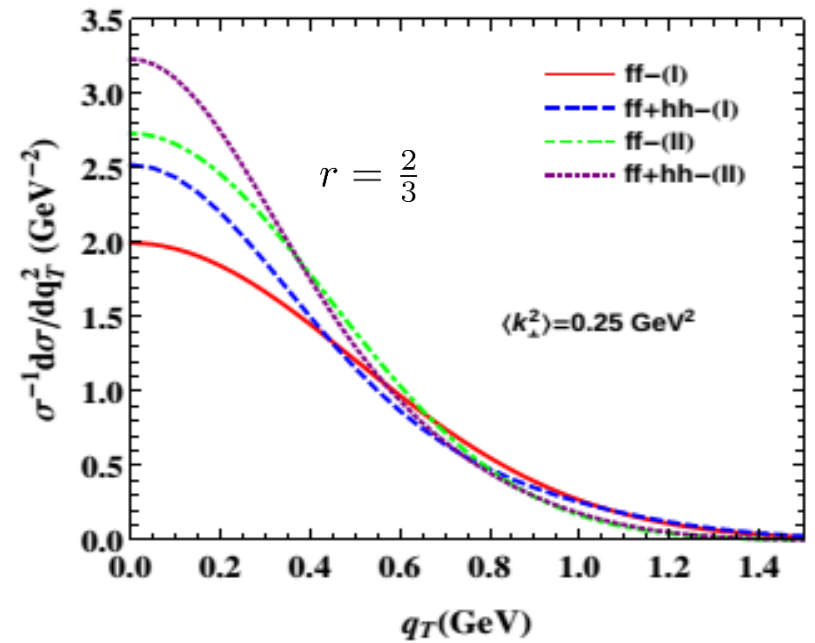
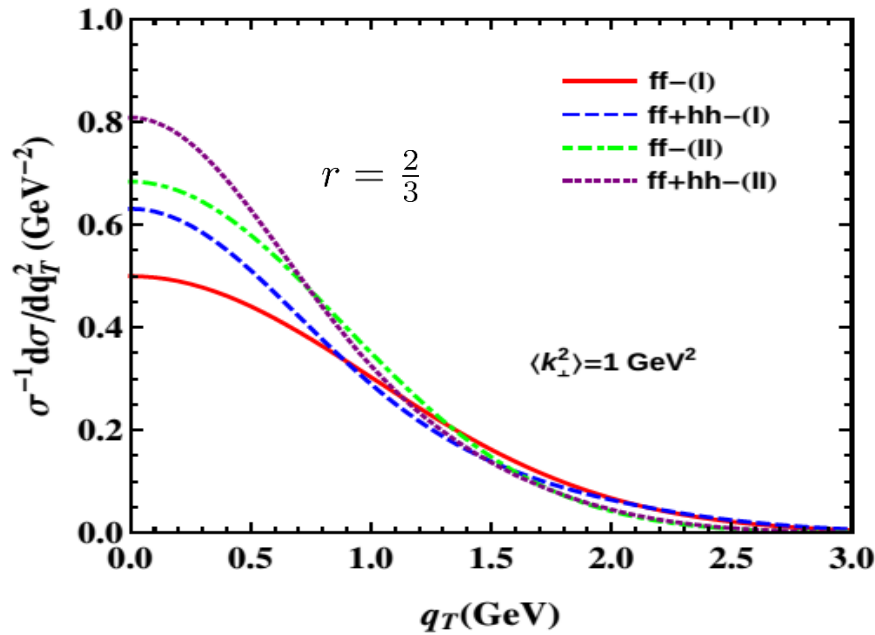
# Differential Cross Section in DGLAP approach

$$\begin{aligned} \frac{d^2\sigma^{ff}}{dydq_T^2} &= \frac{\beta^2\rho}{36s\pi^2} \int_{4m_Q^2}^{4m_{Q\bar{q}}^2} dM^2 \int d\phi_{q_T} \int dk_{\perp a} k_{\perp a} \int d\phi_{k_{\perp a}} e^{-\Delta\beta} \\ &\times \left\{ f_1^g(x_a) f_1^g(x_b) \hat{\sigma}^{gg \rightarrow Q\bar{Q}}(M^2) + \frac{1}{2} \sum_q \left[ f_1^q(x_a) f_1^{\bar{q}}(x_b) + f_1^{\bar{q}}(x_a) f_1^q(x_b) \right] \right. \\ &\quad \left. \times \hat{\sigma}^{q\bar{q} \rightarrow Q\bar{Q}}(M^2) \right\} \end{aligned}$$

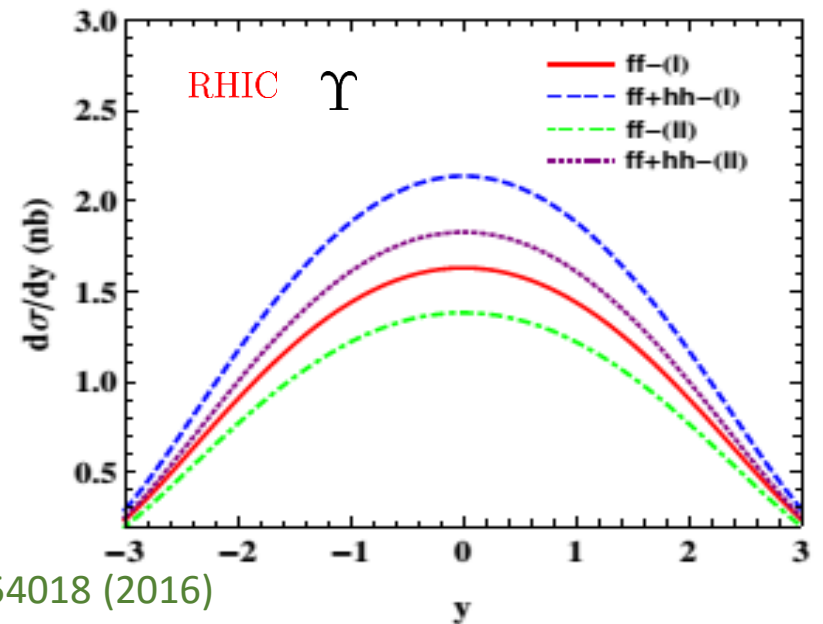
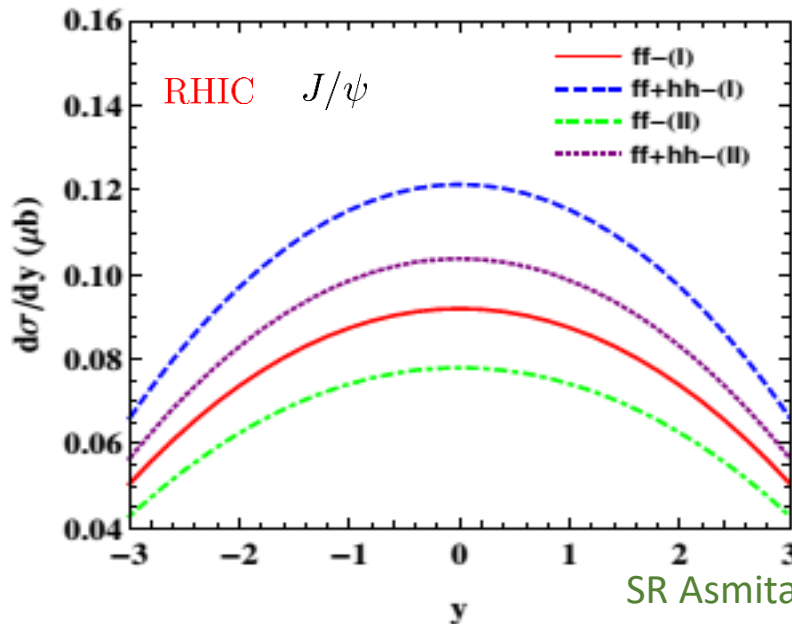
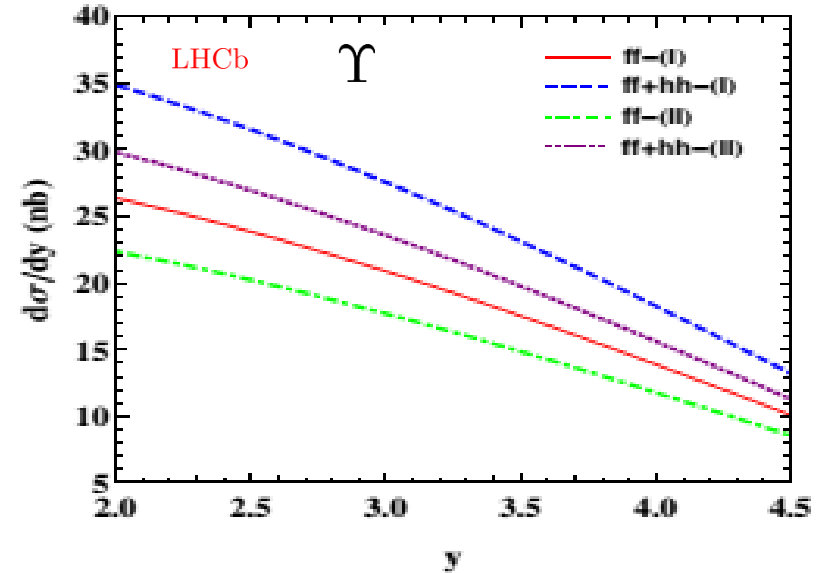
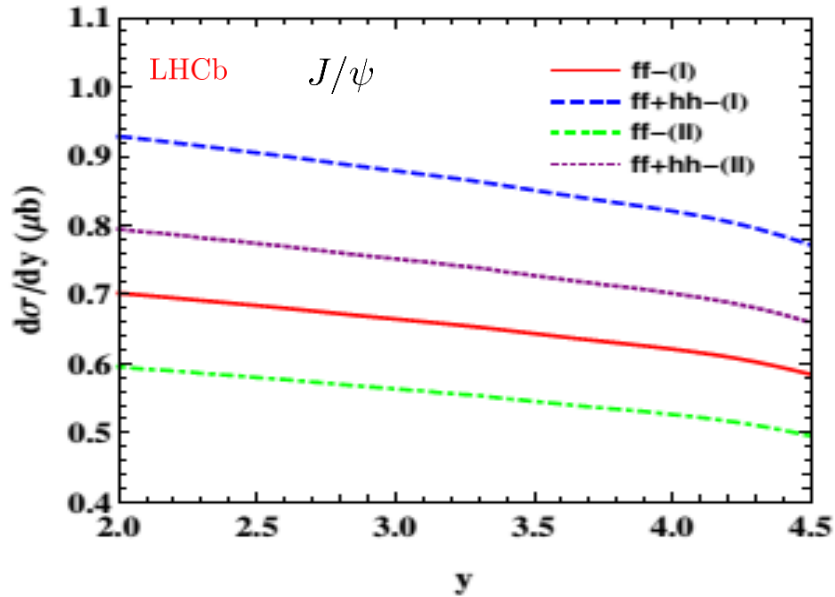
$$\begin{aligned} \frac{d^2\sigma^{hh}}{dydq_T^2} &= \frac{\beta^4\rho(1-r)^2}{18sr^2\pi^2} \int_{4m_Q^2}^{4m_{Q\bar{q}}^2} dM^2 \int d\phi_{q_T} \int dk_{\perp a} k_{\perp a} \int d\phi_{k_{\perp a}} \\ &\times \left[ \frac{1}{2} k_{\perp a}^4 - \frac{1}{2} k_{\perp a}^2 q_T^2 - q_T k_{\perp a}^3 \cos(\phi_{k_{\perp a}} - \phi_{q_T}) + q_T^2 k_{\perp a}^2 \cos^2(\phi_{k_{\perp a}} - \phi_{q_T}) \right] \\ &\times e^{[2 - \frac{\beta}{r} \Delta]} f_1^g(x_a) f_1^g(x_b) \hat{\sigma}^{gg \rightarrow Q\bar{Q}}(M^2) \end{aligned}$$

where  $\Delta = 2k_{\perp a}^2 + q_T^2 - 2q_T k_{\perp a} \cos(\phi_{k_{\perp a}} - \phi_{q_T})$  and  $\beta = \frac{1}{\langle k_{\perp}^2 \rangle}$

# $q_T$ Spectrum in DGLAP Evolution

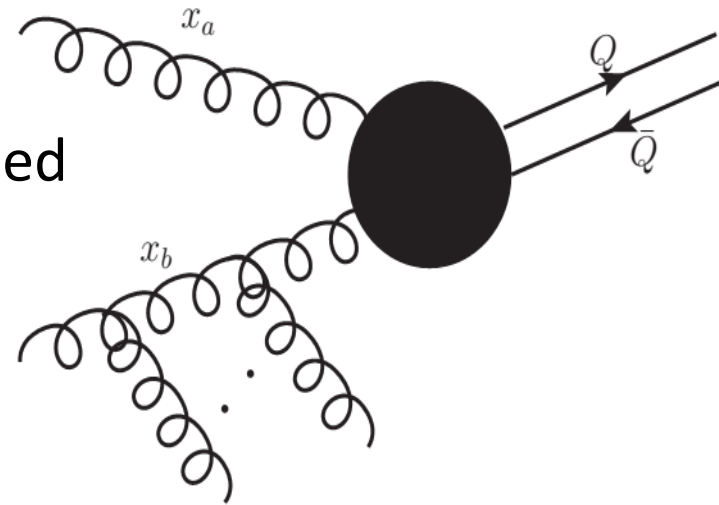


# Rapidity Spectrum in DGLAP Evolution



# Resummation of Sudakov logarithms

The region of low  $q_T$  is strongly influenced by initial state gluon radiation showers



These additional gluon radiation from initial state partons leads to logarithmic corrections for each gluon radiation of the form

$$\alpha_s \log\left(\frac{Q^2}{q_T^2}\right)$$

Collins Soper Sterman (CSS) resummation formalism has been used to resum the large logarithmic terms to all order in  $\alpha_s$

# TMD Evolution

## CSS TMD Evolution

J. Collins, Foundations of Perturbative QCD, 2011

$$f(x, b_{\perp}, Q_f, \zeta) = f(x, b_{\perp}, Q_i, \zeta) R_{pert}(Q_f, Q_i, b_*) R_{NP}(Q_f, Q_i, b_{\perp})$$

$$R_{pert}(Q_f, Q_i, b_*) = \exp \left\{ - \int_{Q_i}^{Q_f} \frac{d\mu}{\mu} \left( A \log \left( \frac{Q_f^2}{\mu^2} \right) + B \right) \right\}$$

$$R_{NP} = \exp \left\{ - \left[ \frac{g_2}{2} \log \frac{Q_f}{2Q_0} + \frac{g_1}{2} \left( 1 + 2g_3 \log \frac{10xx_0}{x_0 + x} \right) \right] b_{\perp}^2 \right\}$$

D. Boer, W. J. den Dunnen, Nucl. Phys. B **886** (2014) 421

where  $Q_i = \frac{2e^{-0.577}}{b_*}$ ,  $Q_f = Q$  and  $\zeta = Q_f^2$

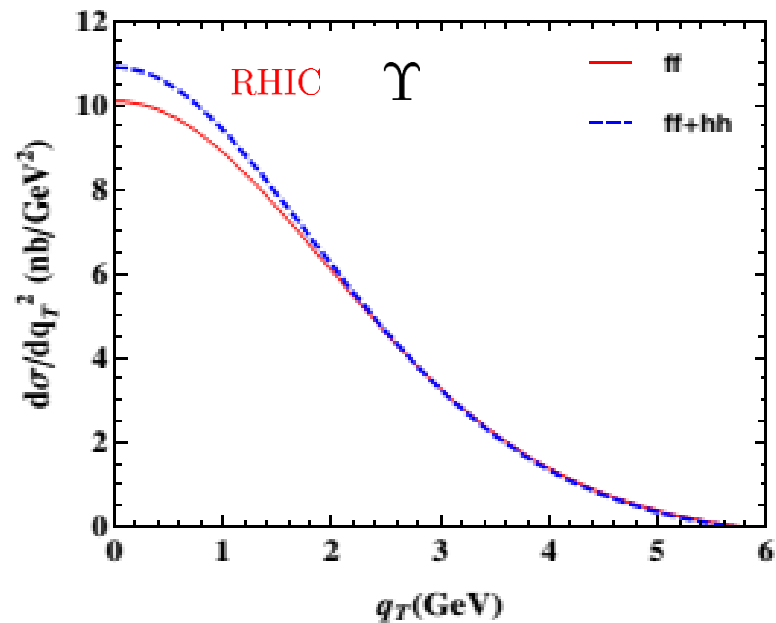
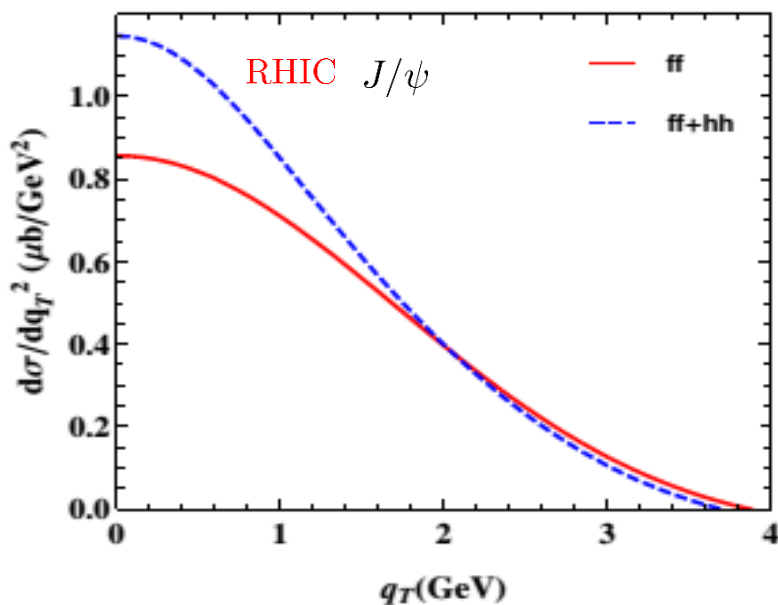
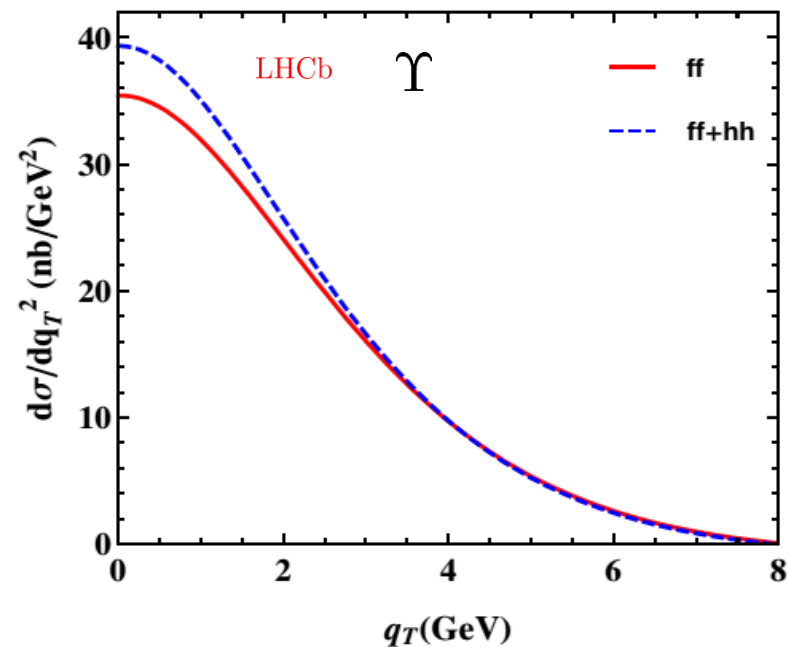
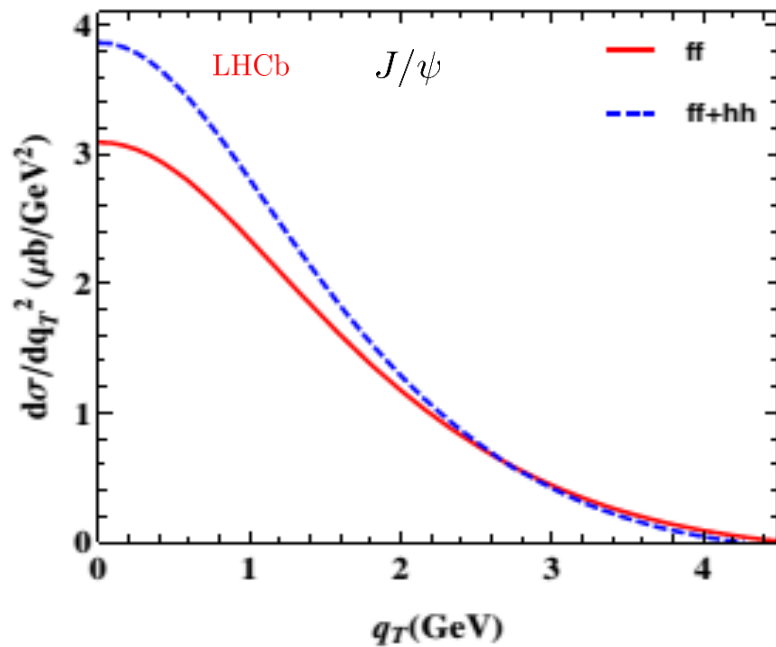
$$b_* = \frac{b_{\perp}}{\sqrt{1 + \left(\frac{b_{\perp}}{b_{\max}}\right)^2}}, \quad A = \sum_{n=1}^{\infty} \left( \frac{\alpha_s(\mu)}{\pi} \right)^n A_n, \quad B = \sum_{n=1}^{\infty} \left( \frac{\alpha_s(\mu)}{\pi} \right)^n B_n$$

# Differential Cross Section in TMD Evolution approach

$$\frac{d^2\sigma^{ff}}{dydq_T^2} = \frac{\rho}{36s} \int_{4m_Q^2}^{4m_{Q\bar{q}}^2} dM^2 \int_0^\infty b_\perp db_\perp J_0(q_T b_\perp) f_1^g(x_a, c/b_*) f_1^g(x_b, c/b_*) \hat{\sigma}^{gg \rightarrow Q\bar{Q}}(M^2) \exp\left\{-2 \int_{c/b_*}^Q \frac{d\mu}{\mu} \left(A \log\left(\frac{Q^2}{\mu^2}\right) + B\right)\right\} \exp\left\{-\left[0.184 \log \frac{Q}{2Q_0} + 0.332\right] b_\perp^2\right\}$$

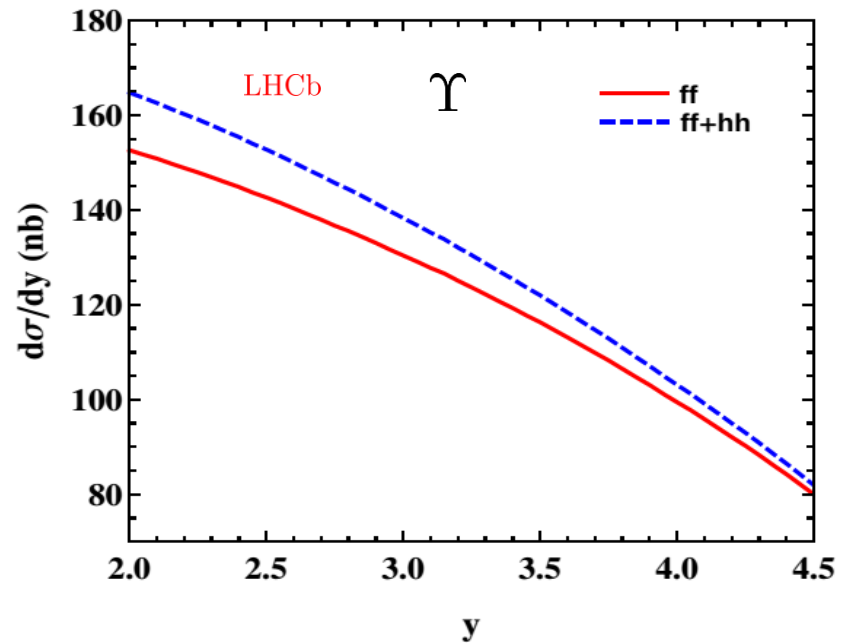
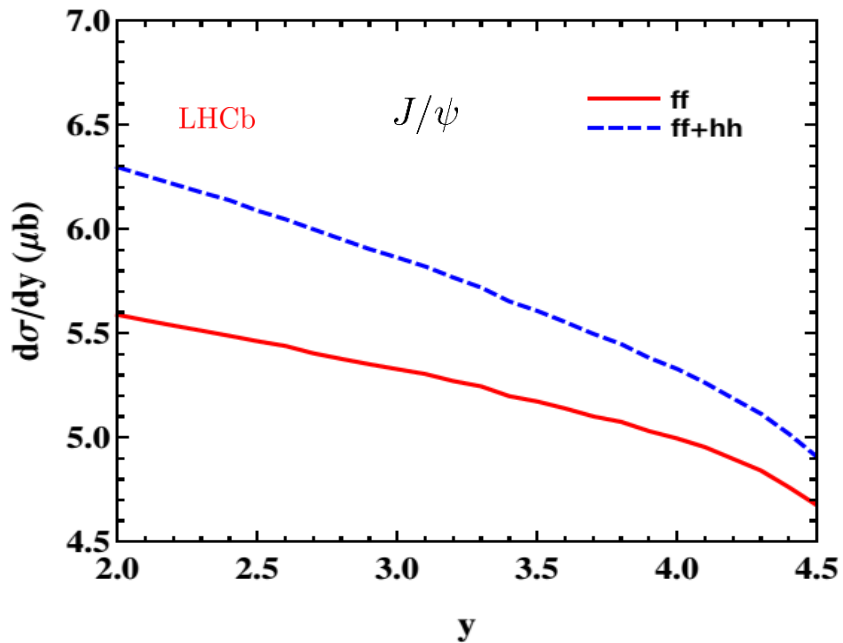
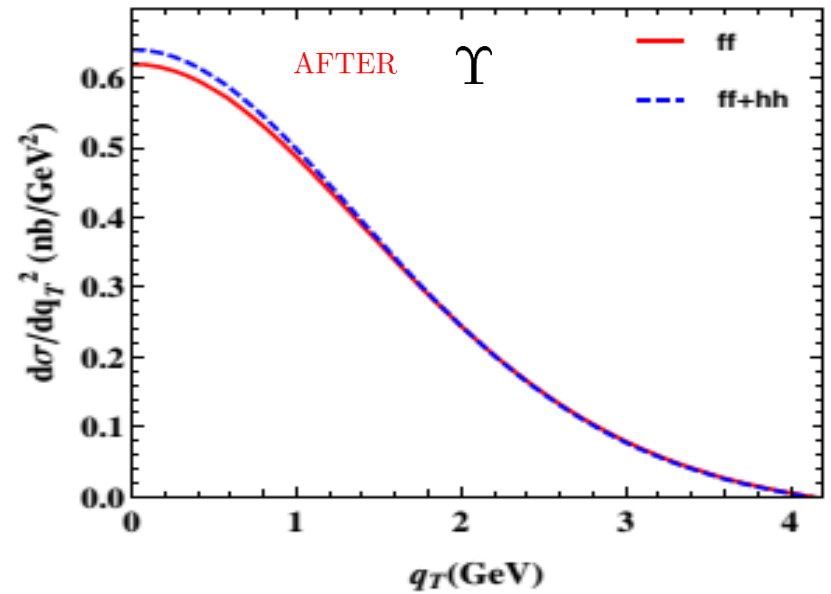
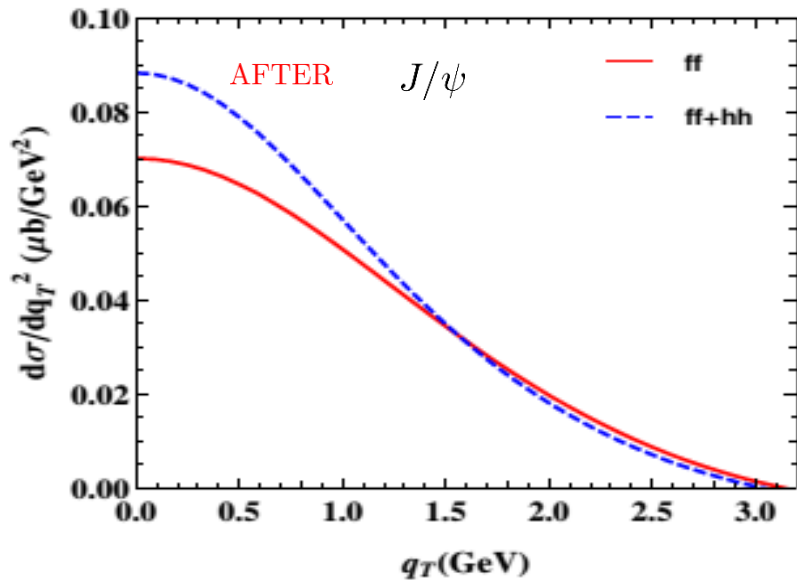
$$\frac{d^2\sigma^{hh}}{dydq_T^2} = \frac{\rho C_A^2}{36s\pi^2} \int_{4m_Q^2}^{4m_{Q\bar{q}}^2} dM^2 \int_0^\infty b_\perp db_\perp J_0(q_T b_\perp) \alpha_s^2(c/b_*) \hat{\sigma}^{gg \rightarrow Q\bar{Q}}(M^2) \int_{x_a}^1 \frac{dx_1}{x_1} \left(\frac{x_1}{x_a} - 1\right) f_1^g(x_1, c/b_*) \int_{x_b}^1 \frac{dx_2}{x_2} \left(\frac{x_2}{x_b} - 1\right) f_1^g(x_2, c/b_*) \exp\left\{-2 \int_{c/b_*}^Q \frac{d\mu}{\mu} \left(A \log\left(\frac{Q^2}{\mu^2}\right) + B\right)\right\} \exp\left\{-\left[0.184 \log \frac{Q}{2Q_0} + 0.332\right] b_\perp^2\right\}$$

# $q_T$ Spectrum in TMD Evolution





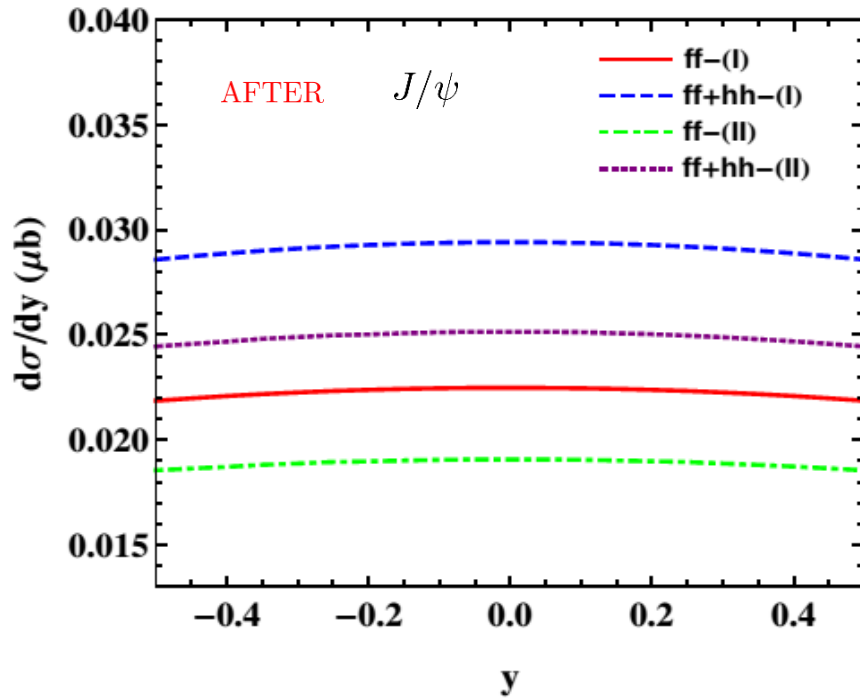
# Results in TMD Evolution



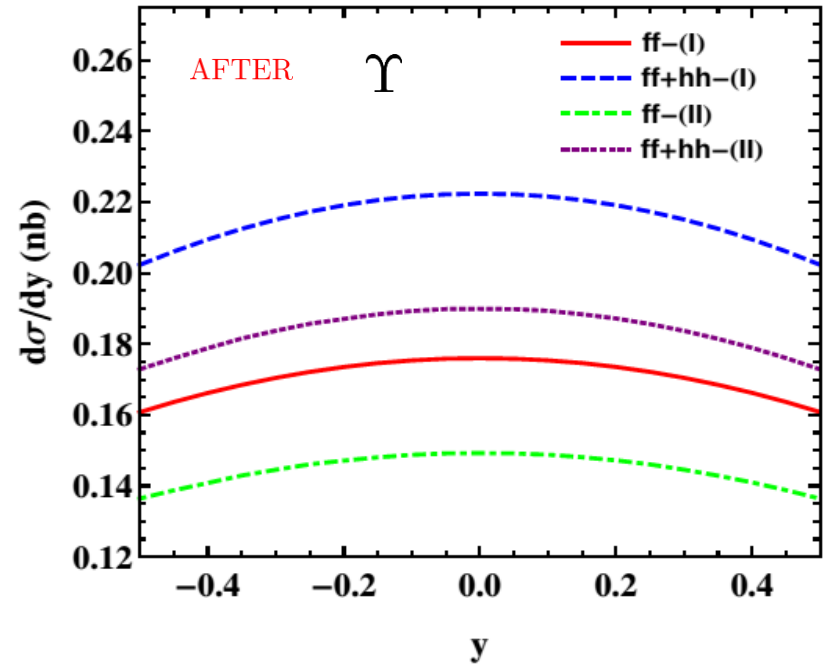
# Conclusion

- ★ The  $q_T$  and  $y$  distributions of quarkonium can be modulated by the presence of linearly polarized gluons in unpolarized proton proton collision.
- ★ Hence, the production of quarkonium is a promising process to probe not only the  $h_1^{\perp g}$  but also unpolarized TMD pdf  $f_1^g$ .
- ★ It would be interesting to include the linearly polarized gluon contribution in the cross section to fit the experimental data to the extent of reasonable accuracy.

Thank you



$$r = \frac{1}{3}$$



$$\langle k_{\perp}^2 \rangle = 1 \text{ GeV}^2$$

