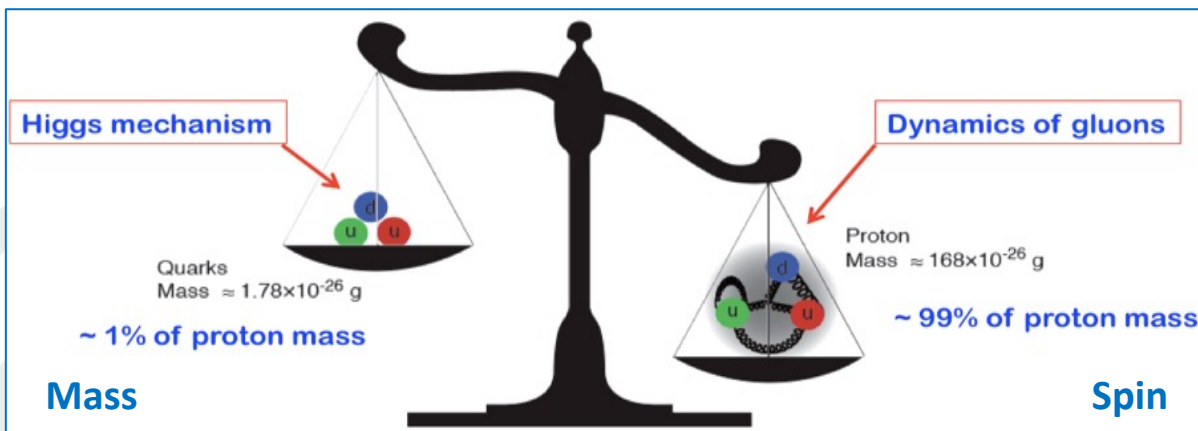
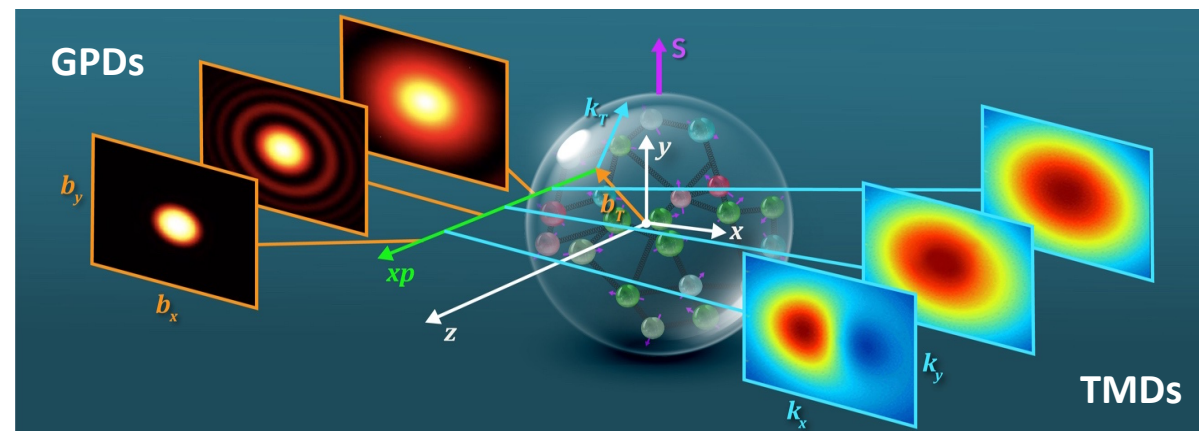


Generalized Parton Distributions from Diffractive Hard Exclusive Processes



Emergent Properties



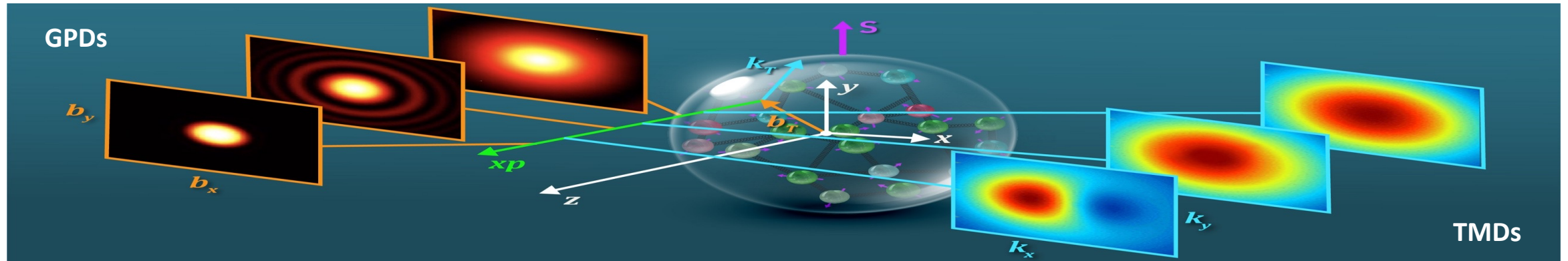
Dynamical Structure

In collaboration with Zhitc Yu (MSU)
JHEP 08 (2022) 103, PRD 107 (2023) 1,
arXiv:2305.15397, and in preparation

Jianwei Qiu
Jefferson Lab, Theory Center

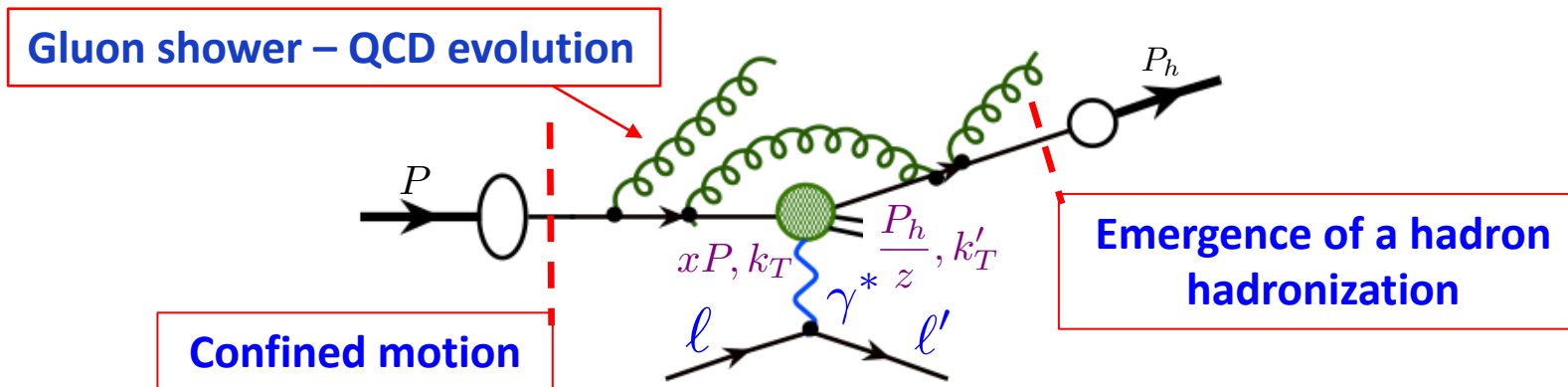
To "See" Internal Structure of Hadron without seeing quarks/gluons?

□ 3D hadron structure:



NO quarks and gluons can be seen in isolation!

□ If the nucleon is broken, e.g., in SIDIS, ...



Transverse momentum broadening:

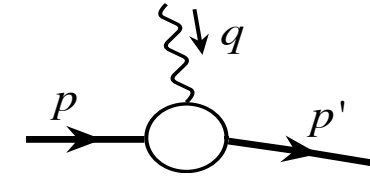
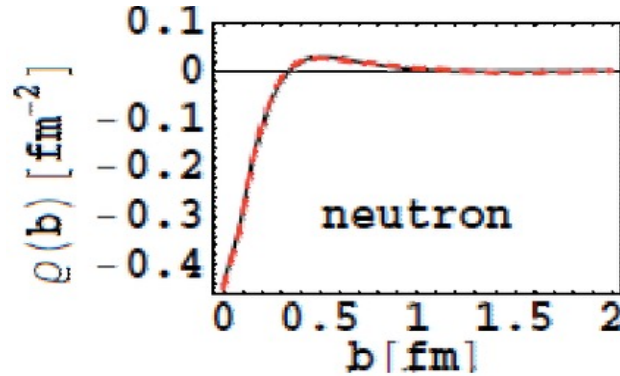
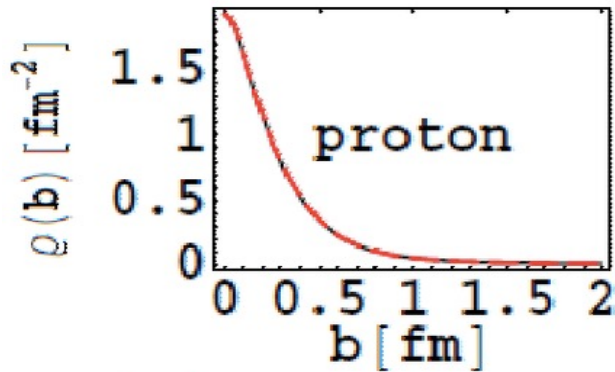
$$\Delta k_T^2 \propto \Lambda_{\text{QCD}}^2 \times \alpha_s(C_F, C_A) \times \log(Q^2/\Lambda_{\text{QCD}}^2) \times \log(s/Q^2) \gtrsim 1$$

Structure information is diluted by the collision induced shower!

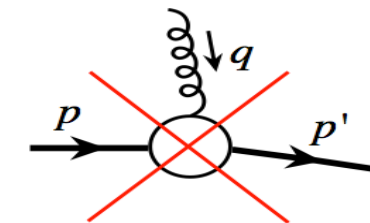
- *Measured k_T is NOT the same as k_T of the confined motion!*
- *Larger Q^2 could weaken our precision to probe the true hadron structure!*

How to Explore Internal Structure of Hadron without Breaking it?

□ **Form factors:** Elastic electric form factor → Charge distributions



Proton "Radius" in EM charge distribution

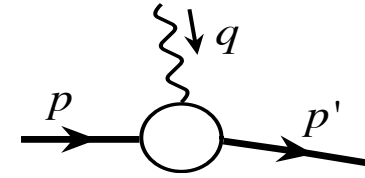
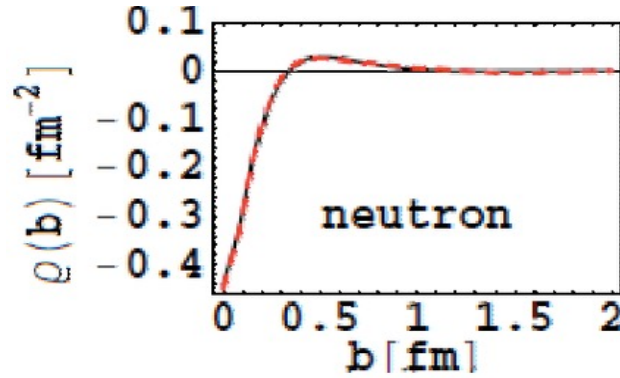
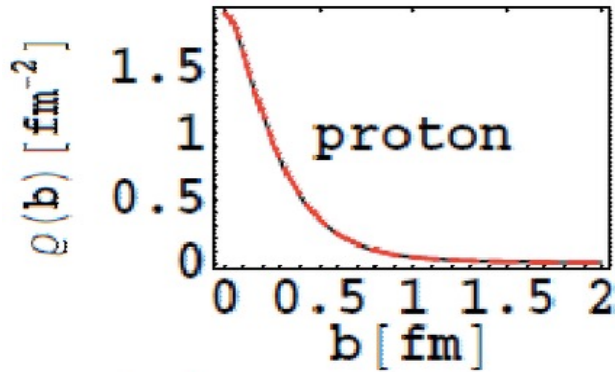


No Proton "Radius" in color charge distribution!

□ **But, there is NO elastic "color" form factor!**

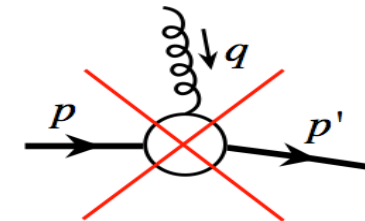
How to Explore Internal Structure of Hadron without Breaking it?

□ **Form factors:** Elastic electric form factor → Charge distributions



Proton "Radius" in EM charge distribution

□ **But, there is NO elastic "color" form factor!**



No Proton "Radius" in color charge distribution!

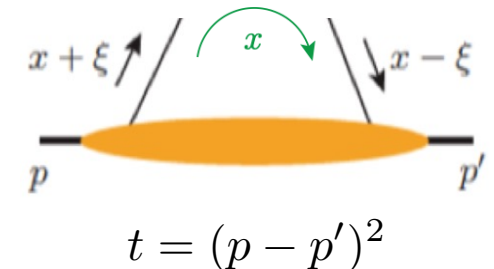
□ **3D hadron tomography:**

Generalized "form factor" for quark and gluon "density" distribution

Generalized PDFs (GPDs) – without breaking the proton

$$F_{q/h}(x, \xi, t) \quad \text{skewness} \quad \xi = \frac{(p - p')^+}{(p + p')^+} \quad t = (p - p')^2$$

Measure p' uniquely fix the t and ξ . It is the x sensitive to partonic structure



GPDs are Fundamental, Responsible for Emergent Hadron Properties

□ “Mass” – QCD energy-momentum tensor:

Ji, PRL78, 1997

$$T^{\mu\nu} = \sum_{i=q,g} T_i^{\mu\nu} \quad \text{with} \quad T_q^{\mu\nu} = \bar{\psi}_q i\gamma^{(\mu} \overleftrightarrow{D}^{\nu)} \psi_q - g^{\mu\nu} \bar{\psi}_q (i\gamma \cdot \overleftrightarrow{D} - m_q) \psi_q \quad \text{and} \quad T_g^{\mu\nu} = F^{a,\mu\eta} F^{a,\eta\nu} + \frac{1}{4} g^{\mu\nu} (F_{\rho\eta}^a)^2$$

□ Gravitational form factors:

$$\langle p' | T_i^{\mu\nu} | p \rangle = \bar{u}(p') \left[A_i(t) \frac{P^\mu P^\nu}{m} + J_i(t) \frac{iP^{(\mu} \sigma^{\nu)\Delta}}{2m} + D_i(t) \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{4m} + m \bar{c}_i(t) g^{\mu\nu} \right] u(p)$$

$$C_i(t) \leftrightarrow D_i(t)/4$$

□ Connection to GPD moments:

$$\int_{-1}^1 dx x F_i(x, \xi, t) \propto \langle p' | T_i^{++} | p \rangle \propto \bar{u}(p') \left[\underbrace{(A_i + \xi^2 D_i)}_{\int_{-1}^1 dx x H_i(x, \xi, t)} \gamma^+ + \underbrace{(B_i - \xi^2 D_i)}_{\int_{-1}^1 dx x E_i(x, \xi, t)} \frac{i\sigma^{+\Delta}}{2m} \right] u(p)$$

Related to pressure & stress force inside h

Polyakov, Schweitzer, *Inntt. J. Mod. Phys.* A33, 1830025 (2018)
Burkert, Elouadrhiri, Girod *Nature* 557, 396 (2018)

□ “Spin” – Angular momentum sum rule:

$$J_i = \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H_i(x, \xi, t) + E_i(x, \xi, t)]$$

$i = q, g$

3D tomography

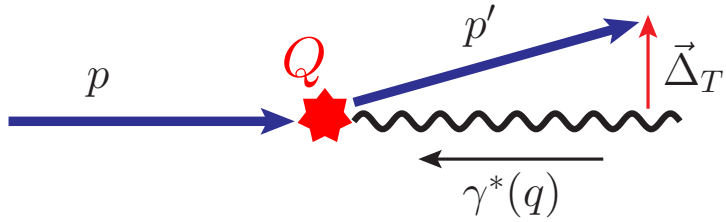
Relation to GFF
Angular Momentum

x-dependence of GPDs!

Need to know the x-dependence of GPDs to construct the proper moments!

Exclusive Diffractive Process for Extracting GPDs

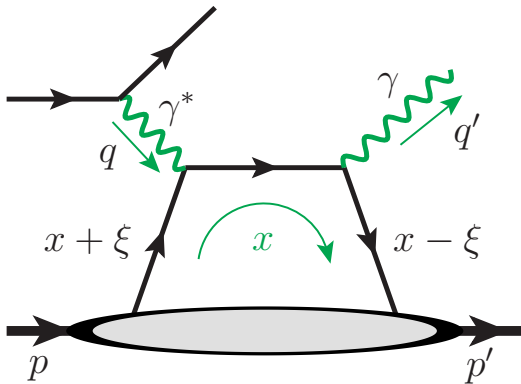
- Hit the proton hard without breaking it \Rightarrow Diffractive scattering to keep proton intact



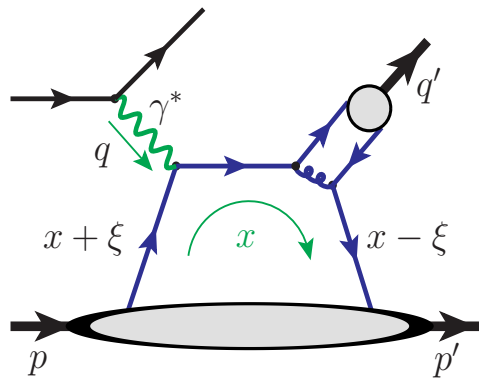
HERA discovery:

\sim 10-15% of HERA events with the Proton stayed intact

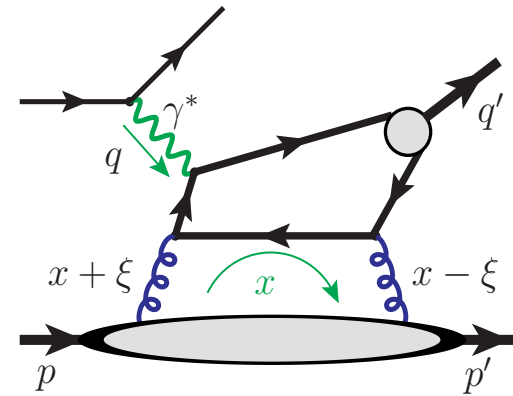
- Known exclusive processes for extracting GPDs:



DVCS: $Q^2 \gg |t|$



DVMP



DVQP

+ DDVCS, ...

Feature: Two-scale observables

$$Q^2 \gg |t|$$

$$t = (p - p')^2$$

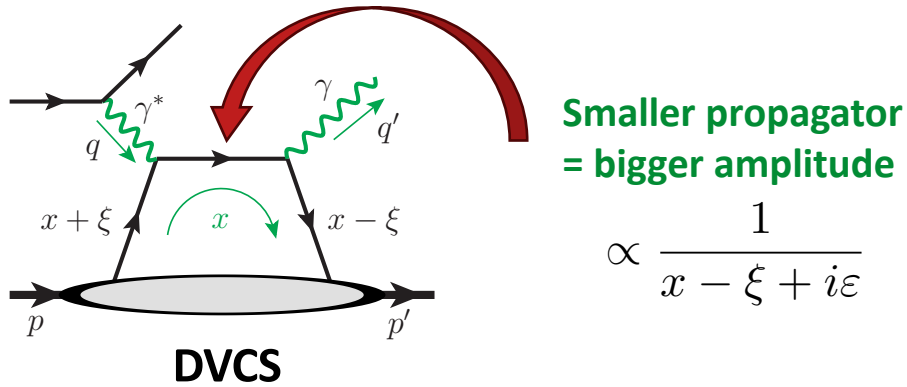
- Hard scale Q : allows pQCD, factorization
- Low scale t : probes non-pert. hadron structure

\rightarrow
Factorization

GPDs: $f_{i/h}(x, \xi, t; \mu)$

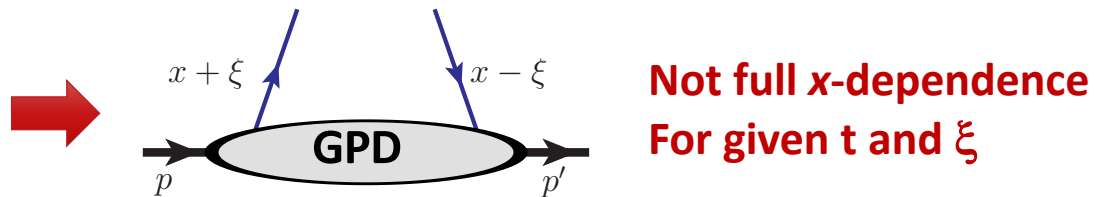
Why it is so Difficult to Extract the x -dependence of GPDs?

□ Amplitude nature: $x \sim$ loop momentum



$$\Rightarrow i\mathcal{M} \propto \int_{-1}^1 dx \frac{F(x, \xi, t)}{x - \xi + i\varepsilon} \equiv "F_0(\xi, t)"$$

- also true for most other processes
- x -dependence is only constrained by a “moment”
- x -integration decouples from external Q^2

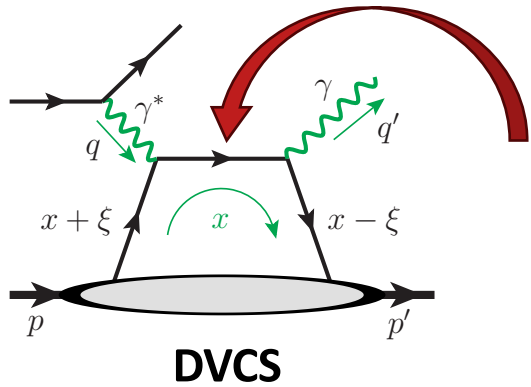


Why it is so Difficult to Extract the x -dependence of GPDs?

□ **Amplitude nature:** $x \sim$ loop momentum

□ **“Shadow GPDs”**

PRD103 (2021) 114019

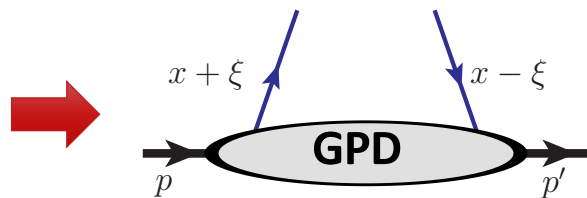


**Smaller propagator
= bigger amplitude**

$$\propto \frac{1}{x - \xi + i\epsilon}$$

➔ $i\mathcal{M} \propto \int_{-1}^1 dx \frac{F(x, \xi, t)}{x - \xi + i\epsilon} \equiv \text{“}F_0(\xi, t)\text{”}$

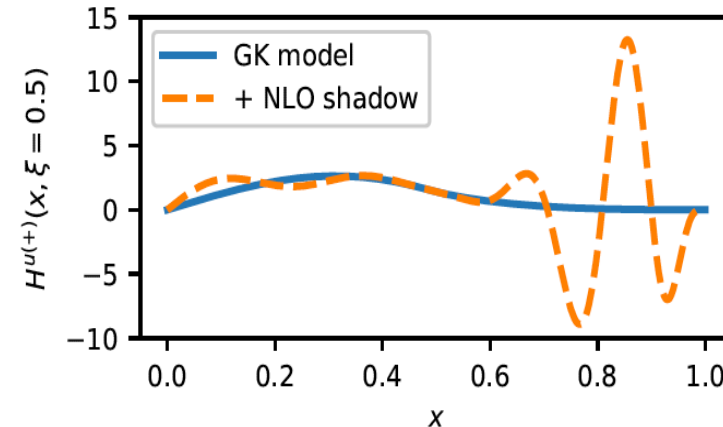
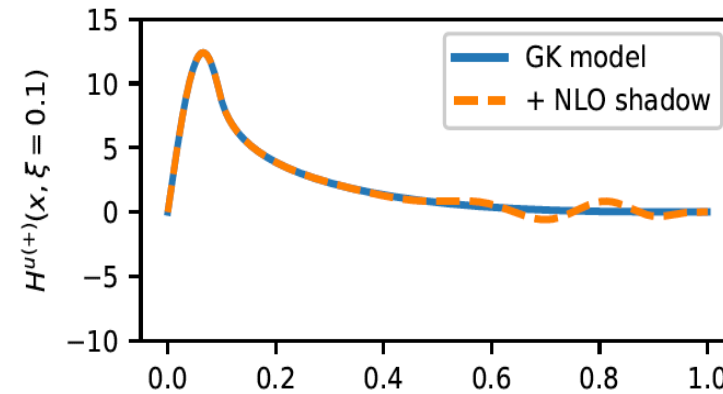
- also true for most other processes
- x -dependence is only constrained by a “moment”
- x -integration decouples from external Q^2



**Not full x -dependence
For given t and ξ**

$$F(x, \xi, t) \rightarrow F(x, \xi, t) + S(x, \xi, t)$$

with $\int_{-1}^1 dx \frac{S(x, \xi, t)}{x - \xi + i\epsilon} = 0$



**Blue and dashed
Fit the same CFFs !**

What kind of process/observable could be sensitive to the x-dependence?

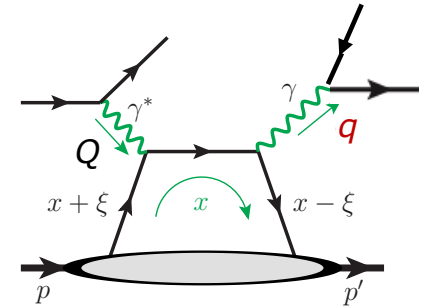
- Create an entanglement between the internal x and an externally measured variable?

$$i\mathcal{M} \propto \int_{-1}^1 dx \frac{F(x, \xi, t)}{x - x_p(\xi, q) + i\varepsilon}$$

Change external q to sample different part of x .

- Double DVCS (two scales):

$$x_p(\xi, q) = \xi \left(\frac{1 - q^2/Q^2}{1 + q^2/Q^2} \right) \rightarrow \xi \text{ same as DVCS if } q \rightarrow 0$$

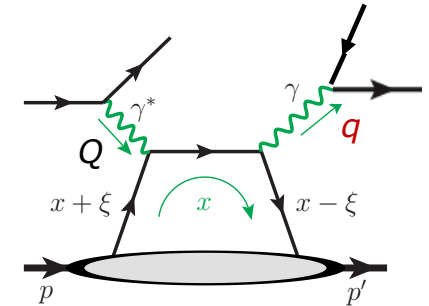


What kind of process/observable could be sensitive to the x -dependence?

- Create an entanglement between the internal x and an externally measured variable?

$$i\mathcal{M} \propto \int_{-1}^1 dx \frac{F(x, \xi, t)}{x - x_p(\xi, q) + i\varepsilon}$$

Change external q to sample different part of x .



- Double DVCS (two scales):

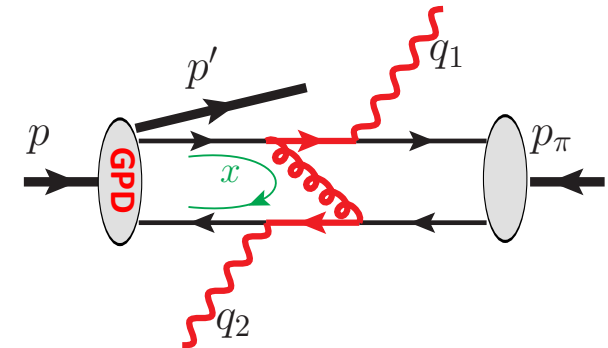
$$x_p(\xi, q) = \xi \left(\frac{1 - q^2/Q^2}{1 + q^2/Q^2} \right) \rightarrow \xi \text{ same as DVCS if } q \rightarrow 0$$

- Production of two back-to-back high p_T particles (say, two photons):

$$\pi^-(p_\pi) + P(p) \rightarrow \gamma(q_1) + \gamma(q_2) + N(p')$$

JHEP 08 (2022) 103

Hard scale: $q_T \gg \Lambda_{\text{QCD}}$ Soft scale: $t \sim \Lambda_{\text{QCD}}^2$



- Factorization:

$$\mathcal{M}(t, \xi, q_T) = \int_{-1}^1 dx F(x, \xi, t; \mu) \cdot C(x, \xi; q_T/\mu) + \mathcal{O}(\Lambda_{\text{QCD}}/q_T)$$

[suppressing pion DA factor]



$$\frac{d\sigma}{dt d\xi dq_T} \sim |\mathcal{M}(t, \xi, q_T)|^2$$

q_T distribution is "conjugate" to x distribution

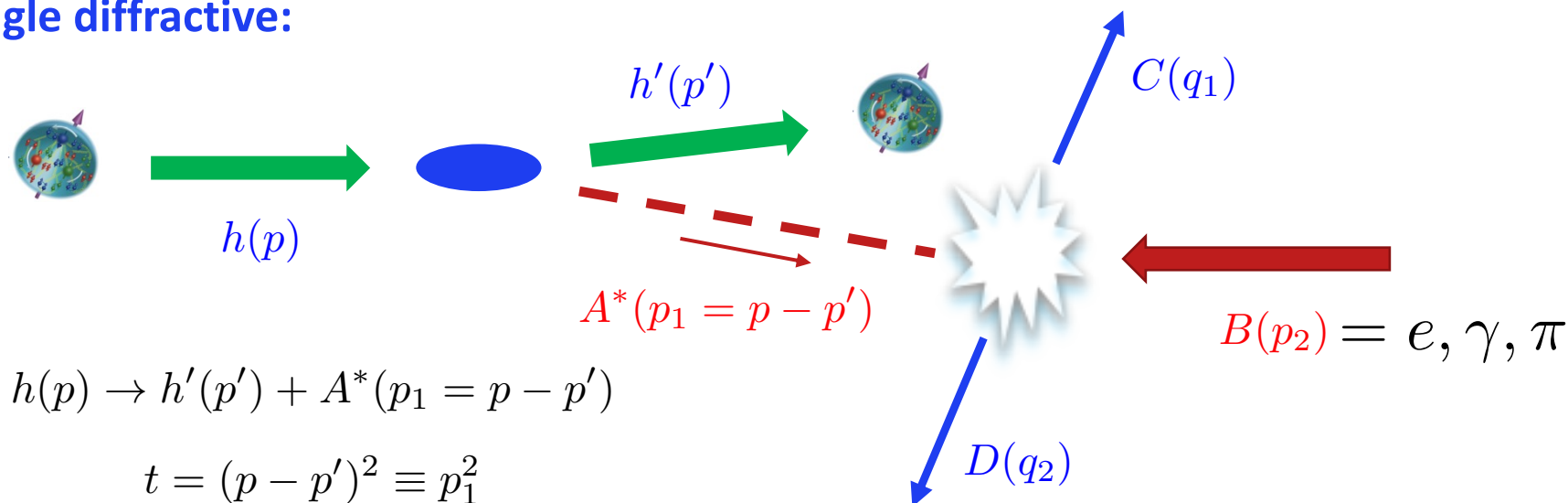
$$x \leftrightarrow q_T$$

Single-Diffractive Hard Exclusive Processes (SDHEP)

Qiu & Yu, JHEP 08 (2022) 103
PRD 107 (2023) 1

□ Two-stage diffractive $2 \rightarrow 3$ hard exclusive processes:

■ Single diffractive:



■ Hard probe: $2 \rightarrow 2$ high q_T exclusive process

$$A^*(p_1) + B(p_2) \rightarrow C(q_1) + D(q_2)$$

$$(p - p') \cdot n \gg \sqrt{|t|} \iff |q_{1T}| = |q_{2T}| \gg \sqrt{-t}$$

■ The single diffractive $2 \rightarrow 3$ exclusive hard processes:

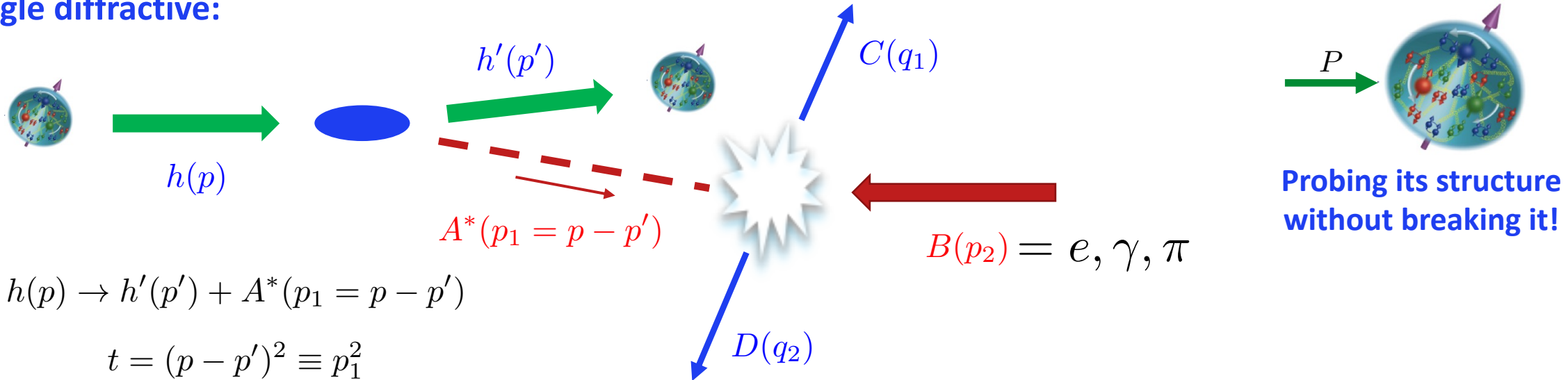
$$h(p) + B(p_2) \rightarrow h'(p') + C(q_1) + D(q_2)$$

Single-Diffractive Hard Exclusive Processes (SDHEP)

Qiu & Yu, JHEP 08 (2022) 103
PRD 107 (2023) 1

□ Two-stage diffractive $2 \rightarrow 3$ hard exclusive processes:

■ Single diffractive:



■ Hard probe: $2 \rightarrow 2$ high q_T exclusive process

■ Necessary condition for QCD factorization:

$$|q_{1T}| = |q_{2T}| \gg \sqrt{-t}$$

The state $A^*(p_1)$ lives much longer than $2 \rightarrow 2$ hard exclusive collision!

■ The single diffractive $2 \rightarrow 3$ exclusive hard processes:

$$h(p) + B(p_2) \rightarrow h'(p') + C(q_1) + D(q_2)$$

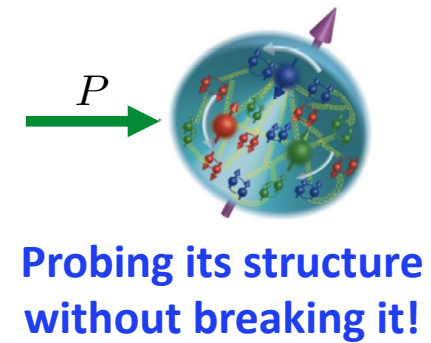
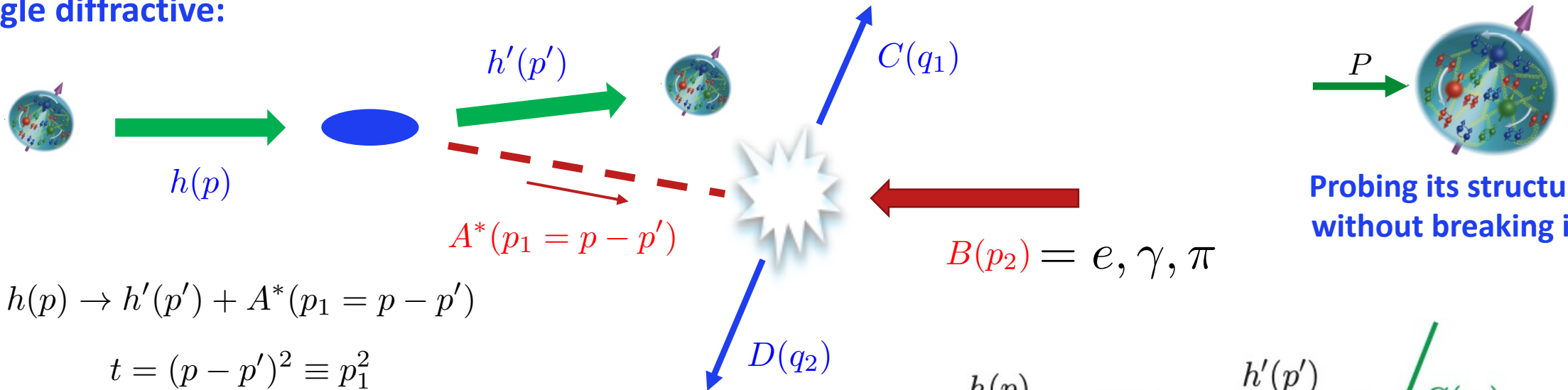
Not necessarily sufficient!

Single-Diffractive Hard Exclusive Processes (SDHEP)

Qiu & Yu, JHEP 08 (2022) 103
PRD 107 (2023) 1

Two-stage diffractive $2 \rightarrow 3$ hard exclusive processes:

Single diffractive:



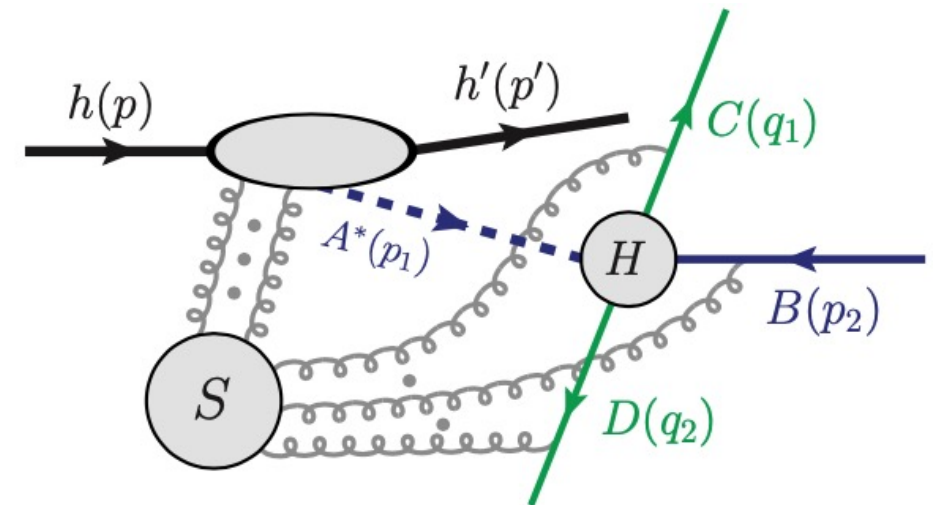
Hard probe: $2 \rightarrow 2$ high q_T exclusive process

$$A^*(p_1) + B(p_2) \rightarrow C(q_1) + D(q_2)$$

$$(p - p') \cdot n \gg \sqrt{|t|} \iff |q_{1T}| = |q_{2T}| \gg \sqrt{-t}$$

The single diffractive $2 \rightarrow 3$ exclusive hard processes:

$$h(p) + B(p_2) \rightarrow h'(p') + C(q_1) + D(q_2)$$

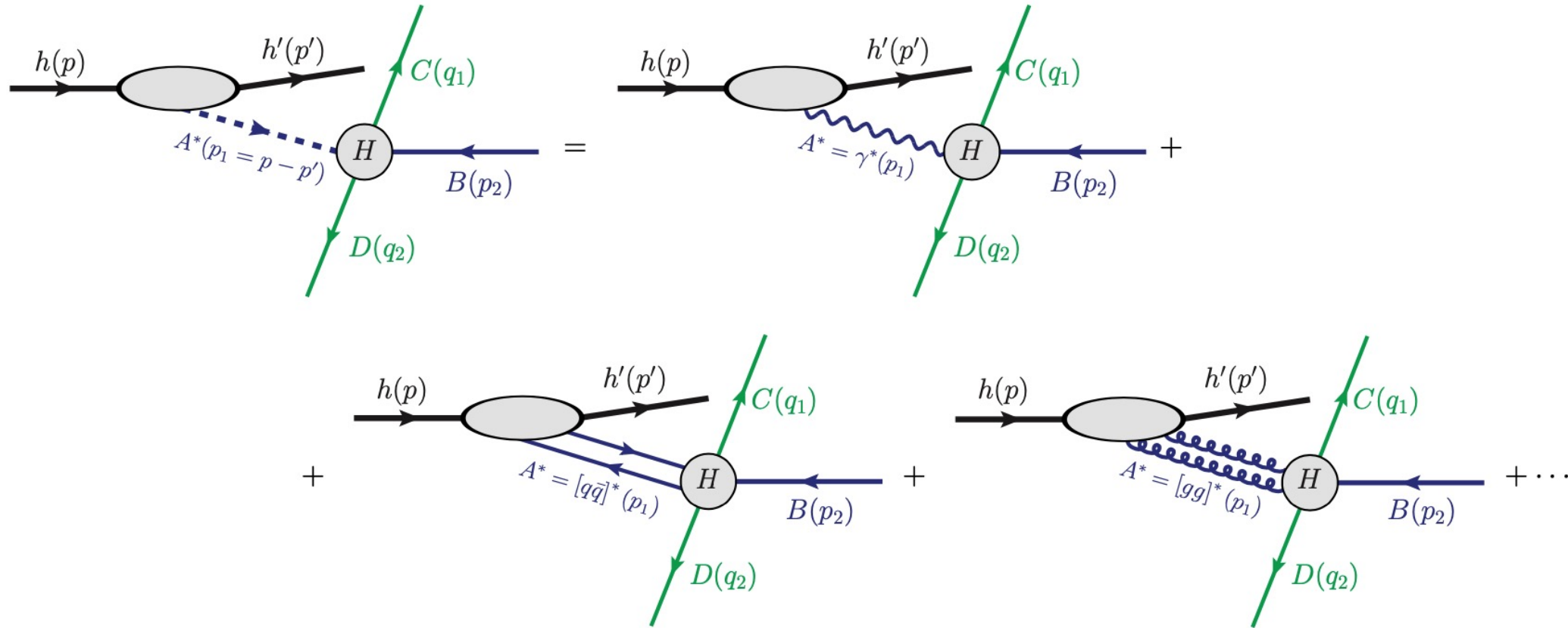


Factorizable if $|q_{1T}| = |q_{2T}| \gg \sqrt{-t}$

Single-Diffractive Hard Exclusive Processes (SDHEP)

Qiu & Yu, JHEP 08 (2022) 103
PRD 107 (2023) 1

□ Two-stage diffractive $2 \rightarrow 3$ hard exclusive processes:



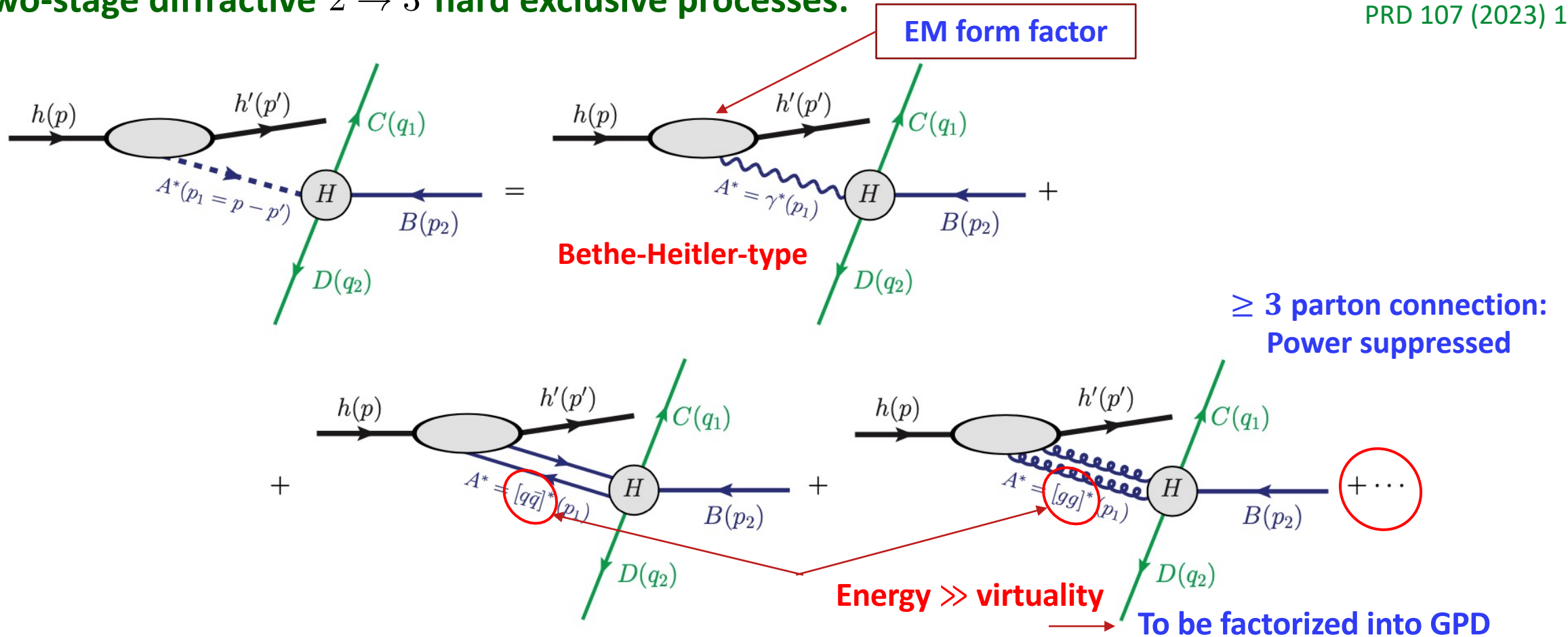
The exchanged state $A^*(p-p')$ is a sum of all possible partonic states, $\sum_{n=1,2,\dots}$, allowed by

- Quantum numbers of $h(p) - h'(p')$
- Symmetry of producing non-vanishing H

Single-Diffractive Hard Exclusive Processes (SDHEP)

Qiu & Yu, JHEP 08 (2022) 103
PRD 107 (2023) 1

□ Two-stage diffractive $2 \rightarrow 3$ hard exclusive processes:

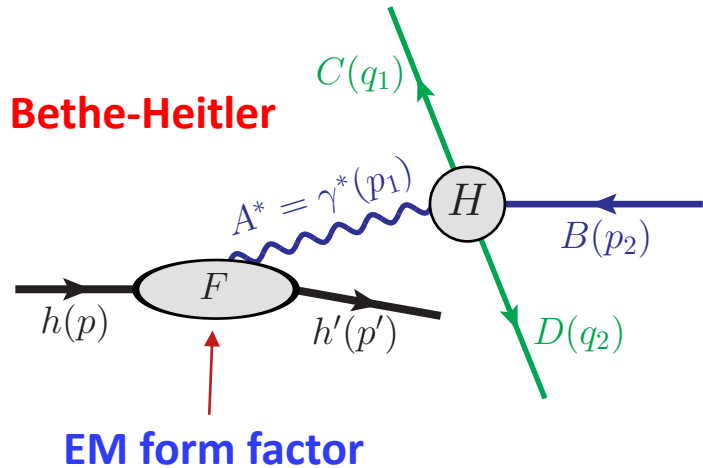


The exchanged state $A^*(p-p')$ is a sum of all possible partonic states, $\sum_{n=1,2,\dots}$, allowed by

- Quantum numbers of $h(p) - h'(p')$
- Symmetry of producing non-vanishing H

General Discussion on n=1 state: γ^*

Exchange of a virtual photon – “GPD background”:



$$\begin{aligned} \mathcal{M}^{(1)} &= \frac{ie^2}{t} \langle h'(p') | J^\mu(0) | h(p) \rangle \langle C(q_1) D(q_2) | J_\mu(0) | B(p_2) \rangle \\ &\equiv \frac{ie^2}{t} F^\mu(p, p') \mathcal{H}_\mu(p_1, p_2, q_1, q_2) \end{aligned}$$

Leading component

$$F^+ \mathcal{H}^- = \frac{1}{p_1^+} F^+ (p_1^+ \mathcal{H}^-) = \frac{1}{p_1^+} F^+ (p_1 \cdot \mathcal{H} + \mathbf{p}_{1\perp} \cdot \mathcal{H}_\perp - p_1^- \mathcal{H}^+) \sim \mathcal{O}(\sqrt{|t|})$$

$$\mathcal{M}^{(1)} \sim \mathcal{O}(1/\sqrt{|t|})$$

$$\mathcal{M}^{(2)} \sim \mathcal{O}(1/Q)$$

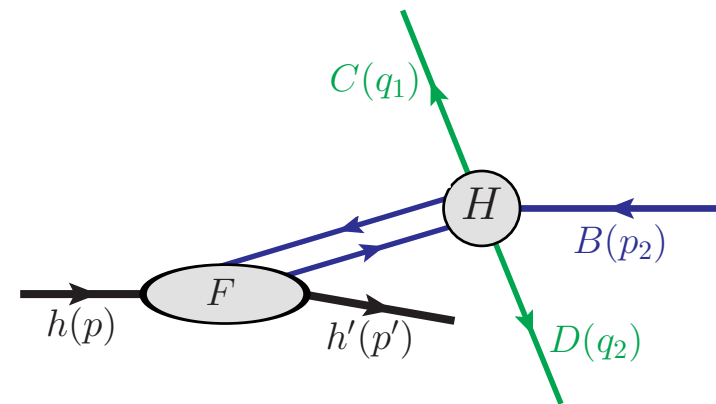


$$\mathcal{M}^{(1)}/\mathcal{M}^{(2)} \sim \mathcal{O}(Q/\sqrt{|t|})$$

γ^* channel is of a **more leading power** than GPD contribution, but higher power in α_{EM}

Generally allowed, except

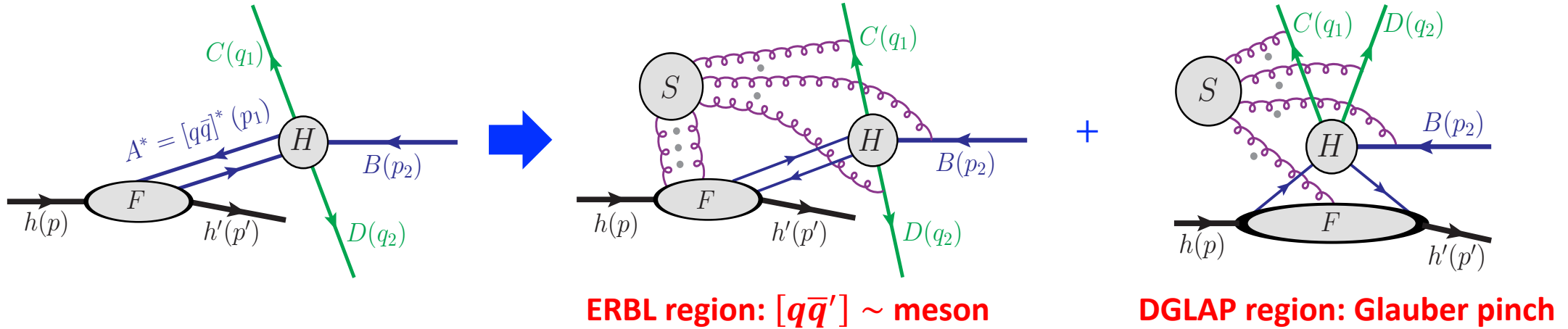
- (1) flavor changing ($p \rightarrow n, n \rightarrow p$, etc.)
- (2) forbidden by symmetry in the hard part



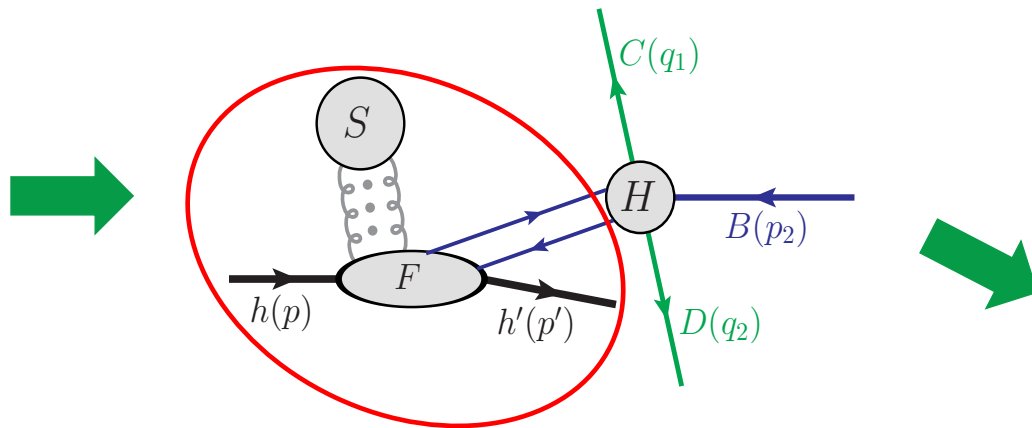
Factorization for SDHEP in the Two-stage Paradigm

Factorization for 2-parton channel factorization:

Qiu & Yu, JHEP 08 (2022) 103,
PRD 107 (2023) 1

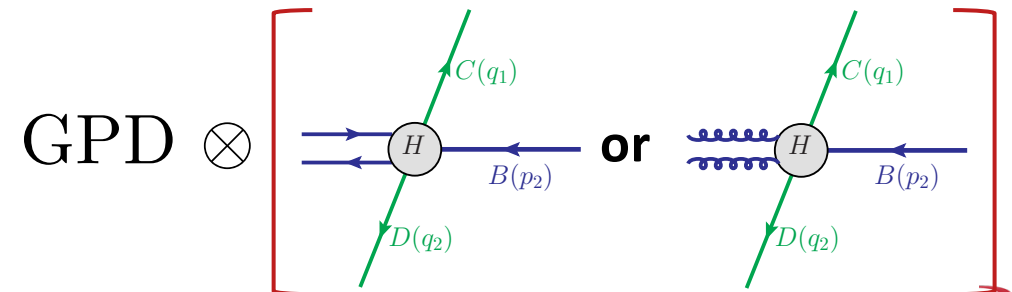


Soft gluons cancel when coupling to color neutral hadrons:



Glauber gluons of SDHEP:

$$k_s = (\lambda^2, \lambda^2, \lambda) \rightarrow (1, \lambda^2, \lambda) \quad \text{Collinear gluons}$$



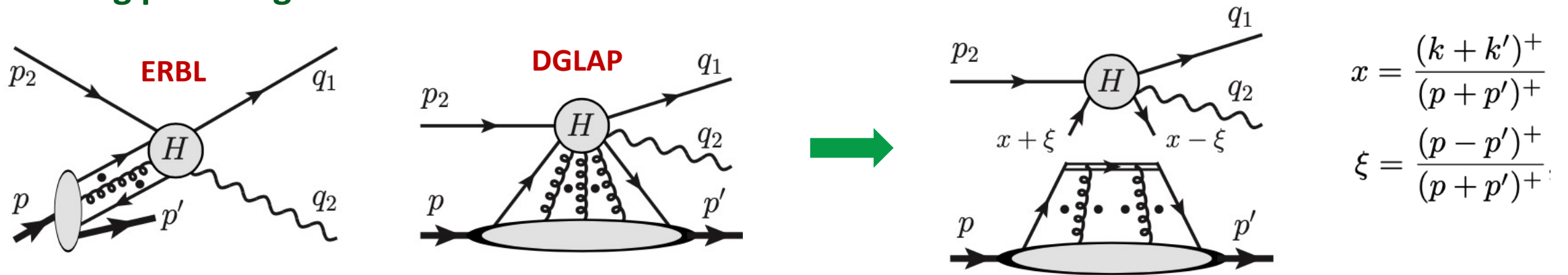
Hard probes Jefferson Lab

SDHEP with a Lepton Beam – JLab, EIC

□ DVCS:

$h(p) = \text{Proton}(p)$, $h'(p') = \text{Proton}(p')$, $B(p_2) = \text{electron}(p_2)$, $C(q_1) = \text{electron}(q_1)$, $D(q_2) = \text{photon}(q_2)$

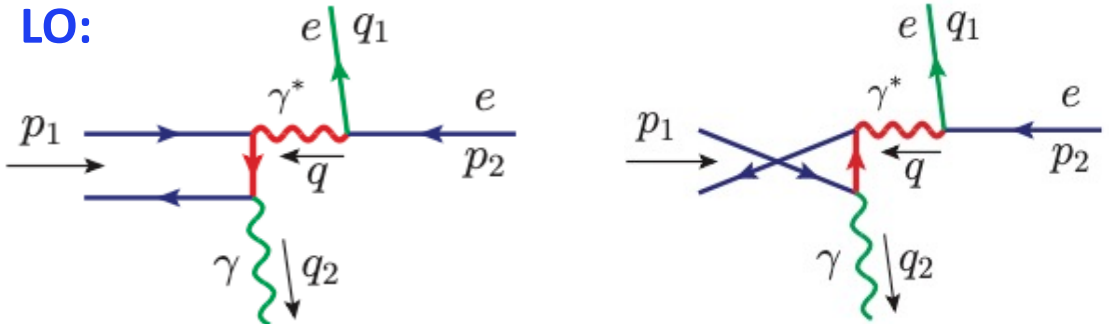
□ Leading pinch region:



□ Factorization formula:

$$\mathcal{M}_{he \rightarrow h'e\gamma}^{(2)} = \sum_i \int_{-1}^1 dx F_i^h(x, \xi, t) C_{ie \rightarrow e\gamma}(x, \xi, q_T),$$

➔ $C^{(0)} \propto \frac{1}{x - \xi + i\epsilon} - \frac{1}{x + \xi - i\epsilon}$



“moment-type” x-dependence

GPD Models for Numerical Results

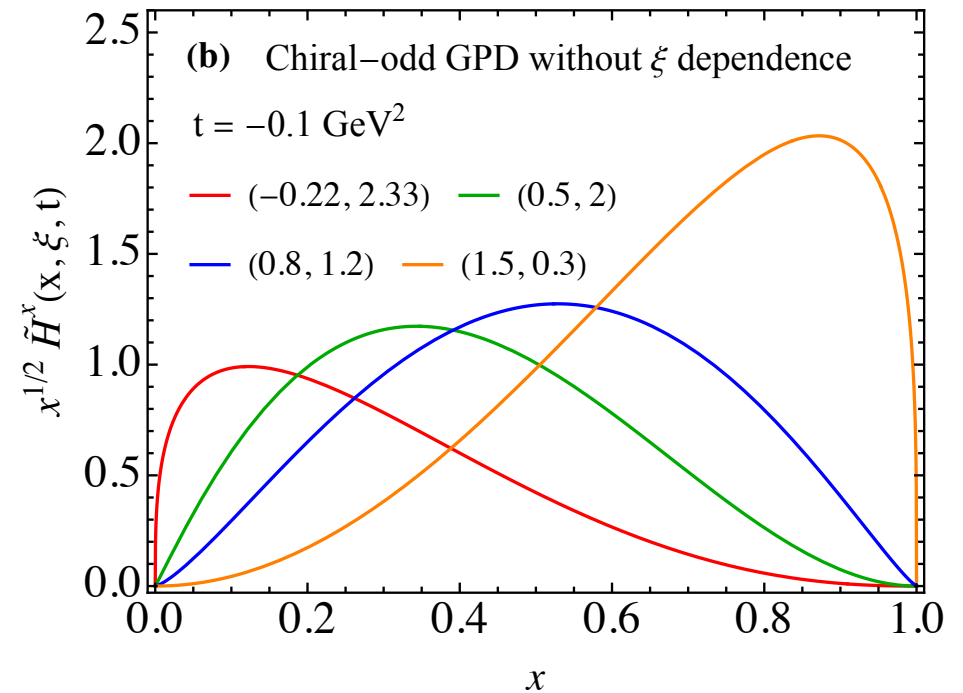
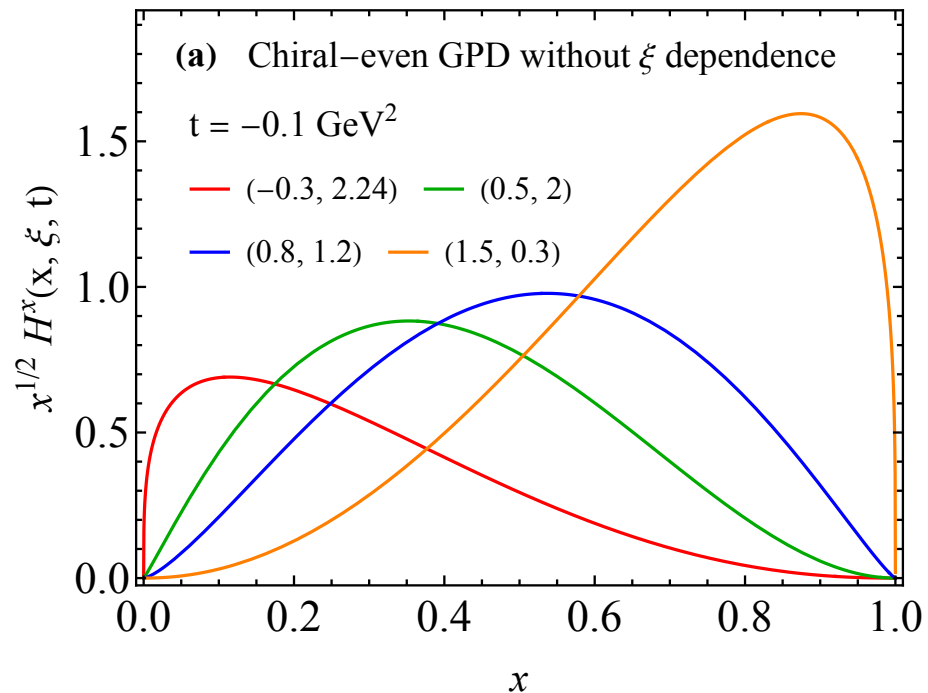
□ Simplified GK models:

$$H_{pn}(x, \xi, t) = \theta(x) x^{-0.9 (t/\text{GeV}^2)} \frac{x^\rho (1-x)^\tau}{B(1+\rho, 1+\tau)}$$

$$\tilde{H}_{pn}(x, \xi, t) = \theta(x) x^{-0.45 (t/\text{GeV}^2)} \frac{1.267 x^\rho (1-x)^\tau}{B(1+\rho, 1+\tau)}$$

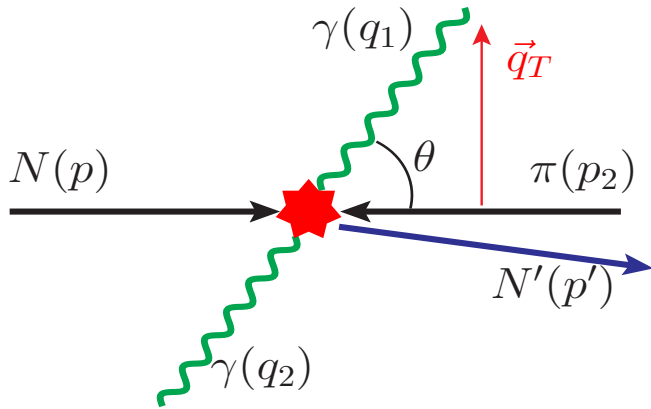
- Neglect E, \tilde{E} . Neglect evolution effect.
- Tune (ρ, τ) to control x shape.
- Fix DA: $D(z) = N z^{0.63} (1-z)^{0.63}$

Goloskokov, Kroll
hep-ph/0501242
arXiv: 0708.3569
arXiv: 0906.0460



Enhanced Sensitivity on x -dependence of GPDs

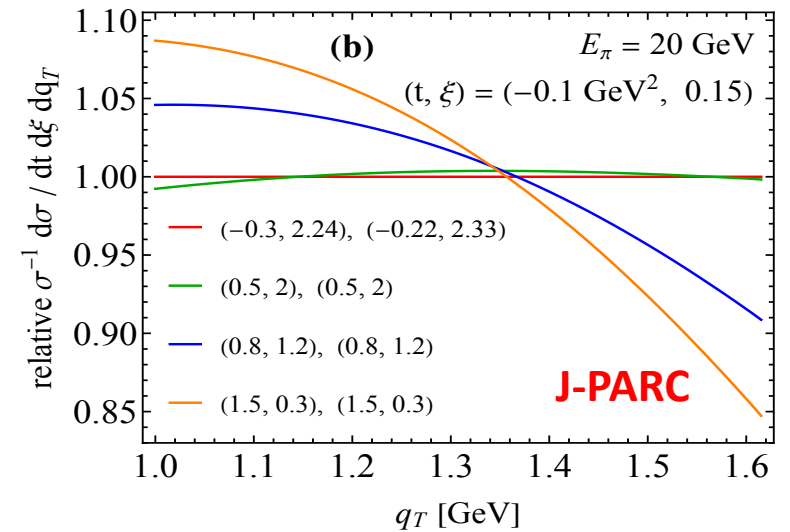
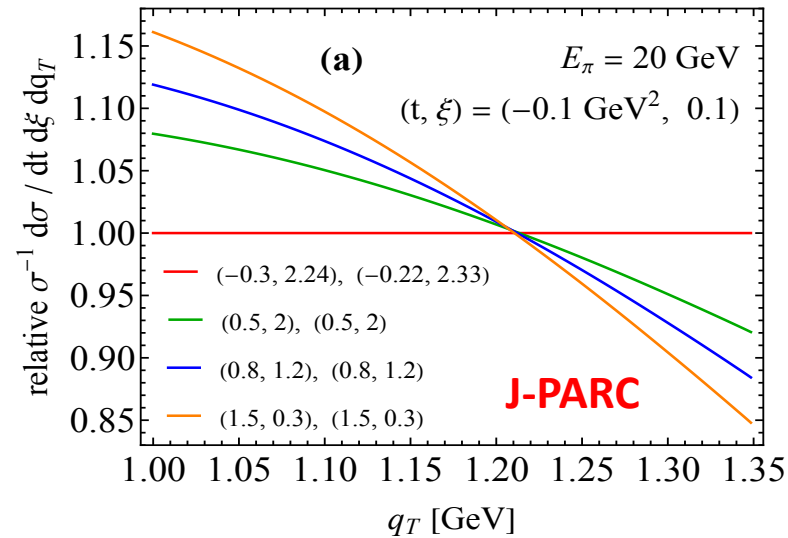
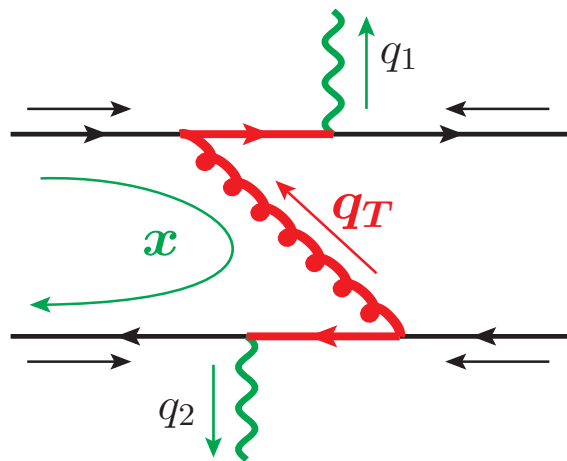
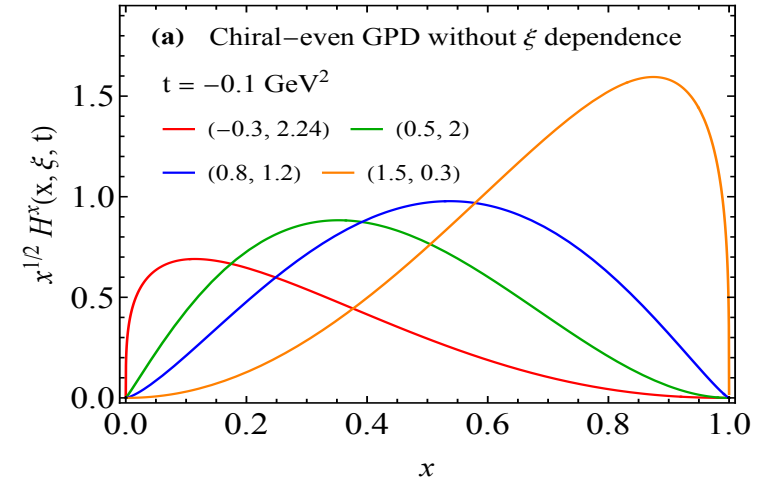
Two-photon production: $\pi^-(p_\pi) + P(p) \rightarrow \gamma(q_1) + \gamma(q_2) + N(p')$ J-PARC, COMPASS Qiu & Yu, JHEP 08 (2022) 103



Vary GPD x shapes

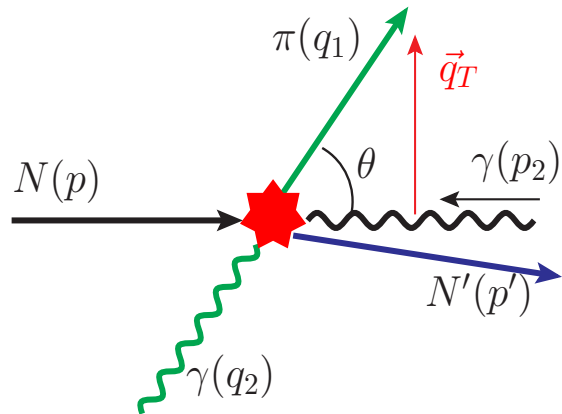


Different q_T shapes



Enhanced Sensitivity on x-dependence of GPDs

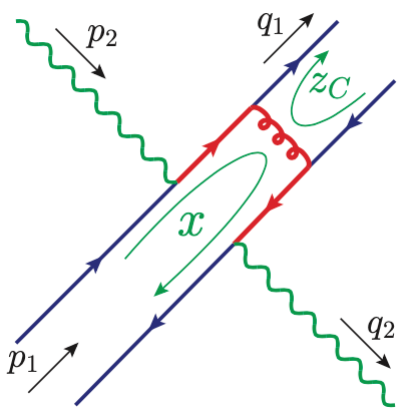
■ **Pion-photon production:** $\gamma(p_\gamma) + h(p) \rightarrow \pi^\pm(q_1) + \gamma(q_2) + h'(p')$
JLab-Hall D, other Halls with quasi-photon beam



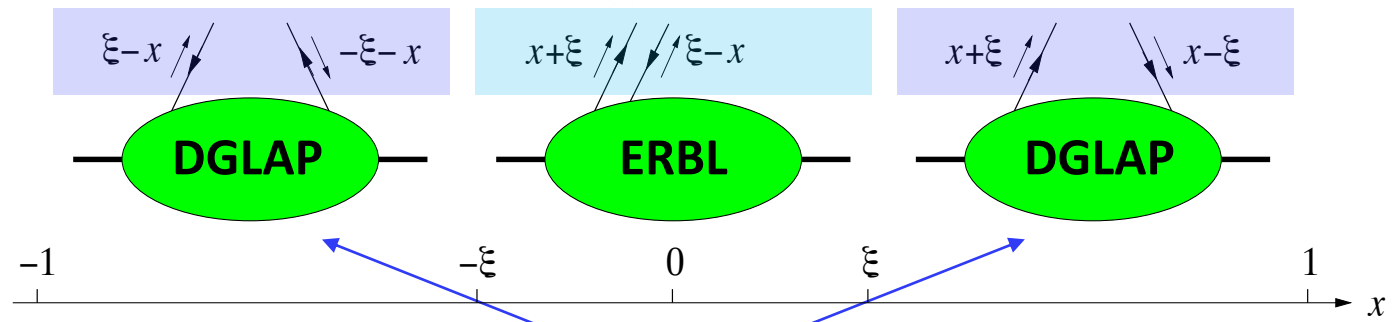
$i\mathcal{M}$ contains the entanglement between x and q_T

$$I'(t, \xi; z, \theta) = \int_{-1}^1 \frac{dx F(x, \xi, t)}{x - \rho'(z; \theta) + i\epsilon}$$

$$\rho'(z; \theta) = \xi \cdot \left[\frac{\cos^2(\theta/2) (1-z) - z}{\cos^2(\theta/2) (1-z) + z} \right] \in [-\xi, \xi]$$



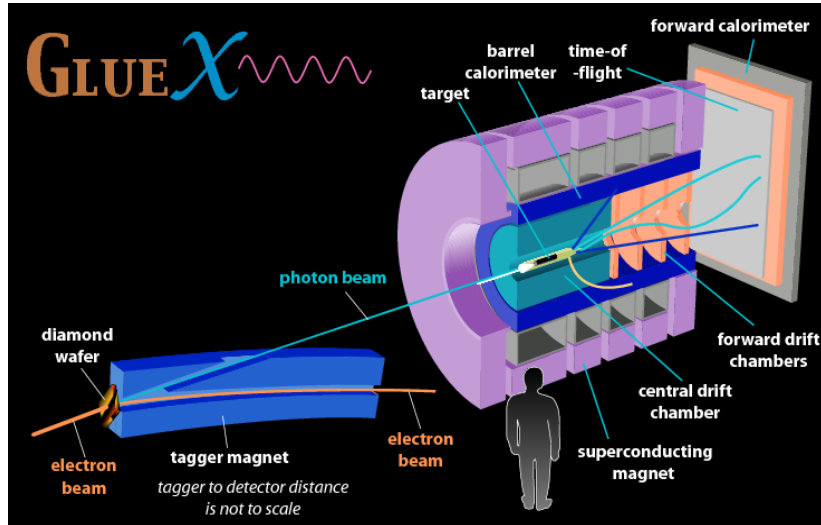
Complementary sensitivity:



$$N \pi \rightarrow N' \gamma \gamma$$

Exclusive Photo-Production of a $\pi\gamma$ Pair – Hall D at JLab

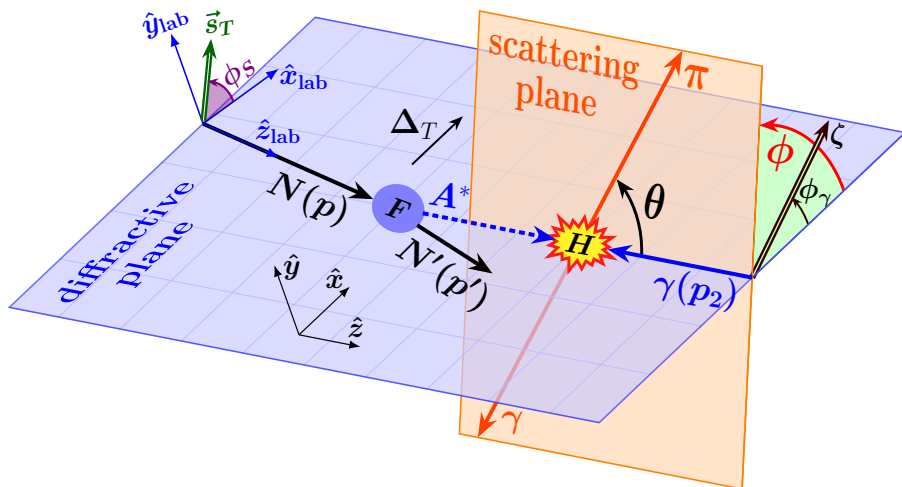
Qiu & Yu, arXiv:2305.15397



□ Polarization asymmetries

$$\frac{d\sigma}{d|t| d\xi d \cos \theta d\phi} = \frac{1}{2\pi} \frac{d\sigma}{d|t| d\xi d \cos \theta} \cdot [1 + \lambda_N \lambda_\gamma A_{LL} + \zeta A_{UT} \cos 2(\phi - \phi_\gamma) + \lambda_N \zeta A_{LT} \sin 2(\phi - \phi_\gamma)]$$

$$\frac{d\sigma}{d|t| d\xi d \cos \theta} = \pi (\alpha_e \alpha_s)^2 \left(\frac{C_F}{N_c} \right)^2 \frac{1 - \xi^2}{\xi^2 s^3} \Sigma_{UU}$$



$$\begin{aligned} \Sigma_{UU} &= |\mathcal{M}_+^{[\tilde{H}]}|^2 + |\mathcal{M}_-^{[\tilde{H}]}|^2 + |\tilde{\mathcal{M}}_+^{[H]}|^2 + |\tilde{\mathcal{M}}_-^{[H]}|^2, \\ A_{LL} &= 2 \Sigma_{UU}^{-1} \text{Re} \left[\mathcal{M}_+^{[\tilde{H}]} \tilde{\mathcal{M}}_+^{[H]*} + \mathcal{M}_-^{[\tilde{H}]} \tilde{\mathcal{M}}_-^{[H]*} \right], \\ A_{UT} &= 2 \Sigma_{UU}^{-1} \text{Re} \left[\tilde{\mathcal{M}}_+^{[H]} \tilde{\mathcal{M}}_-^{[H]*} - \mathcal{M}_+^{[\tilde{H}]} \mathcal{M}_-^{[\tilde{H}]*} \right], \\ A_{LT} &= 2 \Sigma_{UU}^{-1} \text{Im} \left[\mathcal{M}_+^{[\tilde{H}]} \tilde{\mathcal{M}}_-^{[H]*} + \mathcal{M}_-^{[\tilde{H}]} \tilde{\mathcal{M}}_+^{[H]*} \right]. \end{aligned}$$

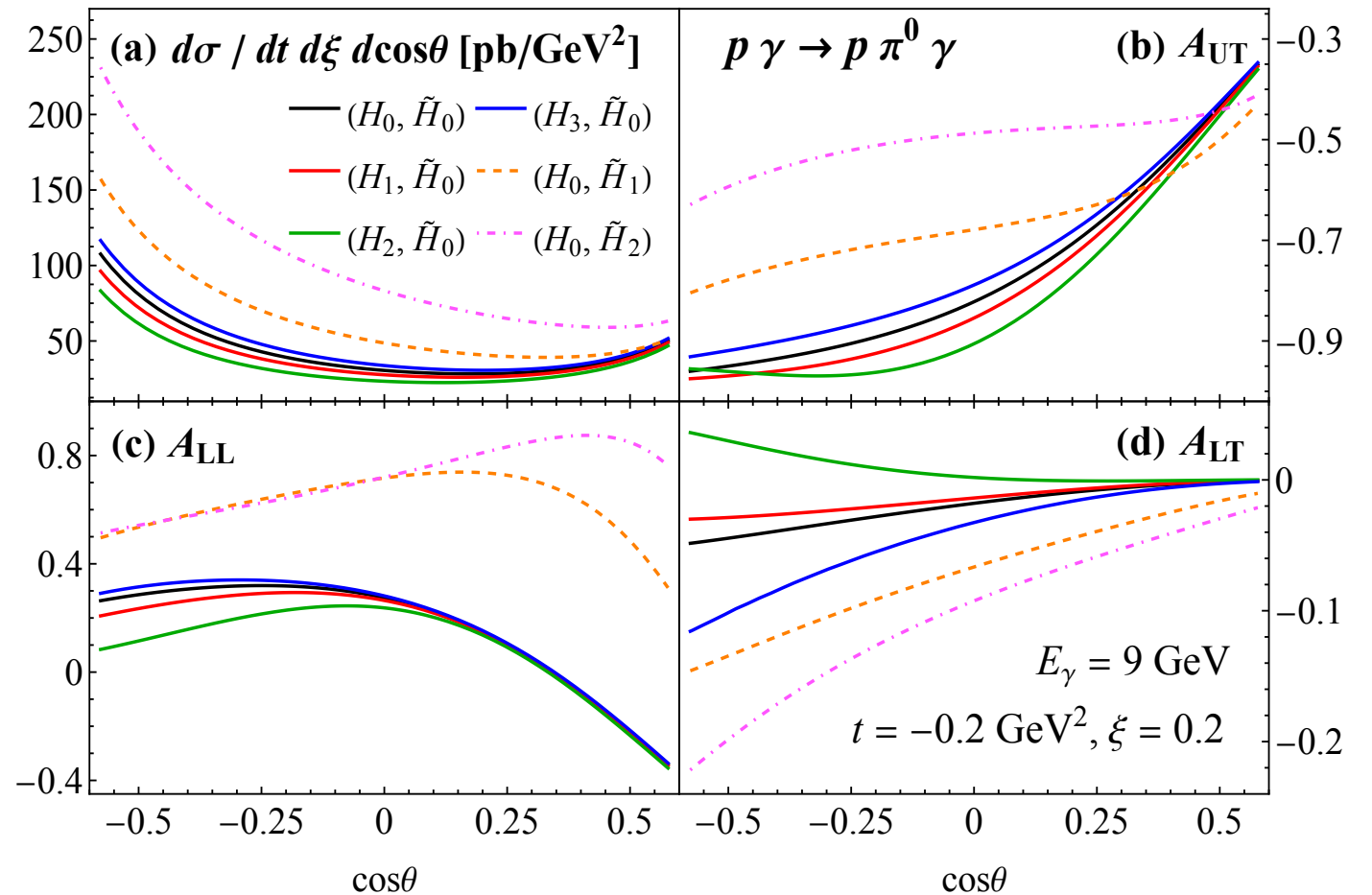
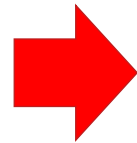
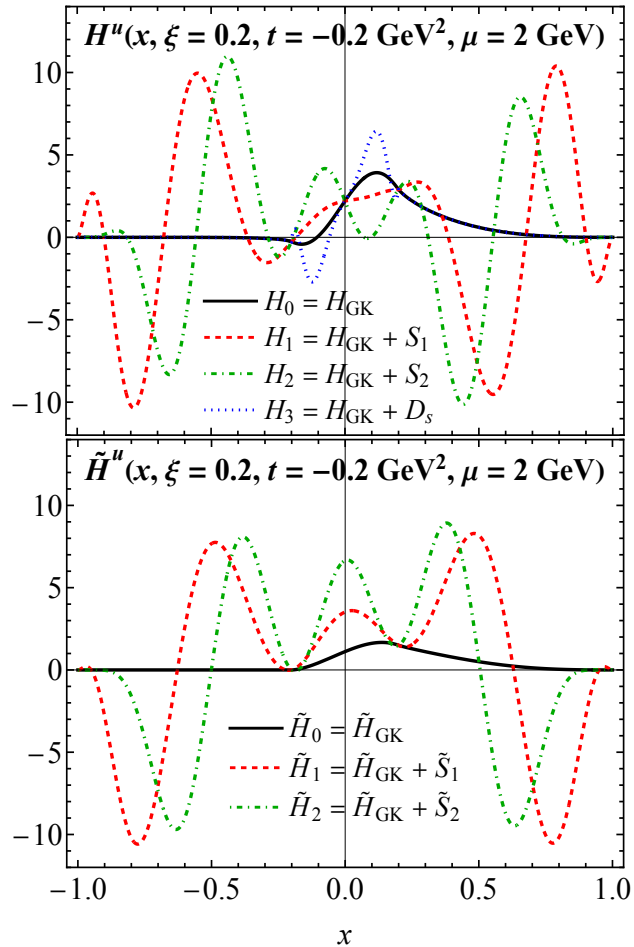
Exclusive Photo-Production of a $\pi\gamma$ Pair – Hall D at JLab

GPD Models:

= GK model + shadow GPDs

$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$

Goloskokov, Kroll, '05, '07, '09
Bertone et al. '21
Moffat et al. '23



Exclusive Photo-Production of a $\pi\gamma$ Pair – Hall D at JLab

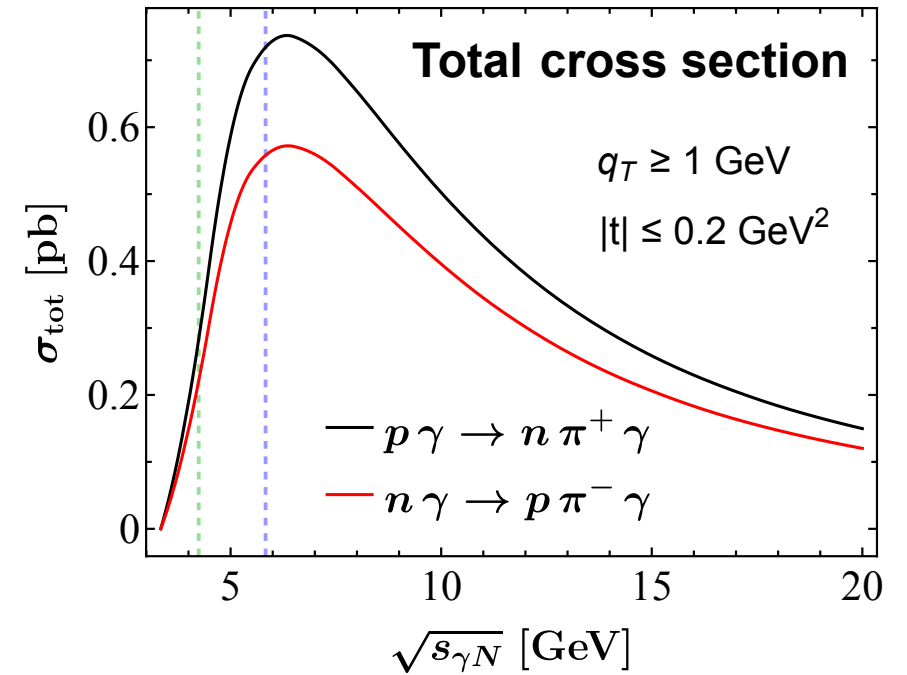
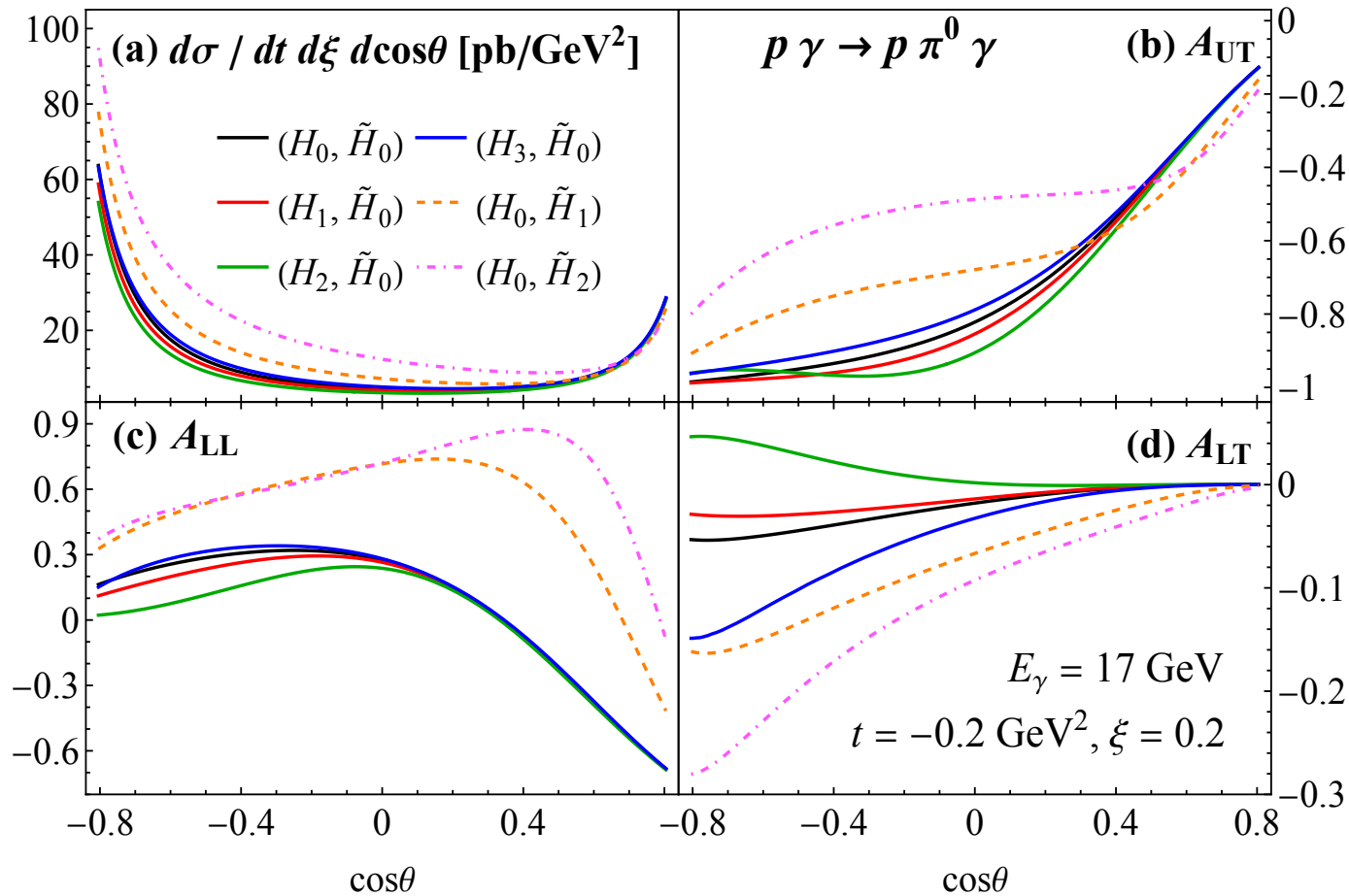
GPD Models:

= GK model + shadow GPDs



$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$

Goloskokov, Kroll, '05, '07, '09
Bertone et al. '21
Moffat et al. '23



Summary and Outlook

□ GPDs are fundamental parton correlation functions:

- Carry rich information on emergent hadron properties (mass, spin, ...) from QCD/parton dynamics
- Are responsible for the tomographic images of confined quarks and gluons inside a bound hadron
- Provide the much needed hints on how confined quarks/gluons respond to the probing scale, ...

Extracting their x -dependence from experimental observable(s) is non-trivial, but, full of opportunities, ...

□ Introduced the single diffractive $2 \rightarrow 3$ hard exclusive processes (SDHEP) for extracting GPDs, ...

- Covered all existing/known processes for extracting GPDs, plus ideas for new observables, ...
- Introduced a path forward to identify new SDHEPs that could be sensitive to x -dependence of GPDs
- Angular modulation between diffractive plane and hard scattering plane could provide unique opportunity to separate various GPDs
- Exclusive photoproduction at JLab provides excellent opportunities for extracting GPDs & x -dependence

Renewed opportunities for exploring the physics of GPDs and the confined phenomena of QCD.

Thanks!