

The three-body problem

looking back at the last 9yrs



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🌐 <http://bit.ly/rbricenoPhD>

🐦 @RaulBriceno12



9+ years

☑ Summer of 2006: W&M / JLab REU

realized that I was useless at the lab 😅

☑ 2012: first three-body paper with Zohreh Davoudi

☑ Fall of 2013: JLab postdoc

"friendliest group!"

☑ Fall of 2015: ODU postdoc

☑ Fall of 2016: Isgur fellow

☑ Summer of 2017: Joint ODU / JLab staff

0 COVID cases reported! 😅

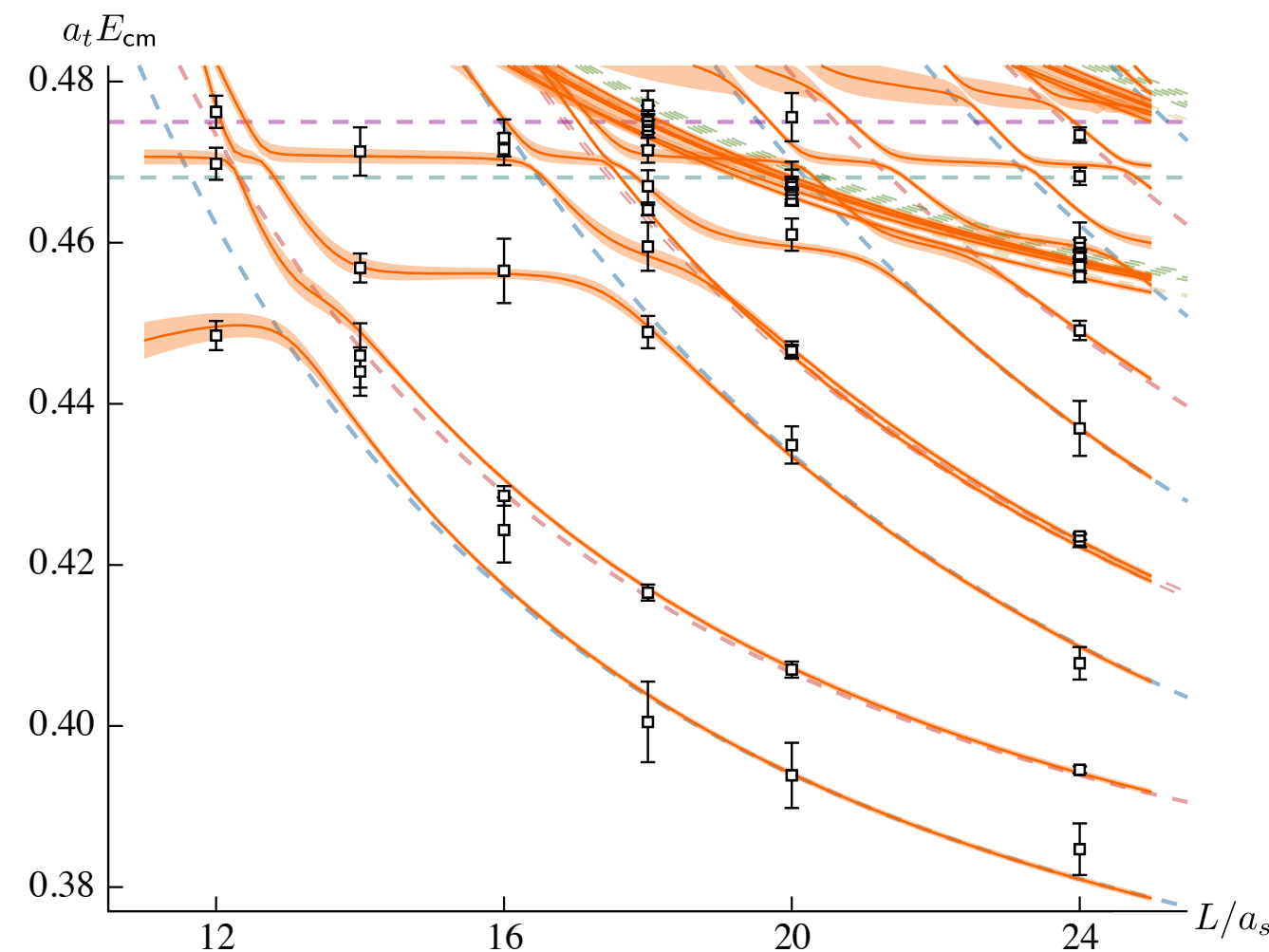


**Raul
Briceno
USER**

Modern-day spectroscopy rooted in QCD

$$\sum_f \bar{\psi}_f (i\gamma^\mu D_\mu - m_f) \psi_f - \frac{1}{4} \text{Tr}[G_{\mu\nu} G^{\mu\nu}]$$

+ electroweak sector

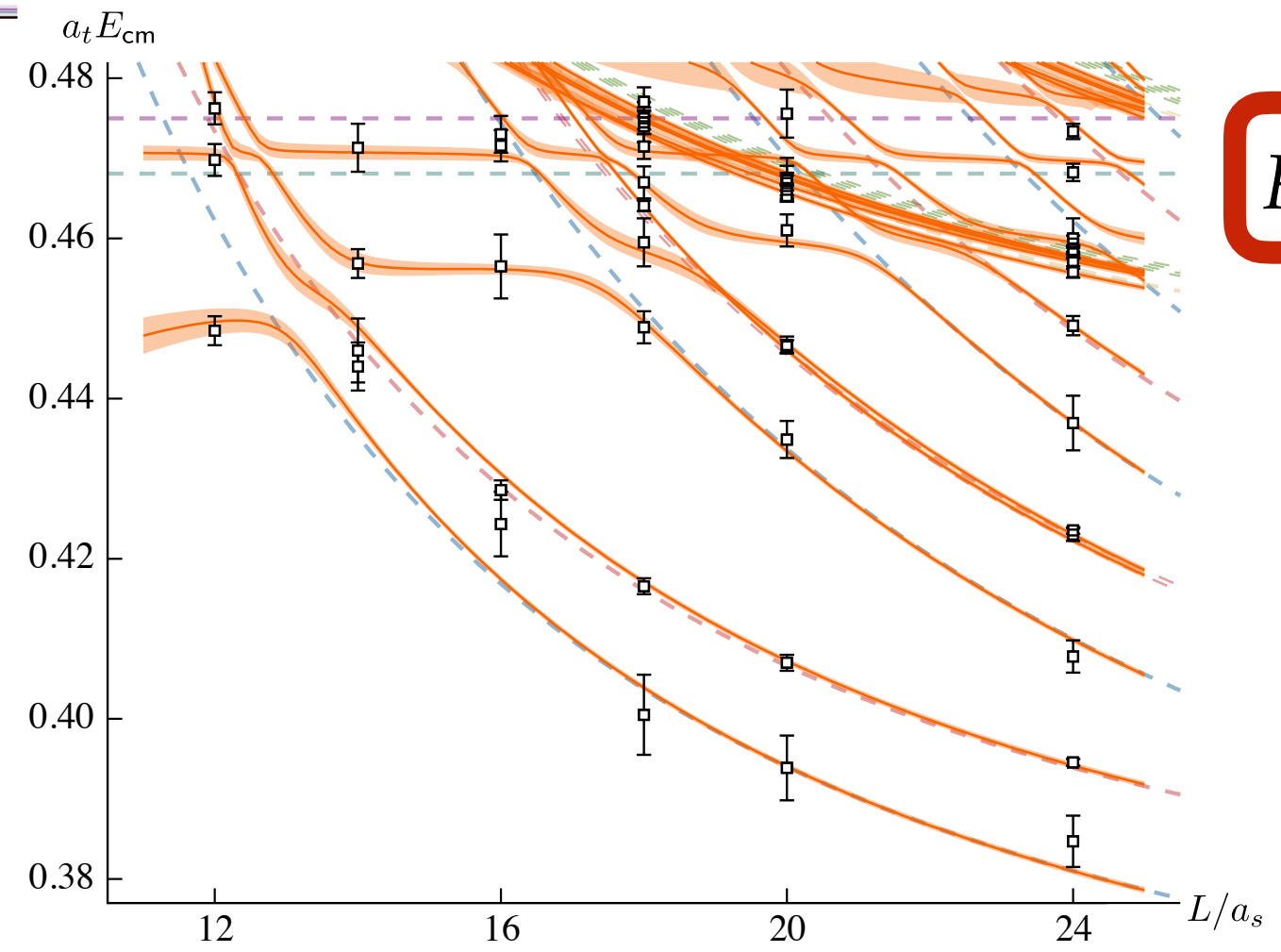
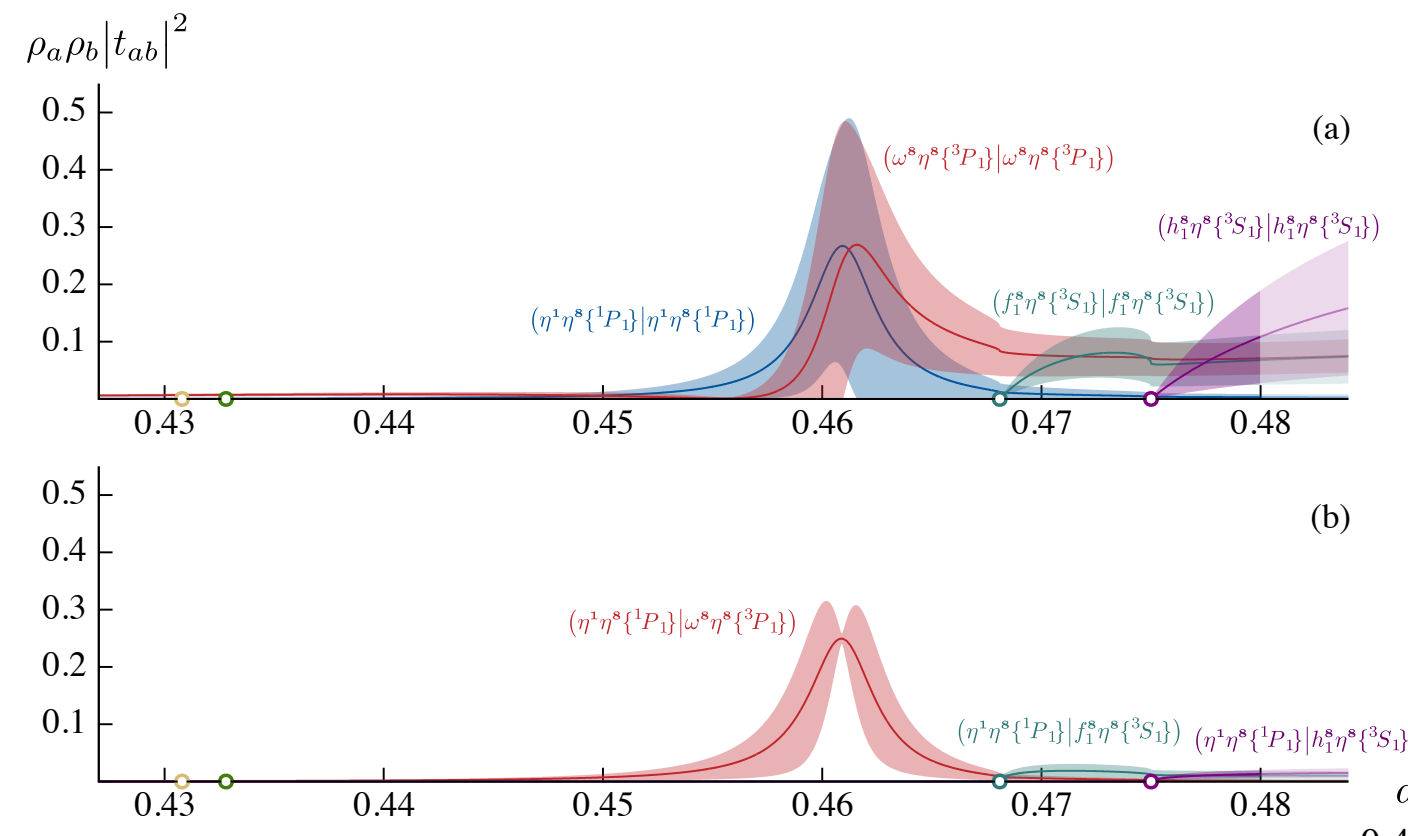


Woss, Dudek,
Edwards, Thomas, Wilson (2020)

QCD Lagrangian &
electroweak probes,
Lattice QCD



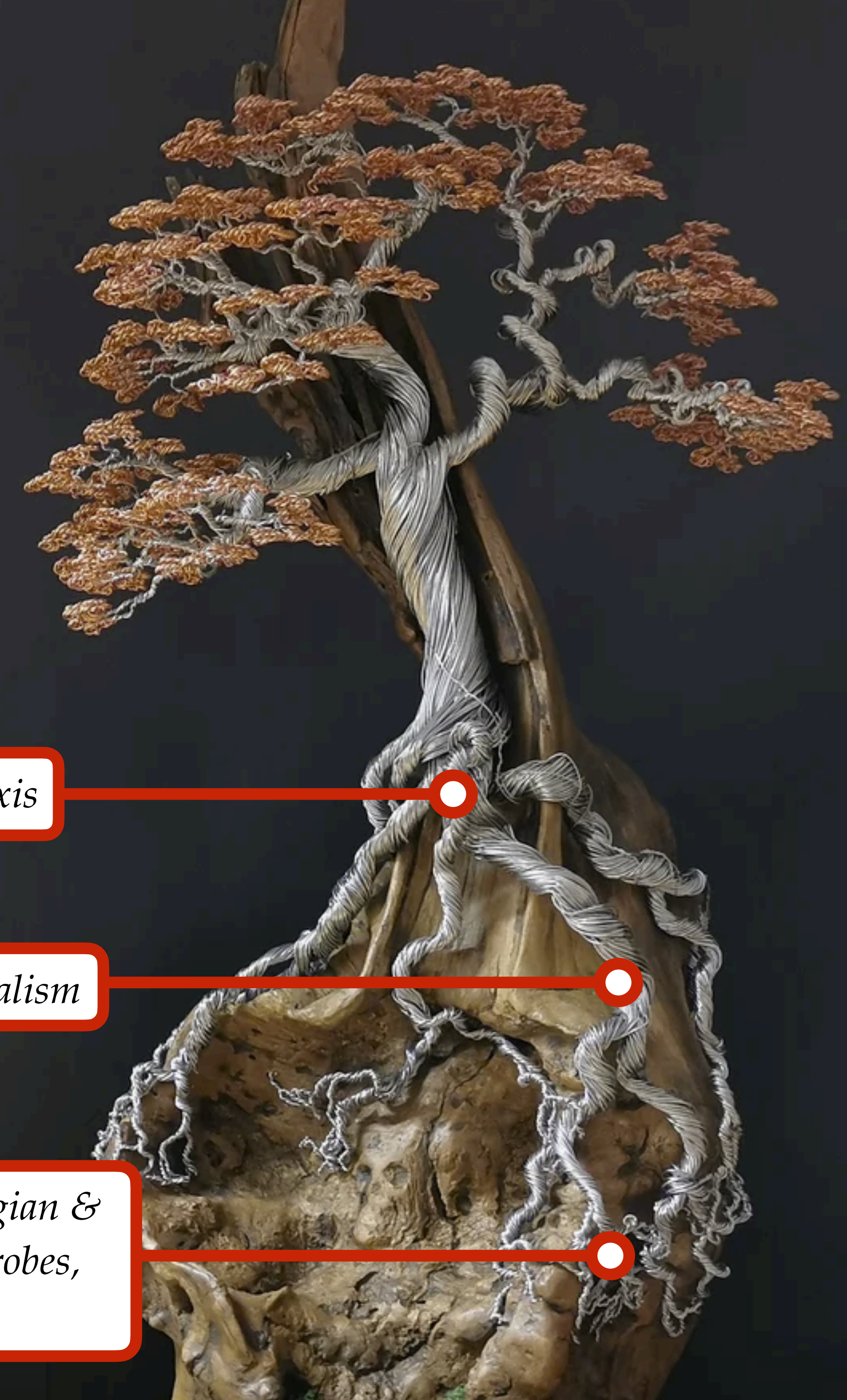
Modern-day spectroscopy rooted in QCD



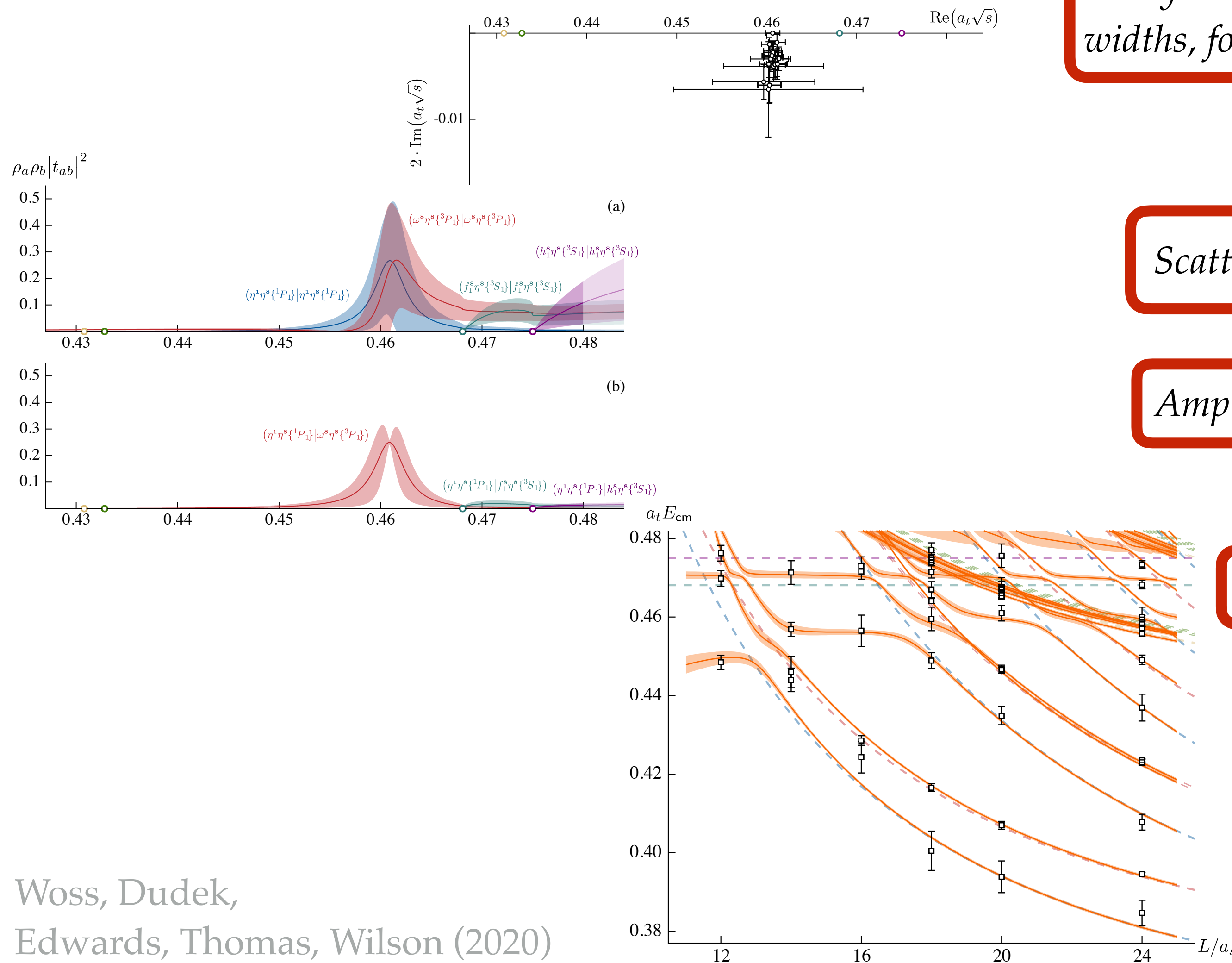
Amplitudes on the real axis

Finite-volume formalism

QCD Lagrangian & electroweak probes, Lattice QCD



Modern-day spectroscopy rooted in QCD



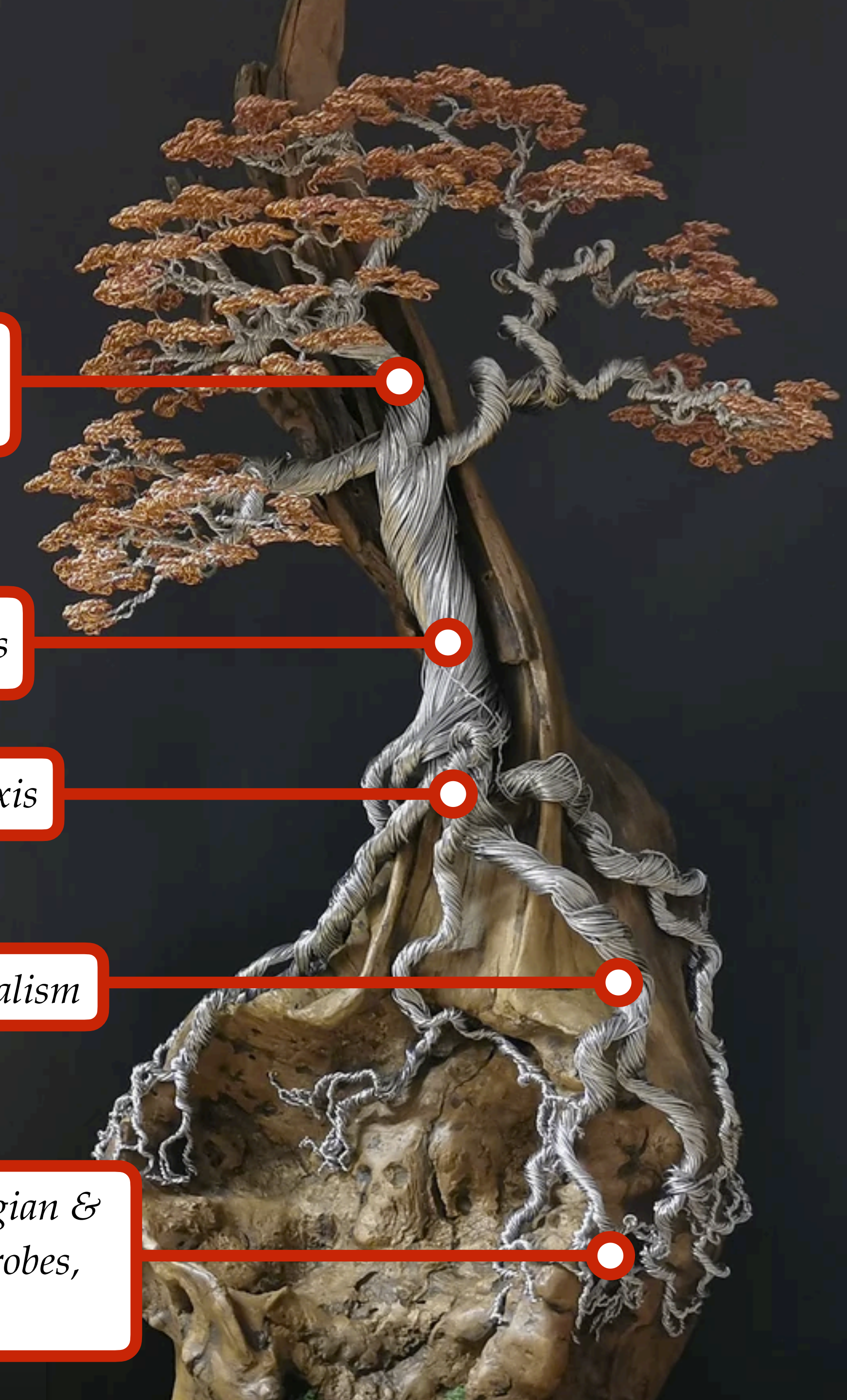
Analytic continuation: poles, widths, form factors,...

Scattering theory & EFTs

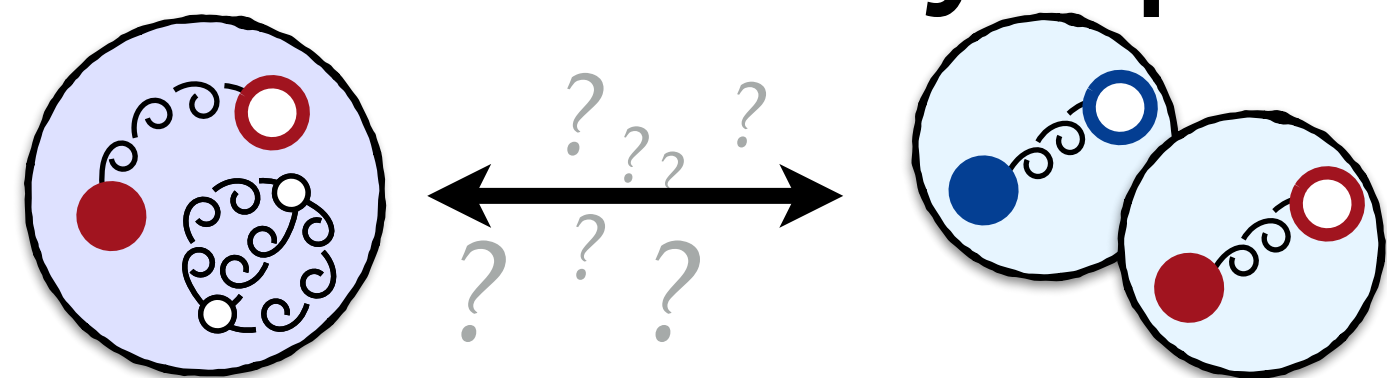
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Modern-day spectroscopy rooted in QCD



Nature of states, patterns in the spectrum, ...

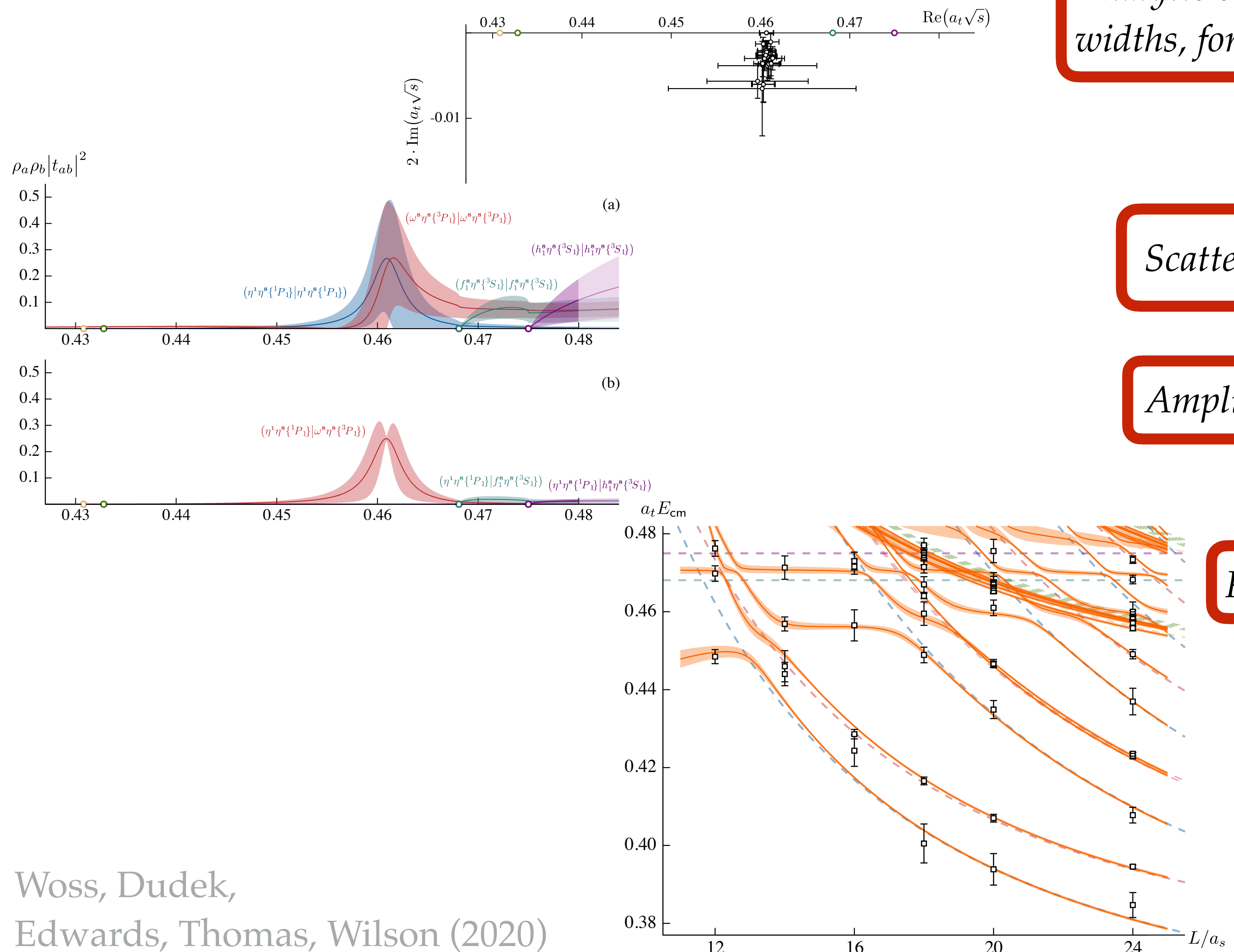
Analytic continuation: poles, widths, form factors, ...

Scattering theory & EFTs

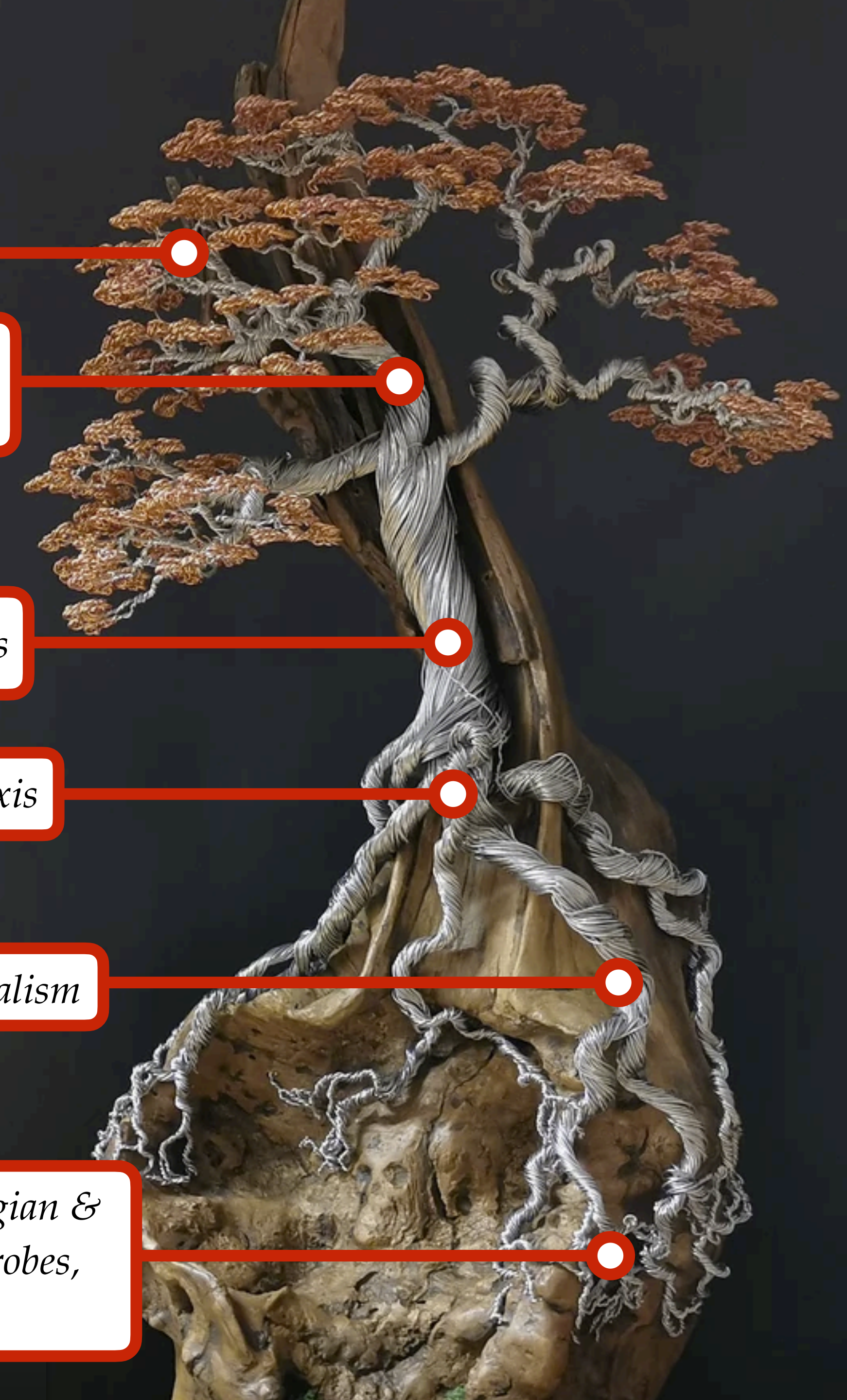
Amplitudes on the real axis

Finite-volume formalism

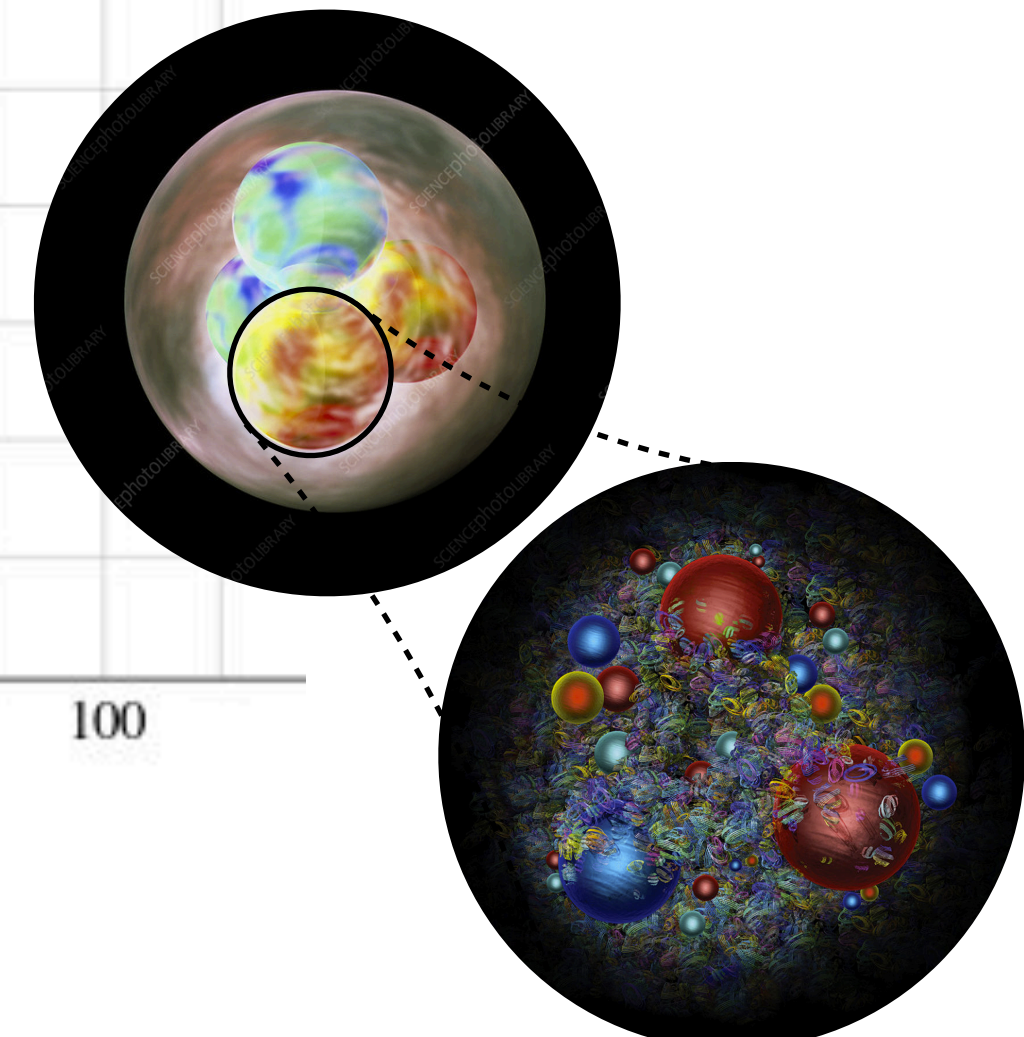
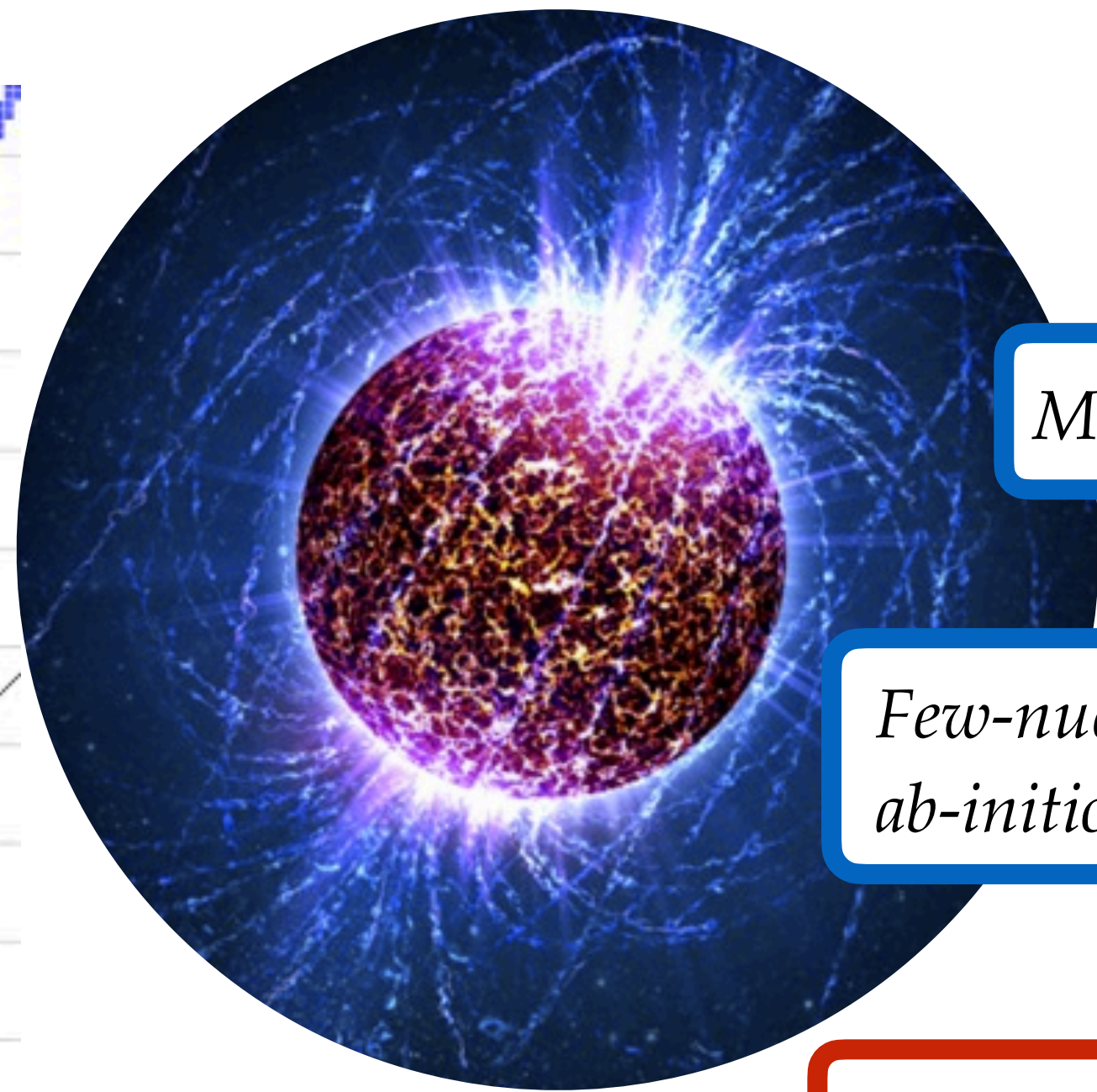
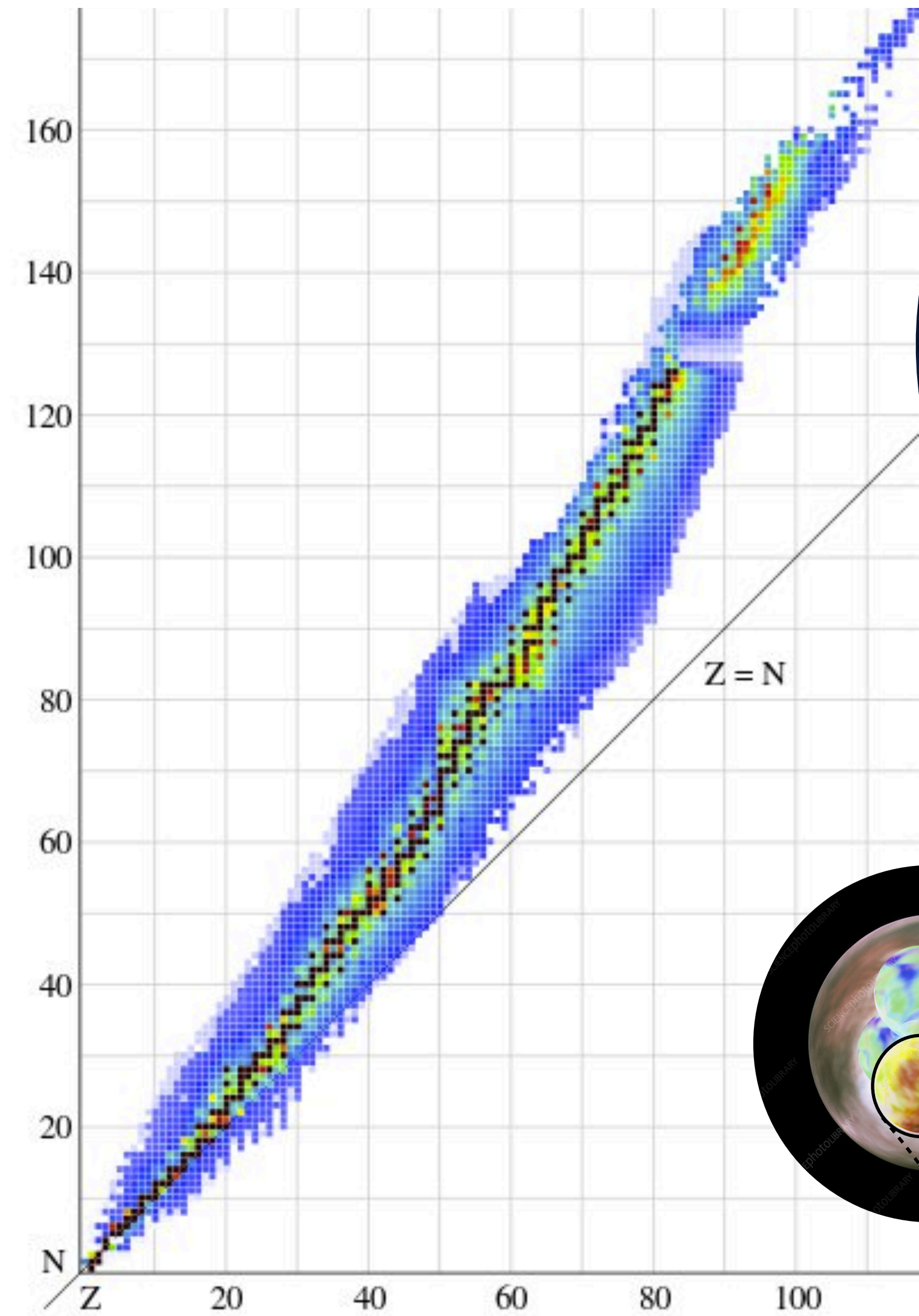
QCD Lagrangian & electroweak probes, Lattice QCD



Woss, Dudek, Edwards, Thomas, Wilson (2020)



Modern-day nuclear structure



Nuclear astro physics

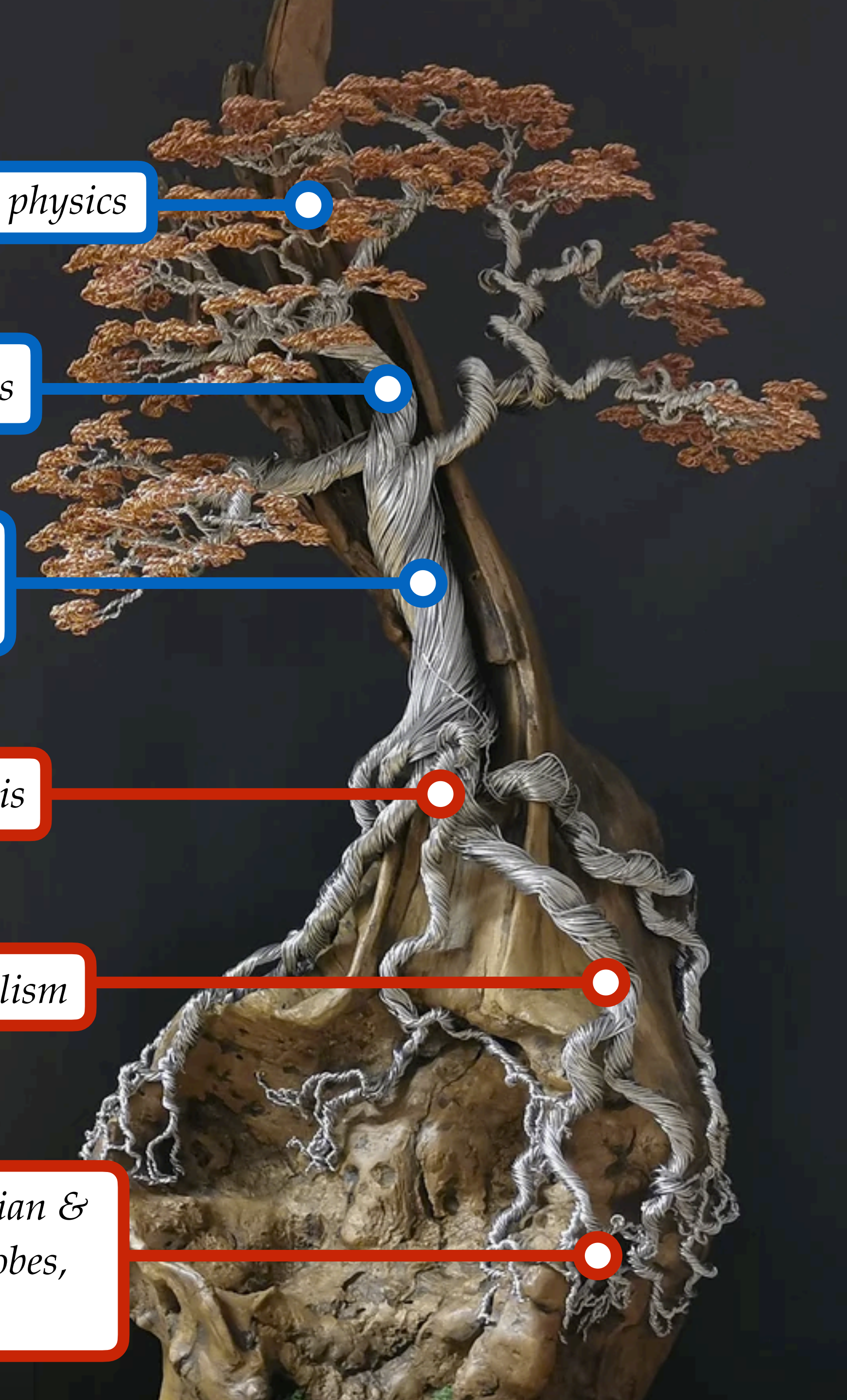
Many-body systems

*Few-nucleon systems:
ab-initio techniques*

Amplitudes on the real axis

Finite-volume formalism

*QCD Lagrangian &
electroweak probes,
Lattice QCD*



The common thread

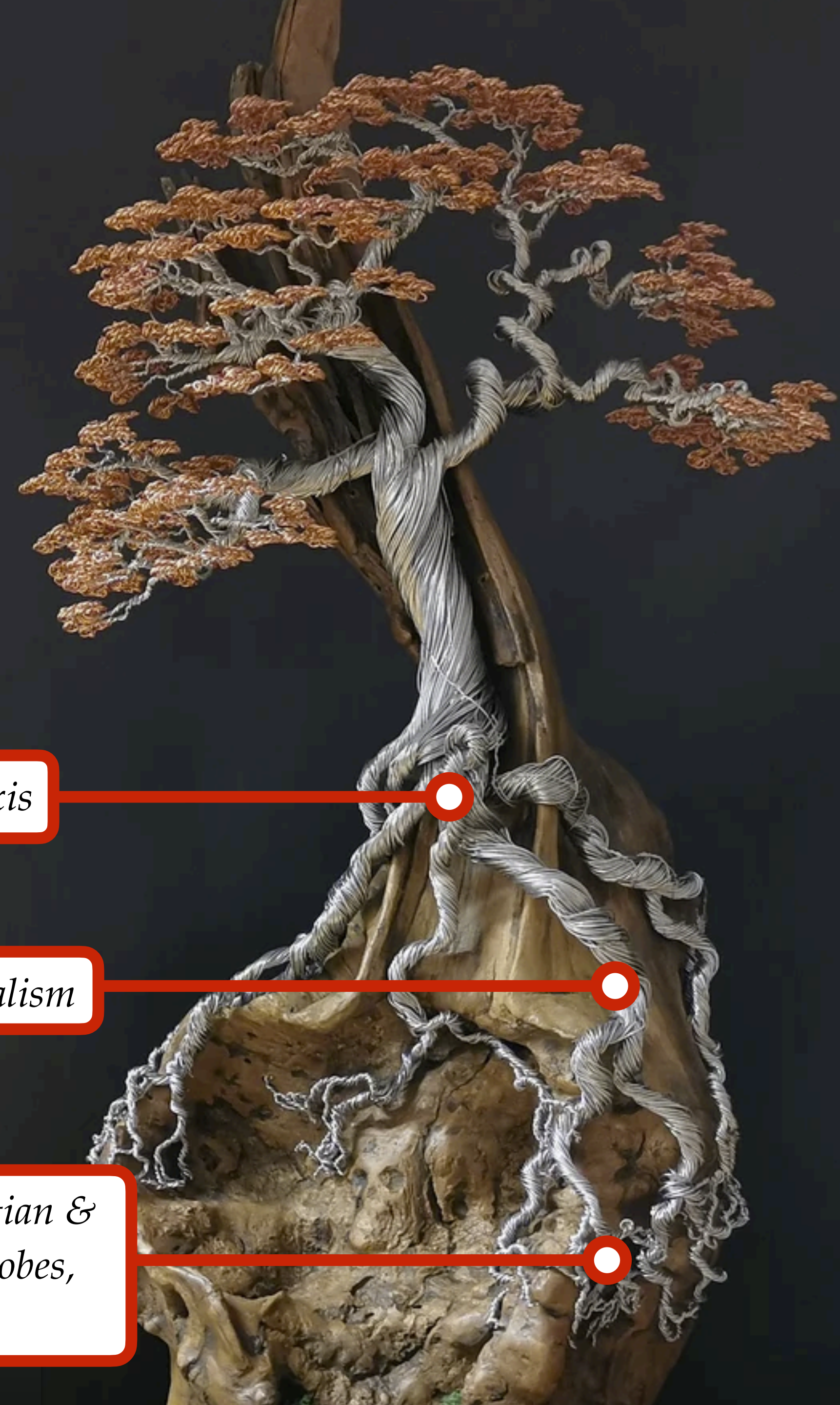
This is at the root of several areas of research:

- hadron spectroscopy,
- nuclear structure,
- hadron structure,
- electroweak precision decays,
- ...

Amplitudes on the real axis

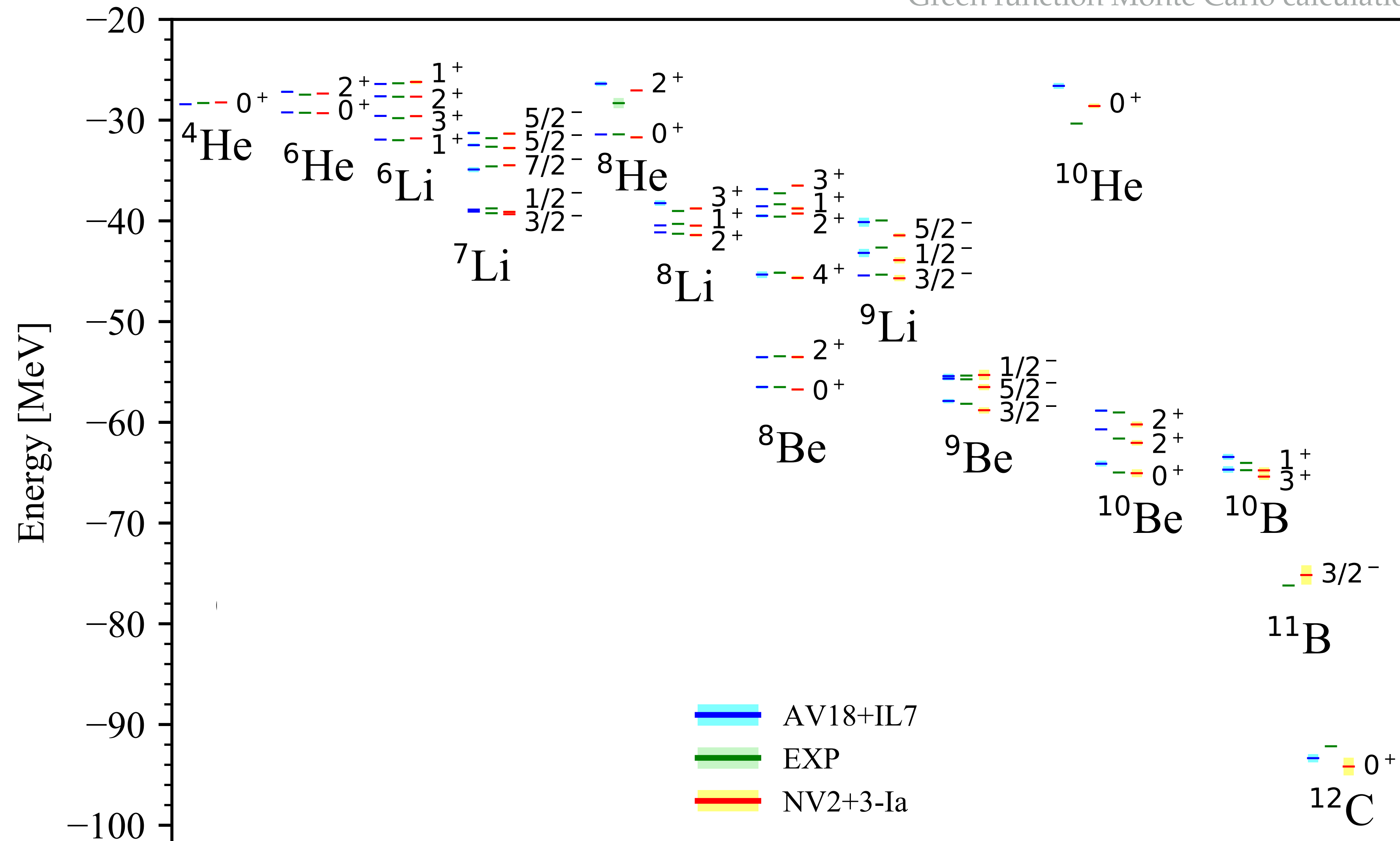
Finite-volume formalism

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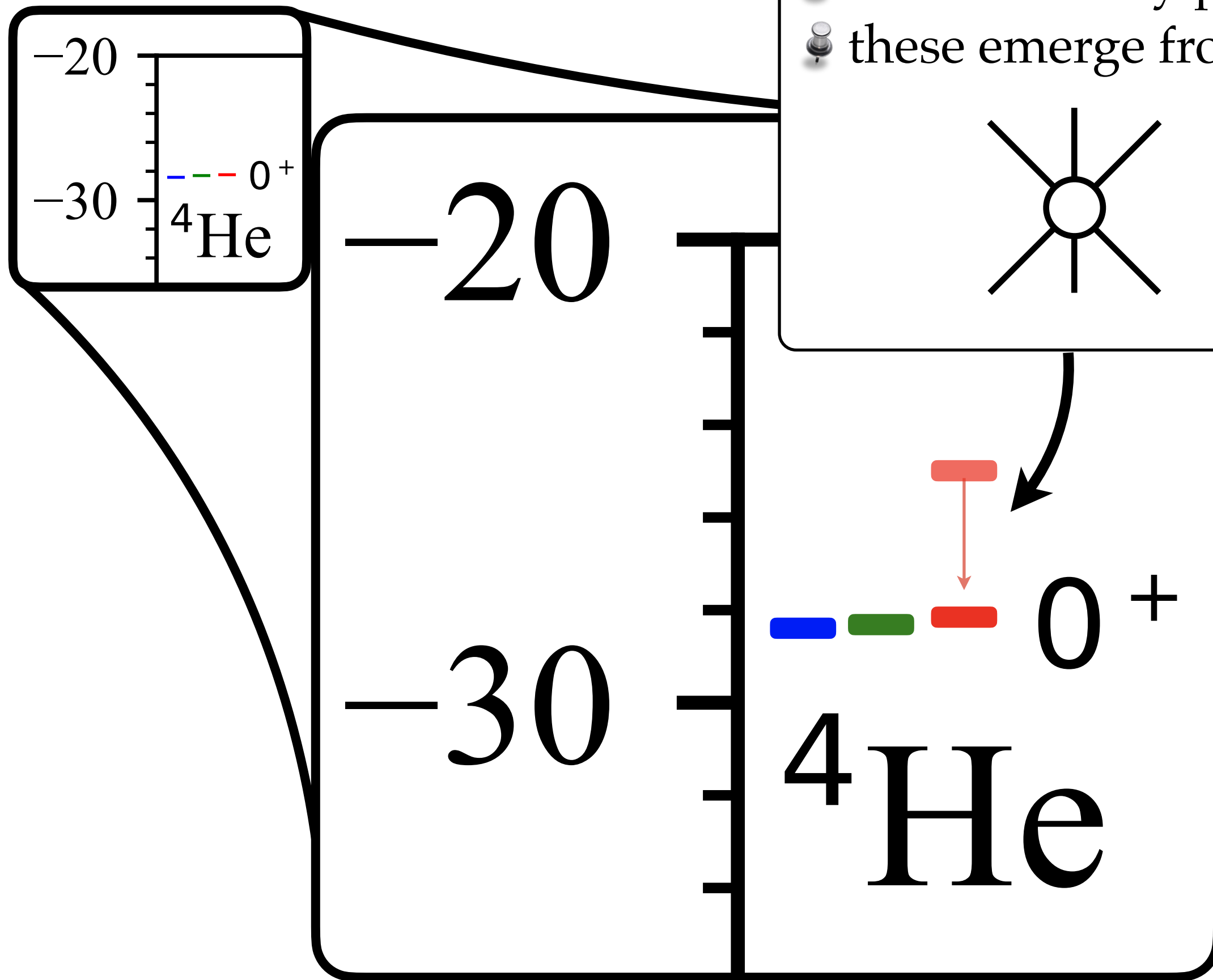


the three-body problem in nuclear physics

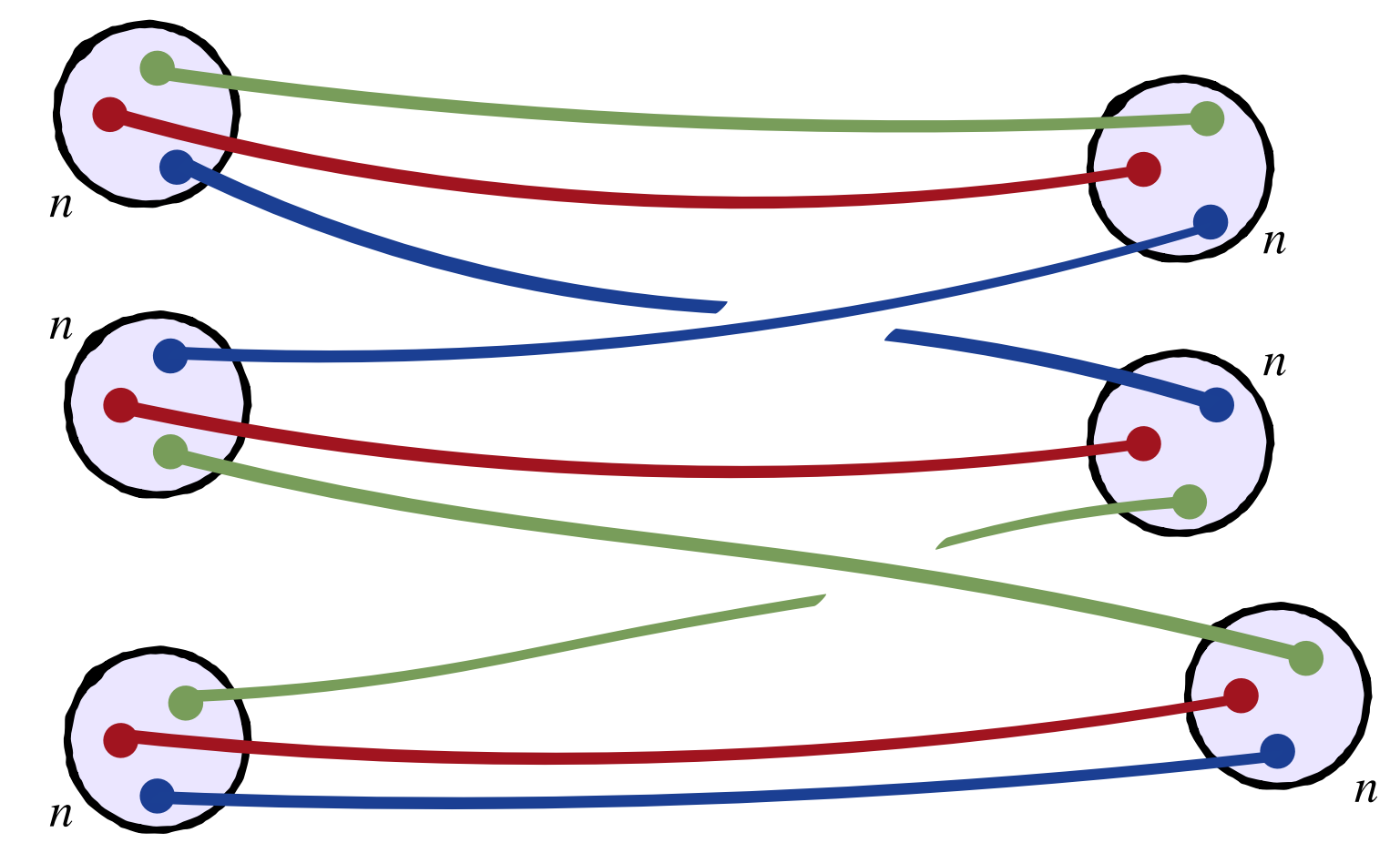
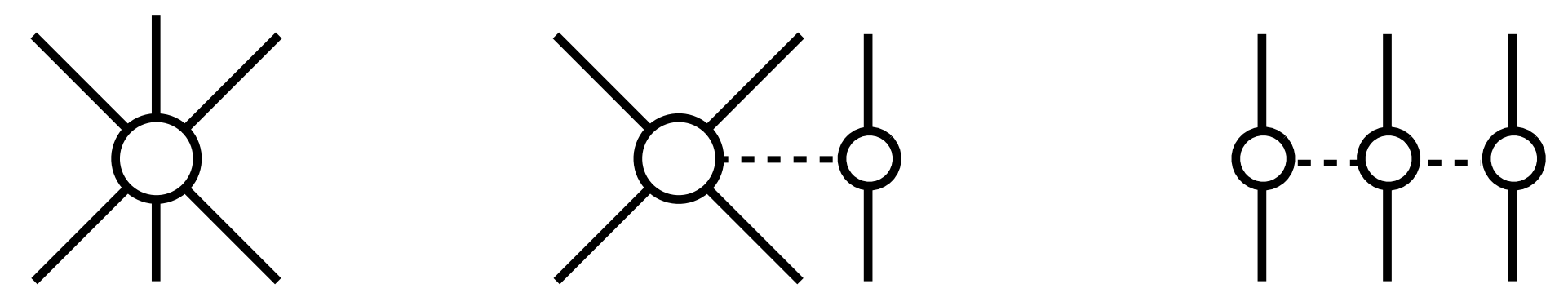
Green function Monte Carlo calculations



the three-body problem



- the three-body potentials shifts spectra by about 10%-20%
- these emerge from "local" & "non-local" interactions



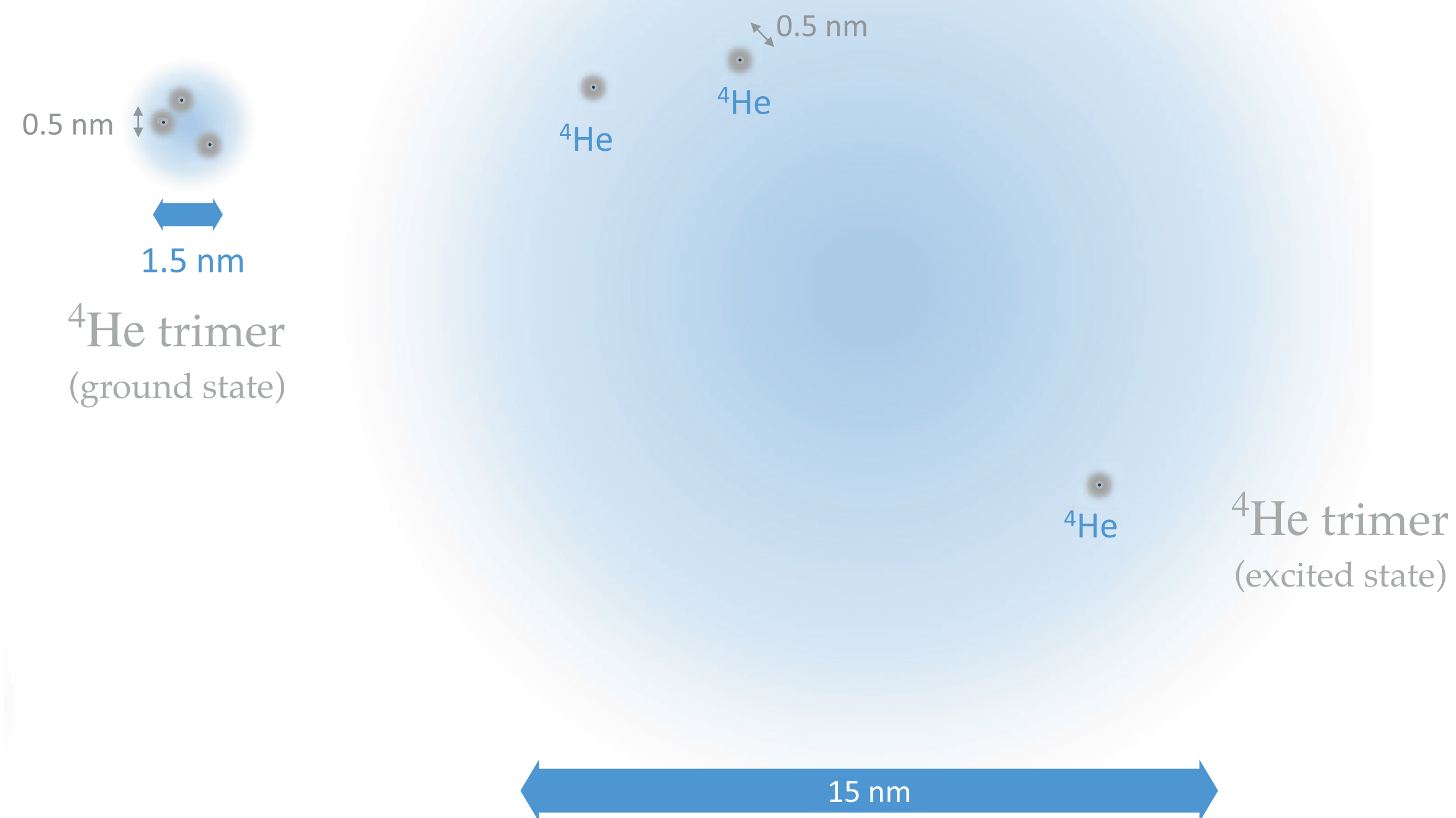
$3n$ scattering

the three-body problem *Efimov physics*

Unitary limit: $p \cot \delta = -\frac{1}{a} + \frac{rp^2}{2} + \dots = 0$

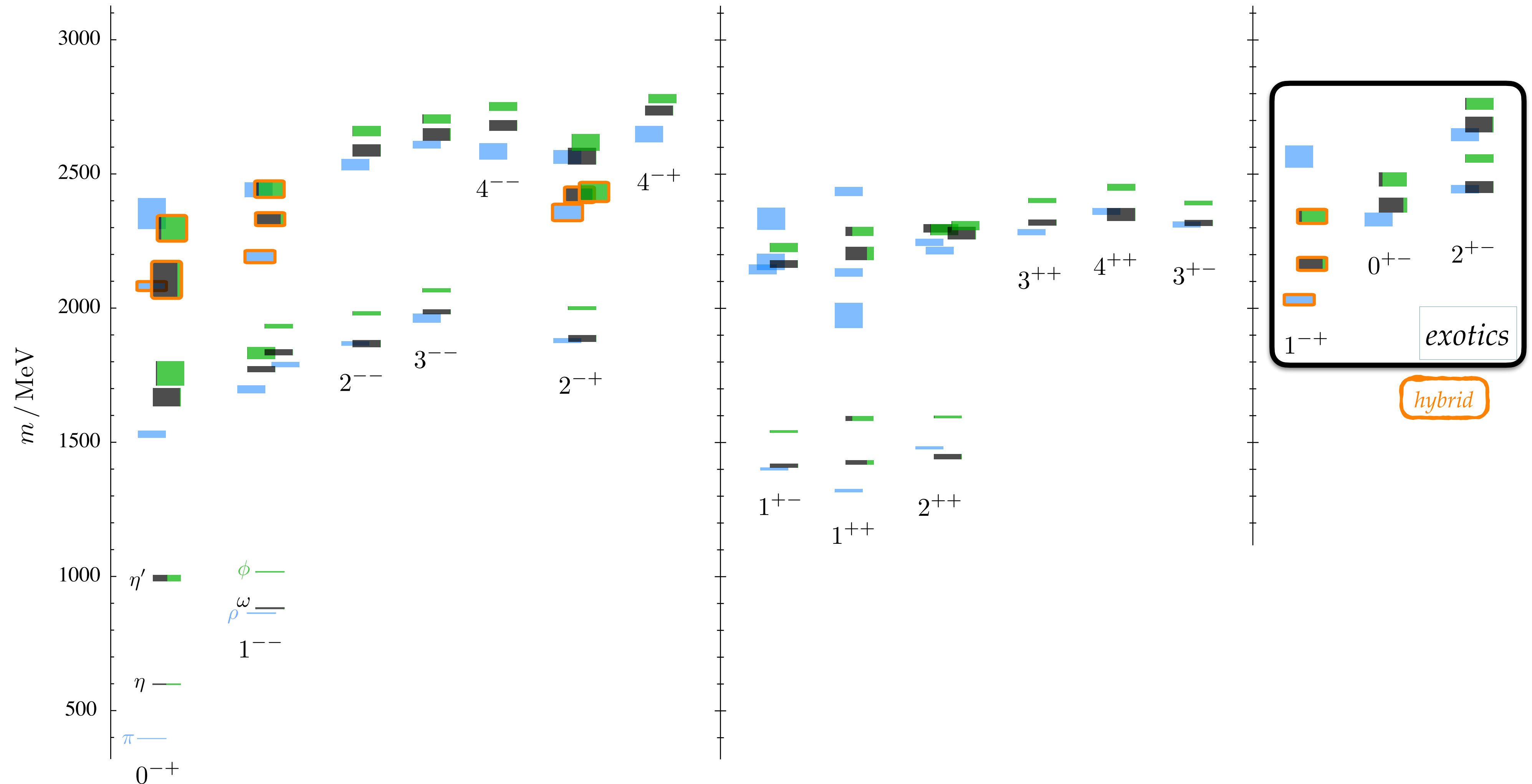
Pole in the two-body scattering amplitude at threshold: $\mathcal{M} \sim \frac{1}{p \cot \delta - ip} = \frac{1}{ip}$

Infinite tower of geometrically-separated three-body bound states: $E_{N+1} = E_N/\lambda$ where $\lambda = 22.69438$



Vitaly Efimov

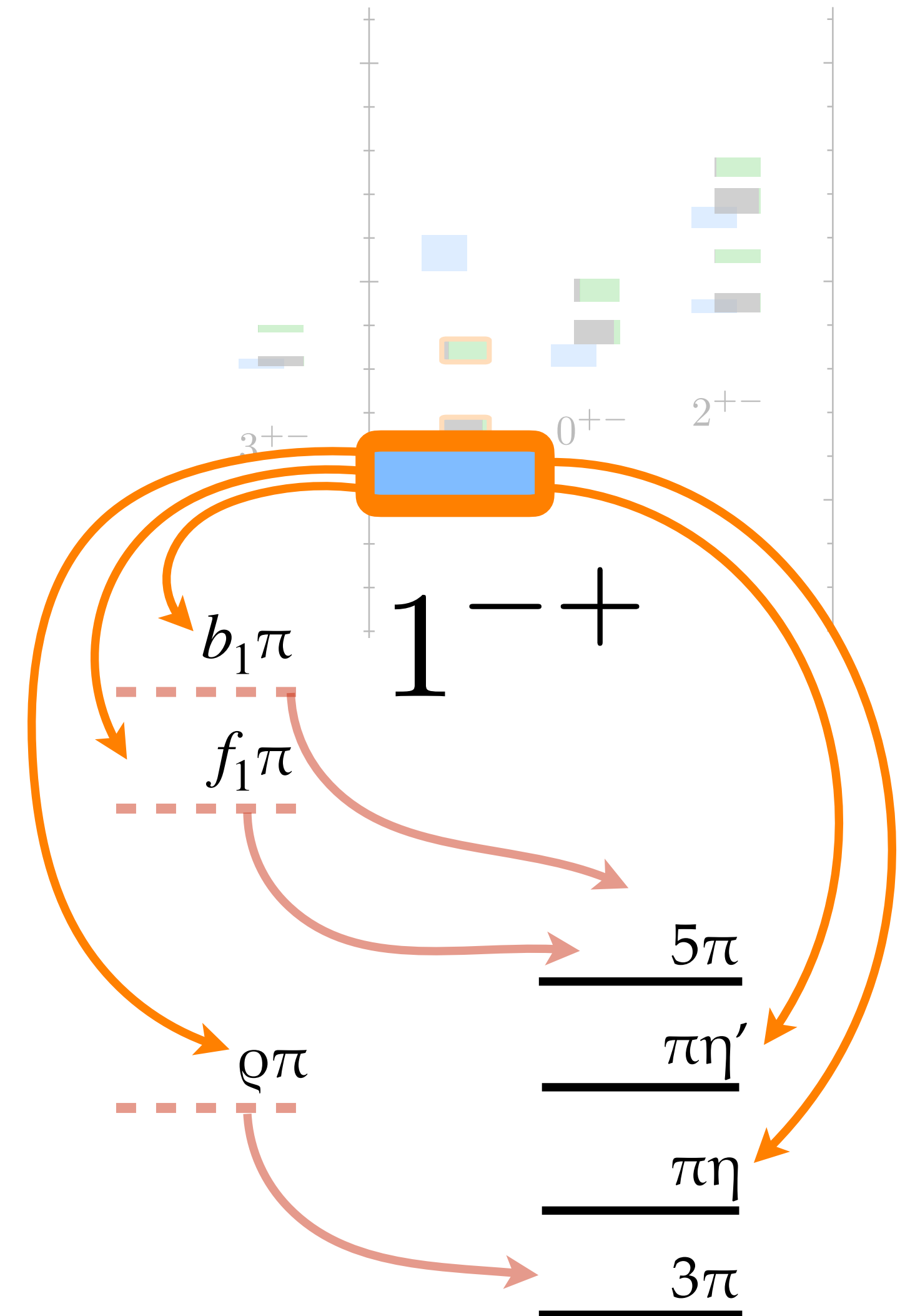
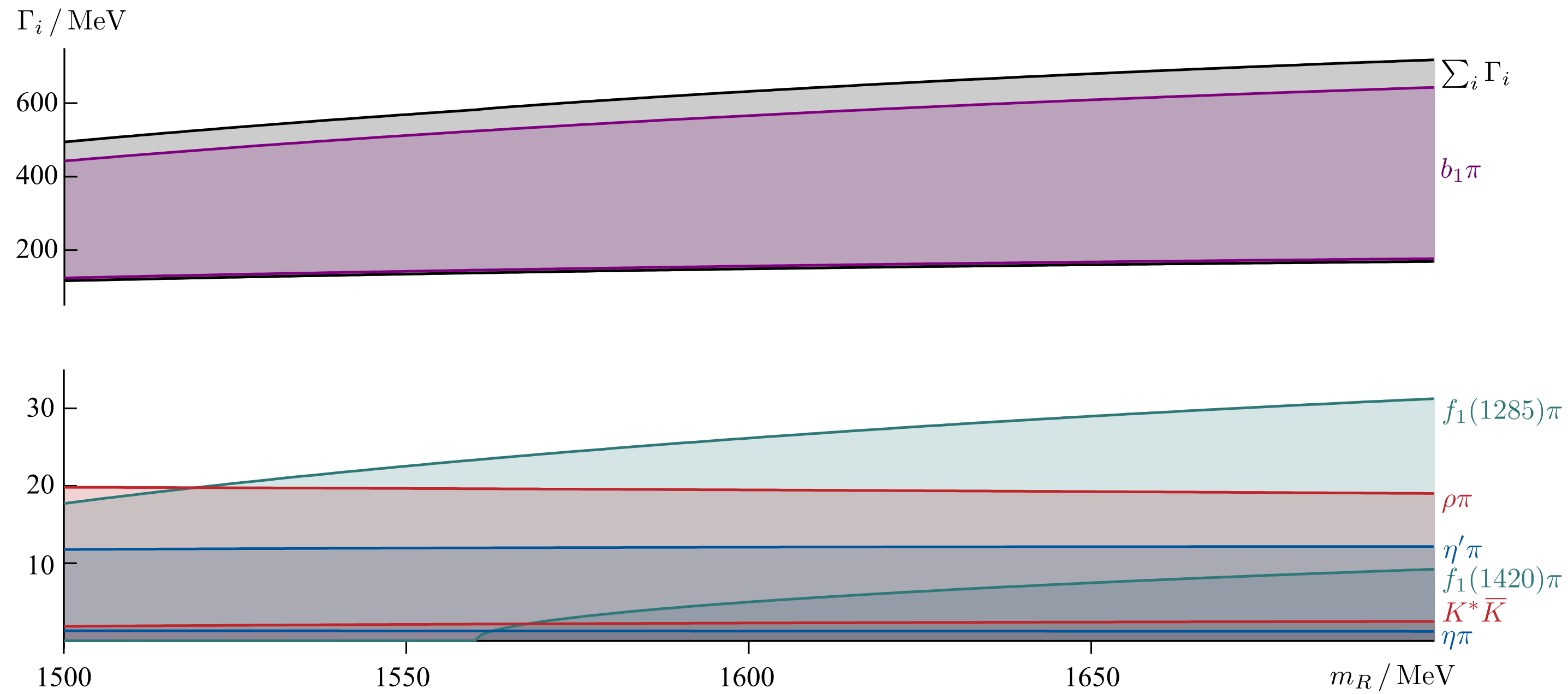
the three-body problem QCD Spectroscopy



$m_{\pi}=391$ MeV

Dudek, Edwards, Guo, Thomas (2013) ~~RB~~

the three-body problem QCD Spectroscopy



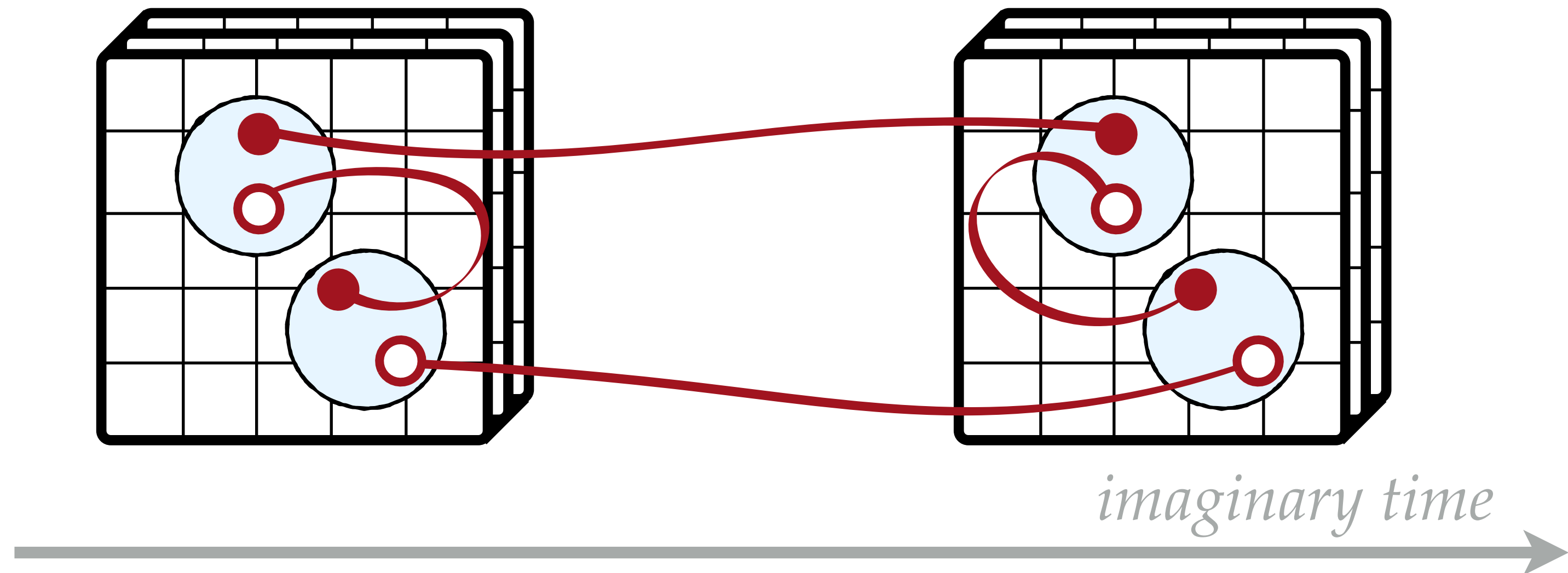
Woss, Dudek, Edwards, Thomas, Wilson (2020)

Dudek, Edwards, Guo, Thomas (2013) ~~RB~~

Lattice QCD in a nutshell

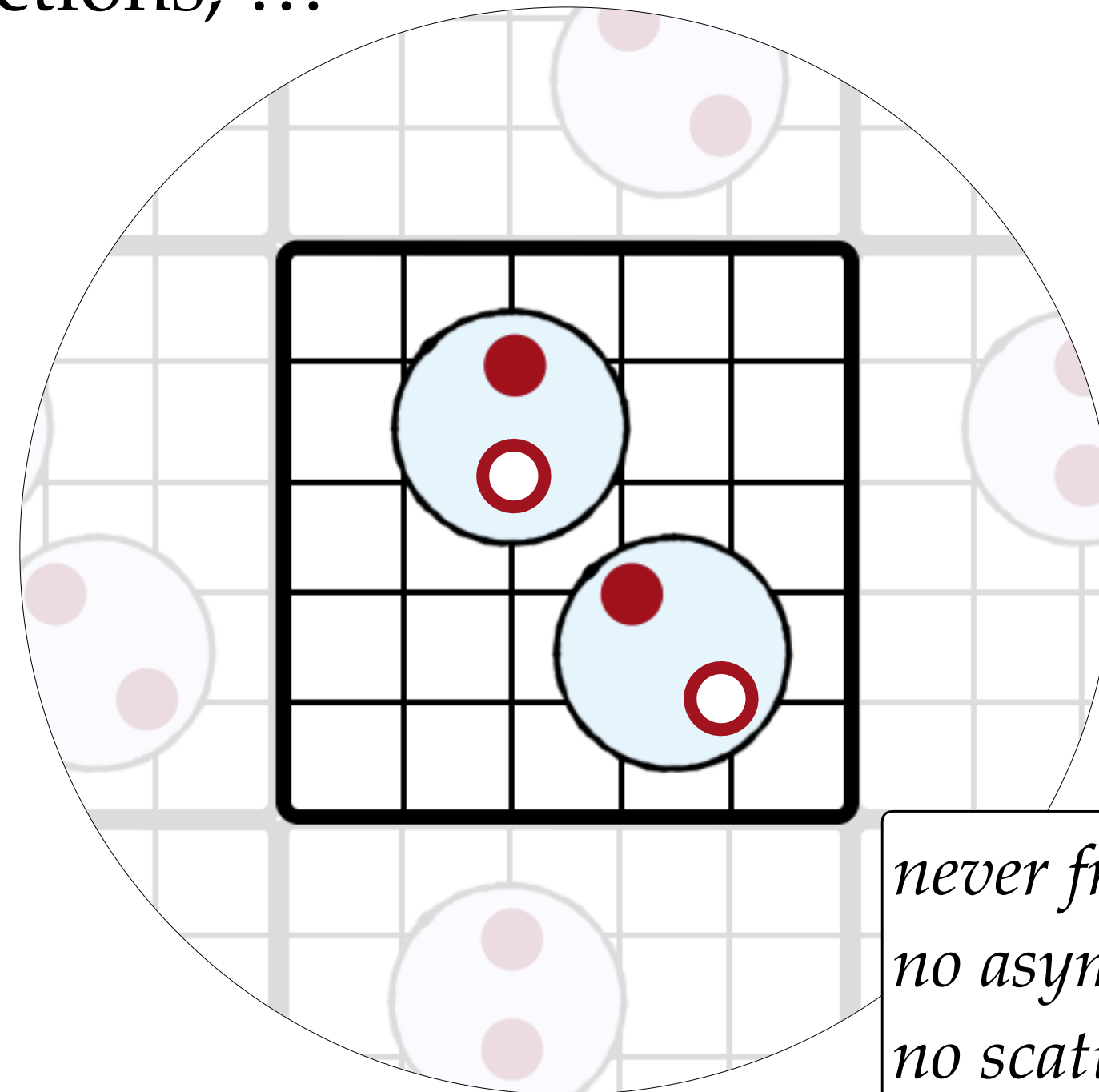
- lattice spacing: $a \sim 0.03 - 0.12$ fm
- finite volume: $L \sim 6 - 12$ fm
- quark masses
- Euclidean spacetime: $t_M \rightarrow -it_E$
- Importance sampling
- Correlation functions, ...

$$C^{2pt.}(t_E) \equiv \langle 0 | \mathcal{O}(t_E) \mathcal{O}^\dagger(0) | 0 \rangle = \sum_n c_n e^{-E_n t_E}$$



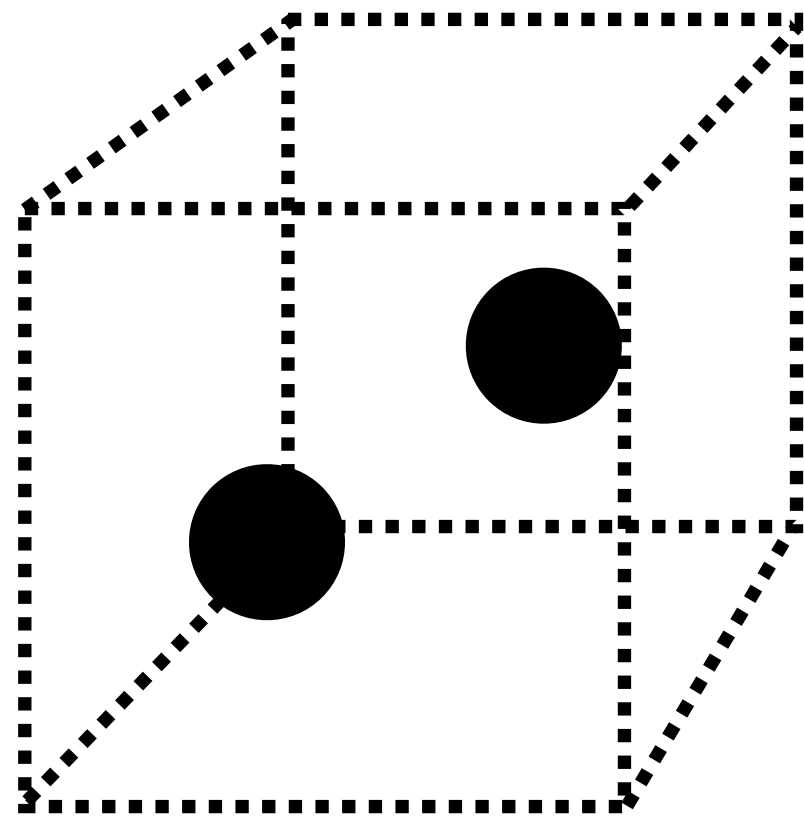
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*never free
no asymptotic states
no scattering*

Two-hadron systems

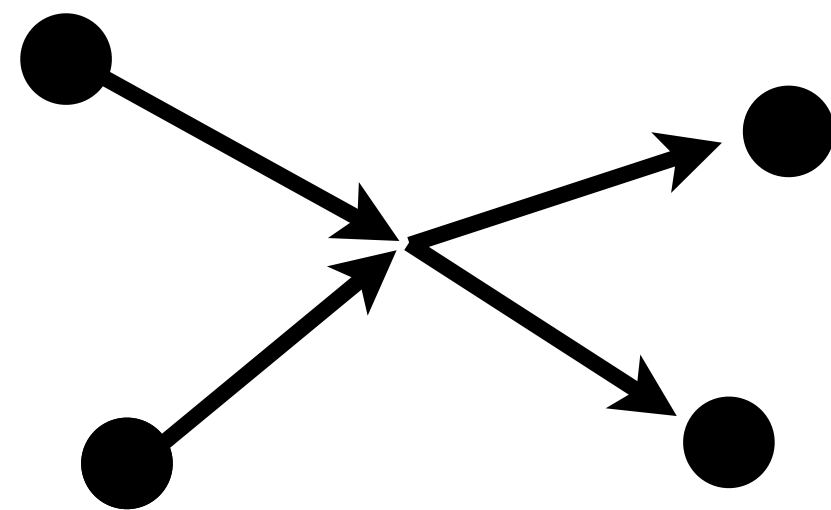


finite-volume
spectroscopy

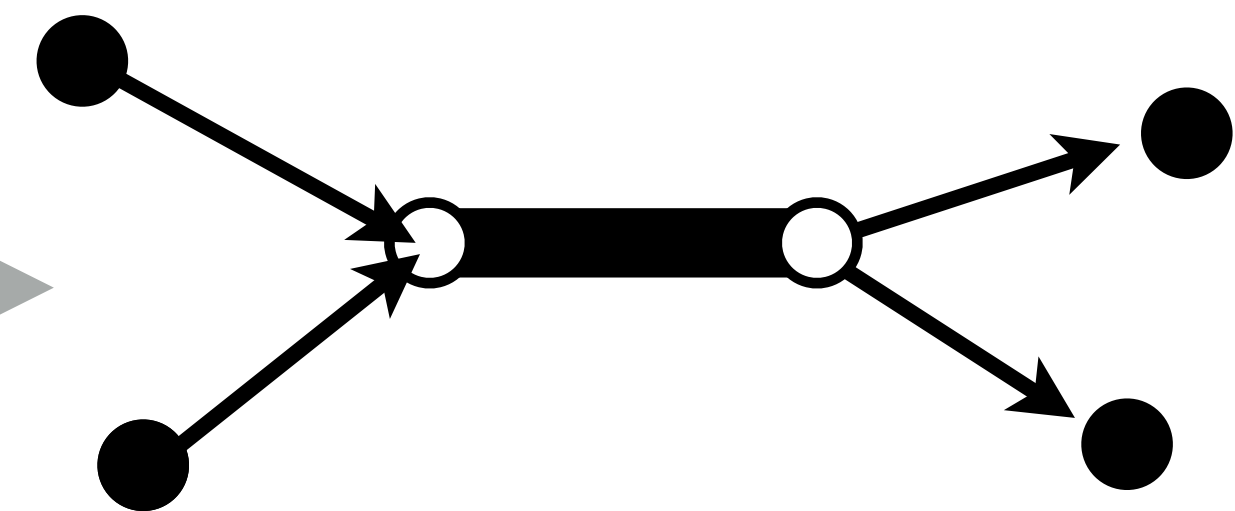
$$\det [F(E_L, L) + \mathcal{M}^{-1}(E_L)] = 0$$



- Lüscher (1986, 1991)
- Rummukainen & Gottlieb (1995)
- Kim, Sachrajda, & Sharpe (2005)
- Christ, Kim & Yamazaki (2005)
- Feng, Li, & Liu (2004)
- Hansen & Sharpe (2021)
- RB & Davoudi (2012)
- RB (2014)



infinite-volume
scattering amplitudes



bound state and
resonance poles

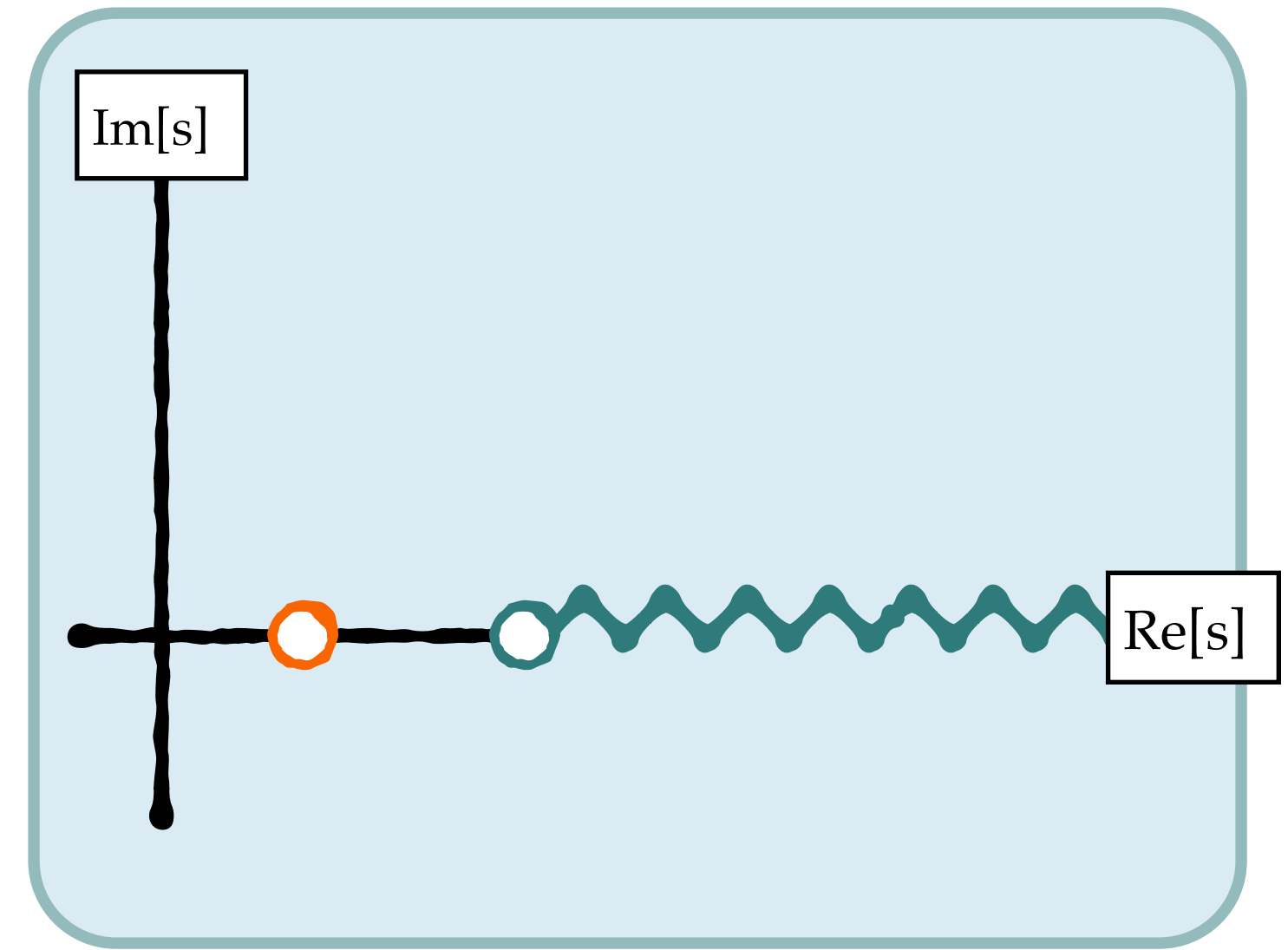
Two-hadron systems

Infinite Minkowski spacetime

$$i\mathcal{M} = \text{blob} = \text{circle} + \text{blob} \circ \text{circle}$$
$$= \text{diamond} + \text{diamond} \circ \text{circle}$$

placing all legs on-shell \longrightarrow $\frac{i}{\mathcal{K}^{-1} - i\rho}$

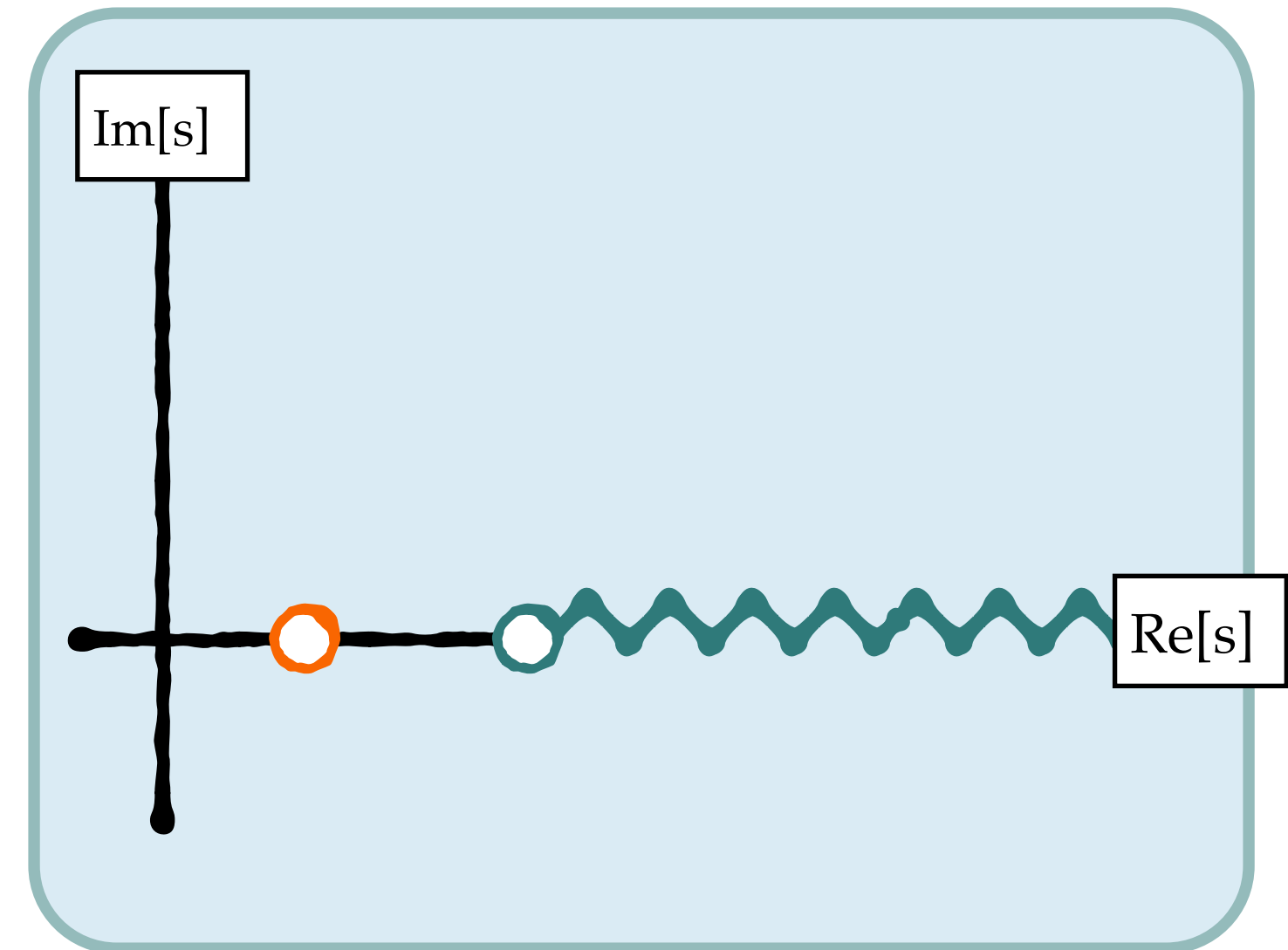
$$\rho \equiv \frac{p}{8\pi E} \sim \sqrt{s - s_{th}}$$



Two-hadron systems

Infinite Minkowski spacetime

$$i\mathcal{M} = \text{blob} = \text{circle} + \text{circle with blob} \\ = \text{diamond} + \text{diamond with dashed line} \xrightarrow{\text{placing all legs on-shell}} \frac{i}{\mathcal{K}^{-1} - i\rho}$$



Finite Minkowski spacetime

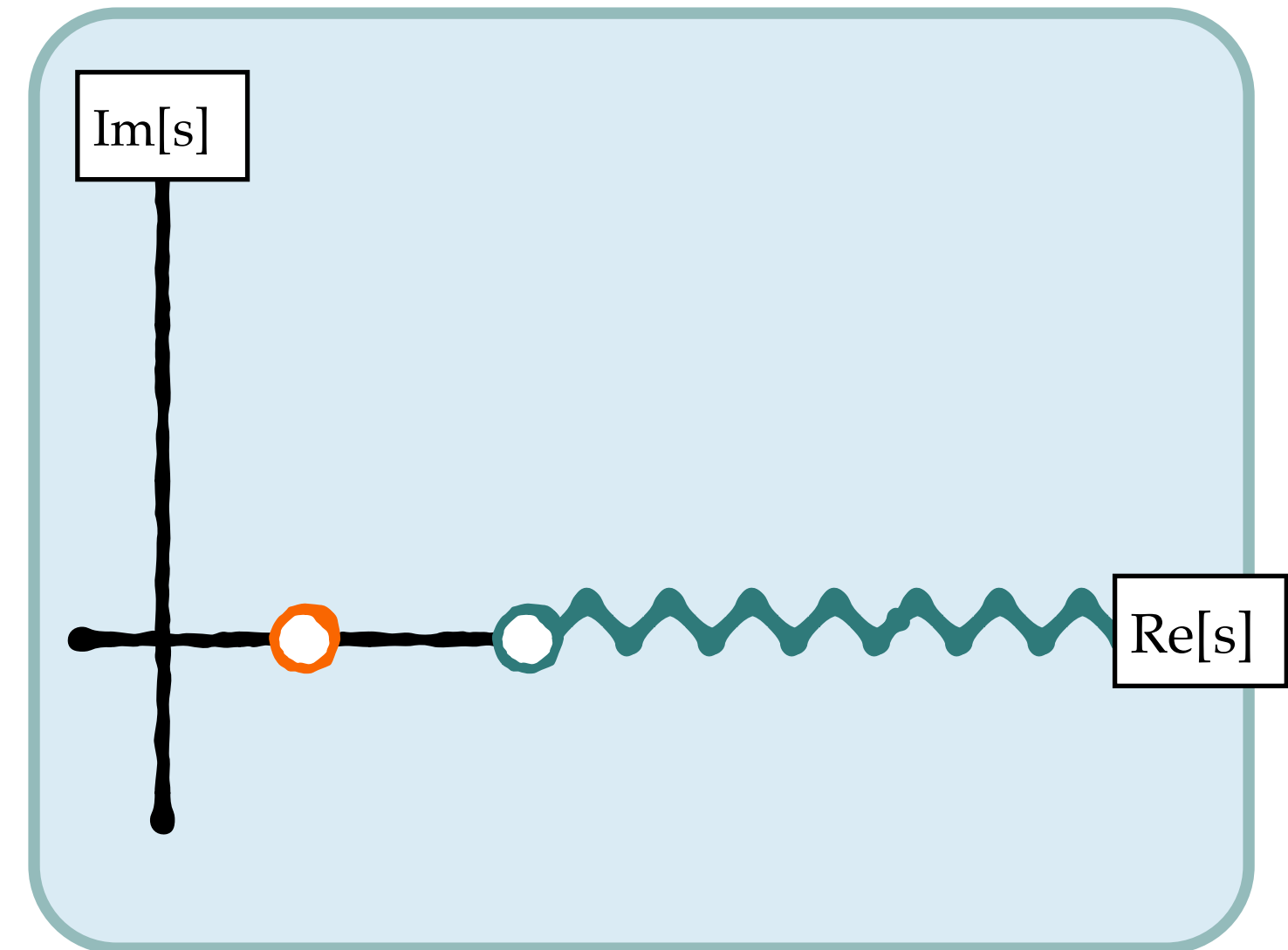
$$i\mathcal{M}_L = \text{square} = \text{circle} + \text{circle with V} \text{ square} \\ = \text{blob} + \text{blob with dashed line} \text{ square} \quad iF$$

$$iF = \left[\frac{1}{L^3} \sum_{\mathbf{k}} - \int \frac{d^3k}{(2\pi)^3} \right] \frac{1}{(2\omega_k)^2} \frac{i}{E - 2\omega_k}$$

Two-hadron systems

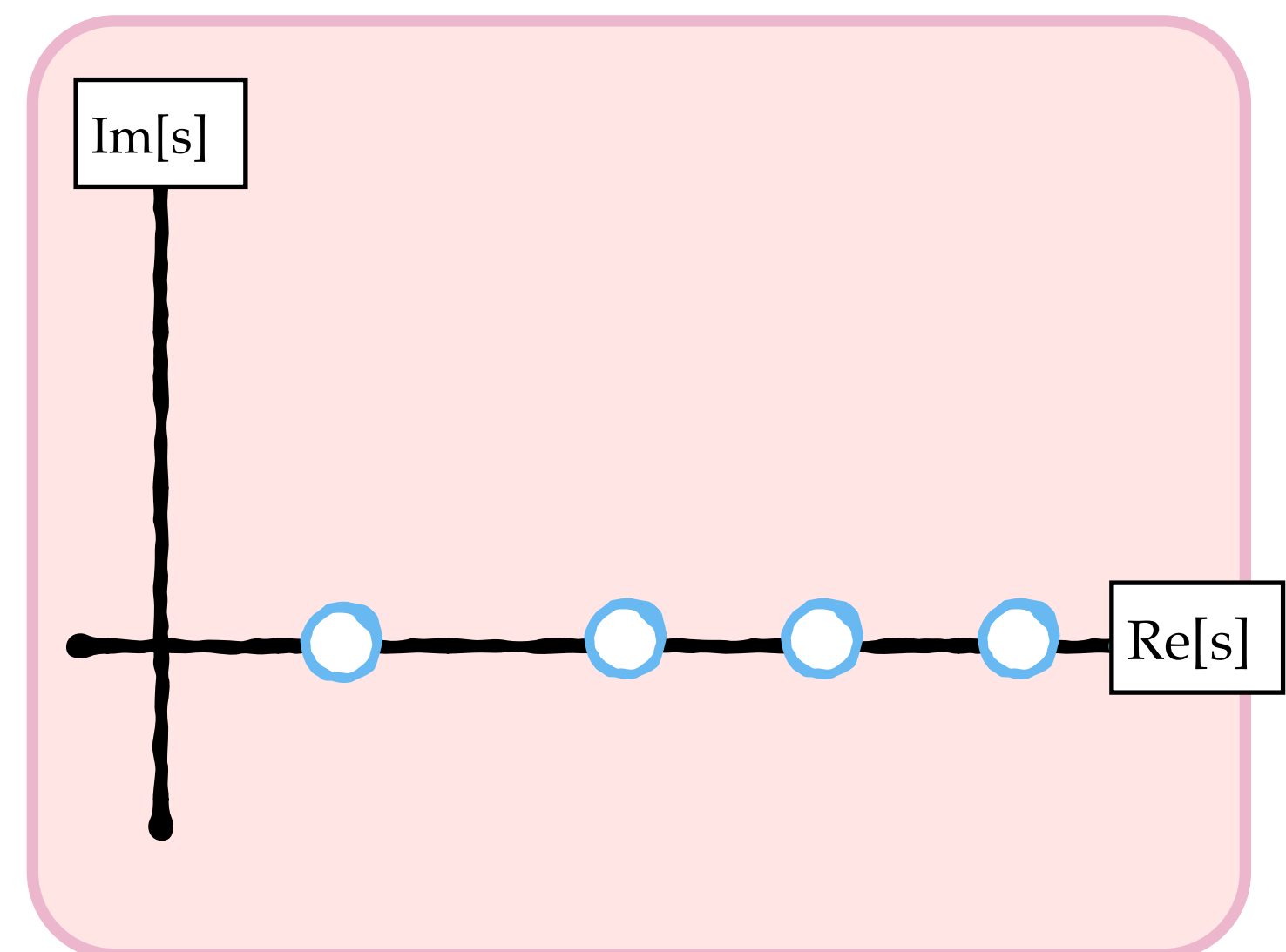
Infinite Minkowski spacetime

$$\begin{aligned}
 i\mathcal{M} &= \text{[Diagram: four external legs meeting at a central black circle]} = \text{[Diagram: four external legs meeting at a central white circle]} + \text{[Diagram: four external legs meeting at a central white circle, with a loop on the right]} \\
 &= \text{[Diagram: four external legs meeting at a central white diamond]} + \text{[Diagram: four external legs meeting at a central white diamond, with a loop on the right and a vertical dashed line through it]} \xrightarrow{\text{placing all legs on-shell}} \frac{i}{\mathcal{K}^{-1} - i\rho}
 \end{aligned}$$



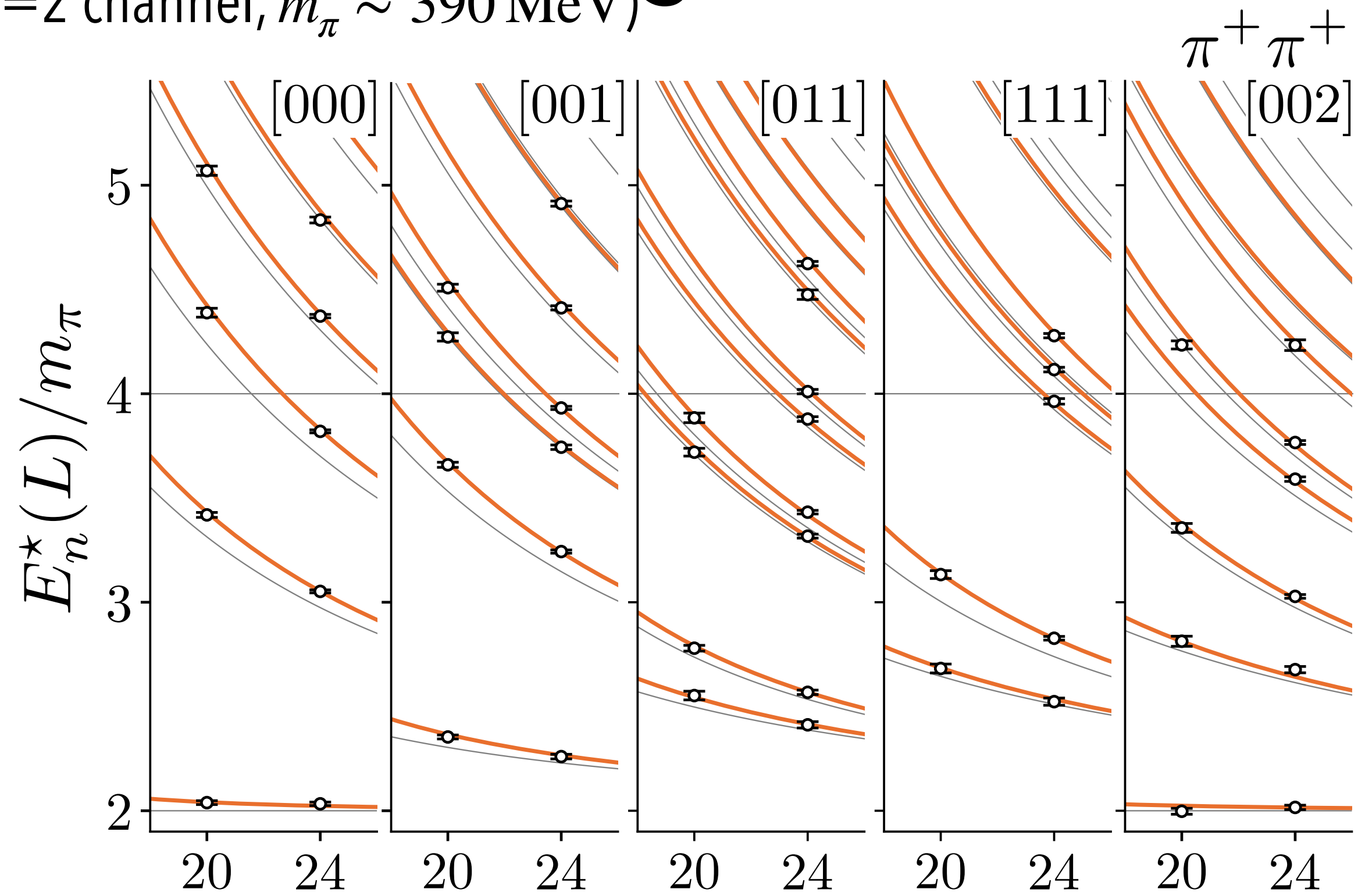
Finite Minkowski spacetime

$$\begin{aligned}
 i\mathcal{M}_L &= \text{[Diagram: four external legs meeting at a central black square]} = \text{[Diagram: four external legs meeting at a central white circle]} + \text{[Diagram: four external legs meeting at a central white circle, with a loop on the right labeled V]} \\
 &= \text{[Diagram: four external legs meeting at a central black circle]} + \text{[Diagram: four external legs meeting at a central black circle, with a loop on the right and a vertical dashed line through it labeled iF]} = \frac{1}{\mathcal{M}^{-1} + F}
 \end{aligned}$$



$\pi\pi$ scattering

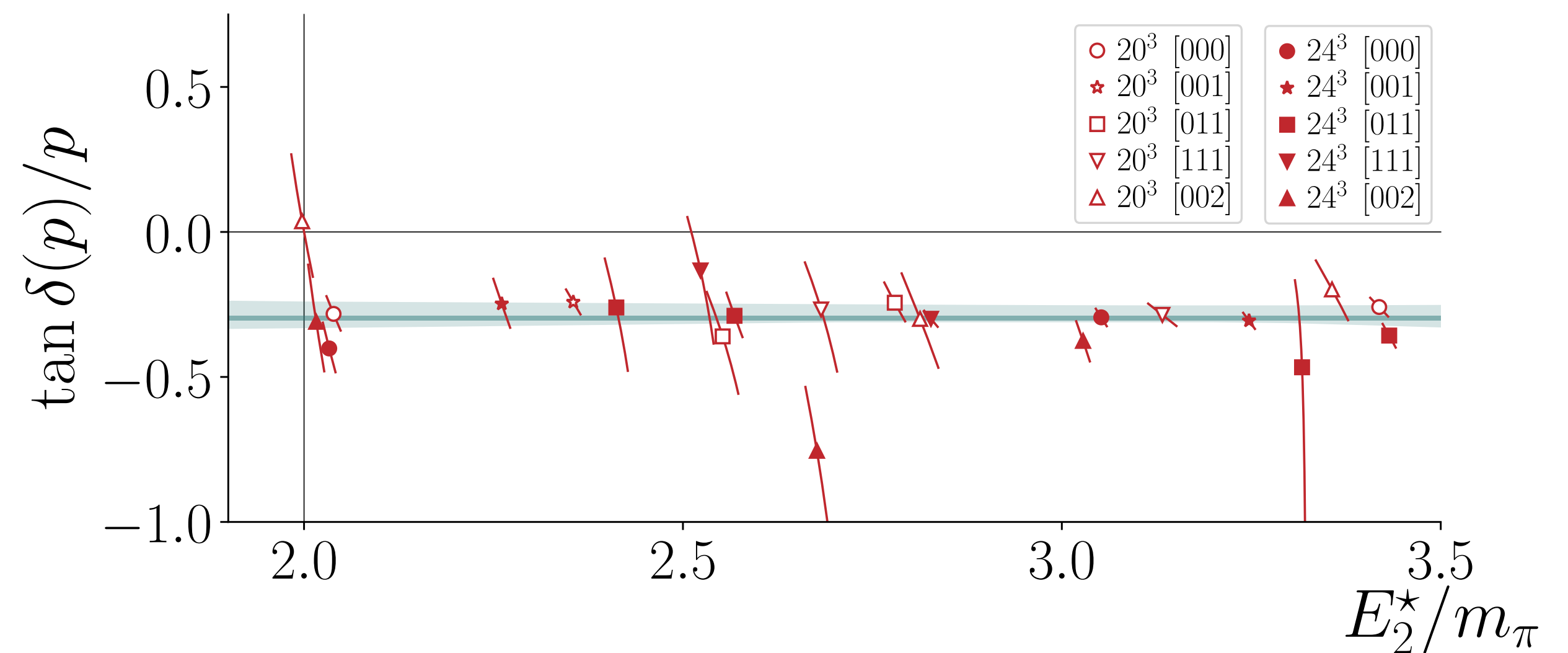
($l=2$ channel, $m_\pi \sim 390$ MeV)



$\pi^+ \pi^+$

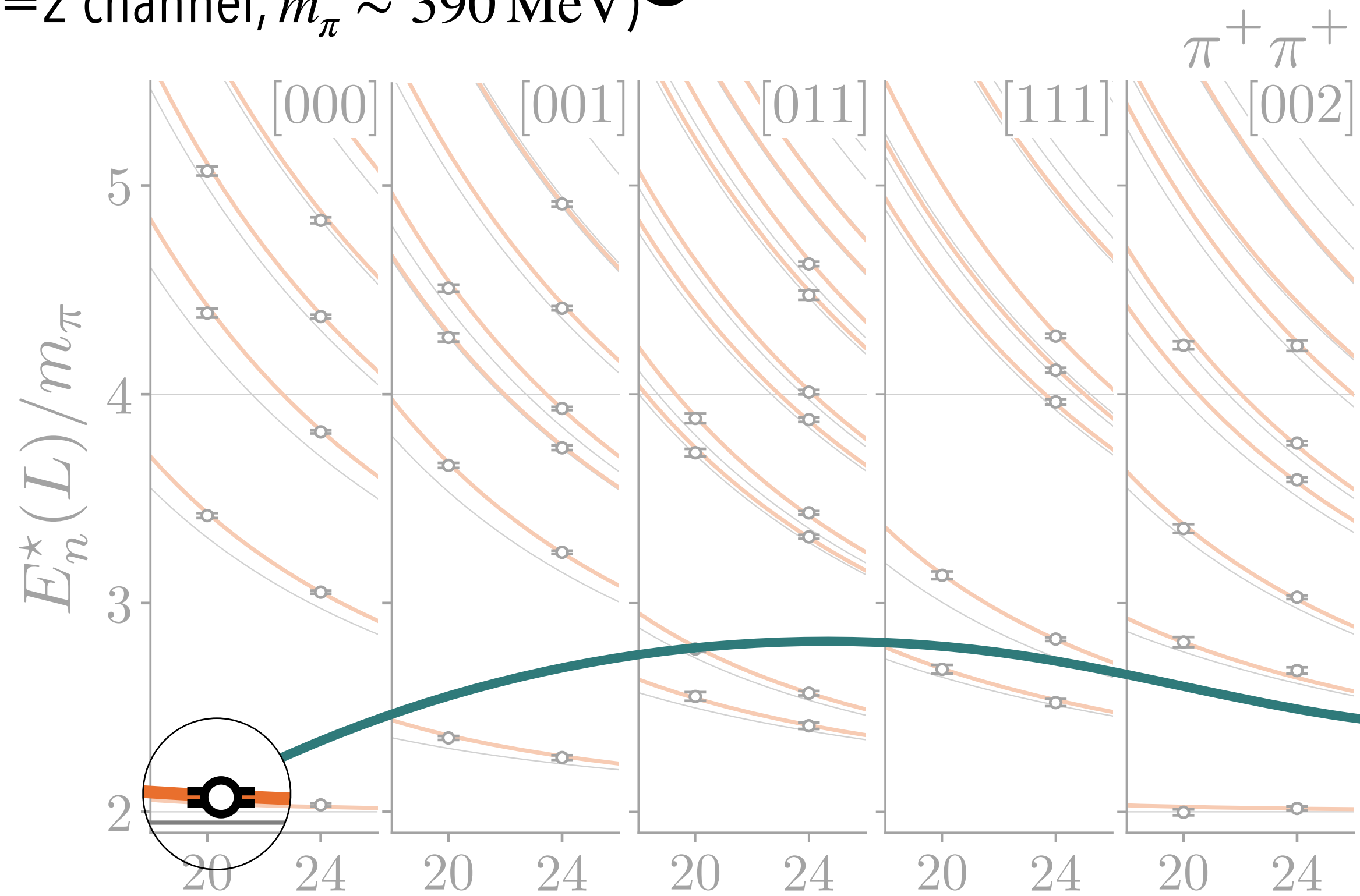
$$\det[F^{-1}(P, L) + \mathcal{M}(P)] = 0$$

$$\mathcal{M} \sim \frac{1}{p \cot \delta - ip}$$



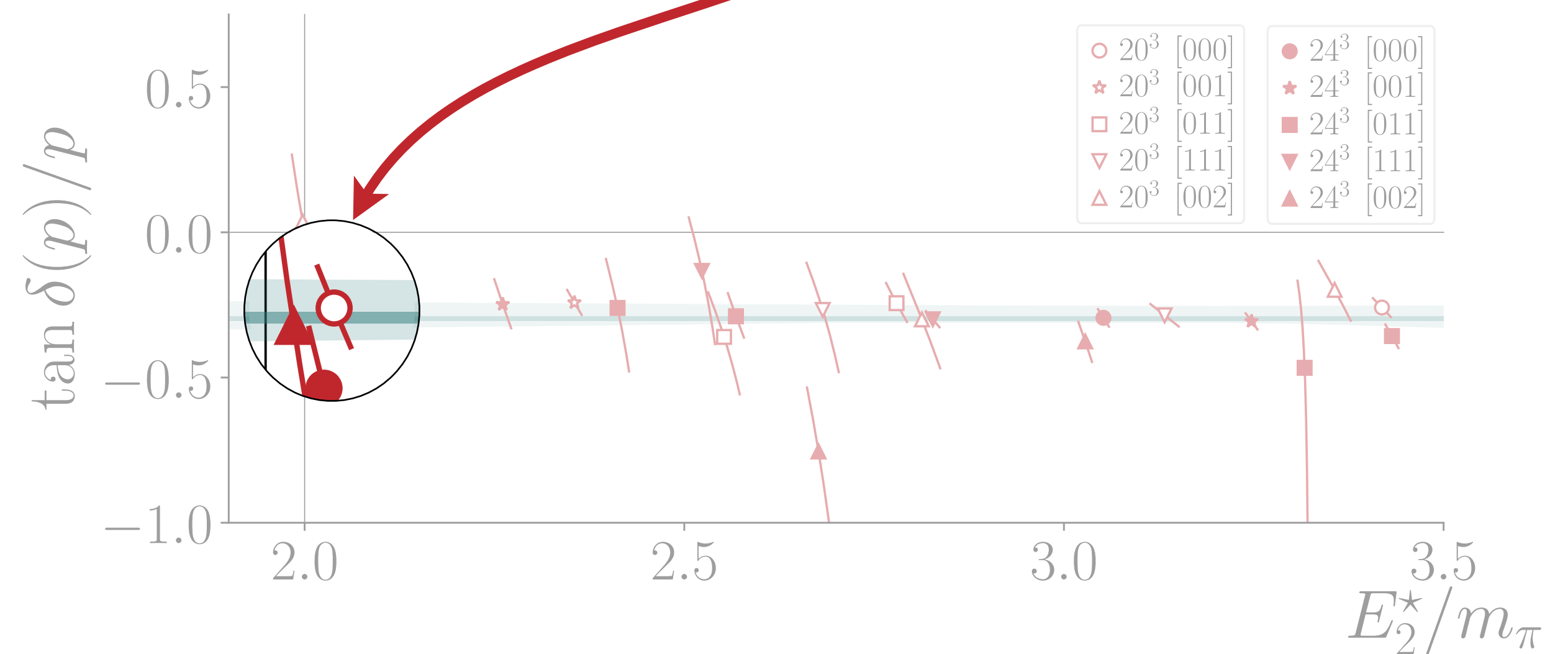
$\pi\pi$ scattering

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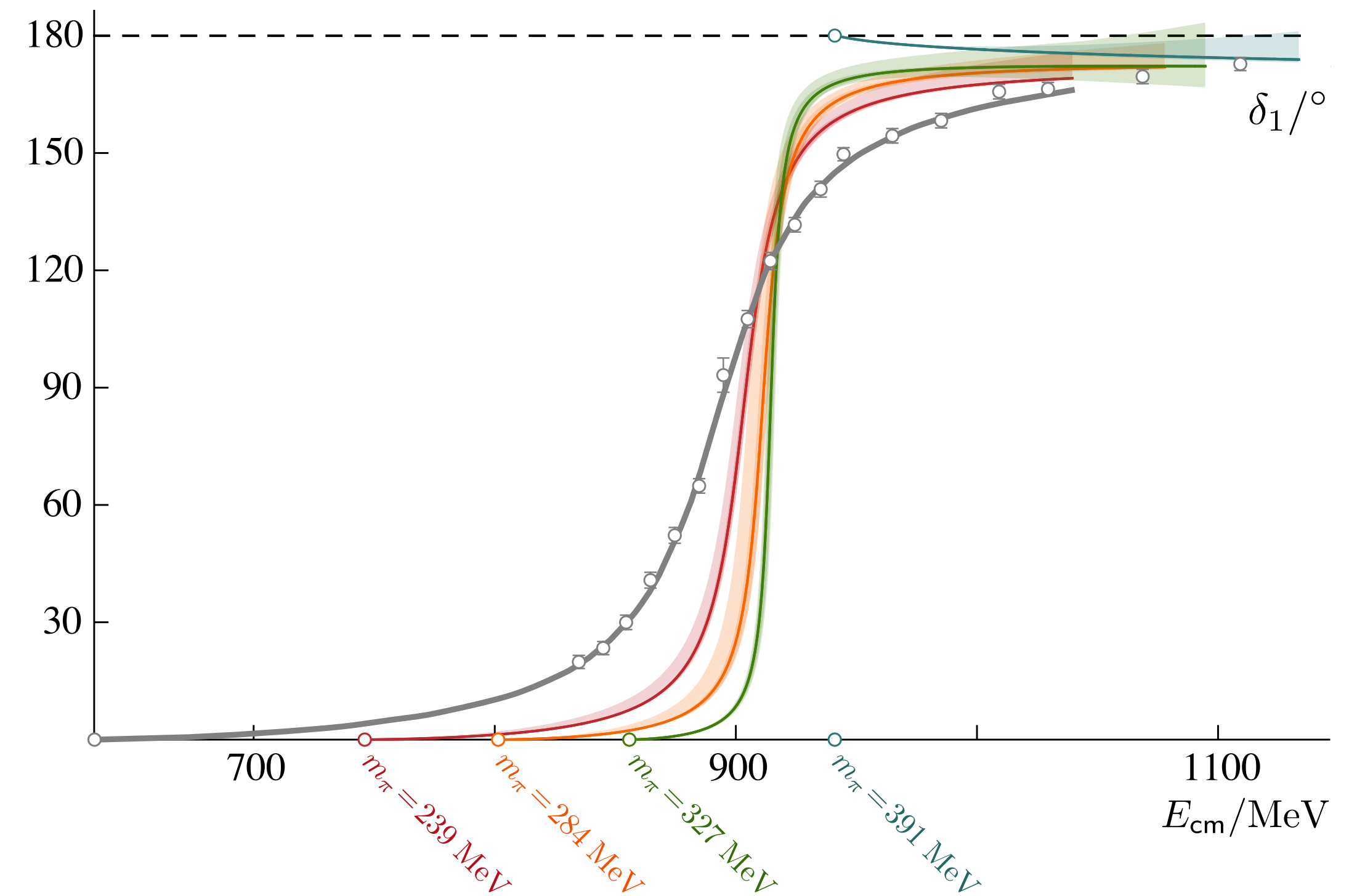
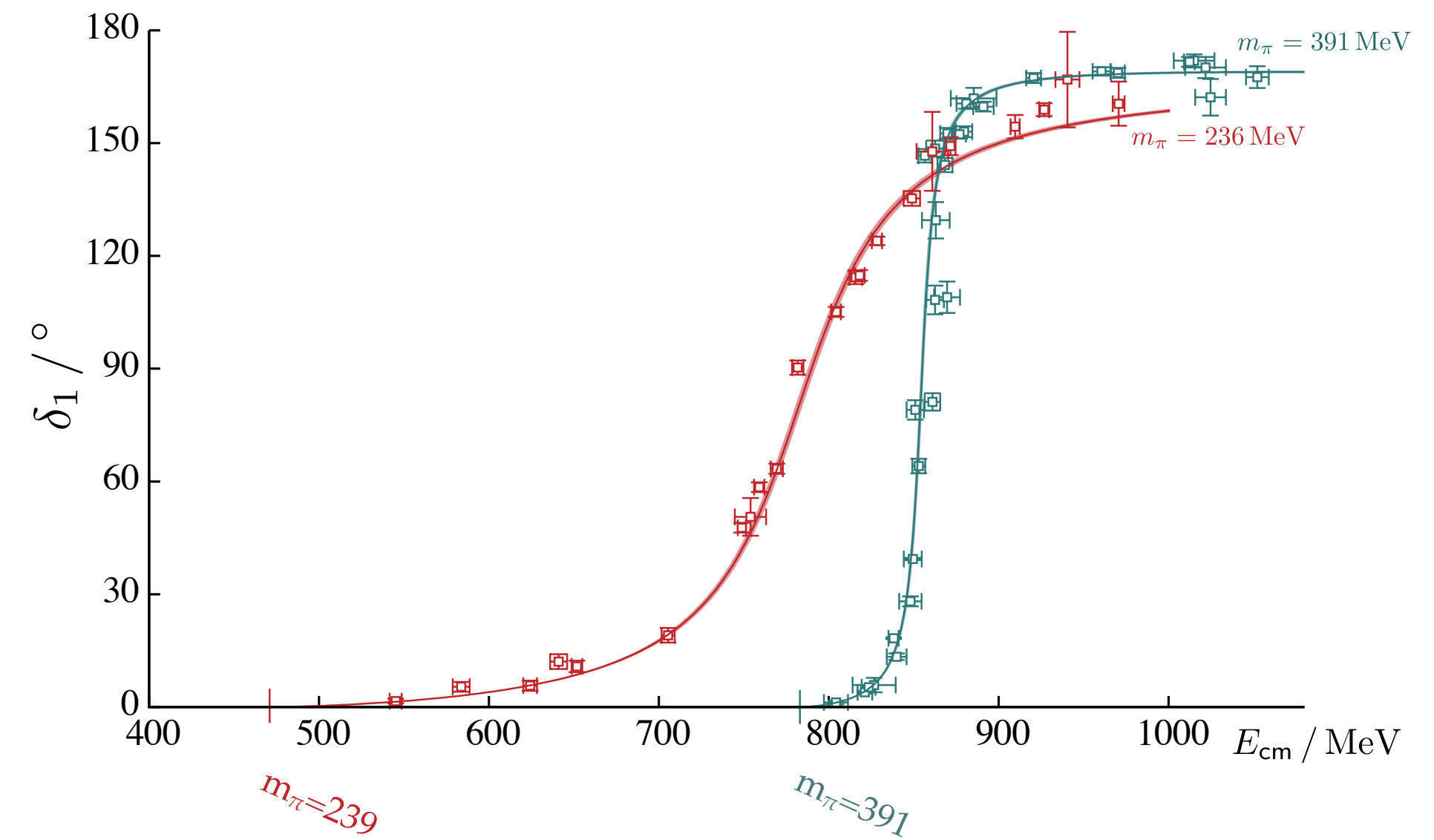
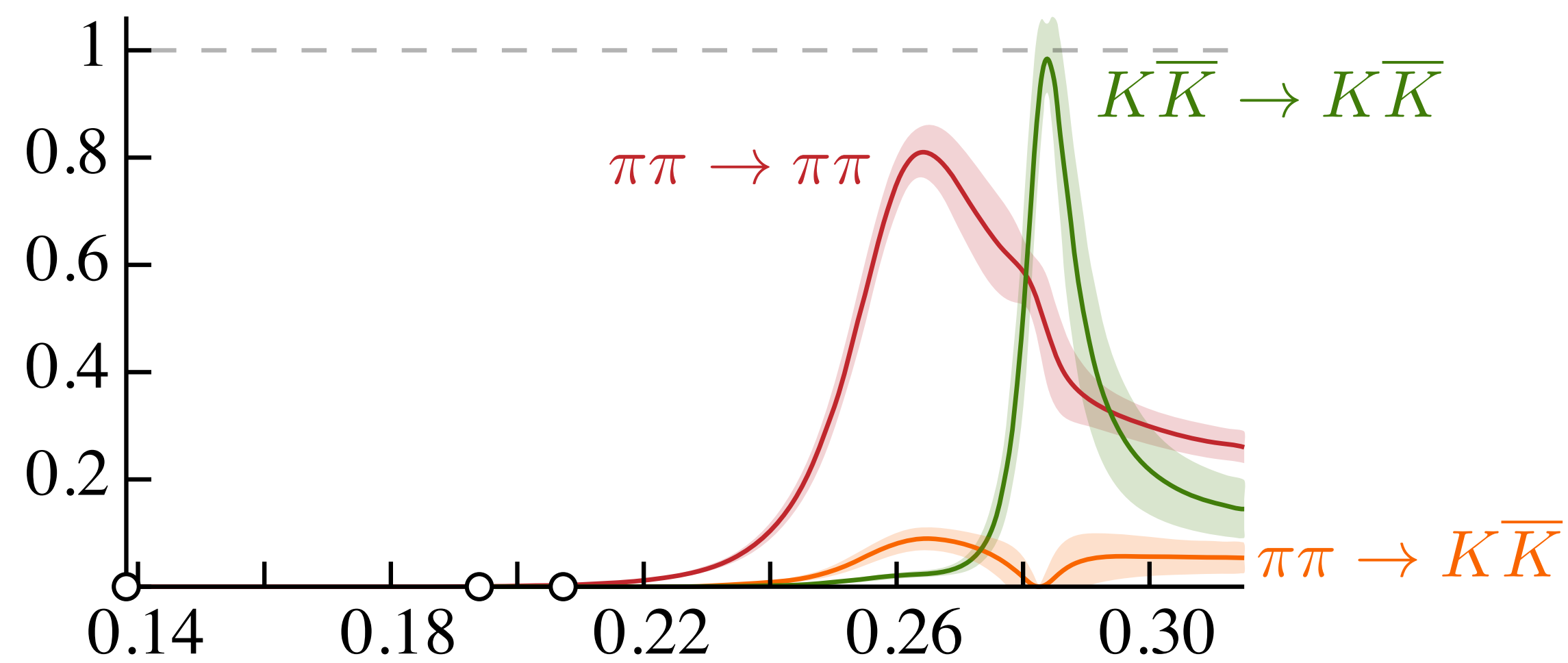
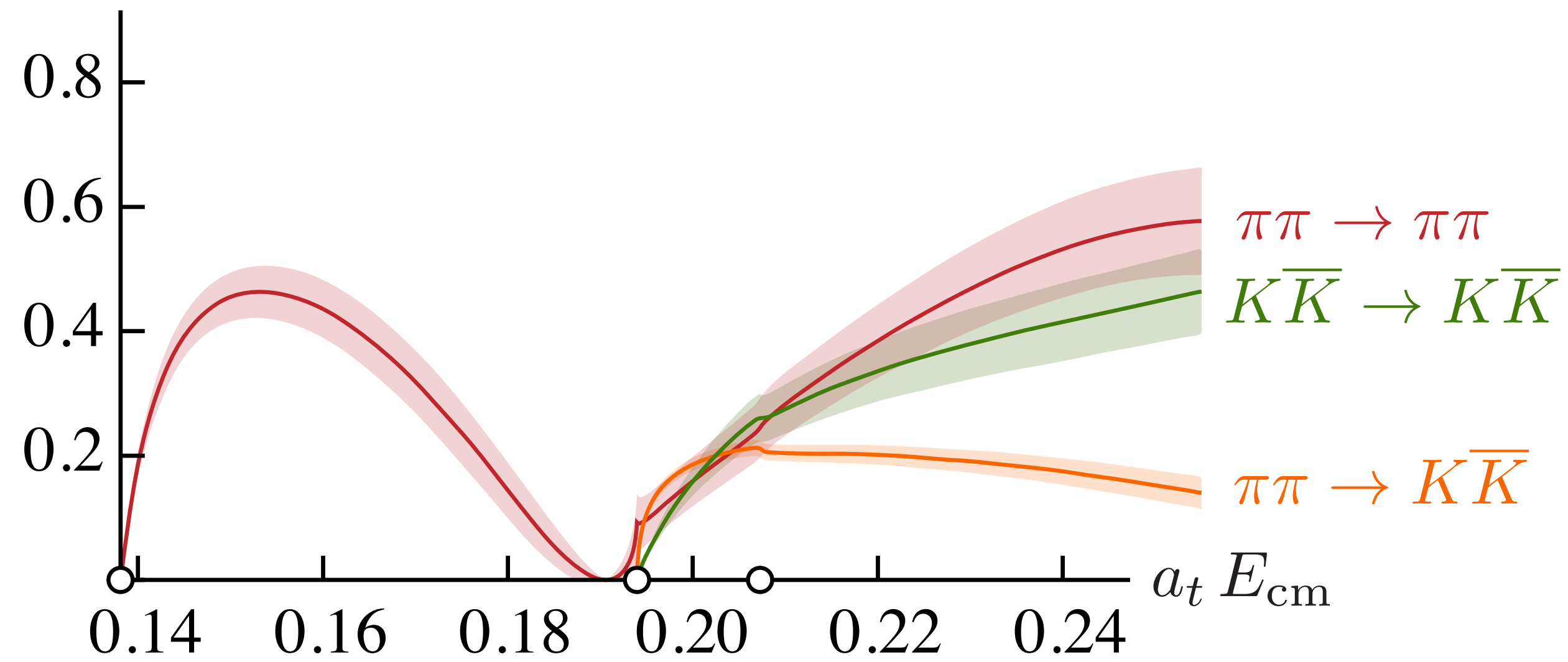
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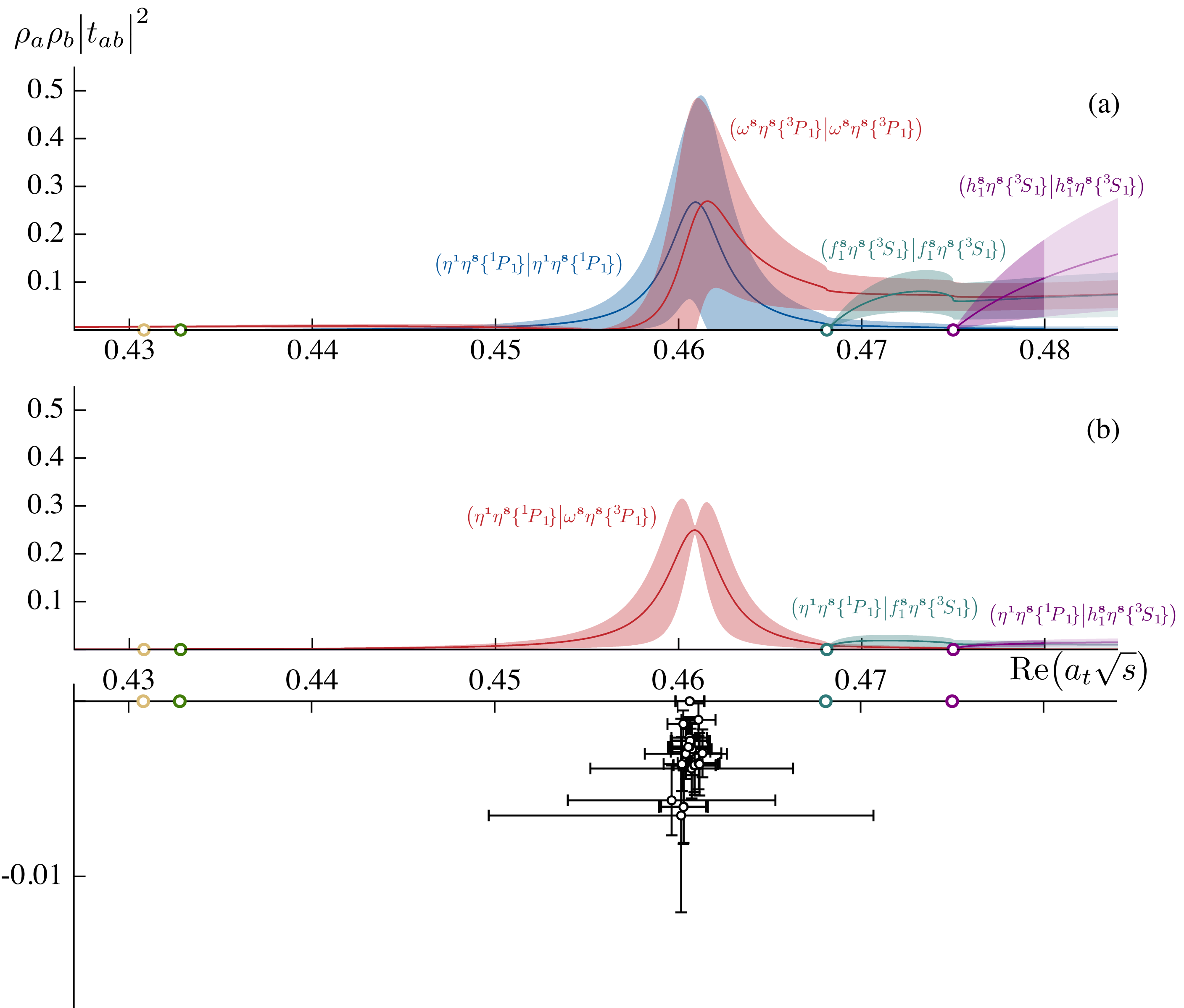
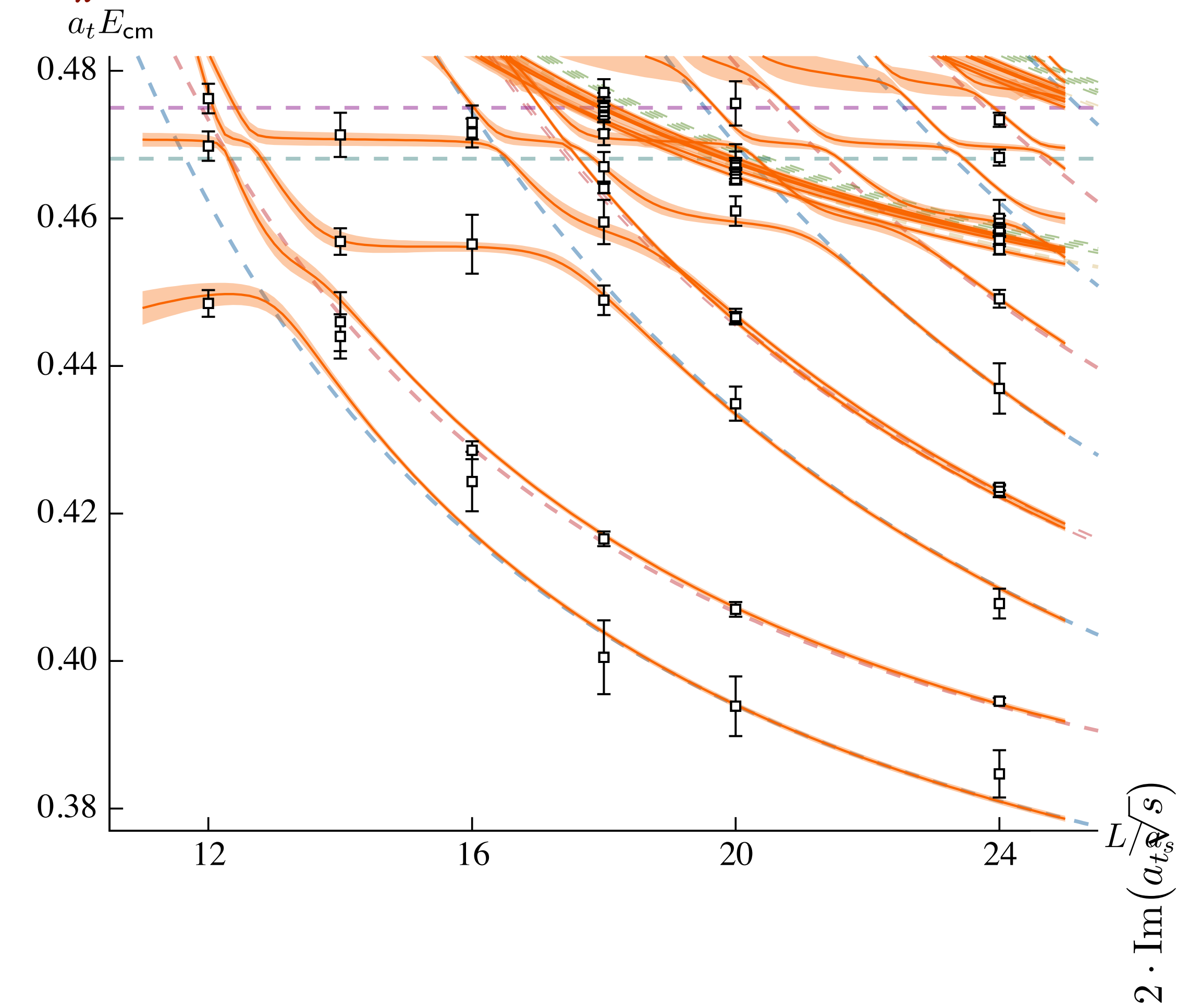
Many other calculations

$$|\mathcal{M}_{ij}|^2 \rho_i \rho_j$$

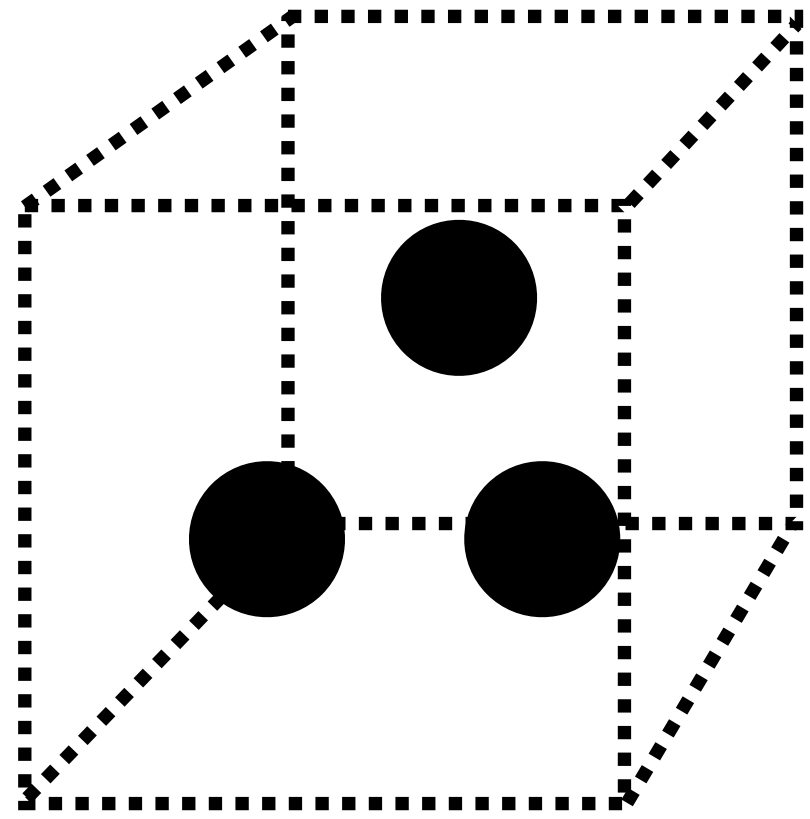


π_1 channel

$(m_\pi \sim 700 \text{ MeV})$



Three-hadron systems

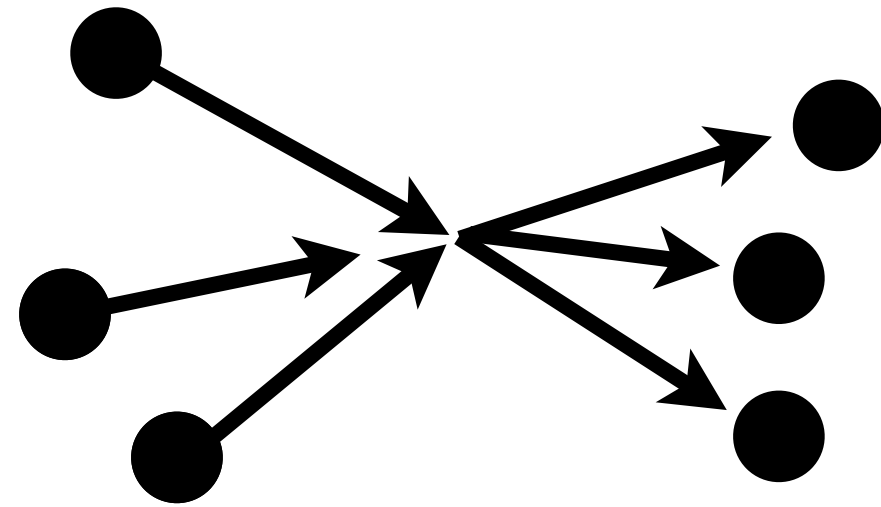


finite-volume
spectroscopy

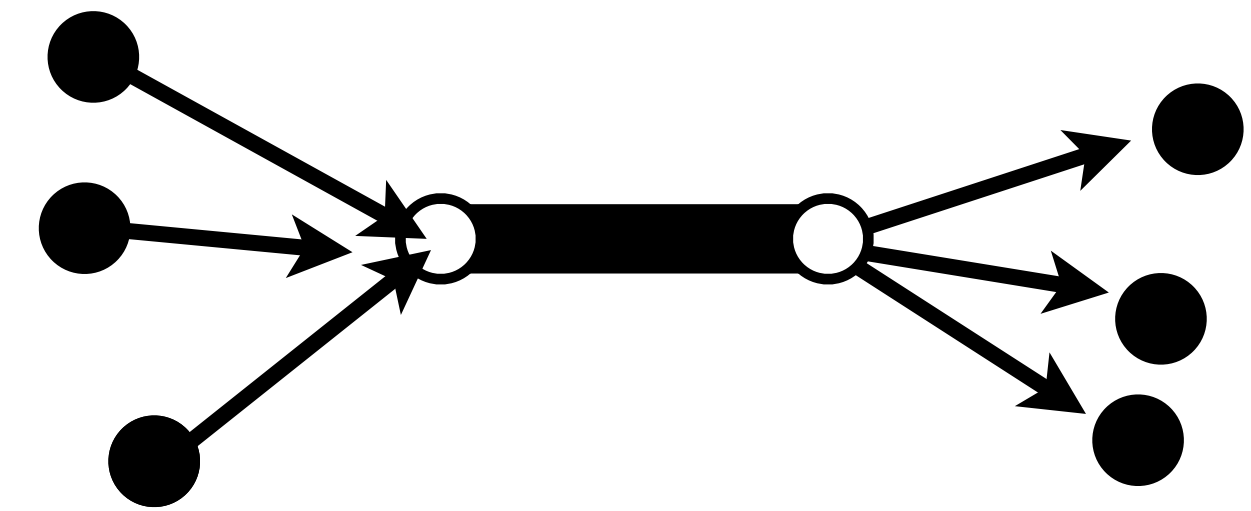
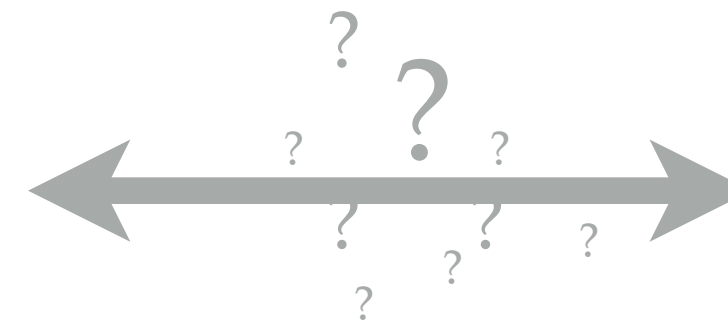
$$\det [F_3(E_L, L) + \mathcal{K}_{\text{df},3}^{-1}(E_L)] = 0$$



Hansen & Sharpe ('14, '15)
Mai & Döring ('17)
RB, Hansen & Sharpe ('18)
Hansen, Romero-Lopez & Sharpe ('20)
Blanton & Sharpe ('20)
Jackura ('22)



infinite-volume
scattering amplitudes



bound state and
resonance poles

Three-hadron systems

Infinite Minkowski spacetime

$$i\mathcal{M}_3 = \text{[Diagram 1]} = \text{[Diagram 2]} + \text{[Diagram 3]} + \dots$$

$G \sim \frac{1}{(P - p - k)^2 - m^2}$

Three-hadron systems

Infinite Minkowski spacetime

$$i\mathcal{M}_3 = \text{[Diagram: 3 legs on a black circle]} = \text{[Diagram: 3 legs on a white circle]} + \text{[Diagram: 3 legs on a white circle, 1 leg on a black circle]} + \text{[Diagram: 3 legs on a white circle, 2 legs on a black circle]} + \text{[Diagram: 3 legs on a black circle, 1 leg on a white circle]} + \text{[Diagram: 3 legs on a black circle, 2 legs on a white circle]}$$

*placing all
legs on-shell*

$$\longrightarrow i\mathcal{D} + i\mathcal{L} \frac{1}{\mathcal{K}_{3,\text{df}}^{-1} + F_3^\infty} \mathcal{L}$$

satisfies an integral equations

$$i\mathcal{D} = i\mathcal{M}_2 iG i\mathcal{M}_2 + \int i\mathcal{M}_2 iG i\mathcal{D}$$

Three-hadron systems

Infinite Minkowski spacetime

$$i\mathcal{M}_3 = \text{[diagram: black circle with 6 legs]} = \text{[diagram: white circle with 6 legs]} + \text{[diagram: white circle with 3 legs connected to black circle with 3 legs]} + \text{[diagram: white circle with 3 legs connected to black circle with 3 legs via a loop]} + \text{[diagram: white circle with 3 legs connected to black circle with 3 legs via a loop with a dot]} + \text{[diagram: white circle with 3 legs connected to black circle with 3 legs via a loop with a dot and a circle}$$

placing all legs on-shell $\longrightarrow i\mathcal{D} + i\mathcal{L} \frac{1}{\mathcal{K}_{3,\text{df}}^{-1} + F_3^\infty} \mathcal{L}$

Finite Minkowski spacetime

$$i\mathcal{M}_{3,L} = \text{[diagram: black square with 6 legs]} = \text{[diagram: white circle with 6 legs]} + \text{[diagram: white circle with 3 legs connected to black square with 3 legs]} + \text{[diagram: white circle with 3 legs connected to black square with 3 legs via a loop labeled V]} + \text{[diagram: white circle with 3 legs connected to black square with 3 legs via a loop labeled V with a dot]} + \text{[diagram: white circle with 3 legs connected to black square with 3 legs via a loop labeled V with a dot and a circle}$$

$\longrightarrow \frac{\#}{\mathcal{K}_{3,\text{df}}^{-1} + F_3}$

spectrum satisfies

$$\det [\mathcal{K}_{3,\text{df}} + F_3^{-1}] = 0$$



Solving integral equations

- Introduce $\mathcal{D} = \mathcal{M}_2 d \mathcal{M}_2$
- Assuming $\mathcal{K}_{\text{df}} = 0$ and partial wave-projecting onto $\ell = 0$ sector,

$$d(p', s, p) = -G(p', s, p) - \int_0^{q_{\text{max}}} \frac{dq q^2}{(2\pi)^2 \omega_q} G(p', s, q) \mathcal{M}_2(q, s) d(q, s, p)$$

- Soften or isolate possible singularities: non-zero epsilon, integrate out poles analytically, ...
- Write as a matrix form:

$$\mathbf{d}(N, \epsilon) = -\mathbf{G}(\epsilon) - \mathbf{G}(\epsilon) \cdot \mathbf{P} \cdot \mathbf{d}(N, \epsilon)$$

- Recover the exact results when taking $\epsilon = \eta/N \rightarrow 0$.
- Including $\mathcal{K}_{\text{df}} \neq 0$ contributions is “easy”

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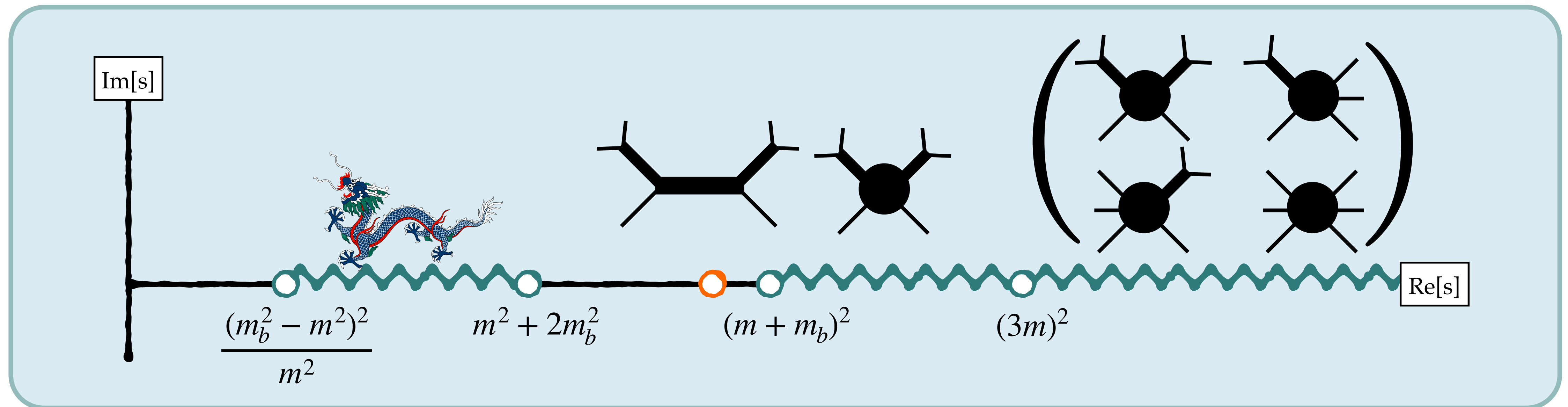
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Consistency checks using toy models

Consider a toy theory with a two-body bound state

□ Simple case $q \cot \delta = -\frac{1}{a}$

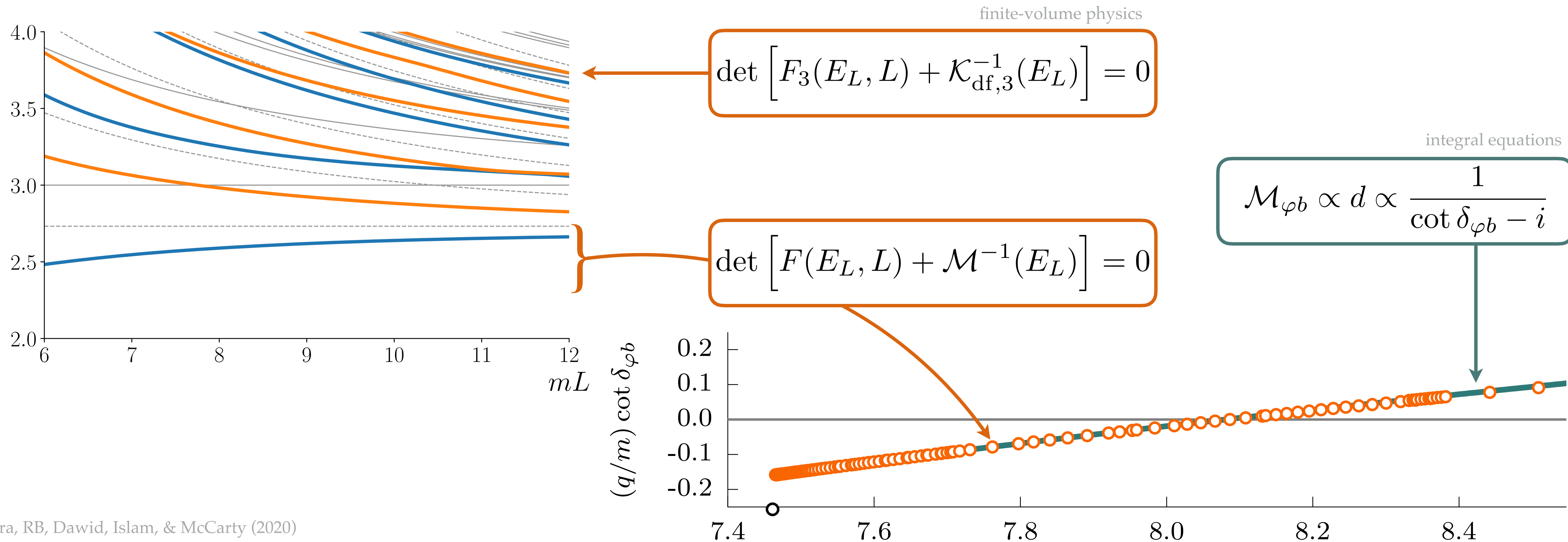
□ if $a > 0$, there is a bound state with binding energy $m_b = 2\sqrt{m^2 - \frac{1}{a^2}}$



Consistency checks for toy model

Alternatively to solving the integral equations, one can instead:

- Fix $ma = 2$ and $\mathcal{K}_{\text{df}} = 0$
- obtain E_L using 3-body quantization by Hansen & Sharpe (2014),
- obtain $q \cot \delta_{\varphi b}$ using the two-body Lüscher formalism.



Trimers

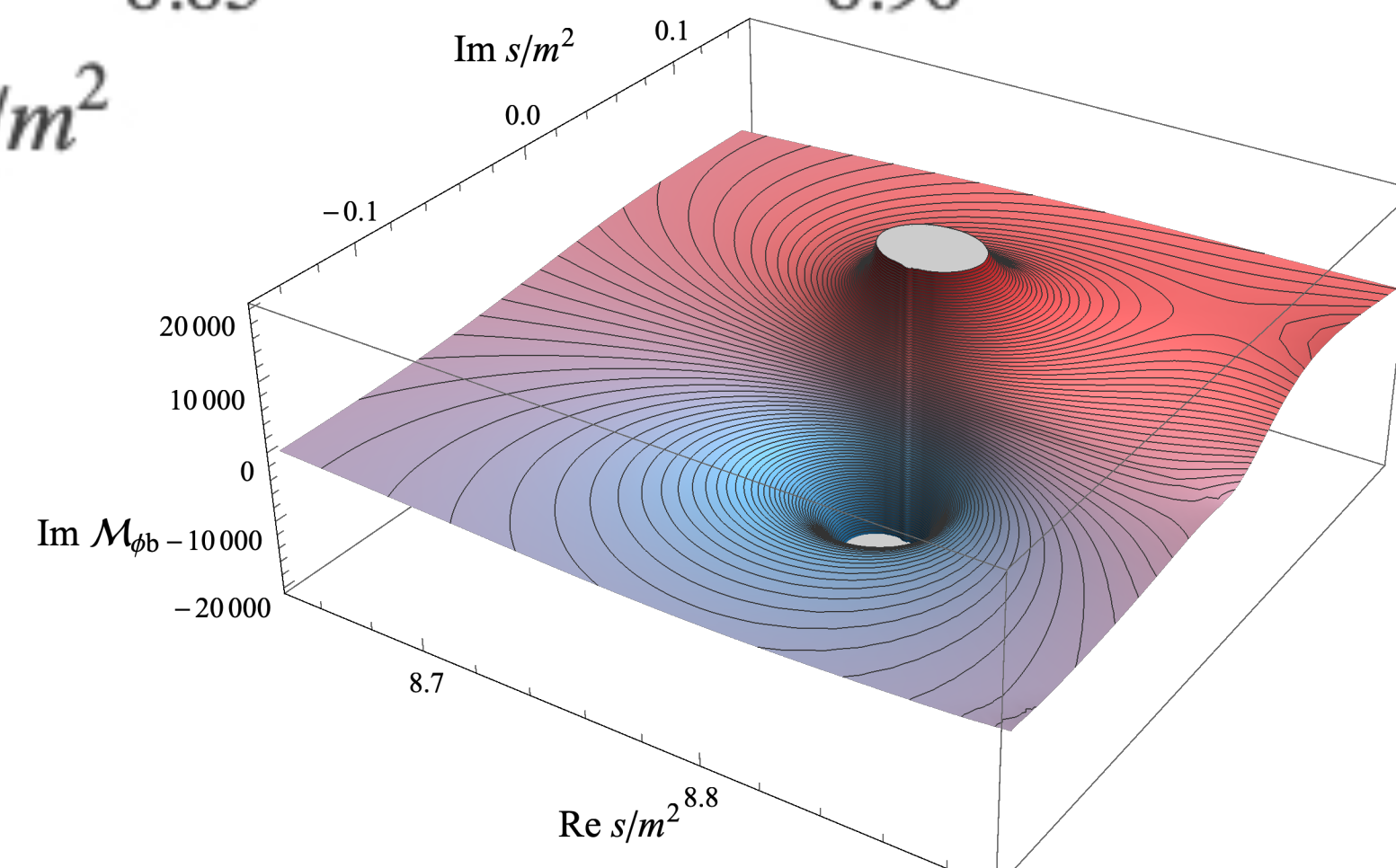
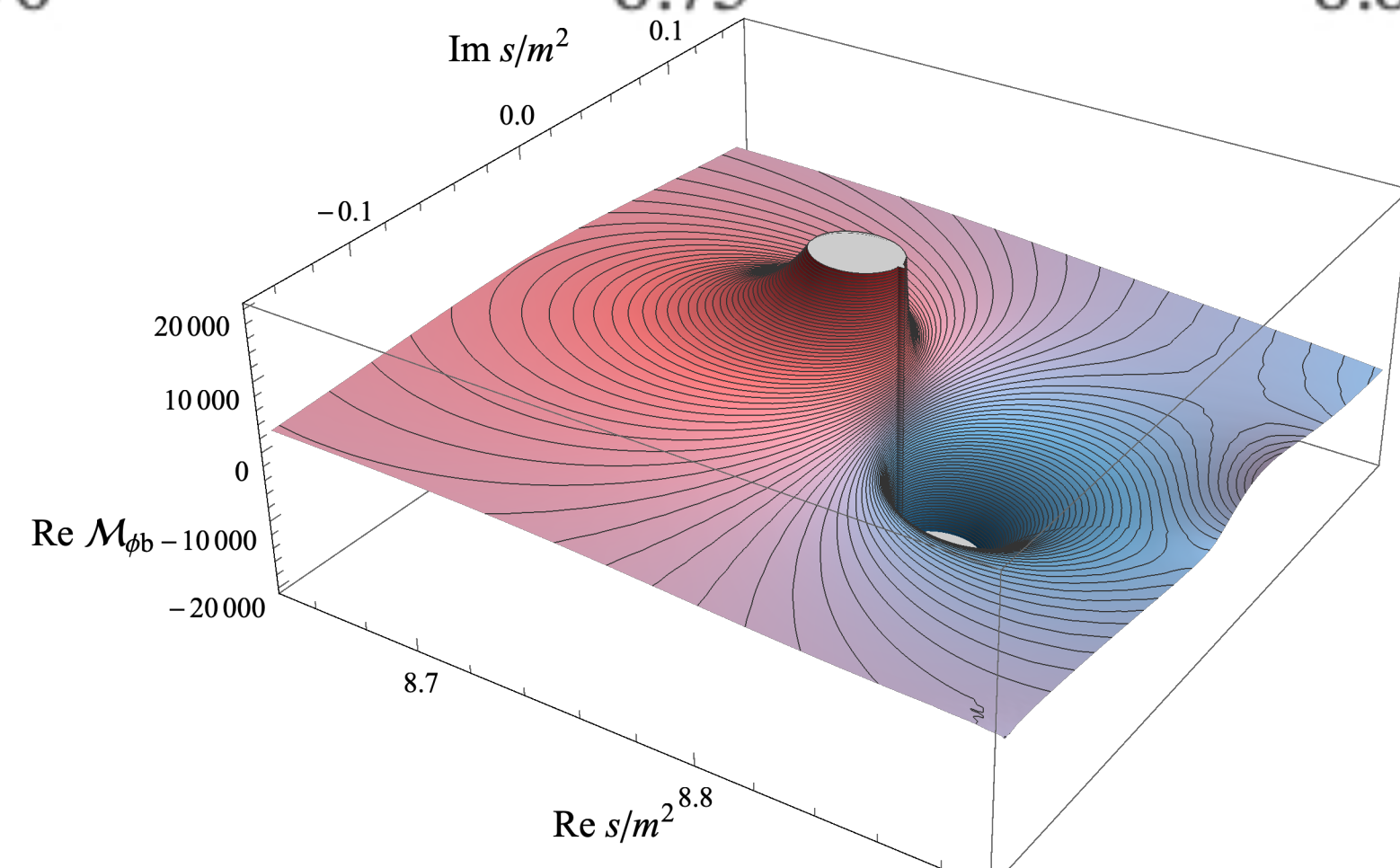
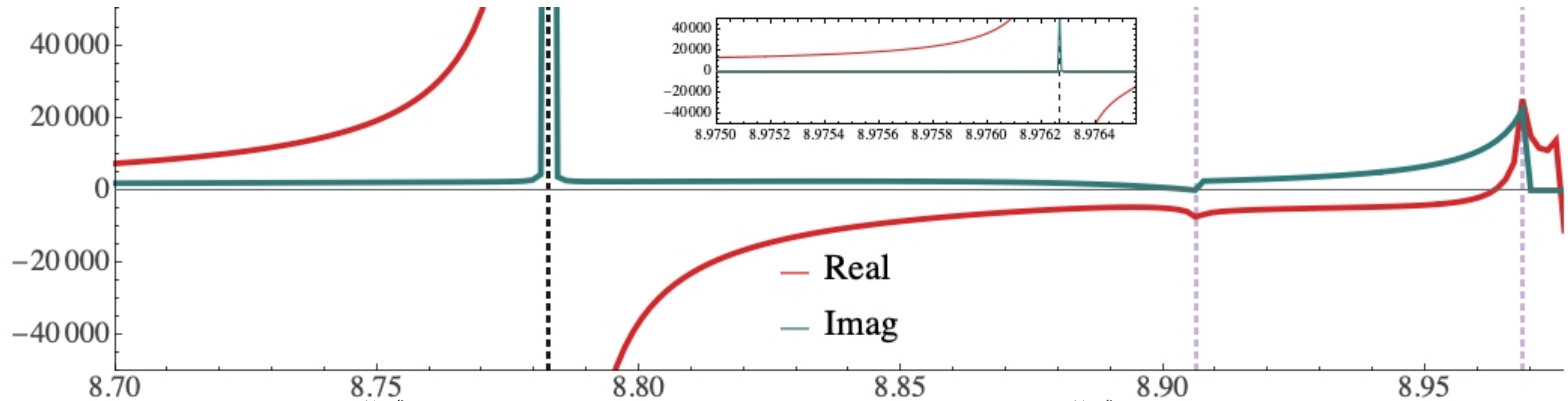
Scan for trimer poles as a function of energy

□ Fix $ma = 16$ and $\mathcal{K}_{df} = 0$



Dawid

Islam

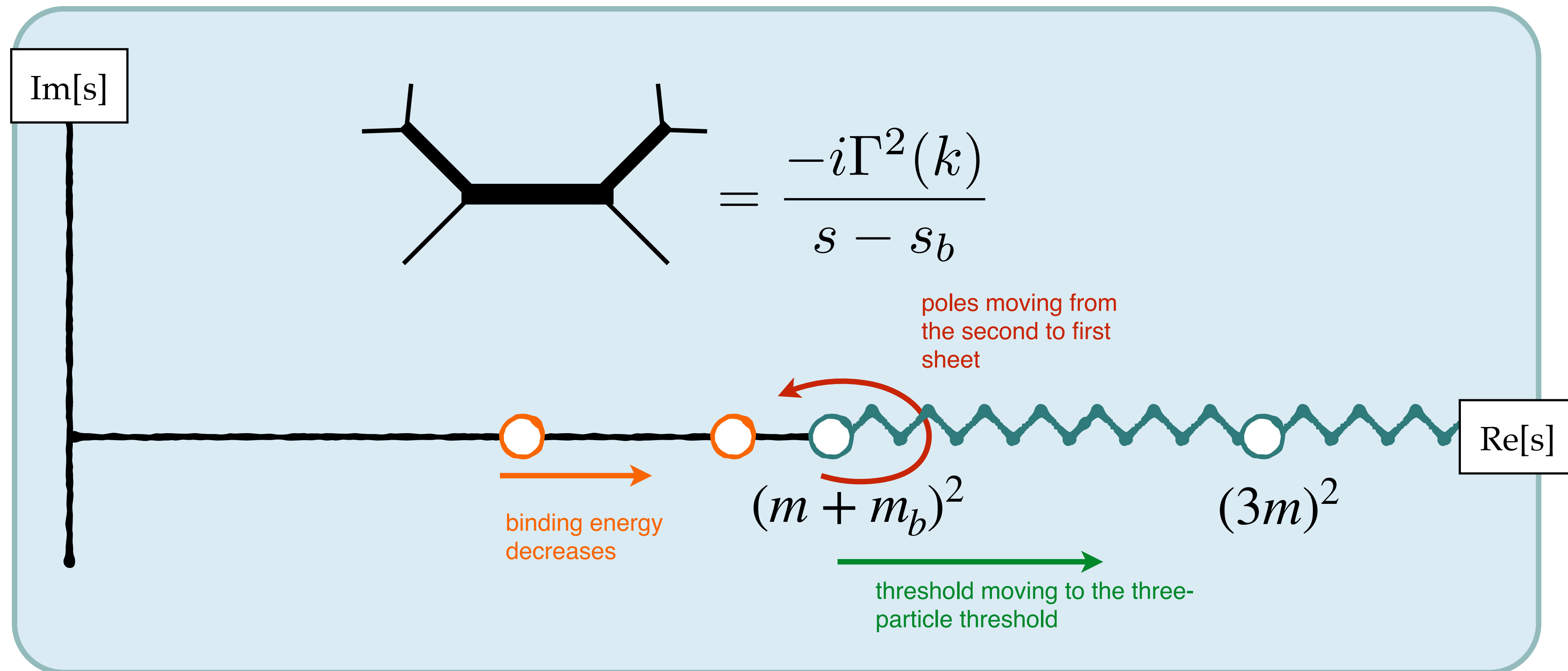


Recovering Efimov Physics

Remember $m_b = 2\sqrt{m^2 - \frac{1}{a^2}}$,

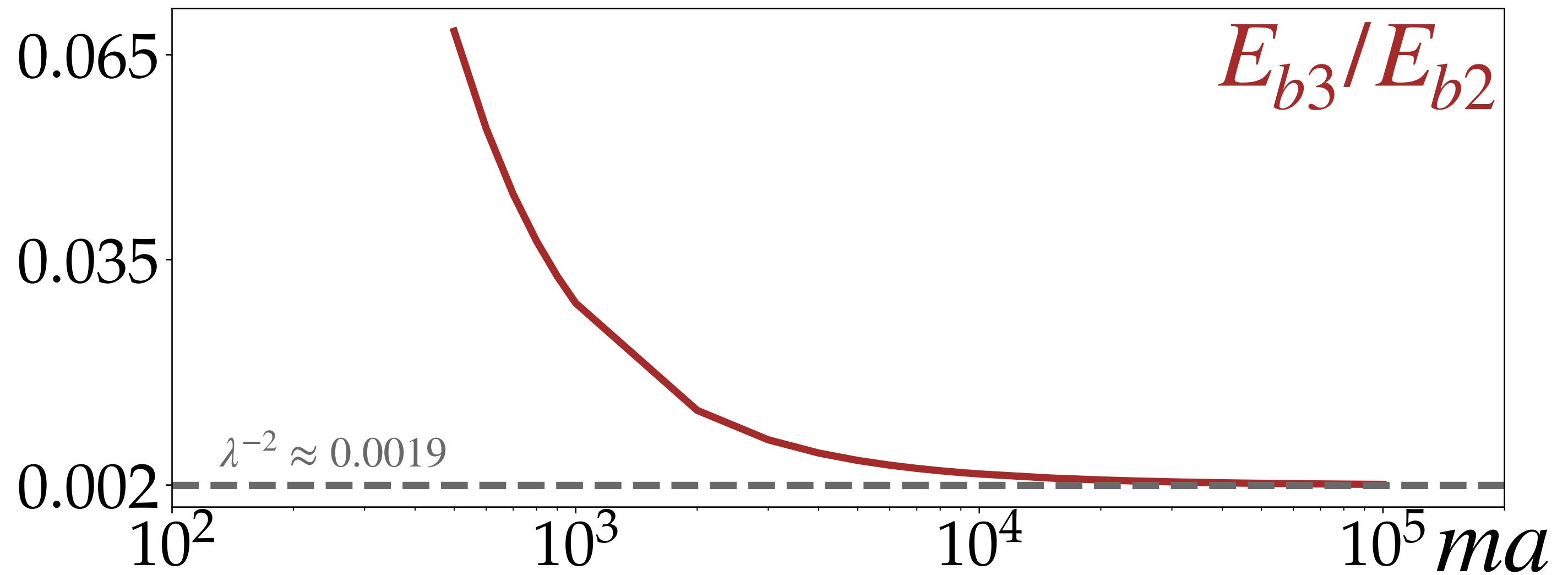
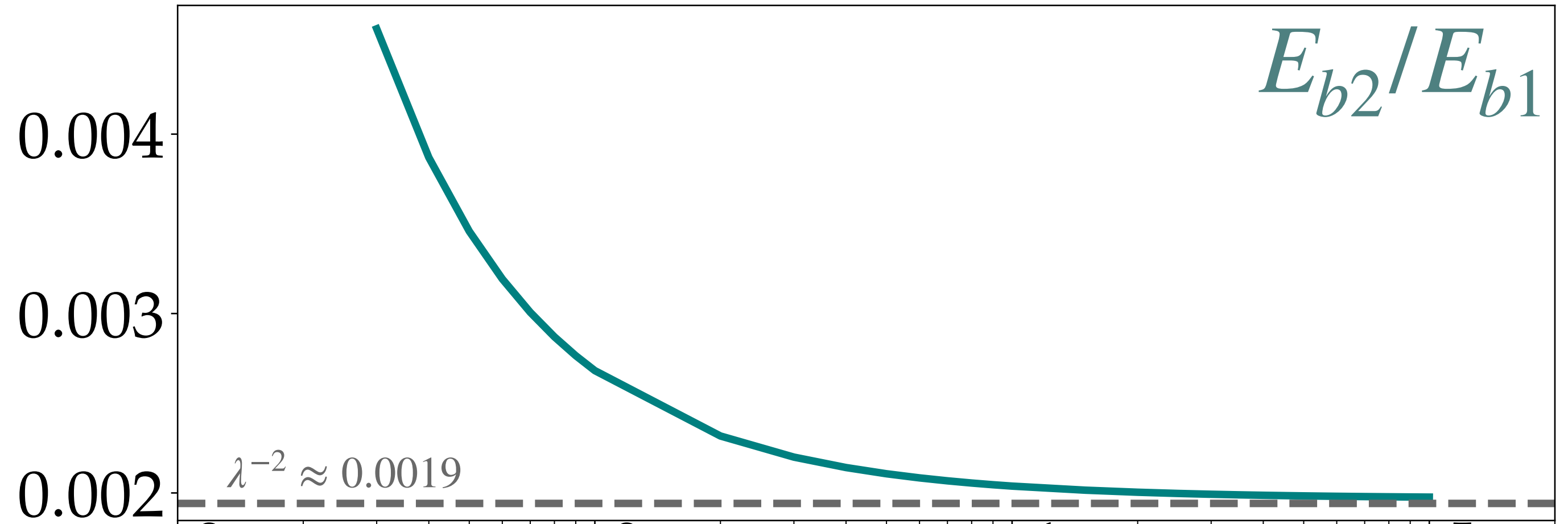
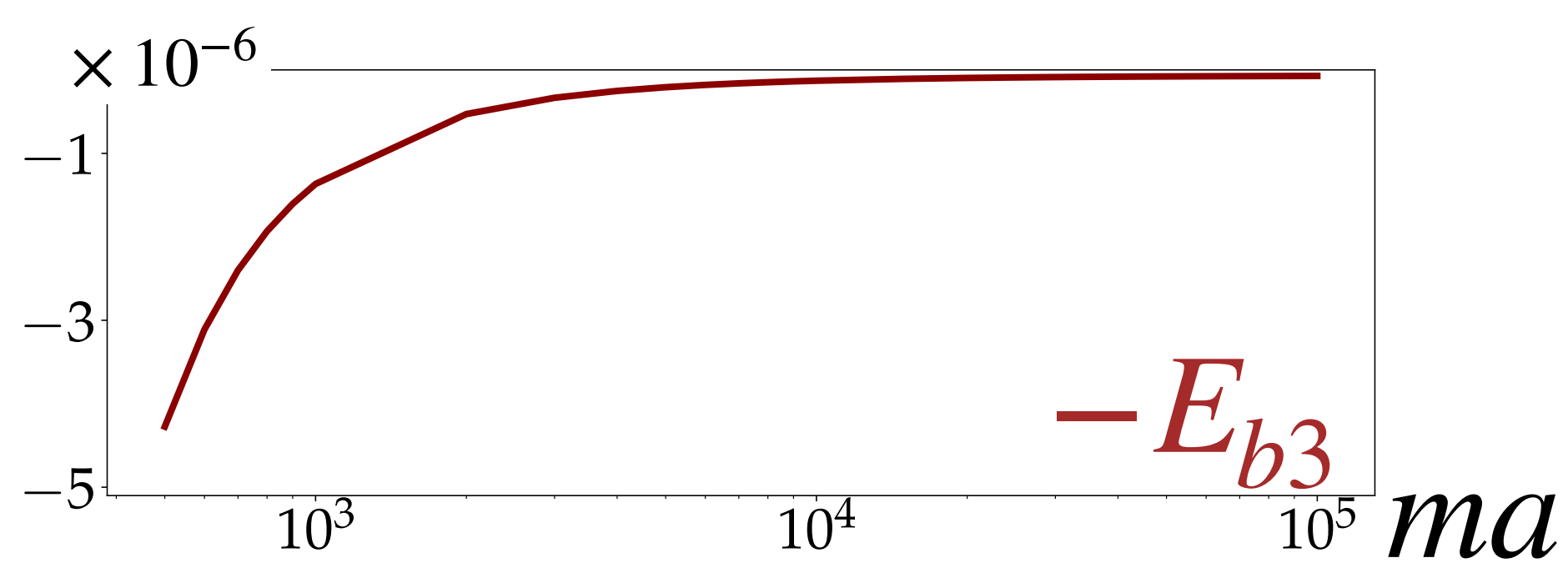
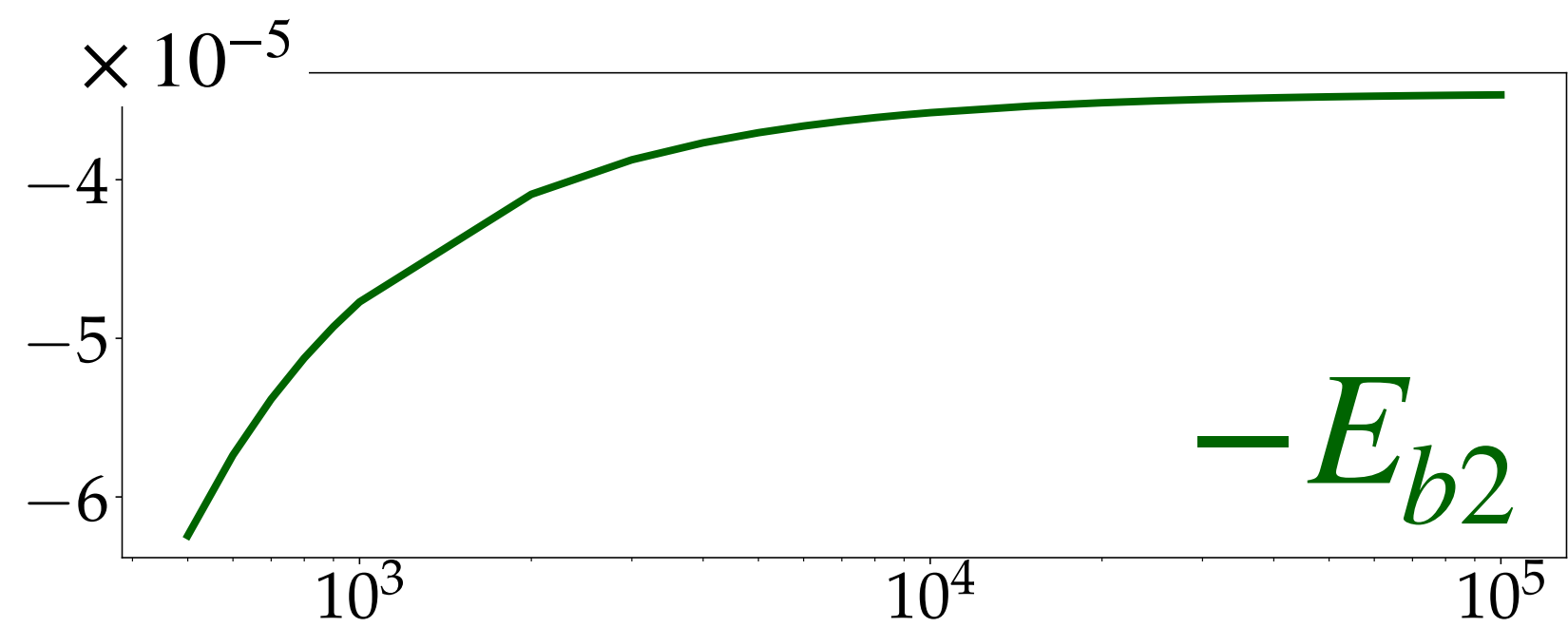
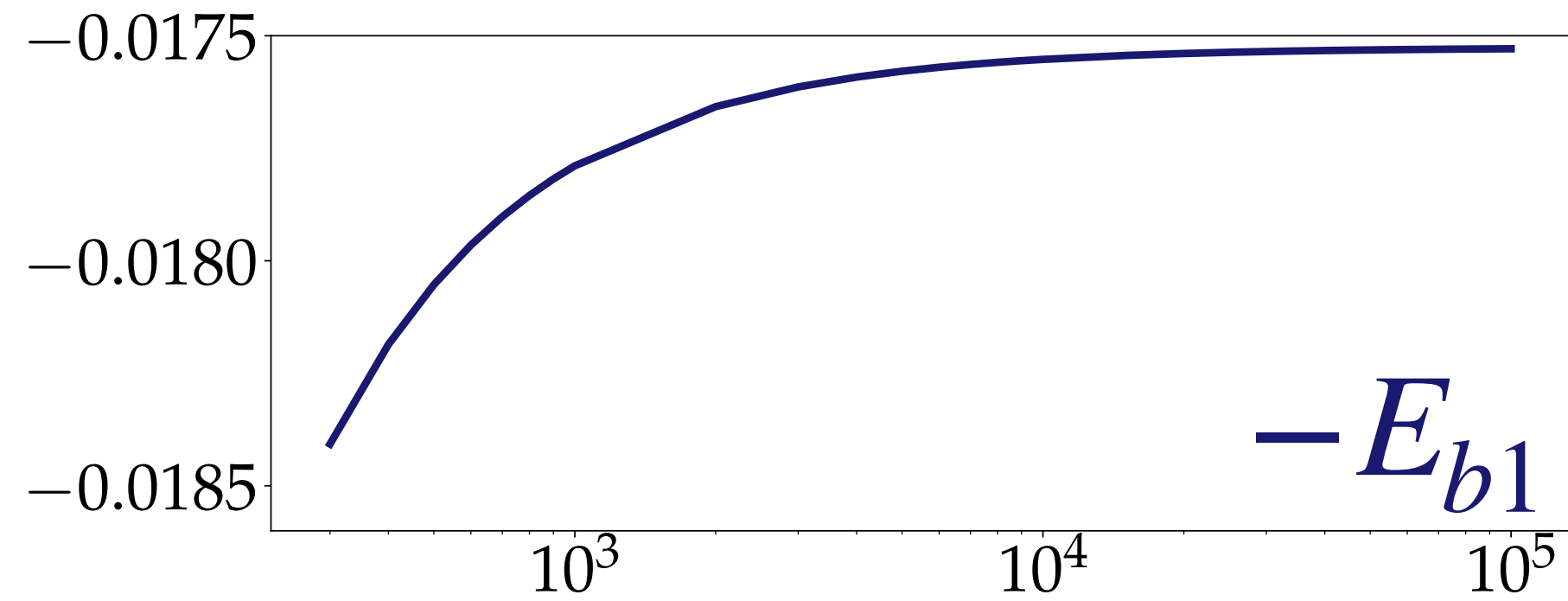
If $a \rightarrow \infty$, then $E_{N+1} = E_N/\lambda^2$ where $\lambda = 22.69438\dots$

But as $a \rightarrow \infty$ life gets more complicated!

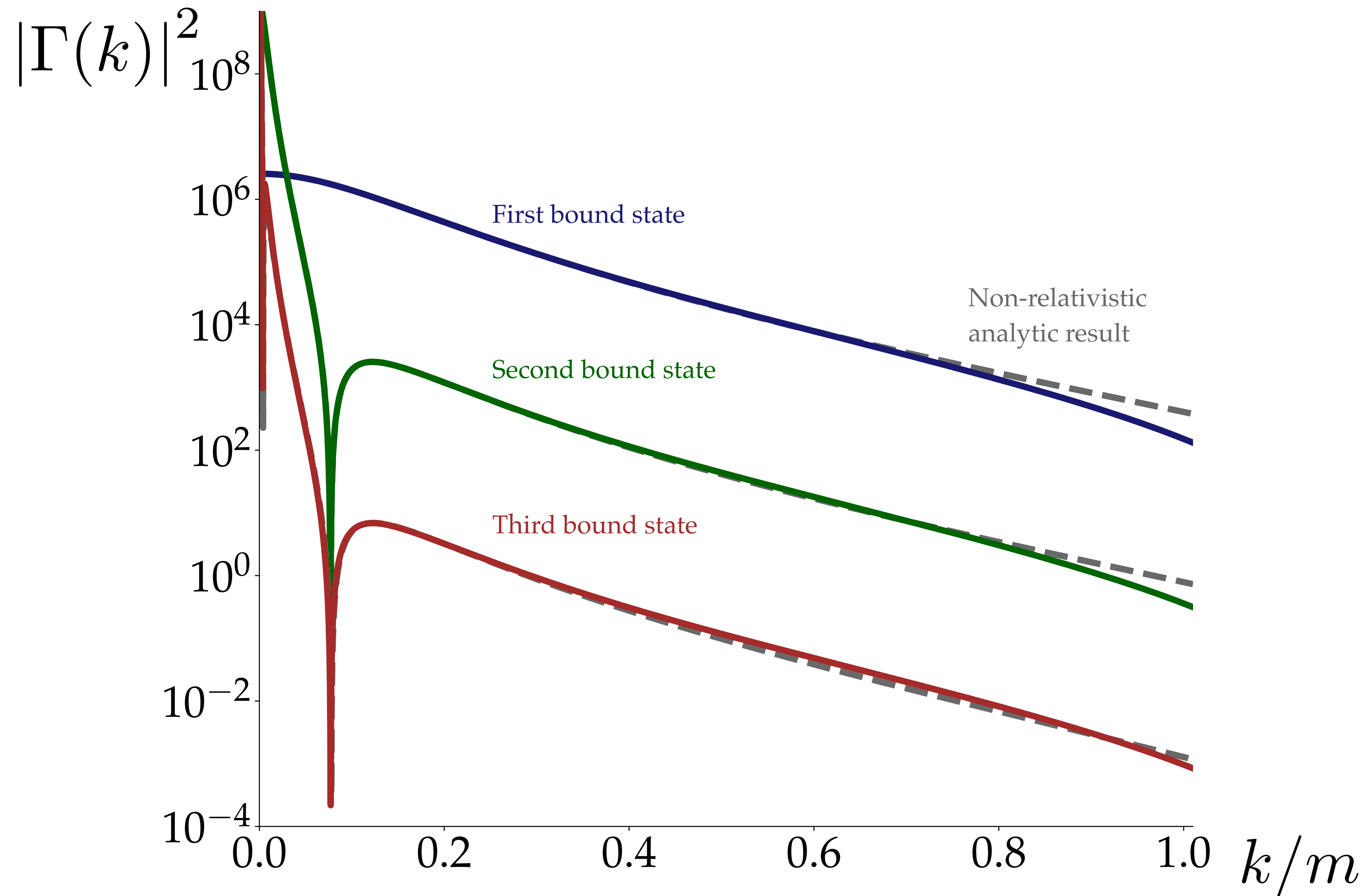


Recovering Efimov Physics

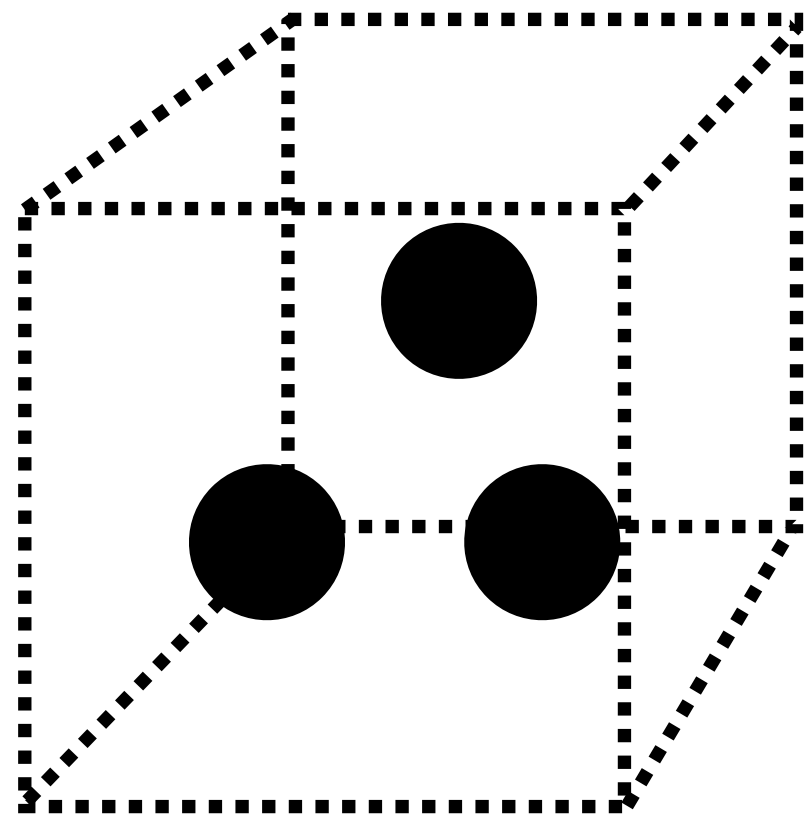
Binding energies



Recovering Efimov Physics *residues*



Back to lattice QCD

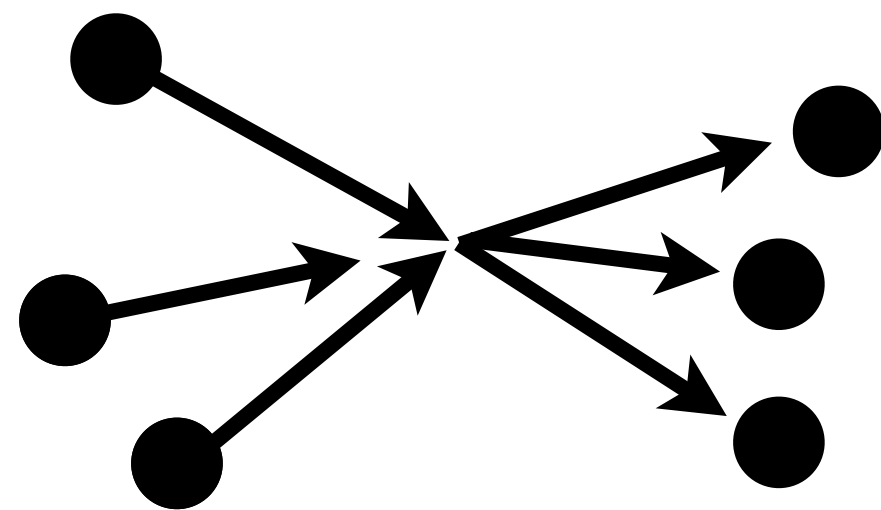


finite-volume
spectroscopy

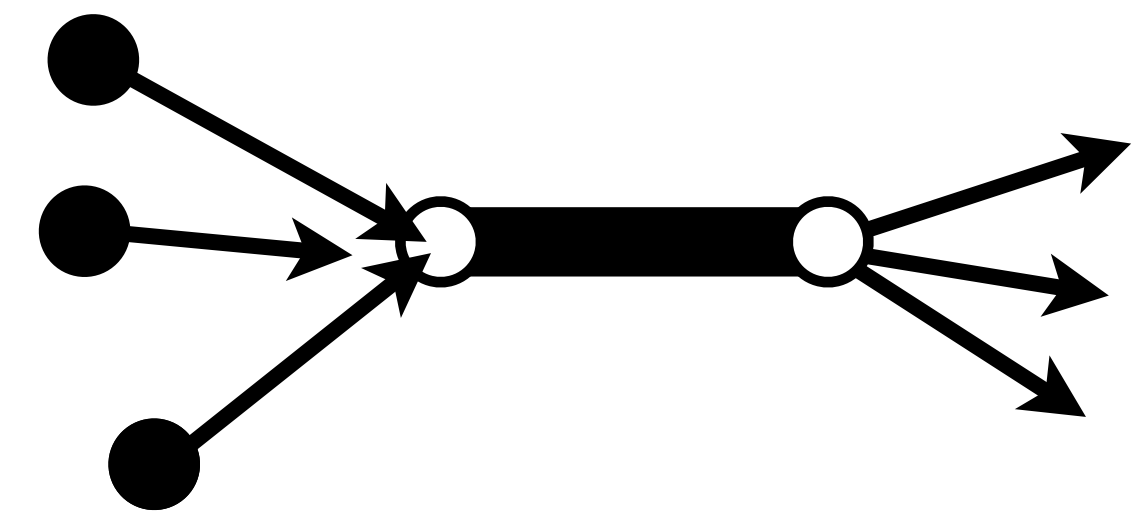
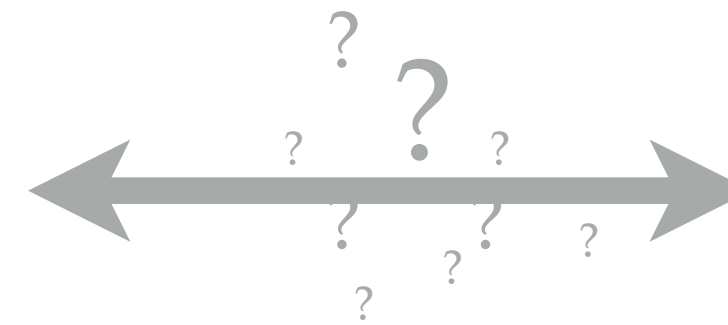
$$\det \left[F_3(E_L, L) + \mathcal{K}_{\text{df},3}^{-1}(E_L) \right] = 0$$



Hansen & Sharpe ('14, '15)
Mai & Döring ('17)
RB, Hansen & Sharpe ('18)
Hansen, Romero-Lopez & Sharpe ('20)
Blanton & Sharpe ('20)
Jackura ('22)



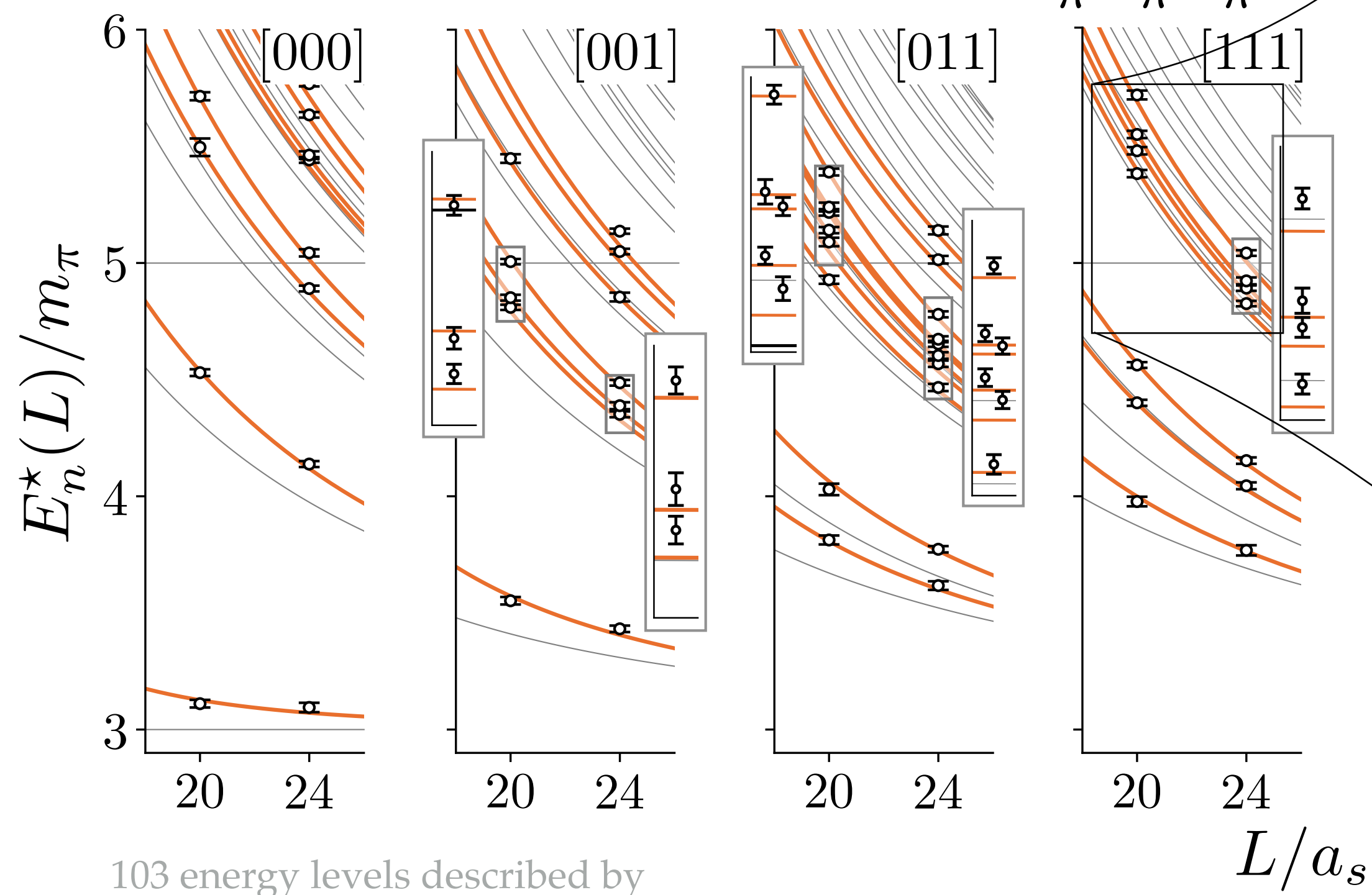
infinite-volume
scattering amplitudes



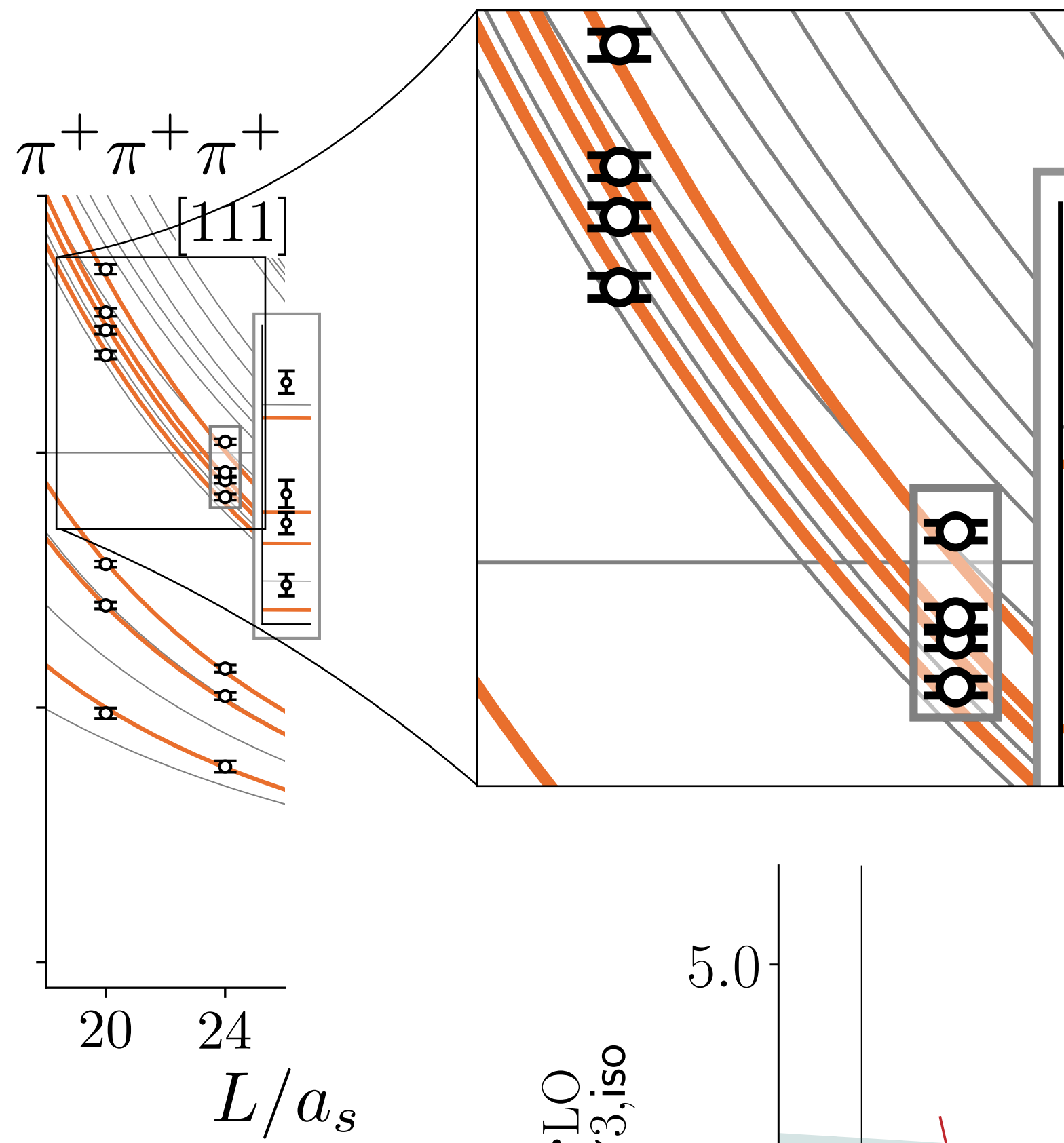
bound state and
resonance poles

$\pi\pi\pi$

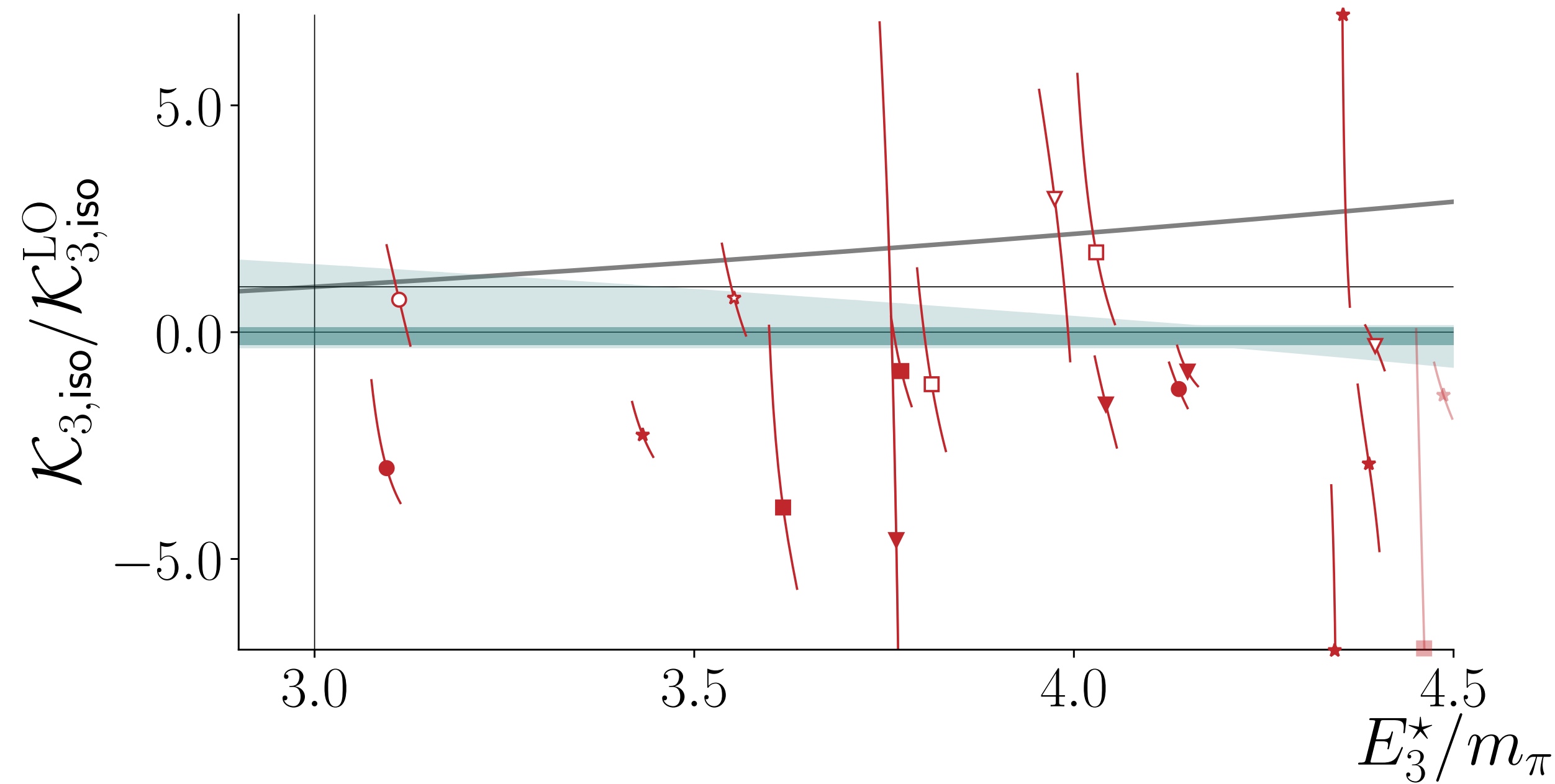
($l=3$ channel, $m_\pi \sim 390$ MeV)



103 energy levels described by three numbers: m_π , $a_{\pi\pi}$, $\mathcal{K}_{3,\text{iso}}$



$$F_{3,\text{iso}}^{-1}(P, L) + \mathcal{K}_{3,\text{iso}}(P^2) = 0$$

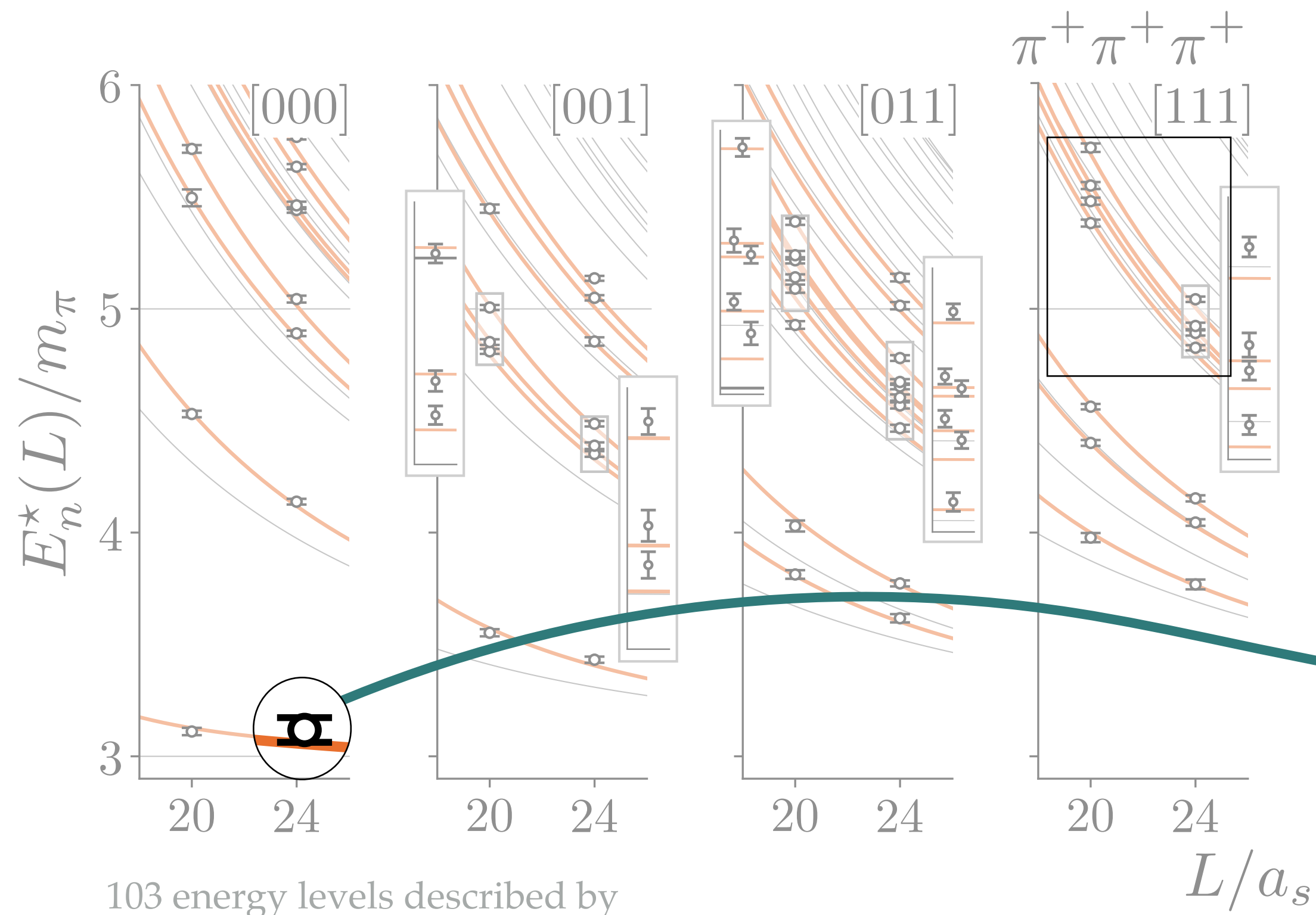


$m_\pi \sim 390$ MeV

Hansen, RB, Edwards, Thomas, & Wilson (2020)

$\pi\pi\pi$

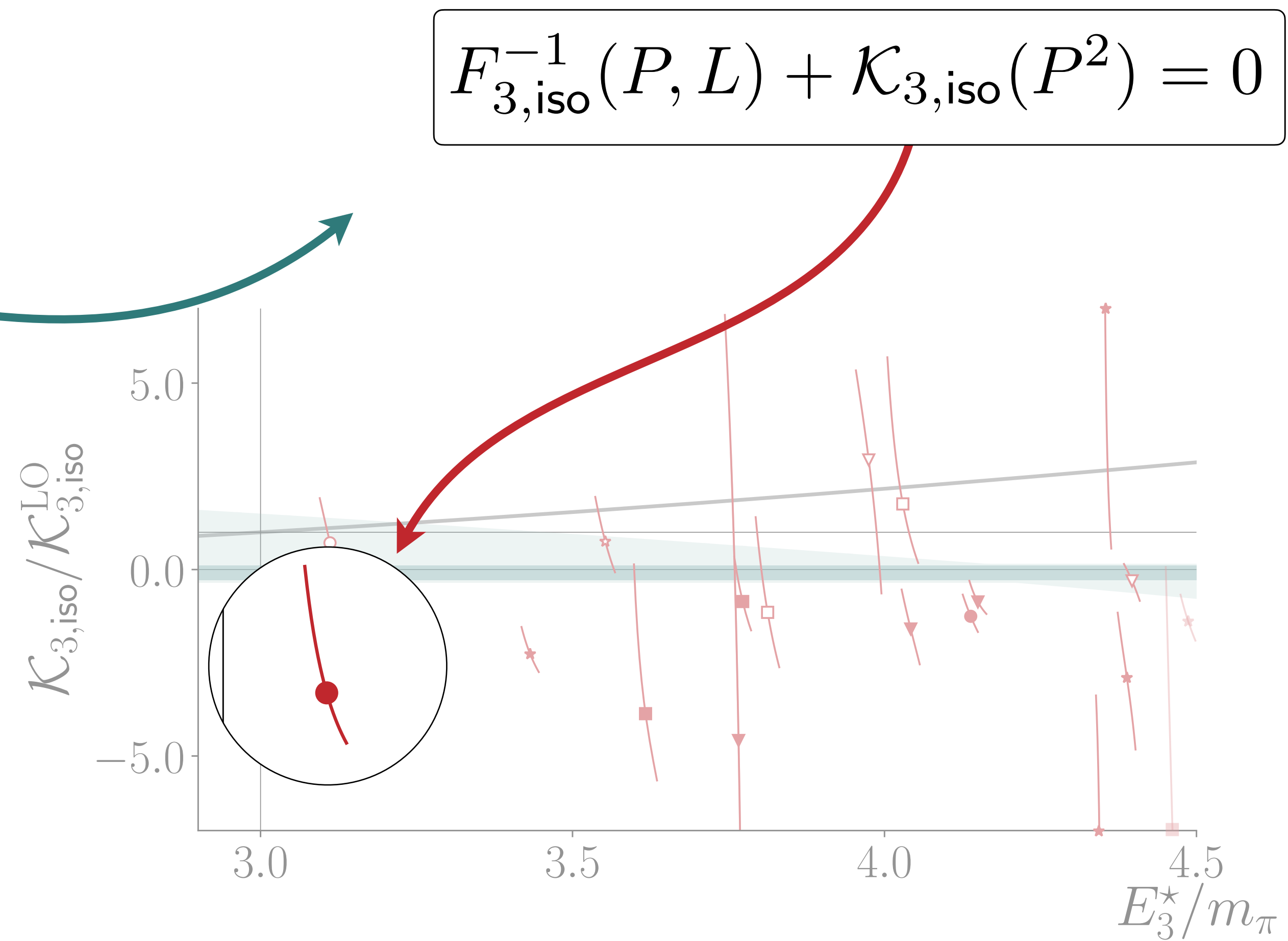
($l=3$ channel, $m_\pi \sim 390$ MeV)



103 energy levels described by three numbers: m_π , $a_{\pi\pi}$, $\mathcal{K}_{3,\text{iso}}$

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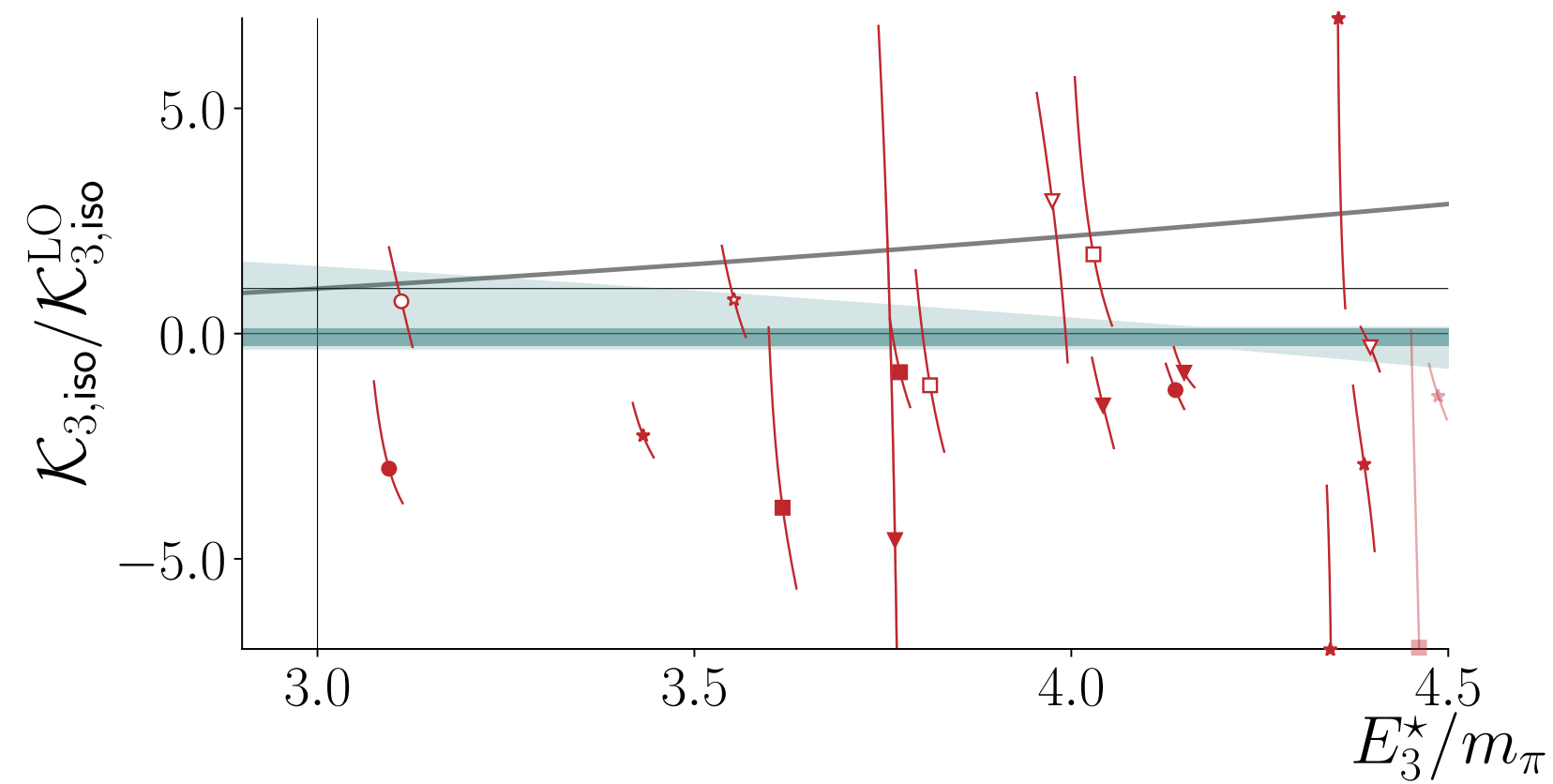
Hansen, RB, Edwards, Thomas, & Wilson (2020)



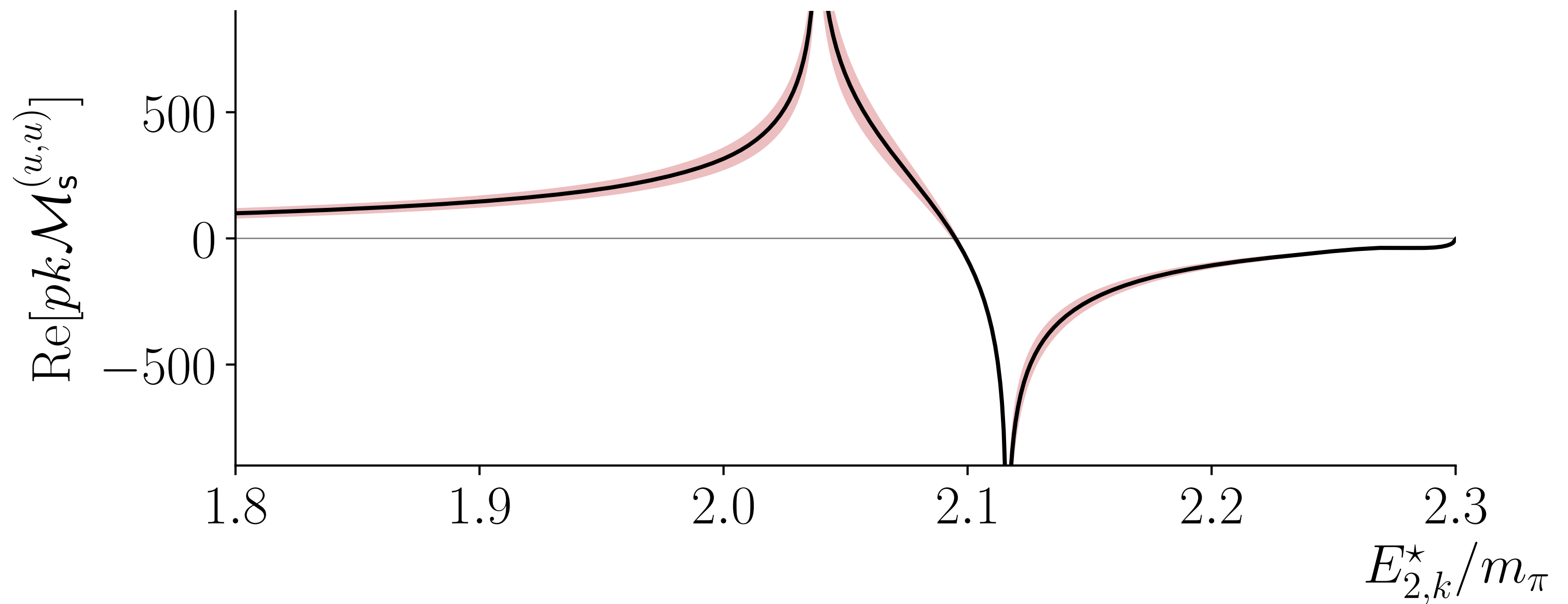
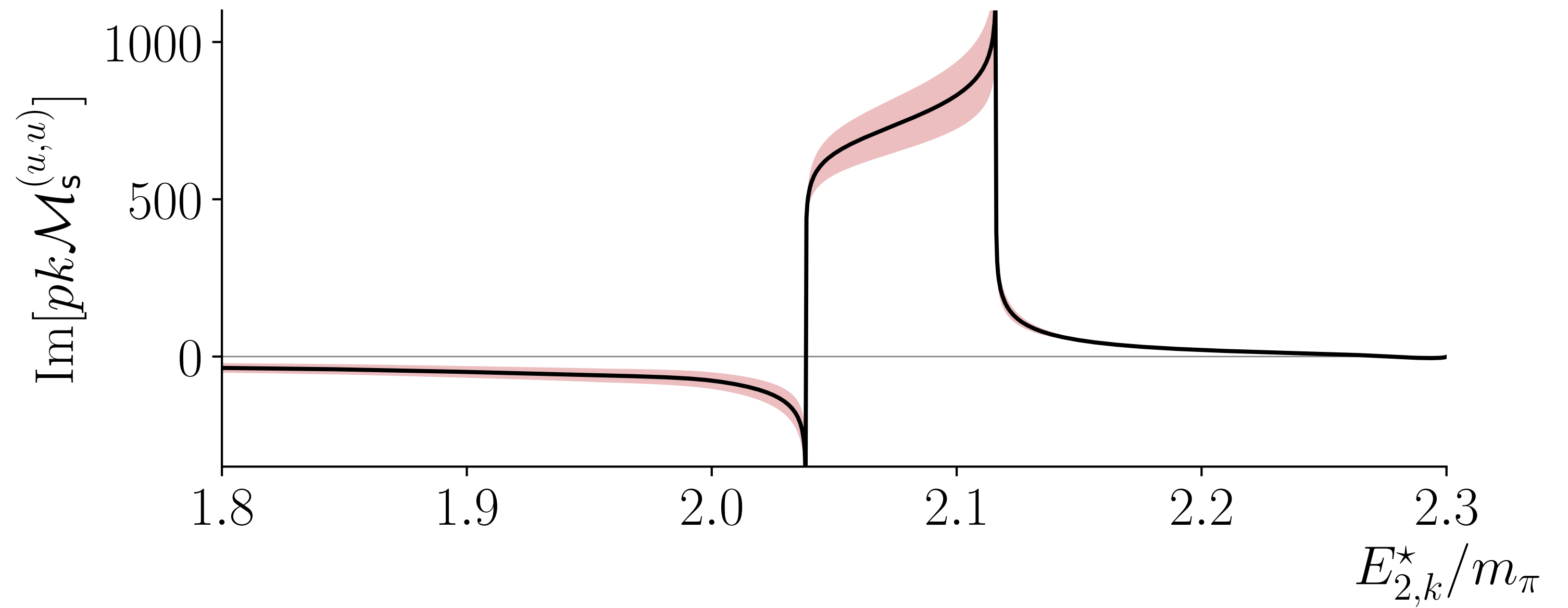
$\pi\pi\pi$ scattering

($l=3$ channel, $m_\pi \sim 390$ MeV)

first 3body scattering amplitude from the lattice QCD!



$$i\mathcal{D} + i\mathcal{L} \frac{1}{\mathcal{K}_{3,\text{df}}^{-1} + F_3^\infty} \mathcal{L}$$

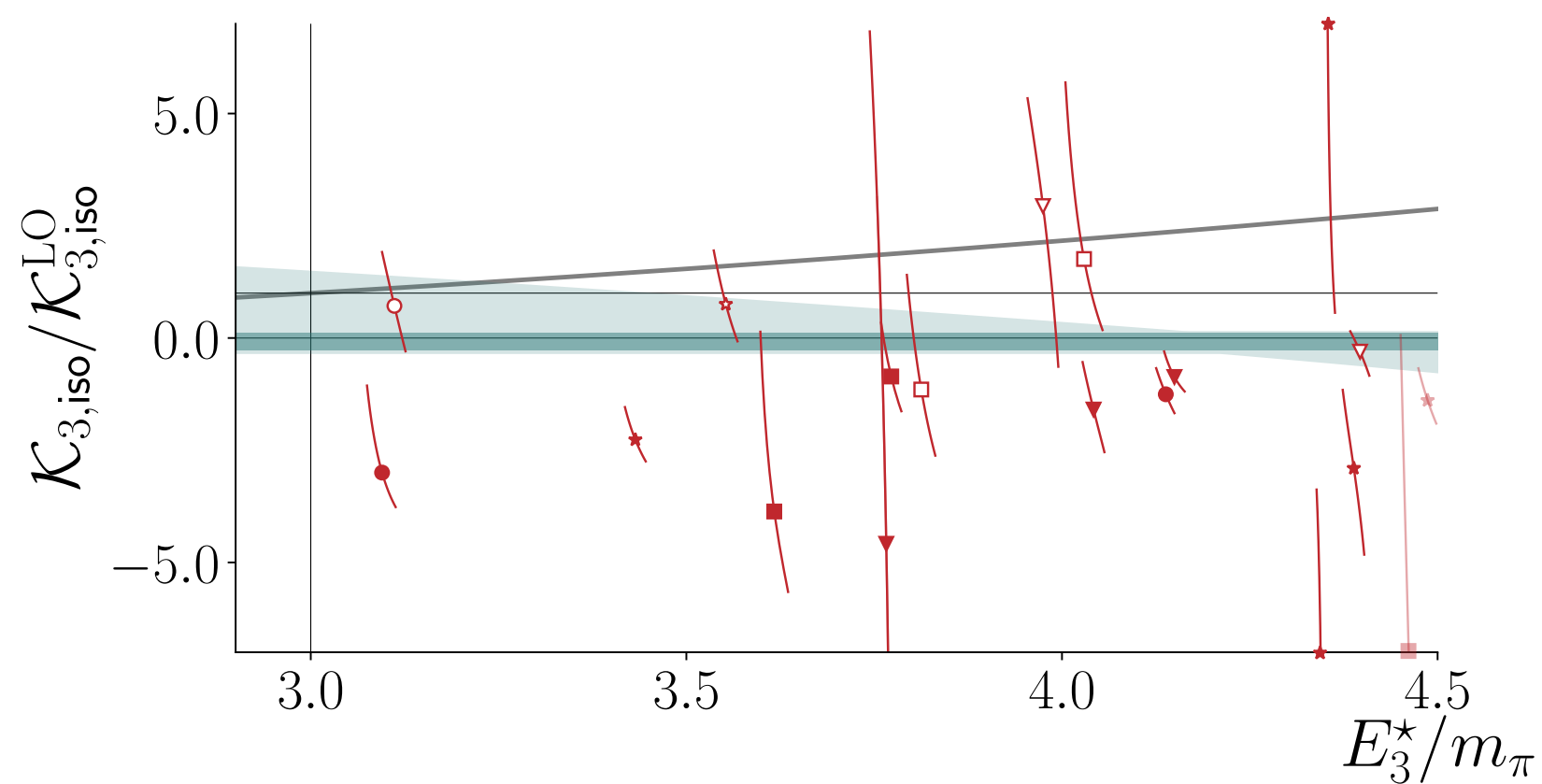


$m_\pi \sim 390$ MeV

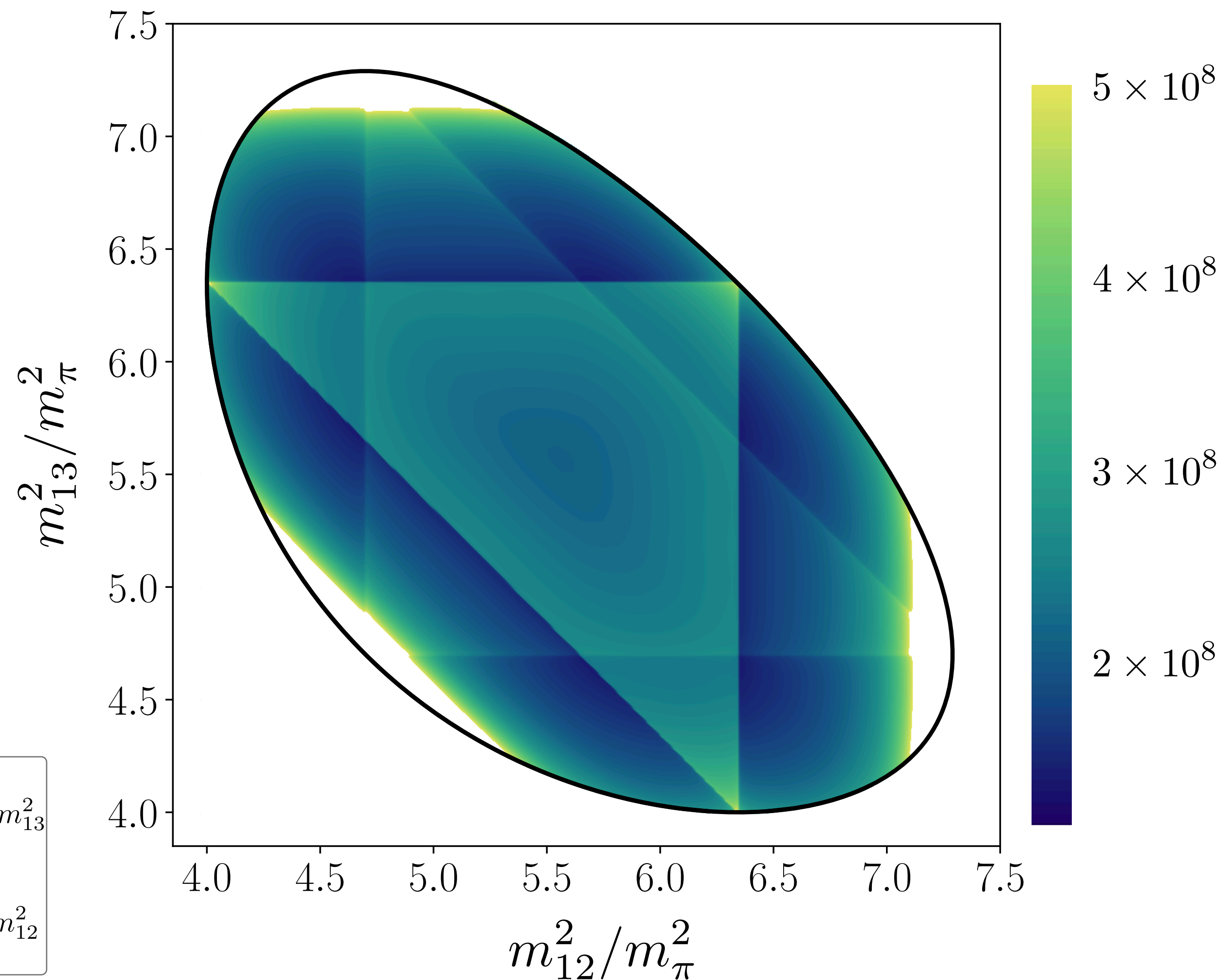
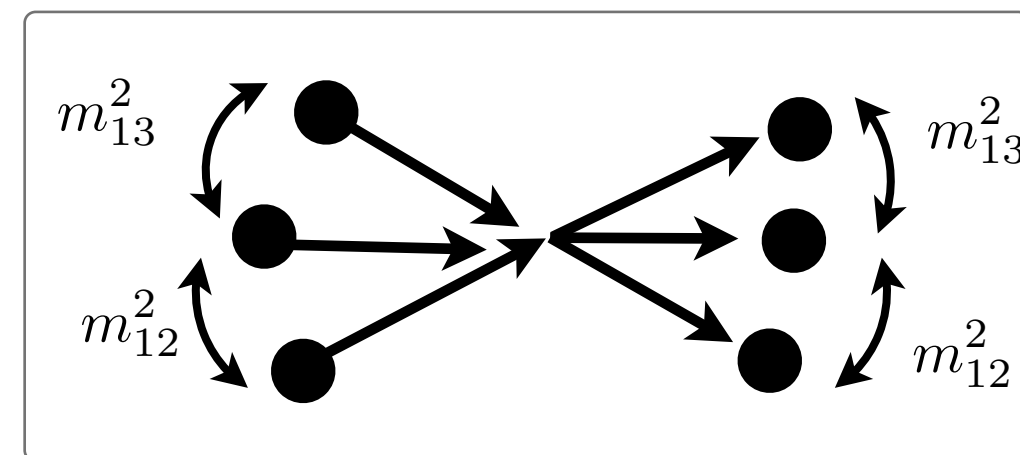
$\pi\pi\pi$ scattering

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Hansen, RB, Edwards, Thomas, & Wilson (2020)

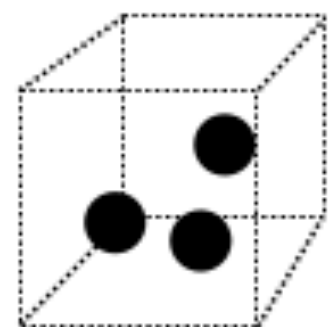
Looking back at 2019!

THE FUTURE IS OURS TO CREATE.

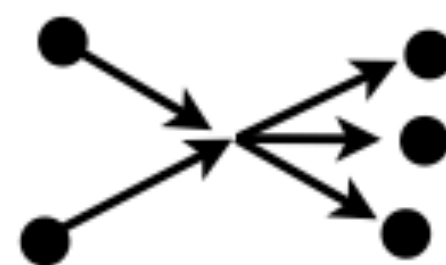
Unitarity in
3Body systems



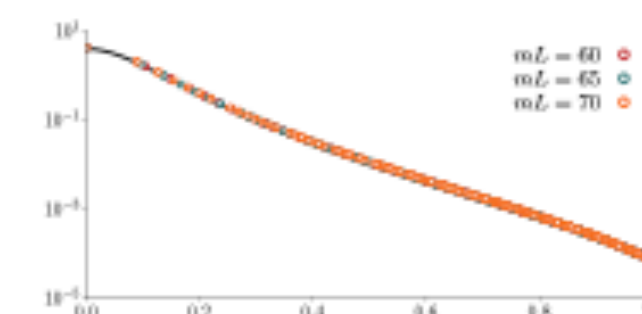
FV formalism for
spinless states



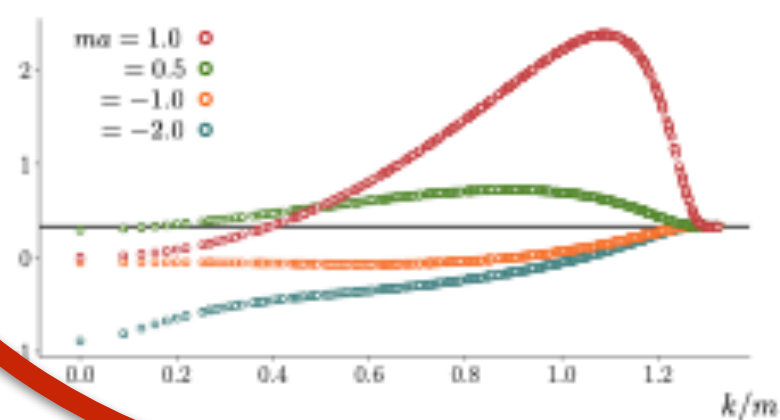
coupled 2/3B
systems



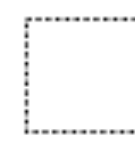
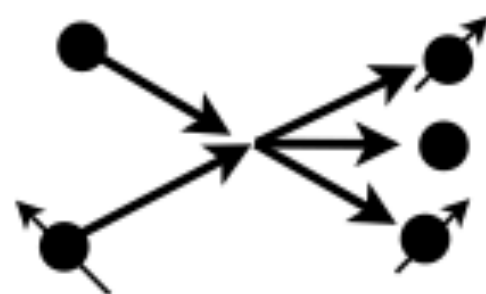
analytic and
numerical checks



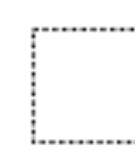
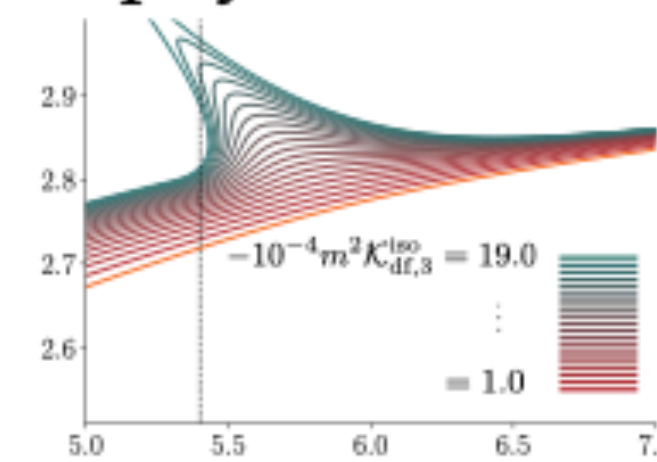
integral equations



Spin, non-degenerate
masses, multichannels



unphysical solutions



do some calculations
already...

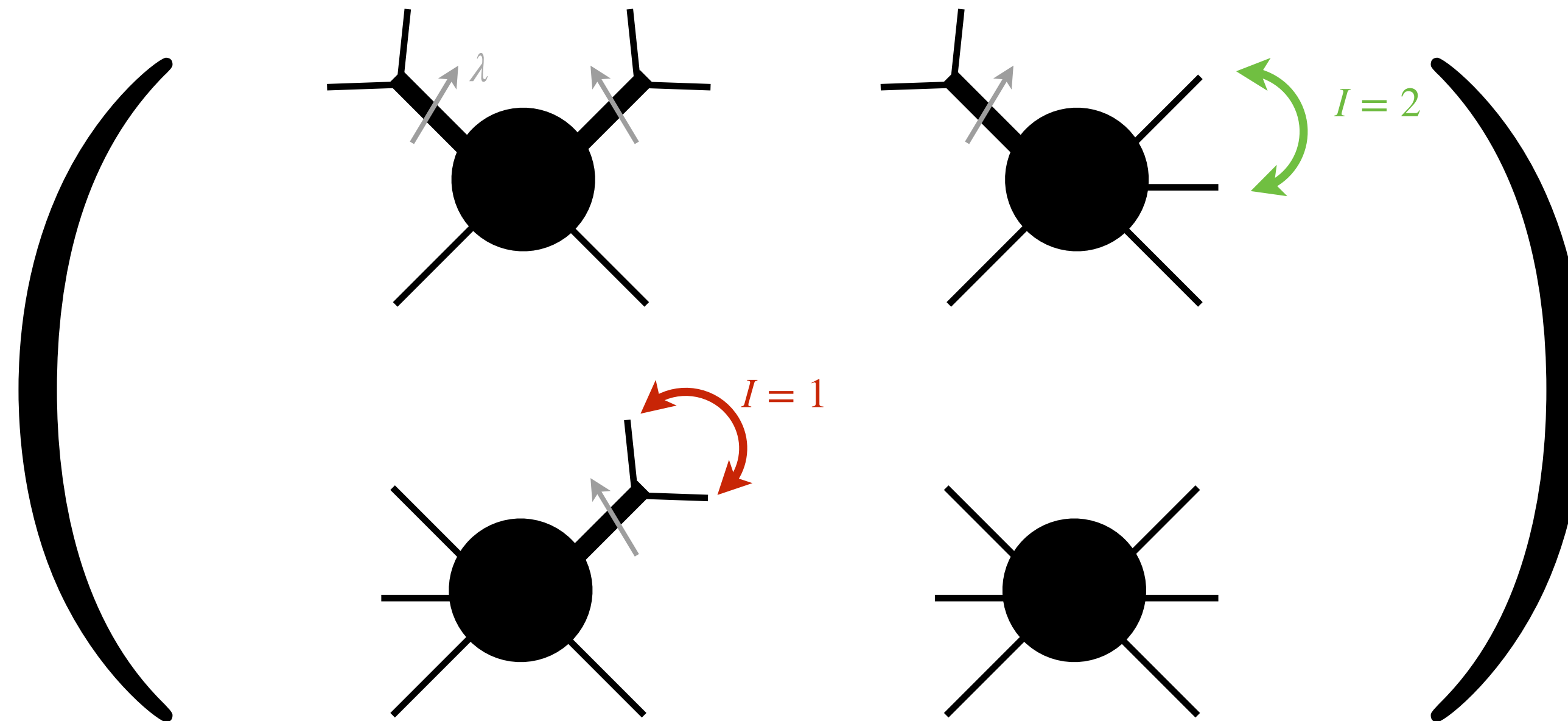


Looking into the future

- Lots of formal developments.
- Lattice QCD calculations are catching up.

Next obvious place to explore is the $I = 2$ channel, which is not so trivial.

- Coupled channels $[\pi(\pi\pi)_{I=2}, \pi(\pi\pi)_{I=1}]$,
- non-zero angular momentum,
- two-body resonances with non-zero helicity.

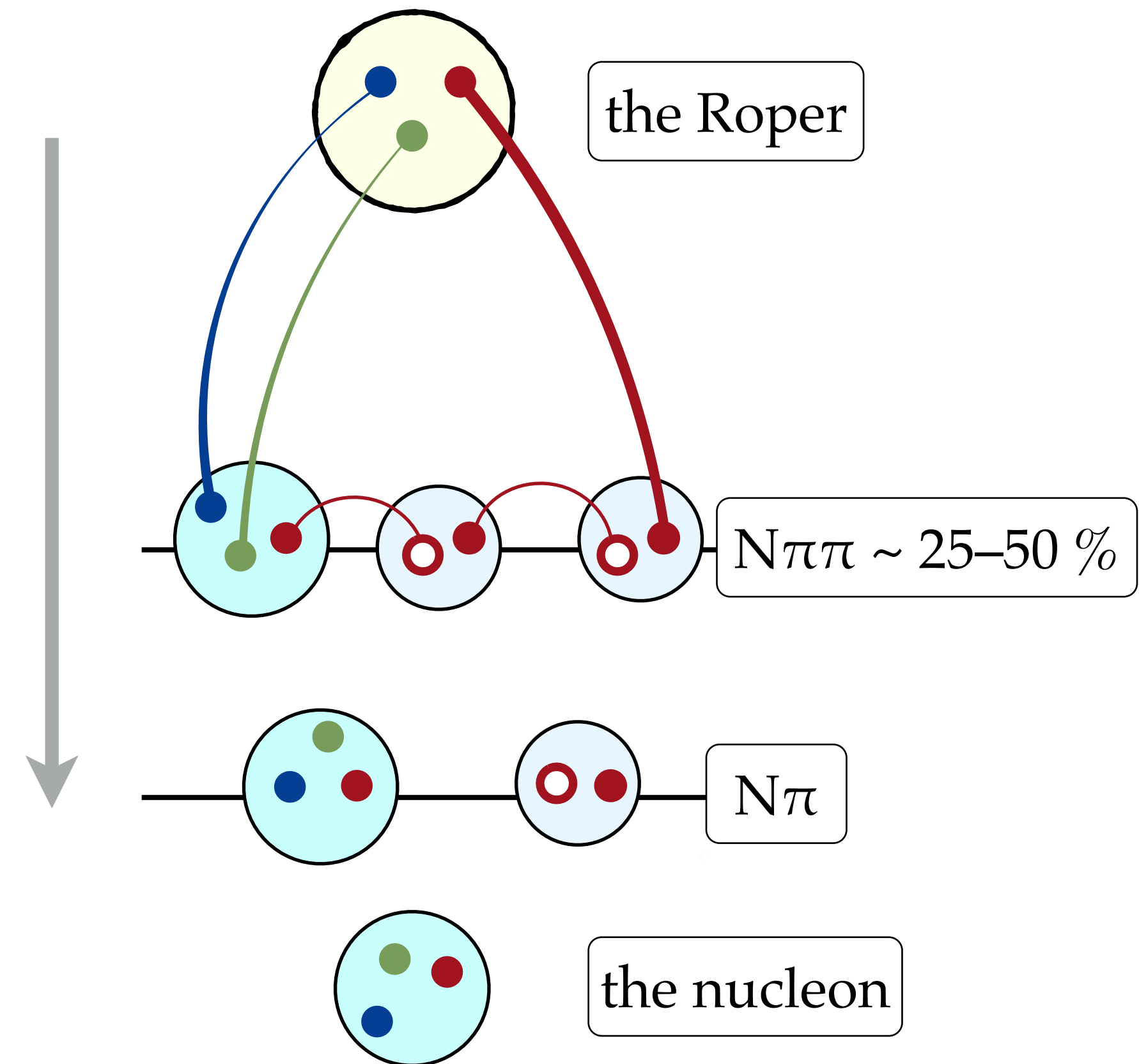


Looking into the future

Outstanding formal questions:

- ! non-zero orbital angular momentum,
- ! analytic continuation into the complex plane,
- ! coupled 2-3 bodies,
- ! non-identical particles,
- ! electroweak production,
- electroweak probes,
- ! long-range processes,
- non-zero intrinsic spin.

$\mathcal{O}(10^{-23}\text{s})$



Modern-day few-body physics

Nature of states, patterns in the spectrum, ...

Analytic continuation: poles, widths, form factors, ...

Scattering theory & EFTs

Amplitudes on the real axis

Finite-volume formalism

QCD Lagrangian & electroweak probes, Lattice QCD

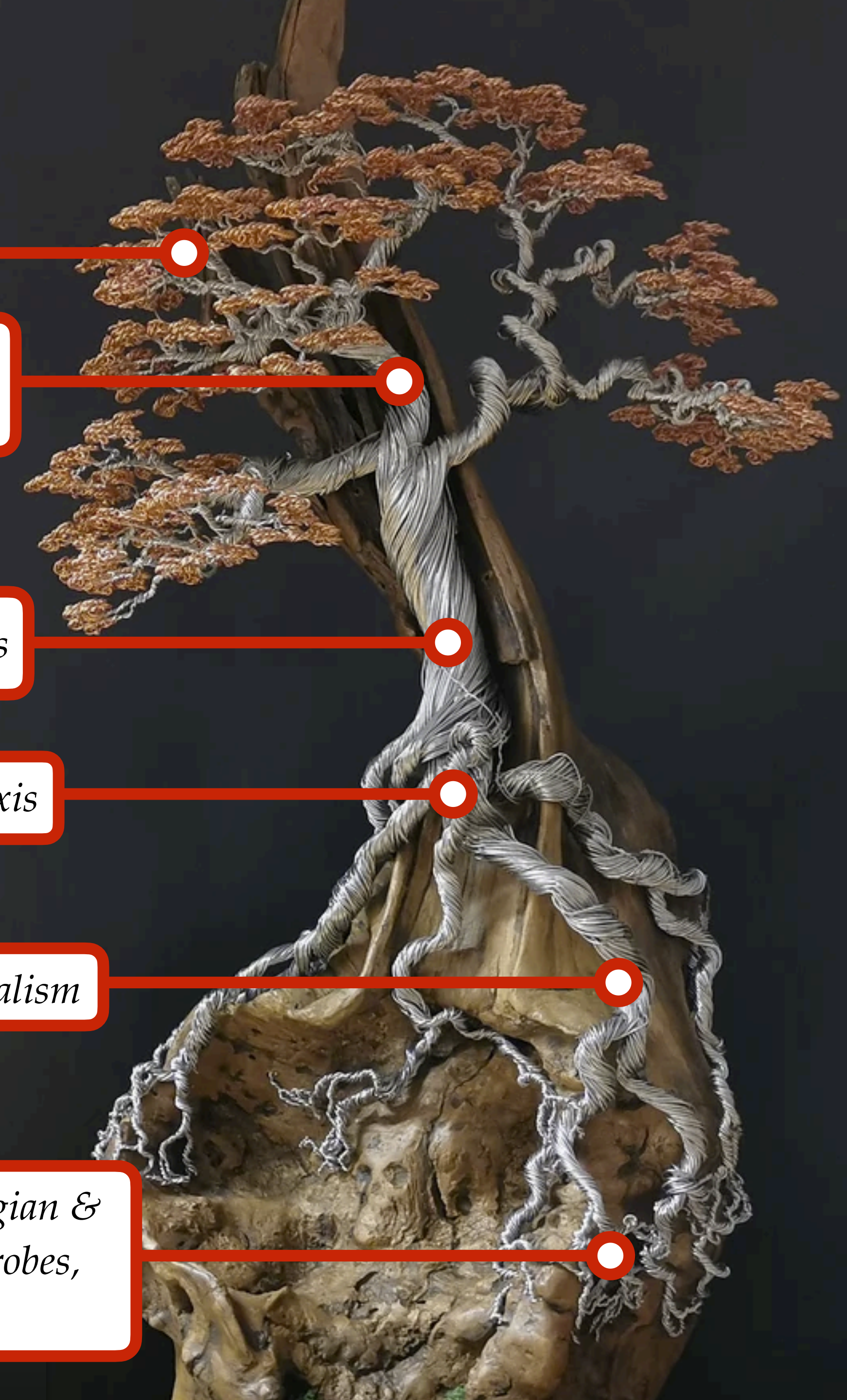
Challenging but not impossible

- Enhancement of operator basis,
 - Tetraquarks, glueballs, pentaquarks, three-particles,

- Contraction costs,
 - Cost grows with the number of external legs...

- Extensions of finite- and infinite-volume formalism,
 - Multi-channel 3-body, spin, etc.
 - Electroweak processes, ...

- Scattering theory,
 - "Earnest amplitudes", analytic continuations, ...



Lots of great memories!

"friendliest group!"

