Extracting Excited States from Lattice Correlation Functions Kimmy Cushman & George Fleming





•—D Yale

Outline

- Introduction and motivation
- Correlation functions
- Prony's method
- Bootstrapping and results
- Clustering
- Future work

Lattice QCD... what is it?

- Gauge field simulation from
 - first principles
- Discretized space-time
- Finite volume with periodic conditions
- Volume ~ size of proton



http://www.physics.adelaide.edu.au/cssm/lattice/

• QCD strongly coupled at <u>low energies</u>

-> not perturbative

-jet fragmentation functions

-parton distribution functions



• QCD strongly coupled at <u>low energies</u>

-> not perturbative

-jet fragmentation functions

-parton distribution functions

 Increased computational power = competitive and supplementary to experiment!





Image credit: Carlos Jones/ORNL

1) > 95% of mass of hadrons come from QCD dynamics



Ab-initio Determination of Light Hadron Masses

S. Dürr¹, Z. Fodor^{1,2,3}, J. Frison⁴, C. Hoelbling^{2,3,4}, R. Hoffmann², S. D. Katz^{2,3}, S. Krieg², T. Kurth², L. Lellouch⁴, T. Lippert^{2,5}, K.K. Szabo², G. Vulvert⁴

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 ⁴Centre de Physique Théorique; Case 907, Campus de Luminy, F-13288 Marseille Cedex 9, France.
 ⁵Jülich Supercomputing Centre, FZ Jülich, D-52425 Jülich, Germany.

Budapest-Marseille-Wuppertal Collaboration

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2) Computation of matrix elements for weak flavor mixing

$$\Gamma(\pi \to l\nu) = \frac{G_F^2 |V_{ud}|^2 f_\pi^2}{8\pi} m_\pi m_l^2 \left(1 - \frac{m_l^2}{m_\pi^2}\right)^2$$

Perturbative Non-perturbative

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

• Predict spectrum of QCD baryons and mesons



New gauge forces to explain...

• Higgs mechanism?



• Dark matter?



https://www.nasa.gov/

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Why composite Higgs?

https://motls.blogspot.com/2015/03/did-homer-simpson-calculate-

1) Hierarchy problem - no longer a *fundamental* scalar



Why composite Higgs?

https://motls.blogspot.com/2015/03/did-homer-simpson-calculate-

Hierarchy problem - no longer a *fundamental* scalar
 no fine tuning necessary if composite



Why composite Higgs?

- Hierarchy problem no longer a *fundamental* scalar
 no fine tuning necessary if composite
- 2) Dynamical symmetry breaking Higgs model describes *effective* potential. Explains where potential comes from

$$\dot{\gamma} V(\phi) = \mu^2 \phi^{\dagger} \phi + \lambda (\phi^{\dagger} \phi)^2 \dot{\gamma}$$



https://motls.blogspot.com/2015/03/did-homer-simpson-calculate-correct.html

• Every confining force has a spectrum of states

 New force of composite Higgs would have spectrum



 $\langle 0|\bar{\pi}(t)\pi(0)|0\rangle = C_{\pi}(t) = \mathcal{O}$

$$\langle 0|\bar{\pi}(t)\pi(0)|0\rangle = C_{\pi}(t) = \mathcal{O}$$

$$\langle \mathcal{O} \rangle = \frac{\int_{\mathcal{D}} O[A] e^{-iS[A]}}{\int_{\mathcal{D}} e^{-iS[A]}}$$





Importance Sampling

 $\langle 0|\bar{\pi}(t)\pi(0)|0\rangle = C_{\pi}(t) = \mathcal{O}$ $\langle \mathcal{O} \rangle = \frac{\int_{\mathcal{D}} \mathcal{O}[A] e^{-iS[A]}}{\int e^{-S[A]}} \longrightarrow \quad \langle \mathcal{O} \rangle = \frac{1}{Z} \sum_{i} e^{-\beta S} \mathcal{O}_{i}$

Configuration 1 \mathcal{O}_1

Configuration 2 \mathcal{O}_2

Configuration 3 \mathcal{O}_3



Configuration 4
$$\mathcal{O}_4$$

$$\langle \mathcal{O} \rangle = \frac{1}{N} \sum_{i}^{N} \mathcal{O}_{i}$$

Importance Sampling

$$\langle 0|\bar{\pi}(t)\pi(0)|0\rangle = C_{\pi}(t) = \mathcal{O}$$

$$\langle \mathcal{O} \rangle = \frac{1}{N} \sum_{i}^{N} \mathcal{O}_{i}$$



 $C(t) = \langle 0 | \bar{\pi}(t) \pi(0) | 0 \rangle$



$$C(t) = \langle 0 | \bar{\pi}(t) \pi(0) | 0 \rangle$$

=
$$\sum_{m}^{\infty} \langle 0 | \pi(0) | E_{m} \rangle e^{-E_{m}t} \langle E_{m} | \pi(0) | 0 \rangle$$



 $C(t) = \langle 0 | \bar{\pi}(t) \pi(0) | 0 \rangle$



http://watersoundimage.yolasite.com/what-is-a-w-s-image.php

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$$\Rightarrow C(t) = \sum_{m=1}^{\infty} a_m \,\mathrm{e}^{-E_m t}$$



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$$C(t) = a_0 e^{-E_0 t} + a_1 e^{-E_1 t} + \cdots$$

 $\approx a_0 e^{-E_0 t}$

$$\Rightarrow C(t) = \sum_{m}^{\infty} a_m \,\mathrm{e}^{-E_m t}$$

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$$C(t) = a_0 e^{-E_0 t} + a_1 e^{-E_1 t} + \cdots$$
$$\approx a_0 e^{-E_0 t}$$

$$\frac{C(t+1)}{C(t)} \approx \frac{a_1 e^{-E_0(t+1)}}{a_1 e^{-E_0 t}}$$
$$= e^{-E_0}$$



$$\Rightarrow E_1 = \log\left(\frac{C(t)}{C(t+1)}\right)$$



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Standard Method for Excited States

Fit more exponentials! but

$$\Rightarrow C(t) = \sum_{m}^{\infty} a_m \,\mathrm{e}^{-E_m t}$$

- Difficult/time consuming with decaying exponentials
- Finding global minimum best a_m , E_m is hard
- 2M dimensional parameter space
- Prone to user bias choosing initial values



Prony, G. R. B. "J. de Lh Ecole Polytechnique." Paris 1 (1795): 24.

G. Fleming, S. Cohen, H. Lin, V. Pereyra (2009)

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$$y_n(t) \equiv C(t+n)$$

= $\sum_m^M a_m e^{-E_m(t+n)}$
= $\sum_m^M a_m e^{-E_m t} e^{-E_m n}$

$$\begin{pmatrix} y_{0} \\ y_{1} \\ \vdots \\ y_{2M-1} \end{pmatrix} = \underbrace{\begin{pmatrix} 1 & 1 & \cdots & 1 \\ z_{1} & z_{2} & \cdots & z_{M} \\ z_{1}^{2} & z_{2}^{2} & \cdots & z_{M}^{2} \\ \vdots & \vdots & \ddots & \vdots \\ z_{1}^{2M-1} & z_{2}^{2M-1} & \cdots & z_{M}^{2M-1} \end{pmatrix}}_{\text{Vandermode matrix}} \begin{pmatrix} A_{1} \\ A_{2} \\ \vdots \\ A_{M} \end{pmatrix}$$

$$\Rightarrow y_n(t) = \sum_m^M A_m(t) z_m^n$$

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$$egin{aligned} y_n(t) &\equiv C(t+n) \ &= \sum_m^M a_m \mathrm{e}^{-E_m(t+n)} \ &= \sum_m^M a_m \mathrm{e}^{-E_m t} \mathrm{e}^{-E_m n} \end{aligned}$$

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Vandermode matrix

$$\Rightarrow y_n(t) = \sum_m^M A_m(t) z_m^n$$

Non-linear

linear

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Prony's Method for M= 3

M = 3 means 6 y's

$$0 = \begin{vmatrix} y_0 & y_1 & y_2 & 1 \\ y_1 & y_2 & y_3 & z \\ y_2 & y_3 & y_4 & z^2 \\ y_3 & y_4 & y_5 & z^3 \end{vmatrix}$$
Prony's Method for M= 3

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Need to solve M^{th} order polynomial

$$0 = p_0 + p_1 z + p_2 z^2 + p_3 z^3 \qquad z_m = e^{-E_m}$$

E. Berkowitz, A. Nicholson, C. Chang et al. (2017)

S. Beane, W. Detmold, T. Luu et al. (2009)

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$$C(t) \rightarrow \begin{pmatrix} C_{11}(t) & C_{12}(t) \\ C_{21}(t) & C_{22}(t) \end{pmatrix}$$

 $C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j(0) | 0 \rangle$

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Effective mass becomes generalized eigenvalue problem

$$0 = C(t+1) - zC(t) \qquad e^{-E_0} = \frac{C(t+1)}{C(t)}$$

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$$C(t+1)v = \lambda C(t)v \qquad \lambda = e^{-E_0}$$

Effective mass becomes generalized eigenvalue problem

$$\lambda = e^{-E} \qquad \qquad C(t+1)v = \lambda C(t)v$$

$$C(t+1)v_0 = e^{-E_0}C(t)v_0$$

$$C(t+1)v_1 = e^{-E_1}C(t)v_1$$

Effective mass becomes generalized eigenvalue problem

$$\lambda = e^{-E} \qquad \qquad C(t+1)v = \lambda C(t)v$$

maximal overlap $C(t+1)v_0 = e^{-E_0}C(t)v_0 \quad \longleftarrow \quad \text{with ground state}$ $C(t+1)v_1 = e^{-E_1}C(t)v_1 \quad \longleftarrow \quad \text{with first excited state}$

Effective mass becomes generalized eigenvalue problem

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maximal overlap

 $C(t+1)v_1 = e^{-E_1}C(t)v_1 \longleftarrow$ with first excited state

Ground state matrix equation

$$\begin{pmatrix} C_{11}(t+1) & C_{12}(t+1) \\ C_{21}(t+1) & C_{22}(t+1) \end{pmatrix} \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = e^{-E_0} \begin{pmatrix} C_{11}(t) & C_{12}(t) \\ C_{21}(t) & C_{22}(t) \end{pmatrix} \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}$$

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$$e^{-E_0(t+1)} \begin{pmatrix} \langle 0|\mathcal{O}_1(b_1\mathcal{O}_1 + b_2\mathcal{O}_2)|0\rangle \\ \langle 0|O_2(b_1\mathcal{O}_1 + b_2\mathcal{O}_2)|0\rangle \end{pmatrix} = e^{-E_0} e^{-E_0t} \begin{pmatrix} \langle 0|\mathcal{O}_1(b_1\mathcal{O}_1 + b_2\mathcal{O}_2)|0\rangle \\ \langle 0|\mathcal{O}_2(b_1\mathcal{O}_1 + b_2\mathcal{O}_2)|0\rangle \end{pmatrix}$$

Hankel Matrix determinant specifies eigenvalue problem

$$0 = \begin{vmatrix} y_0 & y_1 & y_2 & 1 \\ y_1 & y_2 & y_3 & z \\ y_2 & y_3 & y_4 & z^2 \\ y_3 & y_4 & y_5 & z^3 \end{vmatrix}$$

M solutions

$$0 = p_0 + p_1 z + p_2 z^2 + p_3 z^3$$
 $z_m = e^{-E_m}$ Scalar equation

Hankel Matrix determinant specifies eigenvalue problem

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M solutions



6 data points C(t=t_o) for





$$\frac{1}{6}(1+1+2+3+5+5) = C_1(t=t_0)$$

Repeat for all times to obtain $C_1(0), C_1(1), C_1(2), C_1(3), ..., C_1(T)$

$$\frac{1}{6}(1+1+2+3+5+5) = C_1(t=t_0)$$

Repeat for all times to obtain C₁(0), C₁(1), C₁(2), C₁(3), ... C₁(T) $y_n(t) = C(t + n)$

$$\begin{pmatrix} y_0 \\ y_1 \\ y_2 \\ y_3 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ z_1 & z_2 \\ z_1^2 & z_2^2 \end{pmatrix} \begin{pmatrix} A_1 \\ A_2 \end{pmatrix} \qquad 0 = \begin{vmatrix} y_0 & y_1 & 1 \\ y_1 & y_2 & z \\ y_2 & y_3 & z^2 \end{vmatrix}$$

 $C_1(t) \Rightarrow \{ (z_1, a_1), (z_2, a_2) \}$

$$a_m = |\langle 0|\pi |E_m\rangle|^2 \qquad z_m = e^{-E_m}$$



 $M=\!2,\,t=8$













M = 2, t = 8-13.0• -13.2•••• -13.4-13.6 $\log(a)$ -13.8-14.0٠ -14.2 · -14.4-0.4-0.1-0.3-0.2 $\log(z)$

M = 2, t = 8-13.0• -13.2••••• -13.4-13.6 $\log(a)$ -13.8-14.0٠ -14.2 · -14.4-0.4-0.1-0.3-0.2 $\log(z)$



M = 2 States at t = 2













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M = 3 State Extraction



Expectation Maximization Clustering

Dempster, A.P.; Laird, N.M.; Rubin, D.B. (1977). "Maximum Likelihood from Incomplete Data via the EM Algorithm". Journal of the Royal Statistical Society, Series B. 39 (1): 1–38.

$$\vec{\mu} = \left\langle \begin{pmatrix} x \\ y \end{pmatrix} \right\rangle$$
$$\Sigma_{xx} = \sigma_x^2 = \left\langle (x - \mu_x)^2 \right\rangle$$
$$\Sigma_{yy} = \sigma_y^2 = \left\langle (y - \mu_y)^2 \right\rangle$$
$$\Sigma_{xy} = \Sigma_{yx}^{\dagger} = \left\langle (x - \mu_x)(y - \mu_y) \right\rangle$$



$$p(\vec{x}) = \frac{1}{2\pi\sqrt{|\Sigma|}} e^{-\frac{1}{2}(\vec{x}-\vec{\mu})^T \Sigma^{-1} (\vec{x}-\vec{\mu})}$$

Lloyd., S. P. (1982). "Least squares quantization in PCM" (PDF). IEEE Transactions on Information Theory. 28 (2): 129–137.

1. Assign points to initial clusters

- 2. Compute the mean and covariance matrix for each cluster
- 3. For each bootstrap sample, find most probably permutation of points among clusters
- 4. Repeat 2. and 3. until the process converges

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$$p_1(\vec{x}) > p_2(\vec{x})$$

$$\Rightarrow \frac{1}{\sqrt{|\Sigma_1|}} e^{-\frac{1}{2}(\vec{x} - \vec{\mu_1})^T \Sigma_1^{-1}(\vec{x} - \vec{\mu_1})} > \frac{1}{\sqrt{|\Sigma_2|}} e^{-\frac{1}{2}(\vec{x} - \vec{\mu_2})^T \Sigma_2^{-1}(\vec{x} - \vec{\mu_2})},$$

$$\Rightarrow d_1(\vec{x}) + \log|\Sigma_1| < d_2(\vec{x}) + \log|\Sigma_2|$$



$\Rightarrow d_1(\vec{x}) + \log|\Sigma_1| < d_2(\vec{x}) + \log|\Sigma_2|$



 $z_m = \mathrm{e}^{-E_m}$ $0 = p_0 + p_1 z + p_2 z^2$

Clustering M = 2, Initial



Clustering M = 2, 1 iteration



Clustering M = 3, Initial



Clustering M = 3, 1 iteration



Clustering M = 3, 2 iterations



Clustering M = 3, 3 iterations



Clustering M = 3, 4 iterations



Clustering M = 3, 5 iterations



Clustering M = 3, 6 iterations



Clustering M = 3, 7 iterations



Clustering M = 3, 8 iterations



Clustering M = 3, 9 iterations



Clustering M = 3, 10 iterations



Clustering M = 3, 11 iterations



Clustering M = 3, 12 iterations



Clustering M = 3, 13 iterations



Preliminary Results



Preliminary Results



Preliminary Results



Future Work

• Non-Gaussian clusters - try a new distance metric -Don't worry about noisier time slices



Future Work

- Non-Gaussian clusters try a new distance metric -Don't worry about noisier time slices
- All data in "stencil" extract M = T/2 states -Remove all operator bias



Possible solution: Tukey Depth

Multidimensional generalization of *percentiles*



Possible solution: Tukey Depth

Multidimensional generalization of *percentiles*





Possible solution: Tukey Depth

Multidimensional generalization of *percentiles*

Non-parametric statistic!





Tukey Depth Example with non-Gaussian Data

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Tukey Depth Example with non-Gaussian Data



Tukey Depth Example with non-Gaussian Data


Tukey Depth Example with non-Gaussian Data



Conclusion

- Extracting excited states is important for lattice QCD and BSM lattice
- Standard method involves fitting to exponentials and has many know problems
- Prony's method may be a better approach if we can identify clusters
- We will need a better clustering algorithm to account for weirdly shaped clusters

Thanks!

