# Simulating real-time dynamics of hard probes in nuclear matter on a quantum computer

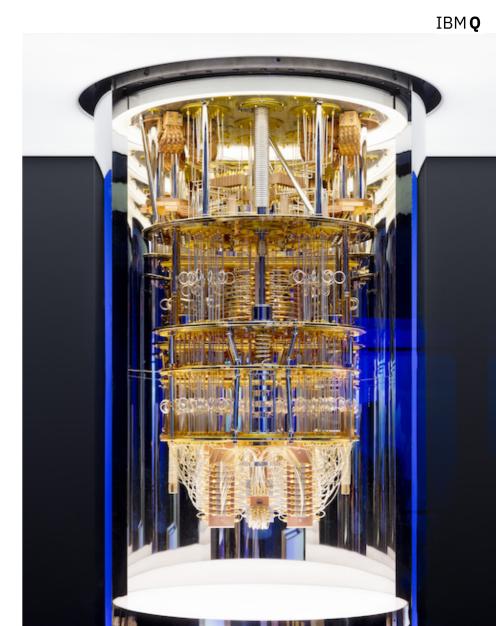
Felix Ringer Lawrence Berkeley National Laboratory

#### arXiv: 2010.03571

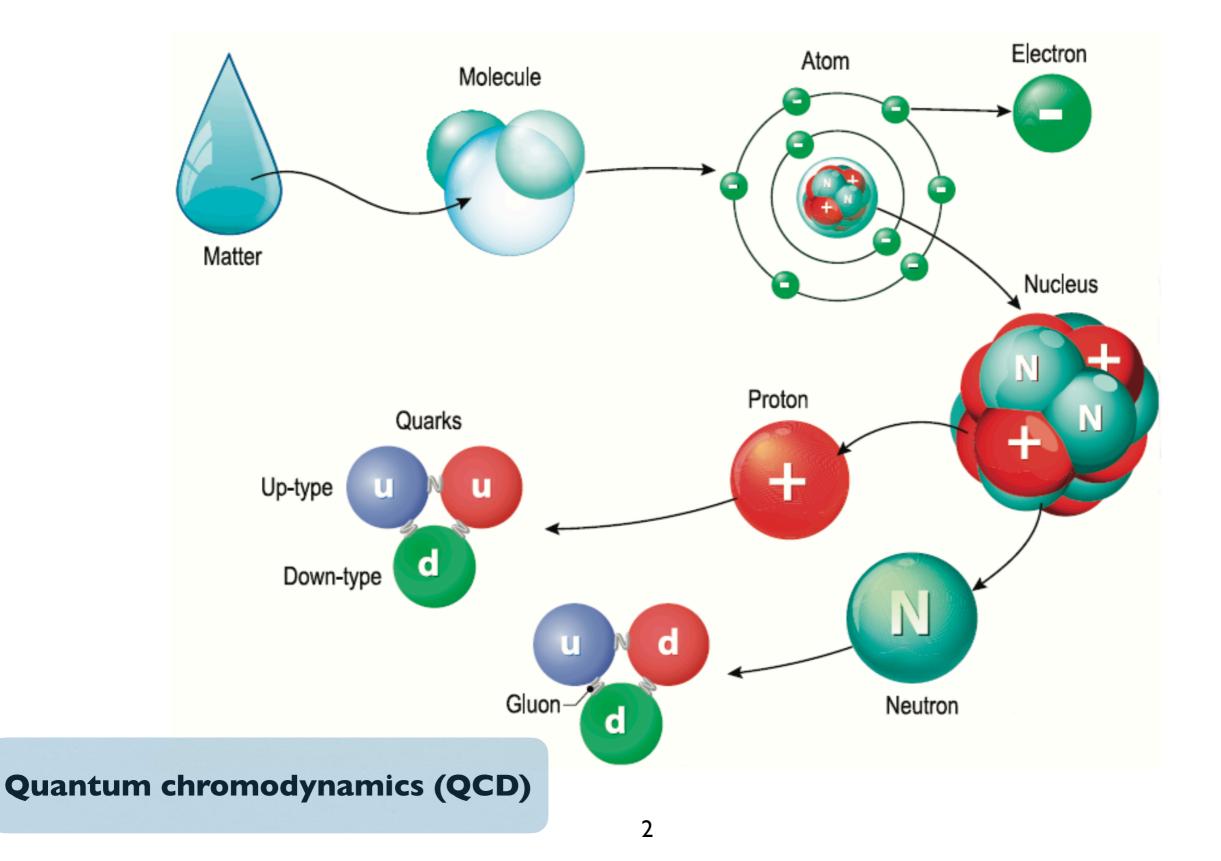
Wibe de JongLBNL (Quantum Information)Mekena MetcalfLBNL (Quantum Information)James MulliganLBNL (Nuclear Science)Mateusz PloskonLBNL (Nuclear Science)Felix RingerMIT (Nuclear Science)



Jefferson Lab January 25, 2021

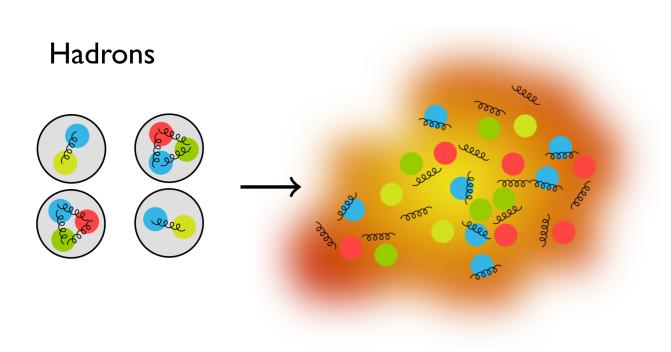


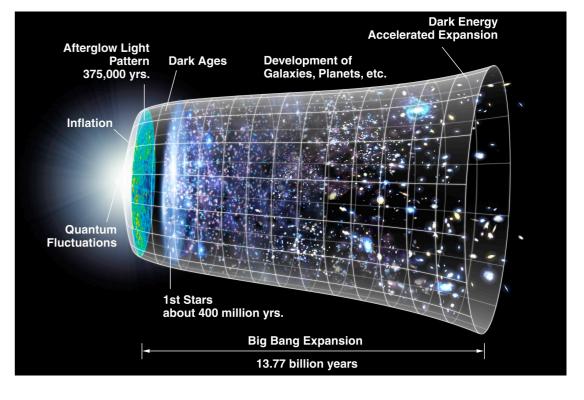
#### Understanding the fundamental structure of matter



### The Quark-Gluon Plasma

If we heat nuclear matter to T = O(100 MeV), quarks and gluons become **deconfined** into a strongly-coupled fluid



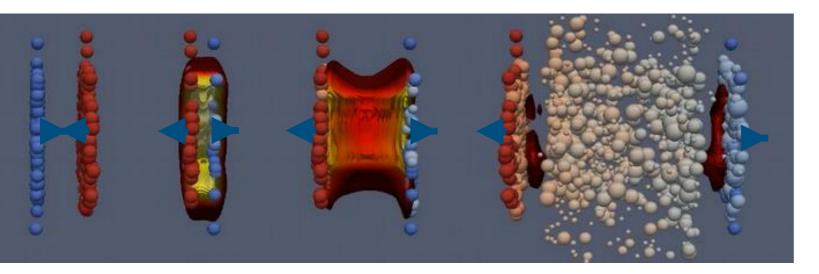


### Heavy-ion collisions



We collide nuclei at

- Large Hadron Collider (LHC)
- Relativistic Heavy Ion Collider (RHIC)



Formation and evolution of the QGP

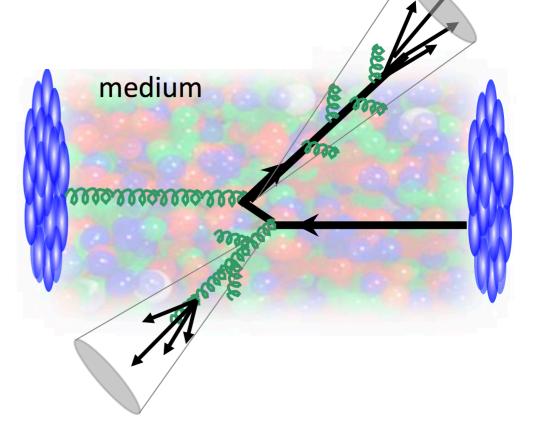
### Hard probes of the Quark Gluon Plasma

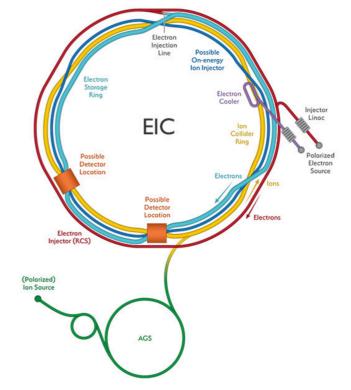
In addition to soft scatterings, there are occasional **hard scatterings** in the collisions

- Highly energetic particles: jets
- Large mass particles: heavy quarks

These "hard probes" interact with the QGP as they traverse it

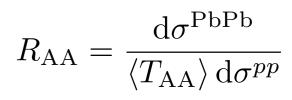
Similar physics relevant in eA collisions at the EIC





### Hard probes — experiment

Experiments measure how cross-sections of hard probes are modified in heavy-ion collisions compared to proton-proton collisions

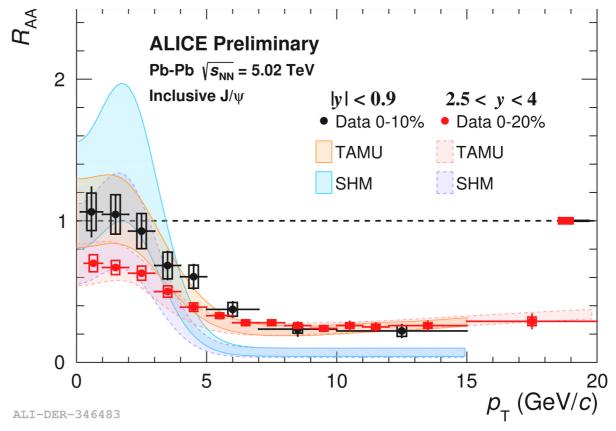


 $\alpha^{\triangleleft}$  1.4 ALICE R = 0.2Pb-Pb 0-10%  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ 1.2 | pp  $\sqrt{s}$  = 5.02 TeV  $|\eta_{iet}| < 0.5 p_{T}^{lead,ch} > 5 \text{ GeV}/c$ ALICE 0-10% SCET Hybrid Model,  $L_{res} = 0$ 0.8 Correlated uncertainty Hybrid Model,  $L_{res} = 2/(\pi T)$ JEWEL, recoils on, 4MomSub Shape uncertainty JEWEL, recoils off 0.6 0.4 0.2 0 50 100 0  $\boldsymbol{p}_{\mathrm{T,jet}}~(\mathrm{GeV/}c)$ PRC 101 034911 (2020)

Jets

#### Heavy quarks





 $c\bar{c}$ 

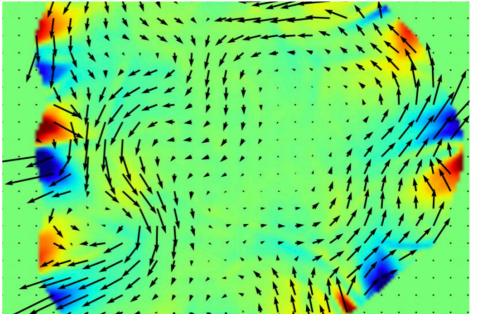
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#### Hard probes — theory

• In vacuum: calculate scattering of asymptotic states using perturbative QCD

Note that there is no sense of "real-time evolution"

• In medium: must combine probe evolution with hydrodynamic evolution of the QGP

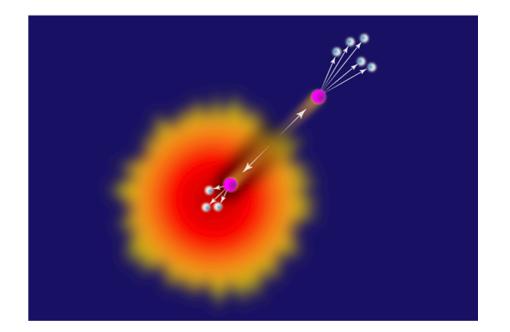


X.N.Wang

Need real-time evolution

Current medium-modified parton shower put in time evolution "by hand"

see e.g. Jetscape

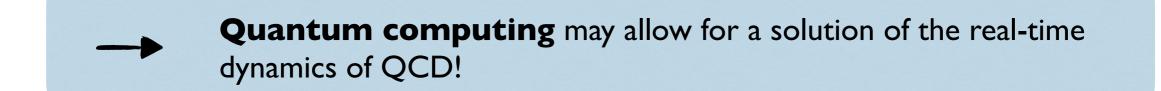


### Can we solve the real-time dynamics of QCD?

 Typical methods in lattice QCD have a sign problem and use imaginary- instead of real time

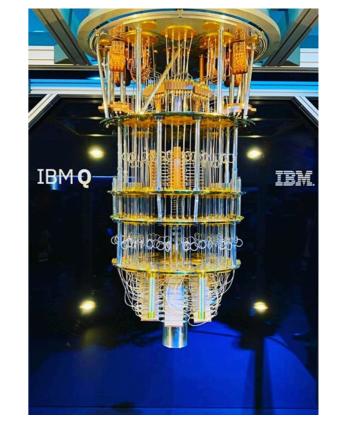
$$\int e^{i\mathscr{L}t} \qquad t \to it$$

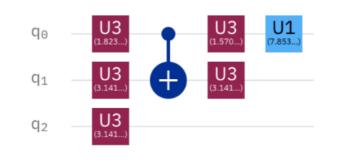
- Can use the Hamiltonian formulation of QCD see e.g. Kogut, Susskind 70s, Preskill `18
  - Theoretical formulation ongoing
  - Gauge, color
  - Large Hilbert space

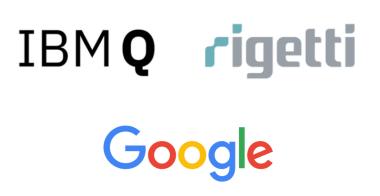


- Significant in process in recent years
- Digital circuit-based e.g. IBMQ, Rigetti
- Noisy Intermediate Scale Quantum (NISQ) era
- Can achieve exponential speedup
- Holds great promise for nuclear physics
- Address computationally expensive problems
- Solve the real-time dynamics of QCD

e.g. Preskill `18, Klco, Savage et al.`18-`20, Cloet, Dietrich et al. `19







## Outline

Open quantum systems in heavy-ion collisions

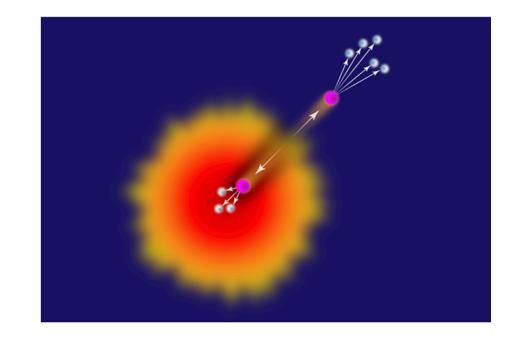
Quantum simulation with IBM Q

Study the real time dynamics of the quantum evolution of probes in the nuclear medium (LHC/RHIC/EIC)

System - Jet/heavy-flavor

**Environment** - Nuclear matter

 $H(t) = H_S(t) + H_E(t) + H_I(t)$ 



Akamatsu, Rothkopf `12-`20, Müller et al `18, Mehen, Yao `18, Qiu, Ringer, Sato, Zurita `19, Vaidya, Yao `20

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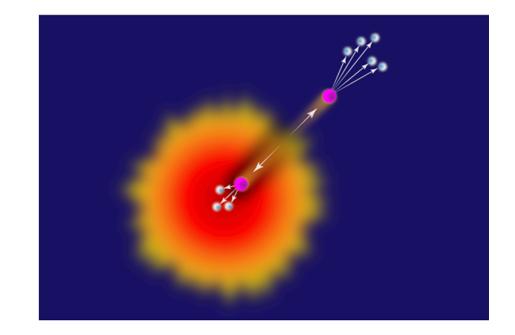
System - Jet/heavy-flavor

**Environment** - Nuclear matter

 $H(t) = H_S(t) + H_E(t) + H_I(t)$ 

The time evolution is governed by the von Neumann equation:

$$\frac{\mathrm{d}}{\mathrm{d}t}\rho^{(\mathrm{int})}(t) = -i\left[H_I^{(\mathrm{int})}(t), \rho^{(\mathrm{int})}(t)\right]$$



where 
$$ho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

Akamatsu, Rothkopf `12-`20, Müller et al `18, Mehen, Yao `18, Qiu, Ringer, Sato, Zurita `19, Vaidya, Yao `20

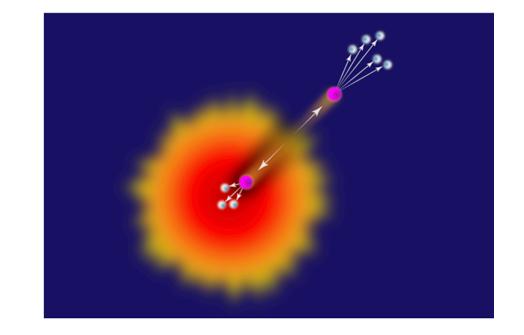
Study the real time dynamics of the quantum evolution of probes in the nuclear medium (LHC/RHIC/EIC)

System - Jet/heavy-flavor

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 $H(t) = H_S(t) + H_E(t) + H_I(t)$ 

In the Markovian limit, the subsystem is described by a **Lindblad equation** 

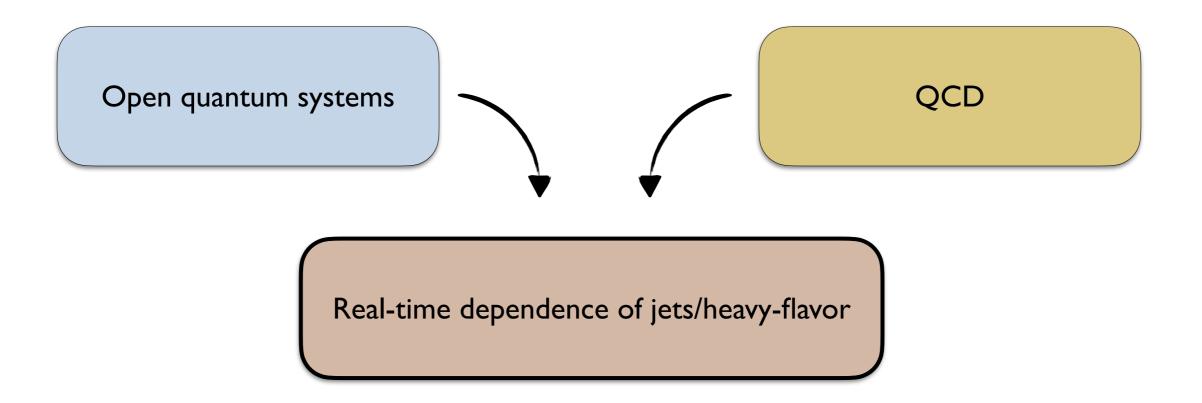


$$\rho_S = \operatorname{tr}_E[\rho]$$
  
$$\frac{\mathrm{d}}{\mathrm{d}t}\rho_S = -i\left[H_S, \rho_S\right] + \sum_{j=1}^m \left(L_j \rho_S L_j^{\dagger} - \frac{1}{2}L_j^{\dagger} L_j \rho_S - \frac{1}{2}\rho_S L_j^{\dagger} L_j\right)$$

#### See also e.g. non-global logarithms and CGC

Neill `15, Armesto et al. `19, Li, Kovner `20

Akamatsu, Rothkopf `12-`20, Brambilla et al. `18, Mehen, Yao `18, Vaidya, Yao `20



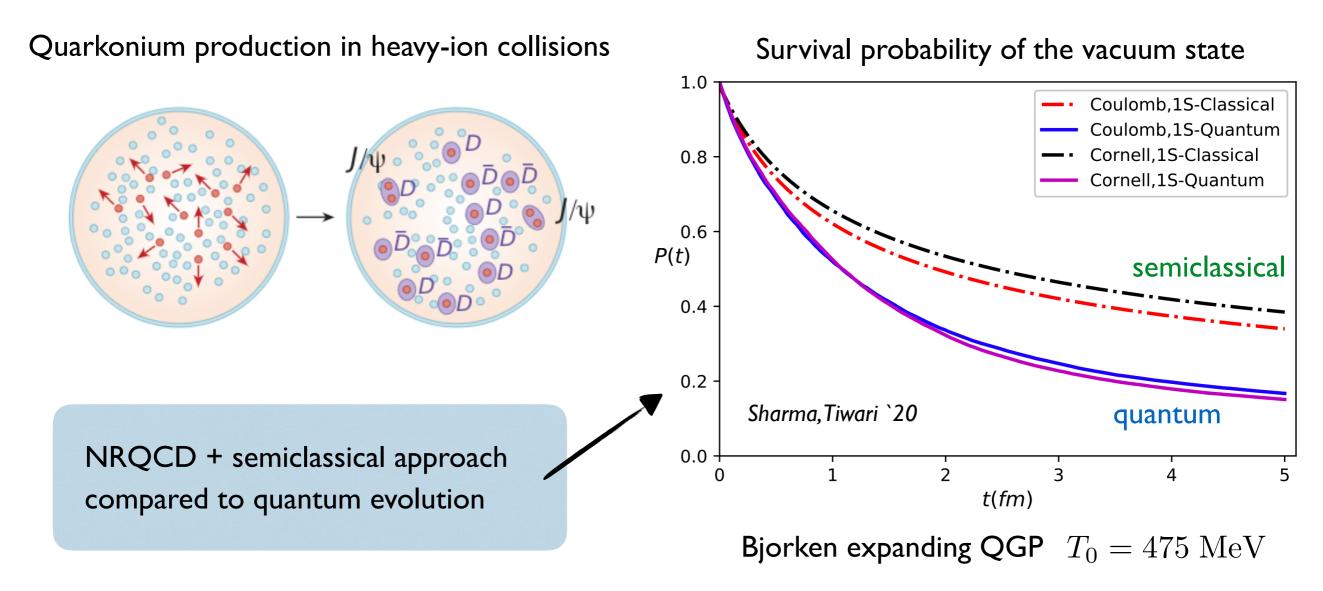
- Currently various approximations are considered Blaizot, Escobedo `18, Yao, Mehen `18, `20
  - Markovian limit
  - Small coupling of system and environment
  - Semi-classical transport

Akamatsu, Rothkopf et al. `12-`20, Brambilla et al. `17-`20 Yao, Mueller, Mehen `18-`20, Sharma, Tiwari `20 Yao, Vaidya `19, Vaidya `20

### Quarkonium suppression

#### **Open quantum system formalism for quarkonia**

Akamatsu, Rothkopf et al. `12-`20, Brambilla et al. `17-`20 Yao, Mueller, Mehen `18-`20, Sharma, Tiwari `20



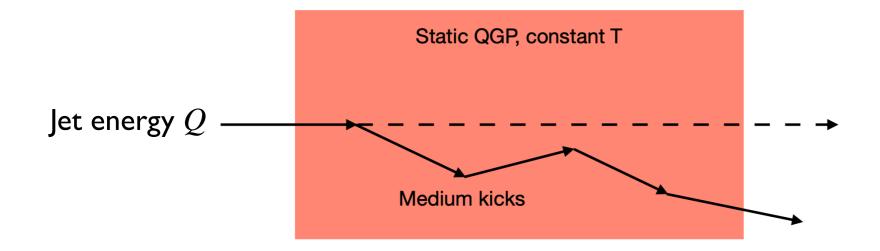
see also Miura, Akamatsu, Asakawa, Rothkopf `19

### Jet broadening

Yao, Vaidya `20

#### **Open quantum system formalism for jets**

First steps in the direction of jet physics



Soft Collinear Effective Theory

 Forward scattering, Glauber gluon exchange

Markovian master equation describes evolution of jet density matrix:

$$\partial_t P(Q,t) = -R(Q)P(Q,t) + \int \widetilde{\mathrm{d}q}K(Q,q)P(q,t)$$

where the probability to be in a given momentum state is:

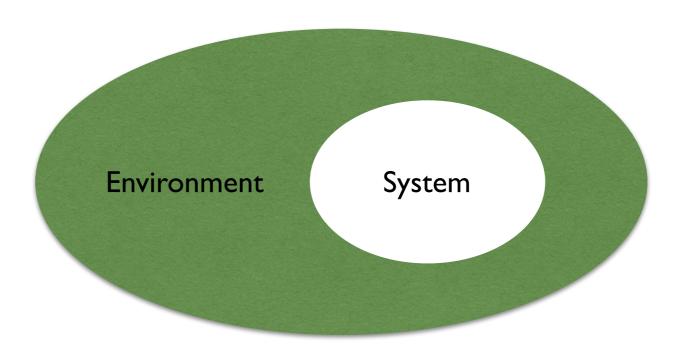
$$P(Q,t) = \langle Q | \rho_S(t) | Q \rangle$$

### Open quantum systems ... more generally

• Basics of Quantum Mechanics/Collapse of the

wave function (measurement theory)

- Cosmology/Inflation
- Qubits



## Outline

Open quantum systems in heavy-ion collisions

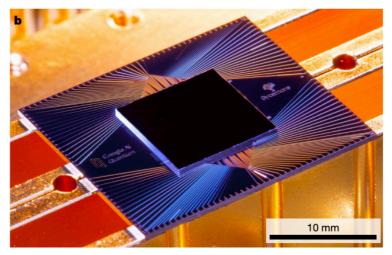
Quantum simulation with IBM Q

### Quantum advantage

Article

Quantum supremacy using a programmable superconducting processor Google

Martinis et al. (2019)



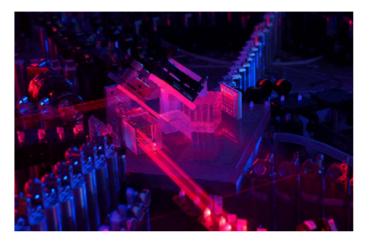
53-qubit superconducting circuit device

Algorithm: sampling of random circuits

## $\mathcal{O}\left(10^3 ight)$ times faster than best classical supercomputers

#### Quantum computational advantage using photons

Han-Sen Zhong<sup>1,2\*</sup>, Hui Wang<sup>1,2\*</sup>, Yu-Hao Deng<sup>1,2\*</sup>, Ming-Cheng Chen<sup>1,2\*</sup>, Li-Chao Peng<sup>1,2</sup>, Yi-Han Luo<sup>1,2</sup>, Jian Qin<sup>1,2</sup>, Dian Wu<sup>1,2</sup>, Xing Ding<sup>1,2</sup>, Yi Hu<sup>1,2</sup>, Peng Hu<sup>3</sup>, Xiao-Yan Yang<sup>3</sup>, Wei-Jun Zhang<sup>3</sup>, Hao Li<sup>3</sup>, Yuxuan Li<sup>4</sup>, Xiao Jiang<sup>1,2</sup>, Lin Gan<sup>4</sup>, Guangwen Yang<sup>4</sup>, Lixing You<sup>3</sup>, Zhen Wang<sup>3</sup>, Li Li<sup>1,2</sup>, Nai-Le Liu<sup>1,2</sup>, Chao-Yang Lu<sup>1,2</sup>, Jian-Wei Pan<sup>1,2†</sup> Science (2020)



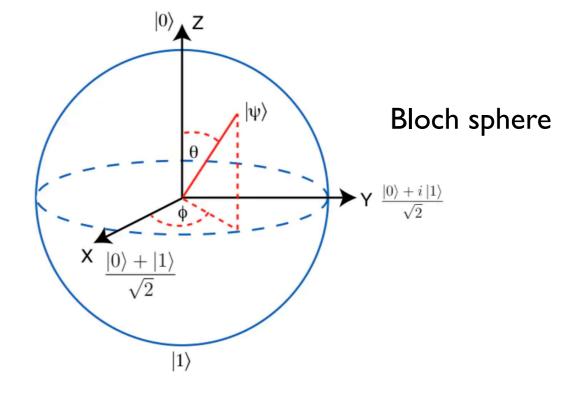
Photonic device — special-purpose

Algorithm: boson sampling

 $\mathcal{O}\left(10^{14}
ight)$  times faster than best classical supercomputers

#### Qubits

- Computational basis |0
  angle, |1
  angle
- Can be in superposition
- Multiple qubits can be entangled



Multi-qubit state vector 
$$\ket{\psi} = \sum_{i=1}^{2^N} a_i \ket{\psi_i}$$

For N qubits, there are  $2^N$  amplitudes

#### 3-qubit example

 $|\psi\rangle = a_1|000\rangle + a_2|001\rangle + a_3|010\rangle + a_4|011\rangle + a_5|100\rangle + a_6|101\rangle + a_7|110\rangle + a_8|111\rangle$ 

#### Potential for exponential speedup

#### Qubits

Multi-qubit state vector

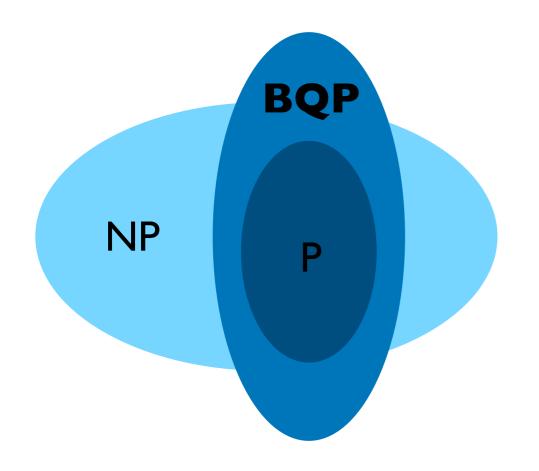
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3-qubit example

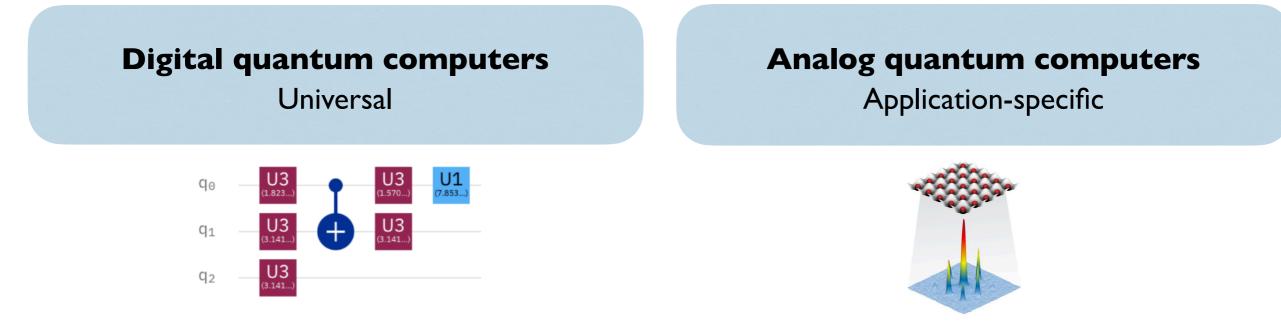
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#### **Complexity classes**



It is expected that quantum computers can solve some classically hard problems with exponential speedup

These include a number of highly impactful problems such as quantum simulation



Both will likely be useful in the "near"-term



Both will likely be useful in the "near"-term

#### • The dream: universal, fault-tolerant digital quantum computer

Shor's and Grover's algorithm, quantum error correction

#### • Noisy Intermediate Scale Quantum (NISQ) era

Decoherence, limited number of qubits, imperfect gates Aim: achieve quantum advantage without full quantum error correction Experimentation and data analysis

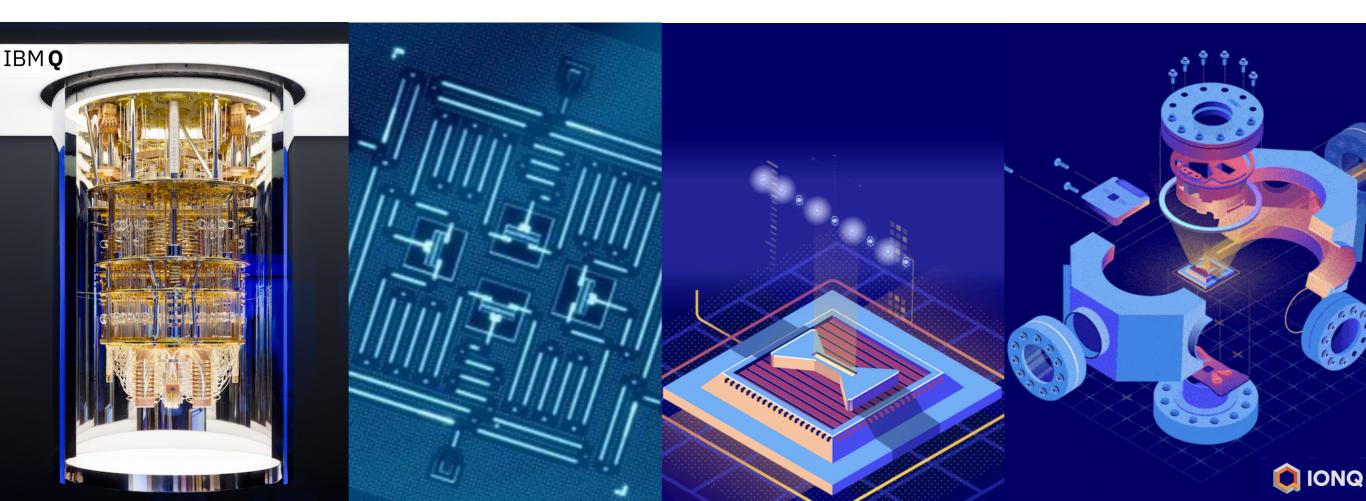
Shor, Preskill, Kitaev, Zoller ...

#### Quantum devices

### Superconducting circuits IBM Q Google rigetti

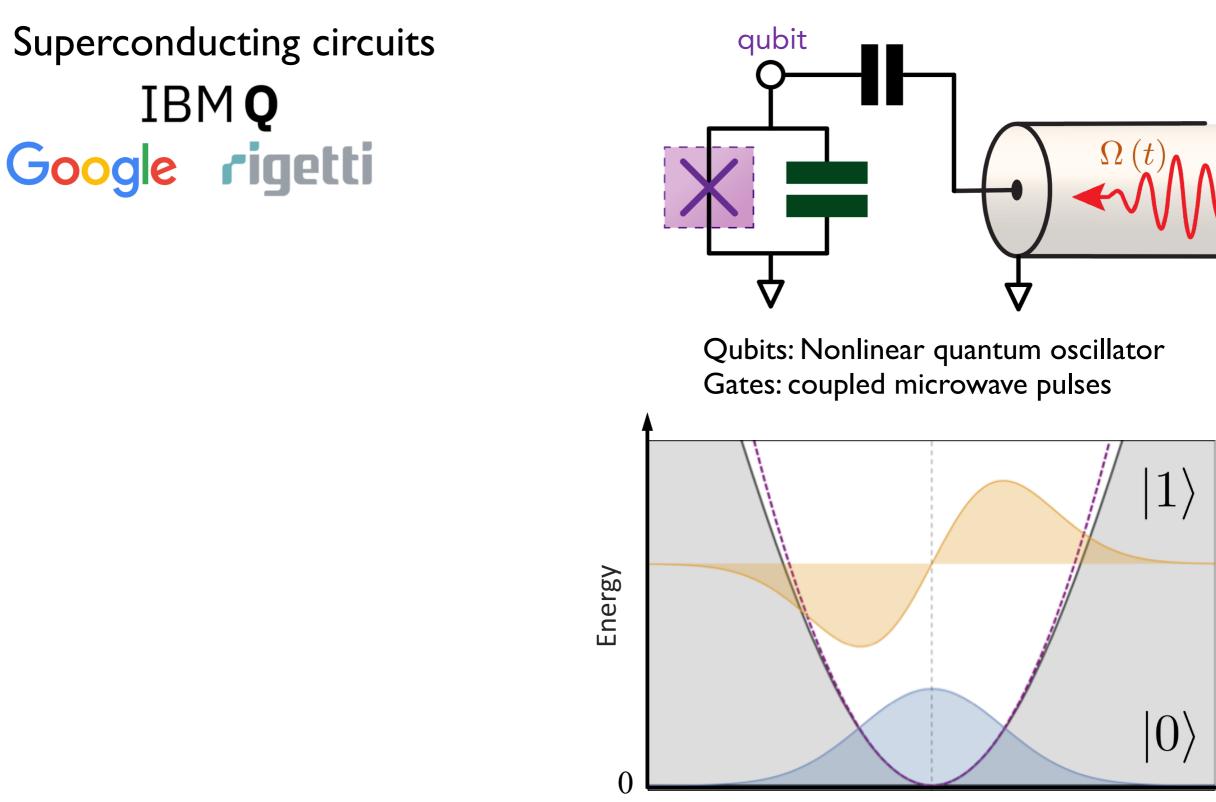
#### And a variety of others...



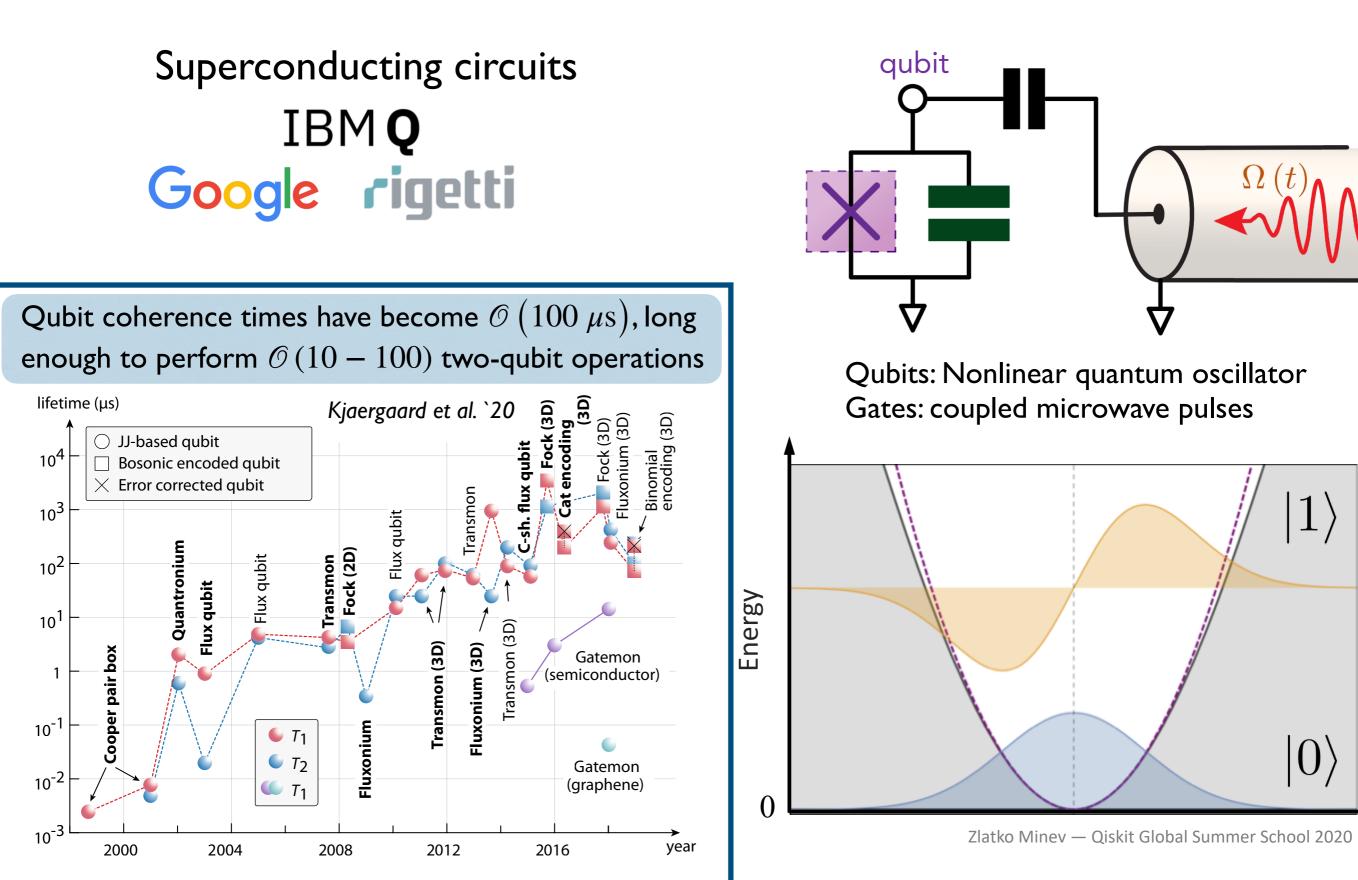


• •

#### Quantum devices

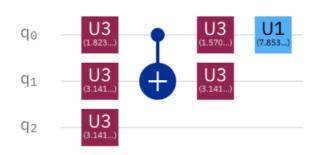


#### Quantum devices

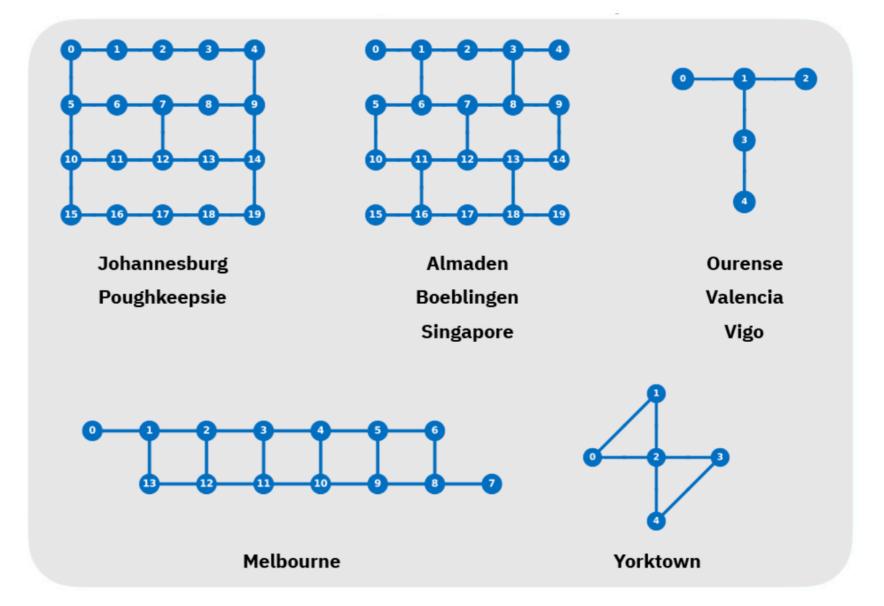


## The IBM Q platform

- up to 65 qubits
- single qubit and CNOT



- Qubits connected to at most 3 others
- Take into account decoherence time, gate errors, read-out error ...





 $|0\rangle_{\mathbf{A}} \mathsf{Z}$ 

 $|\psi\rangle$ 

 $\mathsf{Y} \; \frac{\ket{0} + i \ket{1}}{\sqrt{2}}$ 

### Single- and two-qubit gates

#### • Single qubit rotations

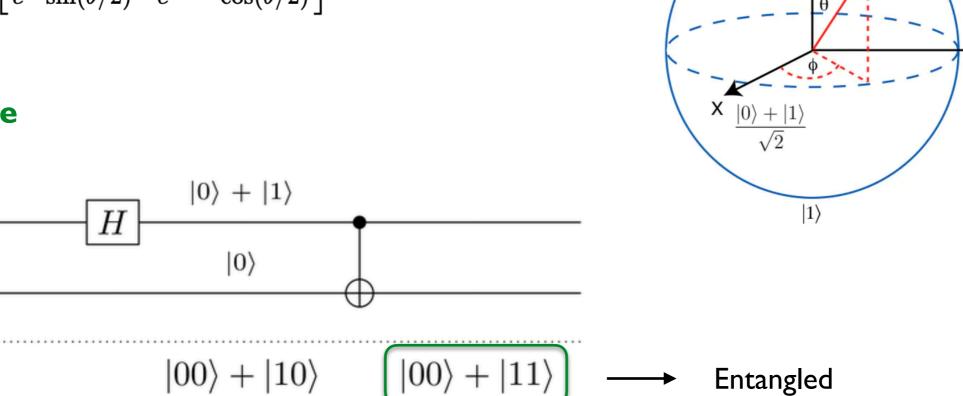
$$U_3( heta,\phi,\lambda) = egin{bmatrix} \cos( heta/2) & -e^{i\lambda}\sin( heta/2) \ e^{i\phi}\sin( heta/2) & e^{i\lambda+i\phi}\cos( heta/2) \end{bmatrix}$$

#### • CNOT gate

 $|0\rangle$ 

 $|0\rangle$ 

 $|00\rangle$ 

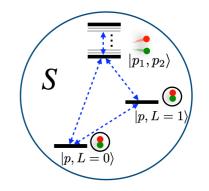


Complete basis of gates

### Closed quantum systems

#### Time evolution of closed systems

• Quantum simulation of the Schrödinger equation





Evolution in time steps  $\Delta t = t/N_{\rm cycle}$ 

• The evolution is unitary and time reversible

For open quantum systems we need to introduce a non-unitarity part

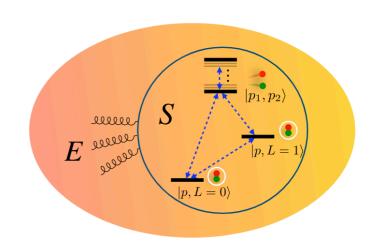
#### Non-unitarity and time irreversible evolution

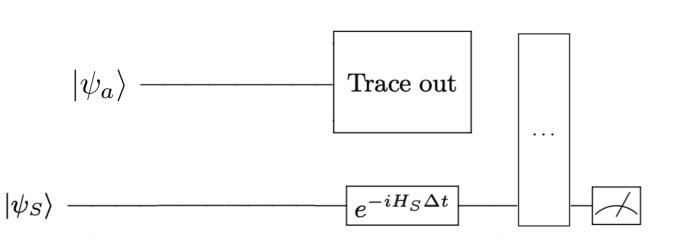
V

$$\frac{\mathrm{d}}{\mathrm{d}t}\rho_S = -i\left[H_S, \rho_S\right] + \sum_{j=1}^m \left(L_j\rho_S L_j^{\dagger} - \frac{1}{2}L_j^{\dagger}L_j\rho_S - \frac{1}{2}\rho_S L_j^{\dagger}L_j\right)$$

• The Stinespring dilation theorem

 $V^{\dagger}$ 





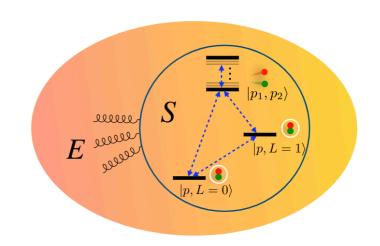
• Introducing and tracing out an ancillary system is not a unitary operation

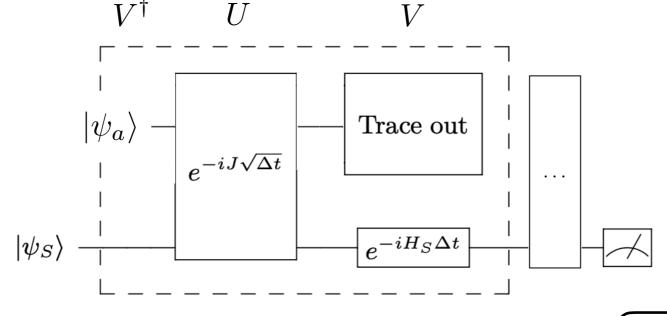
 $V^{\dagger}V = 1 \qquad VV^{\dagger} \neq 1$ 

#### Non-unitarity and time irreversible evolution

$$\frac{\mathrm{d}}{\mathrm{d}t}\rho_S = -i\left[H_S, \rho_S\right] + \sum_{j=1}^m \left(L_j\rho_S L_j^{\dagger} - \frac{1}{2}L_j^{\dagger}L_j\rho_S - \frac{1}{2}\rho_S L_j^{\dagger}L_j\right)$$

#### • The Stinespring dilation theorem





- Introducing and tracing out an ancillary system is not a unitary operation
- Sandwich in between a unitary evolution step
- Evolve in time steps  $\Delta t = t/N_{
  m cycle}$

$$V^{\dagger}V = 1 \qquad VV^{\dagger} \neq 1$$

$$J = \begin{pmatrix} 0 & L_1^{\dagger} & \dots & L_m^{\dagger} \\ L_1 & 0 & \dots & 0 \\ \vdots & \vdots & \ddots & 0 \\ L_m & 0 & \dots & 0 \end{pmatrix}$$

#### Toy model setup

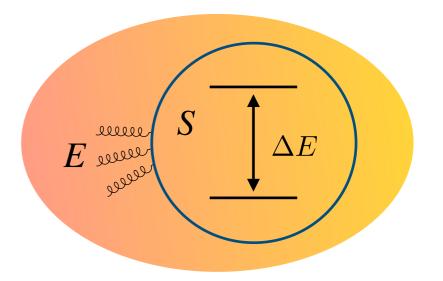
Two-level system in a thermal environment

e.g. bound/unbound  $\,J/\psi,\,c\overline{c}$ 

$$H_S = -\frac{\Delta L}{2} Z$$
$$H_E = \int d^3x \left[ \frac{1}{2} \Pi^2 + \frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} m^2 \phi^2 + \frac{1}{4!} \lambda \phi^4 \right]$$

$$H_I = gX \otimes \phi(x=0)$$

Pauli matrices X, Y, Z, interaction strength g



$$\rho(0) = \rho_S(0) \otimes \rho_E$$
$$\rho_E = \frac{e^{-\beta H_E}}{\operatorname{Tr}(e^{-\beta H_E})}$$

#### Toy model setup

Two-level system in a thermal environment

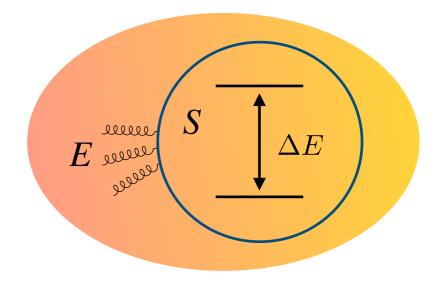
e.g. bound/unbound 
$$J/\psi$$
,  $cc$   
 $H_S = -\frac{\Delta E}{2}Z$   
 $H_E = \int d^3x \left[\frac{1}{2}\Pi^2 + \frac{1}{2}(\nabla\phi)^2 + \frac{1}{2}m^2\phi^2 + \frac{1}{4!}\lambda\phi^4\right]$ 

$$H_I = gX \otimes \phi(x=0)$$

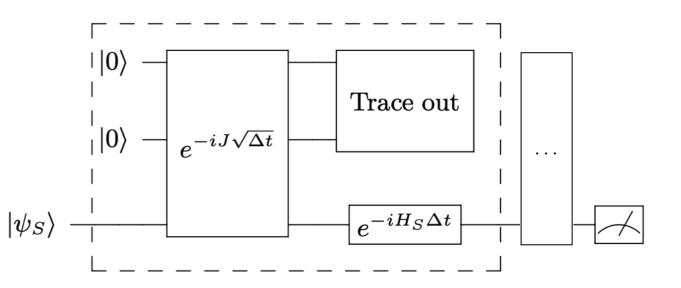
Pauli matrices X, Y, Z, interaction strength g

Lindblad operators

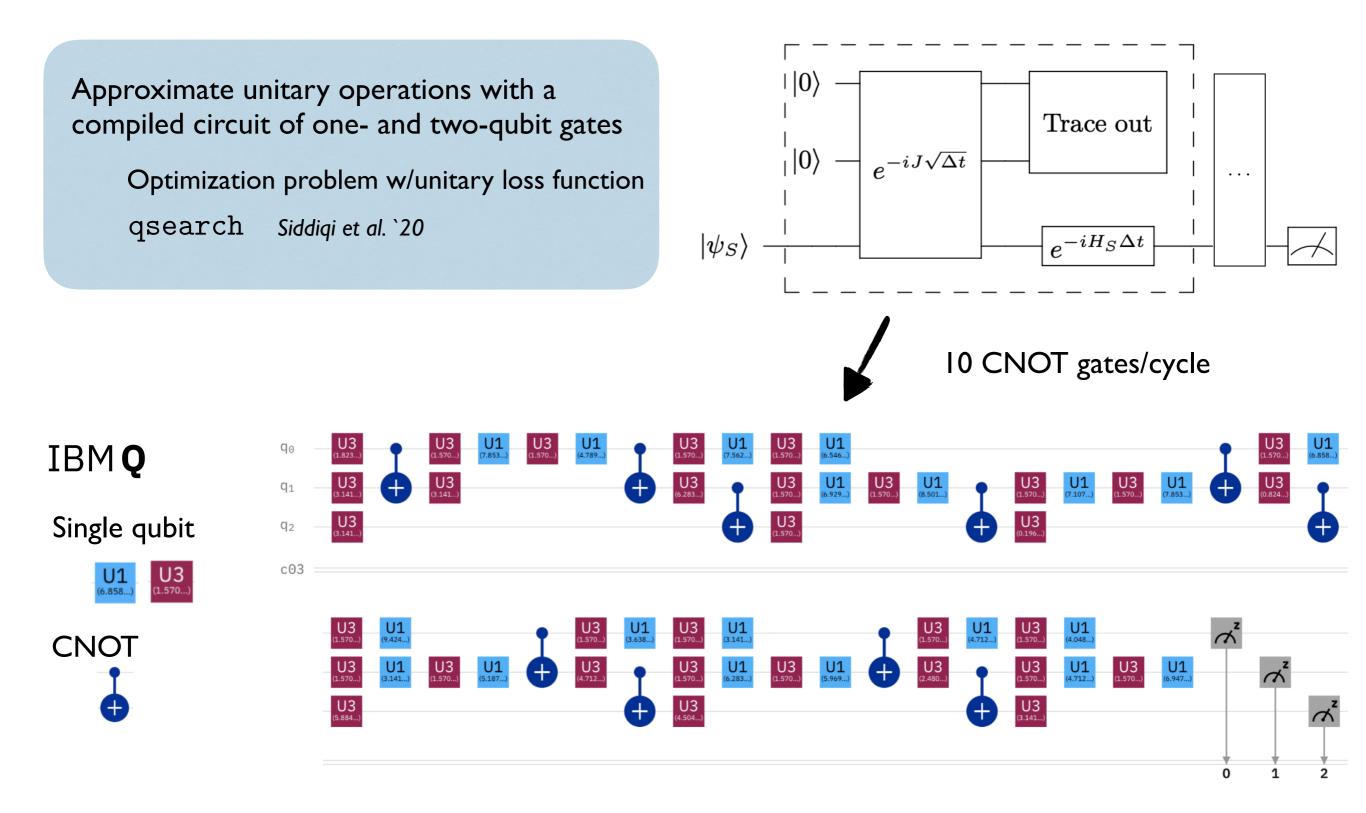
$$L_{j} \sim g(X \mp iY) \quad j = 0, 1$$
$$J = \begin{pmatrix} 0 & L_{0}^{\dagger} & L_{1}^{\dagger} & 0 \\ L_{0} & 0 & 0 & 0 \\ L_{1} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$



$$\rho(0) = \rho_S(0) \otimes \rho_E$$
$$\rho_E = \frac{e^{-\beta H_E}}{\operatorname{Tr}(e^{-\beta H_E})}$$



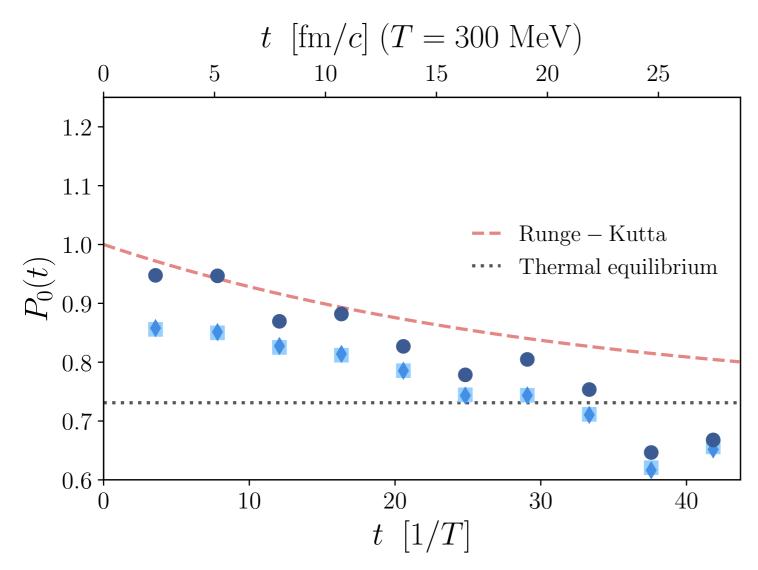
#### Quantum circuit synthesis

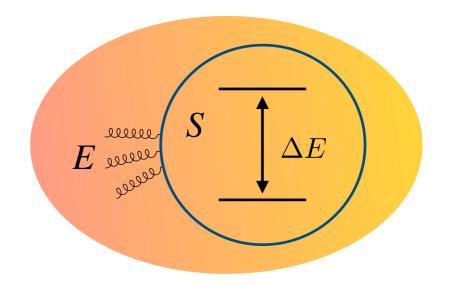


#### **Real-time evolution**

 $P_0(t)$  describes fraction that remains in "bound state"

Similar to *t*-dependent  $R_{AA} = \frac{\mathrm{d}\sigma_{AA}}{\langle N_{\mathrm{coll}} \rangle \,\mathrm{d}\sigma_{pp}}$ 



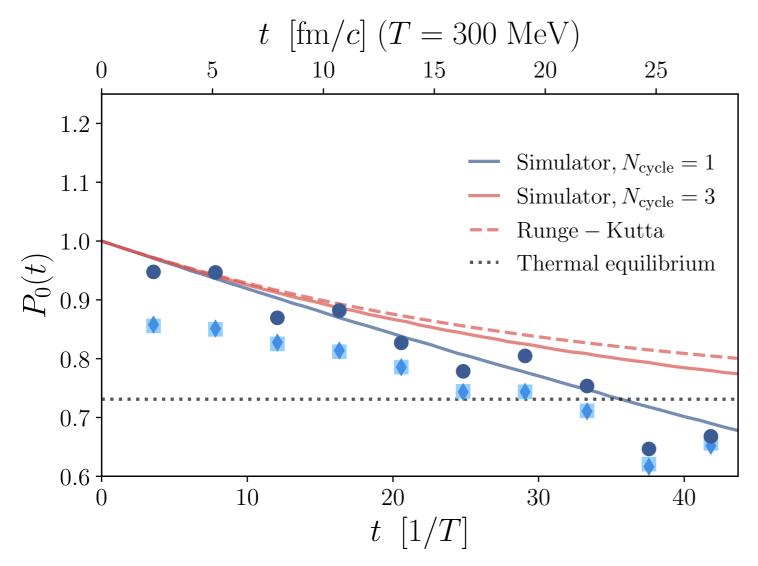


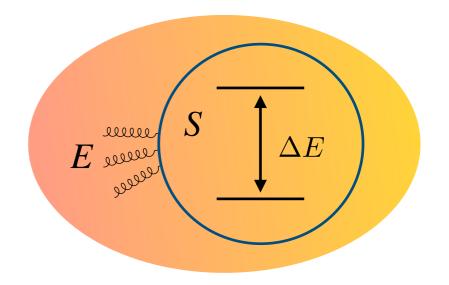
#### arXiv: 2010.03571

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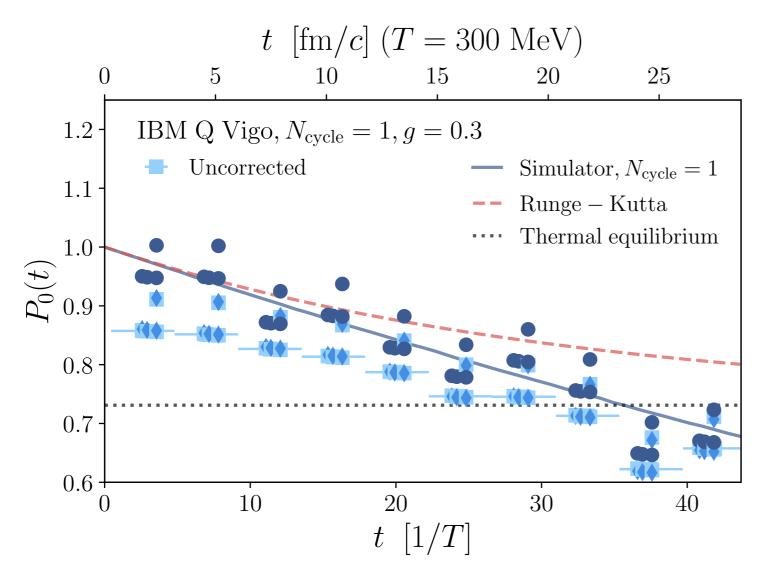
The algorithm converges to Lindblad evolution with a small number of cycles

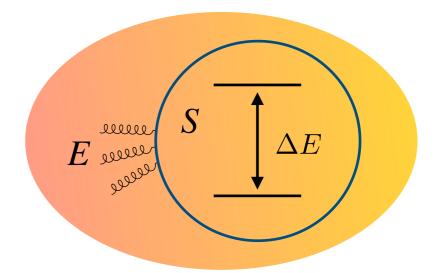
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#### **Real-time evolution**

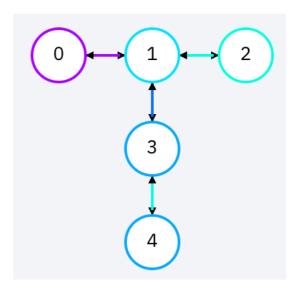
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ibmq\_vigo device



arXiv: 2010.03571

#### Error mitigation

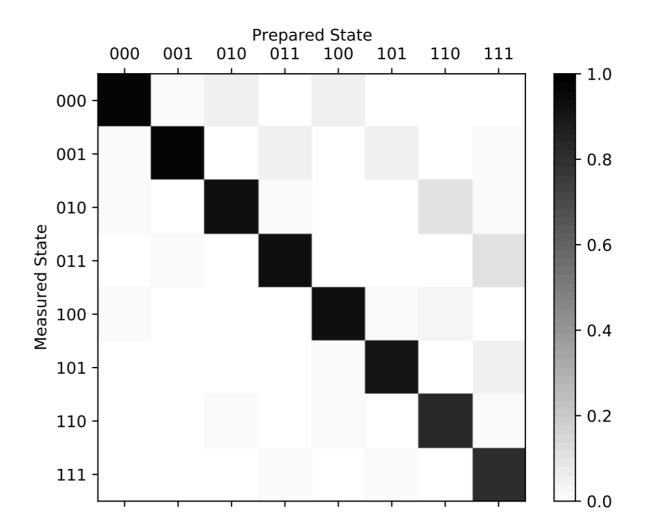
#### **Readout error**

Constrained matrix inversion IBM **Q** qiskit-ignis

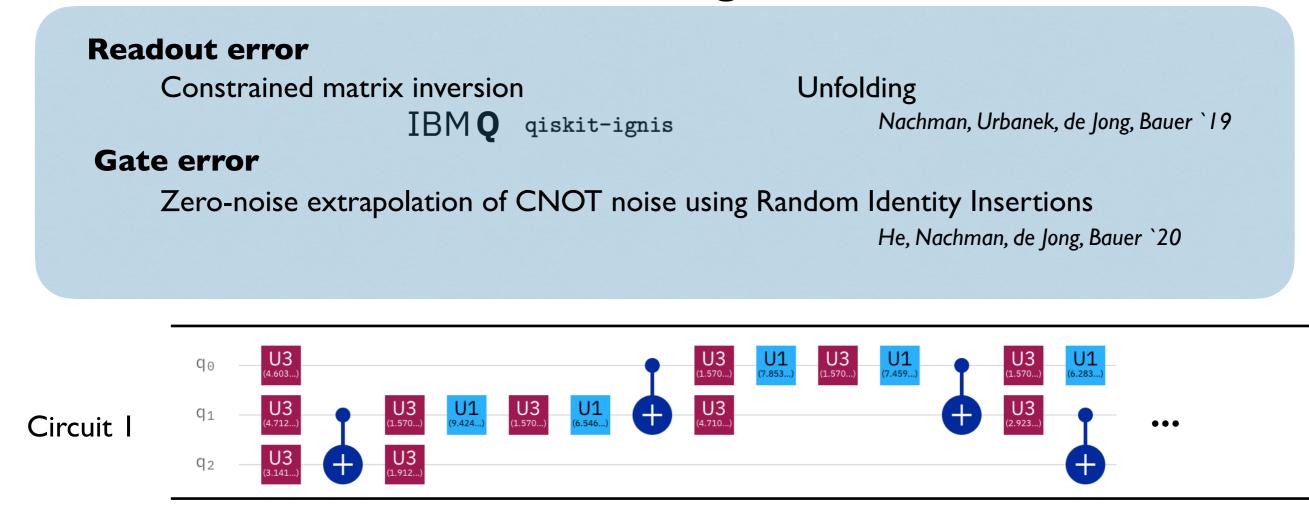
Unfolding Nachman, Urbanek, de Jong, Bauer `19

Prepare states by applying bitflip X gates and read out

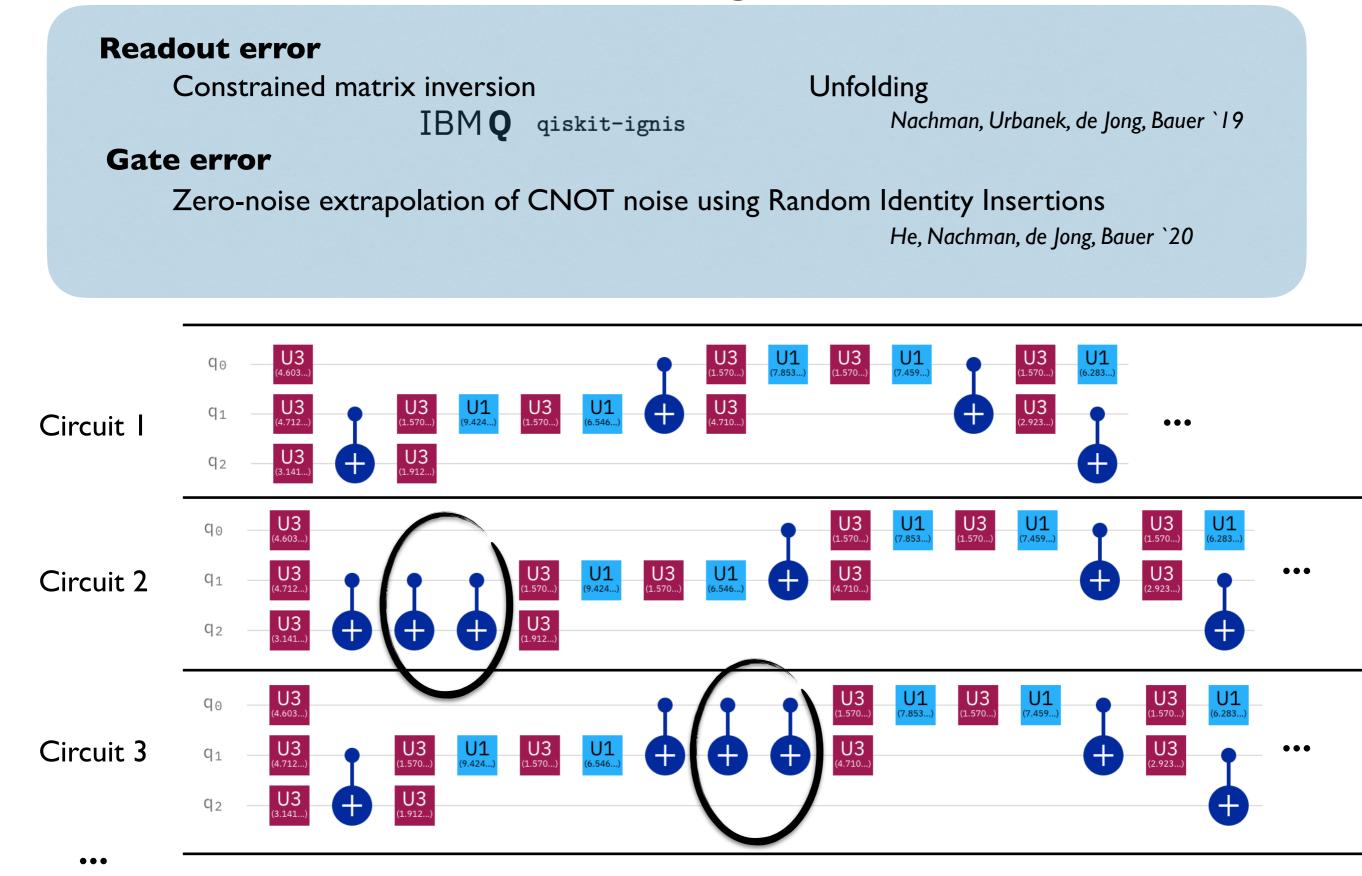
ibmq\_vigo device



#### Error mitigation



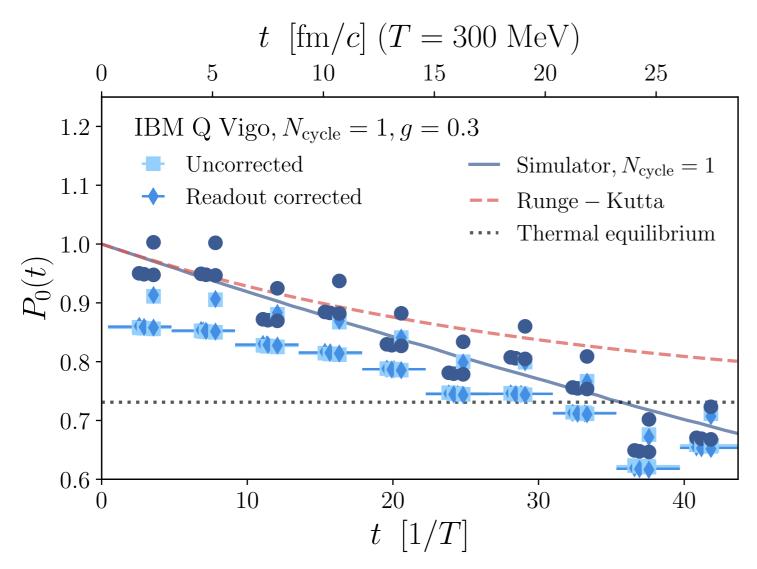
### Error mitigation

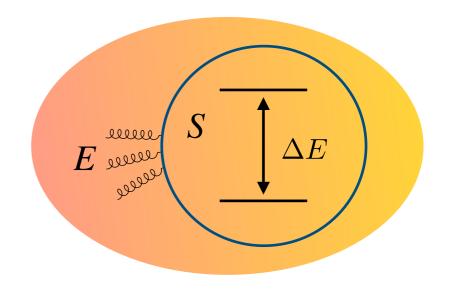


#### **Real-time evolution**

 $P_0(t)$  describes fraction that remains in "bound state"

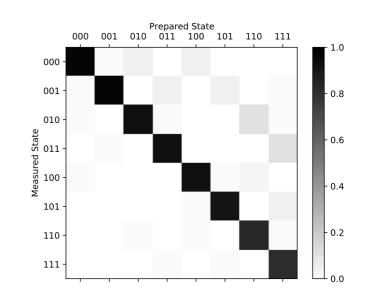
Similar to *t*-dependent  $R_{AA} = \frac{\mathrm{d}\sigma_{AA}}{\langle N_{\mathrm{coll}} \rangle \,\mathrm{d}\sigma_{pp}}$ 





ibmq\_vigo device

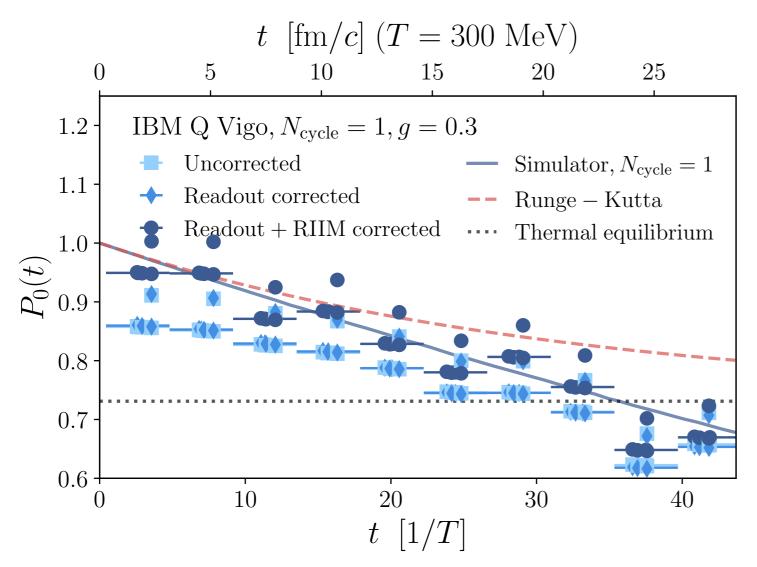
#### Readout correction small



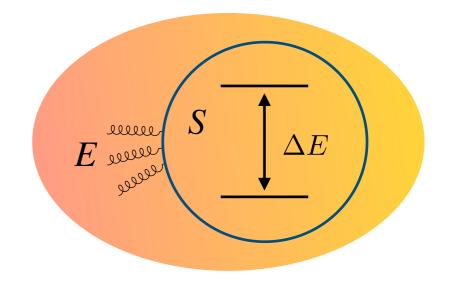
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arXiv: 2010.03571



ibmq\_vigo device

Readout correction small

CNOT gate error correction gives good agreement

Random Identity Insertion Method (RIIM)

Bauer, He, de Jong, Nachman `20

Proof of concept

## Outline

Open quantum systems in heavy-ion collisions

Quantum simulation with IBM Q

### Conclusions and outlook

• Open quantum system formalism describes the real-time evolution of hard probes in heavy-ion collisions

• Allows to go beyond semiclassical approximations in current models

## • Proof of concept that these systems can be simulated on current and near-term quantum computers (IBM Q)

- NISQ era digital quantum computing
- Recently developed error mitigation techniques

#### Future steps

- More efficient quantum algorithms & error mitigation
- Extension toward QCD