# REVISITING THE DISCOVERY POTENTIAL OF THE ISOBAR RUN AT RHIC

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[Jan Hammelmann, ASO, Massimiliano Alvioli, Hannah Elfner, Mark Strikman, arXiv:1908.10231]

#### JLAB Theory Seminar

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# The Chiral Magnetic Effect in a nutshell

QCD axial anomaly in the massless limit



Ingredients

Chiral fermions **2)**  $F\tilde{F} \neq 0$  **3)** Strong magnetic field



 $\mathbf{M}$  Chiral symmetry restoration in the QGP phase  $T_c \sim 156$  MeV

[HotQCD Collab. PLB795 (2019) 15-21]



Chiral symmetry restoration in the QGP phase



3



Chiral symmetry restoration in the QGP phase

#### $\mathbf{M}$ Sources of $F\tilde{F} \neq 0$

Axial charge fluctuations in the Glasma

[T. Lappi, S. Schlichting, PRD 97, 034034 (2018)] [P. Guerrero-Rodriguez, JHEP 1908 (2019) 026]



3



Chiral symmetry restoration in the QGP phase

 $\mathbf{M}$  Several sources of  $F\tilde{F} \neq 0$ 

Skokov et al. Int.J.Mod.Phys. A24 (2009) 5925-5932

$$\vec{j} \propto \mu_5 \overrightarrow{B}$$



# **Experimental determination of the CME**

Charge-dependent azimuthal correlations

$$\gamma_{\alpha\beta} = \langle \cos(\phi_{\alpha} + \phi_{\beta} - 2\Psi_{\rm RP}) \rangle, \alpha, \beta \in \{+, -\}$$

[S.Voloshin PRC 70, 057901 (2004)]

CME prediction:  $\gamma_{++,--} < 0, \gamma_{+-} > 0$ 





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CME prediction:  $\gamma_{++,--} < 0$ ,  $\gamma_{+-} > 0$ 





# **Experimental determination of the CME**

Contributions to the experimental measurement:



Local charge conservation, momentum conservation, resonance decays...

e.g.[S.Schlichting, S.Pratt PRC 83 (2011) 014913] STAR'09 data explained without CME

# Two ways out

• New observables beyond  $\gamma_{\alpha\beta}$  have been proposed.

e.g. [M. Magdy et al. PRC 97, 061901 (2018)]

Isobar run at RHIC

"We have specifically investigated the case for colliding nuclear isobars (nuclei with the same mass but different charge) and find the case compelling. We recommend that a program of nuclear isobar collisions to isolate the chiral magnetic effect from background sources be placed as a high priority item in the strategy for completing the RHIC mission."

> - Chiral Magnetic Effect Task Force Report [Chin.Phys. C41 (2017) no.7, 072001]

 $_{44}$ Ru<sup>96</sup> +  $_{44}$ Ru<sup>96</sup> vs.  $_{40}$ Zr<sup>96</sup> +  $_{40}$ Zr<sup>96</sup> @200GeV Change signal (B), keeping background (v<sub>2</sub>) fixed Check the impact of  ${}_{44}$ Ru<sup>96</sup> and  ${}_{40}$ Zr<sup>96</sup> nuclear structure on the hypothesis "Change signal (B), keeping background (v<sub>2</sub>) fixed"



# Woods-Saxon parameters for isobars

Traditionally nucleons sampled from a Woods-Saxon distribution

 $\rho(r,\theta) = \frac{\rho_0}{e^{(r-R'(\theta,\phi))/d} + 1}$ 

where 
$$R'(\theta) = R_0(1 + \beta_2 Y_2^0(\theta))$$

44**Ru**<sup>96</sup>

 $R_0 = 5.08 \text{ fm}$ 

d = 0.46 fm

#### <u>Option 1</u>: $\beta_2 = 0.158$

[Prytichenko et al. Atom Data Nucl. Data Tabl. 107 1, (2016)]

Option 2:  $\beta_2 \sim 0$ 

[Moller et al. Atom Data Nucl. Data Tabl. 59 185, (1995)]



### **Beyond Woods-Saxon distribution**



X No distinction between protons and neutrons

X Nucleons are considered independent of each other

# Neutron skin within the liquid drop model

 $B(Z,N) = a_v A - a_s A^{2/3} - a_c Z^2 / A^{1/3} - a_A (N-Z)^2 / A + \dots$ [Wikipedia] Surface Coulomb Volume Asymmetry

Neutron-rich nuclei: where to locate extra-neutrons? Surface or core?

Neutron-skin thickness: surface tension vs stiffness of the symmetry energy (L)  $\Delta r_{np}$ 

[X. Roca-Maza et al. PRL106 (2011), 252501]

### Neutron skin: experimental data



[LEAR Collab. Int Jour of Mod Phys E 13 (2004) 343]

#### Neutron skin: experimental data



### Neutron skin: implementation for <sub>40</sub>Zr<sup>96</sup>

Neutron-halo type i.e.  $\Delta r_{np}$ translates into  $d_n > d_p$ 





### Short-range nucleon-nucleon correlations



# Short-range nucleon-nucleon correlations

Metropolis Monte Carlo generator of nuclear configurations using [Alvioli et al. PLB 680 (2009]



Other studies in heavy-ion collisions: [http://inspirehep.net/author/profile/M.Alvioli.1]

# Nuclear configurations for isobar run

[Jan Hammelmann, ASO, Massimiliano Alvioli, Hannah Elfner, Mark Strikman, arXiv:1908.10231]



To be compared with "default" Woods-Saxon distribution

Role of deformation and neutron-skin on eccentricity and magnetic field strength

### SMASH as a tool [Weil et al. PRC 94 (2016) no.5, 054905]

Relativistic hadronic transport approach available at

https://smash-transport.github.io



Binary hadron-hadron interactions proceed via string excitation and decay à la PYTHIA 8 [Mohs et al. arXiv:1909.05586]

# Eccentricity as a proxy for $v_2$

Participant eccentricity at the maximum overlap time  $t = R/\sqrt{\gamma^2 - 1}$ 



Up to 10% difference on background contribution among the isobars in ultra-central collisions

# **Eccentricity probability distribution**

Normalized such that the integral is one



Neat enhancement in ultra-central collisions due to deformation

### "..., keeping background (v<sub>2</sub>) fixed"



Same background from mid-to-peripheral collisions

# **Check with the literature**

Our result is in agreement with other models



Not sufficient to select the same centrality bin in both event samples to ensure identical background

# Magnetic field

Event-by-event calculation via Lienard-Wiechert retarded potentials

$$e\vec{B}(t,\vec{r}) = \alpha_{\text{EM}} \sum_{n} \frac{(1-v_n^2)(\vec{v}_n \otimes \vec{R}_n)}{R_n^3 [1-(\vec{R}_n \otimes \vec{v}_n)/R_n^3]^{3/2}}$$



•  $\vec{R}_n(t) = \vec{r} - \vec{r}_n(t)$ : observation point





# Magnetic field strength

Magnetic field strength computed at the center of the collision



Concentration of protons in the nucleus core leads to a higher magnetic field in peripheral collisions

### **Origin of the enhancement**

$$B^2 = \langle B \rangle^2 + \sigma^2$$



Not an effect of a larger degree of fluctuations

### "Change signal (B), keeping background (v<sub>2</sub>) fixed"

