

Entanglement in pair creation

From string breaking to strong electric fields

Adrien Florio



Plan

“Real-time non-perturbative dynamics of jet production in Schwinger model:
quantum entanglement and vacuum modification”

with D. Frenklakh, K. Ikeda, D. Kharzeev, V. Korepin, S. Shi, K. Yu, arXiv:2301.11991

“Entropy Suppression through Quantum Interference in Electric Pulses”

with G. Dunne, D. Kharzeev, arXiv:2211.13347

Plan

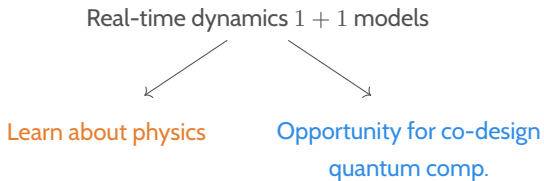
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Motivation



Schwinger model

Electromagnetism in 1 dimension

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Full fledged quantum field theory

Simulable in the near future (?)

Solved in some limit ($m \rightarrow 0$)

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$$\nabla E = \rho$$

Confines $V(x) \propto x$

Highly non-trivial vacuum

Schwinger model

Electromagnetism in 1 dimension

Full fledged quantum field theory

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Simulable in the near future (?)

Confines $V(x) \propto x$

Solved in some limit ($m \rightarrow 0$)

Highly non-trivial vacuum

Use-case/testbed \longleftrightarrow Learn new physics (dynamics)

Word of caution

Not QCD, only toy model (1D, no dynamical gluons)



Need to ask reasonable questions

Only qualitative predictions

Dynamical string breaking

Look at dynamical string breaking



$$E_n \sim m + m + \alpha l_1$$

Dynamical string breaking

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$$E_n \sim m + m + \alpha l_1$$



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$$E_n \sim m + m + \alpha l_3$$

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$$E_n \sim m + m + m + m$$

when $\alpha l_3 > 2m$

Dynamical string breaking

Look at dynamical string breaking



Screen field by creating particles!

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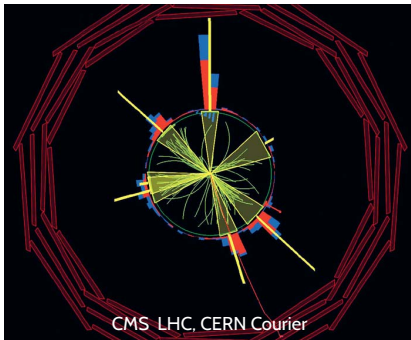


$$E_n \sim m + m + \alpha l_2$$



Screen field by creating particles!

Motivation: QCD jets



Our set-up

$$H(t) = \int dx \left[\frac{1}{2} \mathbf{E}^2 + \hat{\psi} (-i\gamma^1 \partial_1 + \mathbf{g} \mathbf{A}^1 \gamma_1 + m) \hat{\psi} + j_{ext}^1(t) \mathbf{A}_1 \right]$$

Our set-up

Electric field

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Vector potential

$$H(t) = \int dx \left[\frac{1}{2} \mathbf{E}^2 + \hat{\psi} (-i\gamma^1 \partial_1 + g\mathbf{A}^1 \gamma_1 + m) \hat{\psi} + j_{ext}^1(t) \mathbf{A}_1 \right]$$

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Diagram illustrating the components of the Hamiltonian $H(t)$:

- Electric field** (red text) points to $\frac{1}{2} \mathbf{E}^2$.
- Vector potential** (blue text) points to $\mathbf{g} \mathbf{A}^1 \gamma_1$ and $j_{ext}^1(t) \mathbf{A}_1$.
- Fermion** (grey text) points to the fermion bilinear $\hat{\psi} \left(-i\gamma^1 \partial_1 + m \right) \hat{\psi}$.

Our set-up

$$H(t) = \int dx \left[\frac{1}{2} \mathbf{E}^2 + \hat{\psi} \left(-i\gamma^1 \partial_1 + gA^1 \gamma_1 + m \right) \hat{\psi} + \hat{j}_{ext}^1(t) A_1 \right]$$

Electric field

Vector potential

Fermion

External charges: $j_{ext}^1(t) = g (\delta(x+t) + \delta(x-t)) \theta(t)$

2 point charges moving apart at speed of light

Our set-up

Electric field

Vector potential

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- Idea:**
- Find $|\text{vac}\rangle_{t<0}$
 - Compute $|\psi(t)\rangle = e^{-i \int_0^t dt' H(t')} |\text{vac}\rangle_{t<0}$

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see also [74, Casher, Kogut, Susskind], [12,13, Kharzeev, Loshaj], [14, Berges, Hebenstreit]

In practice

- Staggered fermions χ_n
- Integrate out E : $\partial_1 E = \rho + \rho_{ext}$

- (Map to non-local spin chain)

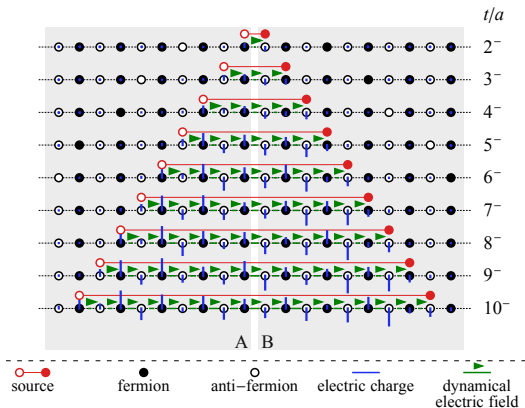
$$H(t) = H_{\pm} + H_{ZZ} + H_Z(t)$$

$$H_{\pm} = \frac{1}{4a} \sum_{n=1}^{N-1} (X_n X_{n+1} + Y_n Y_{n+1})$$

$$H_{ZZ} = \frac{ag^2}{4} \sum_{n=1}^{N-1} \sum_{m=1}^n \sum_{k=1}^{m-1} Z_m Z_k, \quad H_Z = \sum_{n=1}^N f(n) Z_n$$

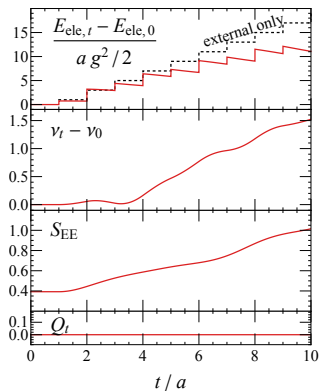
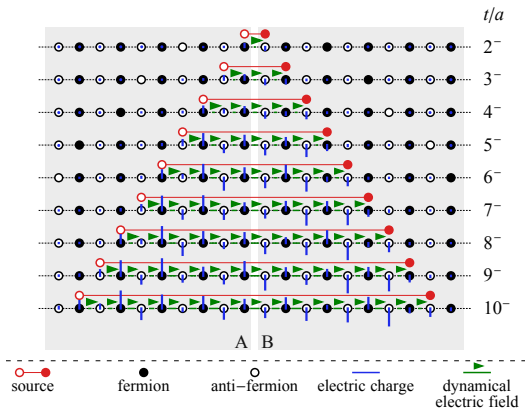
- Use exact diagonalization
- Compute observables $\bar{O}(t) = \langle \psi(t) | O | \psi(t) \rangle$

Results, $m = 0.25, g = 0.5, a = 1$



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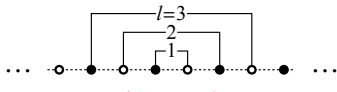
$\nu : \langle \bar{\psi} \psi \rangle, \quad S_{EE} : \text{entanglement entropy A/B}$



Is entanglement manifest in correlations \leftrightarrow measurable?

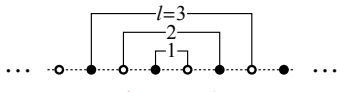
Correlation

1) Look at $\langle \Delta\nu_{N/2+l+1}(t) \Delta\nu_{N/2-l}(t) \rangle$, $\Delta\nu_n = \bar{\psi}\psi|_n(t) - \bar{\nu}$

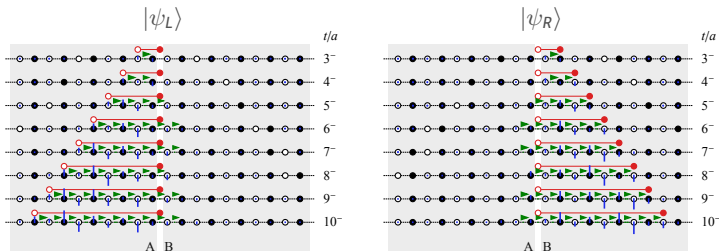


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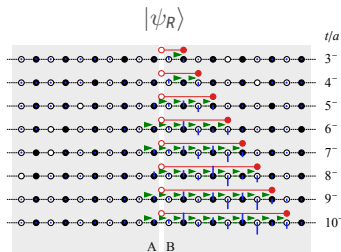
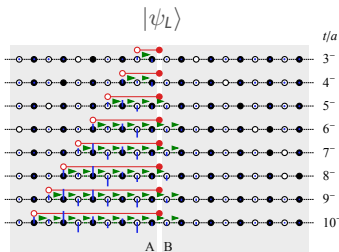


2) Compare to uncorrelated reference case



Correlation

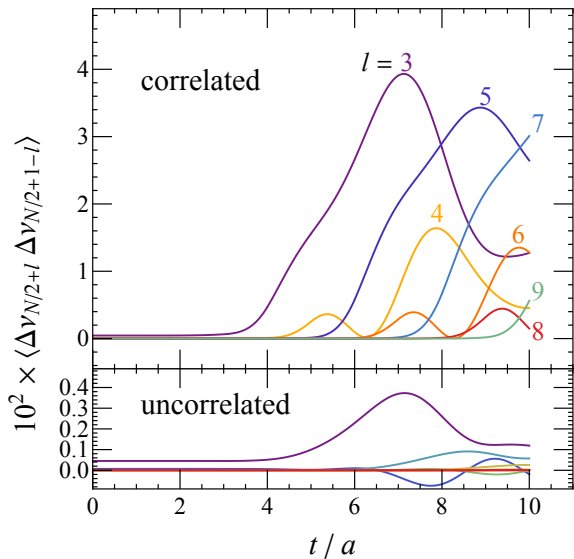
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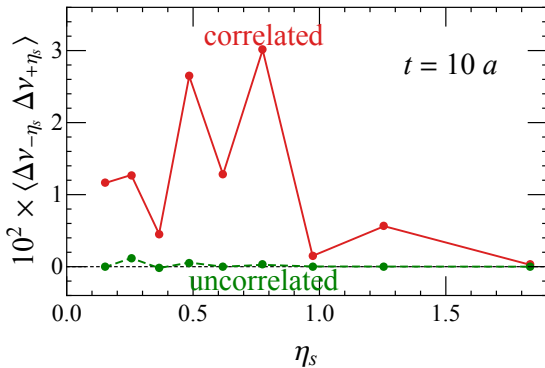
$$|\psi_{ref}\rangle = |\psi_L\rangle + e^{i\phi} |\psi_R\rangle$$

Random uniform phase

$$\langle\langle \psi_{ref} | O | \psi_{ref} \rangle\rangle \equiv \int \langle \psi_{ref} | O | \psi_{ref} \rangle \frac{d\phi}{2\pi} = \frac{\langle \psi_L | O | \psi_L \rangle}{2} + \frac{\langle \psi_R | O | \psi_R \rangle}{2}$$



For exp. \rightarrow spatial rapidity $\eta_s \equiv \operatorname{arctanh} \frac{x}{t}$



Next steps

Finite temperature

Thermalization/ETH

Tensor networks

Summary

- Schwinger model can still teach us some physics
- Direct observation of quantum properties of string breaking
- Suggests enhanced correlations at low/mid rapidities in jet production

Plan

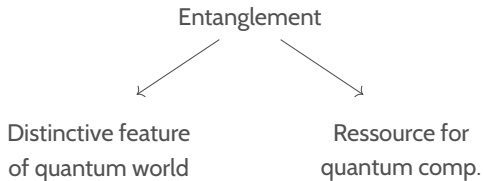
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Motivation



Question: Effect of quantum interferences on entanglement?

Set-up

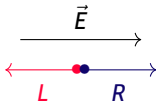
Pair creation in strong electric fields (Schwinger effect)



- Creation of entangled pair of particles
- Typical 2-levels system

Mechanism: Similar to string breaking.

$E \cdot L \approx 2m \rightarrow$ energetically favorable to create particles



Entanglement in momentum space

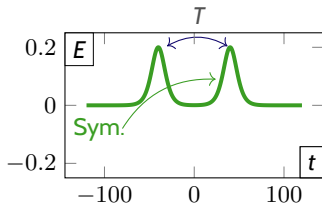
n_k : Probability to create particle with momentum k

Left/right entanglement [AF, Kharzeev, 2021] :

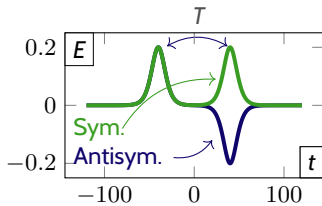
$$S = - \int \frac{dk}{2\pi} [(1 - n_k) \log (1 - n_k) + n_k \log (n_k)]$$
$$\mathcal{N} = - \int \frac{dk}{2\pi} n_k$$

Gibbs entropy!

Interferences

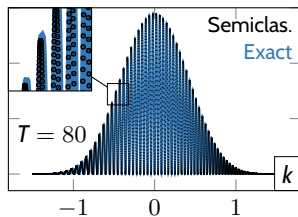
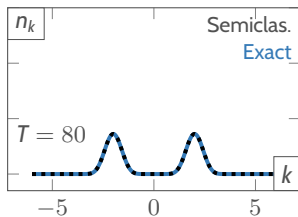
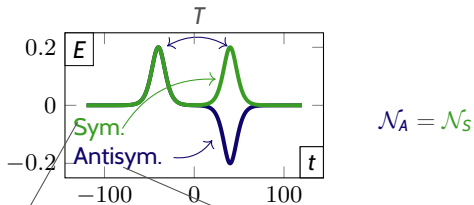


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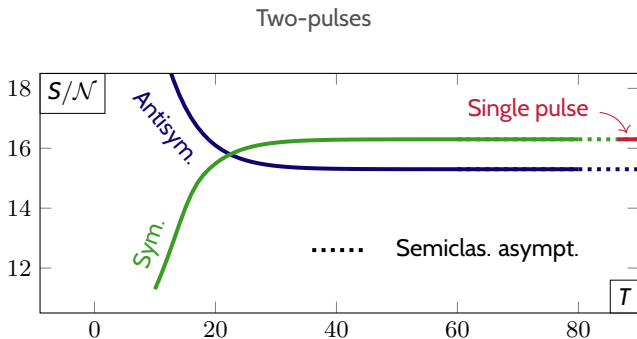


$$\mathcal{N}_A = \mathcal{N}_S$$

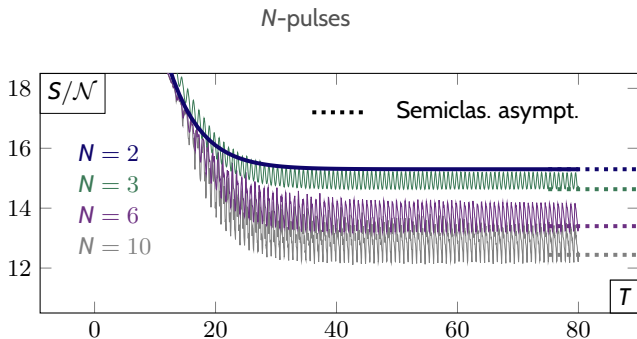
Interferences



Entanglement suppression



Entanglement suppression



Summary # 2

- Interference effects can suppress entropy production
- Potential applications to sensing/hardware?

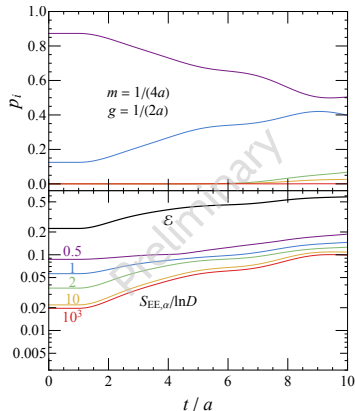
Thank you!

Trailer #1: entanglement spectrum

Entanglement spectrum: $\{p_i\}$, e-values of ρ_A

$$S_{\text{Rényi},\alpha} \equiv \frac{\ln \text{tr}(\rho_A^\alpha)}{1-\alpha}$$

$$\mathcal{E} \equiv \frac{1 - \text{tr} \rho_A^2}{1 - 1/D} = \frac{1 - \sum_{i=1}^D p_i^2}{1 - 1/D}.$$

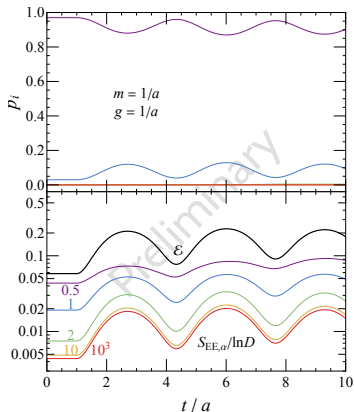
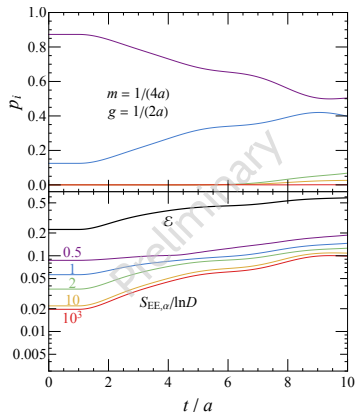


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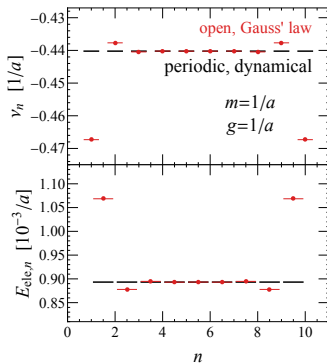
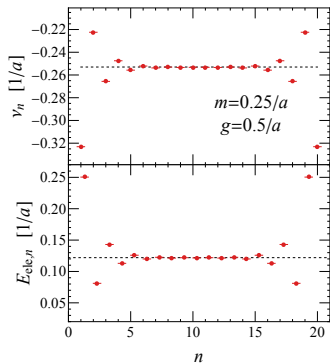
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Boundary effects



Trailer #2: TN

