

Beam-Beam Interaction in Linac-Ring Colliders

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Talk Outline

- Beam-Beam Estimates
- Head-Tail Instability
- Luminosity-Deflection Theorem
- Simulation Methods
 - Coulomb Sum (Beard, Li)
 - PIC (Shi)
- Simulation Status
- Future needs for a complete solution
- Conclusions



Luminosity Factors

$$\mathcal{L} = f \frac{N_e N_i}{2\pi \sqrt{\sigma_{e,x}^2 + \sigma_{i,x}^2} \sqrt{\sigma_{e,y}^2 + \sigma_{i,y}^2}}$$

Ion Tune Shifts:

$$\xi_{i,x} = \beta_{i,x}^* \frac{N_e r_i}{2\pi \gamma_i \sigma_{e,x} (\sigma_{e,x} + \sigma_{e,y})} \quad \xi_{i,y} = \beta_{i,y}^* \frac{N_e r_i}{2\pi \gamma_i \sigma_{e,y} (\sigma_{e,x} + \sigma_{e,y})}$$

“Equivalent” Electron Tune Shifts:

$$\xi_{e,x} = \beta_{e,x}^* \frac{N_i r_e}{2\pi \gamma_e \sigma_{i,x} (\sigma_{i,x} + \sigma_{i,y})} \quad \xi_{e,y} = \beta_{e,y}^* \frac{N_i r_e}{2\pi \gamma_e \sigma_{i,y} (\sigma_{i,x} + \sigma_{i,y})}$$



Gedanken Experiment

For round, equal sized beams, the following scaling applies:

$$\mathcal{L} = \frac{I_e}{e} \frac{\gamma_e \xi_e}{\beta^* r_e}$$

Comparing linac-ring colliders and ring-ring colliders, what can change for the better?

1. Maximum I_e/e is set by ION ring stability. The same in the two cases
2. γ_e set by the physics. The same in the two cases
3. Minimum β^* is set by IR region design issues. Can it be too much better for linac-ring? Should not be any worse than for ring-ring
4. r_e is set by (God, Yahweh, Allah, ...); YOU cannot change it
5. If there are to be luminosity enhancements to be found for linac-ring designs compared to ring-ring designs, they must arise because one is allowed to make the equivalent tune shift ξ_e bigger for linac-ring colliders.
6. Finding the physical phenomena that determine the maximum ξ_e are extremely important for evaluating the linac-ring idea.



Two guesses

1. *Emittance growth generated by a single beam-beam collision. Round Gaussian beam collision integrals can be performed to give*

$$\varepsilon_{n,after}^2 = \varepsilon_{n,before}^2 + (0.194N_p)^2$$

What's the right scaling for circulator ring? If proportional to the number of turns squared, may have a significant problem recovering the beam with small loss. Halo for CR?

2. *Fast Head-Tail instability; Linear Stability Estimate (Lebedev, Yunn, Li). Assume short electron bunch*

$$W_i(s) = \frac{r_e N_e}{\gamma_e \sigma_i^2 \sigma_e^2} s$$



Fast Head-Tail instability

Threshold condition

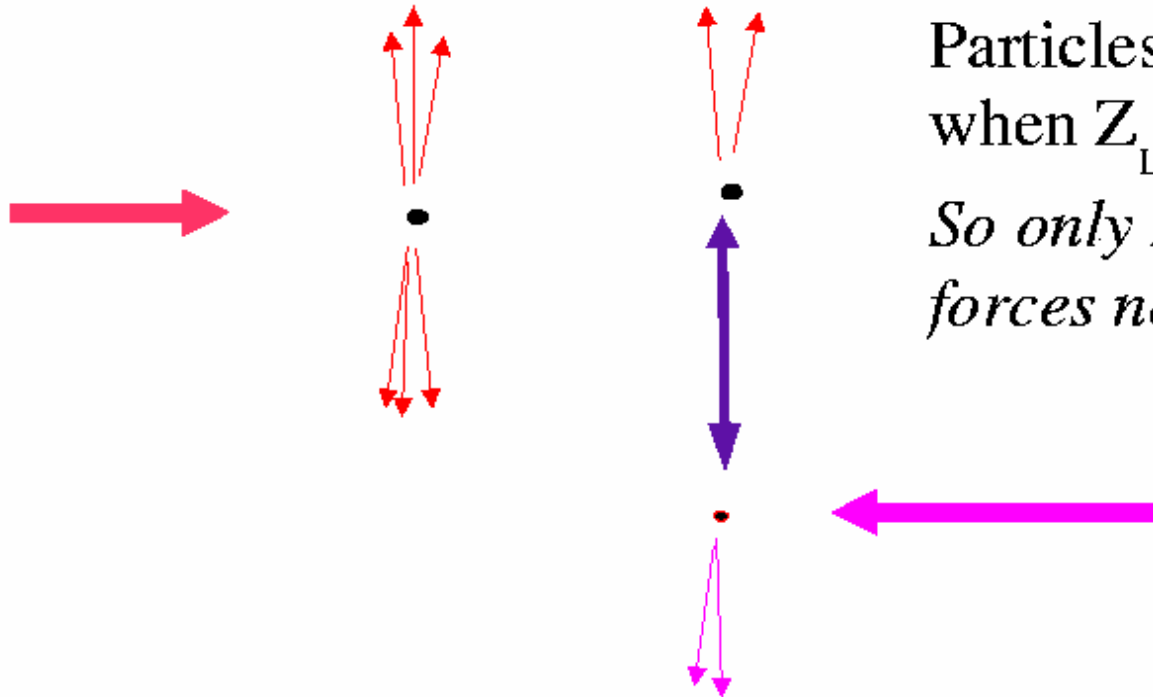
$$\xi_e \xi_p \leq \frac{\beta_e v_s}{\pi^2 \sigma_z}$$

- Larger synchrotron frequency helps
- Different instability than in a ring because the fresh electrons come in at a fixed transverse position, without their own distortion. Makes an “impedance” model reasonable first approximation
- Full calculation needs to account for non-linear effects and the fact that the electron bunch is no longer short
- See Li, et. al, PAC2001, 2014 for a linear “long-bunch” theory, including synchro-betatron coupling



High Energy Simplification

Simplifications... $v \sim c$



Particles only interact
when $Z_L = Z_R$
*So only 2 dimensional
forces need be considered*



Luminosity-Deflection Theorem

Theorem from 2-D electrostatics

$$\vec{F}'_{21}(\vec{b}) = -\vec{F}'_{12} = 2Q'_1Q'_2 \int n_2(\vec{x}_2) \frac{\vec{x}_2 + \vec{b} - \vec{x}_1}{|\vec{x}_2 + \vec{b} - \vec{x}_1|^2} n_1(\vec{x}_1) d^2\vec{x}_1 d^2\vec{x}_2$$

$$\Rightarrow \nabla_{\vec{b}} \cdot \vec{F}'_{21}(\vec{b}) = 4\pi \int \rho_2(\vec{x}) \rho_1(\vec{x}) d^2\vec{x}$$

2-D Gauss's law generalized to transversely extended macroparticles

$$\vec{D}(\vec{b}) = \Delta\gamma_1\vec{\beta}_1 = -\Delta m_2\gamma_2\vec{\beta}_2 / m_1 = \frac{2q_1q_2}{m_1c^2} \int n_1(\vec{x}_1) \frac{\vec{x}_1 - \vec{x}_2 - \vec{b}}{|\vec{x}_1 - \vec{x}_2 - \vec{b}|^2} n_2(\vec{x}_2) d^2\vec{x}_1 d^2\vec{x}_2$$

$$\Rightarrow \nabla_{\vec{b}} \cdot \vec{D}(\vec{b}) = -\frac{4\pi q_1q_2}{m_1c^2} \int n_1(\vec{x}) n_2(\vec{x} - \vec{b}) d^2\vec{x} = \frac{4\pi r_e}{N_1} L(\vec{b})$$



Luminosity-Deflection Transform Pairs

Round Beam Fast Model

$$\vec{D}(\vec{b}) = \frac{2N_2 r_e \vec{b}}{\sigma^2 + b^2} \quad L(\vec{b}) = \frac{N_1 N_2 \sigma^2}{\pi(\sigma^2 + b^2)^2}$$

Gaussian Macroparticles

$$\vec{D}(\vec{b}) = \vec{D}_{\text{Bassetti-Erskine}}(\vec{b}; \sqrt{\sigma_{1x}^2 + \sigma_{2x}^2}; \sqrt{\sigma_{1y}^2 + \sigma_{2y}^2})$$

$$L(\vec{b}) = \frac{N_1 N_2}{2\pi \sqrt{\sigma_{1,x}^2 + \sigma_{2,x}^2} \sqrt{\sigma_{1,y}^2 + \sigma_{2,y}^2}} \exp\left(-\frac{b_x^2}{2\sqrt{\sigma_{1,x}^2 + \sigma_{2,x}^2}}\right) \exp\left(-\frac{b_y^2}{2\sqrt{\sigma_{1,y}^2 + \sigma_{2,y}^2}}\right)$$

Smith-Laslett Model (no complex error function!)

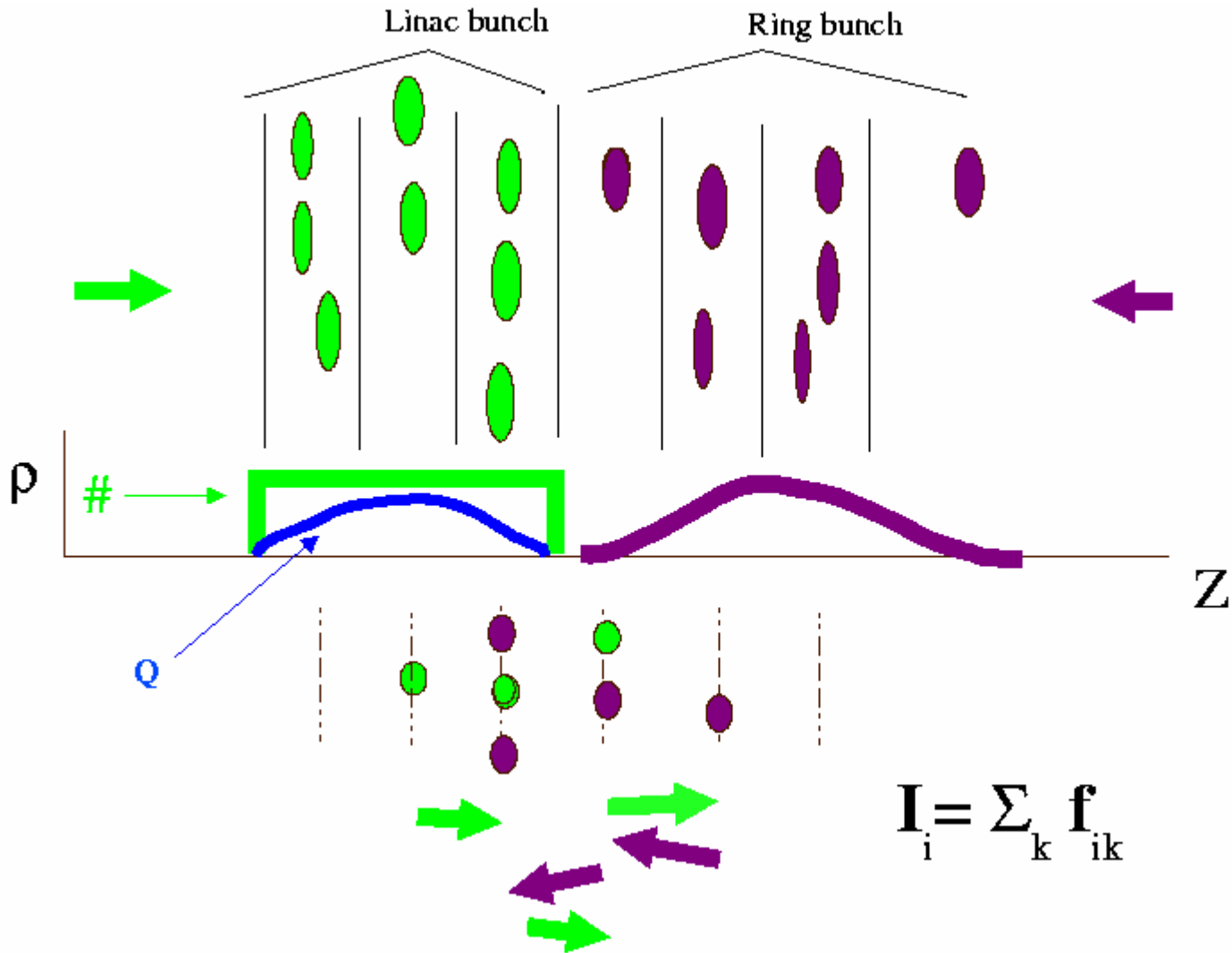
$$\vec{D}(\vec{b}) = \frac{2N_2 r_e \vec{b}}{\hat{b}^2 AB} \left\{ \frac{4\hat{b}^2 + 2\hat{b}^4}{(\hat{b}^4 + 4\hat{b}^2)} - \frac{4\hat{b}^2}{(\hat{b}^4 + 4\hat{b}^2)^{3/2}} \left\{ \sinh^{-1} \left[\frac{\hat{b}^3}{2} + \frac{3\hat{b}}{2} \right] + \sinh^{-1} \left[\frac{\hat{b}}{2} \right] \right\} \right\}$$

$$L(\vec{b}) = \frac{N_1 N_2}{\pi AB} \left\{ \frac{(2\hat{b}^2 - 4)\hat{b}^2}{(\hat{b}^4 + 4\hat{b}^2)^2} + \frac{4\hat{b}^2(1 + \hat{b}^2)}{(\hat{b}^4 + 4\hat{b}^2)^{5/2}} \left\{ \sinh^{-1} \left[\frac{\hat{b}^3}{2} + \frac{3\hat{b}}{2} \right] + \sinh^{-1} \left[\frac{\hat{b}}{2} \right] \right\} \right\}$$

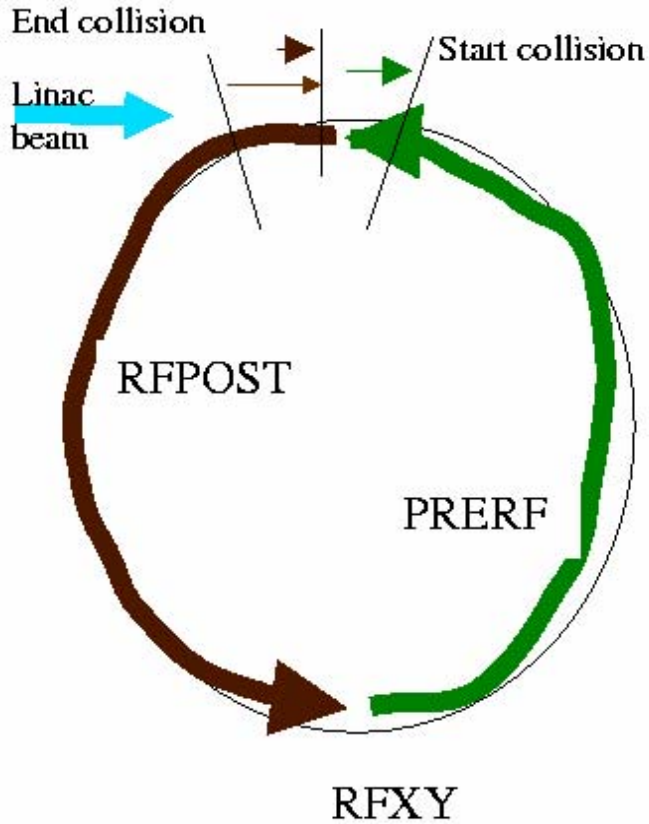
$$\hat{b}^2 = \left(\frac{b_x}{A}\right)^2 + \left(\frac{b_y}{B}\right)^2$$



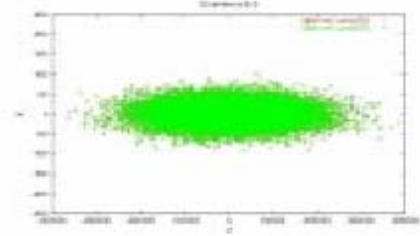
Coulomb Sum Simulation Approach



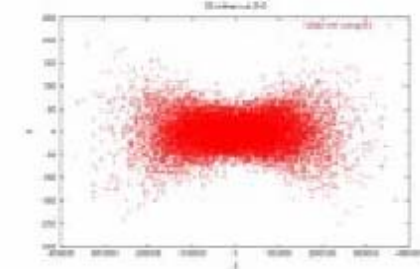
Matching to Reduce Mismatch Oscillations



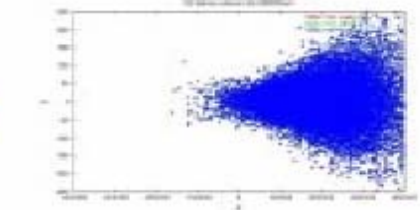
Drifted - in



Drifted - out



Before collision



Transverse motion



Particle-In-Cell (PIC) approaches

Collective Beam-Beam Effects In Hadron Colliders

Jack J. Shi

Department of Physics & Astronomy, University of Kansas

Collaboration:

University of Kansas

Lihui Jin

DESY, Cornell University

Georg Hoffstaetter

Michiko Minty

Thanks: DOE

http://casa.jlab.org/seminars/2004/slides/shi_040213.pdf



2. Direct multi-particle tracking: the beam-beam force is calculated with particles-to-particle individually.

— Precise if N_p is large, but very slow [$O(N_p^2)$],
typical: $N_p \leq 10^4 \implies$ wrong physics in the nonlinear regime.

3. Particle-In-Cell (PIC): evaluate beam-beam force on a mesh.

— Precise, but very slow for separated beams.

Variations:

a. Calculate Beam-Beam Potential Without Boundary

b. Calculate The Potential With Approximated Boundary

c. Directly Calculate Beam-Beam Force on the Mesh

d. With Weighted Functions

Courtesy: Jack Shi



3. Direct Calculation of Beam-Beam Field on the Mesh

The field is calculated with

$$\vec{K}(\vec{r}) = \int d\vec{r}' \rho(\vec{r}') \vec{G}_k(\vec{r} - \vec{r}')$$

where Green's function is

$$\vec{G}_k(\vec{r} - \vec{r}') = \frac{(\vec{r} - \vec{r}')}{(x - x')^2 + (y - y')^2}$$

Comment:

- Accurate — Exact boundary condition
No errors due to numerical derivatives.
- Only a small number of empty cells when using adaptive mesh.
- Slow when a large mesh has to be used (mis-matched beams)
— Computation cost $\sim N_p N_m^2$.

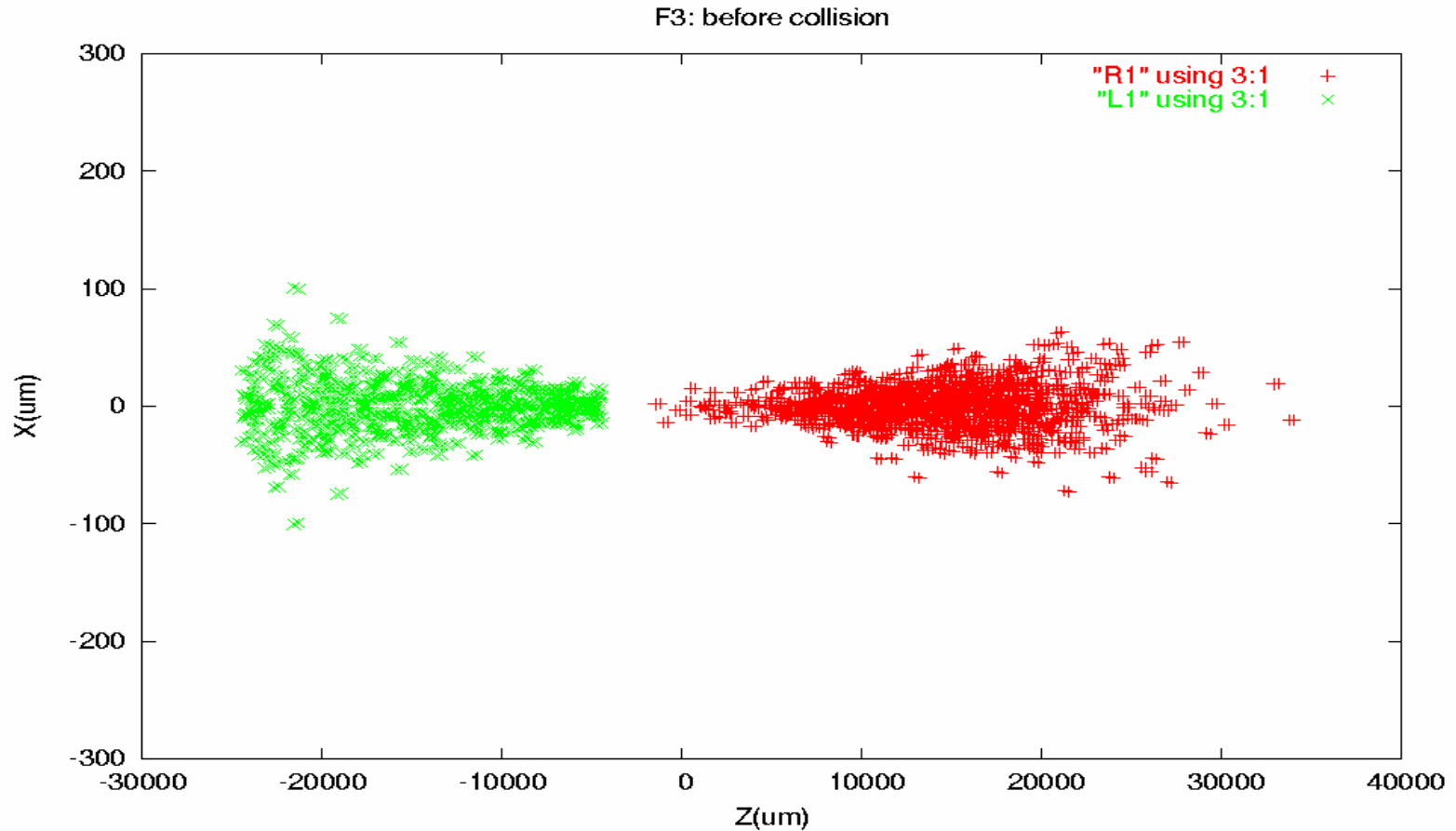


Coulomb Sum Simulation Results (K. Beard)

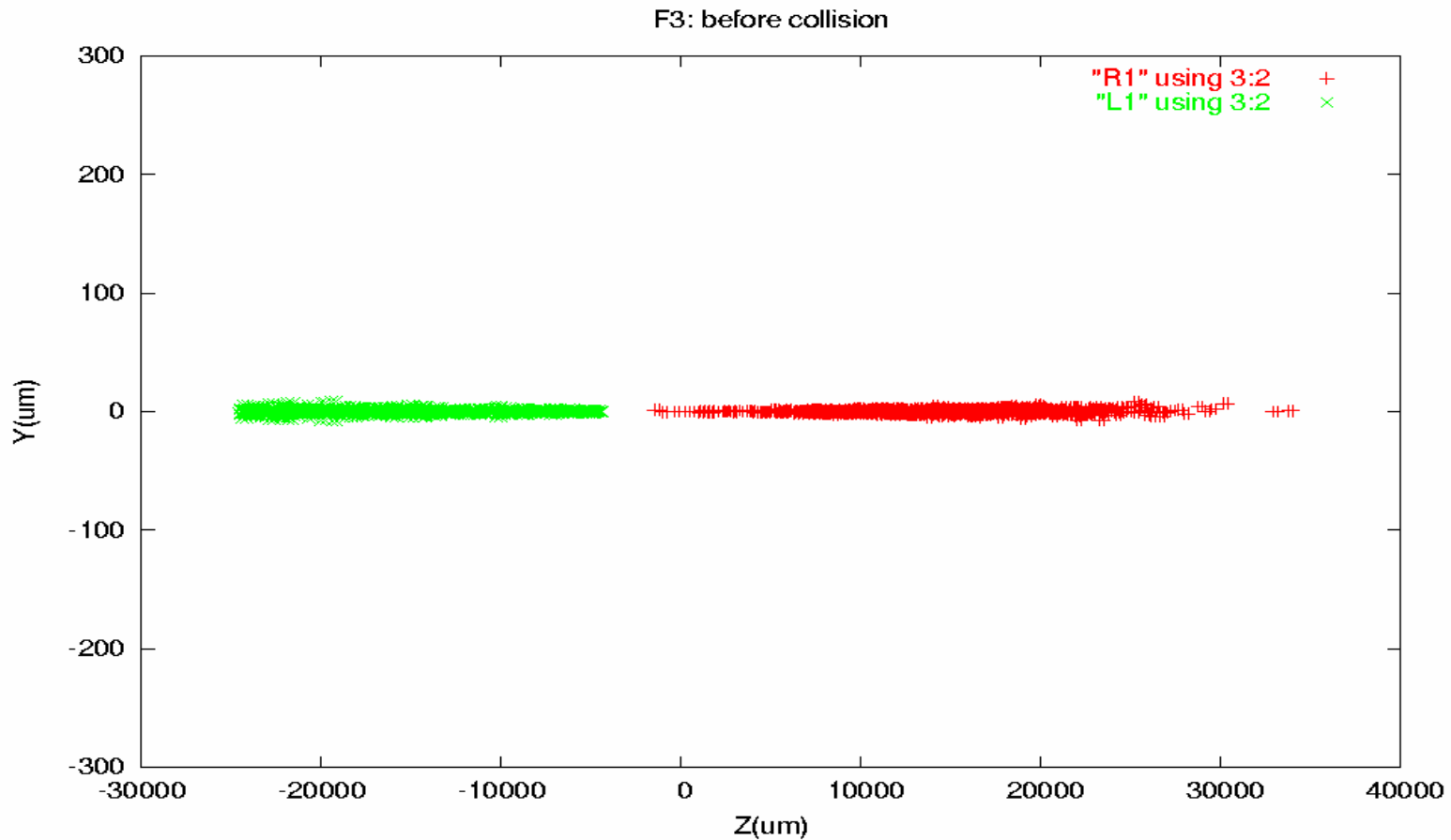
- Number of Beam Macroparticles
 - N_e 400
 - N_p 4000
- Number of turns 20000
- 6 hours on 2 GHz desktop
- Proper synchrotron motion and proton matched properly to ring; electron beam only geometrically matched; dipole offset suppressed
- No circulator ring (new electron bunch for each turn)
- Electrons have varying charge per macroparticle to give longitudinal charge distribution



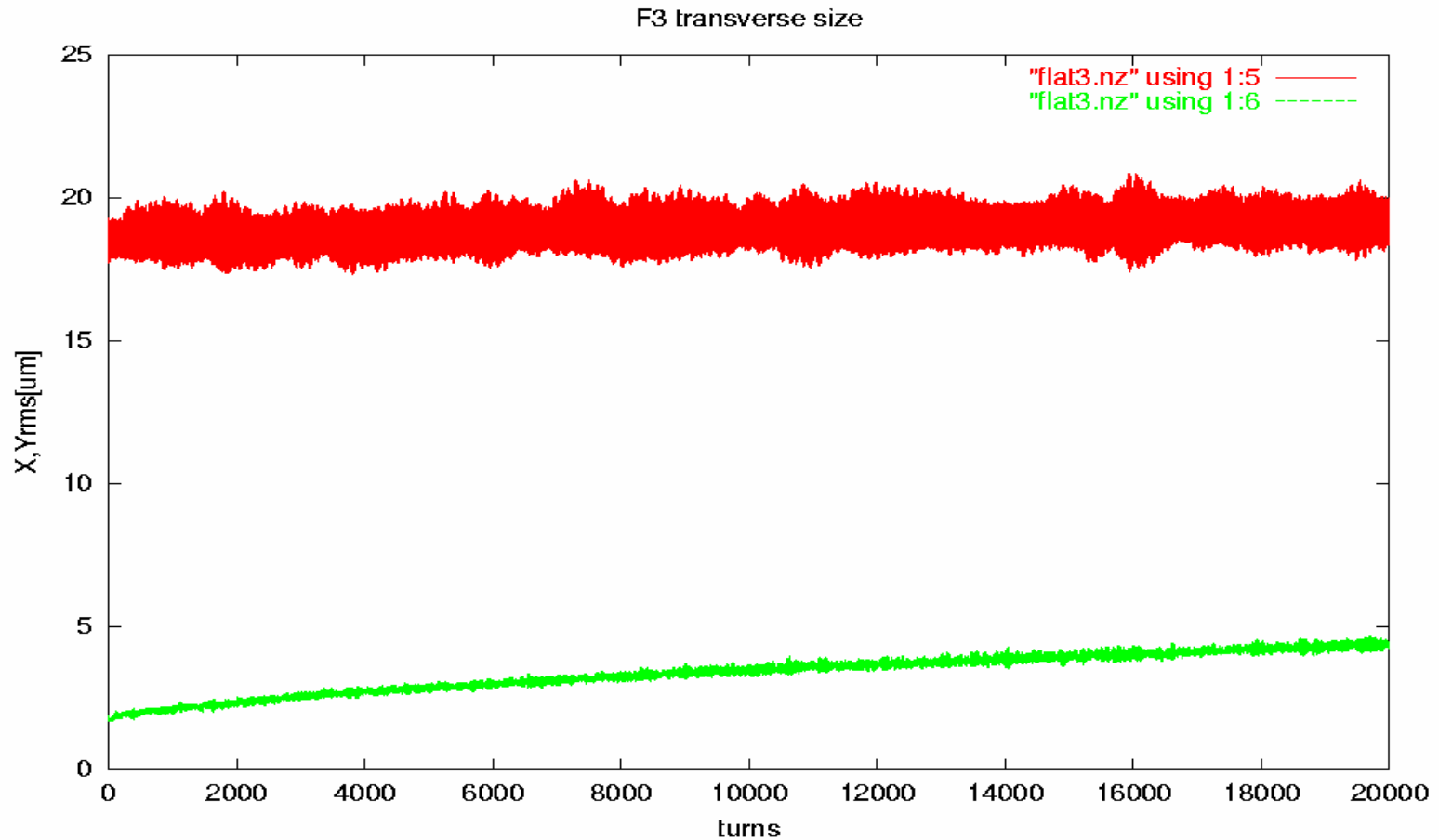
Horizontal Snapshot



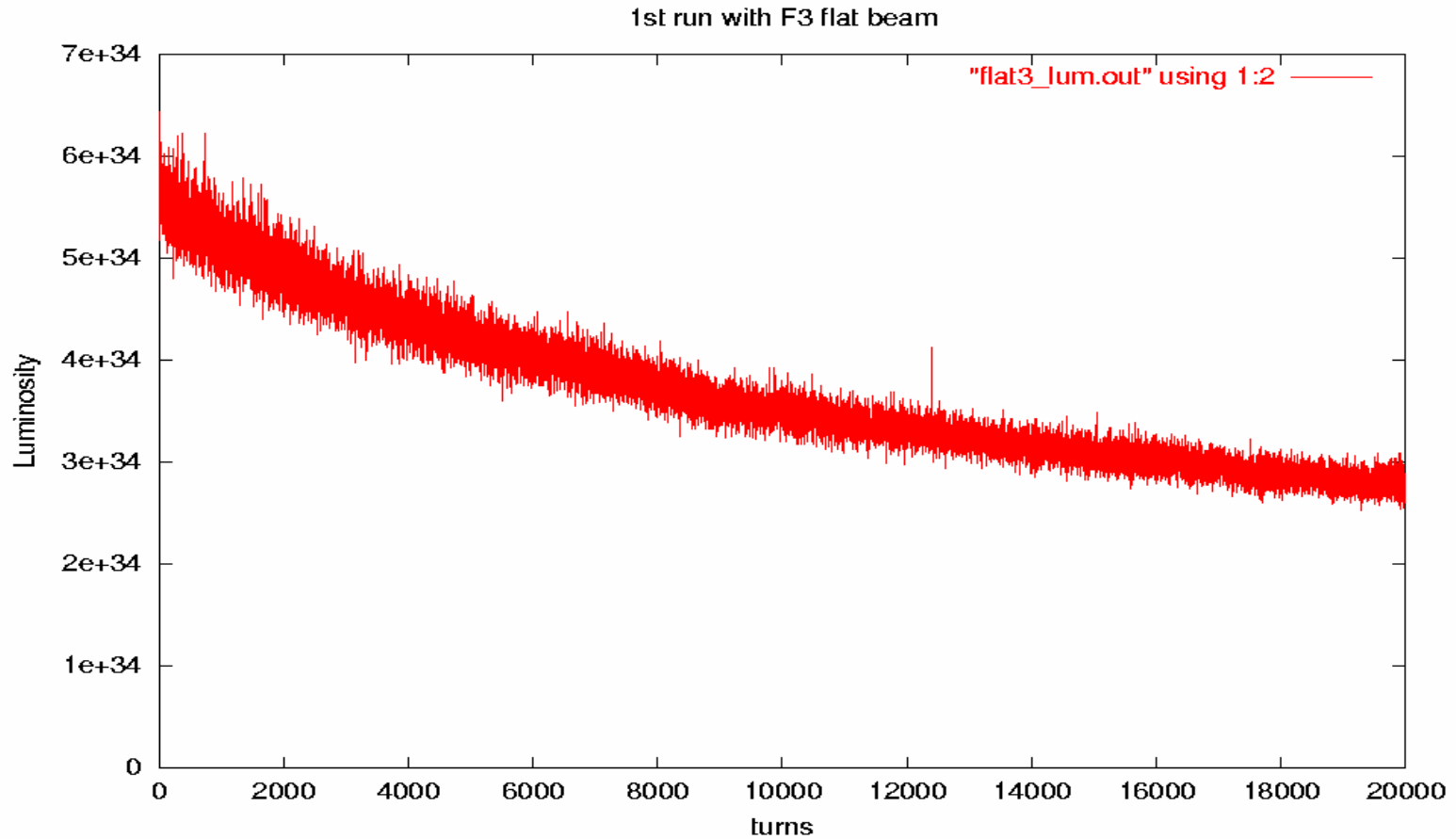
Vertical Snapshot



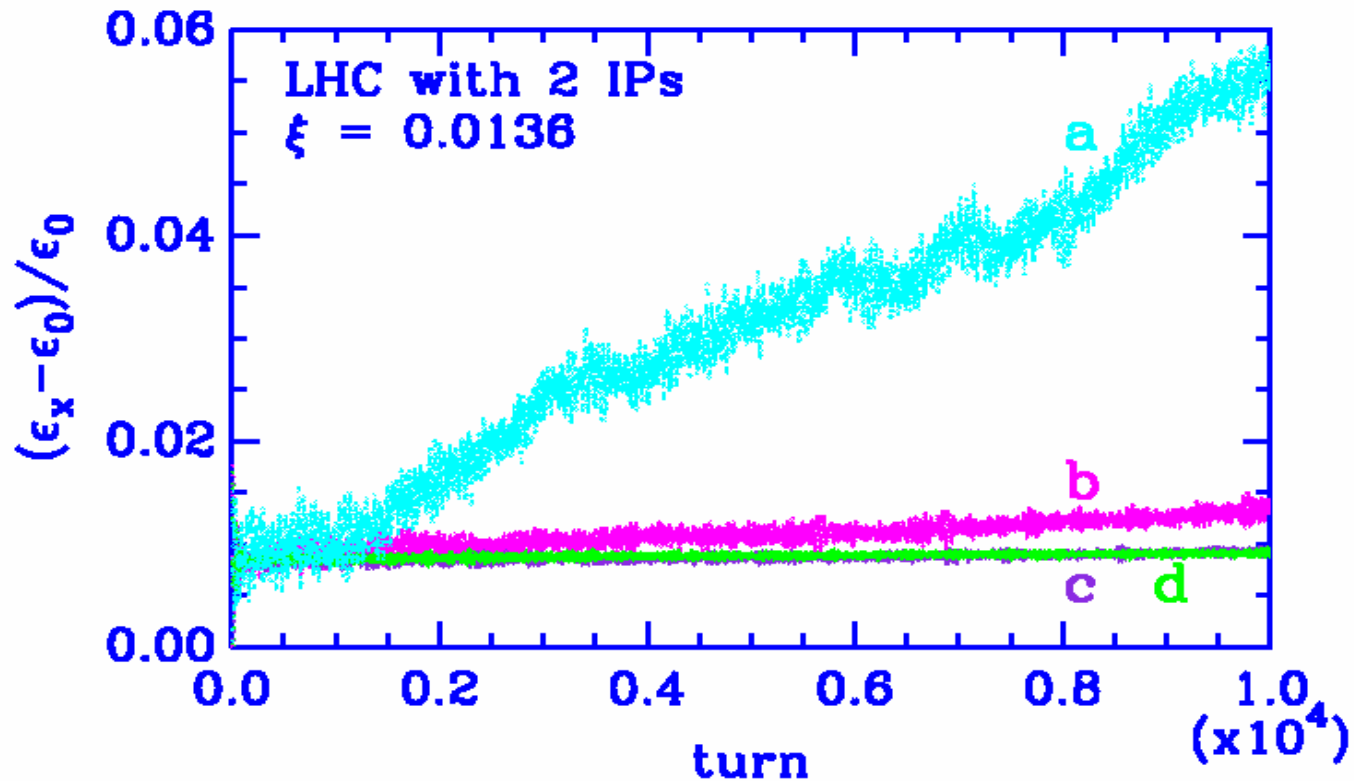
Proton Beam Transverse Sizes



Luminosity



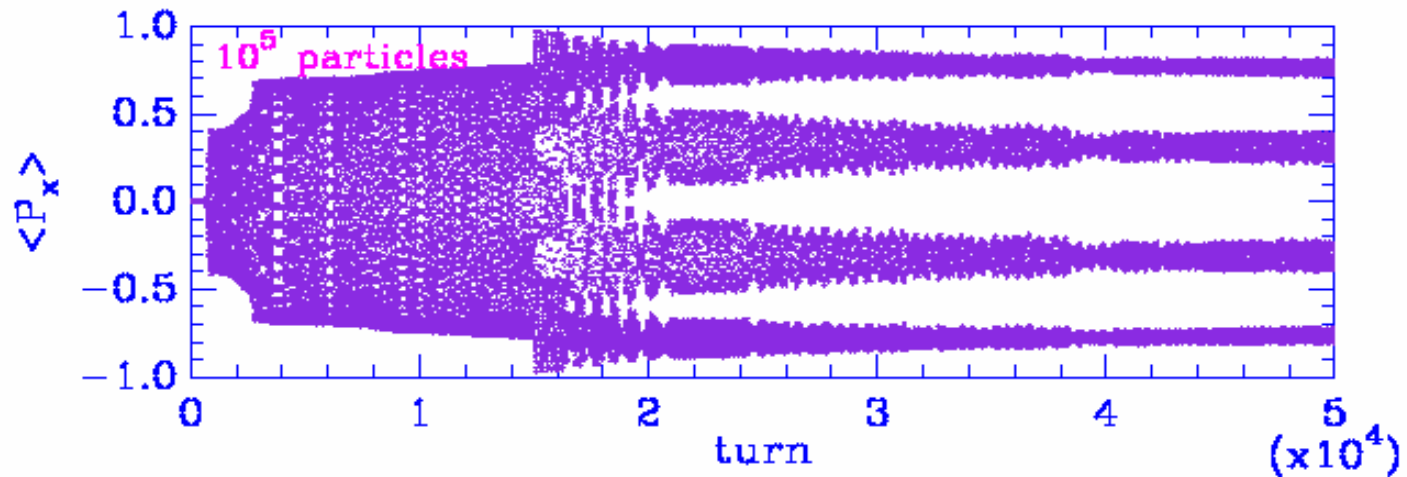
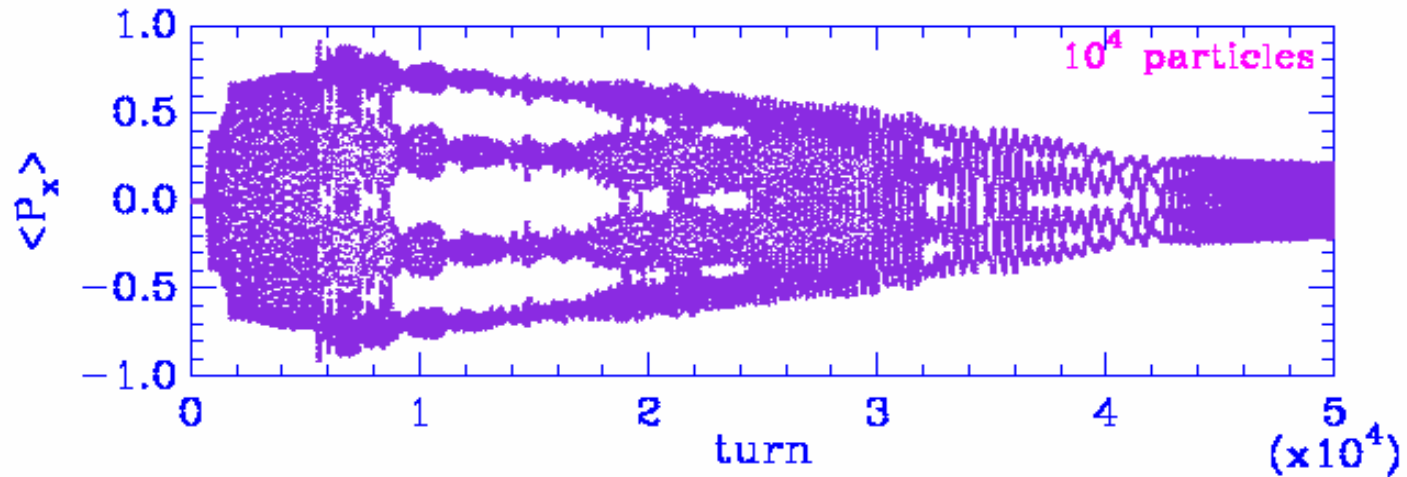
PIC Dependence on Macroparticle Number



- (a) 10^4 particles; (b) 10^5 particles;
(c) 5×10^5 particles; (d) 10^6 particles.

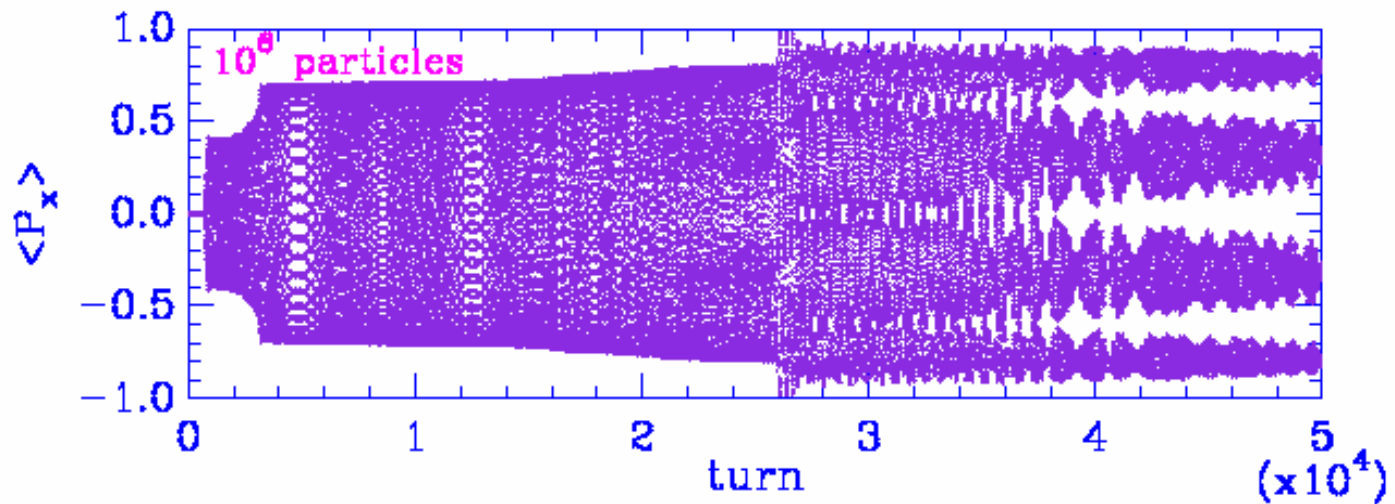
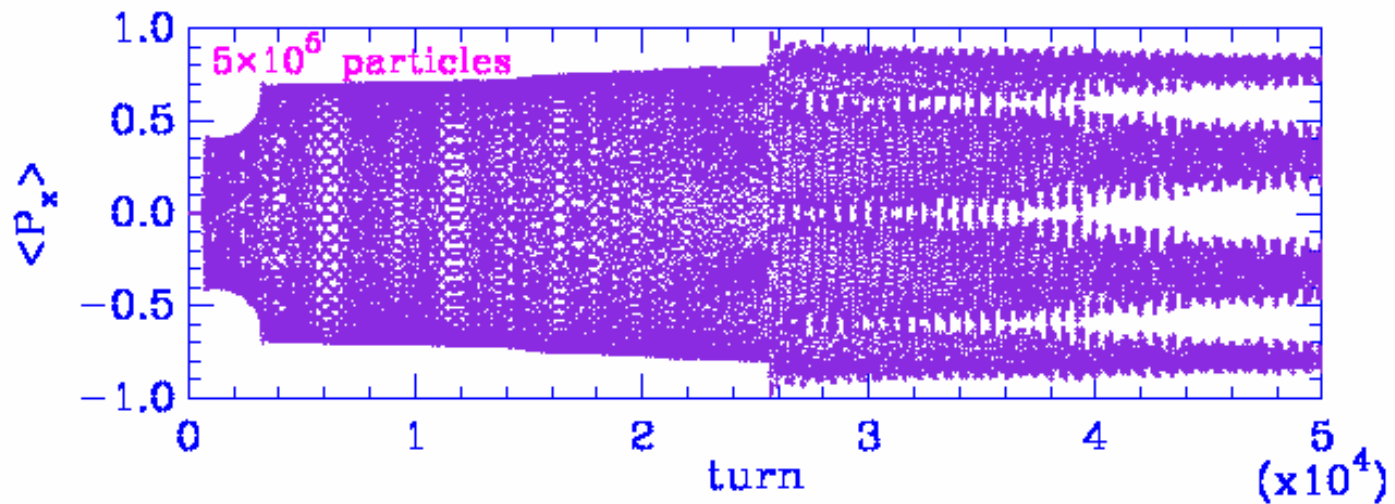


Comparison Between Different Numbers of Macro-Particles



Courtesy: Jack Shi

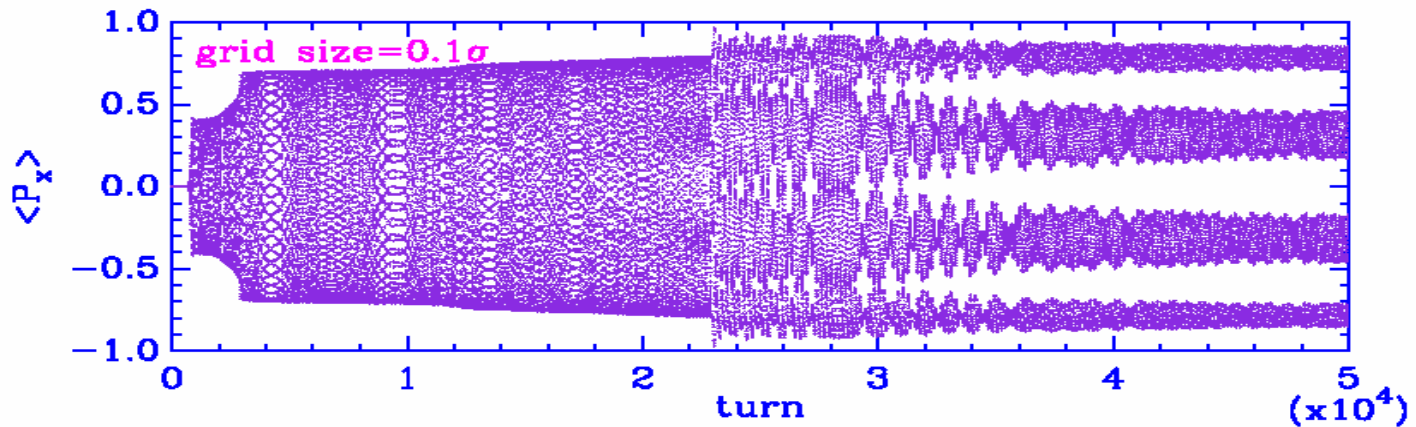
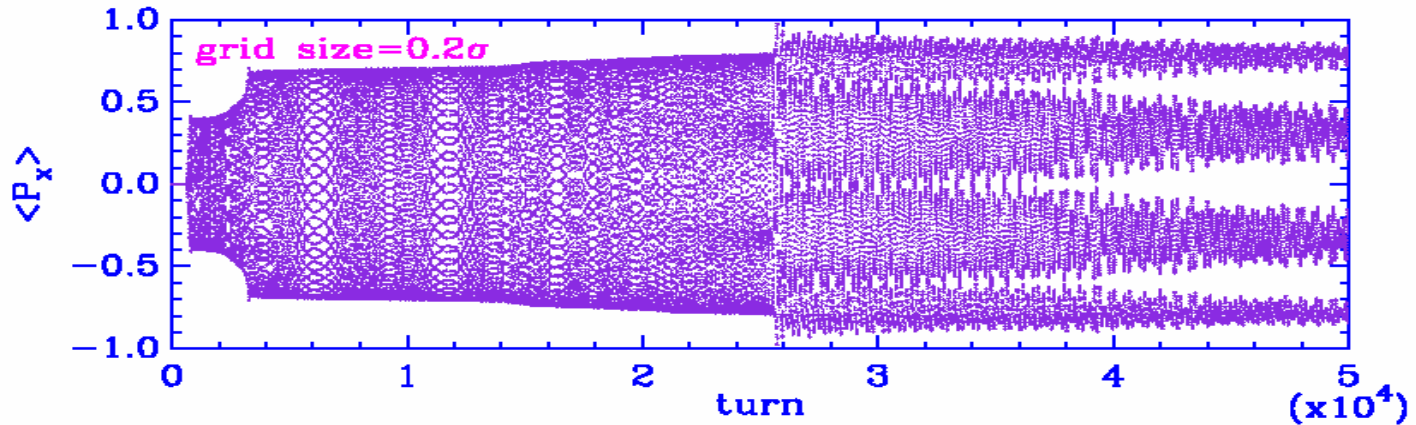




Courtesy: Jack Shi

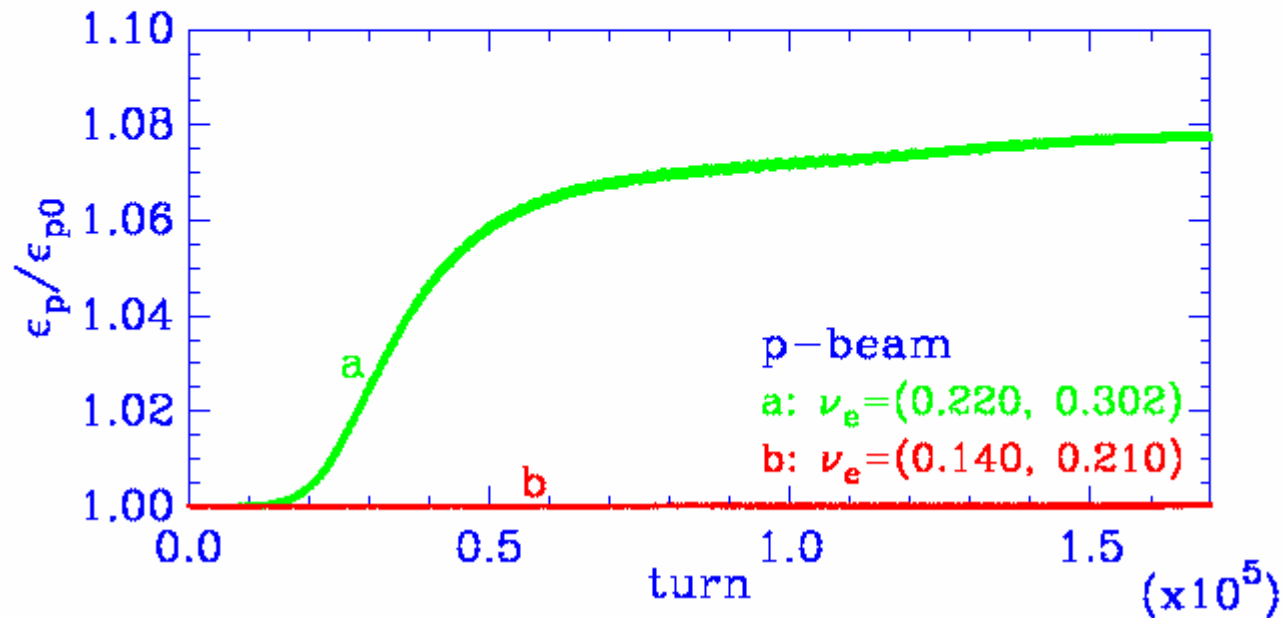


COMPARISON BETWEEN DIFFERENT GRID CONSTANTS



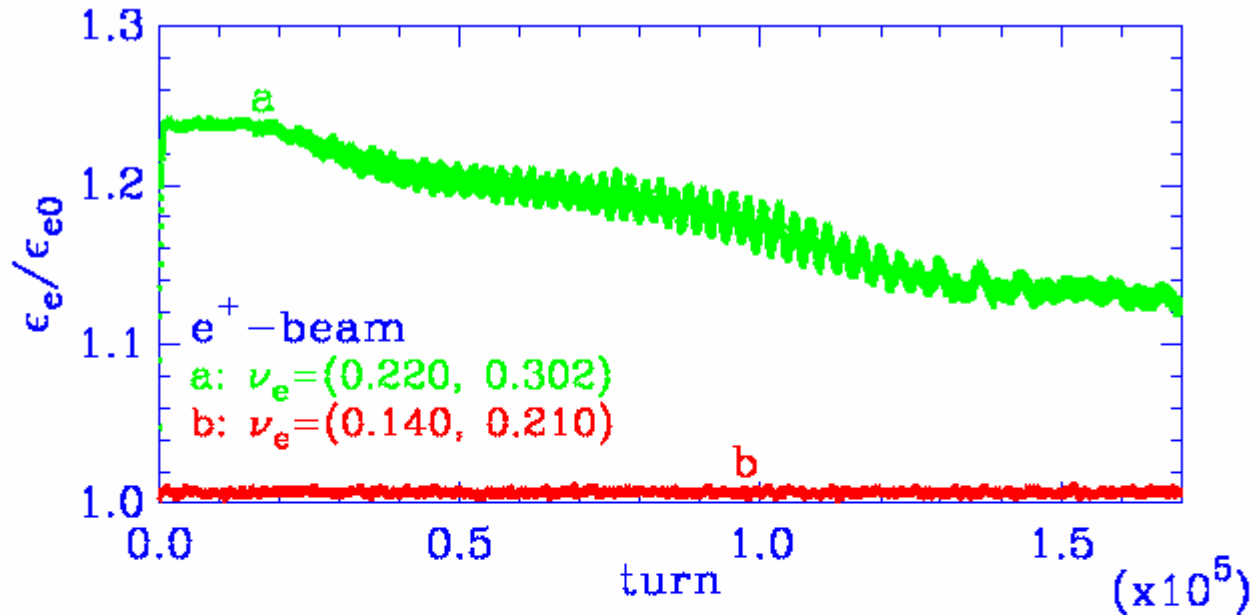
Benchmarking

HERA 2003 High-Luminosity Study With One IP Emittance Growth due to Coherent Beam-Beam Instability



Courtesy: Jack Shi





HERA 2003 Experimental Result:

In case **a**, the proton beam emittance increases $\sim 30\%$ while in case **b**, no emittance increase was observed.

Courtesy: Jack Shi



Simulation Stati (Stata?, Statuses?)

- Coulomb Sums
 - 1000 by 2000 macroparticle simulations, for 100000 turns are possible
 - Simple linear models for the ring transverse optics have been simulated
 - Full synchrotron motion allowed
 - Longitudinal slicing allowed
 - Single IP allowed
- PIC
 - Single slice 10^5 to 10^6 by 10^5 to 10^6 macroparticle simulations, for 100000 turns are possible
 - At least for some physics, such large macroparticle numbers seem necessary
 - At least one such simulation, of LHC, includes many of the important non-linear transverse optics effects in that storage ring
 - As of yet, no longitudinal slicing or synchrotron motion allowed; therefore must have relatively small tune shifts
 - Multiple IPs allowed



Future Needs

- Cooling Model
- Electron matching to the space charge in the ion beam
- Circulator ring simulations, including multiple electron crossings and multiple interactions in multiple IPs
- Crab Crossing
- PIC with slices and synchrotron motion



Conclusions

- We have explored the significant advantages of an energy recovered linac-ring collider
- We have pointed out the similarities and differences between the head-tail instability in such an arrangement and the more conventional ring-ring collider
- Very preliminary simulation studies driven by the CEBAF EIC parameter list, for a single IP configuration, have been undertaken.
- **The next stage for obtaining a still better model may be to make a slice PIC transverse code**

