

HYP2003 at Jefferson Lab

Hypernuclear Excitation
through
Kaon Electroproduction

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Collaboration with

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INTRODUCTION

1) (e,e'K+) experiments started at JLab, confirming the great success of hypernuclear production on ^{12}C :

---> Make clear the impact to hypernuclear physics.

2) Novel characteristics of selectively exciting unnatural-parity high spin states:

---> Complementary role to (K-, π^-) and (π^+ ,K+)

---> Excite various multiplet states which are not accessible before.

3) High resolution reaction spectroscopy (0.3–0.5MeV)
cf. 2.5MeV for (K-, π^-), 1.5MeV for (π^+ ,K+)

---> Helpful tool to investigate dynamical coupling of hyperons to nuclear core excitation.

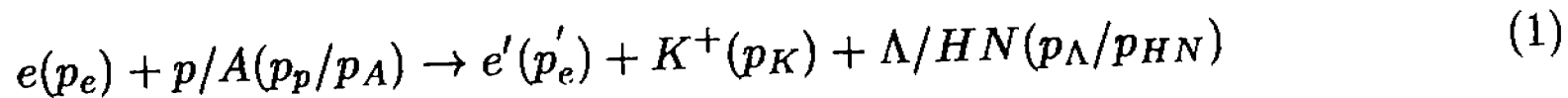
---> Production of p-deficient hypernuclei which are not accessible.

CONTENTS

- 1) Basic properties of the $\gamma p \rightarrow \Lambda K^+$ process
- 2) Make clear the characteristics of photo/electro-production of hypernuclei by taking a typical medium-mass target (^{28}Si) with simplified w.f. ($d_{5/2}$)¹².
- 3) Show a realistic calculation with the full (sd)ⁿ w.f. for $^{28}\text{Si}(\gamma, K^+)_{\Lambda}^{28}\text{Al}$, which should be compared with the coming exp. at JLab.
- 4) Based on the success of the $^{12}\text{C}(e, e' K^+)_{\Lambda}^{12}\text{B}$ experiment, we present theoretical spectra of producing "proton-deficient" p-shell hypernuclei.
- 5) Extend the approach to produce heavier hypernuclei $A \sim 50$, hoping to disclose new areas.
- 6) Summary

Kinematics and amplitudes for electro-production process

The kinematics of the electroproduction reaction



on proton(p)/nuclear(A) target producing Λ hyperon/hypernucleus(HN) is depicted in Fig. 2.

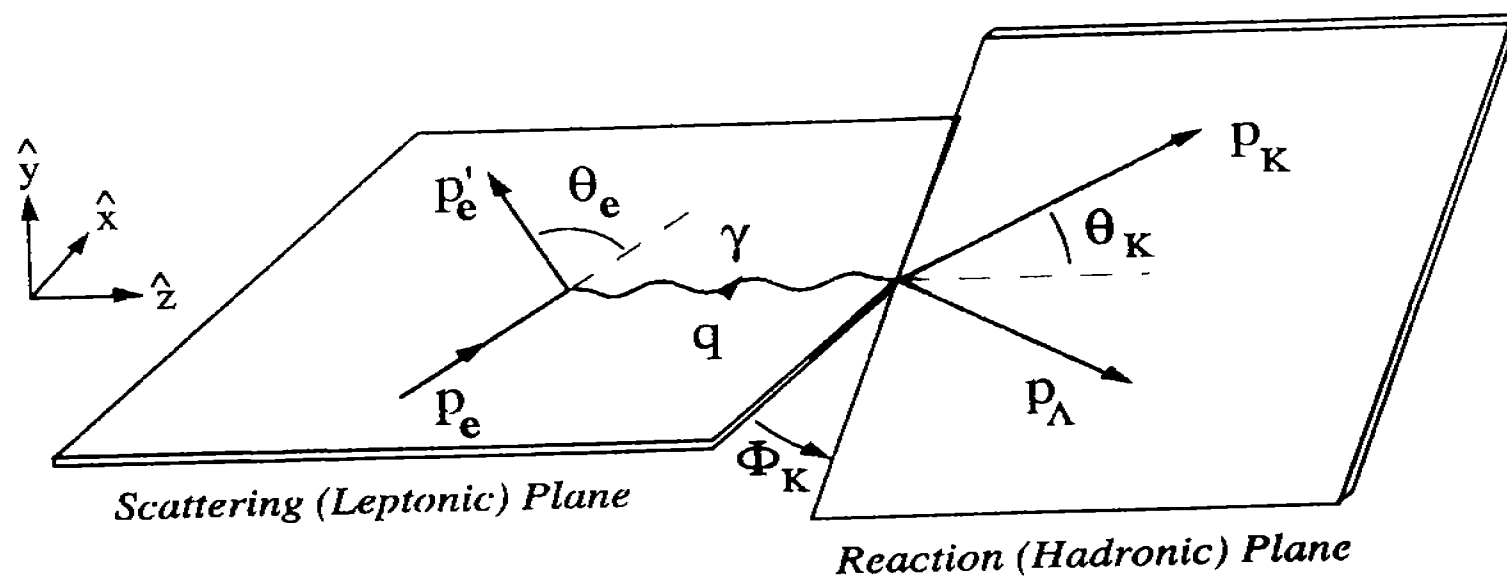
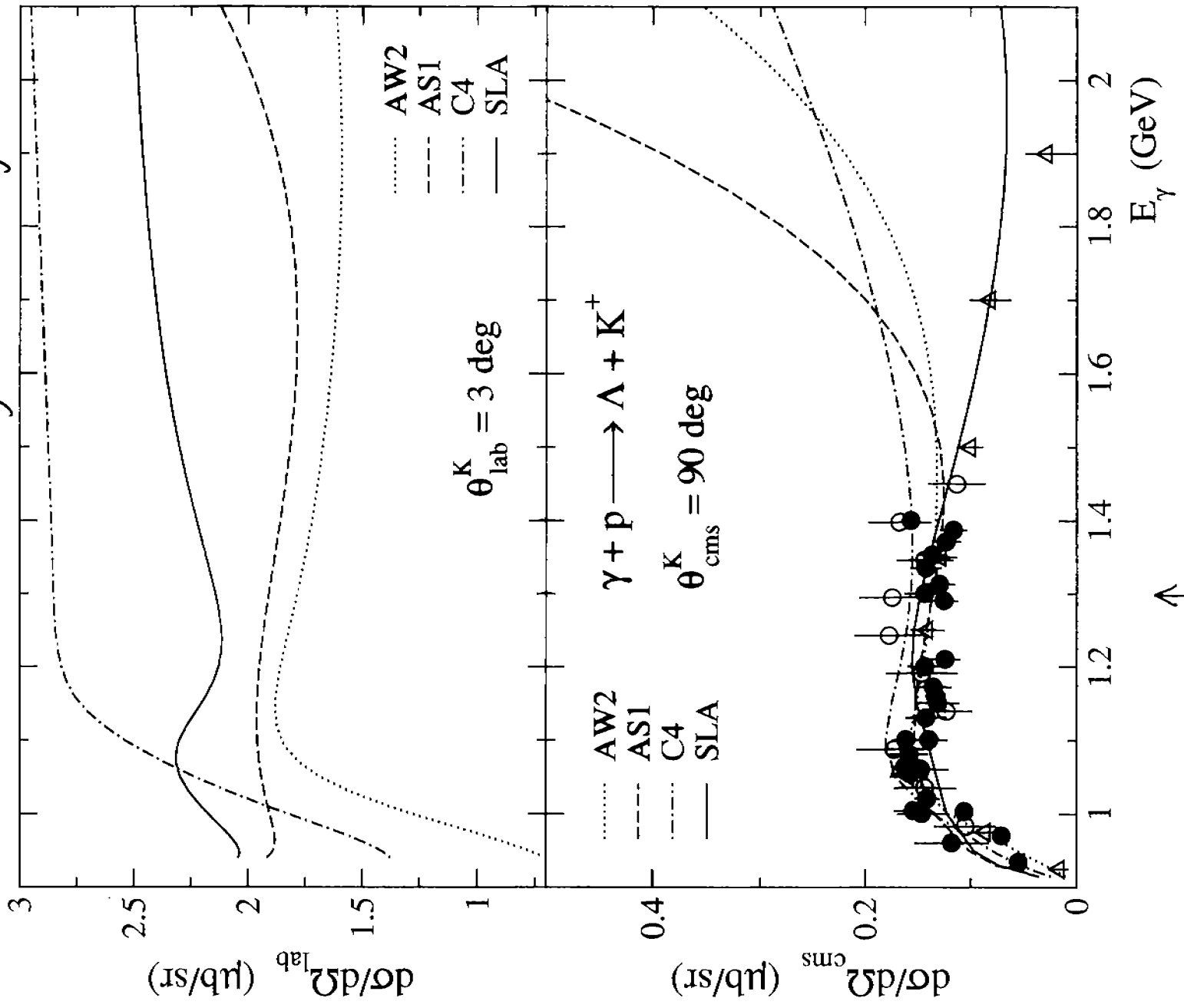


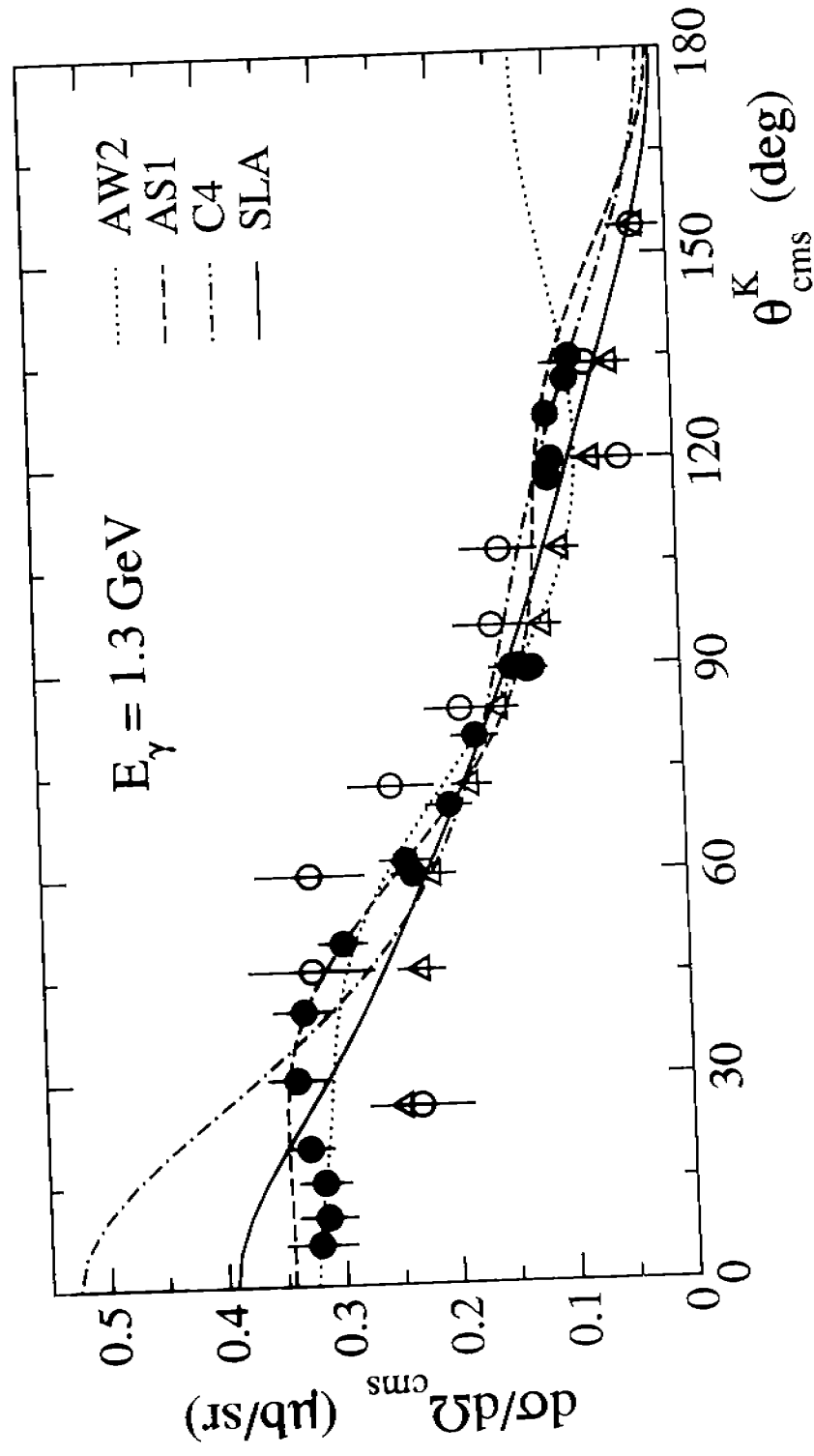
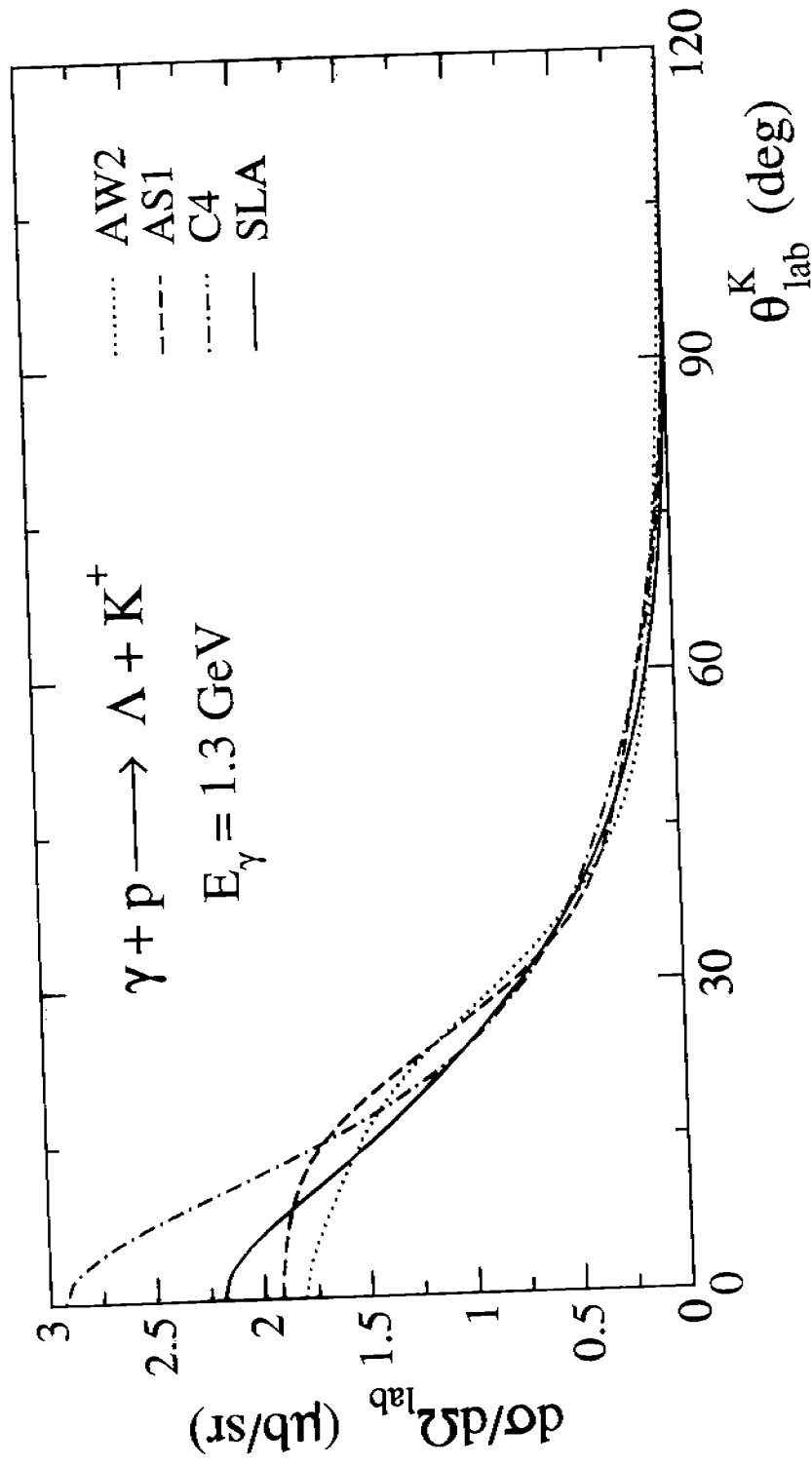
Figure 2: Kinematics of an electroproduction process.

EXP : SAPHIR

- AW4 : Adelseck-Wright (1988)
- AS1 : Adelseck-Saghai (1990)
- C4 : Williams-Ji-Cotanch (1992)
- SLA : Mizutani-Fayard-Lamot-Saghai (1998)



↑
 $E_\gamma = 1.3 \text{ GeV}$



Framework of the calculation

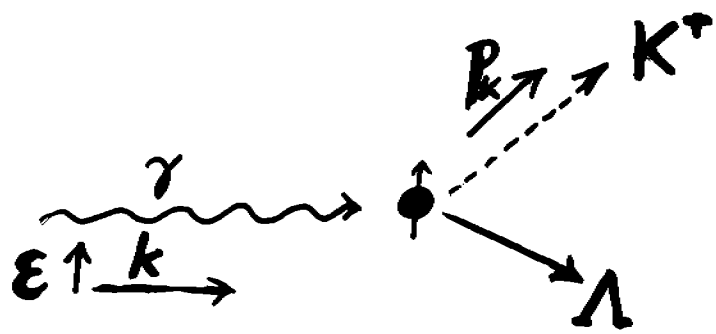
The elementary amplitude for the kaon photoproduction $\gamma + p \rightarrow \Lambda + K^+$ is written in the 2-body center-of-momentum (2CM) system as

$$\begin{aligned} \mathcal{M} &\equiv \frac{(2\pi)^2}{\sqrt{s}} \sqrt{e_\gamma e_p e_\Lambda e_K} \langle \mathbf{q}, -\mathbf{q} | t | \mathbf{k}, -\mathbf{k} \rangle \\ &= iF_1(\boldsymbol{\sigma} \cdot \boldsymbol{\epsilon}) + F_2(\boldsymbol{\sigma} \cdot \hat{\mathbf{q}})[\boldsymbol{\sigma} \cdot (\hat{\mathbf{k}} \times \boldsymbol{\epsilon})] + iF_3(\boldsymbol{\sigma} \cdot \hat{\mathbf{k}})(\hat{\mathbf{q}} \cdot \boldsymbol{\epsilon}) + iF_4(\boldsymbol{\sigma} \cdot \hat{\mathbf{q}})(\hat{\mathbf{q}} \cdot \boldsymbol{\epsilon}). \end{aligned} \quad (2.1)$$

$$\begin{aligned} \left. \frac{d\sigma}{d\Omega} \right|_{2\text{CM}} &= \frac{q}{k} \left\{ |F_1|^2 + |F_2|^2 - 2\cos\theta_{\text{CM}} \text{Re}(F_1^* F_2) \right. \\ &\quad \left. + \frac{1}{2}\sin^2\theta_{\text{CM}}[|F_3|^2 + |F_4|^2 + \text{Re}(F_1^* F_4 + F_2^* F_3 + F_3^* F_4 \cos\theta_{\text{CM}})] \right\}, \end{aligned} \quad (2.2)$$

$$\begin{aligned} \mathcal{P}_Y \left. \frac{d\sigma}{d\Omega} \right|_{2\text{CM}} &= \frac{q}{k} \sin\theta_{\text{CM}} \text{Im} \left\{ -2F_1^* F_2 - F_1^* F_3 + F_2^* F_4 + \sin^2\theta_{\text{CM}} F_3^* F_4 \right. \\ &\quad \left. + \cos\theta_{\text{CM}}[F_2^* F_3 - F_1^* F_4] \right\}, \end{aligned} \quad (2.3)$$

$\{F_i\}$ CGLN type amplitudes
(complex)



ϵ : photon polarization
 σ : Pauli spin of baryon

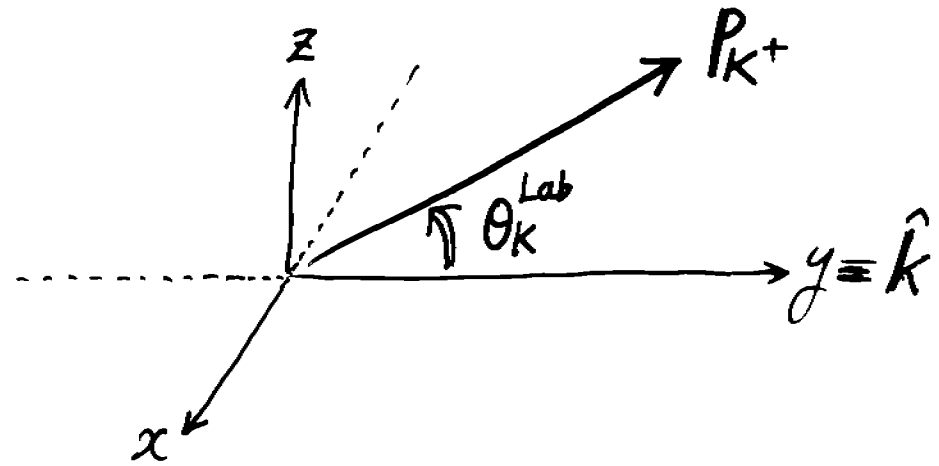
$$\text{Lab: } \mathcal{M} = a_1 (\sigma \cdot \epsilon) + a_2 (\sigma \cdot \hat{k}) (\hat{p}_K \cdot \epsilon) + a_3 (\sigma \cdot \hat{p}_K) (\hat{p}_K \cdot \epsilon) + a_4 (\hat{k} \times \hat{p}_K) \cdot \epsilon$$

$$\{F_i\}_{\text{CM}} \longrightarrow \{a_i\}_{\text{Lab}}$$

New frame

$$\downarrow$$

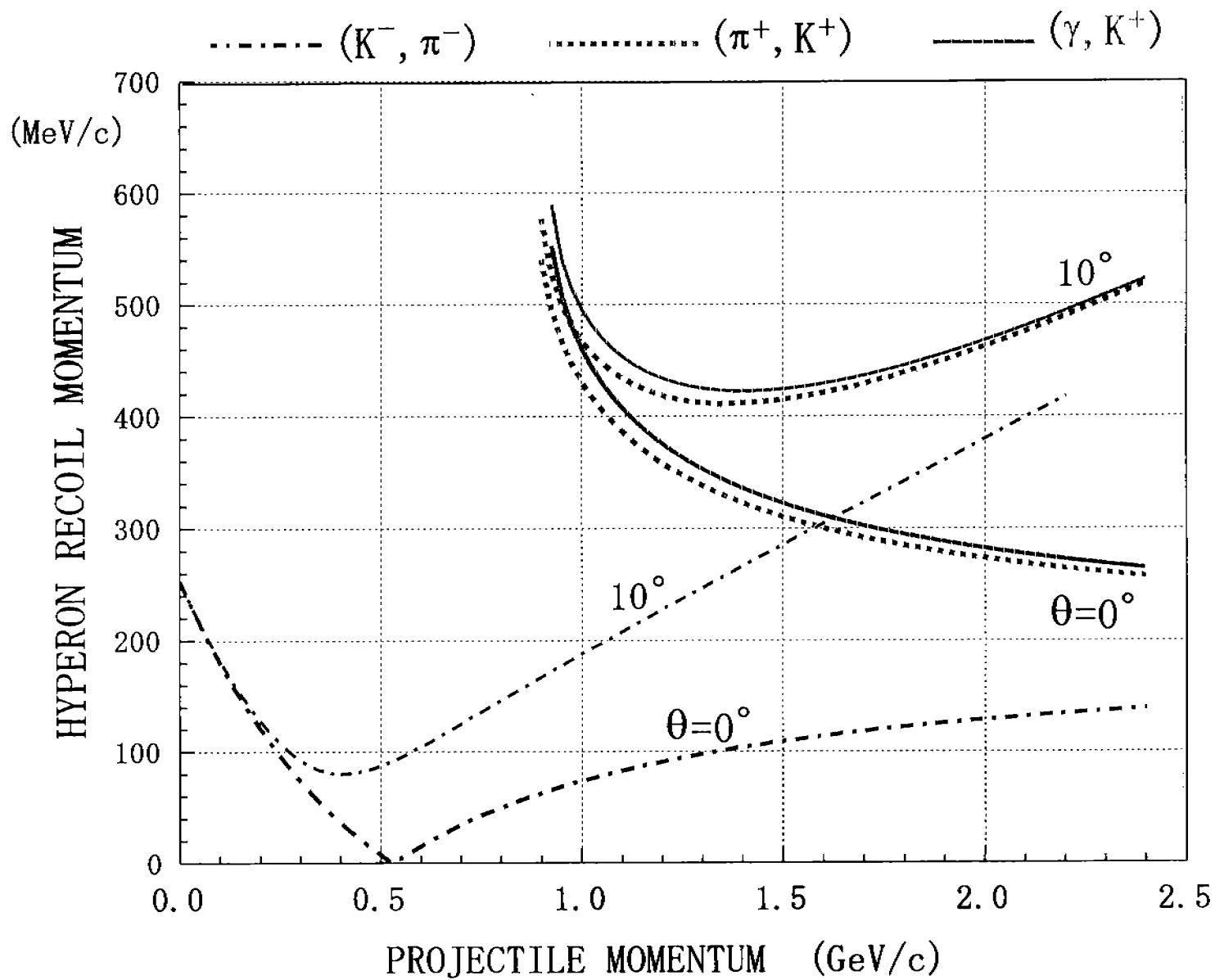
$$\{f_0, g_0, \pm 1\}_{\text{Lab}}$$



$$\mathcal{M} \equiv \epsilon_0 (f_0 + g_0 \sigma_0) + \epsilon_x (g_1 \sigma_{+1} + g_{-1} \sigma_{-1})$$

$$\text{Cross section: } \frac{d\sigma}{d\Omega} = \alpha [|f_0|^2 + |g_0|^2 + |g_1|^2 + |g_{-1}|^2]$$

$$\text{Polarization: } p \frac{d\sigma}{d\Omega} = \alpha [2\text{Re}(f_0 g_0^*) + |g_1|^2 - |g_{-1}|^2]$$



(γ, K^+)

$q_\Lambda = 350-420$ MeV/c at $E_\gamma = 1.3$ GeV

Numerical comparison of f and g 's ($E_\gamma=1.3$ GeV)

MODEL	$ f_0 ^2$	$ g_0 ^2$	$ g_{+1} ^2$	$ g_{-1} ^2$	$d\sigma/d\Omega$	$\text{Pol}(\Lambda)$
$\theta_K=3$ AS1	0.007	0.398	0.191	0.209	1.91	-0.055
$\theta_K=3$ AW2	0.0001	0.369	0.188	0.192	1.78	-0.005
$\theta_K=3$ C4	0.0002	0.609	0.290	0.299	2.85	-0.019
$\theta_K=3$ SLA	0.002	0.451	0.224	0.222	2.14	-0.016
$\theta_K=10$ AS1	0.055	0.375	0.175	0.205	1.84	-0.142
$\theta_K=10$ AW2	0.001	0.296	0.209	0.218	1.65	-0.007
$\theta_K=10$ C4	0.001	0.569	0.190	0.217	2.23	-0.066
$\theta_K=10$ SLA	0.026	0.392	0.186	0.178	1.78	-0.053

($\mu\text{b}/\text{sr}$)

AS1: Adelseck-Saghai (1990)

AW2: Adelseck-Wright (1988)

C4: Williams-Ji-Cotanch (1992)

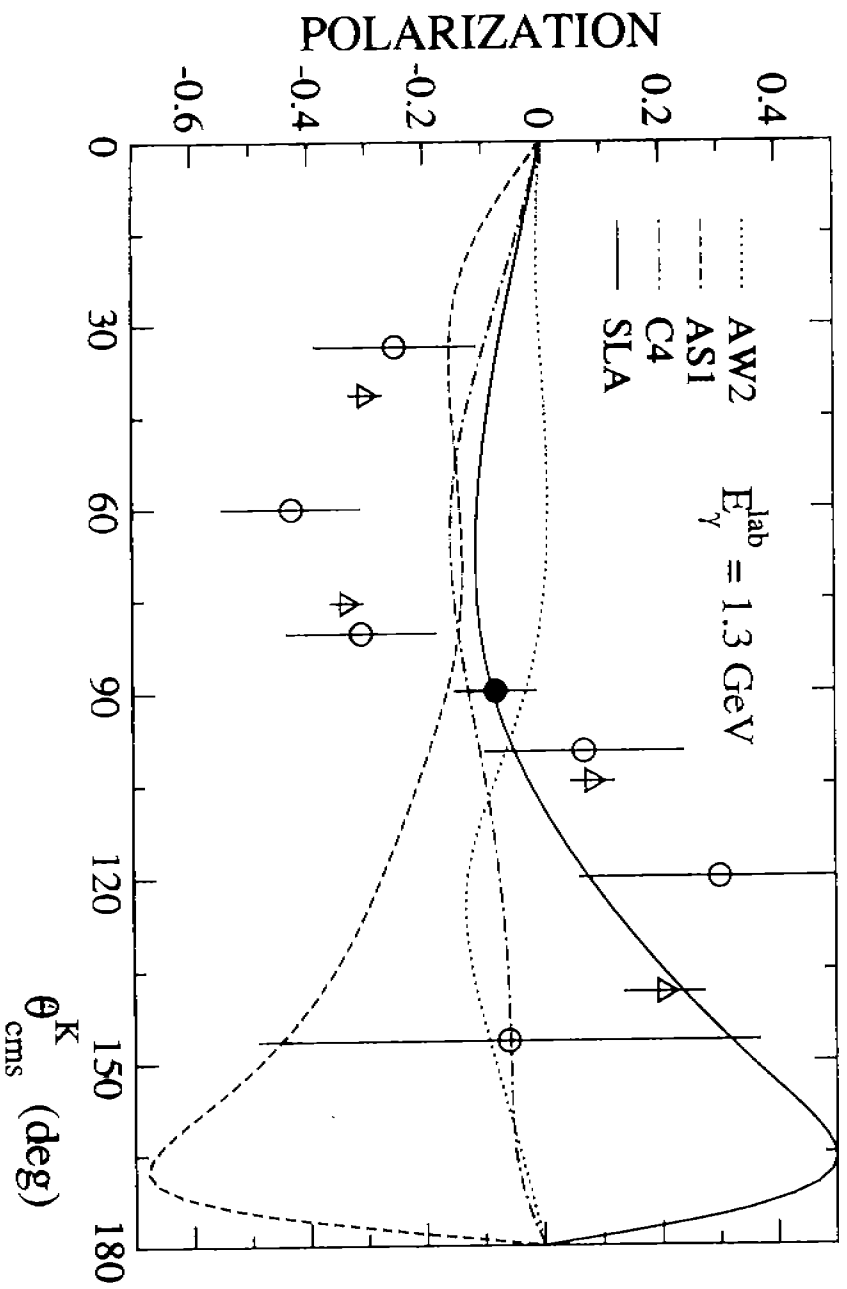
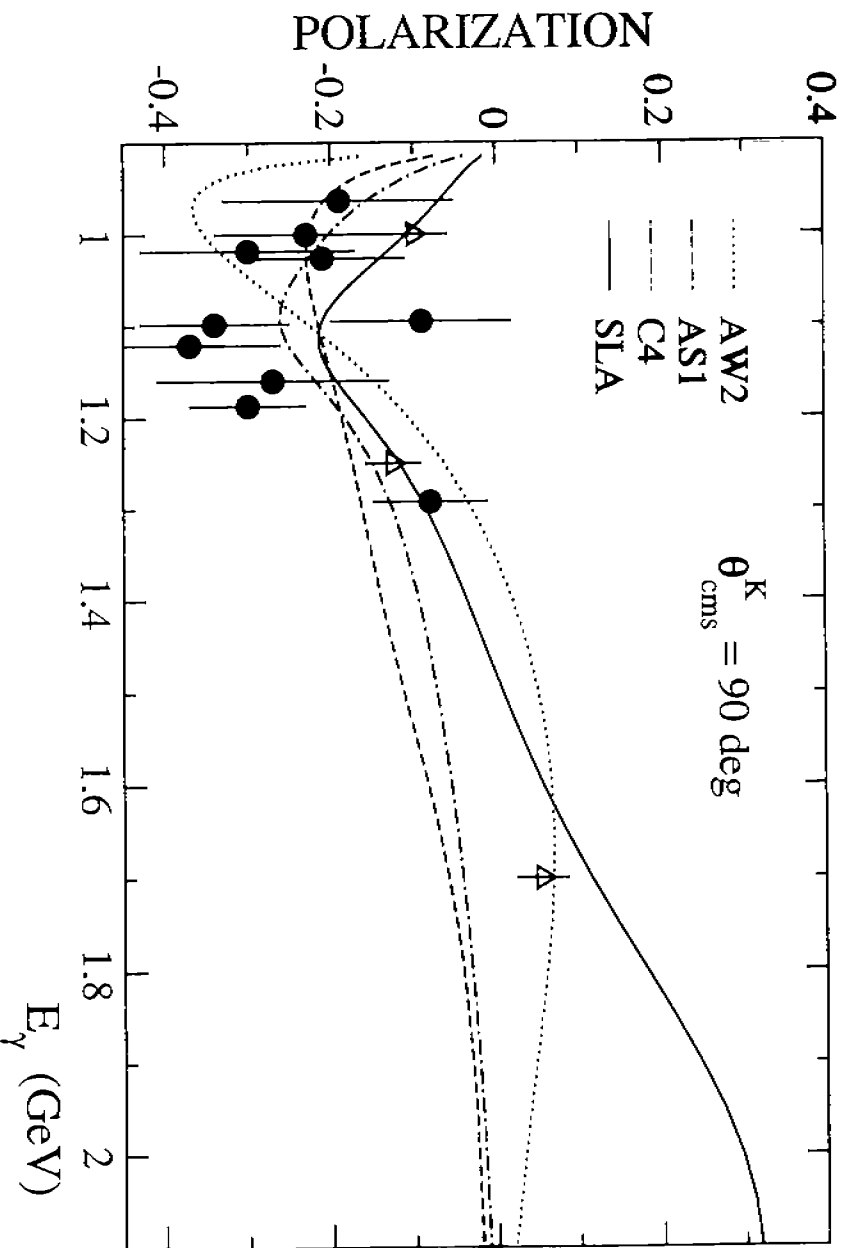
SLA: Mizutani-Fayard-Lamot-Saghai (1998)

(cf. Bennhold-Mart)

@ Spin-nonflip term very small, Spin-flip 3 terms dominate.

@ Momentum transfer: similar to $q(\pi^+, K^+)$. → unnatural parity high-spin states

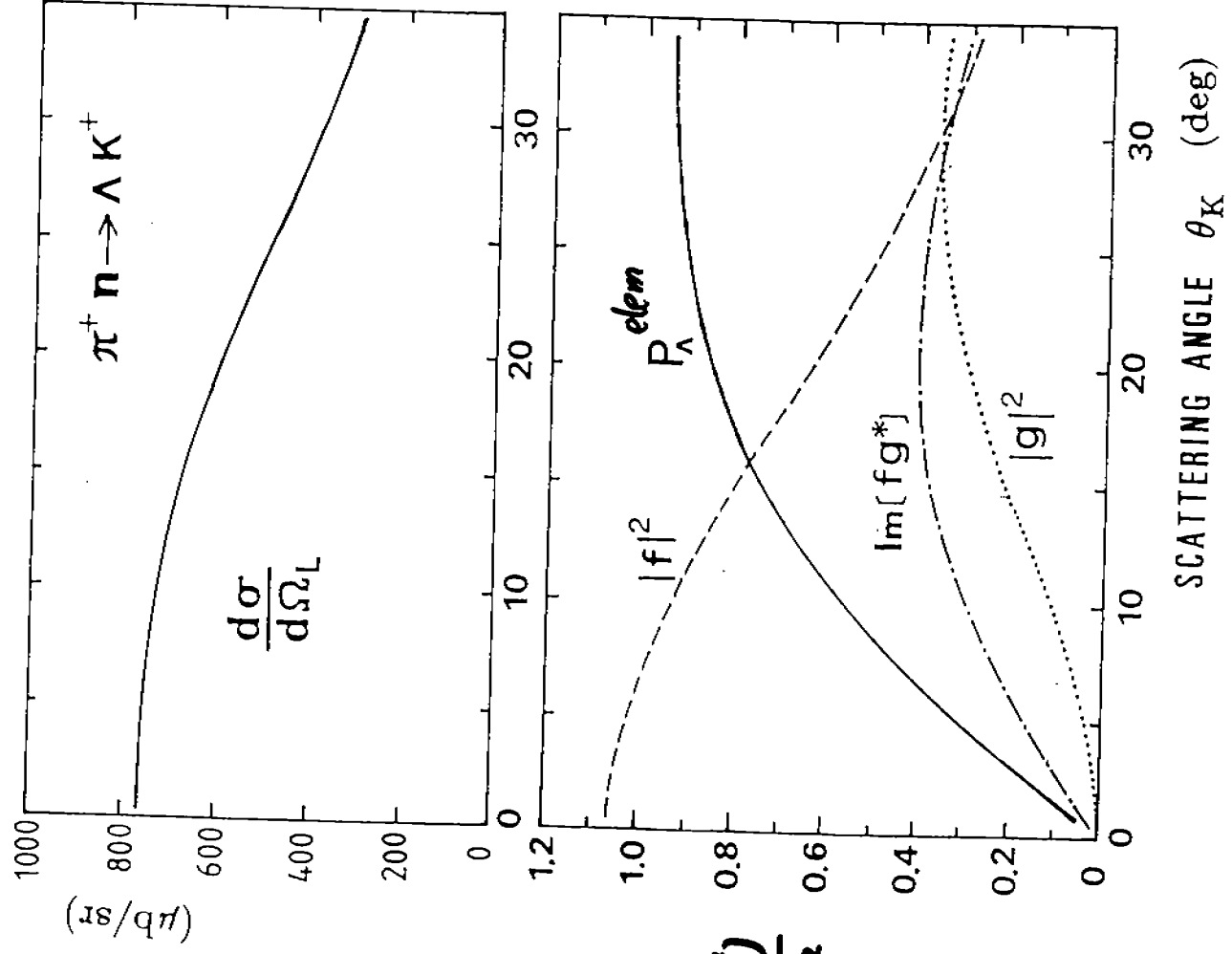
@ Appreciable difference among models
esp. for $d\sigma/d\Omega$ and $\text{pol}(\Lambda)$.





$$f + ig(\sigma \cdot n)$$

$\frac{d\sigma}{d\Omega} = 761 \quad 753 \quad 724 \quad 667 \quad 593 \quad 486 \quad 385 \quad 273$



$$P_\Lambda^{elem} = \frac{2\text{Im}(fg^*)}{|f|^2 + |g|^2}$$

$P_\Lambda = 0.32 \quad 0.57 \quad 0.75 \quad 0.85 \quad 0.91 \quad 0.94$

Fig. 8

For the hypernuclear production reaction on a nuclear target, $\gamma + A(J_i) \rightarrow H(J_f) + K^+$, the differential cross section in the lab frame is expressed analogously by

$$\left. \frac{d\sigma}{d\Omega} \right|_{A\text{-Lab}} = \frac{(2\pi)^4 p^2 E_K E_\gamma E_H}{k \{ p(E_H + E_K) - k E_K \cos \theta_L \}} |\mathbf{T}_{if}^{\text{Lab}}|^2, \quad (2.6)$$

$$|\mathbf{T}_{if}^{\text{Lab}}|^2 = \sum_{M_f} \mathbf{R}(J_i J_f; M_f), \quad (2.7)$$

$$\mathbf{R}(J_i J_f; M_f) = \frac{1}{2J_i + 1} \sum_{\epsilon, M_i} \left| \langle J_f M_f | \sigma_{\gamma K}(\theta) | J_i M_i \rangle \right|^2. \quad (2.8)$$

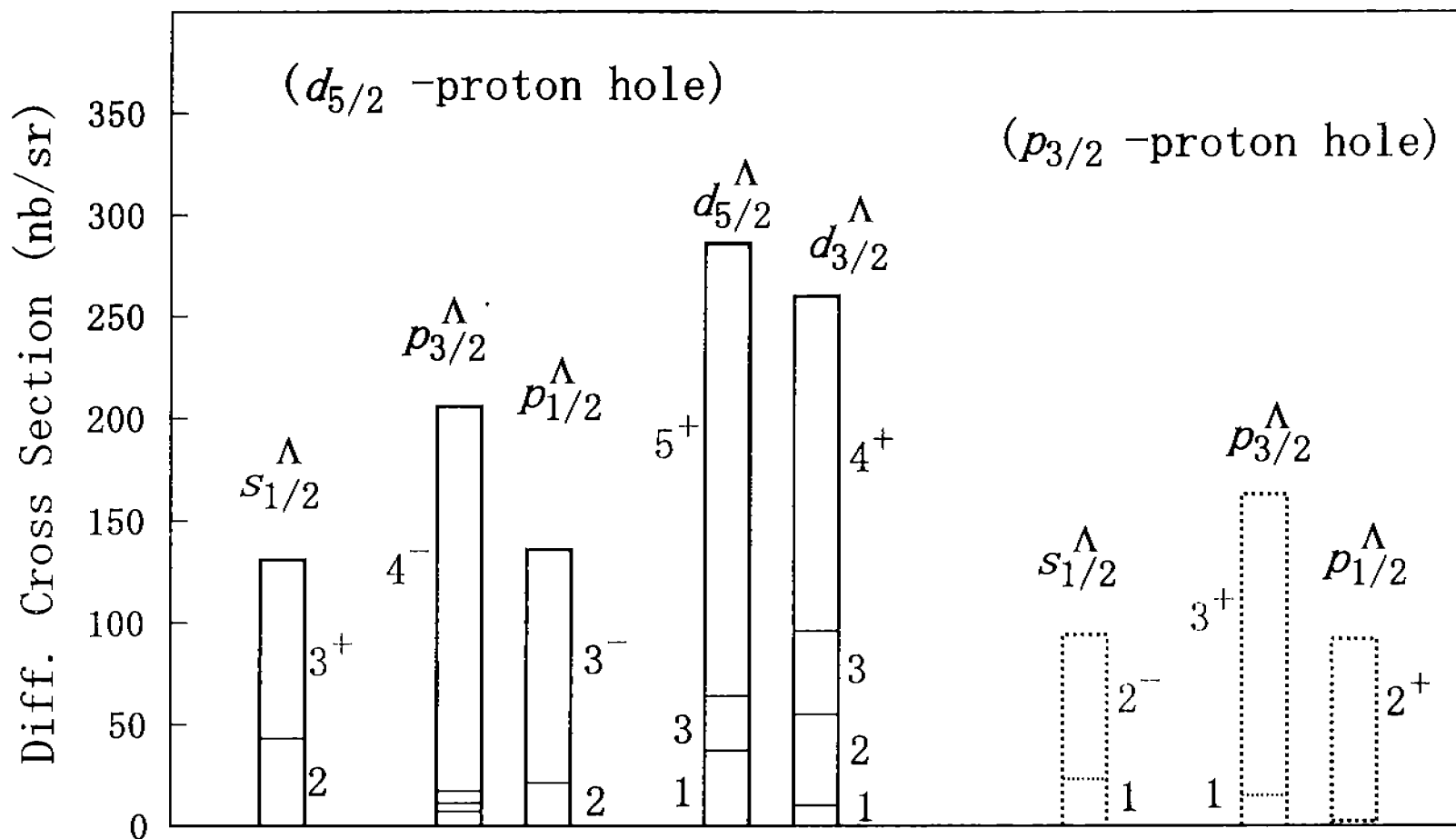
$$\sigma_{\gamma K}(\theta) = \int d^3 r \chi_K^{(-)*}(\mathbf{p}, \xi \mathbf{r}) \chi_\gamma^{(+)}(\mathbf{k}, \mathbf{r}) \sum_{\nu=1}^A V_-(\nu) \delta(\mathbf{r} - \eta \mathbf{r}_\nu) \langle \mathbf{k} - \mathbf{p}, \mathbf{p} | t | \mathbf{k}, 0 \rangle_L. \quad (2.9)$$

$$\mathcal{P}_H(J_f) = \sum_{M_f} \frac{M_f}{J_f} \frac{\mathbf{R}(J_i J_f; M_f)}{\tilde{\mathbf{R}}(J_i J_f)} \quad \text{with} \quad \tilde{\mathbf{R}}(J_i J_f) = \sum_{M_f} \mathbf{R}(J_i J_f; M_f). \quad (2.10)$$

$\chi_K^{(\pm)}$: Distorted wave \leftarrow Klein-Gordon Eq.

Contents of $^{28}\text{Si}(\gamma, \text{K}^+)$ spectrum

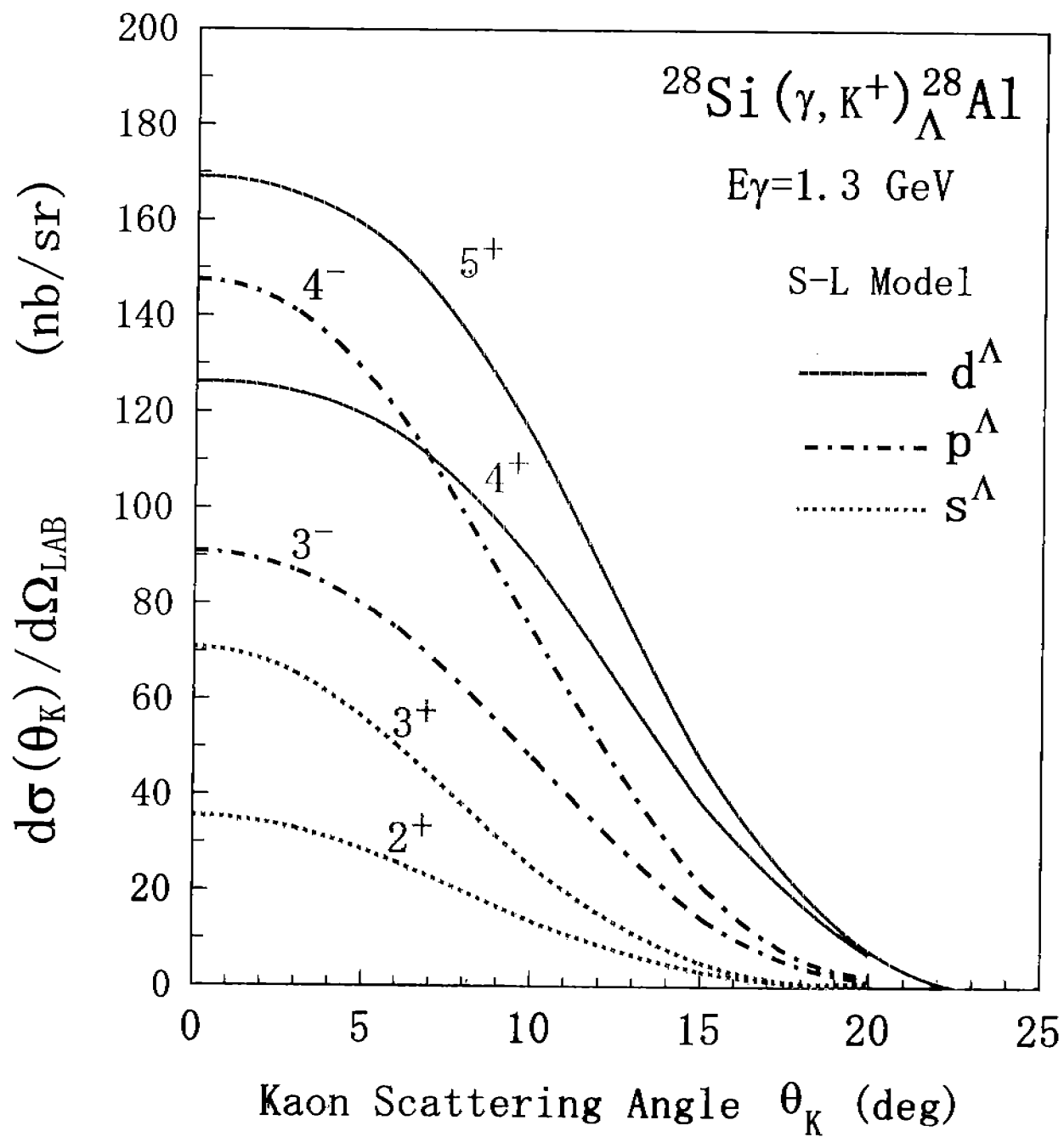
$(d_{5/2})^{12}$ MODEL



Selective Excitation

$$[j_{>}^{-1} j_{>}^{\Lambda}]: J_{\max} = j_{>} + j_{>}^{\Lambda} = l_{\text{N}} + l_{\Lambda} + 1 = L_{\max} + 1 \quad \text{unnatural parity}$$

$$[j_{>}^{-1} j_{<}^{\Lambda}] \text{ or } [j_{<}^{-1} j_{>}^{\Lambda}]: J_{\max} = l_{\text{N}} + l_{\Lambda} = L_{\max} \quad \text{natural parity}$$

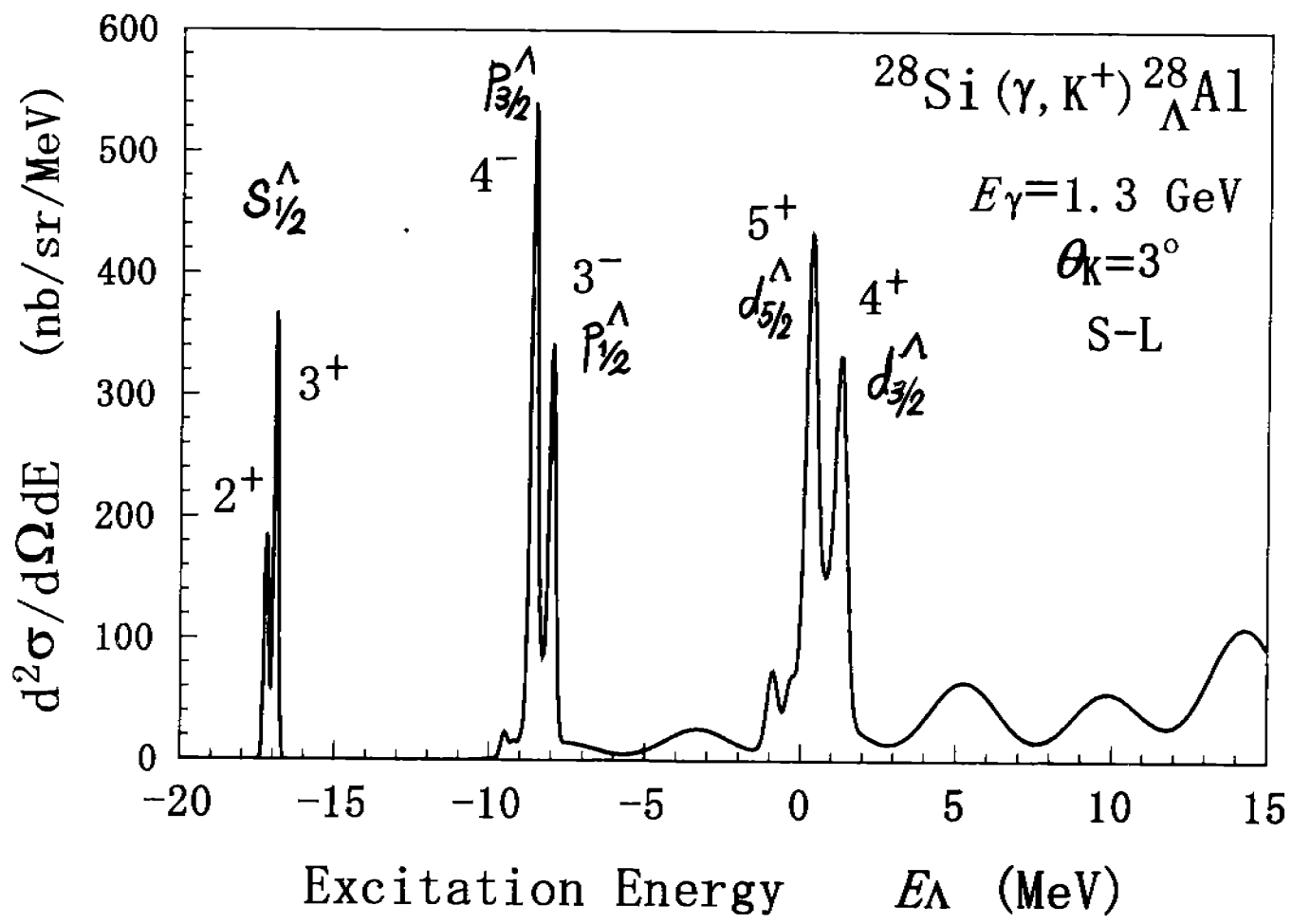


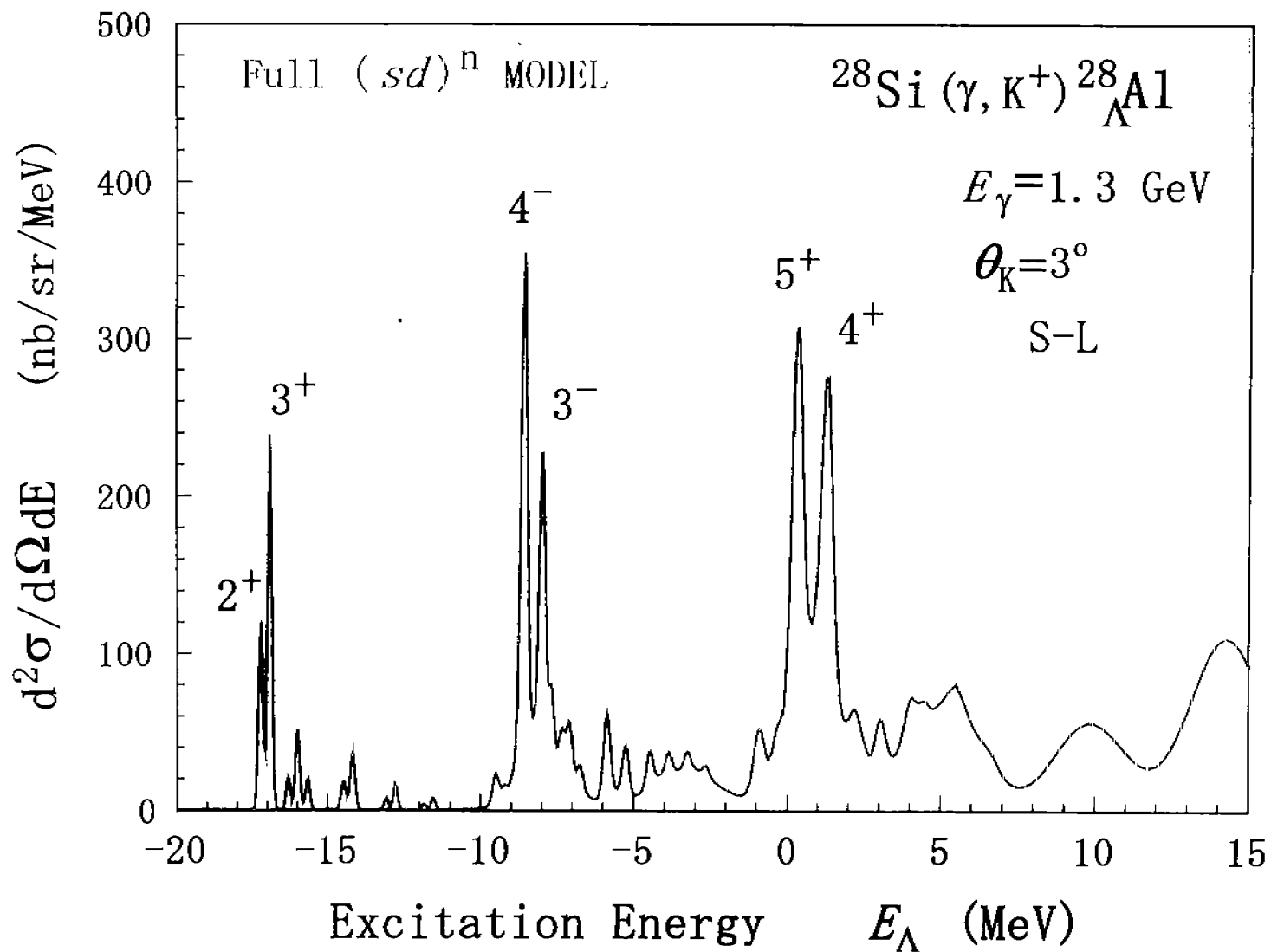
$(sp)^{16} (d_{5/2})^{12}$

Model

S-L Model

— d^{Λ}
 - - - p^{Λ}
 ···· s^{Λ}





- 1) A series of major peaks in the $(d_{5/2})^{12}$ model persist.
- 2) Their cross sections are approx. $0.65 \times$ ($d_{5/2}$ model).
- 3) New substates are expected based on ^{27}Al -core excitation.

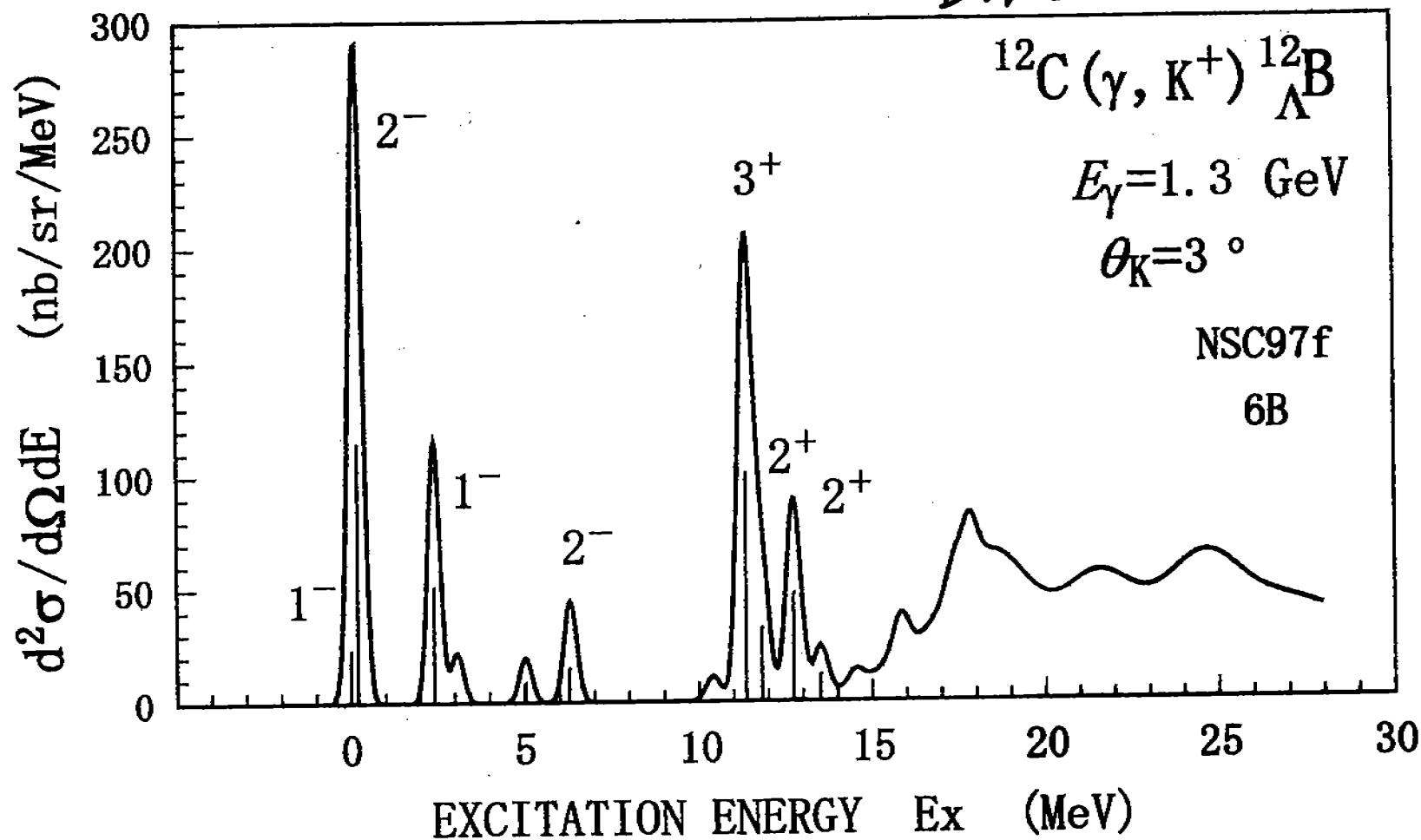
Photo/Electro-production of p-shell hypernuclei:

${}^7\text{Li}(\gamma, K^+) {}^7_{\Lambda}\text{He}$	$T=1$	n-halo
${}^8\text{Be}(\gamma, K^+) {}^8_{\Lambda}\text{Li}$	$T=1$	
${}^{10}\text{B}(\gamma, K^+) {}^{10}_{\Lambda}\text{Be}$		$T=1/2$ mirror of ${}^{10}_{\Lambda}\text{B}$ (π^+K^+)
${}^{11}\text{B}(\gamma, K^+) {}^{11}_{\Lambda}\text{Be}$	$T=1$	
${}^{12}\text{C}(\gamma, K^+) {}^{12}_{\Lambda}\text{B}$	<u>EXP</u>	$T=1/2$ mirror of ${}^{12}_{\Lambda}\text{C}$
${}^{13}\text{C}(\gamma, K^+) {}^{13}_{\Lambda}\text{B}$	$T=1$	
${}^{14}\text{N}(\gamma, K^+) {}^{14}_{\Lambda}\text{C}$		$T=1/2$ mirror of ${}^{14}_{\Lambda}\text{N}$
${}^{16}\text{O}(\gamma, K^+) {}^{16}_{\Lambda}\text{N}$		$T=1/2$ mirror of ${}^{16}_{\Lambda}\text{O}$

Interests:

- proton-deficient hyp. $\{ T=1$ (n-rich)
 $T=1/2$ (mirror)
- Λ dynamical coupling
with core-excited states
- unified knowledge of p-h multiplets
 \rightarrow detailed information of ${}^{21}\text{N}$
- structure change of light systems
(such as shrinkage)

DWIA

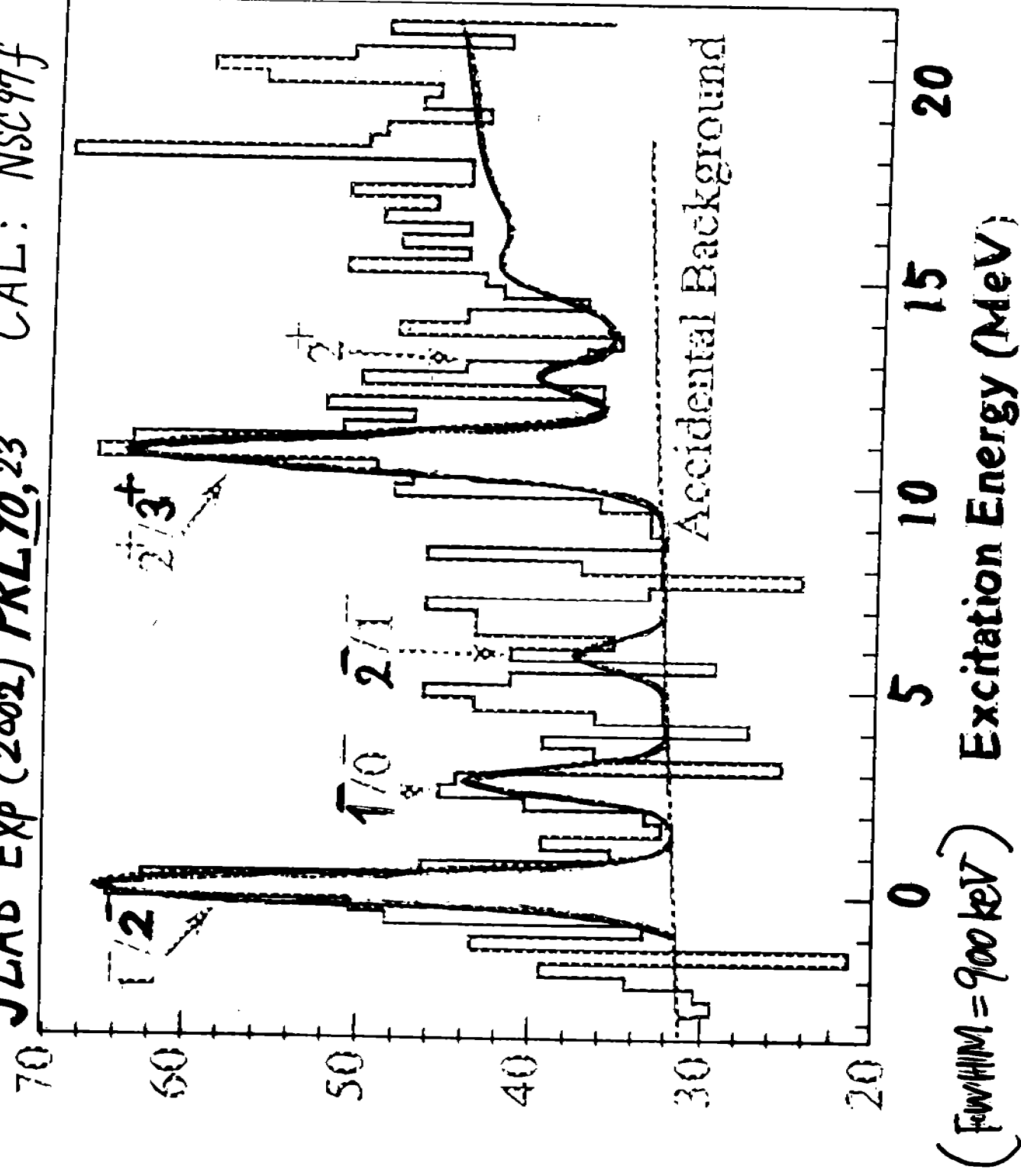


Motoba-Sotona-Itonaga, Prog.Theor.Phys. Suppl. 119, 127 (1994)
Motoba, Mesons & Light Nuclei (2000) updated.

C12gk6B97f

n below ~ 12 MeV, but statistics are insufficient
ly resolve the structure.

JLAB Exp (2002) PRL 90, 23 (2003) CAL: NSC97f



3: The summed ^{12}B missing mass spectrum for bo
ies

$$[p_{3/2}^{-1} p_{3/2}^{\Lambda}]_J$$

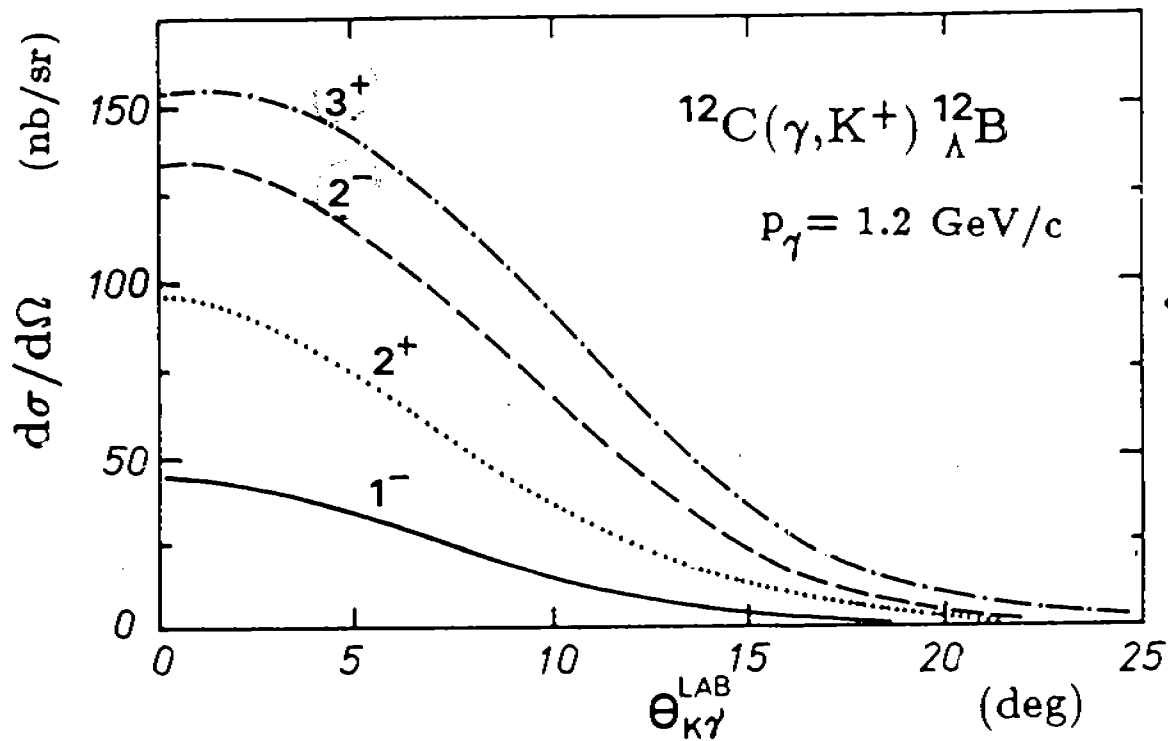
$$J=0^+ (K^-, \pi^-)$$

$$1^+ (K^-, \pi^-) \theta \gg 1$$

$$2^+ (\pi^+, K^+)$$

$$3^+ (\gamma, K^+)$$

Strangeness nuclear physics



$$3^+ [p_{3/2}^{-1} p_{3/2}^{\Lambda}]$$

$$|g_1|^2 - |g_{-1}|^2$$

$$2^- [p_{3/2}^{-1} s_{1/2}^{\Lambda}]$$

Fig. 4.2. Calculated (γ, K^+) cross sections for the $[0p_{3/2}^{-1} 0s_{1/2}^{\Lambda}]_{J=1^-, 2^-}$ and $[0p_{3/2}^{-1} 0p_{3/2}^{\Lambda}]_{J=2^+, 3^+}$ shell model states describing ${}_{\Lambda}^{12}\text{B}$.

neutron
halo

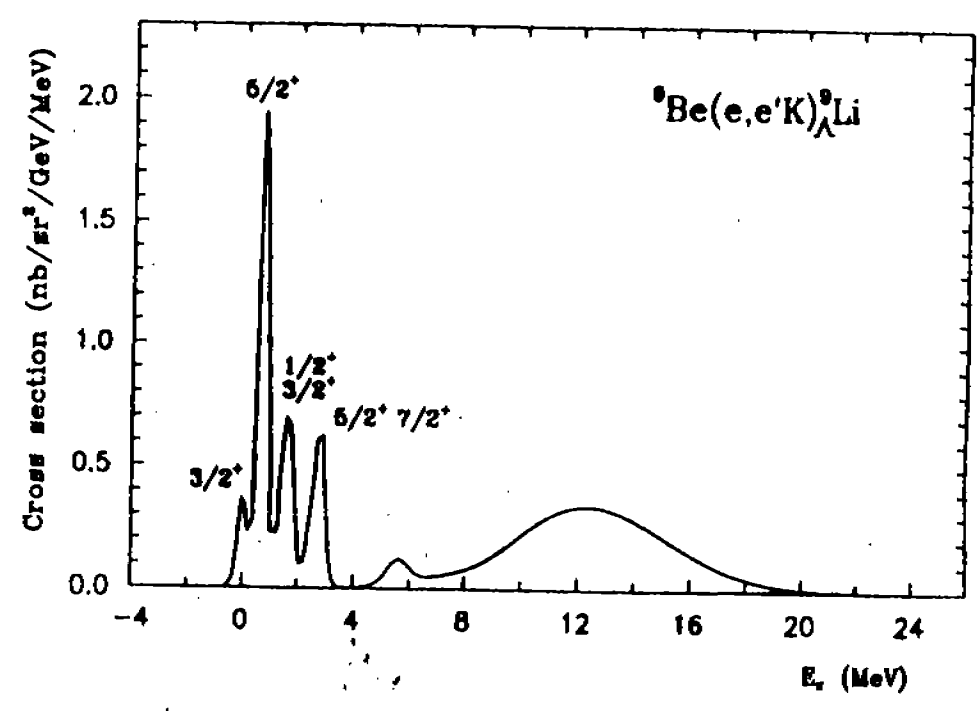
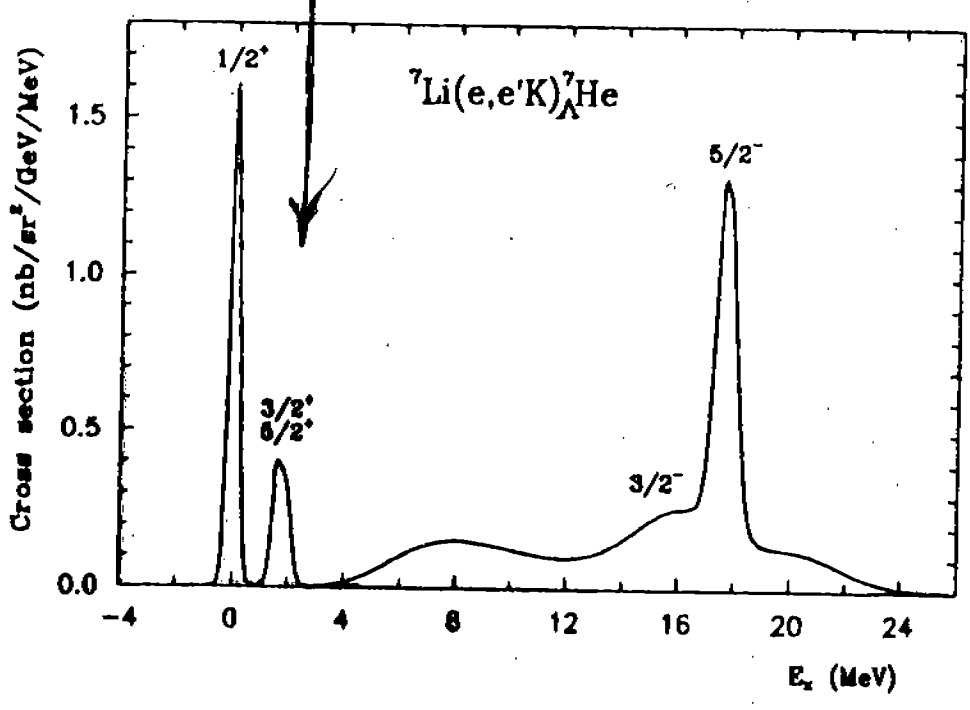
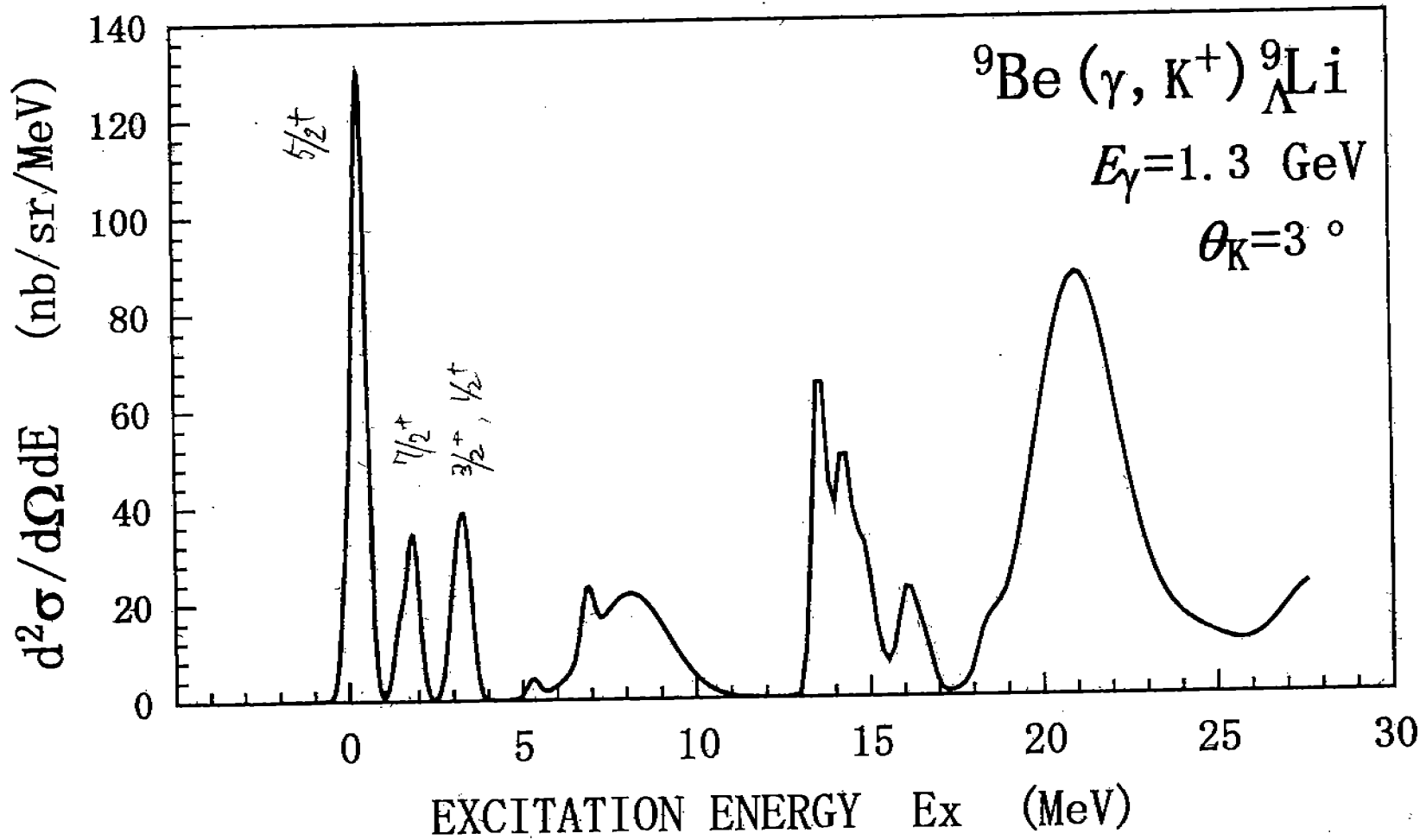


Figure 2: Hypernuclear excitation functions at $E_e = 4.0$ GeV, $E'_e = 1.8$ GeV, $\theta_e = \theta_{Ke} = 6^\circ$ for ${}^7\text{Li}$ target (left) and ${}^9\text{Be}$ target (right).



← Λ

94-02-23

DASHED (J-) FULL (J+)

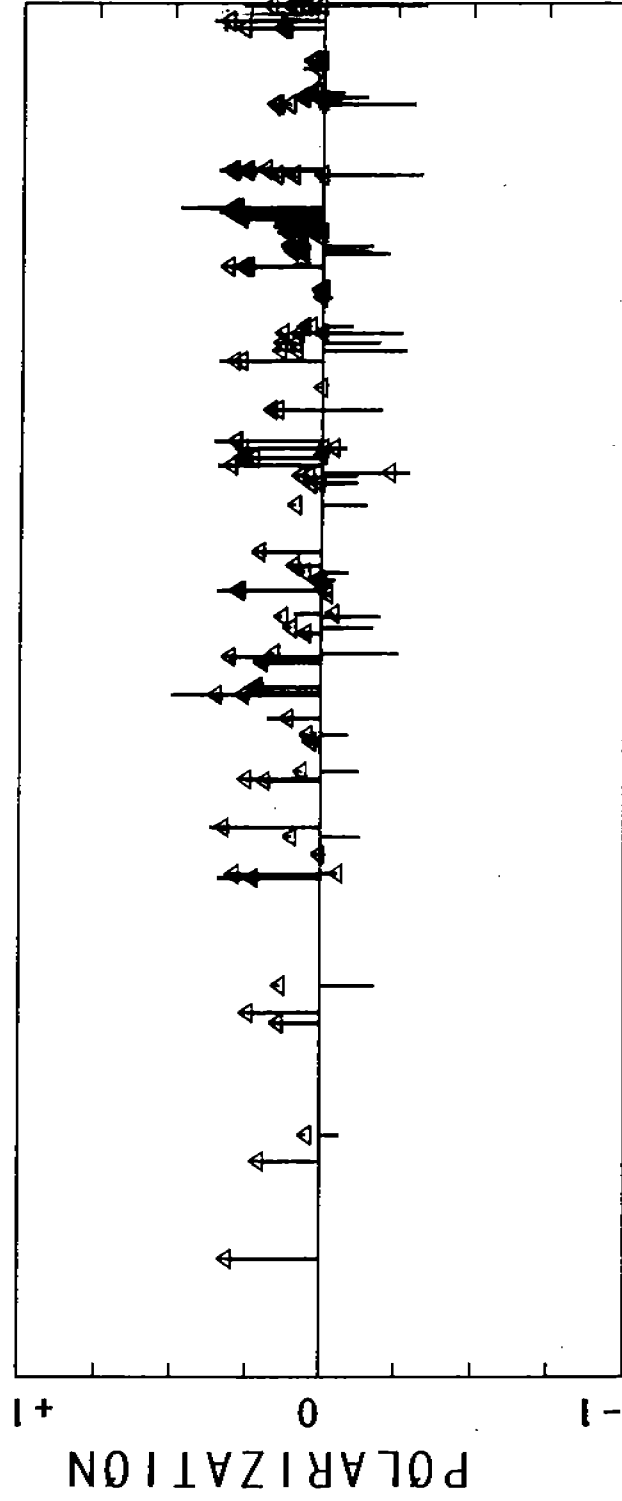
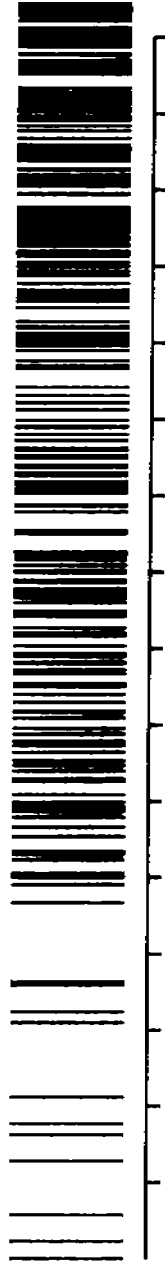
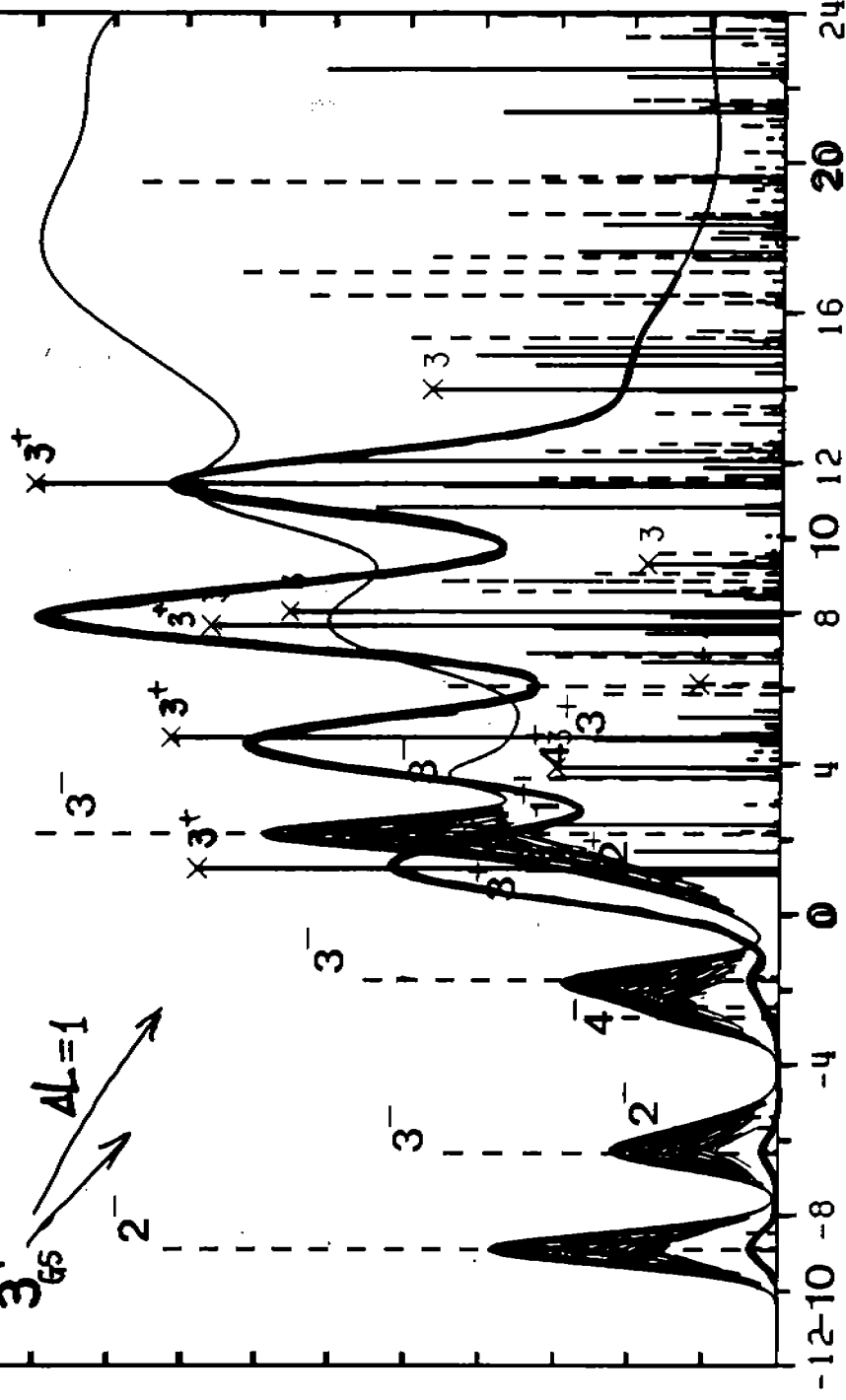
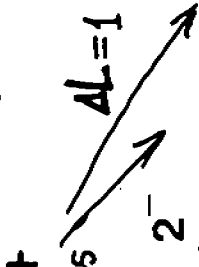
$(K, \pi) \theta \sim 0^\circ$

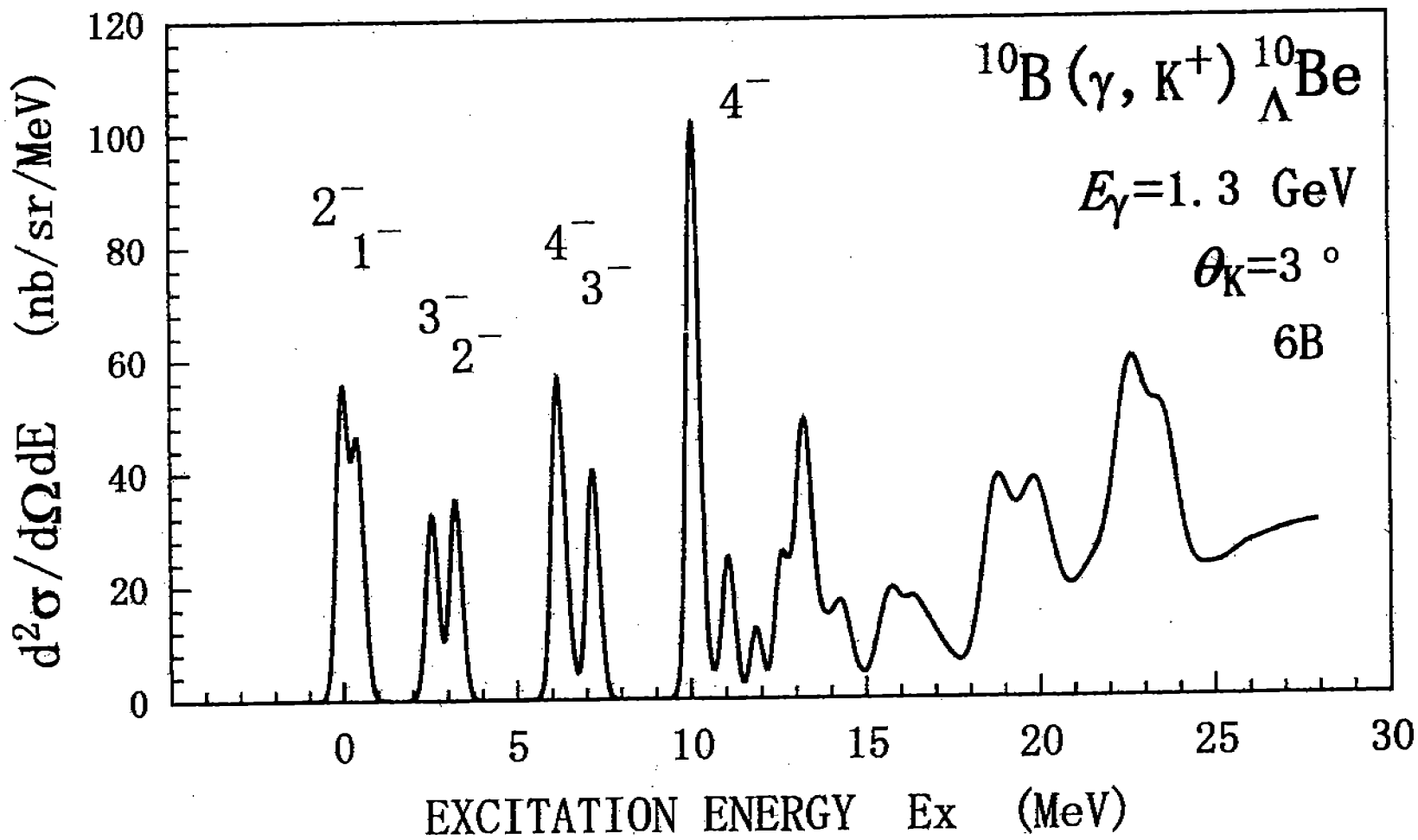
$3^+_{gs} \xrightarrow{\Delta L=0} 3^+$

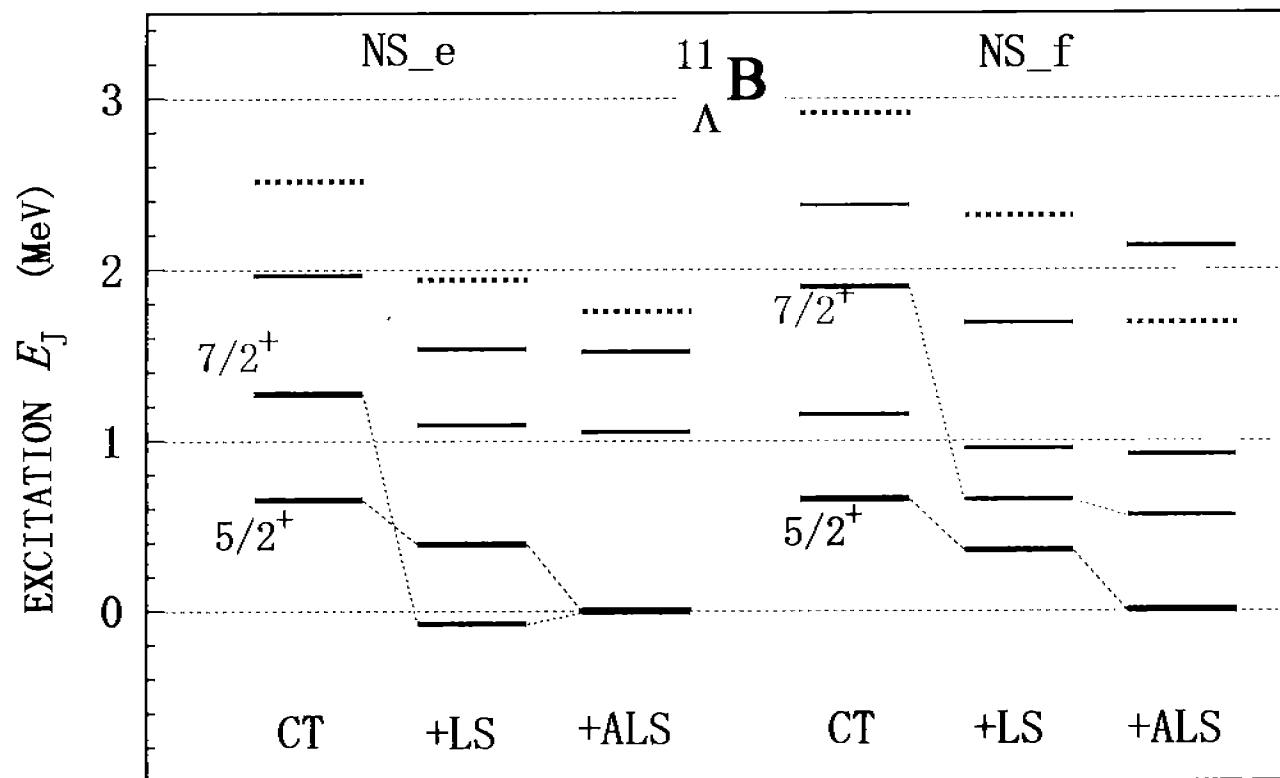
REL. $d\sigma/d\Omega$ (MAX-321144 nb/sr)

$^{10}\text{B} (K^\pi, K^\pi)_{\Lambda}^{10}\text{B}$

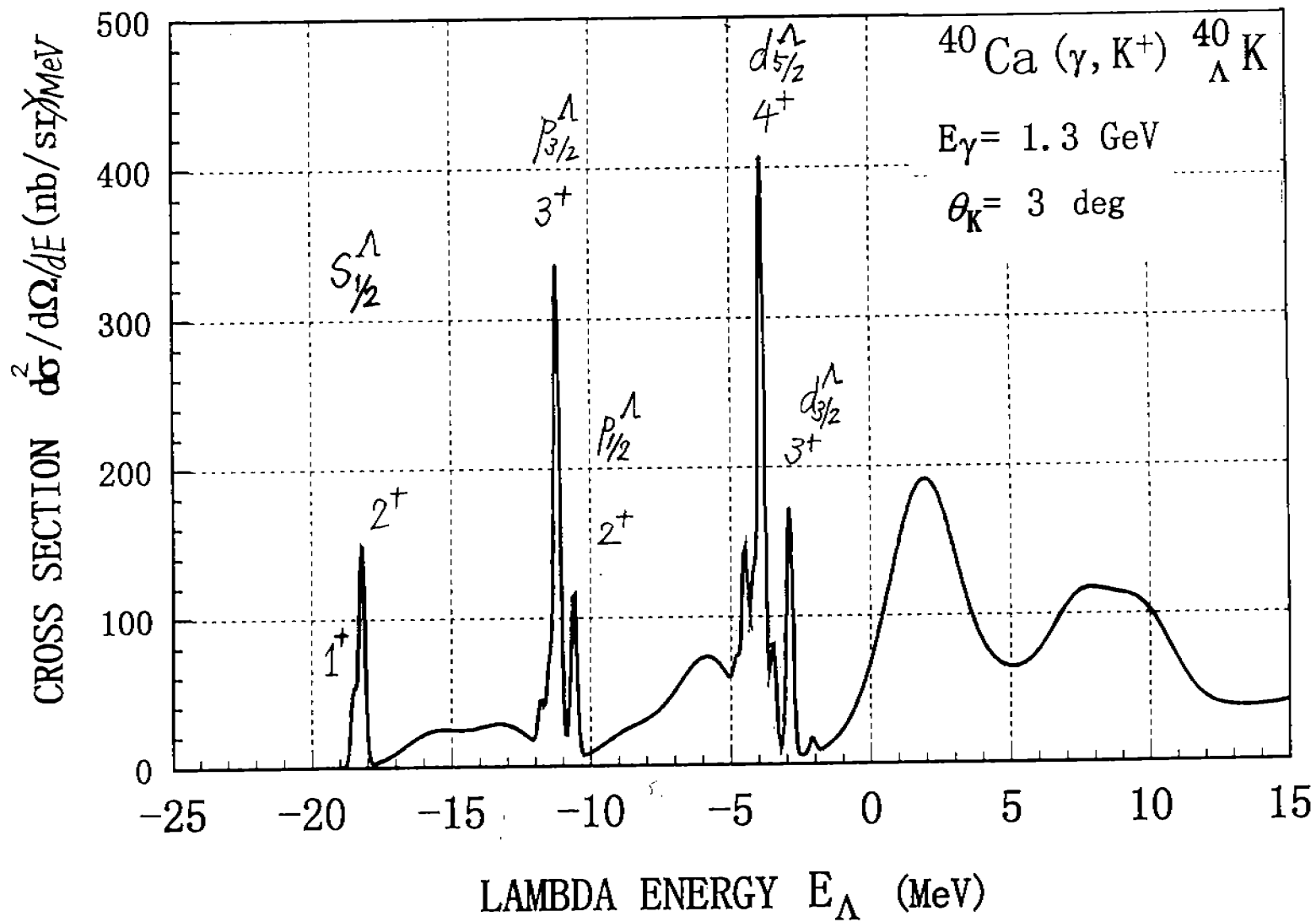
$p_K=0.65 \text{ GeV}/c \quad \theta=0000^\circ$



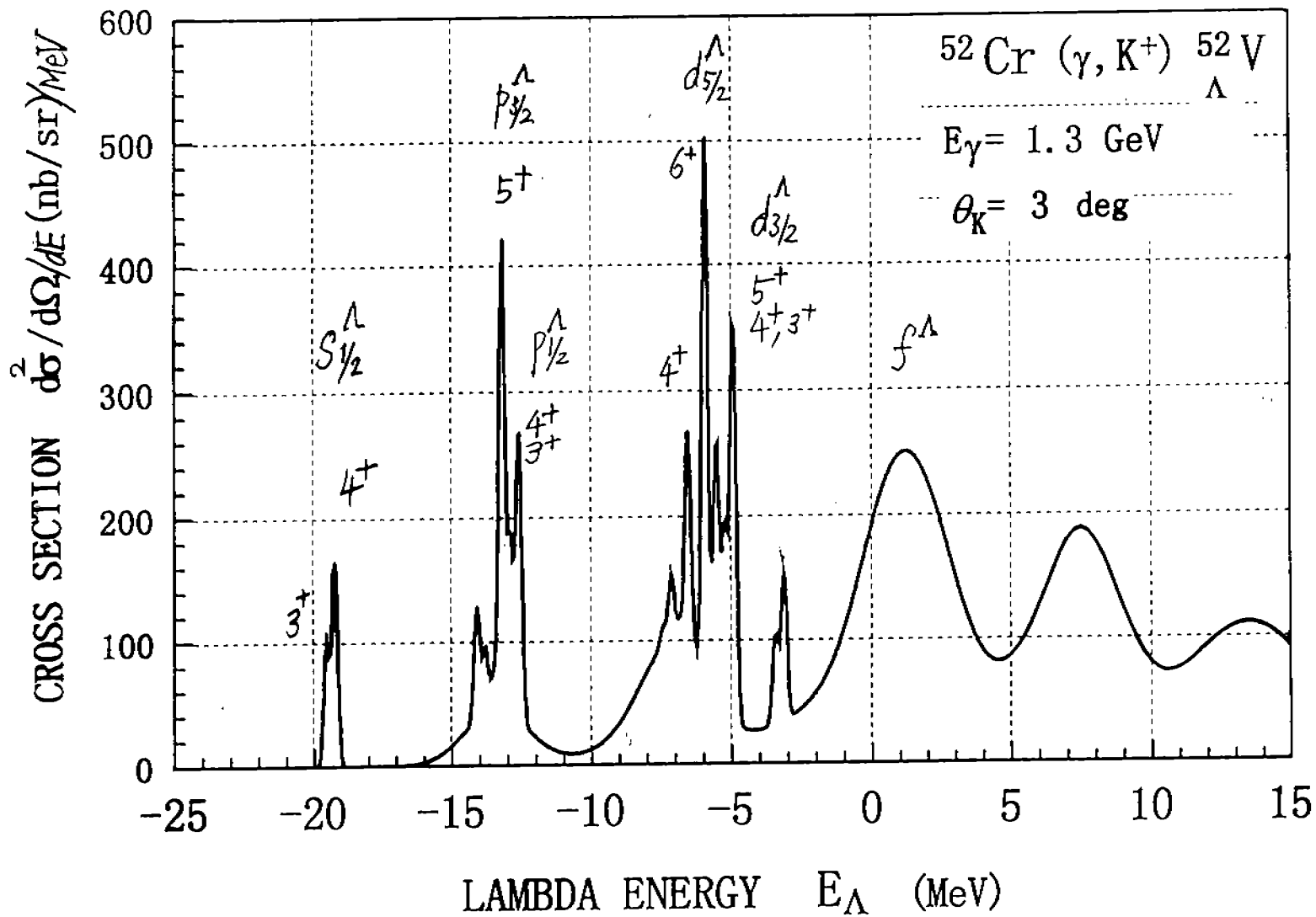




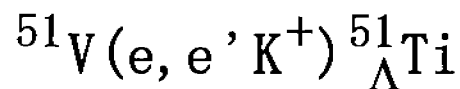
B11_CTL



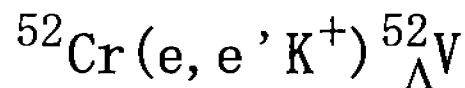
A40gs7



Production of $A \sim 50$ hypernuclei



$$E_e = 4.056 \text{ GeV}, \quad E_{e'} = 1.856 \text{ GeV},$$

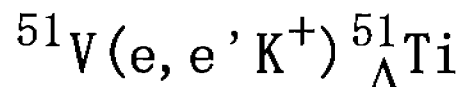


$$\theta_e = 6 \text{ deg.} \quad \rightarrow \text{almost real photon}$$

${}^{48}\text{Ca}$ core assumed.

$$\text{Configuration: } [f_{7/2}^n] + [f_{7/2}^{n-1} p_{3/2}] + [f_{7/2}^{n-1} f_{5/2}]$$

K. Lpis and M. T. McEllistrem, P. R. C1 1009 (1970)



Triple diff. x-section (nb/sr²/GeV)

Λ	E_x	J_f	SL Model	BM Model
0s	0.0	1/2+	1.298	0.849
0p	8.014	3/2-	1.577	1.462
	8.153	1/2-	2.488	0.877
0d	16.575	5/2+	1.378	1.712
	16.840	3/2+	2.468	0.410

dominantly

$${}^{42}\text{Ti}(0^+) \times j_{\Lambda}$$

SUMMARY

- 1) Basic properties of the $\gamma p \rightarrow \Lambda K^+$ process are discussed with typical models for the amplitudes:
 - spin-flip dominance, large q ,
 - appreciable difference among theor. models (P).
- 2) Novel characteristics of photo/electroproduction of hypernuclei are clarified with the ^{28}Si target: $(d_{5/2})^{12}$ Model
 - $[\vec{j}_>^{-1} j_>^{\Lambda}]$: $J_{\max} = j_> + j_>^{\Lambda} = l_N + l_{\Lambda} + 1 = L_{\max} + 1$ unnatural parity
 - $[\vec{j}_>^{-1} j_<^{\Lambda}]$, $[\vec{j}_<^{-1} j_>^{\Lambda}]$: $J'_{\max} = l_N + l_{\Lambda} = L_{\max}$ orbitally stretched

Selective excitation of high-spin states, providing a series of Λ orbits.
- 3) A realistic calculation with the full $(sd)^n$ w.f. has been done for $^{28}\text{Si}(\gamma, K^+)_{\Lambda}^{28}\text{Al}$, which should be compared with the coming exp. at JLab.
- 4) Present theoretical framework has been confirmed by the good agreement with the first hyp. data from Jlab.: $^{12}\text{C}(e, e' K^+)_{\Lambda}^{12}\text{B}$.
 Show possibility of producing "p-deficient" p-shell hypernuclei.
- 5) Extend the approach to produce heavier hypernuclei $A \sim 50$, hoping to disclose new areas.