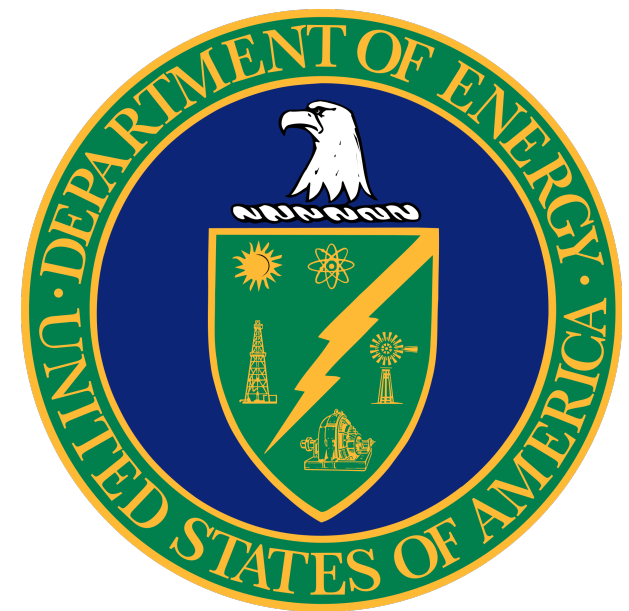
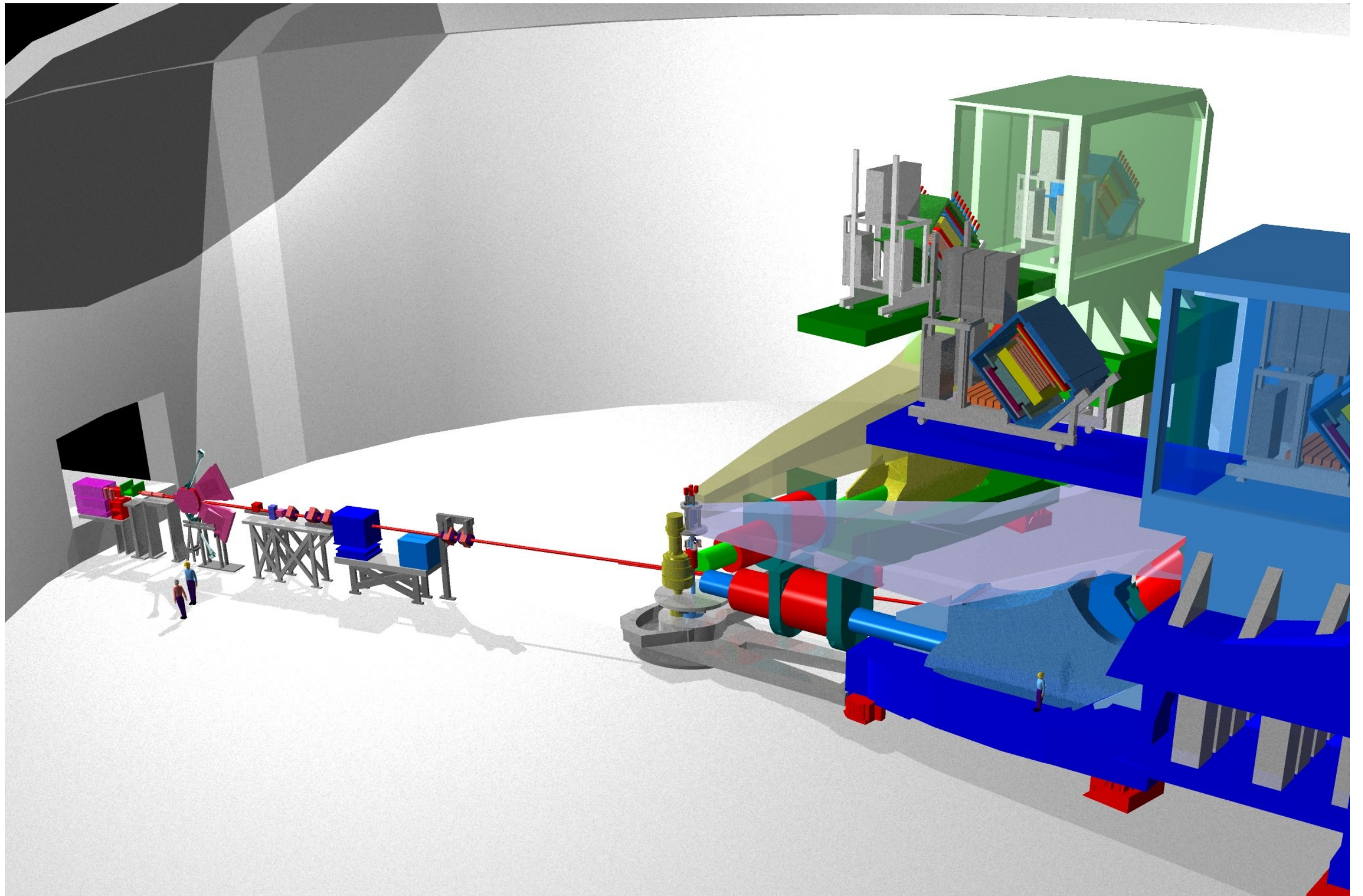


Compton Polarimetry in Hall A

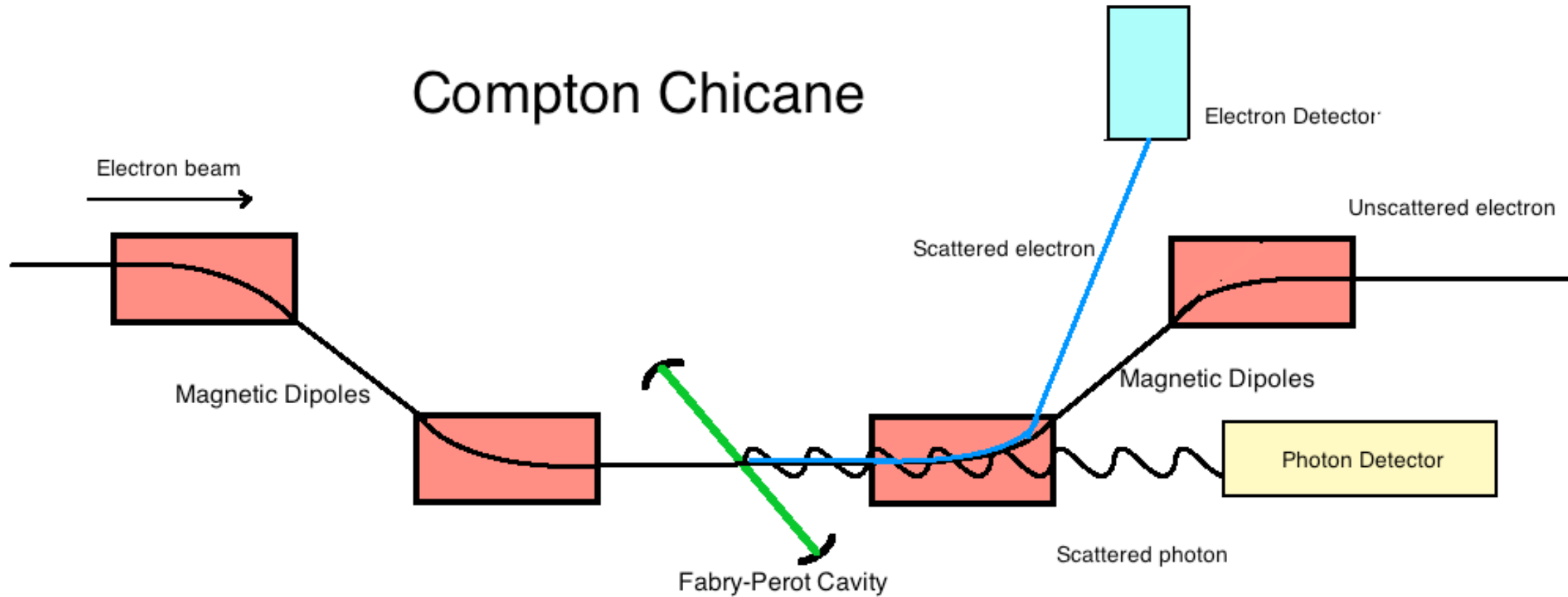
Alexa Johnson



Hall A



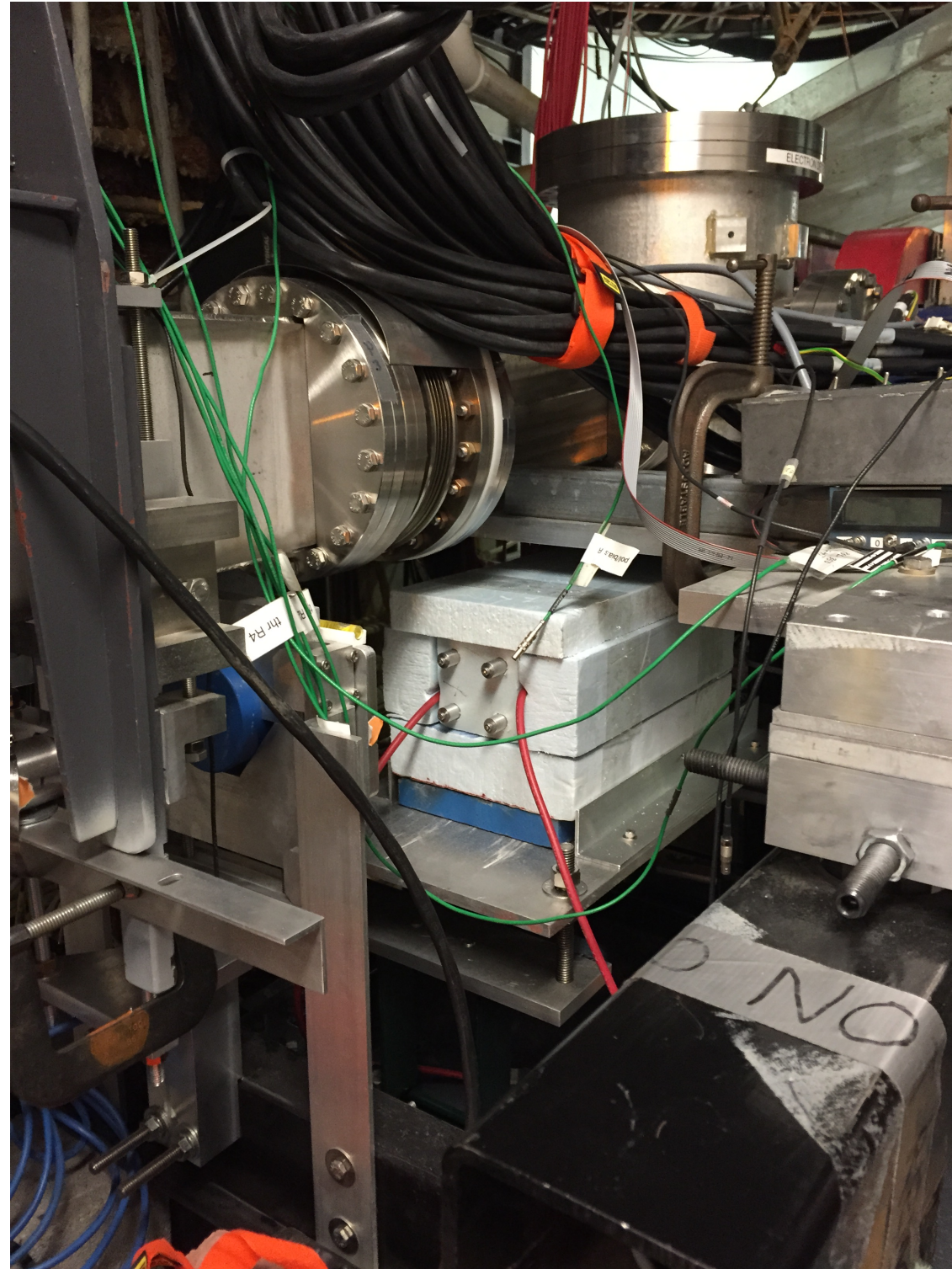
Compton Chicane



- Polarized electrons scatter off polarized photons in a Fabry-Perot cavity
- Spin-dependent Compton scattering

Photon Detector

An array of 4, 3cm x 3cm x 20cm PbWO₄ crystals



Monte Carlo

The polarimeter's response can be modeled using an accurate **Geant4** simulation, containing all relevant, physical beamline components :

- Photons corresponding to Compton-scattered photons are generated in the simulation, and sent through the PbWO4 crystal array.
- The energy deposited in the crystal is recorded.

Monte Carlo

The scattering cross section is:

$$\frac{d\sigma}{d\rho_{in}} = a \left(\frac{(\rho_{in}(1-a))^2}{1-\rho_{in}(1-a)} + \left(\frac{1-\rho_{in}(1+a)}{1-\rho_{in}(1-a)} \right)^2 + 1 \right)$$

$$\rho_{in} = \frac{E_{in}}{E_{max}}$$

$$a = \frac{1}{1 + \frac{4E_{ph}E_{beam}}{M_e^2}}$$

We define the spin-dependent cross sections as:

$$\left(\frac{d\sigma}{d\rho_{in}} \right)_{\pm} = \frac{d\sigma}{d\rho_{in}} (1 \pm A)$$

E_{in} = Scattered photon energy

M_e = Electron mass

E_{ph} = Laser photon energy

E_{max} = Energy of Compton edge

E_{beam} = Electron beam energy

Where:

$$A = \frac{a}{\frac{d\sigma}{d\rho_{in}}} (1 - \rho_{in}(1+a)) \left(1 - \frac{1}{(1 - \rho_{in}(1-a))^2} \right)$$

Monte Carlo

Experimentally, we do not have access to E_{in} ,
and see only the signal of an event in the
detector, defined both aligned and anti-aligned
spin events as:

$$S_{\pm}(\rho_{dep}) = \int d\rho_{in} R(\rho_{in}, \rho_{dep}) \frac{d\sigma}{d\rho_{in\pm}}$$

Define the signal dependent
asymmetry as:

$$A(\rho_{dep}) = \frac{S_{+}(\rho_{dep}) - S_{-}(\rho_{dep})}{S_{+}(\rho_{dep}) + S_{-}(\rho_{dep})}$$

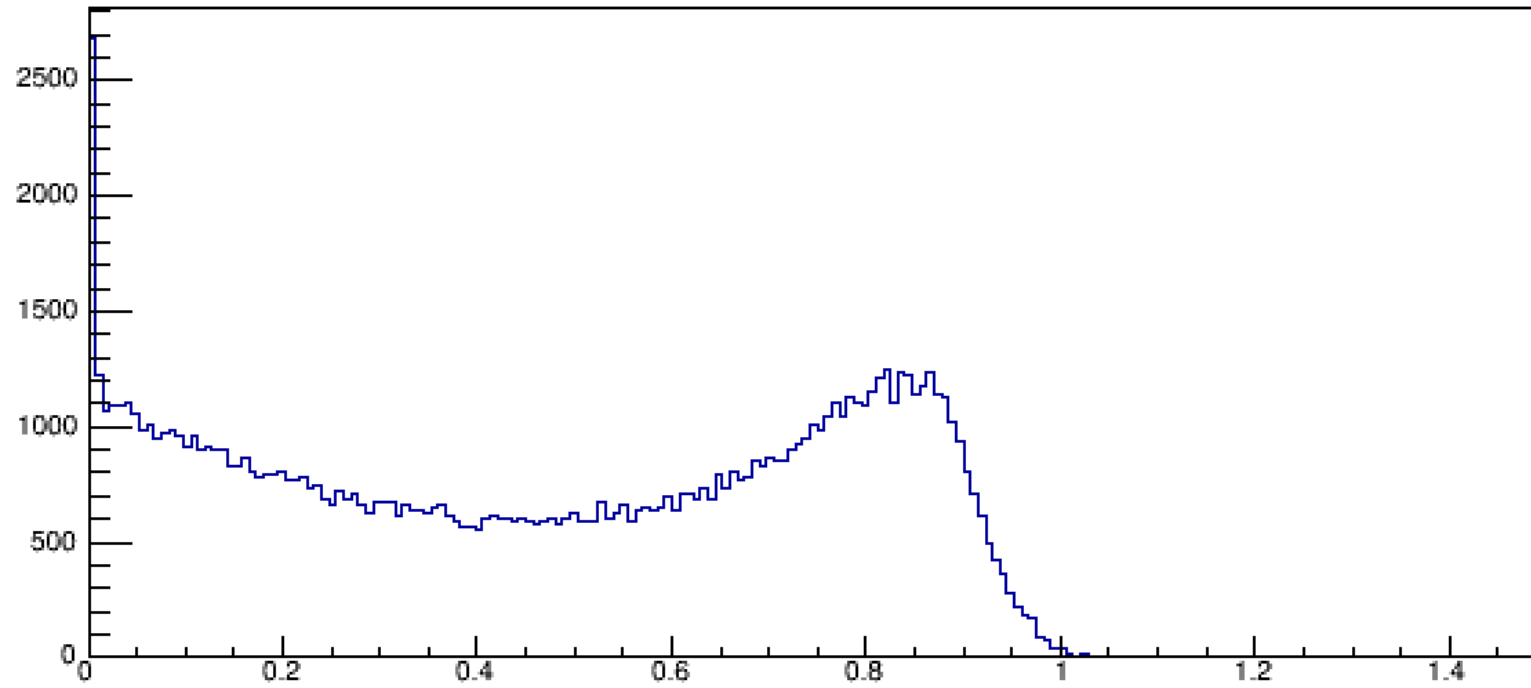
Compare to what we can measure
experimentally:

$$A_{Exp} \equiv \frac{S^{+} - S^{-}}{S^{+} + S^{-}} = P_e P_{\gamma} A$$

$R(\rho_{in}, \rho_{dep})$ is the crystal's resolution function

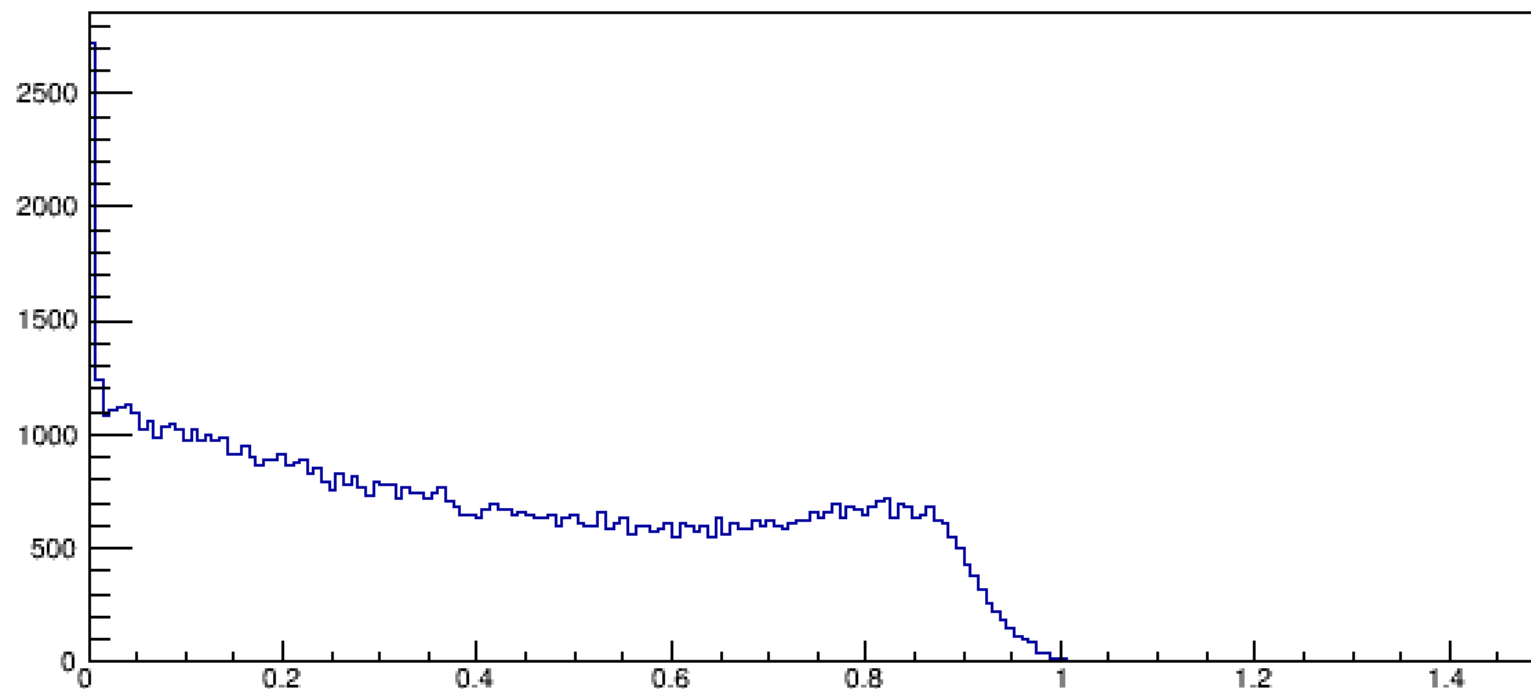
Monte Carlo

Compton spectrum for aligned spin events



$$S_+(\rho_{dep}) \text{ vs. } \rho_{dep}$$

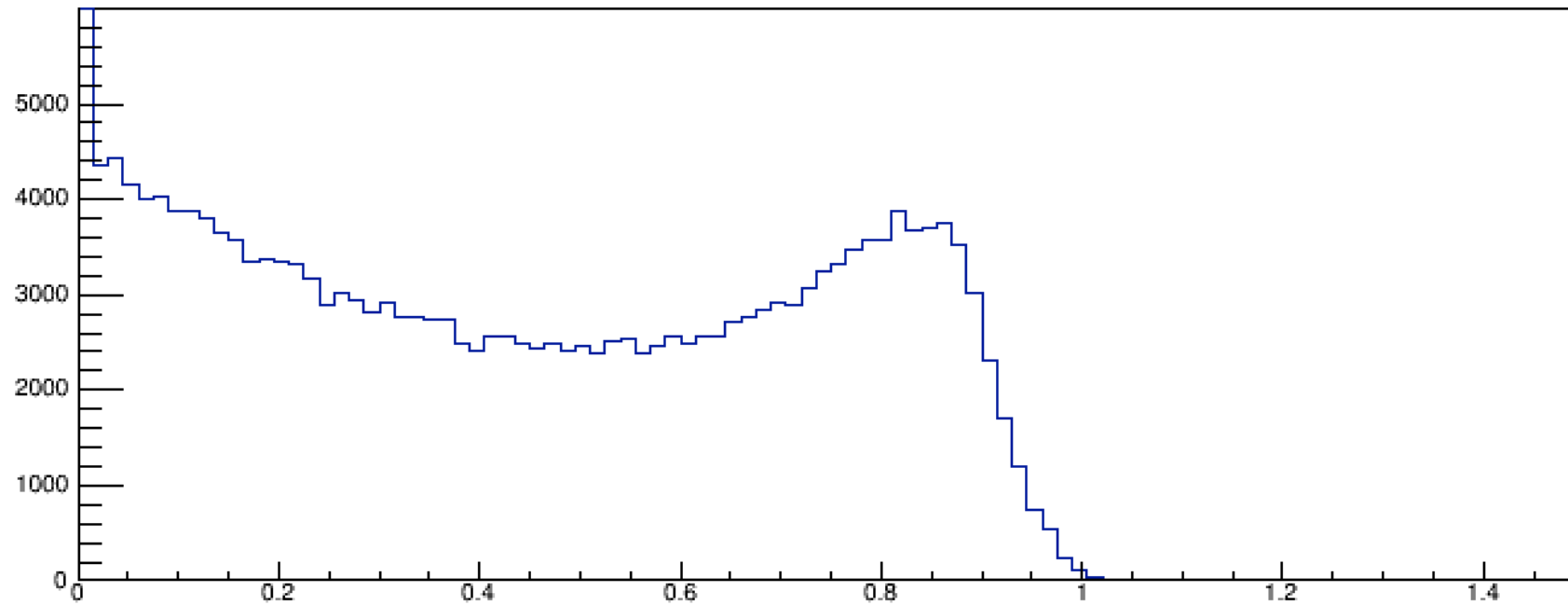
Compton spectrum for anti-aligned spin events



$$S_-(\rho_{dep}) \text{ vs. } \rho_{dep}$$

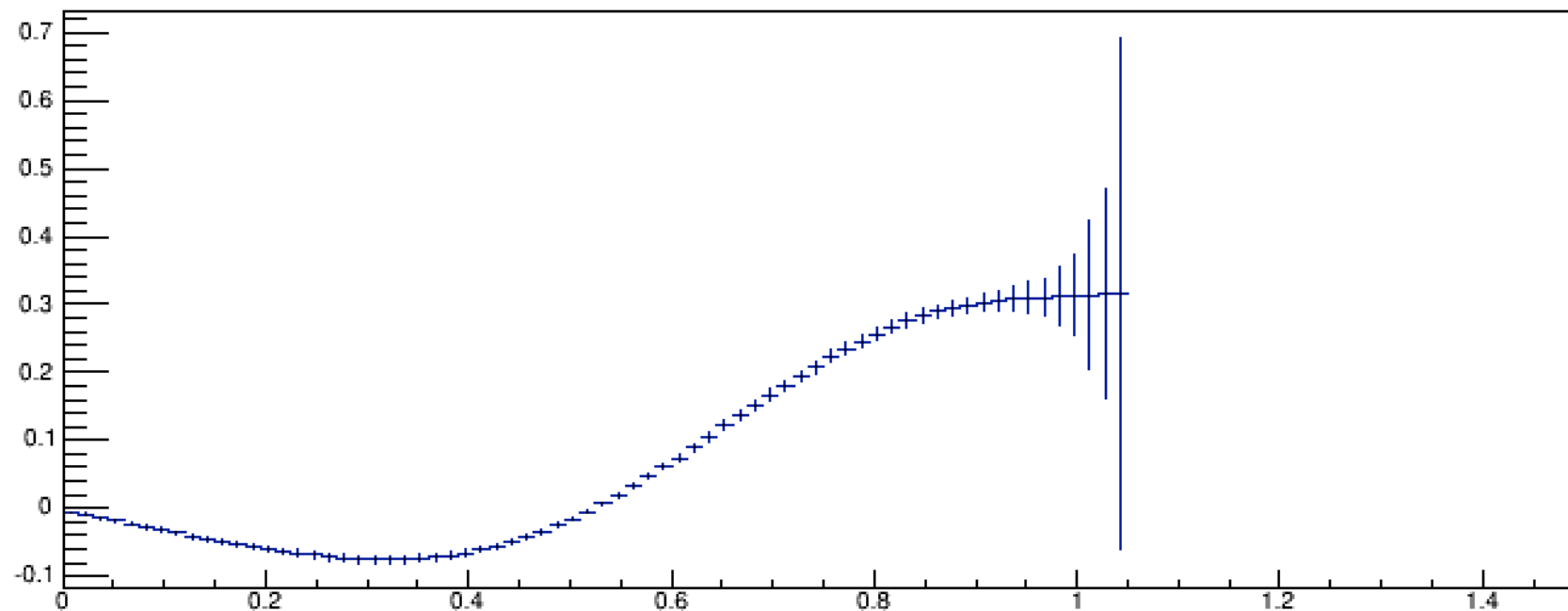
Monte Carlo

Unpolarized Compton Spectrum



$$S_+(\rho_{dep}) + S_-(\rho_{dep})$$

Asymmetry for beam energy 11GeV and shield thickness 3mm



$$A(\rho_{dep}) = \frac{S_+(\rho_{dep}) - S_-(\rho_{dep})}{S_+(\rho_{dep}) + S_-(\rho_{dep})}$$

Measurements

Compton PMT signal leads to FADC, where three measurements are recorded:

My talk



1. All individual PMT events
2. Integrate PMT signal with **no** threshold
3. Integrate PMT signal with **high** threshold

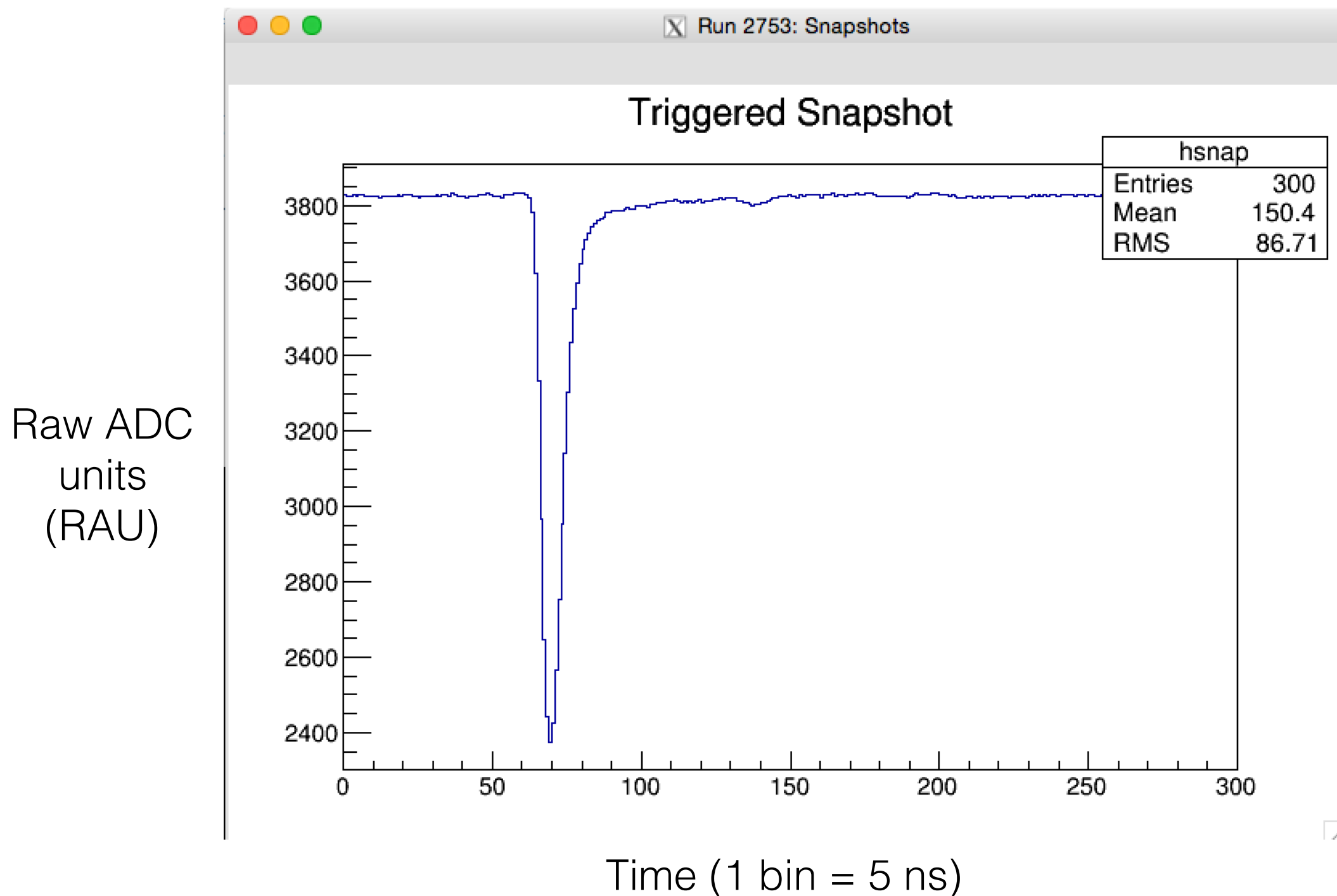


Larisa's talk

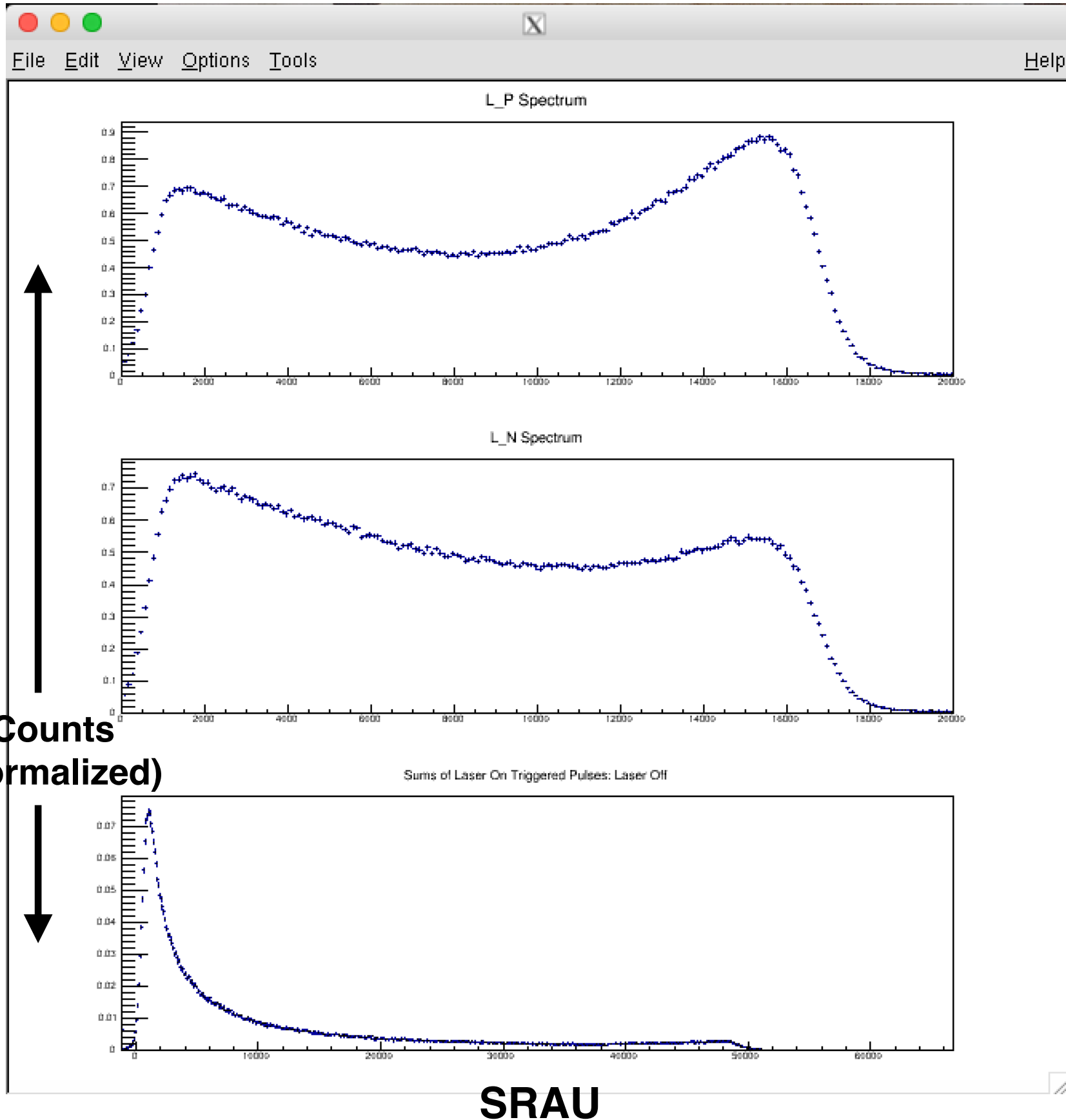
***Each method can be used to determine the electron beam polarization percentage.

Raw Data

PMT signal is integrated for each event in the run. (~200,000)



Making Histograms

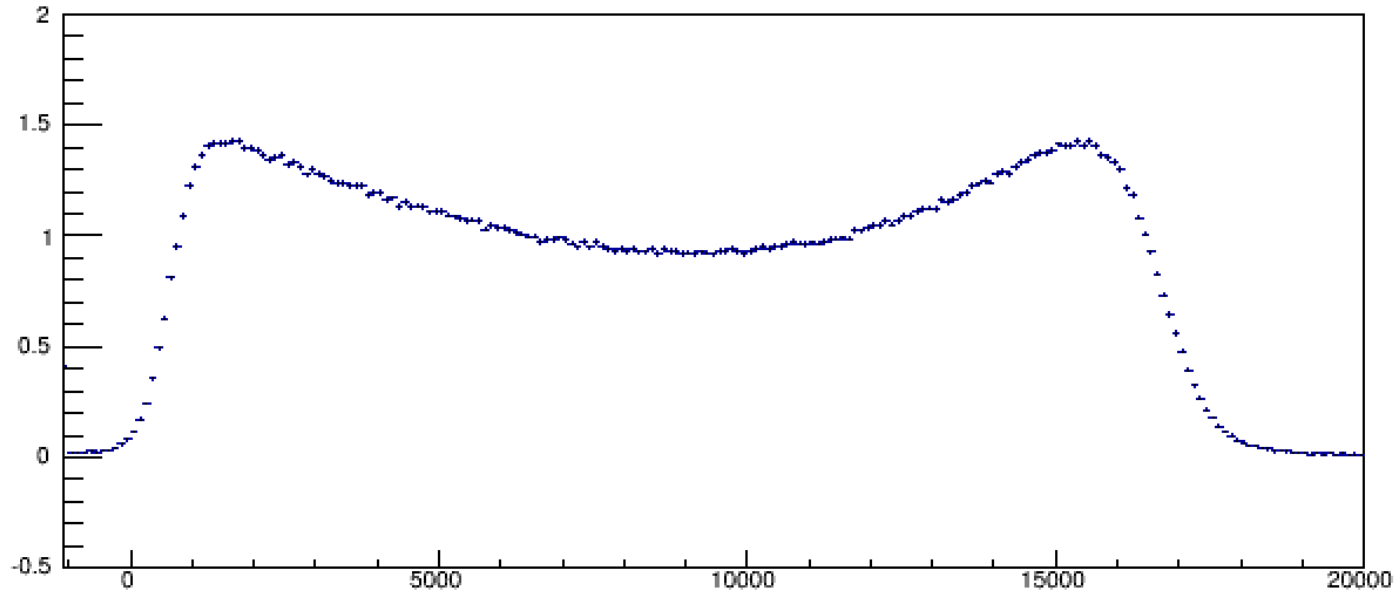


The integrated signal values then fill one of three histograms:

1. Laser **on**, beam helicity **positive**
L_P
2. Laser **on**, beam helicity **negative**
L_N
3. Laser **off**

More Histograms...

Unpolarized Compton Spectrum



From the previous 3 histograms, create 2 more:

1. Unpolarized Compton spectrum

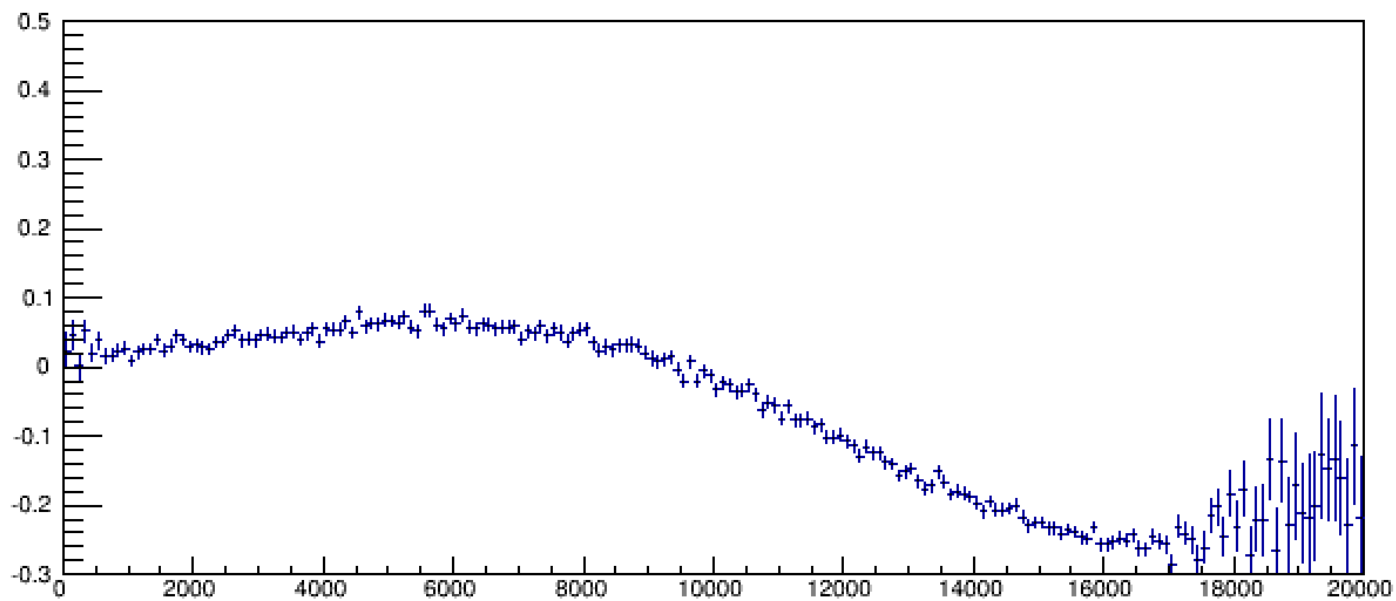
$$\mathbf{L_P + L_N - 2*Off}$$

2. Asymmetry Spectrum

$$\frac{(\mathbf{L_P - L_N})}{(\mathbf{L_P + L_N - 2*Off})}$$

Counts

Compton Asymmetry Spectrum

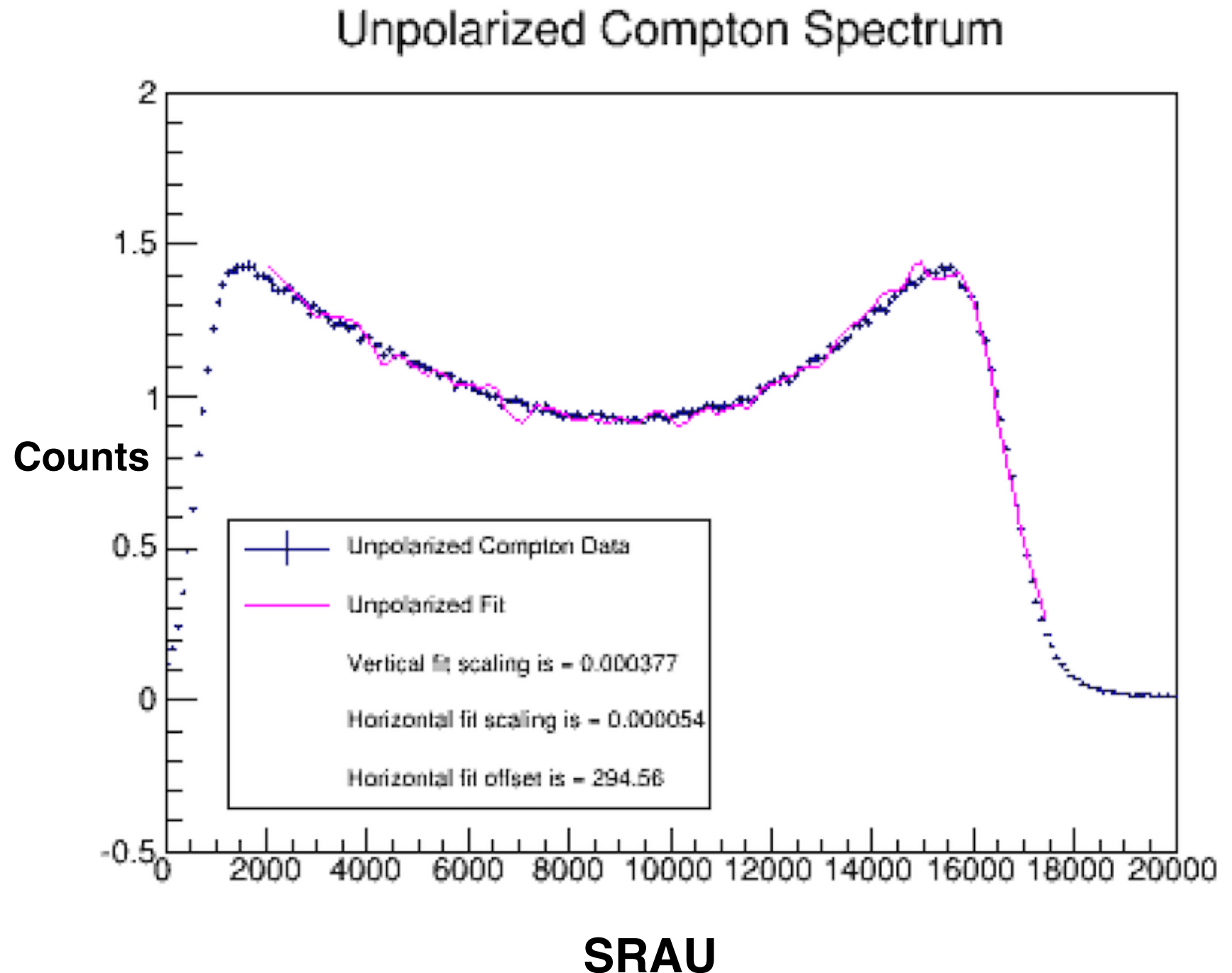


SRAU

Fitting Data...

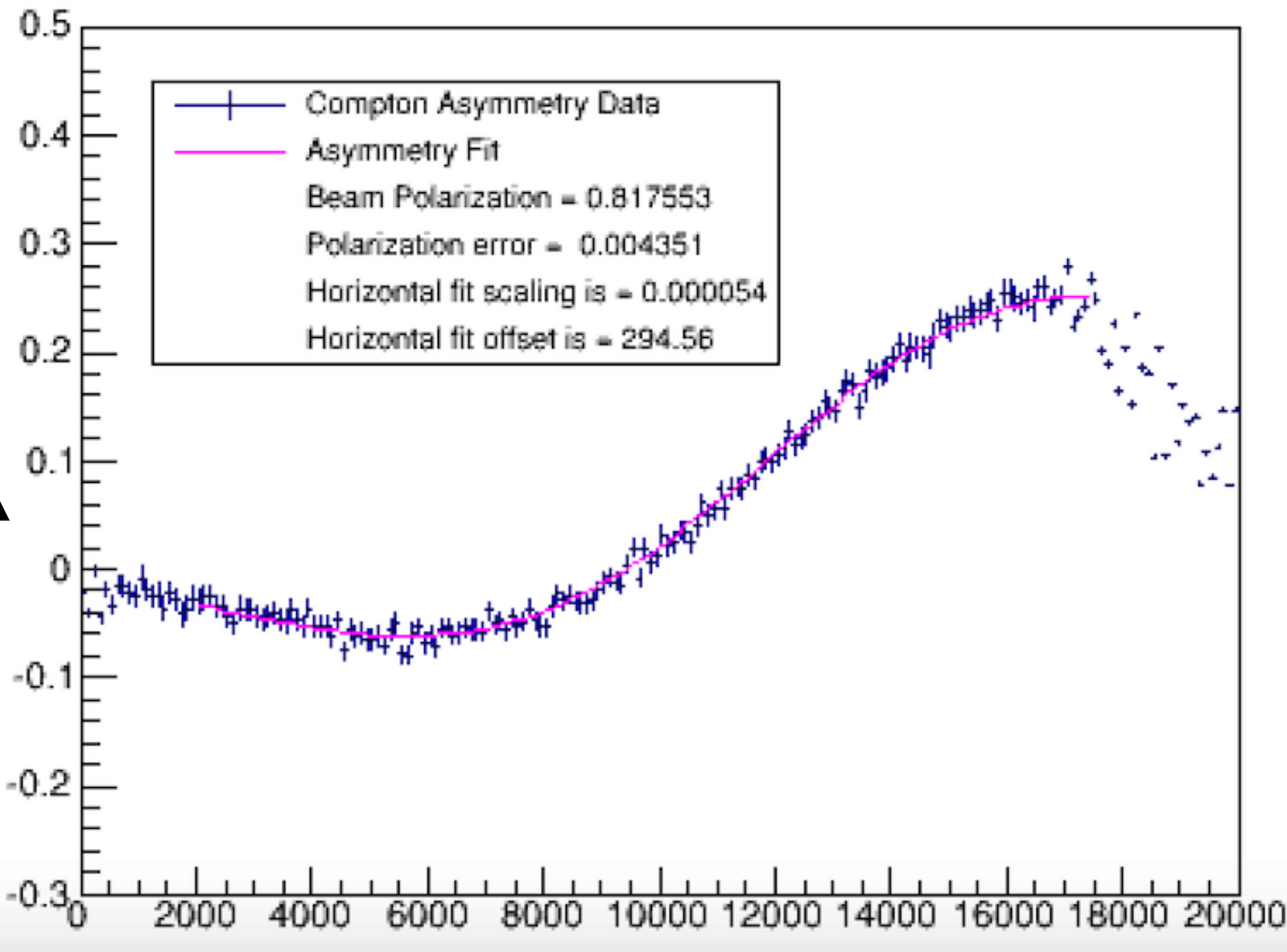
Fit the unpolarized Compton spectrum data to the Monte Carlo with 3 parameters:

1. Vertical scaling
2. Horizontal scaling (proportional to the PMT gain)
3. Horizontal offset



Extracting Beam Polarization

Compton Asymmetry Spectrum



Fit asymmetry spectrum to Monte Carlo, again using 3 parameters, but fixing 2:

Fixed (from unpolarized fit):

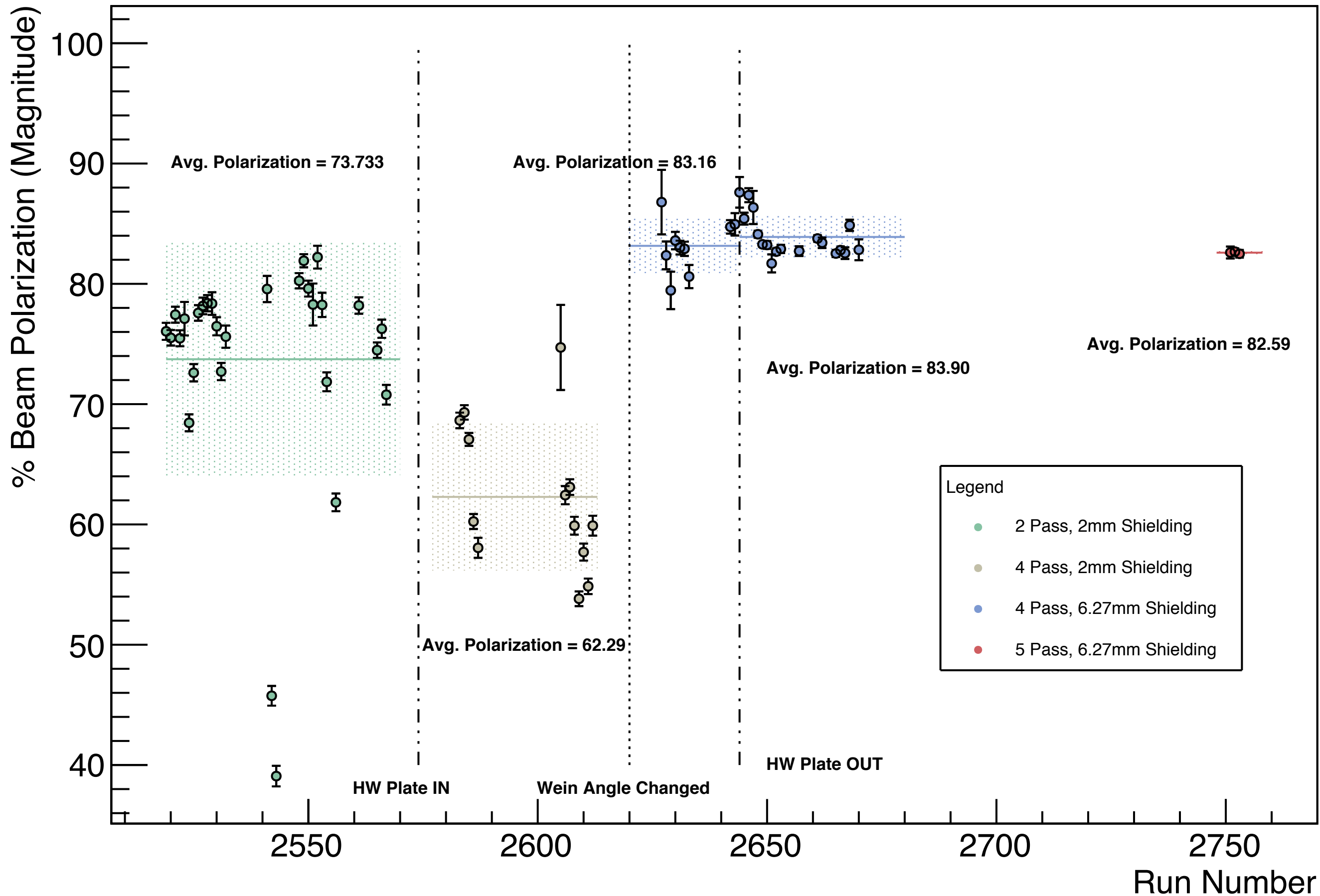
1. Horizontal Scaling
2. Horizontal offset

Free:

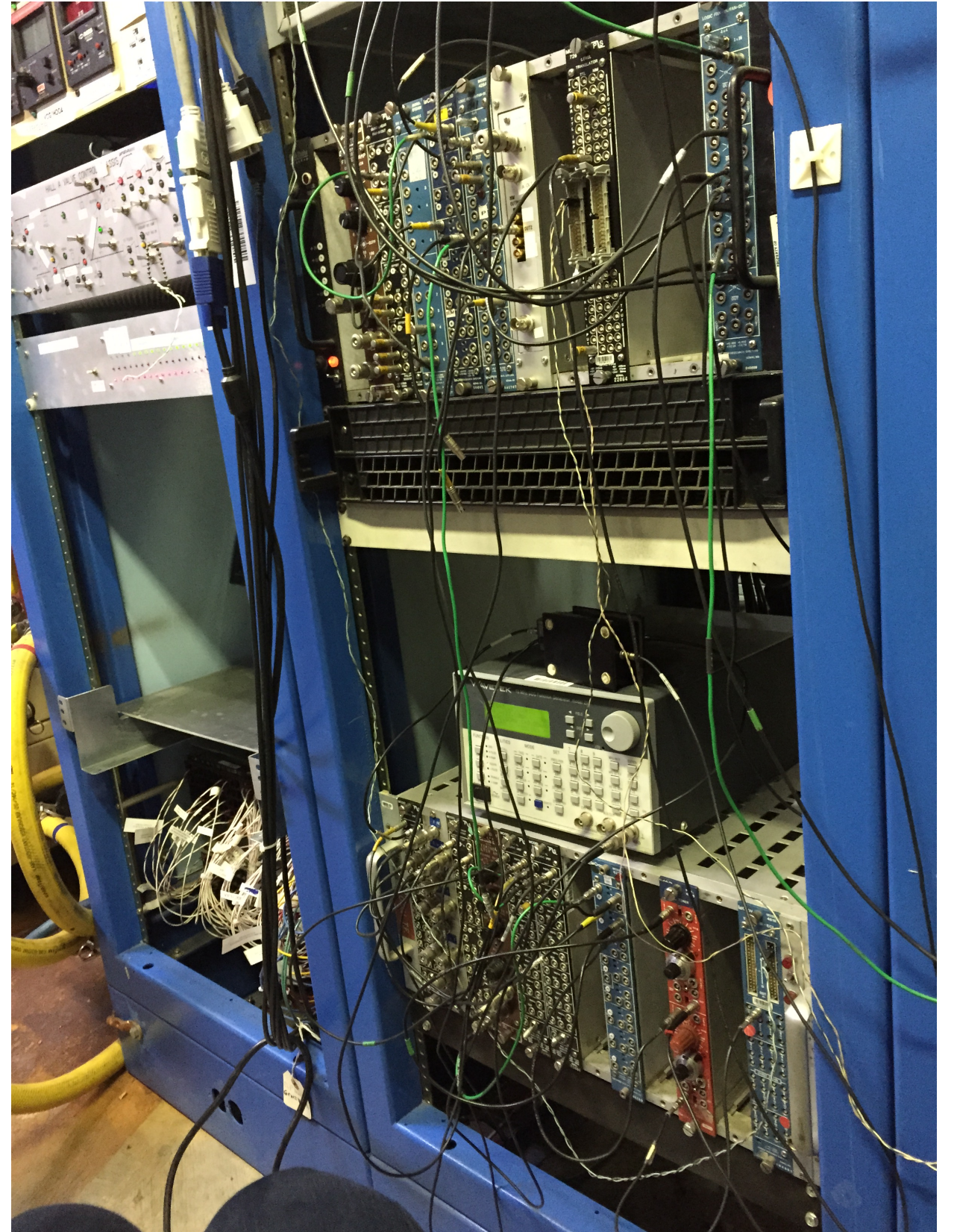
3. Vertical Scaling

Parameter 3 is just the **Beam Polarization**

Beam Polarization vs Run Number



Thank you!



And special thanks to:

Gregg Franklin, Juan Carlos Cornejo, Larisa Thorne, Abel Sun, and Dave Gaskell