## Status of TMD extractions

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Hadronic physics with lepton and hadron beams

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#### Outline of the talk

1) Transverse-Momentum-Dependent distributions (TMDs)

2) formalism

2) extractions : unpolarized

3) extractions : polarized (Sivers)

4) gluon TMDs

5) spin-1 TMDs



#### TMDs

References (intro and reviews) :

- "The 3D structure of the nucleon" EPJ A (2016) 52
- J.C. Collins "Foundations of perturbative QCD"
- P.J. Mulders' lecture notes
- TMD collaboration summer school
- A. Bacchetta's lecture notes



### The hadronic landscape

Manifestation of hadron structure in scattering processes





## The hadronic landscape

Manifestation of hadron structure in scattering processes



Nature is "smooth" : understand the link between TMDs & PDFs



## quark TMD PDFs

#### $\Phi_{ij}(k,P;S) \quad ) \sim \text{F.T.} \ \langle PS \mid \bar{\psi}_j(0) \ U_{[0,\xi]} \ \psi_i(\xi) \ |PS' \rangle_{|_{LF}}$



extraction of a **quark not** collinear with the proton



#### $\Phi_{ij}(k,P;S) \quad ) \sim \text{F.T. } \langle PS \mid \bar{\psi}_j(0) \ U_{[0,\xi]} \ \psi_i(\xi) \ |PS' \rangle_{|_{LF}}$



**bold :** also collinear red : time-reversal odd (universality properties)



extraction of a **quark not** collinear with the proton

encode all the possible **spin-spin** and **spin-orbit correlation** between the proton and its constituents



## Status of TMD phenomenology

Theory, data, fits : we are in a position to start validating the formalism



see, e.g, Bacchetta, Radici, arXiv:1107.5755 Anselmino, Boglione, Melis, PRD86 (12) Echevarria, Idilbi, Kang, Vitev, PRD 89 (14) Anselmino, Boglione, D'Alesio, Murgia, Prokudin, arXiv: 1612.06413 Anselmino et al., PRD87 (13) Kang et al. arXiv:1505.05589

Lu, Ma, Schmidt, arXiv:0912.2031 Lefky, Prokudin arXiv:1411.0580 Barone, Boglione, Gonzalez, Melis, arXiv:1502.04214



## Status of TMD phenomenology

Theory, data, fits : we are in a position to start validating the formalism



Limited data, theory, fits

see, e.g, Bacchetta, Radici, arXiv:1107.5755 Anselmino, Boglione, Melis, PRD86 (12) Echevarria, Idilbi, Kang, Vitev, PRD 89 (14) Anselmino, Boglione, D'Alesio, Murgia, Prokudin, arXiv: 1612.06413 Anselmino et al., PRD87 (13) Kang et al. arXiv:1505.05589

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## The frontier

. . .

Nucleon tomography in momentum space: to understand how hadrons are built in terms of the elementary degrees of freedom of QCD

#### High-energy phenomenology:

to improve our understanding of high-energy scattering experiments and their potential to explore BSM physics

#### A selection of open questions (formalism) :

1) How well do we understand collinear and TMD factorization ?

2) How (well) can we match collinear and TMD factorization ?

3) can we quantify **factorization breaking** effects?

4) how can we investigate gluon TMDs ?

## Jefferson Lab

## The frontier

Nucleon tomography in momentum space: to understand how hadrons are built in terms of the elementary degrees of freedom of QCD High-energy phenomenology:

to improve our understanding of high-energy scattering experiments and their potential to explore BSM physics

More open questions (phenomenology) :

1) what is the functional form of TMDs at low transverse momentum ?

2) what is its kinematic and flavor dependence?

3) can we attempt a global fit of TMDs ?

4) can we test the generalized **universality** of TMDs ?

5) what's the impact of hadron structure on the determination of Standard Model parameters ?



# TMD & collinear factorization

References:

- J.C. Collins "Foundations of perturbative QCD"
- SCET literature



Let's consider a process with three separate scales:

hadronic

(SIDIS, Drell-Yan, e+e- to hadrons, pp to quarkonium, ... )

 $\Lambda_{\rm QCD} \ll q_T \ll Q$ hard scale mass scale

(related to the) transverse momentum of the observed particle

The ratios



select the factorization theorem that we rely on.

According to their values we can access different "projections" of hadron structure



The key of phenomenology :



The key of phenomenology :



The key of phenomenology :



The key of phenomenology :



## Need of TMD evolution





## Need of TMD evolution





25

## Need of TMD evolution



Width of TMDs changes of one order of magnitude: we can we explain this with TMD evolution



# Extraction of unpolarized TMDs

References :

- "The 3D structure of the nucleon" EPJ A (2016) 52
- Bacchetta et al. JHEP 1706 (2017) 081
- A. Signori , PhD thesis
- Angelez-Martinez et al. arXiv:1507.05267
- EIC white paper, JLab 12 GeV white paper, ...

- ...

























**Electron-positron** annihilation data are still **missing** (only some azimuthal asymmetries are available)





**Electron-positron** annihilation data are still **missing** (only some azimuthal asymmetries are available)

crucial for analyses of TMD FFs !!



## Comparison with collinear PDF fits





## Comparison with collinear PDF fits





## Comparison with collinear PDF fits



On top of extending **data set**, many improvements are needed: higher **perturbative** orders, **matching** with high transverse momentum, **flavor** dependence, flexible functional forms...



## What do we know ?

(only a selection of results!)

	Framework	HERMES	COMPASS	DY	Z production	N of points
KN 2006 <u>hep-ph/0506225</u>	LO-NLL	×	×		<ul> <li>✓</li> </ul>	98
Pavia 2013 (+Amsterdam, Bilbao) <u>arXiv:1309.3507</u>	No evo (QPM)	•	×	*	×	1538
Torino 2014 (+JLab) <u>arXiv:1312.6261</u>	No evo (QPM)	✓ (separately)	<ul><li>(separately)</li></ul>	×	×	576 (H) 6284 (C)
DEMS 2014 <u>arXiv:1407.3311</u>	NLO-NNLL	×	×		•	223
EIKV 2014 <u>arXiv:1401.5078</u>	LO-NLL	1 (x,Q <sup>2</sup> ) bin	1 (x,Q <sup>2</sup> ) bin		•	500 (?)
Pavia/JLab 2017 <u>arXiv:1703.10157</u>	LO-NLL	~		~	<ul> <li></li> </ul>	8059
SV 2017 arXiv:1706.01473	NNLO- NNLL	×	×	•	~	309

( courtesy A. Bacchetta )



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Torino 2014 (+JLab) <u>arXiv:1312.6261</u>	No evo (QPM)	<ul> <li>(separately)</li> </ul>	<ul><li>(separately)</li></ul>	×	×	576 (H) 6284 (C)
DEMS 2014 <u>arXiv:1407.3311</u>	NLO-NNLL	×	×		•	223
EIKV 2014 <u>arXiv:1401.5078</u>	LO-NLL	1 (x,Q <sup>2</sup> ) bin	1 (x,Q <sup>2</sup> ) bin		•	500 (?)
Pavia/JLab 2017 <u>arXiv:1703.10157</u>	LO-NLL	~	~		~	8059
SV 2017 arXiv:1706.01473	NNLO- NNLL	×	×		~	309

( courtesy A. Bacchetta )



#### Features

	Framework	HERMES	COMPASS	DY	Z production	N of points
Pavia/JLab 2017 <u>arXiv:1703.10157</u>	LO-NLL	~	~	~	~	8059

#### **PRO**s

almost a **global fit** of quark unpolarized TMDs

includes TMD evolution

**replica (bootstrap)** fitting methodology

#### **kinematic dependence** in intrinsic part of TMDs

intrinsic momentum: **beyond the Gaussian** assumption

#### CONs

no "pure" info on TMD FFs

accuracy of TMD evolution : not the state of the art

only "low" transverse momentum (no fixed order and Y-term)

> flavor separation in the transverse plane : problematic



#### TMD PDFs at 1 GeV

$$\widetilde{f}_{1\text{NP}}^{a}(x, b_{T}^{2}) = \frac{1}{2\pi} e^{-g_{1a}\frac{b_{T}^{2}}{4}} \left(1 - \frac{\lambda g_{1a}^{2}}{1 + \lambda g_{1a}}\frac{b_{T}^{2}}{4}\right)$$

$$f^{a}_{1\text{NP}}(x, \boldsymbol{k}_{\perp}^{2}) = \frac{1}{\pi} \ \frac{\left(e^{-\frac{\boldsymbol{k}_{\perp}^{2}}{g_{1a}}} + \lambda \boldsymbol{k}_{\perp}^{2} e^{-\frac{\boldsymbol{k}_{\perp}^{2}}{g_{1a}}}\right)}{g_{1a} + \lambda \ g^{2}_{1a}}$$



## TMD PDFs at 1 GeV




# TMD PDFs at 1 GeV

$$\begin{split} \widetilde{f}_{1\mathrm{NP}}^{a}(x,b_{T}^{2}) &= \frac{1}{2\pi}e^{-g_{1a}\frac{b_{T}^{2}}{4}}\left(1 - \frac{\lambda g_{1a}^{2}}{1 + \lambda g_{1a}}\frac{b_{T}^{2}}{4}\right) \\ f_{1\mathrm{NP}}^{a}(x,\boldsymbol{k}_{\perp}^{2}) &= \frac{1}{\pi} \frac{\left(e^{-\frac{\boldsymbol{k}_{\perp}^{2}}{g_{1a}}} + \lambda \boldsymbol{k}_{\perp}^{2}e^{-\frac{\boldsymbol{k}_{\perp}^{2}}{g_{1a}}}\right)}{g_{1a} + \lambda g_{1a}^{2}} \end{split}$$

x-dependent width

$$g_1(x) = N_1 \frac{(1-x)^{\alpha} x^{\sigma}}{(1-\hat{x})^{\alpha} \hat{x}^{\sigma}}$$



# TMD PDFs at 1 GeV



x-dependent width

 $g_1(x) = N_1 \frac{(1-x)^{\alpha} x^{\sigma}}{(1-\hat{x})^{\alpha} \hat{x}^{\sigma}}$ 

Fragmentation function is similar Including TMD PDFs and FFs, in total: 11 free parameters (4 for TMD PDFs, 6 for TMD FFs, 1 for TMD evolution)



# Agreement data-theory

Flavor independent scenario

Flavor independent configuration | 11 parameters

	HERMES	HERMES	HERMES	HERMES
	$p \to \pi^+$	$p \to \pi^-$	$p \to K^+$	$p \to K^-$
Points	190	190	189	187
$\chi^2$ /points (	4.83	2.47	0.91	0.82

Points	Parameters	$\chi^2$	$\chi^2/{ m d.o.f.}$
8059	11	$12629\pm363$	$1.55\pm0.05$

**Hermes** P/D into π+: problems at low z

	HERMES	HERMES	HERMES	HERMES	COMPASS	COMPASS
	$D \to \pi^+$	$D \to \pi^-$	$D \to K^+$	$D \to K^-$	$D \rightarrow h^+$	$D \rightarrow h^{-}$
Points	190	190	189	189	3125	3127
$\chi^2$ /points	3.46	2.00	1.31	2.54	1.11	1.61

	E288 [200]	E288 [300]	E288 [400]	E605
Points	45	45	78	35
$\chi^2$ /points	0.99	0.84	0.32	1.12

	CDF Run I	D0 Run I	CDF Run II	D0 Run II
Points	31	14	37	8
$\chi^2$ /points	1.36	1.11	2.00	1.73

**Hermes** kaons better than pions: larger uncertainties from FFs

**Compass** : better agreement due to #points and normalization



### COMPASS, selected bins



#### Deuteron h<sup>-</sup> $\chi^2$ /dof = 1.58



### COMPASS, selected bins



Deuteron h<sup>-</sup>  $\chi^2/dof = 1.58$ 



# HERMES, selected bins



Contributions to chi2 mainly from **normalization**, not shape (also in Z-boson production)



# HERMES, selected bins



$$\chi^{2}/dof = 4.80$$

The worst of all channels...

Contributions to chi2 mainly from **normalization**, not shape (also in Z-boson production)



# HERMES, selected bins



 $\chi^{2}/dof = 4.80$ 

The worst of all channels...

However **normalizing** the theory curves to the first bin, without changing the parameters of the fit,  $\chi^2$ /dof becomes good

Contributions to chi2 mainly from **normalization**, not shape (also in Z-boson production)



# Best-fit values



#### **Caveat for comparisons** :

NP effects (as the intrinsic momentum) always depend on the accuracy of the perturbative part ;

#### determined as observed - calculable

Bacchetta, Delcarro, Pisano, Radici, Signori,
Signori, Bacchetta, Radici, Schnell arXiv:1309.3507
Schweitzer, Teckentrup, Metz, arXiv:1003.2190
Anselmino et al. arXiv:1312.6261 [HERMES]
Anselmino et al. arXiv:1312.6261 [HERMES, high z]
Anselmino et al. arXiv:1312.6261 [COMPASS, norm.]
Anselmino et al. arXiv:1312.6261 [COMPASS, high z, norm.]
Echevarria, Idilbi, Kang, Vitev arXiv:1401.5078 (Q = 1.5 GeV)

Red/orange regions : 68% CL from replica method

Inclusion of  $\ensuremath{\text{DY/Z}}$  diminishes the correlation

Inclusion of Compass increases the  $\langle P_{\perp}^2 \rangle$  and reduces its spread

**e+e-** would further reduce the correlation



# Polarized TMDs (Sivers)

References :

- ...

- "The 3D structure of the nucleon" EPJ A (2016) 52
- STAR arXiv:1511.06003
- Compass: arxiv:1704.00488



Gauge invariance and T-reversal invariance generate a sign change between the Sivers TMD PDF in Drell-Yan and Semi-Inclusive DIS



Gauge invariance and T-reversal invariance generate a sign change between the Sivers TMD PDF in Drell-Yan and Semi-Inclusive DIS

Collins, PLB 536 (02)



Gauge invariance and T-reversal invariance generate a sign change between the Sivers TMD PDF in Drell-Yan and Semi-Inclusive DIS

Collins, PLB 536 (02)





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Collins, PLB 536 (02)







#### Gluon TMDs



# Gluon TMDs

 $e \ p \to e \ \text{jet jet } X$ 

 $p p \to J/\psi \gamma X$ 

gluon TMD

gluon TMD

 $p \ p \to \eta_c \ X$ 



see, e.g.,

 $\pm 1$ 

 $\pm 1$ 

- Boer, den Dunnen, Pisano, Schlegel, Vogelsang, PRL108 (12)
- den Dunnen, Lansberg, Pisano, Schlegel, PRL 112 (14)
- AS: PhD thesis , arXiv:1602.03405

gluon TMD

- AFTER@LHC working group: arXiv:1702.01546 , arXiv:1610.05228 , ....
- Echevarria et al. arXiv:1502.05354

÷ ...



### Higgs transverse momentum

G. Ferrera, talk at REF 2014, Antwerp, <u>https://indico.cern.ch/event/330428/</u>





#### Higgs transverse momentum

G. Ferrera, talk at REF 2014, Antwerp, <u>https://indico.cern.ch/event/330428/</u>





# Spin 1 TMDs

References :

- quark TMDs : Phys.Rev. D62 (2000) 114004

- gluon TMDs : JHEP 1610 (2016) 013

- ...



#### $\Phi_{ij}(k,P;S,T) \sim \text{F.T.} \langle PST | \ \overline{\psi}_j(0) \ U_{[0,\xi]} \ \psi_i(\xi) \ |PST\rangle_{|_{LF}}$

Quarks	$\gamma^+$	$\gamma^+\gamma^5$	$i\sigma^{i+}\gamma^5$
U	$f_1$		$h_1^\perp$
L		$g_1$	$h_{1L}^{\perp}$
Т	$f_{1T}^{\perp}$	$g_{1T}$	$oldsymbol{h_1},h_{1T}^\perp$
LL	$f_{1LL}$		$h_{1LL}^{\perp}$
$\operatorname{LT}$	$f_{1LT}$	$g_{1LT}$	$h_{1LT},h_{1LT}^{\perp}$
TT	$f_{1TT}$	$g_{1TT}$	$h_{1TT},h_{1TT}^{\perp}$



extraction of a **quark not** collinear with the proton

a similar scheme holds for TMD FFs and gluons



**bold :** also collinear red : time-reversal odd (universality properties)

# quark TMD PDFs



**bold :** also collinear red : time-reversal odd (universality properties)



# Conclusions : a path to move forward

**1)** Phenomenology of TMDs is well underway ...

**2)** ... but there are a lot of theoretical challenges to be addressed: definition of kinematic regions in SIDIS, matching, perturbative accuracy, a better understanding of hadronization, context for gluon TMDs , ...

3) we definitely need more data, at the moment especially for e+e-

**4)** Working with some approximations, we are getting closer to a global fit analysis of TMDs

**5)** polarized structure functions unexplored from the point of view of QCD, but we have guidance from parton model studies (see JLab activities)



# Backup



#### Beware of different notations...

#### Amsterdam Seattle (arXiv:1108.1713)

p	k	momentum of parton in distribution function
$p_T$	$m{k}_{\perp}$	parton transverse momentum in distribution function
k	p	momentum of fragmenting parton
$oldsymbol{k}_T$	$p_{\perp}$	trans. momentum of fragmenting parton w.r.t. final hadron
$oldsymbol{K}_T$	$P_{\!\perp}$	trans. momentum of final hadron w.r.t. fragmenting parton
$P_{h\perp}$	$oldsymbol{P}_{hT}$	transverse momentum of final hadron w.r.t. virtual photon



#### TMDs and their evolution

FT of TMDs :



Nonperturbative parts defined in a "negative" way : **observed-calculable** 



#### TMDs and their evolution

Distribution for intrinsic transverse momentum (and its FT):

 $(\tilde{F}_{i,NP}(x,\bar{b}_T;\{\lambda\}))$ a Gaussian ?

Soft gluon emission





#### TMDs and their evolution

Distribution for intrinsic transverse momentum (and its FT):

$$\tilde{F}_{i,NP}(x,\bar{b}_T;\{\lambda\})$$
a Gaussian ?

Soft gluon emission

$$g_K(\overline{b}_T; \{\lambda\})$$

Separation of **b**<sub>T</sub> regions

$$\hat{b}_T(b_T; b_{\min}, b_{\max}) \xrightarrow{b_{\max}}, \begin{array}{c} b_T \rightarrow +\infty \\ \sim & b_T \\ \sim & b_T \\ b_{\min} \\ \end{array}, \begin{array}{c} b_T \rightarrow +\infty \\ b_{\max} \\ b_T \rightarrow 0 \end{array}$$

High b<sub>T</sub> limit : avoid Landau pole

Low **b**<sub>T</sub> limit : recover fixed order expression



#### Models - evolution and $b_{\mathsf{T}}$ regions

$$g_K(b_T; g_2) = -g_2 \frac{b_T^2}{2}$$

$$\hat{b}(b_T; b_{\min}, b_{\max}) = b_{\max} \left( \frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}} \right) \xrightarrow{b_{\max}, b_T \to +\infty} b_{\min}, b_T \to 0$$

$$b_{\max} = 2e^{-\gamma_E}$$

$$\mu b_{\min} \frac{C_1}{\overline{b_*}} 2e^{-\gamma_E \overline{b_*}} \not \Rightarrow b_{\max} \left(\frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}}\right)^{\frac{1}{4}} \qquad b_T$$
These choicesegorarantee that for  $C_1$   $b_{\min} = \frac{C_1}{Q}$ 

$$Q=2 \text{ GeV}$$

$$Q=2 \text{ GeV}$$

$$Q=5 \text{ GeV}$$

$$Q=20 \text{ GeV}$$

$$Q=20 \text{ GeV}$$

$$D=1 \text{ GeV the TMD coincides with the NP model}$$

#### Models - evolution and $b_{\mathsf{T}}$ regions

$$g_K(b_T; g_2) = -g_2 \frac{b_T^2}{2}$$

$$\hat{b}(b_T; b_{\min}, b_{\max}) = b_{\max} \left( \frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}} \right) \xrightarrow{b_{\max}, b_T \to +\infty} b_{\min}, b_T \to 0$$

$$\underbrace{b_{\min} \sim 1/Q, \ \mu_{\hat{b}} < Q}_{b_{\min} \sim 1/Q, \mu_{\hat{b}} < Q}$$
The phenomenological importance of b\_{\min} is a signal that -especially in SIDIS data at low Q- we are exiting the proper 2 MD region and x = approaching the region of collinear factorization
$$C_1 \quad b_{\min} = \frac{C_1}{Q}$$

$$\frac{d_0}{d_0} = 2 \text{ GeV} \quad b_T \text{ (GeV-1)}$$

# Kinematic dependence

Comparison with other extractions :














## Data sets and selections

	HERMES	HERMES	HERMES	6 HERMI	ES		
	$p \to \pi^+$	$p \to \pi^-$	$p \to K^+$	$p \to K$	-	IND factorization (Pht/ z << U <sup>2</sup> )	
Reference			[61]			avoid target fragmentation (low z	
₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩₩	ENER STREETS CONTRACTOR	$Q^2 > 1.4 \ { m GeV}^2$				and exclusive contributions (high z)	
Cuts		0.2 < z < 0.7					
uter a constant and a constant and	$P_{hT}$	$< Min[0.2 \ Q$	$[0, 0.7 \ Qz] + 0$	$0.5  { m GeV}$			
Points	190	190	189	187		1 I I I I I I I I I I I I I I I I I I I	
Max. $Q^2$	$9.2 \text{ GeV}^2$					IN ORDER TO AVOID THE PROBLEMS	
x range		0.06 <	< x < 0.4			(see Compass coll., Erratum)	
	HERMES	HERMES	HERMES	HERMES	COMPASS	COMPASS	
	$D \to \pi^+$	$D \to \pi^-$	$D \to K^+$	$D \to K^-$	$D \to h^+$	$D \rightarrow h^-$	
Reference		[6]	[]			[62]	
	$Q^2 > 1.4 \ { m GeV}^2$						
Cuts	0.2 < z < 0.7						
	an a		$P_{hT} < M$	$\operatorname{in}[0.2\;Q,0.]{}$	$7 \ Qz] + 0.5 \ Ge$	eV	
Points	190	190	189	189	3125	3127	
Max. $Q^2$	$9.2 \text{ GeV}^2$					$10 \text{ GeV}^2$	
x range	0.06 < x < 0.4					0.006 < x < 0.12	
					Obcorruphle	$m = (m \sim \mathbf{P}^2 \cap \mathbf{C}^2)  \text{and}  (29)$	

## Data sets and selections

	E288 200	E288 300	E288 400	E605	
Reference	[65]	[65]	[65]	[66]	
Cuts	$q_T < 0.2 \ Q + 0.5 \ { m GeV}$				
Points	45	45	78	$\overline{35}$	
$\sqrt{s}$	$19.4 \mathrm{GeV}$	$23.8 \mathrm{GeV}$	$27.4 \mathrm{GeV}$	$38.8 \mathrm{GeV}$	
Q range	4-9 GeV	4-9 GeV	5-9, 11-14 GeV	7-9, 10.5-18 GeV	
Kin. var.	y = 0.4	y = 0.21	y = 0.03	$-0.1 < x_F < 0.2$	

TMD factorization  $(q_T \ll Q^2)$ 

#### **Drell-Yan**

	CDF Run I	D0 Run I	CDF Run II	D0 Run II		
Reference	[67]	[68]	[69]	[70]		
Cuts	$q_T < 0.2 \ Q + 0.5 \ \text{GeV} = 18.7 \ \text{GeV}$					
Points	31	14	37	8		
$\sqrt{s}$	1.8 TeV	$1.8 \mathrm{TeV}$	$1.96 { m ~TeV}$	1.96 TeV		

#### Ζ

normalization : fixed from DEMS fit, different from exp. (not really relevant for TMD parametrizations)



### Target vs current vs central regions





### Target vs current vs central regions



## Evolution at work







Overview of the terminology

 $C_{i/j}$  Wilson coefficients : expansion of the TMD distribution on a basis of collinear PDFs







Anomalous dimension of the TMD and logarithmic expansion

$$\begin{split} \gamma_F[\alpha_s(\mu), \zeta/\mu^2] &\sim \underbrace{\alpha_s L}_{\text{LL}} + \underbrace{(\alpha_s + \alpha_s^2 L)}_{\text{NLL}} + \underbrace{(\alpha_s^2 + \alpha_s^3 L)}_{\text{NNLL}} + \cdots \\ &\sim 1 + \alpha_s + \alpha_s^2 + \cdots \end{split} \qquad L = \ln \frac{Q^2}{\mu} , \quad \alpha_s L \sim 1 \end{split}$$





Anomalous dimension of the TMD and logarithmic expansion

$$\begin{split} \gamma_F[\alpha_s(\mu), \zeta/\mu^2] &\sim \underbrace{\alpha_s L}_{\text{LL}} + \underbrace{(\alpha_s + \alpha_s^2 L)}_{\text{NLL}} + \underbrace{(\alpha_s^2 + \alpha_s^3 L)}_{\text{NNLL}} + \cdots \\ &\sim 1 + \alpha_s + \alpha_s^2 + \cdots \end{split} \qquad \begin{array}{l} L &= \ln \frac{Q^2}{\mu} \ , \ \alpha_s L \sim 1 \end{split}$$

**Collins-Soper kernel** : a power series in the coupling

$$K(b_T;\mu_b) \sim 1 + \alpha_s + \alpha_s^2 \cdots$$

accuracy chosen consistently with Wilson coefficients and anomalous dimension



 $C_{i/i}$  Wilson coefficients : expansion of the TMD distribution on a basis of collinear PDFs



Anomalous dimension of the TMD and logarithmic expansion

$$\mu_{\hat{b}} = 2e^{-\gamma_E}/\bar{b}_{\star}$$

$$\gamma_F[\alpha_s(\mu), \zeta/\mu^2] \sim \underbrace{\alpha_s L}_{\text{LL}} + \underbrace{(\alpha_s + \alpha_s^2 L)}_{\text{NLL}} + \cdots$$
$$\sim 1 + \alpha_s + \cdots$$

**Collins-Soper kernel** : a power series in the coupling

 $K(b_T;\mu_b) \sim 1 + \alpha_s + \cdots$ 

$C_{i/j}$	$\gamma_{ m nc}$	$\Gamma_{\rm cusp}$	K	accuracy
0	0	0	0	$\mathbf{QPM}$
0	0	1	0	LO-LL
0	1	2	1	LO-NLL
0	2	3	2	LO-NNLL
1	1	2	1	NLO-NLL
1	2	3	2	NLO-NNLL
2	2	3	2	NNLO-NNLL

Jefferson Lab

## Pavia / Amsterdam / Bilbao 2013

### proton target global $\chi^2$ / d.o.f. = 1.63 ± 0.12 no flavor dep. 1.72 ± 0.11





 $Q^2~({
m GeV}^2)$ COMPASS  $M_D^{h^+}$  $10^{\circ}$ <*z*>=0.23  $Q^2 = 7.36 \,\, {
m GeV}^2$  $Q^2$ =7.57 GeV<sup>2</sup>  $10^{-1}$ <z>=0.28 y = 0.45y = 0.27<z>=0.33 TILII <z>=0.38 IIII IIIIII <z>=0.45 10<sup>4</sup> <z>=0.55 •  $Q^2$ =4.07 GeV<sup>2</sup>  $Q^2$ =4.47 GeV<sup>2</sup>  $Q^2\!\!=\!\!4.57~{
m GeV}^2$  $Q^2$ =4.62 GeV<sup>2</sup>  $10^{-1}$ y = 0.63*y* =0.46 y = 0.28y = 0.17 $0.25 \quad 0.50 \quad 0.75$ IIIII IIII TITI  $10^{0}$ .........  $Q^2$ =2.94 GeV<sup>2</sup>  $Q^2$ =2.95 GeV<sup>2</sup>  $Q^2 = 2.90 \,\,{
m GeV}^2$  $Q^2 = 2.95 \,\,\mathrm{GeV}^2$  $10^{-1}$ y = 0.64y = 0.46y = 0.31y = 0.181111 IIIII IIII IIIII IIIII IIIIIT. 10  $Q^2$ =1.92 GeV<sup>2</sup>  $Q^2 = 1.76 \,\, {
m GeV}^2$  $Q^2 = 1.92 \,\,{
m GeV}^2$  $Q^2 = 1.92 \,\,{
m GeV}^2$  $Q^2 = 1.93 {
m ~GeV}^2$  $10^{-1}$ y = 0.43y = 0.20y = 0.59y = 0.30y = 0.14 $0.25 \quad 0.50 \quad 0.75 \quad 0.25 \quad 0.50 \quad 0.75 \quad 0.50 \quad$  $P_T (\text{GeV})$  $x_B$ Anselmino, Boglione, Gonzalez, Melis, Prokudin, JHEP 1404 (14) Jefferson Lab

COMPASS  $M_D^{h^+}$ 

COMPASS  $M_D^{h^+}$ 

 $\chi^2$  / dof = 3.79 with ad-hoc

normalization



COMPASS  $M_D^{h^+}$ 

 $\chi^2/dof = 3.79$ with ad-hoc normalization

see Compass coll. Erratum



Anselmino, Boglione, Gonzalez, Melis, Prokudin, JHEP 1404 (14) Jefferson Lab

COMPASS  $M_D^{h^+}$ 

simple Gaussian ansatz

 $\chi^2$  / dof = 3.79 with ad-hoc normalization

see Compass coll. Erratum





#### D'Alesio, Echevarria, Melis, Scimemi, JHEP 1411 (14)





NLO-NNLL analysis with evaluation of theoretical uncertainties

very good









#### D'Alesio, Echevarria, Melis, Scimemi, JHEP 1411 (14)



NLO-NNLL analysis with evaluation of theoretical uncertainties

very good



# KN 2006



≈100 data points Q²>4 GeV



# KN 2006



Brock, Landry, Nadolsky, Yuan, PRD67 (03)



# KN 2006



65

## EIKV 2014

Parametrizations for intrinsic momenta and soft gluon emission :

 $F_{NP}(b_T, Q)^{\text{pdf}} = \exp\left[-b_T^2 \left(g_1^{\text{pdf}} + \frac{g_2}{2}\ln(Q/Q_0)\right)\right]$  $F_{NP}(b_T, Q)^{\text{ff}} = \exp\left[-b_T^2 \left(g_1^{\text{ff}} + \frac{g_2}{2}\ln(Q/Q_0)\right)\right]$ 

#### **Pros and Cons** :

1) a global analysis of SIDIS and DY/Z/W data

2) TMD evolution at LO-NLL

3) multidimensionality not exploited

4) chi-square not provided

5) can't be considered as a "complete" fit

Jefferson Lab





p<sub>T</sub> (GeV)

#### SIDIS











20



 $V^{-2}$ 





 $V^{-2}$ )

Echevarria et al. arXiv:1401.5078

## Other studies

. . .

...

CSS formalism on DY/Z/W data:

1) Davies-Webber-Stirling (DOI: <u>10.1016/0550-3213(85)90402-X</u>)

2) Ladinsky-Yuan (DOI: <u>10.1103/PhysRevD.50.R4239</u>)

3) BLNY [DOI: <u>10.1103/PhysRevD.63.013004</u>]

4) Hirai, Kawamura, Tanaka (DOI: <u>10.3204/DESY-PROC-2012-02/136</u>) - complexb prescription

combined SIDIS/DY/W/Z :

5) Sun, Yuan (arXiv:1308.5003)

6) Isaacson, Sun, Yuan, Yuan (arXiv:1406.3073)



## ... and the next challenges

#### The goal is not only to fit data, but to answer fundamental questions in QCD in the best possible way

**11** identification of the current fragmentation region in SIDIS ?

**12)** rise the accuracy of transverse momentum resummation

**13)** match TMD and collinear factorization : fixed-order description of the high transverse momentum region and its matching to the low transverse momentum one

**14)** order the hadronic tensor in terms of definite rank

15) include electron-positron annihilation, LHC and JLab data
16) address the flavor decomposition in transverse momentum
17) address the polarized structure functions **18) Monte Carlo generators and TMDs**19) what about spin 1 targets ?



20) ...

## Monte Carlo generators



### Mapping the hadronization description in the Pythia MCEG to the correlation functions of TMD factorization



see the talk by M. Diefenthaler



