

**“Gluon distribution,
fluctuations and saturation at
an Electron Ion Collider”**

Jefferson Lab

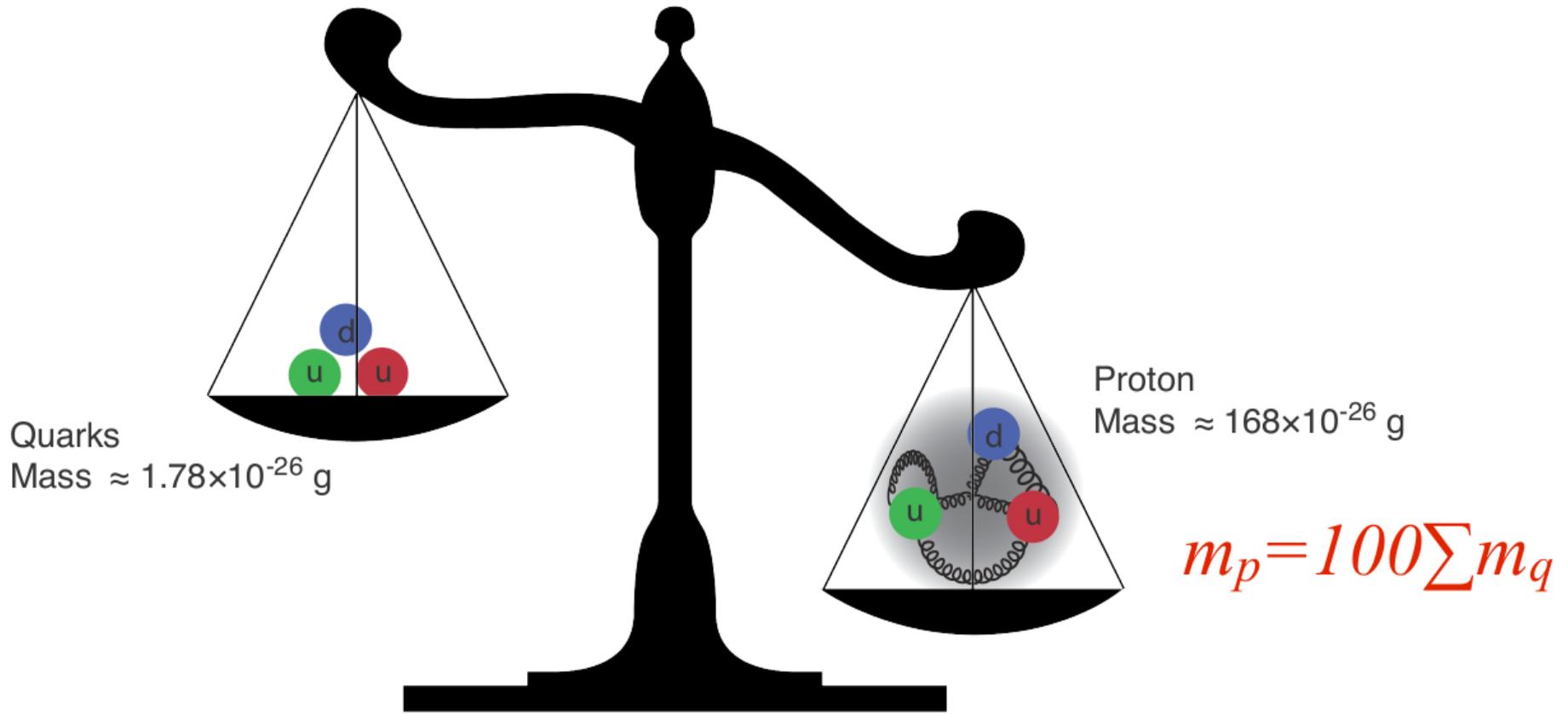
June 16, 2017

Tobias Toll

SHIV NADAR UNIVERSITY



The Hadron Mass Puzzle



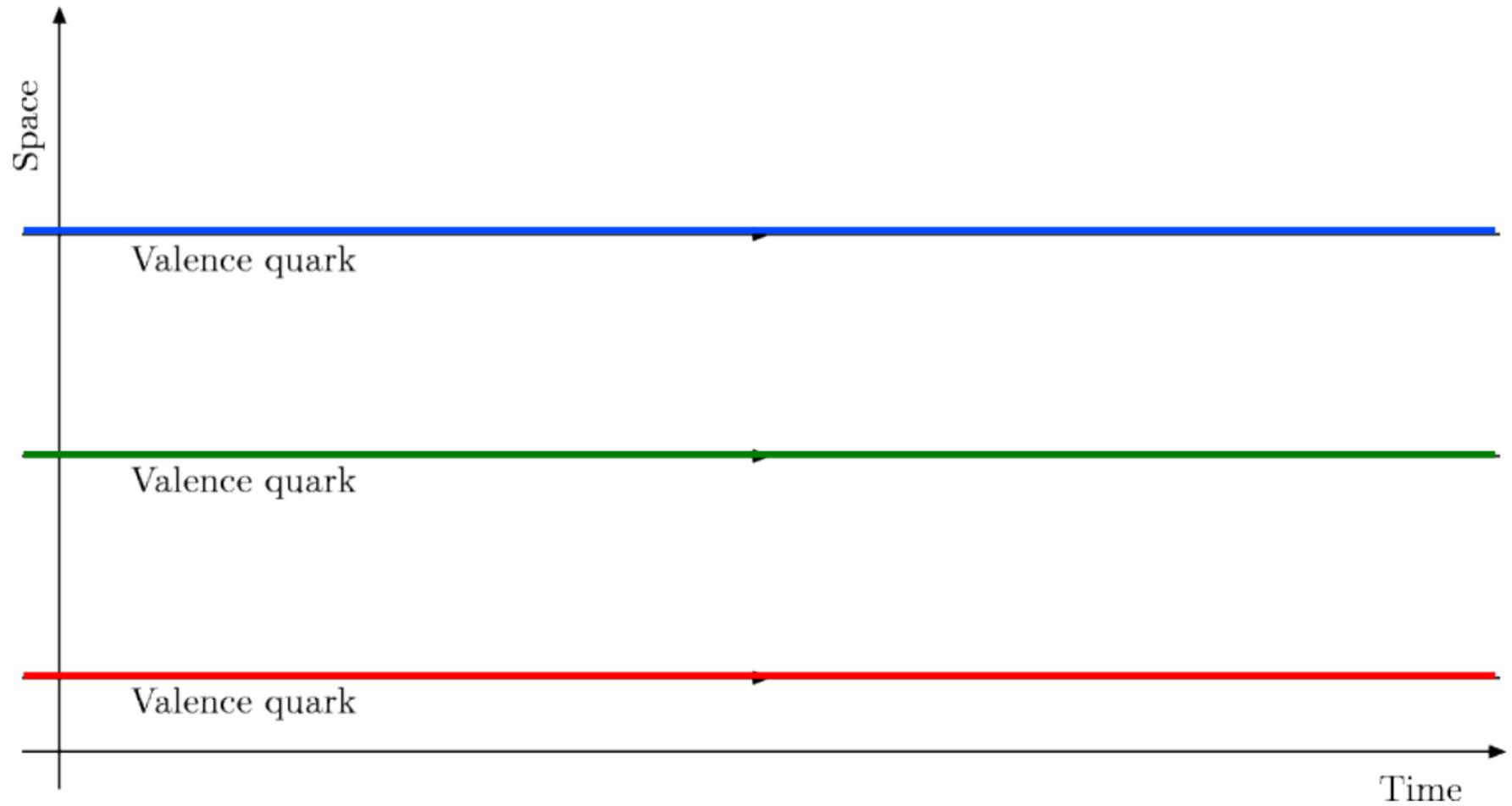
Quarks carry only $\sim 1\%$ of proton mass

Answer can come only from **Quantum Chromodynamics**

Gluons make up 98% of the visible mass of the universe!

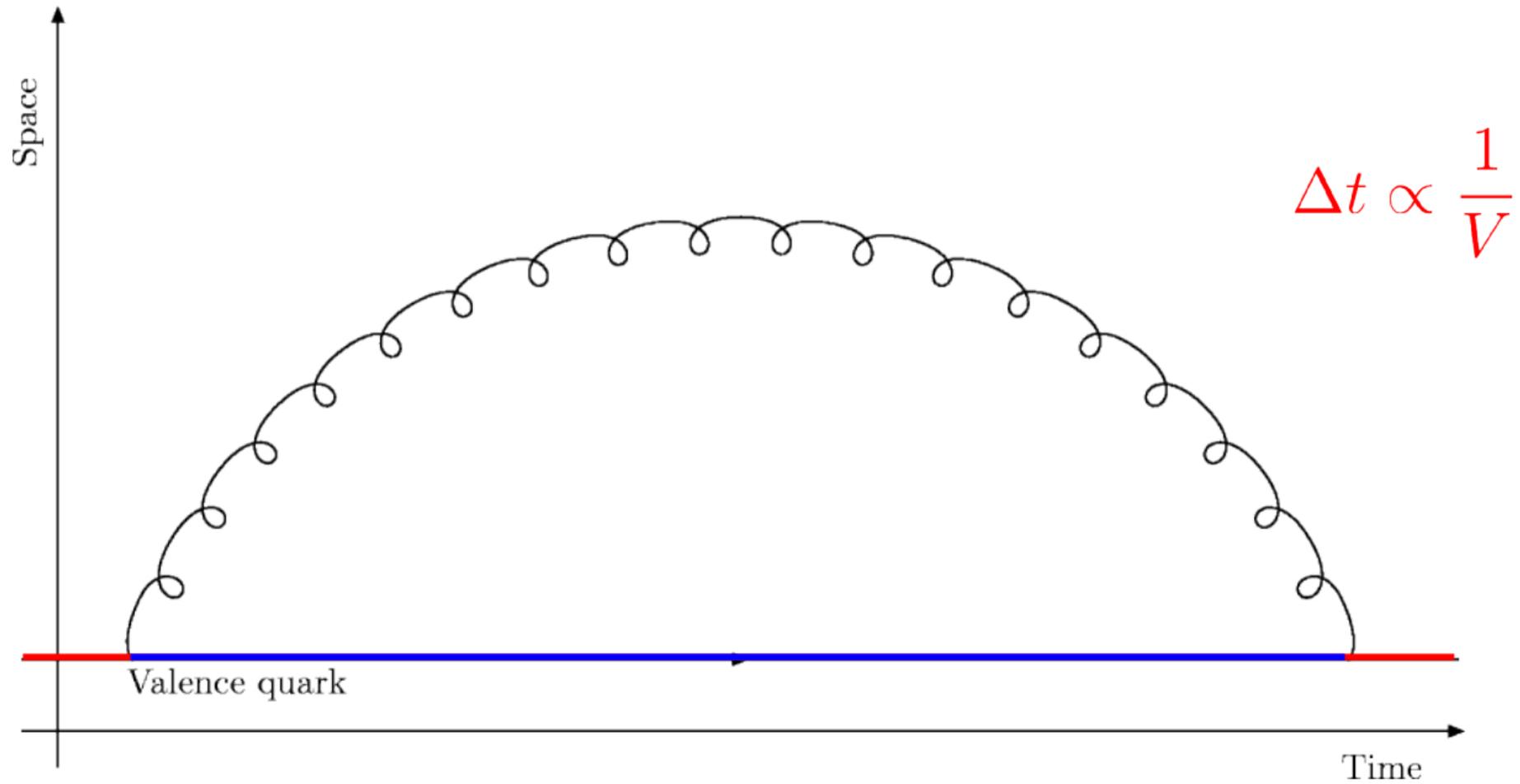
Our Understanding of a Proton

$$x \approx 0.33$$



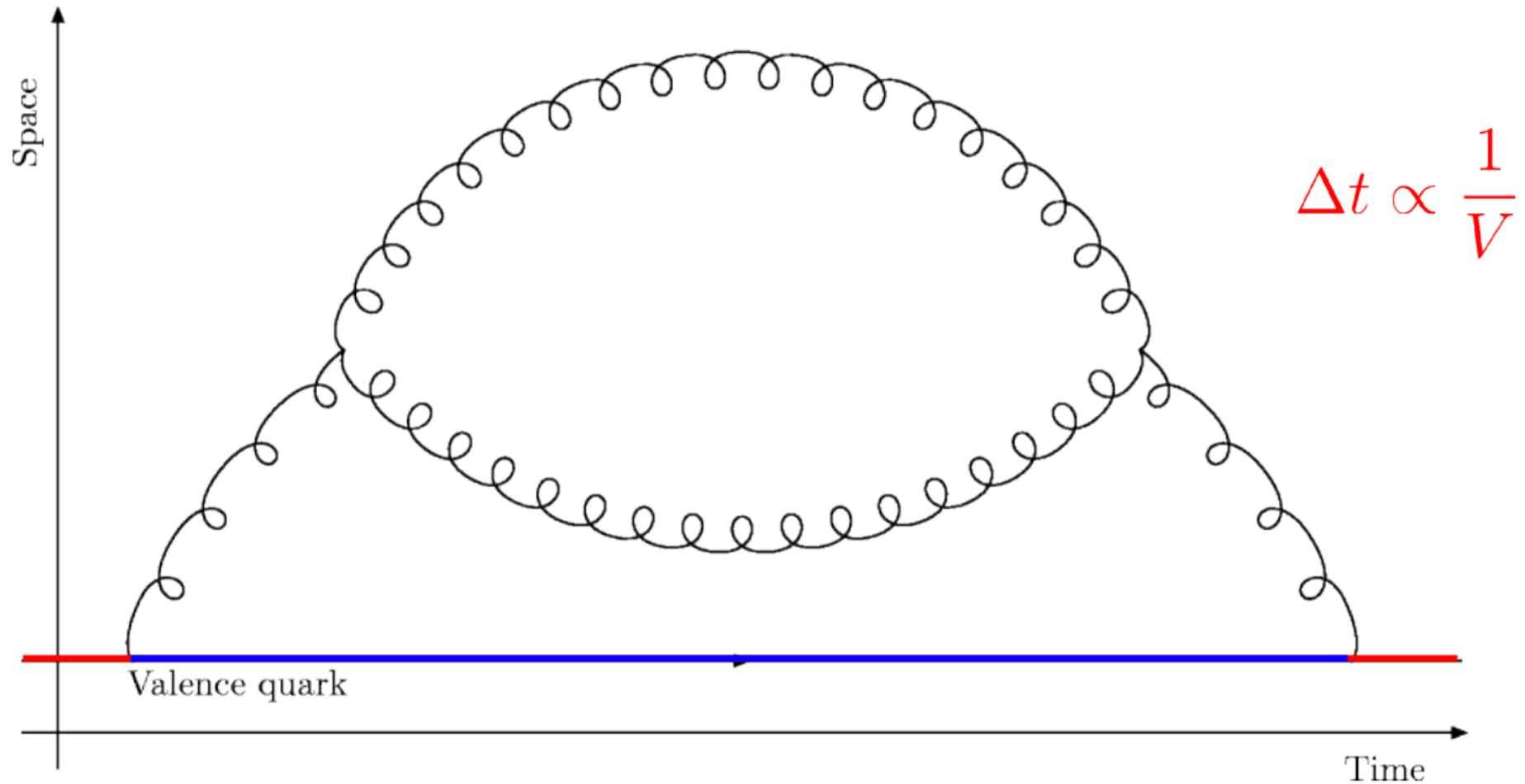
Our Understanding of a Proton

$$x \approx 0.17$$



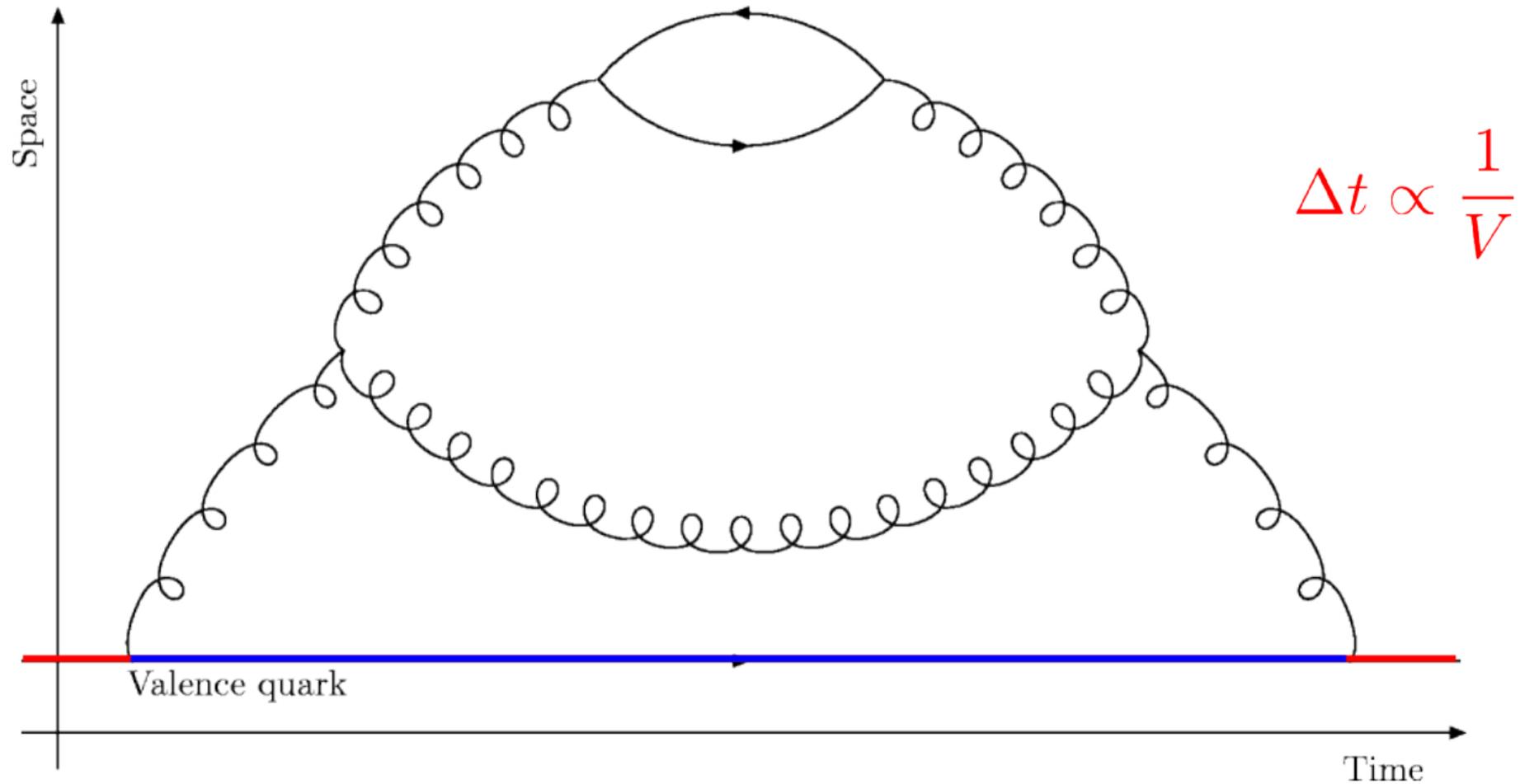
Our Understanding of a Proton

$$x \approx 0.11$$

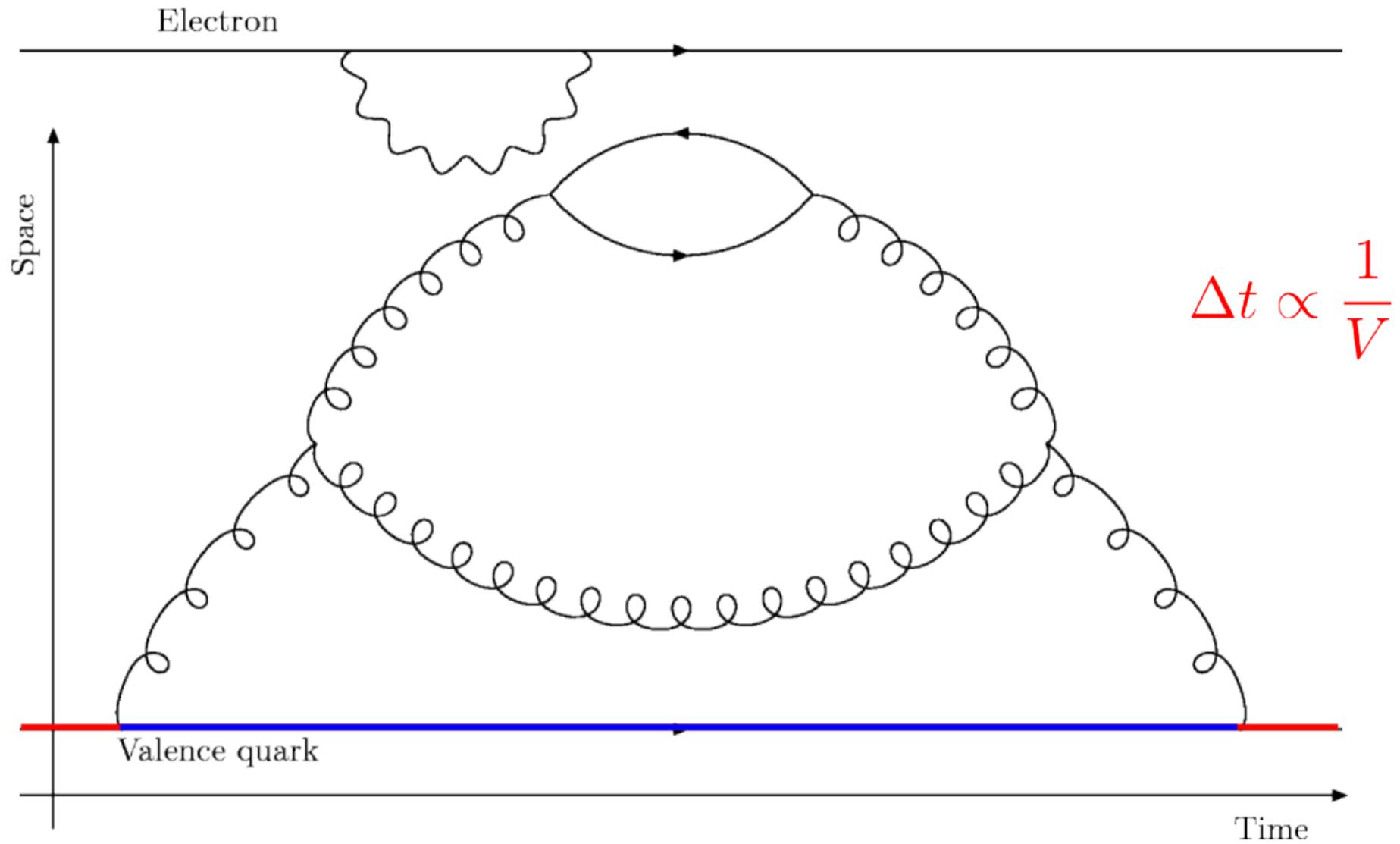


Our Understanding of a Proton

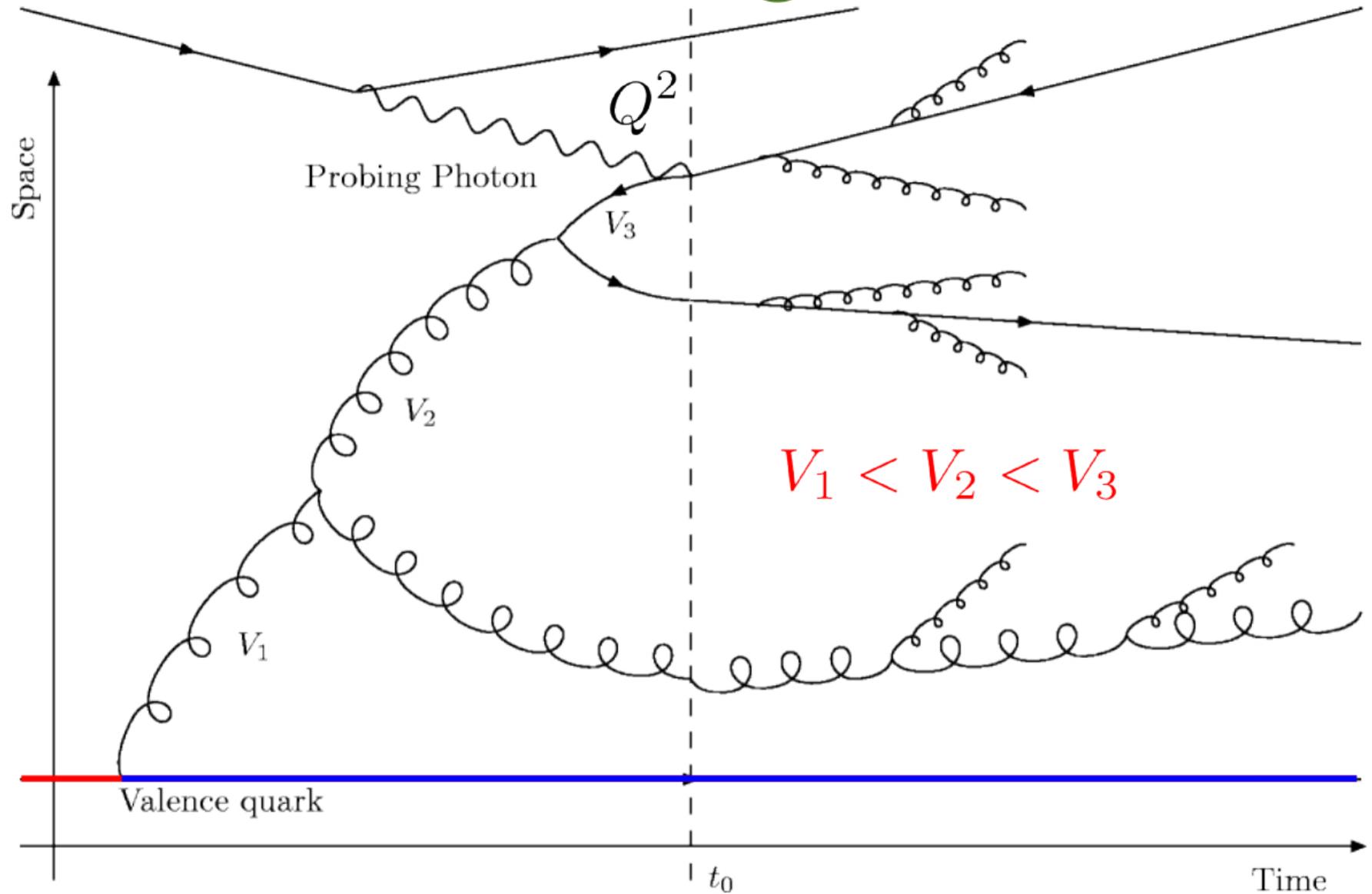
$$x \approx 0.08$$



Our Understanding of a Proton

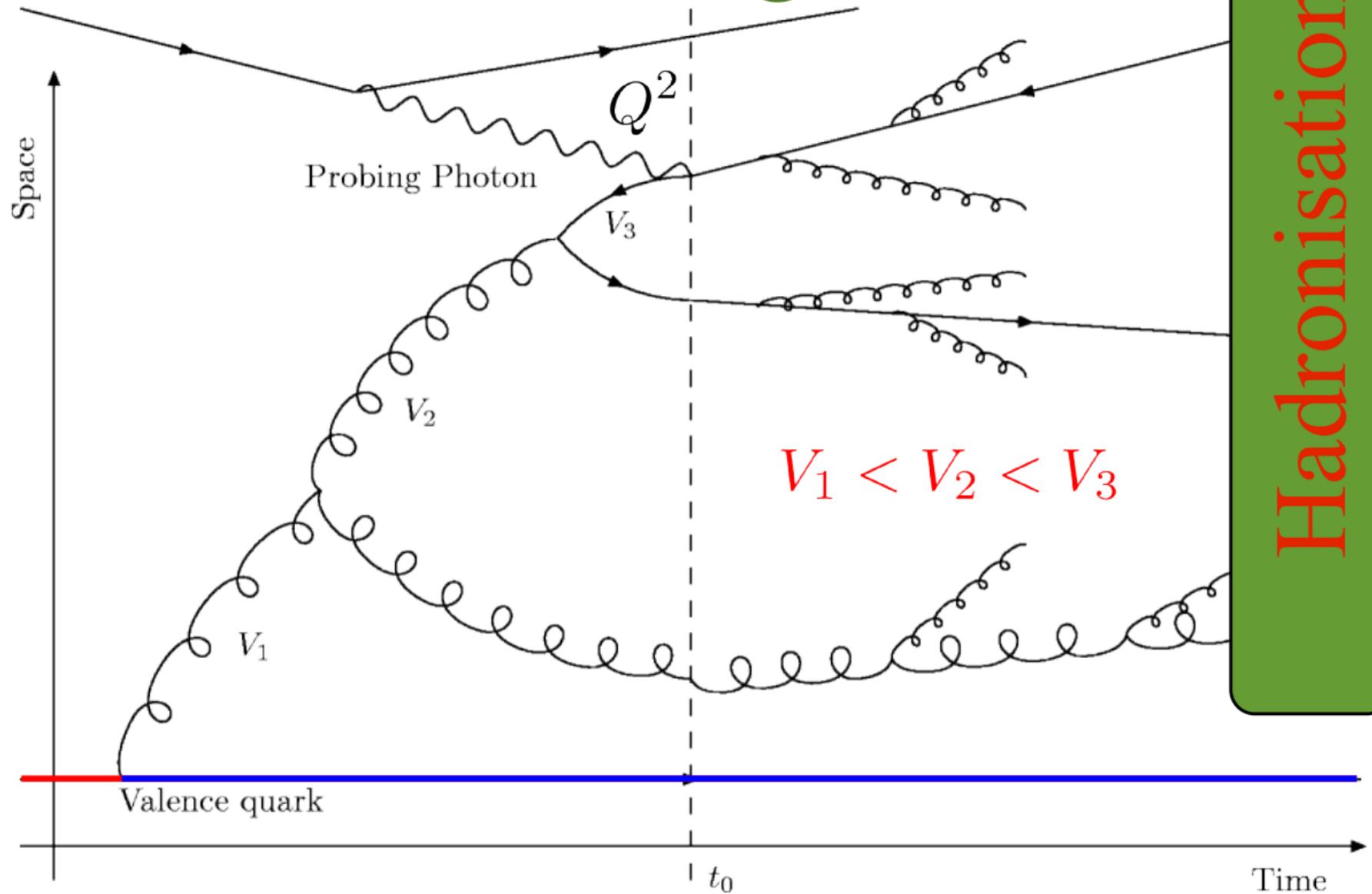


Our Understanding of a Proton



Dokshitzer **Gribov** **Lipatov** **Altarelli** **Parisi** DGLAP

Our Understanding of a Proton

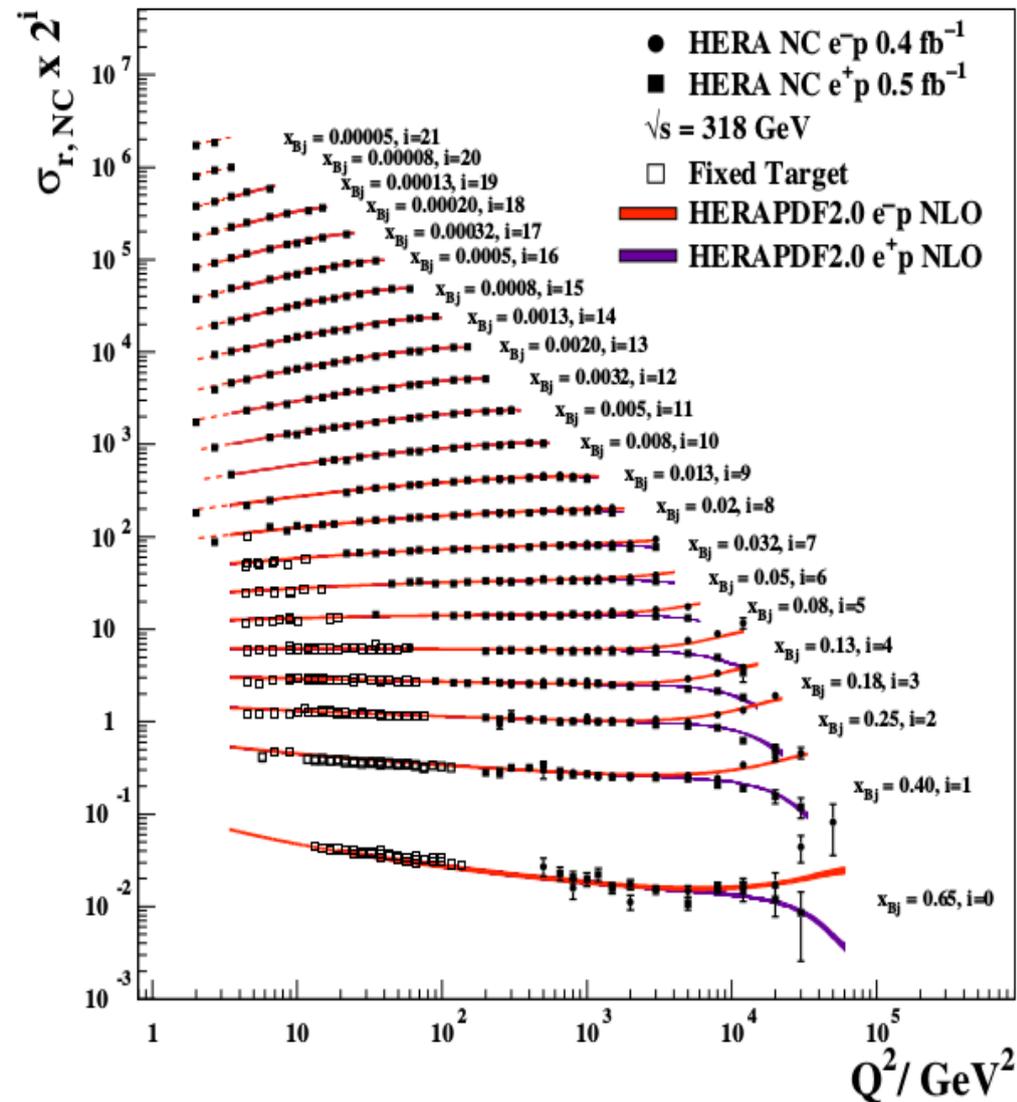


Dokshitzer **Gribov** **Lipatov** **Altarelli** **Parisi** DGLAP

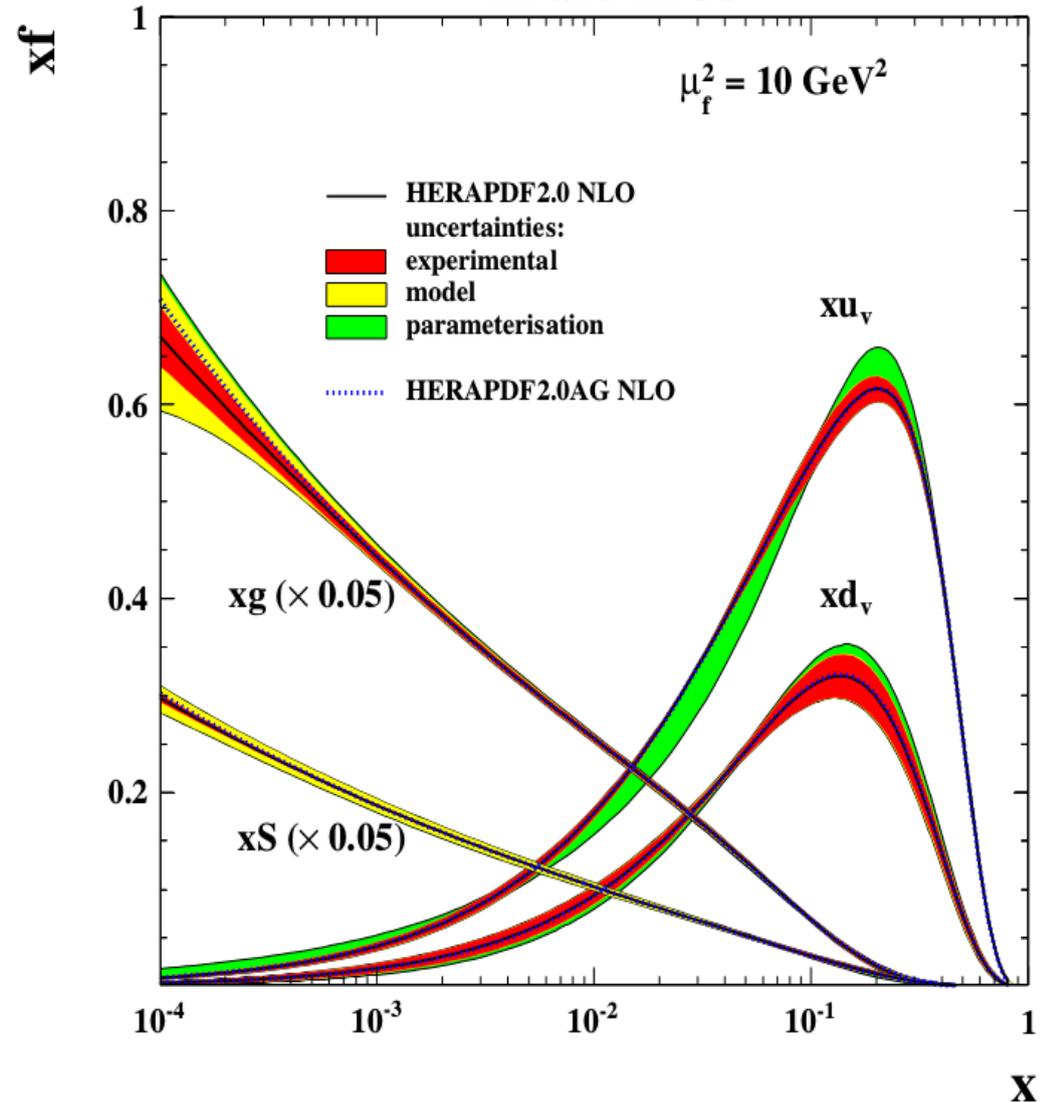
DGLAP in e^+p collisions at HERA

$$\sigma_r(x, Q^2) = F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2)$$

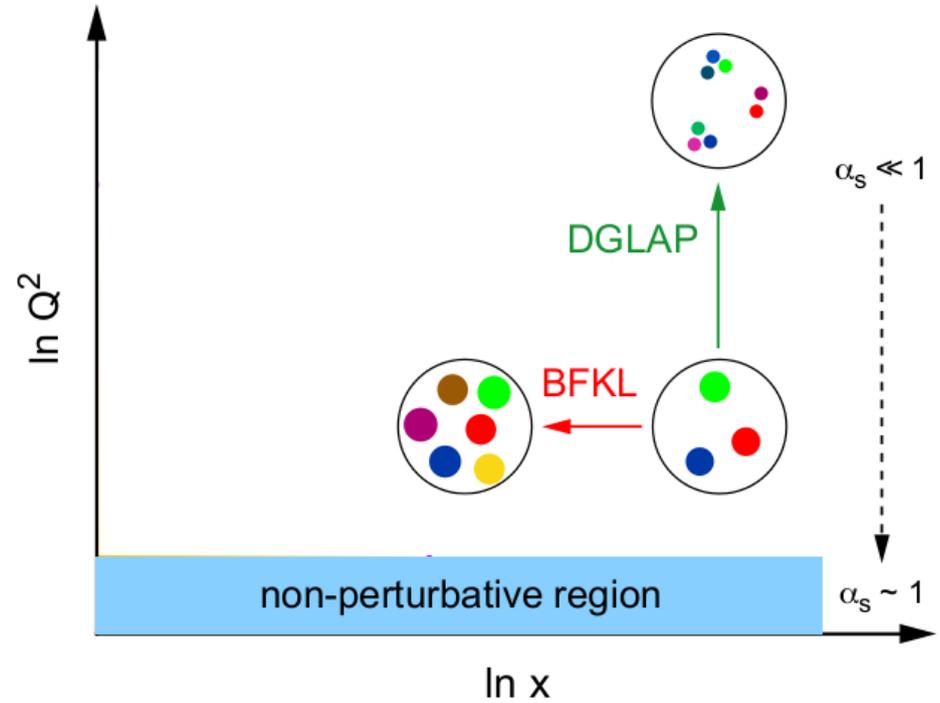
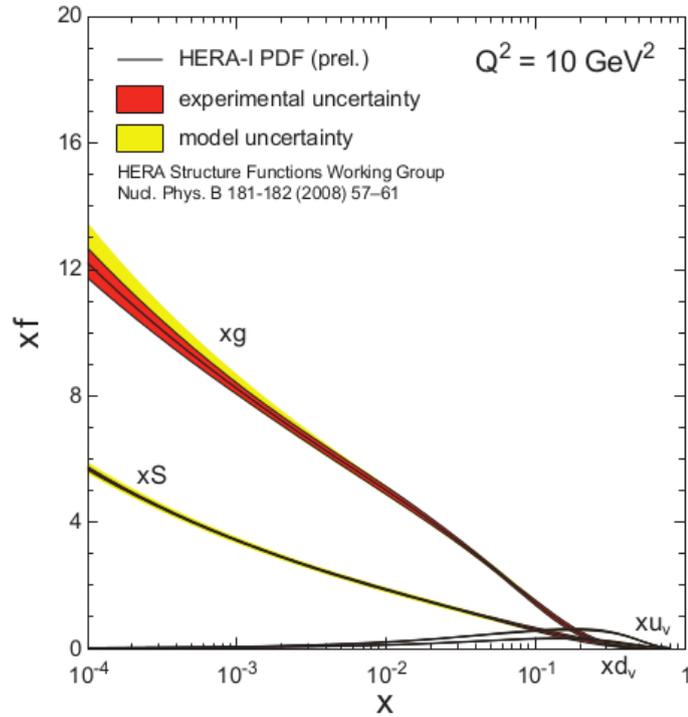
H1 and ZEUS



H1 and ZEUS



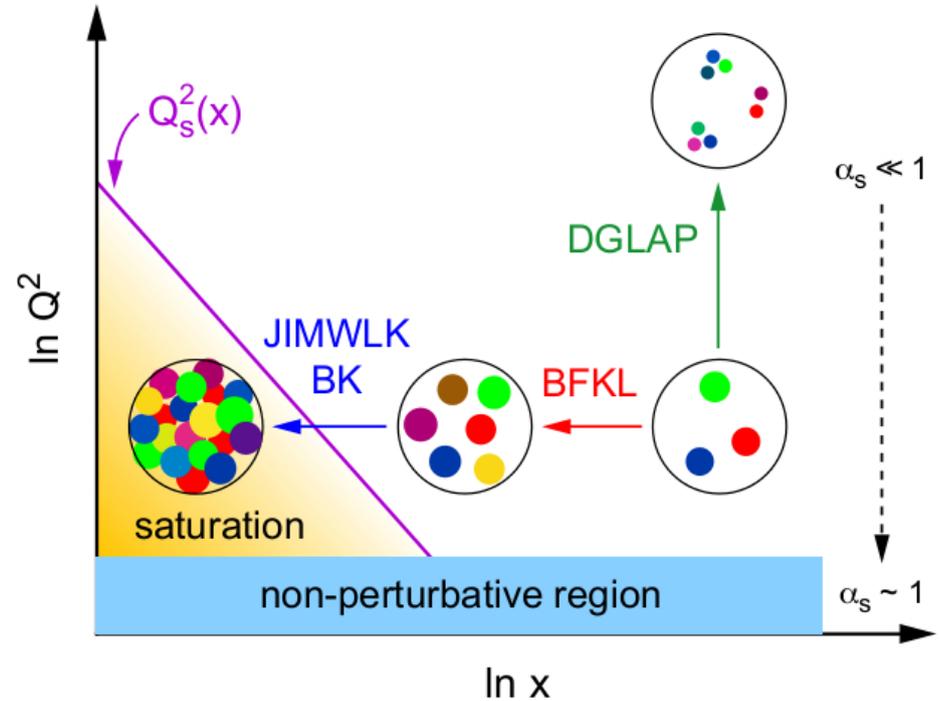
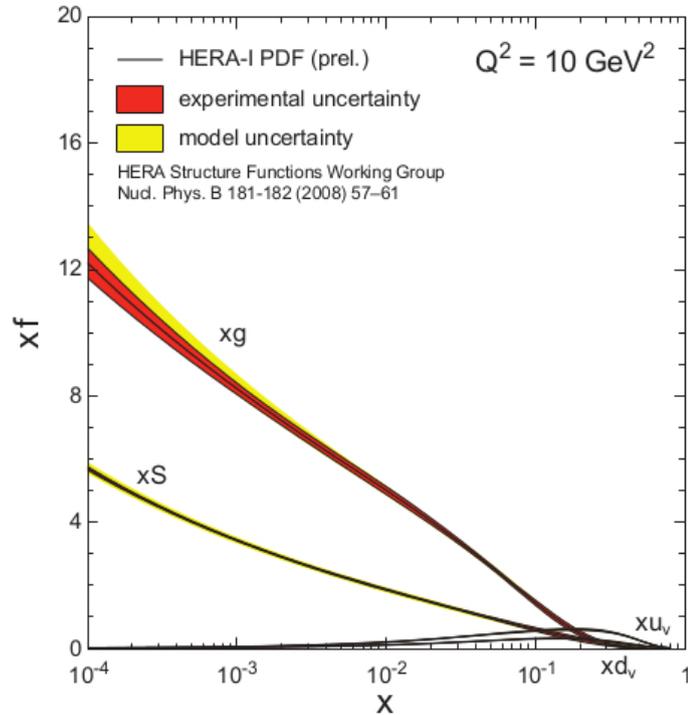
The small- x Puzzle



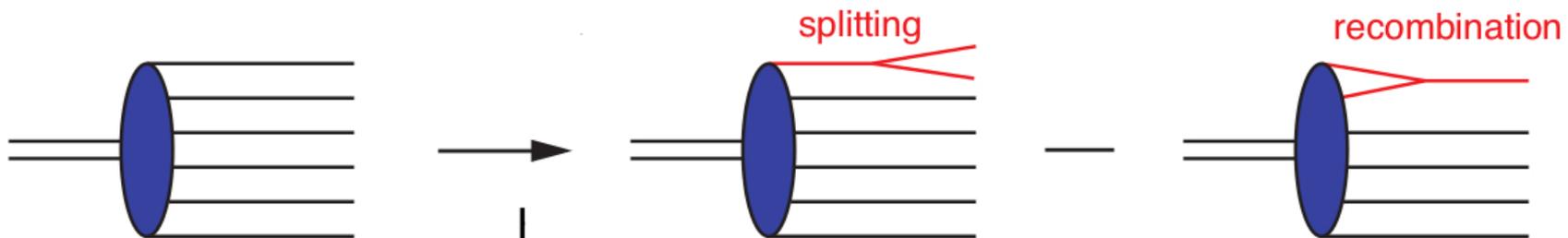
- Gluons dominate the PDFs at small- to intermediate- x ($x < 0.1$)
- ➔ Rapid rise in gluons described naturally by linear DGLAP evolution equations



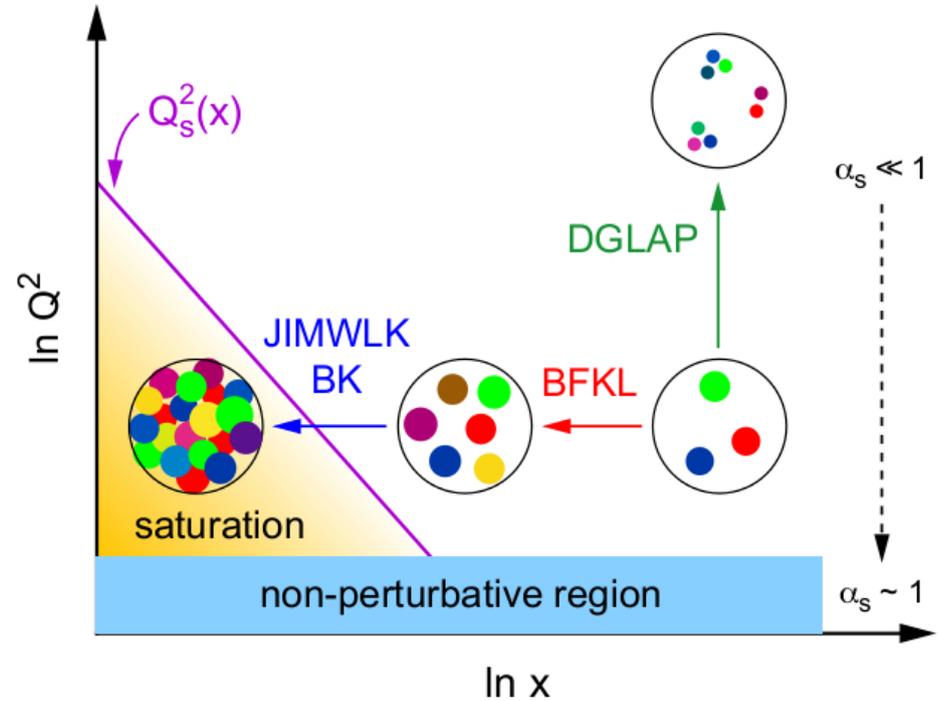
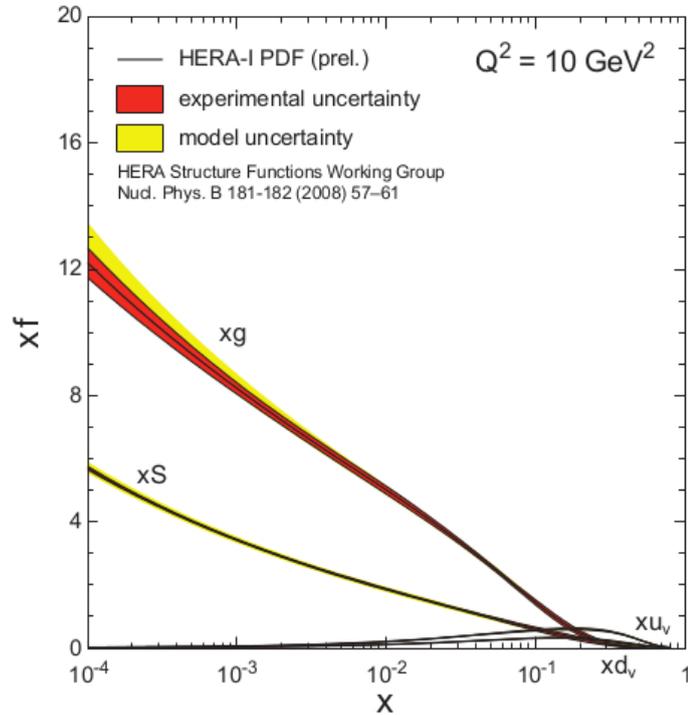
The small- x Puzzle



- Gluons dominate the PDFs at small- to intermediate- x ($x < 0.1$)
 - ➔ Rapid rise in gluons described naturally by linear DGLAP evolution equations
 - ➔ This rise cannot increase forever - gluon cross-section > proton cross-section!
 - non-linear pQCD evolution equations provide a natural way to tame this growth and lead to a **saturation** of gluons, characterised by the saturation scale $Q_s^2(x)$



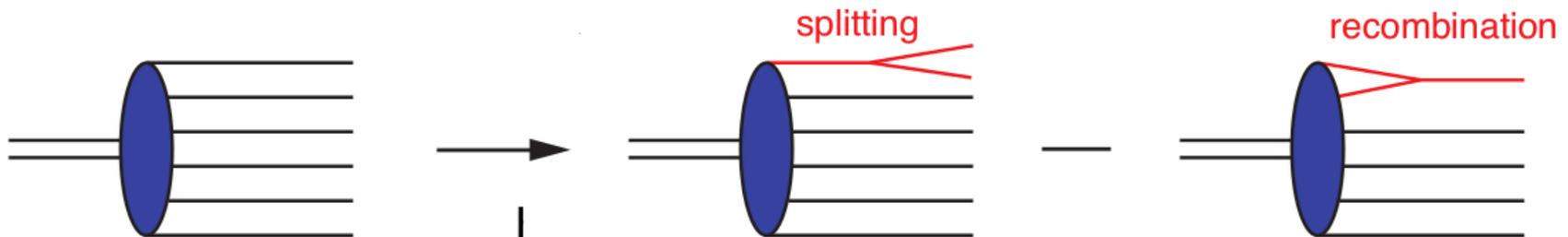
The small- x Puzzle



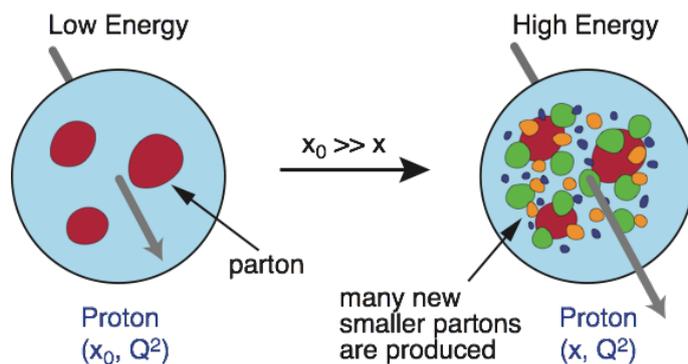
- Gluons dominate the PDFs at small- to intermediate- x ($x < 0.1$)
 - ➔ Rapid rise in gluons described naturally by linear DGLAP evolution equations

Saturation in the gluon density is not observed in the gluon distribution at HERA, too high x

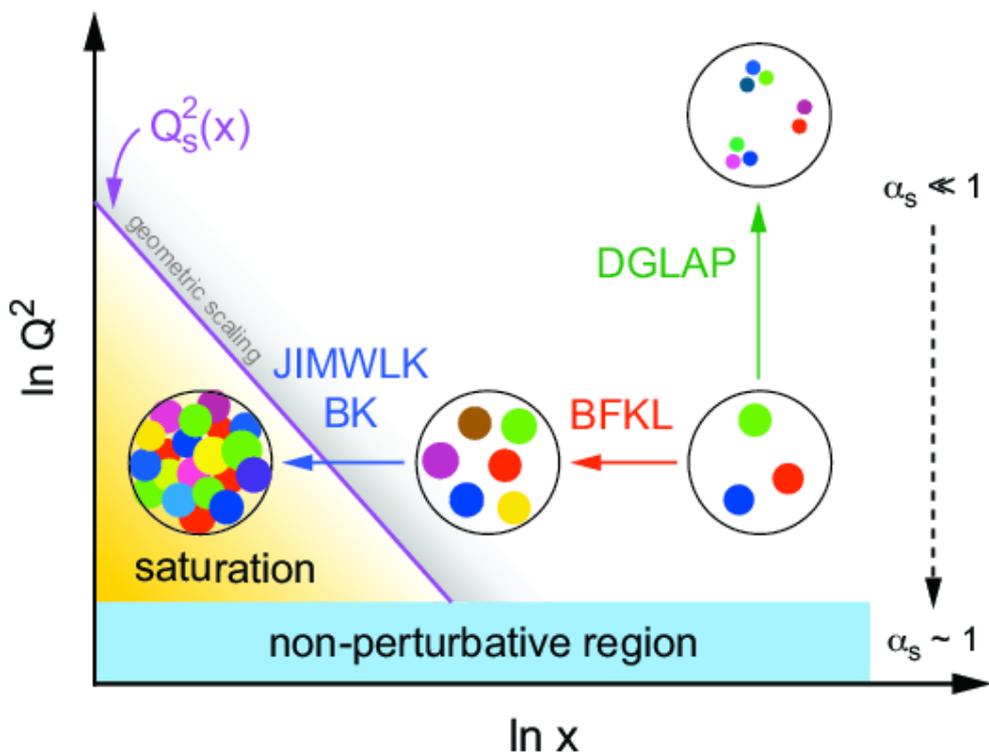
lead to a saturation of gluons, characterised by the saturation scale $Q_s(x)$



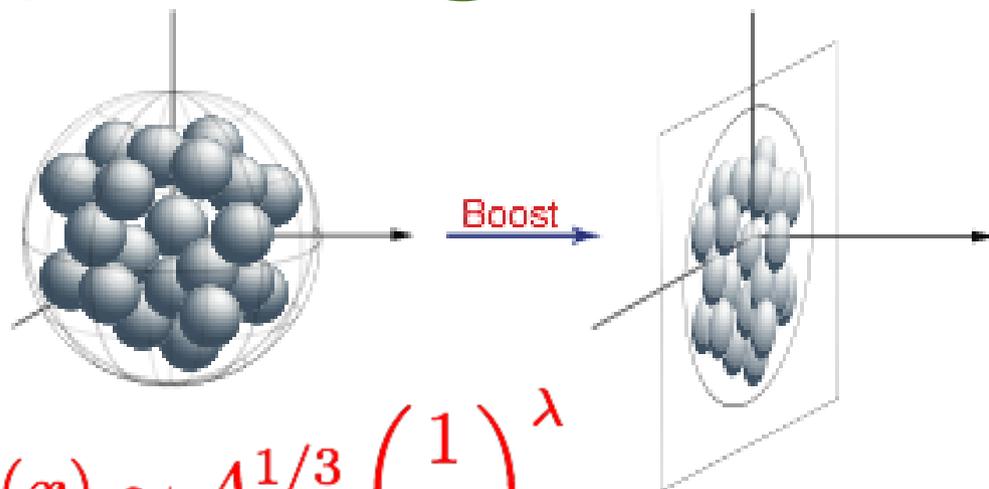
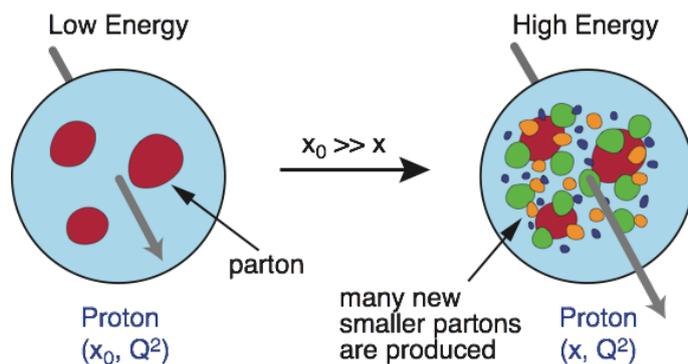
Saturation at EIC



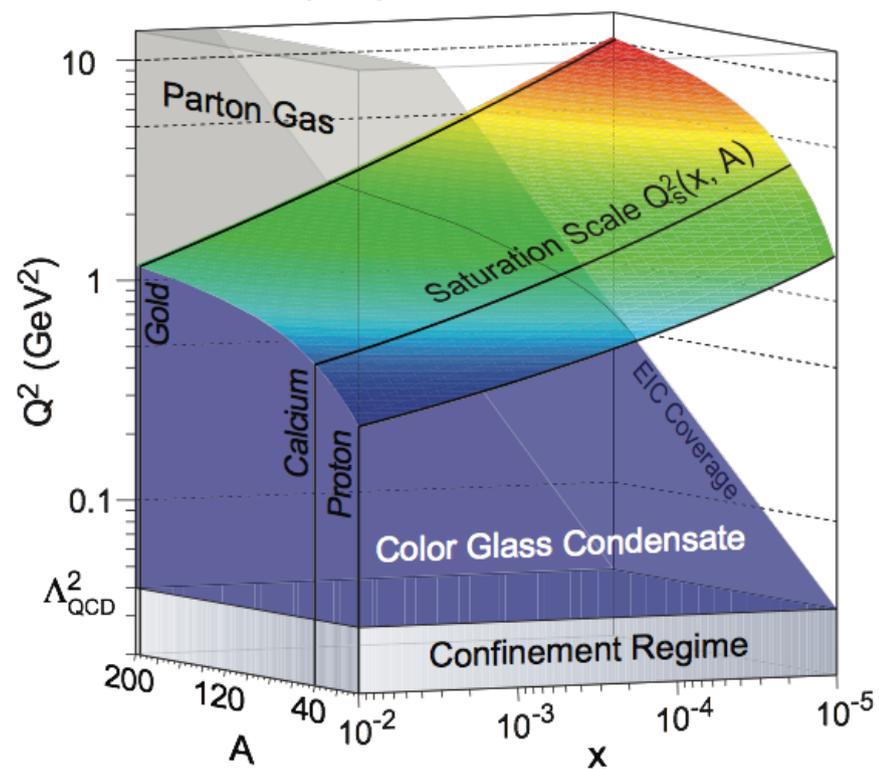
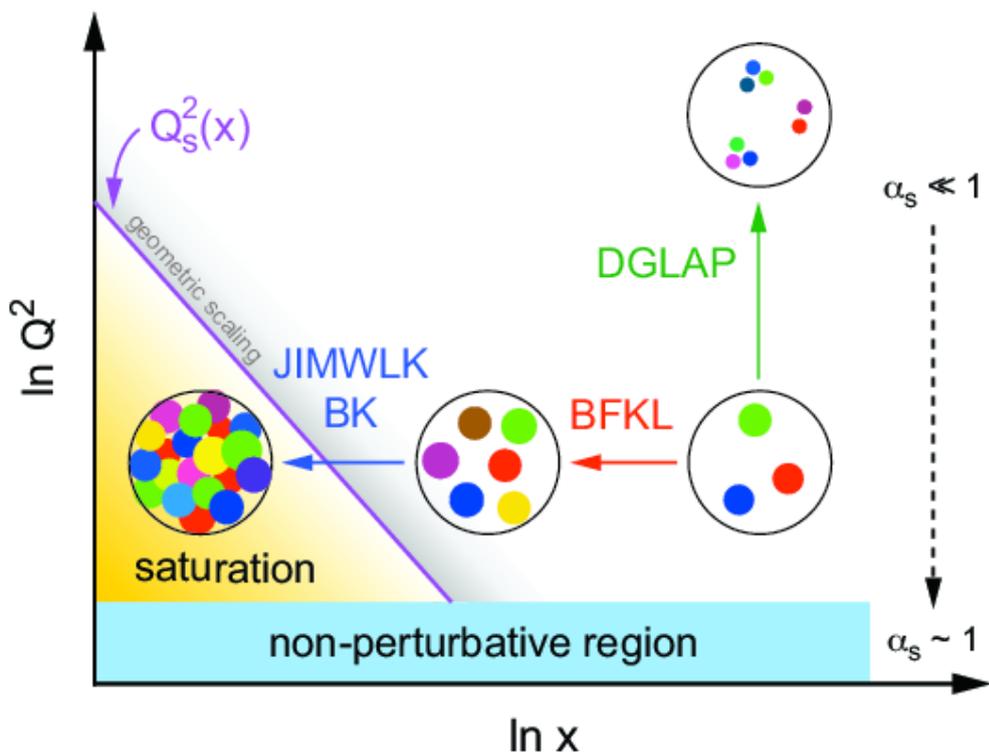
$$Q_s^2(x) \sim \left(\frac{1}{x} \right)^\lambda$$



Saturation at EIC



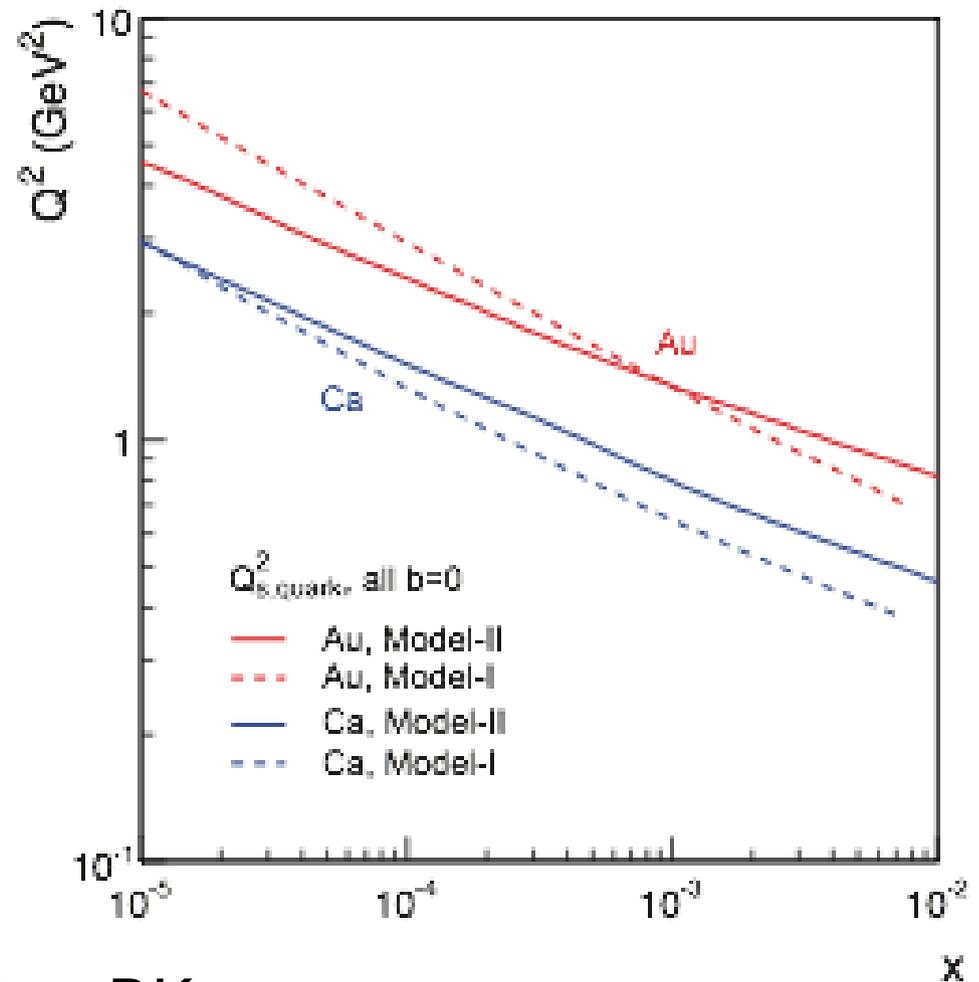
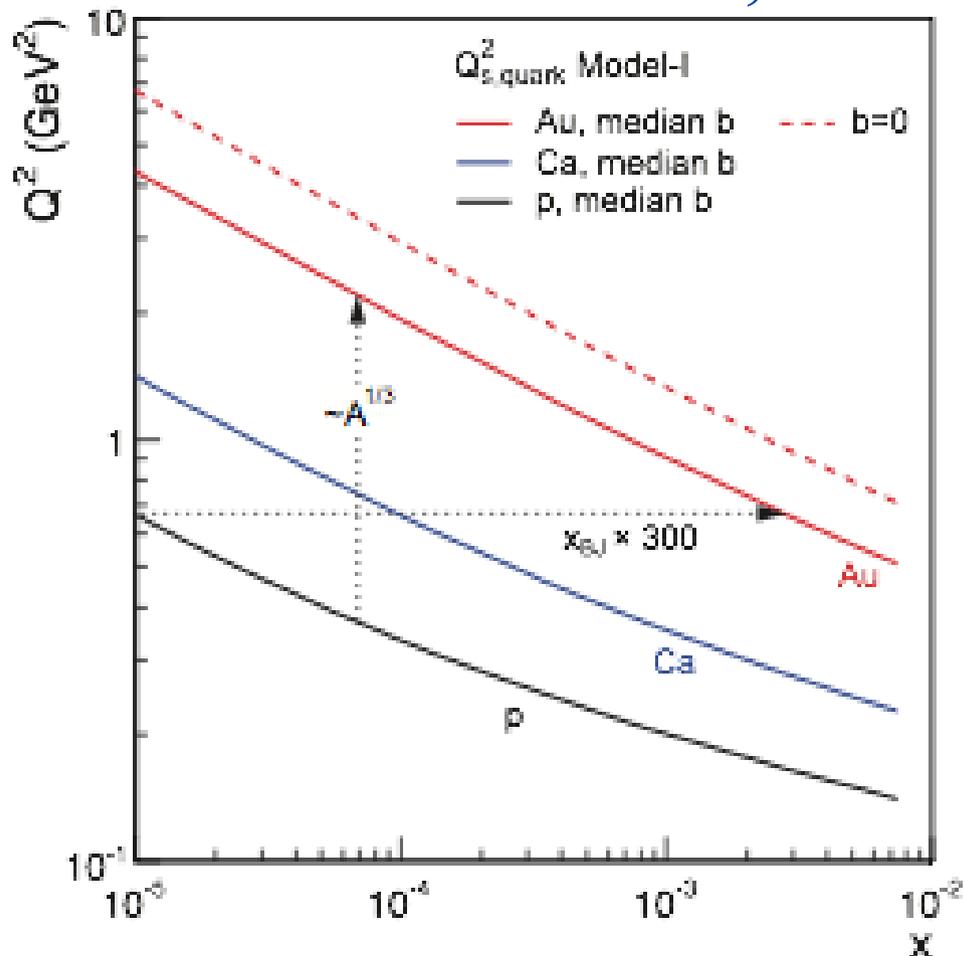
$$Q_s^2(x) \sim A^{1/3} \left(\frac{1}{x} \right)^\lambda$$



Saturation at EIC

Pocket formula: $Q_s^2(x) \sim A^{1/3} \left(\frac{1}{x}\right)^\lambda \sim \left(\frac{A}{x}\right)^{1/3}$

Gold: $A=197$, x 197 times smaller!

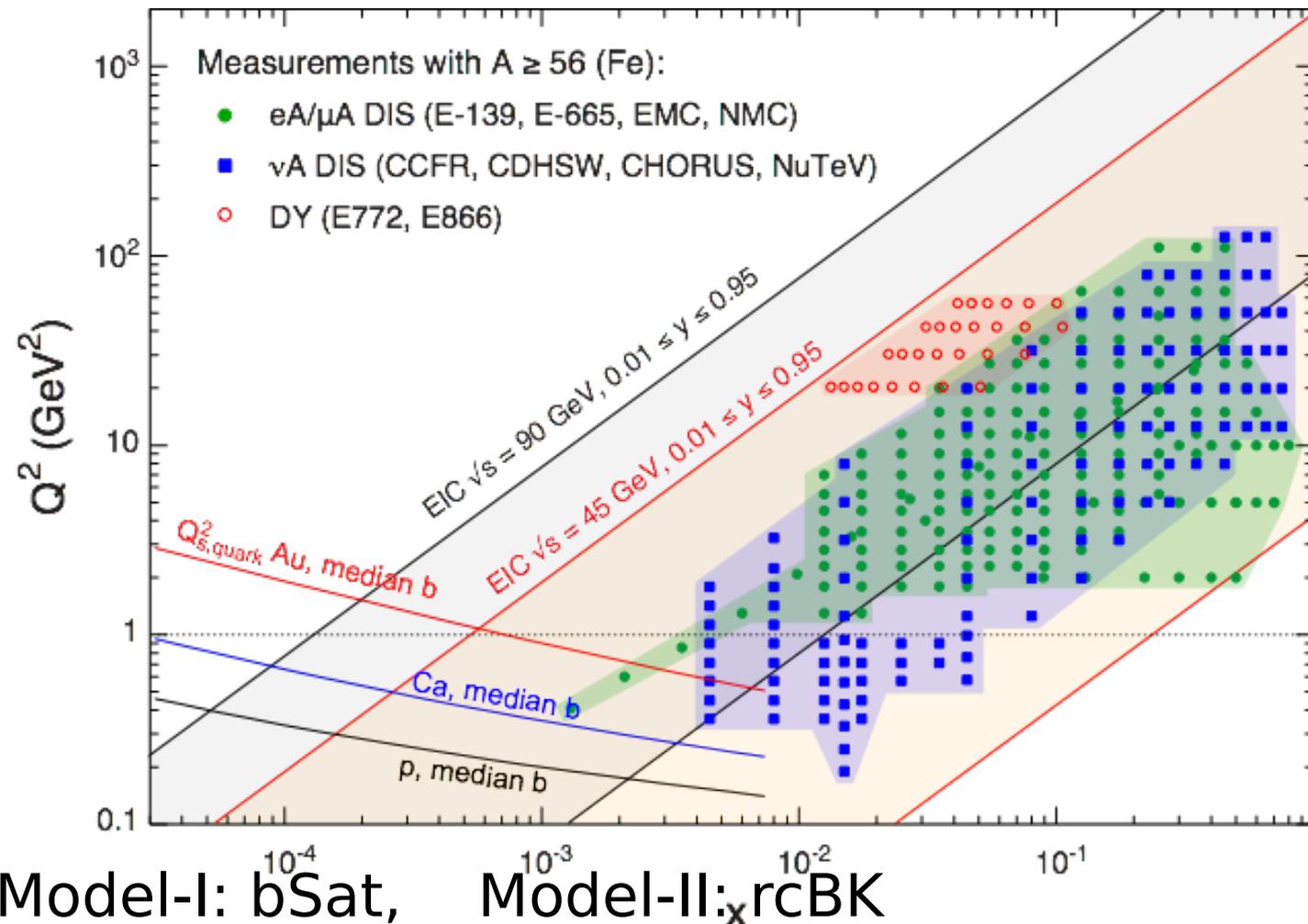


Model-I: bSat, Model-II: rcBK

Saturation at EIC

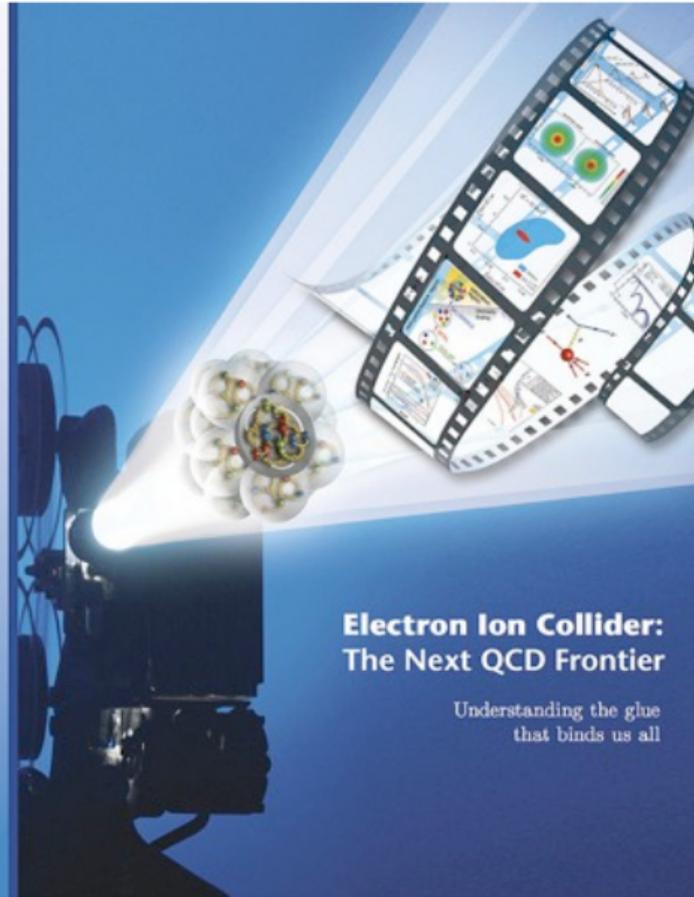
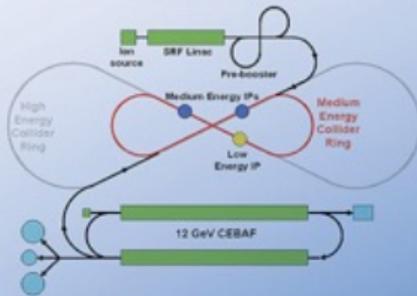
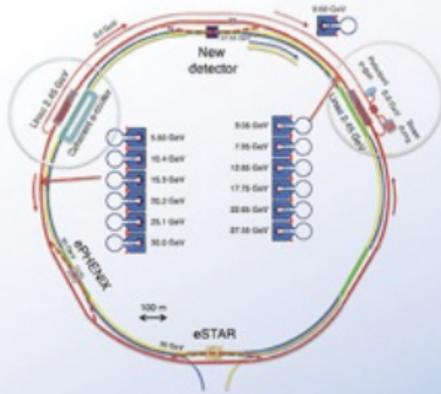
Pocket formula: $Q_s^2(x) \sim A^{1/3} \left(\frac{1}{x}\right)^\lambda \sim \left(\frac{A}{x}\right)^{1/3}$

Gold: $A=197$, x 197 times smaller!



The EIC Physics Programme

Eur.Phys.J. A52 (2016) no.9, 268



Physics of Strong
Colour Fields

Spin Physics

3D Imaging

Hadronisation

Electroweak

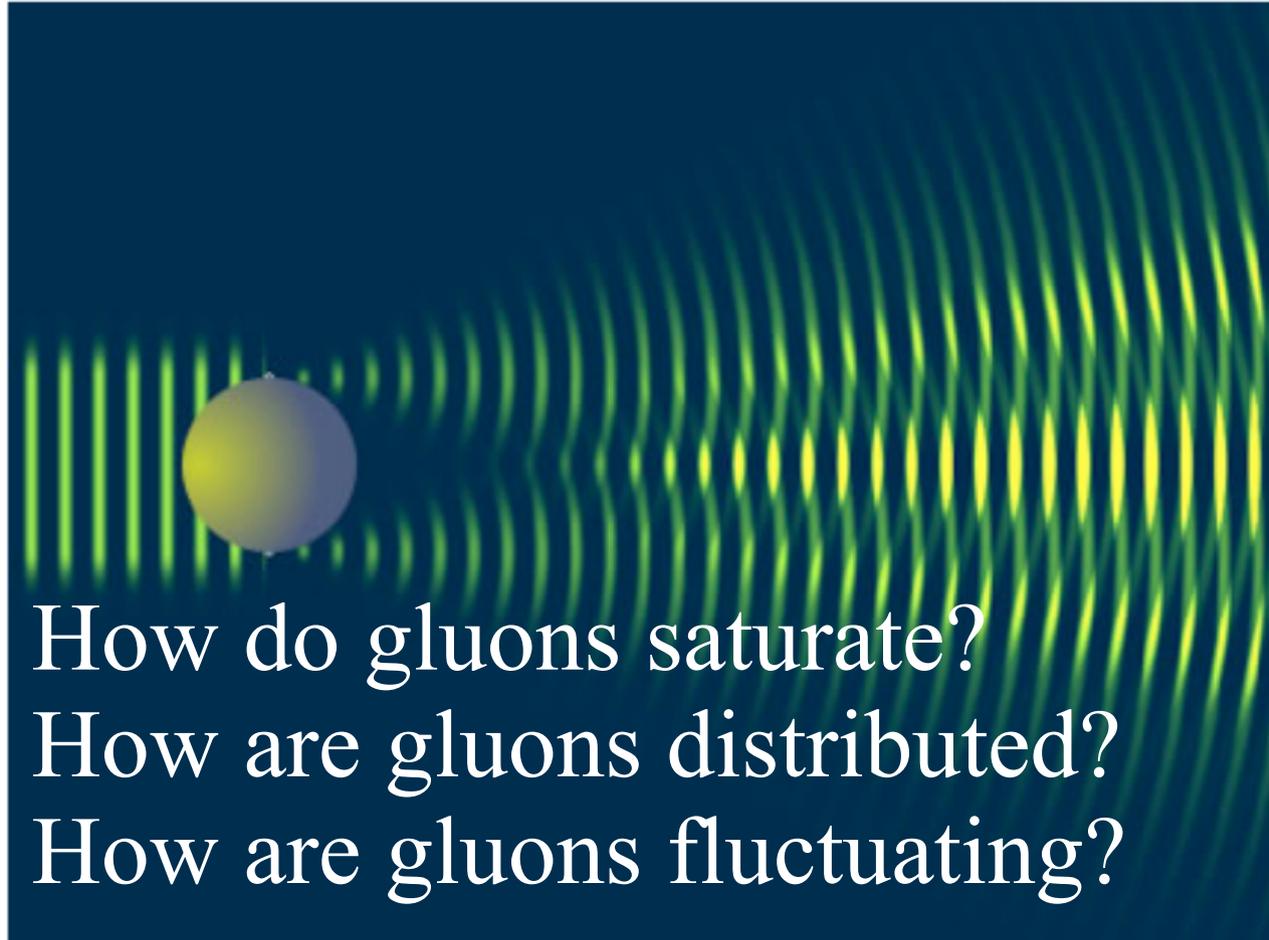
[arXiv:1212.1701](https://arxiv.org/abs/1212.1701)

Jefferson Lab
Thomas Jefferson National Accelerator Facility

BROOKHAVEN
NATIONAL LABORATORY

U.S. DEPARTMENT OF
ENERGY
Office of
Science

Key Measurements: Diffractive Processes in eA

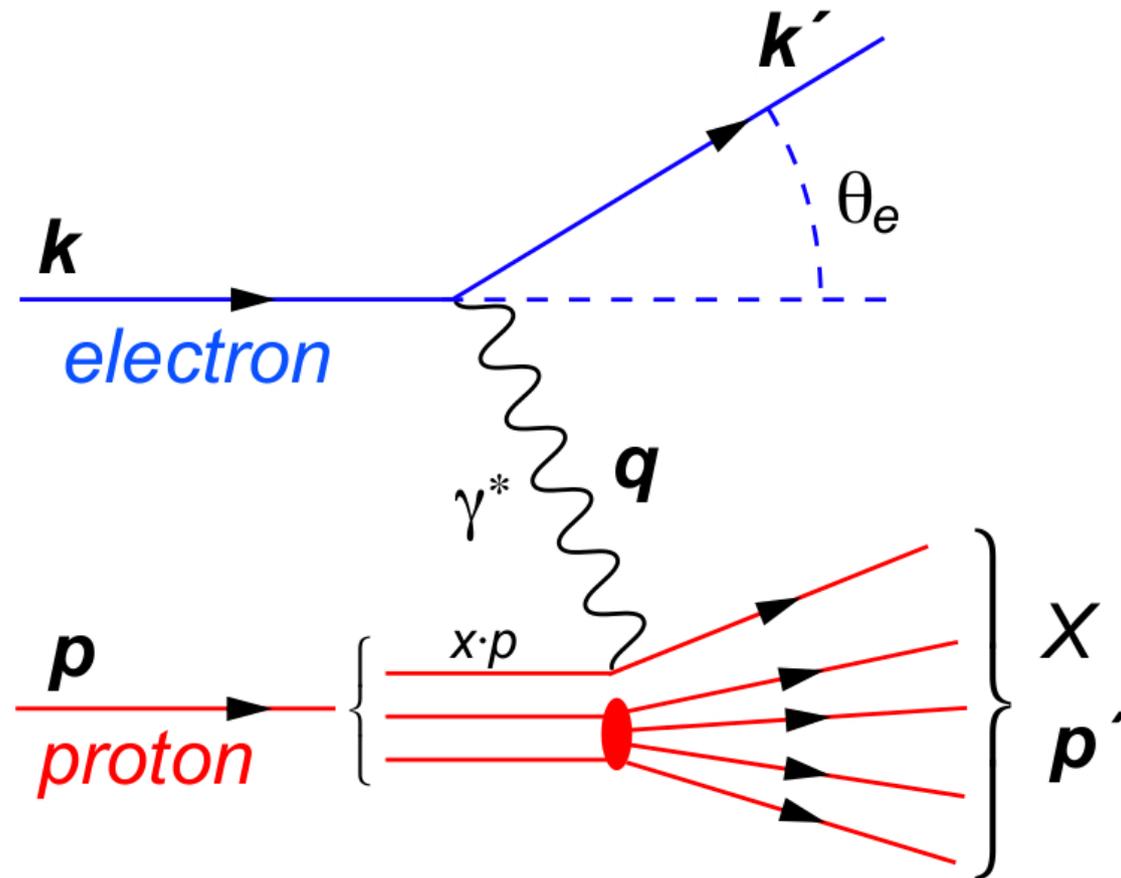


How do gluons saturate?
How are gluons distributed?
How are gluons fluctuating?

In **Nuclei** and **Protons**

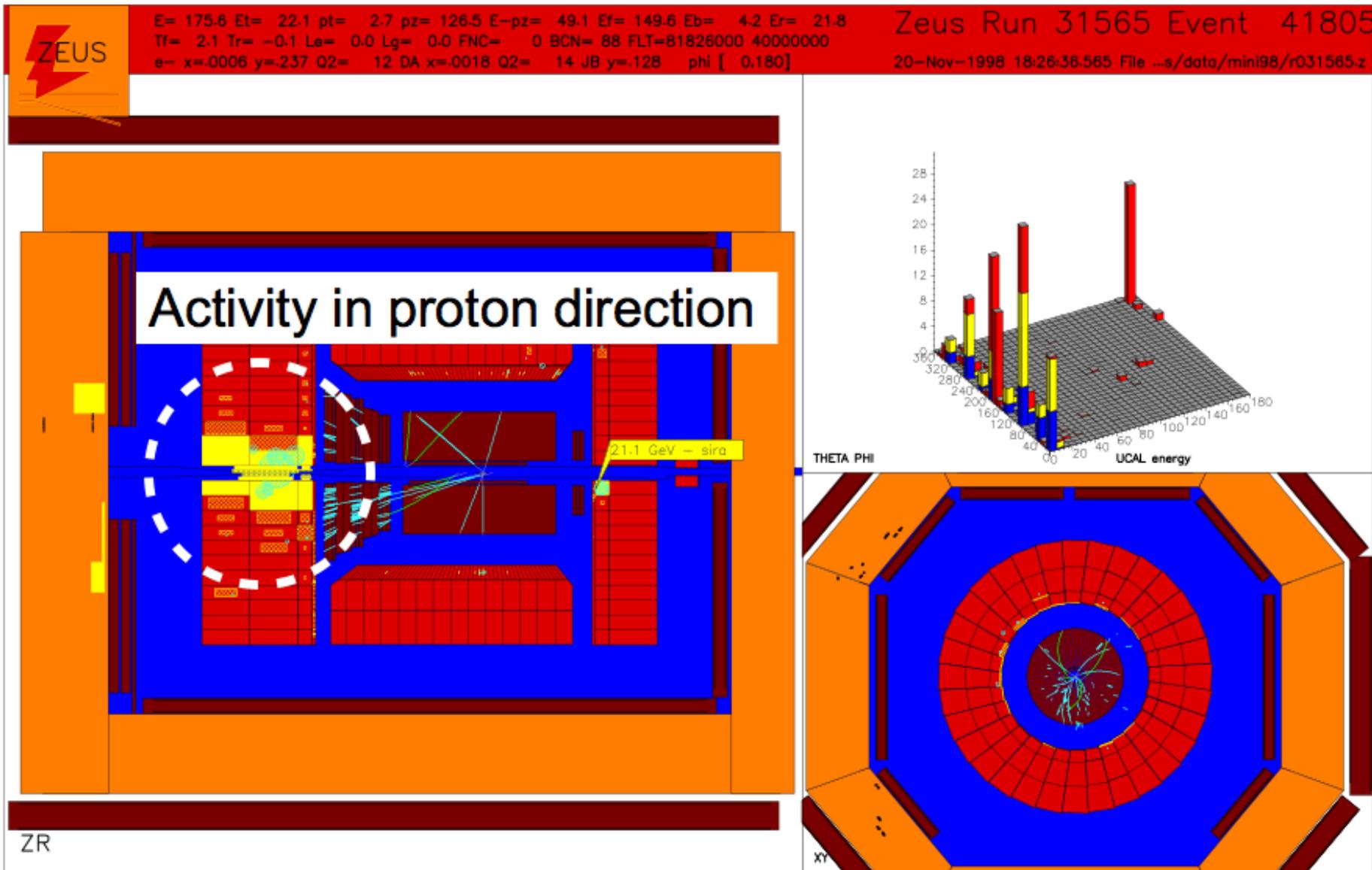
Deep Inelastic Scattering

ep and eA

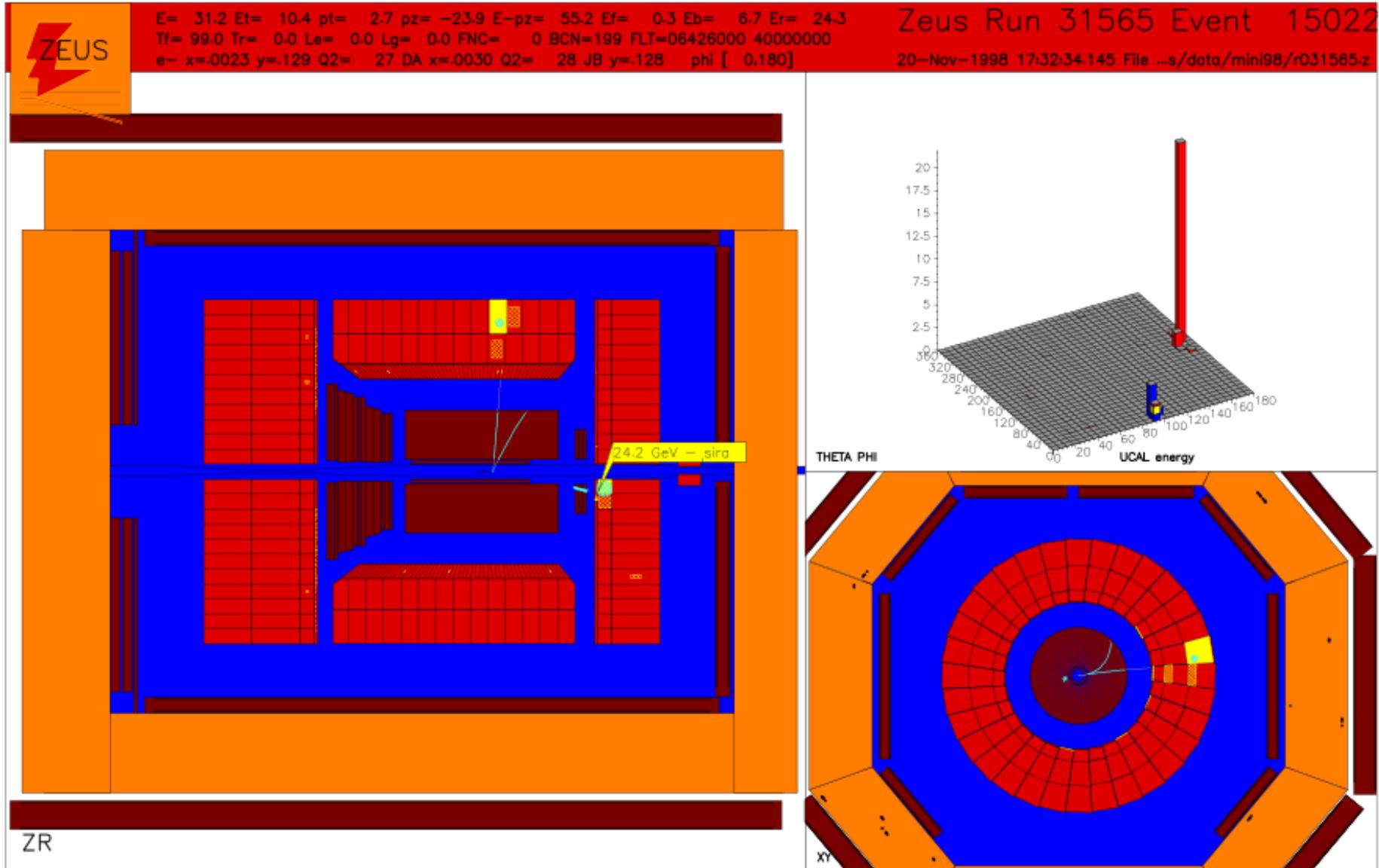


Theoretical view

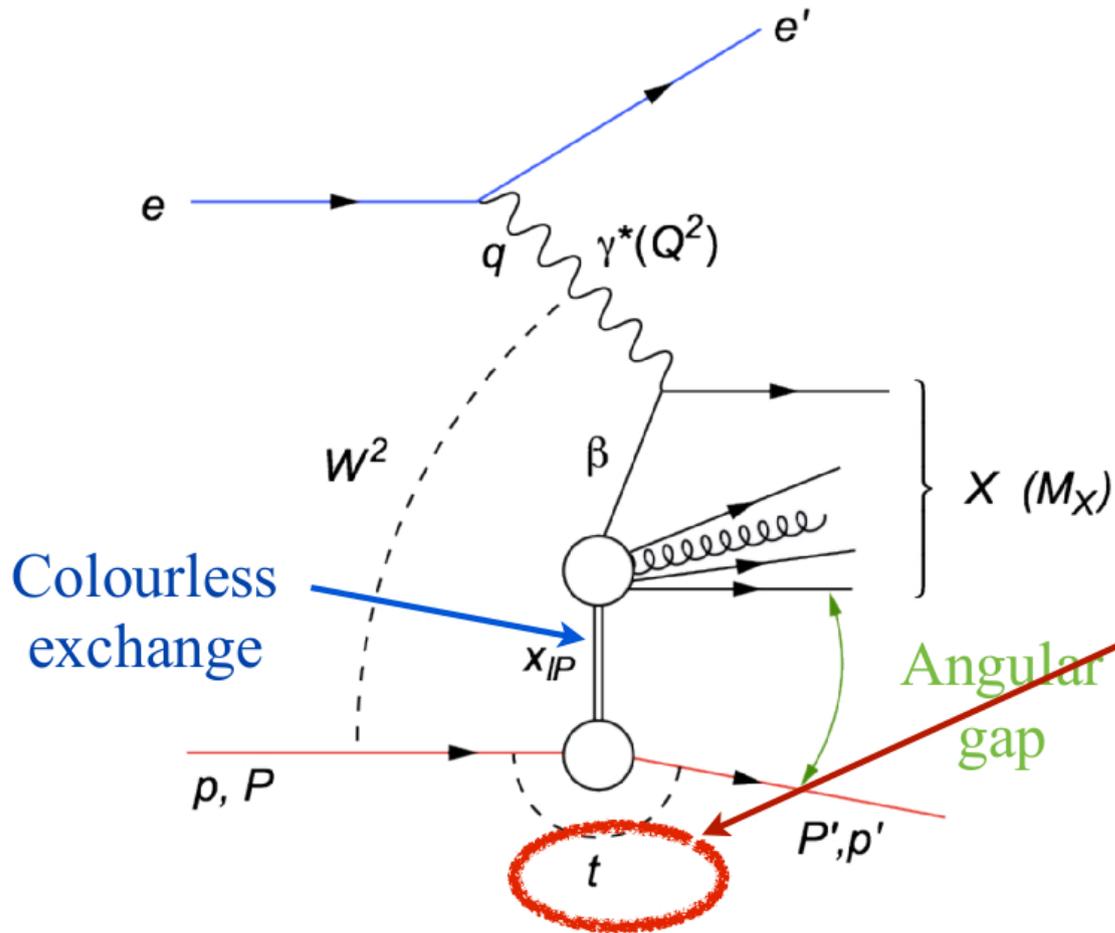
DIS ep and eA



diffraction ep and eA



Diffraction ep and eA



Another kinematic handle

$$t = (p - p')^2$$

HERA:

Proton collides with electron at
CMS energy $\sim 300 m_p$.

In $\sim 15\%$ of measured collisions
proton stays intact!

EIC $e+A$:

Ion predicted to stay intact in
 $25\%-40\%$ of events w.
saturation!

how to measure $t = (P_A - P_{A'})^2$

need to measure $P_{A'}$

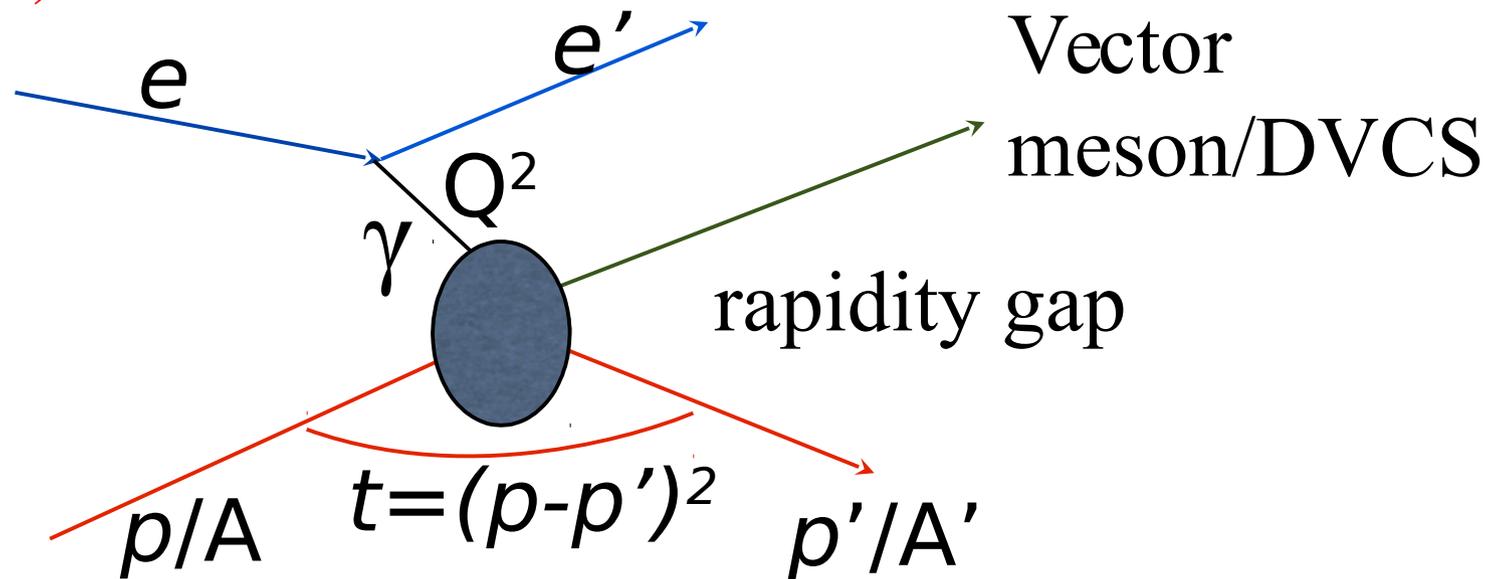
coherent case: A' disappears down beam pipe

incoherent case: cannot measure all beam remnants

only possibility: Exclusive diffraction

$e + A \rightarrow e' + VM + A'$

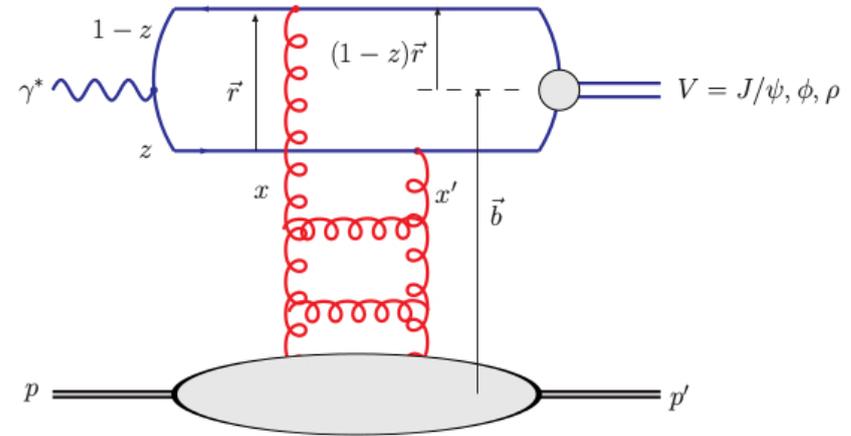
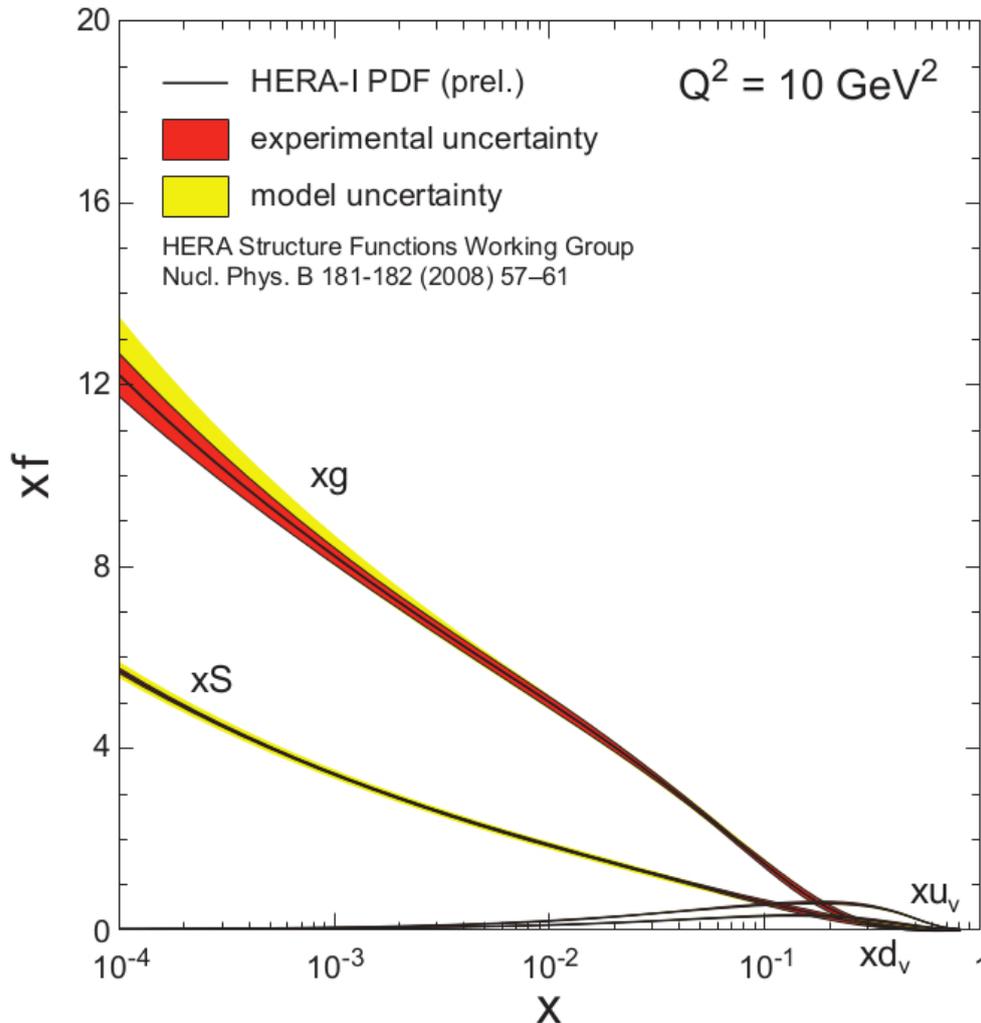
$t = (P_{VM} + P_{e'} - P_e)^2$



Why diffraction is so great

Diffraction sensitive to gluon **momentum** distributions²:

$$\sigma \propto g(x, Q^2)^2$$



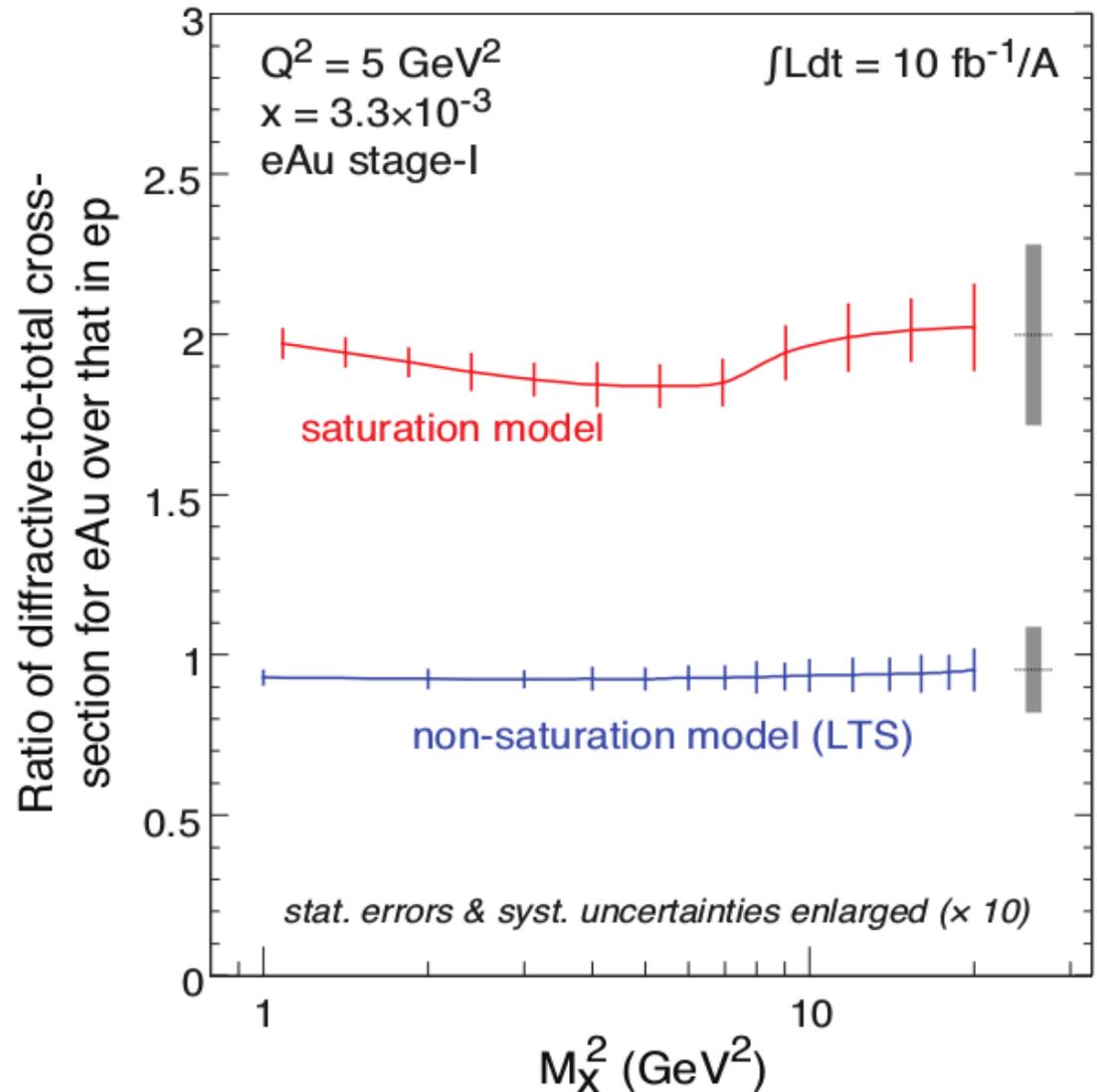
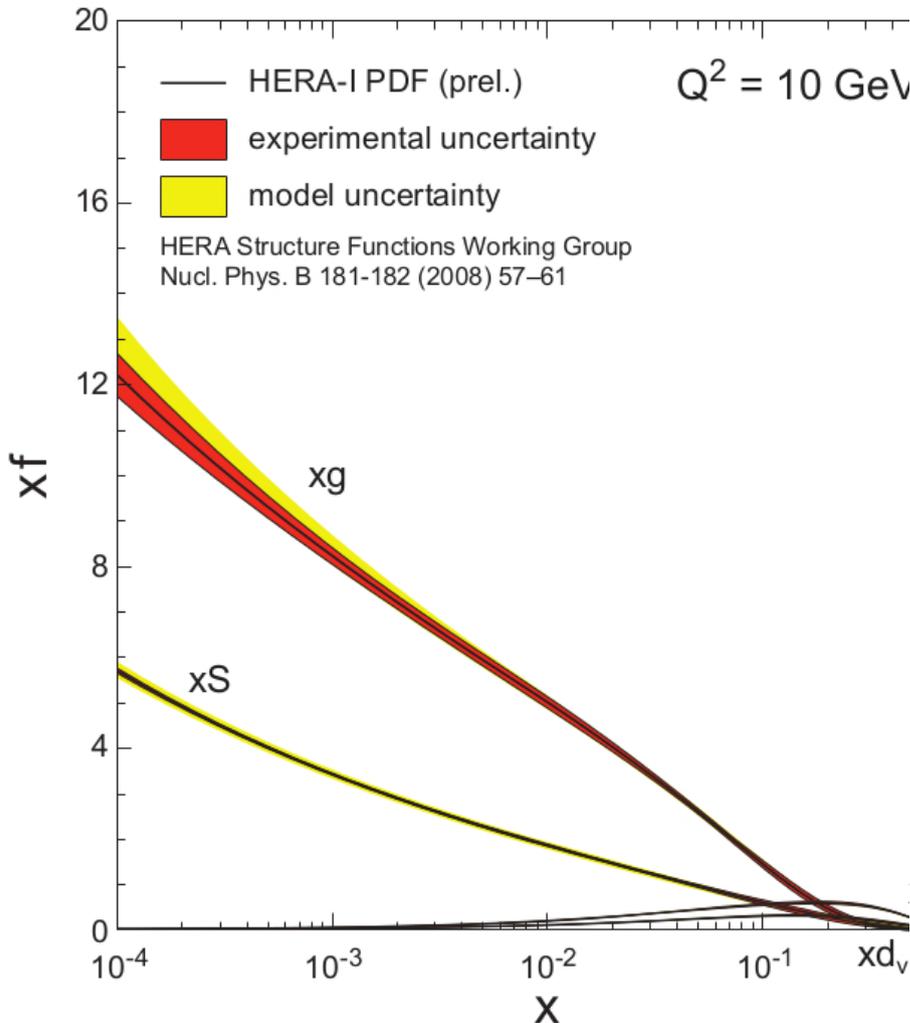
The colourless exchange can be understood as **two** gluons screening each other's colour

Sensitive to how the gluon distribution saturates at small x

Why diffraction is so great

Diffraction sensitive to gluon **momentum** distributions²:

$$\sigma \propto g(x, Q^2)^2$$



Why diffraction is so great

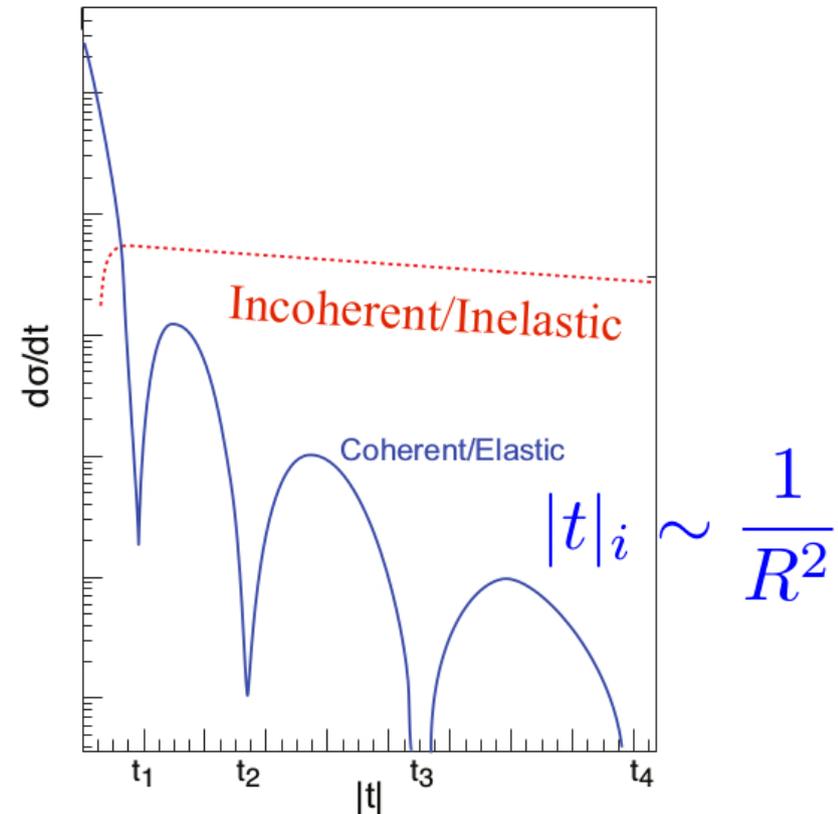
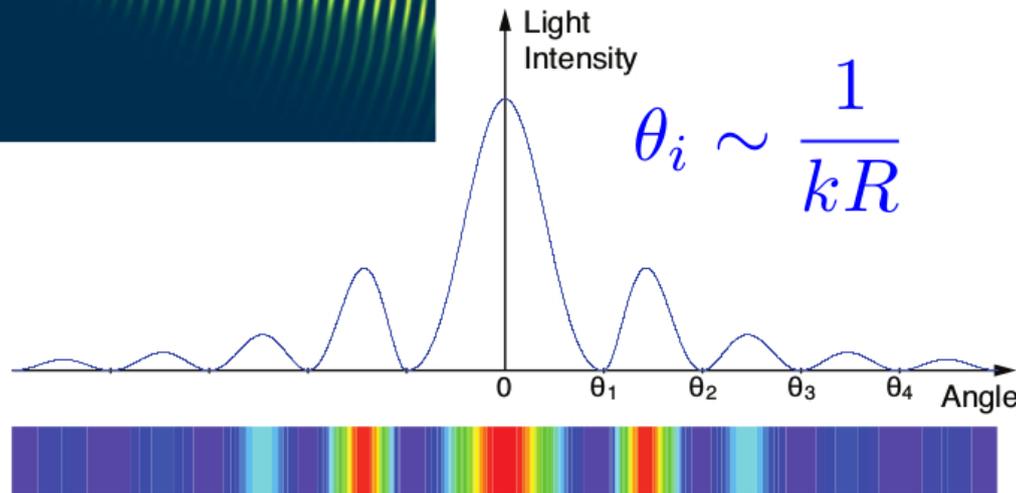
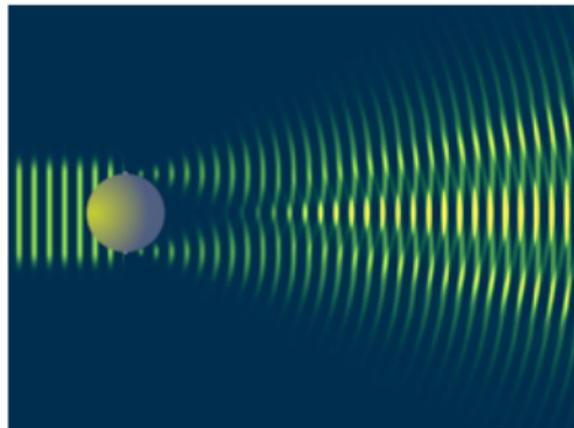
Sensitive to **spatial** gluon distributions

a projectile scattering off a nucleus of radius R

-not a 'black disk', edge effects
-inelastic scattering

Light scattering elastically off a circular screen of radius R

t : Fourier conjugate to ion shape



Incoherent Scattering

Good, Walker:

nucleus dissociates ($f \neq i$):

$$\sigma_{\text{incoherent}} \propto \sum_{f \neq i} \langle i | \mathcal{A} | f \rangle^\dagger \langle f | \mathcal{A} | i \rangle \quad \text{complete set}$$

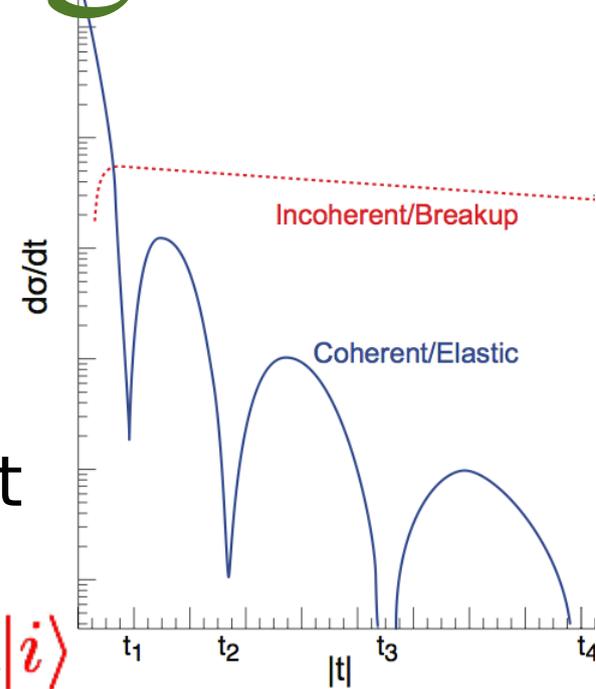
$$= \sum_f \langle i | \mathcal{A} | f \rangle^\dagger \langle f | \mathcal{A} | i \rangle - \langle i | \mathcal{A} | i \rangle^\dagger \langle i | \mathcal{A} | i \rangle$$

$$= \langle i | |\mathcal{A}|^2 | i \rangle - |\langle i | \mathcal{A} | i \rangle|^2 = \langle |\mathcal{A}|^2 \rangle - |\langle \mathcal{A} \rangle|^2$$

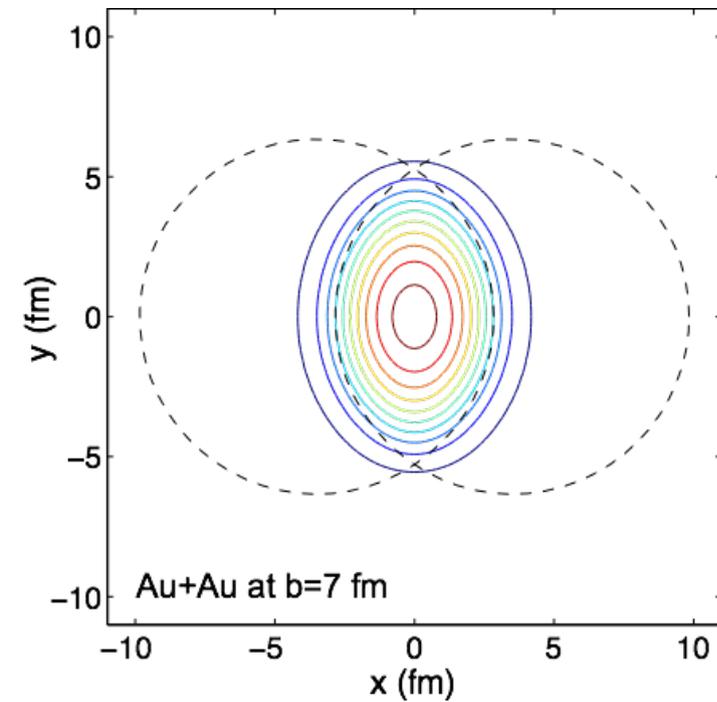
the incoherent CS is the variance of the amplitude!!

$$\frac{d\sigma_{\text{total}}}{dt} = \frac{1}{16\pi} \langle |\mathcal{A}|^2 \rangle$$

$$\frac{d\sigma_{\text{coherent}}}{dt} = \frac{1}{16\pi} |\langle \mathcal{A} \rangle|^2$$



Eccentricity and the spatial distribution



$$\frac{dN}{d\varphi} \propto 1 + 2v_2 \cos[2(\varphi - \psi_R)] + \dots$$

$$v_2 = \langle \cos[2(\varphi - \psi_R)] \rangle$$

Sensitive to **early interactions** and **pressure gradients**

Ideal hydrodynamics, $v_2 \propto$ spatial eccentricity ϵ : $\epsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$

v_2/ϵ versus particle density is sensitive test if system approaches ideal hydrodynamic

$$\frac{v_2}{\epsilon} = \frac{h}{1 + B / \left(\frac{1}{S} \frac{dN}{dy} \right)}$$

S transverse area, h hydro limit of v_2/ϵ and

$B \propto \eta/s$ **Shear Viscosity/Entropy Density**

Bhalerao, Blaizot, Borghini and Ollitrault,
Phys. Lett. B 627 (2005) 49

Luzum and Romatschke, Phys. Rev. C 78, 034915

Eccentricity and the spatial distribution

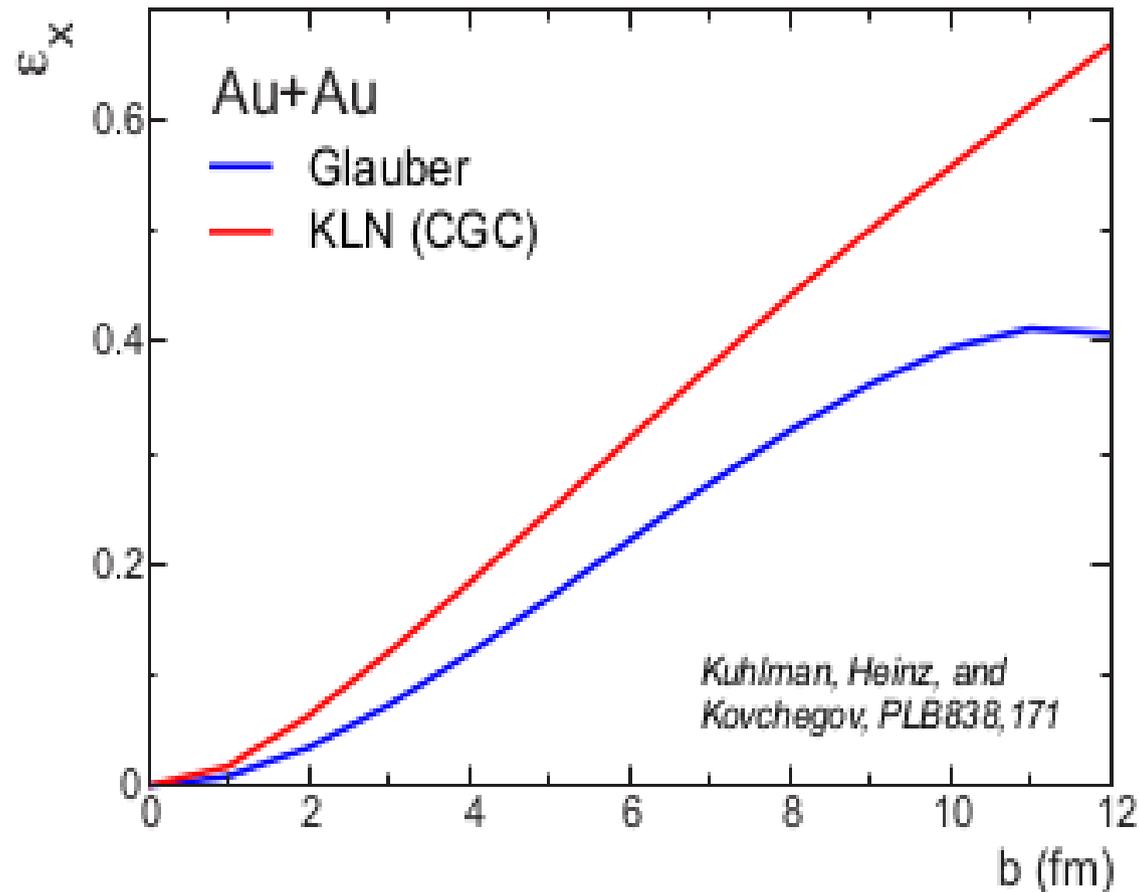
The question is what is ε ?

RHIC & LHC: low- p_T realm driven almost entirely by glue

\Rightarrow spatial distribution of glue in nuclei?

Two methods for ε :

- ▶ Glauber (non-saturated)?
- ▶ CGC (saturated)?



Eccentricity and the spatial distribution

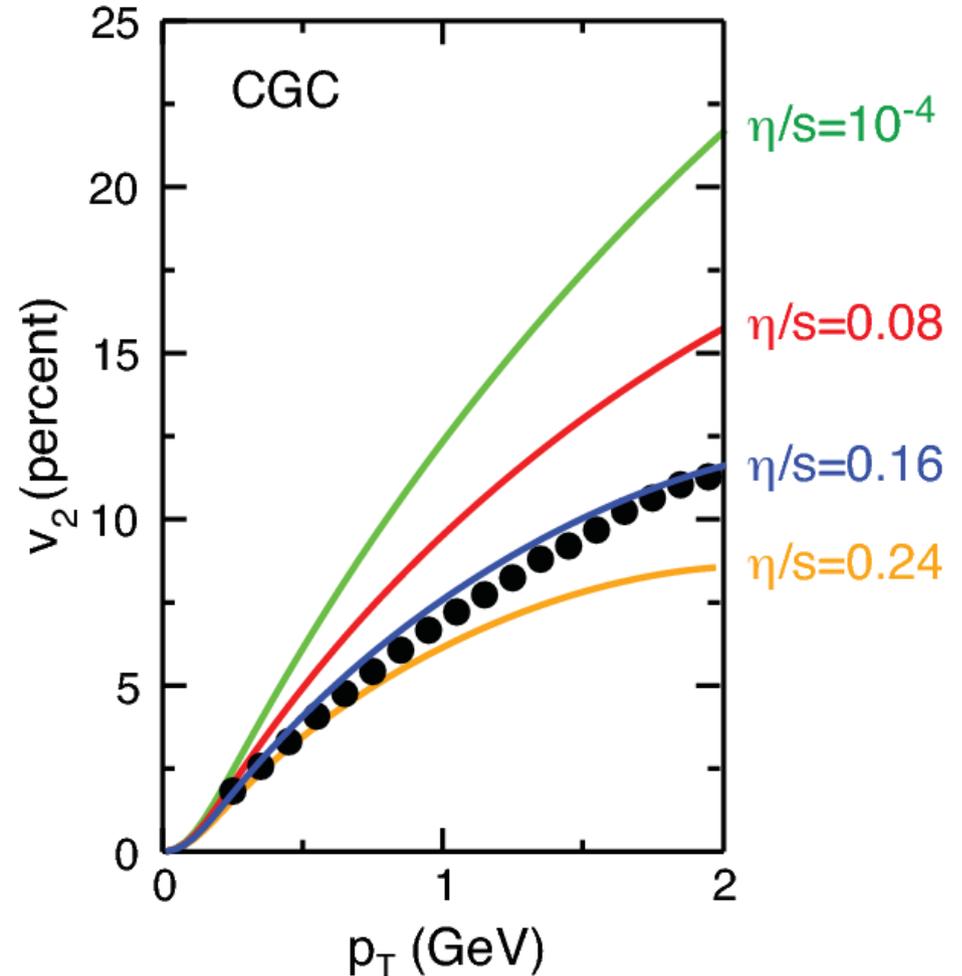
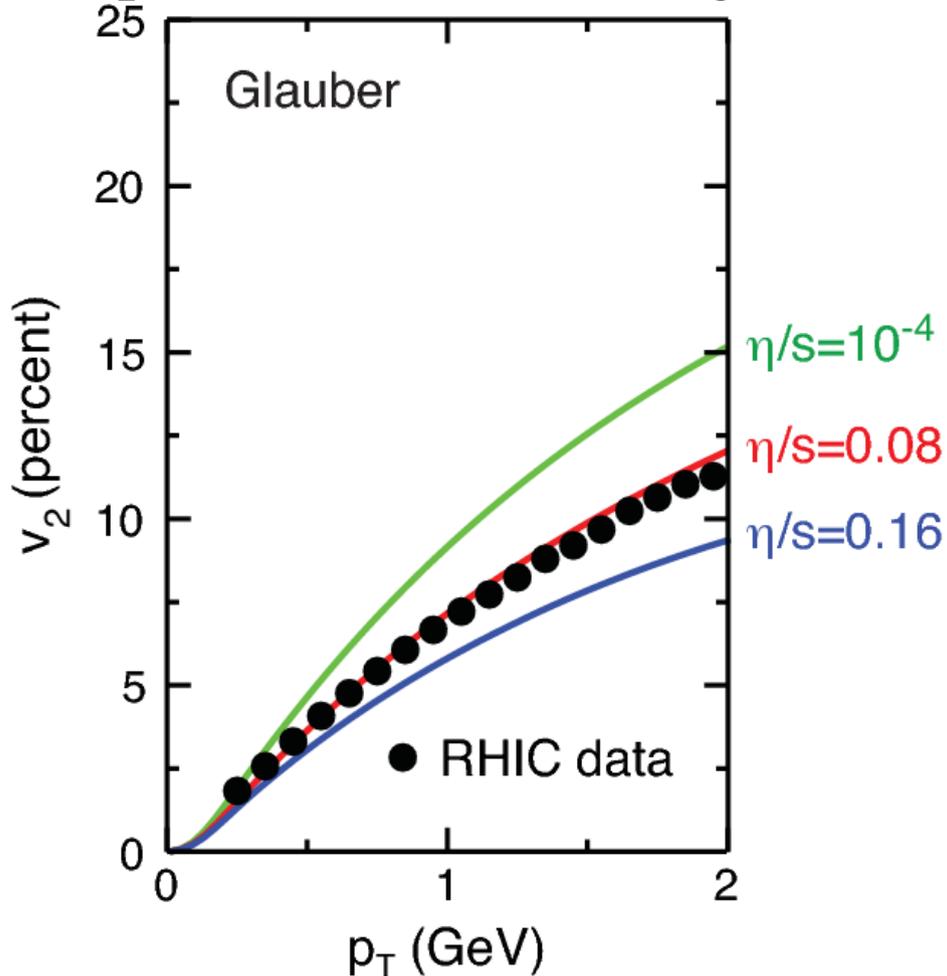
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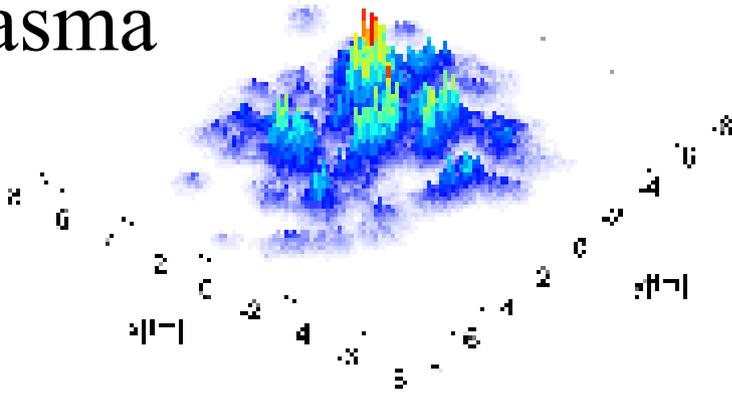
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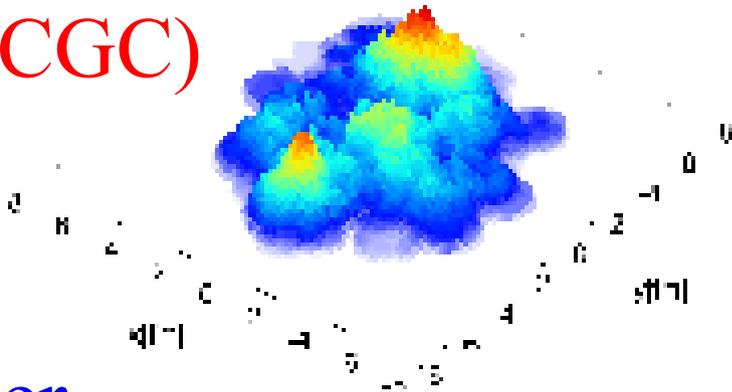


Fluctuations

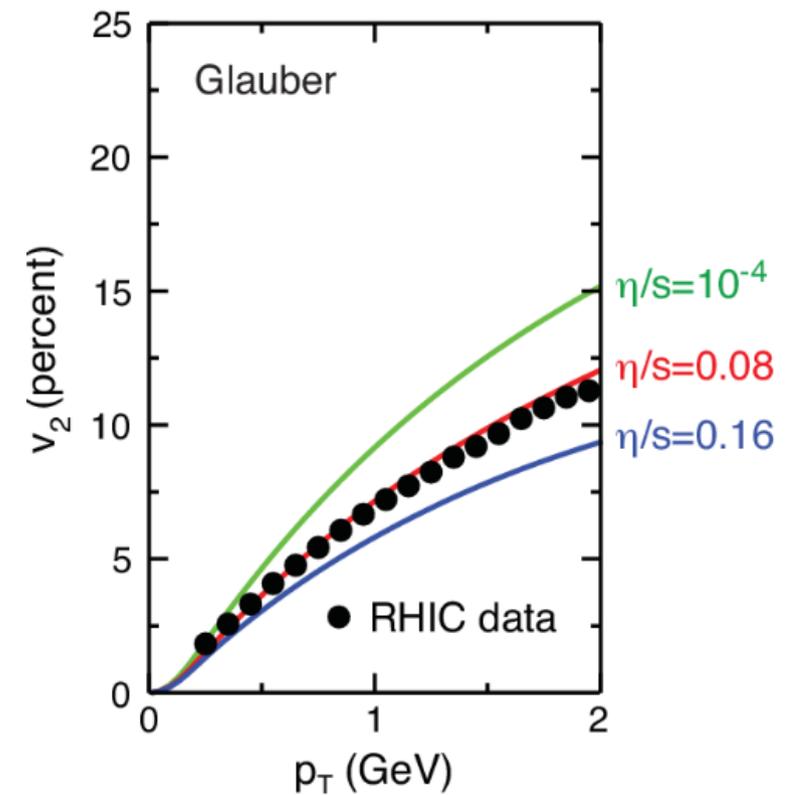
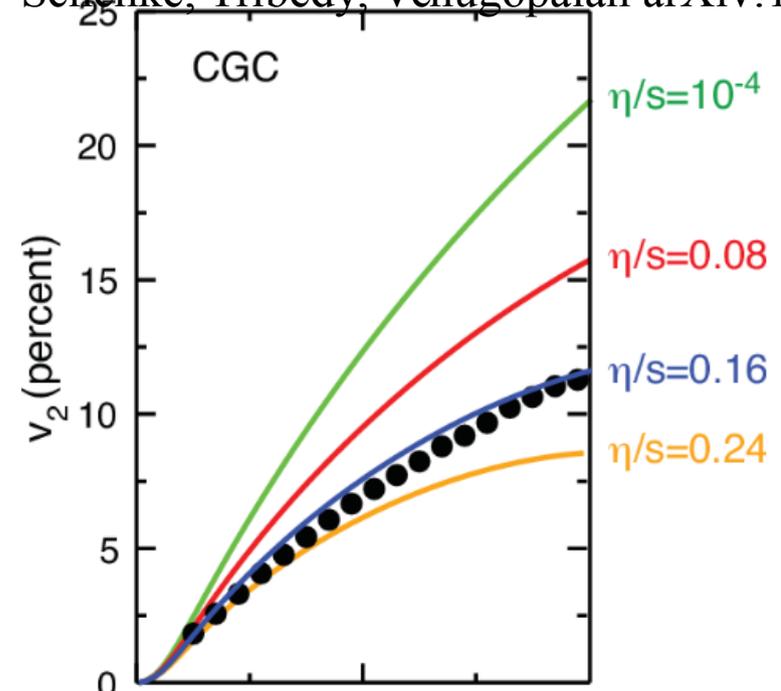
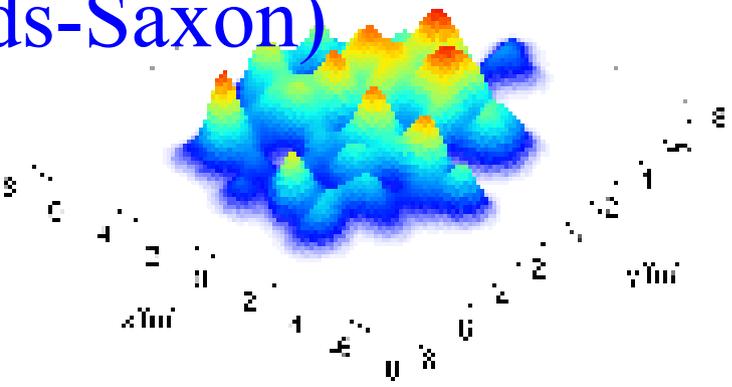
IP-Glasma



KLN(CGC)



Glauber
(Woods-Saxon)

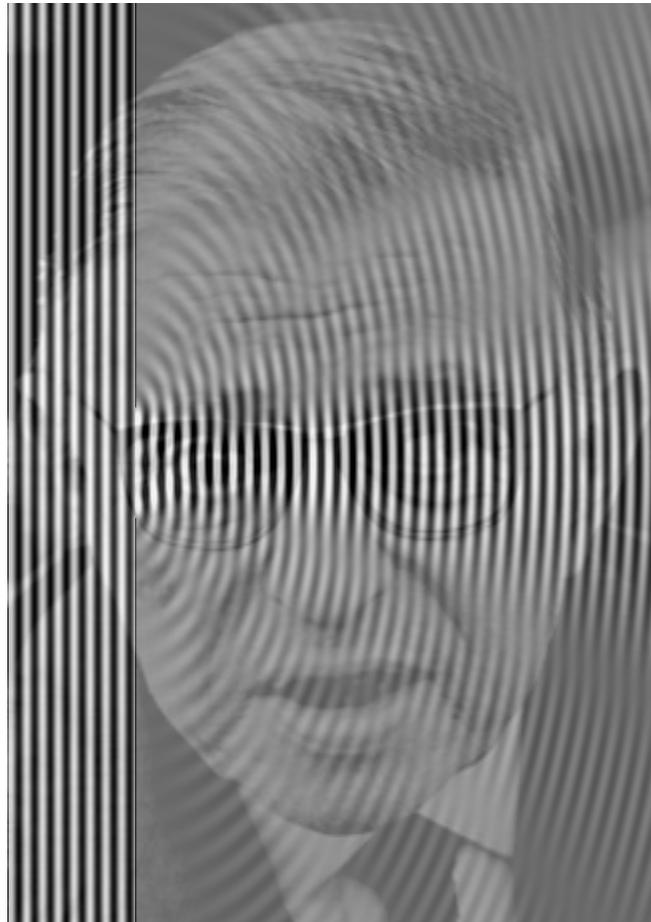


EIC predictions:

Physics Event Generator

Sartre

Exclusive Diffractive
Vector Meson and
DVCS production in
ep and *eA*



TT and T. Ullrich,
Exclusive diffractive processes in electron-ion collisions

Phys. Rev. C 87, 024913
(2013)

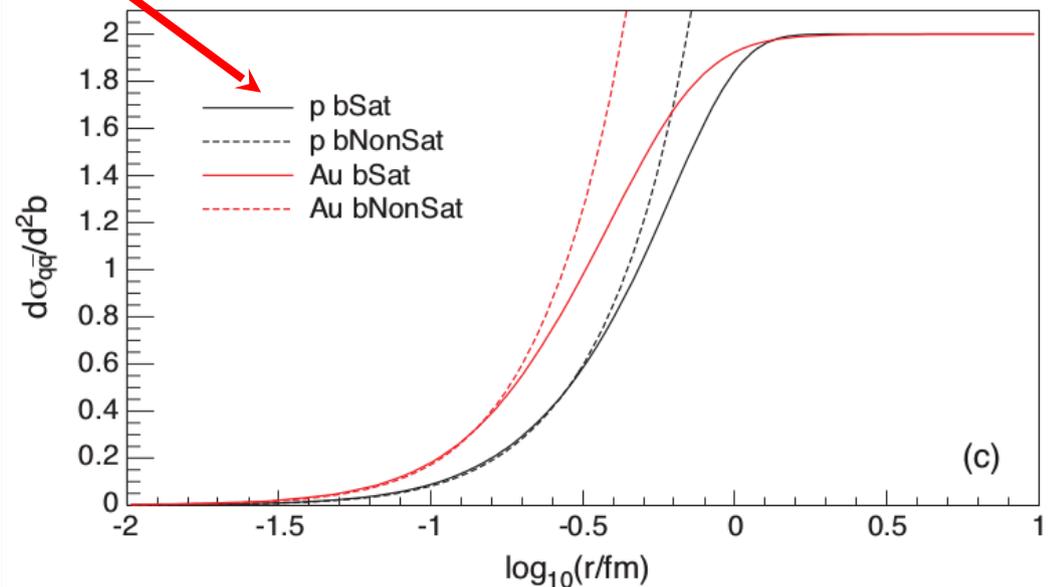
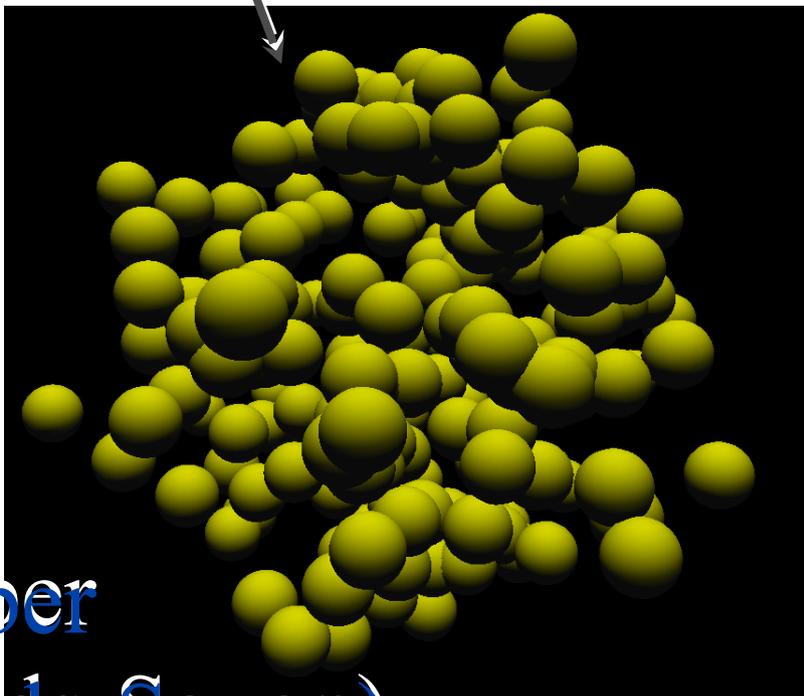
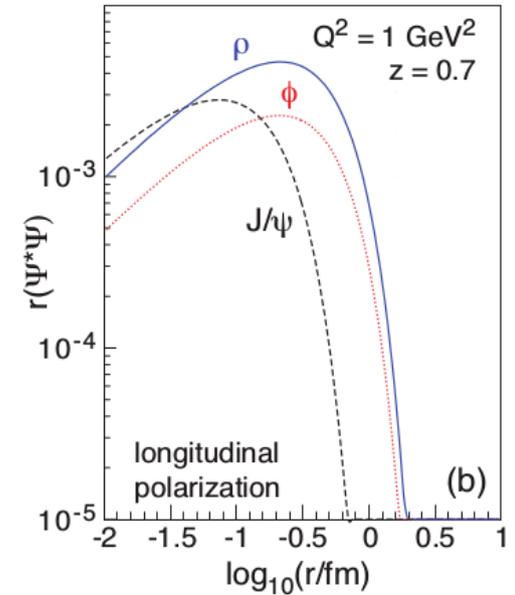
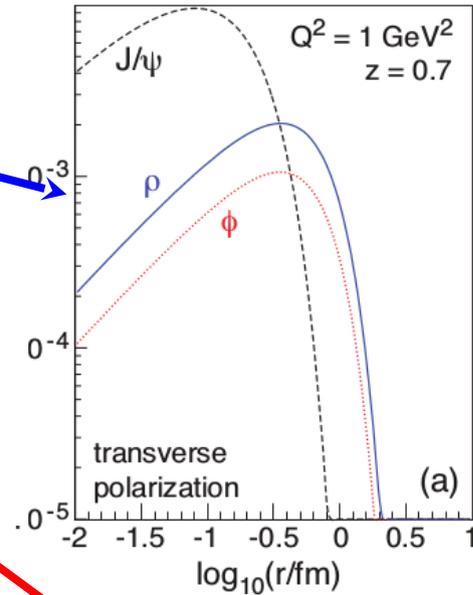
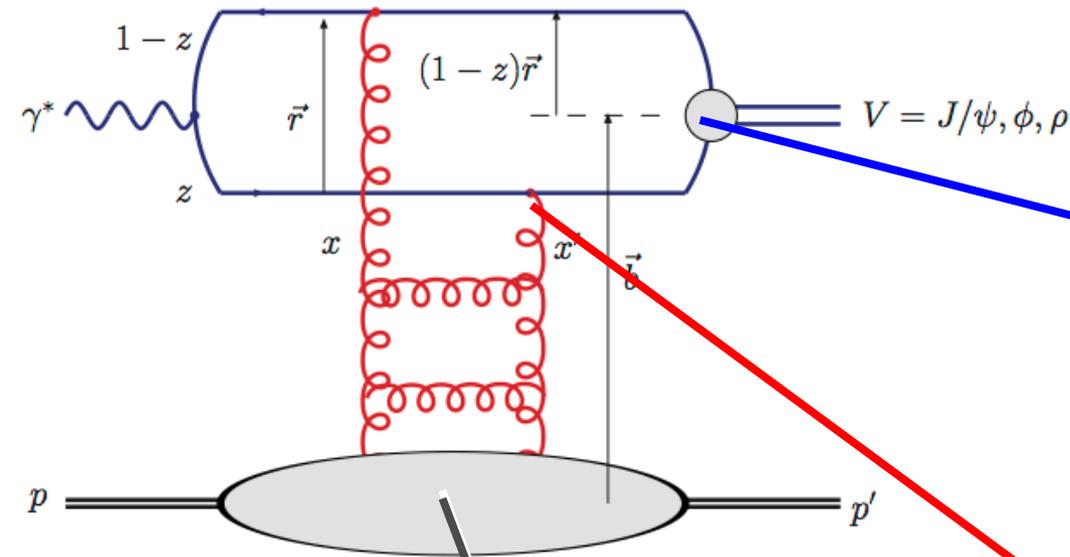
[arXiv:1211.3048](https://arxiv.org/abs/1211.3048)

TT and T. Ullrich,
The dipole model Monte Carlo generator Sartre 1

Comput.Phys.Commun. 185
(2014) 1835-1853

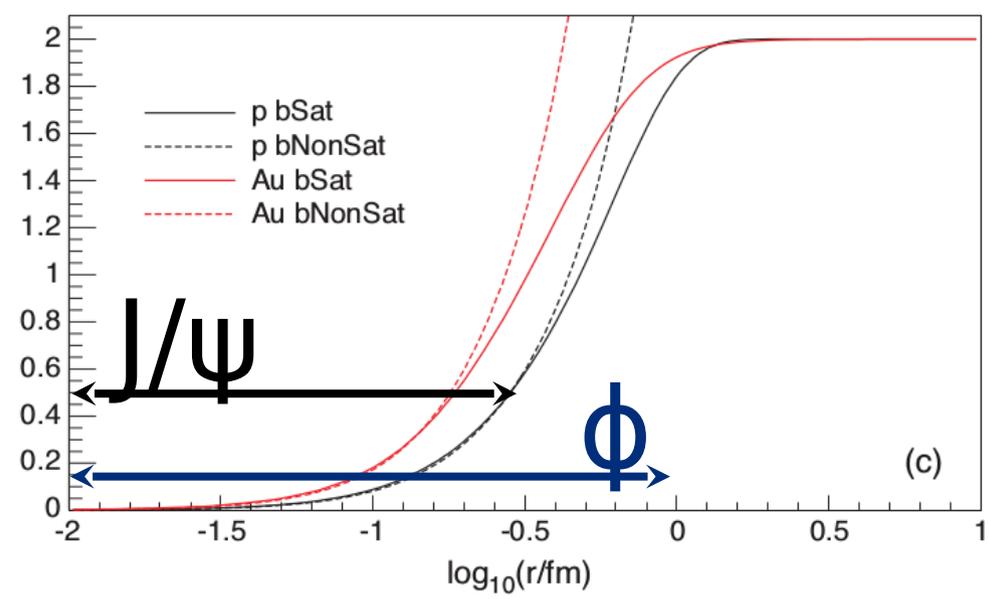
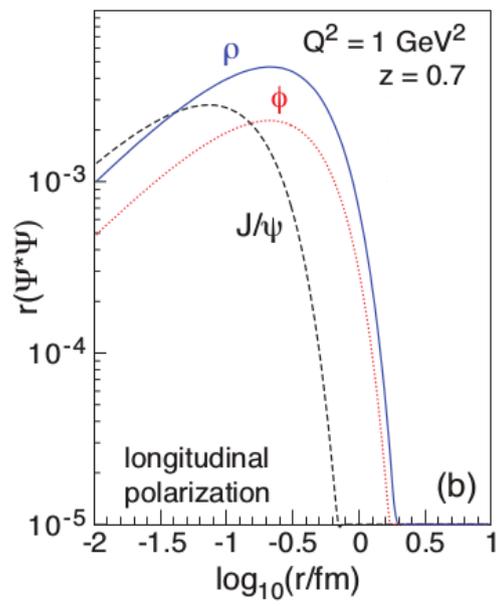
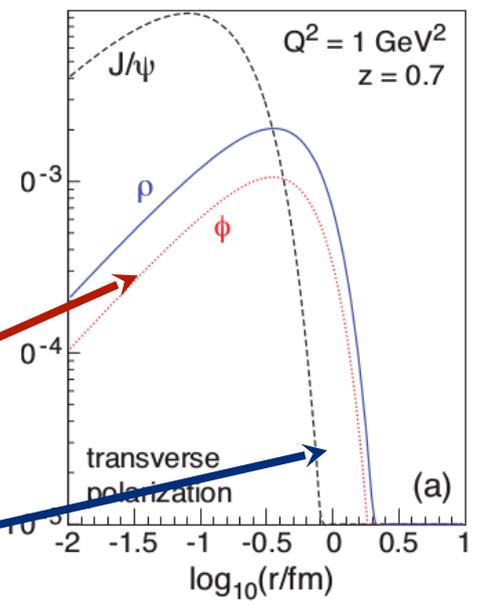
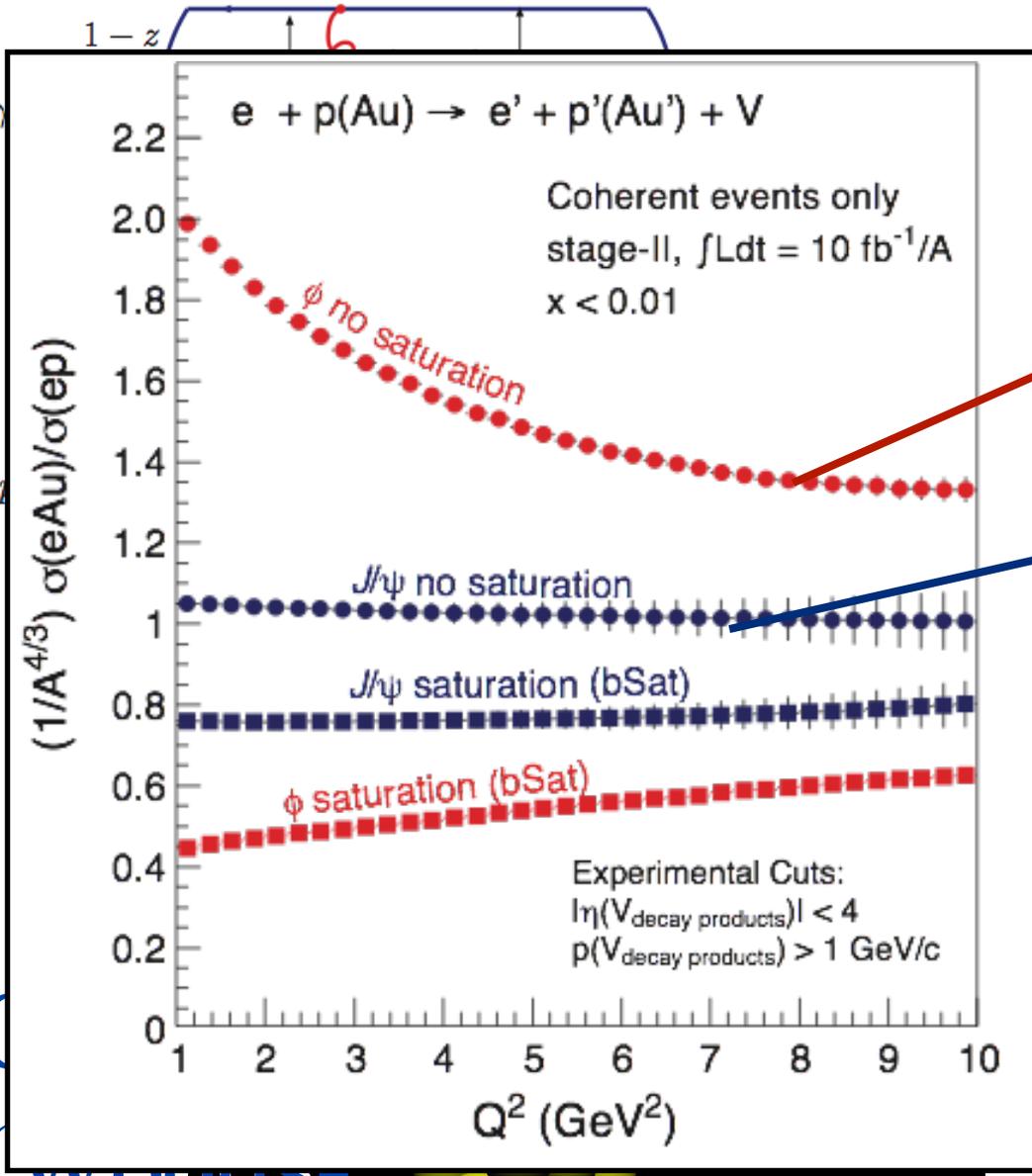
[arXiv:1307.8059](https://arxiv.org/abs/1307.8059)

Sartre dipole model with Glauber bSat and bNonSat



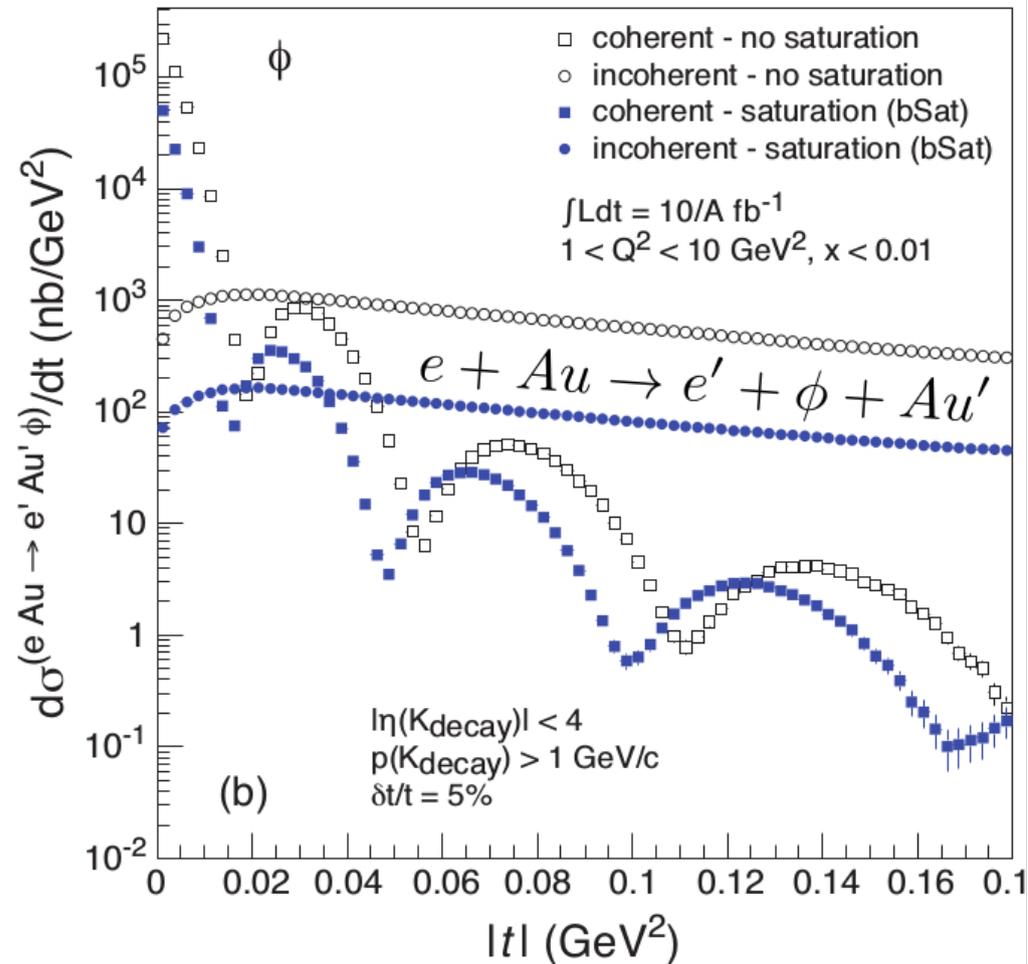
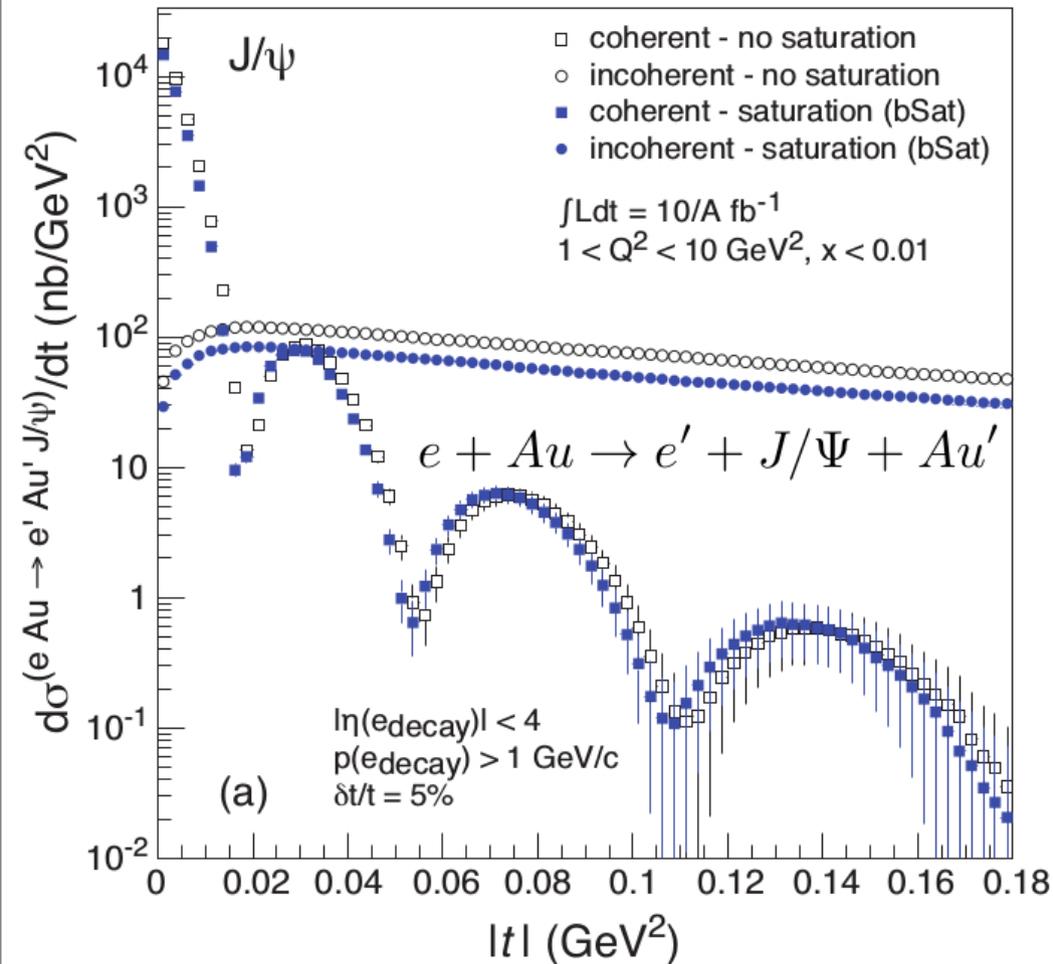
Glauber
(Woods-Saxon)

Sartre dipole model with Glauber bSat and bNonSat



EIC predictions: exclusive diffraction

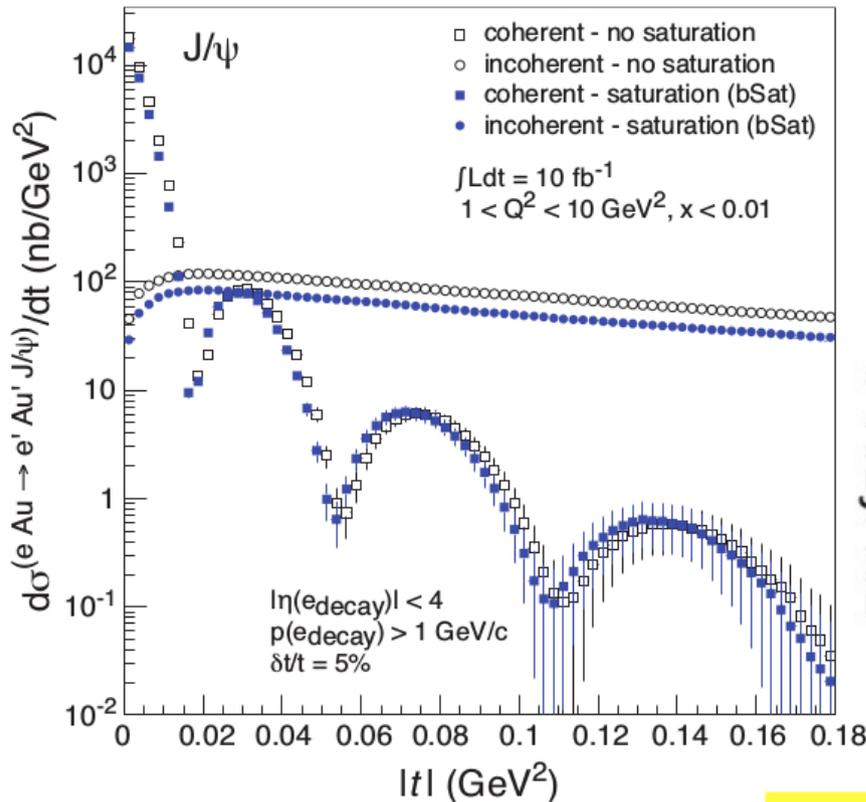
T.T. & Thomas Ullrich
PRC 87 (2013) 024913



Can constrain models **a lot** with a few months of running!
 First 4 dips obtainable.

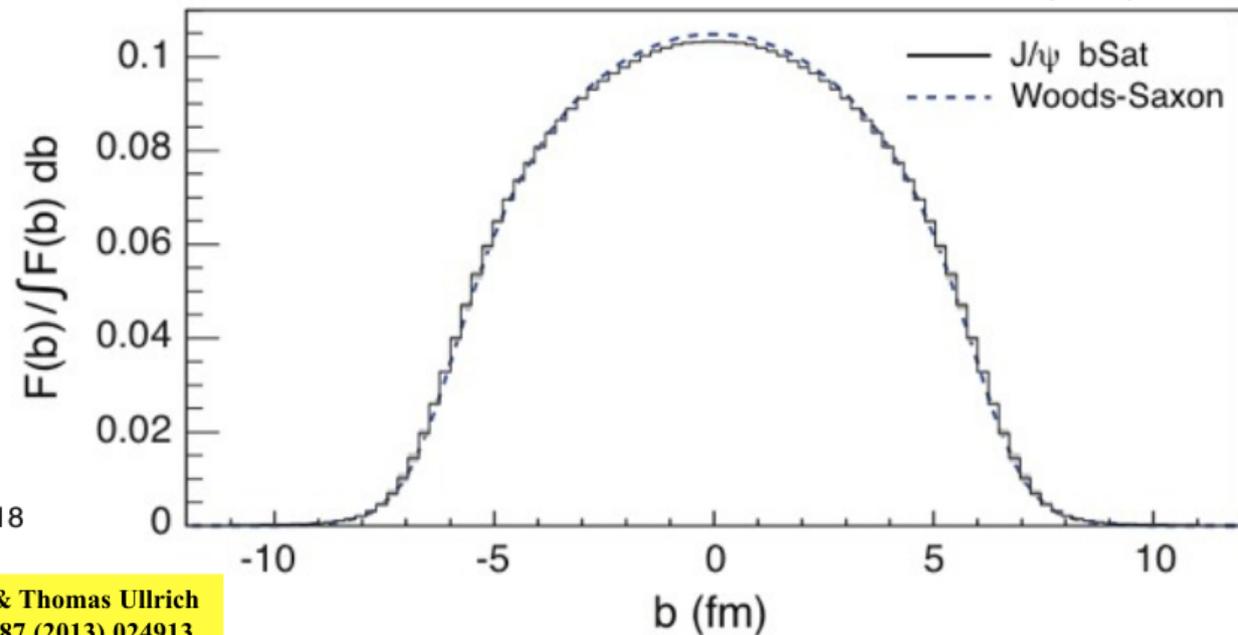
Probing the **spatial** gluon distribution at EIC

Momentum transfer t conjugate to transverse coordinate b



$$F(b) \propto \int d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$

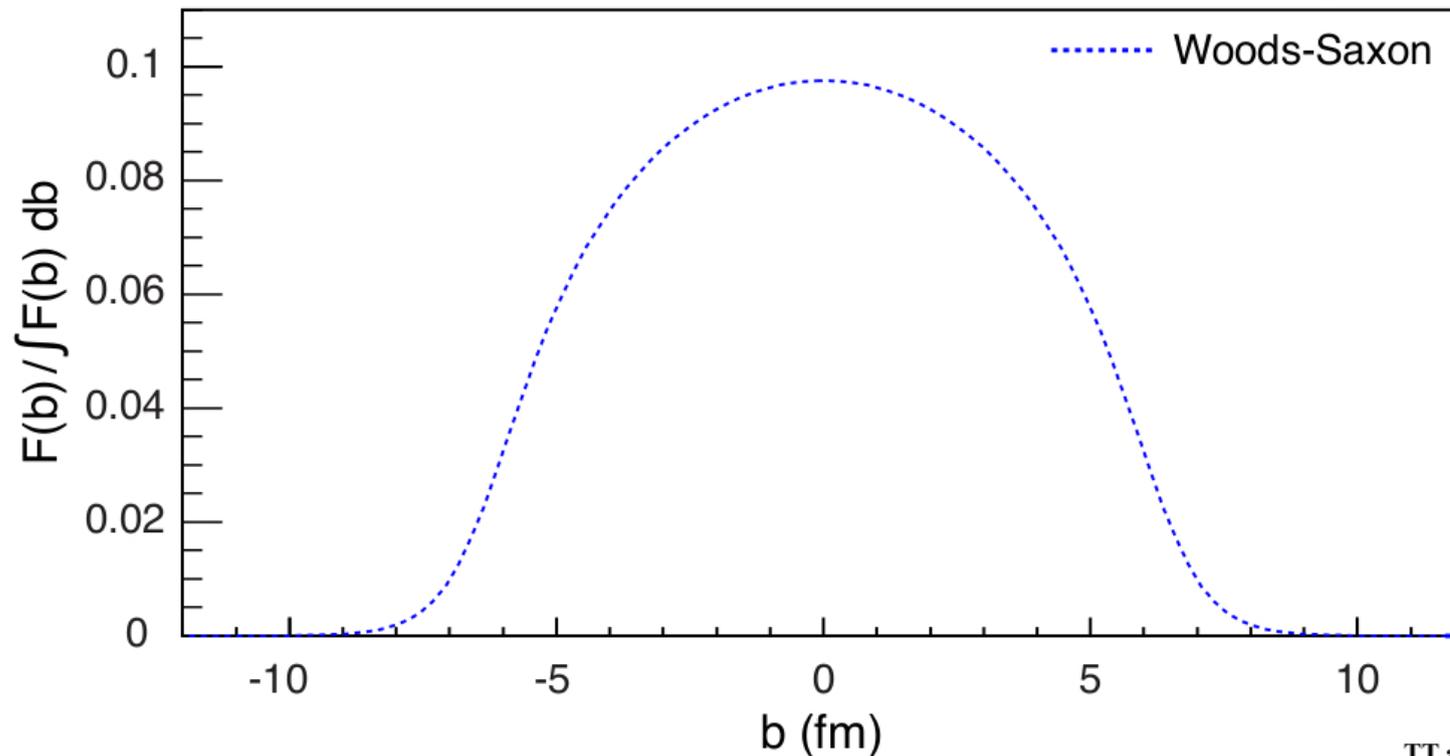
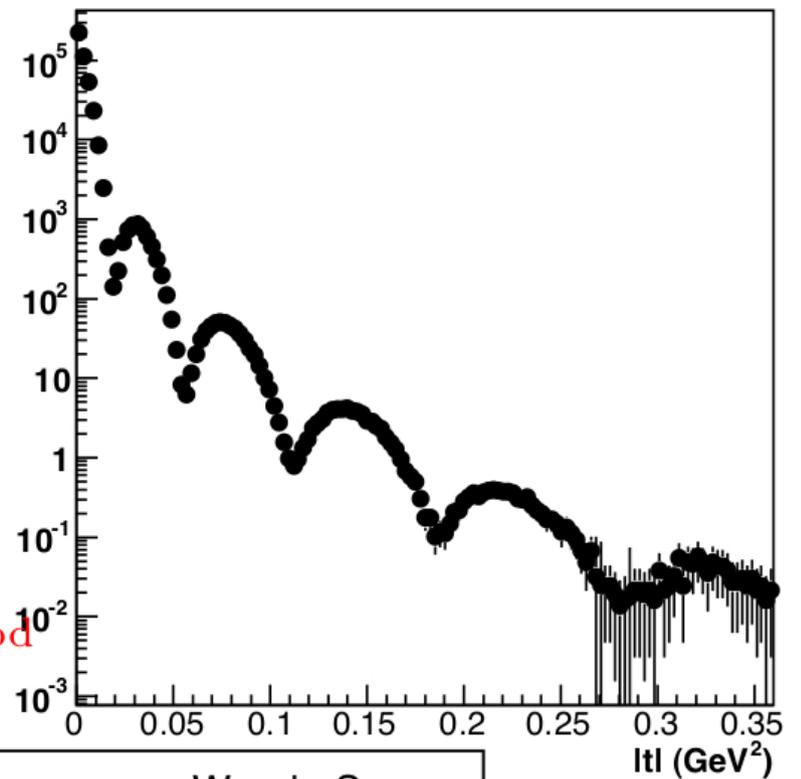


T.T. & Thomas Ullrich
PRC 87 (2013) 024913

EIC will be able to retrieve the spatial gluon distribution with high precision.

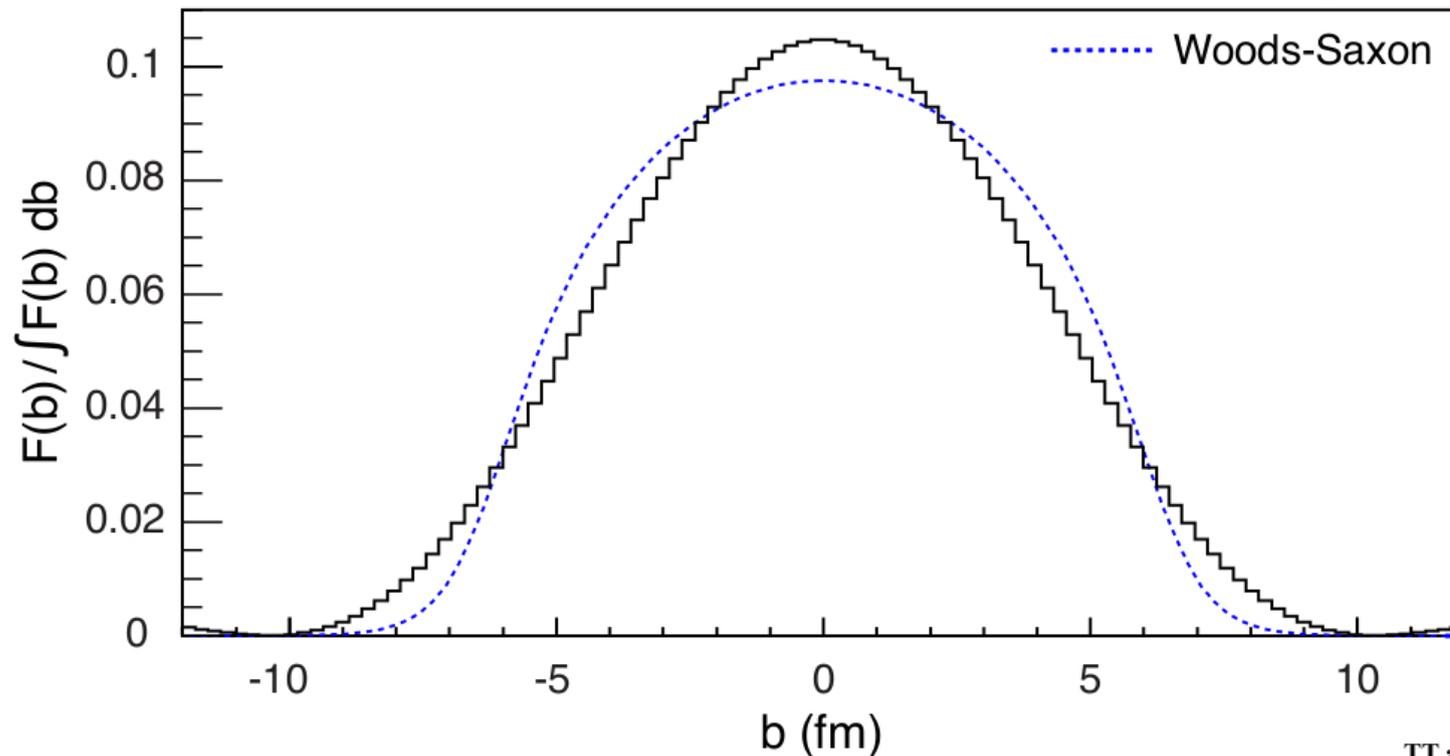
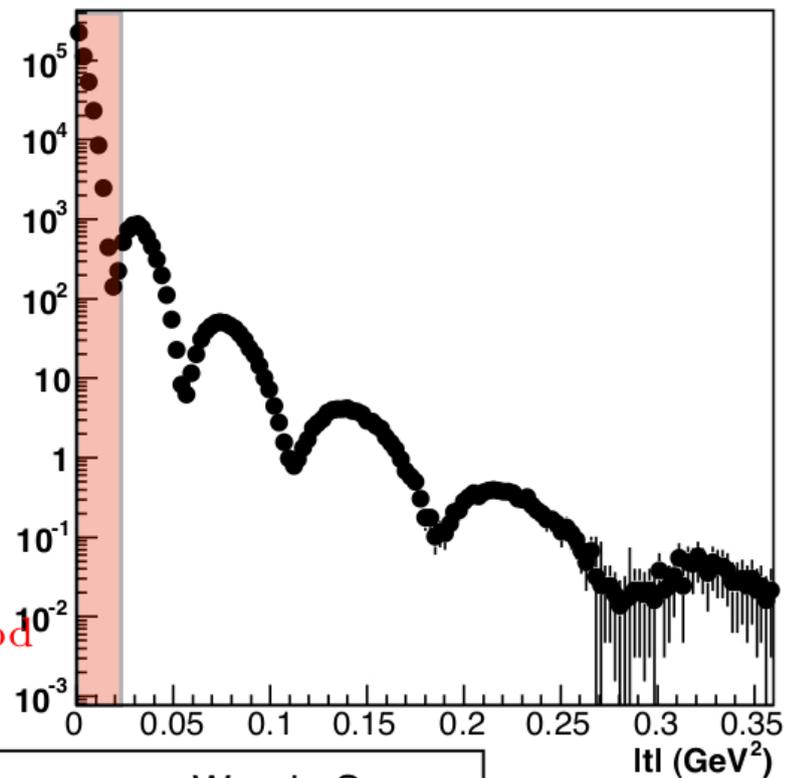
Probing the **spatial** gluon distribution at EIC

$$F(b) = \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\left. \frac{d\sigma_{\text{coherent}}}{dt}(\Delta) \right|_{\text{mod}}}$$



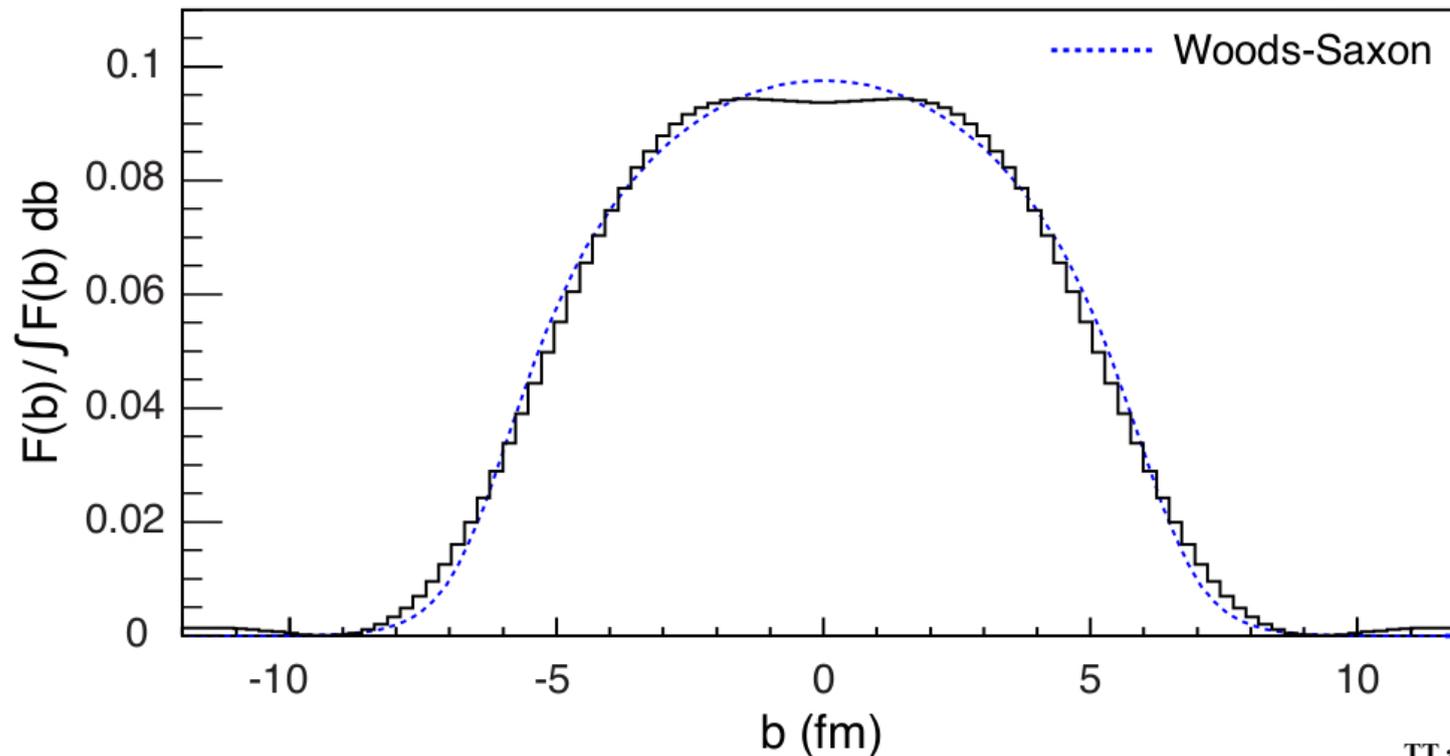
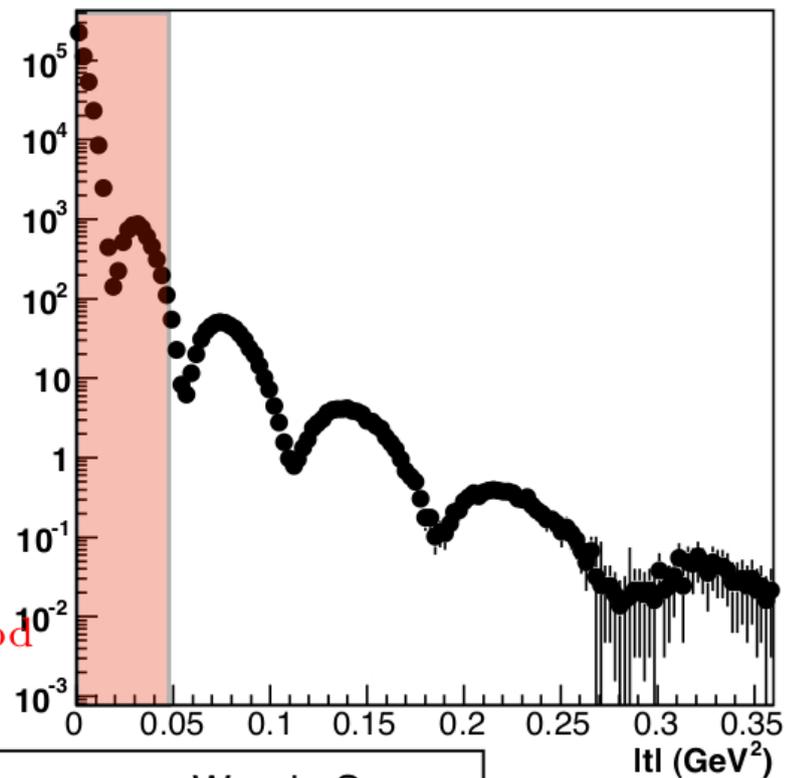
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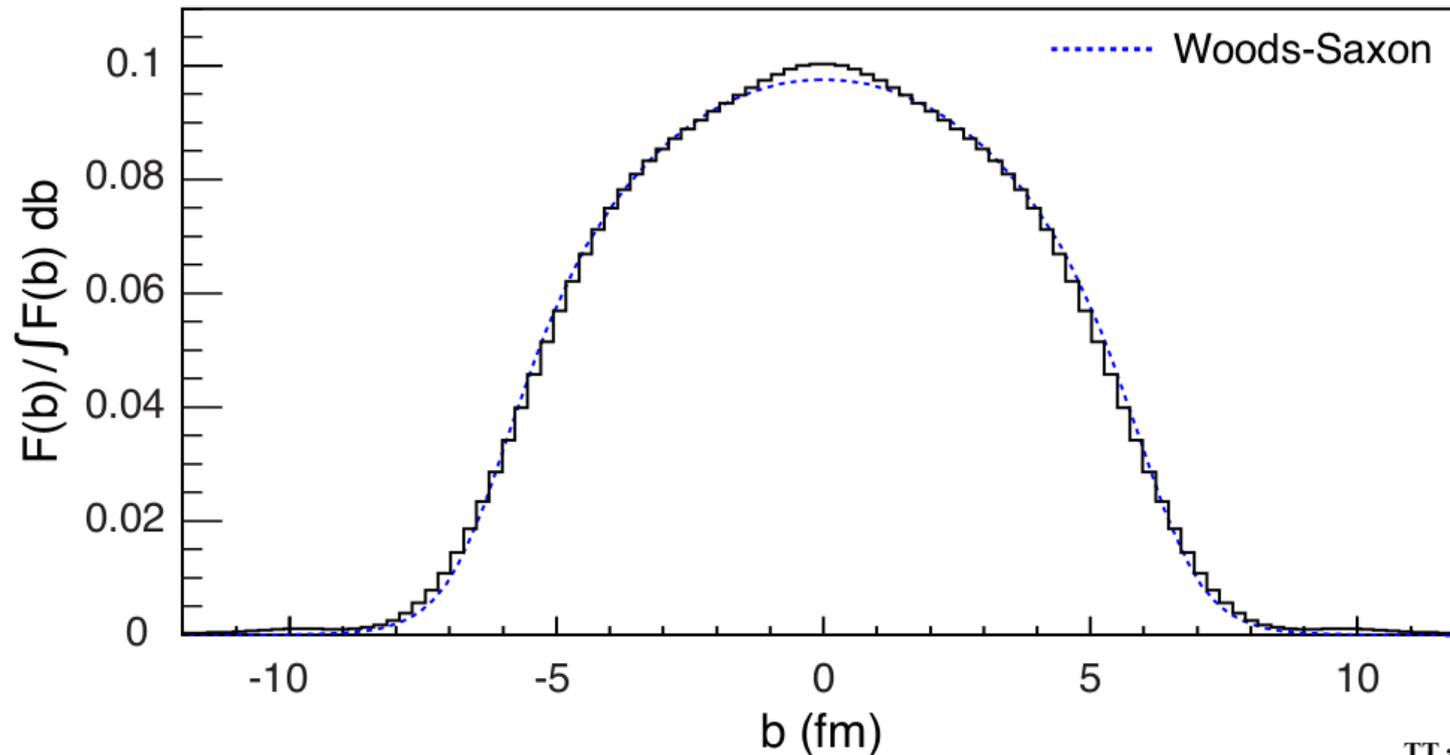
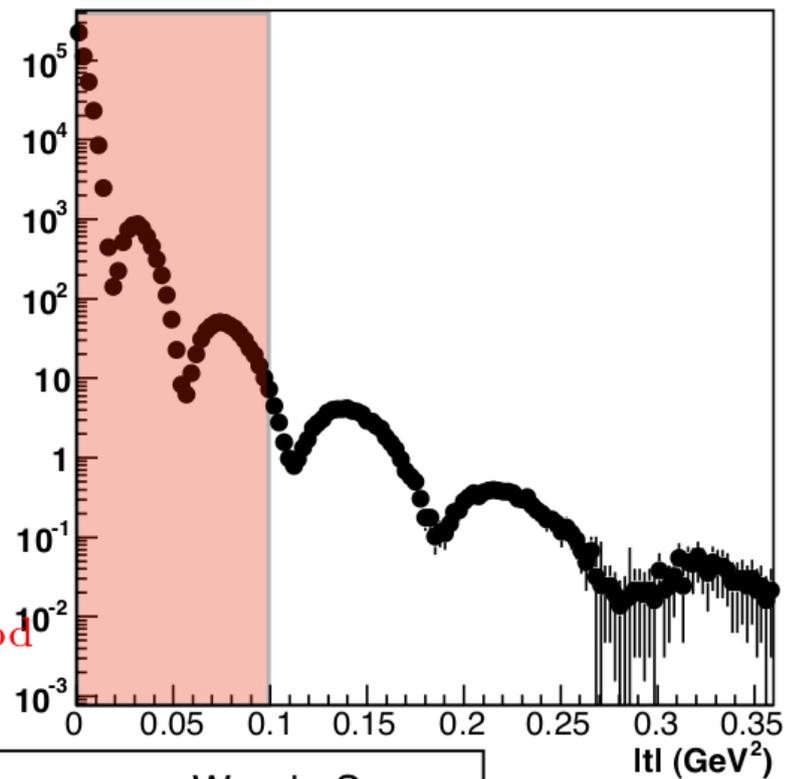
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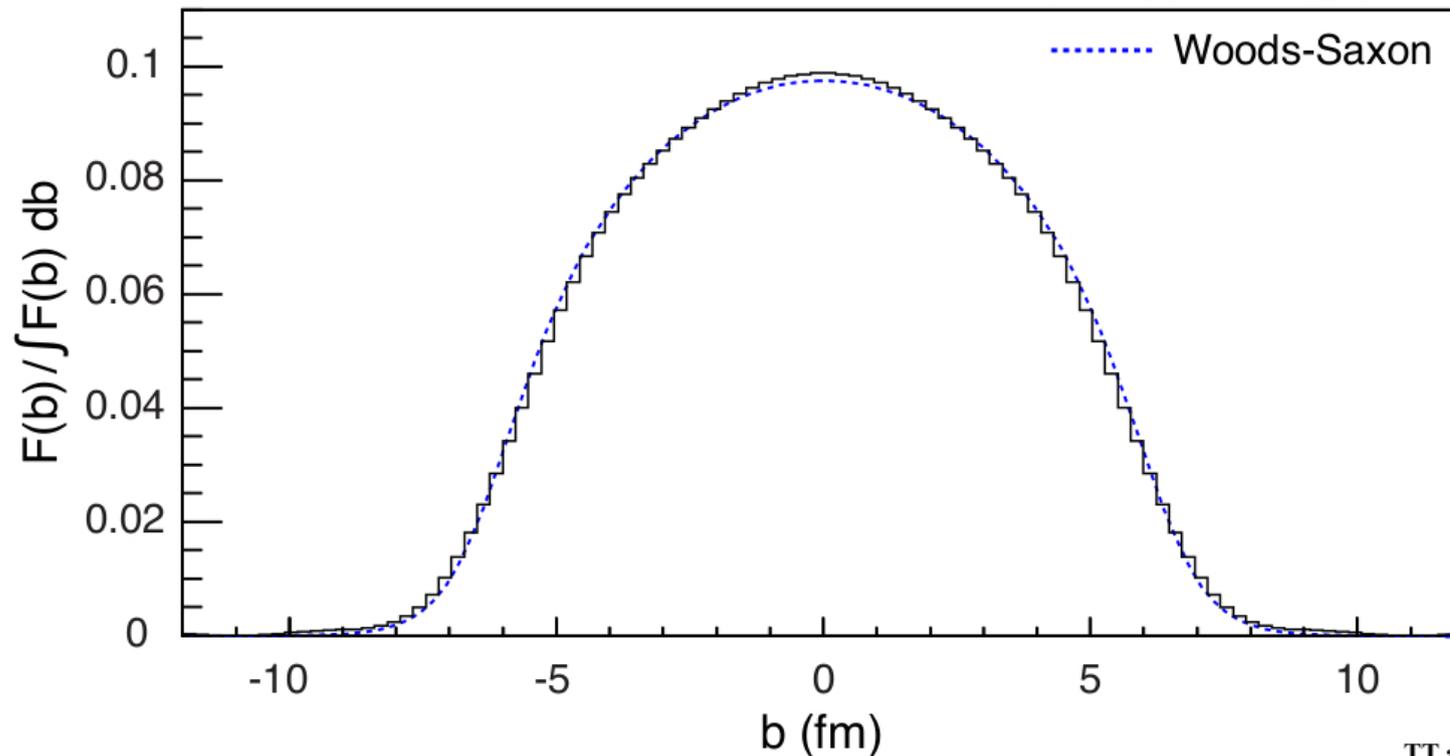
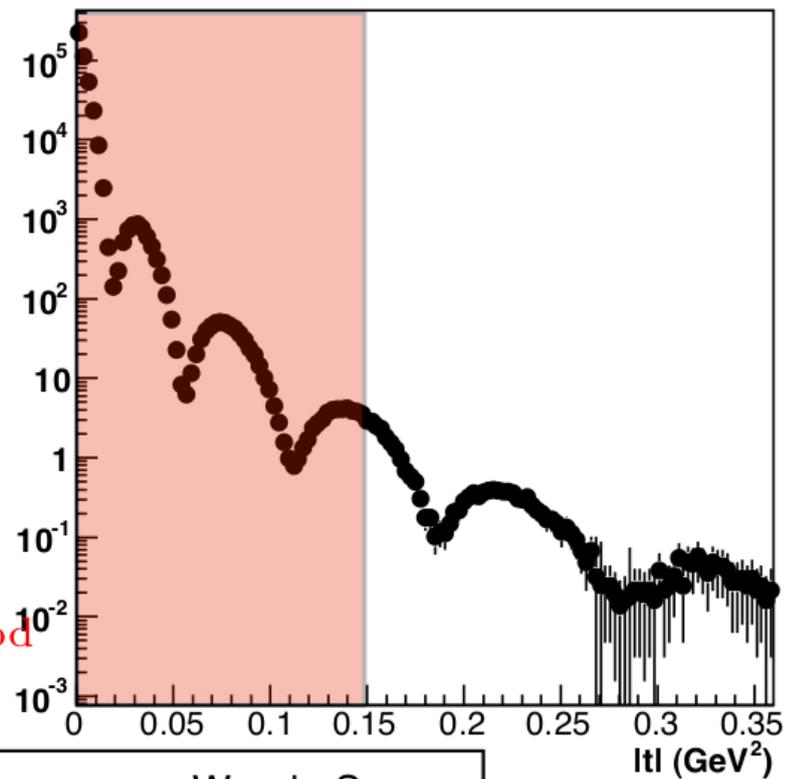
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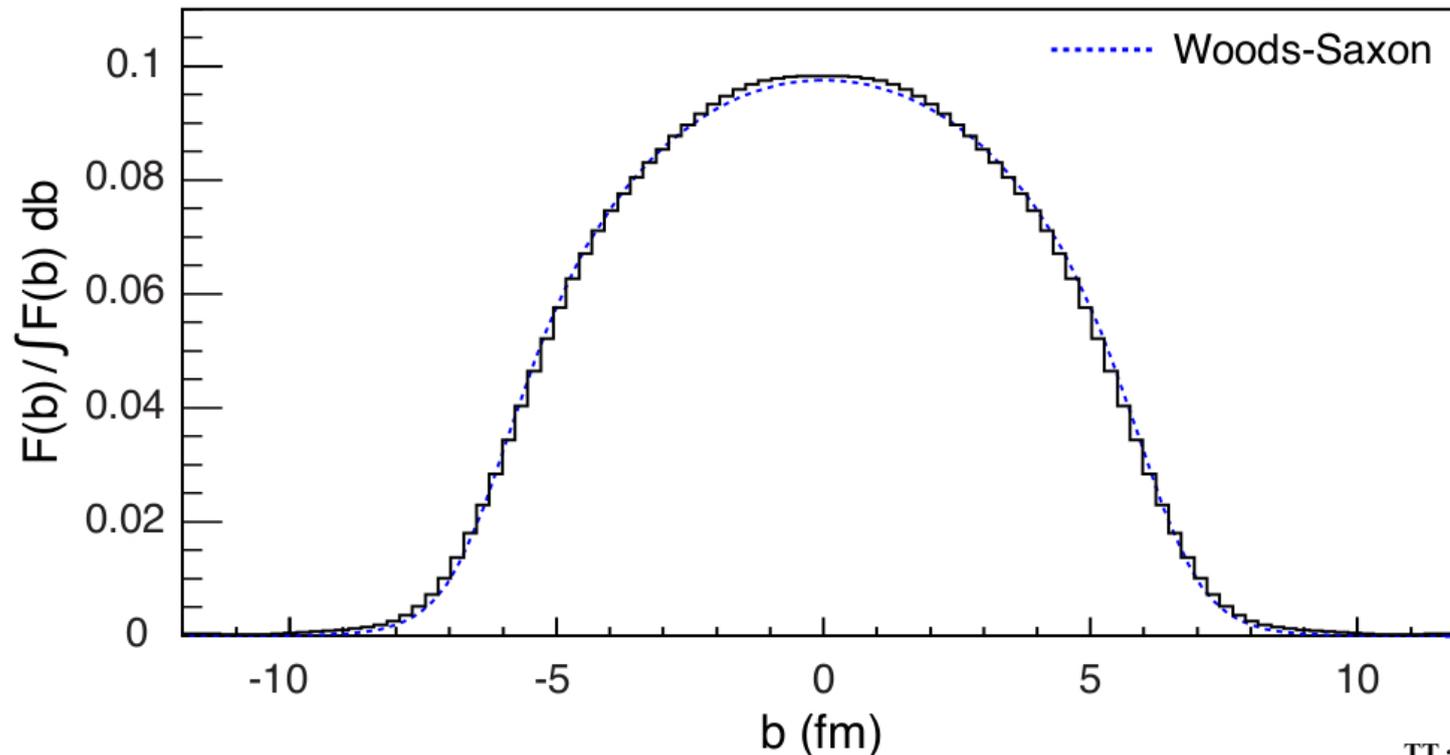
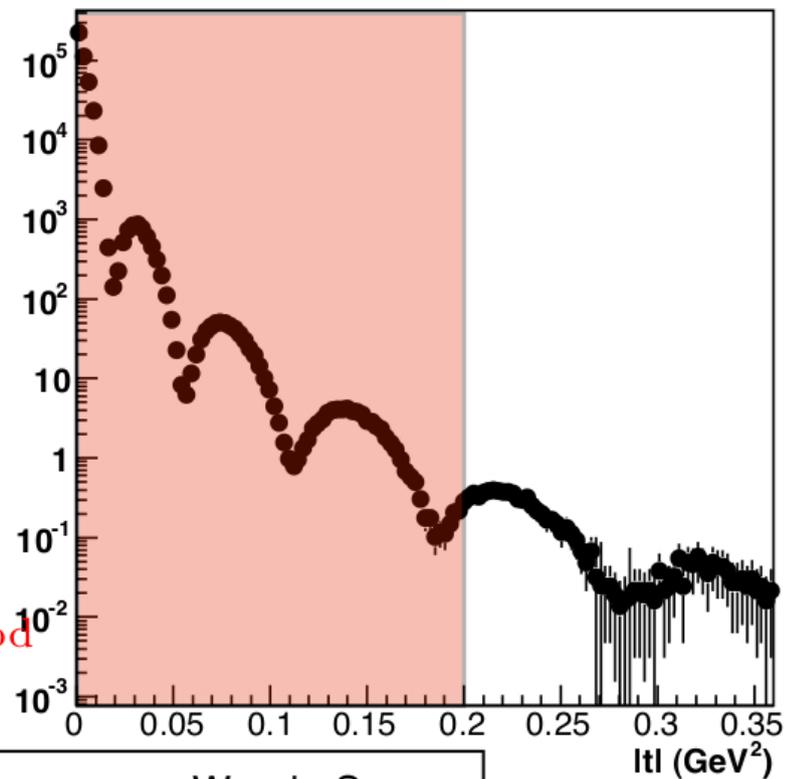
Probing the **spatial** gluon distribution at EIC

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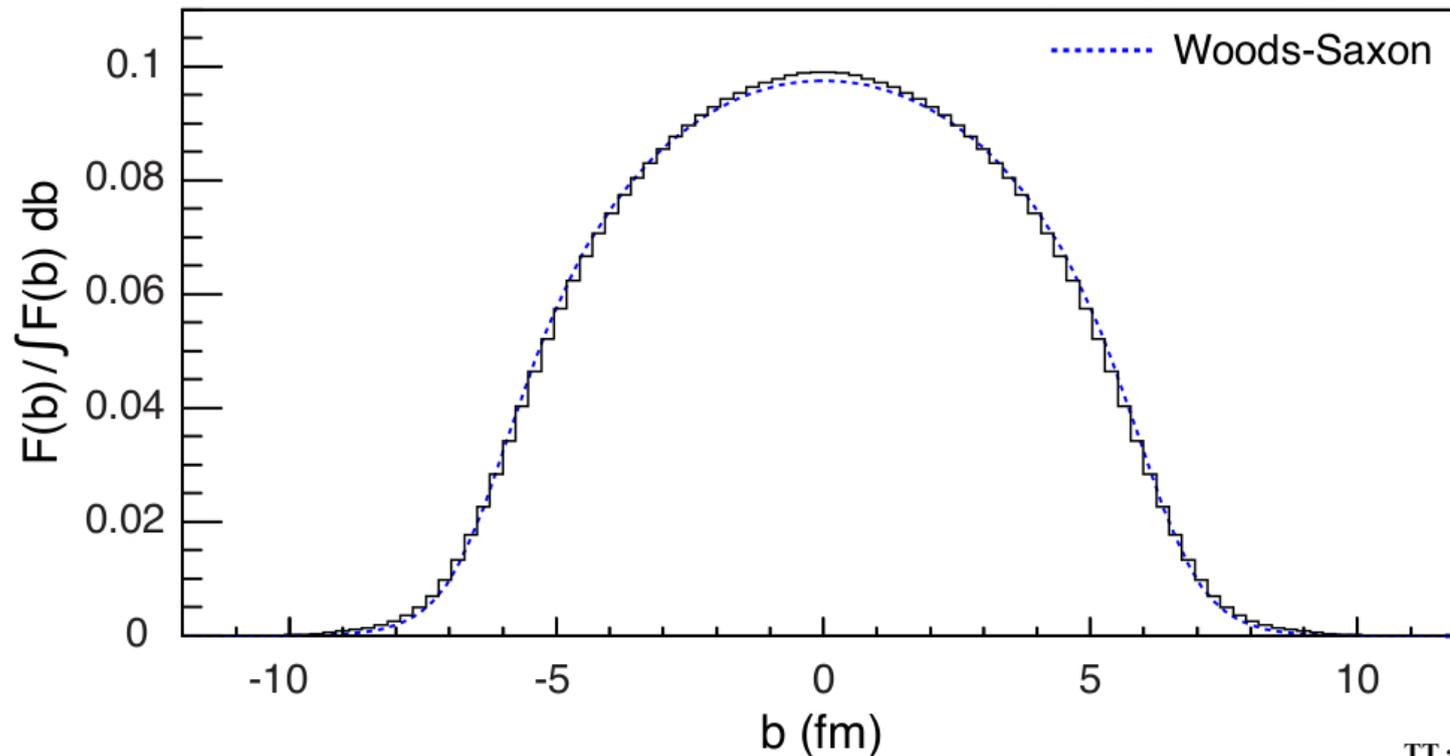
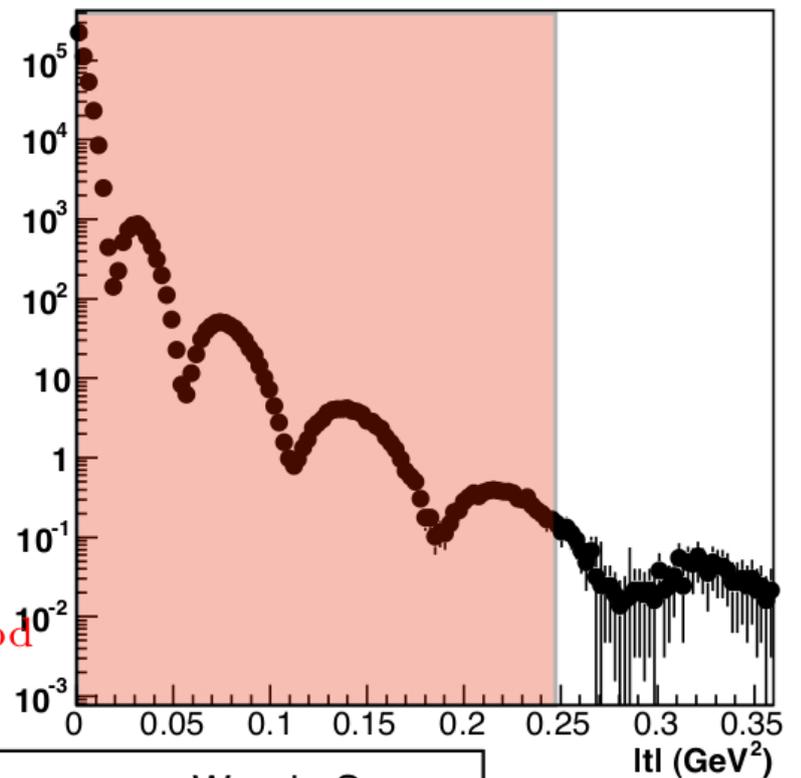
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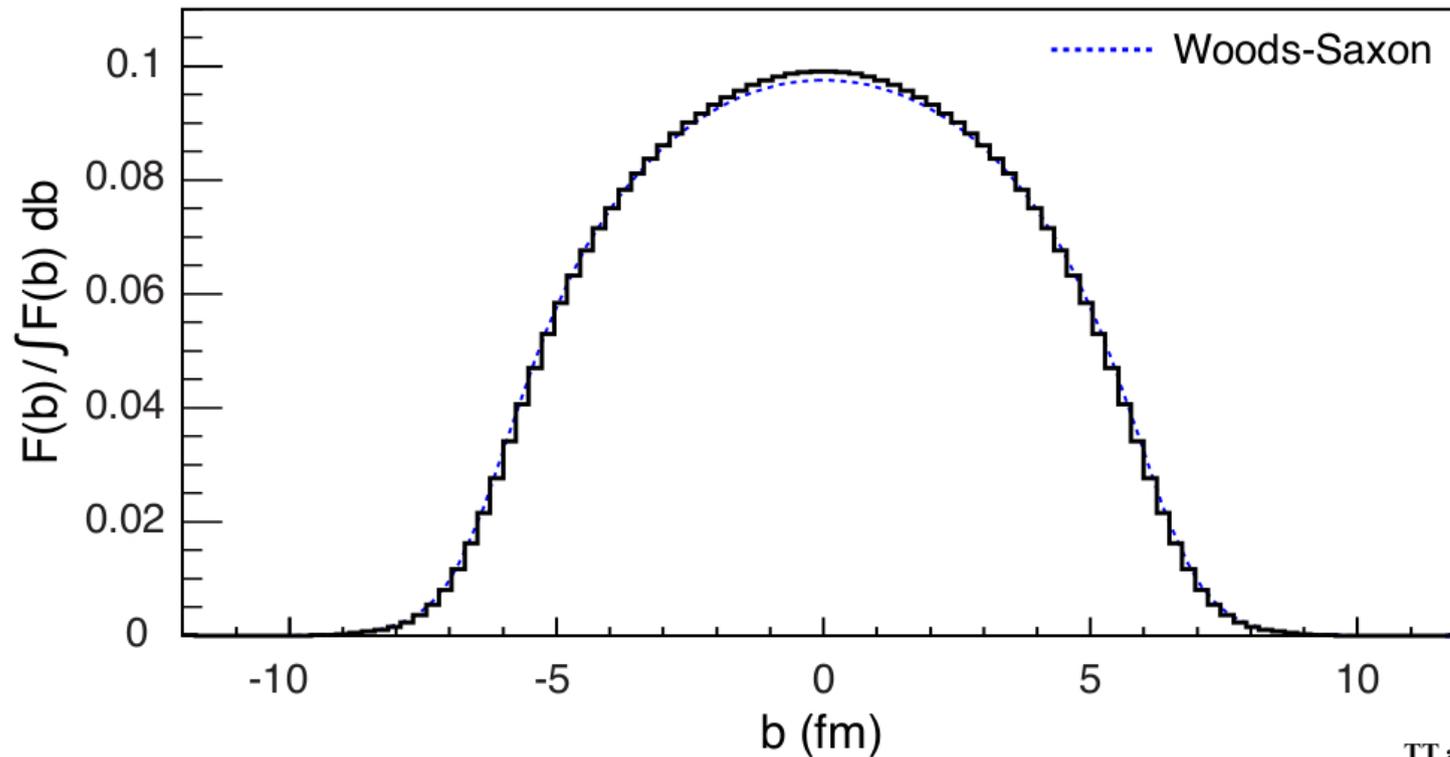
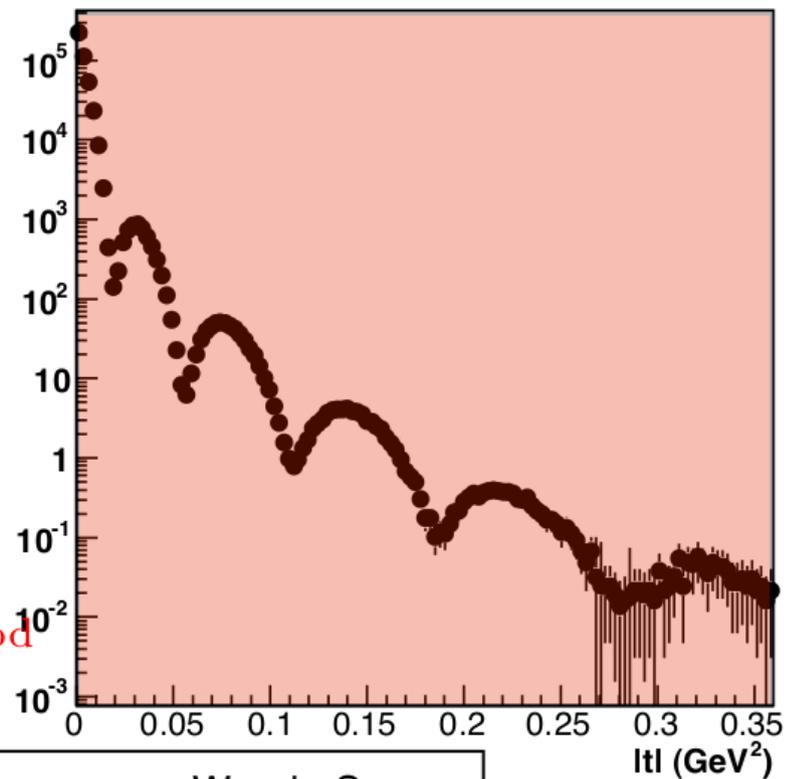
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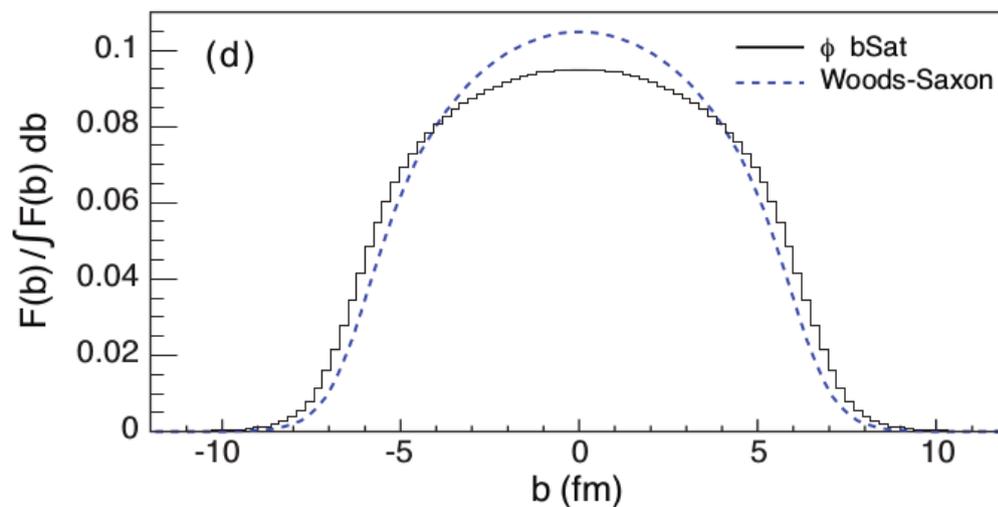
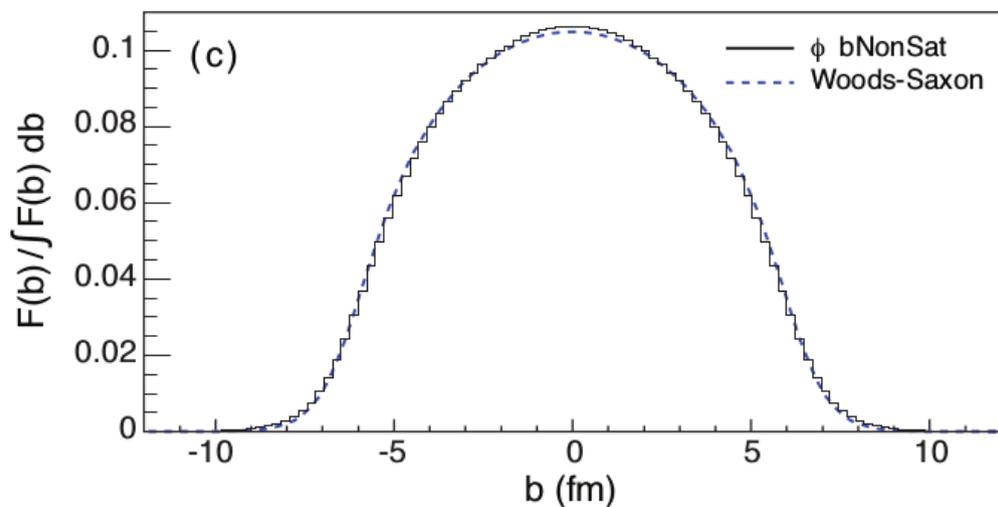
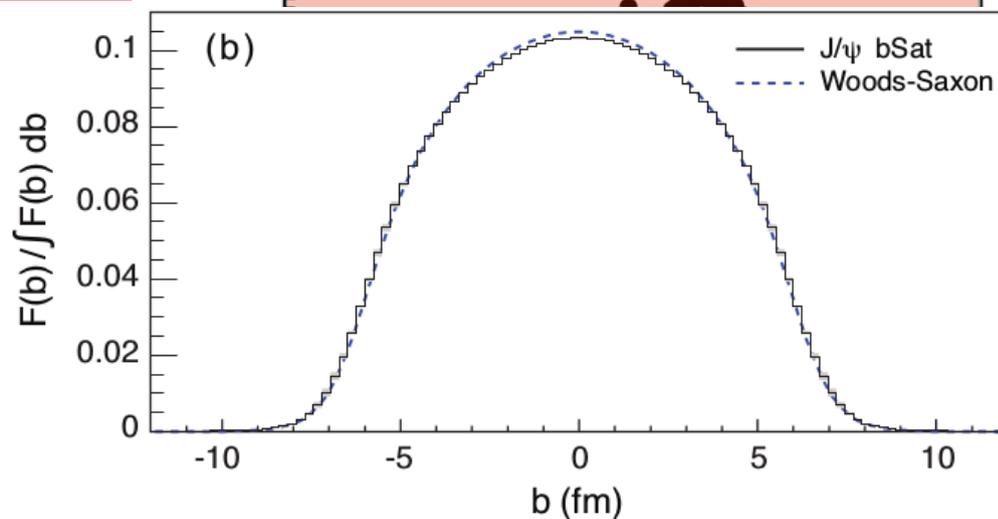
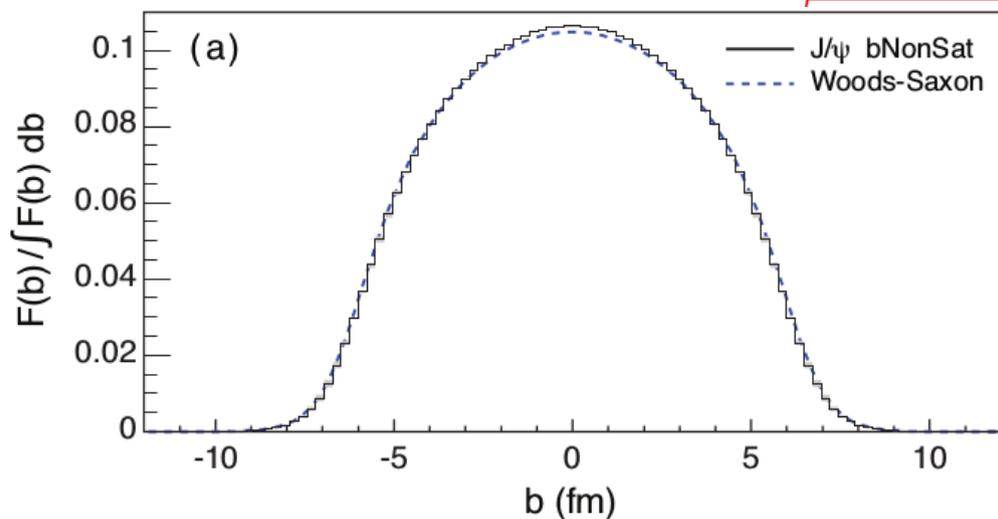


Probing the **spatial** gluon distribution at EIC

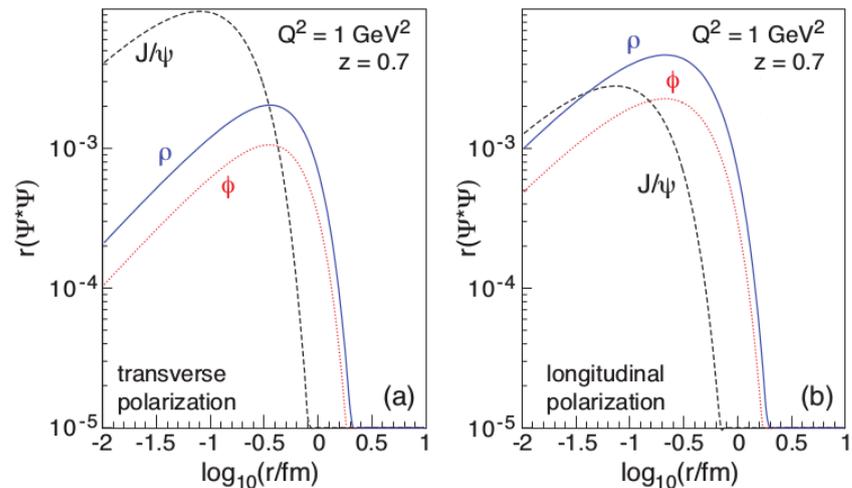
$$F(b) = \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\left. \frac{d\sigma_{\text{coherent}}}{dt}(\Delta) \right|_{\text{mod}}}$$



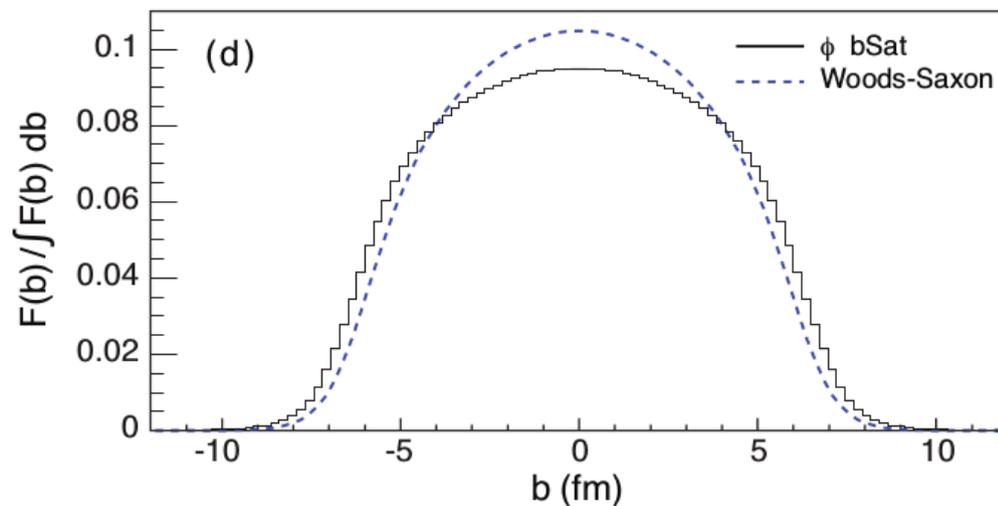
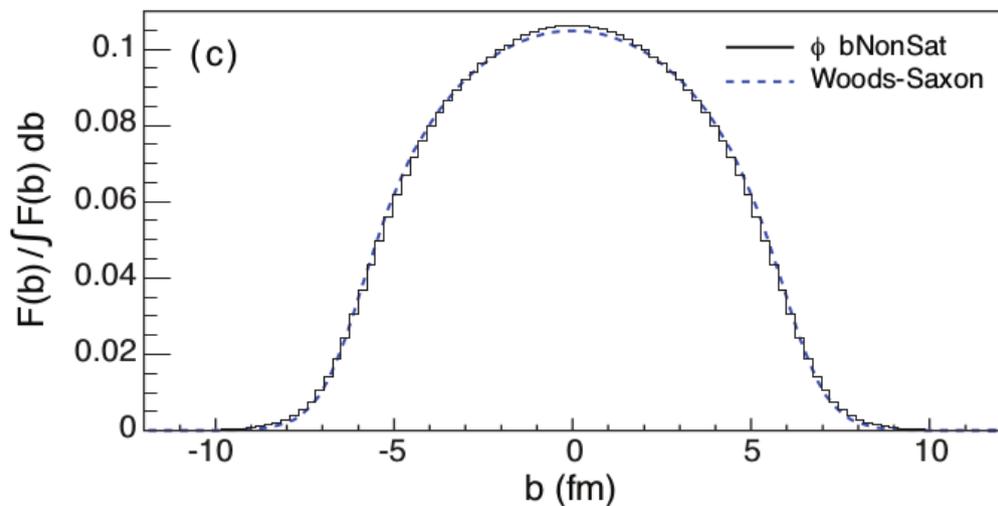
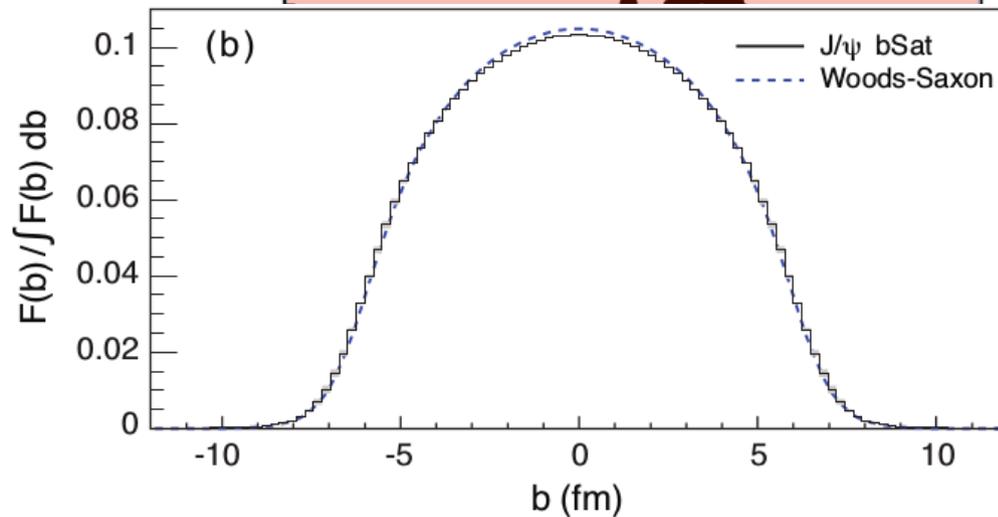
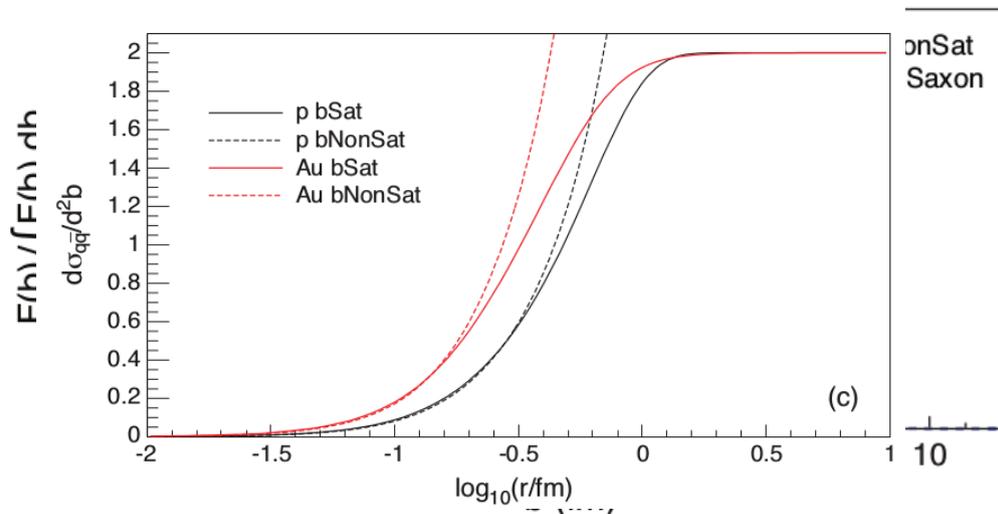
Probing the **spatial** gluon distribution at EIC



F



gluon
EIC

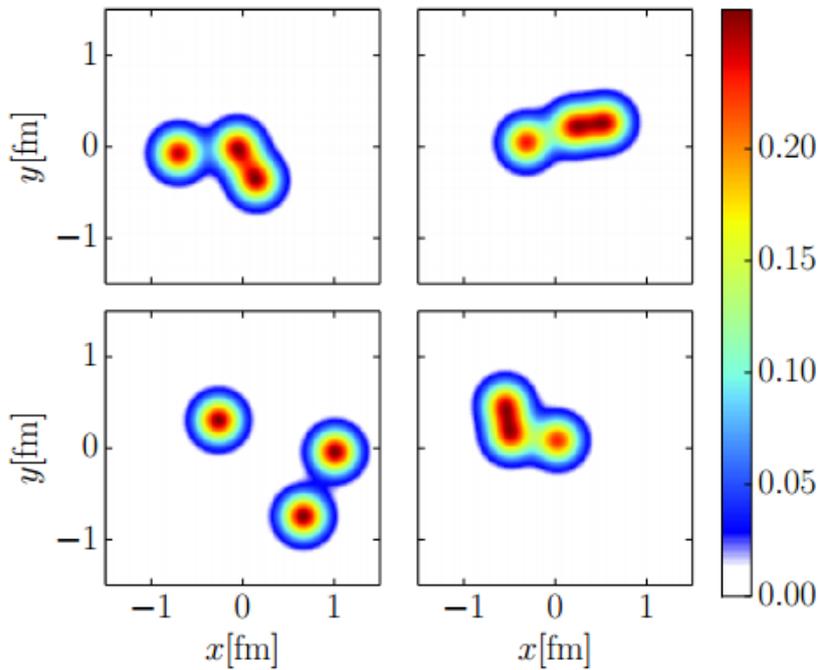


Recent work on Fluctuations

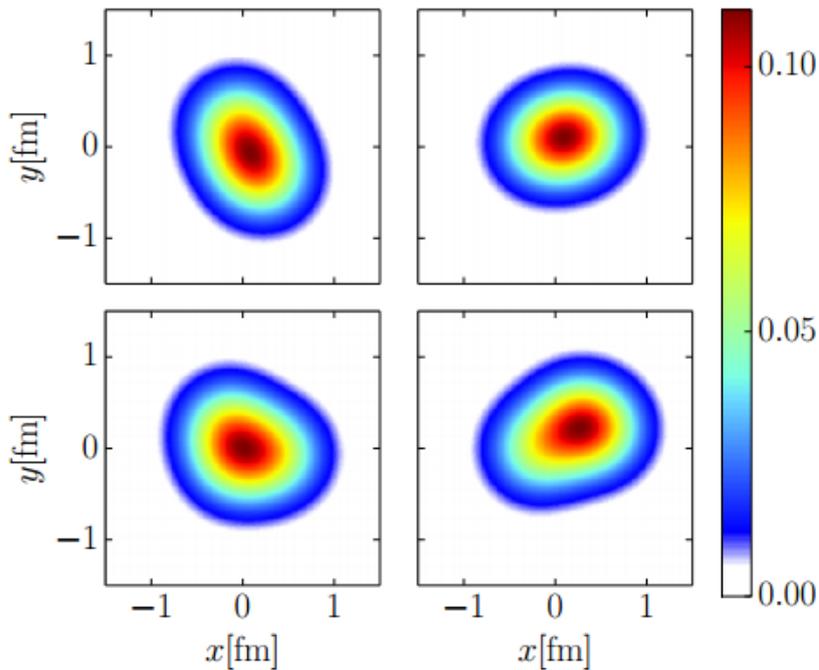
Two different fluctuations:
 Around low scale partons
 + Saturation Scale

In Nuclei:

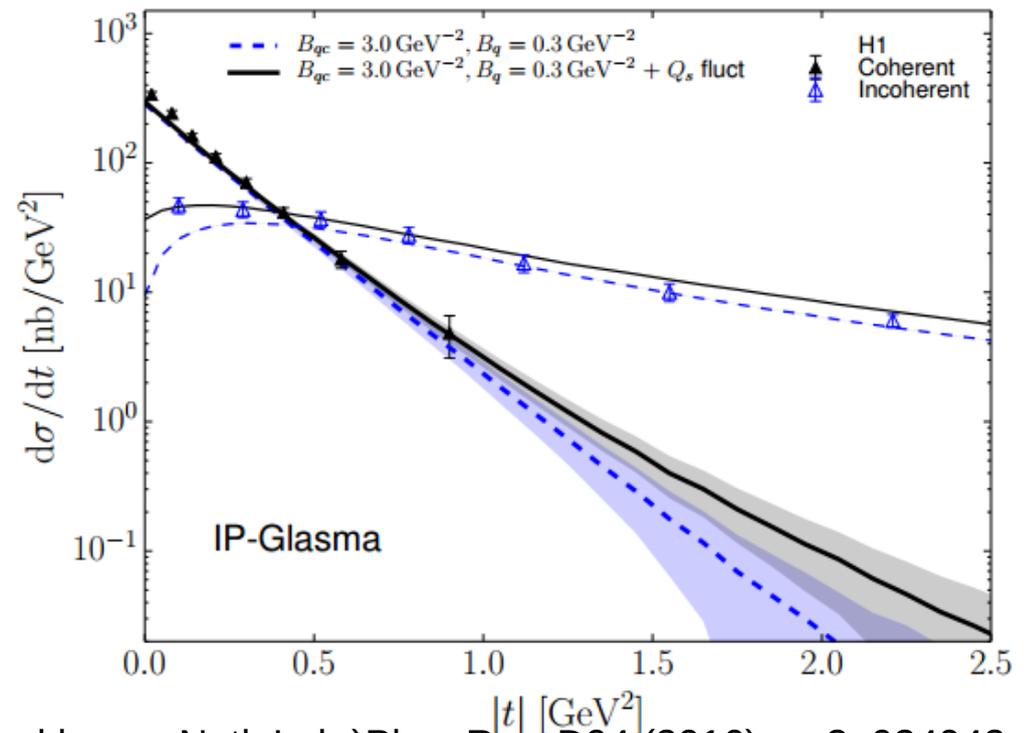
Nucleon Configuration Fluctuations



(a) $B_{qc} = 3.3 \text{ GeV}^{-2}, B_q = 0.7 \text{ GeV}^{-2}$



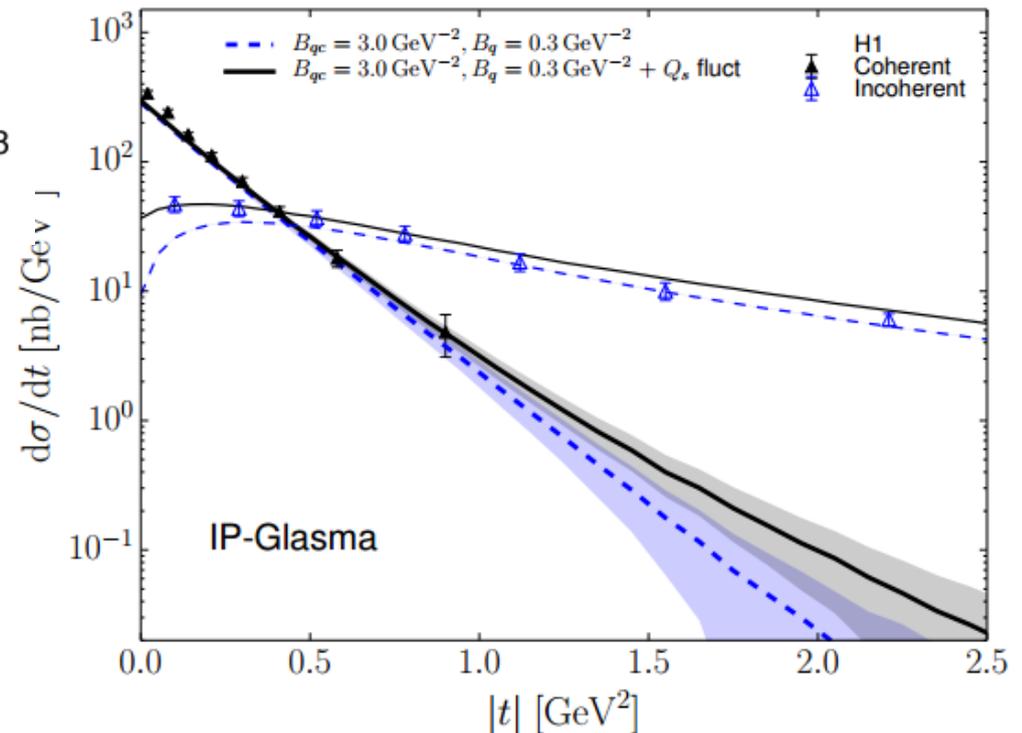
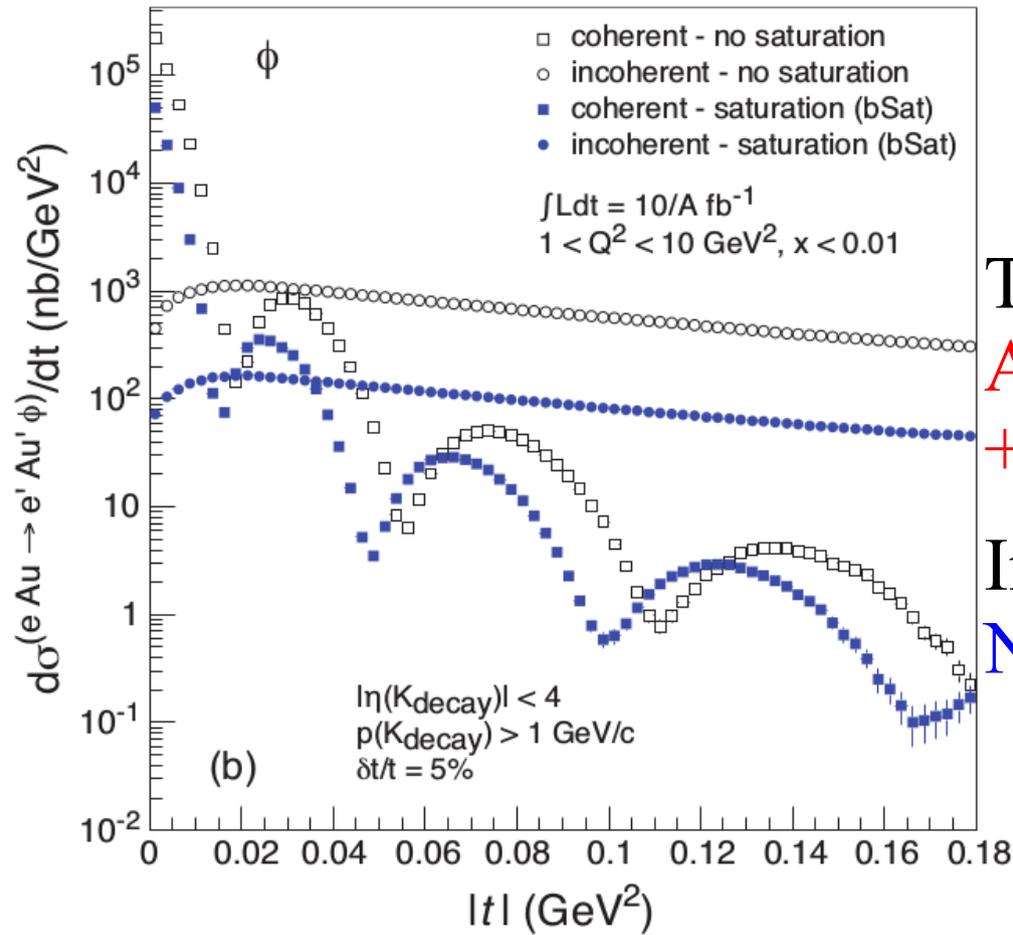
(b) $B_{qc} = 1.0 \text{ GeV}^{-2}, B_q = 3.0 \text{ GeV}^{-2}$



Recent work on Fluctuations

Two different fluctuations:
 Around low scale partons
 + Saturation Scale

In Nuclei:
 Nucleon Configuration Fluctuations



Ongoing Work: **Inclusive Diffraction**

Investigating the Proton Size & Shape in the Dipole Model

Making a new fit for the dipole model parameters to combined HERA data for:

Inclusive DIS

Inclusive Diffraction

F_2 charm

Sartre 2: Inclusive and Exclusive Diffraction

Ultrapерipheral Collisions with Sartre

At first glance, change the photon flux of the electron:

$$\frac{d^2 N_{\gamma,T}^e}{dQ^2 dW^2} = \frac{\alpha_{\text{EM}}(1-y)}{2\pi Q^2 s y} \cdot \left(1 + (1-y)^2 - 2(1-y) \frac{Q_{\text{min}}^2}{Q^2} \right)$$

into that of a proton or ion:

$$\frac{d^2 N_{\gamma,T}^A}{dQ^2 dW^2} = \frac{\alpha_{\text{EM}} Z^2 Q^2}{\pi E_\gamma} \left(\frac{F(Q^2 + E_\gamma^2/\gamma_{\text{Lorentz}}^2)}{Q^2 + E_\gamma^2/\gamma_{\text{Lorentz}}^2} \right)^2 \frac{dE_\gamma}{dW^2}$$

$$F^{(p)}(Q^2) = \frac{1}{(1 + Q^2/0.71 \text{ GeV}^2)^2} \quad F^{(A)}(Q^2) = \frac{4\pi\rho_0}{AQ^3} \frac{\sin(QR) - qR \cos(QR)}{1 + a_{\text{Yukawa}}^2 Q^2}$$

Ongoing work: **Sartre** can soon be used for UPC at RHIC and LHC

Summary

The EIC will profoundly impact our understanding of QCD, with its high energy, its polarised beams, its versatility in accelerating different nuclear species, and its high luminosity.

ep: Precision studies of PDFs, TMDs, and GPDs will lead to the most comprehensive picture of the nucleon ever: its flavor, spin, **spatial and momentum structure and parton fluctuations**

eA: Unprecedented study of matter in a new regime where physics is not described by “ordinary” QCD: non-linear QCD (**saturation**) effects, **properties of gluon distribution (momentum and space)**, in nucleons and nuclei, as well as their **fluctuations** due to different effects.

Only way to measure nuclear initial state.

Sartre: Simulates exclusive diffractive *ep* and *eA*, as well as UPC in pA and AA.

Sartre 2 will be an event generator for **inclusive** diffractive DIS as well.

This physics is tied to a future high energy polarised Electron Ion Collider and cannot be done without it.