

Evidence for intrinsic charm quarks in the proton [Nature608.483]

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INTRINSIC CEEVROLETS AT THE SSC

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Summary

The possibility of the production at high energy of heavy quarks, supersymmetric particles and other large mass colored systems via the intrunic twist-six components in the proton wave function is discussed. While the existing data do not rule out the possible relevance of intrinsic charm production at present energies, the extrapolation of such intrinsic contributions to very high masses and energies suggests that they will not play an important role at the SSC. sufficiently large. The data from the EMC collaboration⁴ on deep-inelastic muon scattering could also be intepreted as suggesting an unexpectedly large charm structure function in the region z > 0.3.

The possible existence of such a new production mechanism is of great importance for design considerations at the SSC^{5,4}. An example of the importance of this issue is that, if intrinsic large x production is dominant, experiments and, perhaps, even the machine should be designed to focus on the forward "difficrative" regime⁶. The que-

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ature 608, 483–487 (2022) <u>Cite this article</u> 4k Accesses 11 Citations 361 Altmetric <u>Metrics</u>	Do protons have intrinsic charm? New evidence suggests yes Benjamin Thompson & Nick Petrić Howe Nature Nature Podcast 17 Aug 2022		
he theory of the strong force, quantum chromodynamics, describes the proton in terms of uarks and gluons. The proton is a state of two up quarks and one down quark bound by luons, but quantum theory predicts that in addition there is an infinite number of quark- ntiquark pairs. Both light and heavy quarks, whose mass is respectively smaller or bigger	Evidence at last that the proton has intrinsic charm Ramona Vogt Nature News & Views 17 Aug 2022		
an the mass of the proton, are revealed inside the proton in high-energy collisions. owever, it is unclear whether heavy quarks also exist as a part of the proton wavefunction, hich is determined by non-perturbative dynamics and accordingly unknown: so called trinsic heavy quarks ¹ . It has been argued for a long time that the proton could have a sizable trinsic heavy courts of the light heavy end the heavy mark of the trinsic heavy shows the size t	Sections Figures References Abstract Main Methods		

Nature 2022 in the media



0. Introduction

- 1. NNPDF4.0 [EPJC82.428]
- 2. Intrinsic Charm [Nature608.483]

3. Summary

1. NNPDF4.0 [EPJC82.428]



Data: History



6

Methodology: Replicas





- Data is given by central values and covariance matrix
- generate Monte Carlo data replicas which as an ensemble represent the experiment
- fit one PDF replica to each data replica
- \Rightarrow ensemble of PDF replica



Theory: pineline [2302.12124]



https://nnpdf.github.io/pineline

- "Industrialization of High-Energy Theory Predictions":
 - · collect diverse generators in an "assembly line"
 - NNPDF4.0: > 4.5k datapoints + > 10 generators
- be reproducible (i.e. track data and metadata) and open source!
- not yet in NNPDF4.0 but any future release

 \Rightarrow please provide new calculations in an "interfaceable" way

PineAPPL is a fast interpolation grid library that

- extends to arbitrary orders in QCD and EW coupling
- provides a very good Command Line Interface
- provides several interfaces: C, C++, Fortran, Rust, Python
- can convert APPLgrid [EPJC66.503] and FastNLO [DIS12.217]
- interfaces to Mg5 [JHEP07.079], yadism, Vrap [PRD69.094008] soon MATRIX [EPJC78.537]

https://github.com/NNPDF/pineappl
 E https://nnpdf.github.io/pineappl/

Theory: EKO [EPJC82.976]



DGLAP:

$$\mu_F^2 \frac{\mathrm{d}}{\mathrm{d}\mu_F^2} \mathbf{E}(\mu_F^2 \leftarrow \mu_{F,0}^2) = \mathbf{P}(a_s(\mu_R^2), \mu_F^2) \otimes \mathbf{E}(\mu_F^2 \leftarrow \mu_{F,0}^2)$$

with

$$\mathbf{f}(\mu_F^2) = \mathbf{E}(\mu_F^2 \leftarrow \mu_{F,0}^2) \otimes \mathbf{f}(\mu_{F,0}^2)$$

- compute in Mellin space, but deliver in x space
- correct treatment of intrinsic PDFs

Theory: EKO Project Management

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Oldto

Interpolation

implementation: etc.interpolation

order to obtain the operators in an PDF independent way we use approximation theory rerefere, we define the basis grid

 $G = \{x_j : 0 < x_j <= 1, j = 0, ..., N_{grid} - 1\}$

from which we define our interpolatio

$$f(x) \sim \bar{f}(x) = \sum_{i=0}^{N_{FU}-1} f(x_i)p_i$$

This each grid paint x_j has an isoscitated hidepotation polynomia $p_j(x)$ (represented by in interpolation, maintenetian). We hidepotate in $\ln(x_j)$ using Lagrange interpolation among the event $(\Lambda_{max} + 1 \text{ points}, which renders the <math>p_i(x)$ (polynomials of order $O(2e^{N_{max}}(x))$).

Algorithm

First, we split the interpolation region into several areas (represented by size interpolation drug which are bound by the grid points:

$A_j = \{x_j, x_{j+1}\}, \text{ for } j = 0, ..., N_{grid} - 2$

Note, that we include the right border point into the definition, but not the left which keeps all areas objoint. This assumption is based on the physical fact, that PDFs do have a fixed upper bound to = 20, but no fixed invert bound.

econd, we define the interpolation blocks, which will build the interpolation polynomials and outsin the mended space of enterty.

- Fully open source: **()** https://github.com/NNPDF/eko
- Written in Python
- Fully documented: E https://eko.readthedocs.io/



https://github.com/NNPDF/yadism
 E https://yadism.readthedocs.io

• DIS coefficient function database:



• implemented flavor number schemes: ZM-VFNS, FFNS, FONLL

Checks: Future Tests [Acta Phys.Polon.B52.243]

Go to the past and look into the (back then) future!



χ /N (only exp. covmat)					
(dataset)	NNPDF4.0	pre-LHC	pre-Hera		
pre-HERA	1.09	1.01	0.90		
pre-LHC	1.21	1.20	23.1		
NNPDF4.0	1.29	3.30	23.1		



Checks: Future Tests [Acta Phys.Polon.B52.243]

Go to the past and look into the (back then) future!



- without data PDF errors have to be big
- with PDF errors the total uncertainty increases, and accommodates for difference between predictions and new data

Fake a universe with known input assumptions

- 1. Assume a "true" underlying PDF (e.g. a single PDF replica)
- 2. Produce fake data distributed accordingly
- 3. Perform a fit to this fake data

Observe statistical estimators (e.g. bias and variance) \rightarrow Is the truth within one sigma in 68% of cases?

$\sqrt{\text{bias/variance}}$	$\xi_{1\sigma}^{(\mathrm{data})}$
1.03 ± 0.05	0.68 ± 0.02

NNPDF4.0: PDF plot



NNPDF4.0: Uncertainties



- typical uncertainties in data region: singlet $\sim 1\%$, nonsinglet $\sim 2-3\%$
- data region: $10 \lesssim M_X \lesssim 3 imes 10^3 \, {\rm GeV}, -4 \lesssim y \lesssim 4$

2. Intrinsic Charm [Nature608.483]

• perturbative charm

- is fully perturbative, i.e., predictable at all scales
- generated by matching conditions and evolution
- always present for $\mu_F > \mu_h = m_h$
- it is $g
 ightarrow c ar{c}$, so (mostly) no asymmetry possible $(c
 eq ar{c})$
- default for CTEQ and MSHT

• intrinsic charm

- is non-perturbative
- charm in 3 light flavor scheme
- CTEQ: use a model (e.g. [BHPS])

• fitted charm

- default for NNPDF
- don't assume anything just fit charm in 4 flavor scheme!
- is an arbitrary mixture of intrinsic and perturbative charm

Mass Dependency on PDFs



Mass Dependency by OMEs



slow perturbative convergence of OME: NNLO and N3LO differ significantly

For (forward) evolution across a matching scale μ_h^2 :

 $\tilde{\mathbf{f}}^{(n_f+1)}(\mu_{F,1}^2) = \tilde{\mathbf{E}}^{(n_f+1)}(\mu_{F,1}^2 \leftarrow \mu_h^2) \mathbf{R}^{(n_f)} \tilde{\mathbf{A}}^{(n_f)}(\mu_h^2) \tilde{\mathbf{E}}^{(n_f)}(\mu_h^2 \leftarrow \mu_{F,0}^2) \tilde{\mathbf{f}}^{(n_f)}(\mu_{F,0}^2)$

with $\mathbf{R}^{(n_f)}$ a flavor rotation matrix and $\tilde{\mathbf{A}}^{(n_f)}(\mu_h^2)$ the operator matrix elements (partially known up to N³LO)

For (forward) evolution across a matching scale μ_h^2 :

 $\tilde{\mathbf{f}}^{(n_f+1)}(\mu_{F,1}^2) = \tilde{\mathbf{E}}^{(n_f+1)}(\mu_{F,1}^2 \leftarrow \mu_h^2) \mathbf{R}^{(n_f)} \tilde{\mathbf{A}}^{(n_f)}(\mu_h^2) \tilde{\mathbf{E}}^{(n_f)}(\mu_h^2 \leftarrow \mu_{F,0}^2) \tilde{\mathbf{f}}^{(n_f)}(\mu_{F,0}^2)$ with $\mathbf{R}^{(n_f)}$ a flavor rotation matrix and $\tilde{\mathbf{A}}^{(n_f)}(\mu_h^2)$ the operator matrix elements (partially known up to N³LO)

for backward evolution:

- invert $\tilde{\mathbf{E}}^{(n_f)}$: simple (invert RGE flow) \checkmark
- invert $\mathbf{R}^{(n_f)}$: simple (static matrix) \checkmark
- invert $\tilde{\mathbf{A}}^{(n_f)}$: expanded or exact

based on NNPDF4.0 [EPJC82.428]



The PDF Plot



- in **3FNS** a valence-like peak is present
- for $x \le 0.2$ the perturbative uncertainties are quite large
- the carried momentum fraction is within 1%

The PDF Plot with Model Comparison



[BHPS] or [Meson/Baryon Cloud Model]

- in **3FNS** a valence-like peak is present
- for $x \le 0.2$ the perturbative uncertainties are quite large
- the carried momentum fraction is within 1%



• predict better recent measurement

• reweighting is consistent

EMC [NPB461.181]



- direct measurement of F_2^c
- · evidence for intrinsic charm claimed, but experiment disputed
- adding EMC data is consistent



- direct measurement of F_2^c
- can distinguish intrinsic charm scenarios

Significance



- we find a 3σ evidence of intrinsic charm
- result is stable with mass variation, dataset variation

Charm Asymmetry (PRELIMINARY!)



- also parametrize $c^- = c \bar{c} \Rightarrow$ intrinsic!
- significance for baseline now $> 3\sigma$
- $\sim 1.5\sigma$ evidence for $c^-
 eq 0$

3. Summary

Summary

We fit the charm PDF in order to get

- realistic error estimate
- no strong dependence on charm mass
- no sensitivity to MHOU in matching condition

We find

- $\bullet\,$ large uncertainties and charm compatible with zero at small $\times\,$
- 3σ evidence for an intrinsic charm valence-like peak

The road ahead:

- more data $ightarrow 5\sigma$ evidence
- $c \bar{c}$ asymmetry phenomenology

Thank you!

4. Backup slides

Data: Kinematic Plot



- about 50 new datasets & 400 extra datapoints
- DIS/FTDY: dataset as in NNPDF3.1 + nomad neutrino + SeaQuest DY
- LHC: full 7 TeV and 8 TeV dataset & extensive use of 13 TeV data
- several new processes: prompt photon; single top; dijets; Hera jets

Methodology: Neural Network and Hyperoptimization

$$f_j(x, Q_0) = x^{-\alpha_j} (1-x)^{\beta_j} \mathrm{NN}_j(x)$$



- functional form: neural network (corresponds to many "effective parametrizations")
- choose model parameters? hyperoptimization! (i.e. scan parameter space)
- prevent overfitting!

Produce FastKernel (FK) tables!



The workhorse in the background: PineAPPL

IC - uncertainties splitted



IC - dataset variation



IC - mass dependency

