

HUGS 2018
Jefferson Lab, Newport News, VA
May 29- June 15 2018

Fundamental Symmetries - 3

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Los Alamos National Laboratory



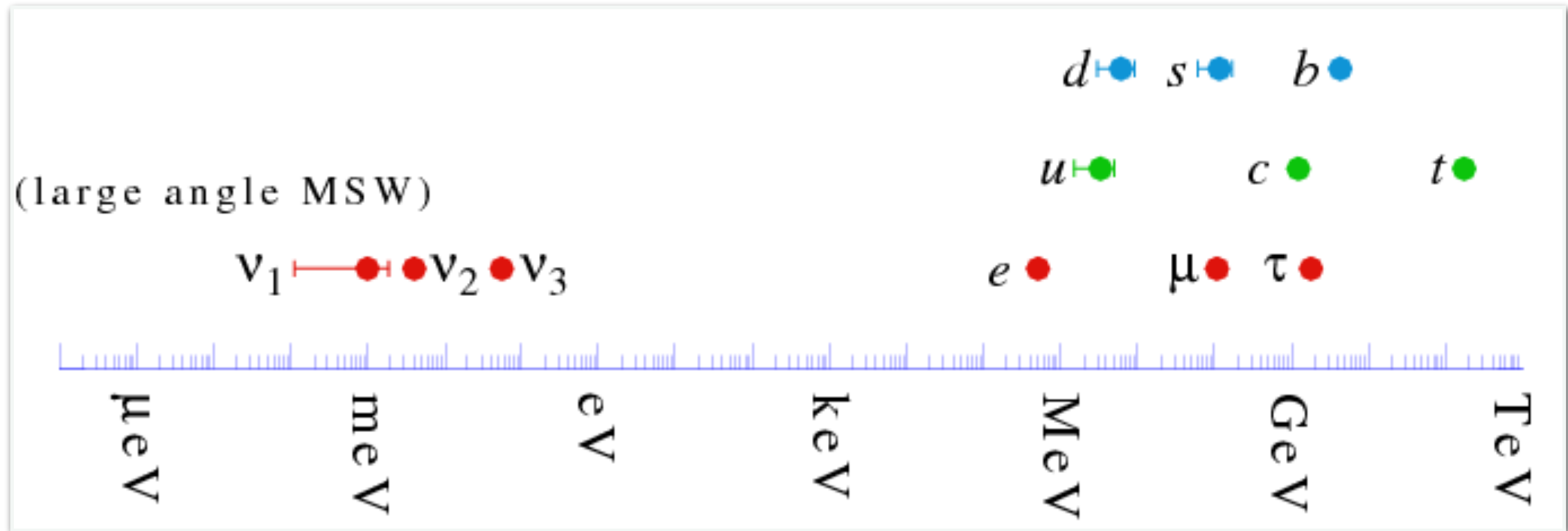
Plan of the lectures

- Review symmetry and symmetry breaking
- Introduce the Standard Model and its symmetries
- Beyond the SM:
 - hints from current discrepancies?
 - effective theory perspective
- Discuss a number of “worked examples”
 - Precision measurements: charged current (beta decays); neutral current (Parity Violating Electron Scattering).
 - Symmetry tests: CP (T) violation and EDMs; Lepton Number violation and neutrino-less double beta decay.

Status of the Standard Model

- Tested at the loop level in both electroweak and flavor sector
- Experimental hints pointing to new physics:
 - **Neutrino mass**
 - Sterile neutrinos?
 - Few-sigma discrepancies in precision physics
 - Muon $g-2$
 - Lepton universality in B meson decays
 - Neutron lifetime: beam vs bottle
 - ...

What about neutrino masses?



arXiv:1010.4131

What about neutrino masses?

- Neutrino mass requires new degrees of freedom

- Lorentz invariant “mass terms” for fermions

(1)

$$m_D \overline{\psi}_L \psi_R + \text{h.c.}$$

(2)

$$m_M \psi_L^T C \psi_L + \text{h.c.}$$

- In the case of neutrinos:
 - option (1) requires introducing ν_R and using Higgs to make it SU(2) gauge invariant (as for other fermions)
 - option (2) is not SU(2) gauge invariant

What about neutrino masses?

- Neutrino mass requires new degrees of freedom
- Simple / natural option: three R-handed neutrinos ν_{Ri} (gauge singlets)

$$\mathcal{L}_{\nu SM} = \mathcal{L}_{SM} + i\bar{\nu}_R \not{\partial} \nu_R - \left(\frac{1}{2} \nu_R^T C M_R \nu_R + \bar{\ell} Y_\nu \nu_R \tilde{\varphi} + \text{h.c.} \right)$$

Both allowed by gauge symmetry
Mass term breaks $U(1)_L$

$$l \rightarrow e^{i\alpha} l \quad e \rightarrow e^{i\alpha} e \quad \nu_R \rightarrow e^{i\alpha} \nu_R$$

What about neutrino masses?

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$$\mathcal{L}_{\nu SM} = \mathcal{L}_{SM} + i\bar{\nu}_R \not{\partial} \nu_R - \left(\frac{1}{2} \nu_R^T \mathbf{M}_R \nu_R + \bar{\ell} Y_\nu \nu_R \tilde{\varphi} + \text{h.c.} \right)$$

- Dirac neutrinos: $\mathbf{M}_R = 0$. Complete analogy to quark sector ($B \rightarrow L$), except for tiny ($O(10^{-10})$) Yukawa couplings

$$Y_e = V_{eL}^\dagger Y_e^{\text{diag}} V_{eR}$$

 \Rightarrow

$$\frac{g}{\sqrt{2}} W_\mu^- \bar{e}_L^\alpha \gamma^\mu U^{\alpha i} \nu_L^i$$

Unitary mixing in CC vertex: 3 angles, 1 phase

$$Y_\nu = V_{\nu L}^\dagger Y_\nu^{\text{diag}} V_{\nu R}$$

$$U = V_{eL} V_{\nu L}^\dagger$$

- Majorana neutrinos: $M_R \neq 0$. L not conserved

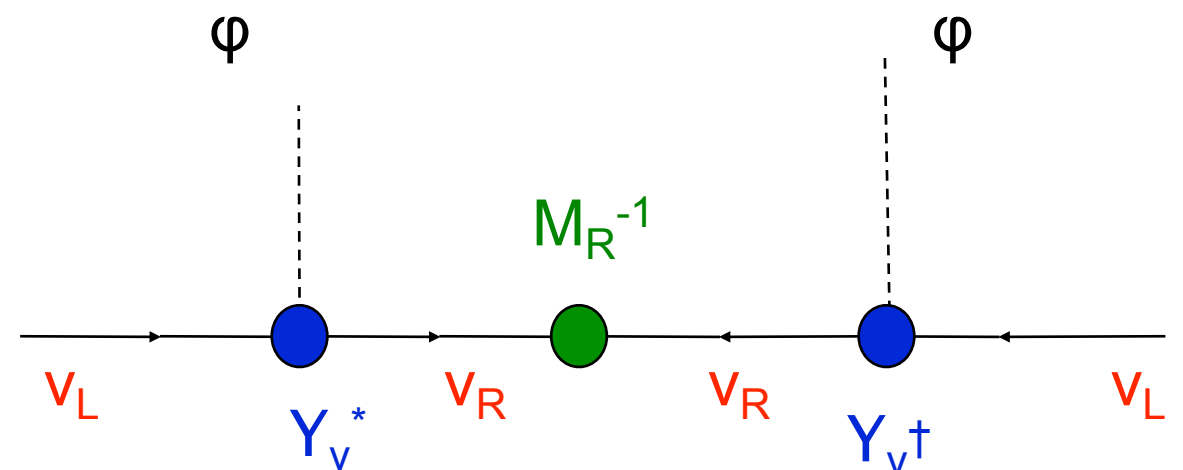
$$\mathcal{L}_{\nu SM} = \mathcal{L}_{SM} + i\bar{\nu}_R \not{\partial} \nu_R - \left(\frac{1}{2} \nu_R^T C M_R \nu_R + \bar{\ell} Y_\nu \nu_R \tilde{\varphi} + \text{h.c.} \right)$$

- In general 6x6 mass matrix for $\begin{pmatrix} \nu_L \\ \nu_R \end{pmatrix}$: six Majorana ($\nu = \nu^c$) eigenstates
- If $M_R \gg v Y_\nu$: 3 light ($\nu_L \rightarrow \nu_i$) and 3 heavy ($\nu_R \rightarrow N_i$) eigenstates

$$\mathcal{L}_{\nu SM} \supset -\frac{1}{2} \nu_L^T C m_\nu \nu_L$$

$$m_\nu = v^2 Y_\nu^* M_R^{-1} Y_\nu^\dagger$$

We could have written this term without reference to ν_R and in SU(2) gauge-invariant form (more later)



- Majorana neutrinos: $M_R \neq 0$. L not conserved

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- If $M_R \gg v Y_\nu$: 3 light ($\nu_L \rightarrow \nu_i$) and 3 heavy ($\nu_R \rightarrow N_i$) eigenstates

- Mixing of 3 light Majorana neutrinos:

$$\mathcal{L}_{\nu SM} \supset -\frac{1}{2} \nu_L^T C m_\nu \nu_L$$

$$m_\nu = V_{\nu L}^T m_\nu^{\text{diag}} V_{\nu L}$$

 \Rightarrow

$$\frac{g}{\sqrt{2}} W_\mu^- \bar{e}_L^\alpha \gamma^\mu U^{\alpha i} \nu_L^i$$

$$Y_e = V_{eL}^\dagger Y_e^{\text{diag}} V_{eR}$$

$$U = V_{eL} V_{\nu L}^\dagger$$

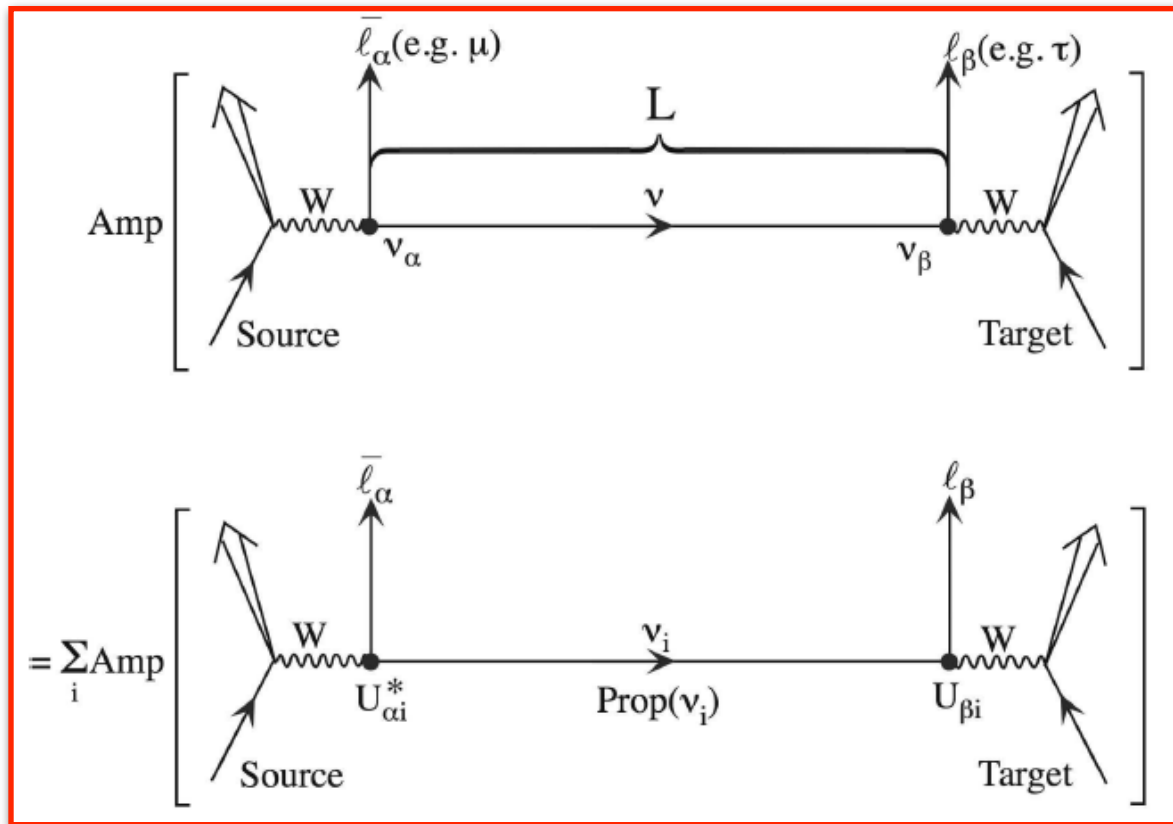
Unitary mixing in CC vertex: 3 angles, 1+2 phases

$$U = U_{\text{Dirac}} \times \begin{pmatrix} 1 & & \\ & e^{i\alpha_1} & \\ & & e^{i\alpha_2} \end{pmatrix}$$

Neutrino phenomenology

- \mathcal{L}_{VSM} probed at the **Intensity Frontier** (accelerator, reactor) and **Cosmic Frontier** (solar, atmospheric, astro)
- Oscillation experiments sensitive to mass splittings and mixing angles

Image credit: B. Kayser



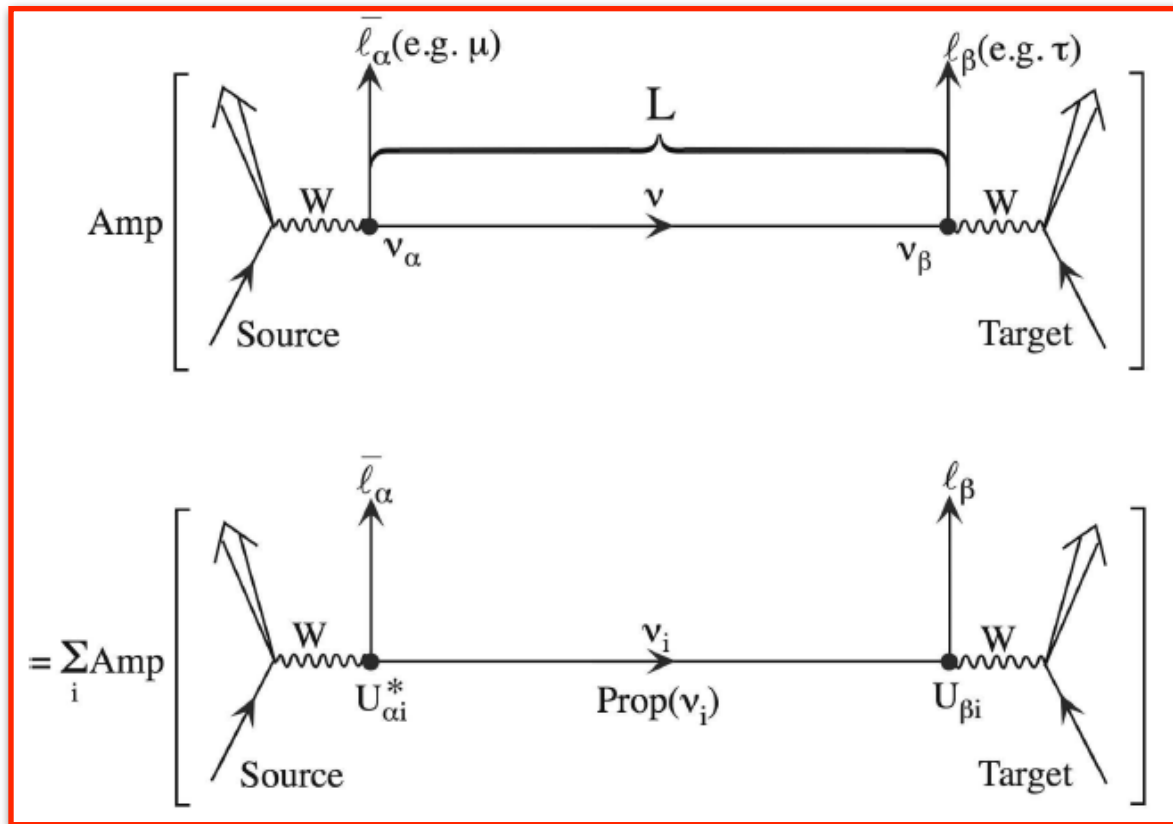
$$U = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix},$$

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \sin^2 2\theta \times \frac{1}{2} \left(1 - \cos \frac{\Delta m^2 L}{2E} \right)$$

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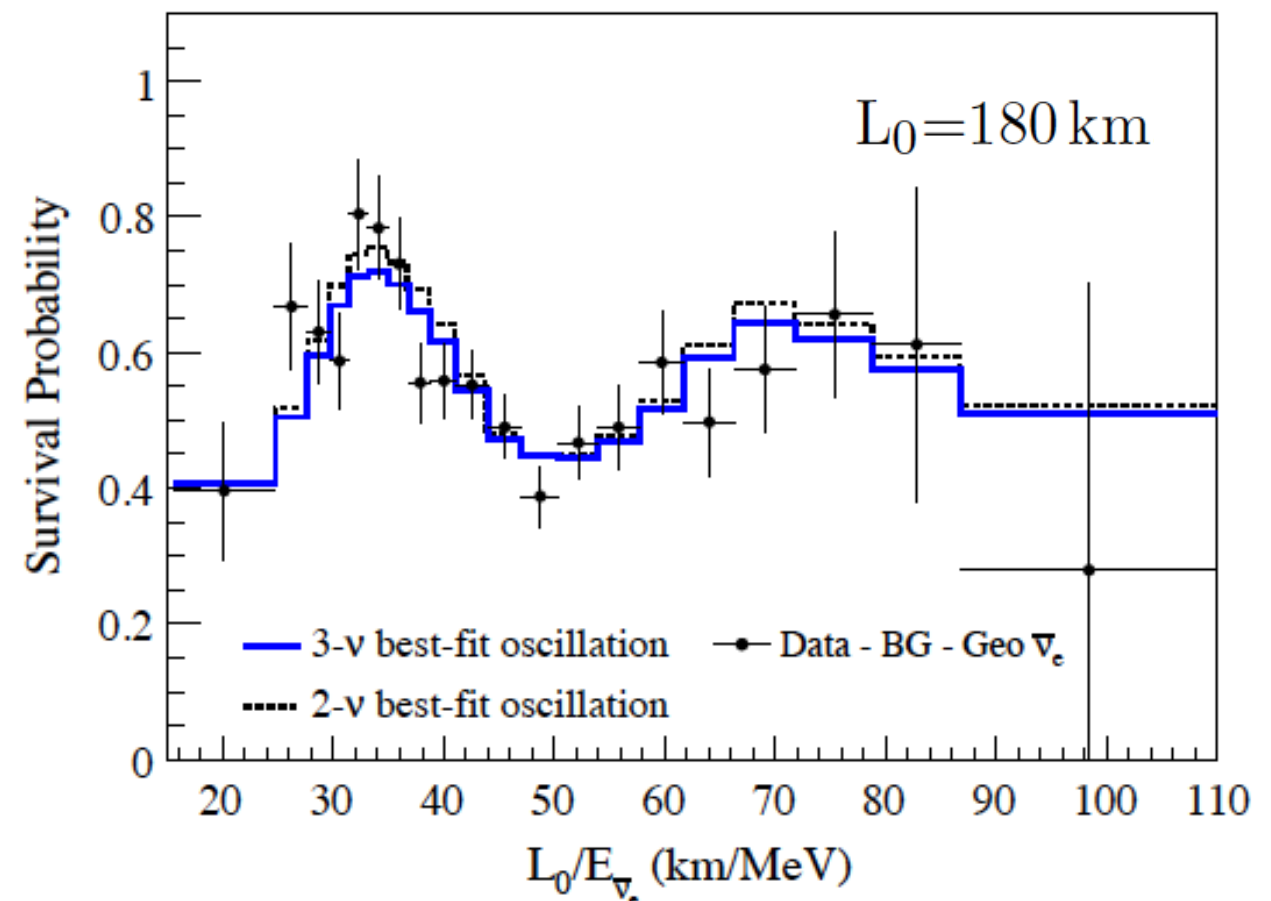
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KAMLAND 2011



Reactor electron anti-neutrino survival probability

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- Oscillation experiments sensitive to mass splittings and mixing angles

World data consistent with 3 light states,
but other light ν not excluded

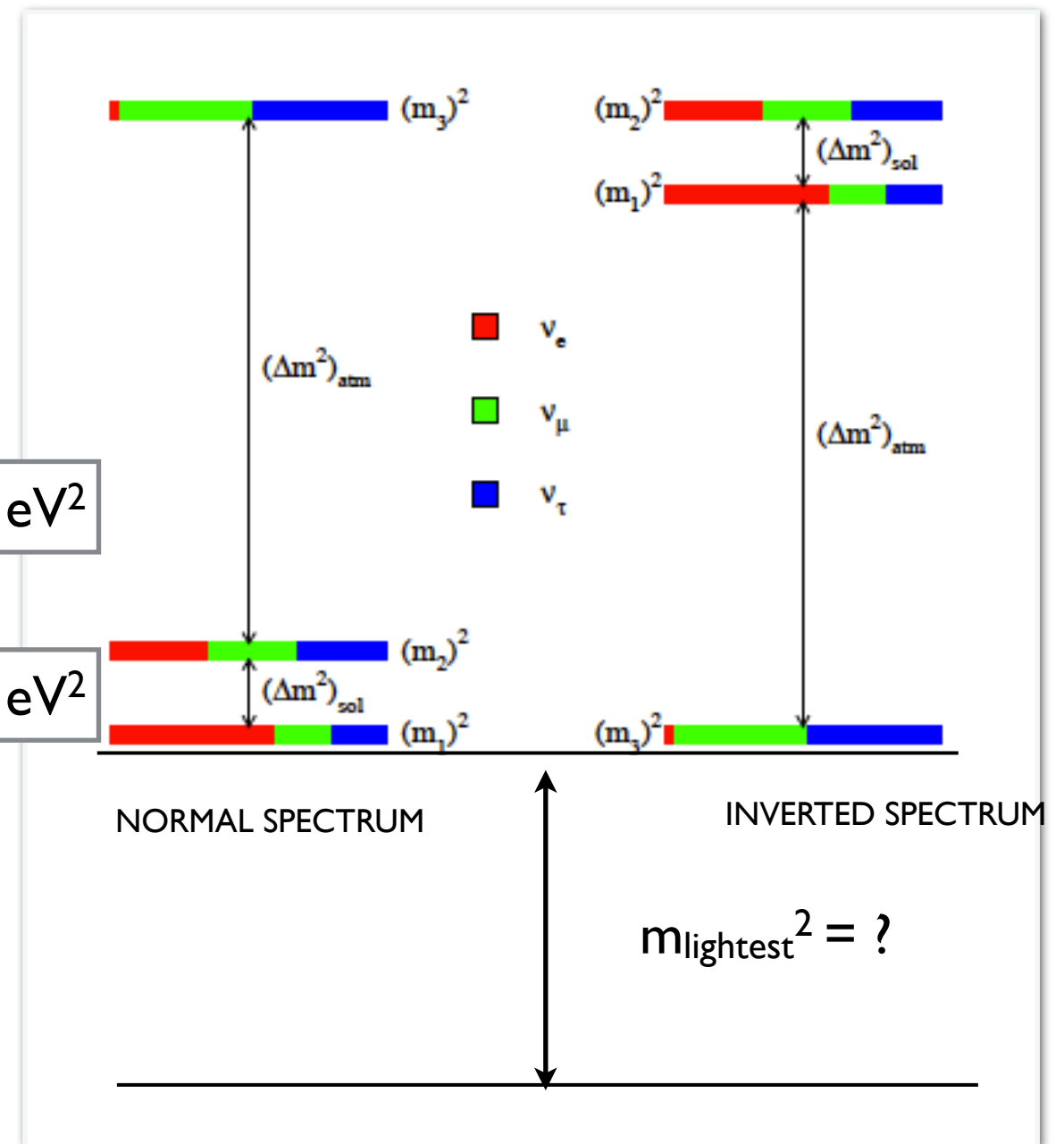
$$U = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix}$$

Parameter	best-fit ($\pm 1\sigma$)
Δm_{21}^2 [10^{-5} eV ²]	$7.54^{+0.26}_{-0.22}$
$ \Delta m^2 $ [10^{-3} eV ²]	2.43 ± 0.06 (2.38 ± 0.06)
$\sin^2 \theta_{12}$	0.308 ± 0.017
$\sin^2 \theta_{23}, \Delta m^2 > 0$	$0.437^{+0.033}_{-0.023}$
$\sin^2 \theta_{23}, \Delta m^2 < 0$	$0.455^{+0.039}_{-0.031}$
$\sin^2 \theta_{13}, \Delta m^2 > 0$	$0.0234^{+0.0020}_{-0.0019}$
$\sin^2 \theta_{13}, \Delta m^2 < 0$	$0.0240^{+0.0019}_{-0.0022}$
δ/π (2σ range quoted)	$1.39^{+0.38}_{-0.27}$ ($1.31^{+0.29}_{-0.33}$)

PDG 2014

$$\sim 2.4 \cdot 10^{-3} \text{ eV}^2$$

$$\sim 7.5 \cdot 10^{-5} \text{ eV}^2$$



Neutrino phenomenology

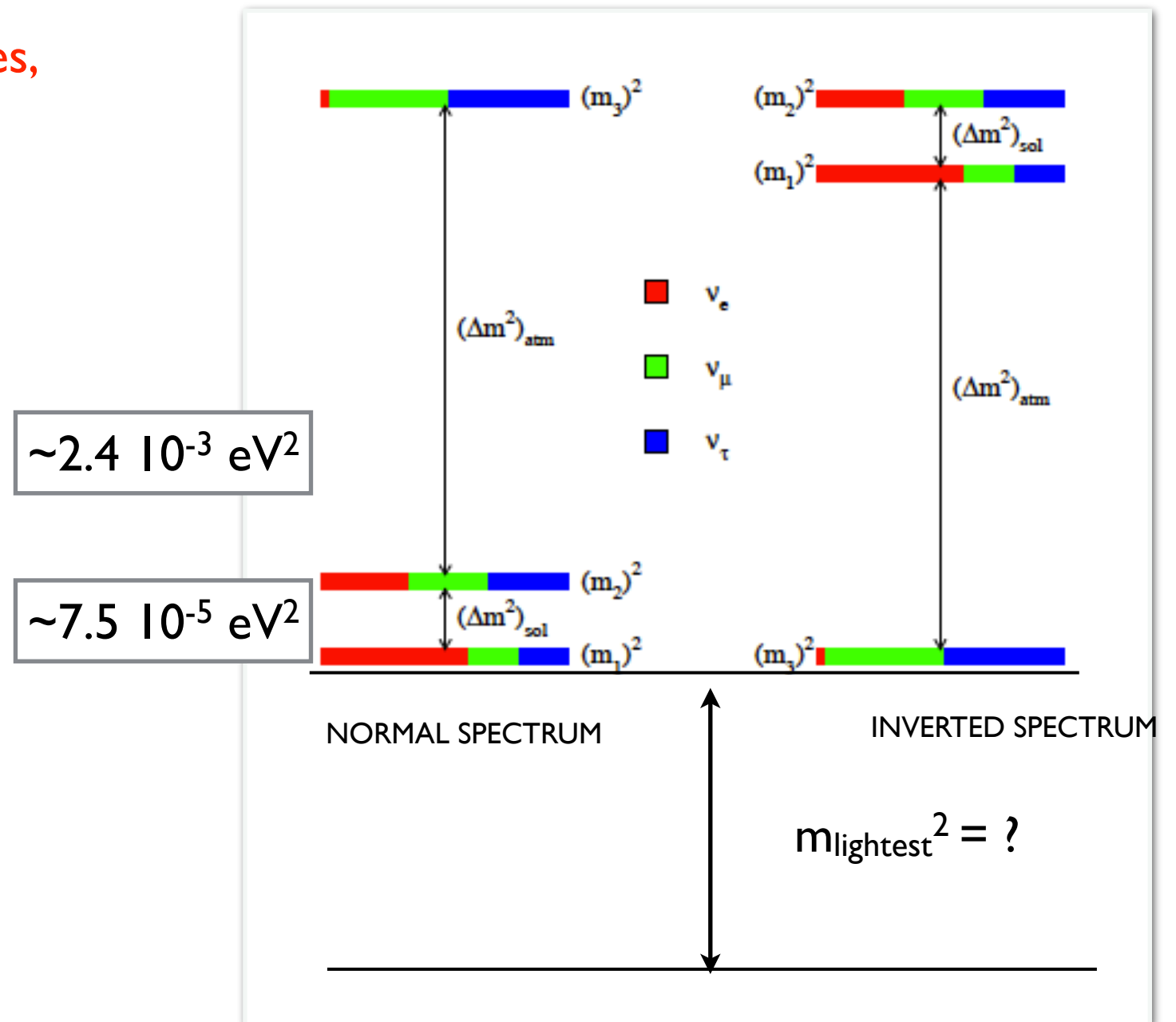
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World data consistent with 3 light states,
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$$U \sim \begin{pmatrix} 0.8 & 0.5 & 0.2 \\ 0.4 & 0.6 & 0.7 \\ 0.4 & 0.6 & 0.7 \end{pmatrix}$$

A. de Gouvea

$$V_{CKM} \sim \begin{pmatrix} 1 & 0.2 & 0.001 \\ 0.2 & 1 & 0.01 \\ 0.001 & 0.01 & 1 \end{pmatrix}$$

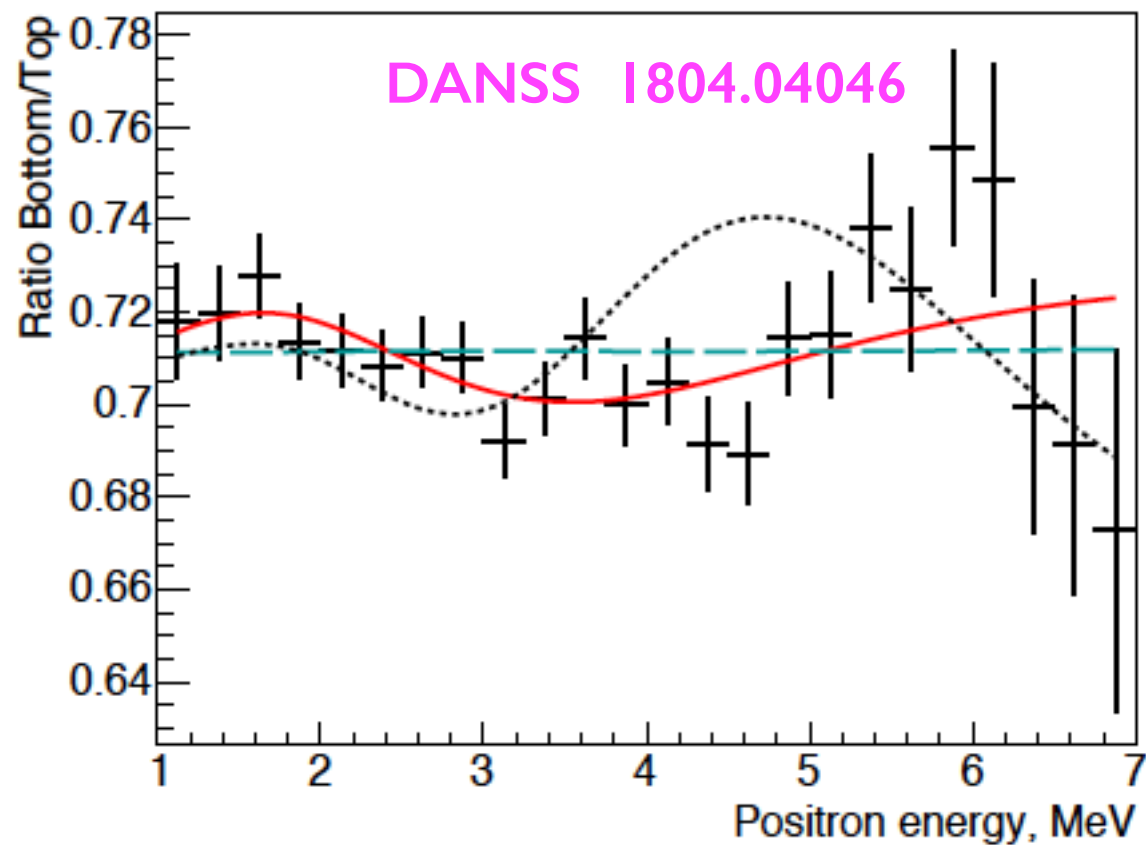


Open questions

- Many key aspects of ν dynamics remain unknown, and will be explored by experiments in the next decade
- **Symmetries / particle content:**
 - Is lepton number (L) broken? (Dirac vs Majorana) ($0\nu\beta\beta$)
 - Are there light sterile ν 's? (short-baseline anomalies, cosmo)
- Determine **parameters of mass matrix** (regardless its origin):
 - Absolute mass scale (beta decay, $0\nu\beta\beta^*$, cosmology*)
 - Mass ordering (oscillation experiments)
 - Mixing angles (✓), CPV phase

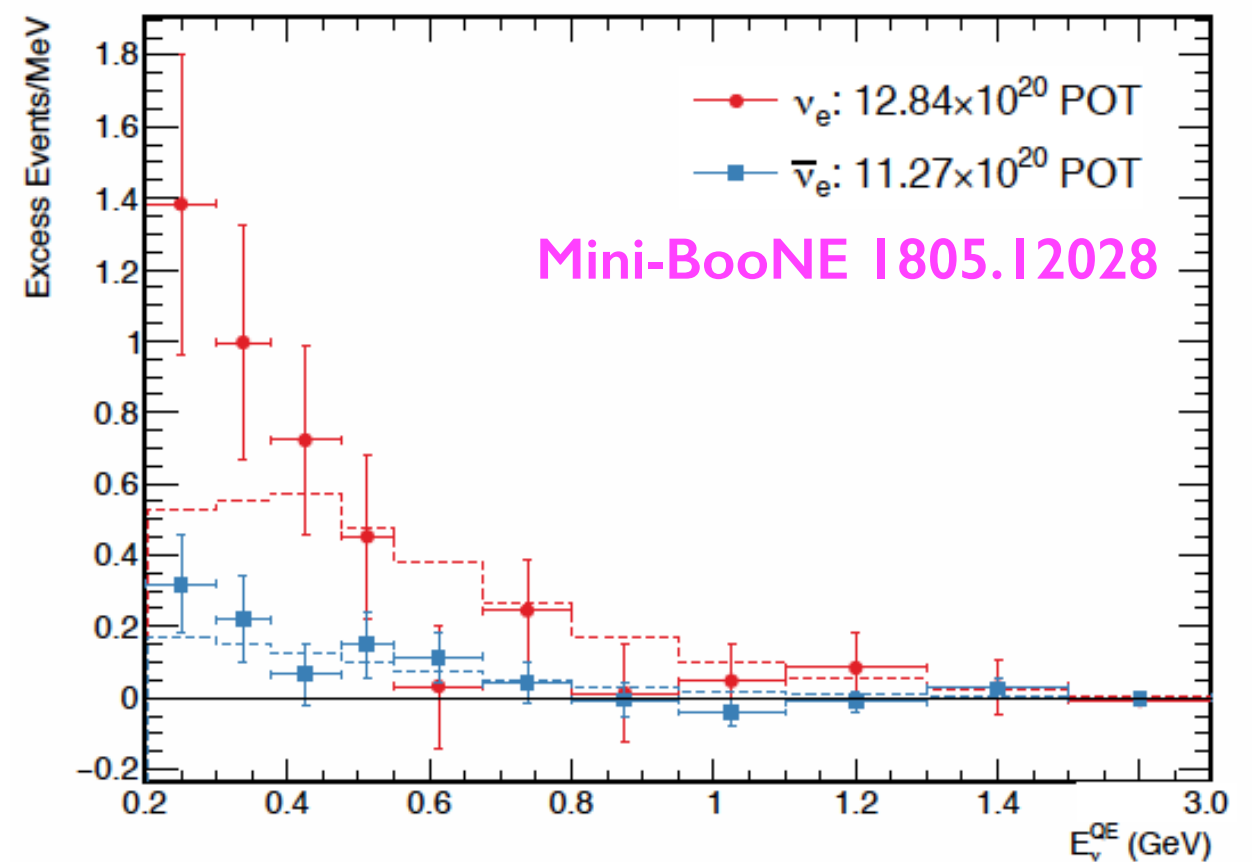
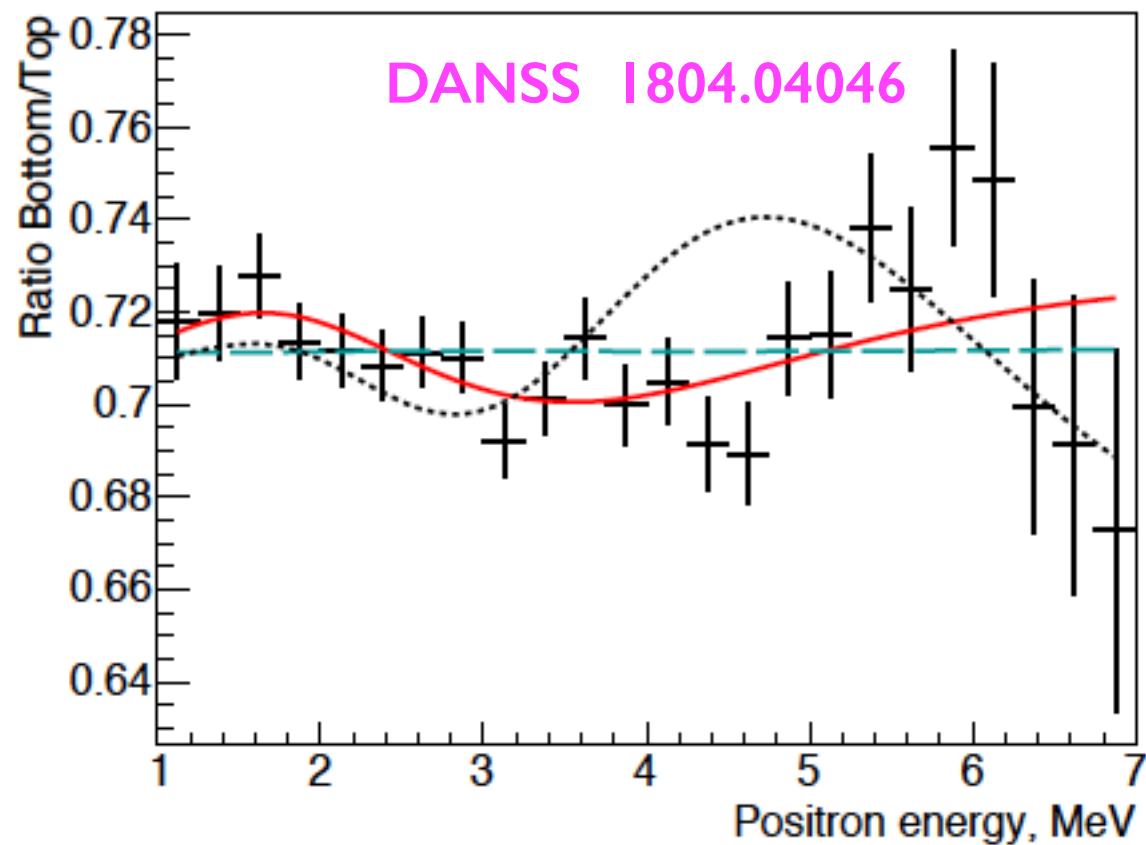
Sterile neutrinos?

- Anomalies in ν_e disappearance ν_e to ν_μ appearance data
- Reactor anomaly: $\bar{\nu}_e$ to $\bar{\nu}_e$ ($\sim 3\sigma$), now independent of calculated flux



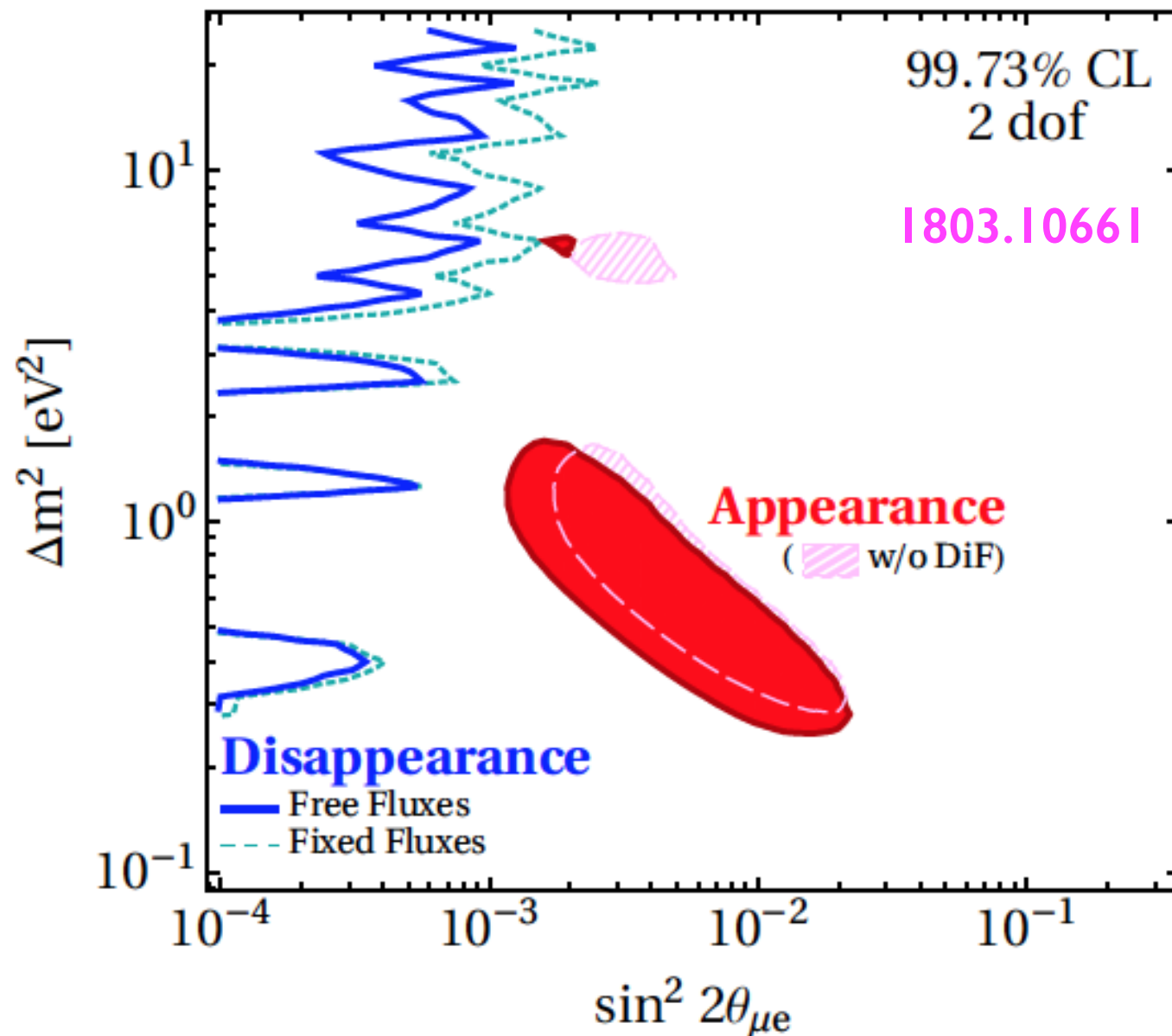
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- Anomalies in ν_e disappearance ν_e to ν_μ appearance data
- Reactor anomaly: $\bar{\nu}_e$ to $\bar{\nu}_e$ ($\sim 3\sigma$), now independent of calculated flux
- Short-baseline anomaly: ν_μ to ν_e excess (4.8σ)



Sterile neutrinos?

- Anomalies in ν_e disappearance ν_e to ν_μ appearance data
- Global analysis in 3+1 scheme, using increasingly strong bounds on ν_μ disappearance (MINOS+, ...)



$$P(\nu_e \text{ to } \nu_e) \sim 1 - 4|U_{e4}|^2 \sin^2(\text{phase})$$

$$P(\nu_\mu \text{ to } \nu_\mu) \sim 1 - 4|U_{\mu 4}|^2 \sin^2(\text{phase})$$

$$P(\nu_\mu \text{ to } \nu_e) \sim 1 - 4 \underbrace{|U_{e4} U_{\mu 4}|^2}_{\sin^2(2\theta_{\mu e})} \sin^2(\text{phase})$$

$$\text{phase} = \left(\frac{\Delta m_{41}^2 L}{4E} \right)$$

4.7 σ tension between
different data sets!

MicroBooNE at FNAL will
shed light on MiniBooNE
excess (e vs γ discrimination)

Status of the Standard Model

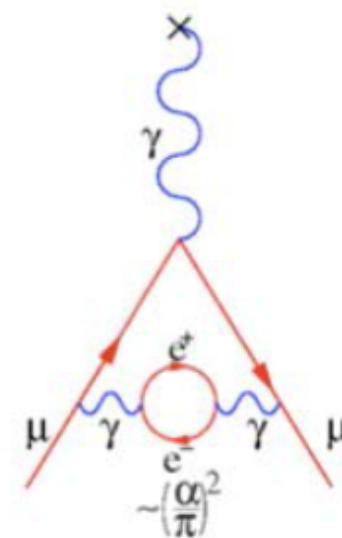
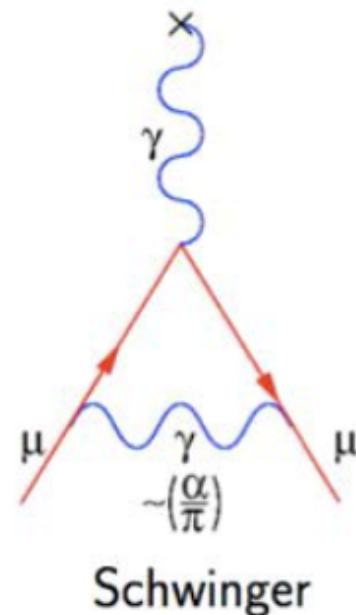
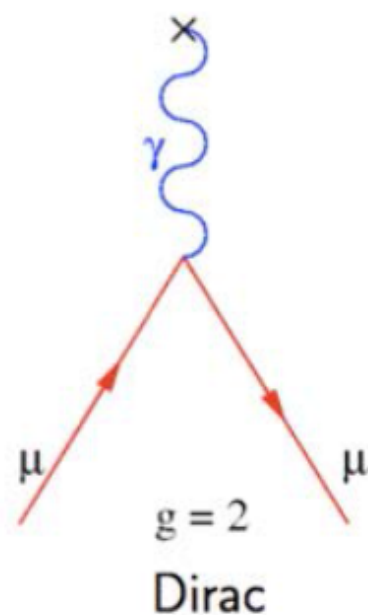
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Muon anomalous magnetic moment

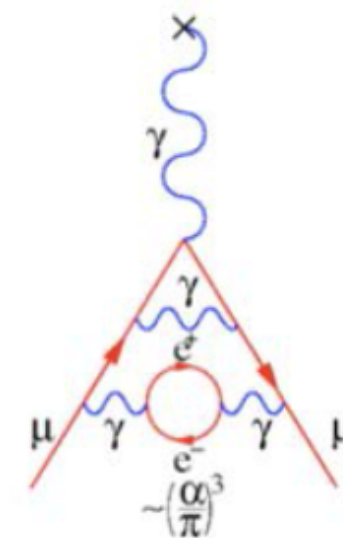
$$\vec{\mu} = g \frac{e}{2mc} \vec{s}, \quad \vec{s} = \frac{\hbar}{2} \vec{\sigma}$$

- Dirac predicts $g=2$ in 1928
- 1947: Measurements find $g_e \neq 2$

- Schwinger calculated $g_e = 2(1 + a_e)$ $a_e = \frac{(g_e - 2)}{2} = \frac{\alpha}{2\pi} \approx 0.00116$



Kinoshita and others

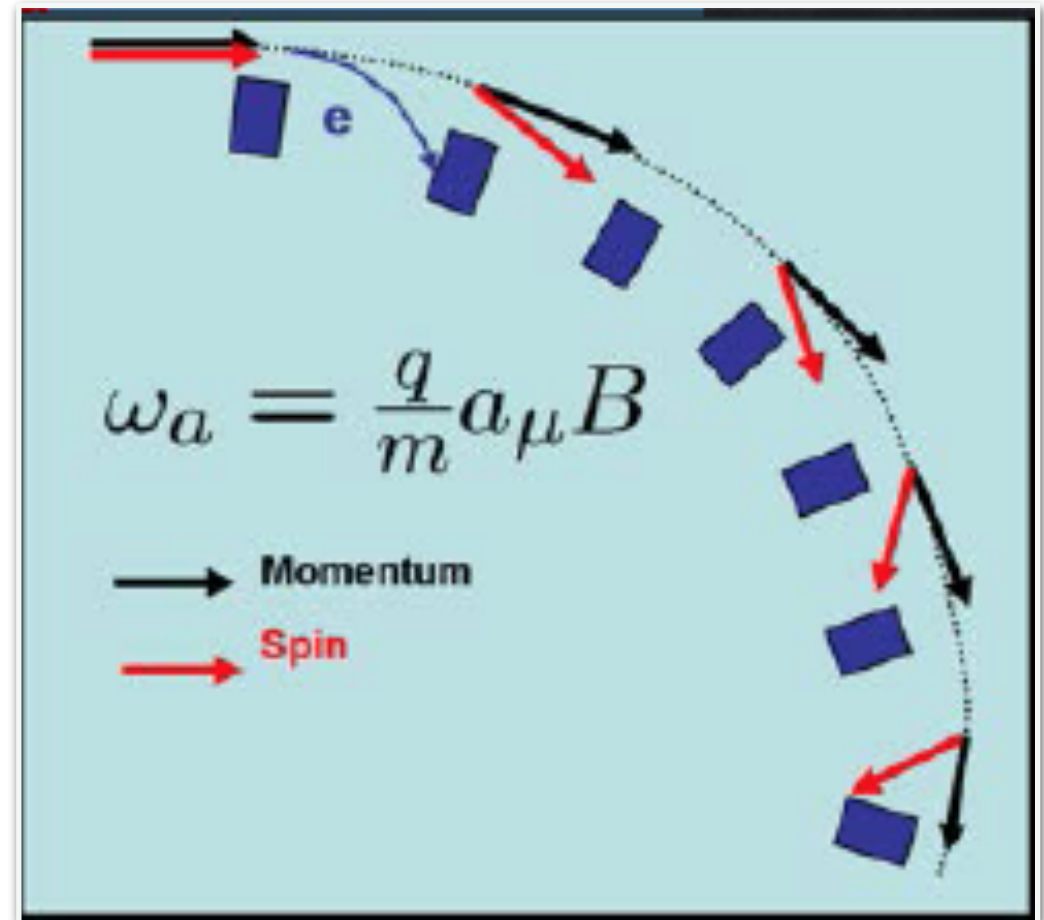


Great success of QED

- Current experimental precision: $\Delta g_e = 5.2 \times 10^{-13}$ and $\Delta g_\mu = 1.2 \times 10^{-9}$
 - g_e used to determine the electromagnetic coupling
 - g_μ used to challenge the SM!

- How is g_μ (a_μ) measured?
 - Exploit the fact that momentum and spin do not precess in the same way in a B field
 - Relative frequency ω_a proportional to $(g-2)*B$

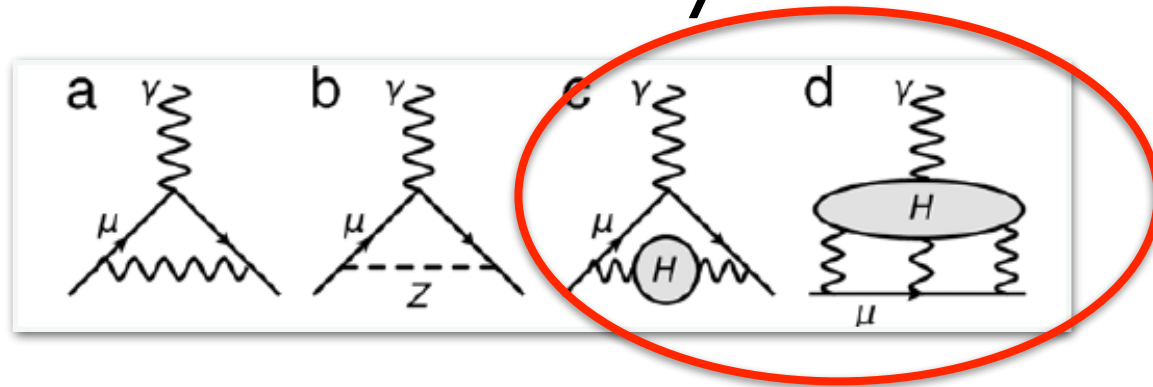
$$\omega_C = \frac{eB}{mc\gamma}$$



$$\omega_S = \frac{geB}{2mc} + (1 - \gamma) \frac{eB}{\gamma mc}$$

- Where are we?

SM theory



VS

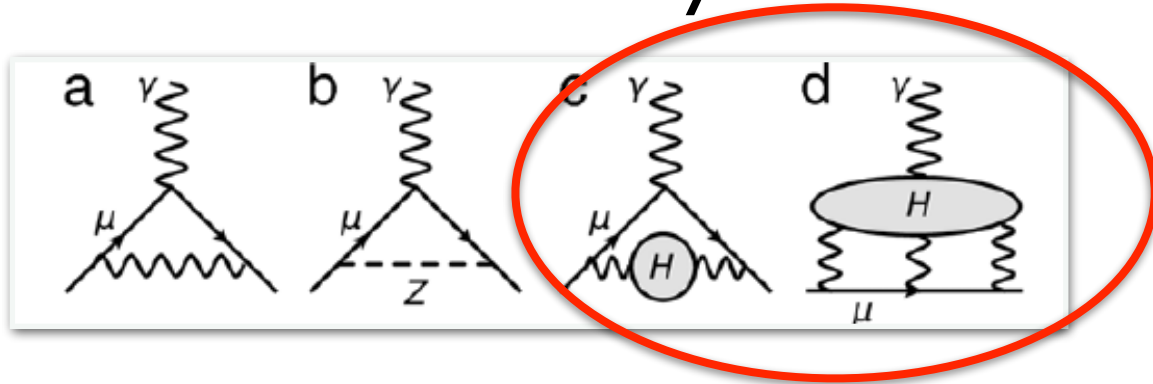
Experiment



Dominant uncertainties: ongoing efforts to improve these results using Lattice QCD

- Where are we?

SM theory



VS

Experiment



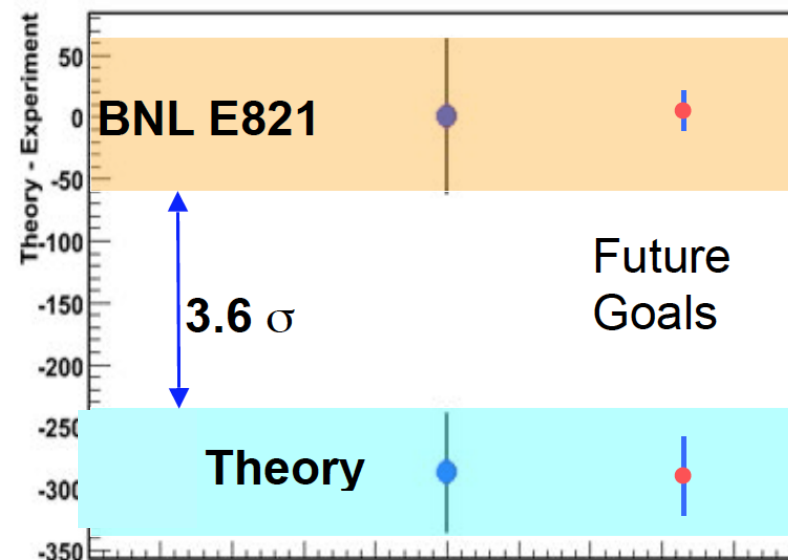
Dominant uncertainties: ongoing efforts to improve these results using Lattice QCD

$$\vec{\mu} = g \frac{e}{2mc} \vec{s}$$

$$a_\mu = (g_\mu - 2)/2$$

$a_\mu(\text{Expt})$	$=$	$116\,592\,089(54)(33) \times 10^{-11}$	BNL E821 (2006)
$a_\mu(\text{SM})$	$=$	$116\,591\,802(42)(26)(02) \times 10^{-11}$	
		$\Rightarrow \Delta a_\mu = 287(80) \times 10^{-11}$	3.6 σ discrepancy

D. Hertzog



Goal: 0.14 ppm

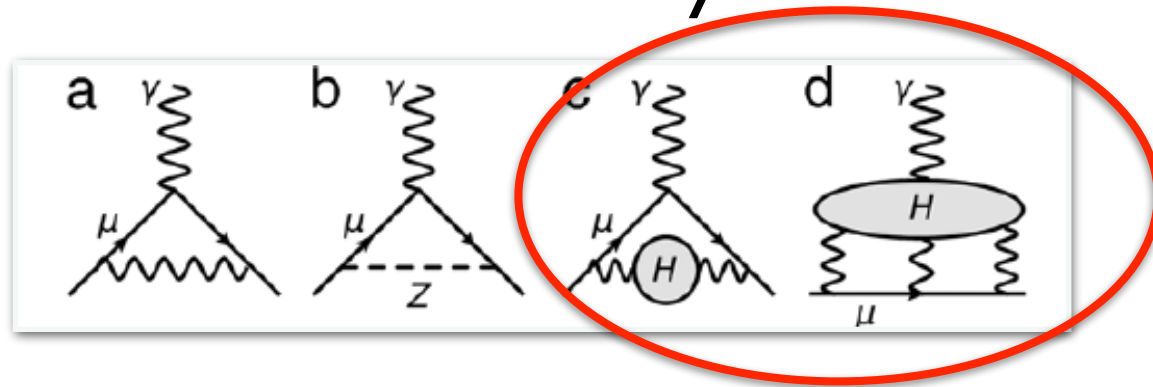
New g-2 at Fermilab (FNALE989) and J-PARCE34 will improve uncertainty factor of 4

Establish confidence in error bar

- Where are we?

Experiment

SM theory



VS

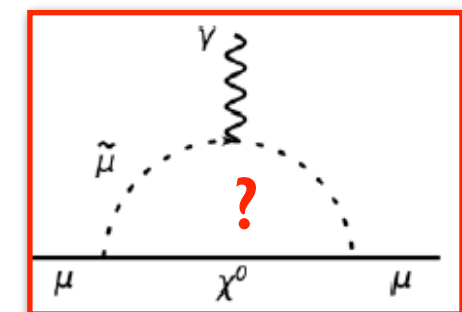
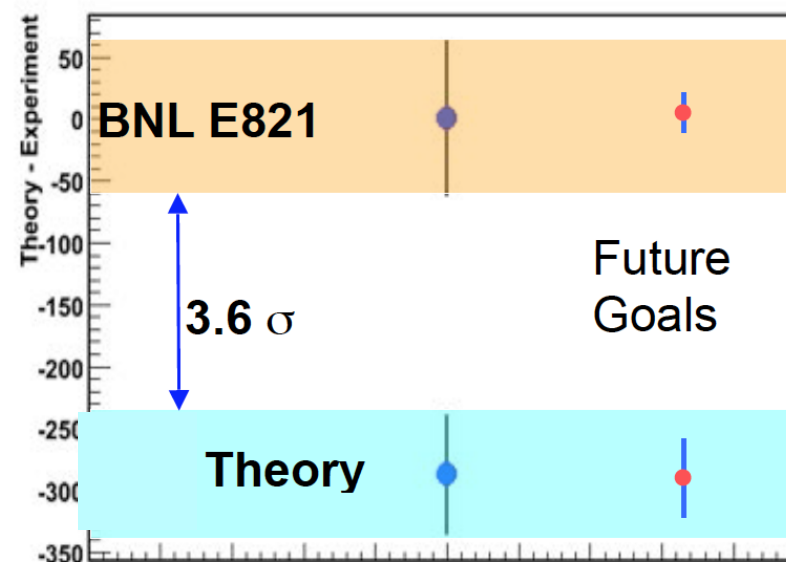


Dominant uncertainties: ongoing efforts to improve these results using Lattice QCD

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New physics?

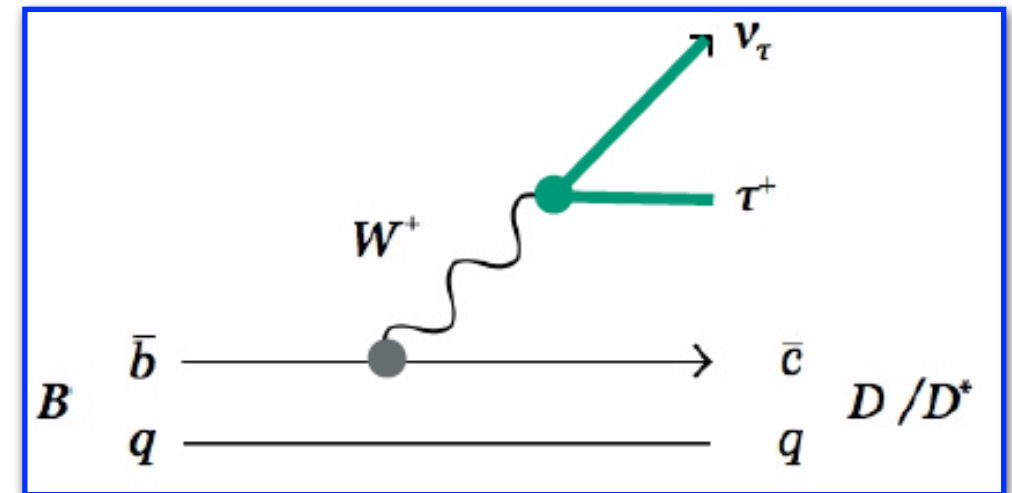
Lepton universality in B decays

Tree level, $b \rightarrow c \ell \nu$

abundant

well known in SM

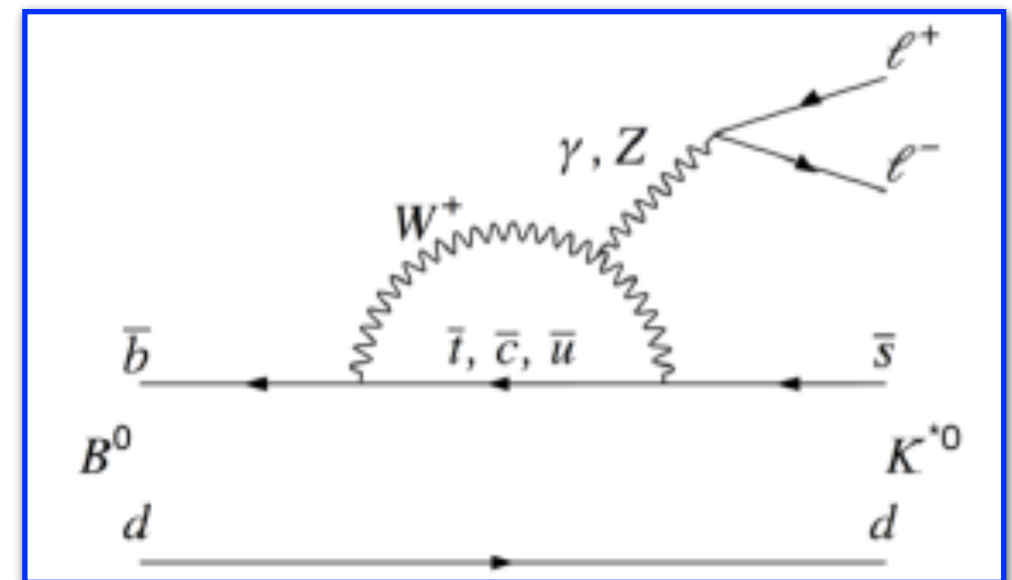
possible new physics in 3rd gen.



Loop level, $b \rightarrow s \ell \ell$

FCNC forbidden at tree-level in SM

sensitive to new physics in loops



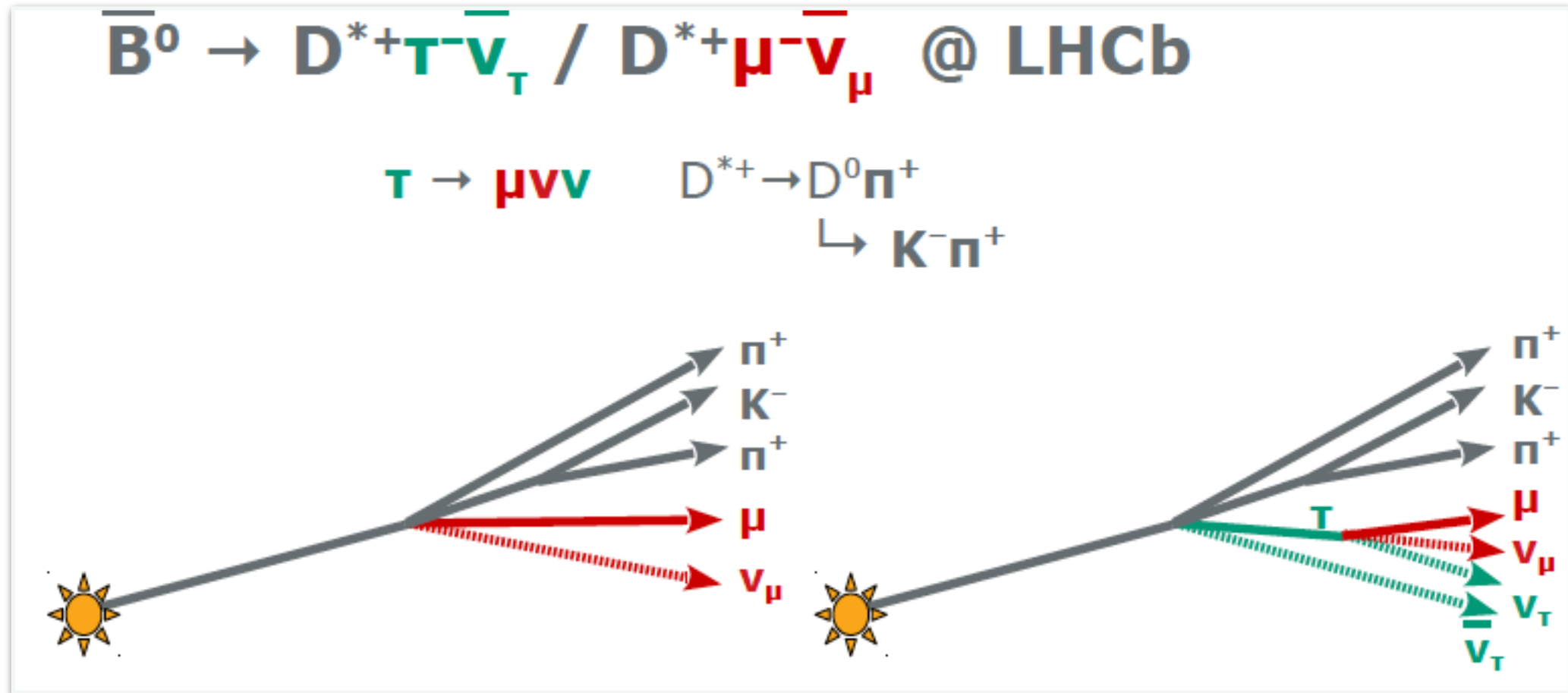
Approach : determine ratio of branching fractions

experimentally clean \rightarrow many systematics cancel

theoretically clean \rightarrow many QCD effects cancel

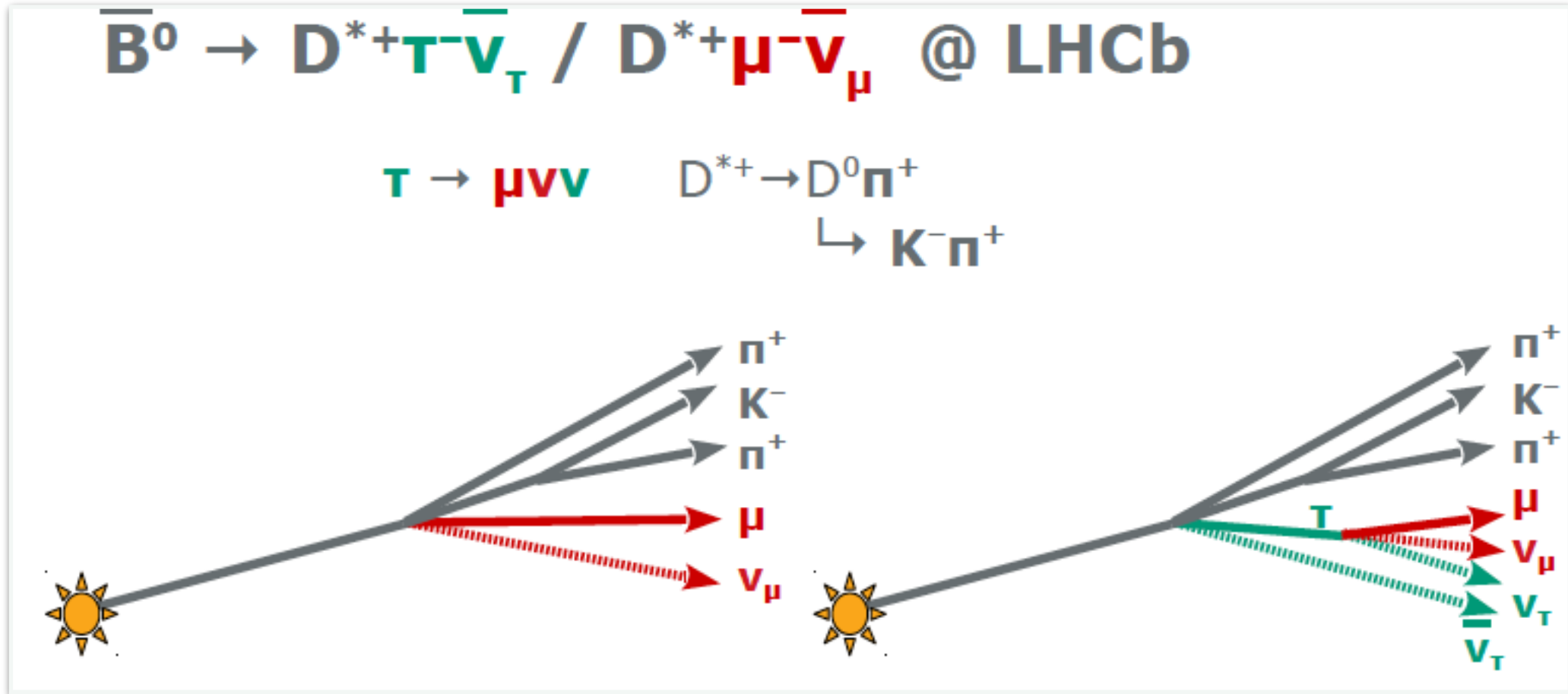
Lepton universality in B decays

R(D*)



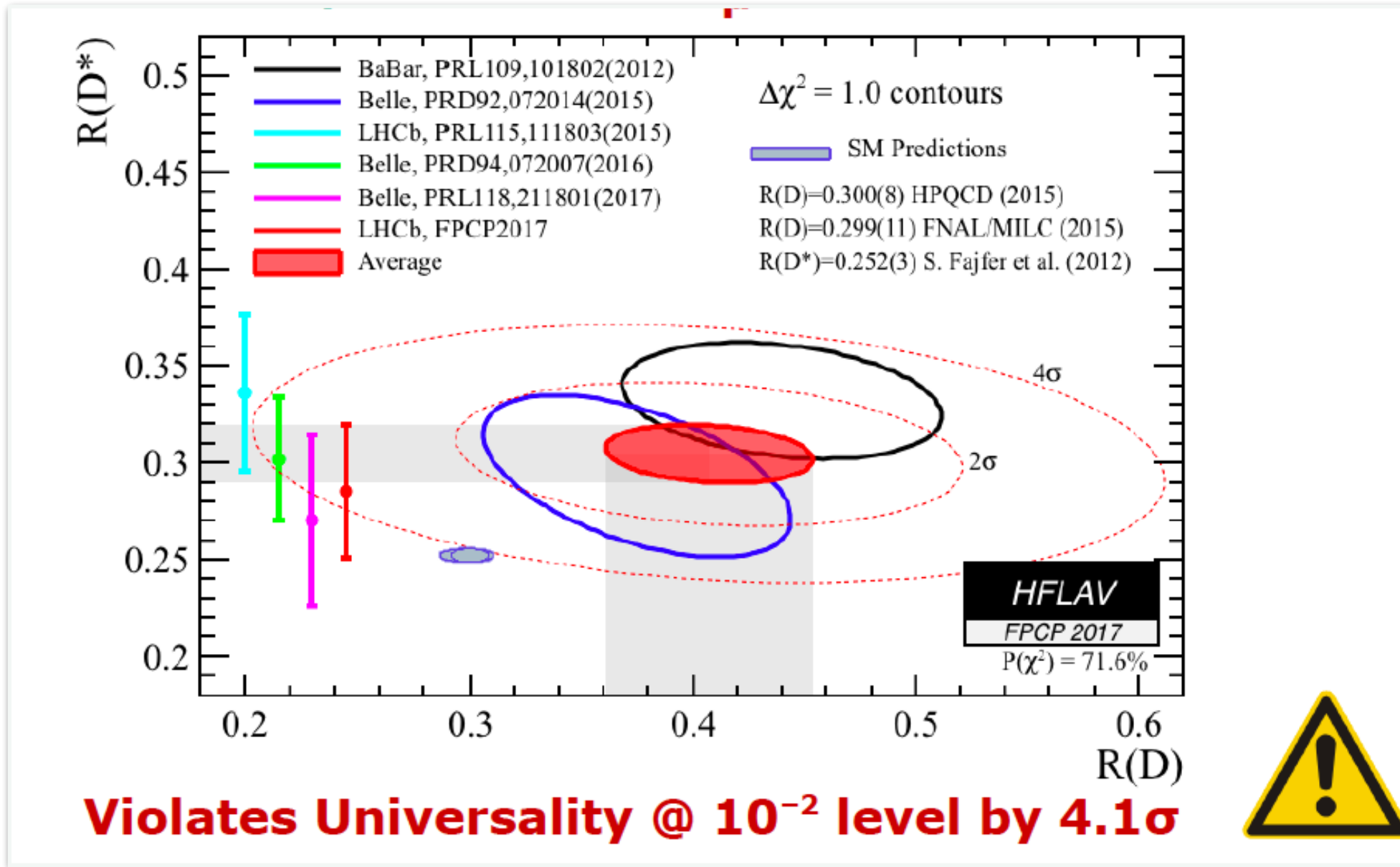
Lepton universality in B decays

$R(D^*)$



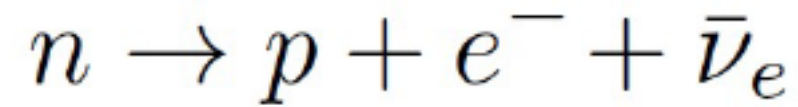
$$R_{K^{*0}} = \frac{\mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{\mathcal{B}(B^0 \rightarrow K^{*0} J/\psi (\rightarrow \mu^+ \mu^-))} / \frac{\mathcal{B}(B^0 \rightarrow K^{*0} e^+ e^-)}{\mathcal{B}(B^0 \rightarrow K^{*0} J/\psi (\rightarrow e^+ e^-))}$$

Lepton universality in B decays

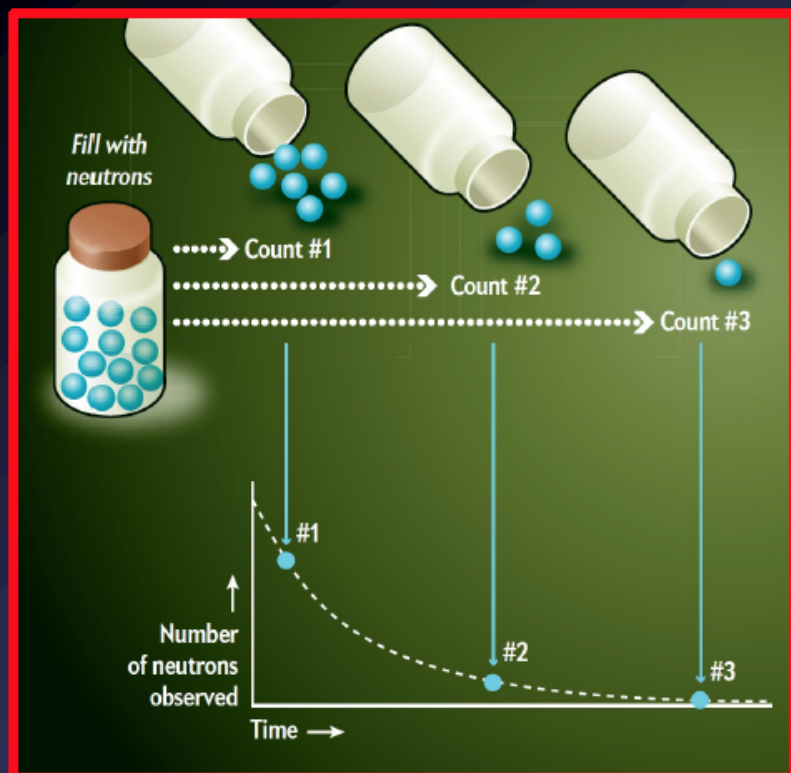


$R_{D^*}(\mu, \tau) > \text{prediction}$: **4.1σ tension**
 $R_{K^*}(\mu) < \text{prediction}$: **$\sim 4\sigma$ tension**
 (not shown here)

Neutron lifetime puzzle



Bottle experiments



Data points fit to an exponential decay

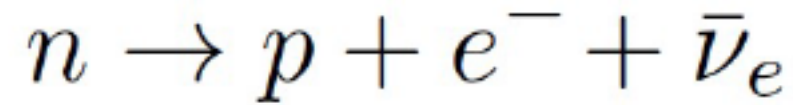
$$N = N_0 e^{-\lambda t}$$

Lifetime

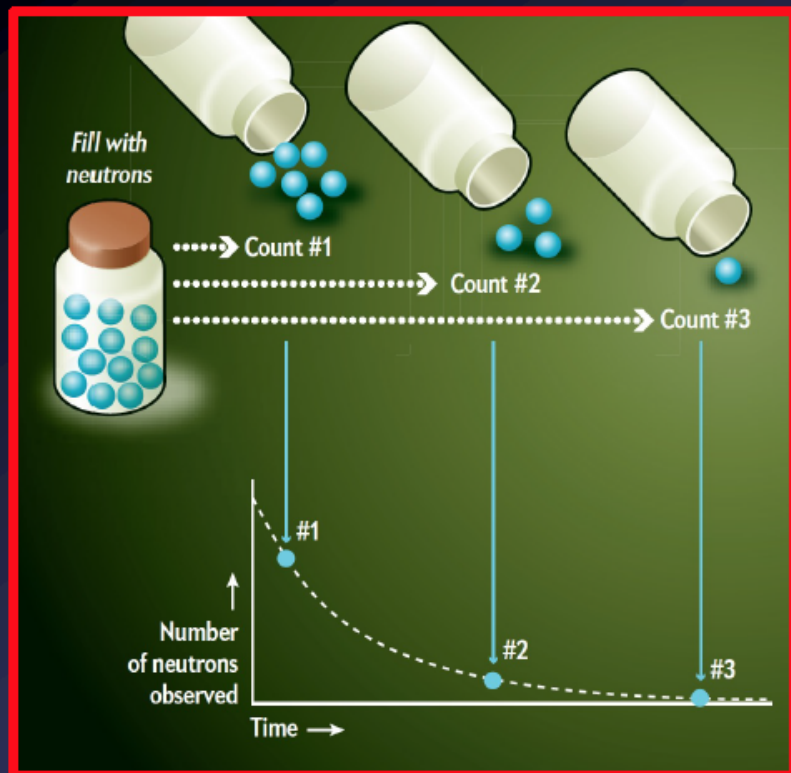
$$\tau = \frac{1}{\lambda}$$

Count the living

Neutron lifetime puzzle



Bottle experiments



Data points fit to an exponential decay

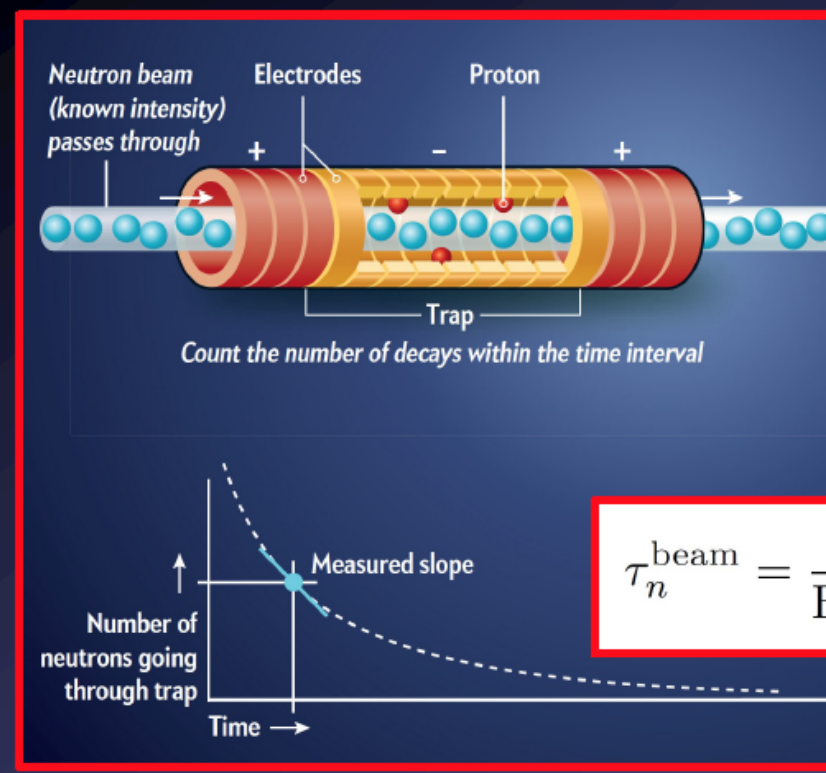
$$N = N_0 e^{-\lambda t}$$

Lifetime

$$\tau = \frac{1}{\lambda}$$

Source: <https://www.scientificamerican.com>

Beam experiments



Only the decay rate to protons is measured

$$\frac{dN}{dt} = -\lambda N$$

$$\tau_n^{\text{beam}} = \frac{\tau_n}{\text{Br}(n \rightarrow p + \text{anything})}$$

Source: <https://www.scientificamerican.com>

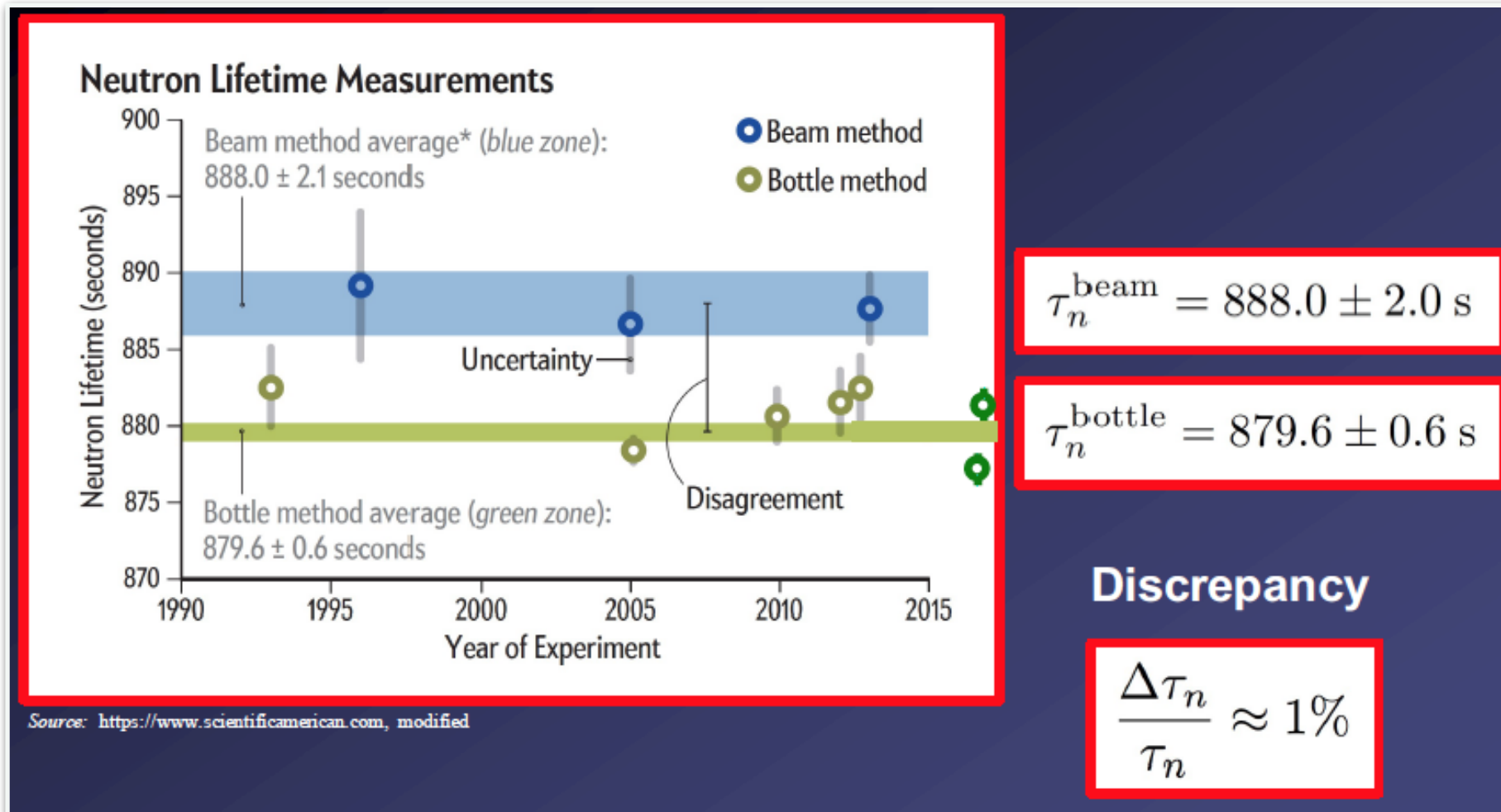
Count the living

Count the dead

Neutron lifetime puzzle



Neutron lifetime puzzle



$$\text{Br}(n \rightarrow p + \text{anything}) \approx 99\%$$

Remaining 1% :

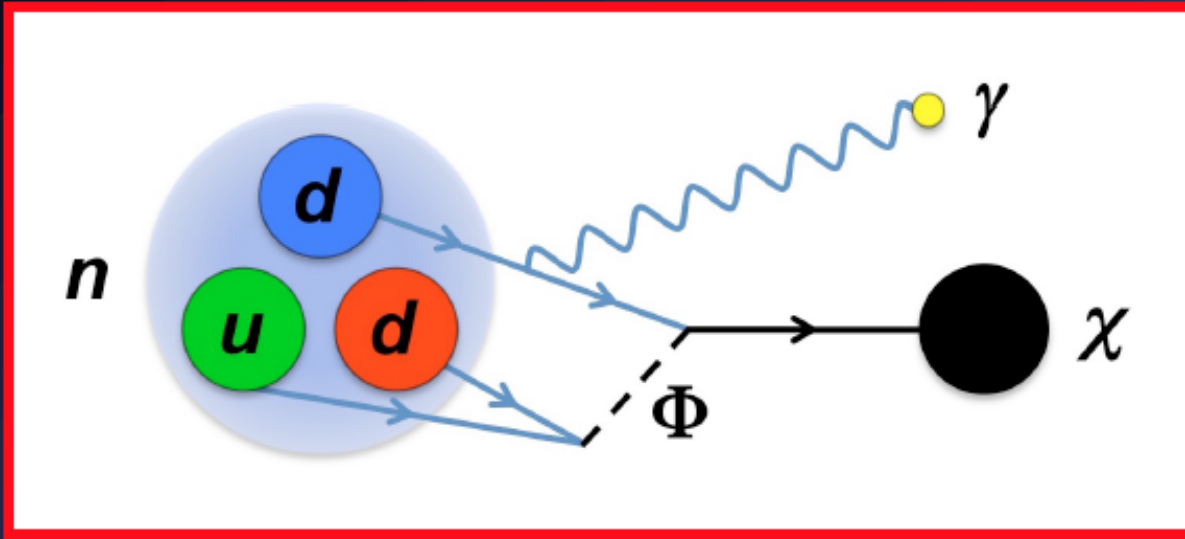
→ $n \rightarrow \text{SM particles (other than } p)$ ✗

→ $n \rightarrow \text{dark particle(s) + SM particle(s)}$ ✓

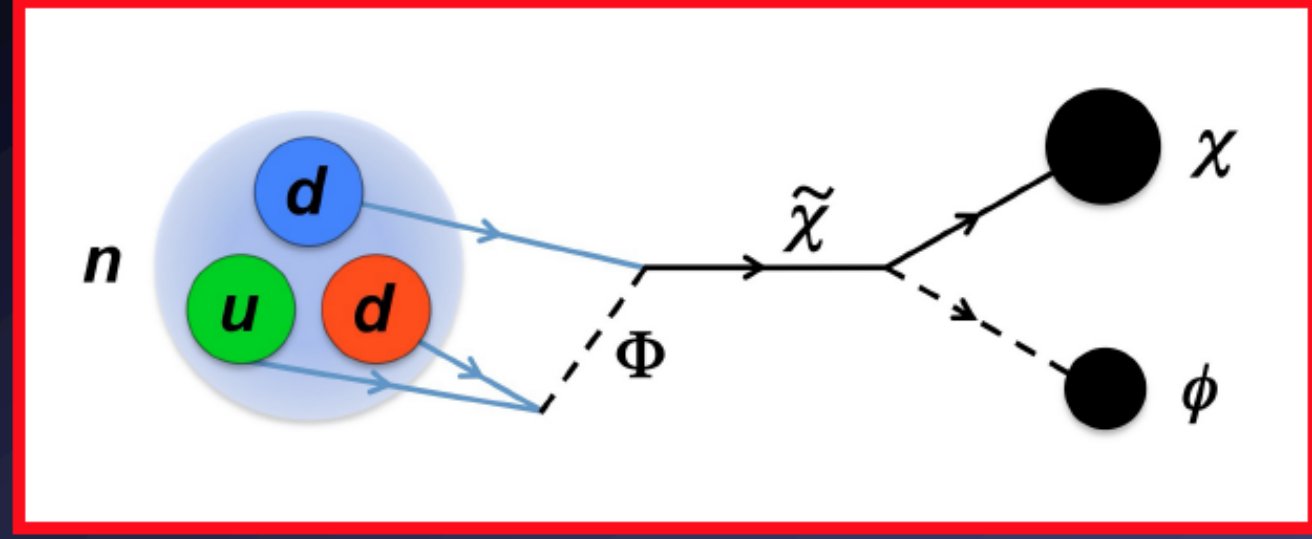
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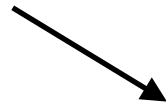
Neutron \longrightarrow dark particle + photon



Neutron \longrightarrow two dark particles



From ${}^9\text{Be}$ stability



Dark particle mass

$$937.900 \text{ MeV} < m_\chi < 939.565 \text{ MeV}$$

Dark decay photon energy

$$0 < E_\gamma < 1.664 \text{ MeV}$$

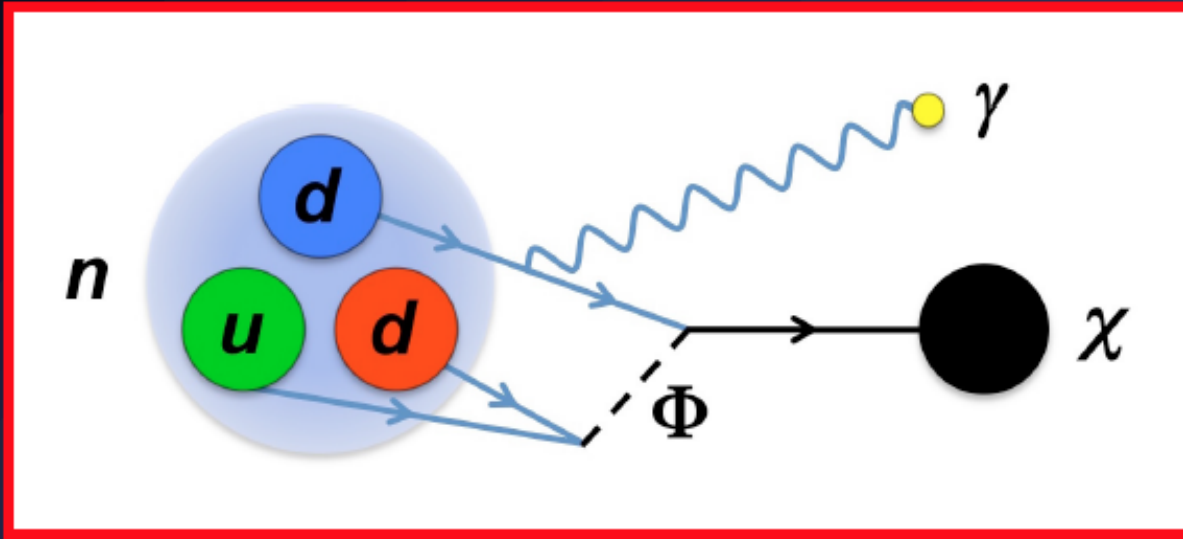
Constraints on masses

$$937.900 \text{ MeV} < m_{\tilde{\chi}}$$

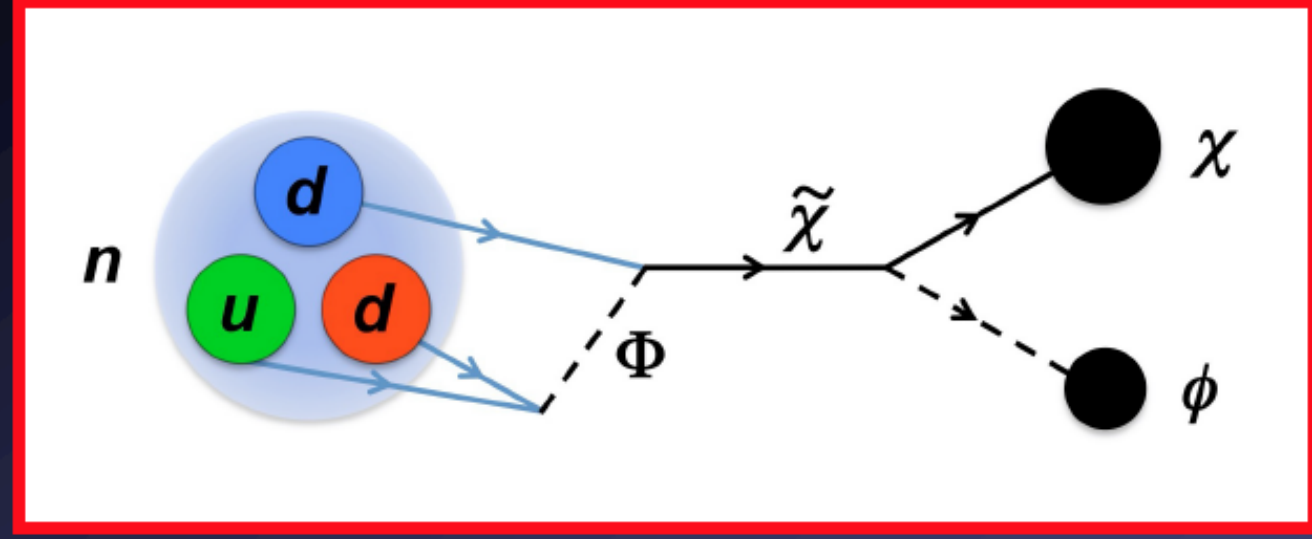
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Neutron lifetime puzzle

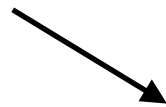
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$$0 < E_\gamma < 1.664 \text{ MeV}$$

Constraints on masses

$$937.900 \text{ MeV} < m_{\tilde{\chi}}$$

$$937.900 \text{ MeV} < m_\chi + m_\phi < 939.565 \text{ MeV}$$



“Missing Neutrons May Lead a Secret Life as Dark Matter”,

C. Moskowitz, Scientific American (January 29, 2018)